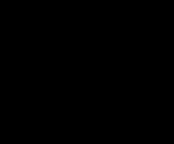
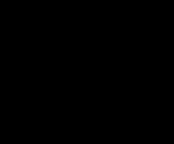


# *Neutrinos as Pathfinders*

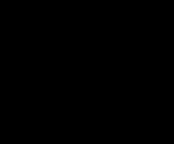
José W F Valle



VNIVERSITAT  
DE VALÈNCIA

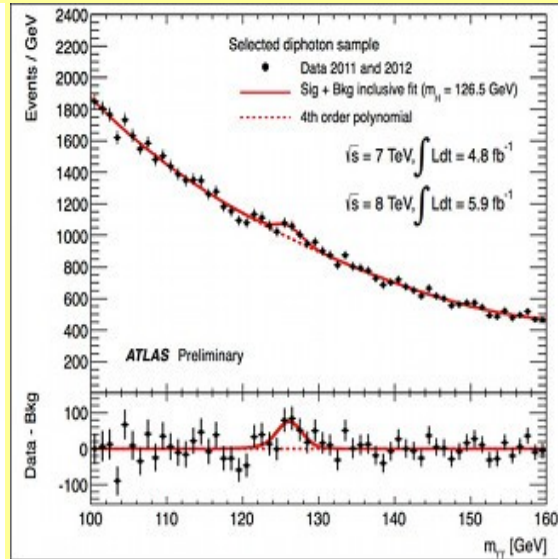
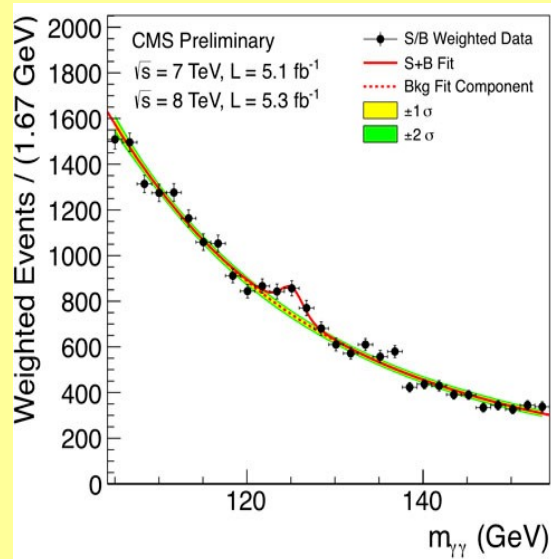


IFIC AHEP on facebook



November 2015

# HISTORIC DISCOVERY 1



# HISTORIC DISCOVERY 1



## Physics 2013



Photo: Pnicolet via Wikimedia Commons  
**François Englert**

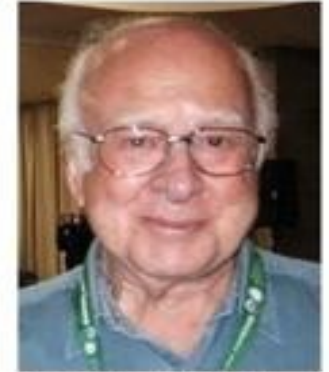
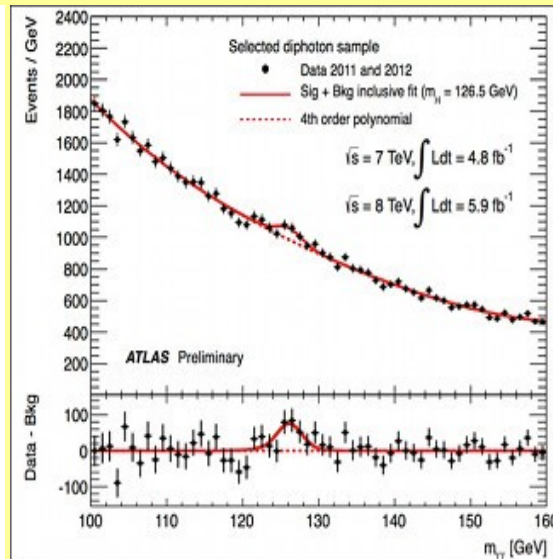
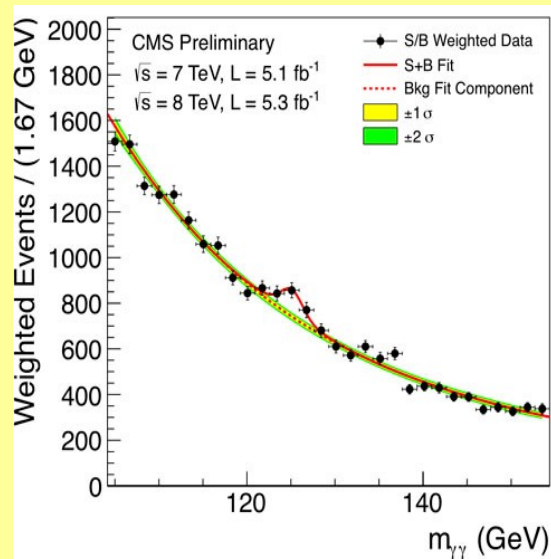


Photo: G-M Greuel via Wikimedia Commons  
**Peter W. Higgs**



Last stone ...

# HISTORIC DISCOVERY 2

PHYSICAL REVIEW D 90, 093006 (2014)

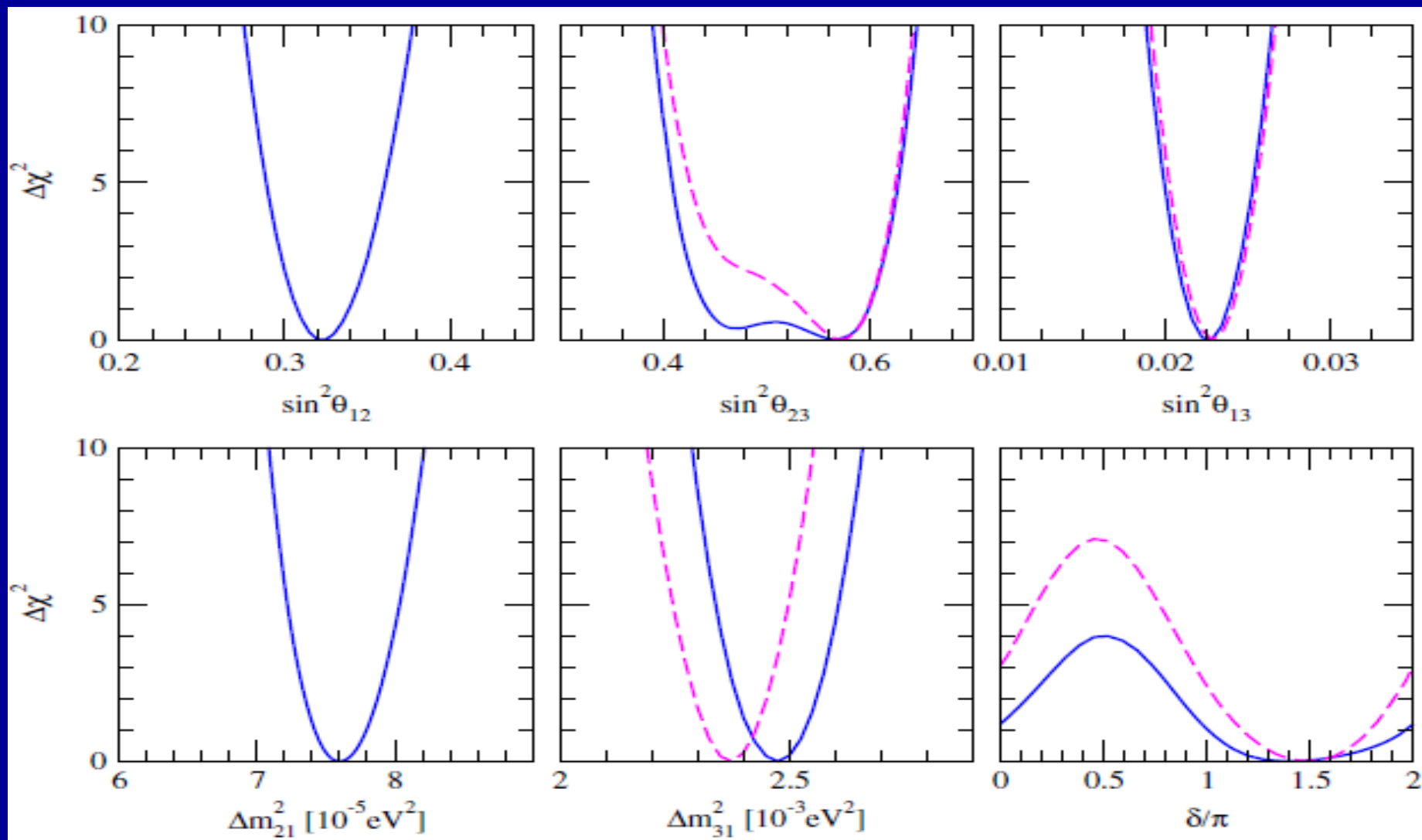
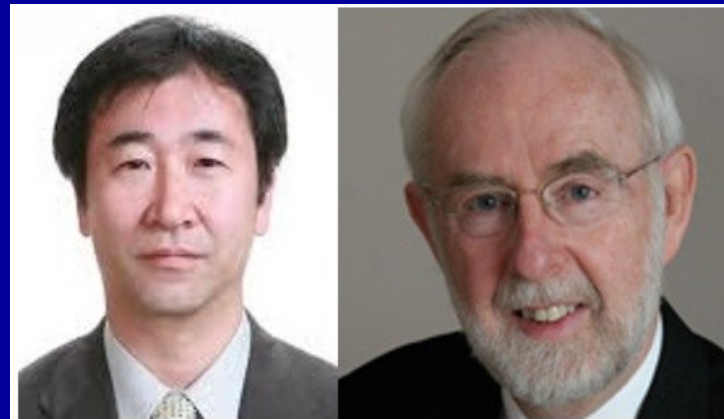
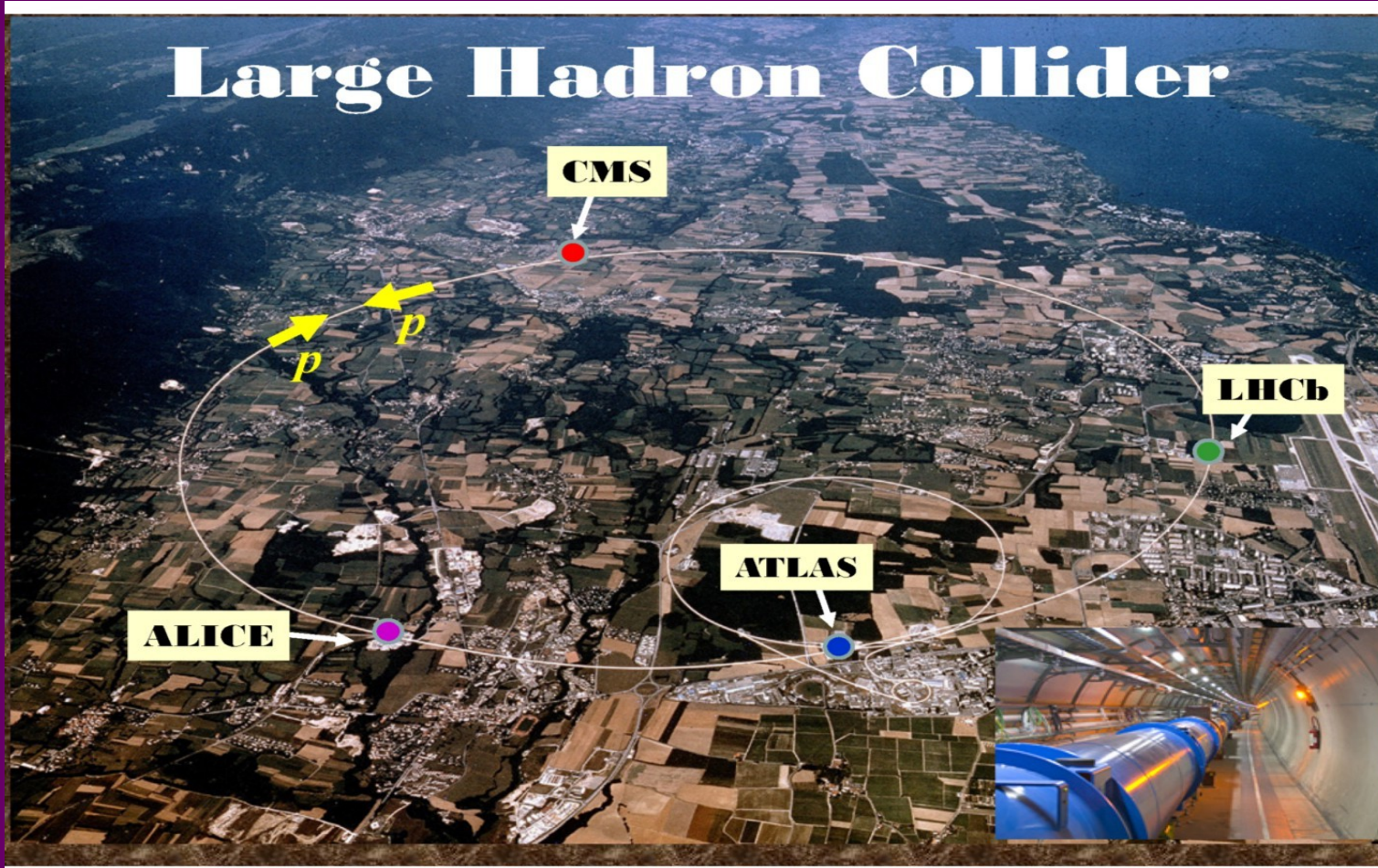


TABLE II. Neutrino oscillation parameters summary from the global analysis updated after Neutrino 2014 conference.

Parameter	Best fit $\pm 1\sigma$	$2\sigma$ range	$3\sigma$ range
$\Delta m_{21}^2 [10^{-5} \text{ eV}^2]$	$7.60^{+0.19}_{-0.18}$	7.26–7.99	7.11–8.18
$ \Delta m_{31}^2  [10^{-3} \text{ eV}^2]$ (NH)	$2.48^{+0.05}_{-0.07}$	2.35–2.59	2.30–2.65
$ \Delta m_{31}^2  [10^{-3} \text{ eV}^2]$ (IH)	$2.38^{+0.05}_{-0.06}$	2.26–2.48	2.20–2.54
$\sin^2 \theta_{12} / 10^{-1}$	$3.23 \pm 0.16$	2.92–3.57	2.78–3.75
$\theta_{12} / ^\circ$	$34.6 \pm 1.0$	32.7–36.7	31.8–37.8
$\sin^2 \theta_{23} / 10^{-1}$ (NH)	$5.67^{+0.32\text{a}}_{-1.24}$	4.14–6.23	3.93–6.43
$\theta_{23} / ^\circ$	$48.9^{+1.8}_{-7.2}$	40.0–52.1	38.8–53.3
$\sin^2 \theta_{23} / 10^{-1}$ (IH)	$5.73^{+0.25}_{-0.39}$	4.35–6.21	4.03–6.40
$\theta_{23} / ^\circ$	$49.2^{+1.5}_{-2.3}$	41.3–52.0	39.4–53.1
$\sin^2 \theta_{13} / 10^{-2}$ (NH)	$2.26 \pm 0.12$	2.02–2.50	1.90–2.62
$\theta_{13} / ^\circ$	$8.6^{+0.3}_{-0.2}$	8.2–9.1	7.9–9.3
$\sin^2 \theta_{13} / 10^{-2}$ (IH)	$2.29 \pm 0.12$	2.05–2.52	1.93–2.65
$\theta_{13} / ^\circ$	$8.7 \pm 0.2$	8.2–9.1	8.0–9.4
$\delta / \pi$ (NH)	$1.41^{+0.55}_{-0.40}$	0.0–0.2.0	0.0–2.0
$\delta / ^\circ$	$254^{+99}_{-72}$	0–360	0–360
$\delta / \pi$ (IH)	$1.48 \pm 0.31$	0.00–0.09 & 0.86–2.0	0.0–2.0
$\delta / ^\circ$	$266 \pm 56$	0–16 & 155–360	0–360

<sup>a</sup>There is a local minimum in the first octant, at  $\sin^2 \theta_{23} = 0.473$  with  $\Delta\chi^2 = 0.36$  with respect to the global minimum

