

The Energy Frontier in Nature: Highest Energy Cosmic Radiation

Günter Sigl

1. Introduction and Overview
2. Astrophysics
3. Particle Physics at High Energies

DFG Deutsche
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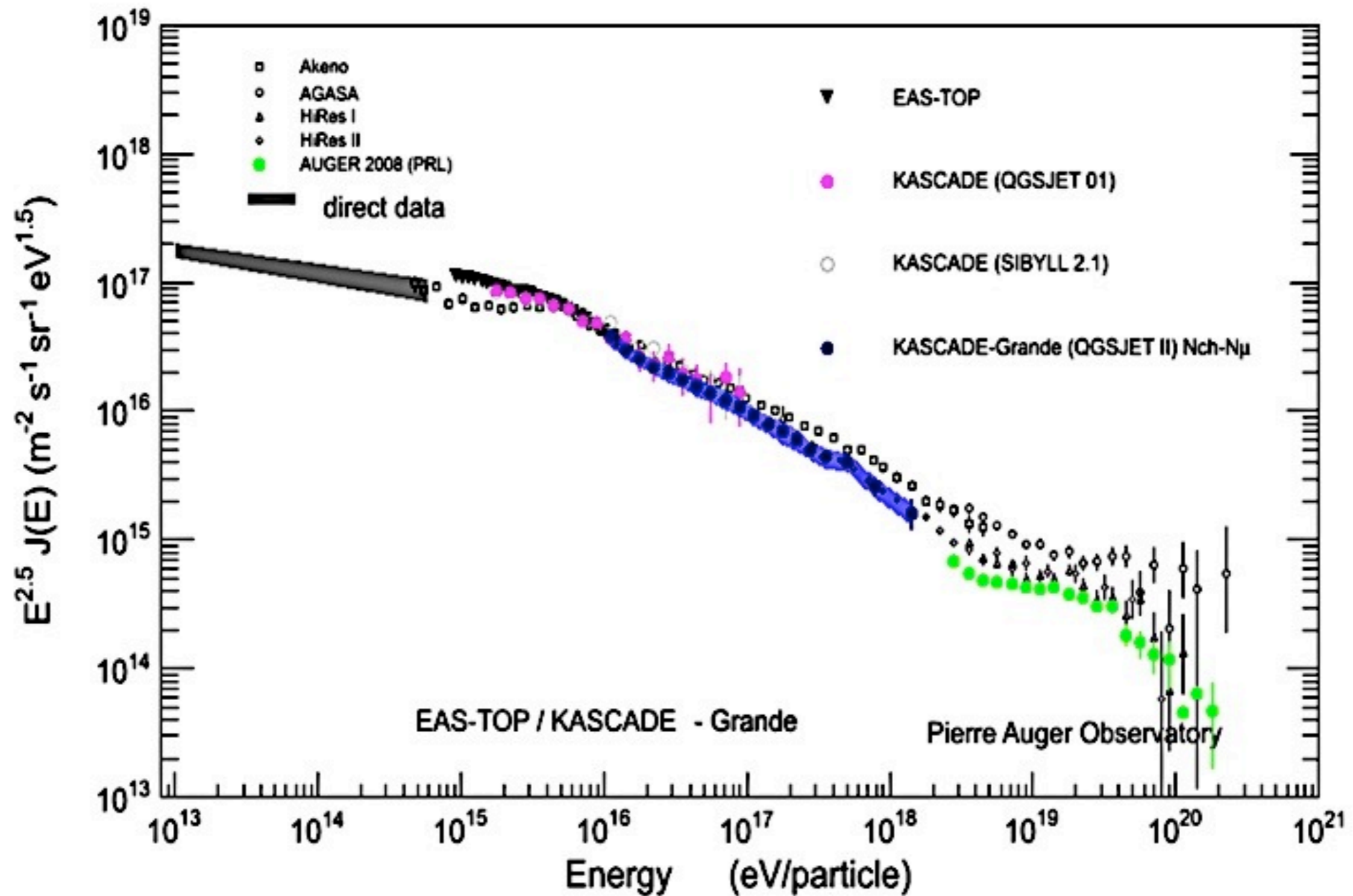


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II. Institut theoretische Physik, Universität Hamburg

<http://www2.iap.fr/users/sigl/homepage.html>

The All Particle Cosmic Ray Spectrum

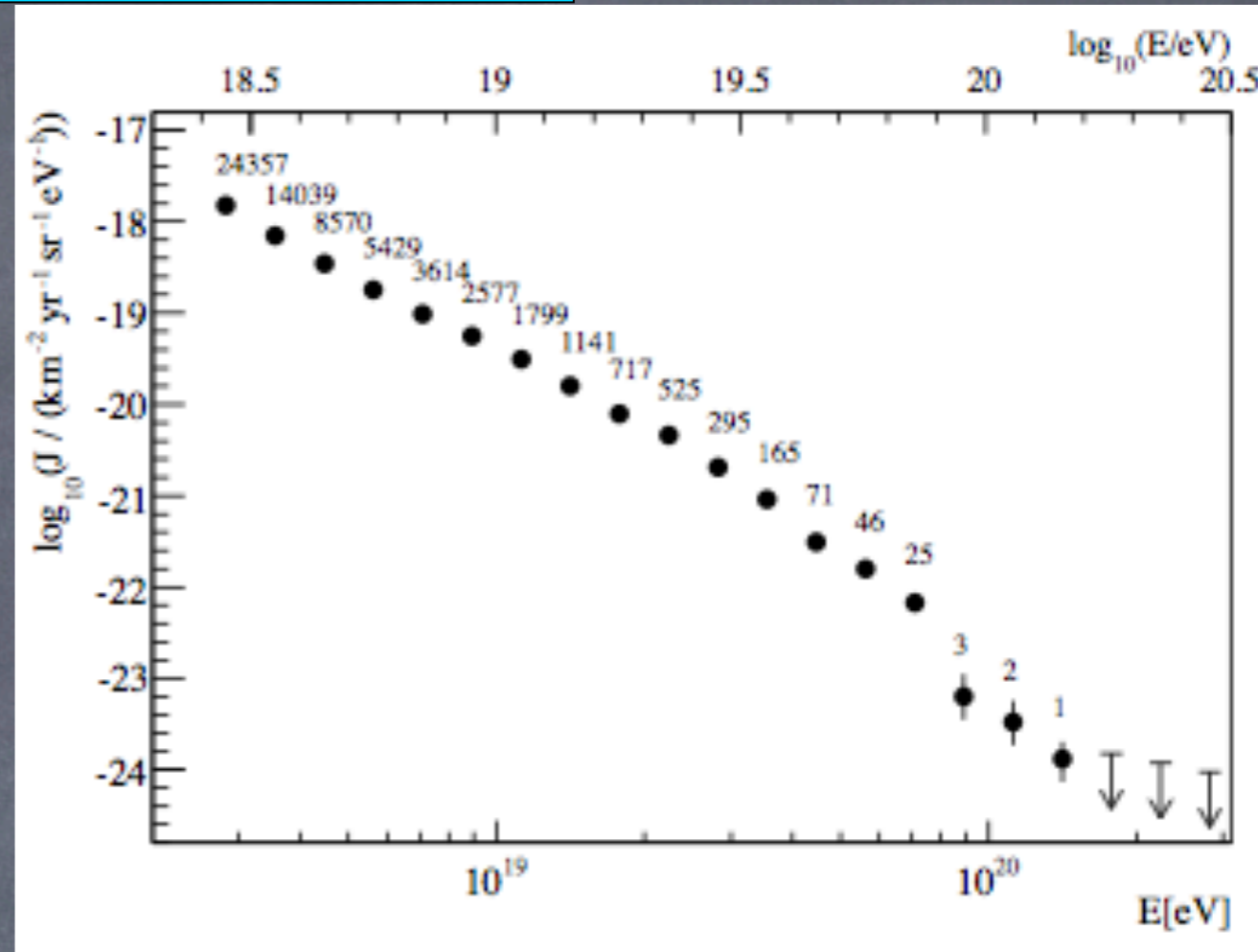


KASCADE-Grande collaboration, arXiv:1009.4716

Auger and HiRes Spectra

Auger exposure = 20905 km² sr yr
up to December 2010

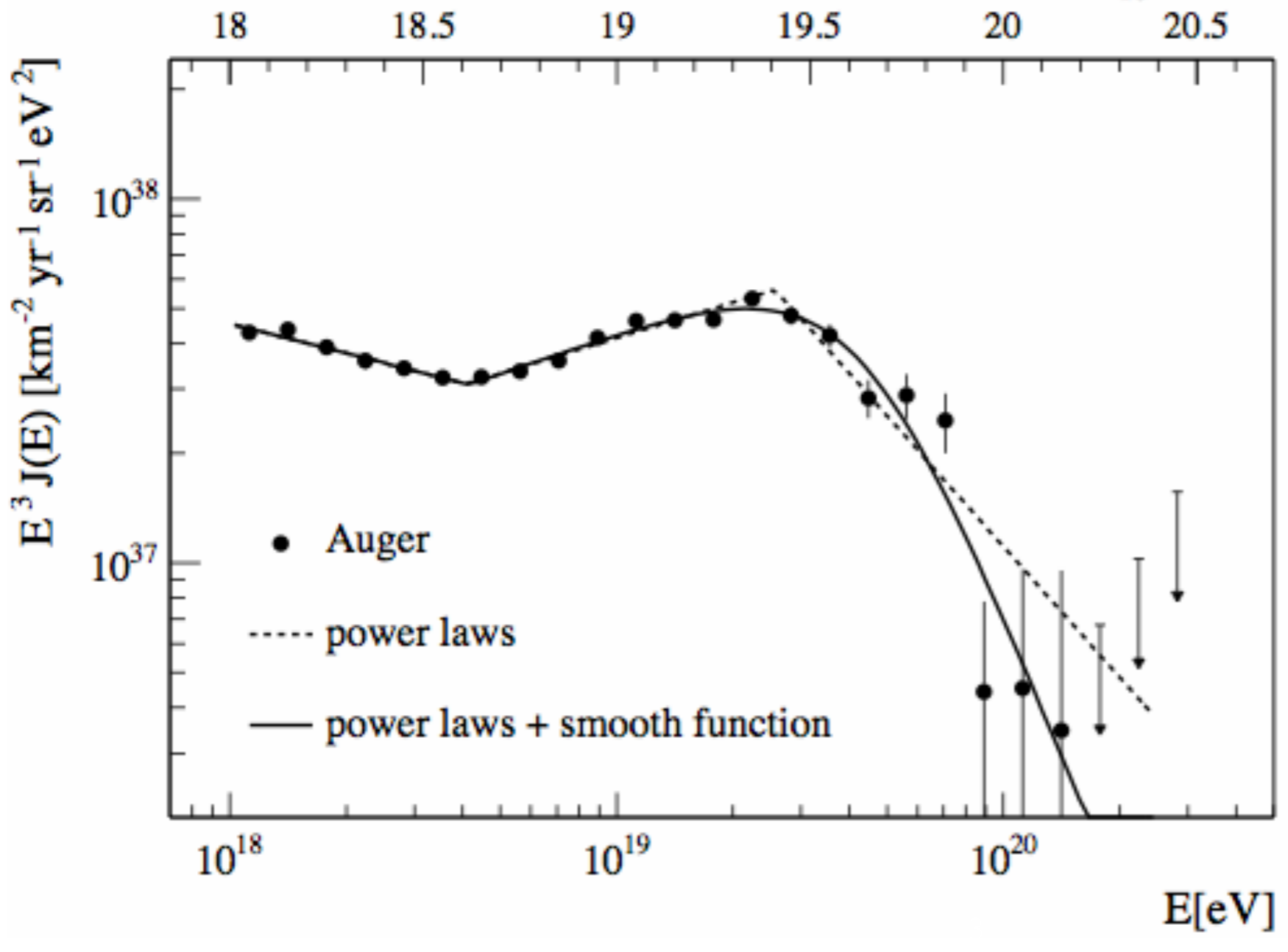
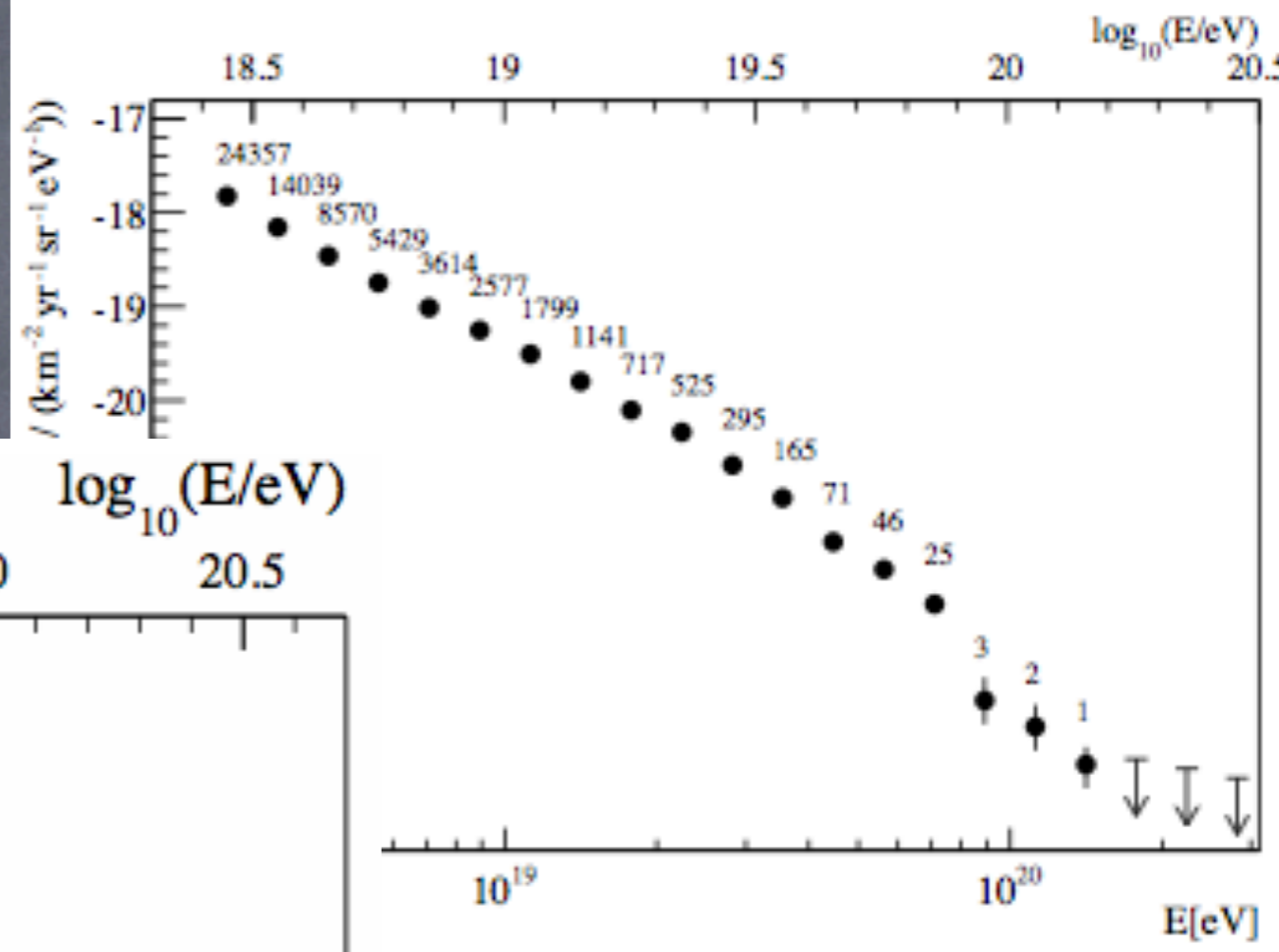
Pierre Auger Collaboration, PRL 101, 061101 (2008)
and Phys.Lett.B 685 (2010) 239
and ICRC 2011, arXiv:1107.4809



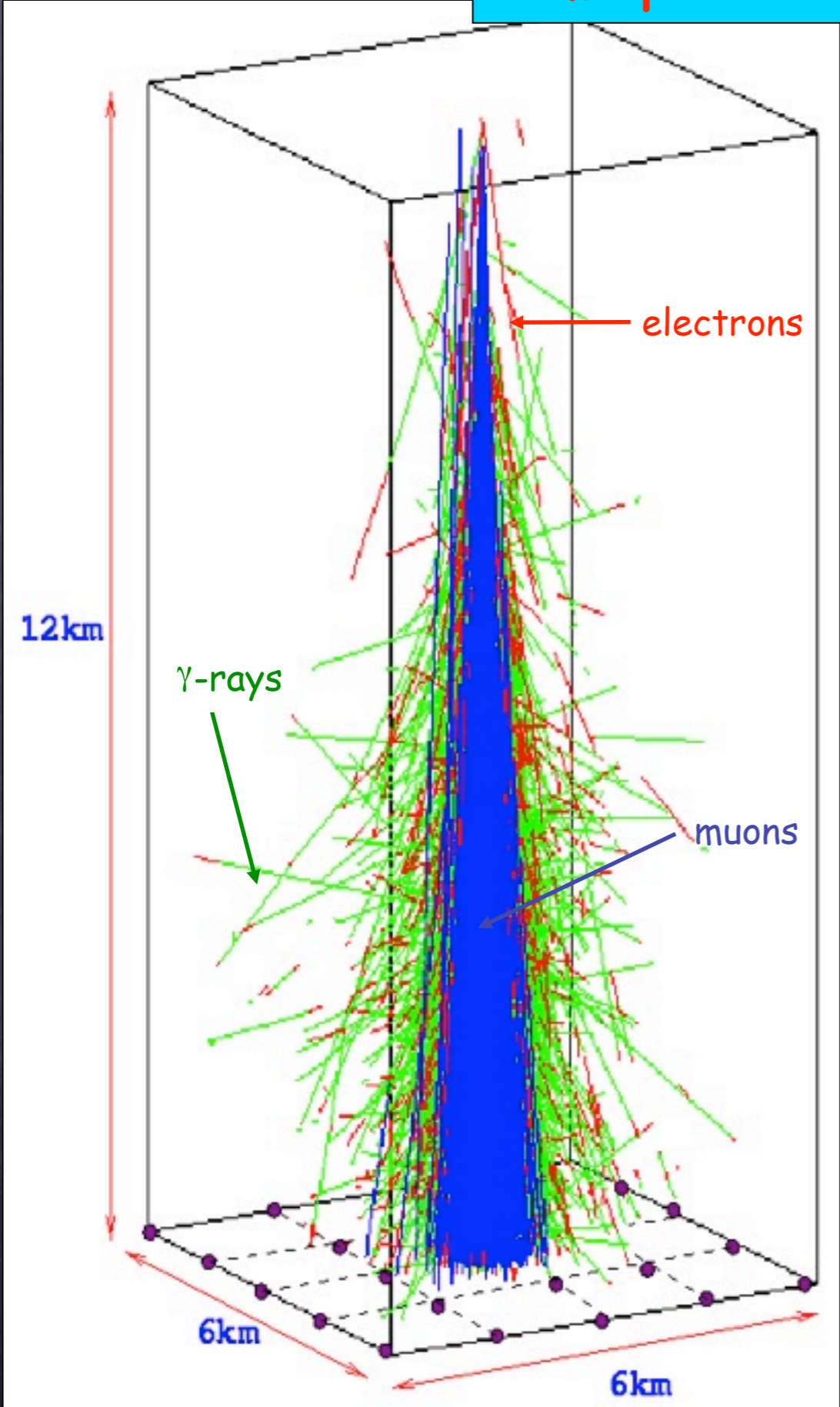
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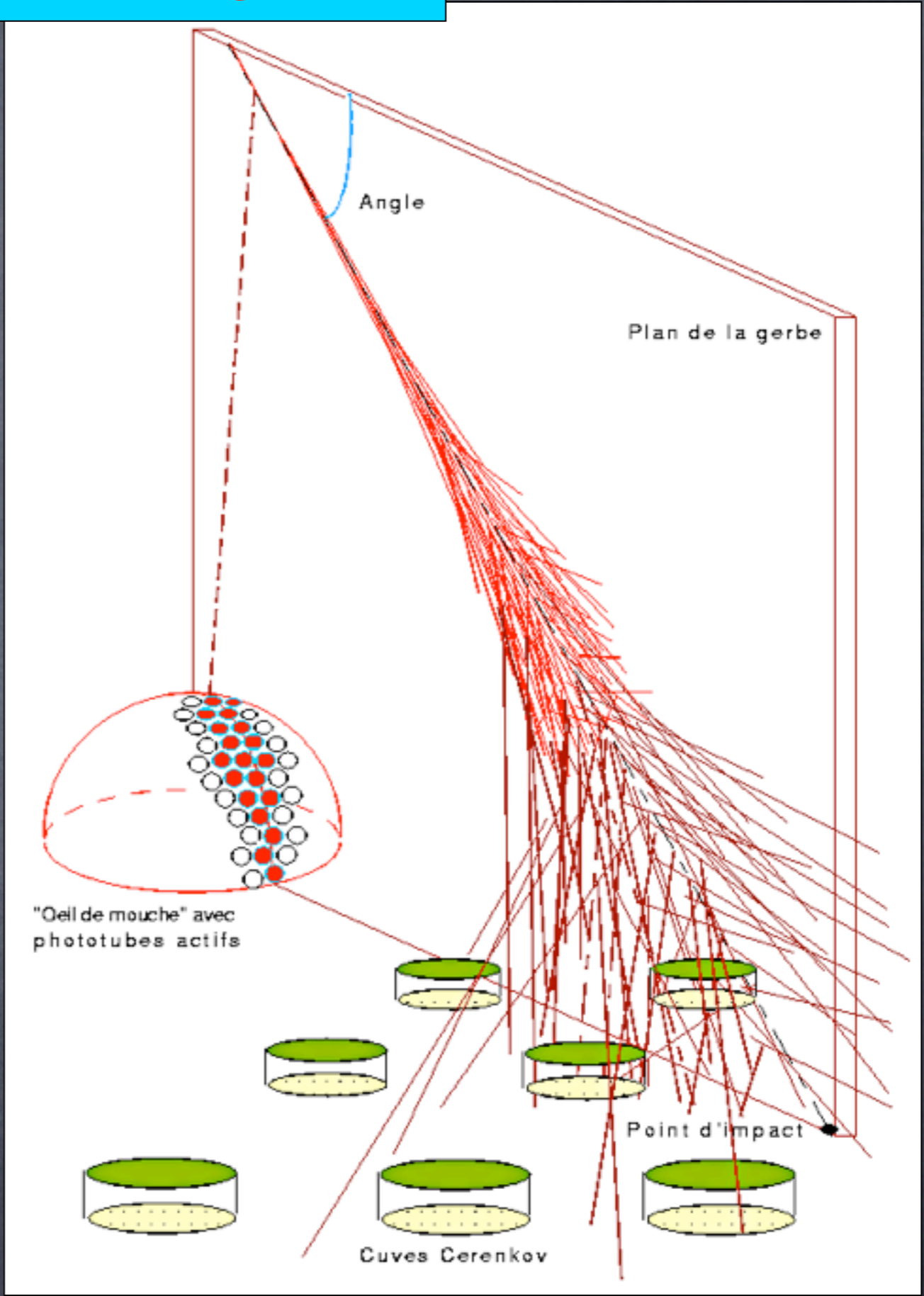
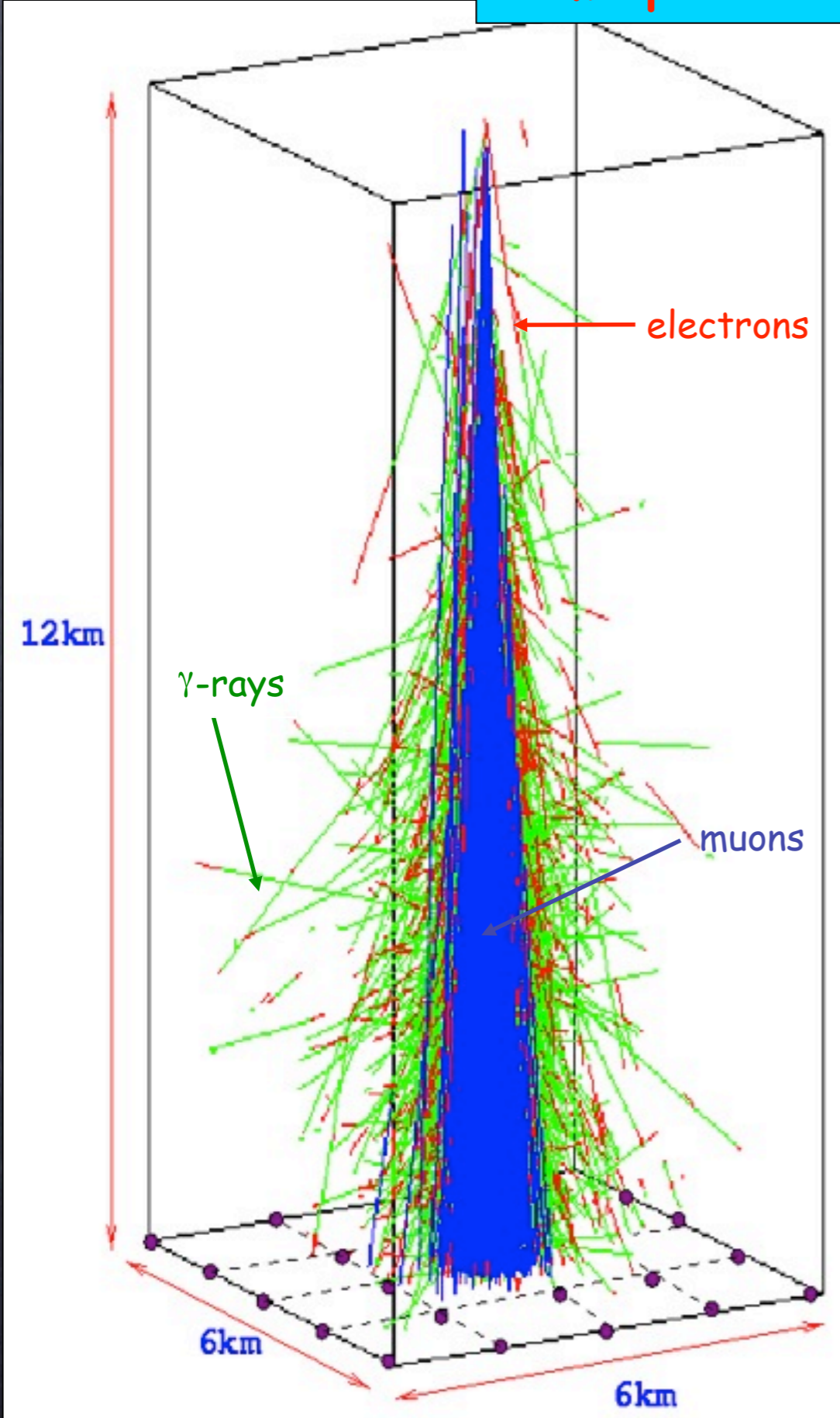
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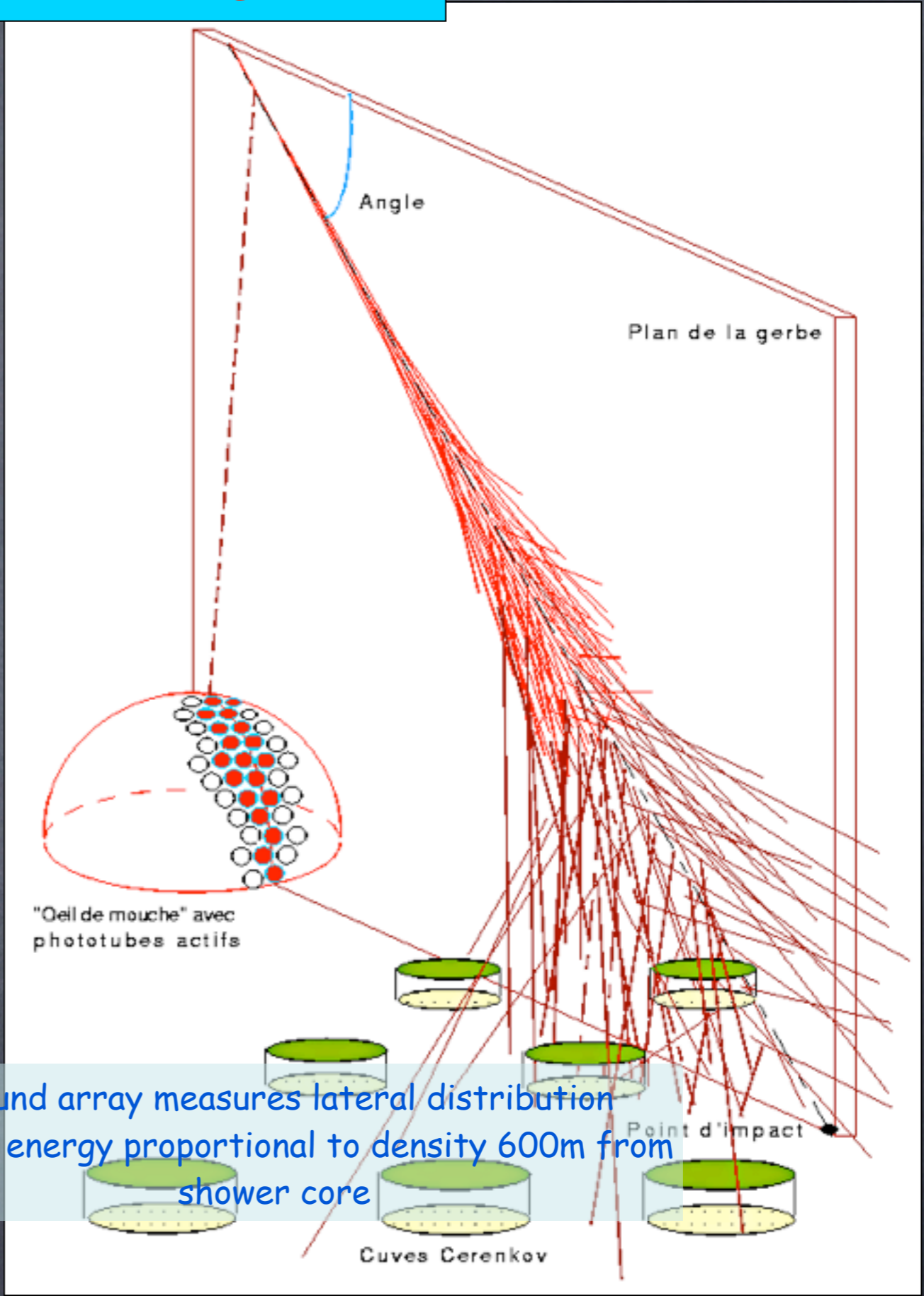
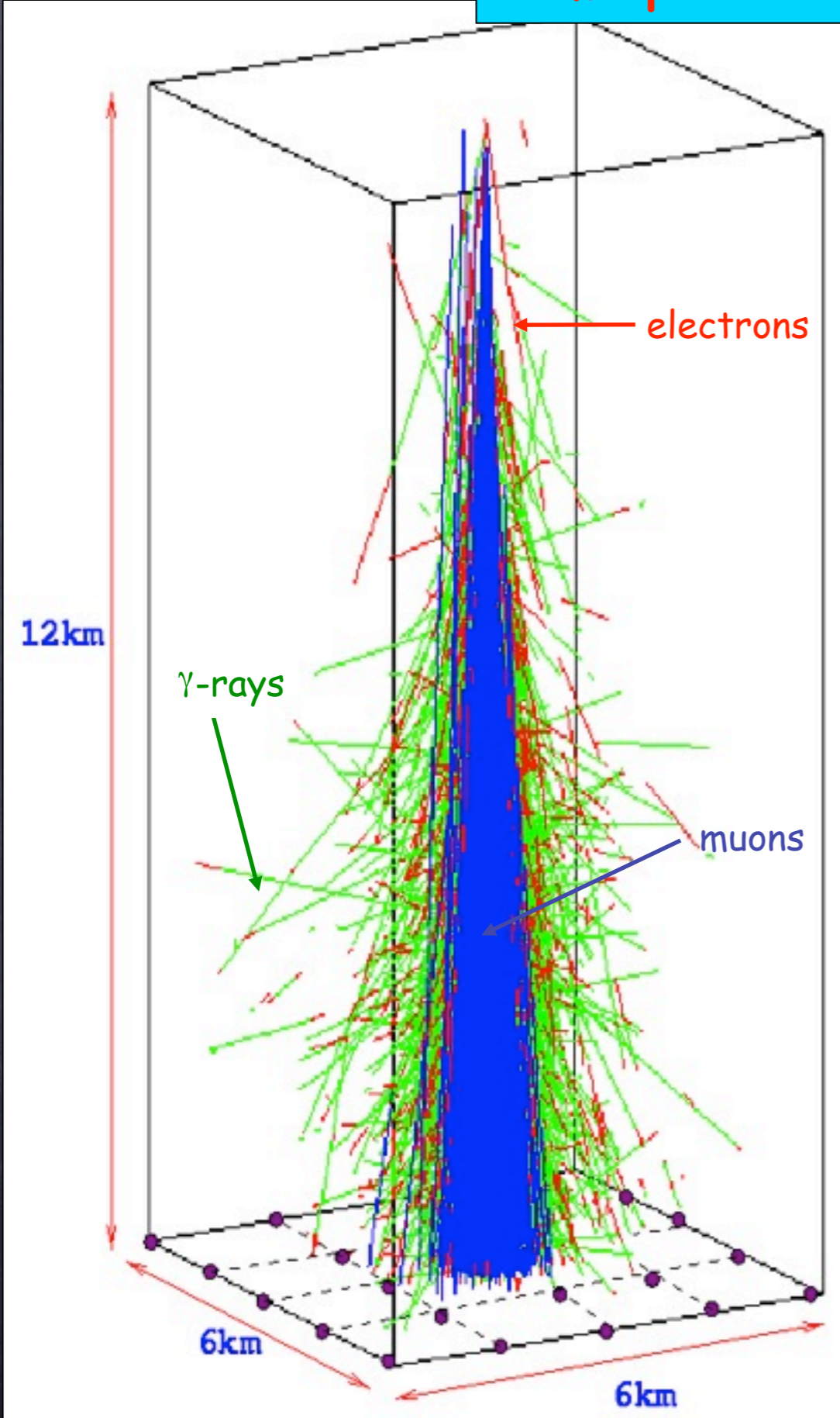
Atmospheric Showers and their Detection



