

Density Perturbations from Curvatons Revisited

Takeshi Kobayashi (CITA)

based on: arXiv:1107.6011 w/ M. Kawasaki, F. Takahashi
arXiv:1203.3011 w/ T. Takahashi

PASCOS 2012

The Curvaton Paradigm

Enqvist, Sloth '01 Lyth, Wands '01 Moroi, Takahashi '01

generates cosmological perturbations from
field fluctuations of a “curvaton” field

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however...

many studies have been limited to rather trivial
curvaton energy potentials, e.g., **quadratic** ones

Why Consider Non-Quadratic Curvatons?

- microscopic realizations can give complicated energy potentials
- observational constraints on the spectral index requires a **tachyonic** potential, or rather large $|\dot{H}/H^2|$

$$n_s - 1 = \frac{2}{3} \frac{V''}{H^2} + 2 \frac{\dot{H}}{H^2} = -0.032 \pm 0.012 \quad (\text{WMAP7, 68\%CL})$$

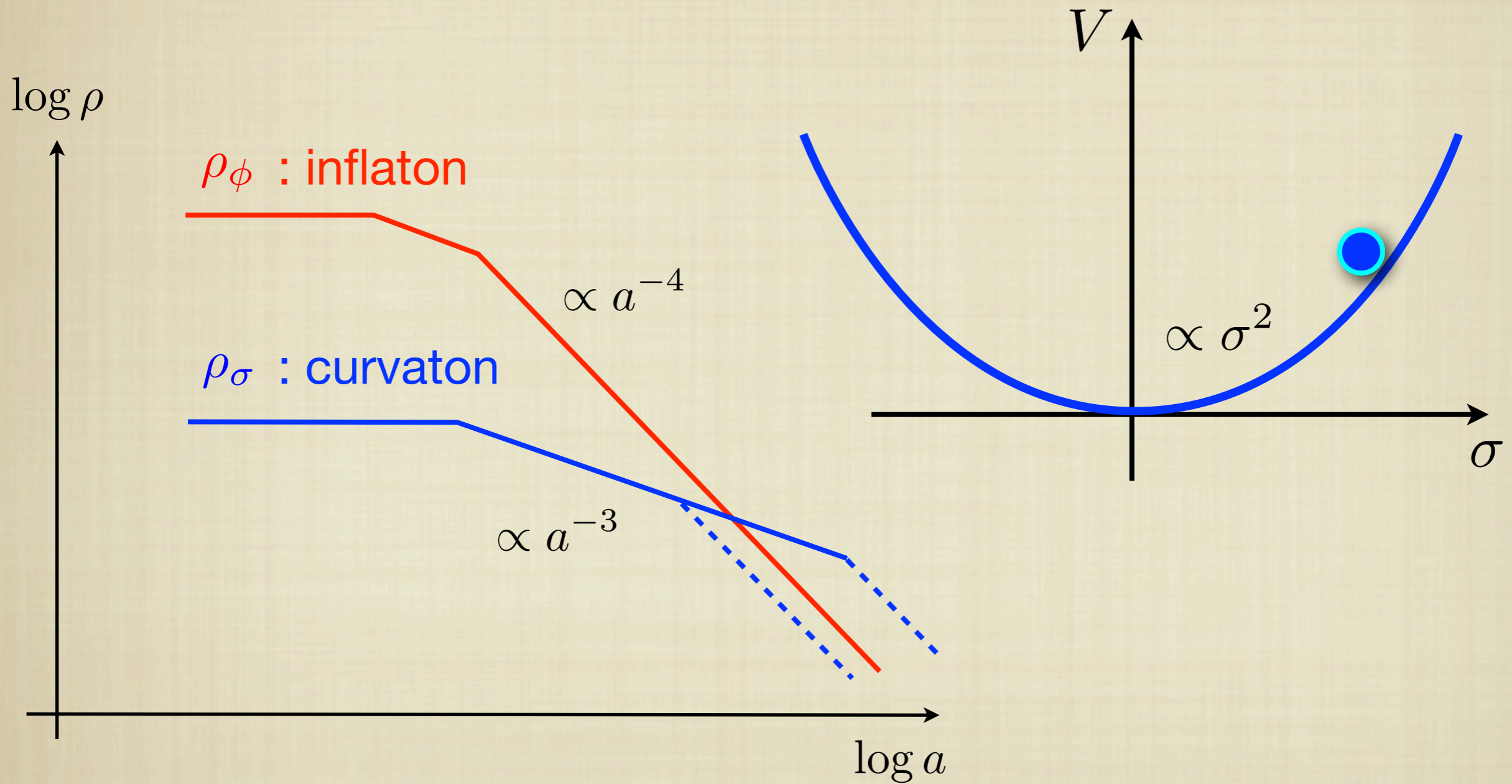
This Work :

Curvatons with Generic Potentials

- non-quadratic curvatons behave quite differently from quadratic ones
- interesting features in cosmological observables, especially in non-Gaussianity
- Our analyses are analytic!

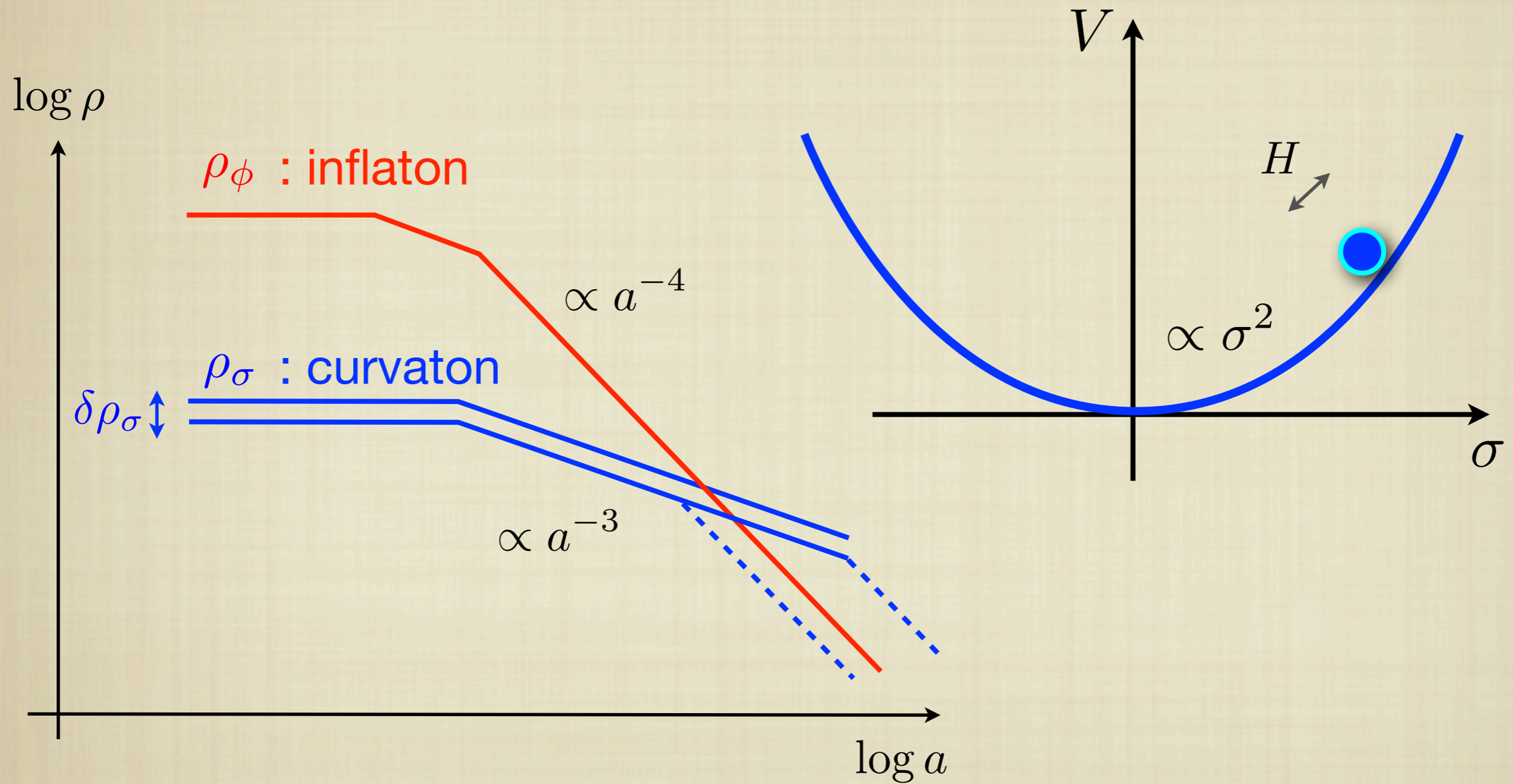
The Quadratic Curvaton Scenario

: a light field with negligible energy density during inflation



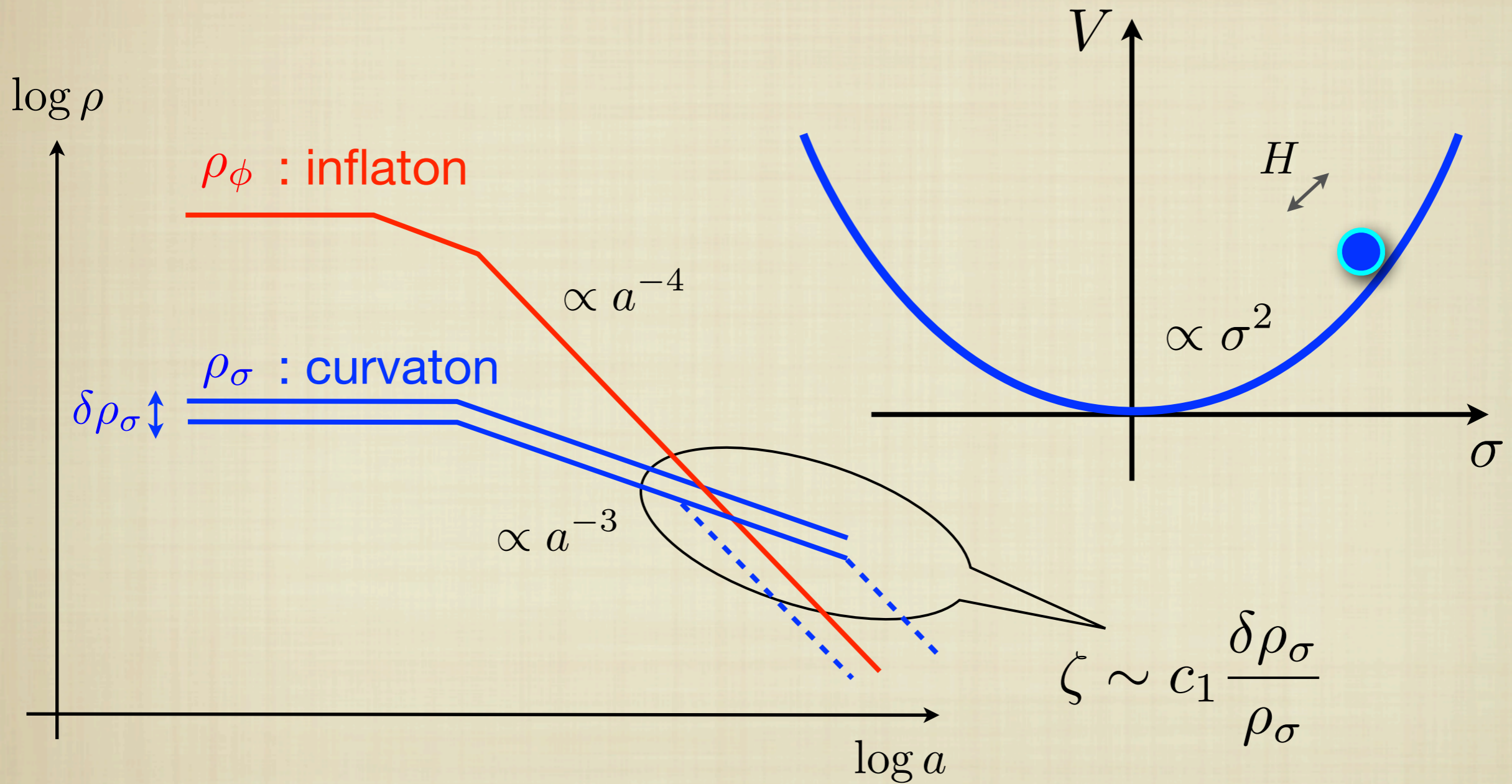
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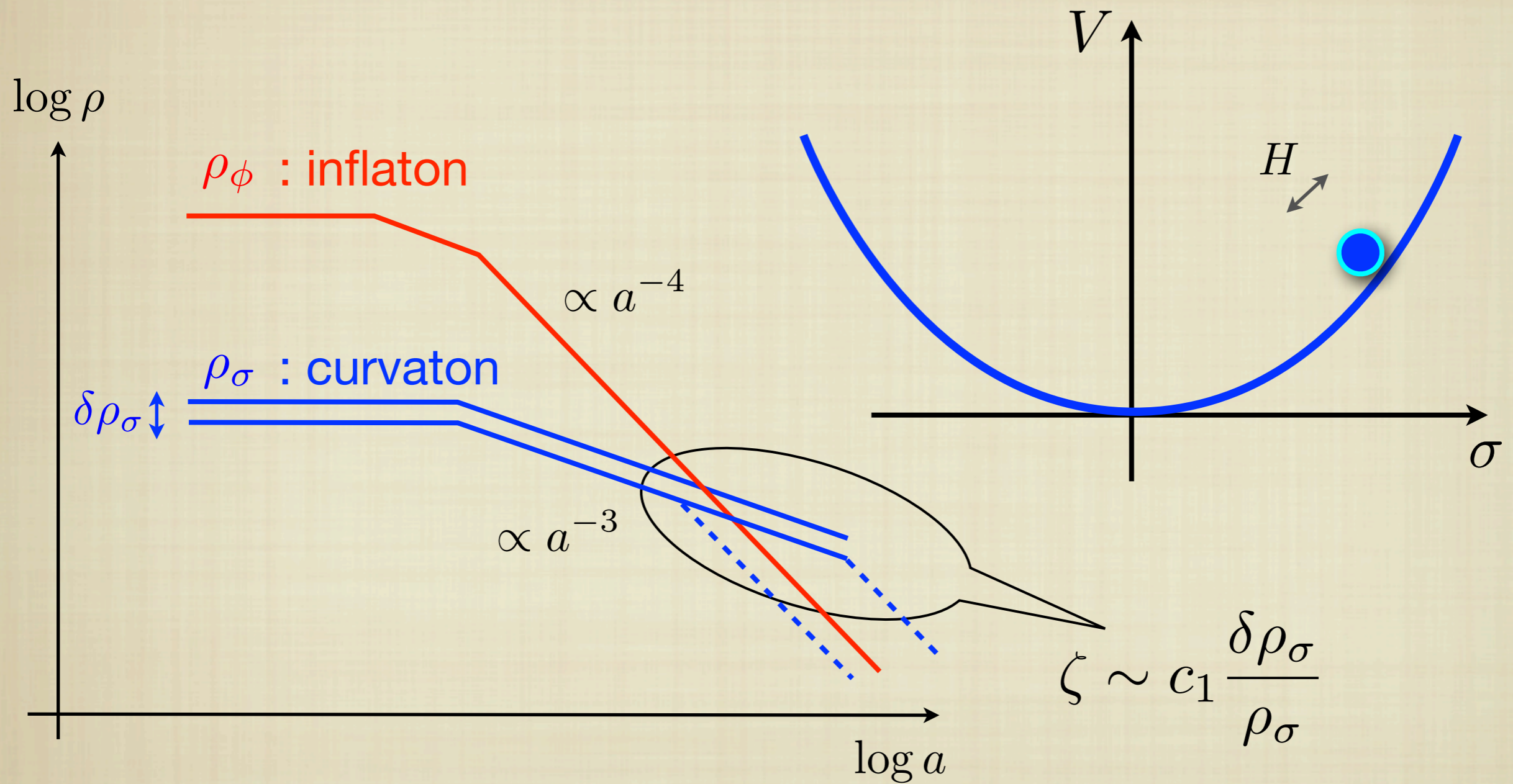


