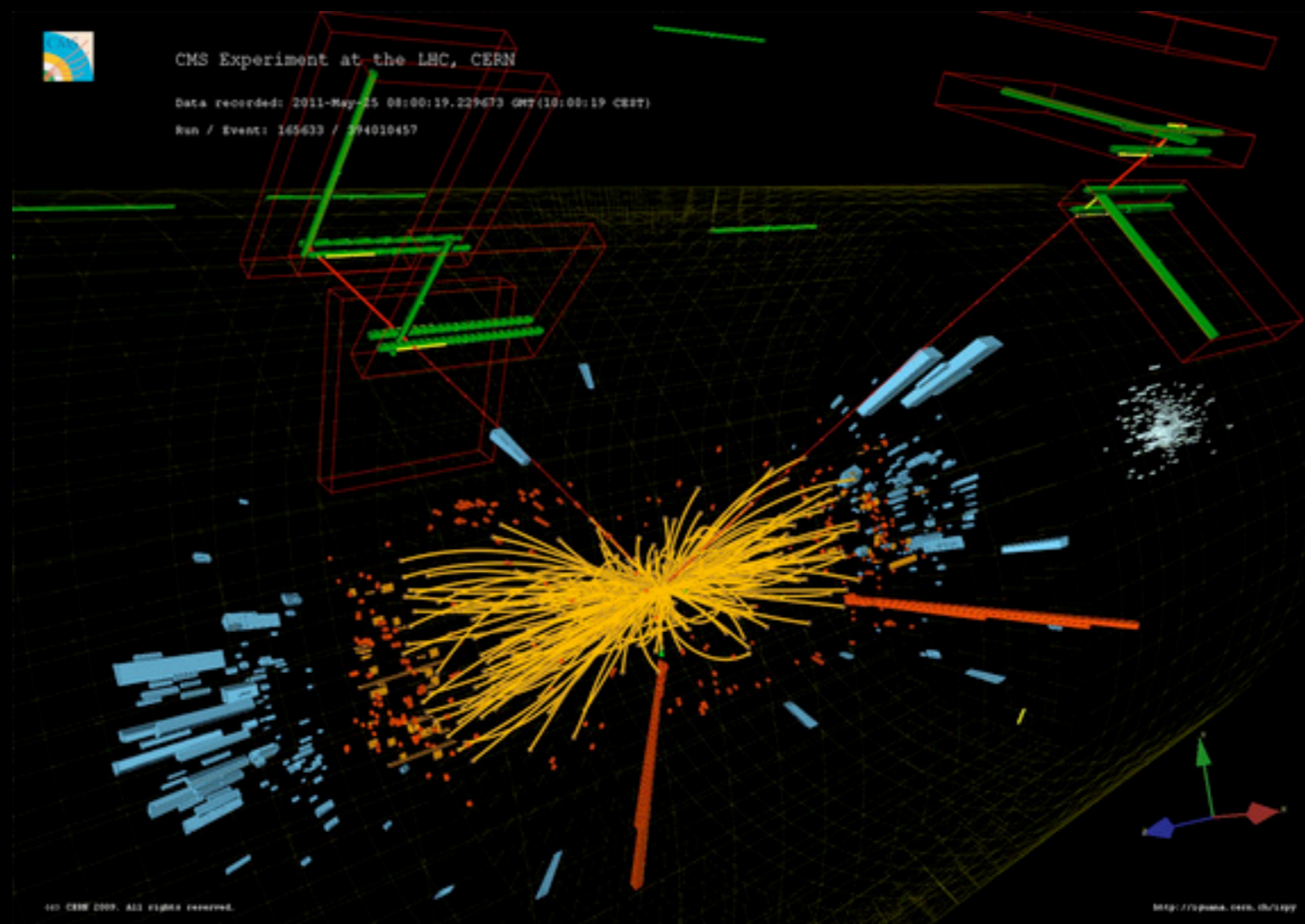




El bosón Higgs en teoría y fenomenología



Jens Erler

Seminario de Altas Energías (ICN y IF-UNAM)

25 de Abril 2012

Table of the Elementary Particles

ν_τ $s=1/2$ ~ 0	τ^- $s=1/2$ 1.9075	τ^+ $s=1/2$ 1.9075	t $s=1/2$ 176	t $s=1/2$ 176	t $s=1/2$ 176	\bar{t} $s=1/2$ 176	\bar{t} $s=1/2$ 176	\bar{t} $s=1/2$ 176	b $s=1/2$ 4.5	b $s=1/2$ 4.5	b $s=1/2$ 4.5	\bar{b} $s=1/2$ 4.5	\bar{b} $s=1/2$ 4.5	\bar{b} $s=1/2$ 4.5	
ν_μ $s=1/2$ ~ 0	μ^- $s=1/2$ 0.11343	μ^+ $s=1/2$ 0.11343	c $s=1/2$ 1.4	c $s=1/2$ 1.4	c $s=1/2$ 1.4	\bar{c} $s=1/2$ 1.4	\bar{c} $s=1/2$ 1.4	\bar{c} $s=1/2$ 1.4	s $s=1/2$ 0.1	s $s=1/2$ 0.1	s $s=1/2$ 0.1	\bar{s} $s=1/2$ 0.1	\bar{s} $s=1/2$ 0.1	\bar{s} $s=1/2$ 0.1	
ν_e $s=1/2$ ~ 0	e^- $s=1/2$ 0.00055	e^+ $s=1/2$ 0.00055	u $s=1/2$ 0.003	u $s=1/2$ 0.003	u $s=1/2$ 0.003	\bar{u} $s=1/2$ 0.003	\bar{u} $s=1/2$ 0.003	\bar{u} $s=1/2$ 0.003	d $s=1/2$ 0.005	d $s=1/2$ 0.005	d $s=1/2$ 0.005	\bar{d} $s=1/2$ 0.005	\bar{d} $s=1/2$ 0.005	\bar{d} $s=1/2$ 0.005	
H $s=0$ 134	H^\pm $s=0$ 86.3 ξ	Z $s=1$ 97.9	W^- $s=1$ 86.3	W^+ $s=1$ 86.3	g $s=1$ 0	g $s=1$ 0	g $s=1$ 0	g $s=1$ 0	g $s=1$ 0	g $s=1$ 0	g $s=1$ 0	g $s=1$ 0	g $s=1$ 0	γ $s=1$ 0	G $s=2$ 0

Table of the Elementary Particles

ν_τ $s=1/2$ ~ 0	τ^- $s=1/2$ 1.9075	τ^+ $s=1/2$ 1.9075	t $s=1/2$ 176	t $s=1/2$ 176	t $s=1/2$ 176	\bar{t} $s=1/2$ 176	\bar{t} $s=1/2$ 176	\bar{t} $s=1/2$ 176	b $s=1/2$ 4.5	b $s=1/2$ 4.5	b $s=1/2$ 4.5	\bar{b} $s=1/2$ 4.5	\bar{b} $s=1/2$ 4.5	\bar{b} $s=1/2$ 4.5	
ν_μ $s=1/2$ ~ 0	μ^- $s=1/2$ 0.11343	μ^+ $s=1/2$ 0.11343	c $s=1/2$ 1.4	c $s=1/2$ 1.4	c $s=1/2$ 1.4	\bar{c} $s=1/2$ 1.4	\bar{c} $s=1/2$ 1.4	\bar{c} $s=1/2$ 1.4	s $s=1/2$ 0.1	s $s=1/2$ 0.1	s $s=1/2$ 0.1	\bar{s} $s=1/2$ 0.1	\bar{s} $s=1/2$ 0.1	\bar{s} $s=1/2$ 0.1	
ν_e $s=1/2$ ~ 0	e^- $s=1/2$ 0.00055	e^+ $s=1/2$ 0.00055	u $s=1/2$ 0.003	u $s=1/2$ 0.003	u $s=1/2$ 0.003	\bar{u} $s=1/2$ 0.003	\bar{u} $s=1/2$ 0.003	\bar{u} $s=1/2$ 0.003	d $s=1/2$ 0.005	d $s=1/2$ 0.005	d $s=1/2$ 0.005	\bar{d} $s=1/2$ 0.005	\bar{d} $s=1/2$ 0.005	\bar{d} $s=1/2$ 0.005	
H $s=0$ 125	H^\pm $s=0$ 86.3 ξ	Z $s=1$ 97.9	W^- $s=1$ 86.3	W^+ $s=1$ 86.3	g $s=1$ 0	g $s=1$ 0	g $s=1$ 0	g $s=1$ 0	g $s=1$ 0	g $s=1$ 0	g $s=1$ 0	g $s=1$ 0	g $s=1$ 0	γ $s=1$ 0	G $s=2$ 0

Outline

Outline

- the weak interaction

Outline

- the weak interaction
- spontaneous symmetry breaking

Outline

- the weak interaction
- spontaneous symmetry breaking
- theoretical constraints

Outline

- the weak interaction
- spontaneous symmetry breaking
- theoretical constraints
- indirect constraints

Outline

- the weak interaction
- spontaneous symmetry breaking
- theoretical constraints
- indirect constraints
- direct constraints

Outline

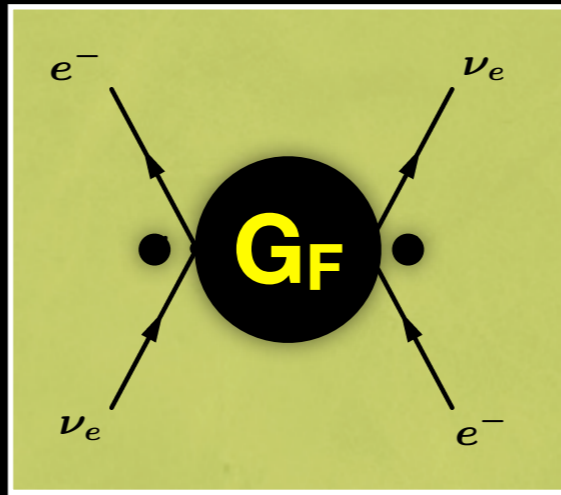
- the weak interaction
- spontaneous symmetry breaking
- theoretical constraints
- indirect constraints
- direct constraints
- synthesis

Outline

- the weak interaction
- spontaneous symmetry breaking
- theoretical constraints
- indirect constraints
- direct constraints
- synthesis
- conclusions

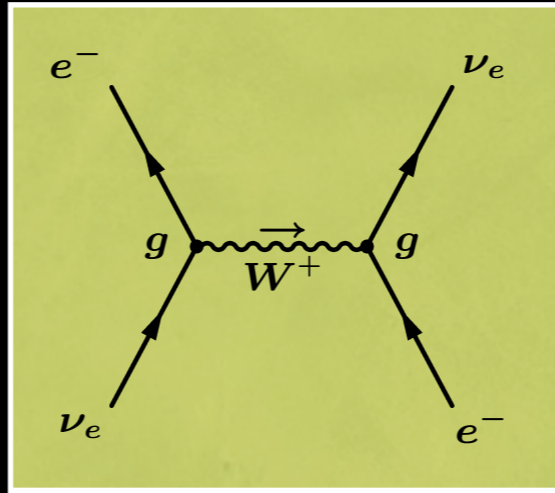
the weak interaction

Problems with Fermi theory



- $\sigma(\nu_e e^- \rightarrow e^- \nu_e) \rightarrow E^2 / (2\pi v^4) \leq 16\pi / E^2 \Rightarrow$
tree level unitarity violation for $E \gtrsim 1 \text{ TeV}$

Problems with Fermi theory

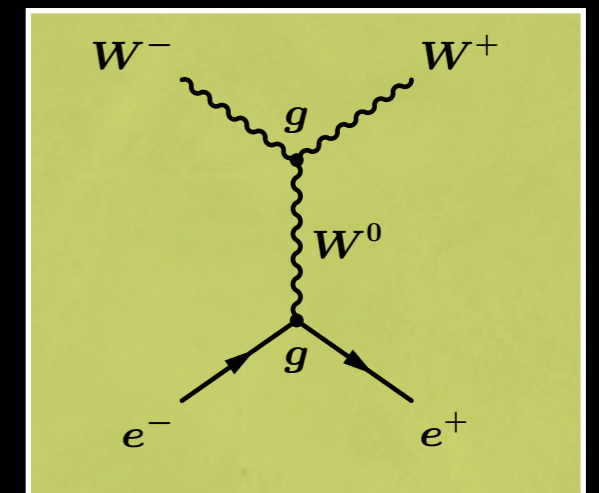
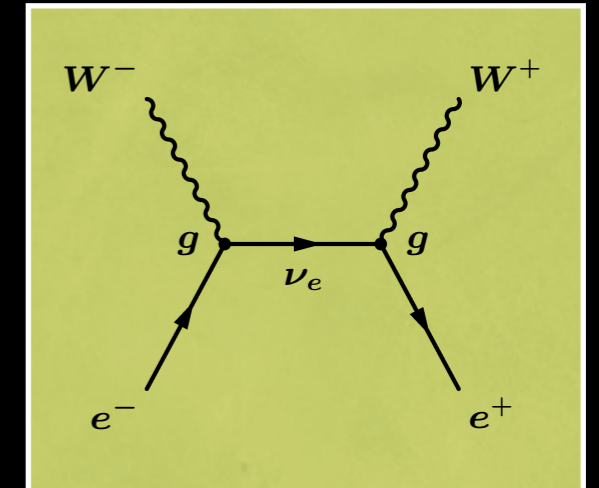


- $\sigma(\nu_e e^- \rightarrow e^- \nu_e) \rightarrow E^2 / (2\pi v^4) \leq 16\pi / E^2 \Rightarrow$
tree level unitarity violation for $E \gtrsim 1 \text{ TeV}$

➔ intermediate vector bosons *Yukawa 1935, Schwinger 1957*

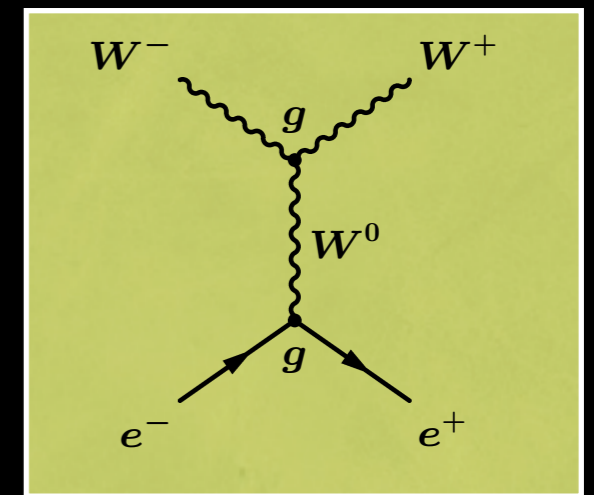
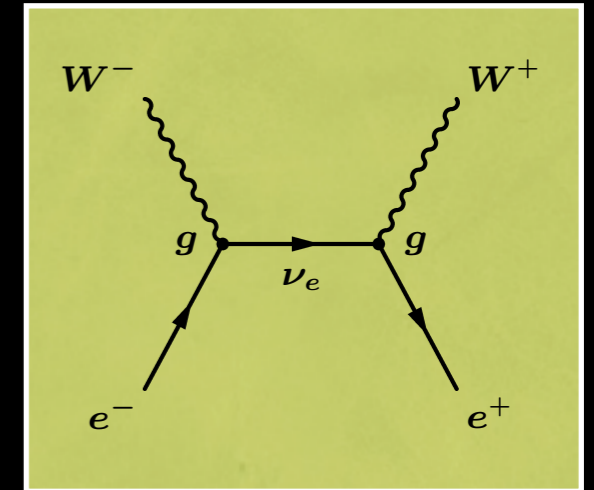
- but now trouble computing $\sigma(e^+ e^- \rightarrow W^+ W^-)$

Gauge invariance I



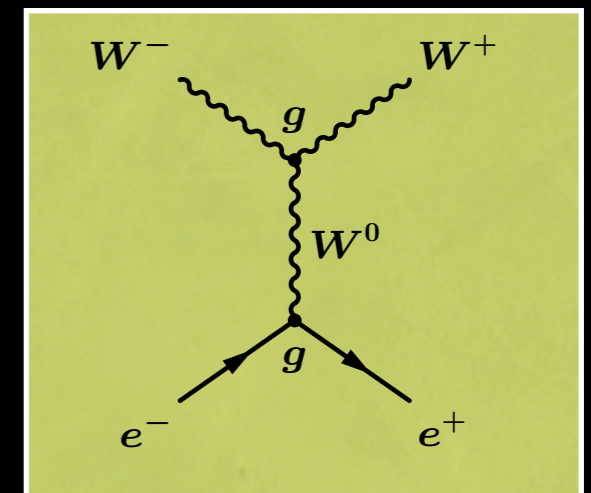
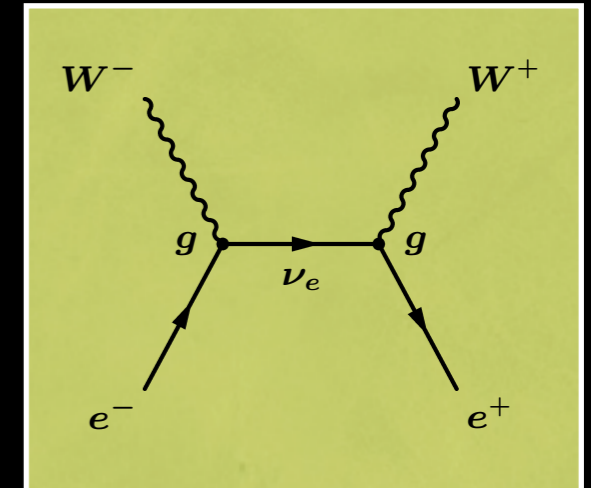
Gauge invariance I

- Bad behavior in $\sigma(e^+ e^- \rightarrow W^+ W^-)$ and $\sigma(\bar{\nu}_e e^- \rightarrow W^- W^0)$ cancels if the currents satisfy $[J, J^\dagger] \propto J^0 \Rightarrow$



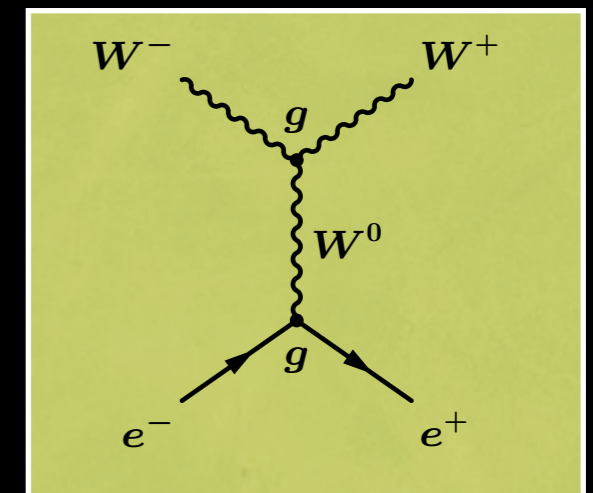
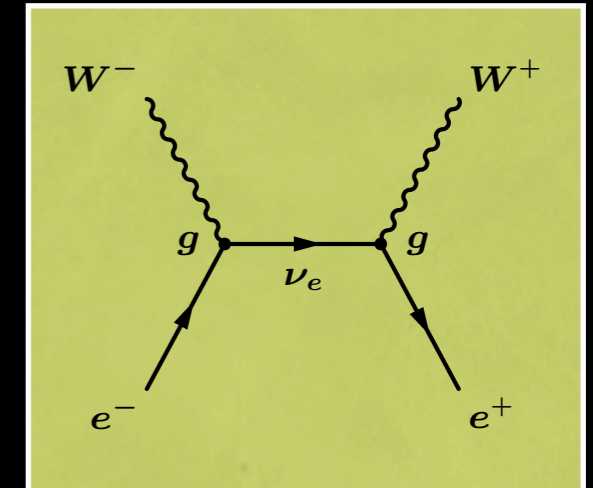
Gauge invariance I

- Bad behavior in $\sigma(e^+ e^- \rightarrow W^+ W^-)$ and $\sigma(\bar{\nu}_e e^- \rightarrow W^- W^0)$ cancels if the currents satisfy $[J, J^\dagger] \propto J^0 \Rightarrow$
- SU(2) gauge invariance



Gauge invariance I

- Bad behavior in $\sigma(e^+ e^- \rightarrow W^+ W^-)$ and $\sigma(\bar{\nu}_e e^- \rightarrow W^- W^0)$ cancels if the currents satisfy $[J, J^\dagger] \propto J^0 \Rightarrow$
- SU(2) gauge invariance
- Need to add photons
 $\Rightarrow SU(2)_L \times U(1)_Y$ *Glashow 1961*



Gauge invariance II

Gauge invariance II

- **Theorem** (*Weinberg 1964*): no 4-vector field $A_\mu(x)$ can be built from the a and a^\dagger for $m = 0$ and $h = \pm 1$ particles:

Gauge invariance II

- **Theorem** (*Weinberg 1964*): no 4-vector field $A_\mu(x)$ can be built from the a and a^\dagger for $m = 0$ and $h = \pm 1$ particles:
- $e_\mu(\mathbf{p}, \pm 1) \rightarrow \Lambda_{\mu\nu} e^\nu(\mathbf{p}, \pm 1) + p_\mu \Omega_\pm(\mathbf{p}, \Lambda)$