# HISTORIC DISCOVERY 3 ???



**Barring confirmation  
or other surprises ...**

JWF Valle



**neutrino  
masses**

**consistency of vacuum  
flavour  
coupling unification  
gravity**

**dark matter  
baryon asymmetry  
Inflation**

**If so ...**



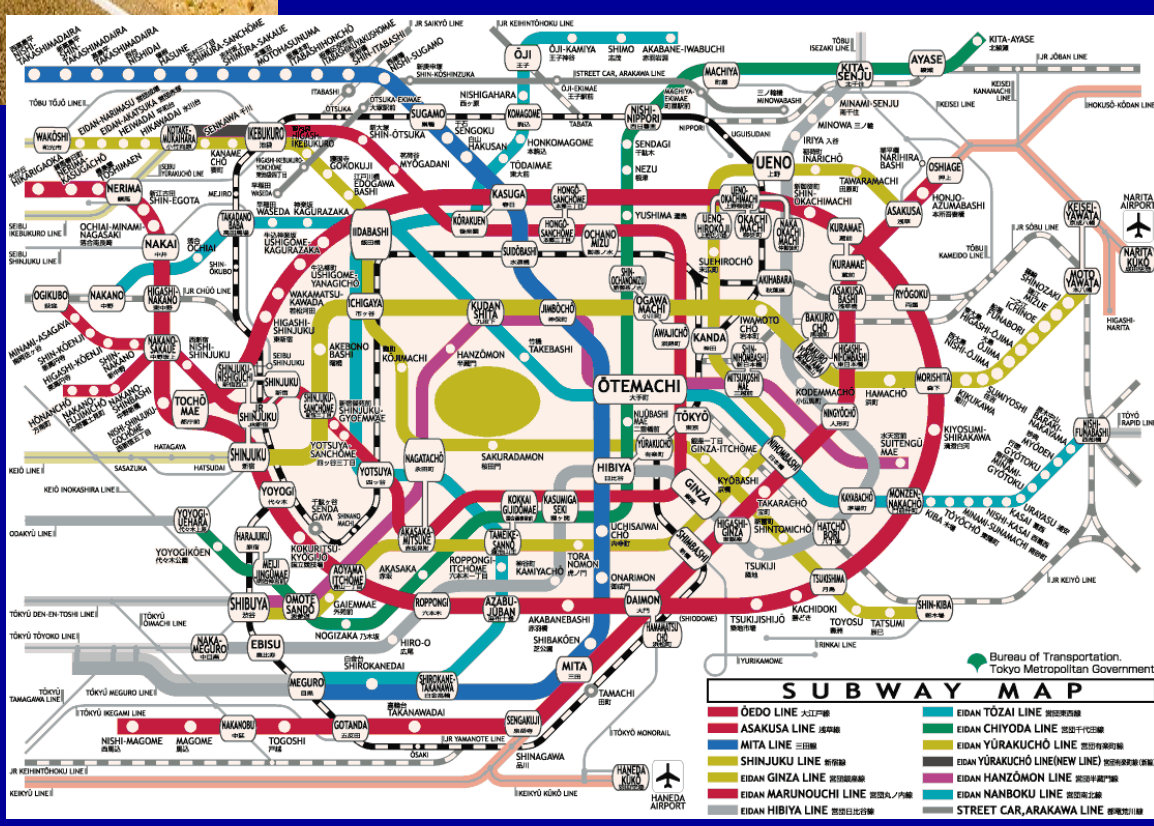
*It's a loooong*





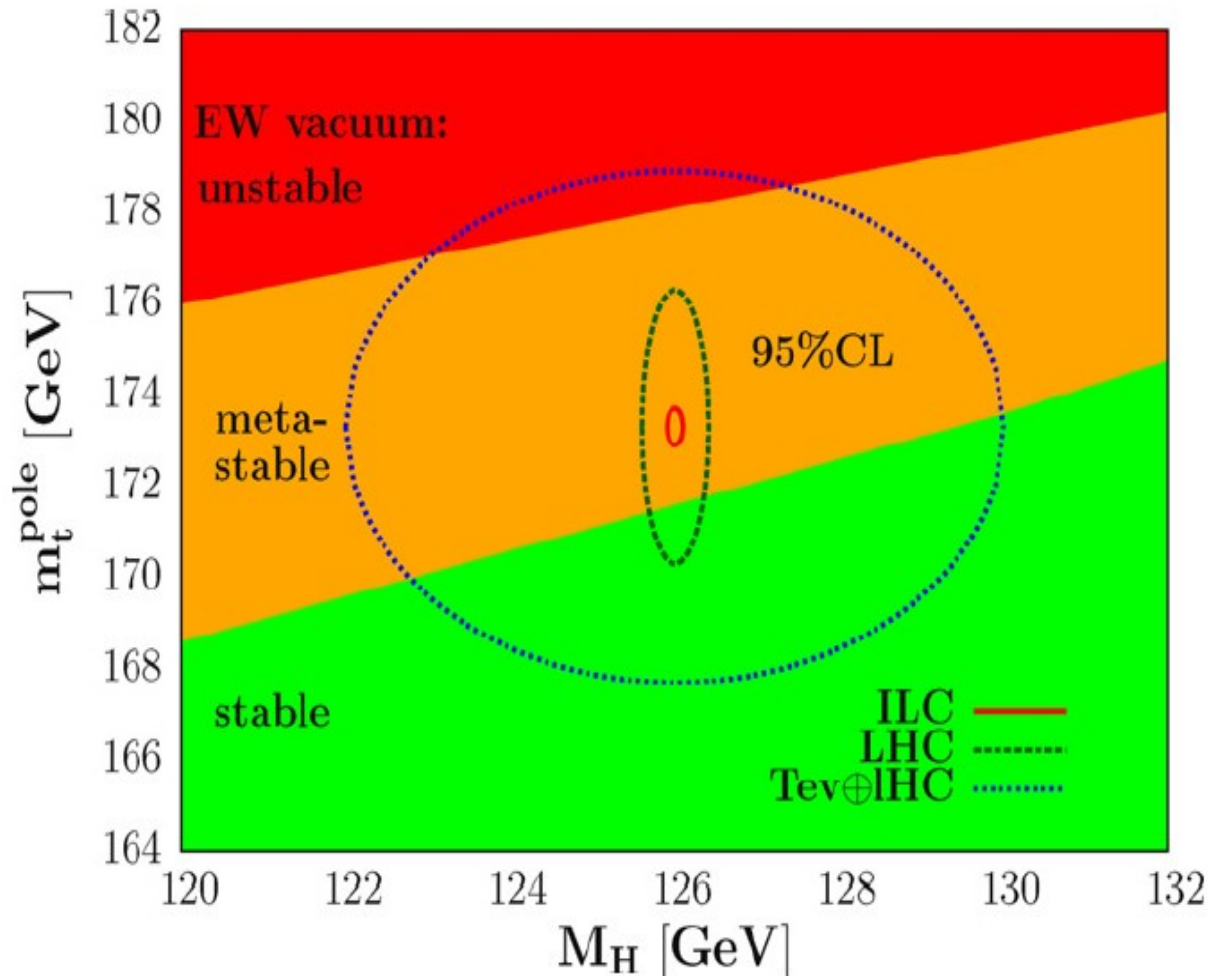
*It's a loooong*

*And winding  
Road...*



# SM vacuum ...

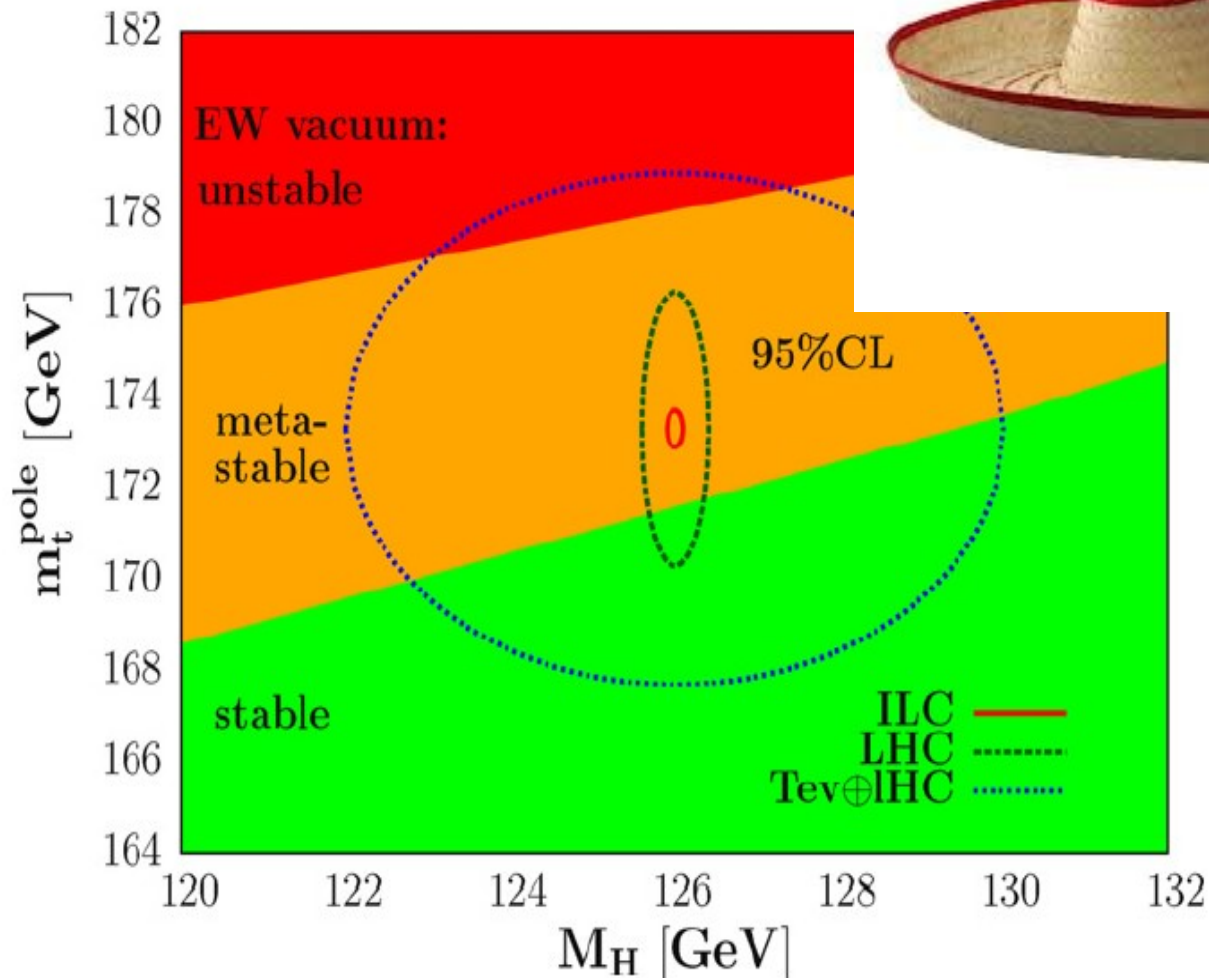
Physics Letters B 716 (2012) 214–219



**Fig. 1.** The  $2\sigma$  ellipses in the  $[M_H, m_t^{\text{pole}}]$  plane that one obtains from the current top quark and Higgs mass measurements at the Tevatron and LHC and which can be expected in future measurements at the LHC and at the ILC, when confronted with the areas in which the SM vacuum is absolutely stable, metastable and unstable up to the Planck scale.

# SM vacuum ...

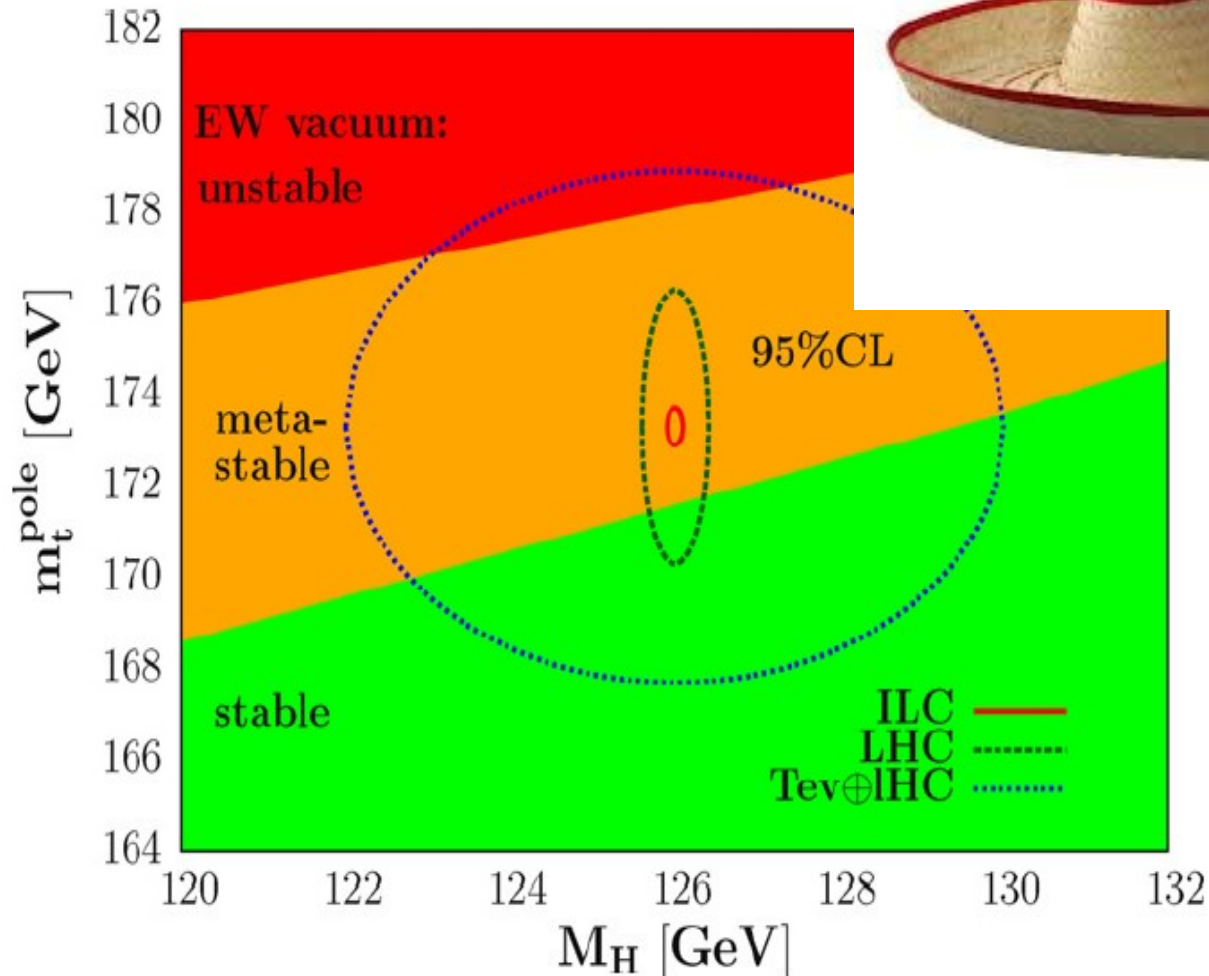
Physics Letters B 716 (2012) 214–219



**Fig. 1.** The  $2\sigma$  ellipses in the  $[M_H, m_t^{\text{pole}}]$  plane that one obtains from the current top quark and Higgs mass measurements at the Tevatron and LHC and which can be expected in future measurements at the LHC and at the ILC, when confronted with the areas in which the SM vacuum is absolutely stable, metastable and unstable up to the Planck scale.

# SM vacuum ...

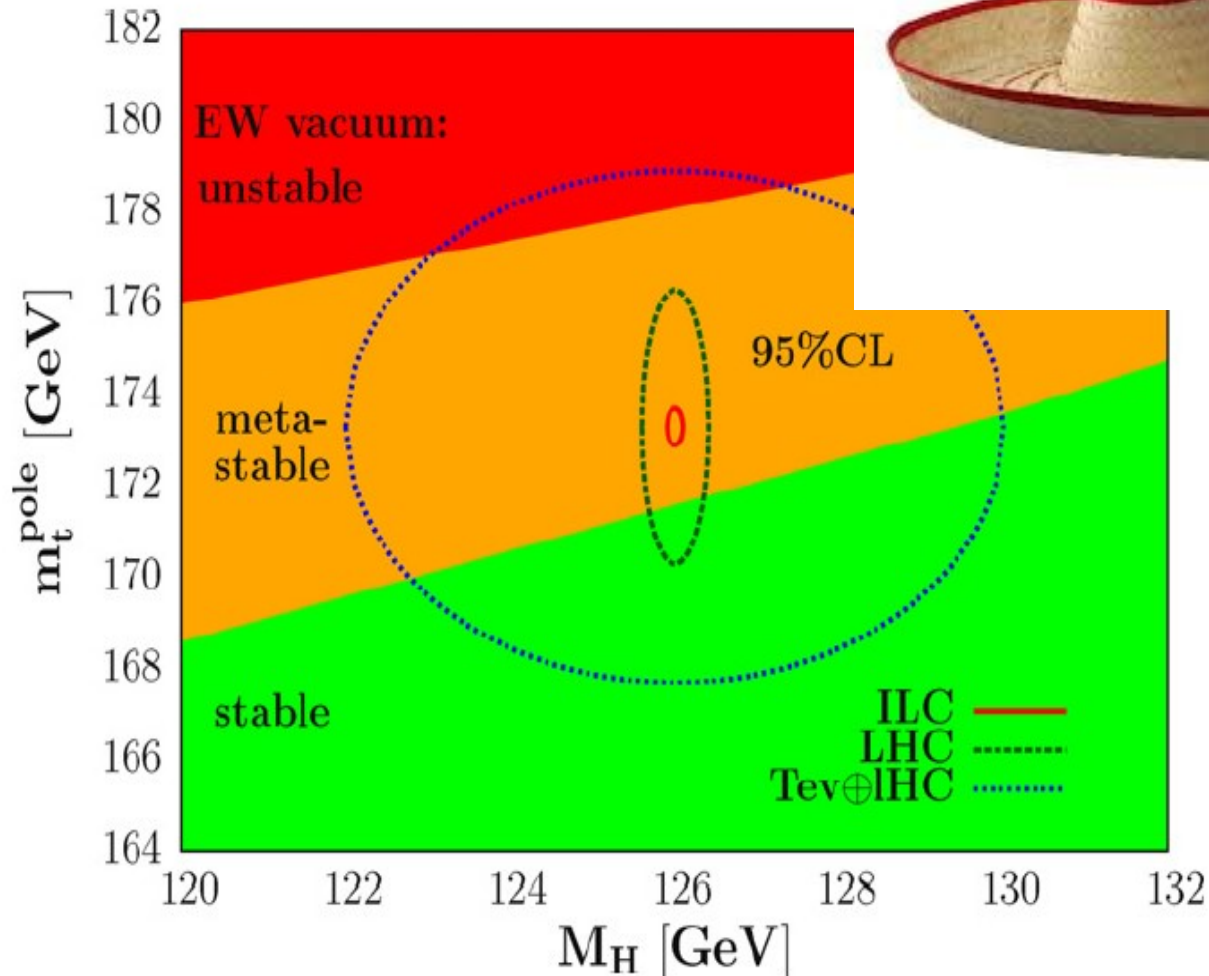
Physics Letters B 716 (2012) 214–219



**Fig. 1.** The  $2\sigma$  ellipses in the  $[M_H, m_t^{\text{pole}}]$  plane that one obtains from the current top quark and Higgs mass measurements at the Tevatron and LHC and which can be expected in future measurements at the LHC and at the ILC, when confronted with the areas in which the SM vacuum is absolutely stable, metastable and unstable up to the Planck scale.

# SM vacuum ...

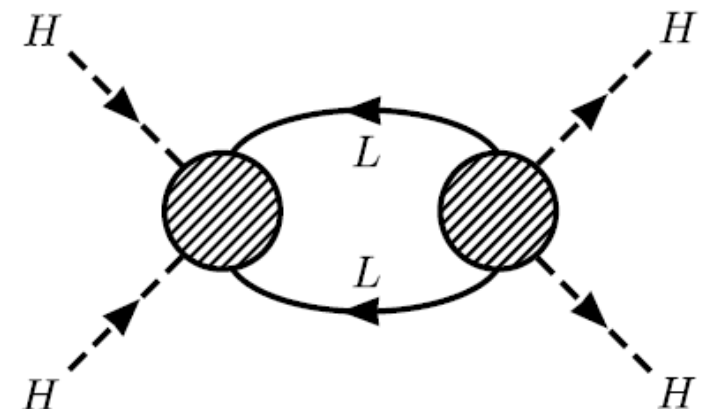
Physics Letters B 716 (2012) 214–219



## Neutrinos & Stability?

arXiv:1506.04031

Fig. 1. The  $2\sigma$  ellipses in the  $[M_H, m_t^{\text{pole}}]$  plane that one obtains from the current top quark and Higgs mass measurements at the Tevatron and LHC and which can be expected in future measurements at the LHC and at the ILC, when confronted with the areas in which the SM vacuum is absolutely stable, metastable and unstable up to the Planck scale.

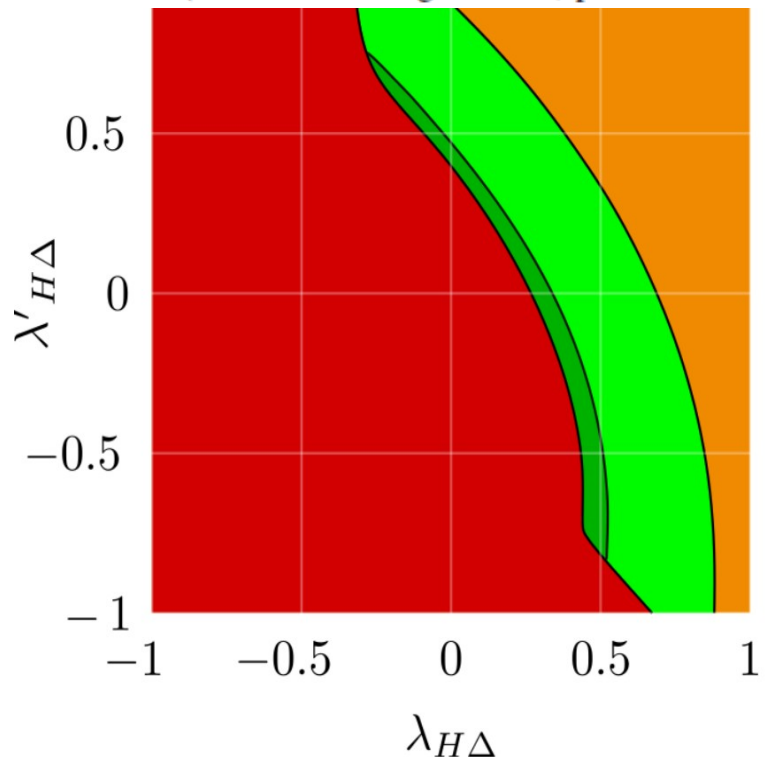


## Consistency of the triplet seesaw model revisited

Cesar Bonilla,<sup>\*</sup> Renato M. Fonseca,<sup>†</sup> and J. W. F. Valle<sup>‡</sup>

*AHEP Group, Instituto de Física Corpuscular, C.S.I.C./Universitat de València,  
Edificio Institutos de Investigación, Apartado 22085, E-46071 Valencia, Spain*

(Received 17 August 2015; published 21 October 2015)



arXiv:1508.02323

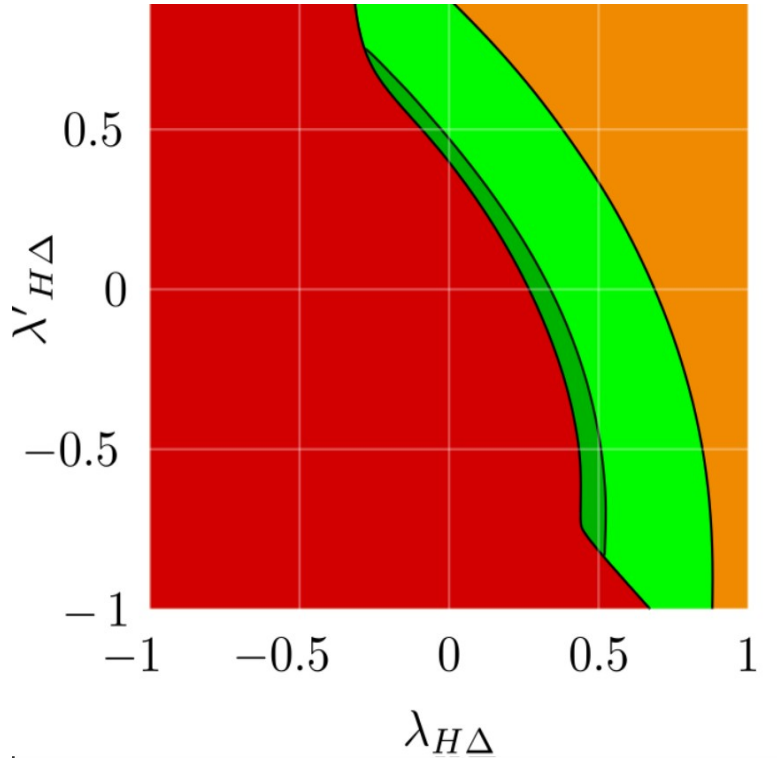
In addition to SM gauge invariance  
must break lepton number to give  
Masses to neutrinos ..