Atmospheric Showers and their Detection

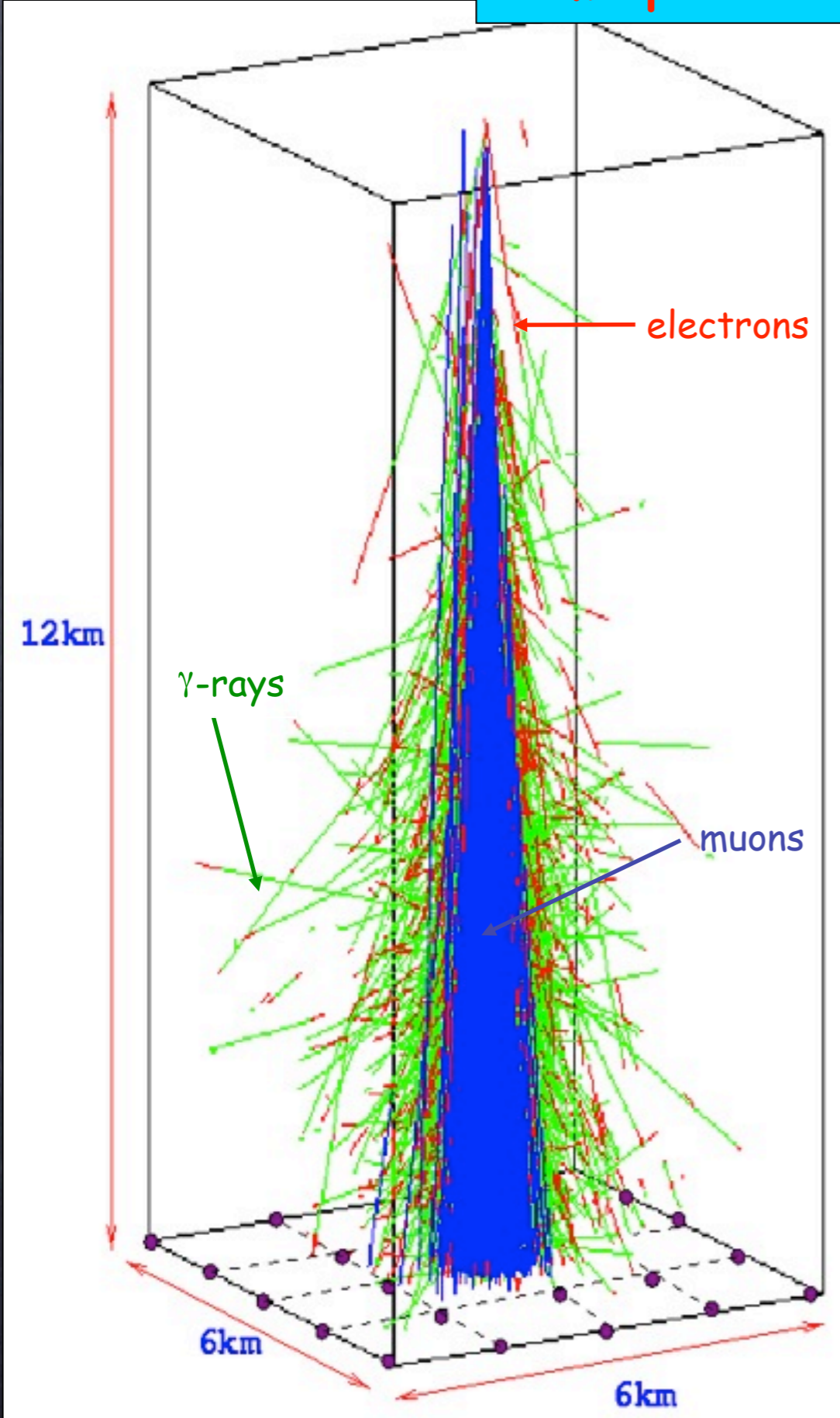


Atmospheric Showers and their Detection



Ground array measures lateral distribution
Primary energy proportional to density 600m from
shower core

Atmospheric Showers and their Detection



Fly's Eye technique measures fluorescence emission

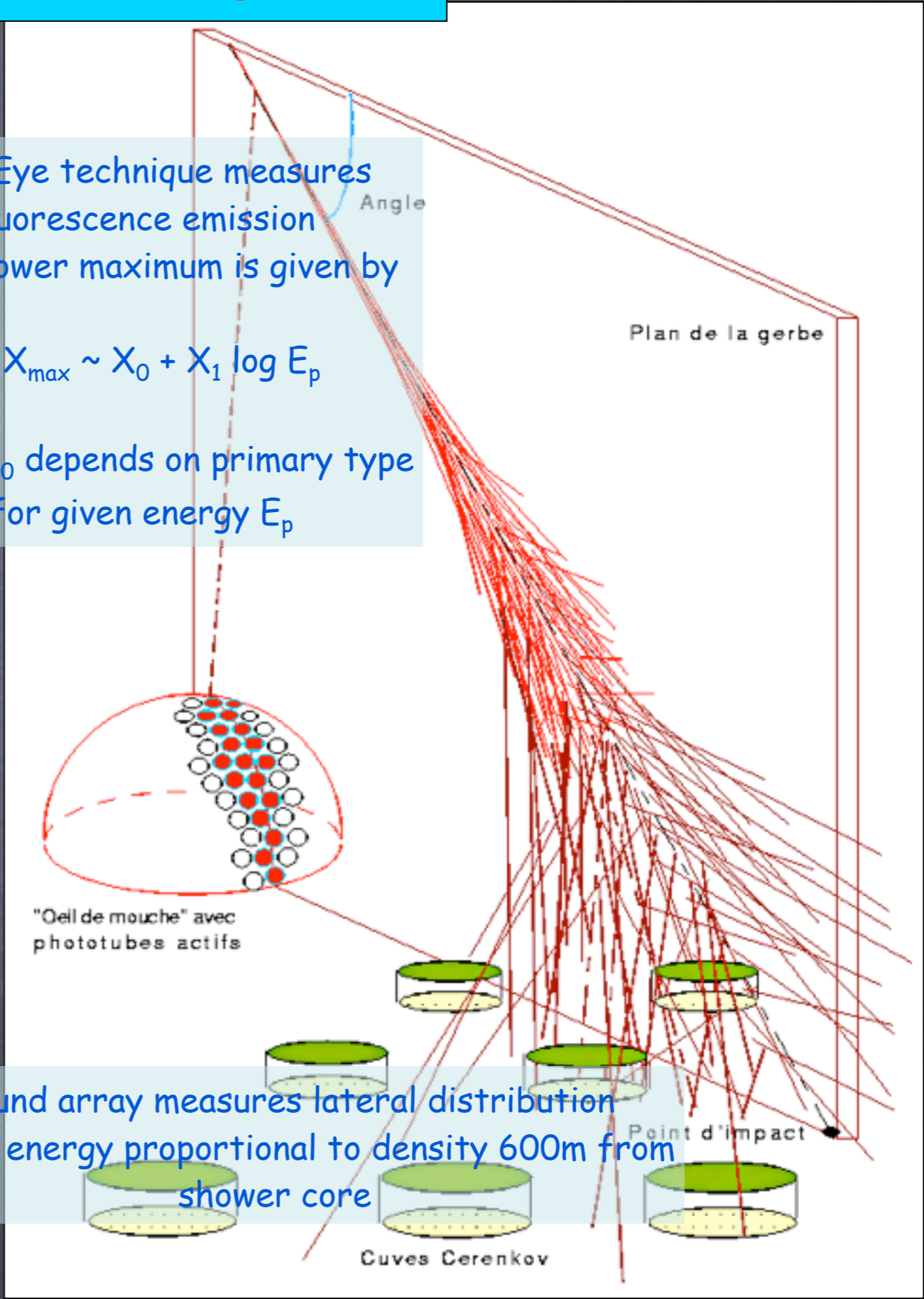
The shower maximum is given by

$$X_{\max} \sim X_0 + X_1 \log E_p$$

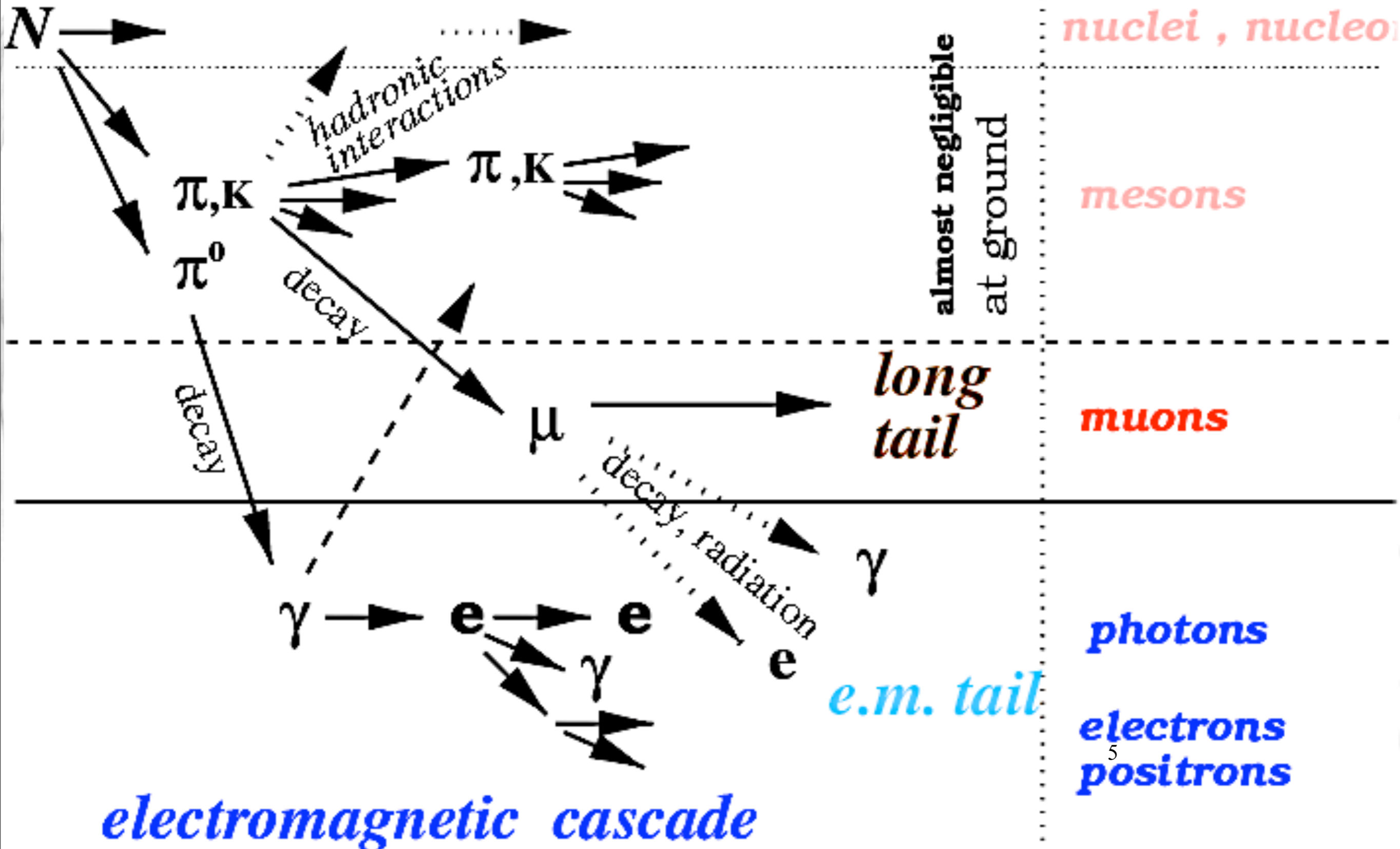
where X_0 depends on primary type for given energy E_p

Ground array measures lateral distribution

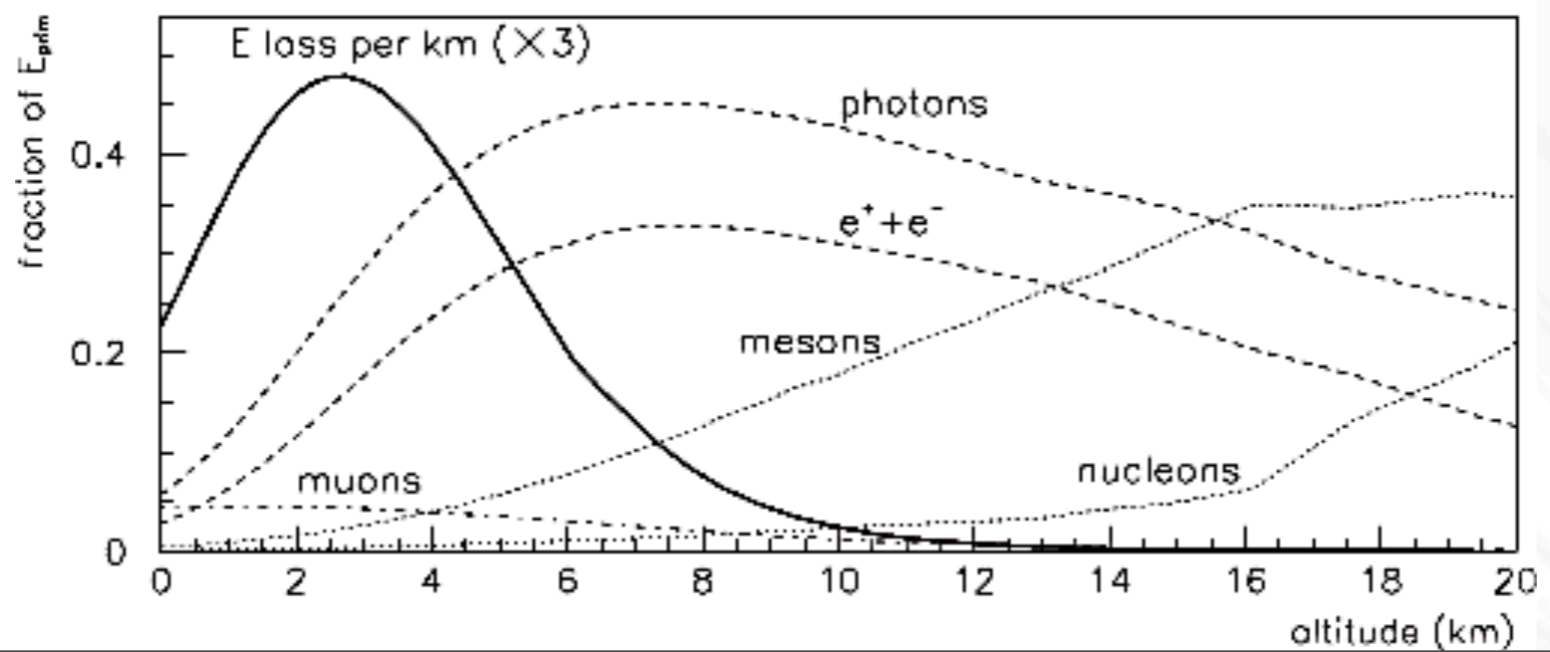
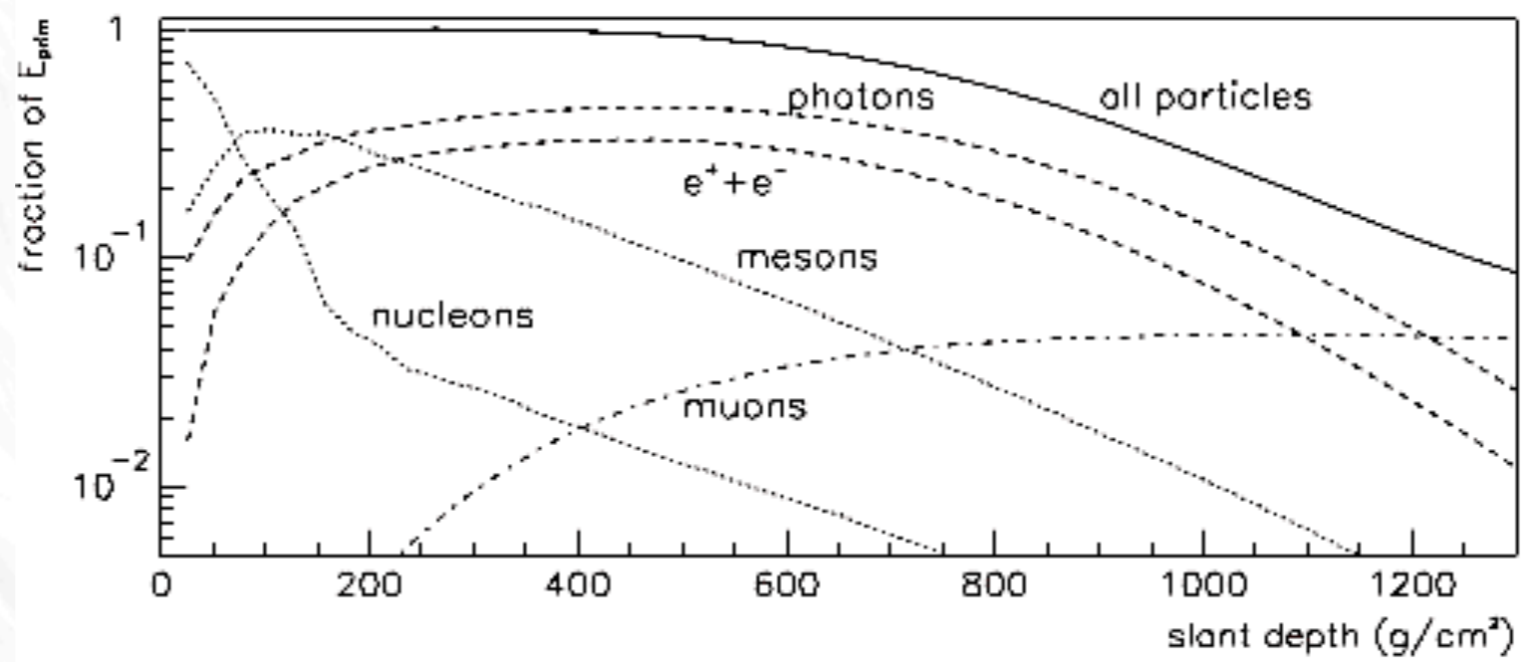
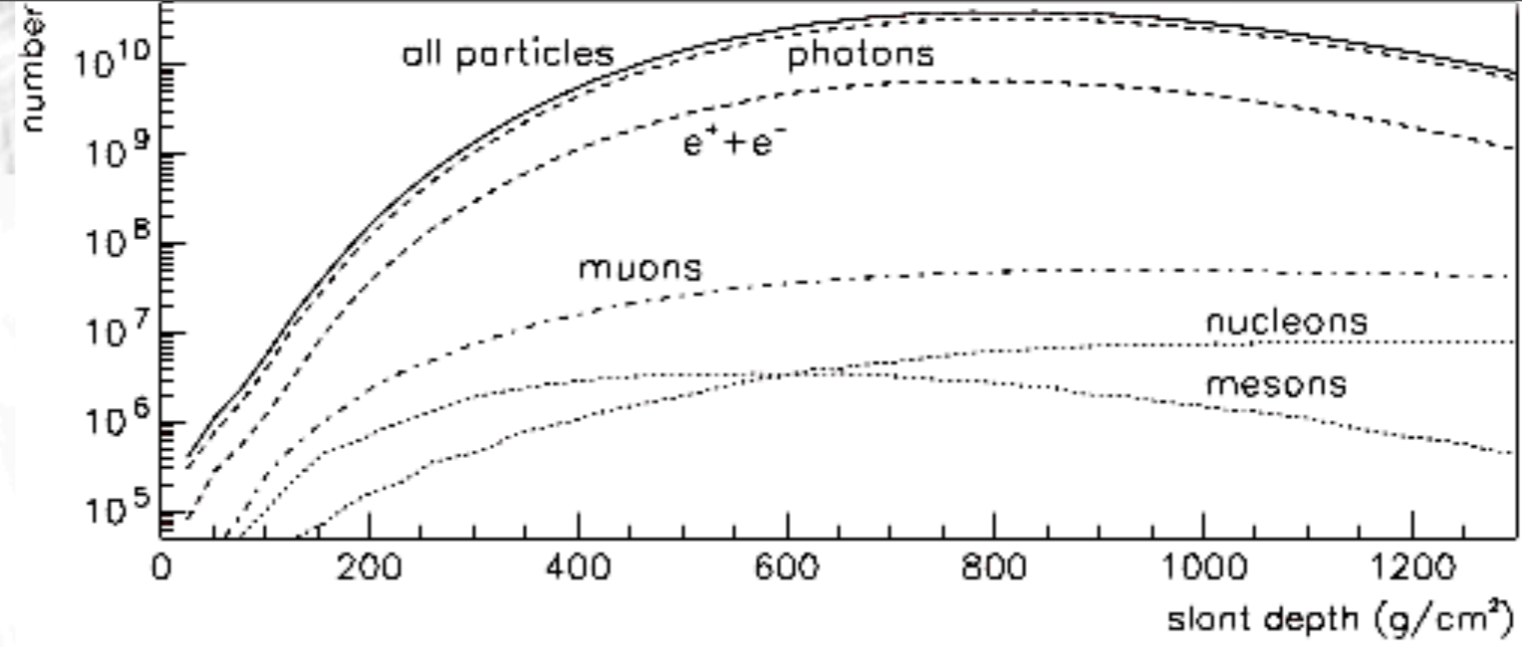
Primary energy proportional to density 600m from shower core