Gauge invariance II

- **Theorem** (*Weinberg 1964*): no 4-vector field $A_\mu(x)$ can be built from the a and a^\dagger for $m = 0$ and $h = \pm 1$ particles:
- $e_\mu(\mathbf{p}, \pm 1) \rightarrow \Lambda_{\mu\nu} e^\nu(\mathbf{p}, \pm 1) + p_\mu \Omega_\pm(\mathbf{p}, \Lambda)$
- $U(\Lambda) A_\mu(x) U^{-1}(\Lambda) = \Lambda_{\mu\nu} A^\nu(\Lambda x) + \partial_\mu \Omega(x, \Lambda)$, where Ω depends on a and a^\dagger

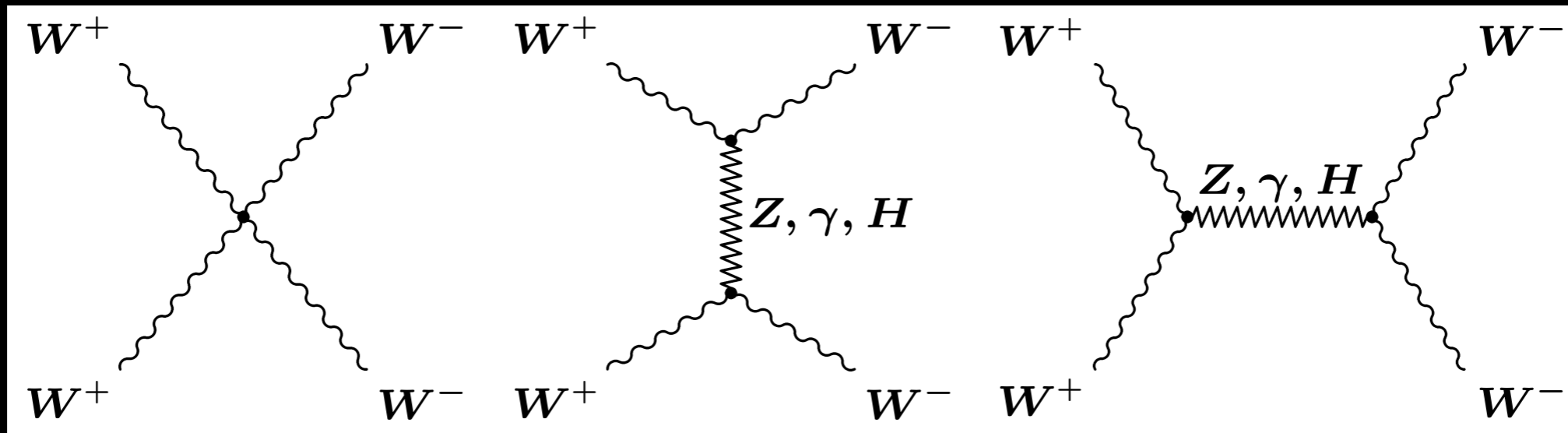
Gauge invariance II

- **Theorem** (*Weinberg 1964*): no 4-vector field $A_\mu(x)$ can be built from the a and a^\dagger for $m = 0$ and $h = \pm 1$ particles:
- $e_\mu(\mathbf{p}, \pm 1) \rightarrow \Lambda_{\mu\nu} e^\nu(\mathbf{p}, \pm 1) + p_\mu \Omega_\pm(\mathbf{p}, \Lambda)$
- $U(\Lambda) A_\mu(x) U^{-1}(\Lambda) = \Lambda_{\mu\nu} A^\nu(\Lambda x) + \partial_\mu \Omega(x, \Lambda)$, where Ω depends on a and a^\dagger
- **Solution**: take couplings of A_μ as $A_\mu J^\mu$ with $\partial_\mu J^\mu = 0$

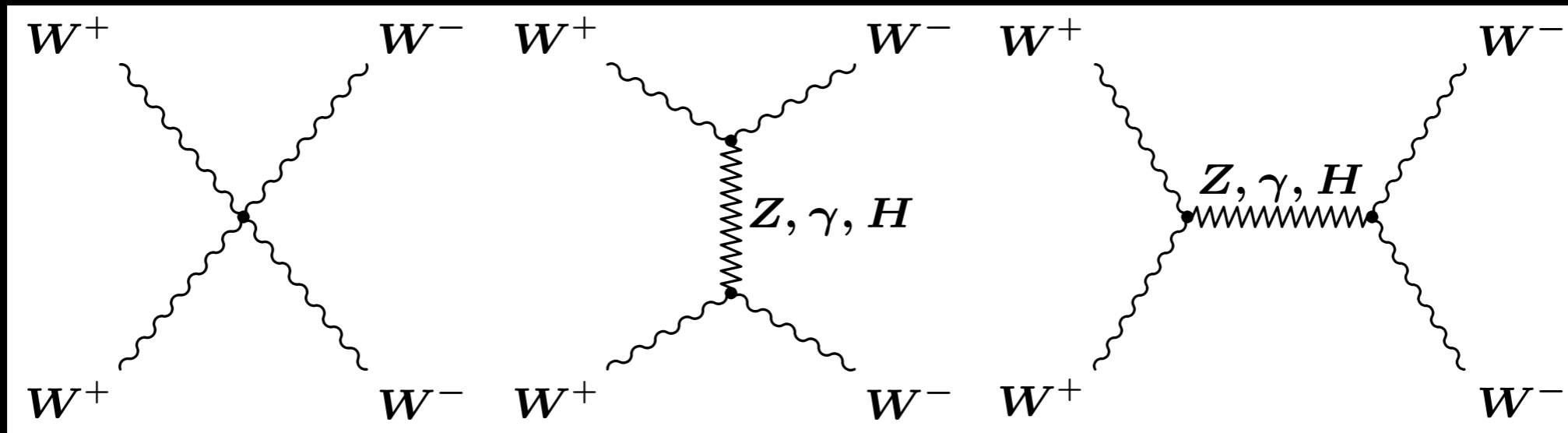
Gauge invariance II

- **Theorem** (*Weinberg 1964*): no 4-vector field $A_\mu(x)$ can be built from the a and a^\dagger for $m = 0$ and $h = \pm 1$ particles:
- $e_\mu(\mathbf{p}, \pm 1) \rightarrow \Lambda_{\mu\nu} e^\nu(\mathbf{p}, \pm 1) + p_\mu \Omega_\pm(\mathbf{p}, \Lambda)$
- $U(\Lambda) A_\mu(x) U^{-1}(\Lambda) = \Lambda_{\mu\nu} A^\nu(\Lambda x) + \partial_\mu \Omega(x, \Lambda)$, where Ω depends on a and a^\dagger
- **Solution:** take couplings of A_μ as $A_\mu J^\mu$ with $\partial_\mu J^\mu = 0$
- **Implications:**
 - breaking gauge invariance breaks Lorentz invariance
 - $m_A = 0$

Need for Higgs particles

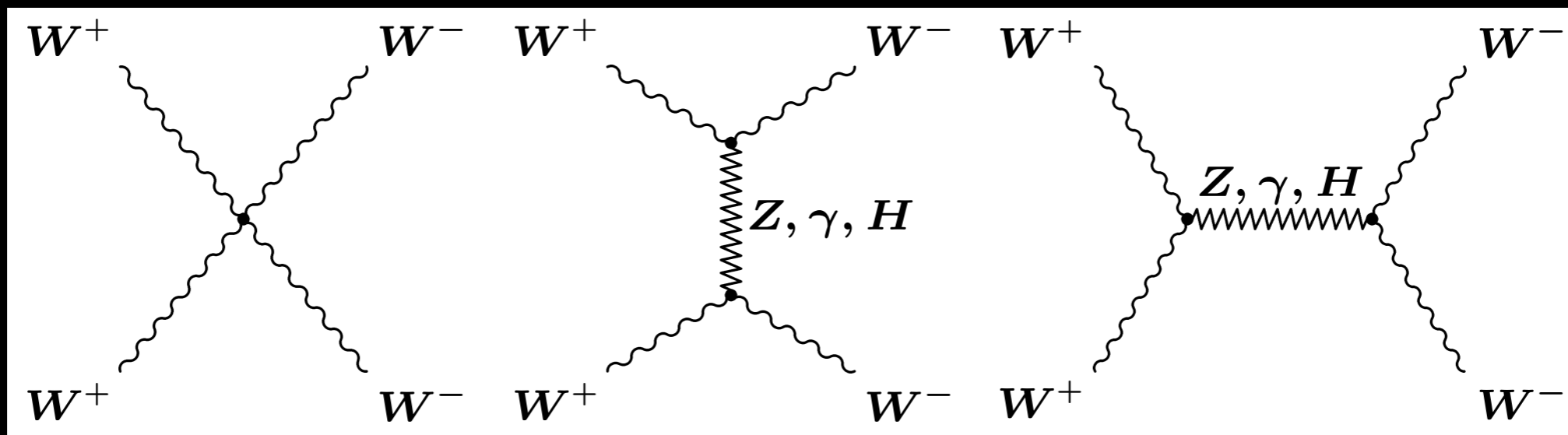


Need for Higgs particles



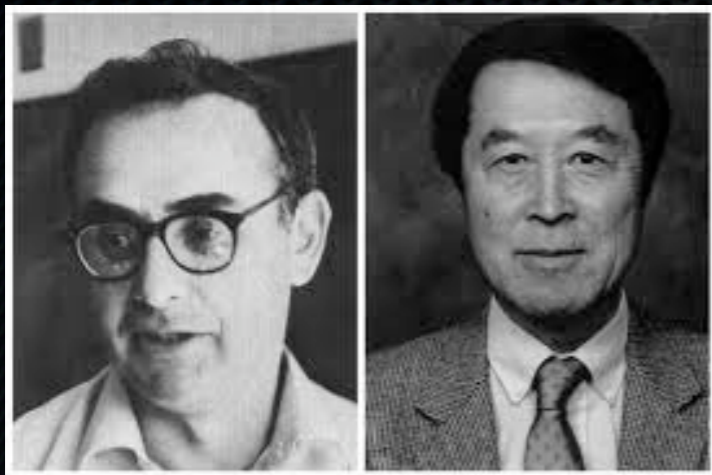
- After fixing $\sigma(e^+ e^- \rightarrow W^+ W^-)$ & $\sigma(\bar{\nu}_e e^- \rightarrow W^- W^0)$, now tree level unitarity violated in $\sigma(W^+ W^- \rightarrow W^+ W^-)$, unless the Higgs is introduced with

Need for Higgs particles



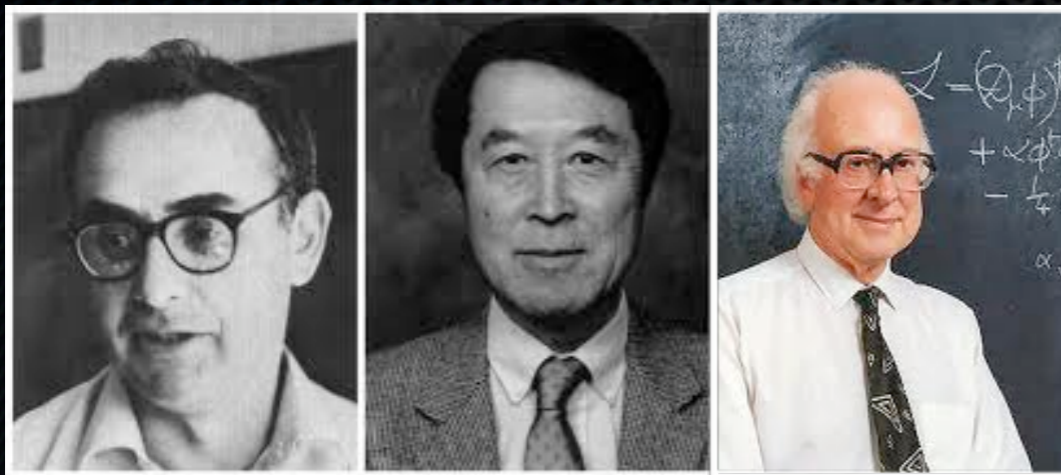
- After fixing $\sigma(e^+ e^- \rightarrow W^+ W^-)$ & $\sigma(\bar{\nu}_e e^- \rightarrow W^- W^0)$, now tree level unitarity violated in $\sigma(W^+ W^- \rightarrow W^+ W^-)$, unless the Higgs is introduced with
- $M_H^2 \lesssim 16\pi/5 v^2 \approx (780 \text{ GeV})^2$ *Lüscher, Weisz 1988*

spontaneous symmetry breaking

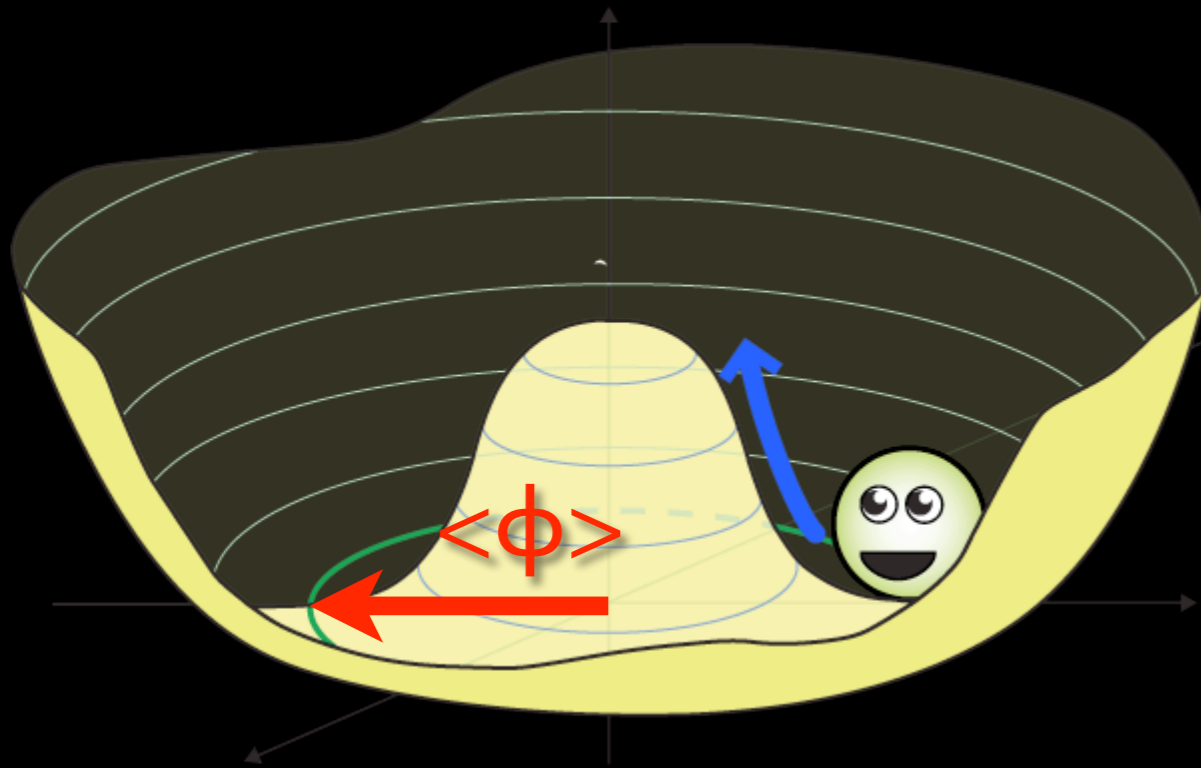


mechanisms of *Nambu-Goldstone* (Nobel Prize 2008) and
Brout-Englert-Higgs-Guralnik-Hagen-Kibble (Sakurai Prize 2010)

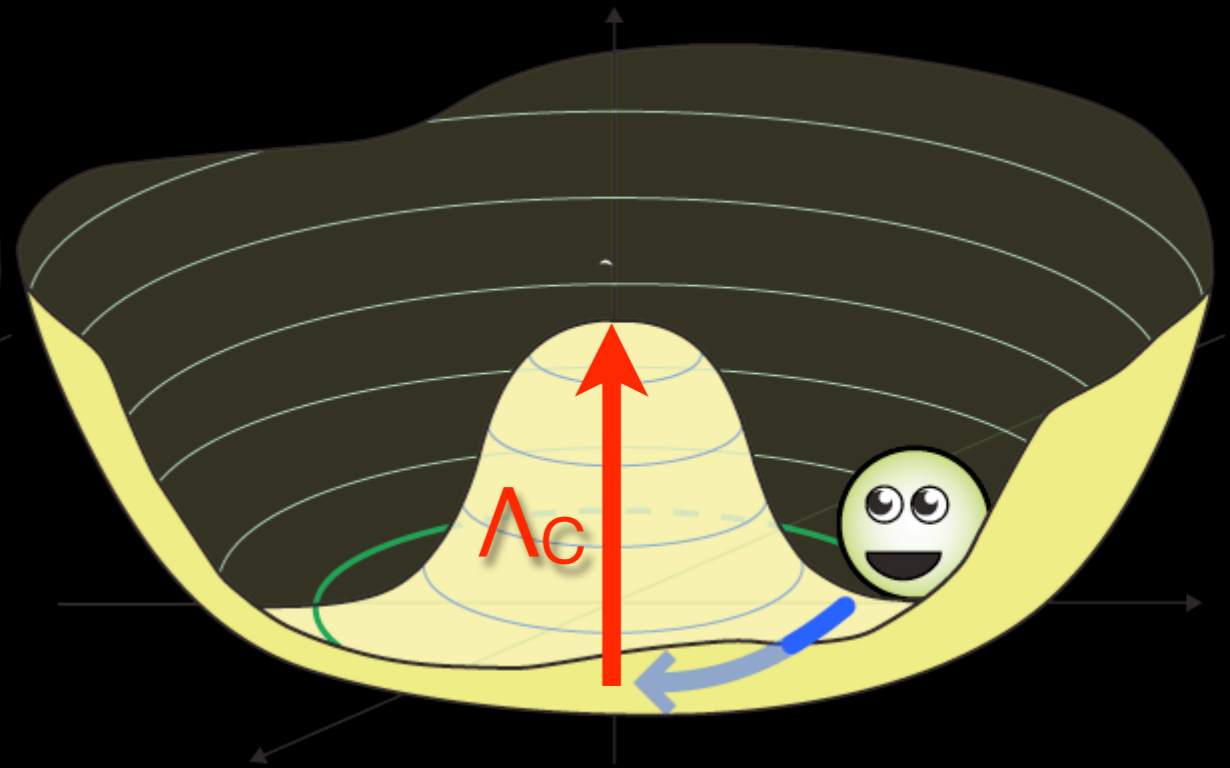
spontaneous symmetry breaking



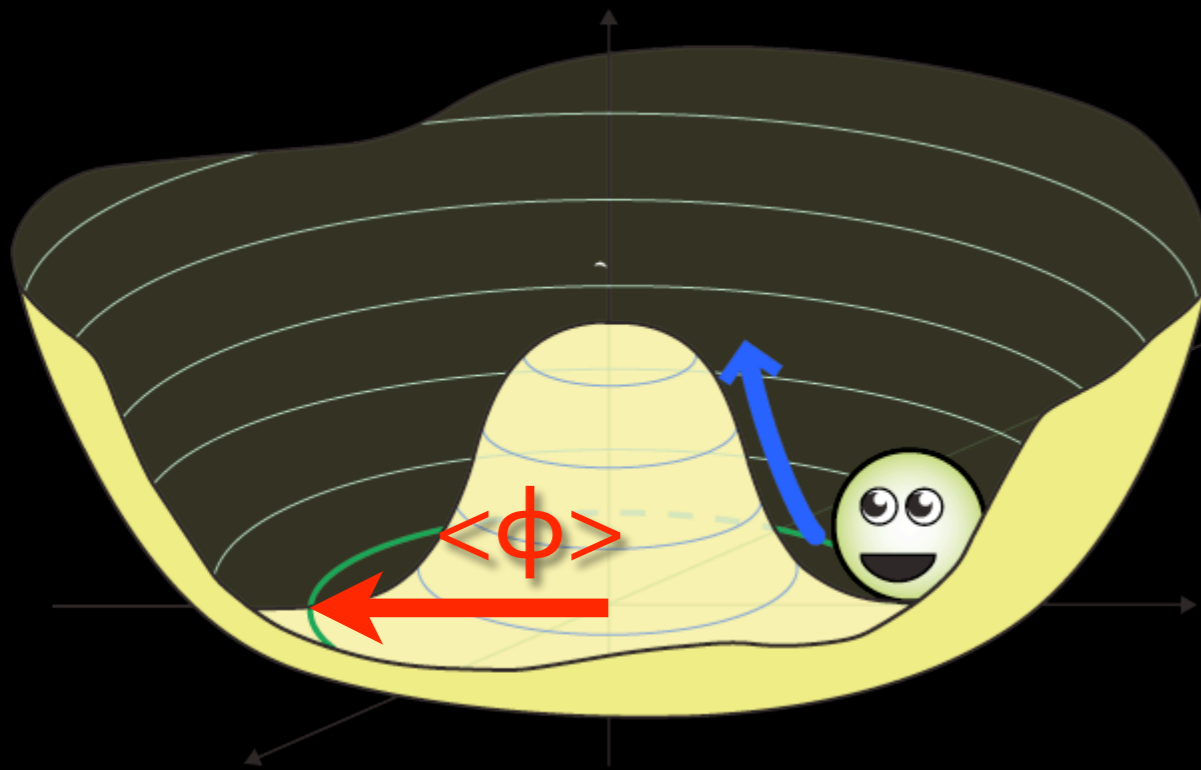
mechanisms of *Nambu-Goldstone* (Nobel Prize 2008) and
Brout-Englert-Higgs-Guralnik-Hagen-Kibble (Sakurai Prize 2010)