# Consistency of the triplet seesaw model revisited

Cesar Bonilla,<sup>\*</sup> Renato M. Fonseca,<sup>†</sup> and J. W. F. Valle<sup>‡</sup>

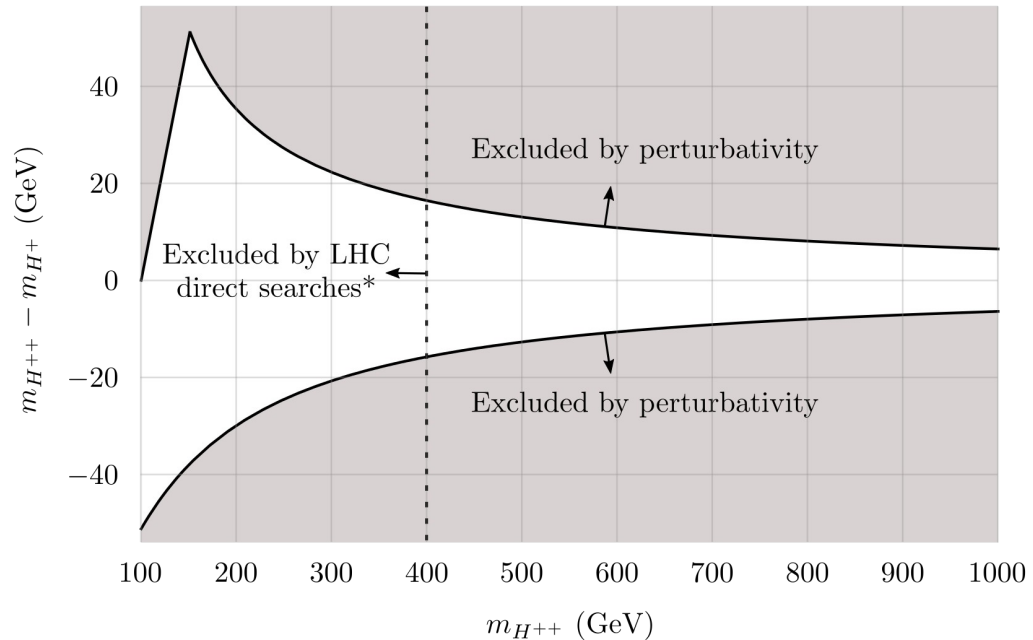
*AHEP Group, Instituto de Física Corpuscular, C.S.I.C./Universitat de València, Edificio Institutos de Investigación, Apartado 22085, E-46071 Valencia, Spain*

(Received 17 August 2015; published 21 October 2015)



arXiv:1508.02323

In addition to SM gauge invariance must break lepton number to give Masses to neutrinos ..

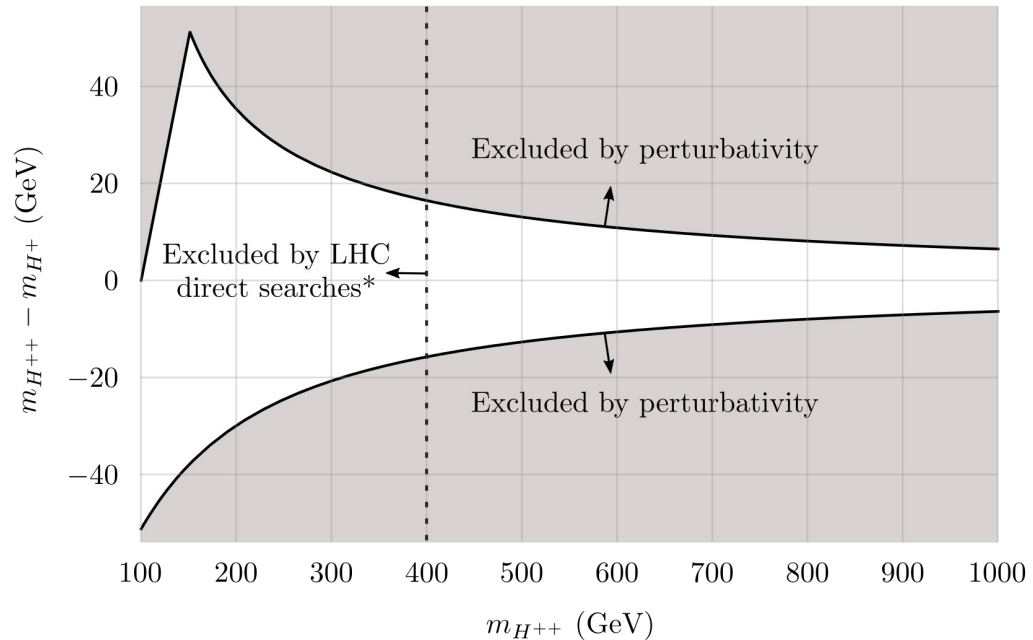
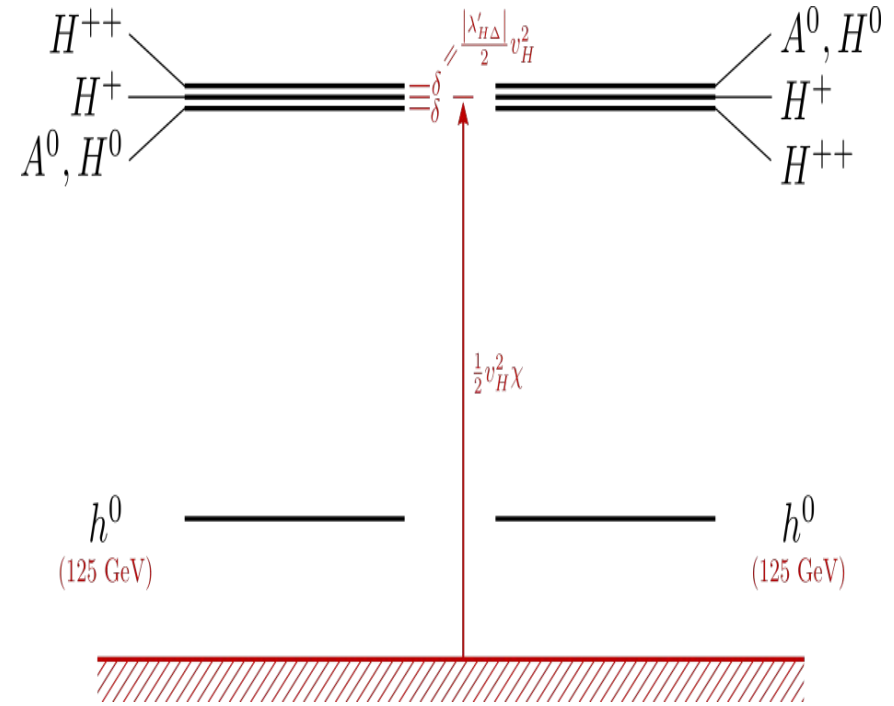
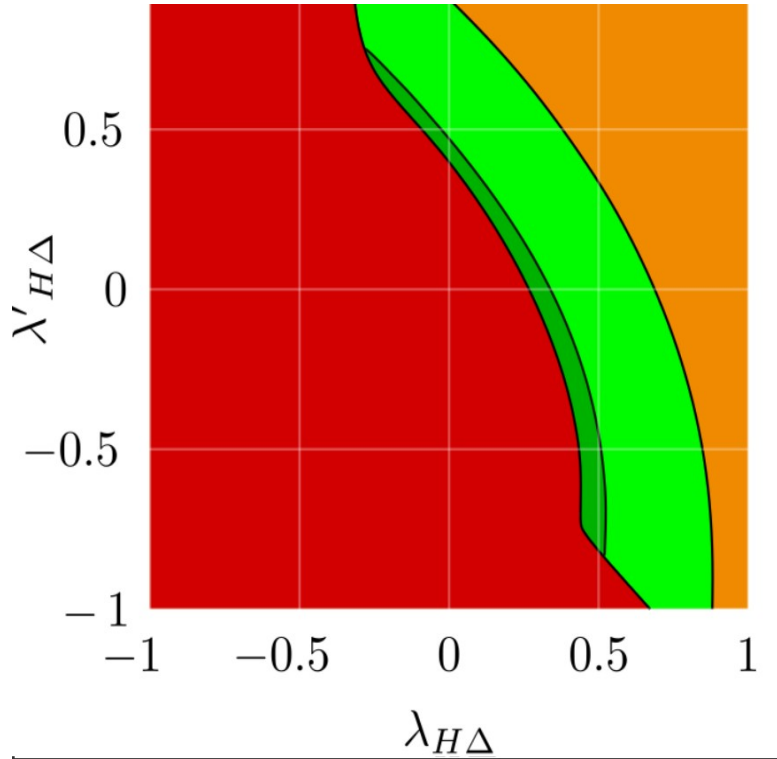


# Consistency of the triplet seesaw model revisited

Cesar Bonilla,\* Renato M. Fonseca,† and J. W. F. Valle‡

AHEP Group, Instituto de Física Corpuscular, C.S.I.C./Universitat de València, Edificio Institutos de Investigación, Apartado 22085, E-46071 Valencia, Spain

(Received 17 August 2015; published 21 October 2015)



arXiv:1508.02323

In addition to SM gauge invariance must break lepton number to give Masses to neutrinos ..



# *Neutrinos & Stability*

arXiv:1506.04031

$$V(\sigma, H) = \mu_1^2 |\sigma|^2 + \mu_2^2 H^\dagger H + \lambda_1 |\sigma|^4 \\ + \lambda_2 (H^\dagger H)^2 + \lambda_{12} (H^\dagger H) |\sigma|^2 .$$

In addition to SM gauge invariance must break lepton number to give Masses to neutrinos ..

# Neutrinos & Stability

arXiv:1506.04031

$$V(\sigma, H) = \mu_1^2 |\sigma|^2 + \mu_2^2 H^\dagger H + \lambda_1 |\sigma|^4 \\ + \lambda_2 (H^\dagger H)^2 + \lambda_{12} (H^\dagger H) |\sigma|^2 .$$

Vacuum stability with spontaneous violation of lepton number

Cesar Bonilla, Renato M. Fonseca and José W. F. Valle

*AHEP Group, Institut de Física Corpuscular – C.S.I.C./Universitat de València, Parc Científic de Paterna.*

*C/ Catedrático José Beltrán, 2 E-46980 Paterna (Valencia) - SPAIN*

(Dated: June 15, 2015)

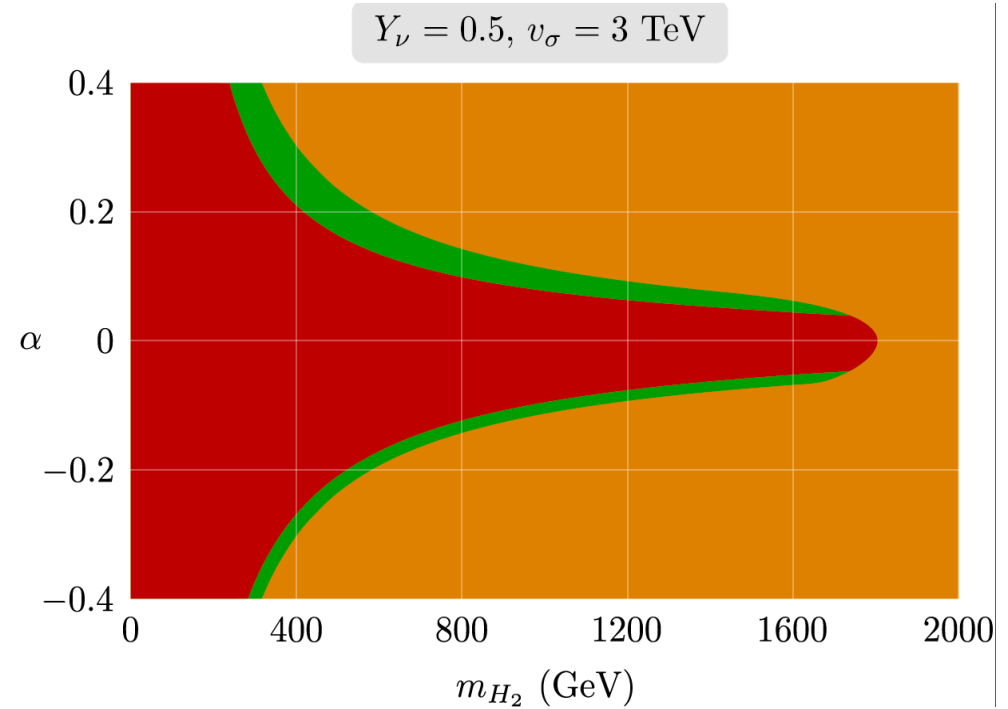
In addition to SM gauge invariance must break lepton number to give Masses to neutrinos ..

JWF Valle

# Neutrinos & Stability

arXiv:1506.04031

$$V(\sigma, H) = \mu_1^2 |\sigma|^2 + \mu_2^2 H^\dagger H + \lambda_1 |\sigma|^4 + \lambda_2 (H^\dagger H)^2 + \lambda_{12} (H^\dagger H) |\sigma|^2 .$$



Vacuum stability with spontaneous violation of lepton number

Cesar Bonilla, Renato M. Fonseca and José W. F. Valle

*AHEP Group, Institut de Física Corpuscular – C.S.I.C./Universitat de València, Parc Científic de Paterna.*

*C/ Catedrático Jose Beltrán, 2 E-46980 Paterna (Valencia) - SPAIN*

(Dated: June 15, 2015)

In addition to SM gauge invariance must break lepton number to give Masses to neutrinos ..

# *Neutrinos & Invisible Higgs*

PHYSICAL REVIEW D **91**, 113015 (2015)

## **Neutrino mass and invisible Higgs decays at the LHC**

Cesar Bonilla,<sup>1,\*</sup> Jorge C. Romão,<sup>2,†</sup> and José W. F. Valle<sup>1,‡</sup>

$$\Gamma(H_i \rightarrow JJ) = \frac{1}{32\pi} \frac{g_{H_i JJ}^2}{m_{H_i}}$$

arXiv:1502.01649

# Neutrinos & Invisible Higgs

PHYSICAL REVIEW D **91**, 113015 (2015)

## Neutrino mass and invisible Higgs decays at the LHC

Cesar Bonilla,<sup>1,\*</sup> Jorge C. Romão,<sup>2,†</sup> and José W. F. Valle<sup>1,‡</sup>

$$\Gamma(H_i \rightarrow JJ) = \frac{1}{32\pi} \frac{g_{H_i JJ}^2}{m_{H_i}}$$

channel	ATLAS	CMS
$\mu_{\gamma\gamma}$	$1.17 \pm 0.27$	$1.14_{-0.23}^{+0.26}$
$\mu_{WW}$	$1.00_{-0.29}^{+0.32}$	$0.83 \pm 0.21$
$\mu_{ZZ}$	$1.44_{-0.35}^{+0.40}$	$1.00 \pm 0.29$
$\mu_{\tau^+\tau^-}$	$1.4_{-0.4}^{+0.5}$	$0.91 \pm 0.27$
$\mu_{b\bar{b}}$	$0.2_{-0.6}^{+0.7}$	$0.93 \pm 0.49$

[arXiv:1502.01649](https://arxiv.org/abs/1502.01649)

# Neutrinos & Invisible Higgs

PHYSICAL REVIEW D **91**, 113015 (2015)

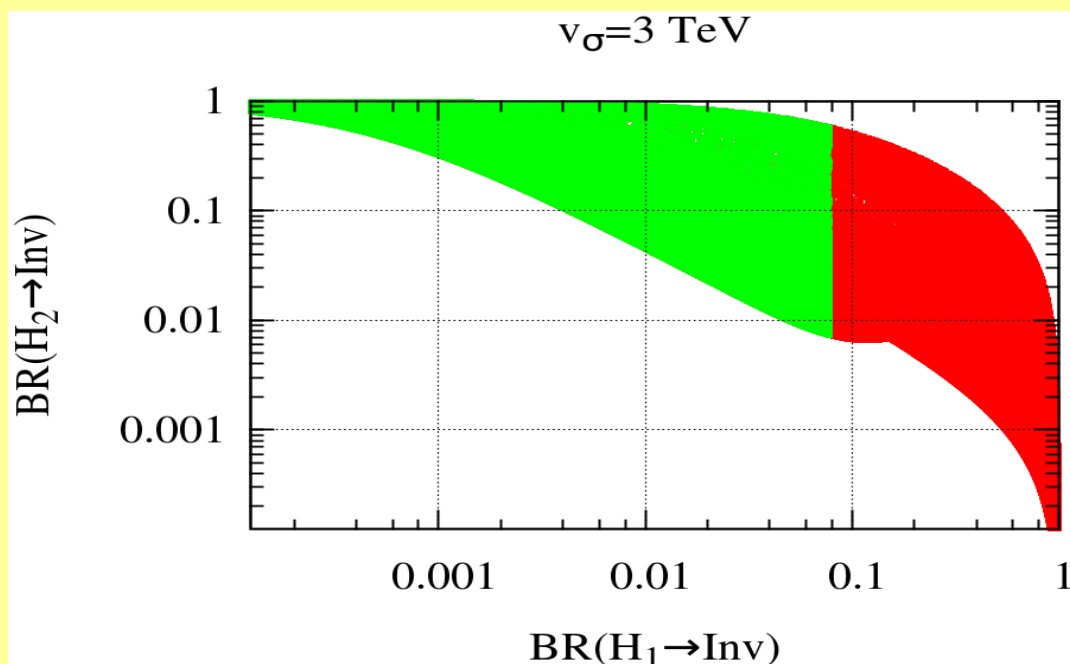
## Neutrino mass and invisible Higgs decays at the LHC