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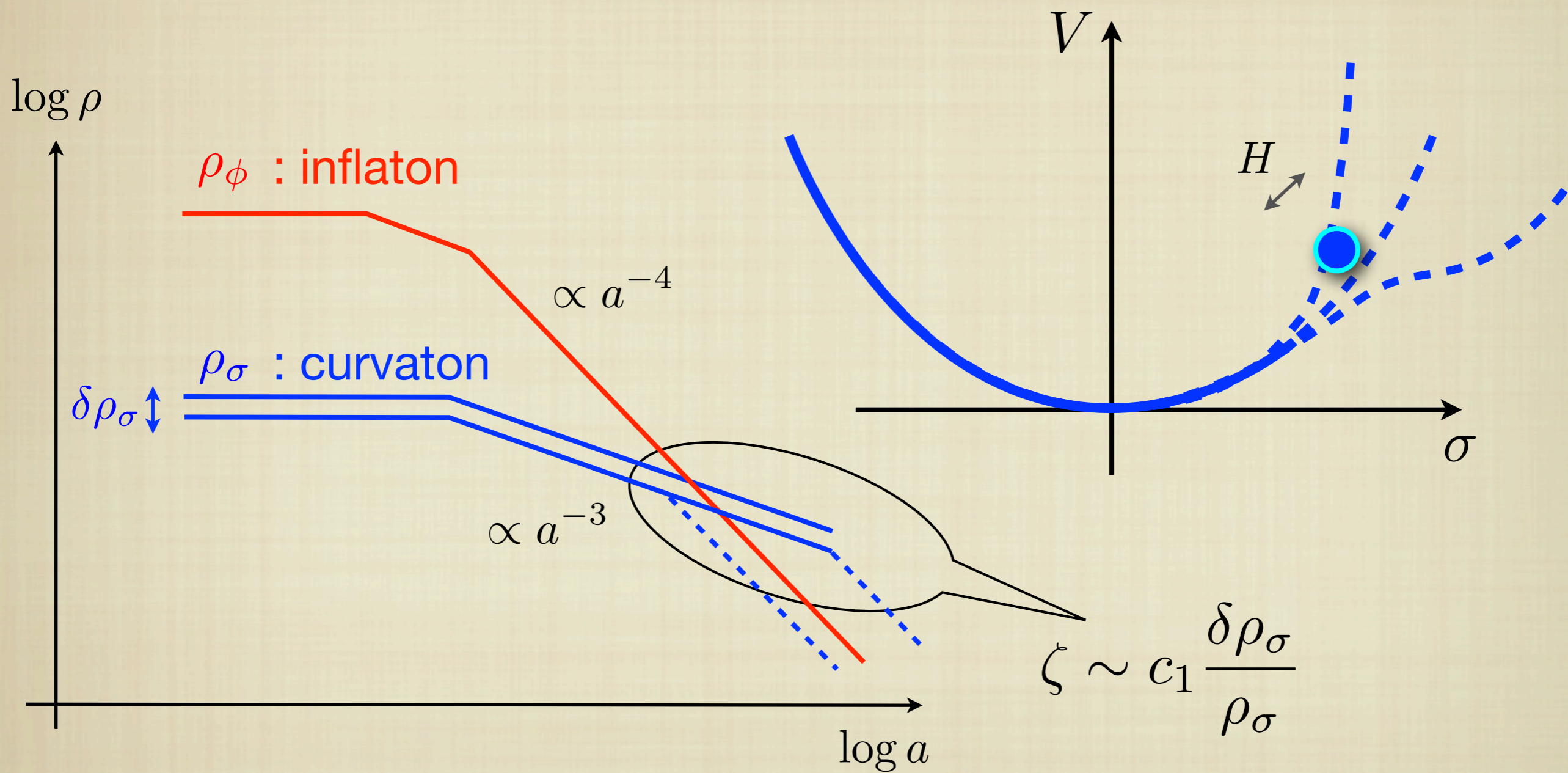
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Curvatons with Arbitrary Potentials

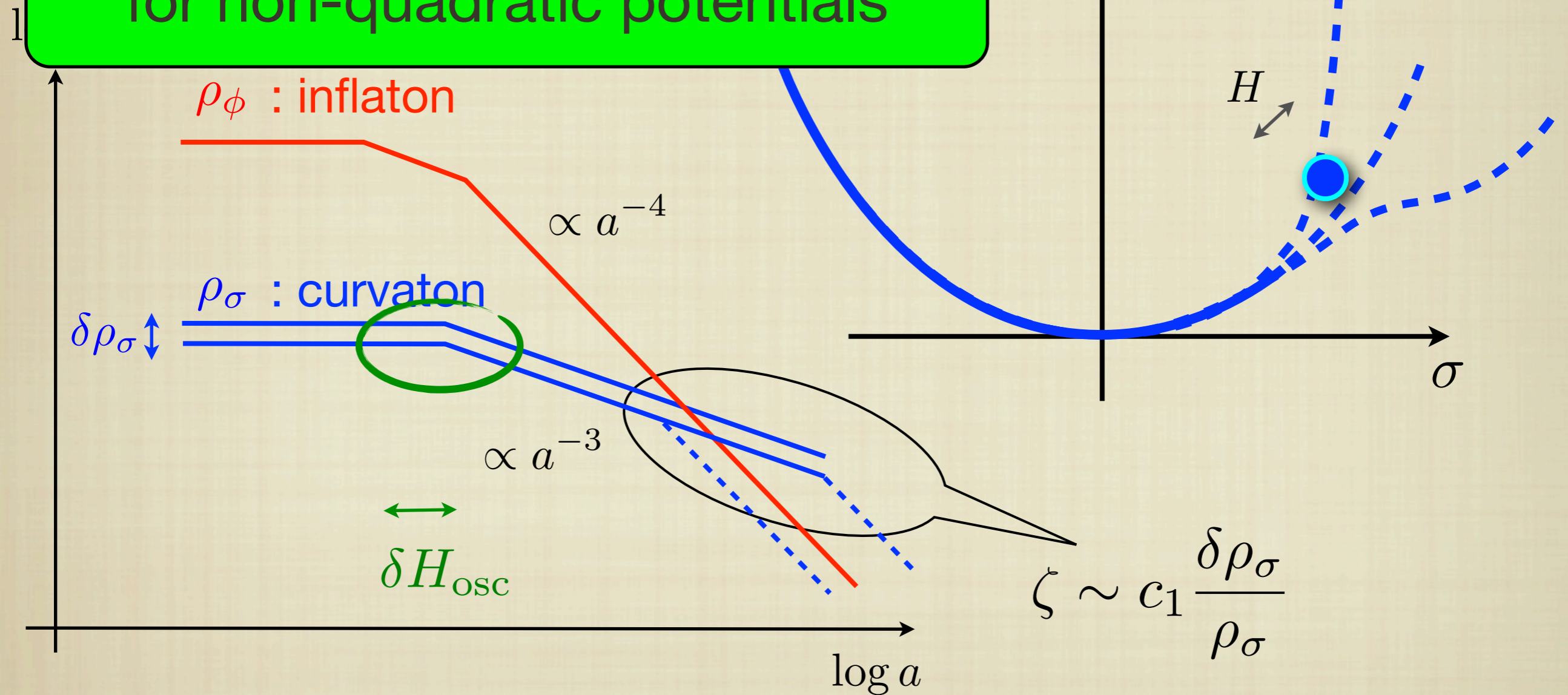


Curvatons with Arbitrary Potentials



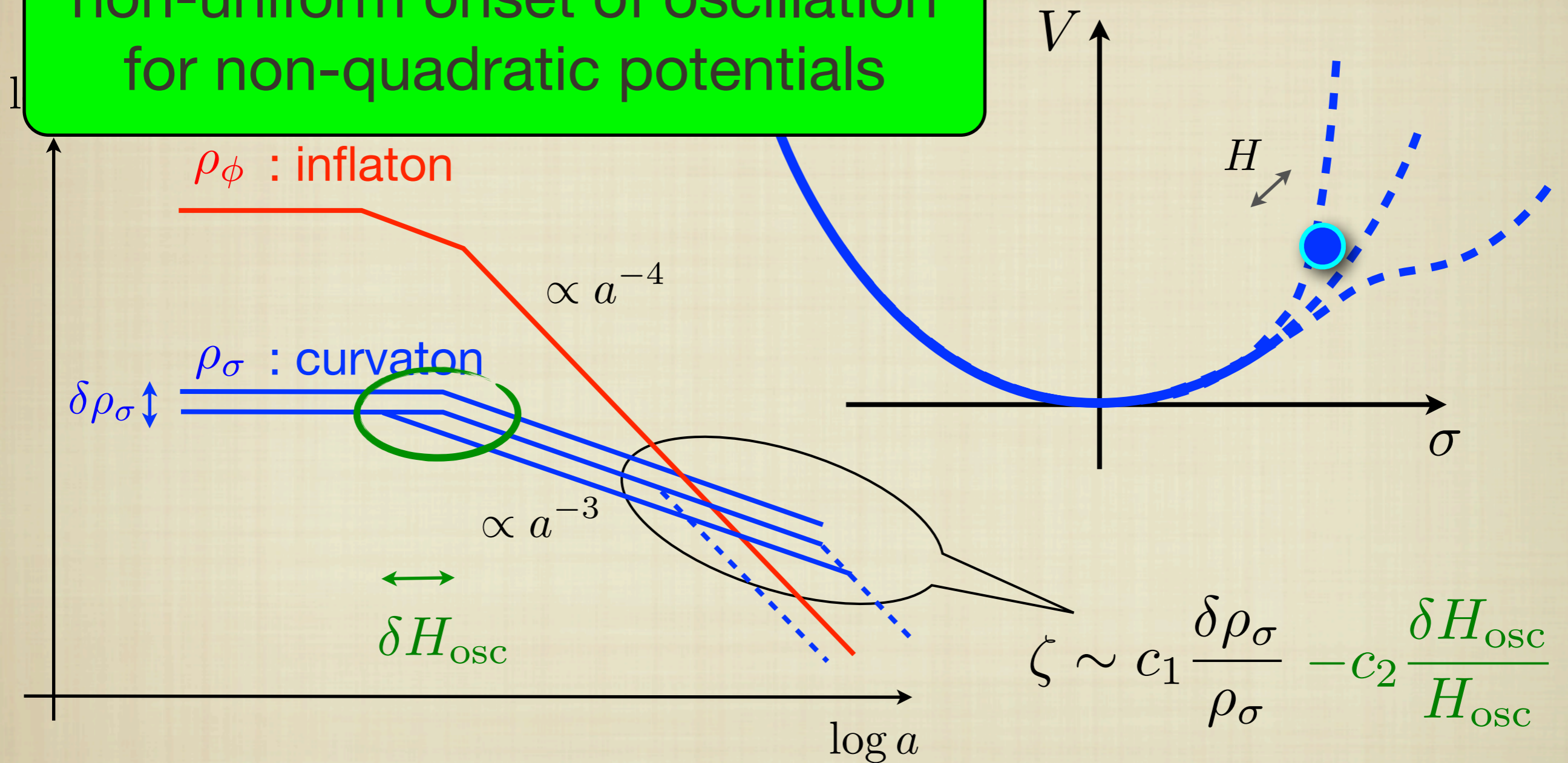
Curvaton with Arbitrary Potentials

non-uniform onset of oscillation
for non-quadratic potentials



Curvaton with Arbitrary Potentials

non-uniform onset of oscillation
for non-quadratic potentials



Additional contributions to the density perturbations!

Density Perturbations

$$\mathcal{P}_\zeta = \left(\frac{\partial \mathcal{N}}{\partial \sigma_*} \frac{H_*}{2\pi} \right)^2$$

$$\frac{\partial \mathcal{N}}{\partial \sigma_*} = \frac{r}{4 + 3r} (1 - X(\sigma_{\text{osc}}))^{-1} \left\{ \frac{V'(\sigma_{\text{osc}})}{V(\sigma_{\text{osc}})} - \frac{3X(\sigma_{\text{osc}})}{\sigma_{\text{osc}}} \right\} \frac{V'(\sigma_{\text{osc}})}{V'(\sigma_*)}$$

$$r \equiv \frac{\rho_\sigma}{\rho_r} \quad @ \text{ curvaton decay}$$

* : @ horizon exit

osc : @ onset of curvaton oscillation

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$$X(\sigma_{\text{osc}}) \equiv \frac{1}{2(c-3)} \left(\frac{\sigma_{\text{osc}} V''(\sigma_{\text{osc}})}{V'(\sigma_{\text{osc}})} - 1 \right)$$

: effects due to non-uniform onset of oscillation

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$$\text{spectral index} \quad n_s - 1 \equiv \frac{d \ln \mathcal{P}_\zeta}{d \ln k} = \frac{2}{3} \frac{V''(\sigma_*)}{H_*^2} + 2 \frac{\dot{H}_*}{H_*^2}$$

observational data = -0.032 ± 0.012 (WMAP7, 68%CL)

requires a tachyonic curvaton, or rather large \dot{H}

Non-Gaussianity

$$\begin{aligned}
 f_{\text{NL}} = & \frac{40(1+r)}{3r(4+3r)} + \frac{5(4+3r)}{6r} \left\{ \frac{V'(\sigma_{\text{osc}})}{V(\sigma_{\text{osc}})} - \frac{3X(\sigma_{\text{osc}})}{\sigma_{\text{osc}}} \right\}^{-1} \left[(1-X(\sigma_{\text{osc}}))^{-1} X'(\sigma_{\text{osc}}) \right. \\
 & + \left. \left\{ \frac{V'(\sigma_{\text{osc}})}{V(\sigma_{\text{osc}})} - \frac{3X(\sigma_{\text{osc}})}{\sigma_{\text{osc}}} \right\}^{-1} \left\{ \frac{V''(\sigma_{\text{osc}})}{V(\sigma_{\text{osc}})} - \frac{V'(\sigma_{\text{osc}})^2}{V(\sigma_{\text{osc}})^2} - \frac{3X'(\sigma_{\text{osc}})}{\sigma_{\text{osc}}} + \frac{3X(\sigma_{\text{osc}})}{\sigma_{\text{osc}}^2} \right\} \right. \\
 & \left. \left. + \frac{V''(\sigma_{\text{osc}})}{V'(\sigma_{\text{osc}})} - (1-X(\sigma_{\text{osc}})) \frac{V''(\sigma_*)}{V'(\sigma_{\text{osc}})} \right]
 \end{aligned}$$

$$r \equiv \frac{\rho_\sigma}{\rho_r} \quad @ \text{ curvaton decay}$$

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cf. quadratic curvatons : $f_{\text{NL}} \sim \frac{1}{r}$

$f_{\text{NL}} \gg 1$ only for curvatons decaying when subdominant ($r \ll 1$)

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Large f_{NL} (with either sign) possible for both dominant/subdominant curvatons!