massive scalar
 also for discrete symmetry
 as in *Landau-Ginzburg* theory

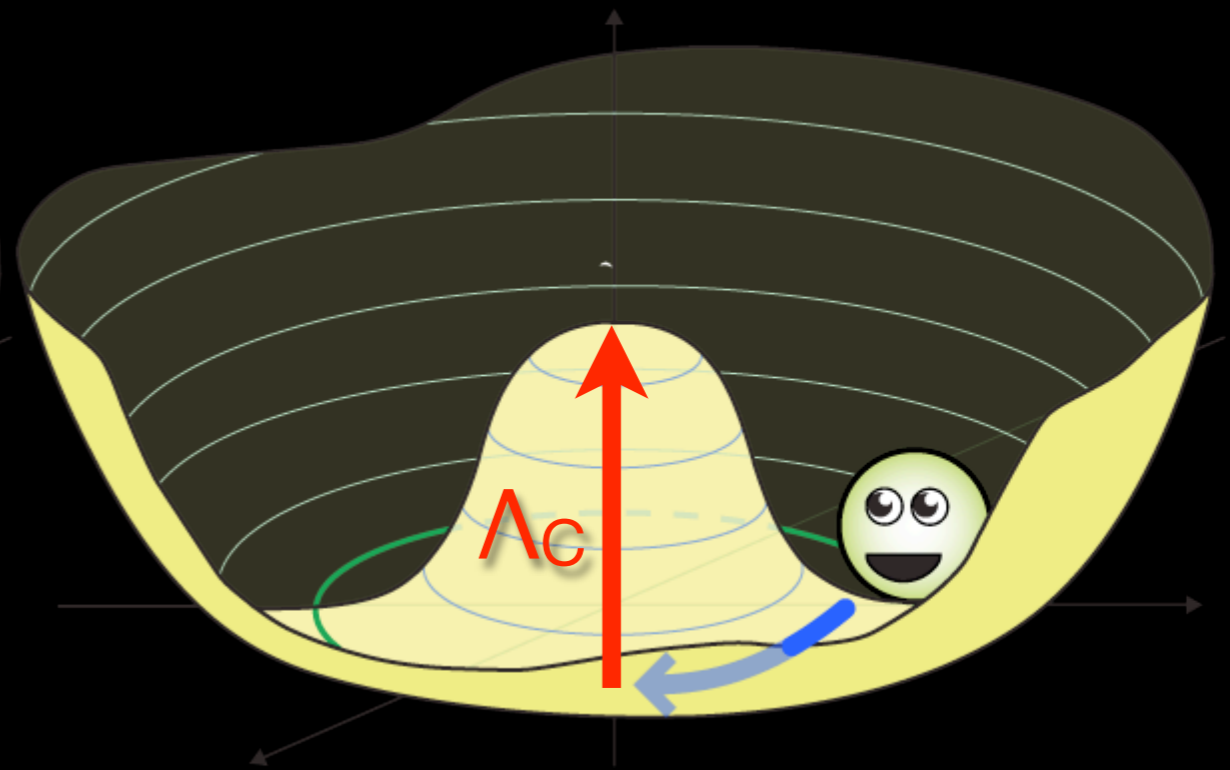


massless
Nambu-Goldstone mode
 ($V = 0$ exactly)



massive scalar

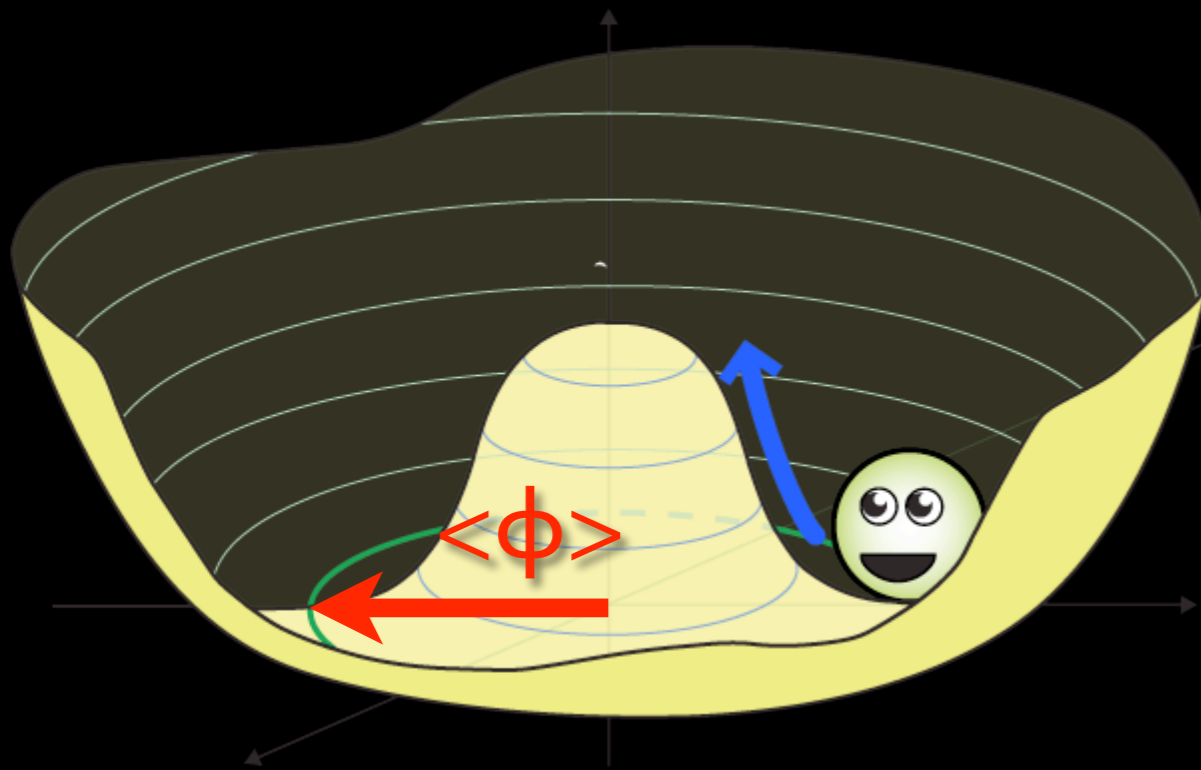
also for discrete symmetry
as in *Landau-Ginzburg* theory



massless

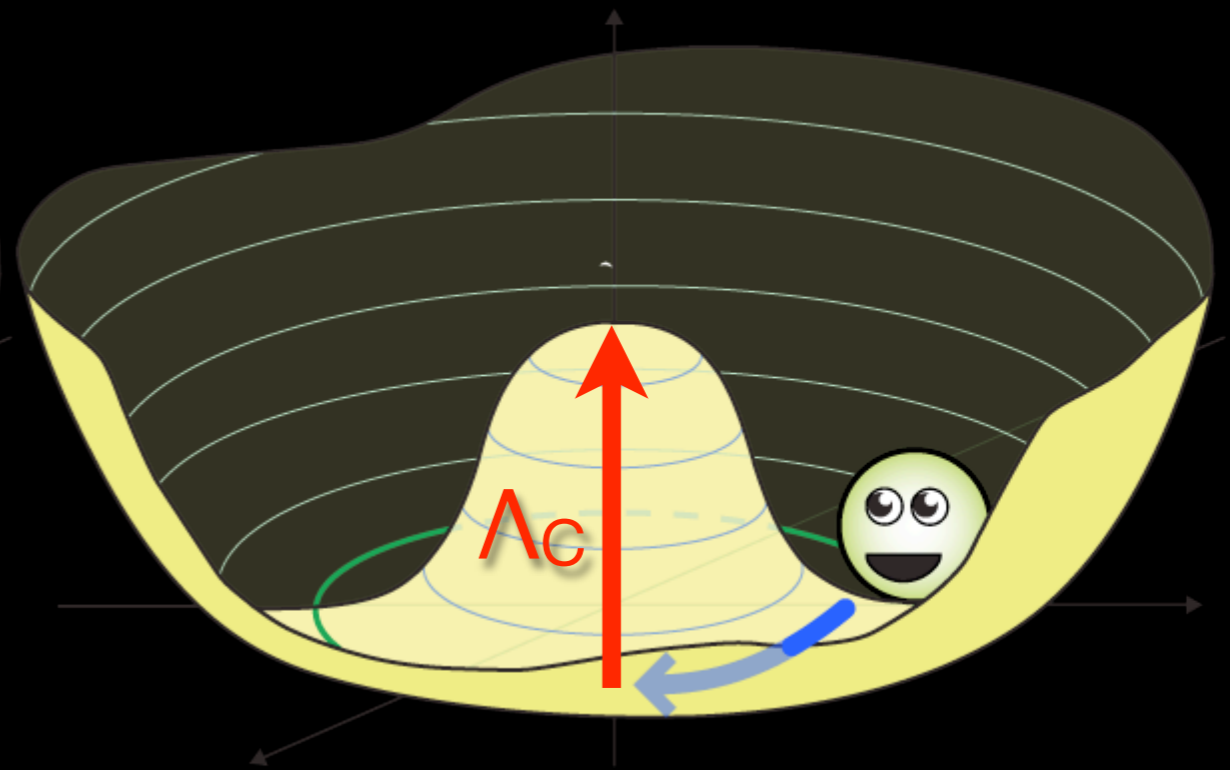
Nambu-Goldstone mode
($V = 0$ exactly)

- field gets vacuum expectation value $\langle \phi \rangle$ (order parameter)



massive scalar

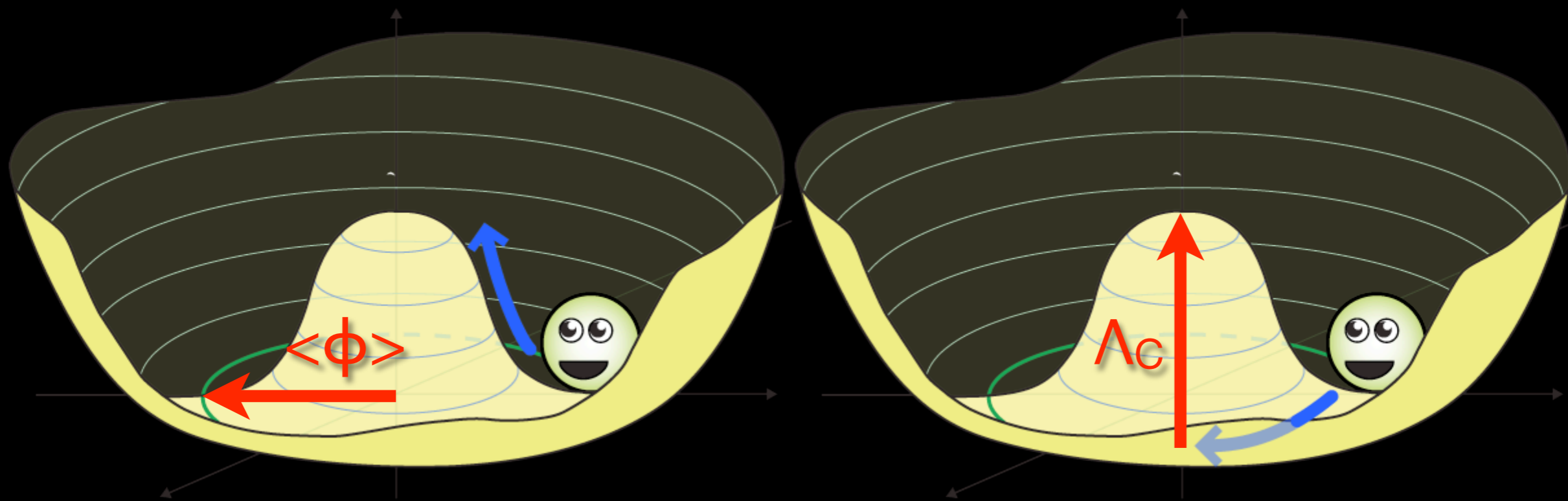
also for discrete symmetry
as in *Landau-Ginzburg* theory

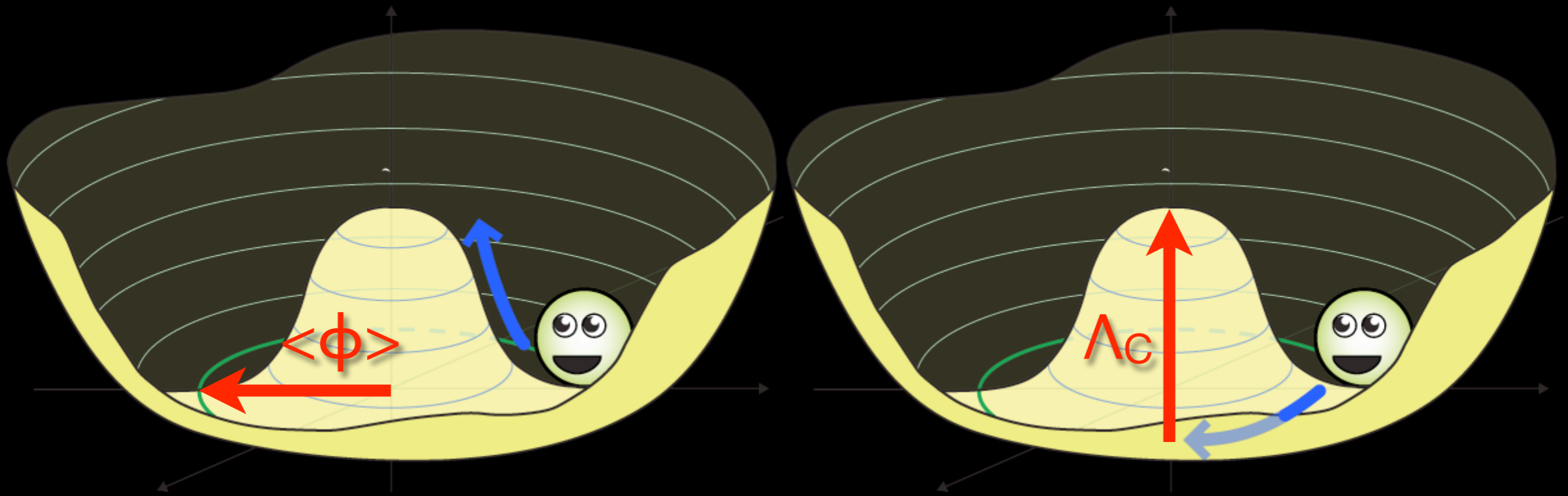


massless

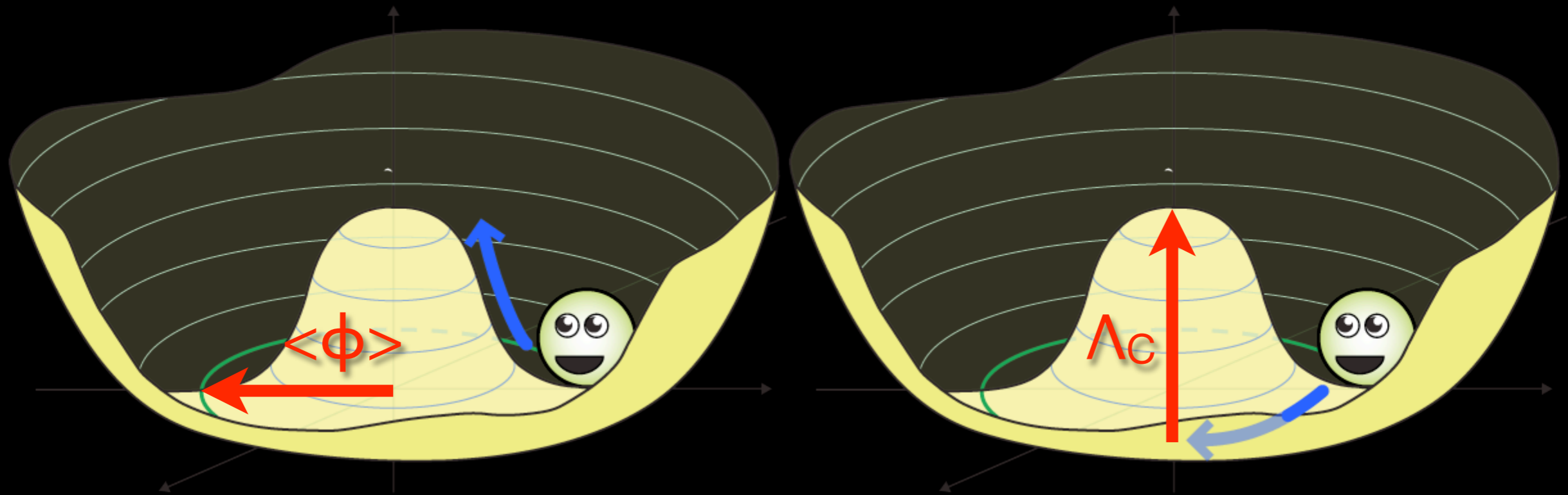
Nambu-Goldstone mode
($V = 0$ exactly)

- field gets vacuum expectation value $\langle \phi \rangle$ (order parameter)
- but may be elementary or composite (e.g. quark condensate breaking chiral symmetry of strong interaction, or *Cooper* pairs in *BCS theory* of superconductivity)



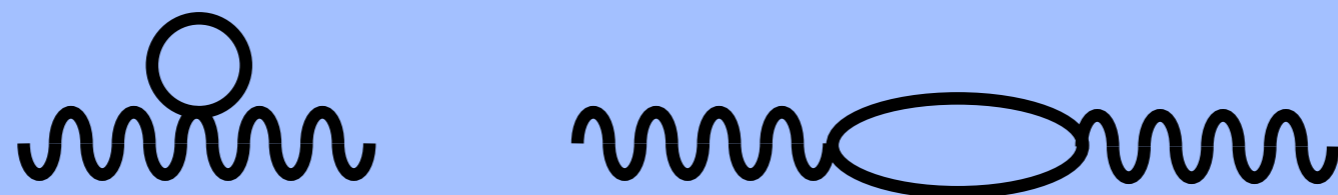


- **BCS theory:** approximate model; exact properties (like zero resistance and flux quantization) follow from symmetry breaking

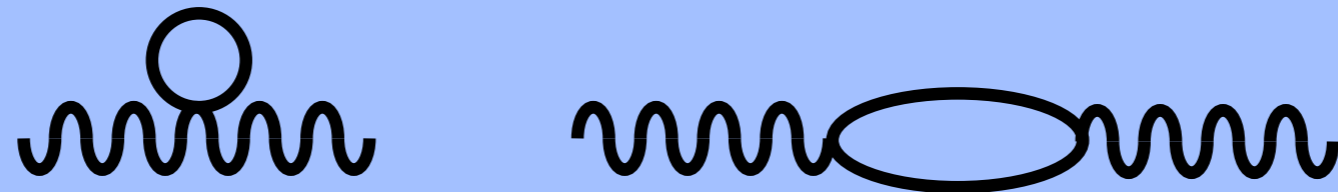


- **BCS theory:** approximate model; exact properties (like zero resistance and flux quantization) follow from symmetry breaking
- **electroweak theory:** need **new field** (the **Higgs**) or **new strong interaction** (e.g. **technicolor**); if **elementary** it must be **scalar** so as not to break **Lorentz invariance**

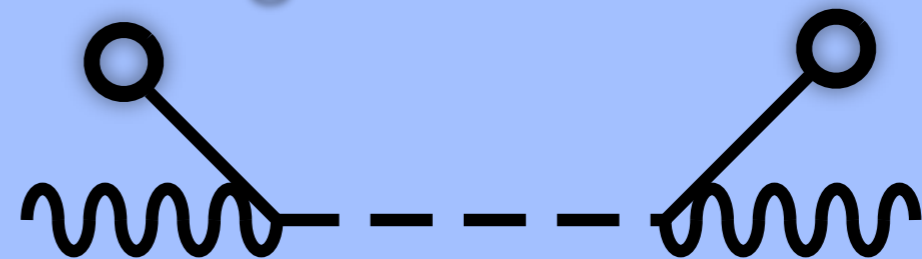
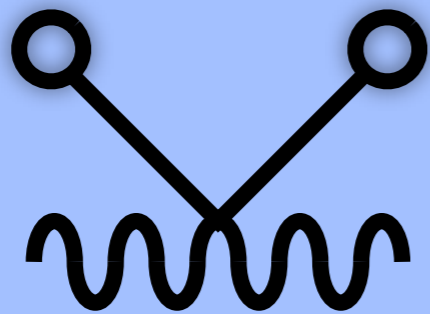
One-loop corrections to gauge boson propagator.