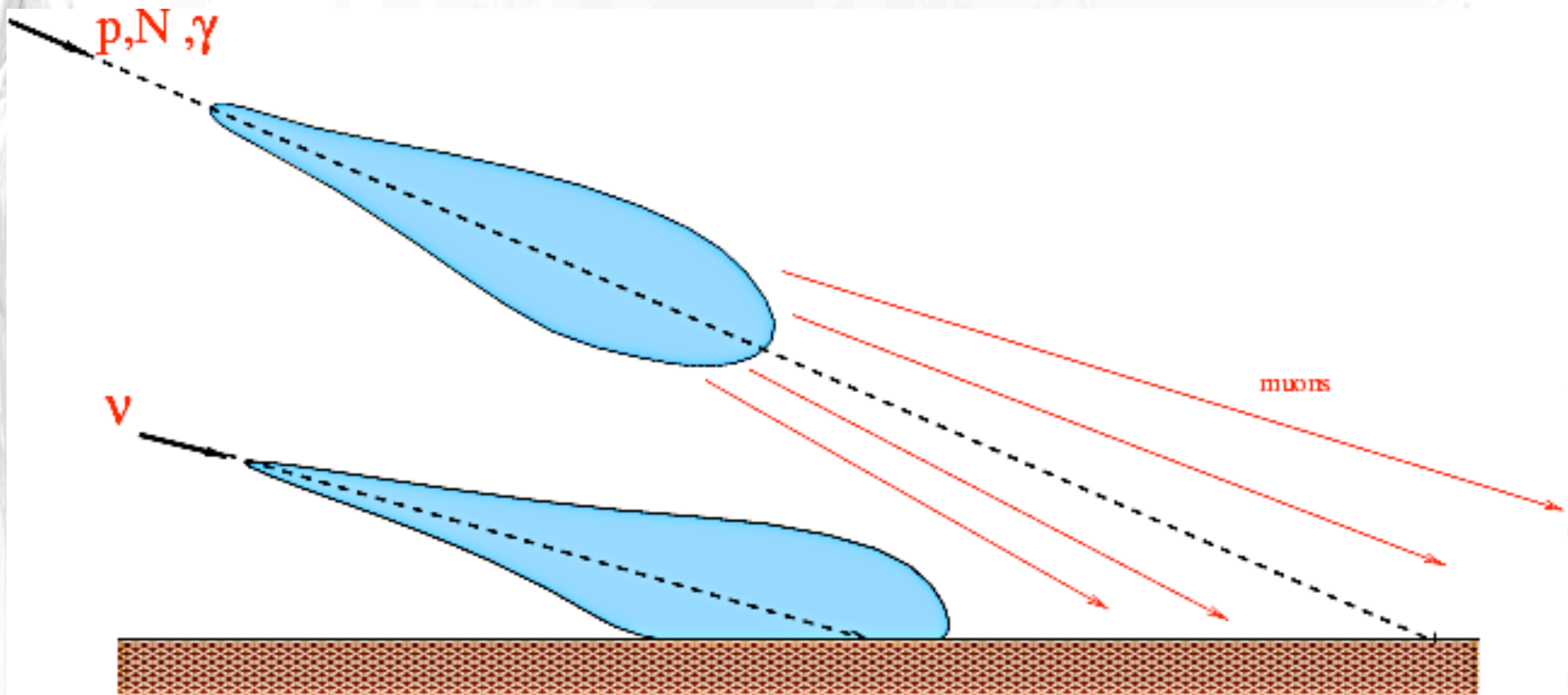
hadronic cascade

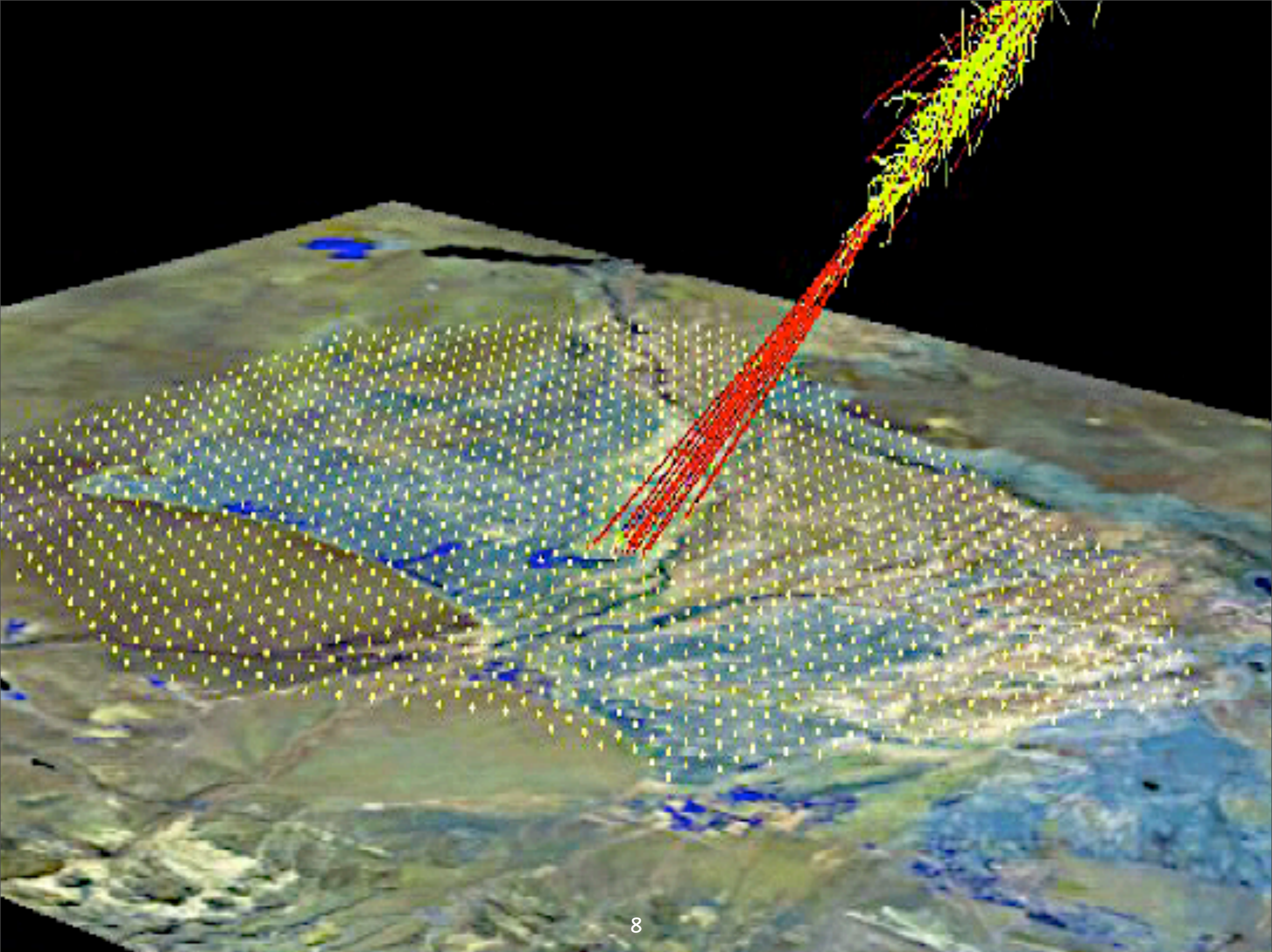


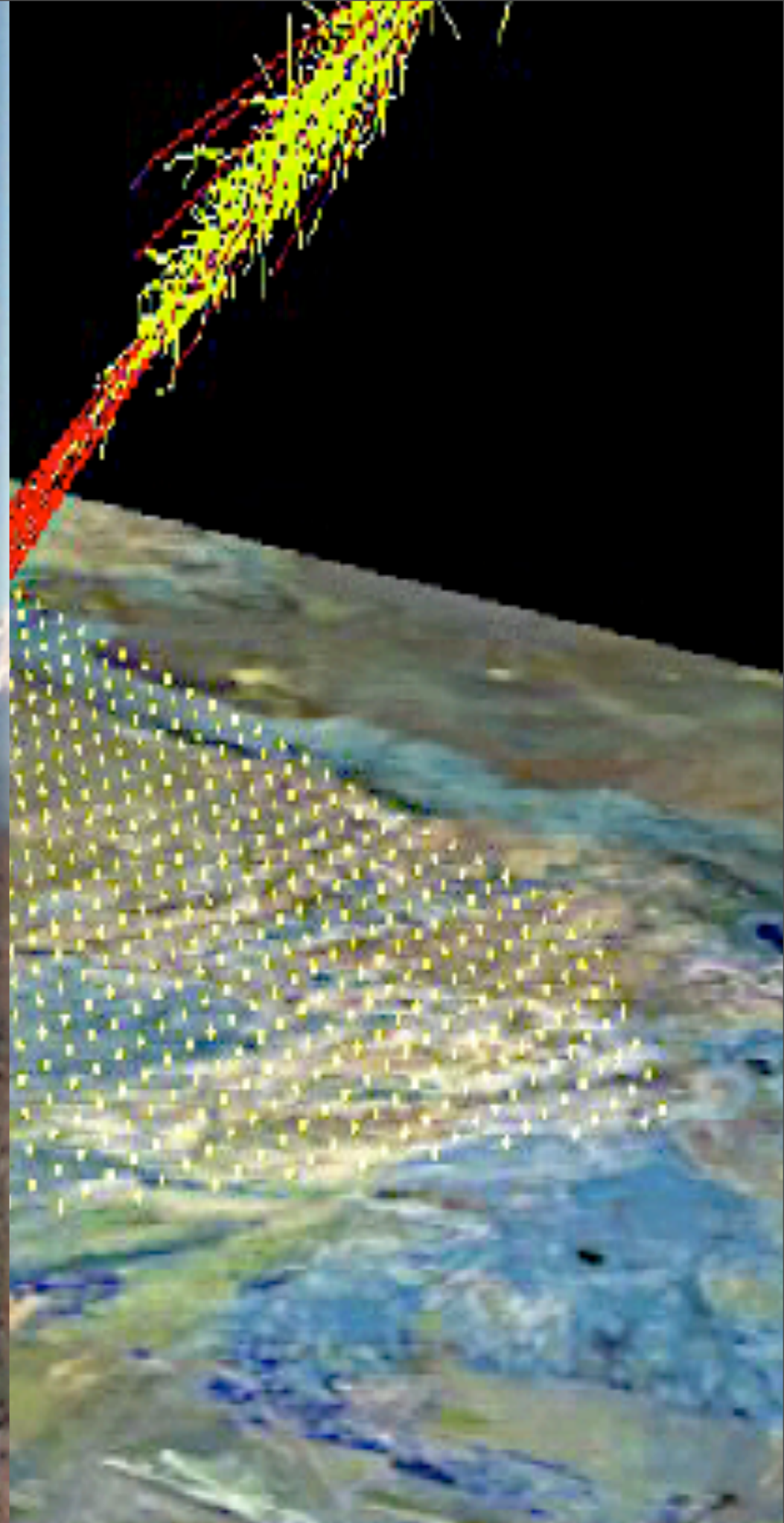
electromagnetic cascade

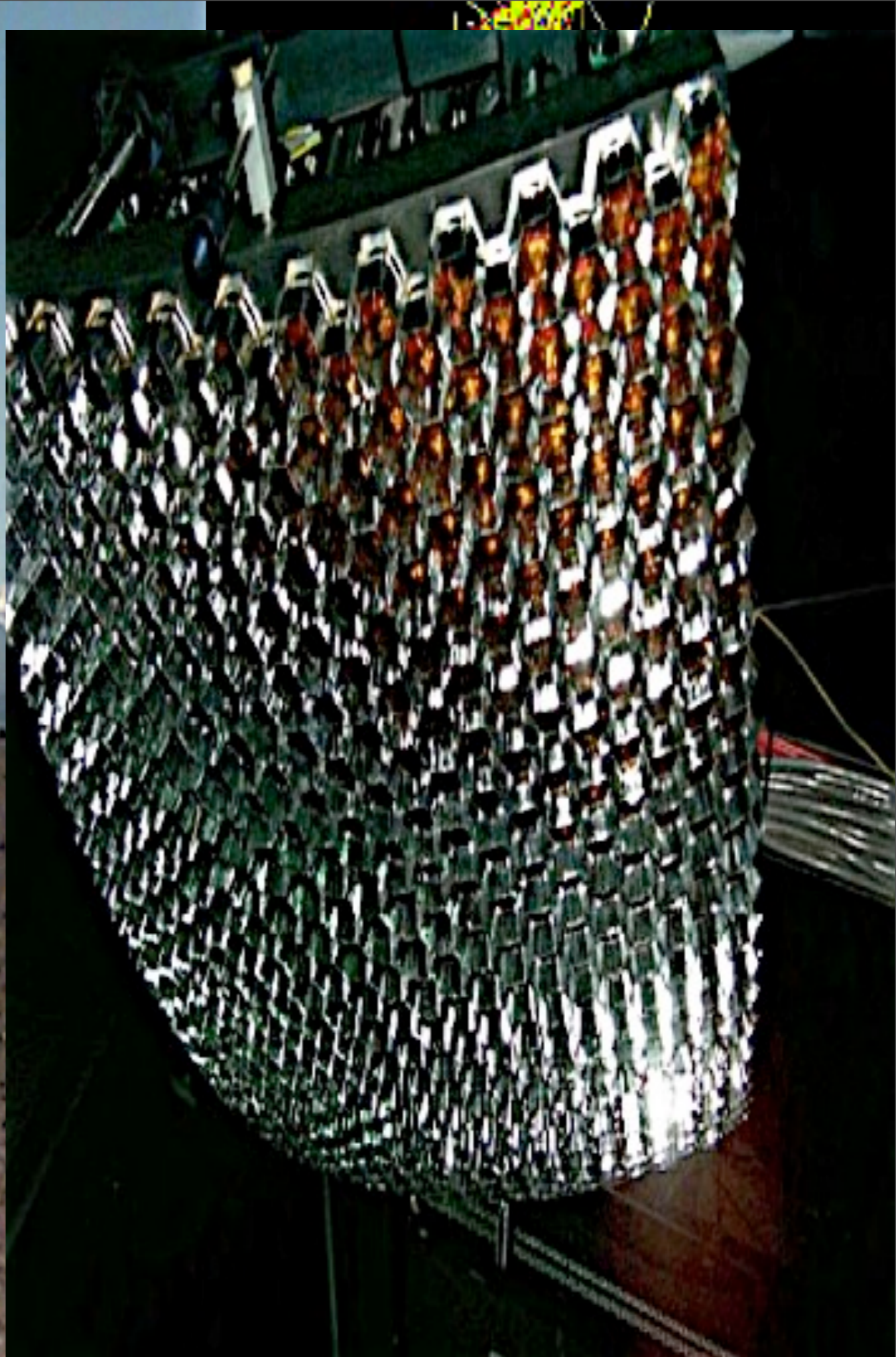
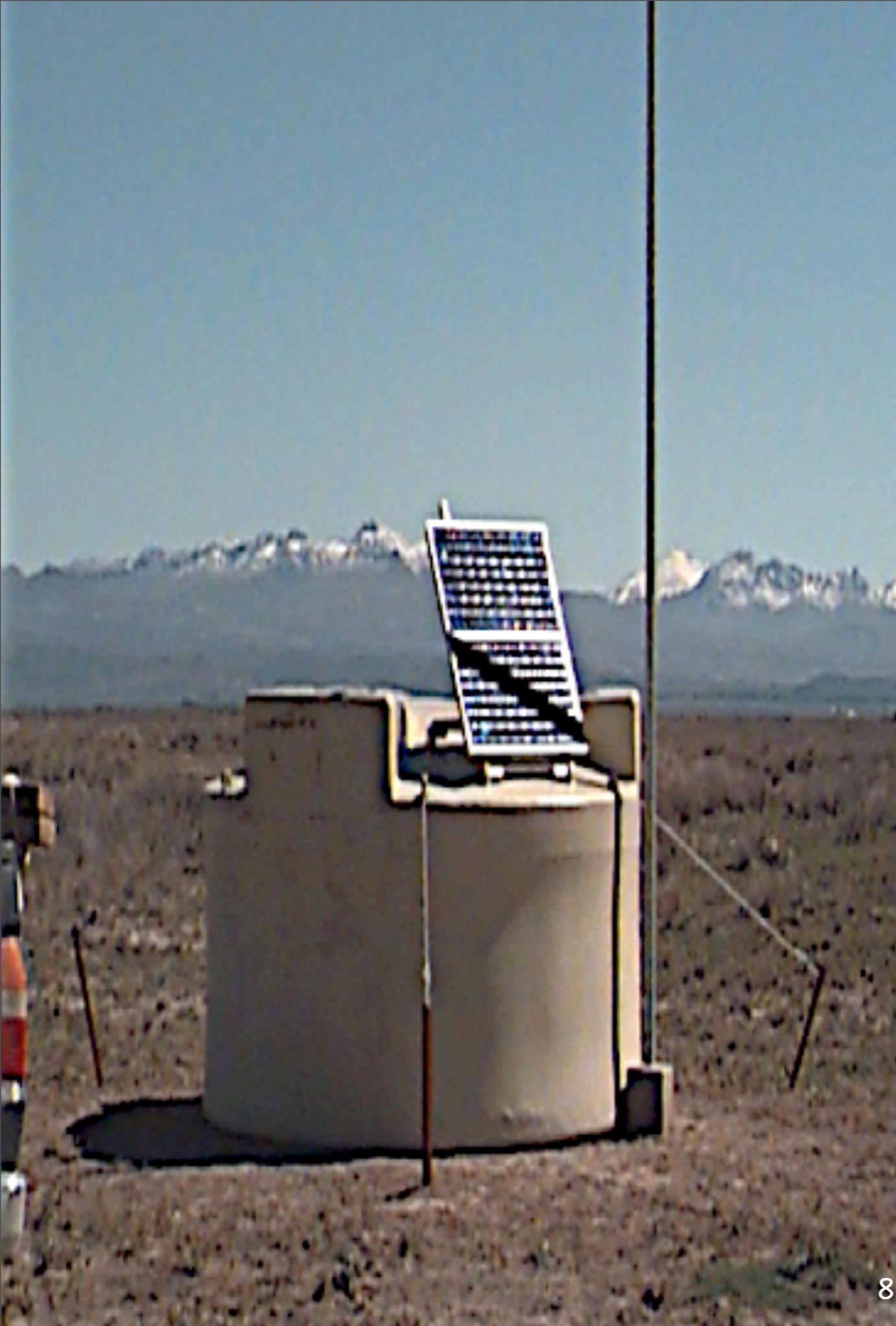


Cosmic ray versus neutrino induced air showers

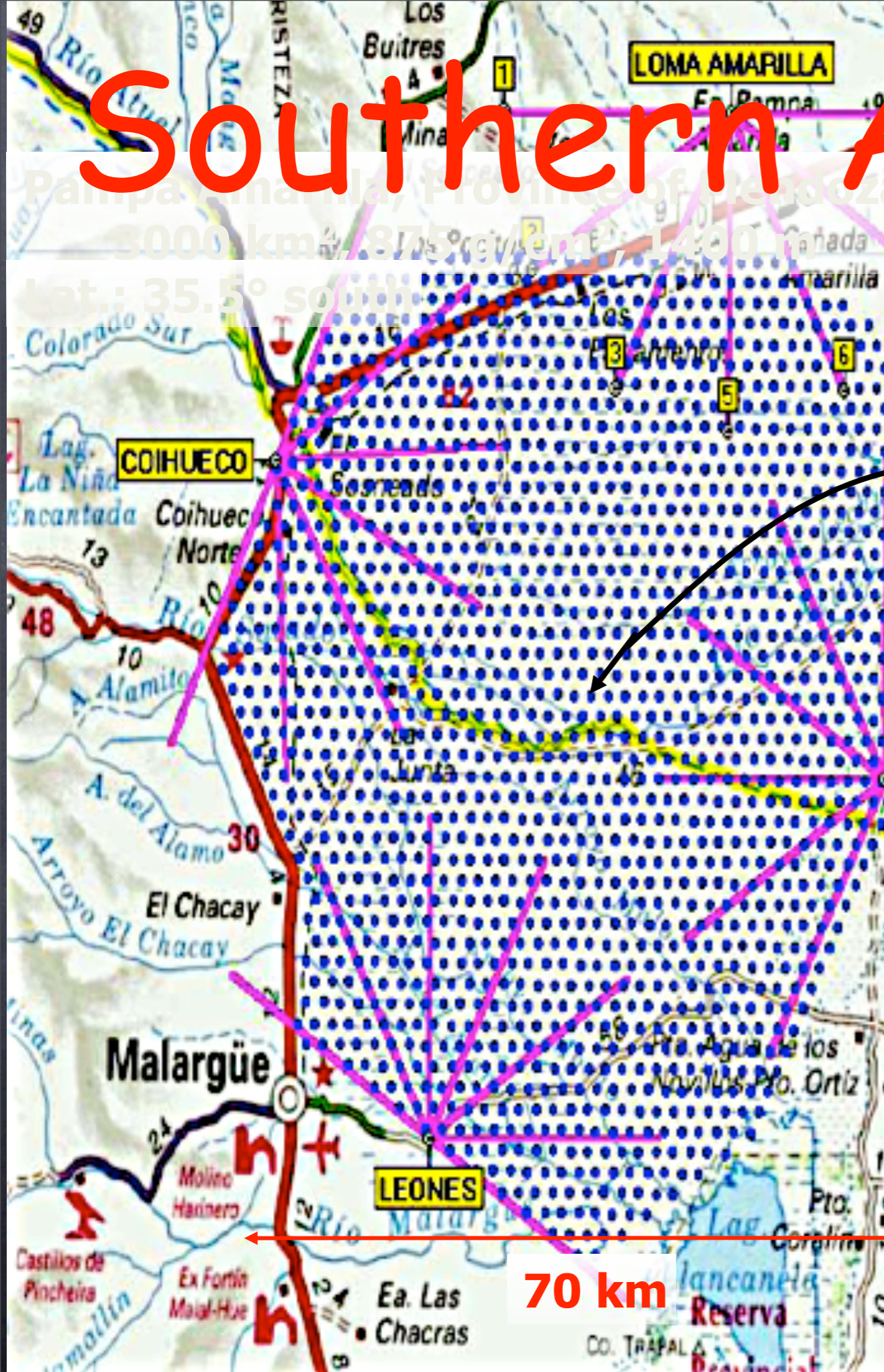




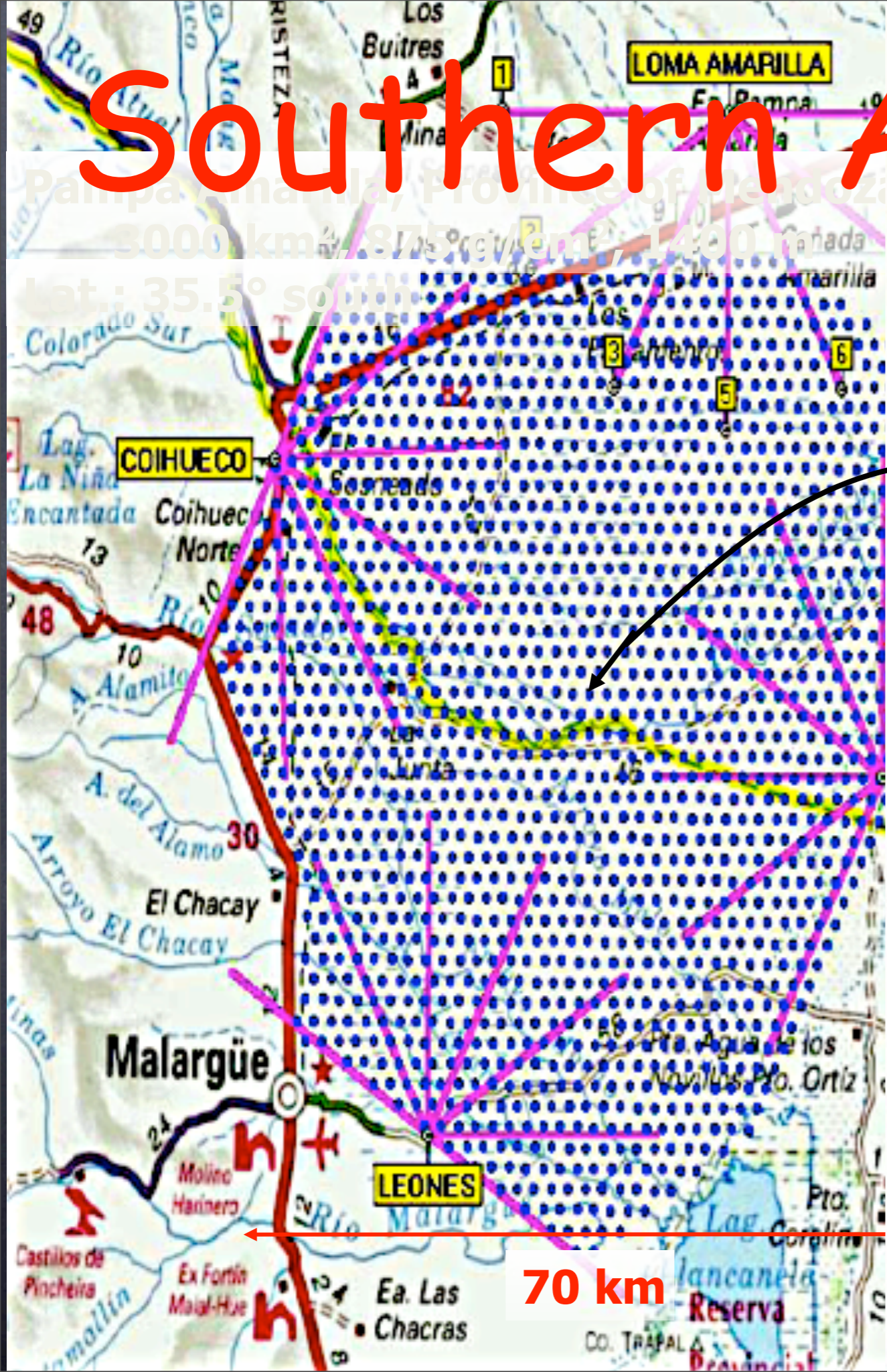




Southern Auger Site



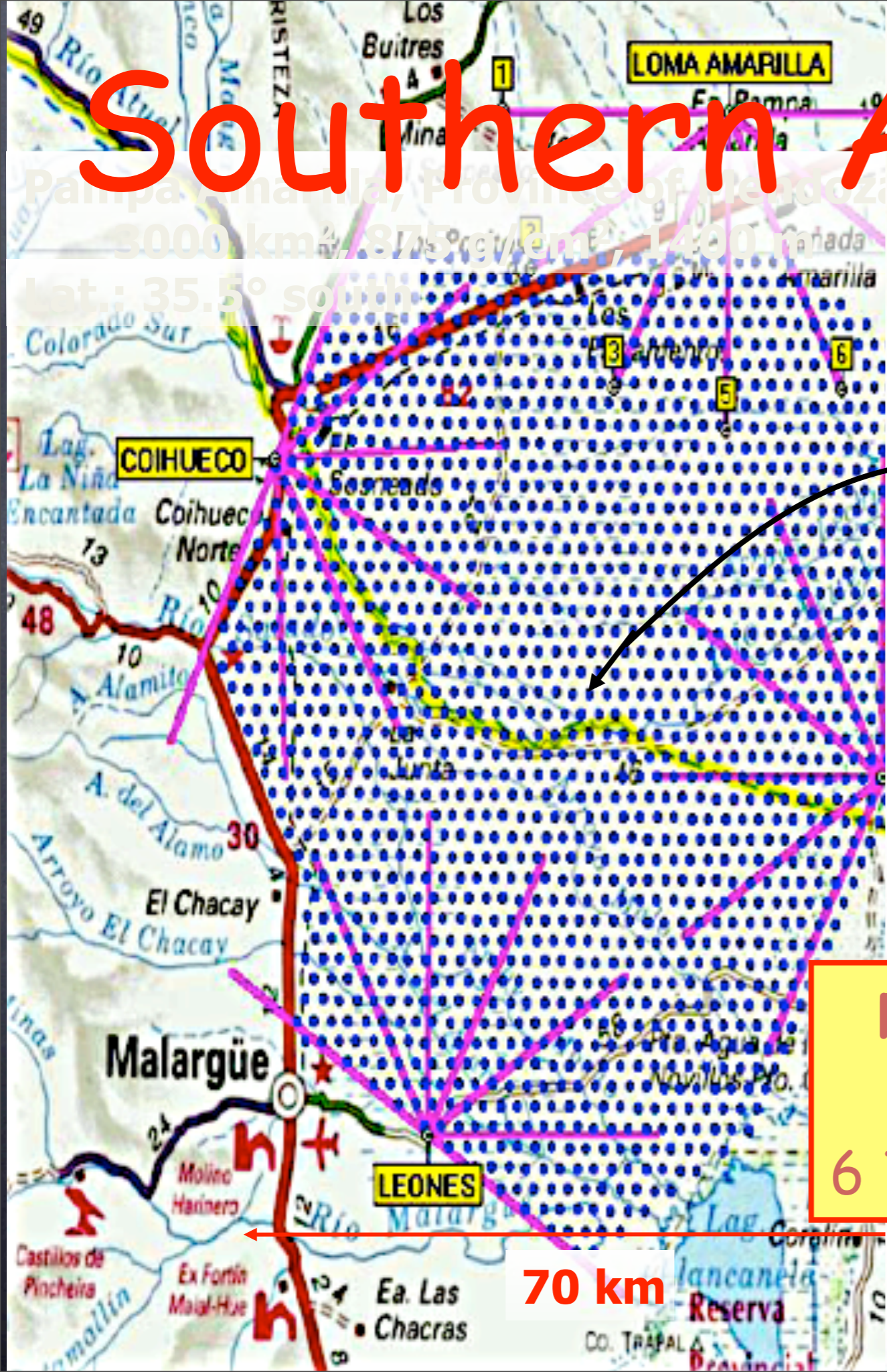
Southern Auger Site



Surface Array (SD):
1600 Water Tanks
1.5 km spacing
3000 km²

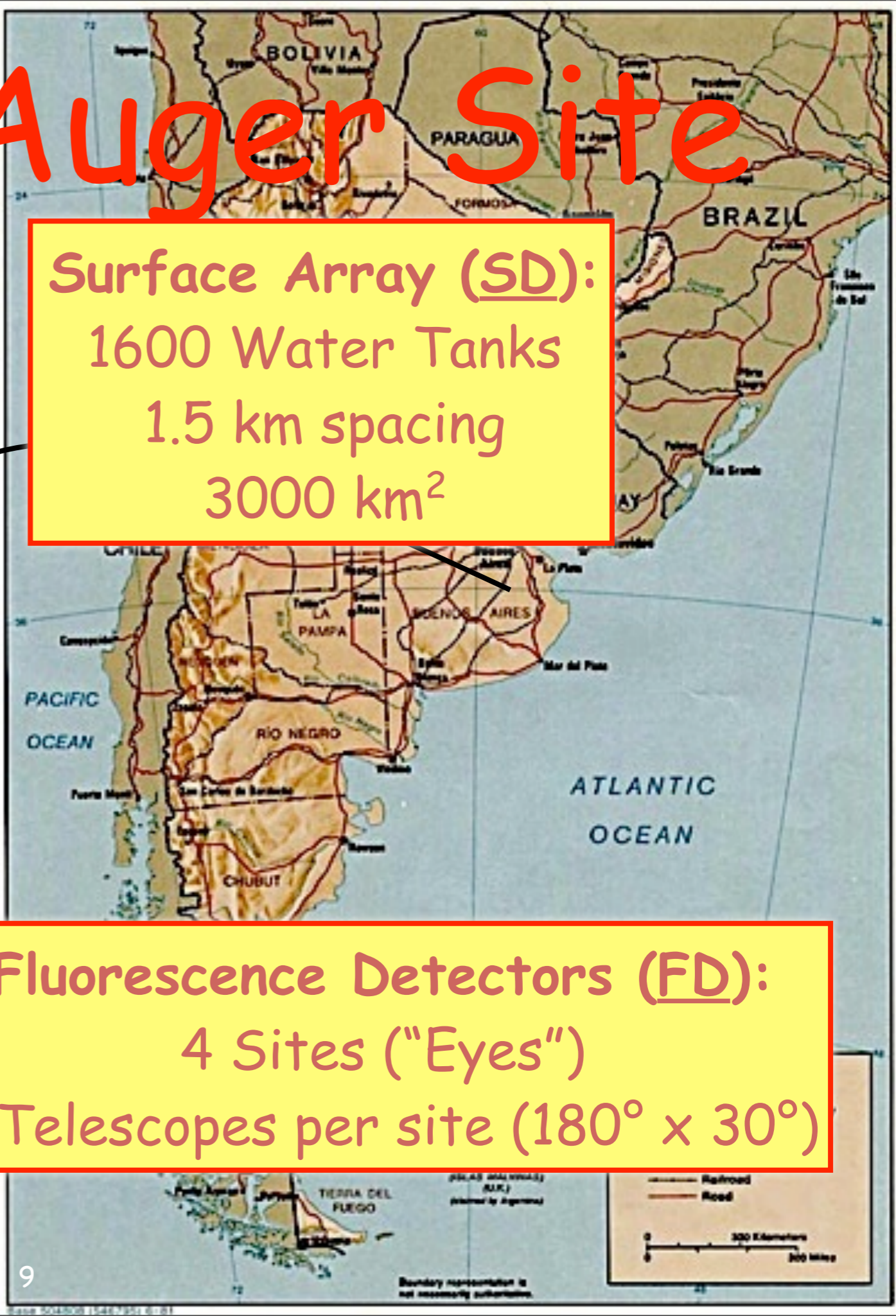


Southern Auger Site



Surface Array (SD):
1600 Water Tanks
1.5 km spacing
3000 km²

Fluorescence Detectors (FD):
4 Sites ("Eyes")
6 Telescopes per site (180° x 30°)



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2.) in most conventional scenarios exceptionally powerful acceleration sources within that distance are needed.