$$r \equiv \frac{\rho_\sigma}{\rho_r} \quad @ \text{ curvaton decay}$$

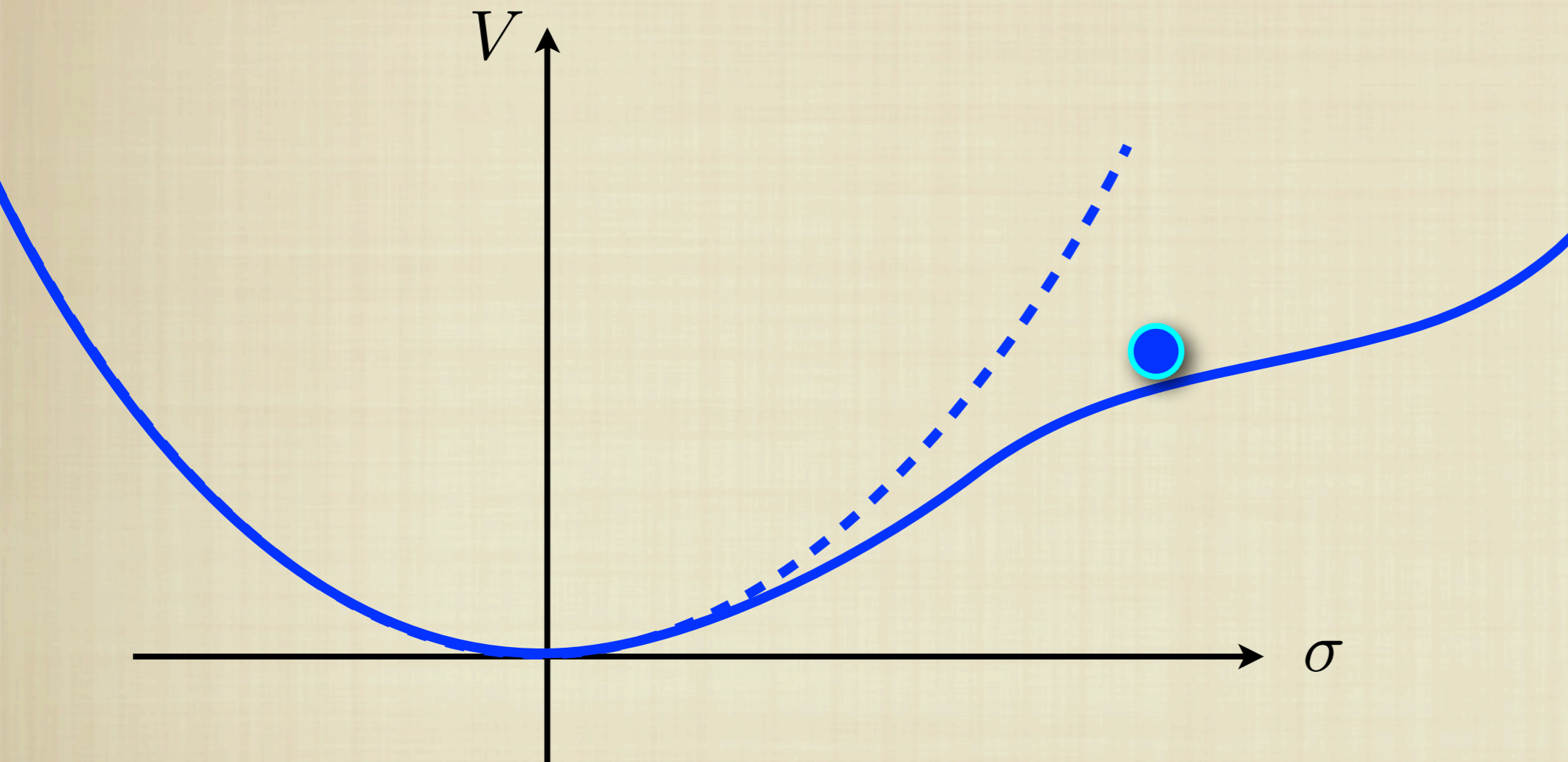
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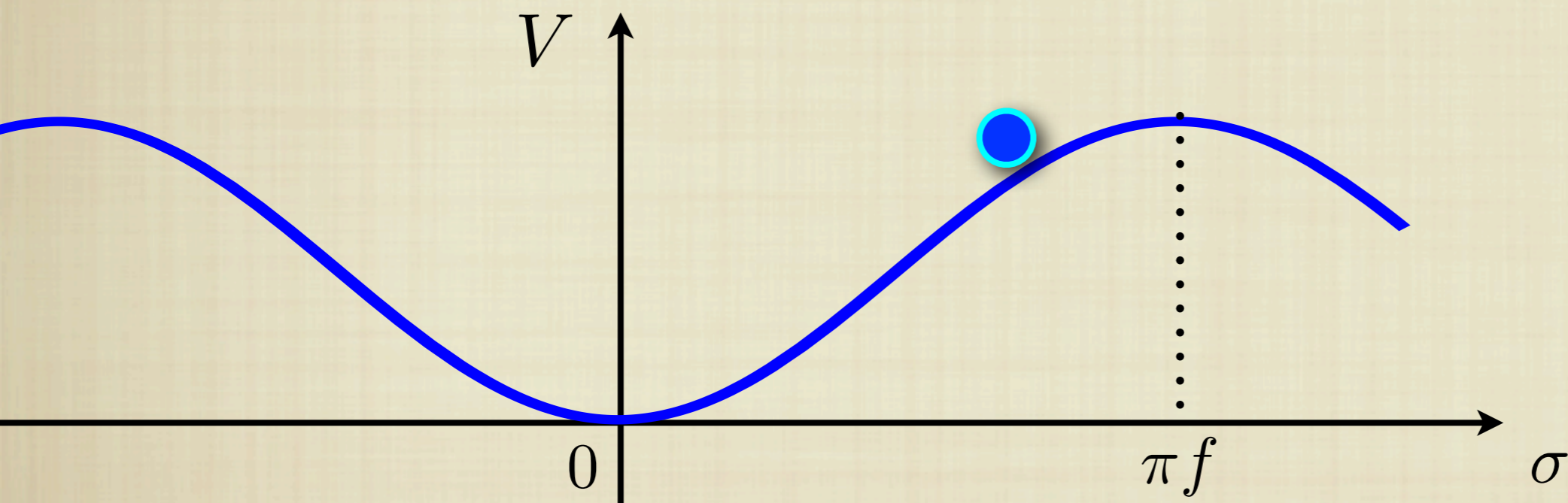
1. Flat Potentials



Effects due to non-uniform onset of oscillation can become significant, leading to **strong enhancement of linear-order perturbations & non-Gaussianity.**

case study : Curvaton =
pseudo-NG boson of a broken U(1)
symmetry

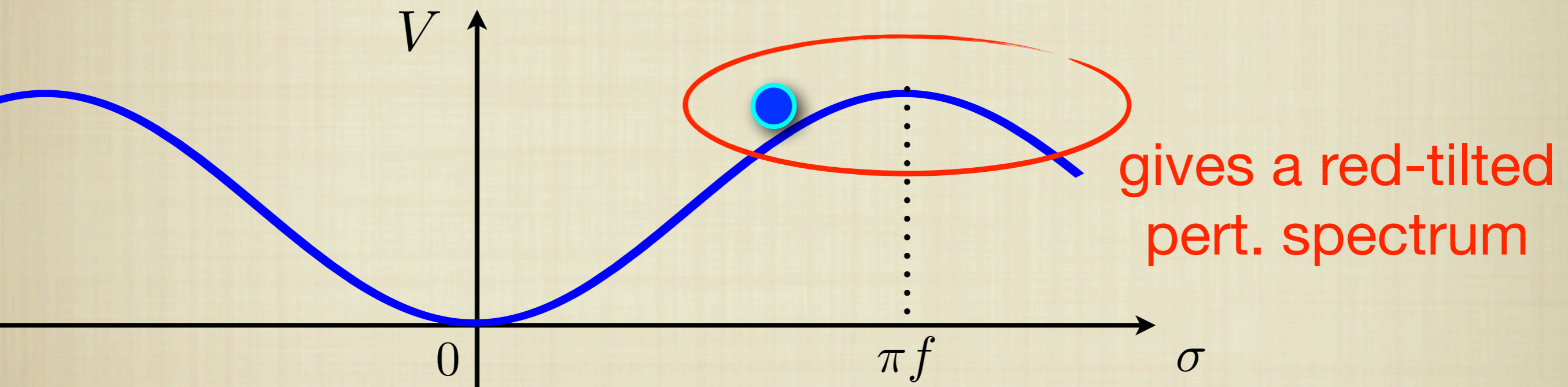
$$V(\sigma) = \Lambda^4 \left[1 - \cos \left(\frac{\sigma}{f} \right) \right]$$



curvaton decay rate : $\Gamma_\sigma \sim \frac{1}{16\pi} \frac{m^3}{f^2} = \frac{1}{16\pi} \frac{\Lambda^6}{f^5}$

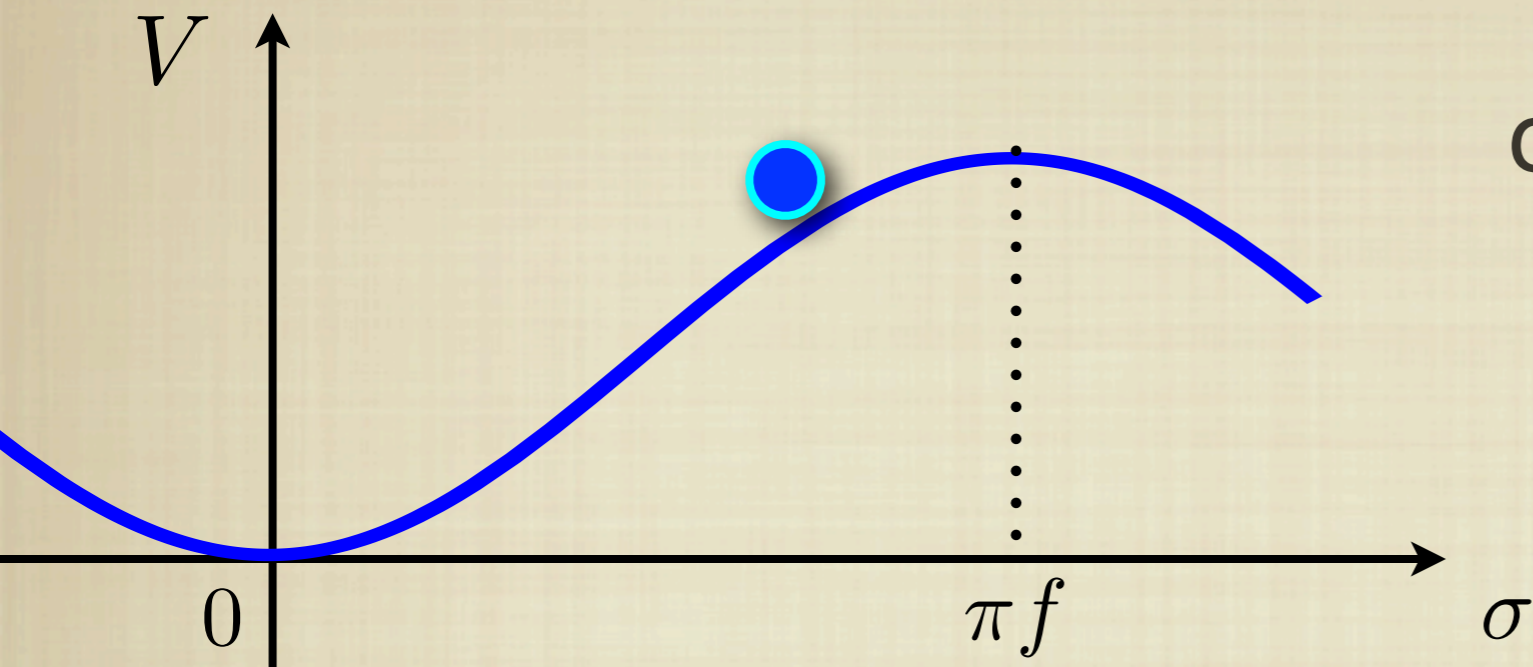
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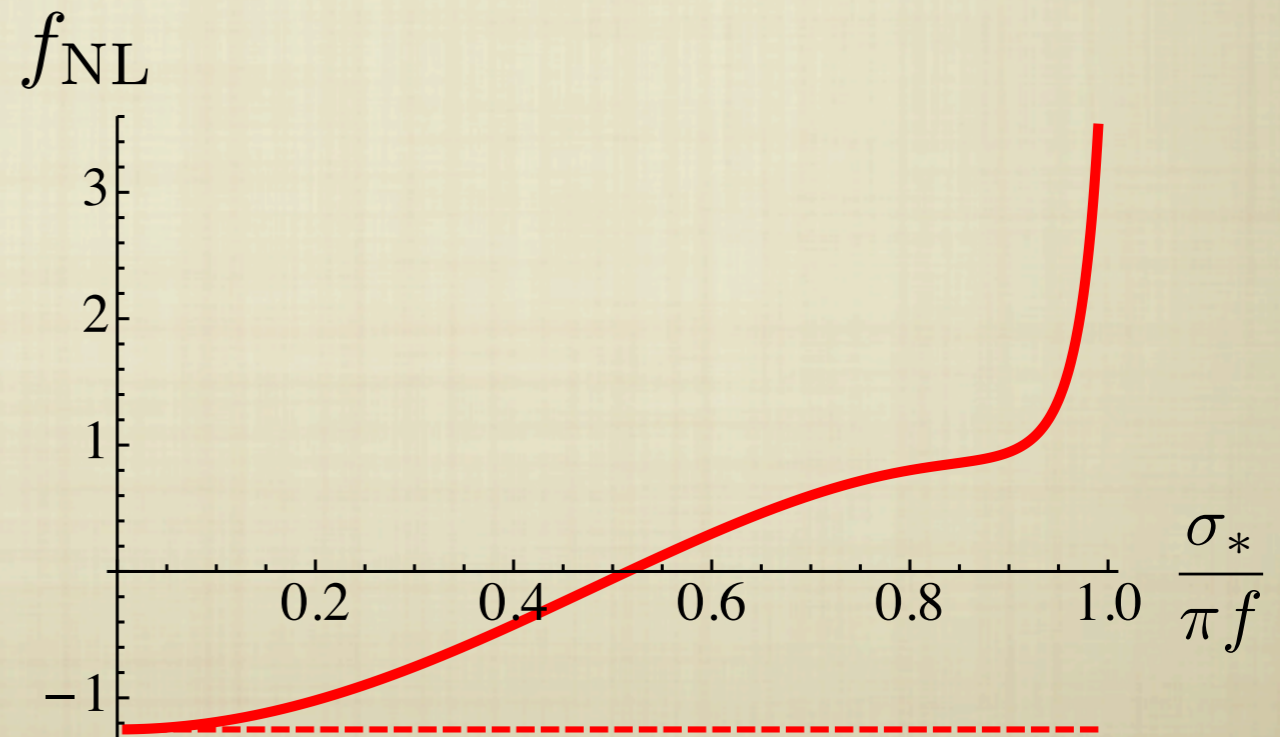
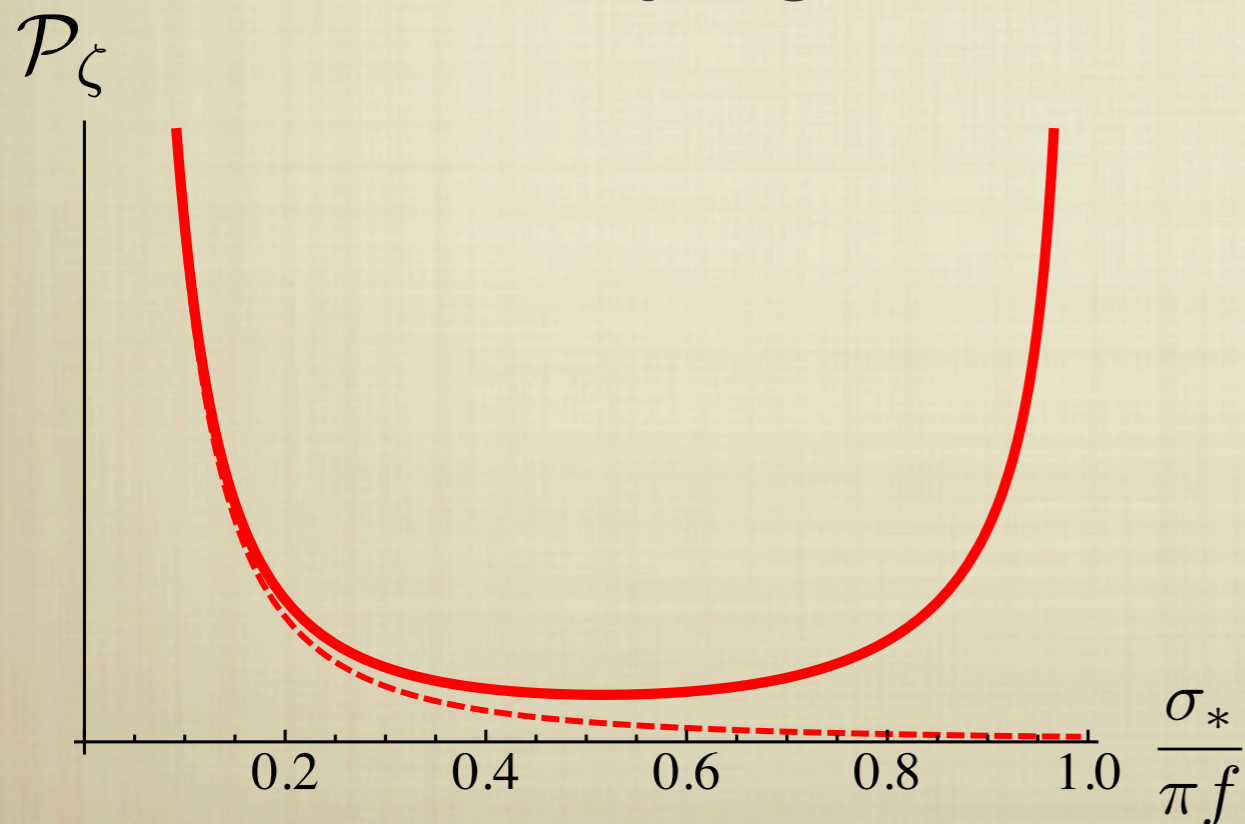
Density Pert. from a NG-Curvaton