One-loop corrections to gauge boson propagator.



Additional tadpole graph corrections after spontaneous symmetry breaking.



effective mass term
(London's penetration depth)
renormalizability!

Nambu-Goldstone boson
provides longitudinal
degree of freedom
unitarity!

theoretical constraints

One-loop RGE

One-loop RGE

- $d M_H^2 / d \ln \mu^2 = 1 / (8 \pi^2 v^2)$
[$3 M_H^4 + 3 M_Z^4 + 6 M_W^4 - 12 m_t^4 - M_H^2 (3 M_Z^2 + 6 M_W^2 - 6 m_t^2)$]
- $\text{Str} = \sum (-)^{2S} (2S + 1) N_C$
- scalar field wave function renormalization (γ)
- the masses in the RGE are running masses (e.g. MS-bar)
- $v = [\sqrt{2} G_F]^{-1/2} = 246.22 \text{ GeV}$ (slightly modified definition of G_F)

One-loop RGE

- $d M_H^2 / d \ln \mu^2 = 1 / (8 \pi^2 v^2)$
[$3 M_H^4 + 3 M_Z^4 + 6 M_W^4 - 12 m_t^4 - M_H^2 (3 M_Z^2 + 6 M_W^2 - 6 m_t^2)$]
- $\text{Str} = \sum (-)^{2S} (2S + 1) N_C$
- scalar field wave function renormalization (γ)
- the masses in the RGE are running masses (e.g. MS-bar)
- $v = [\sqrt{2} G_F]^{-1/2} = 246.22 \text{ GeV}$ (slightly modified definition of G_F)
- $M_H \gg v \Rightarrow 1 / M_H^2(v) - 1 / M_H^2(\Lambda) = 3 / (8 \pi^2 v^2) \ln \Lambda^2 / v^2$
- $M_H^2 < 8 \pi^2 v^2 / (3 \ln \Lambda^2 / v^2) = [147 (144) \text{ GeV}]^2$
for $\Lambda = 2.4 \times 10^{18} (1.2 \times 10^{19}) \text{ GeV}$
- $M_H < 816 \text{ GeV}$ for $\Lambda = M_H$

Lower bound

Lower bound

- rewrite the RGE

$$\begin{aligned} d M_H^2 / d \ln \mu^2 &= 3 / (8 \pi^2 v^2) (M_H^2 - M_+^2) (M_H^2 - M_-^2) \text{ with } M_{\pm}^2 = \\ & \frac{1}{2} M_Z^2 + M_W^2 - m_t^2 \pm [(\frac{1}{2} M_Z^2 + M_W^2 - m_t^2)^2 - M_Z^4 - 2 M_W^4 + 4 m_t^4]^{1/2} \\ & \approx (-1 \pm \sqrt{5}) m_t^2 \Rightarrow M_+ \sim 200 \text{ GeV}, M_-^2 \sim - (300 \text{ GeV})^2 \\ & \Rightarrow \text{for } M_H \gtrsim 200 \text{ GeV, Landau pole persists (potential bounded below)} \end{aligned}$$

Lower bound

- rewrite the RGE

$$d M_H^2 / d \ln \mu^2 = 3 / (8 \pi^2 v^2) (M_H^2 - M_+^2) (M_H^2 - M_-^2) \text{ with } M_{\pm}^2 = \\ \frac{1}{2} M_Z^2 + M_W^2 - m_t^2 \pm [(\frac{1}{2} M_Z^2 + M_W^2 - m_t^2)^2 - M_Z^4 - 2 M_W^4 + 4 m_t^4]^{1/2} \\ \approx (-1 \pm \sqrt{5}) m_t^2 \Rightarrow M_+ \sim 200 \text{ GeV}, M_-^2 \sim - (300 \text{ GeV})^2$$

\Rightarrow for $M_H \gtrsim 200 \text{ GeV}$, Landau pole persists (potential bounded below)

- we have worked with constant masses, but $d m_t^2 / d \ln \mu^2 =$

$$3 m_t^2 / (16 \pi^2 v^2) [(1_L + 2_R) m_t^2 - (2_L + 0_R) \dim \text{SU}(2) / N_D M_W^2 - \\ ((1/3)^2_L + (4/3)^2_R) (M_Z^2 - M_W^2) - (2_L + 2_R) \dim \text{SU}(3) / N_C (\pi \alpha_s v^2)] = \\ m_t^2 / (16 \pi^2 v^2) [9 m_t^2 - 10/3 M_W^2 - 17/3 M_Z^2 - 32 (\pi \alpha_s v^2)] < 0$$

summary of bounds

summary of bounds

- “triviality”:

$$M_H^2/v^2 \ln M_H^2/v^2 < 8 \pi^2/3 \Rightarrow M_H < 816 \text{ GeV}$$

summary of bounds

- “triviality”:

$$M_H^2/v^2 \ln M_H^2/v^2 < 8 \pi^2/3 \Rightarrow M_H < 816 \text{ GeV}$$

- tree-level unitarity of the partial S-wave amplitude of elastic Goldstone boson scattering:

$$M_H^2/v^2 < 16 \pi/5 \Rightarrow M_H < 781 \text{ GeV}$$

summary of bounds

- “triviality”:

$$M_H^2/v^2 \ln M_H^2/v^2 < 8 \pi^2/3 \Rightarrow M_H < 816 \text{ GeV}$$

- tree-level **unitarity** of the partial S-wave amplitude of elastic Goldstone boson scattering:

$$M_H^2/v^2 < 16 \pi/5 \Rightarrow M_H < 781 \text{ GeV}$$

- absence of **Landau pole**:

$$M_H^2/v^2 < 4 \pi^2/3 \ln^{-1} \kappa_P/v \Rightarrow M_H < 147 \text{ GeV}$$

summary of bounds

- “triviality”:

$$M_H^2/v^2 \ln M_H^2/v^2 < 8 \pi^2/3 \Rightarrow M_H < 816 \text{ GeV}$$

- tree-level **unitarity** of the partial S-wave amplitude of elastic Goldstone boson scattering:

$$M_H^2/v^2 < 16 \pi/5 \Rightarrow M_H < 781 \text{ GeV}$$

- absence of **Landau pole**:

$$M_H^2/v^2 < 4 \pi^2/3 \ln^{-1} \kappa_P/v \Rightarrow M_H < 147 \text{ GeV}$$

- **vacuum stability**: $M_H \gtrsim 130 \text{ GeV}$ *Casas, Espinosa, Quiros 1995*

summary of bounds

- “triviality”:

$$M_H^2/v^2 \ln M_H^2/v^2 < 8 \pi^2/3 \Rightarrow M_H < 816 \text{ GeV}$$

- tree-level **unitarity** of the partial S-wave amplitude of elastic Goldstone boson scattering:

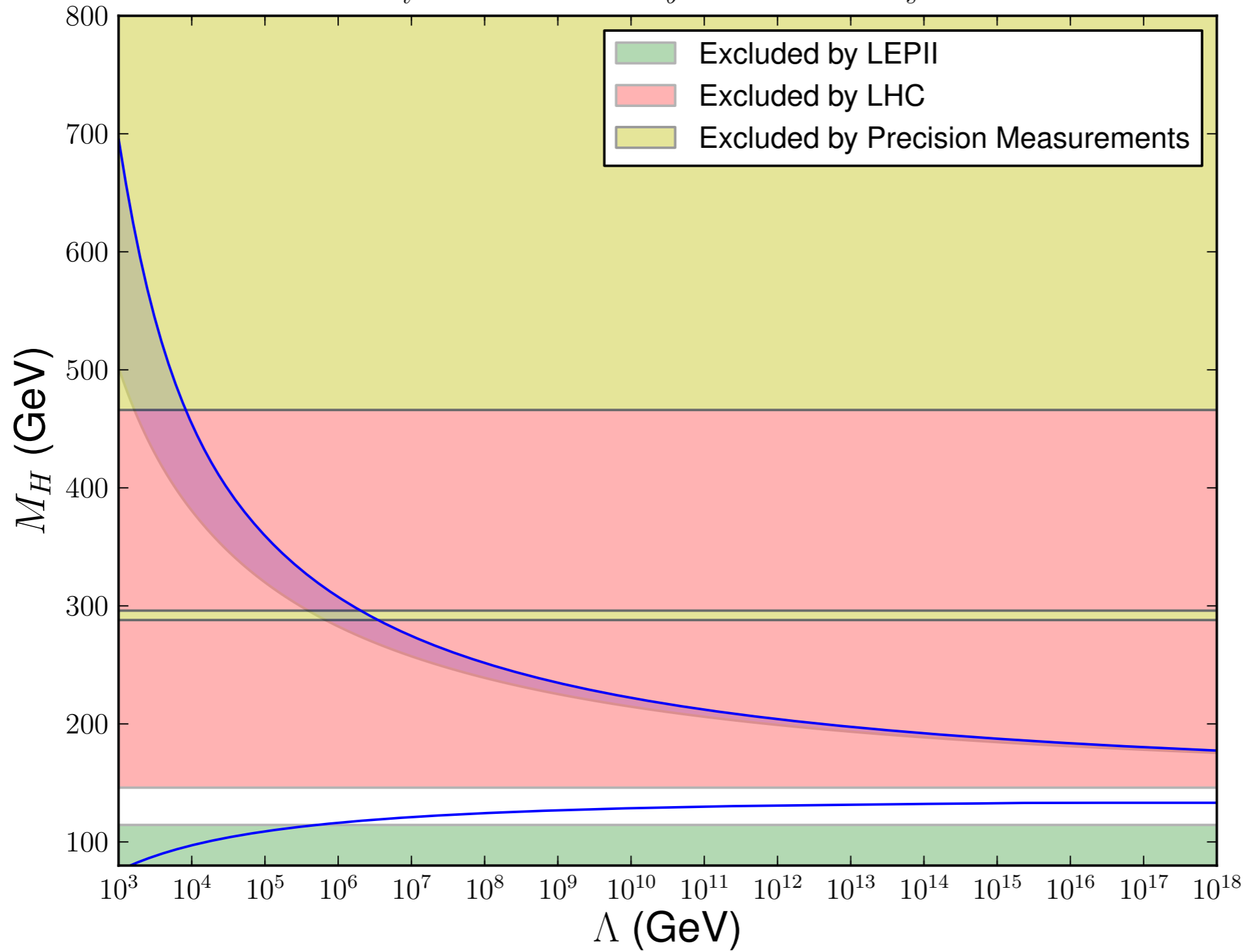
$$M_H^2/v^2 < 16 \pi/5 \Rightarrow M_H < 781 \text{ GeV}$$

- absence of **Landau pole**:

$$M_H^2/v^2 < 4 \pi^2/3 \ln^{-1} \kappa_P/v \Rightarrow M_H < 147 \text{ GeV}$$

- **vacuum stability**: $M_H \gtrsim 130 \text{ GeV}$ *Casas, Espinosa, Quiros 1995*
- **vacuum meta-stability**: $M_H \gtrsim 115 \text{ GeV}$ *Isidori, Ridolfi, Strumia 2001*

SM3: $m_t = 173.3 \text{ GeV}$ $\bar{m}_b = 4.19 \text{ GeV}$ $\alpha_s = 0.1184$



*originally by
Hambye
Riesselmann
1996*

*figure from
Wingerter
2011*

indirect constraints

M_w

M_W

- $M_W = 80.387 \pm 0.016 \text{ GeV}$ *CDF & D0 2012* ($\pm 19 \text{ MeV}$ *CDF* 2.2 fb^{-1})

M_W

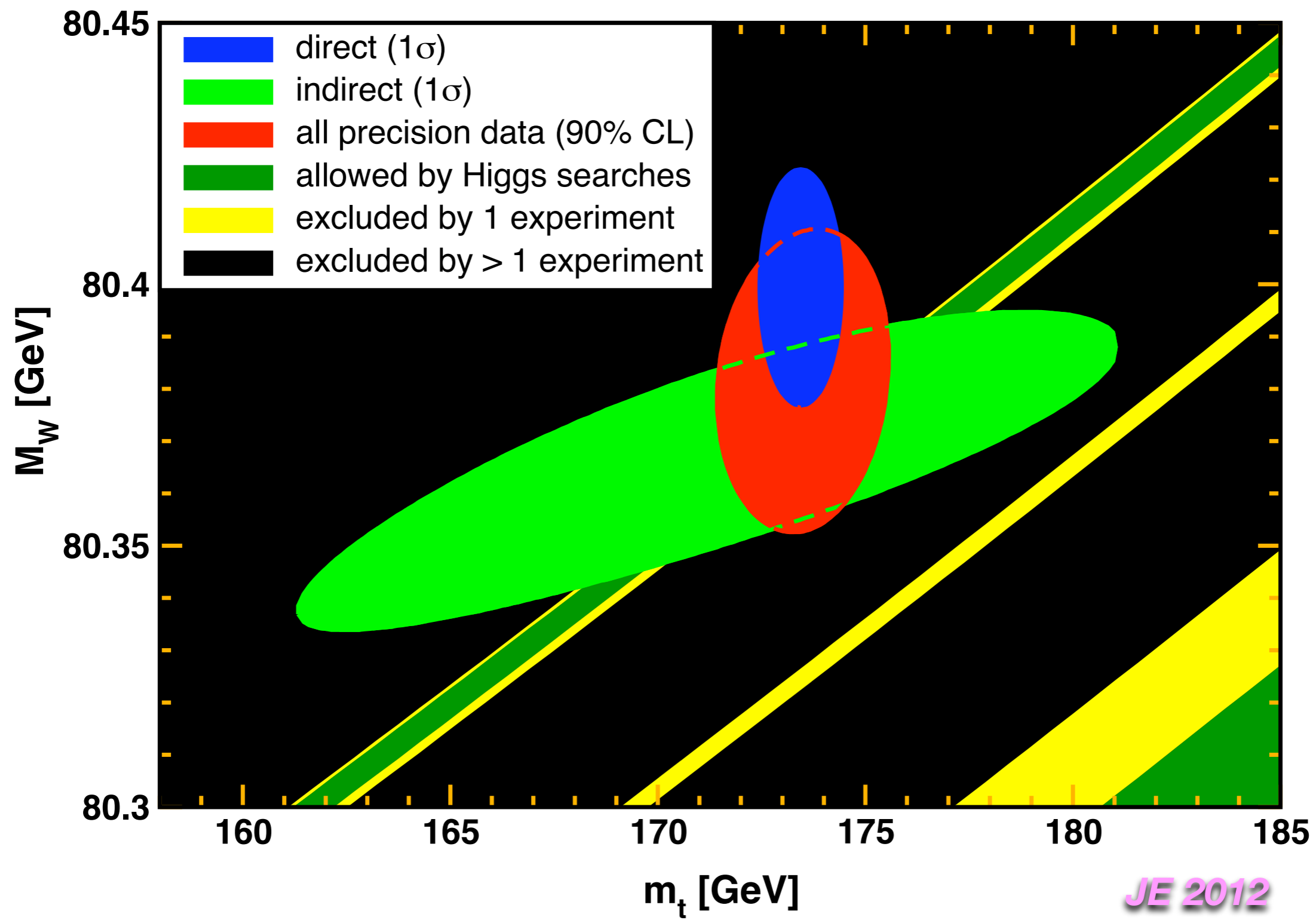
- $M_W = 80.387 \pm 0.016$ GeV *CDF & D0 2012* (± 19 MeV *CDF* 2.2 fb⁻¹)
- $M_W = 80.376 \pm 0.033$ GeV *LEP 2*
 - $\Rightarrow \sin^2 \theta_W^{\text{on-shell}} \equiv 1 - M_W^2/M_Z^2 = 0.22290 \pm 0.00028$
 - $\Rightarrow \sin^2 \theta_W^{\text{eff}} = 0.23141 \pm 0.00013$ and $M_H = 96^{+29}_{-25}$ GeV

M_W

- $M_W = 80.387 \pm 0.016$ GeV *CDF & D0 2012* (± 19 MeV *CDF* 2.2 fb^{-1})
- $M_W = 80.376 \pm 0.033$ GeV *LEP 2*
 - $\Rightarrow \sin^2 \theta_W^{\text{on-shell}} \equiv 1 - M_W^2/M_Z^2 = 0.22290 \pm 0.00028$
 - $\Rightarrow \sin^2 \theta_W^{\text{eff}} = 0.23141 \pm 0.00013$ and $M_H = 96^{+29}_{-25}$ GeV
- new global electroweak fit: $M_H = 102^{+24}_{-20}$ GeV *JE 2012*

M_W

- $M_W = 80.387 \pm 0.016$ GeV *CDF & D0 2012* (± 19 MeV *CDF* 2.2 fb⁻¹)
- $M_W = 80.376 \pm 0.033$ GeV *LEP 2*
 - $\Rightarrow \sin^2 \theta_W^{\text{on-shell}} \equiv 1 - M_W^2/M_Z^2 = 0.22290 \pm 0.00028$
 - $\Rightarrow \sin^2 \theta_W^{\text{eff}} = 0.23141 \pm 0.00013$ and $M_H = 96^{+29}_{-25}$ GeV
- new global electroweak fit: $M_H = 102^{+24}_{-20}$ GeV *JE 2012*
- prospects:
 - no PDF (± 10 MeV) & QED (± 4 MeV) improvement $\Rightarrow \pm 13$ MeV *CDF* 10 fb⁻¹
 - ± 7 MeV_{PDF} $\Rightarrow \pm 11$ MeV *CDF* 10 fb⁻¹
 - ± 5 MeV_{PDF} & lepton energy scale $\pm 6 \rightarrow \pm 3$ MeV $\Rightarrow \pm 10$ MeV *CDF* 10 fb⁻¹
 - ILC threshold scan: ± 6 MeV



m_t

m_t

- $m_t = 173.4 \pm 0.9_{\text{exp}} \pm 0.5_{\text{th}} \text{ GeV}$

m_t

- $m_t = 173.4 \pm 0.9_{\text{exp}} \pm 0.5_{\text{th}} \text{ GeV}$
- Question: What is the definition of m_t ?

Correct but useless answer: $m_t \equiv m_t^{\text{Pythia}}$ (“Pythia tuning parameter”)

We assume $m_t^{\text{Pythia}} = m_t^{\text{pole}} \pm \Lambda_{\text{QCD}}$ where

$$m_t^{\text{pole}} \equiv \bar{m}_t(\bar{m}_t) \left[1 + \frac{4}{3} \alpha_s(\bar{m}_t) / \pi + O(\alpha_s^2) + O(\alpha_s^3) \right]$$

and $\Lambda_{\text{QCD}} \equiv$ the $O(\alpha_s^3)$ term above (see also *Skands, Wicke 2007*)

m_t

- $m_t = 173.4 \pm 0.9_{\text{exp}} \pm 0.5_{\text{th}}$ GeV

- Question: What is the definition of m_t ?

Correct but useless answer: $m_t \equiv m_t^{\text{Pythia}}$ (“Pythia tuning parameter”)

We assume $m_t^{\text{Pythia}} = m_t^{\text{pole}} \pm \Lambda_{\text{QCD}}$ where

$$m_t^{\text{pole}} \equiv \bar{m}_t(\bar{m}_t) [1 + 4/3 \alpha_s(\bar{m}_t)/\pi + O(\alpha_s^2) + O(\alpha_s^3)]$$

and $\Lambda_{\text{QCD}} \equiv$ the $O(\alpha_s^3)$ term above (see also *Skands, Wicke 2007*)

- Alternative I: SCET + HQET \rightarrow “jet mass” *Fleming, Hoang, Mantry, Stewart 2008*

m_t

- $m_t = 173.4 \pm 0.9_{\text{exp}} \pm 0.5_{\text{th}} \text{ GeV}$

- Question: What is the definition of m_t ?