3.) The observed distribution does not yet reveal unambiguously the sources, although there is some correlation with local large scale structure

4.) The observed mass composition may become heavy toward highest energies, but no consistent picture yet between experiments and air shower models

The Greisen-Zatsepin-Kuzmin (GZK) effect

Nucleons can produce pions on the cosmic microwave background

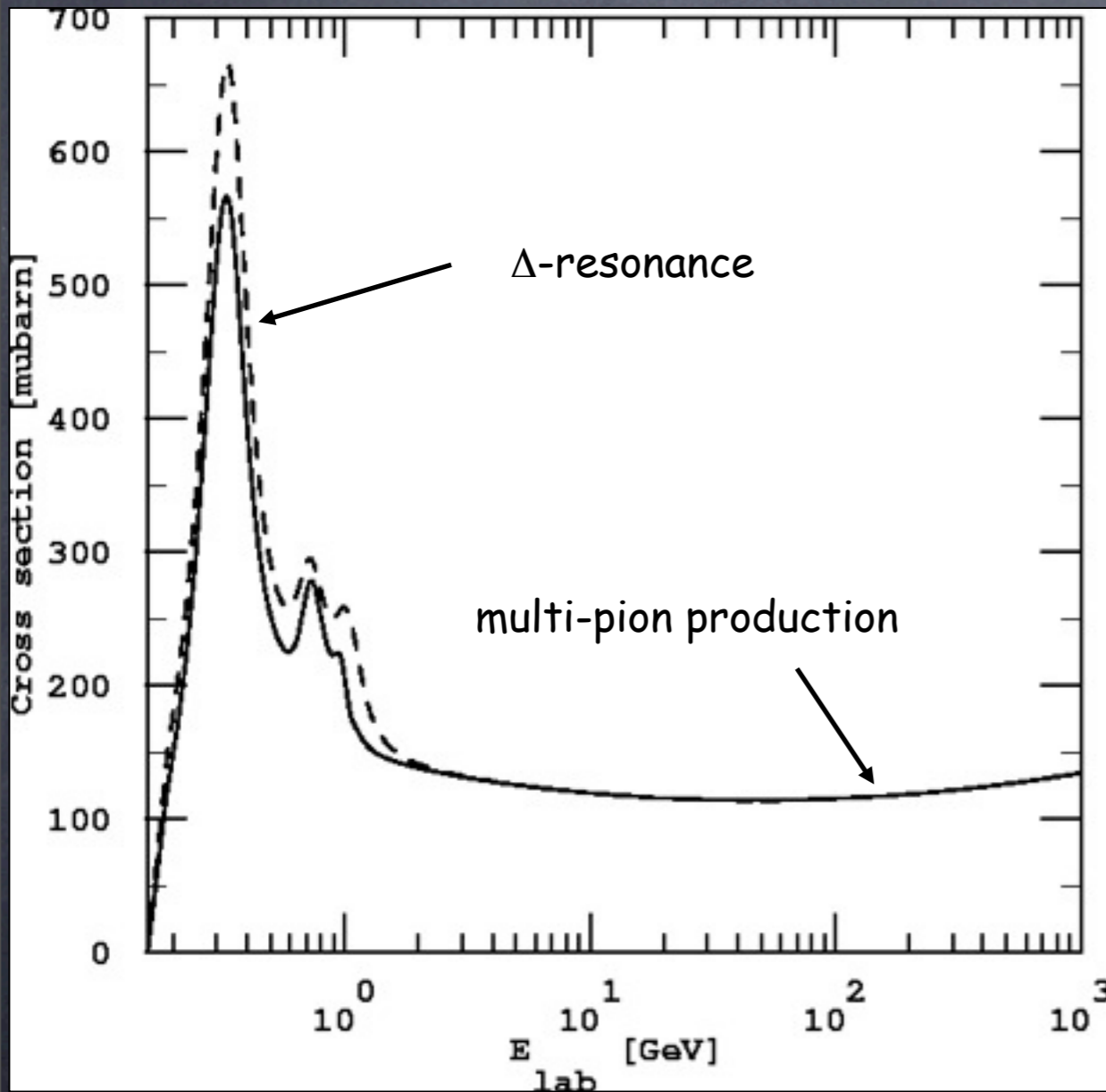


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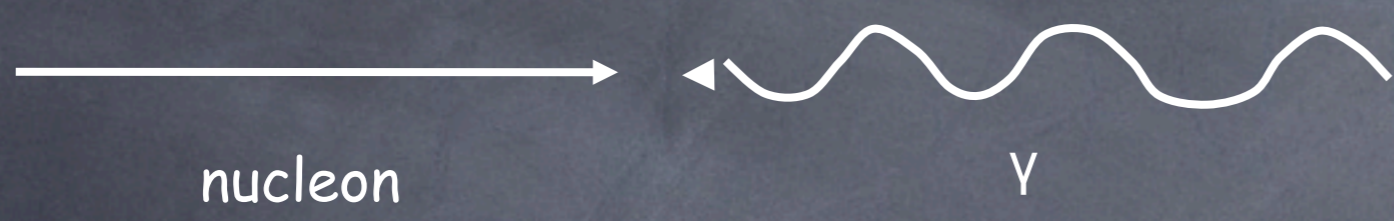


$$E_{\text{th}} = \frac{2m_N m_\pi + m_\pi^2}{4\epsilon} \simeq 4 \times 10^{19} \text{ eV}$$

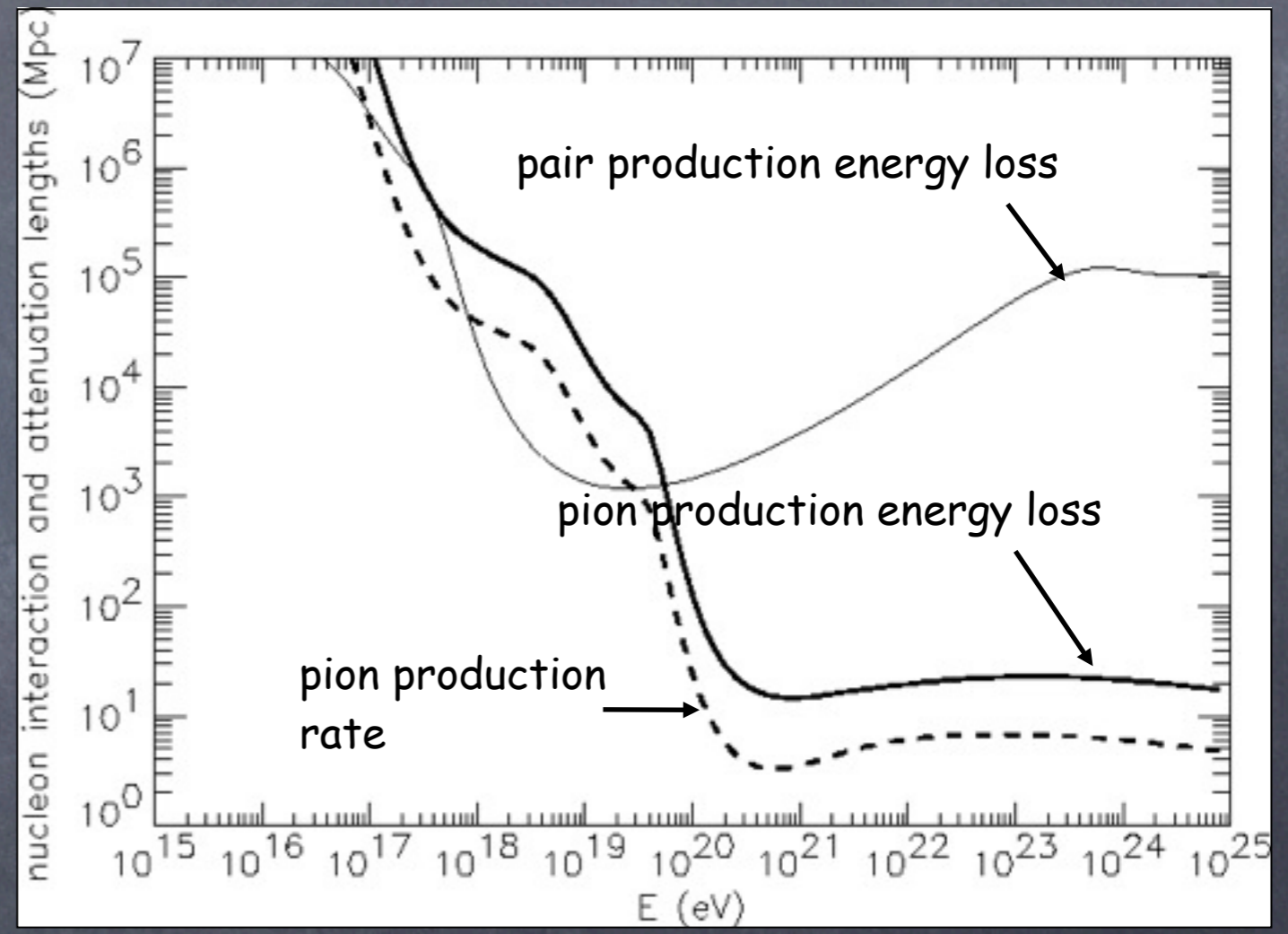
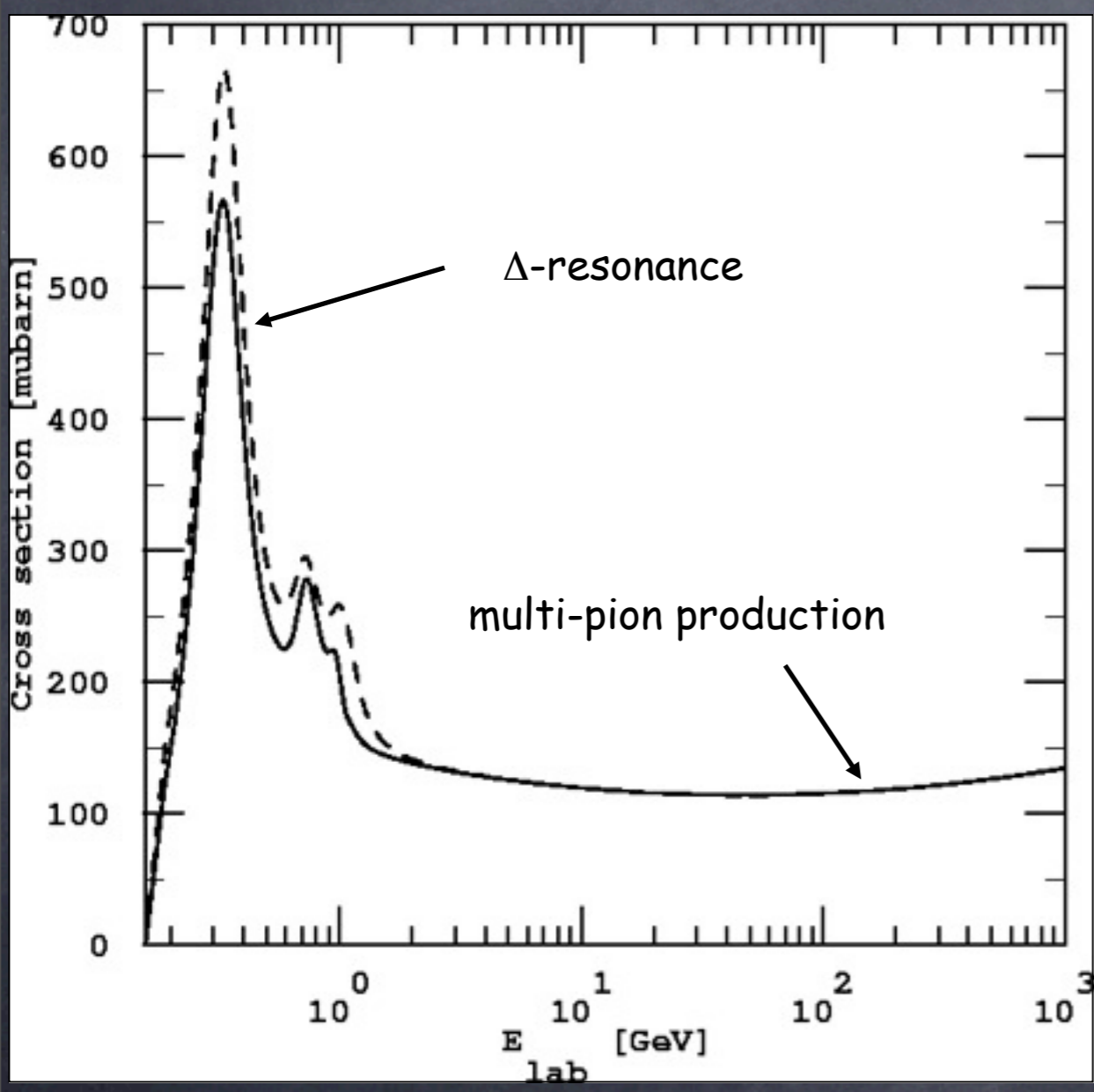


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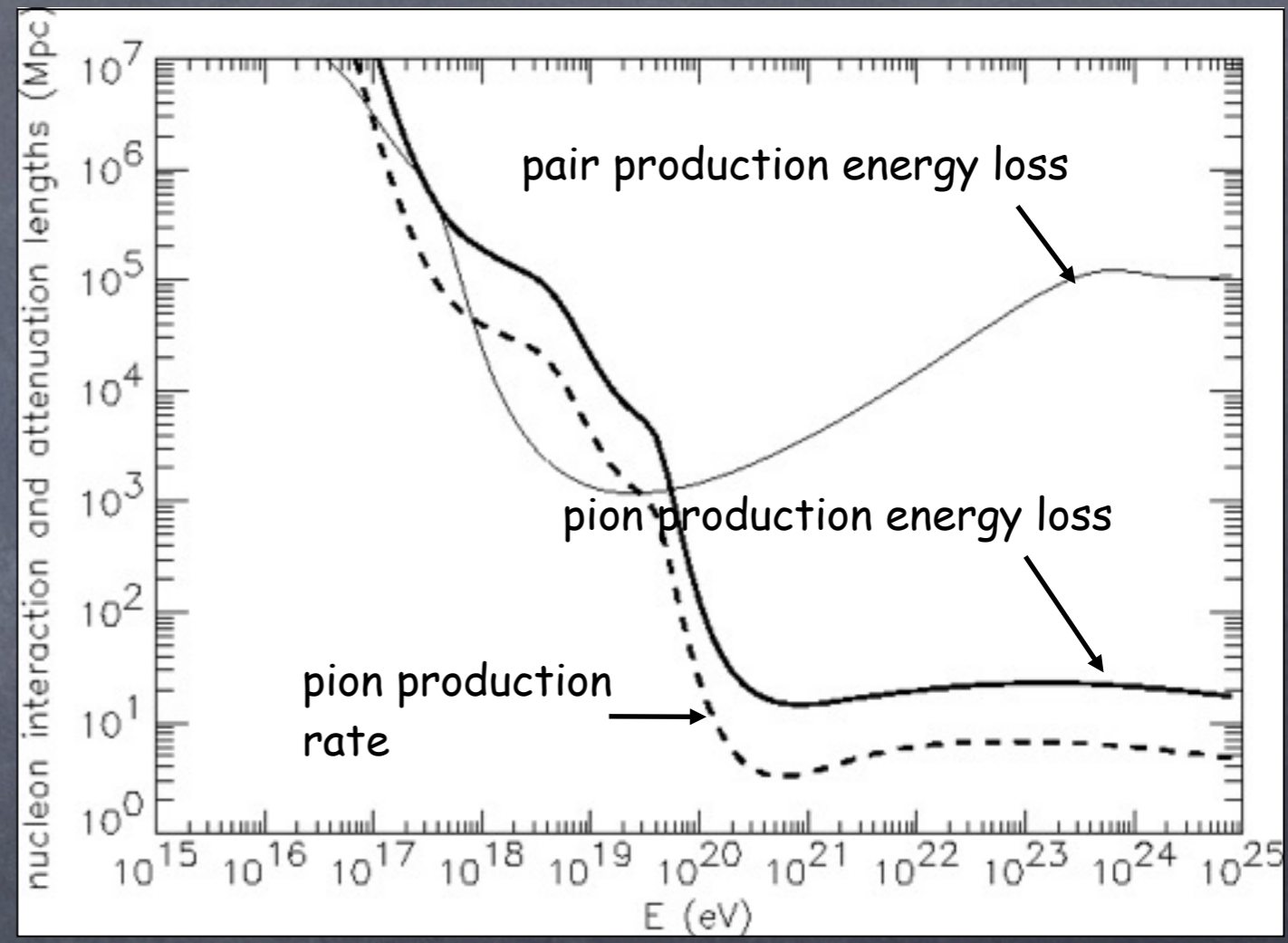
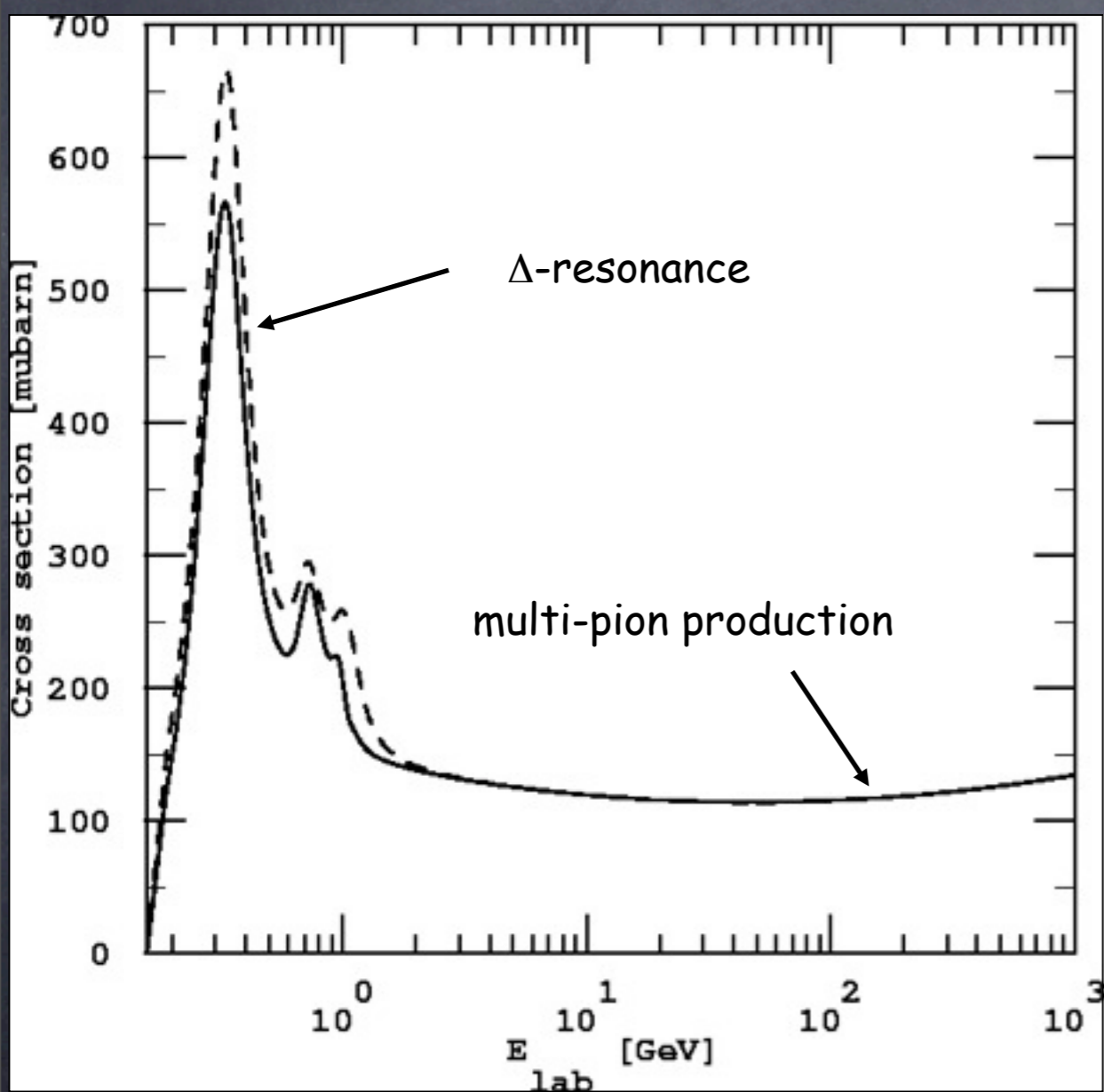


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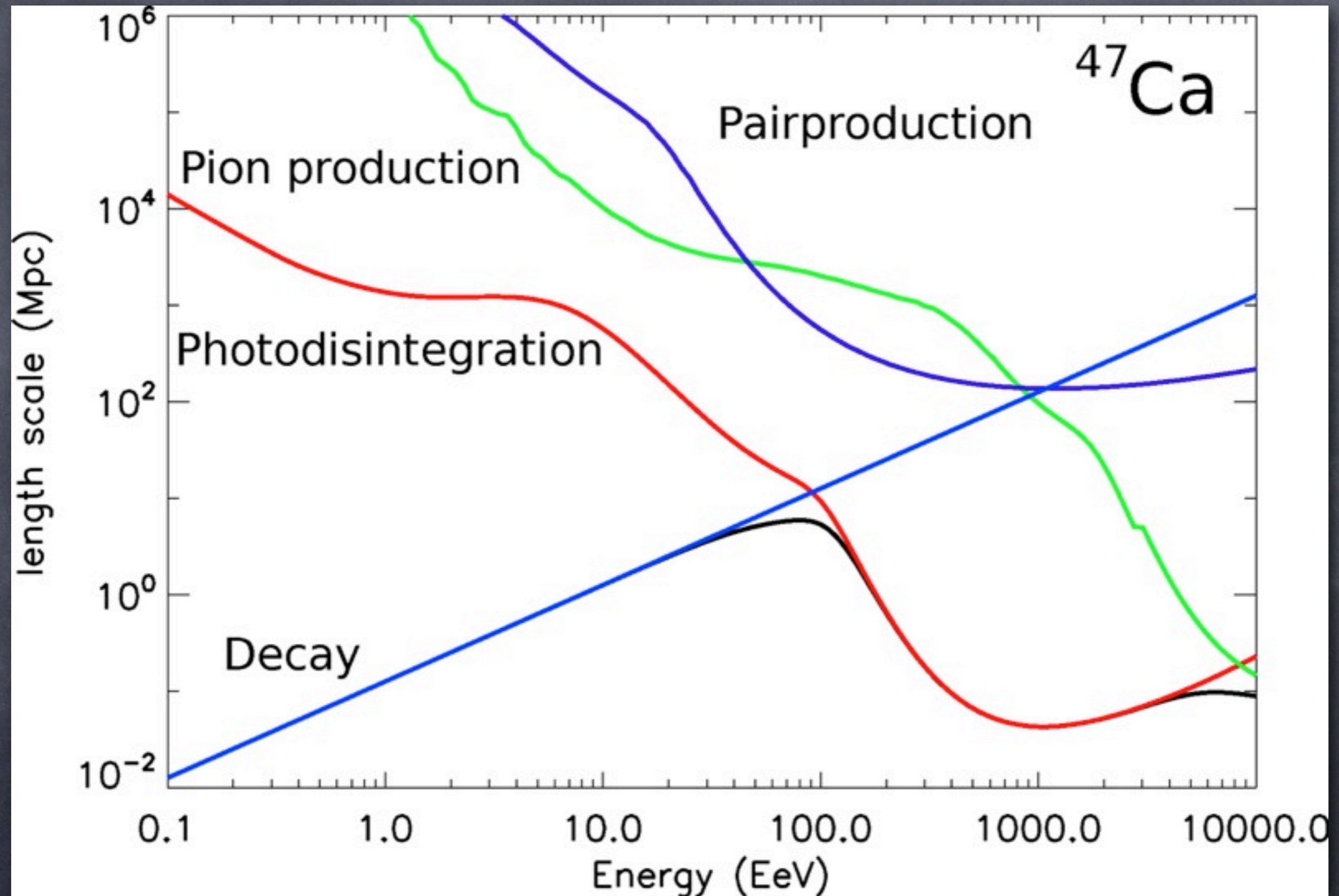
The Greisen-Zatsepin-Kuzmin (GZK) effect

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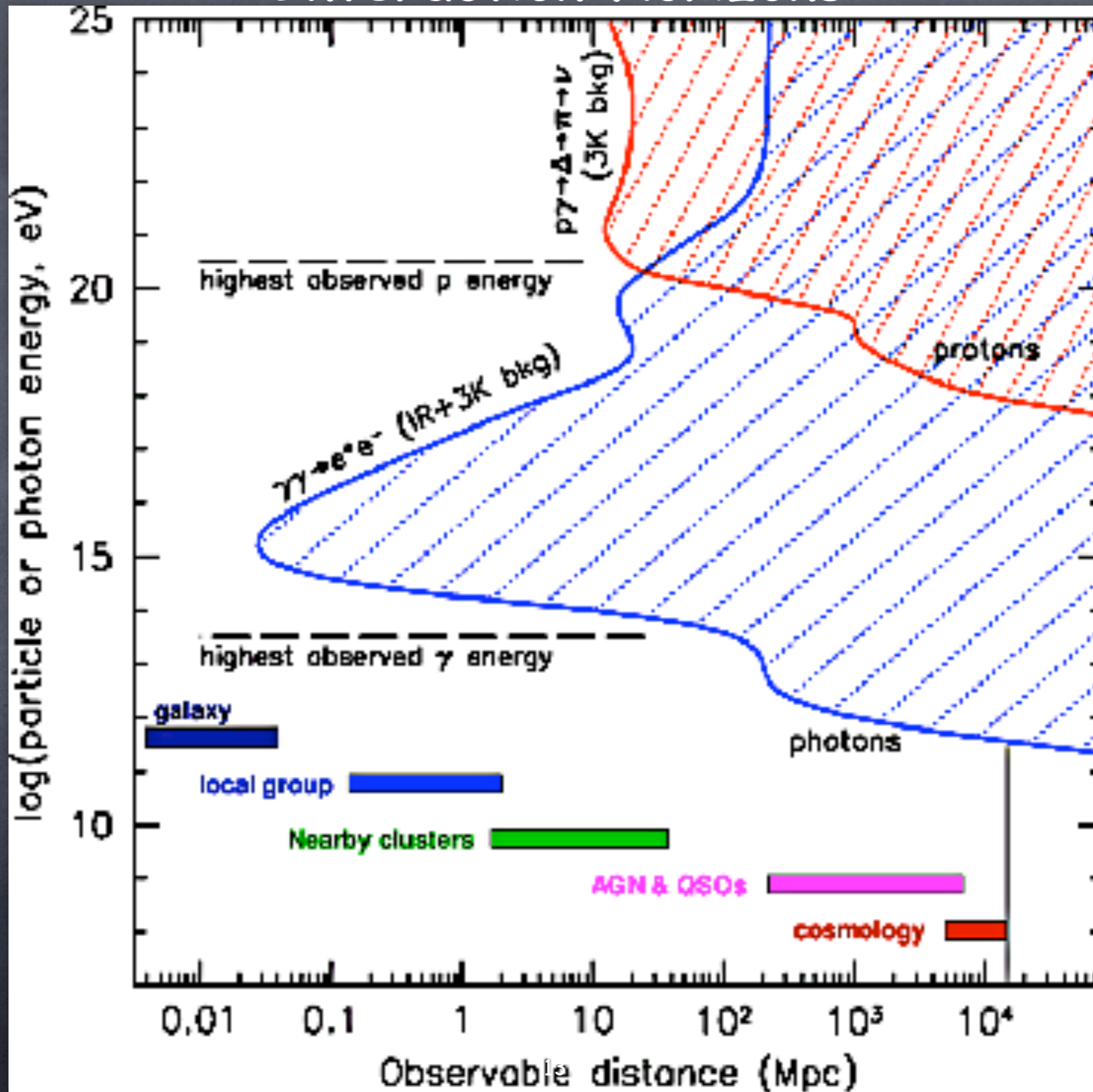


sources must be in cosmological backyard
 Only Lorentz symmetry breaking at $\Gamma > 10^{11}$
 could avoid this conclusion.

