Closing in on mass-degenerate dark matter scenarios with antiprotons and direct detection

On the complementarity of direct and indirect detection

Stefan Vogl

with Mathias Garny (DESY), Alejandro Ibarra and Miguel Pato (TUM) to appear soon



# Outline

- Introduction
- Particle Physics Framework
- Relic Density
- Indirect Detection
- Direct Detection
- Results
- Conclusion

Lets consider the case when the dark matter particle  $\chi$  and the next to lightest beyond the Standard Model particle  $\eta$  have a similar mass

$$\Delta m = m_{\chi} - m_{\eta} \lesssim m_{\chi}.$$

#### Colliders

- minimal transverse momentum p<sub>T</sub> is required to distinguish jet
- $p_T \approx \Delta m$
- low sensitivity to compressed mass spectra

Lets consider the case when the dark matter particle  $\chi$  and the next to lightest beyond the Standard Model particle  $\eta$  have a similar mass

$$\Delta m = m_{\chi} - m_{\eta} \lesssim m_{\chi}.$$

thermal production

• for  $\frac{m_{\eta}}{m_{\gamma}} \approx 1.2$  coannihilations become important

Lets consider the case when the dark matter particle  $\chi$  and the next to lightest beyond the Standard Model particle  $\eta$  have a similar mass

$$\Delta m = m_{\chi} - m_{\eta} \lesssim m_{\chi}.$$

#### **Indirect Detection**

- compressed mass spectra exhibit very characteristic features
- annihilation rates are enhanced for small  $\Delta m$
- huge astrophysical uncertainties

Lets consider the case when the dark matter particle  $\chi$  and the next to lightest beyond the Standard Model particle  $\eta$  have a similar mass

$$\Delta m = m_{\chi} - m_{\eta} \lesssim m_{\chi}.$$

#### **Direct Detection**

- scattering rates are enhanced for small  $\Delta m$
- less astrophysical uncertainties than in Indirect Detection
- good experimental limits

Lets consider the case when the dark matter particle  $\chi$  and the next to lightest beyond the Standard Model particle  $\eta$  have a similar mass

$$\Delta m = m_{\chi} - m_{\eta} \lesssim m_{\chi}.$$

#### **Direct Detection**

- scattering rates are enhanced for small  $\Delta m$
- less astrophysical uncertainties than in Indirect Detection
- good experimental limits

But: We need to specify the model in order to compare observables.

# Particle Physics Framework

Begin with the SM and add news physics

#### Particles

- Majorana fermion  $\chi$  as dark matter
- a scalar  $\eta$  as the next to lightest beyond the Standard Model particle

### Assign charges

- $\chi$  is a singlet under  $SU(3) \times SU(2) \times U(1)$
- $\eta$  is a triplet under SU(3) and (for simplicity) a singlet under SU(2)
- u,d,s or b flavor quantum number for  $\eta$

#### Interactions

• a Yukawa interaction with the quarks:  $\mathcal{L}_{int} = f \bar{\chi} q_R \eta$ 

(I) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1))

# Particle Physics Framework

Begin with the SM and add news physics

#### **Particles**

- Majorana fermion  $\chi$  as dark matter
- a scalar  $\eta$  as the next to lightest beyond the Standard Model particle

### Assign charges

- $\chi$  is a singlet under  $SU(3) \times SU(2) \times U(1)$
- $\eta$  is a triplet under SU(3) and (for simplicity) a singlet under SU(2)
- u,d,s or b flavor quantum number for η

#### Interactions

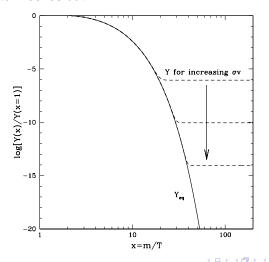
• a Yukawa interaction with the quarks:  $\mathcal{L}_{int} = f \bar{\chi} q_R \eta$ 

Notice: similar to SUSY with light squarks

A D F A B F A B F A B

### thermal freeze out

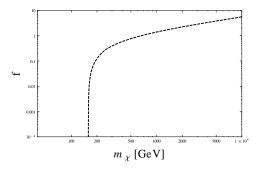
- all particles are in thermal equilibrium in the early Universe
- when temperature  $T \ll m_{\chi}$  dark matter can't be produced anymore  $\rightarrow$  dark matter freezes out



S. Vogl (TU München)

# Coannihilations

- for  $\frac{\Delta m}{m_{\chi}} \lesssim 1.2$  more particles need to be included in the Boltzmann equation
- we use MicrOMEGAS for the calculation of the relic density
- specifying  $m_{\chi}$  and  $\Delta m$  yields constraint on f
- Example: Coupling to u and  $m_{\chi}/m_{\eta} = 1.1$



• for  $m_{\chi}$  smaller that a certain scale the relic density can not be obtained