Correct but useless answer: $m_t \equiv m_t^{\text{Pythia}}$ (“Pythia tuning parameter”)

We assume $m_t^{\text{Pythia}} = m_t^{\text{pole}} \pm \Lambda_{\text{QCD}}$ where

$$m_t^{\text{pole}} \equiv \bar{m}_t(\bar{m}_t) [1 + 4/3 \alpha_s(\bar{m}_t)/\pi + O(\alpha_s^2) + O(\alpha_s^3)]$$

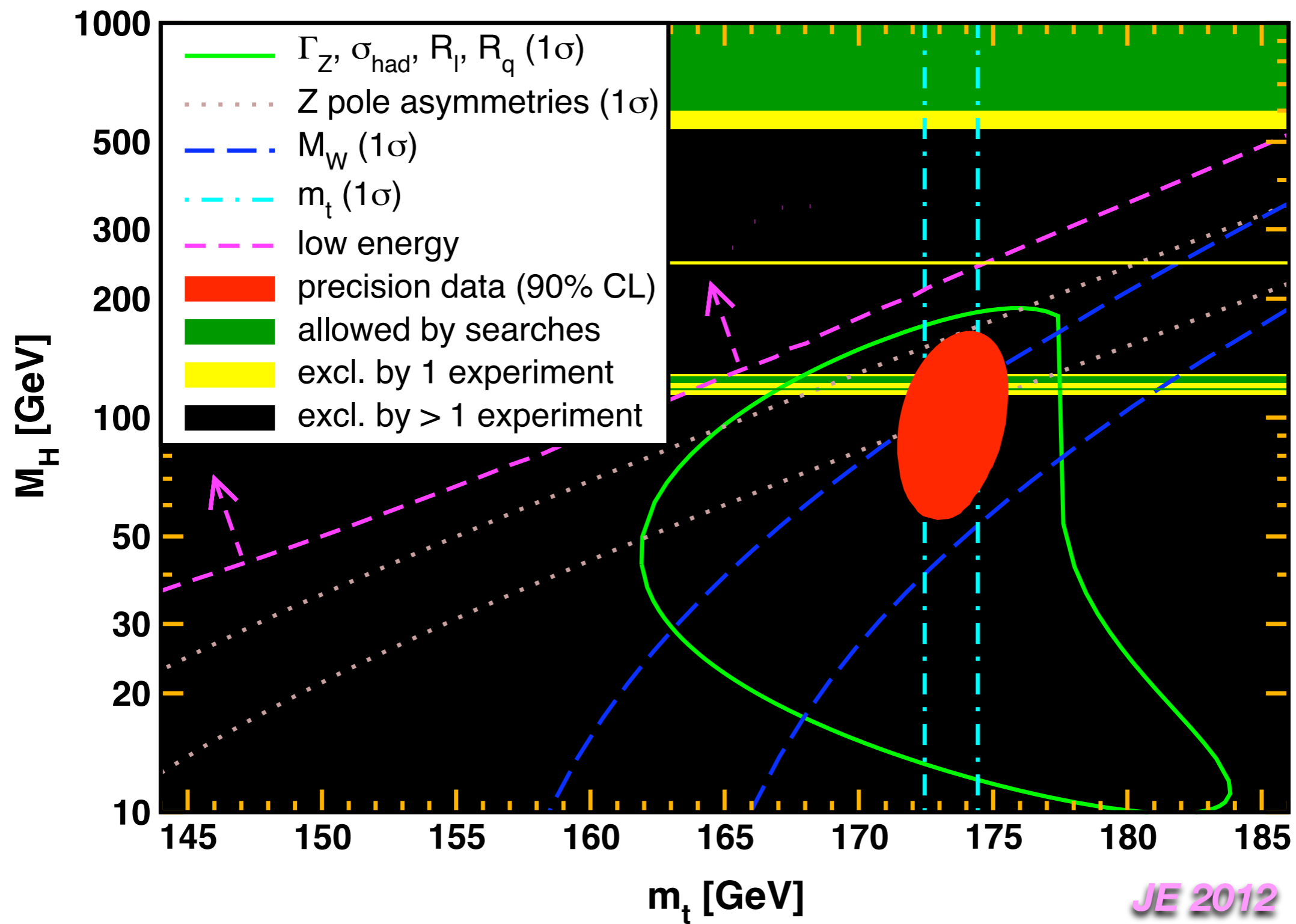
and $\Lambda_{\text{QCD}} \equiv$ the $O(\alpha_s^3)$ term above (see also *Skands, Wicke 2007*)

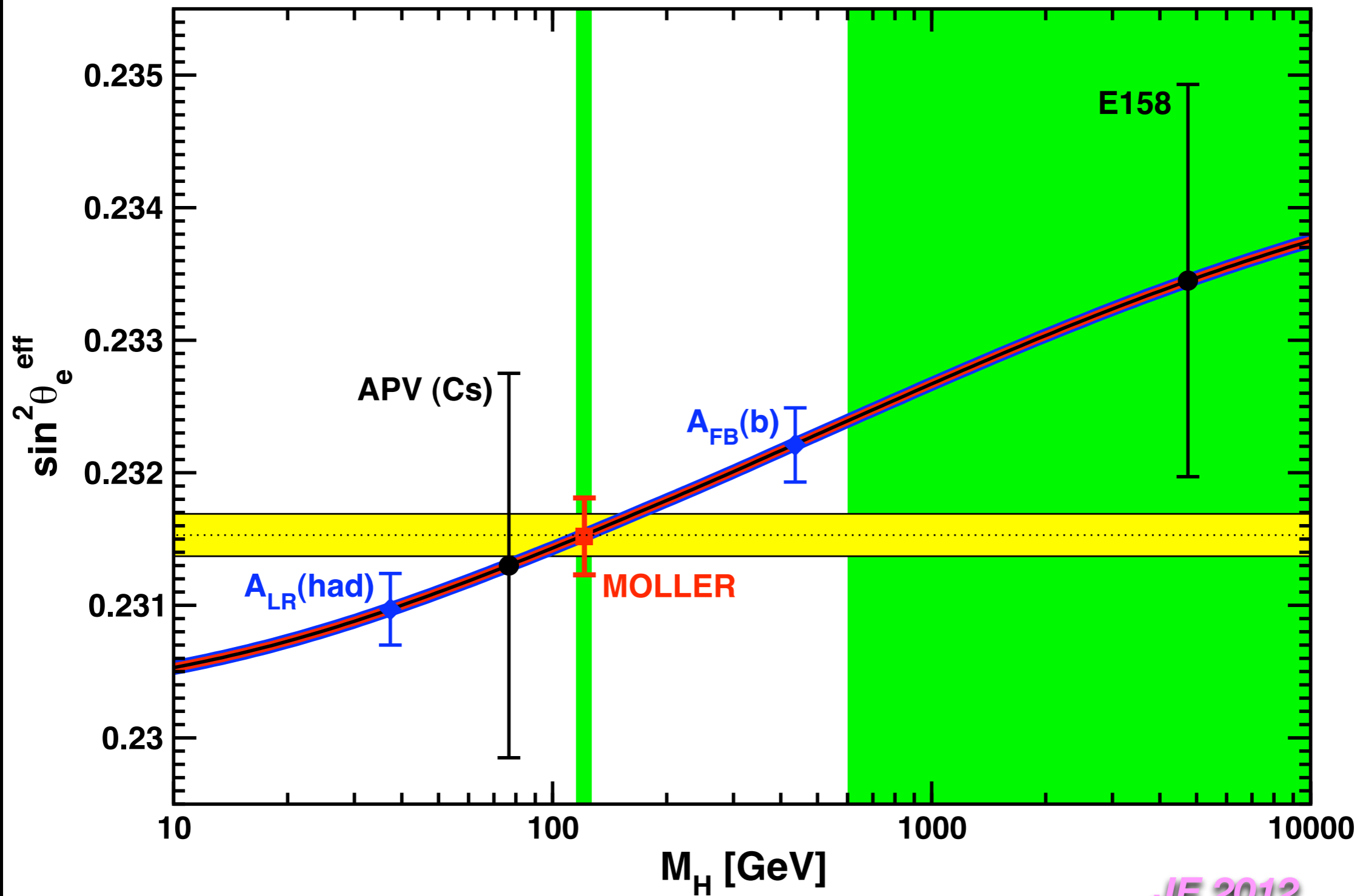
- Alternative I: SCET + HQET \rightarrow “jet mass” *Fleming, Hoang, Mantry, Stewart 2008*

- Alternative II: get $\bar{m}_t(\bar{m}_t)$ directly from $t \bar{t}$ cross-section \Rightarrow

$$\bar{m}_t(\bar{m}_t) = 160.0 \pm 3.3 \text{ GeV} \text{ *Langenfeld, Moch, Uwer 2008*}$$

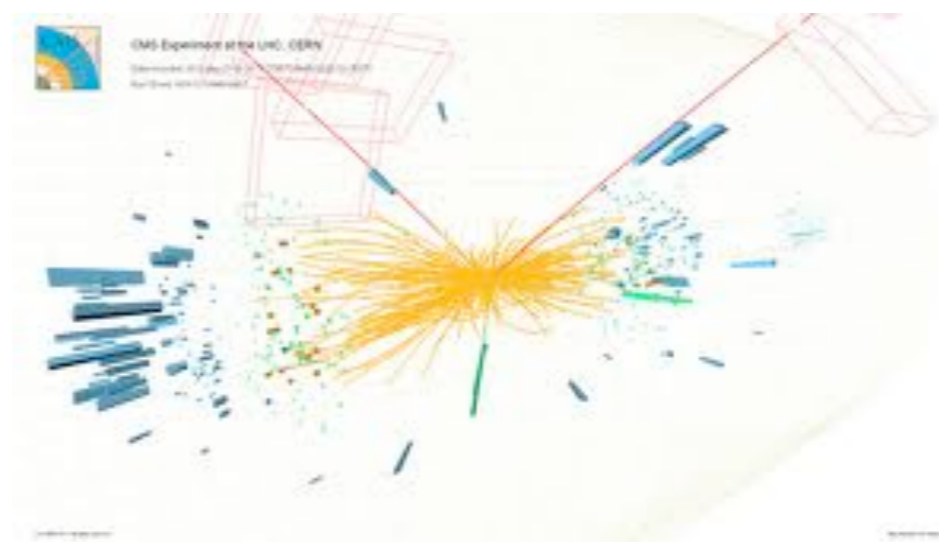
$$\Rightarrow M_H = 81^{+32}_{-24} \text{ GeV} \text{ (} m_t^{\text{pole}} = 169.6 \pm 3.5 \text{ GeV)}$$



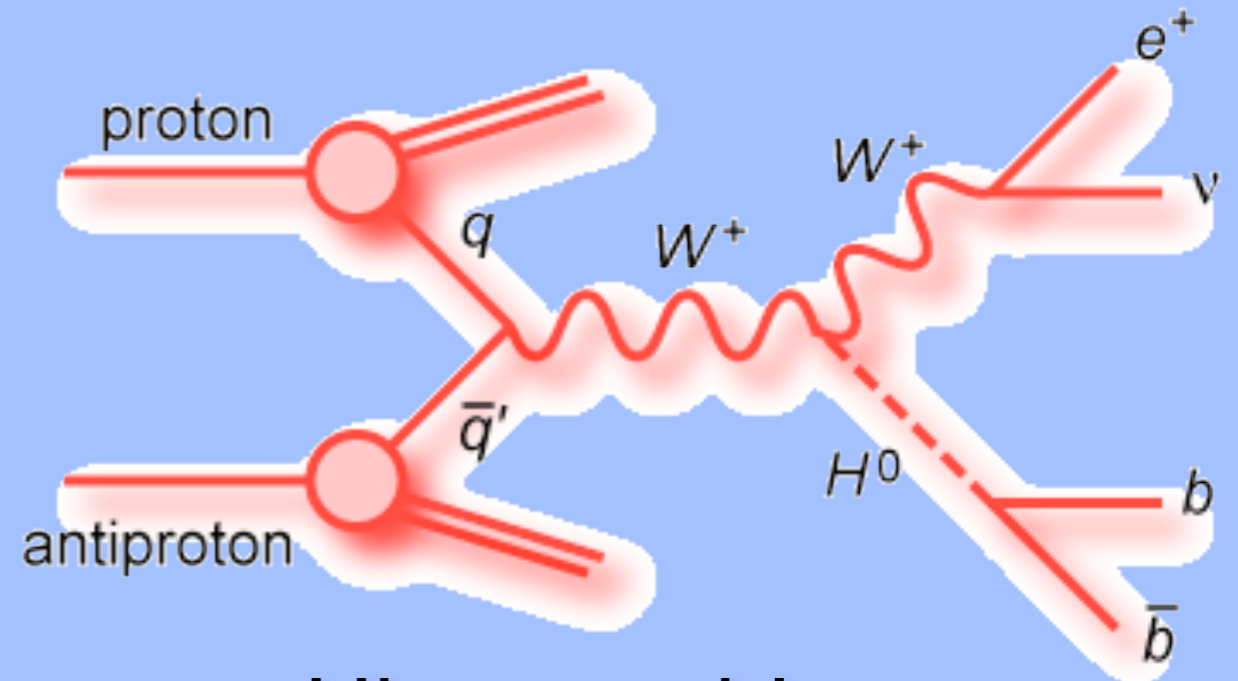


JE 2012

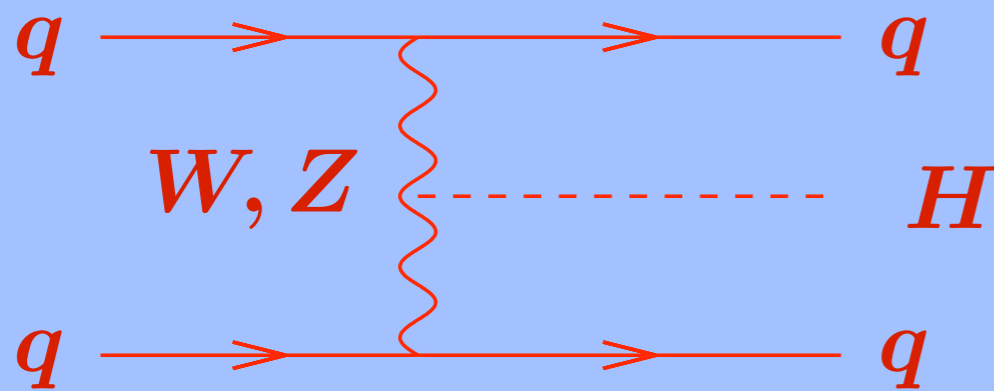
Constraints from the Higgs hunt



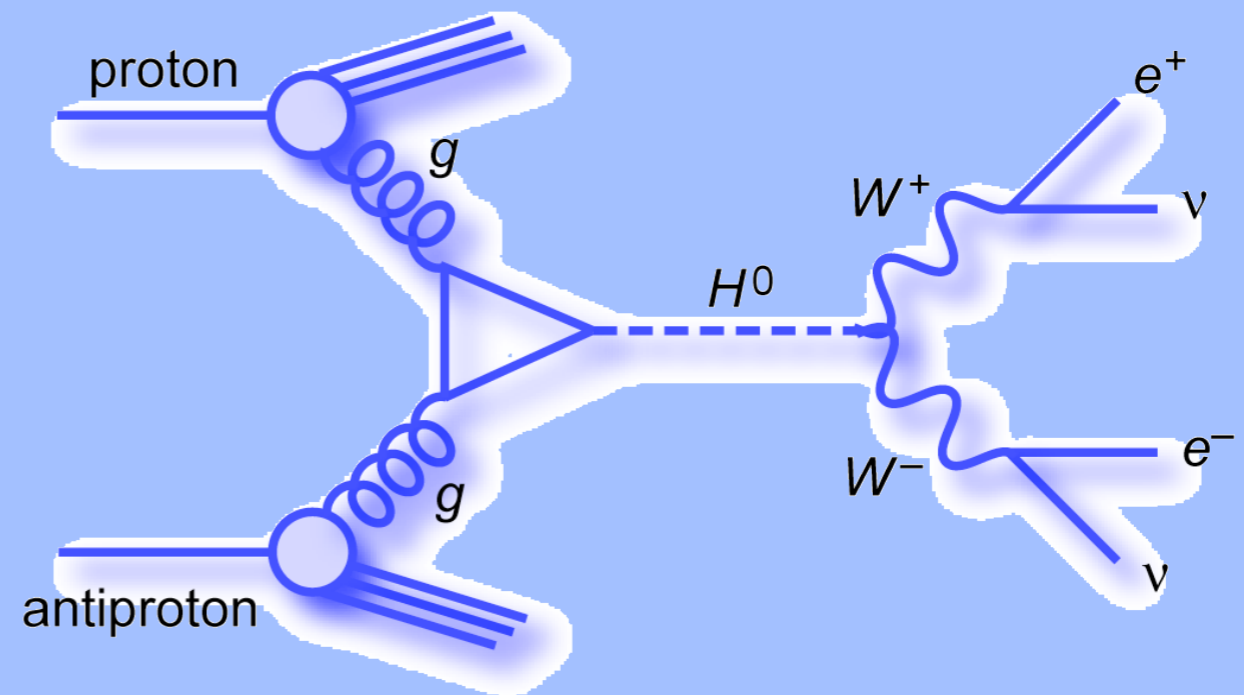
How to make a Higgs



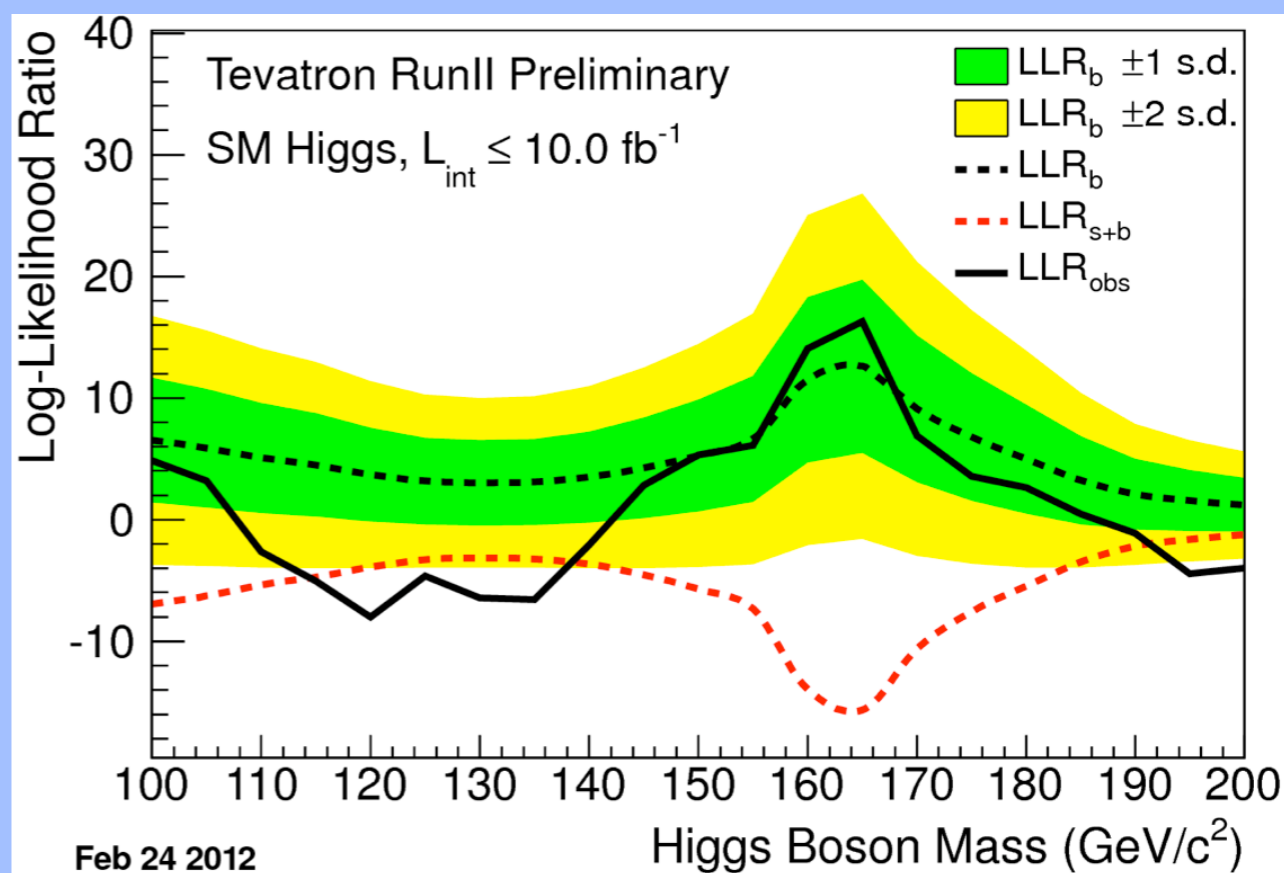
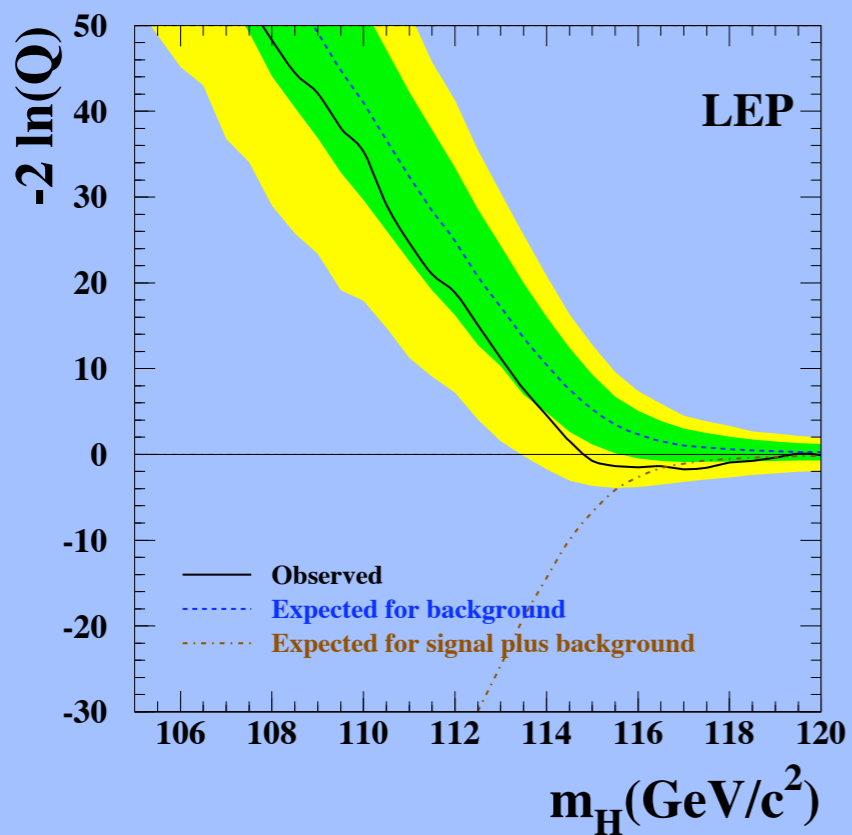
Higgsstrahlung