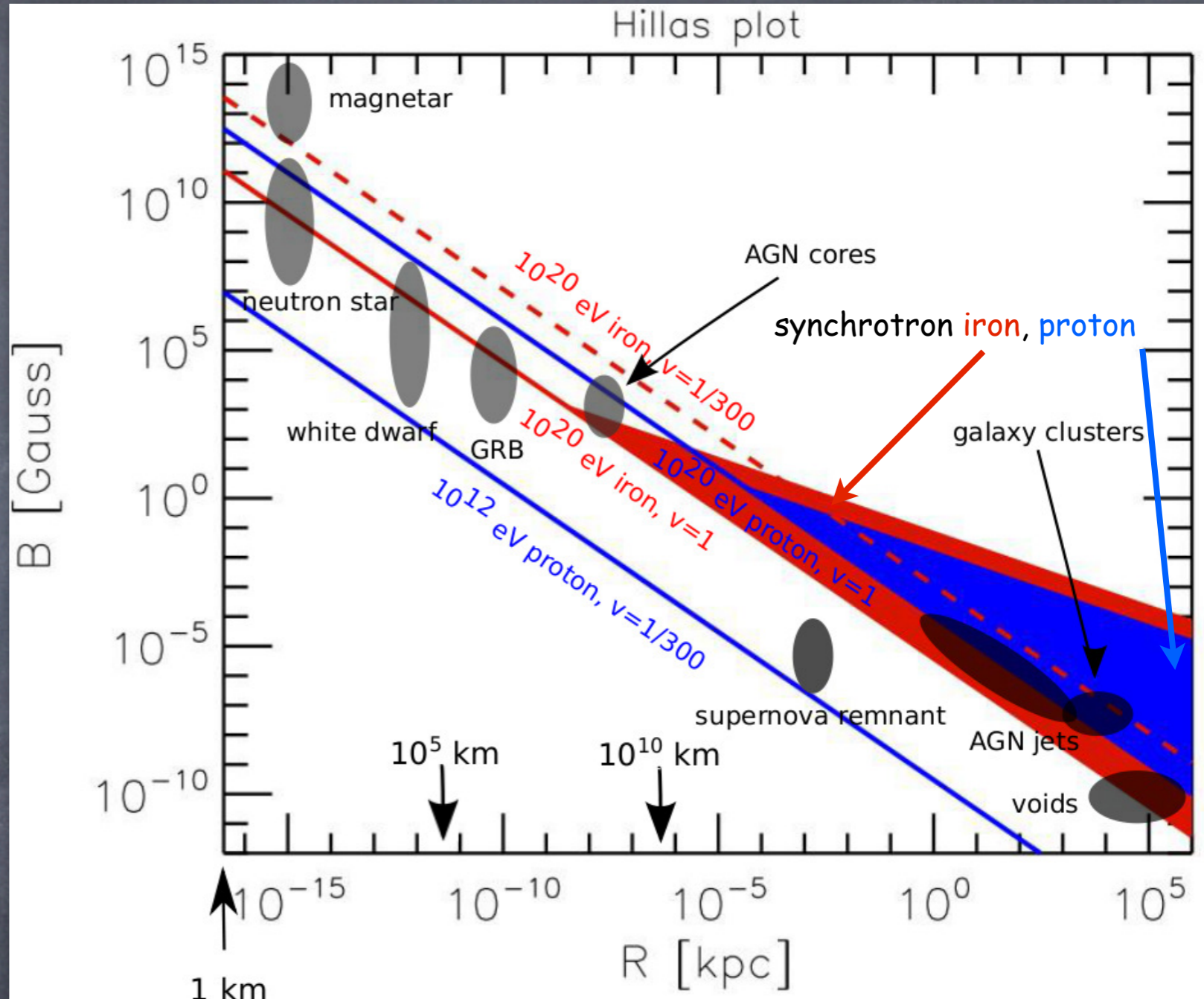
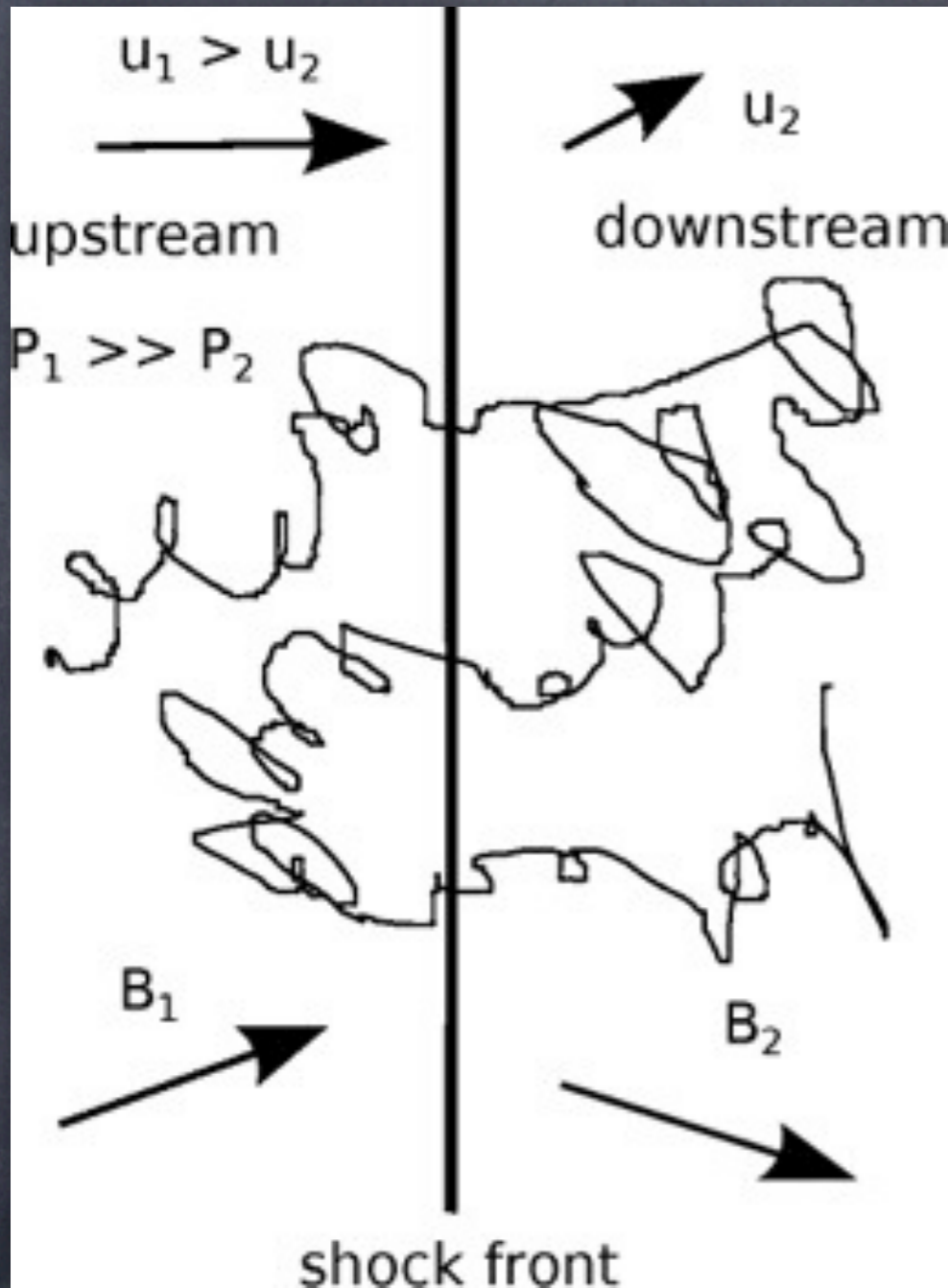
Length scales for relevant processes of a typical heavy nucleus



Interaction Horizons



1st Order Fermi Shock Acceleration



Fractional energy gain per shock crossing $\sim u_1 - u_2$ on a time scale r_L/u_2 .

Together with downstream losses this leads to a spectrum E^{-q} with $q > 2$ typically.

Confinement, gyroradius $<$ shock size, and energy loss times define maximal energy

Some general Requirements for Sources

Accelerating particles of charge eZ to energy E_{\max} requires induction $\epsilon > E_{\max}/eZ$. With $Z_0 \sim 100\Omega$ the vacuum impedance, this requires dissipation of minimum power of

$$L_{\min} \sim \frac{\epsilon^2}{Z_0} \simeq 10^{45} Z^{-2} \left(\frac{E_{\max}}{10^{20} \text{ eV}} \right)^2 \text{ erg s}^{-1}$$

This „Poynting“ luminosity can also be obtained from $L_{\min} \sim (BR)^2$ where BR is given by the „Hillas criterium“:

$$BR > 3 \times 10^{17} \Gamma^{-1} \left(\frac{E_{\max}/Z}{10^{20} \text{ eV}} \right) \text{ Gauss cm}$$

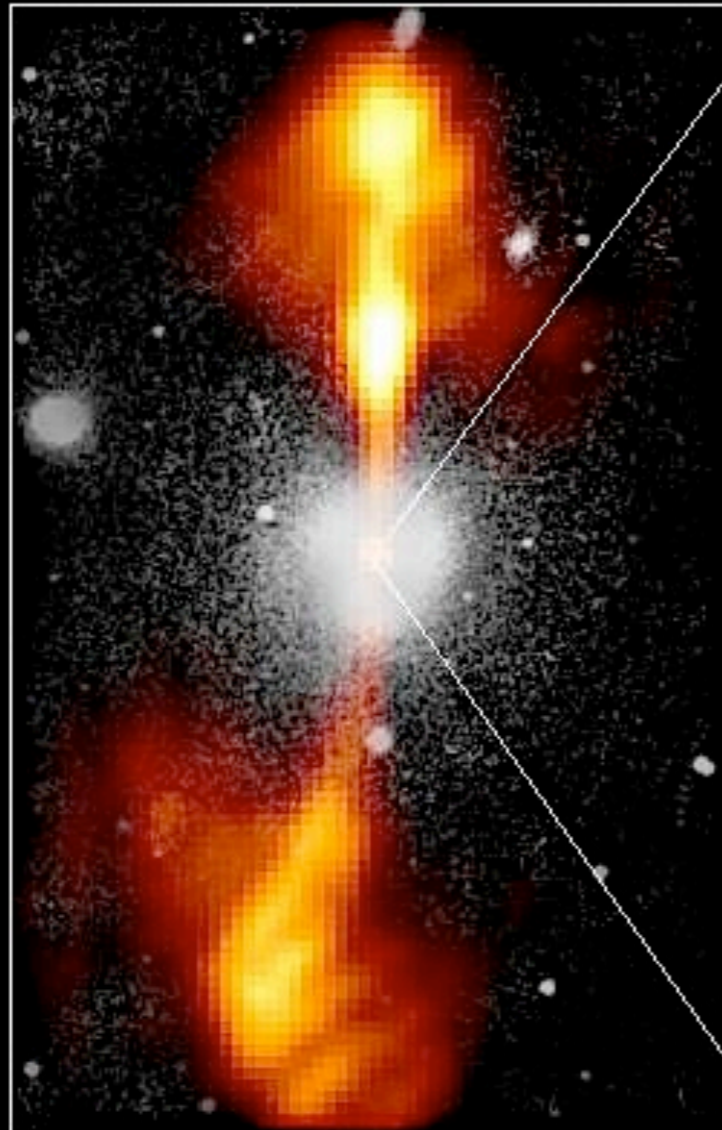
where Γ is a possible beaming factor.

If most of this goes into electromagnetic channel, only AGNs and maybe gamma-ray bursts could be consistent with this.

Core of Galaxy NGC 4261

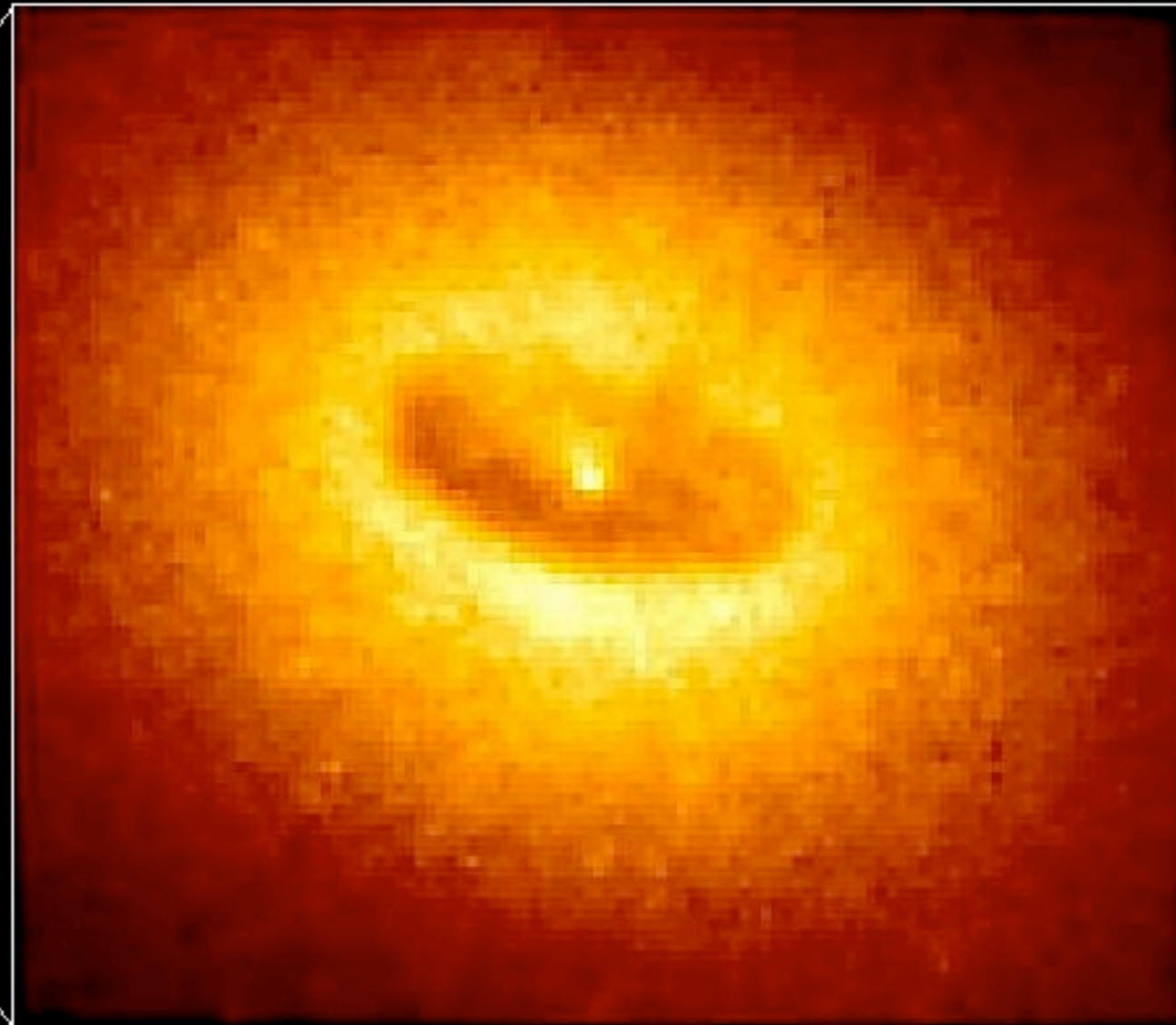
Hubble Space Telescope
Wide Field / Planetary Camera

Ground-Based Optical/Radio Image



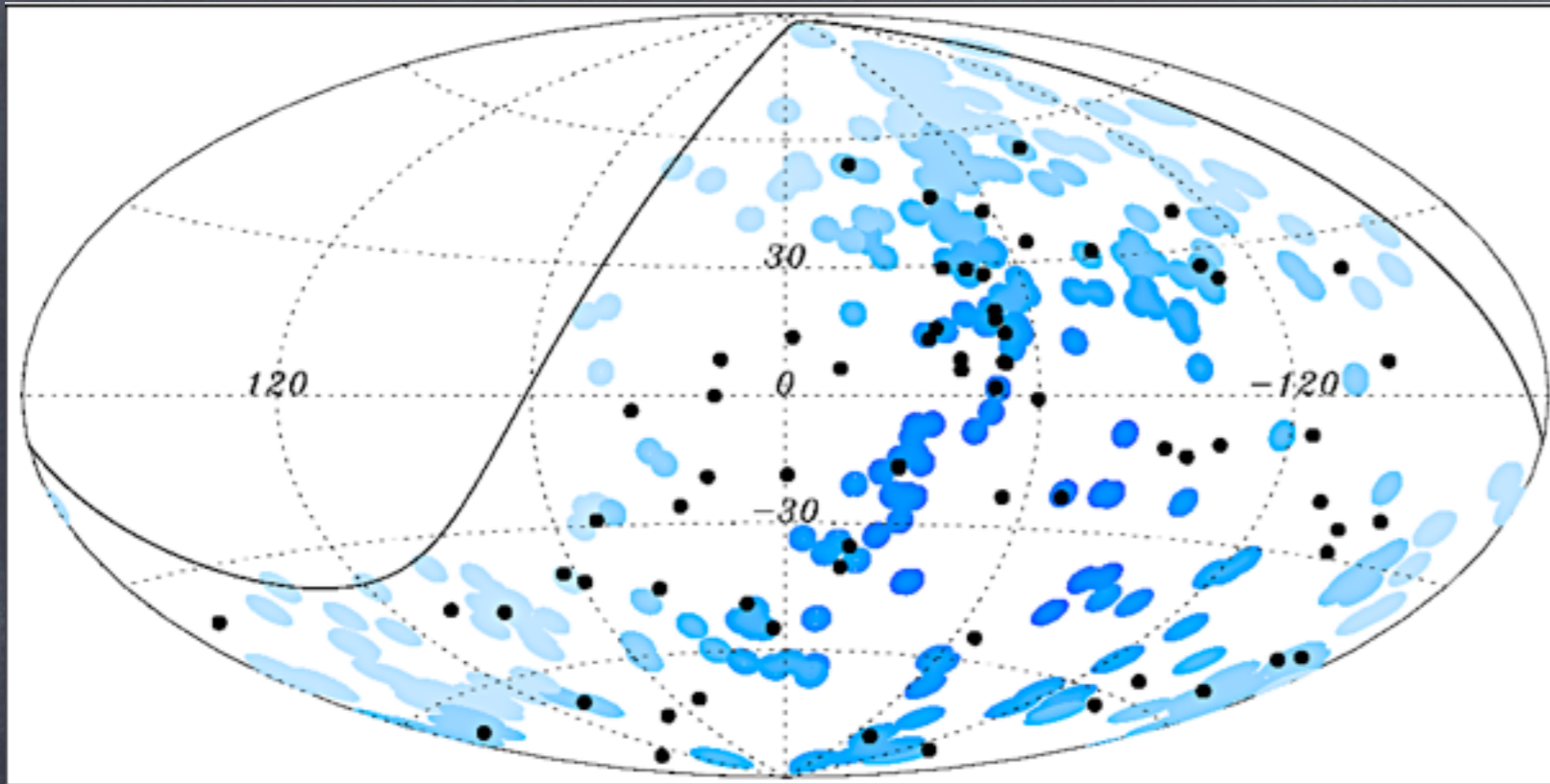
380 Arc Seconds
88,000 LIGHT-YEARS

HST Image of a Gas and Dust Disk

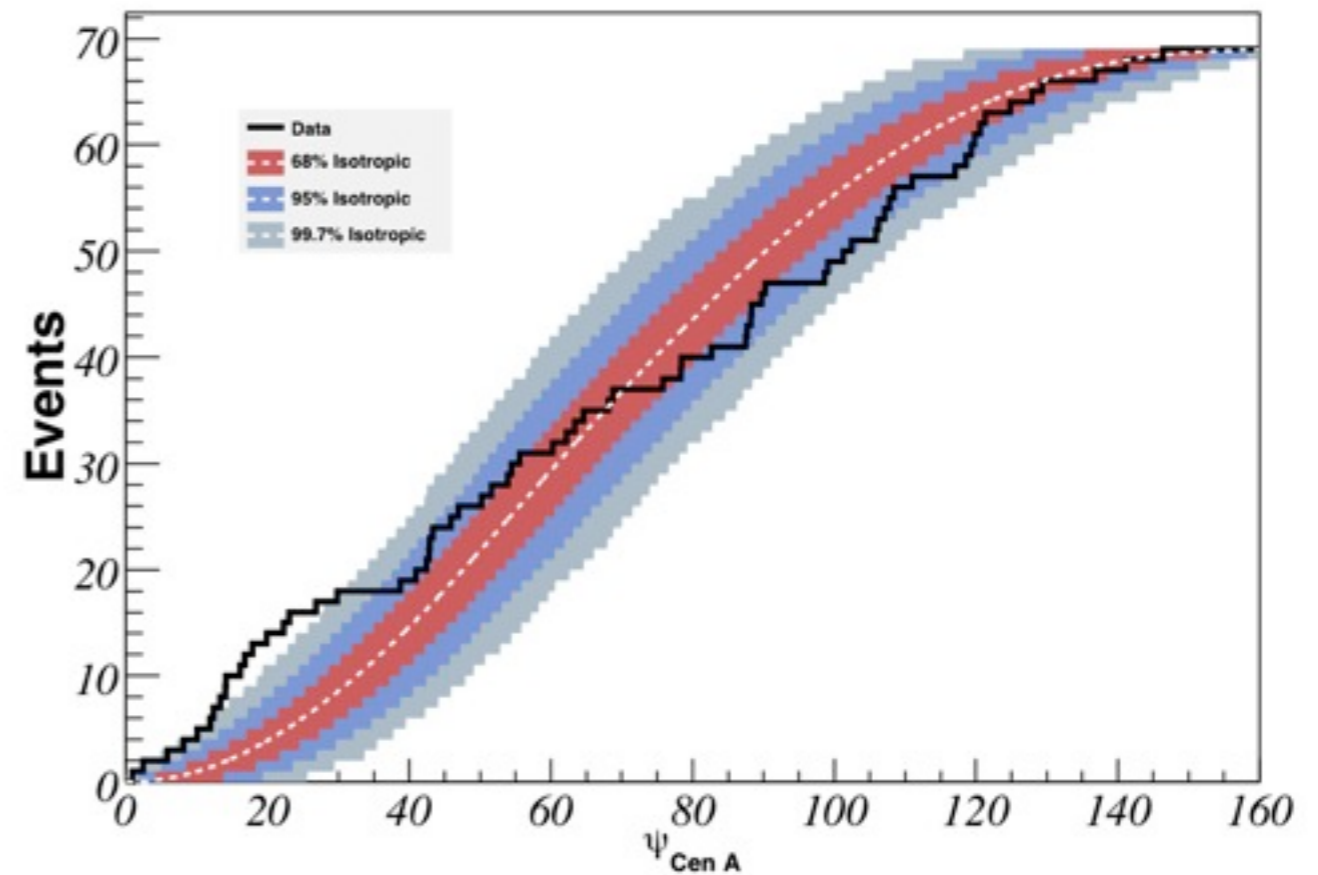


17 Arc Seconds
400 LIGHT-YEARS

Centaurus A is a UHECR source candidate

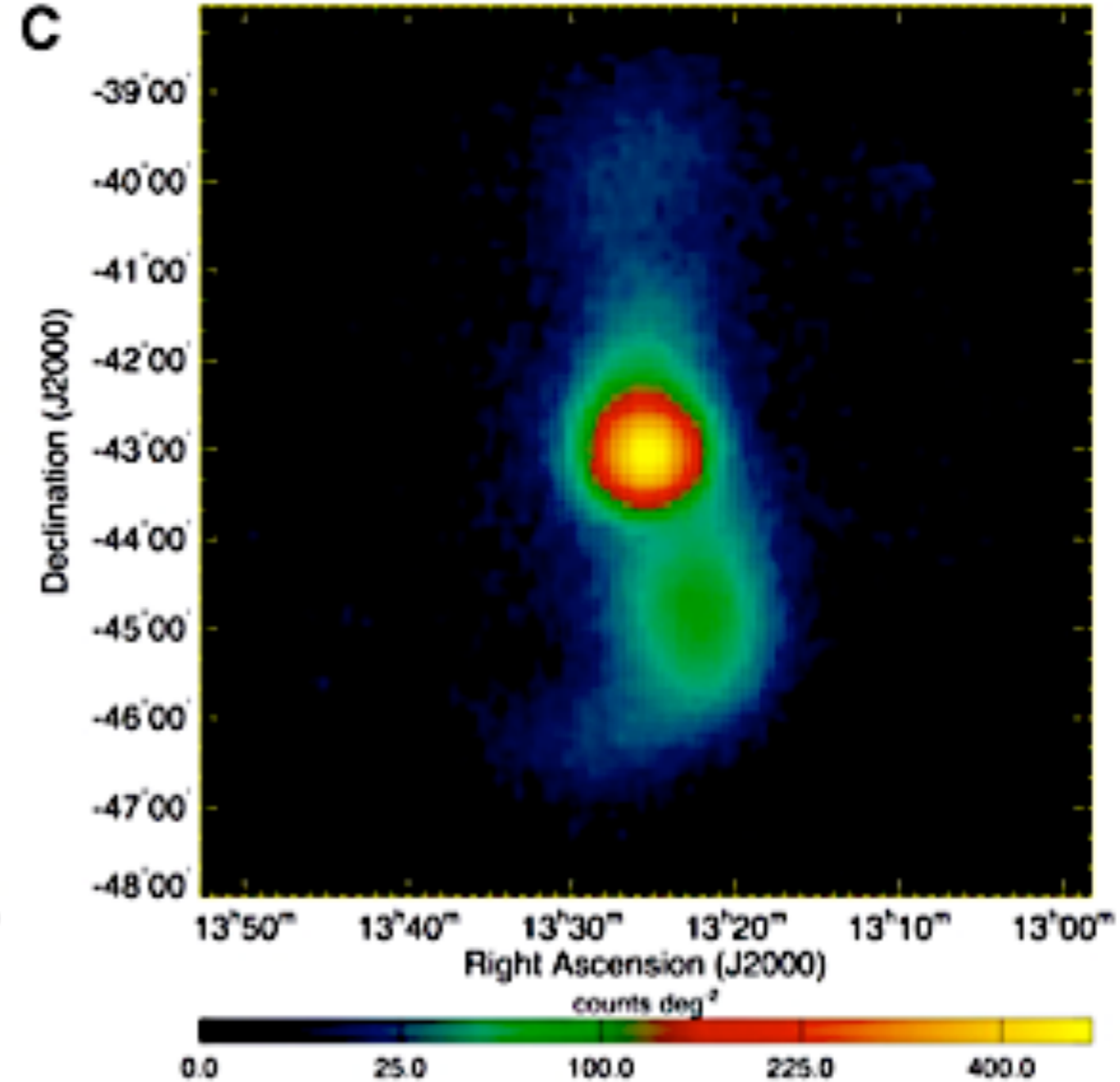
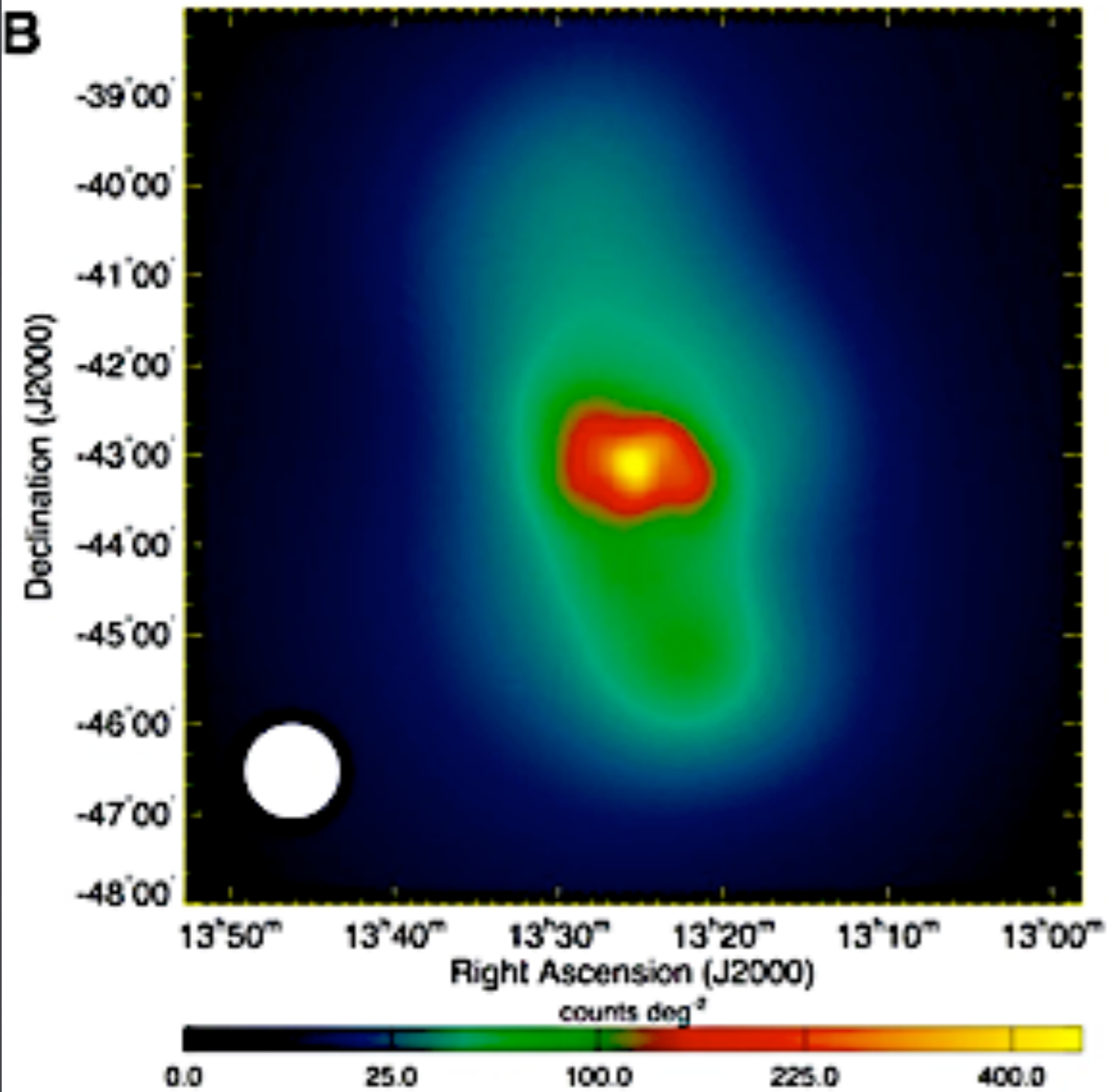


Pierre Auger sees an excess
in the direction of Centaurus A



Pierre Auger Collaboration, *Astropart.Phys.* 34 (2010) 314

Lobes of Centaurus A seen by Fermi-LAT

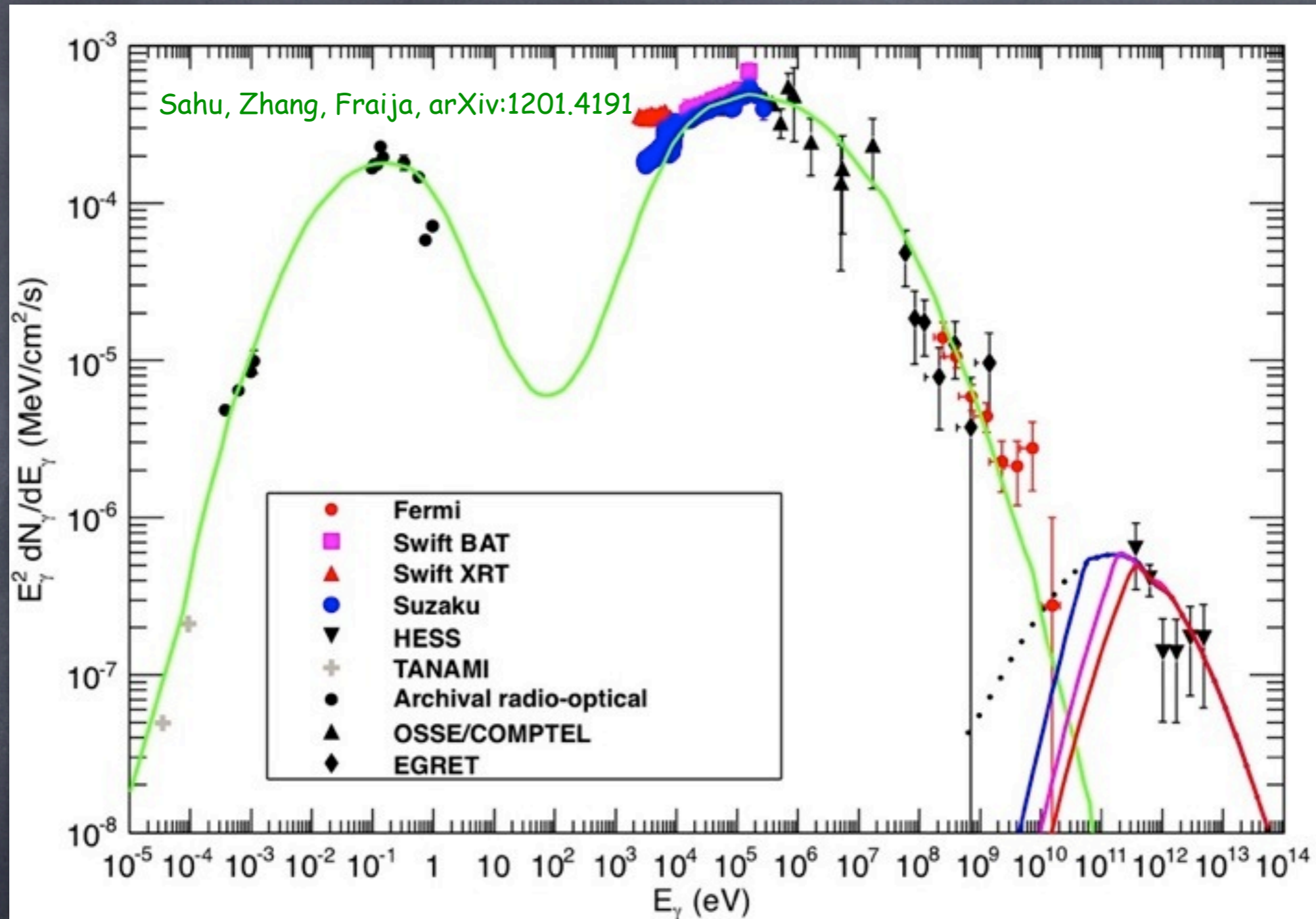


> 200 MeV γ -rays

Radio observations

Abdo et al., Science Express 1184656, April 1, 2010

Centaurus A as Multimessenger Source: A Mixed hadronic+leptonic Model



Low energy bump = synchrotron
high energy bump = synchrotron self-Compton
TeV- γ -rays: $p\gamma$ interactions of shock-accelerated protons