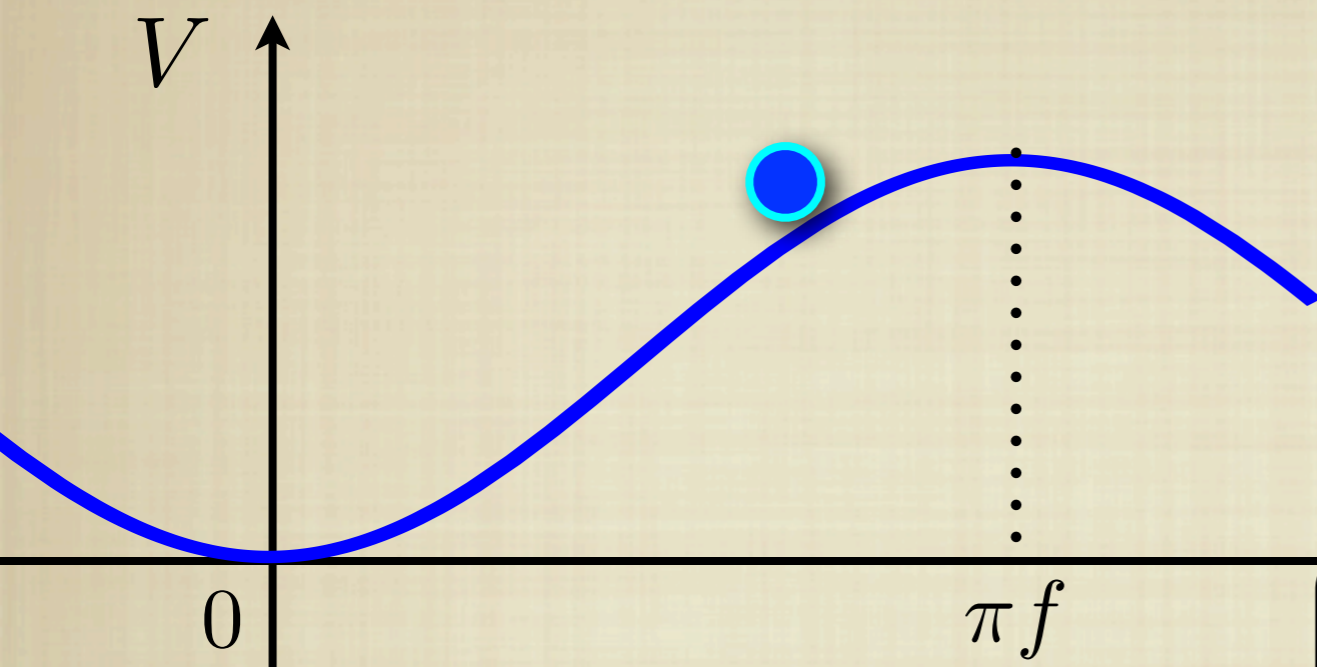
curvaton dominant case,

$$\text{i.e. } r \equiv \left. \frac{\rho_\sigma}{\rho_r} \right|_{\text{dec}} \gg 1$$

When varying σ_* :



Density Pert. from a NG-Curvaton

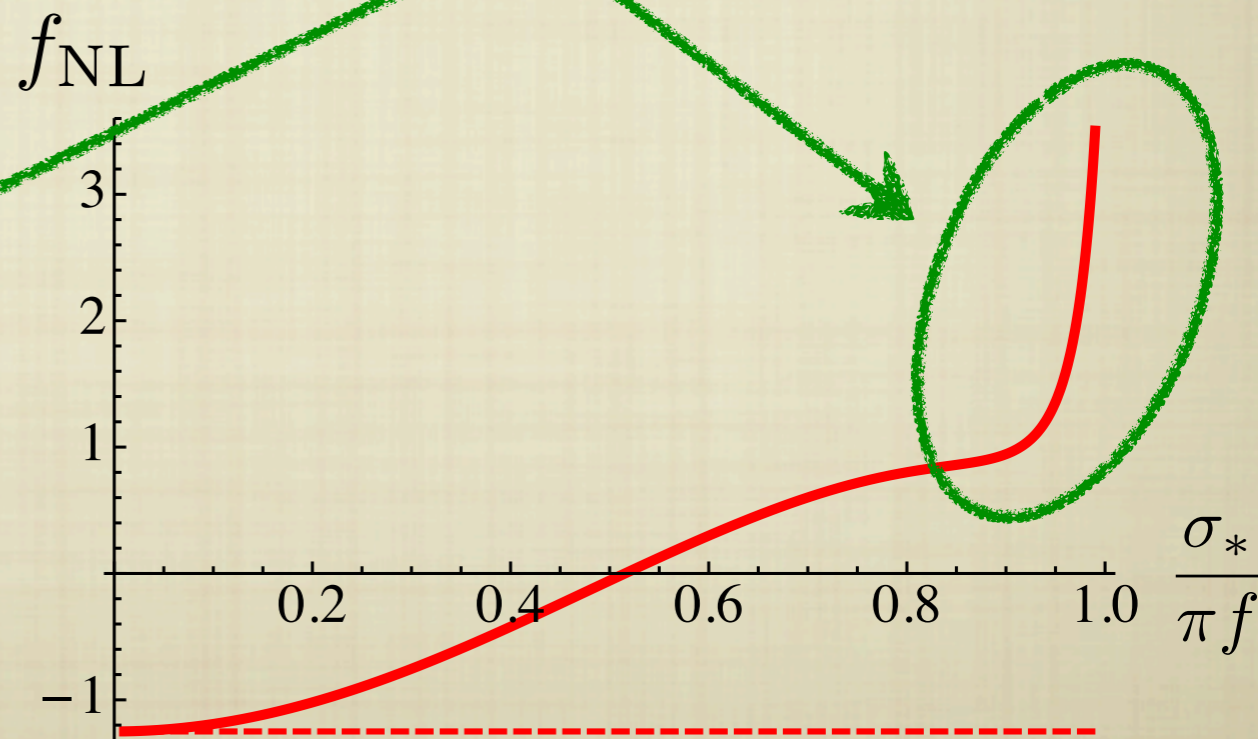
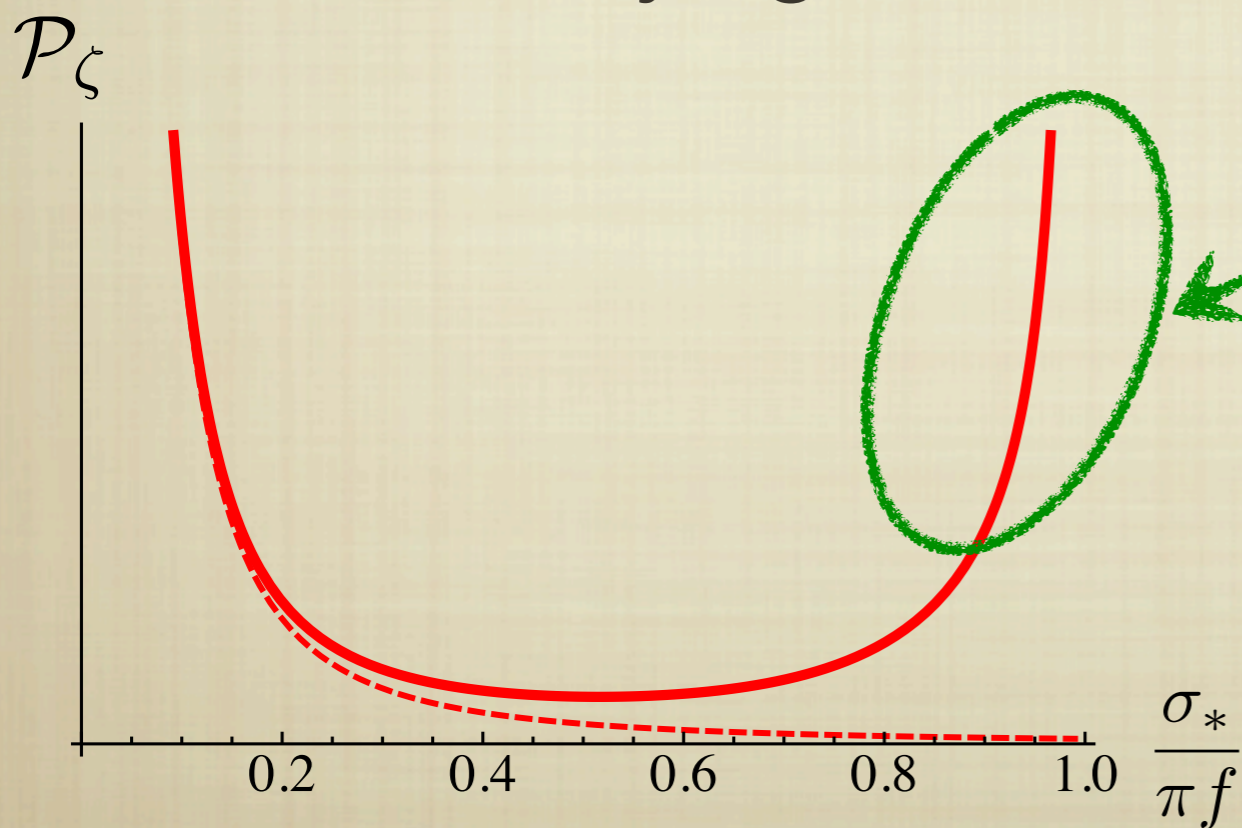


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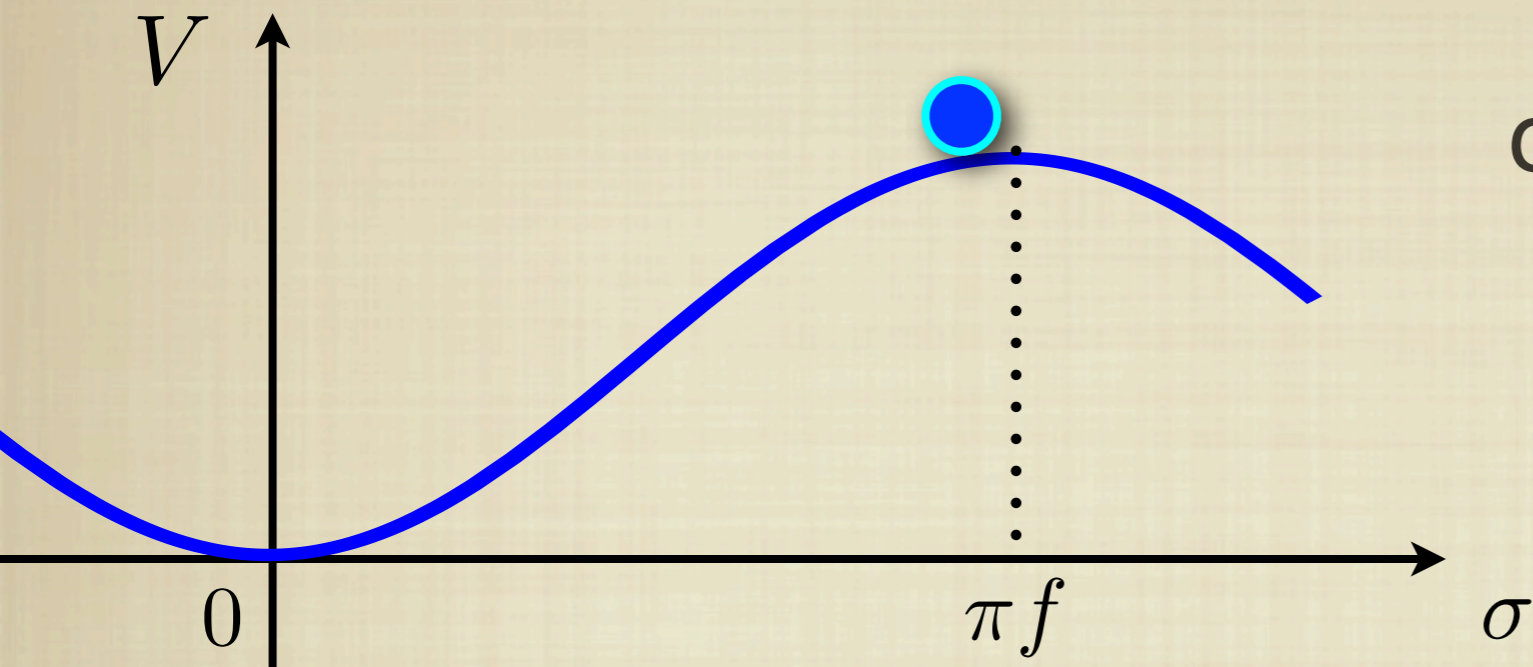
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effects due to non-uniform onset of oscillation

When varying σ_* :