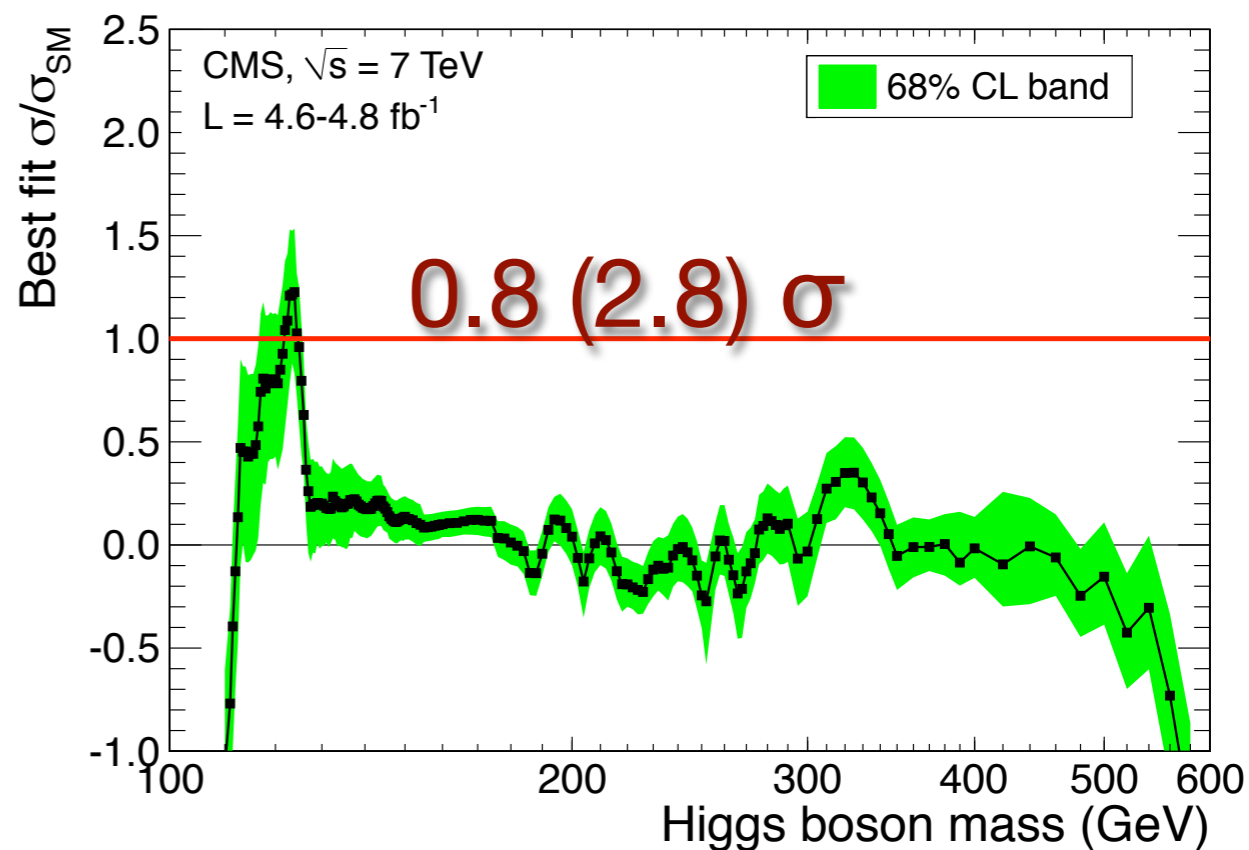
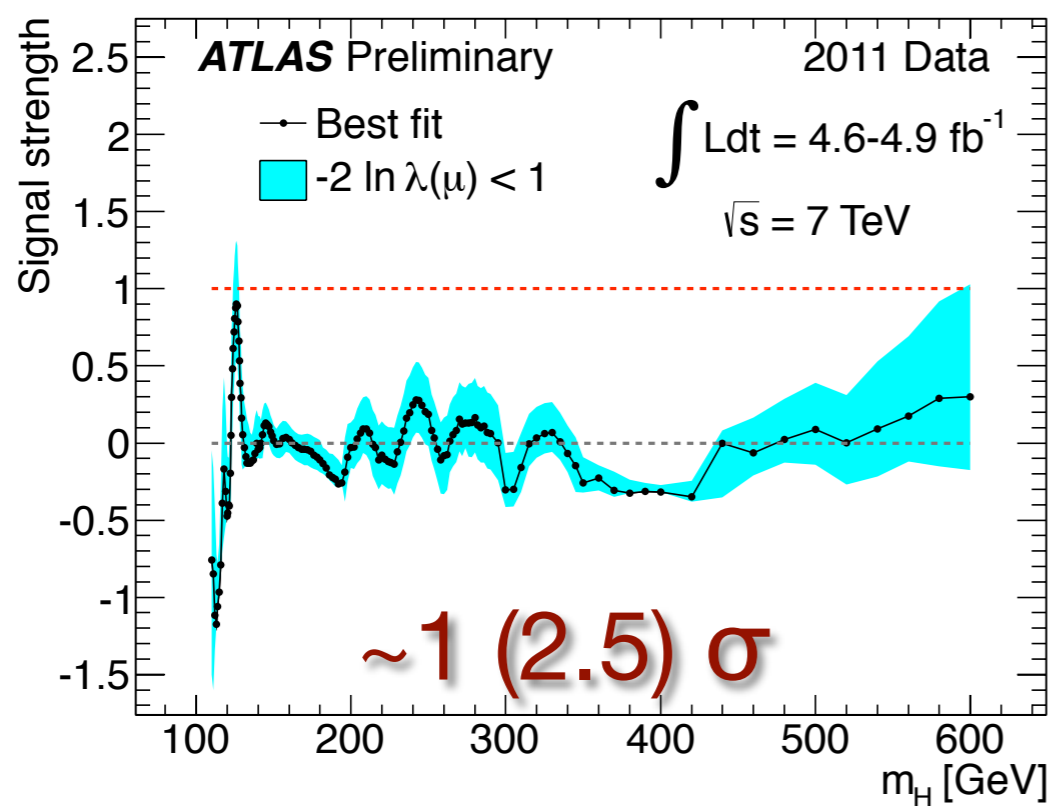
vector boson fusion

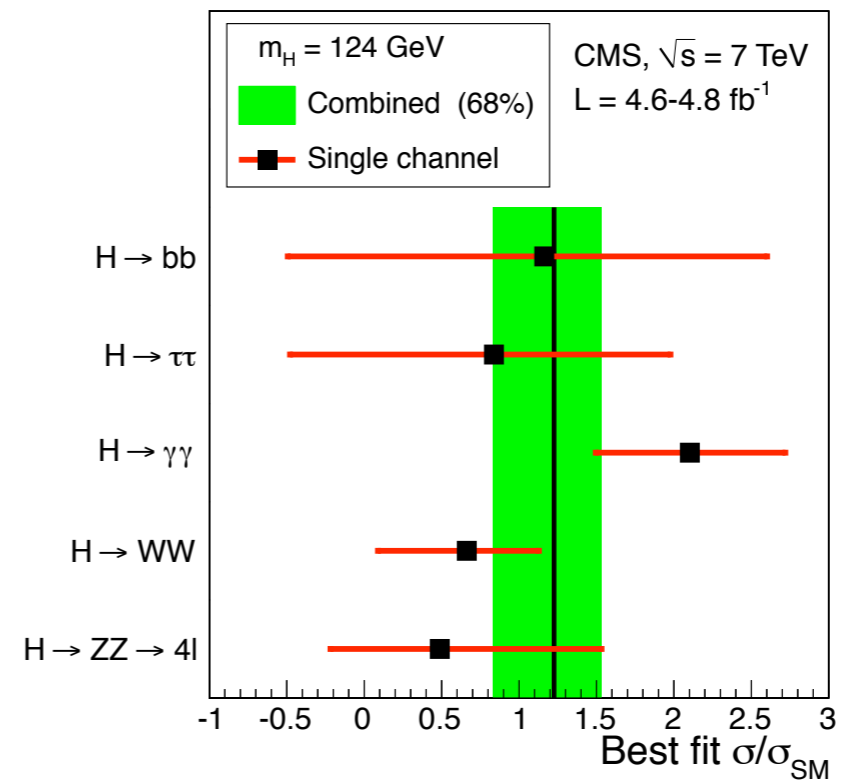
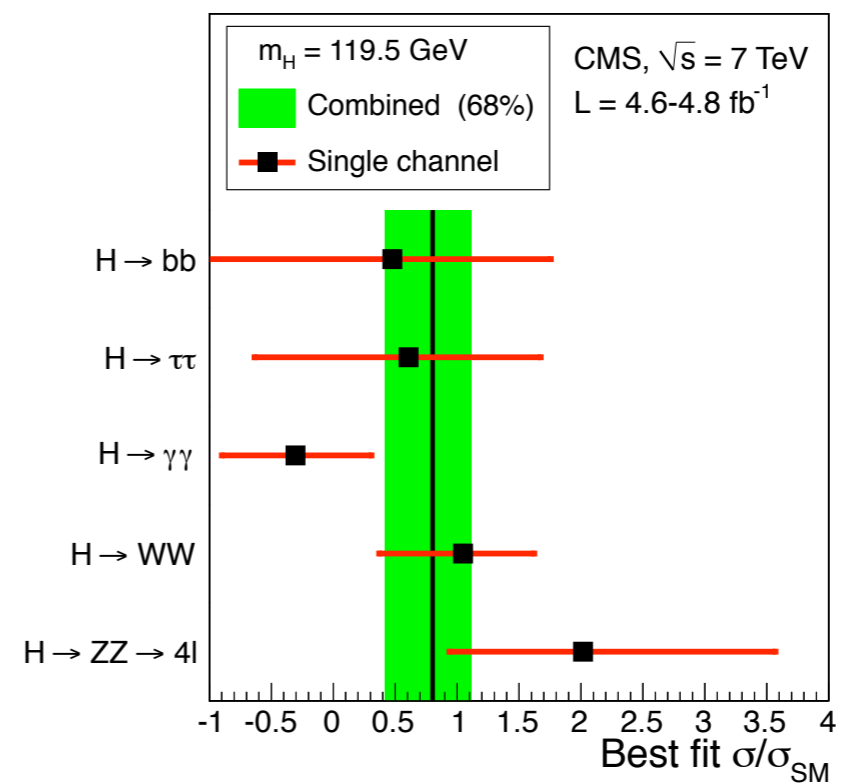
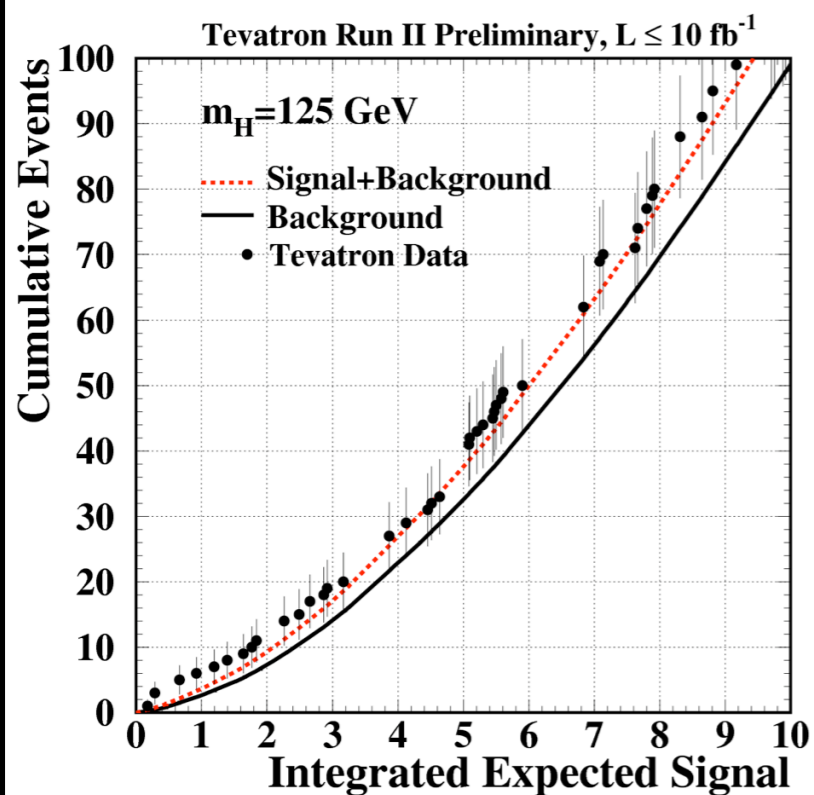
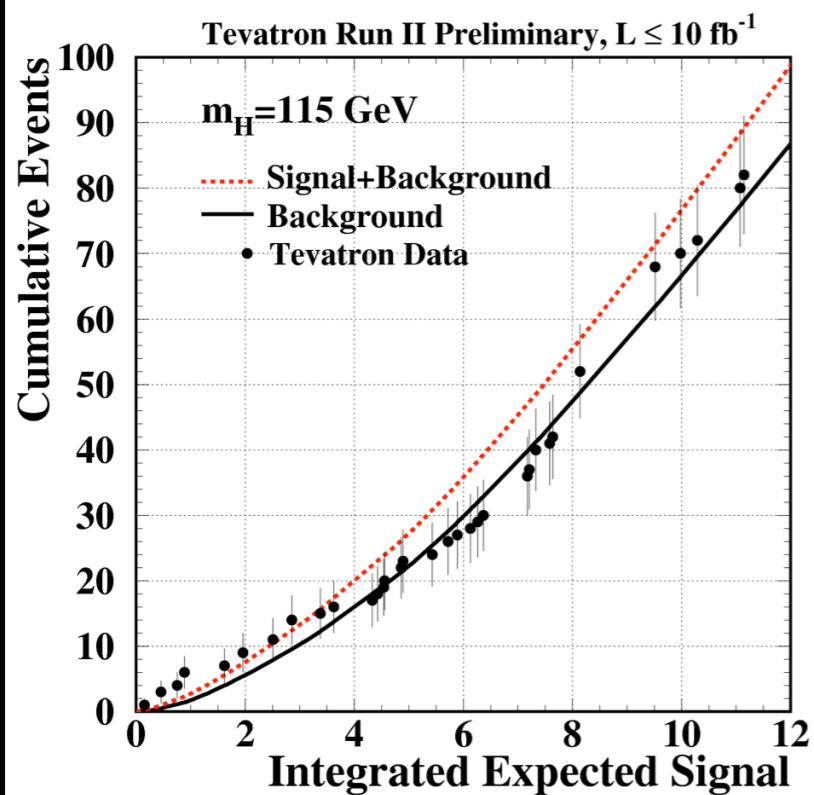


gluon fusion



Feb 24 2012





synthesis

M_H probability density

M_H probability density

- $p(M_H) \equiv \exp[-\chi^2_{EW}(M_H)/2] Q_{LEP} Q_{Tevatron} Q_{LHC} M_H^{-1}$

factorized form: neglect of correlations

M_H probability density

- $p(M_H) \equiv \exp[-\chi^2_{EW}(M_H)/2] Q_{LEP} Q_{Tevatron} Q_{LHC} M_H^{-1}$

factorized form: neglect of correlations

- $Q_{LEP}(M_H), Q_{Tevatron}(M_H)$: likelihood ratios H/H+B

M_H probability density

- $p(M_H) \equiv \exp[-\chi^2_{EW}(M_H)/2] Q_{LEP} Q_{Tevatron} Q_{LHC} M_H^{-1}$

factorized form: neglect of correlations

- $Q_{LEP}(M_H), Q_{Tevatron}(M_H)$: likelihood ratios H/H+B

- $Q_{LHC}(M_H) = Q_{ATLAS}(M_H) Q_{CMS}(M_H)$ (but not available)

instead: $2 \ln Q \equiv \chi^2_{H+B}(M_H) - \chi^2_B(M_H) \equiv$

$$(1 - \bar{\sigma}_{obs})^2 / \Delta \bar{\sigma}_+^2 - \bar{\sigma}_{obs}^2 / \Delta \bar{\sigma}_-^2$$

M_H probability density

- $p(M_H) \equiv \exp[-\chi^2_{EW}(M_H)/2] Q_{LEP} Q_{Tevatron} Q_{LHC} M_H^{-1}$

factorized form: neglect of correlations

- $Q_{LEP}(M_H)$, $Q_{Tevatron}(M_H)$: likelihood ratios H/H+B
- $Q_{LHC}(M_H) = Q_{ATLAS}(M_H) Q_{CMS}(M_H)$ (but not available)

instead: $2 \ln Q \equiv \chi^2_{H+B}(M_H) - \chi^2_B(M_H) \equiv$

$$(1 - \bar{\sigma}_{obs})^2 / \Delta\bar{\sigma}_+^2 - \bar{\sigma}_{obs}^2 / \Delta\bar{\sigma}_-^2$$

- $\bar{\sigma}_{obs}$: effective observed X-section combining all channels

M_H probability density

- $p(M_H) \equiv \exp[-\chi^2_{EW}(M_H)/2] Q_{LEP} Q_{Tevatron} Q_{LHC} M_H^{-1}$

factorized form: neglect of correlations

- $Q_{LEP}(M_H), Q_{Tevatron}(M_H)$: likelihood ratios H/H+B

- $Q_{LHC}(M_H) = Q_{ATLAS}(M_H) Q_{CMS}(M_H)$ (but not available)

instead: $2 \ln Q \equiv \chi^2_{H+B}(M_H) - \chi^2_B(M_H) \equiv$

$$(1 - \bar{\sigma}_{obs})^2 / \Delta\bar{\sigma}_+^2 - \bar{\sigma}_{obs}^2 / \Delta\bar{\sigma}_-^2$$

- $\bar{\sigma}_{obs}$: effective observed X-section combining all channels

- $\Delta\bar{\sigma}_{\pm}$: error pointing in signal (+) and background (-) direction

M_H probability density

- $p(M_H) \equiv \exp[-\chi^2_{EW}(M_H)/2] Q_{LEP} Q_{Tevatron} Q_{LHC} M_H^{-1}$

factorized form: neglect of correlations

- $Q_{LEP}(M_H), Q_{Tevatron}(M_H)$: likelihood ratios H/H+B

- $Q_{LHC}(M_H) = Q_{ATLAS}(M_H) Q_{CMS}(M_H)$ (but not available)