Mass Composition

Depth of shower maximum and its distribution contain information on primary mass composition

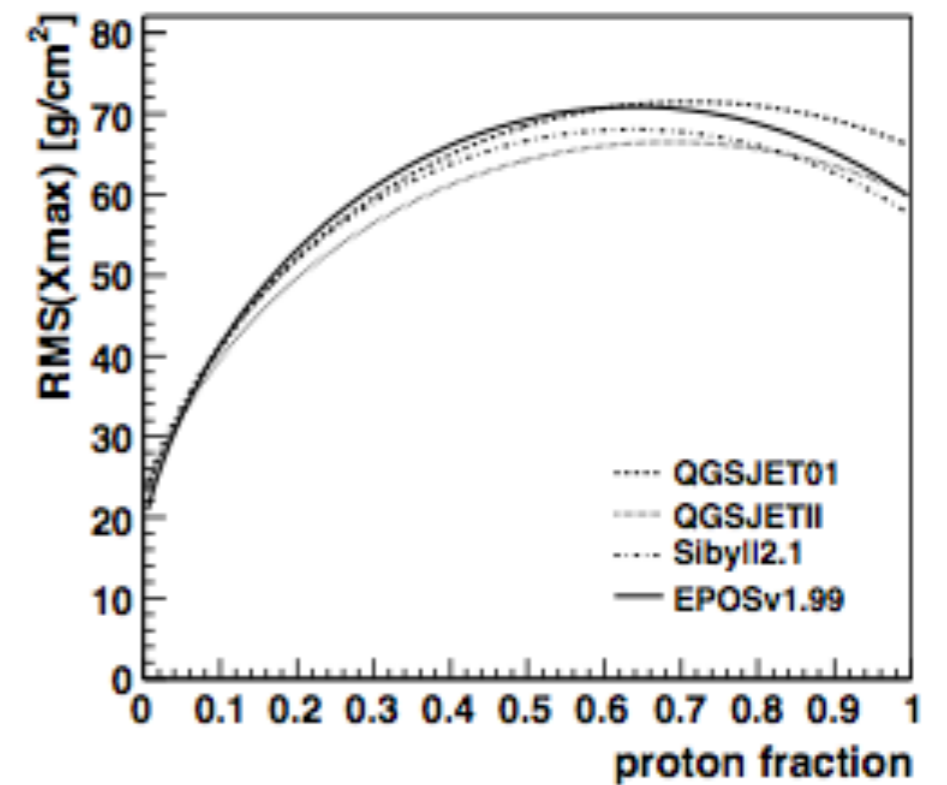
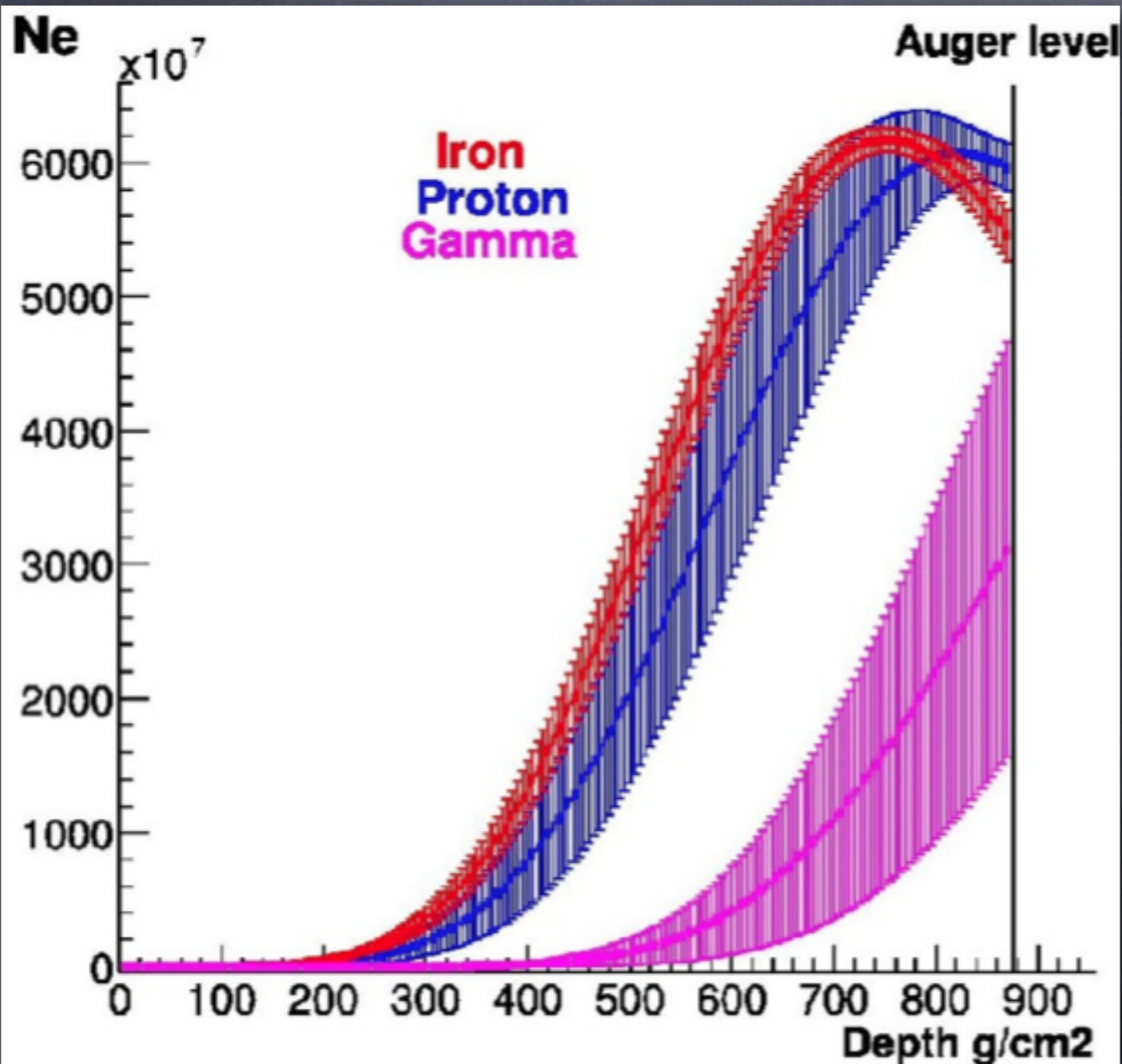
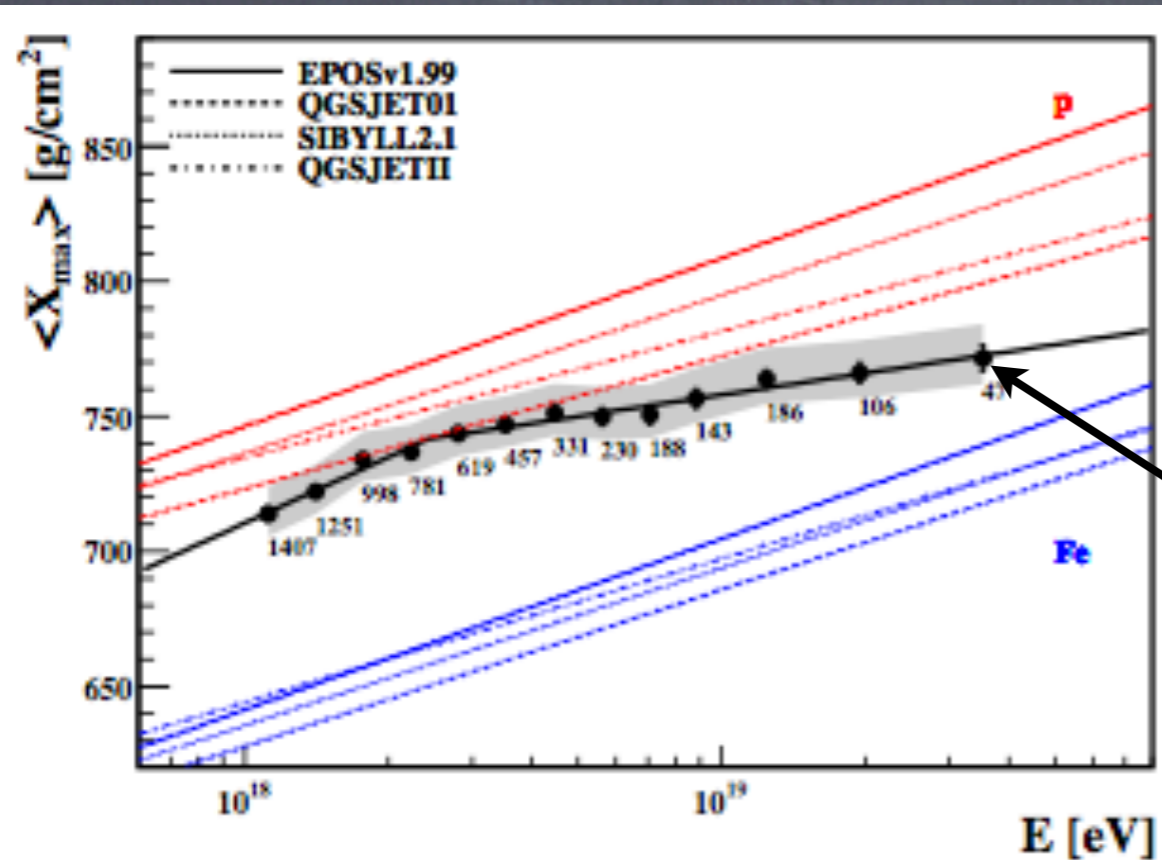


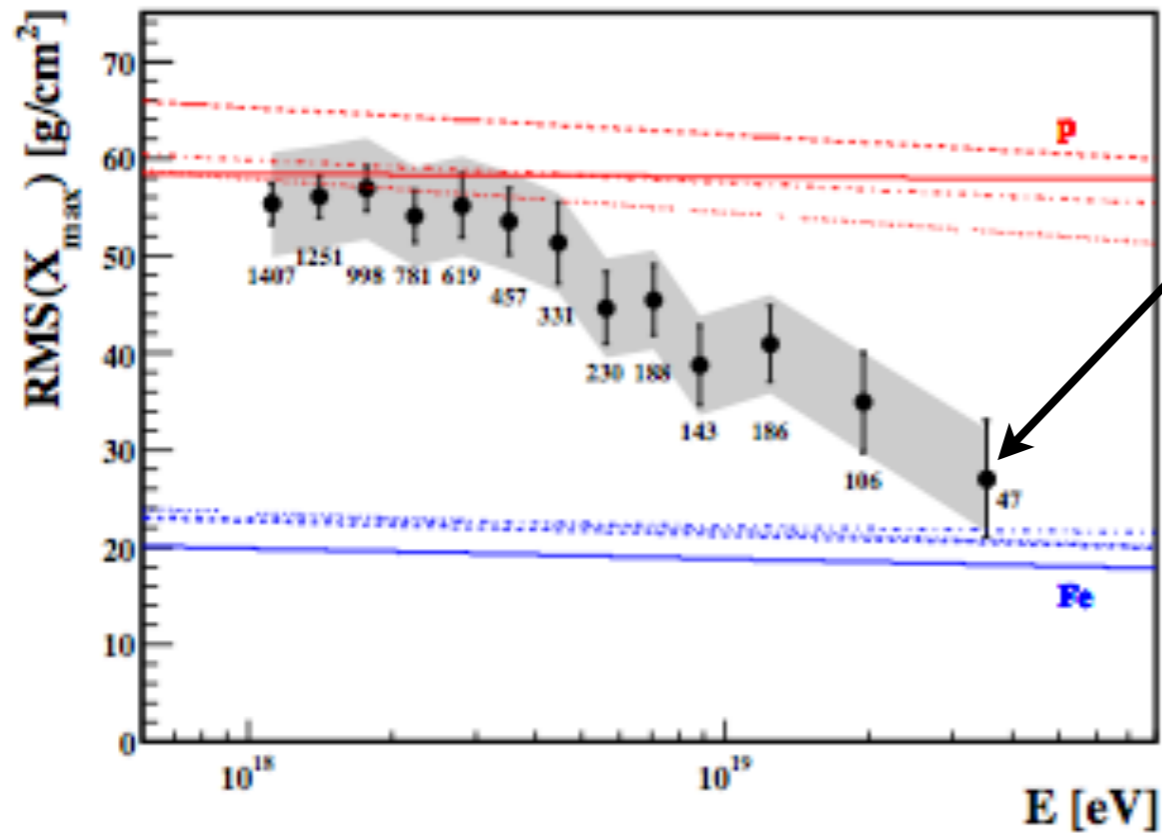
FIGURE 1. $RMS(X_{max})$ from different hadronic interaction models [23] and a two-component p/Fe composition model ($E = 10^{18}$ eV).

Pierre Auger data suggest a heavier composition toward highest energies:

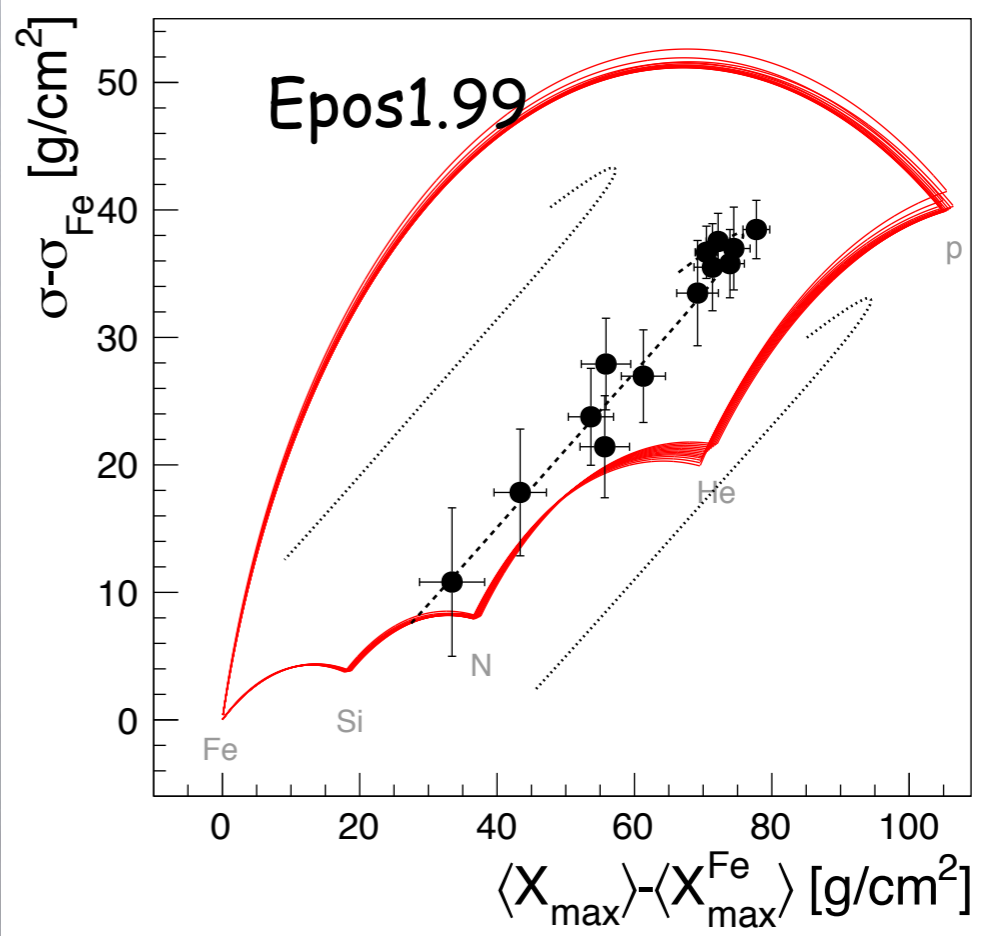
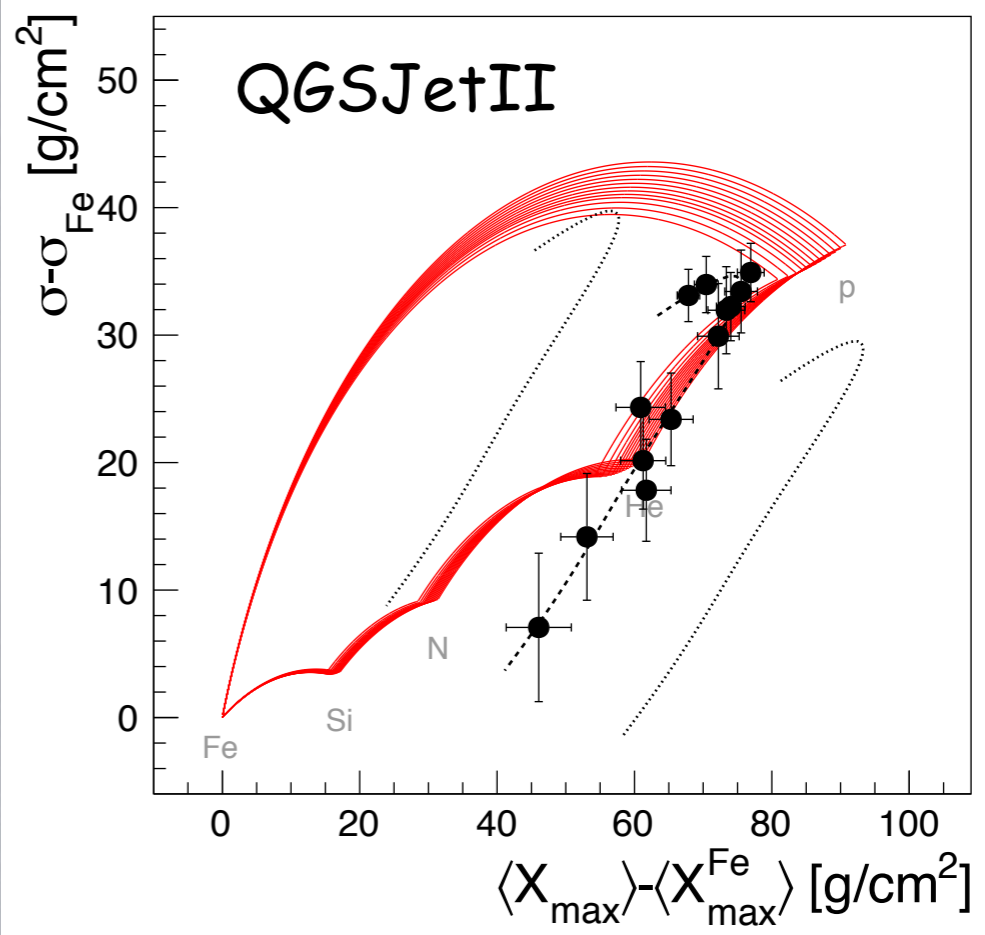
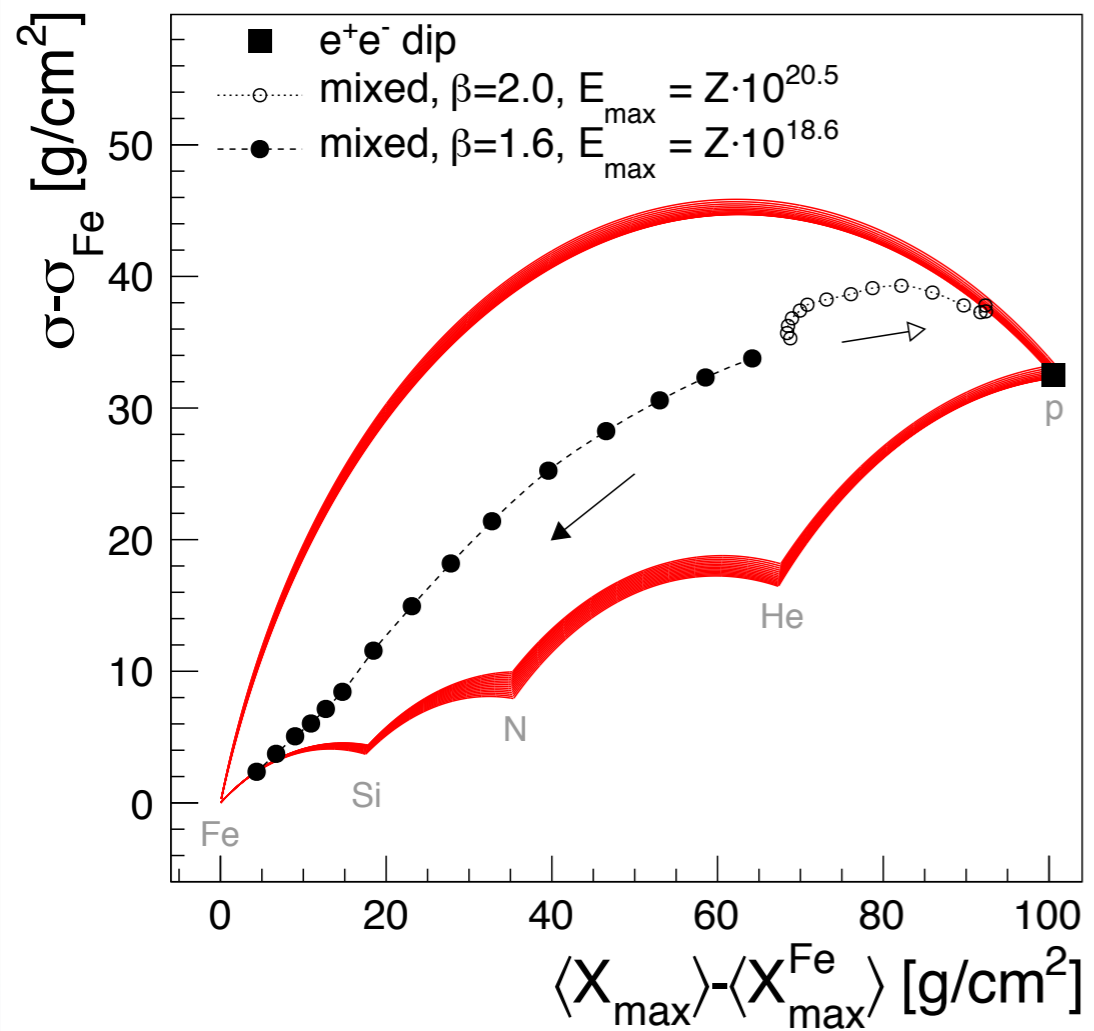


but not confirmed on the northern hemisphere by HiRes and Telescope Array which are consistent with protons

potential tension with air shower simulations and some hadronic interaction models because a mixed composition would predict larger $\text{RMS}(X_{\text{max}})$



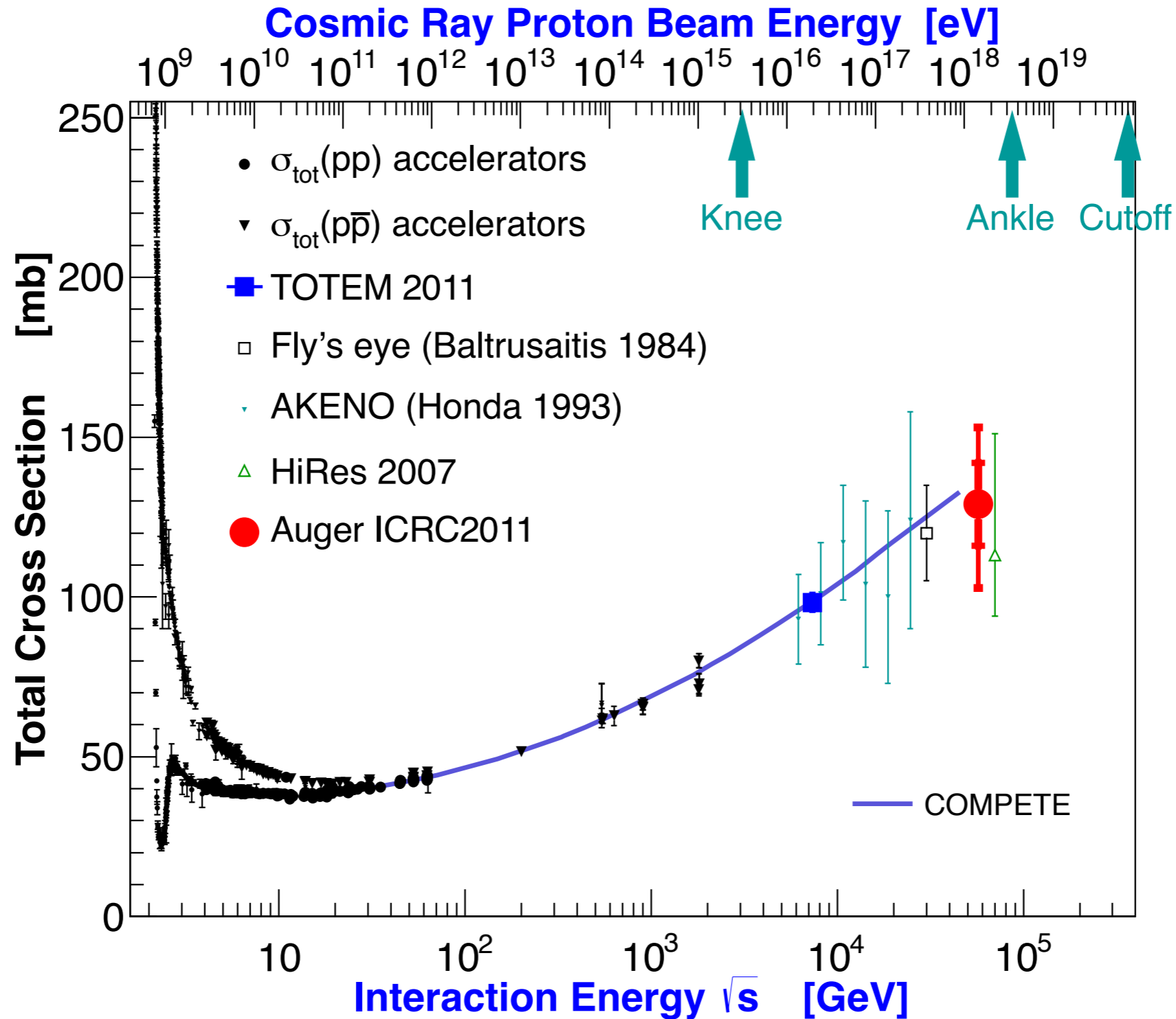
Pierre Auger Collaboration,
Phys.Rev.Lett., 104 (2010) 091101,
and ICRC 2011, arXiv:1107.4804



combined measurement of X_{\max} and its fluctuation σ constrains composition within a given hadronic interaction model

Kampert and Unger, arXiv:1201.0018

Hadronic Cross-Sections

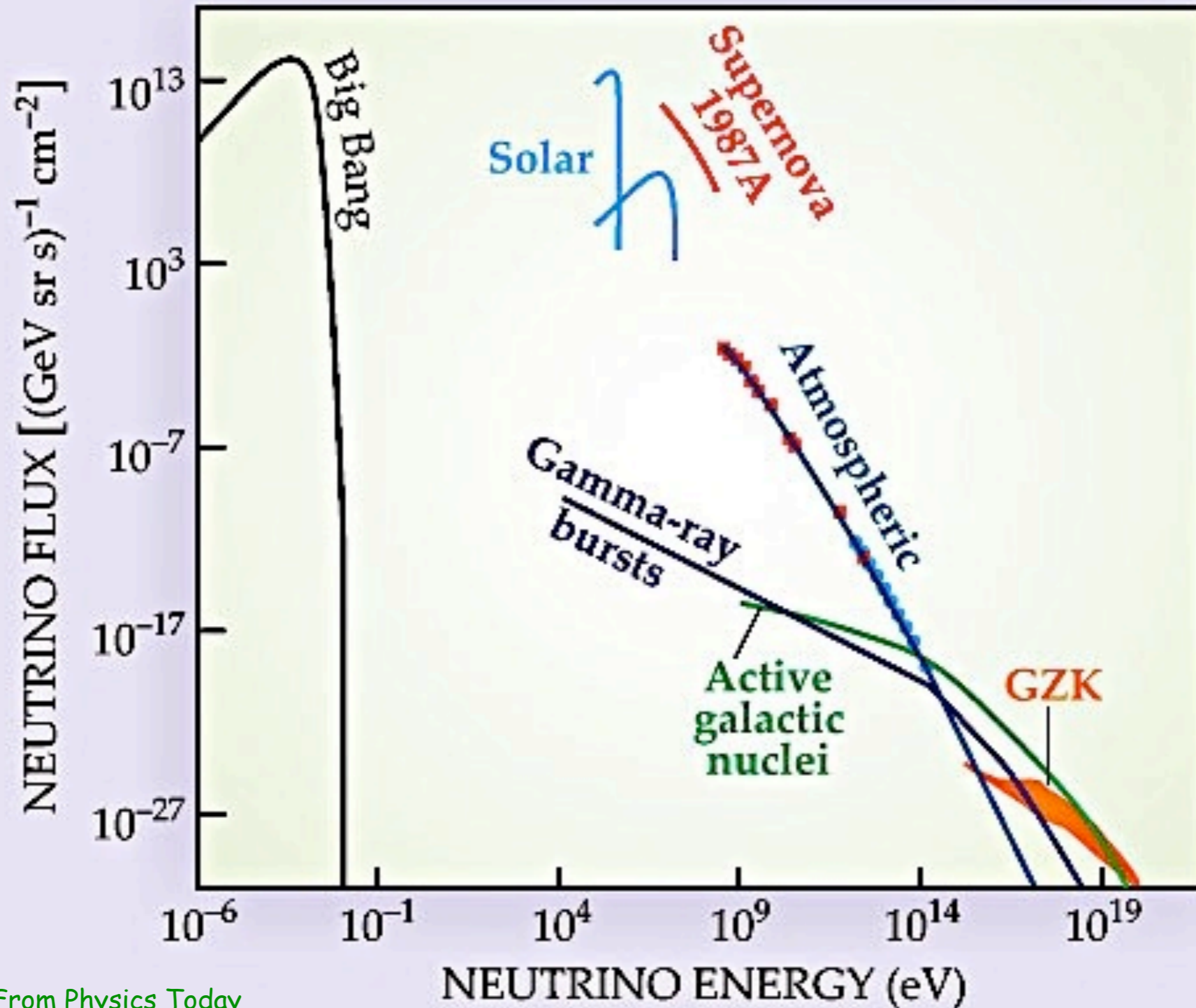


Ralf Ulrich

The high-energy frontier. No indications of missing physics.