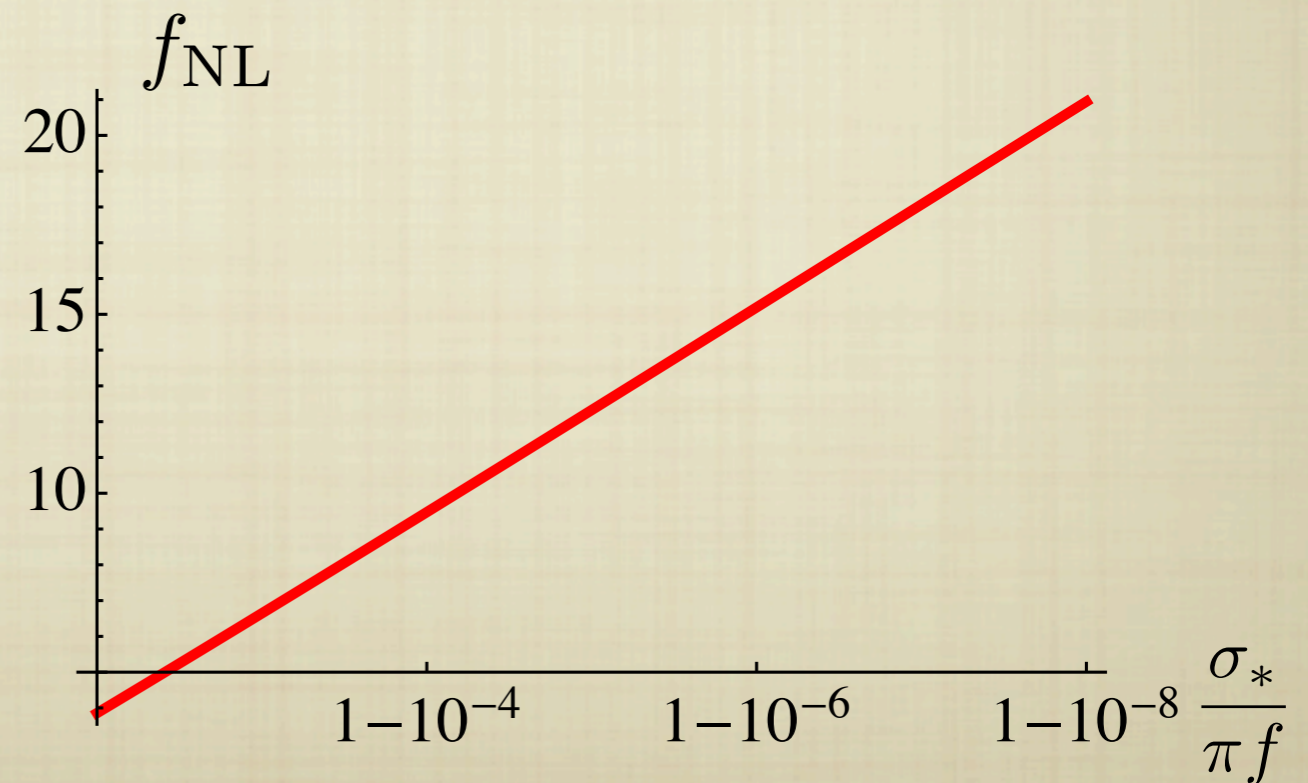
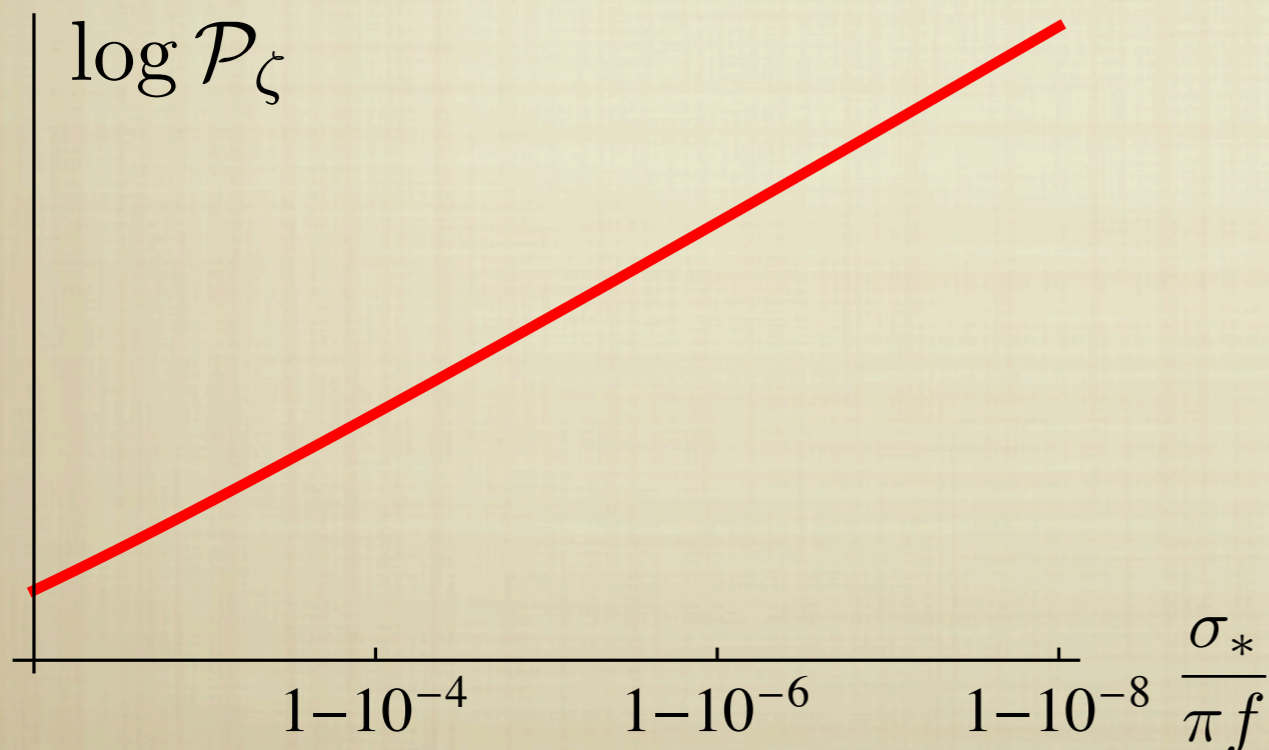
Density Pert. from the Hilltop



curvaton dominant case,

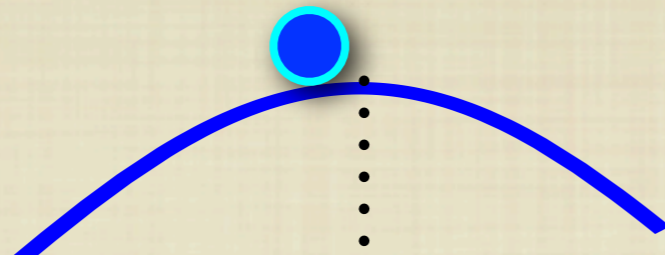
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When varying σ_* :



Density Pert. from the Hilltop

V



curvaton dominant case,

ρ_{σ}

Strong enhancement of linear-order density pert. with mildly increasing f_{NL} of $O(10)$ towards the hilltop.

$\log \mathcal{P}_{\zeta}$

$1-10^{-4}$

$1-10^{-6}$

$1-10^{-8}$

$\frac{\sigma_*}{\pi f}$

f_{NL}

20

15

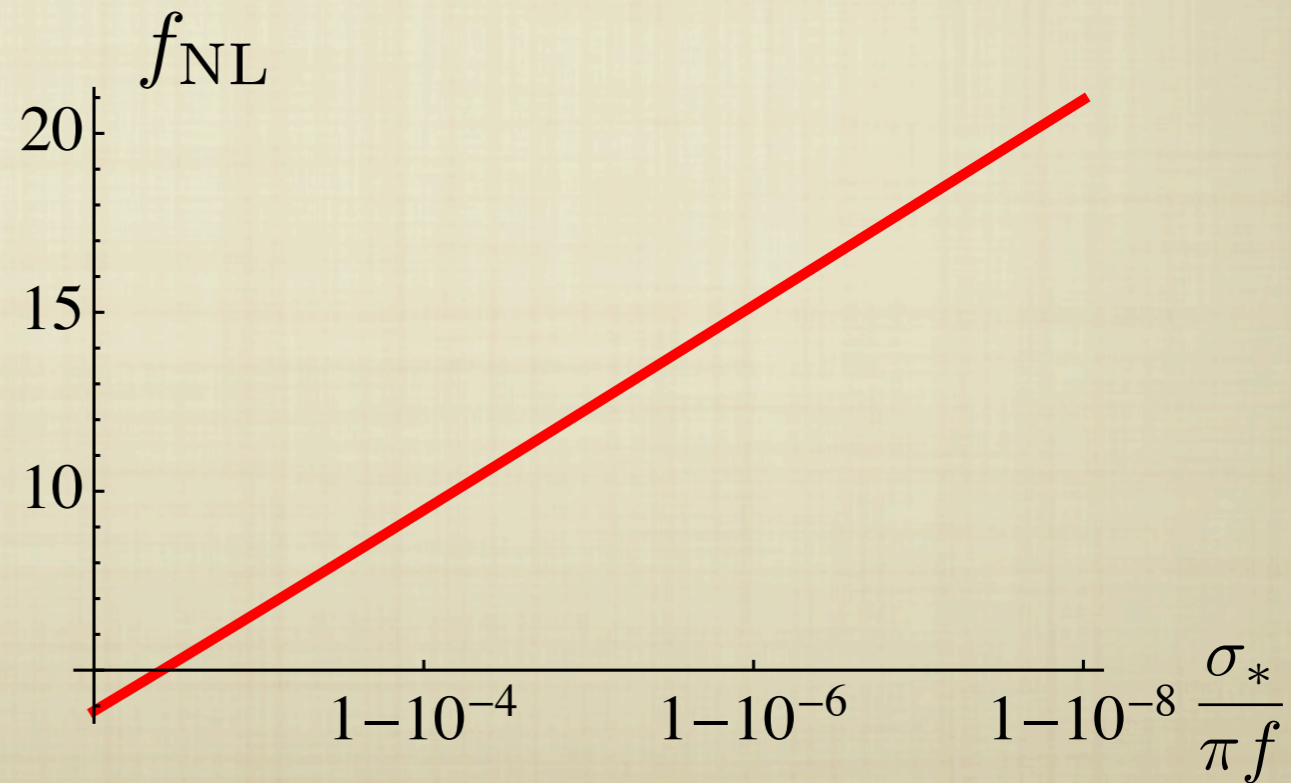
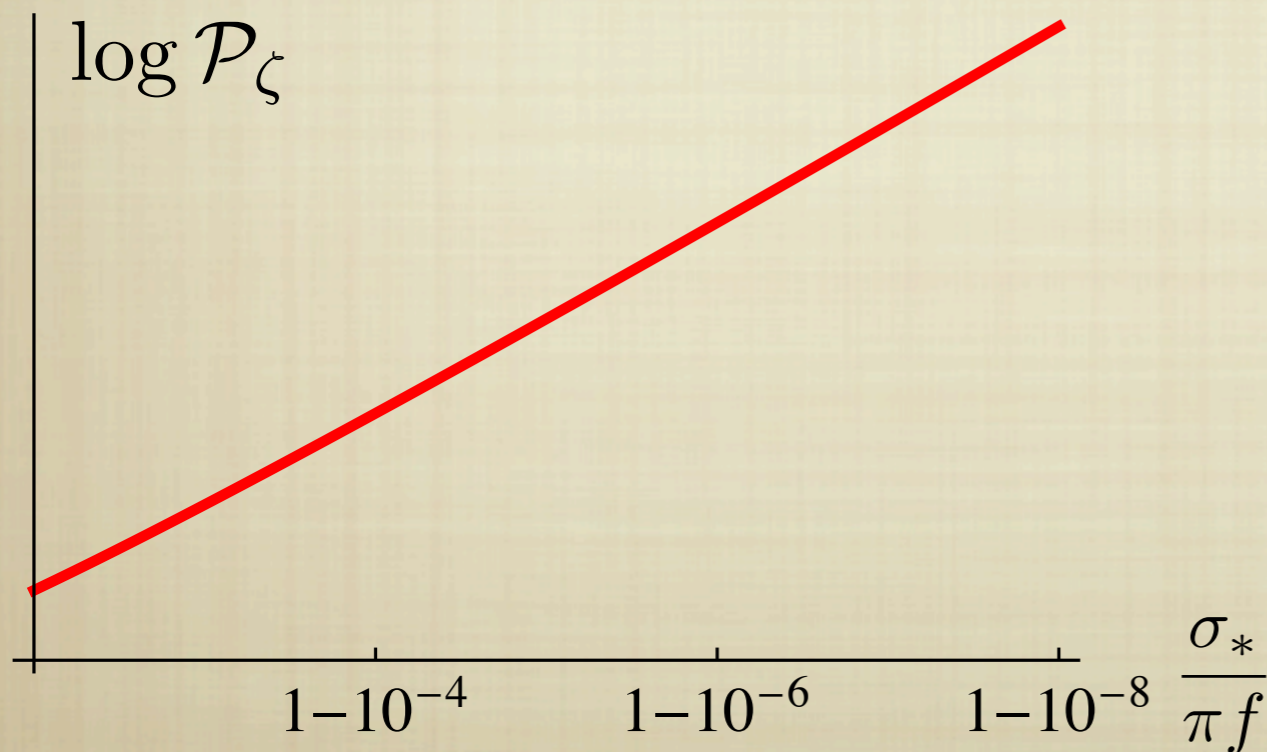
10