S. Vogl	(TU München)	

# Majorana fermions annihilating into light quarks

• thermally averaged cross section  $\langle \sigma_{ann} v \rangle$  can be expanded as  $\langle \sigma_{ann} v \rangle = a + bv^2 + O(v^4)$ 

consider annihilations into quarks

s-wave annihilation is suppressed by chirality

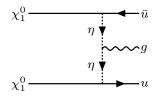
$$\langle \sigma_{ann} m{v} 
angle pprox m{a} pprox rac{m_{f}^{2}}{m_{{\scriptscriptstyle DM}}^{2}}$$

p-wave suppressed by velocity

$$\langle \sigma_{ann} v 
angle pprox v^2 pprox 10^{-6}$$

# Lifting the chirality suppression

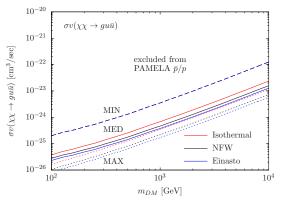
 the suppression can be lifted by the emission of a boson, i.e. γ, W<sup>±</sup>, Z or a gluon



the fragmentation of the gluon increases the production of antiprotons

## Constraints from Antiprotons

- the  $\bar{p}/p$  ratio measured by Pamela constrains  $\sigma v$
- main uncertainty: halo model and cosmic ray propagation
- Example:  $m_{\eta}/m_{\chi} = 1.1$



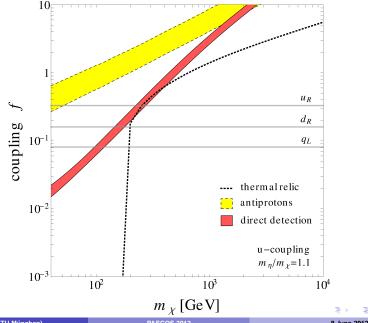
# Dark Matter Nucleon Scattering

 dark matter nucleon scattering is induced microscopical by scattering of quarks and gluons in the nucleus



- interactions can be described in terms of effective Lagrangian
- suppression scale  $\Lambda = m_\eta^2 (m_\chi + m_q)^2$
- compressed spectrum → small Λ
- recoil rate is enhanced
- uncertainties: astrophysics (mainly neglected here) and composition of the nucleon

## Putting everything together

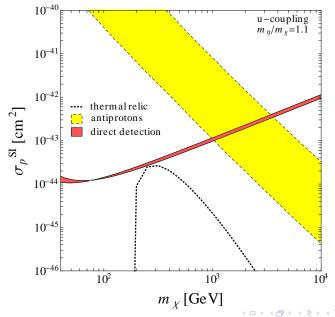


S. Vogl (TU München)

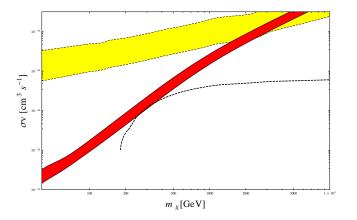
PASCOS 2012

8 June 2012 11/15

### The Direct Detection Plane

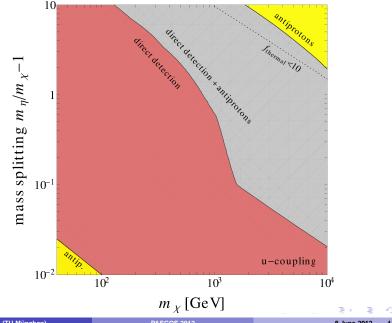


# The Indirect Detection Plane



A B A B A B A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A

# Which constraint is strongest?



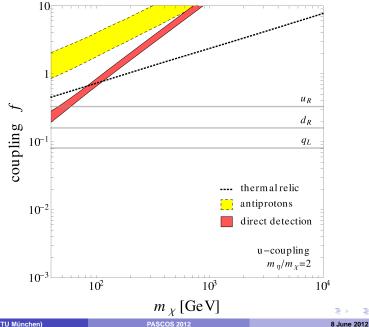
S. Vogl (TU München)

PASCOS 2012

## Conclusions

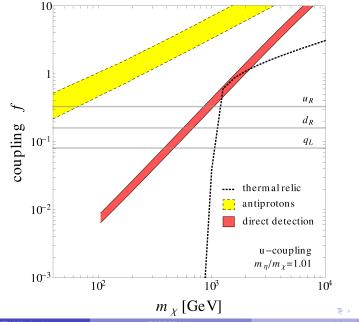
- compressed mass spectra lead to enhanced signals for dark matter detection experiments
- probes region of parameter space inaccessible at colliders
- direct detection experiments are cutting into the parameter space allowed by thermal production

# **Backup**



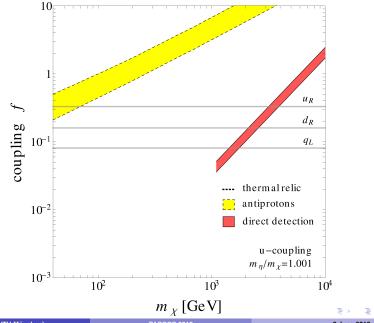


# Backup





# Backup



S. Vogl (TU München)

PASCOS 2012

8 June 2012 15 / 15