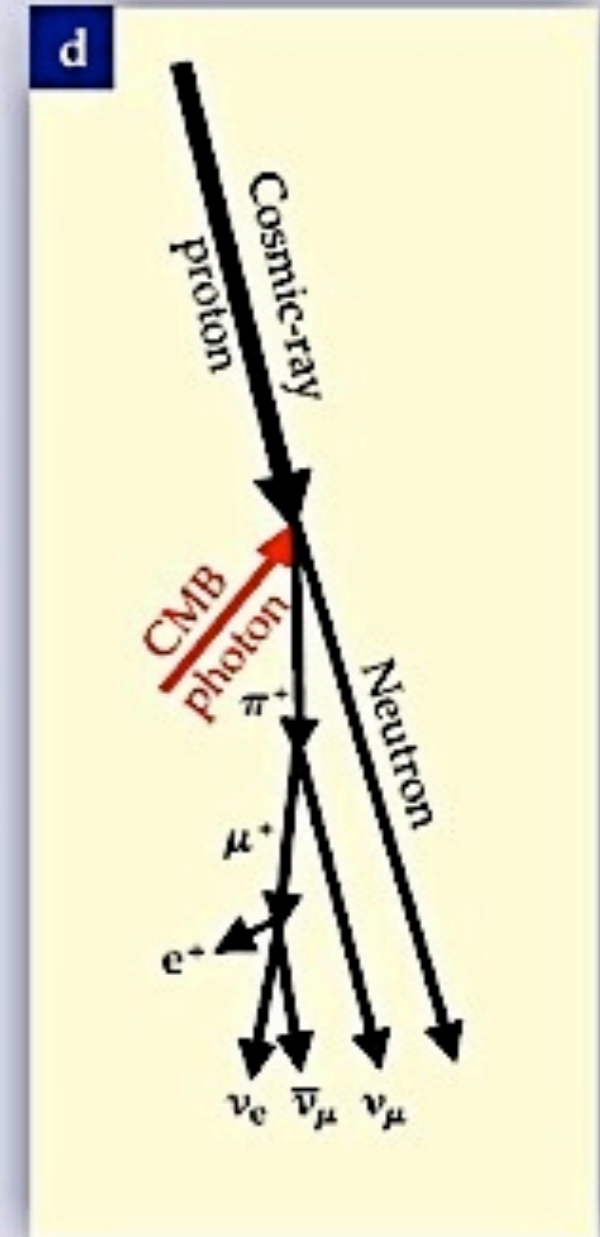
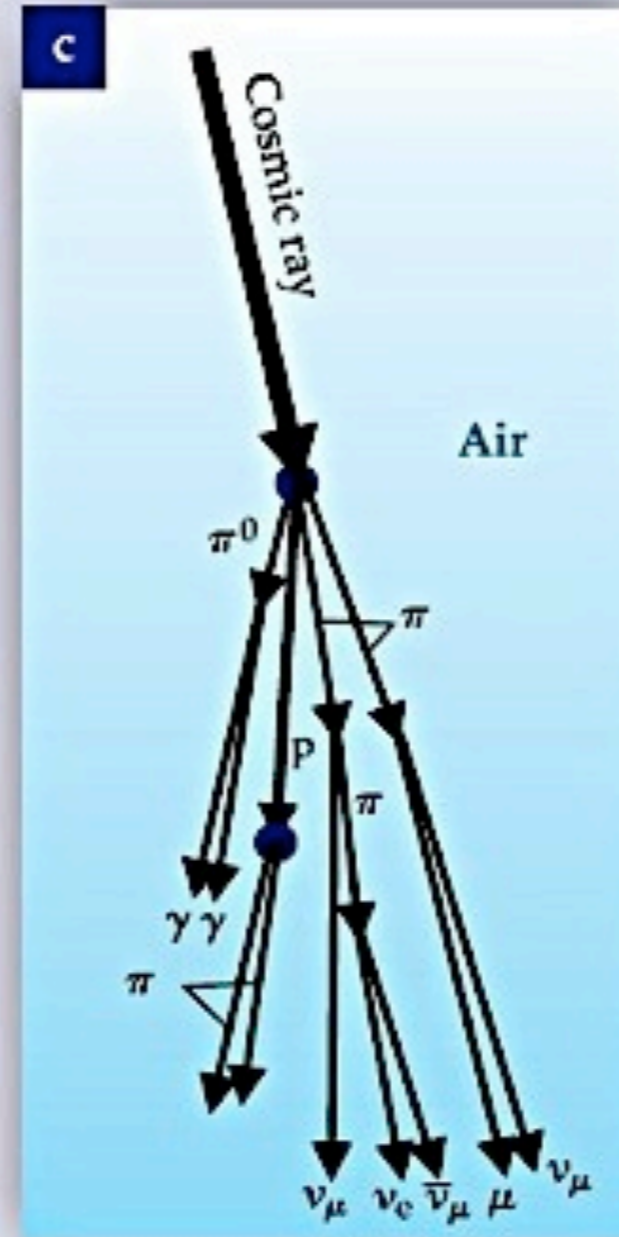
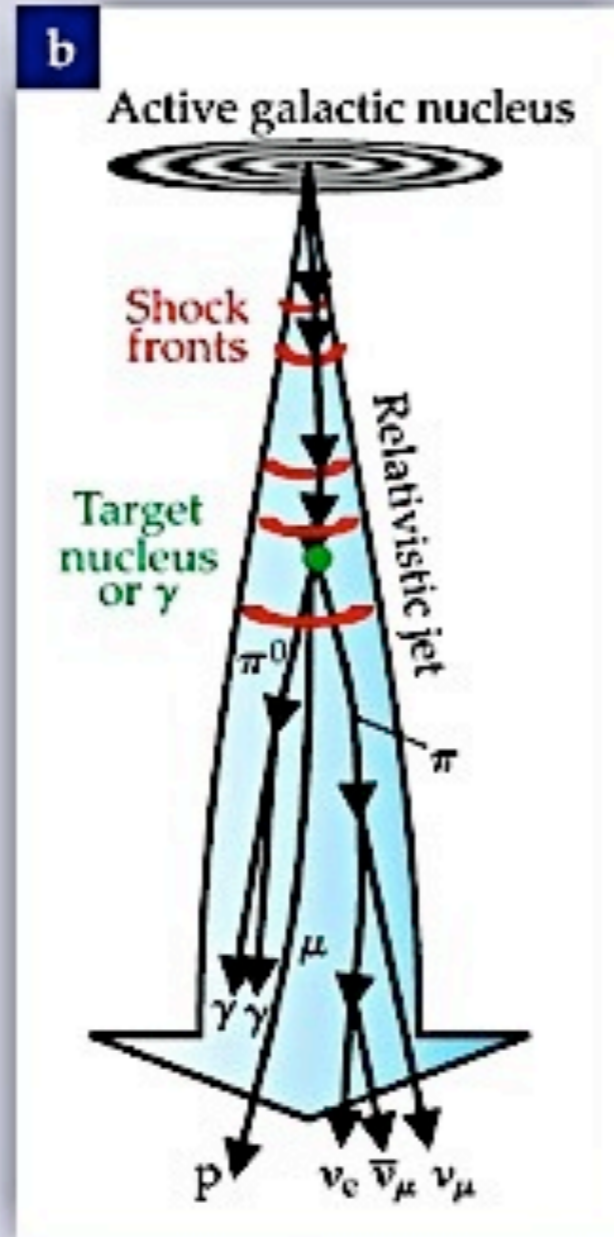
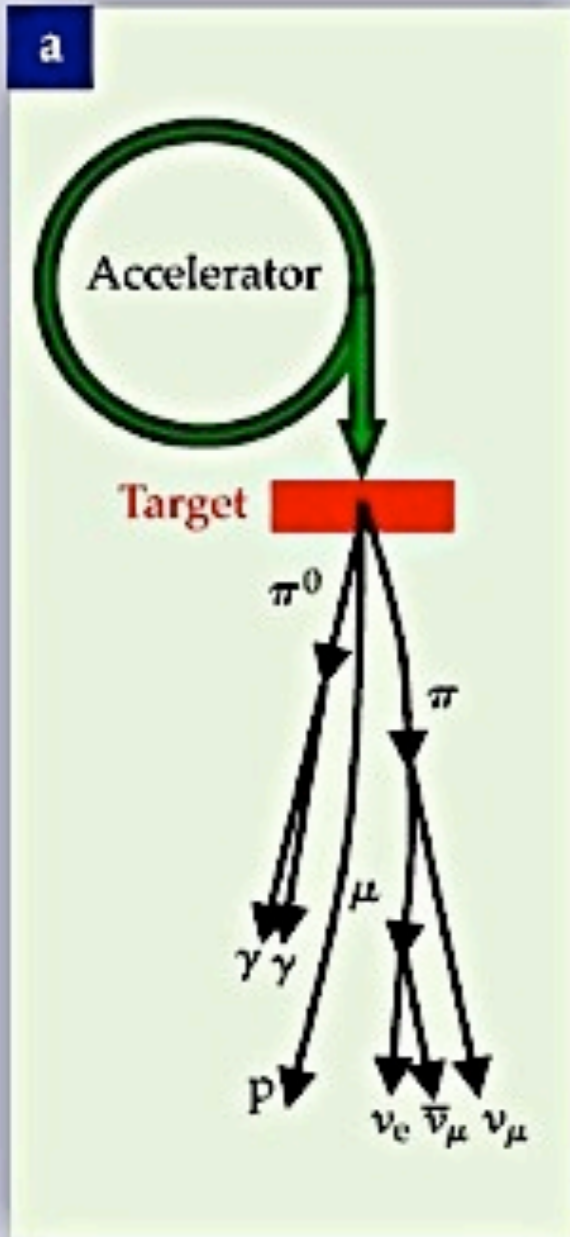
Very High High Energy Neutrinos

The „grand unified“ differential neutrino number spectrum



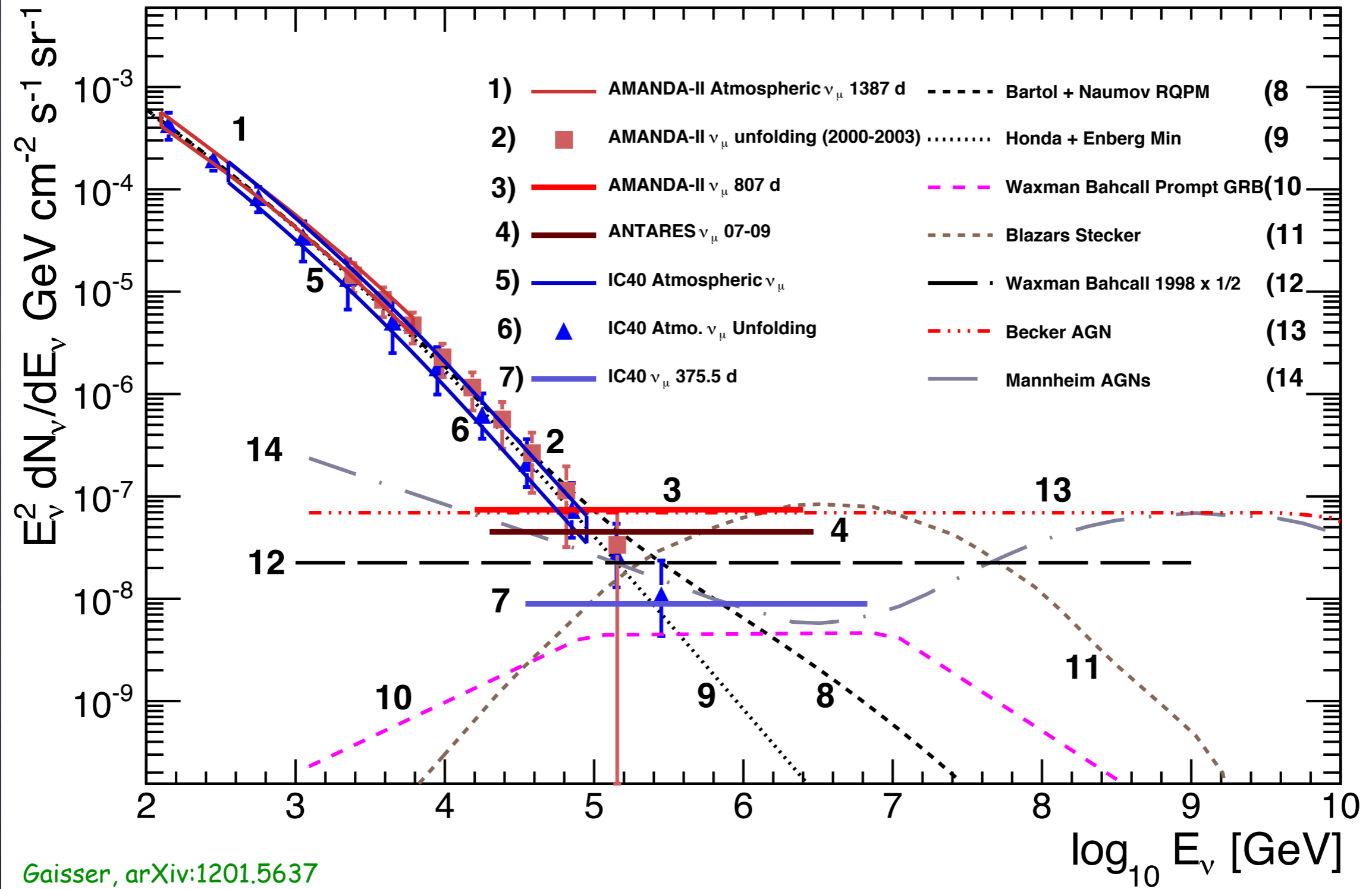
From Physics Today

Summary of neutrino production modes

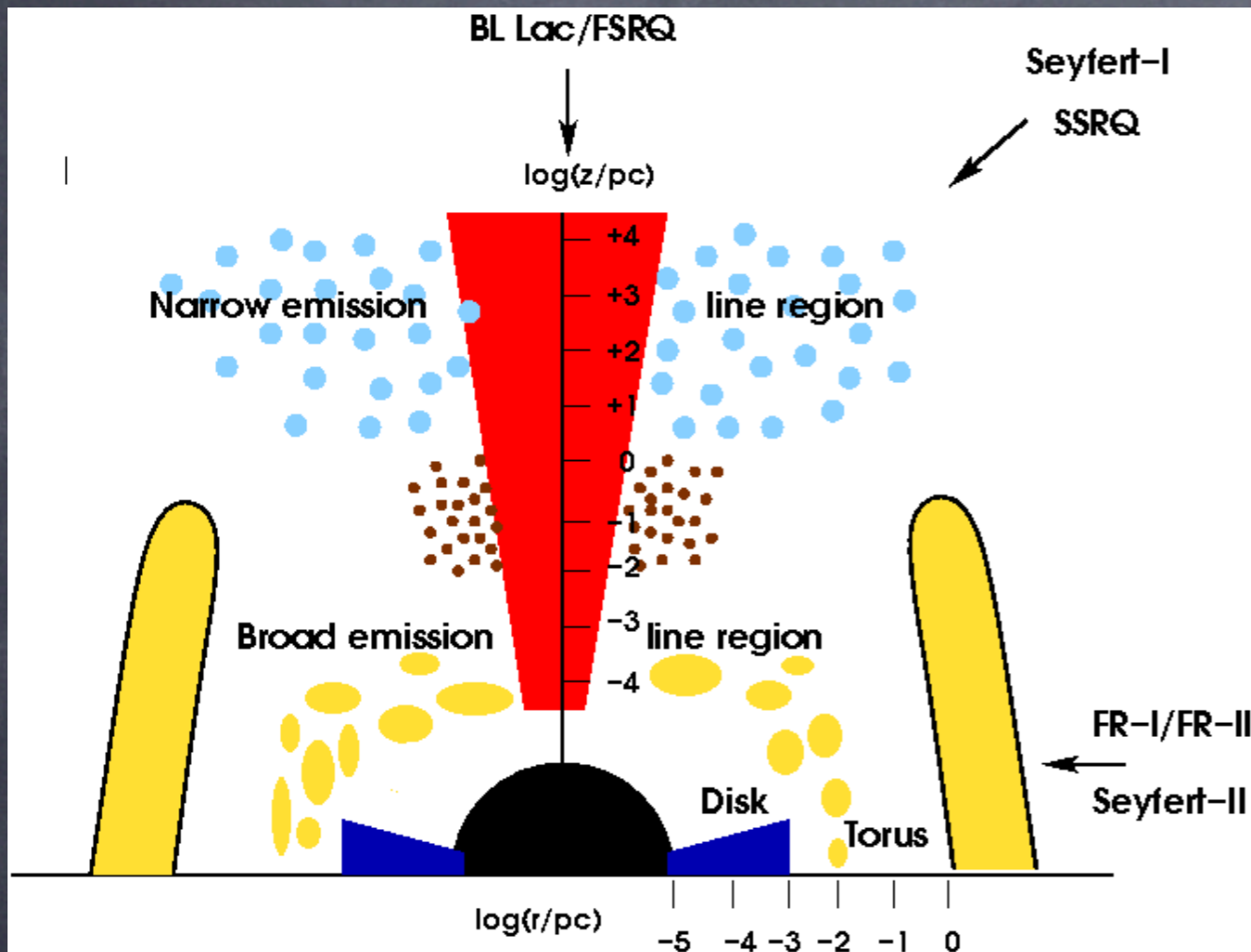


From Physics Today

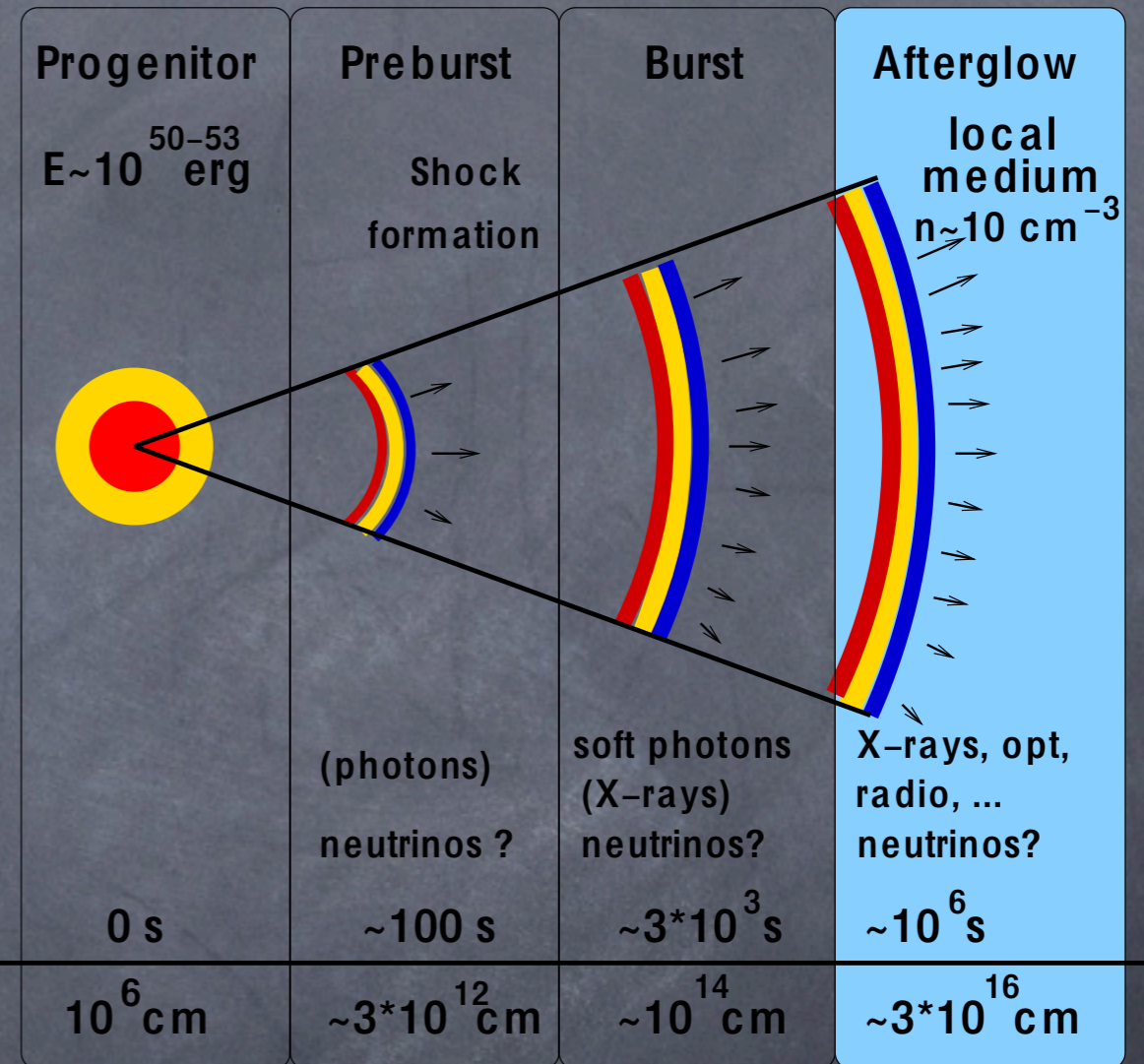
Current Neutrino Flux Upper Limits at TeV-EeV energies



Discrete Extragalactic High Energy Neutrino Sources



active galaxies



gamma ray bursts

Figures from J. Becker, Phys.Rep. 458 (2008) 173

Neutrino Fluxes from Gamma-Ray Bursts

GRBs are optically thick to charged cosmic rays and nuclei are disintegrated
=> only neutrons escape and contribute to the UHECR flux by decaying back
into protons

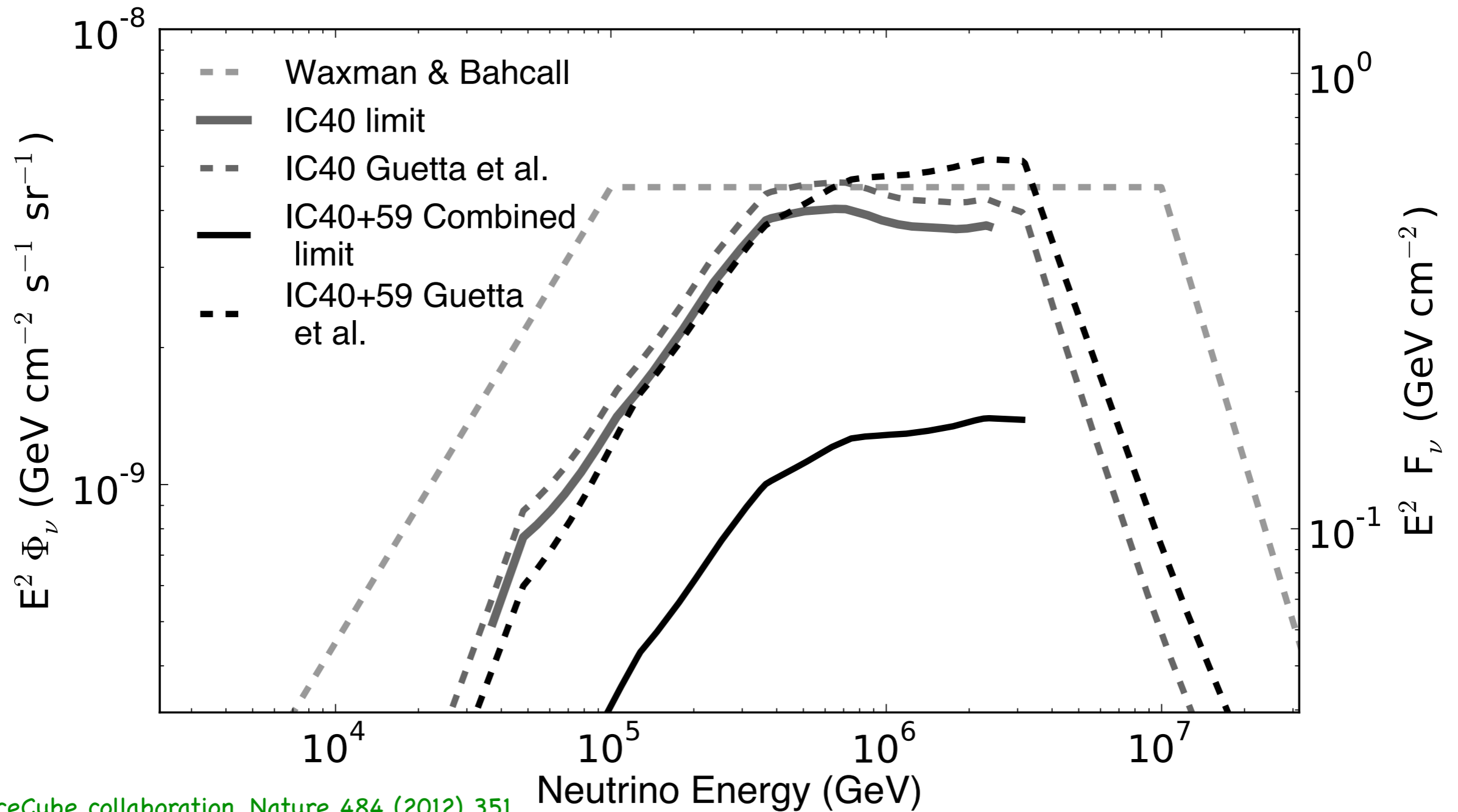
Diffuse neutrino flux from GRBs can thus be linked to UHECR flux (if it is
dominantly produced by GRBs)

$$\Phi_\nu(E_\nu) \sim \frac{1}{\eta_\nu} \Phi_p \left(\frac{E}{\eta_\nu} \right),$$

where $\eta_\nu \simeq 0.1$ is average neutrino energy in units of the parent proton energy.