Cesar Bonilla,<sup>1,\*</sup> Jorge C. Romão,<sup>2,†</sup> and José W. F. Valle<sup>1,‡</sup>

$$\Gamma(H_i \rightarrow JJ) = \frac{1}{32\pi} \frac{g_{H_i JJ}^2}{m_{H_i}}$$

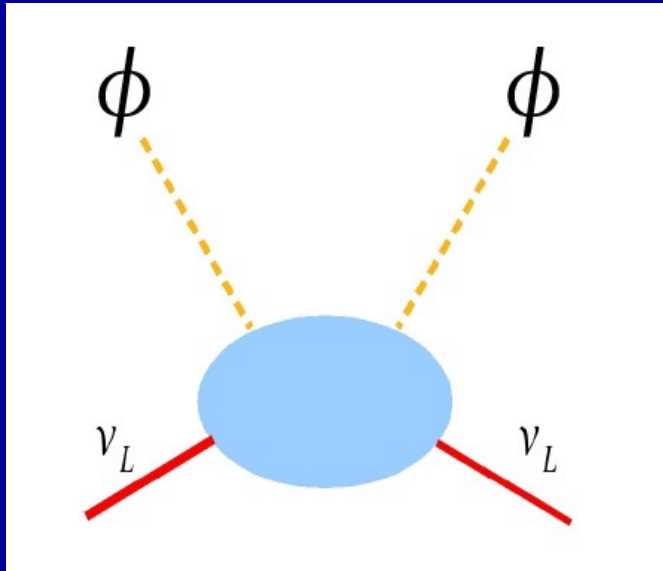
channel	ATLAS	CMS
$\mu_{\gamma\gamma}$	$1.17 \pm 0.27$	$1.14_{-0.23}^{+0.26}$
$\mu_{WW}$	$1.00_{-0.29}^{+0.32}$	$0.83 \pm 0.21$
$\mu_{ZZ}$	$1.44_{-0.35}^{+0.40}$	$1.00 \pm 0.29$
$\mu_{\tau+\tau^-}$	$1.4_{-0.4}^{+0.5}$	$0.91 \pm 0.27$
$\mu_{b\bar{b}}$	$0.2_{-0.6}^{+0.7}$	$0.93 \pm 0.49$

arXiv:1502.01649

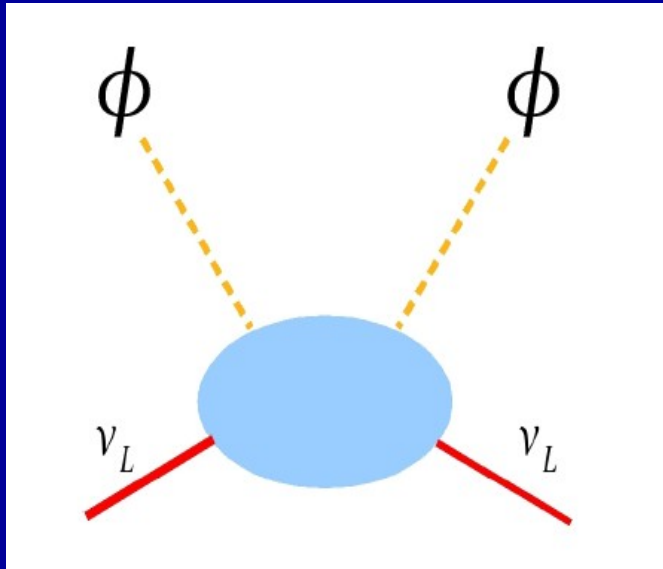


JWF Valle

# *origin of neutrino mass and seesaw*



# *origin of neutrino mass and seesaw*



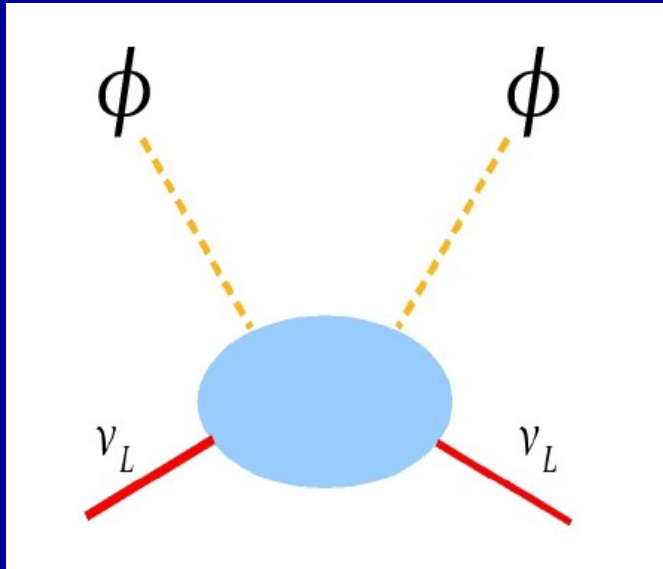
**SCALE**

**MECHANISM**

**FLAVOR STRUCTURE**



# origin of neutrino mass *and seesaw*



$$v_3 v_1 \sim v_2^2 \text{ with } v_1 \gg v_2 \gg v_3$$

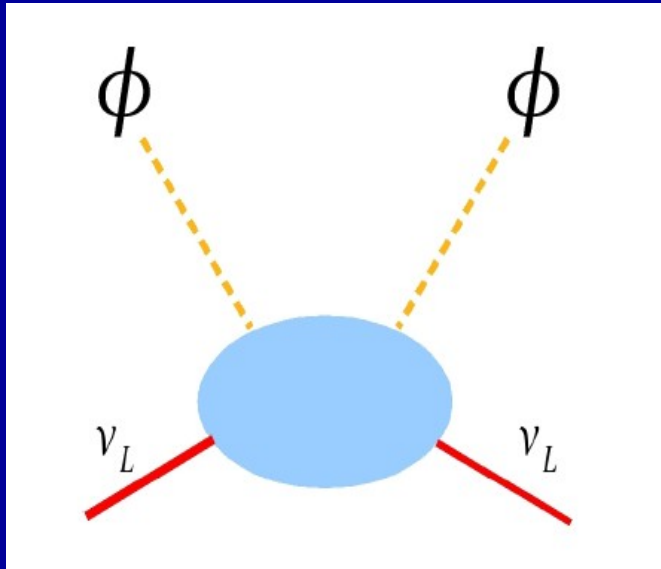


**SCALE**

**MECHANISM**

**FLAVOR STRUCTURE**

# origin of neutrino mass *and seesaw*



fermion exchange

## TYPE I

Minkowski 77  
Gellman Ramond Slansky 80  
Glashow, Yanagida 79  
Mohapatra Senjanovic 80  
Lazarides Shafi Weterrich 81  
Schechter-Valle, 80 & 82

Scalar-exchange

## TYPE II

Schechter-Valle 80/82



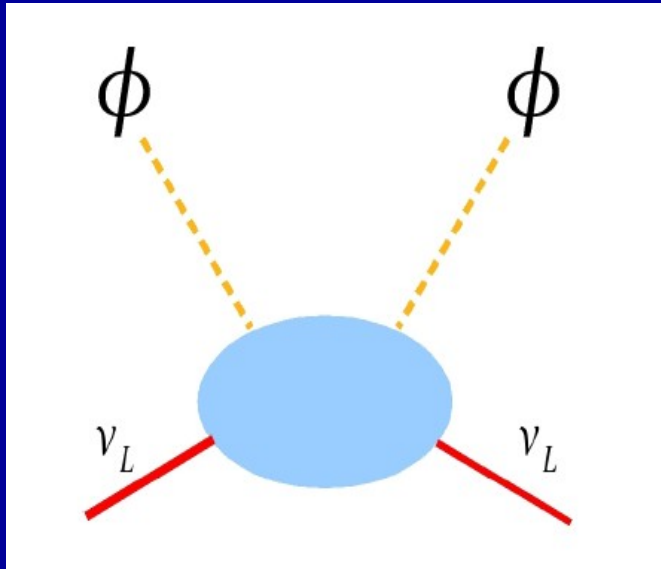
$$v_3 v_1 \sim v_2^2 \text{ with } v_1 \gg v_2 \gg v_3$$

**SCALE**

**MECHANISM**

**FLAVOR STRUCTURE**

# origin of neutrino mass *and seesaw*



fermion exchange

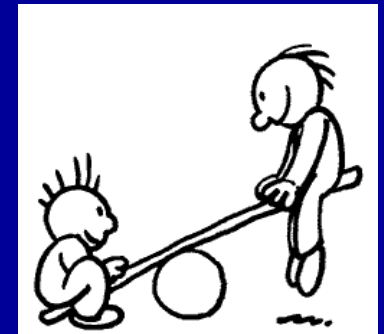
## TYPE I

Minkowski 77  
Gellman Ramond Slansky 80  
Glashow, Yanagida 79  
Mohapatra Senjanovic 80  
Lazarides Shafi Weterrich 81  
Schechter-Valle, 80 & 82

Scalar-exchange

## TYPE II

Schechter-Valle 80/82



$$v_3 v_1 \sim v_2^2 \text{ with } v_1 \gg v_2 \gg v_3$$

**SCALE**

**MECHANISM**

**FLAVOR STRUCTURE**

Number & properties of messengers

## LOW-SCALE SEESAW

Mohapatra-Valle 86  
Akhmedov et al PRD53 (1996) 2752  
Malinsky et al PRL95(2005)161801  
Bazzocchi et al, PRD81 (2010) 051701

# *Radiative neutrino mass*

many low-scale neutrino mass schemes ...

arXiv:1404.3751

*3<sub>L</sub>3<sub>c</sub>1 scheme* # generations = # colours

Singer, Valle, Schechter, Phys.Rev. D22 (1980) 738

# Radiative neutrino mass

many low-scale neutrino mass schemes ...

arXiv:1404.3751

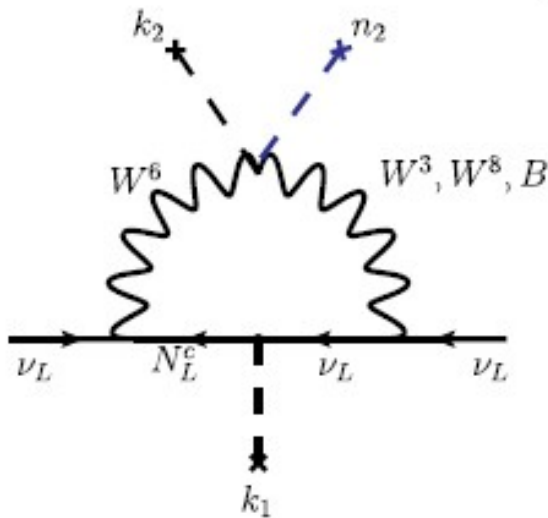
**3L3c1 scheme** # generations = # colours

Singer, Valle, Schechter, Phys.Rev. D22 (1980) 738

PHYSICAL REVIEW D **90**, 013005 (2014)

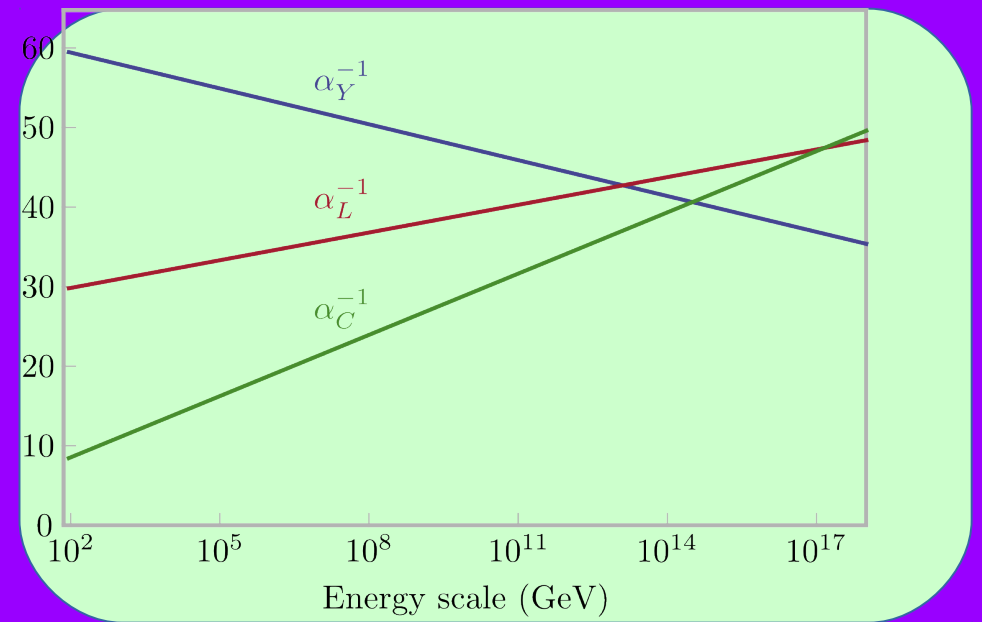
**Radiative neutrino mass in 3-3-1 scheme**

PHYSICAL REVIEW D **90**, 013005 (2014)



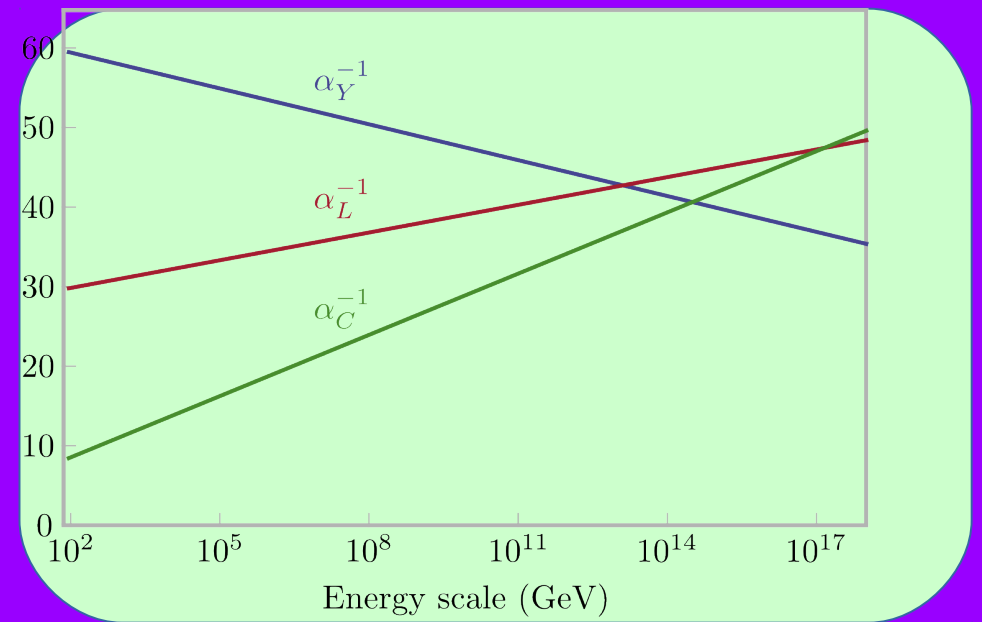
Gauge vs Higgs

*gauge coupling  
unification :  
a near miss ...*



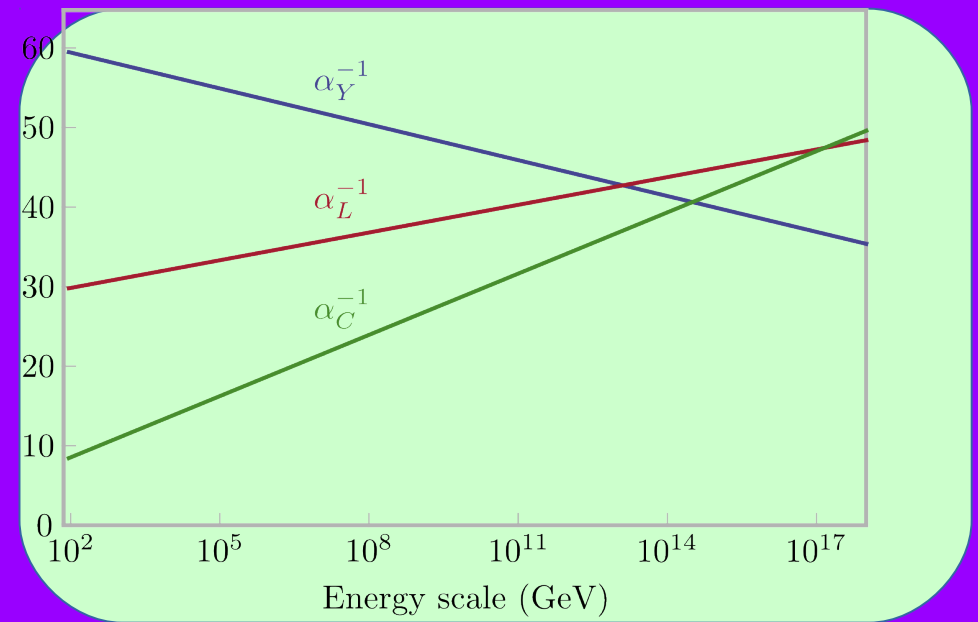
What makes the gauge couplings unify?

*gauge coupling  
unification :  
a near miss ...*

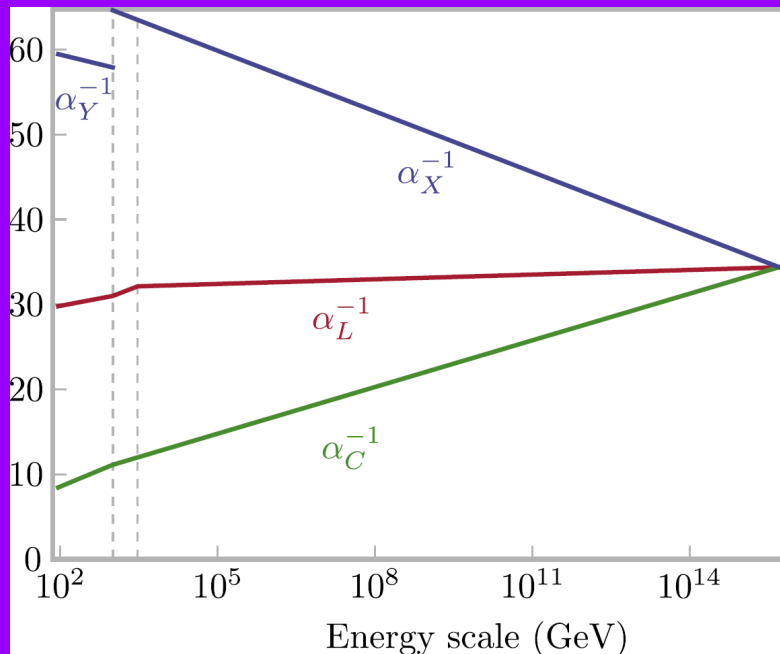


What makes the gauge couplings unify? GUT (p decay)  
SUSY

*gauge coupling unification :  
a near miss ...*



What makes the gauge couplings unify? **GUT (p decay)**  
**SUSY**  
**NEUTRINO**

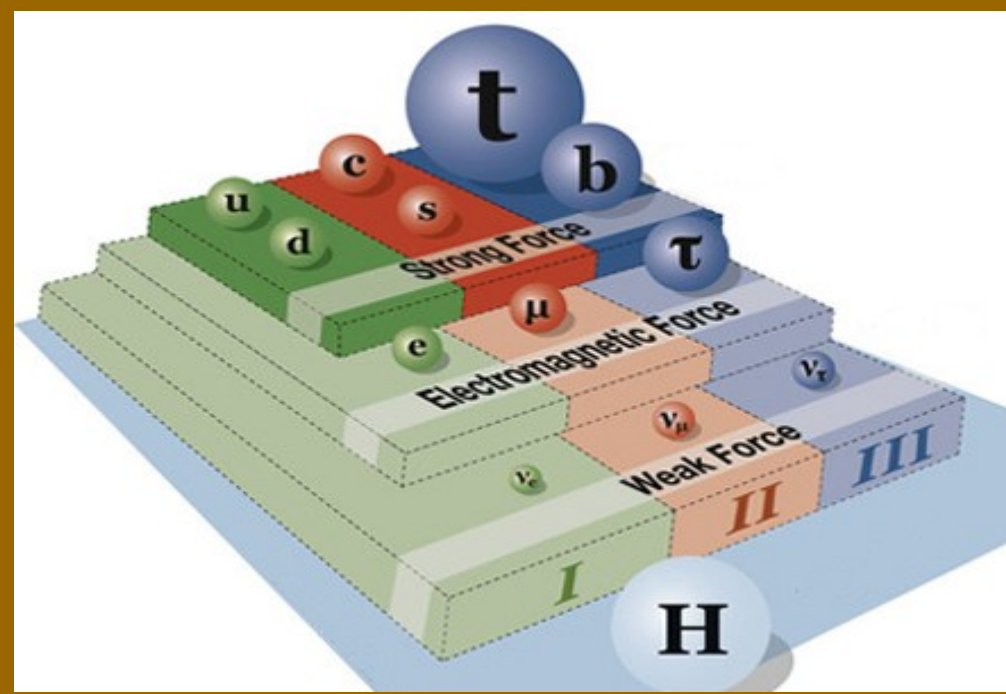


The physics responsible for gauge coupling unification may also induce small neutrino masses

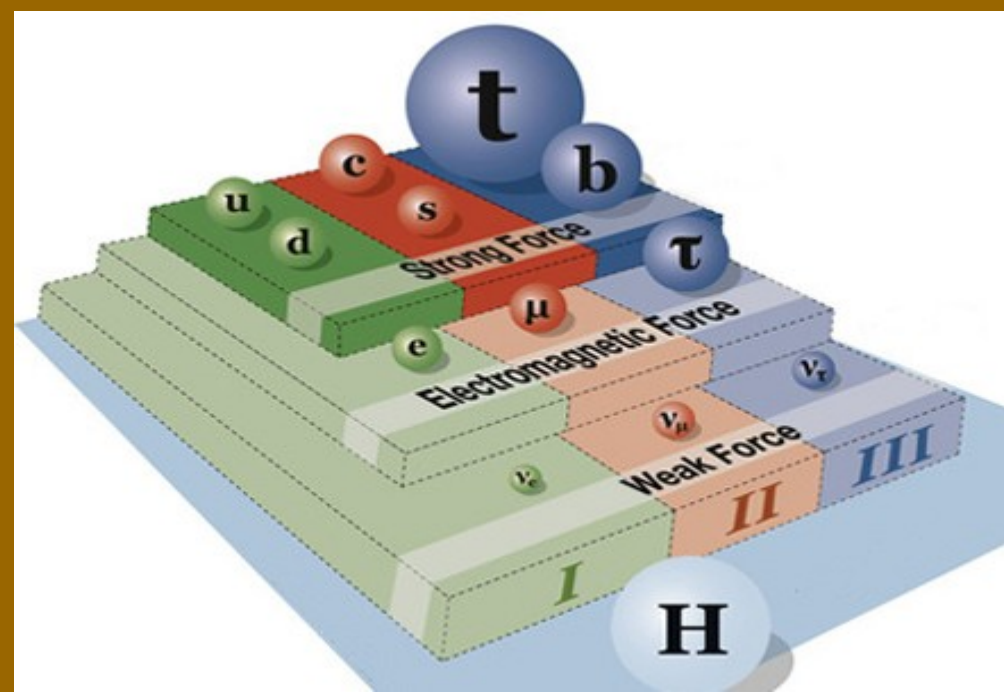
Phys. Rev. D 91, 031702 (2015)  
Boucenna, Fonseca, Gonzalez-Canales, JV