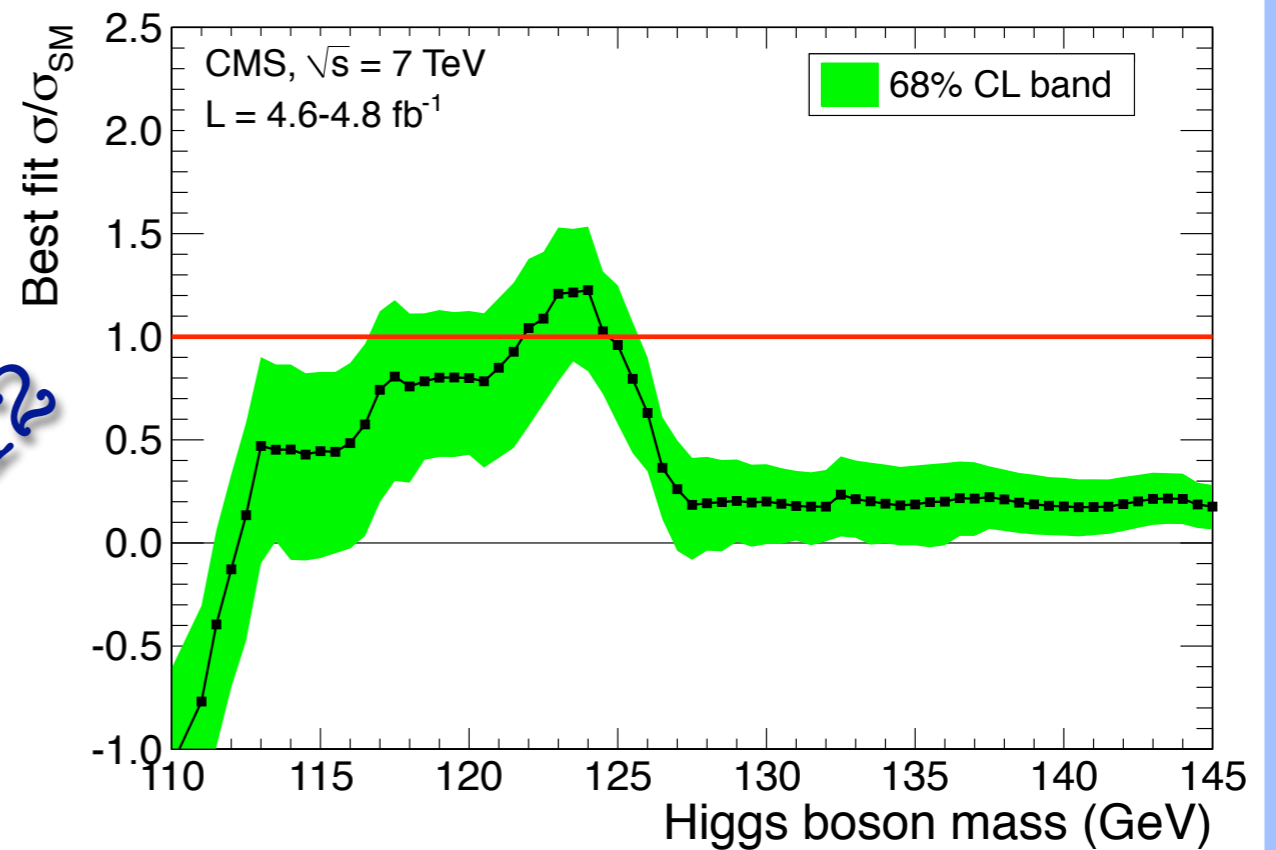
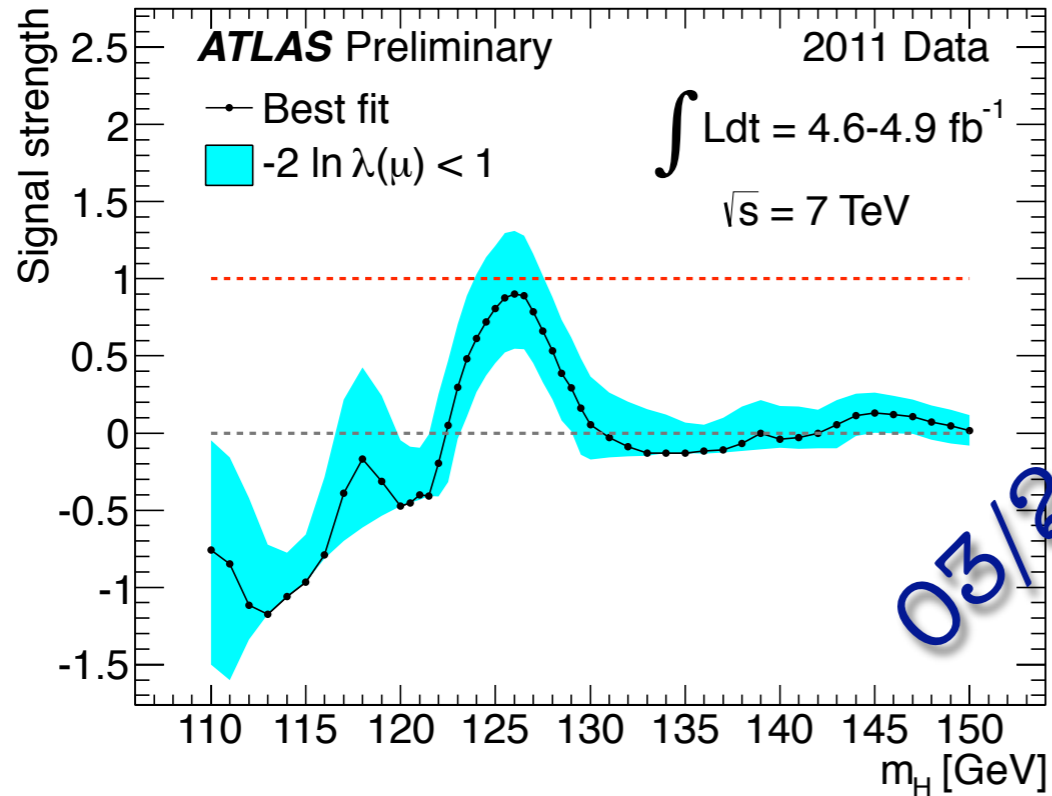
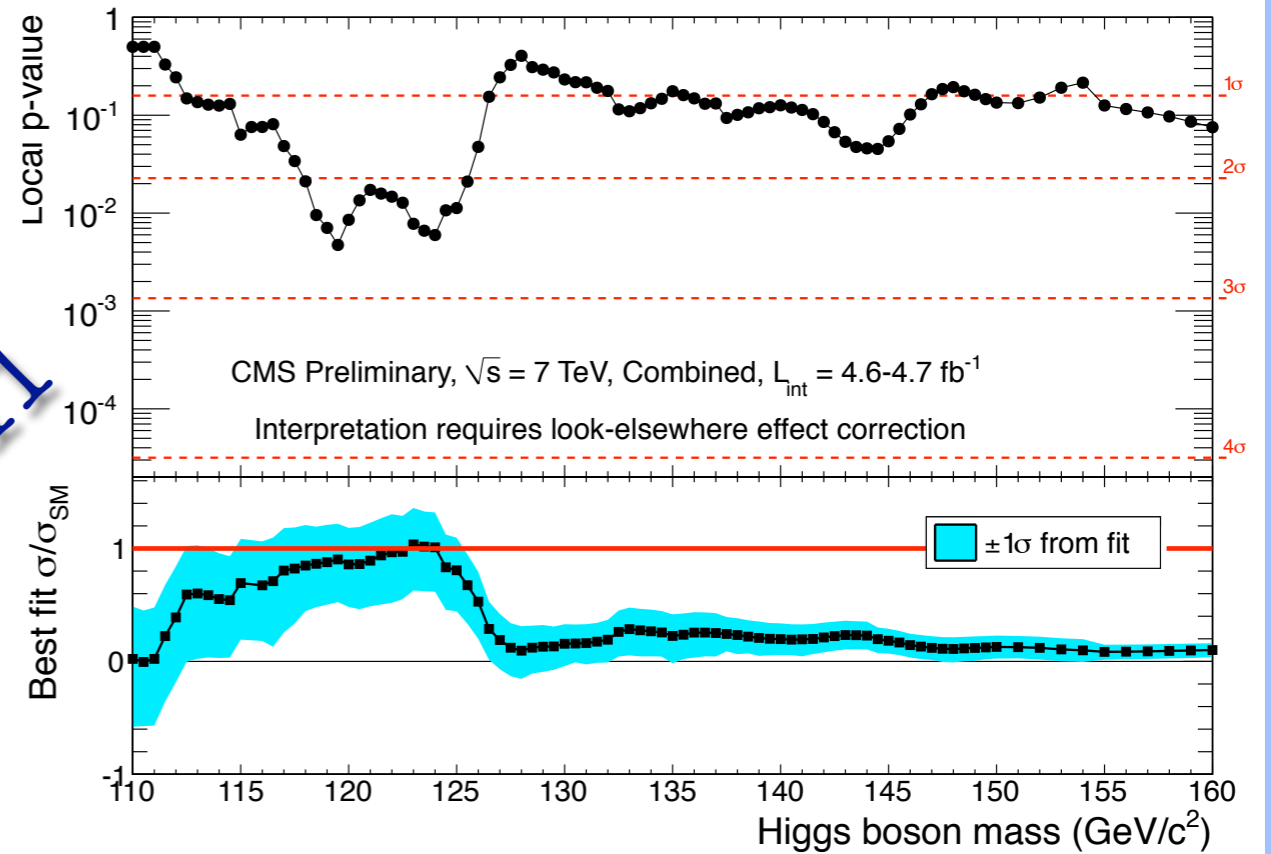
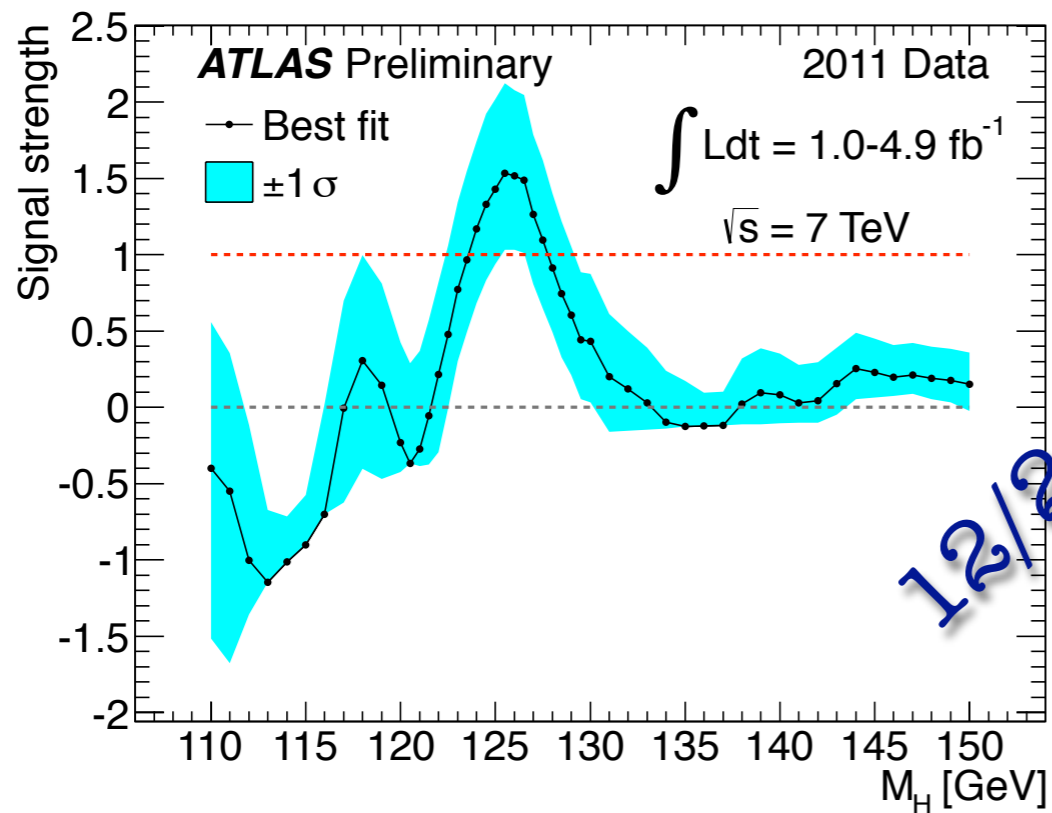
instead: $2 \ln Q \equiv \chi^2_{H+B}(M_H) - \chi^2_B(M_H) \equiv$

$$(1 - \bar{\sigma}_{obs})^2 / \Delta\bar{\sigma}_+^2 - \bar{\sigma}_{obs}^2 / \Delta\bar{\sigma}_-^2$$

- $\bar{\sigma}_{obs}$: effective observed X-section combining all channels

- $\Delta\bar{\sigma}_\pm$: error pointing in signal (+) and background (-) direction

- Poisson statistics $\Rightarrow \Delta\bar{\sigma}_+ > \Delta\bar{\sigma}_-$ but often also $\Delta\bar{\sigma}_+ < \Delta\bar{\sigma}_-$



Examples

Examples

- $2 \ln Q_{\text{ATLAS}}(126 \text{ GeV}) = 9.8 - 1.1 = -8.7$ ($H \rightarrow \gamma\gamma, ZZ^*$)
 $\sqrt{8.7} = 2.9$ while quoted local significance of excess = 3.6σ

Examples

- $2 \ln Q_{\text{ATLAS}}(126 \text{ GeV}) = 9.8 - 1.1 = -8.7$ ($H \rightarrow \gamma\gamma, ZZ^*$)
 $\sqrt{8.7} = 2.9$ while quoted local significance of excess = 3.6σ
- $2 \ln Q_{\text{ATLAS}}(244 \text{ GeV}) \approx 2 \ln Q_{\text{ATLAS}}(560 \text{ GeV}) \approx -3$ ($H \rightarrow ZZ$)

Examples

- $2 \ln Q_{\text{ATLAS}}(126 \text{ GeV}) = 9.8 - 1.1 = -8.7$ ($H \rightarrow \gamma\gamma, ZZ^*$)
 $\sqrt{8.7} = 2.9$ while quoted local significance of excess = 3.6σ
- $2 \ln Q_{\text{ATLAS}}(244 \text{ GeV}) \approx 2 \ln Q_{\text{ATLAS}}(560 \text{ GeV}) \approx -3$ ($H \rightarrow ZZ$)
- $2 \ln Q_{\text{CMS}}(119.5 \text{ GeV}) = -5.6$ ($H \rightarrow ZZ^*, WW^*, b \bar{b}, \tau^+ \tau^-$)

Examples

- $2 \ln Q_{\text{ATLAS}}(126 \text{ GeV}) = 9.8 - 1.1 = -8.7$ ($H \rightarrow \gamma\gamma, ZZ^*$)
 $\sqrt{8.7} = 2.9$ while quoted local significance of excess = 3.6σ
- $2 \ln Q_{\text{ATLAS}}(244 \text{ GeV}) \approx 2 \ln Q_{\text{ATLAS}}(560 \text{ GeV}) \approx -3$ ($H \rightarrow ZZ$)
- $2 \ln Q_{\text{CMS}}(119.5 \text{ GeV}) = -5.6$ ($H \rightarrow ZZ^*, WW^*, b\bar{b}, \tau^+\tau^-$)
- $2 \ln Q_{\text{CMS}}(124 \text{ GeV}) = -6.6$ (mostly $H \rightarrow \gamma\gamma$)

Examples

- $2 \ln Q_{\text{ATLAS}}(126 \text{ GeV}) = 9.8 - 1.1 = -8.7$ ($H \rightarrow \gamma\gamma, ZZ^*$)
 $\sqrt{8.7} = 2.9$ while quoted local significance of excess = 3.6σ
- $2 \ln Q_{\text{ATLAS}}(244 \text{ GeV}) \approx 2 \ln Q_{\text{ATLAS}}(560 \text{ GeV}) \approx -3$ ($H \rightarrow ZZ$)
- $2 \ln Q_{\text{CMS}}(119.5 \text{ GeV}) = -5.6$ ($H \rightarrow ZZ^*, WW^*, b\bar{b}, \tau^+\tau^-$)
- $2 \ln Q_{\text{CMS}}(124 \text{ GeV}) = -6.6$ (mostly $H \rightarrow \gamma\gamma$)
- $2 \ln Q_{\text{Tevatron}}(130 \text{ GeV}) = -1.9$ ($H \rightarrow 165$ different channels)

Examples

- $2 \ln Q_{\text{ATLAS}}(126 \text{ GeV}) = 9.8 - 1.1 = -8.7$ ($H \rightarrow \gamma\gamma, ZZ^*$)
 $\sqrt{8.7} = 2.9$ while quoted local significance of excess = 3.6σ
- $2 \ln Q_{\text{ATLAS}}(244 \text{ GeV}) \approx 2 \ln Q_{\text{ATLAS}}(560 \text{ GeV}) \approx -3$ ($H \rightarrow ZZ$)
- $2 \ln Q_{\text{CMS}}(119.5 \text{ GeV}) = -5.6$ ($H \rightarrow ZZ^*, WW^*, b \bar{b}, \tau^+ \tau^-$)
- $2 \ln Q_{\text{CMS}}(124 \text{ GeV}) = -6.6$ (mostly $H \rightarrow \gamma\gamma$)
- $2 \ln Q_{\text{Tevatron}}(130 \text{ GeV}) = -1.9$ ($H \rightarrow 165$ different channels)
- $2 \ln Q_{\text{Tevatron 2012}}(120 \text{ GeV}) = -8.0$ (mostly $H \rightarrow b \bar{b}$, not yet included)

Examples

- $2 \ln Q_{\text{ATLAS}}(126 \text{ GeV}) = 9.8 - 1.1 = -8.7$ ($H \rightarrow \gamma\gamma, ZZ^*$)
 $\sqrt{8.7} = 2.9$ while quoted local significance of excess = 3.6σ
- $2 \ln Q_{\text{ATLAS}}(244 \text{ GeV}) \approx 2 \ln Q_{\text{ATLAS}}(560 \text{ GeV}) \approx -3$ ($H \rightarrow ZZ$)
- $2 \ln Q_{\text{CMS}}(119.5 \text{ GeV}) = -5.6$ ($H \rightarrow ZZ^*, WW^*, b \bar{b}, \tau^+ \tau^-$)
- $2 \ln Q_{\text{CMS}}(124 \text{ GeV}) = -6.6$ (mostly $H \rightarrow \gamma\gamma$)
- $2 \ln Q_{\text{Tevatron}}(130 \text{ GeV}) = -1.9$ ($H \rightarrow 165$ different channels)
- $2 \ln Q_{\text{Tevatron 2012}}(120 \text{ GeV}) = -8.0$ (mostly $H \rightarrow b \bar{b}$, not yet included)
- $2 \ln Q_{\text{LEP}}(117 \text{ GeV}) = -1.7$ ($H \rightarrow 4$ jets *ALEPH*)

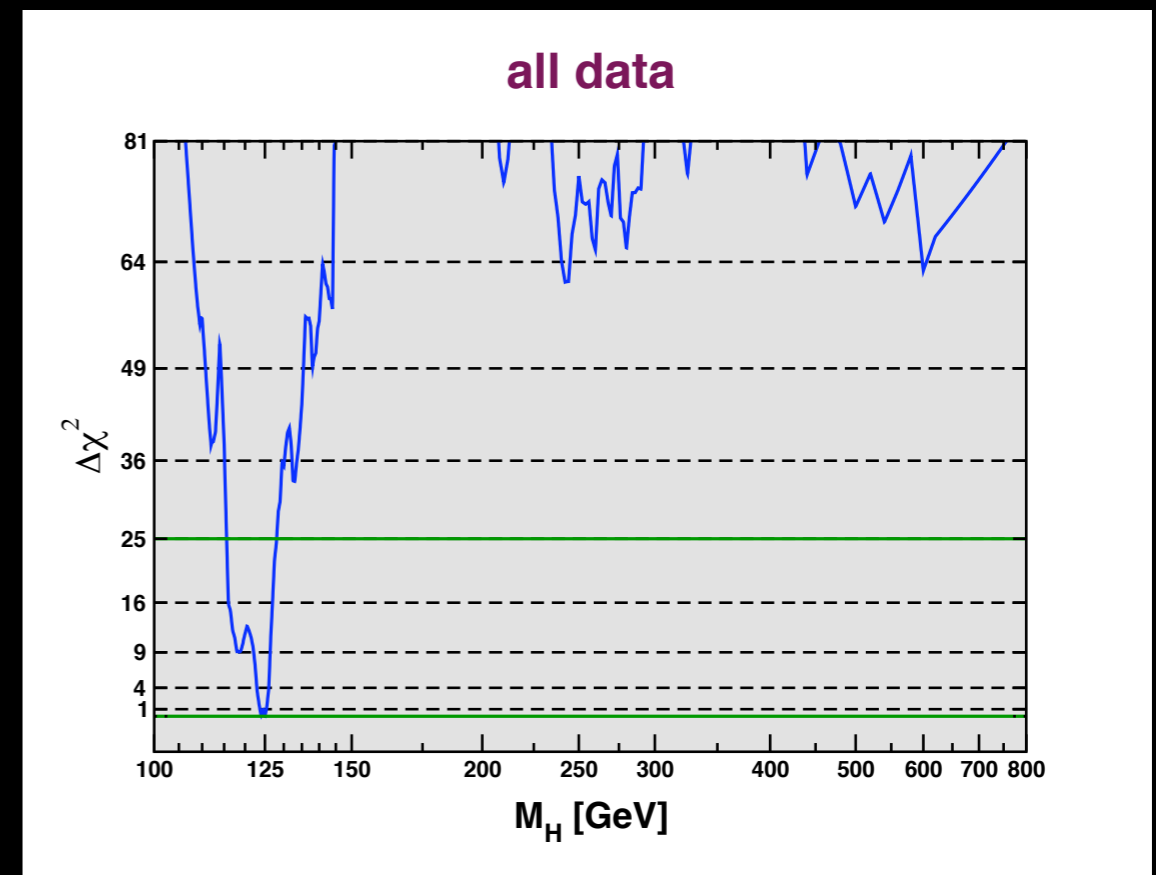
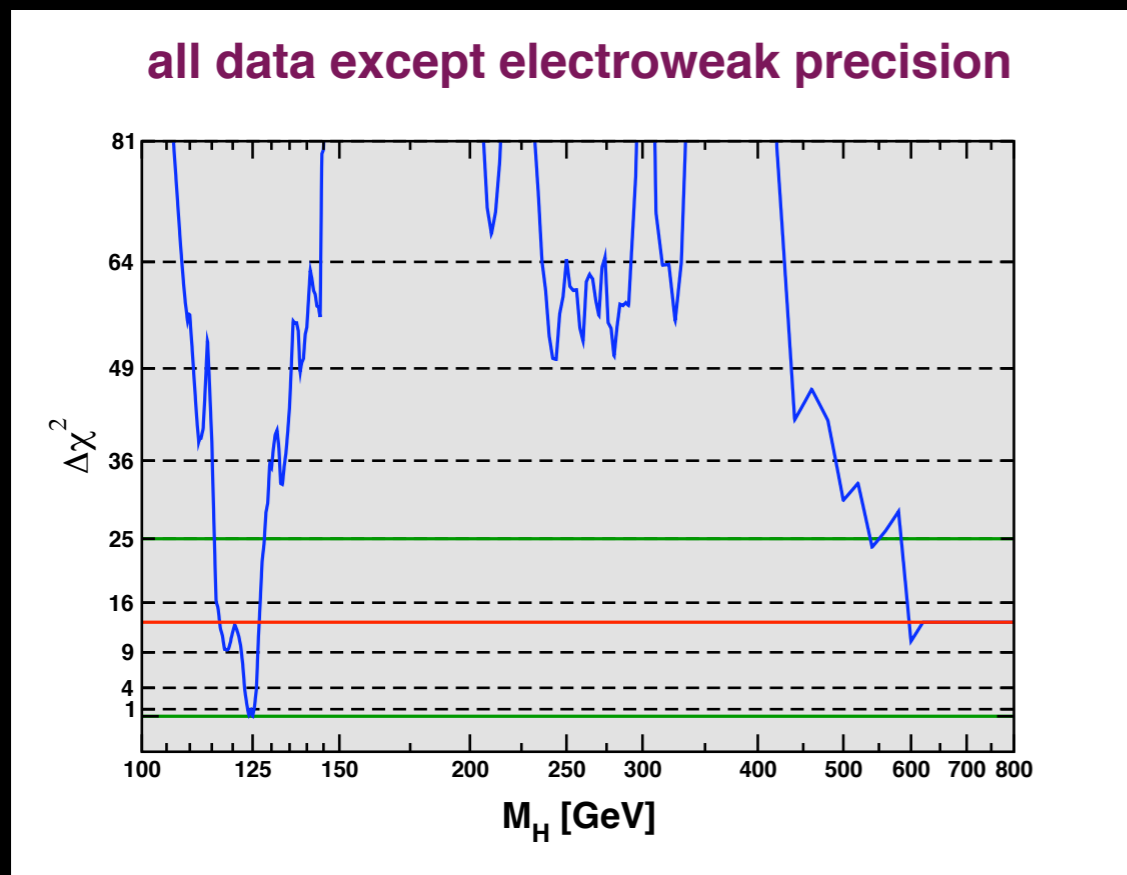
Examples