$1-10^{-4}$

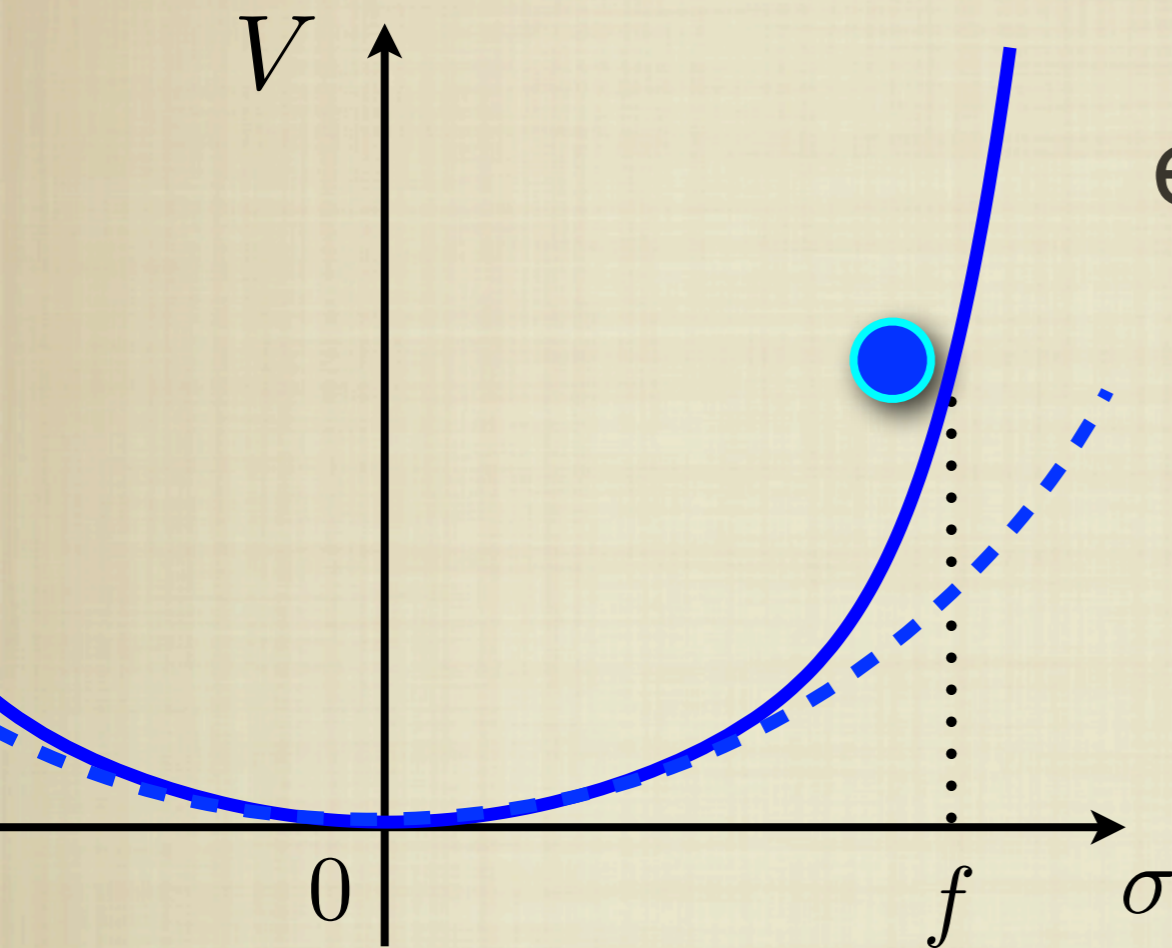
$1-10^{-6}$

$1-10^{-8}$

$\frac{\sigma_*}{\pi f}$



2. Steep Potentials

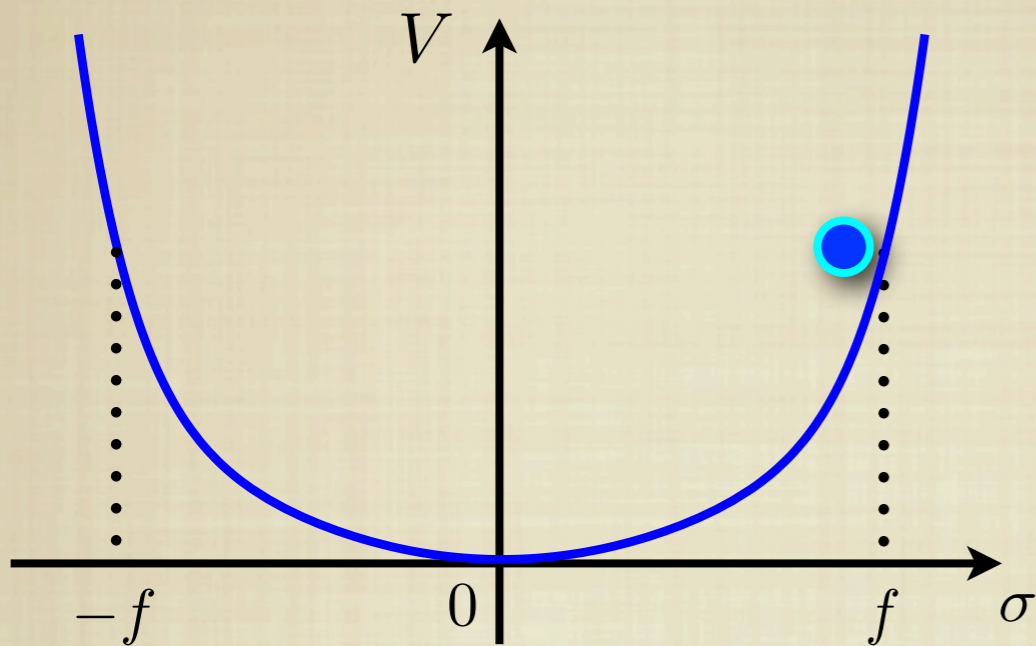


e.g. Self-Interacting Curvatons

$$V(\sigma) = \Lambda^4 \left[\left(\frac{\sigma}{f} \right)^2 + \left(\frac{\sigma}{f} \right)^m \right]$$

Curvaton rolling along the steep potential can lead to **strongly scale-dependent non-Gaussianity.**

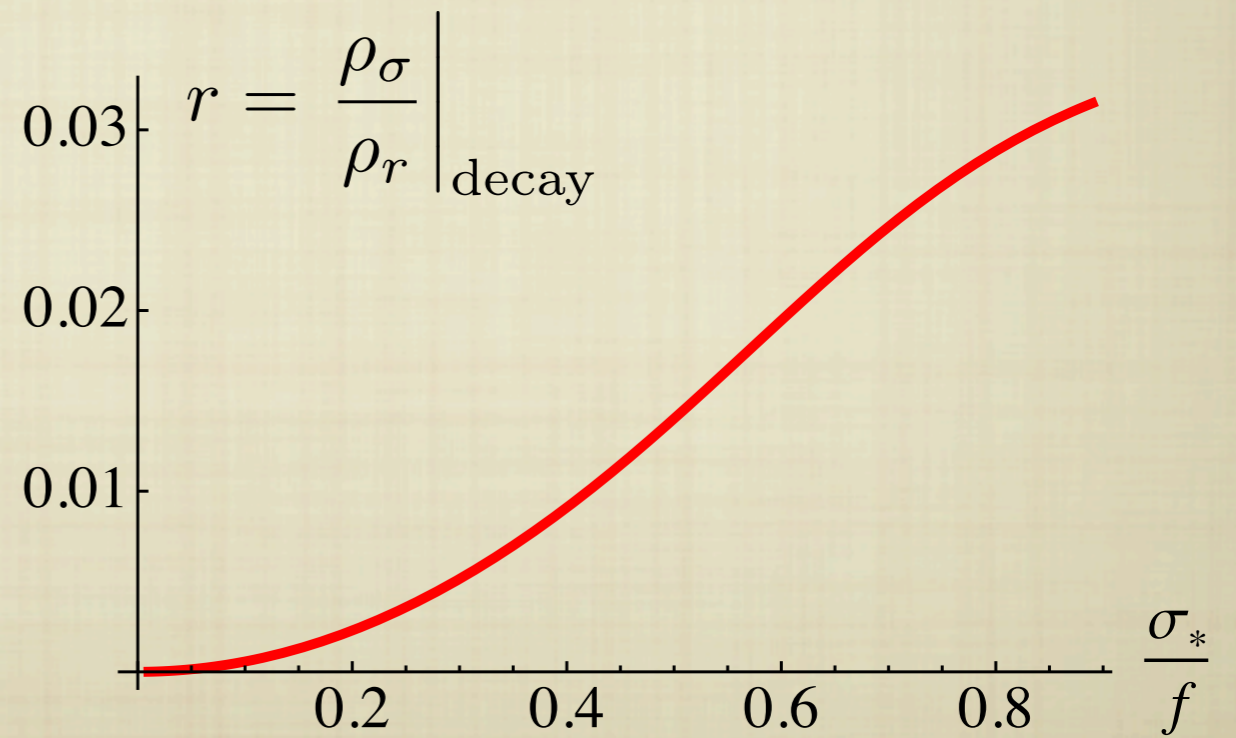
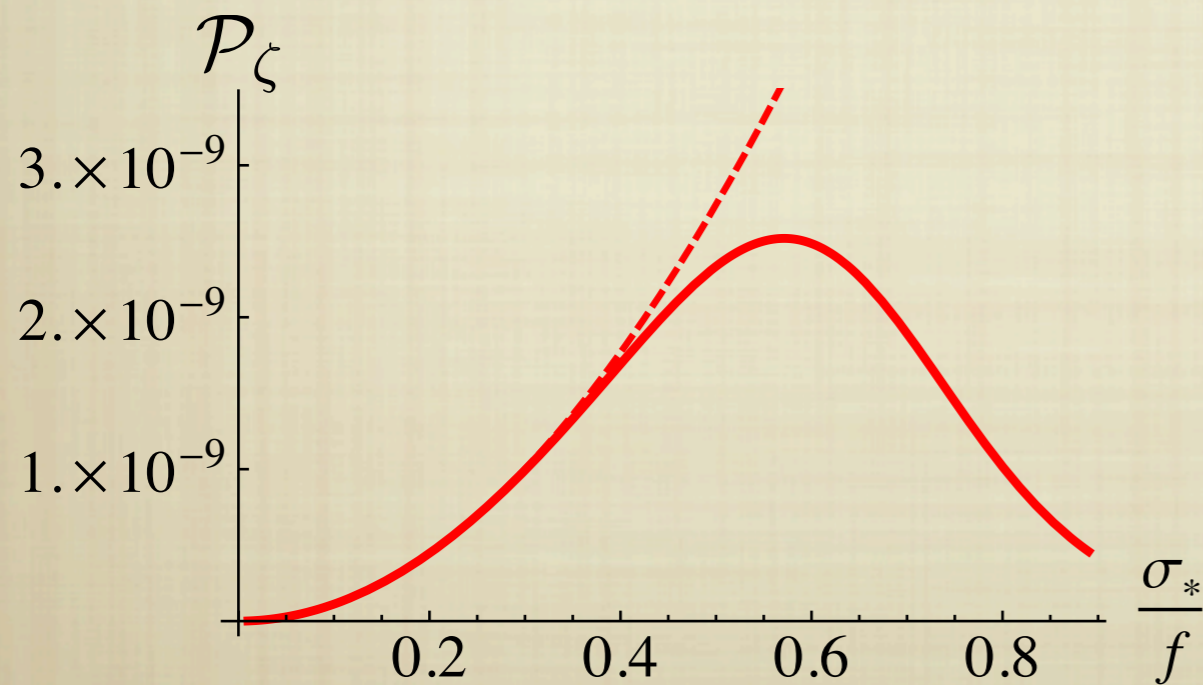
Self-Interacting Curvaton



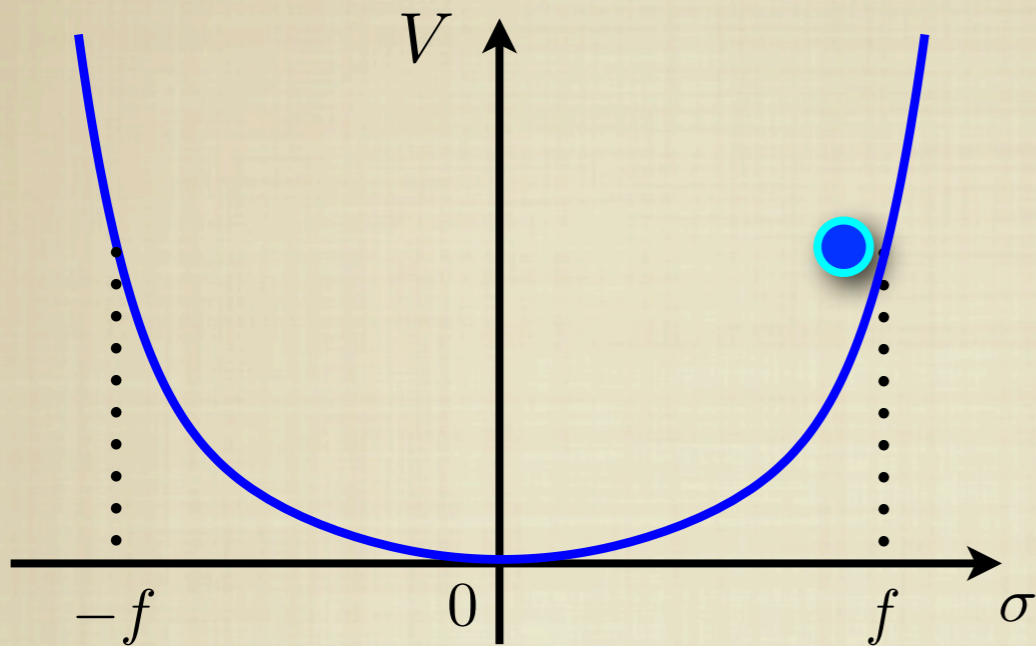
$$V(\sigma) = \Lambda^4 \left[\left(\frac{\sigma}{f} \right)^2 + \left(\frac{\sigma}{f} \right)^8 \right]$$

ex.) $\Lambda \sim 10^{12} \text{ GeV}$ $f \sim 10^{13} \text{ GeV}$
 $H_{\text{inf}} \sim 10^{12} \text{ GeV}$ $T_{\text{reh}} \sim 10^{11} \text{ GeV}$
 $T_{\text{dec}} \sim 100 \text{ GeV}$

When varying σ_* :



Self-Interacting Curvatons

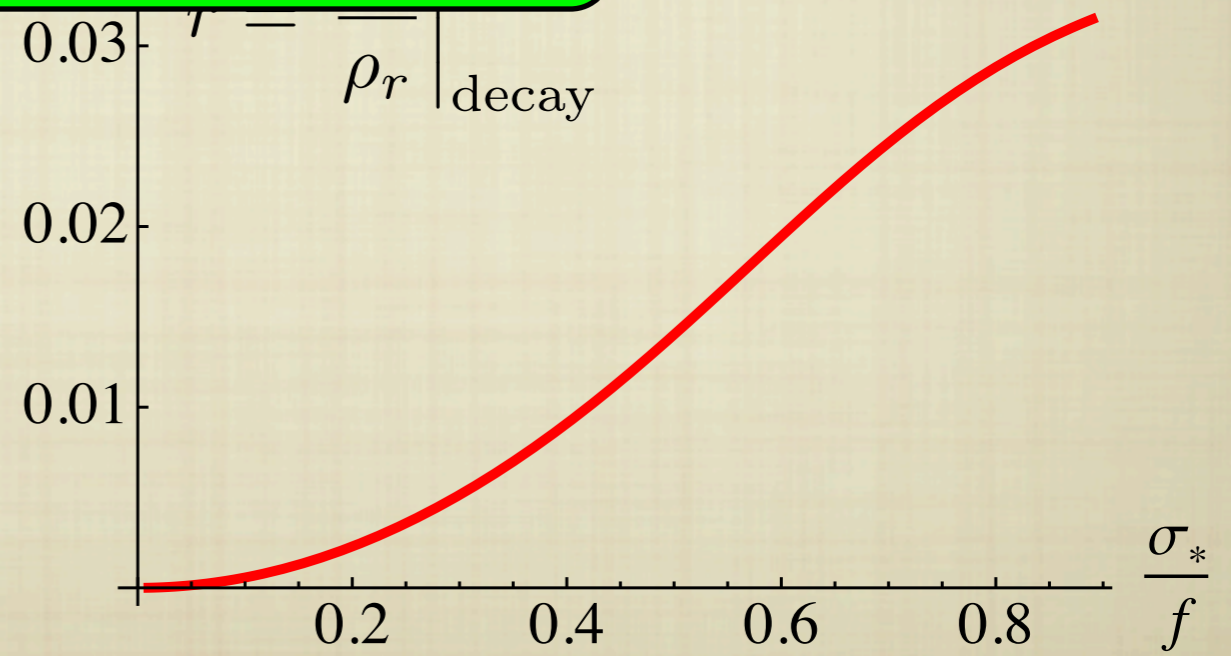
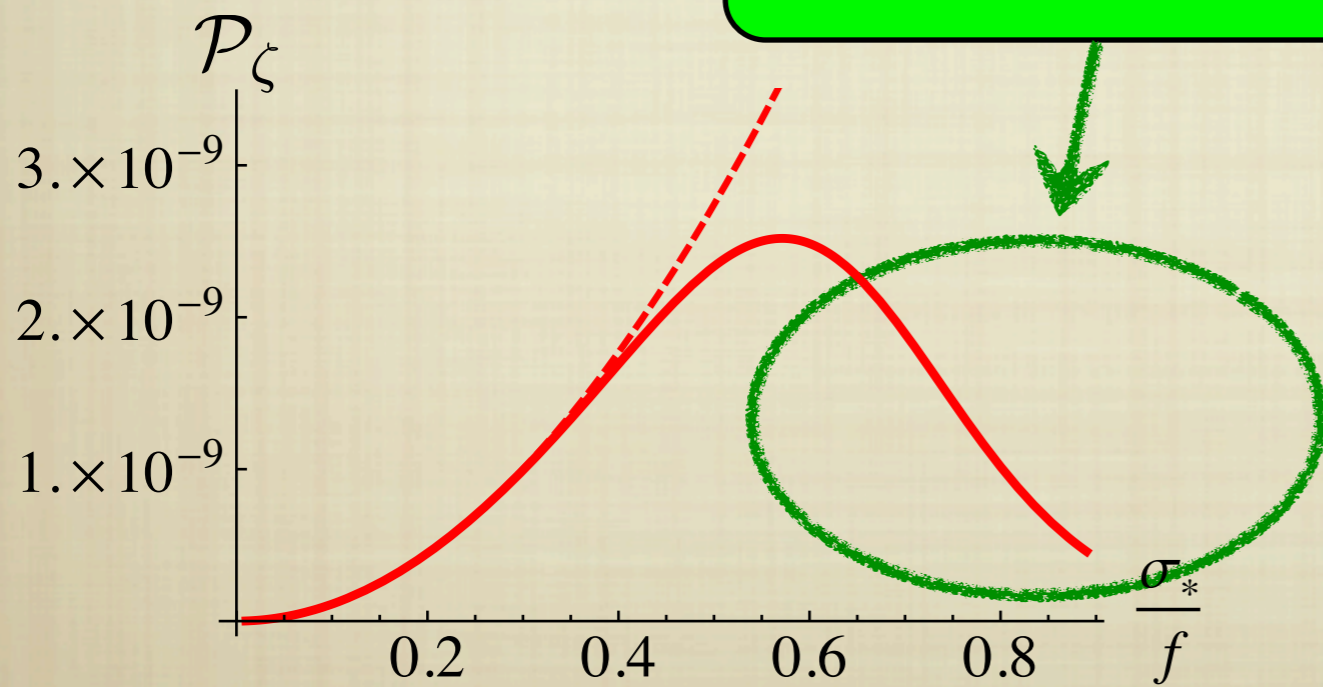