# *Flavor problem*



# Flavor problem



*pattern of masses...*

$$\frac{m_\tau}{\sqrt{m_e m_\mu}} \approx \frac{m_b}{\sqrt{m_d m_s}}$$

*b-tau unification without GUTS...*

Morisi et al Phys.Rev. D84 (2011) 036003

King et al Phys. Lett. B 724 (2013) 68

Morisi et al Phys.Rev. D88 (2013) 036001

Bonilla et al Phys.Lett. B742 (2015) 99

JWF Valle

# Flavor problem

pattern of masses...

$$\frac{m_\tau}{\sqrt{m_e m_\mu}} \approx \frac{m_b}{\sqrt{m_d m_s}}$$

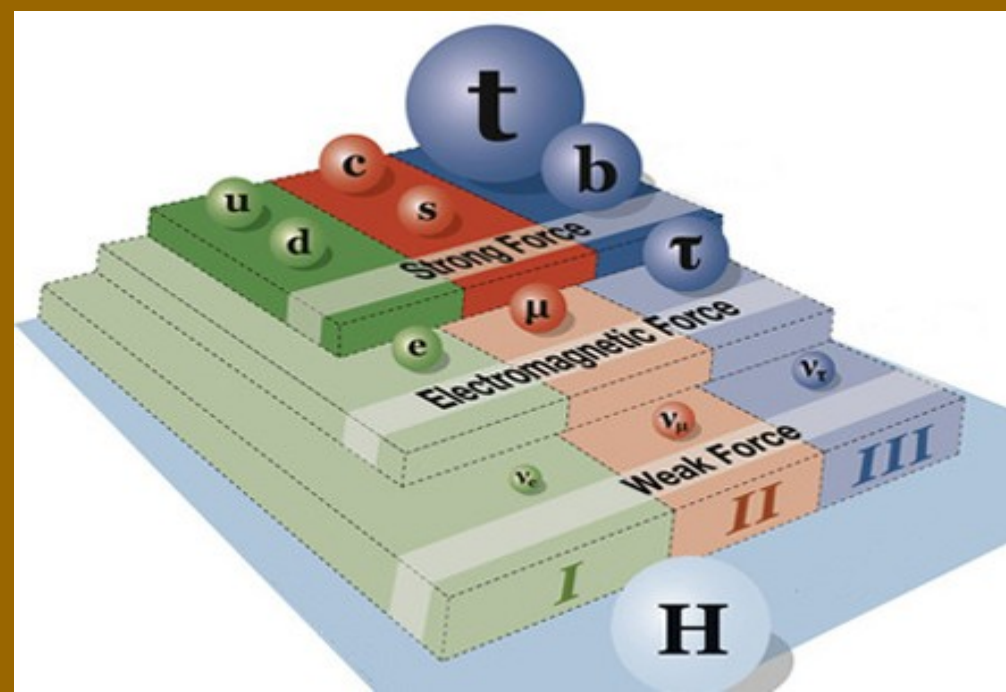
*b-tau unification without GUTS...*

Morisi et al Phys.Rev. D84 (2011) 036003

King et al Phys. Lett. B 724 (2013) 68

Morisi et al Phys.Rev. D88 (2013) 036001



Bonilla et al Phys.Lett. B742 (2015) 99



pattern of mixing parameters ...

*Anarchy ?*

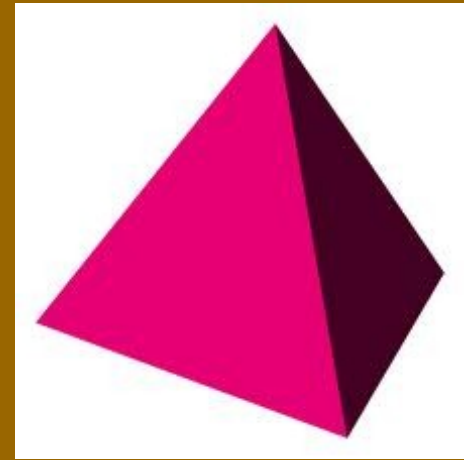
Donoghue et al PRD73  
Hall, Murayama, Weiner, PRL  
Altarelli, Feruglio, Masina, JHEP

$\begin{pmatrix} \nu_e \\ e \\ e_R \end{pmatrix}_L$	$\begin{pmatrix} \nu_\mu \\ \mu \\ \mu_R \end{pmatrix}_L$	$\begin{pmatrix} \nu_\tau \\ \tau \\ \tau_R \end{pmatrix}_L$
$\begin{pmatrix} u \\ d \\ u_R \\ d_R \end{pmatrix}_L$	$\begin{pmatrix} c \\ s \\ c_R \\ s_R \end{pmatrix}_L$	$\begin{pmatrix} t \\ b \\ t_R \\ b_R \end{pmatrix}_L$

# FLAVOR SYMMETRY

A4

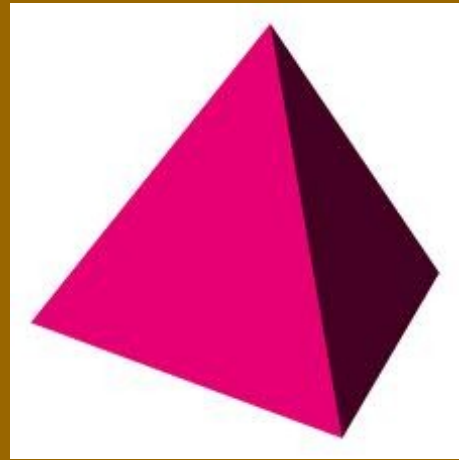


Babu-Ma-Valle PLB552 (2003) 207  
 Hirsch et al PRD69 (2004) 093006

$$\sin^2 \theta_{23} = 0.5$$

$$\sin^2 \theta_{13} = 0$$

# FLAVOR SYMMETRY



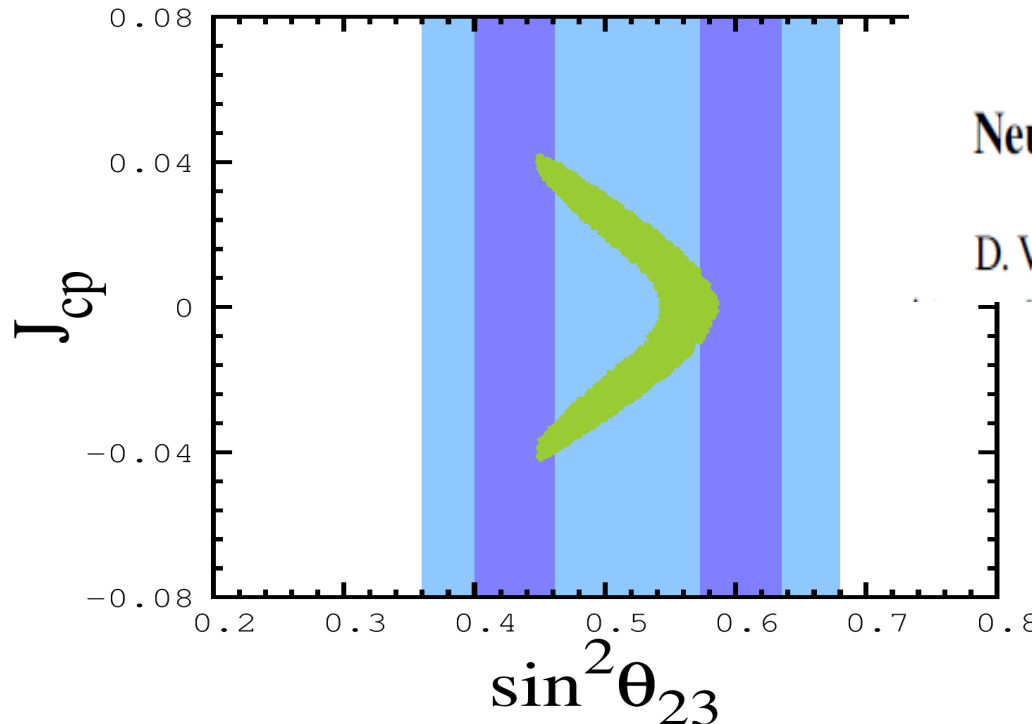
A4

$$\begin{array}{ccc}
 \begin{pmatrix} \nu_e \\ e \\ e_R \end{pmatrix}_L & \begin{pmatrix} \nu_\mu \\ \mu \\ \mu_R \end{pmatrix}_L & \begin{pmatrix} \nu_\tau \\ \tau \\ \tau_R \end{pmatrix}_L \\
 \begin{pmatrix} u \\ d \\ u_R \\ d_R \end{pmatrix}_L & \begin{pmatrix} c \\ s \\ c_R \\ s_R \end{pmatrix}_L & \begin{pmatrix} t \\ b \\ t_R \\ b_R \end{pmatrix}_L
 \end{array}$$

Babu-Ma-Valle PLB552 (2003) 207  
 Hirsch et al PRD69 (2004) 093006

$$\sin^2 \theta_{23} = 0.5$$

$$\sin^2 \theta_{13} = 0$$



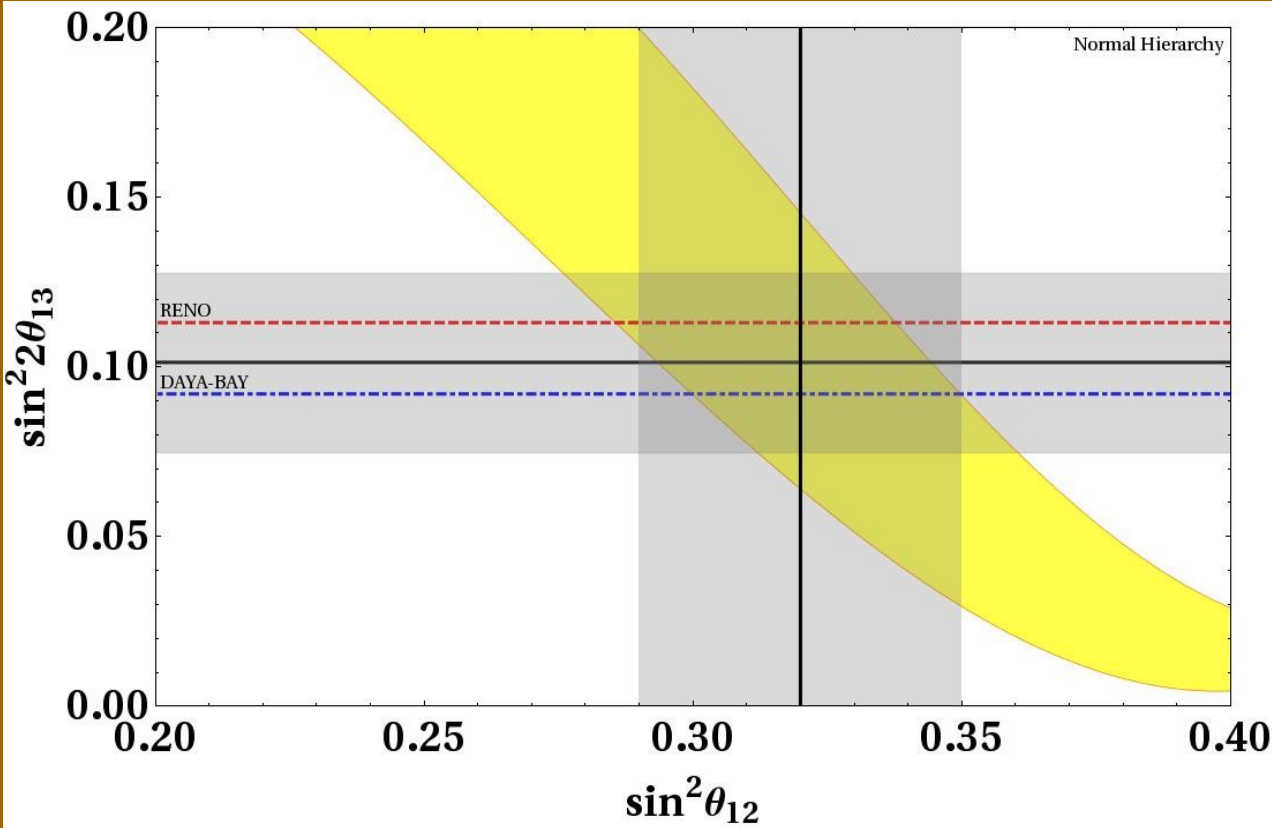
PHYSICAL REVIEW D 88, 016003 (2013)

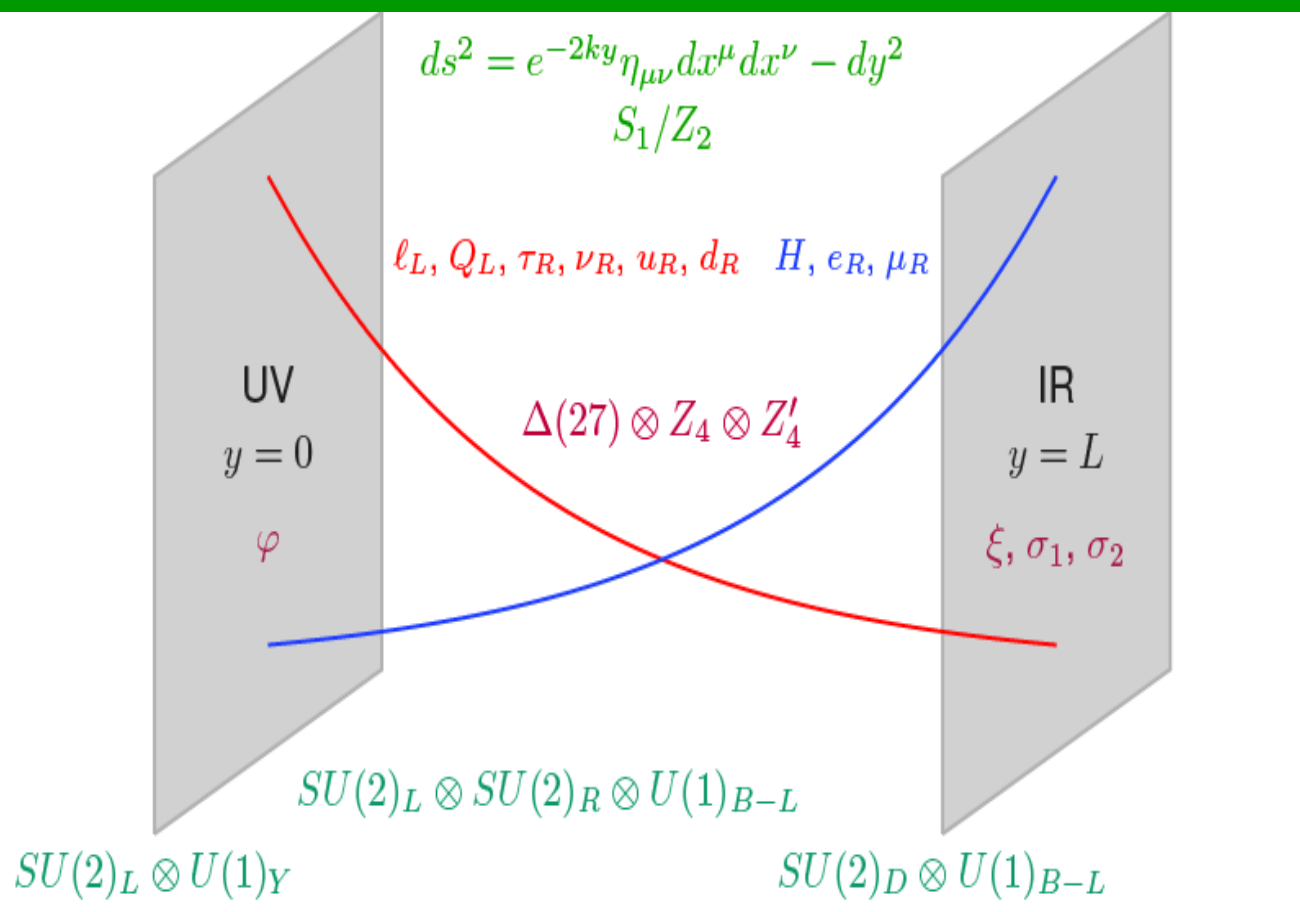
## Neutrino mixing with revamped $A_4$ flavor symmetry