- $2 \ln Q_{\text{ATLAS}}(126 \text{ GeV}) = 9.8 - 1.1 = -8.7$ ($H \rightarrow \gamma\gamma, ZZ^*$)
 $\sqrt{8.7} = 2.9$ while quoted local significance of excess = 3.6σ
- $2 \ln Q_{\text{ATLAS}}(244 \text{ GeV}) \approx 2 \ln Q_{\text{ATLAS}}(560 \text{ GeV}) \approx -3$ ($H \rightarrow ZZ$)
- $2 \ln Q_{\text{CMS}}(119.5 \text{ GeV}) = -5.6$ ($H \rightarrow ZZ^*, WW^*, b\bar{b}, \tau^+\tau^-$)
- $2 \ln Q_{\text{CMS}}(124 \text{ GeV}) = -6.6$ (mostly $H \rightarrow \gamma\gamma$)
- $2 \ln Q_{\text{Tevatron}}(130 \text{ GeV}) = -1.9$ ($H \rightarrow 165$ different channels)
- $2 \ln Q_{\text{Tevatron 2012}}(120 \text{ GeV}) = -8.0$ (mostly $H \rightarrow b\bar{b}$, not yet included)
- $2 \ln Q_{\text{LEP}}(117 \text{ GeV}) = -1.7$ ($H \rightarrow 4$ jets *ALEPH*)
- $\chi^2_{\text{EW}}(127 \text{ GeV}) - \chi^2_{\text{EW}}(115.5 \text{ GeV}) = 0.63$

Examples

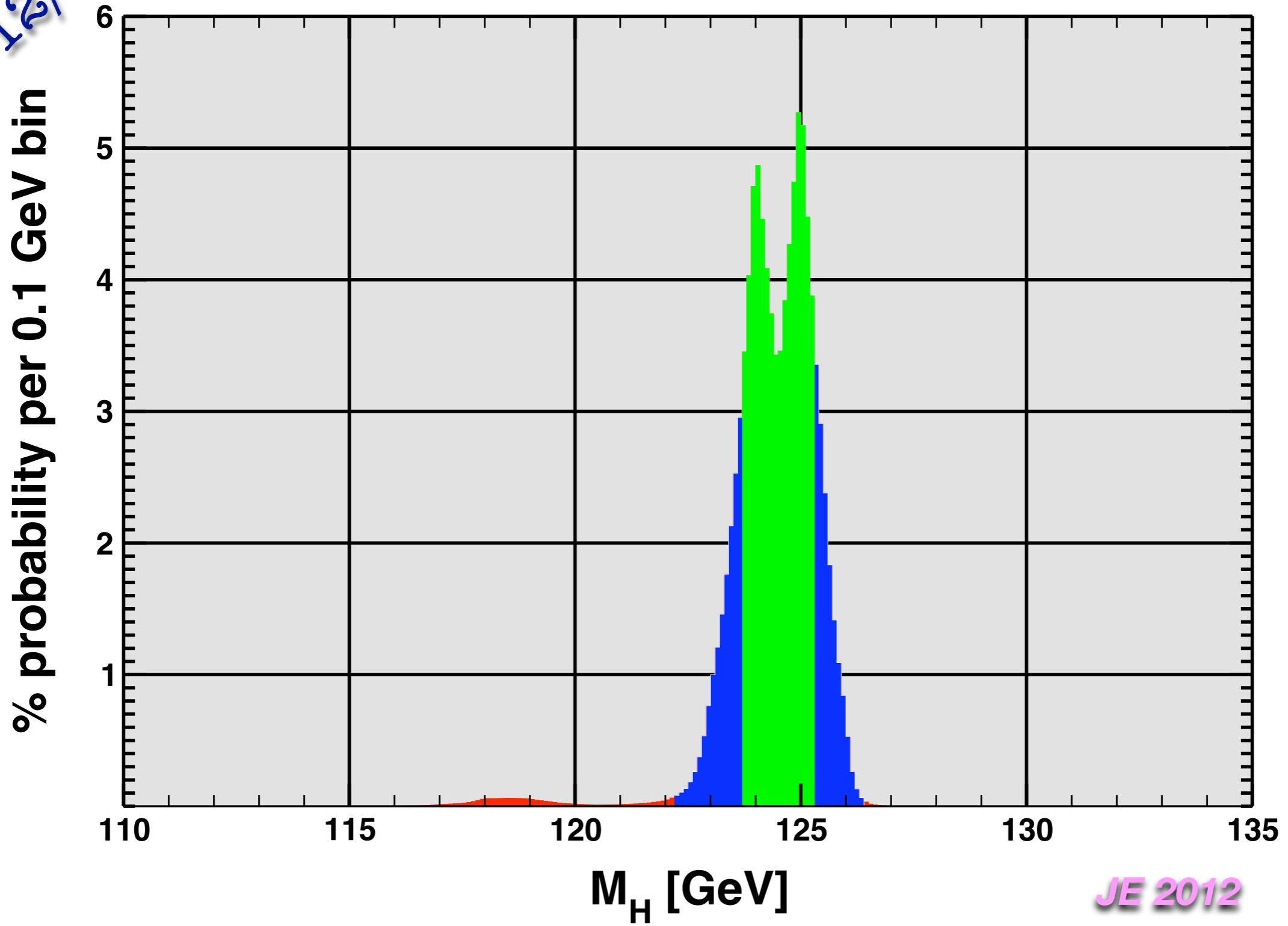
- $2 \ln Q_{\text{ATLAS}}(126 \text{ GeV}) = 9.8 - 1.1 = -8.7$ ($H \rightarrow \gamma\gamma, ZZ^*$)
 $\sqrt{8.7} = 2.9$ while quoted local significance of excess = 3.6σ
- $2 \ln Q_{\text{ATLAS}}(244 \text{ GeV}) \approx 2 \ln Q_{\text{ATLAS}}(560 \text{ GeV}) \approx -3$ ($H \rightarrow ZZ$)
- $2 \ln Q_{\text{CMS}}(119.5 \text{ GeV}) = -5.6$ ($H \rightarrow ZZ^*, WW^*, b \bar{b}, \tau^+ \tau^-$)
- $2 \ln Q_{\text{CMS}}(124 \text{ GeV}) = -6.6$ (mostly $H \rightarrow \gamma\gamma$)
- $2 \ln Q_{\text{Tevatron}}(130 \text{ GeV}) = -1.9$ ($H \rightarrow 165$ different channels)
- $2 \ln Q_{\text{Tevatron } 2012}(120 \text{ GeV}) = -8.0$ (mostly $H \rightarrow b \bar{b}$, not yet included)
- $2 \ln Q_{\text{LEP}}(117 \text{ GeV}) = -1.7$ ($H \rightarrow 4$ jets *ALEPH*)
- $\chi^2_{\text{EW}}(127 \text{ GeV}) - \chi^2_{\text{EW}}(115.5 \text{ GeV}) = 0.63$
- $2 \ln p_{\text{direct}}(125 \text{ GeV}) = -13.2$

- LHC data require “look elsewhere effect correction”
- Can be avoided when combined with electroweak precision data *JE 2012*



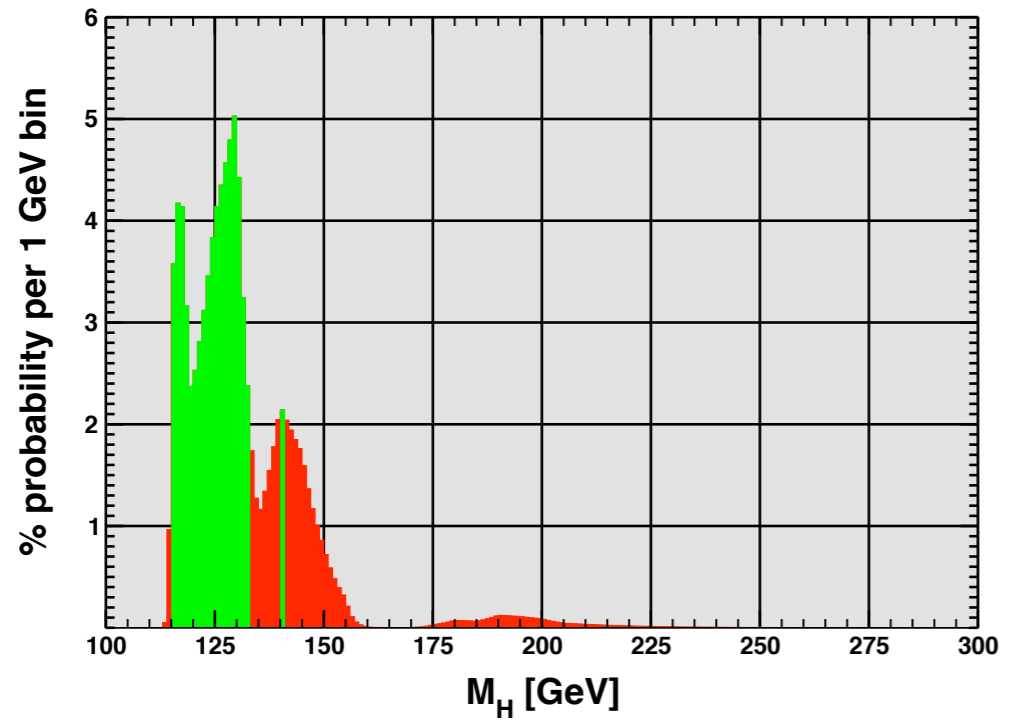
12/2011

all data

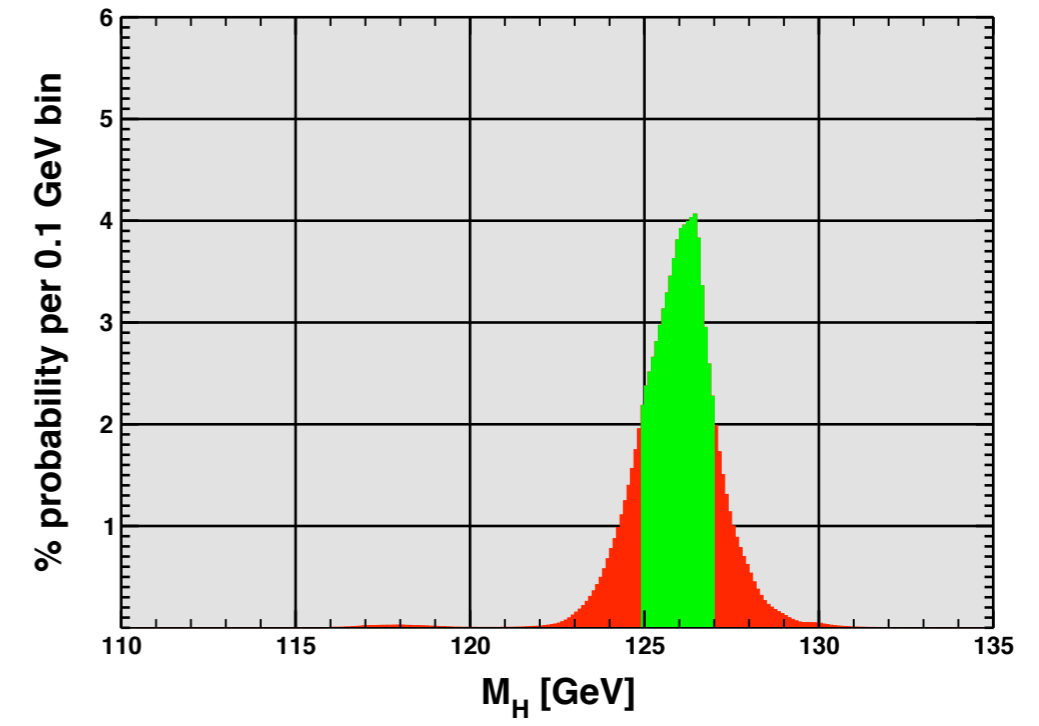


JE 2012

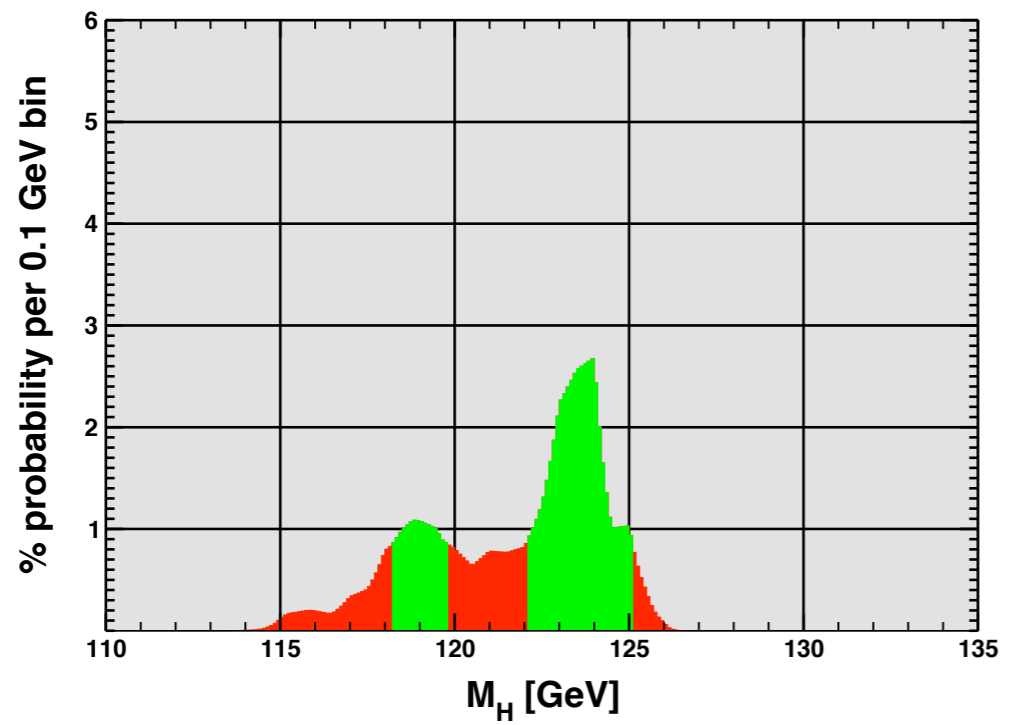
all data except LHC



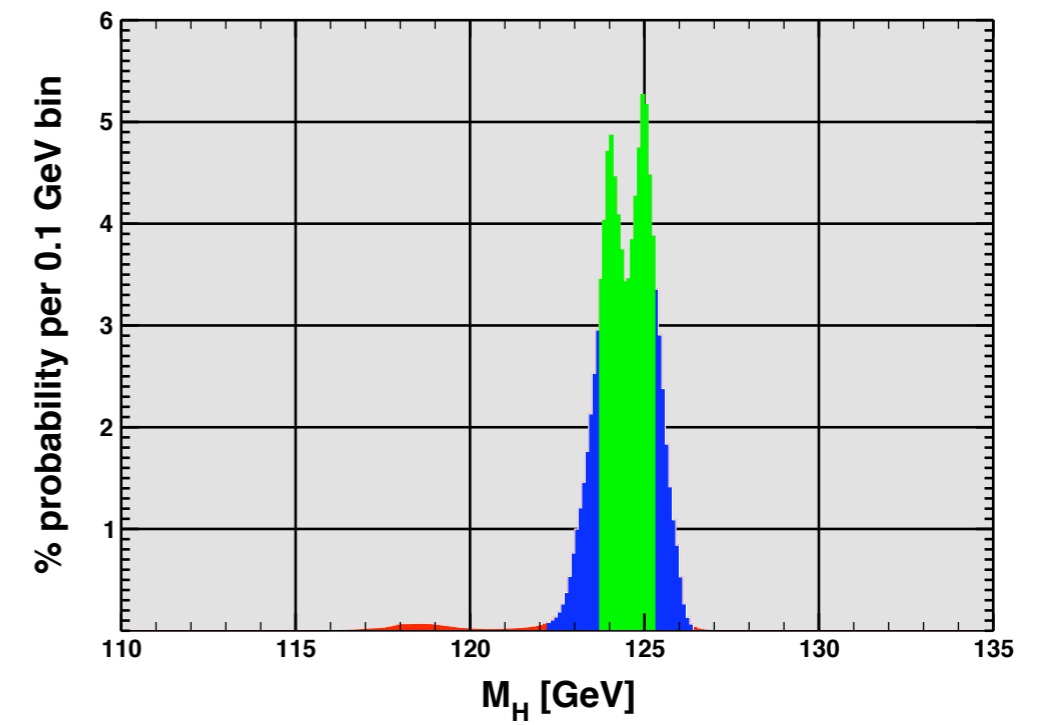
all data except CMS



all data except ATLAS



all data



conclusions

Summary and outlook

Summary and outlook



Summary and outlook

- Remarkably consistent picture between **QFT** constraints, electroweak **precision tests** and **direct search** results.



Summary and outlook

- Remarkably consistent picture between QFT constraints, electroweak precision tests and direct search results.
- It walks, quacks and looks like a duck!



Summary and outlook

- Remarkably consistent picture between QFT constraints, electroweak precision tests and direct search results.
 - It walks, quacks and looks like a duck!
- Highly likely that the discovery (“observation”) of the Higgs particle will be announced in the course of this year, with an “evidence” claim in the summer (unless lucky or unlucky fluctuations occur).



Summary and outlook

- Remarkably consistent picture between QFT constraints, electroweak precision tests and direct search results.
 - It walks, quacks and looks like a duck!
- Highly likely that the discovery (“observation”) of the Higgs particle will be announced in the course of this year, with an “evidence” claim in the summer (unless lucky or unlucky fluctuations occur).
- The attention will move rapidly (or is already moving) to the question whether the Higgs is Standard Model-like, or more to the point how Standard Model-like.



Summary and outlook

- Remarkably consistent picture between QFT constraints, electroweak precision tests and direct search results.
 - It walks, quacks and looks like a duck!
- Highly likely that the discovery (“observation”) of the Higgs particle will be announced in the course of this year, with an “evidence” claim in the summer (unless lucky or unlucky fluctuations occur).
- The attention will move rapidly (or is already moving) to the question whether the Higgs is Standard Model-like, or more to the point how Standard Model-like.
 - Is it a mutated duck?



Summary and outlook

- Remarkably consistent picture between QFT constraints, electroweak precision tests and direct search results.
 - It walks, quacks and looks like a duck!
- Highly likely that the discovery (“observation”) of the Higgs particle will be announced in the course of this year, with an “evidence” claim in the summer (unless lucky or unlucky fluctuations occur).
- The attention will move rapidly (or is already moving) to the question whether the Higgs is Standard Model-like, or more to the point how Standard Model-like.
 - Is it a mutated duck?