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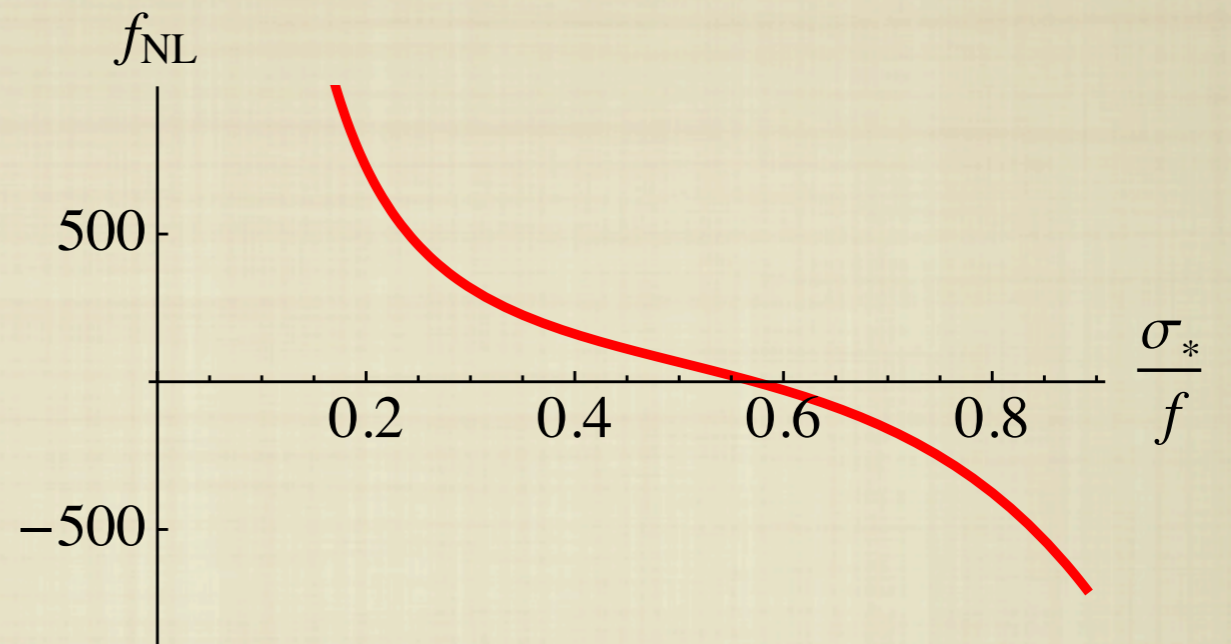
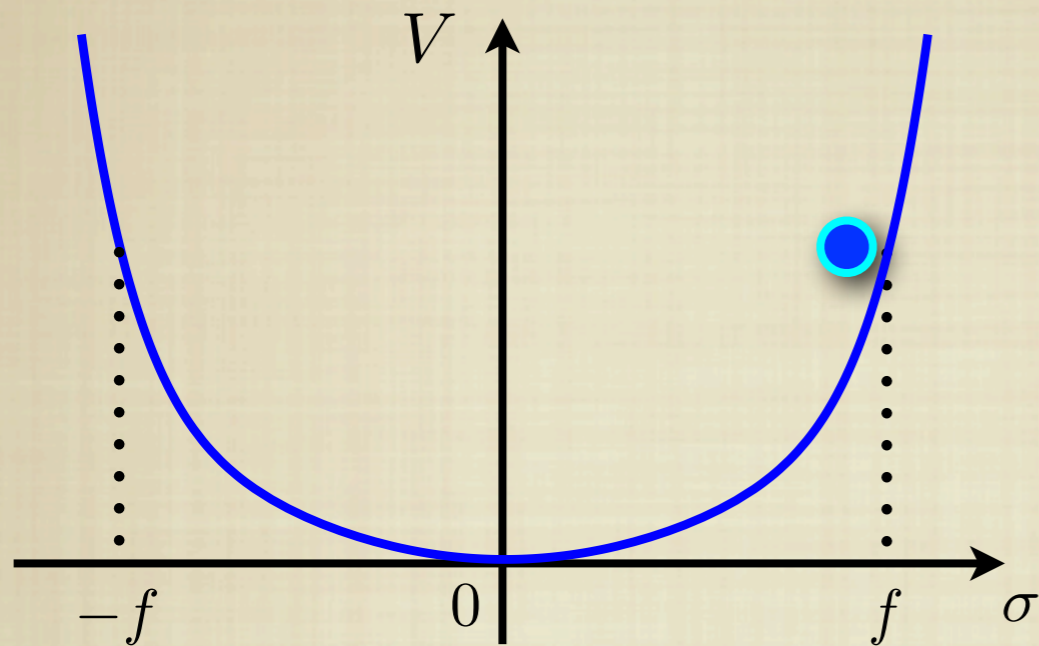
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When varying

suppression due to curvaton rolling

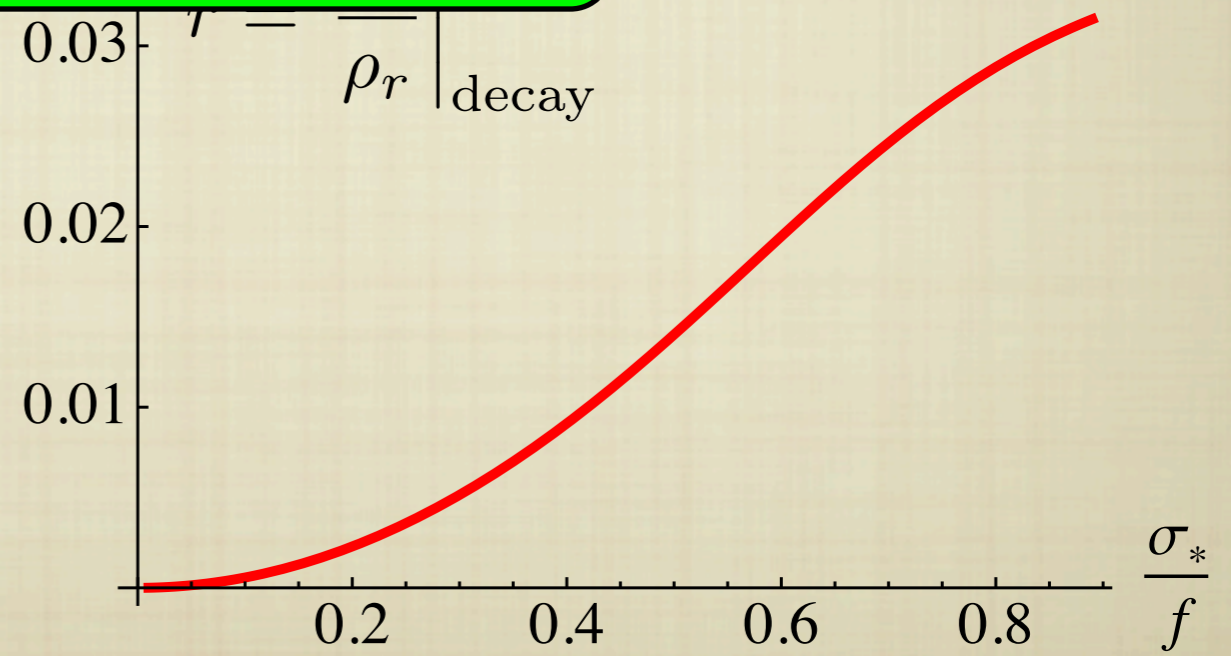
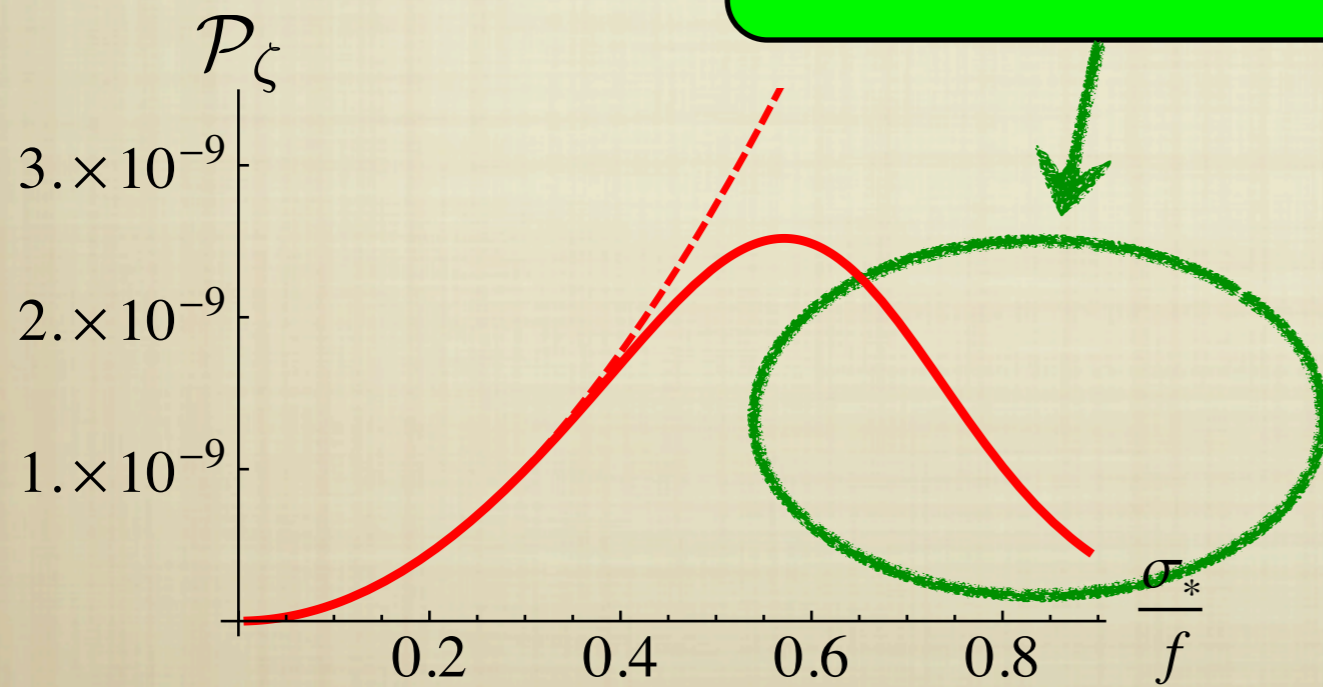


Self-Interacting Curvaton

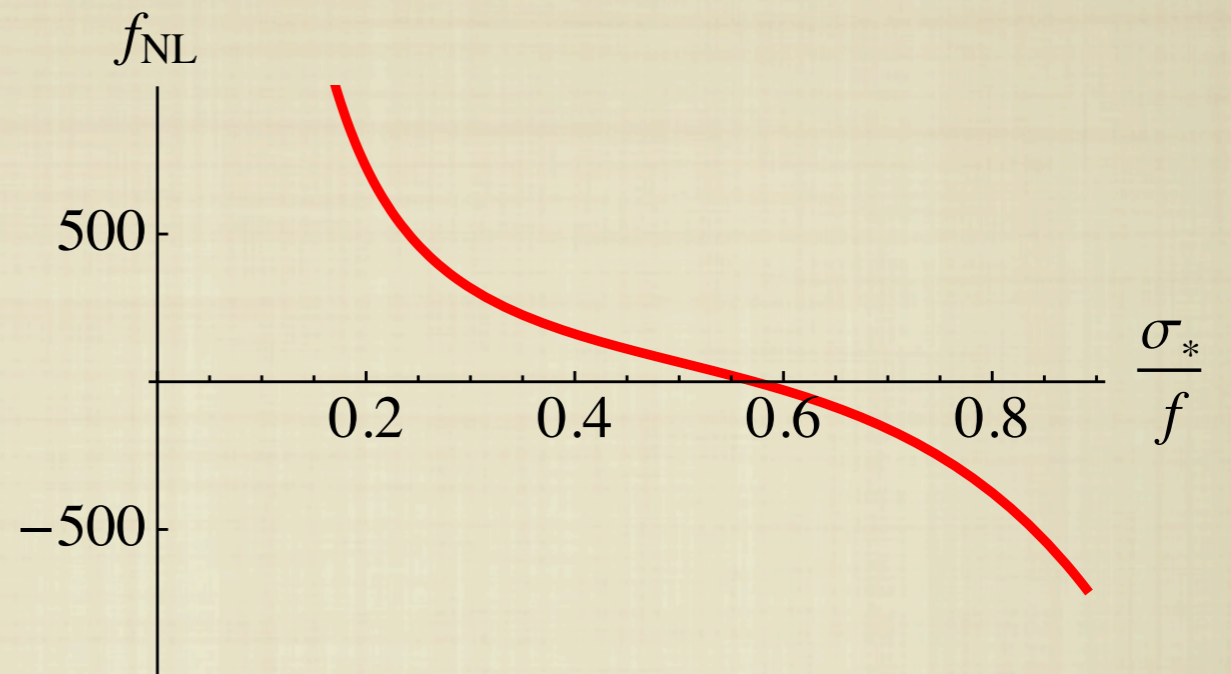
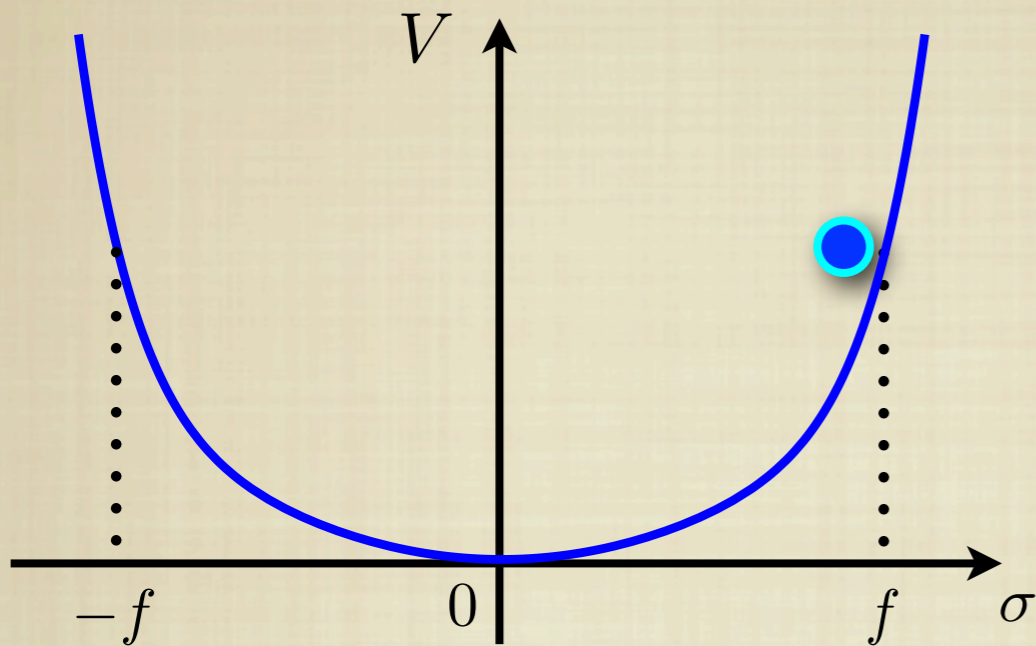