D. V. Forero,<sup>1,2,\*</sup> S. Morisi,<sup>3,†</sup> J. C. Romão,<sup>1,‡</sup> and J. W. F. Valle<sup>2,§</sup>

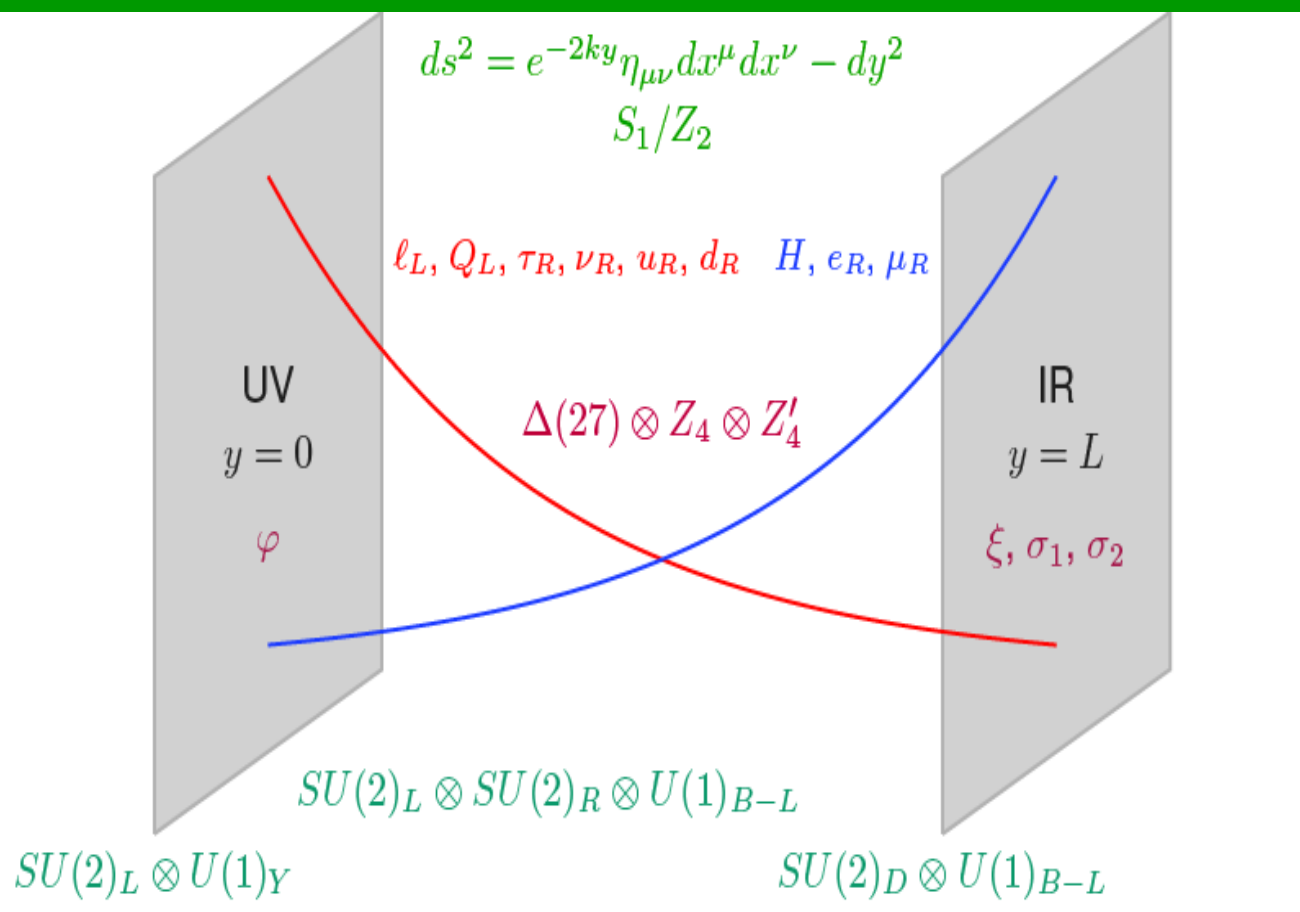
# OSCILLATION PARAMETER CORRELATIONS

Boucenna et al  
PhysRevD.86.073008





**Mass hierarchies in principle be accounted for by judicious choices of the bulk mass parameters**

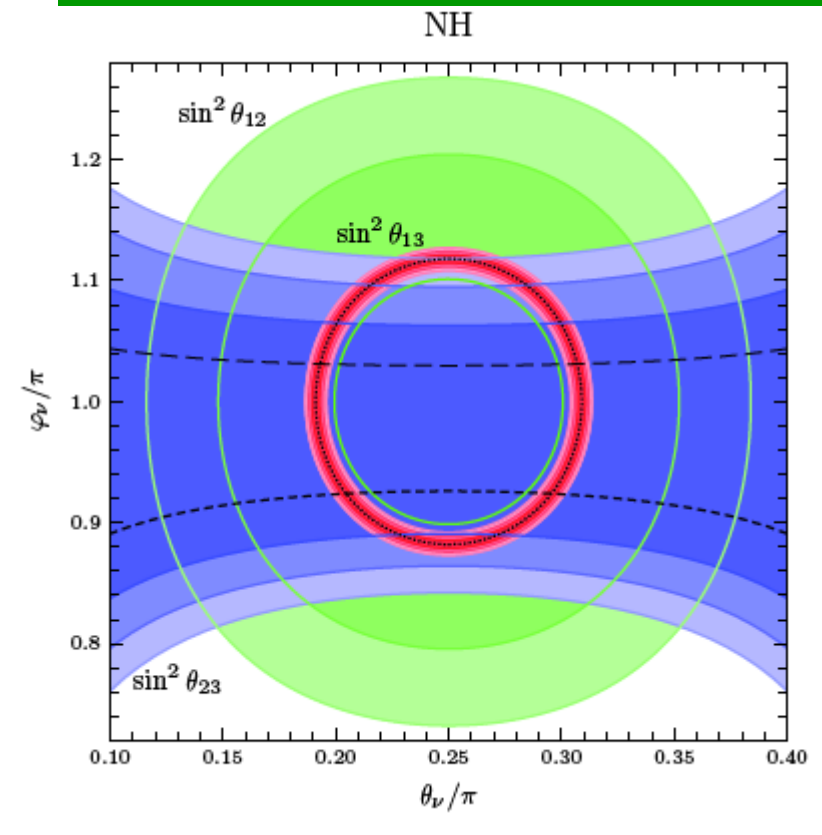
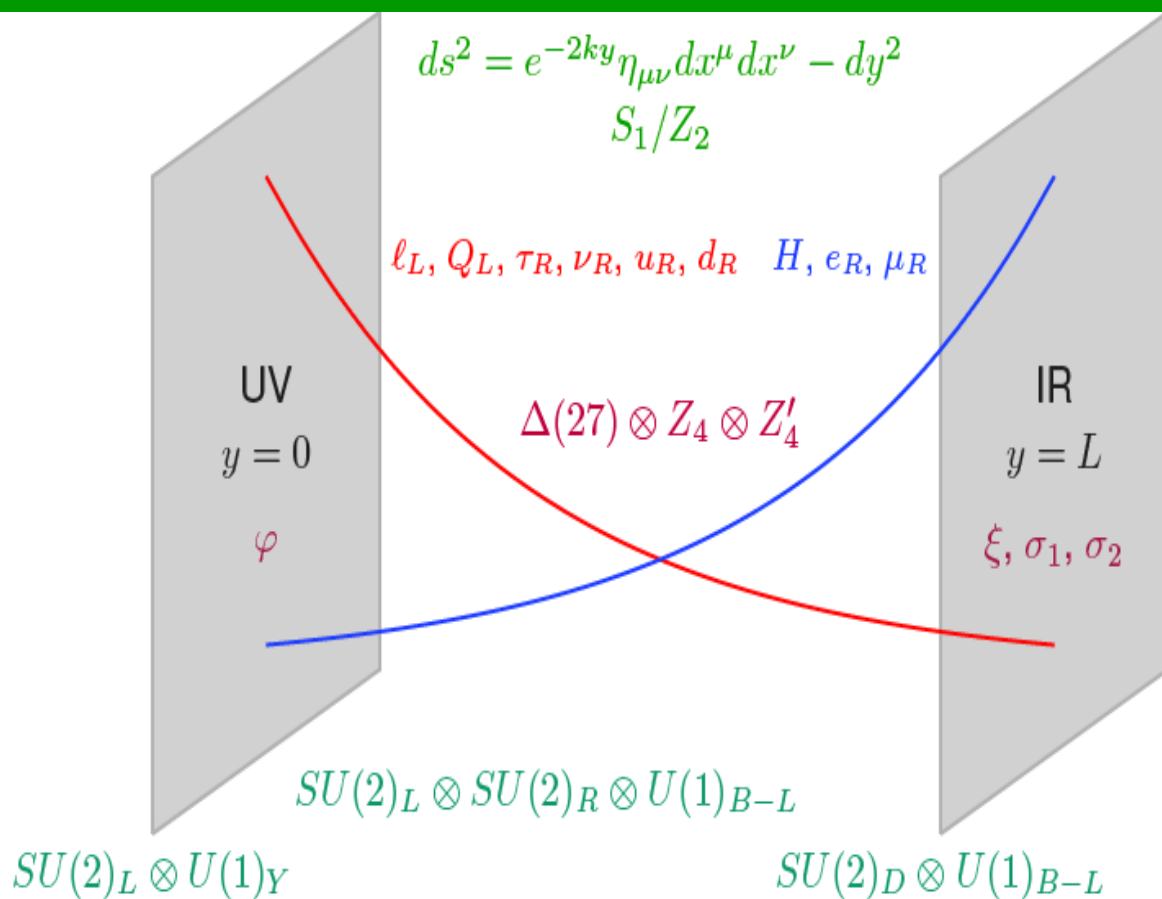


**Mass hierarchies in principle be accounted for by judicious choices of the bulk mass parameters**

### Particle content and transformation properties

Field	$\Psi_\ell$	$\Psi_e$	$\Psi_\mu$	$\Psi_\tau$	$\Psi_{\nu_1}$	$\Psi_{\nu_2}$	$\Psi_{\nu_3}$	$H$	$\varphi$	$\xi$	$\sigma_1$	$\sigma_2$
$\Delta(27)$	$\mathbf{3}$	$\mathbf{1}_{0,0}$	$\mathbf{1}_{1,0}$	$\mathbf{1}_{2,0}$	$\mathbf{1}_{0,0}$	$\mathbf{1}_{0,0}$	$\mathbf{1}_{0,0}$	$\mathbf{1}_{0,0}$	$\mathbf{3}$	$\mathbf{3}$	$\mathbf{1}_{0,1}$	$\mathbf{1}_{0,0}$
$Z_4$	1	1	1	1	-1	$i$	-1	1	1	-1	1	$i$
$Z'_4$	1	$i$	$i$	$i$	-1	-1	-1	1	$-i$	1	-1	-1

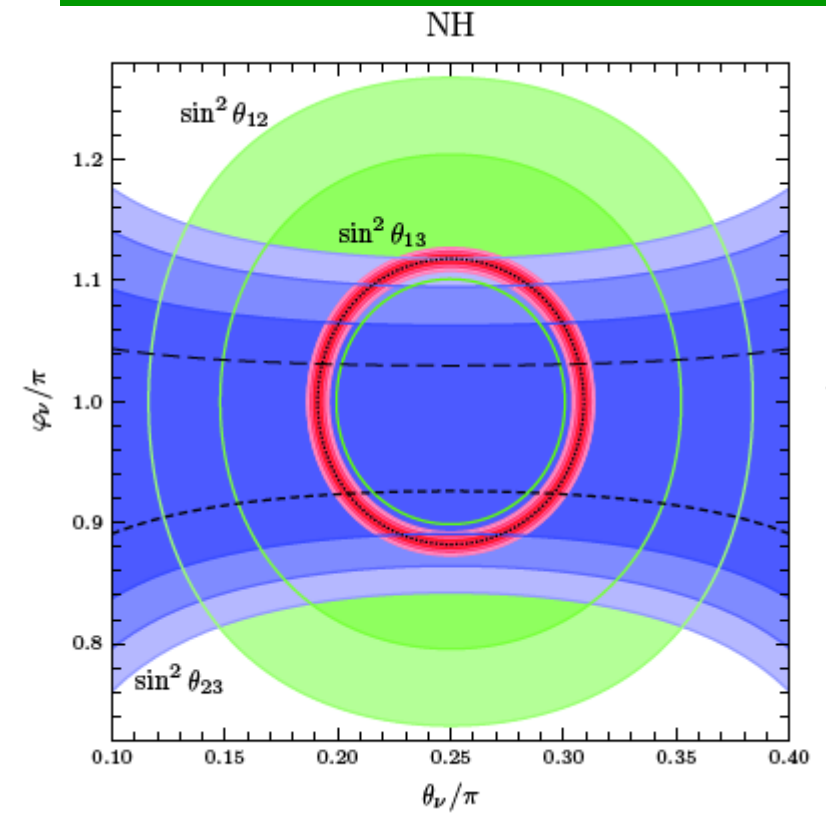
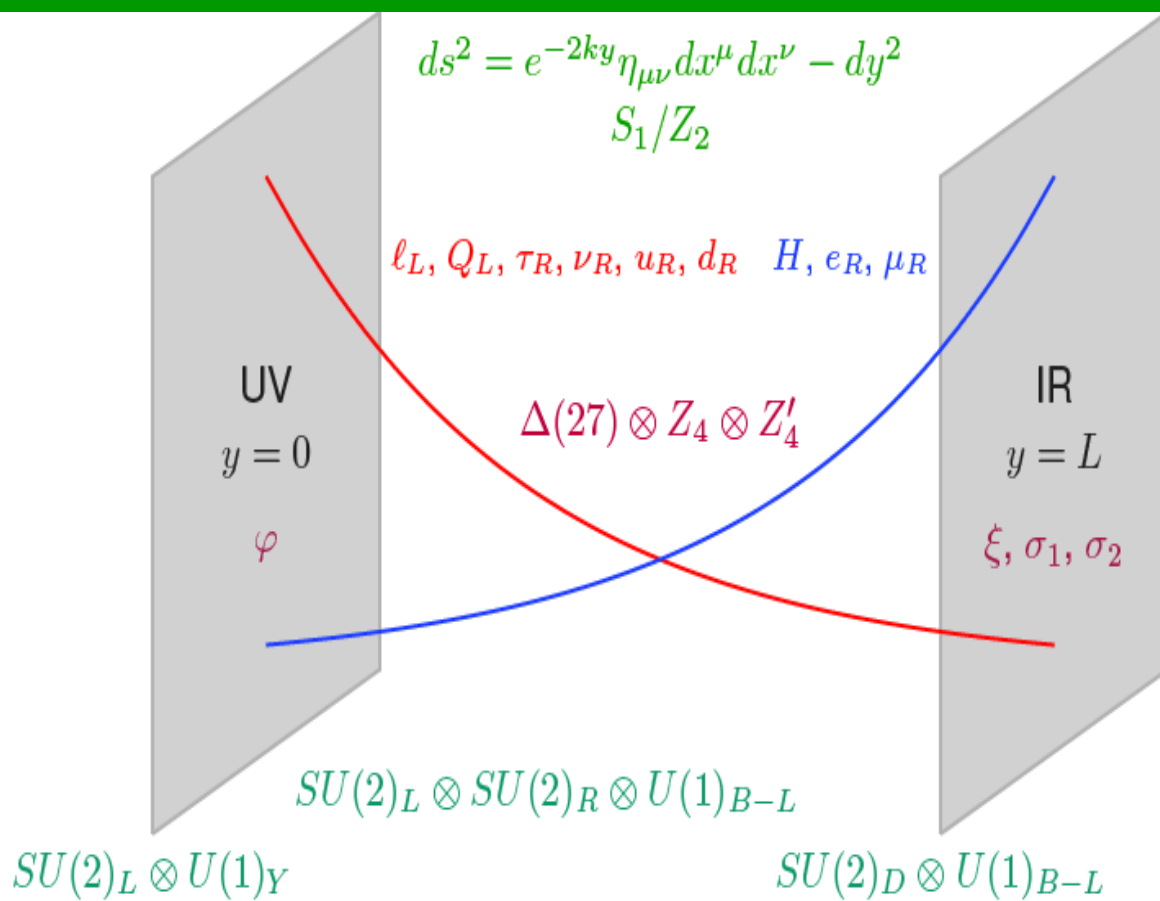




Mass hierarchies in principle be accounted for by judicious choices of the bulk mass parameters

### Particle content and transformation properties

Field	$\Psi_\ell$	$\Psi_e$	$\Psi_\mu$	$\Psi_\tau$	$\Psi_{\nu_1}$	$\Psi_{\nu_2}$	$\Psi_{\nu_3}$	$H$	$\varphi$	$\xi$	$\sigma_1$	$\sigma_2$
$\Delta(27)$	$\mathbf{3}$	$\mathbf{1}_{0,0}$	$\mathbf{1}_{1,0}$	$\mathbf{1}_{2,0}$	$\mathbf{1}_{0,0}$	$\mathbf{1}_{0,0}$	$\mathbf{1}_{0,0}$	$\mathbf{1}_{0,0}$	$\mathbf{3}$	$\mathbf{3}$	$\mathbf{1}_{0,1}$	$\mathbf{1}_{0,0}$
$Z_4$	1	1	1	1	-1	$i$	-1	1	1	-1	1	$i$
$Z'_4$	1	$i$	$i$	$i$	-1	-1	-1	1	$-i$	1	-1	-1



**Mass hierarchies in principle be accounted for by judicious choices of the bulk mass parameters**

**Particle content and transformation properties**

Field	$\Psi_\ell$	$\Psi_e$	$\Psi_\mu$	$\Psi_\tau$	$\Psi_{\nu_1}$	$\Psi_{\nu_2}$	$\Psi_{\nu_3}$	$H$	$\varphi$	$\xi$	$\sigma_1$	$\sigma_2$
$\Delta(27)$	<b>3</b>	$1_{0,0}$	$1_{1,0}$	$1_{2,0}$	$1_{0,0}$	$1_{0,0}$	$1_{0,0}$	$1_{0,0}$	<b>3</b>	<b>3</b>	$1_{0,1}$	$1_{0,0}$
$Z_4$	1	1	1	1	-1	$i$	-1	1	1	-1	1	$i$
$Z'_4$	1	$i$	$i$	$i$	-1	-1	-1	1	$-i$	1	-1	-1

$\sin^2 \theta_{12} \cos^2 \theta_{13} = \frac{1}{3}$

# B ANOMALIES

$$R_K = \frac{\text{BR}(B \rightarrow K \mu^+ \mu^-)}{\text{BR}(B \rightarrow K e^+ e^-)} = 0.745_{-0.074}^{+0.090} \pm 0.036$$

# B ANOMALIES

$$R_K = \frac{\text{BR}(B \rightarrow K \mu^+ \mu^-)}{\text{BR}(B \rightarrow K e^+ e^-)} = 0.745_{-0.074}^{+0.090} \pm 0.036$$

$$\mathcal{O}^{ij} = \frac{1}{\Lambda^2} J_\alpha^d J_{\ell ij}^\alpha,$$

where

$$J_\alpha^d = C_{bs}^Q \bar{b} \gamma_\alpha P_L s,$$

$$J_{\ell ij}^\alpha = C_{ij}^L \bar{\ell}_i \gamma^\alpha P_L \ell_j,$$

$$-\frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \frac{e^2}{16\pi^2} \sum_i (C_i \mathcal{O}_i + C'_i \mathcal{O}'_i)$$

# B ANOMALIES

$$R_K = \frac{\text{BR}(B \rightarrow K \mu^+ \mu^-)}{\text{BR}(B \rightarrow K e^+ e^-)} = 0.745_{-0.074}^{+0.090} \pm 0.036$$

$$\mathcal{O}^{ij} = \frac{1}{\Lambda^2} J_\alpha^d J_{\ell ij}^\alpha,$$

where

$$J_\alpha^d = C_{bs}^Q \bar{b} \gamma_\alpha P_L s,$$

$$J_{\ell ij}^\alpha = C_{ij}^L \bar{\ell}_i \gamma^\alpha P_L \ell_j,$$

The violation of lepton universality usually comes together with the violation of lepton flavor. Based on symmetry arguments, Glashow, Guadagnoli and Lane [19] recently argued that the observation of universality violation in the lepton flavor conserving (LFC)  $B \rightarrow K \ell_i^+ \ell_i^-$  decays implies the existence of the lepton flavor violating (LFV) processes  $B \rightarrow K \ell_i^+ \ell_j^-$  (with  $i \neq j$ ).

$$-\frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \frac{e^2}{16\pi^2} \sum_i (C_i \mathcal{O}_i + C'_i \mathcal{O}'_i)$$

# B ANOMALIES

$$R_K = \frac{\text{BR}(B \rightarrow K \mu^+ \mu^-)}{\text{BR}(B \rightarrow K e^+ e^-)} = 0.745_{-0.074}^{+0.090} \pm 0.036$$

$$\mathcal{O}^{ij} = \frac{1}{\Lambda^2} J_\alpha^d J_{\ell ij}^\alpha,$$

where

$$J_\alpha^d = C_{bs}^Q \bar{b} \gamma_\alpha P_L s,$$

$$J_{\ell ij}^\alpha = C_{ij}^L \bar{\ell}_i \gamma^\alpha P_L \ell_j,$$

The violation of lepton universality usually comes together with the violation of lepton flavor. Based on symmetry arguments, Glashow, Guadagnoli and Lane [19] recently argued that the observation of universality violation in the lepton flavor conserving (LFC)  $B \rightarrow K \ell_i^+ \ell_i^-$  decays implies the existence of the lepton flavor violating (LFV) processes  $B \rightarrow K \ell_i^+ \ell_j^-$  (with  $i \neq j$ ).

$$-\frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \frac{e^2}{16\pi^2} \sum_i (C_i \mathcal{O}_i + C'_i \mathcal{O}'_i)$$

Physics Letters B 750 (2015) 367–371

Here we raise the following question: can the leptonic mixing matrix provide the required lepton flavor structure in  $\mathcal{O}_9$  and  $\mathcal{O}_{10}$ ?

# B ANOMALIES

$$R_K = \frac{\text{BR}(B \rightarrow K \mu^+ \mu^-)}{\text{BR}(B \rightarrow K e^+ e^-)} = 0.745_{-0.074}^{+0.090} \pm 0.036$$

$$\mathcal{O}^{ij} = \frac{1}{\Lambda^2} J_\alpha^d J_{\ell ij}^\alpha,$$

where

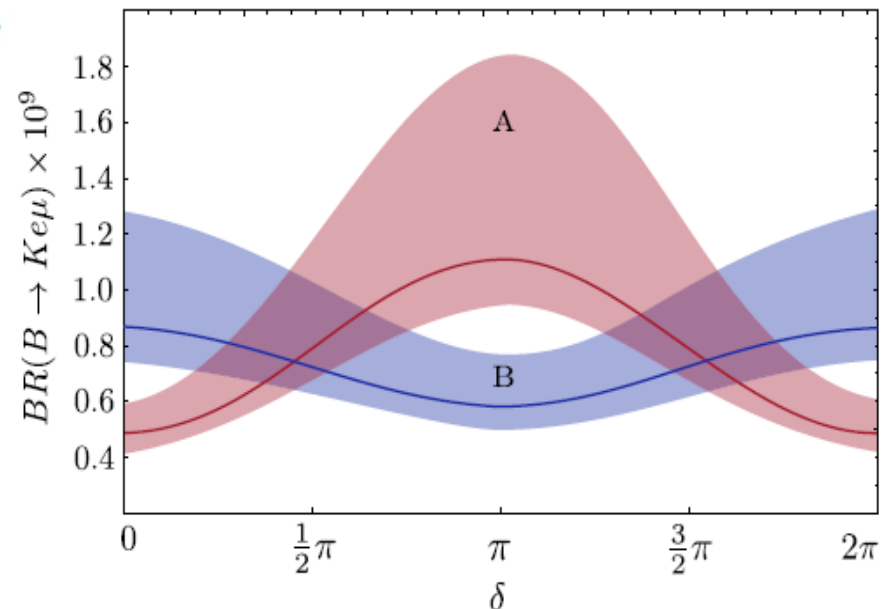
$$J_\alpha^d = C_{bs}^Q \bar{b} \gamma_\alpha P_L s,$$

$$J_{\ell ij}^\alpha = C_{ij}^L \bar{\ell}_i \gamma^\alpha P_L \ell_j,$$

The violation of lepton universality usually comes together with the violation of lepton flavor. Based on symmetry arguments, Glashow, Guadagnoli and Lane [19] recently argued that the observation of universality violation in the lepton flavor conserving (LFC)  $B \rightarrow K \ell_i^+ \ell_i^-$  decays implies the existence of the lepton flavor violating (LFV) processes  $B \rightarrow K \ell_i^+ \ell_j^-$  (with  $i \neq j$ ).

$$-\frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \frac{e^2}{16\pi^2} \sum_i (C_i \mathcal{O}_i + C'_i \mathcal{O}'_i)$$

Physics Letters B 750 (2015) 367–371

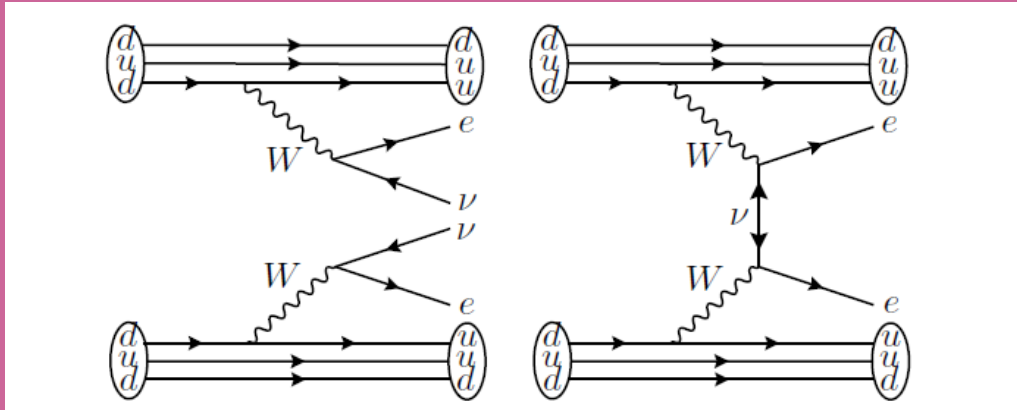


Boucenna, Vicente JWF Valle

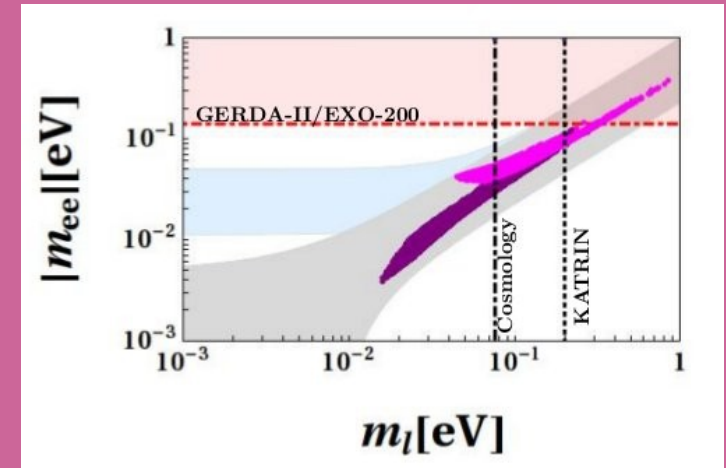
Here we raise the following question: can the leptonic mixing matrix provide the required lepton flavor structure in  $\mathcal{O}_9$  and  $\mathcal{O}_{10}$ ?

# Neutrinoless Double Beta Decay and flavor

A.S. Barabash arXiv:1104.2714



Family symmetry dependent lower bound

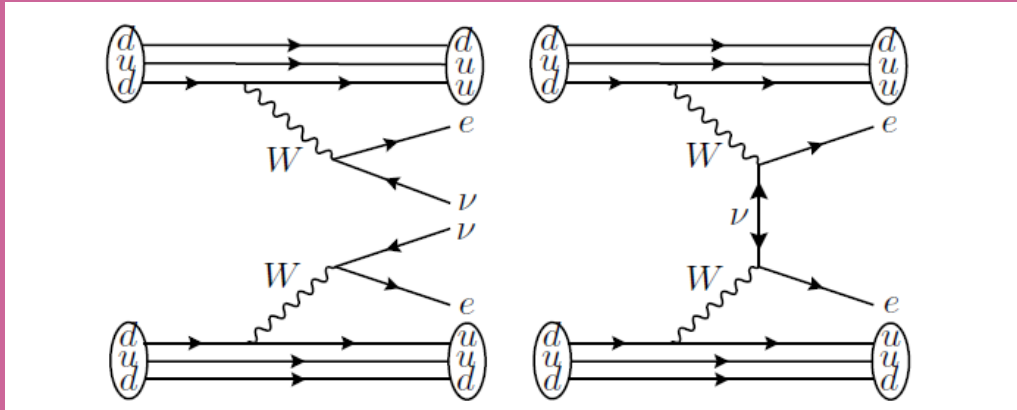


Bonilla et al arXiv:1411.4883

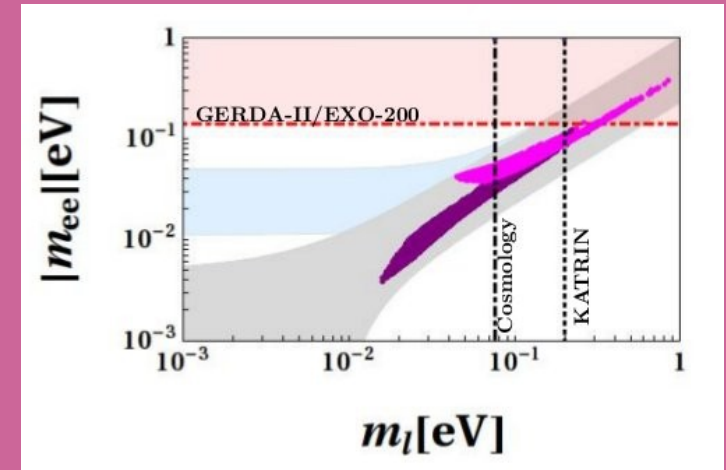


# Neutrinoless Double Beta Decay and flavor

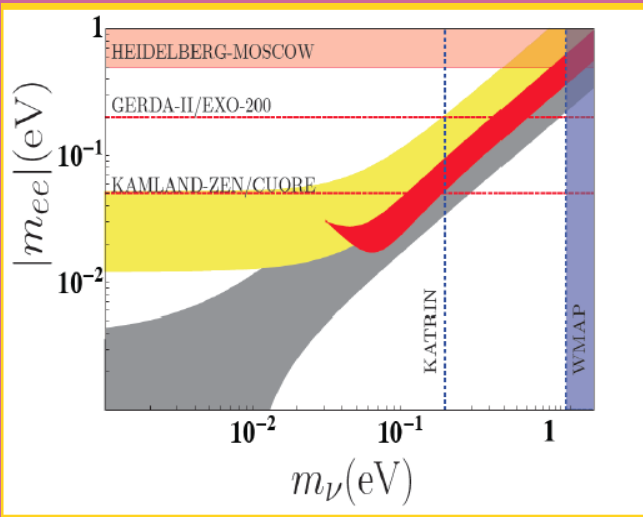
A.S. Barabash arXiv:1104.2714



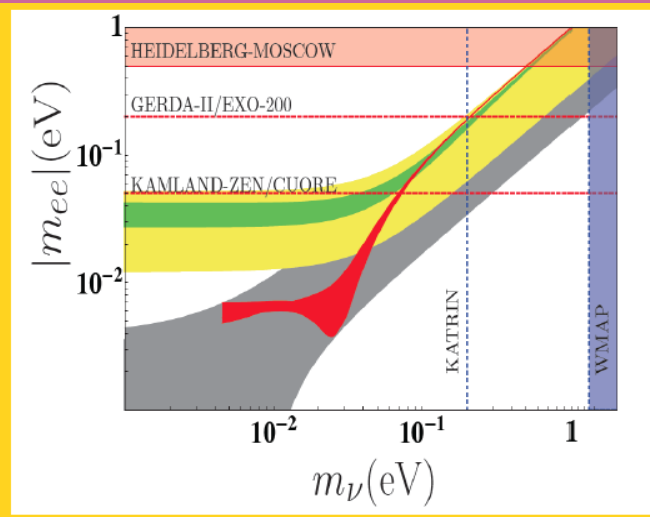
Family symmetry dependent lower bound



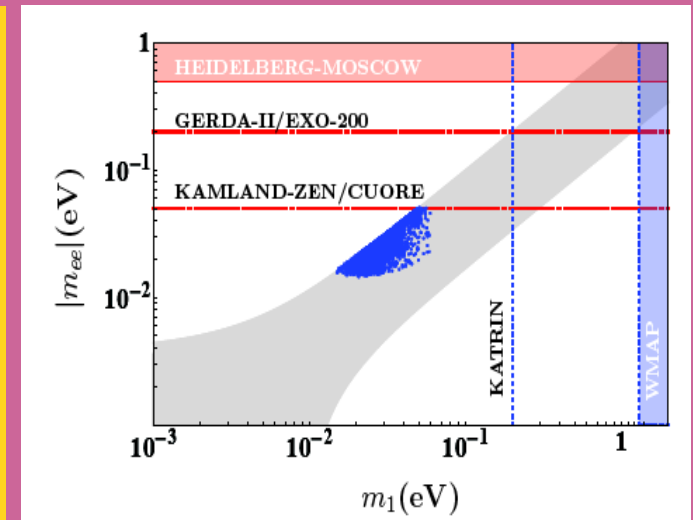
Bonilla et al arXiv:1411.4883



Dorame et al  
NPB861 (2012) 259-270

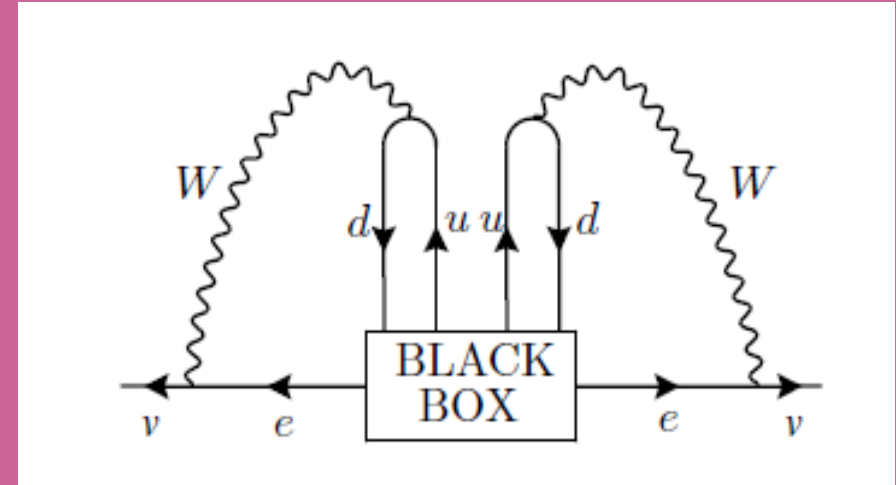


PhysRevD.86.056001  
JWF Valle



King et al Phys. Lett. B 724 (2013) 68

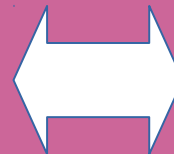
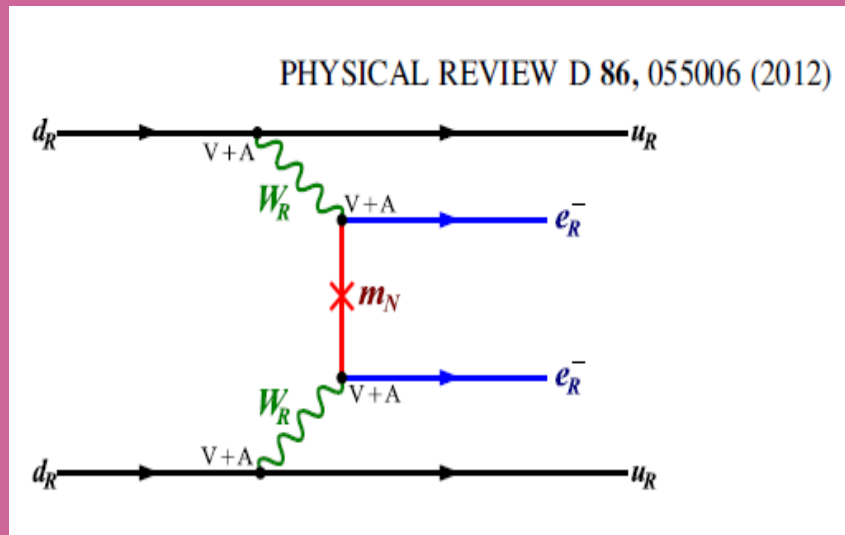
# Neutrinoless Double Beta Decay and COLLIDERS



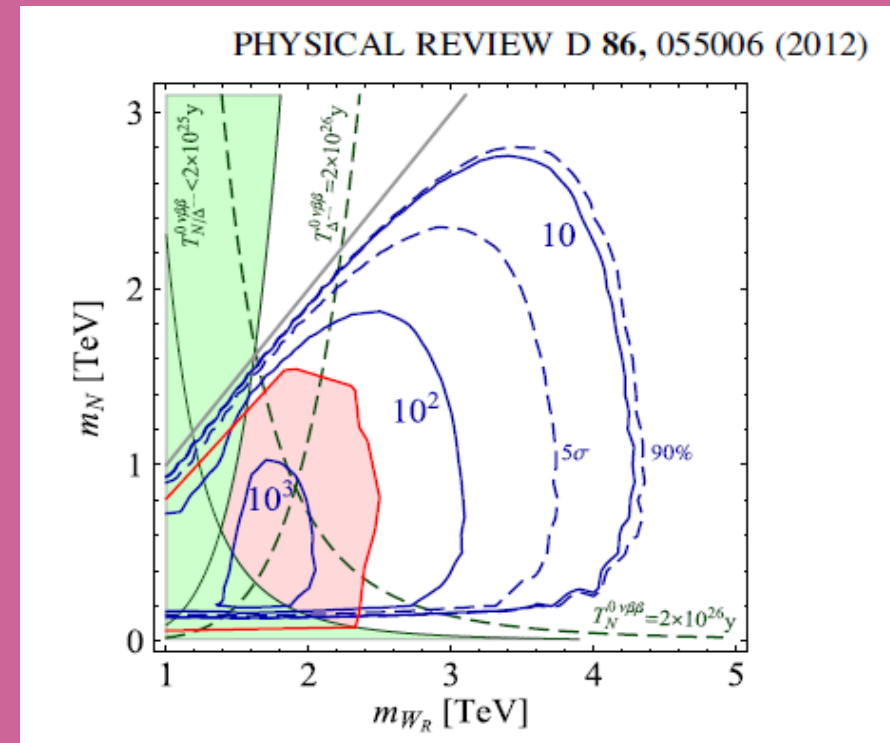
Schechter, JWFV 82

Lindner et al JHEP 1106 (2011) 091

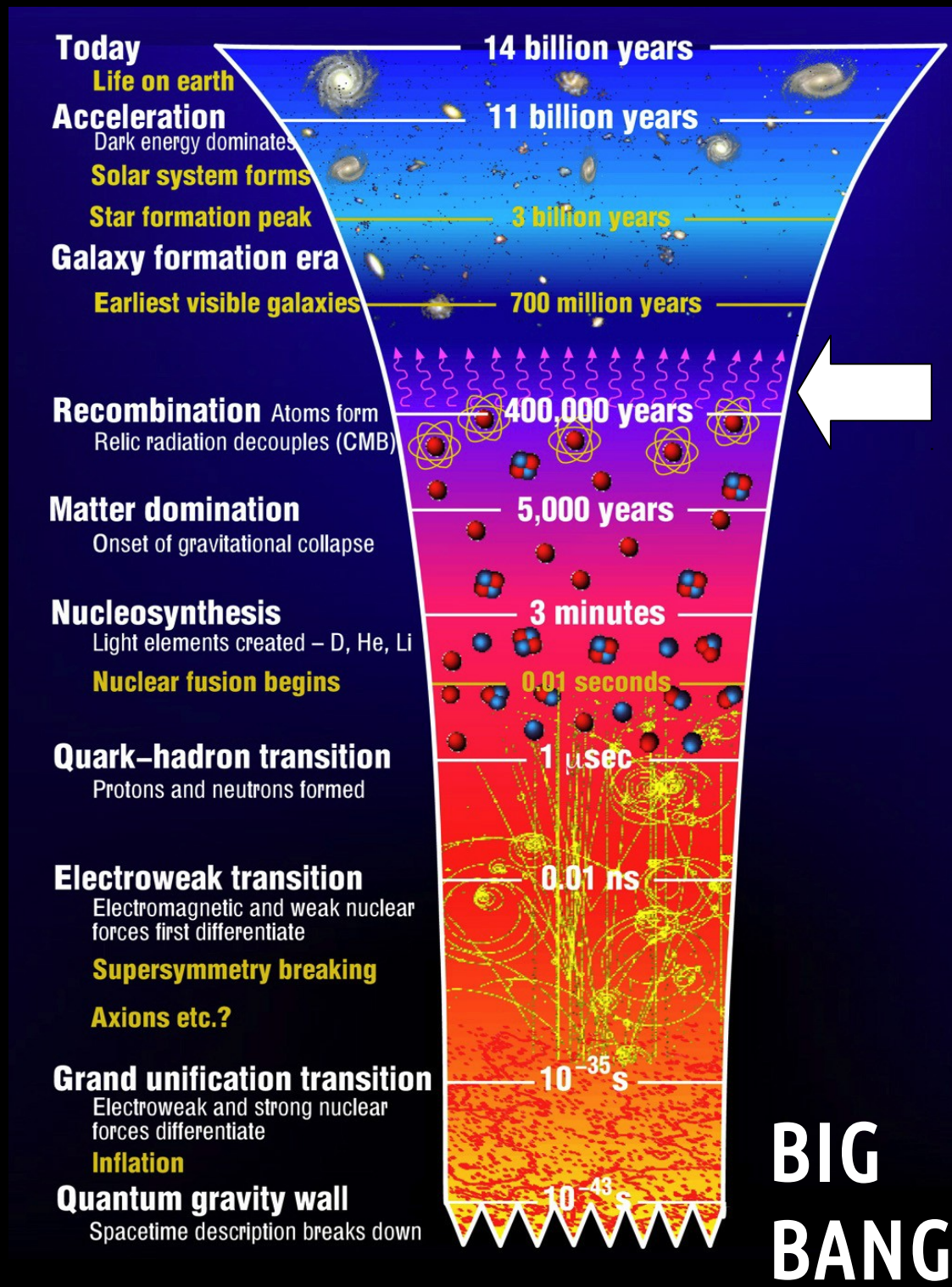
## Short versus long-range and the LHC



JWF Valle

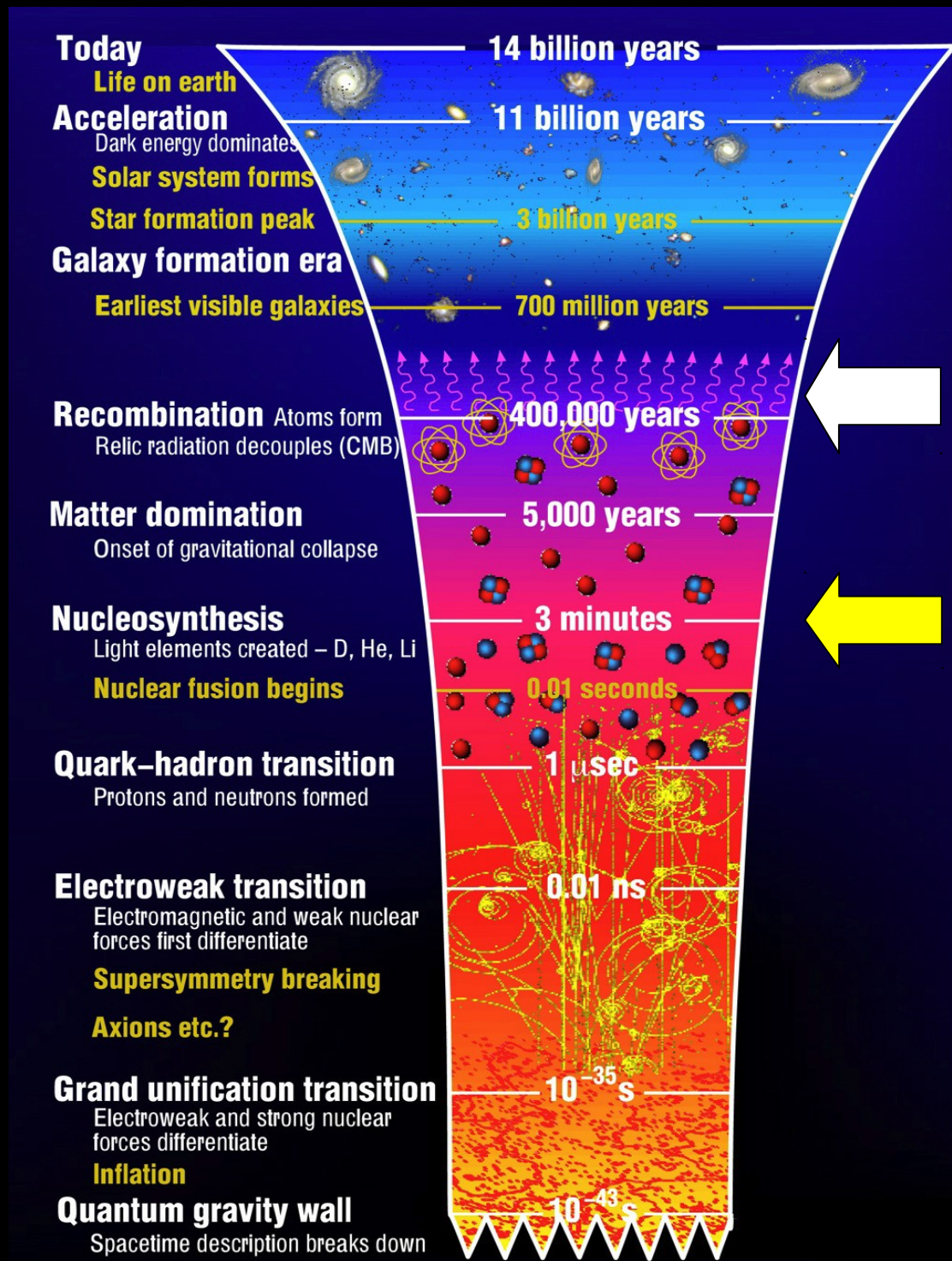


Neutrinos affect the CMB  
and large scale structure  
in the Universe ...