When varying

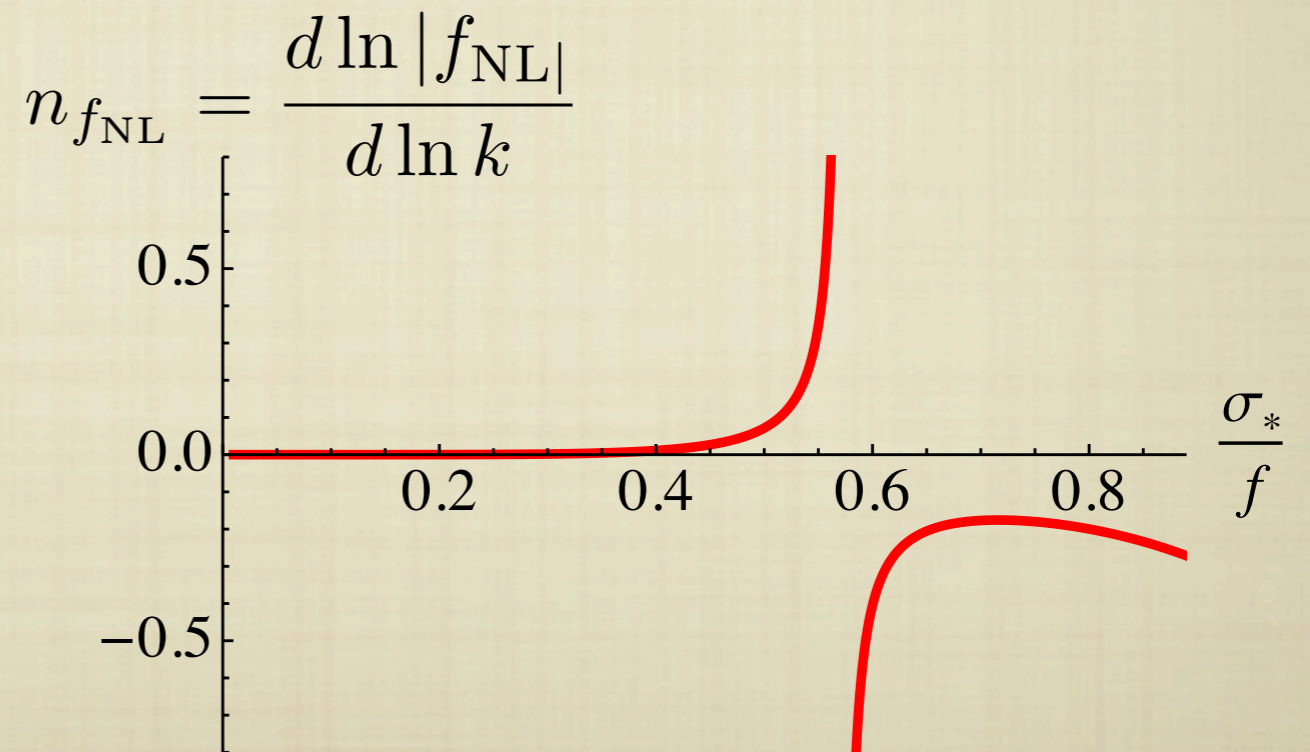
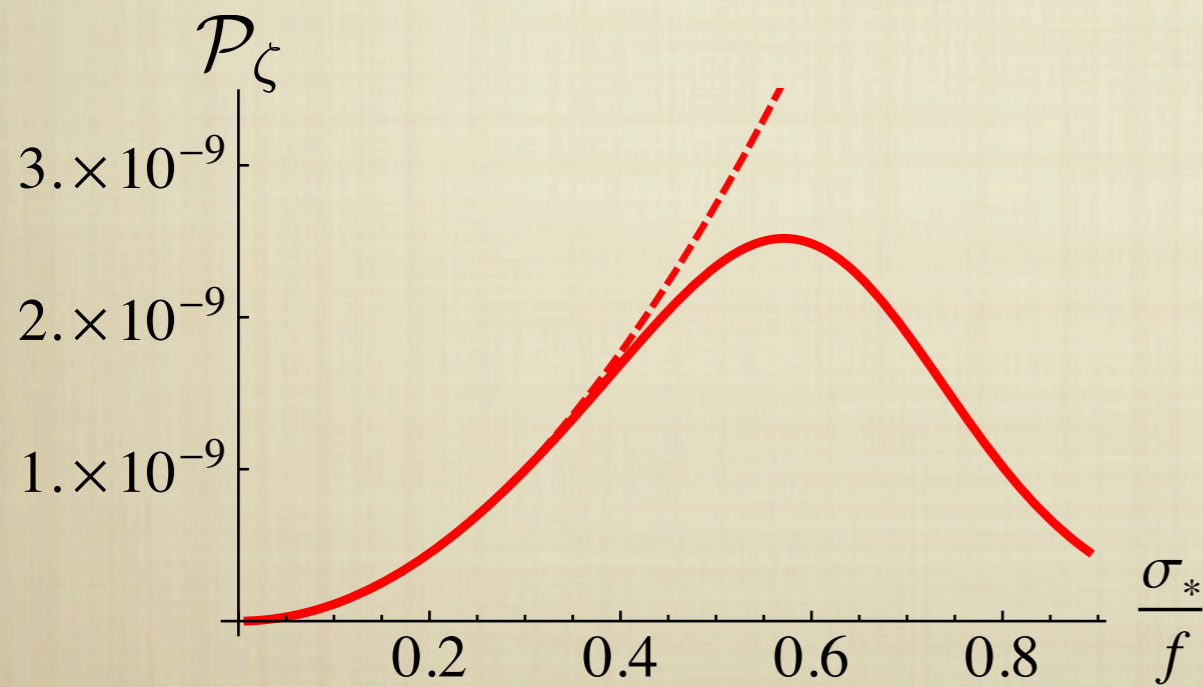
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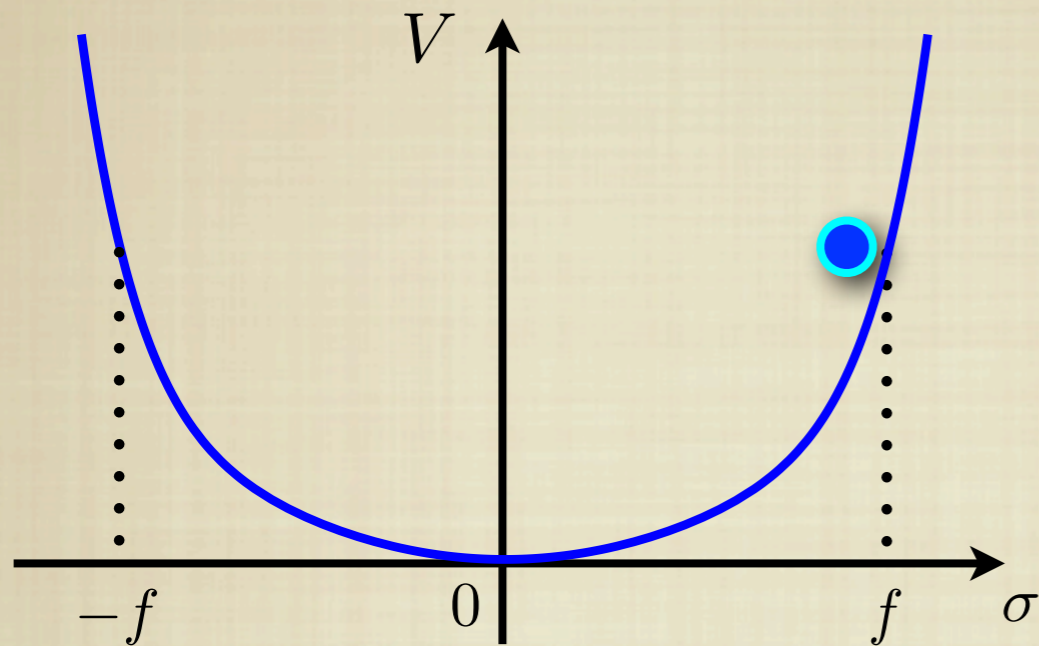
Self-Interacting Curvaton



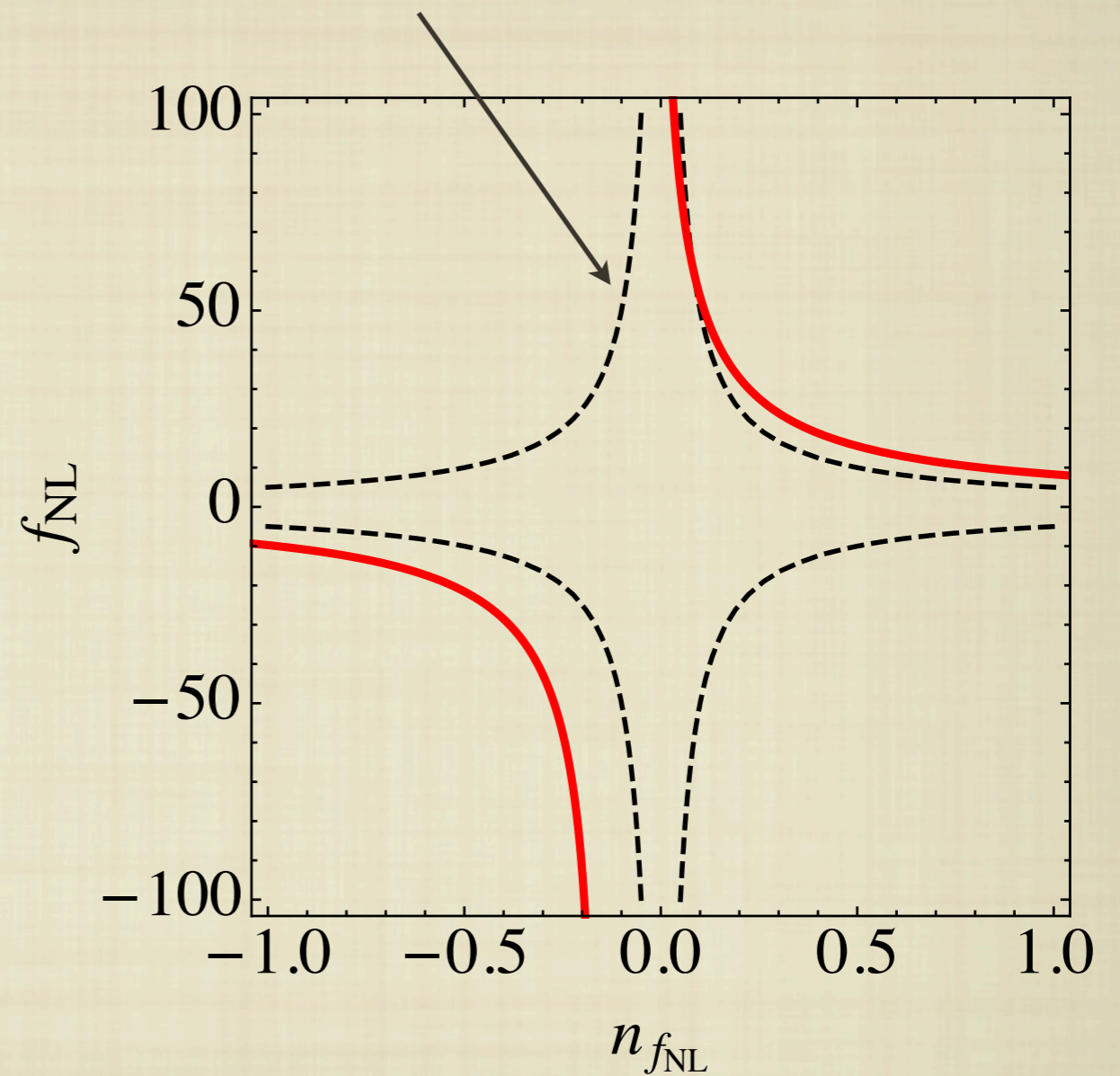
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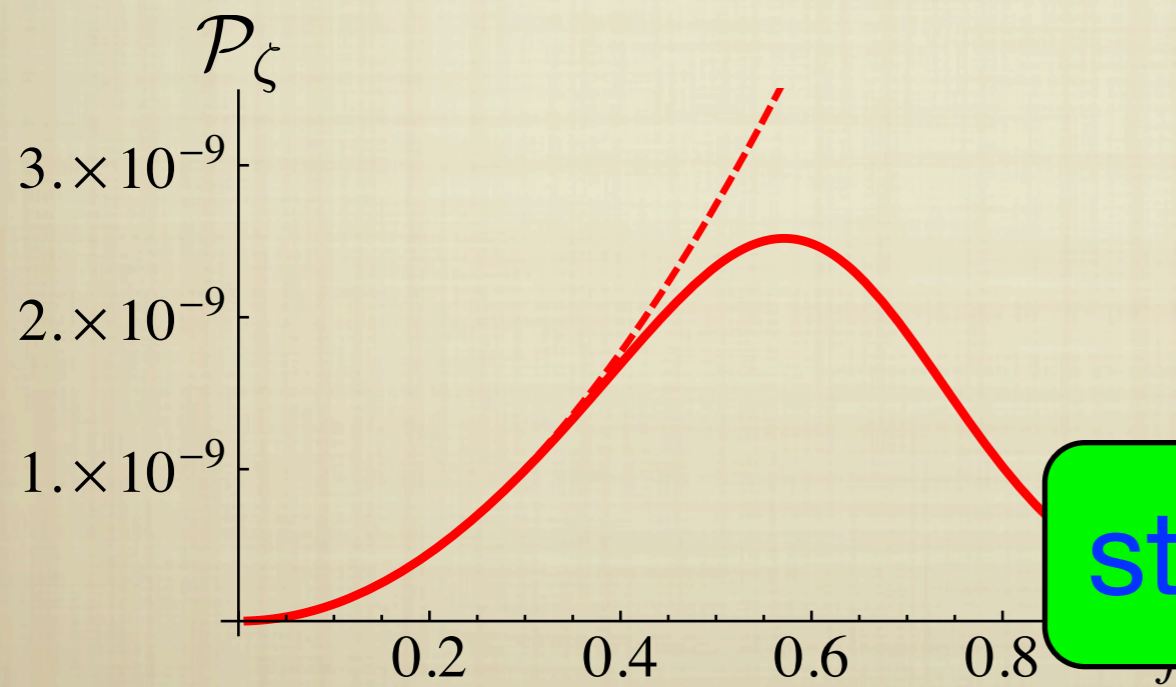
Self-Interacting Curvaton



PLANCK detection limit (Sefusatti et al. '09)



When varying σ_* :



strongly scale-dependent f_{NL}

Summary

- We analytically investigated density perturbations from a curvaton with a generic energy potential.
- Non-quadratic curvatons exhibit new behaviors, such as inhomogeneous onset of oscillations.
- Rich phenomenology : **Flattened** (compared to a quadratic) **potentials can enhance linear & second order perturbations (large f_{NL} even for dominant curvatons!),** **steepened potentials can source running f_{NL} ,** and more.
- Future work : applications to microscopic models.