Neutrinos affect the CMB and large scale structure in the Universe ...

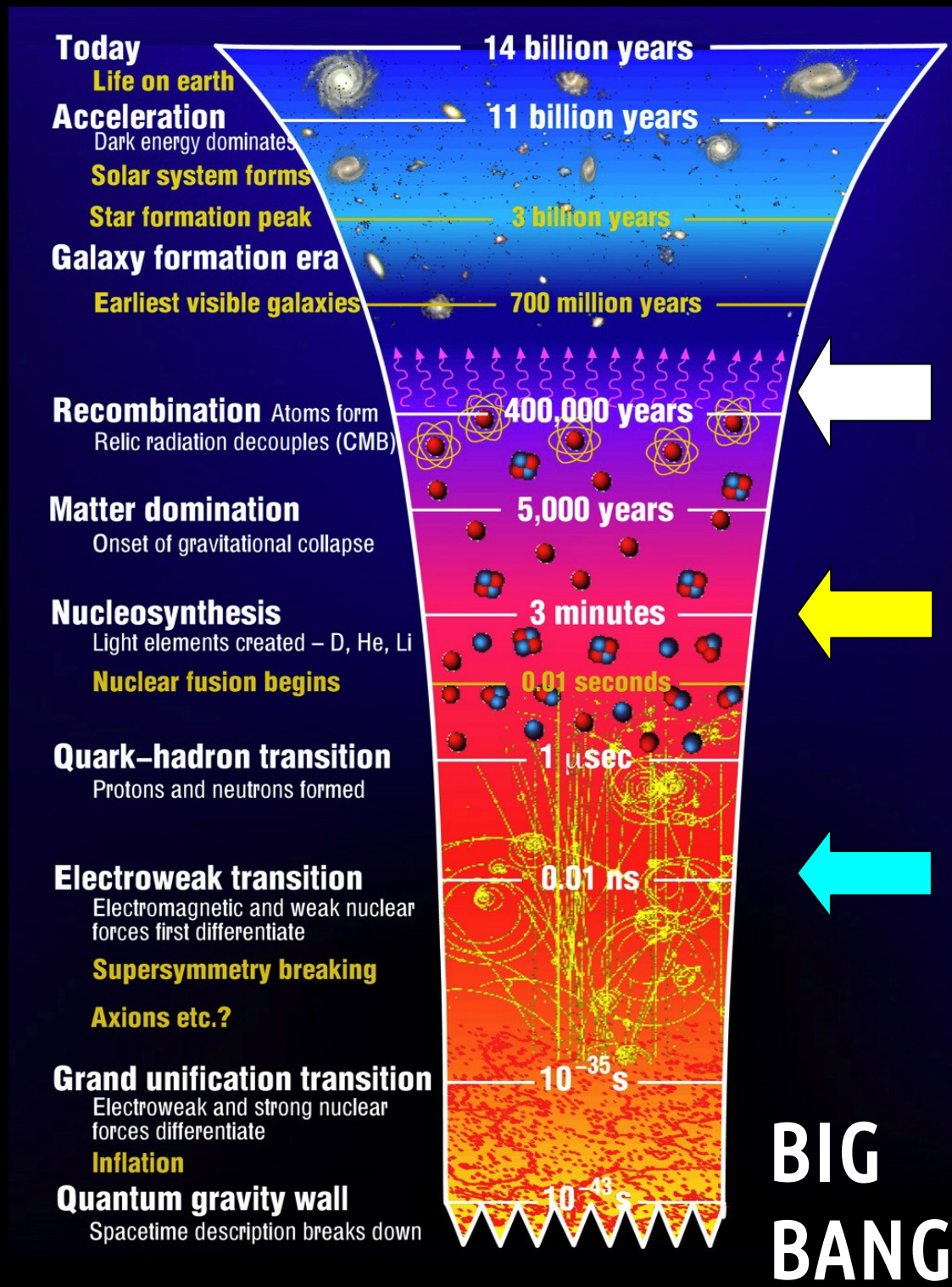
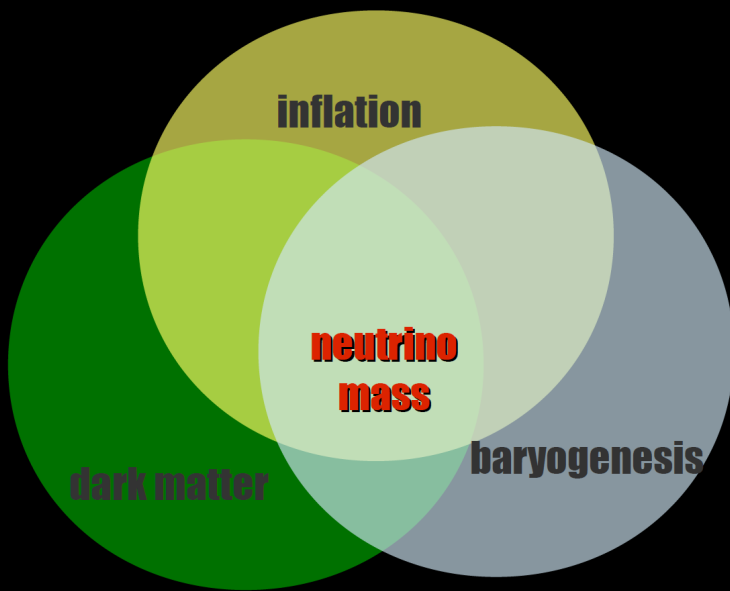
are key in the synthesis of light elements



Neutrinos affect the CMB and large scale structure in the Universe ...

are key in the synthesis of light elements

can “probe” the Universe earlier than photons ...



# SEESAW INFLATION & MAJORON DARK MATTER

$$\sigma = \frac{1}{\sqrt{2}} (\langle \sigma \rangle + \rho + iJ)$$

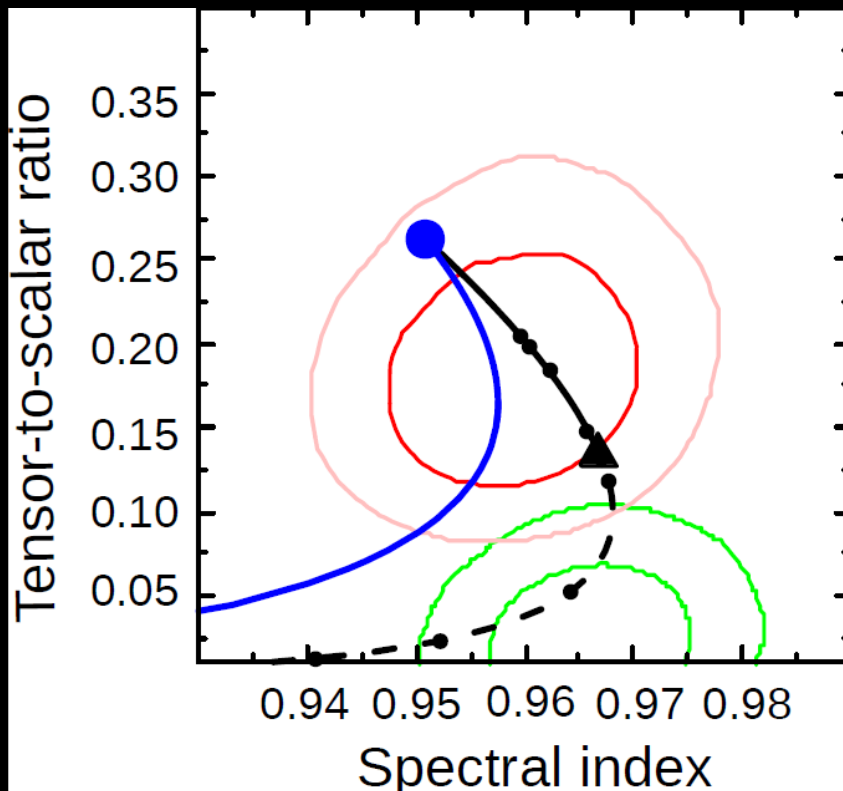
NEUTRINO MASSES

DARK MATTER

INFLATON

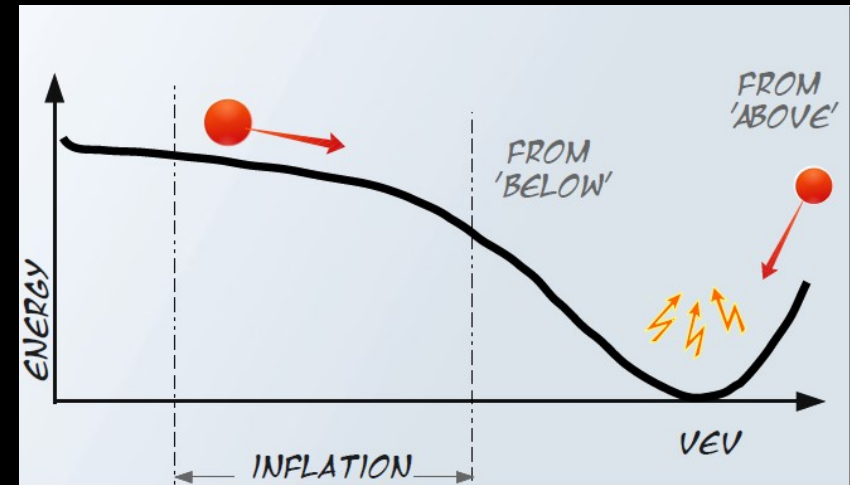
Boucenna et al arXiv:1405.2332 PRD90 (2014) 055023

Quartic versus Higgs Inflation



## type-I seesaw Leptogenesis

Aristizabal et al arXiv:1405.4706



<http://arxiv.org/pdf/1502.00612v1>

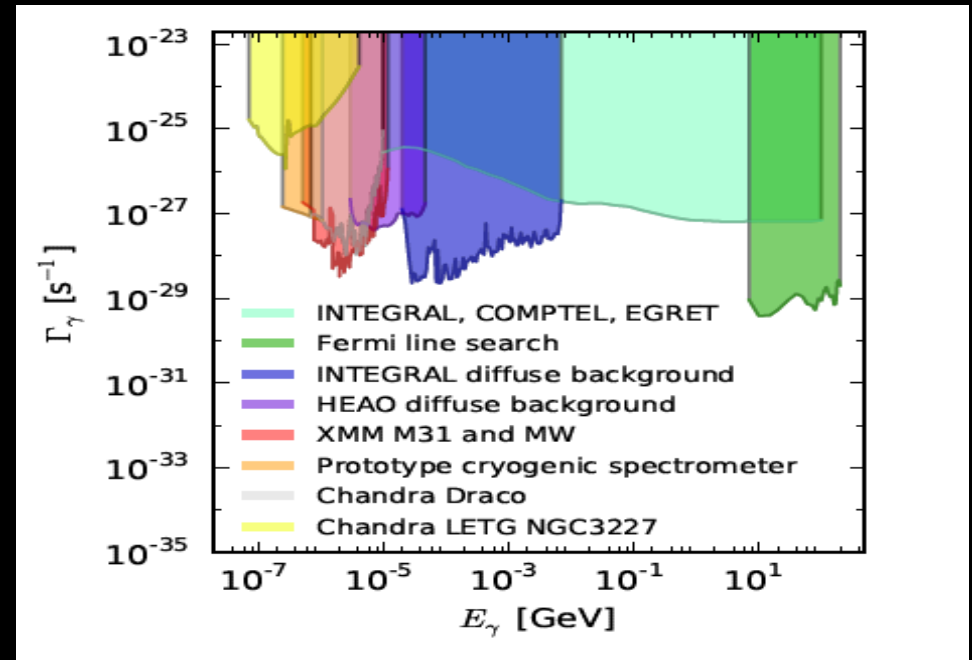
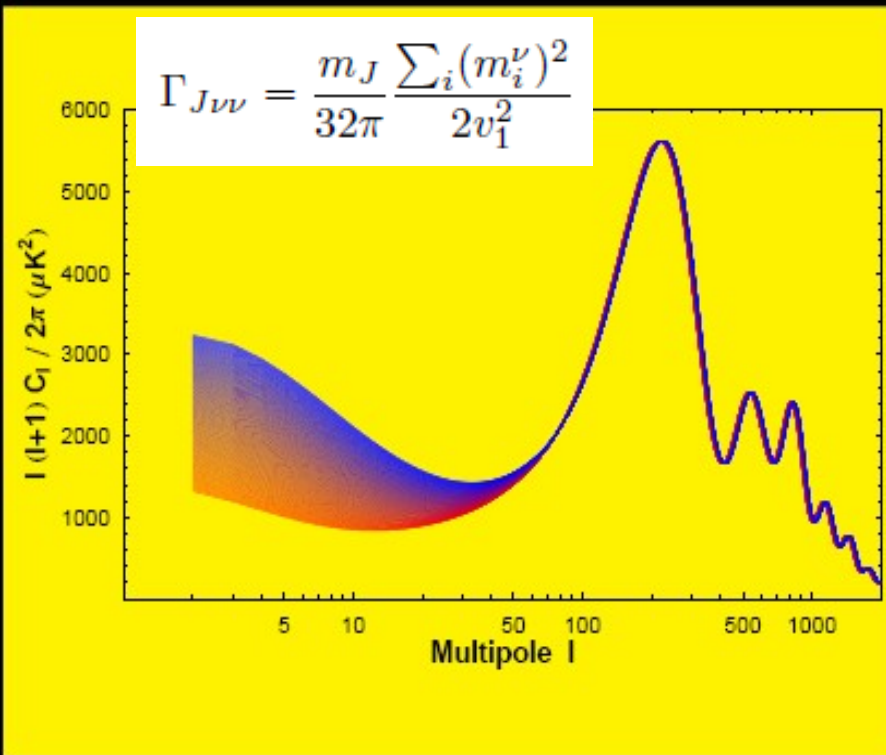
# DARK MATTER MAJORONS

Berezinsky, Valle PLB318 (1993) 360

## Consistency with CMB

Lattanzi & Valle, PRL99 (2007) 121301

$J \rightarrow \gamma\gamma$



Lattanzi et al PRD88 (2013) 063528

Esteves et al, PRD 82, 073008 (2010)

Bazzocchi & al JCAP 0808 (2008) 013

**neutrino  
masses**

**consistency of vacuum  
flavour  
coupling unification  
gravity**

**dark matter  
baryon asymmetry  
Inflation**

**First nail ???**

**neutrino as pathfinder**



# *Thank you*

Apart from **Miranda** and his  
Group at CINVESTAV ...

**Alfredo Aranda**

**Cesar Bonilla**

**Luis Dorame**

Felix Gonzalez

Eduardo Peinado

Alma Rojas

Carlos Vaquera

<http://t.co/yUSLiF1cGX>

<fb.me/2ZAD7khZf>

*Thank you*

Apart from [Miranda](#) and his Group at CINVESTAV ...

[Alfredo Aranda](#)

[Cesar Bonilla](#)

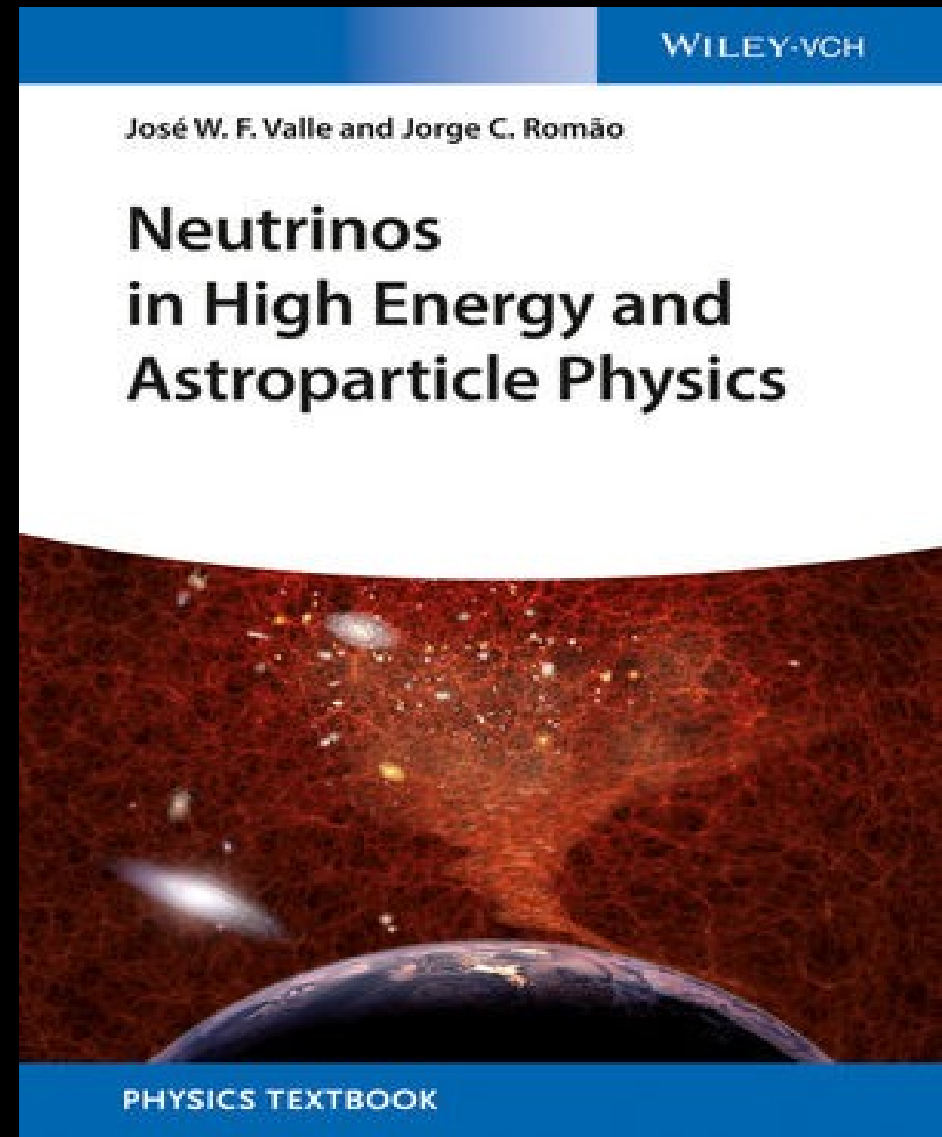
[Luis Dorame](#)

Felix Gonzalez

Eduardo Peinado

Alma Rojas

Carlos Vaquera



ISBN: 978-3-527-41197-9

458 pages 2015