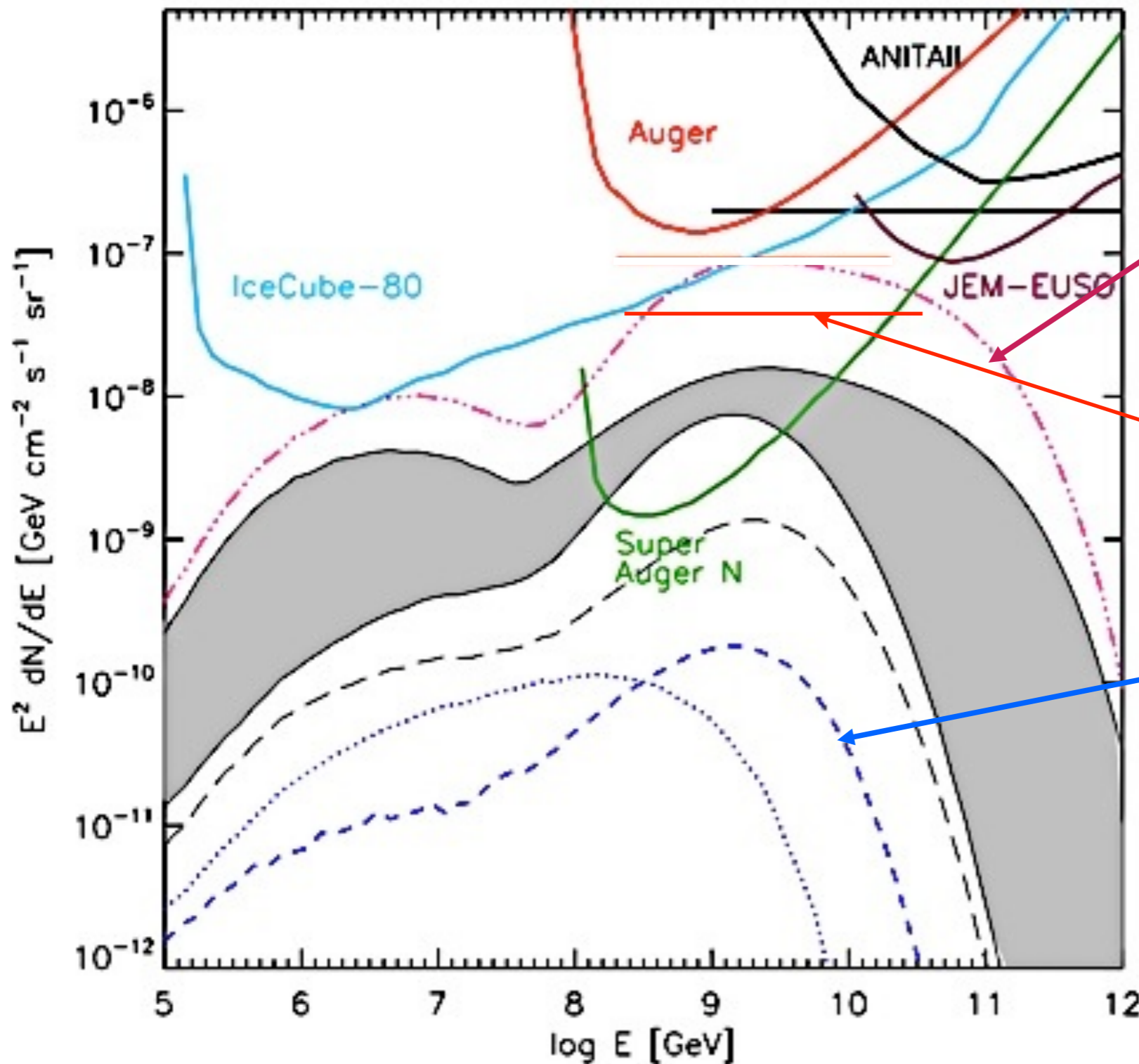
Above $\sim 10^{17}$ eV neutrino spectrum is steepened by one power of E_ν because pions/
muons interact before decaying

GRBs as UHECR sources now strongly challenged by non-observation of neutrinos by IceCube



Physics with Diffuse Cosmogenic Neutrino Fluxes

Cosmogenic neutrino fluxes depend on number of nucleons produced above GZK threshold which is proportional to E_{\max}/A
Further suppressed for heavy nuclei due to increased pair production



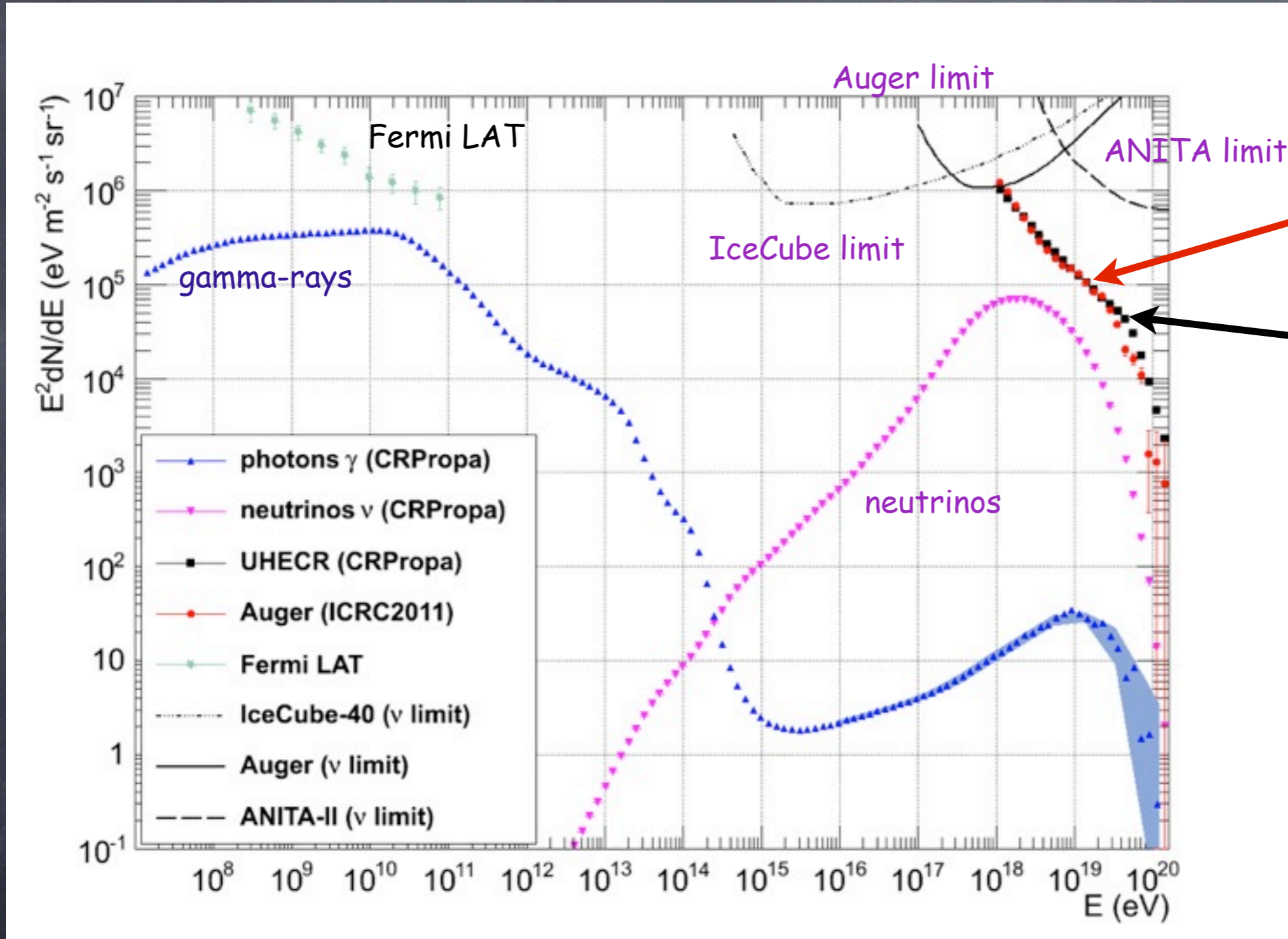
Pure protons, $E_{\max} = 3 \cdot 10^{21}$ eV,
strong evolution

current Auger limit on
Earth-skimming neutrinos

Pure iron, $E_{\max} = 10^{20}/26$ eV,
no evolution

Kotera, Allard, Olinto, JCAP 1010 (2010) 013

TeV γ -ray fluxes also constrain cosmogenic neutrino fluxes
sensitive to redshift evolution; complementary to UHE γ -ray fluxes

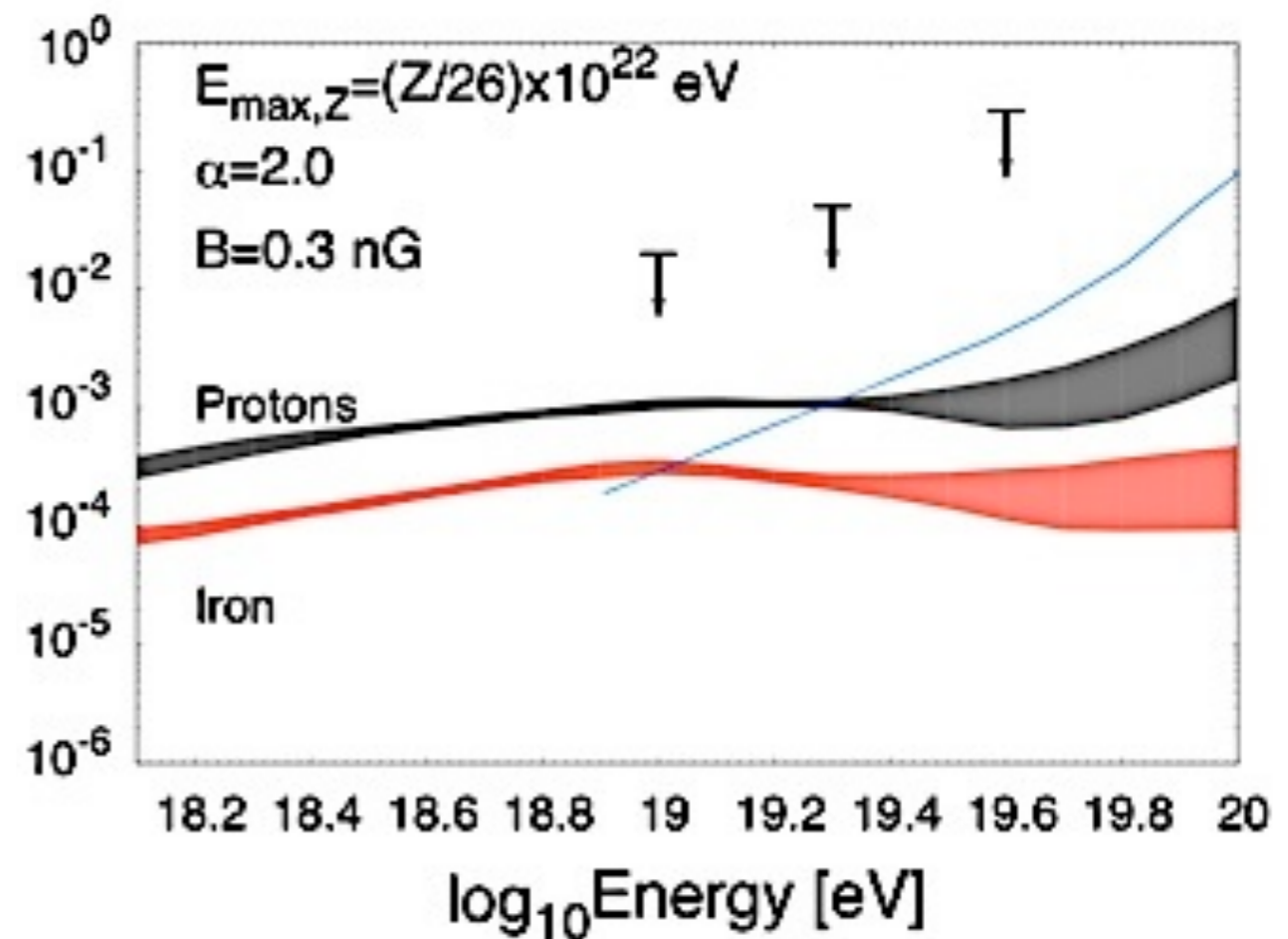
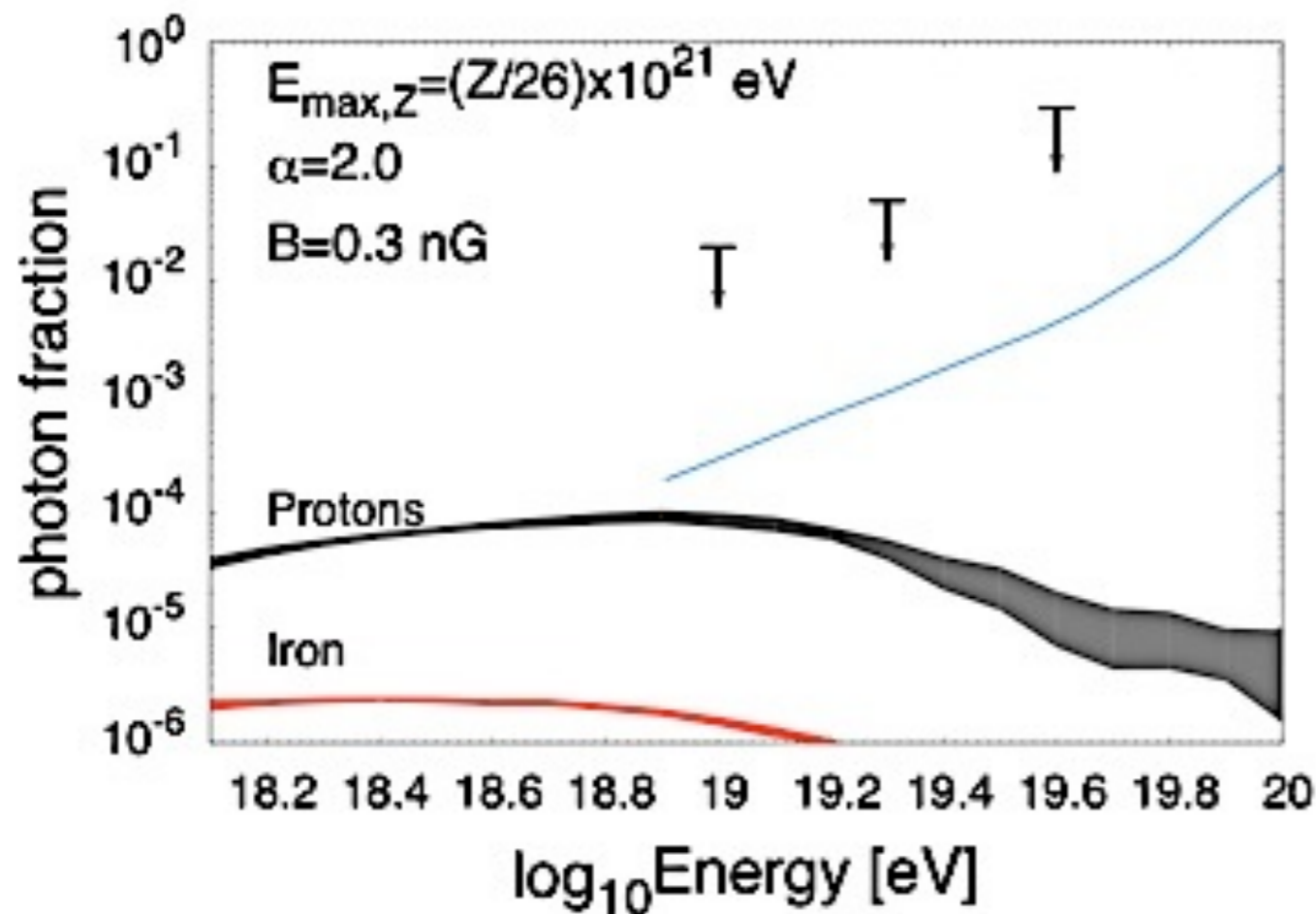


comoving injection rate scaling as $(1+z)^4$ up to $z=2$ injection spectral slope $\alpha=2.2$
up to $E_{\text{max}} = Z \times 3.86 \times 10^{20}$ eV for a galactic composition at the sources

Physics with Diffuse Secondary Gamma-Ray Fluxes

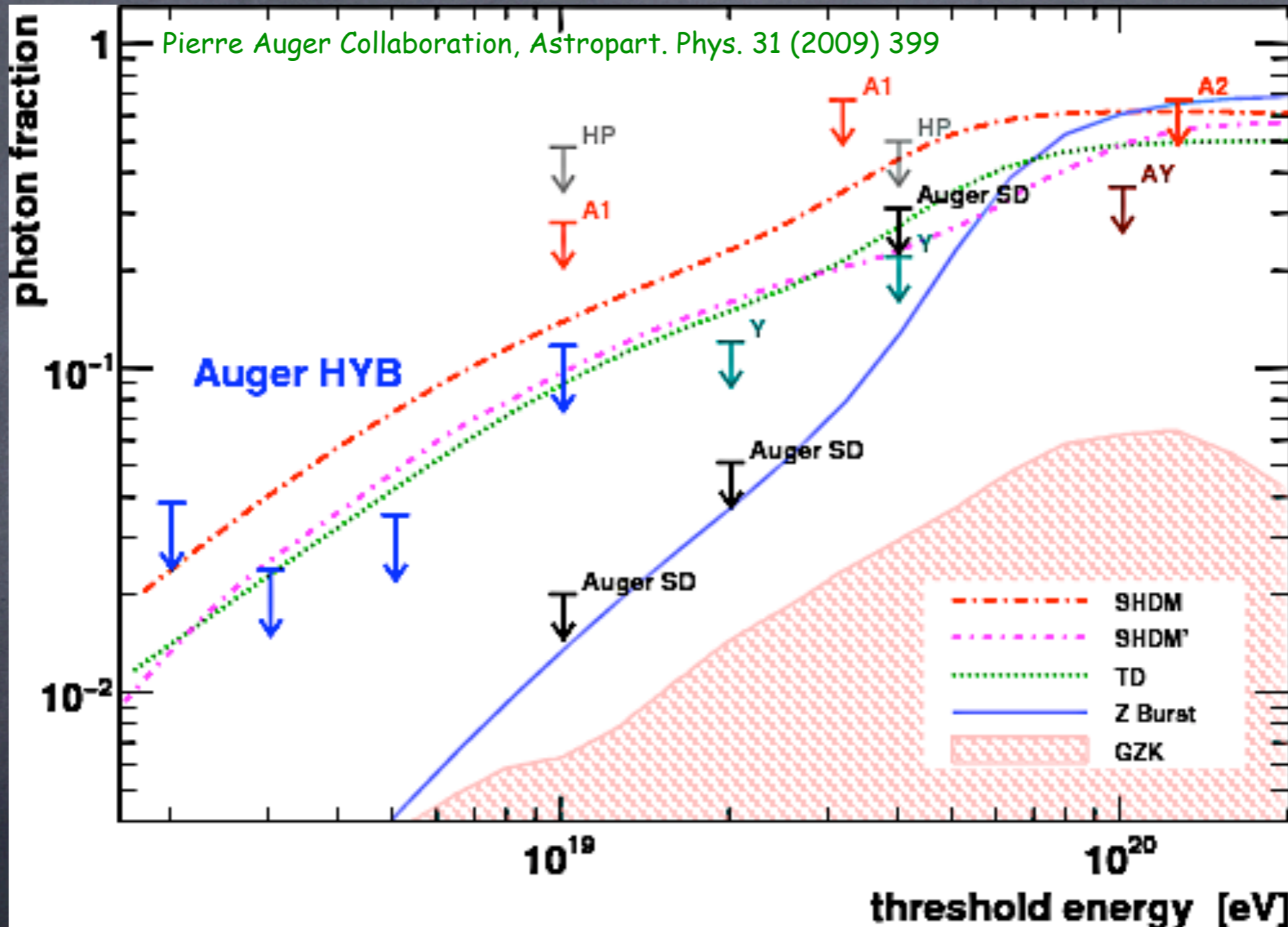
UHE gamma-ray fluxes depend on number of nucleons locally produced above GZK threshold which is proportional to E_{\max}/A

Further suppressed for heavy nuclei due to increased pair production
complementary to cosmogenic neutrinos: does not depend on redshift evolution

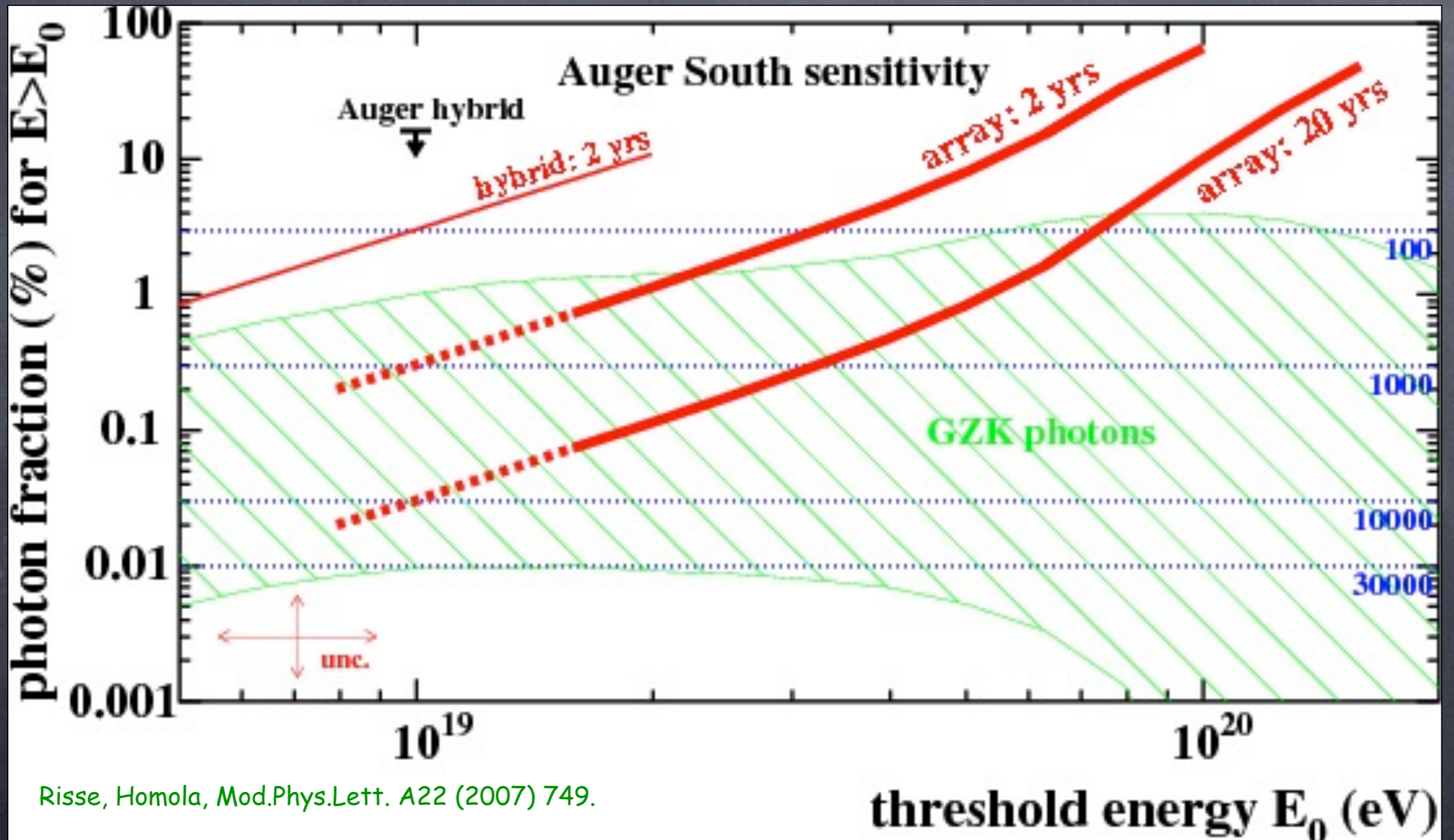


Hooper, Taylor, Sarkar, *Astropart.Phys.* 34 (2011) 340

Current upper limits on the photon fraction are of order 2% above 10^{19} eV from latest results of the Pierre Auger experiments (ICRC) and order 30% above 10^{20} eV.



Future data will allow to probe smaller photon fractions and the GZK photons

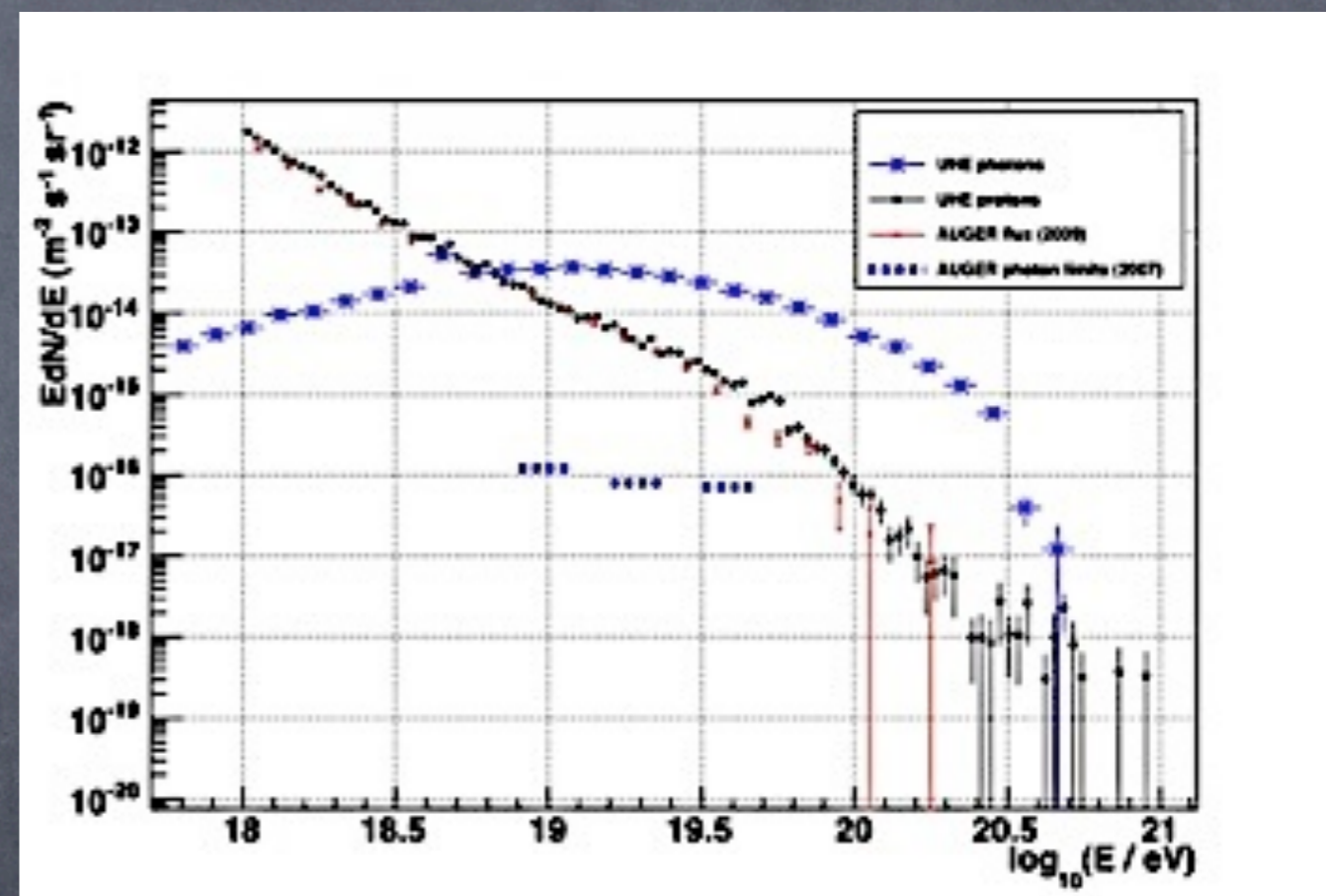
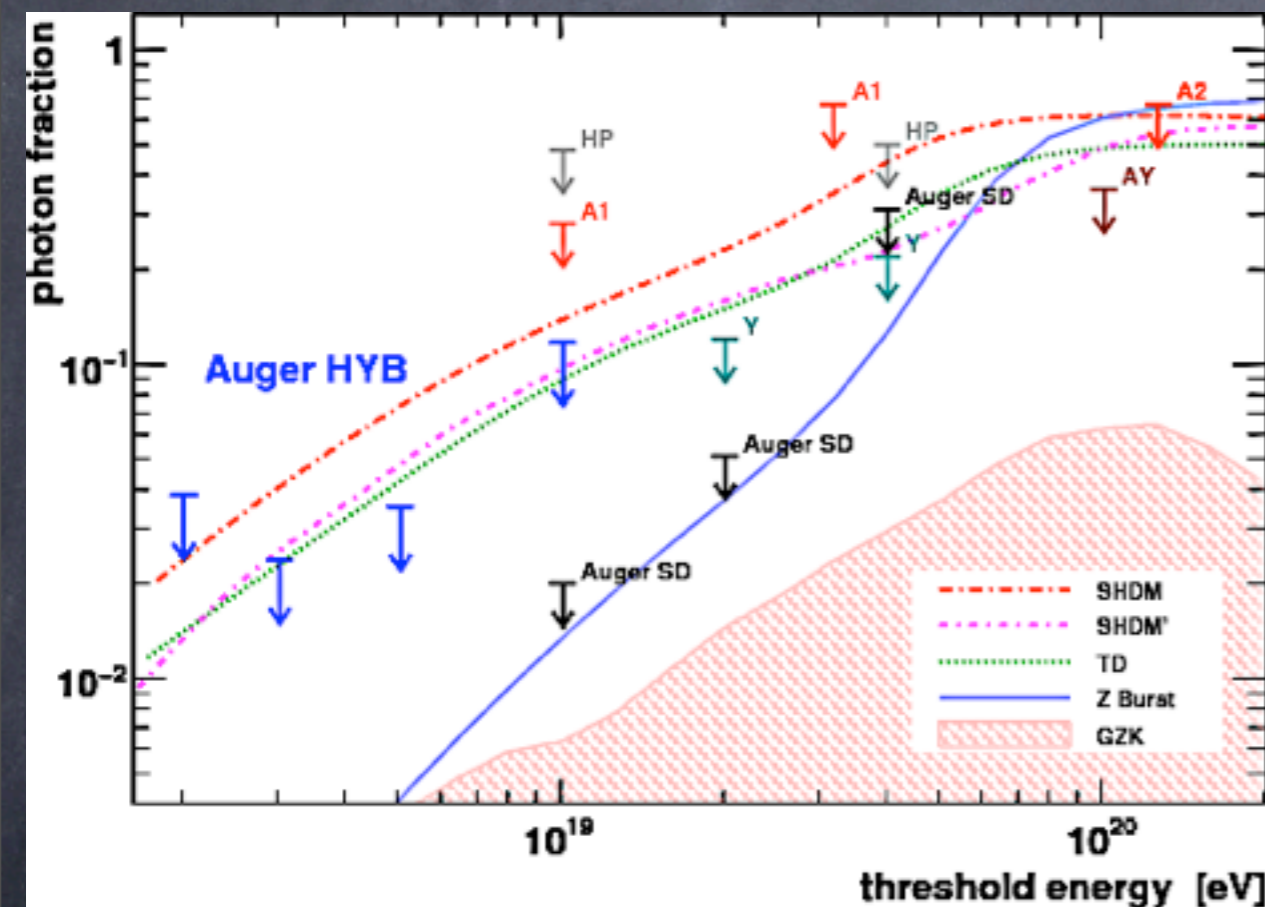


Lorentz Symmetry Violation in the Electromagnetic Sector

The idea:

Experimental upper limits on UHE photon fraction

Contradict predictions if pair production is absent



Pierre Auger Collaboration,
Astropart. Phys. 31 (2009) 399

Maccione, Liberati, Sigl,
PRL 105 (2010) 021101

For photons we assume the dispersion relation

$$\omega_{\pm}^2 = k^2 + \xi_n^{\pm} k^2 \left(\frac{k}{M_{\text{Pl}}} \right)^n, n \geq 1,$$

and for electrons

$$E_{e,\pm}^2 = p_e^2 + m_e^2 + \eta_n^{e,\pm} p_e^2 \left(\frac{p_e}{M_{\text{Pl}}} \right)^n, n \geq 1,$$

with only one term present. Polarizations denoted with \pm . For positrons, effective field theory implies $\eta_n^{p,\pm} = (-1)^n \eta_n^{e,\pm}$. Furthermore, $\xi_n^+ = (-1)^n \xi_n^-$ so that the problem depends on three parameters which in the following we denote by

$$\xi_n, \eta_n^+, \eta_n^-$$

for each n .

Consider pair production on a background photon of energy k_b and assume kinematics with ordinary energy-momentum conservation, with $p_- = (1-y)k$, $p_+ = yk$. Using $x = 4y(1-y)k/k_{LI}$ with the threshold in absence of Lorentz invariance (LI) violation, $k_{LI} = m_e^2/\omega_b$, the condition for pair production is then

$$\alpha_n x^{n+2} + x - 1 \geq 0$$

where

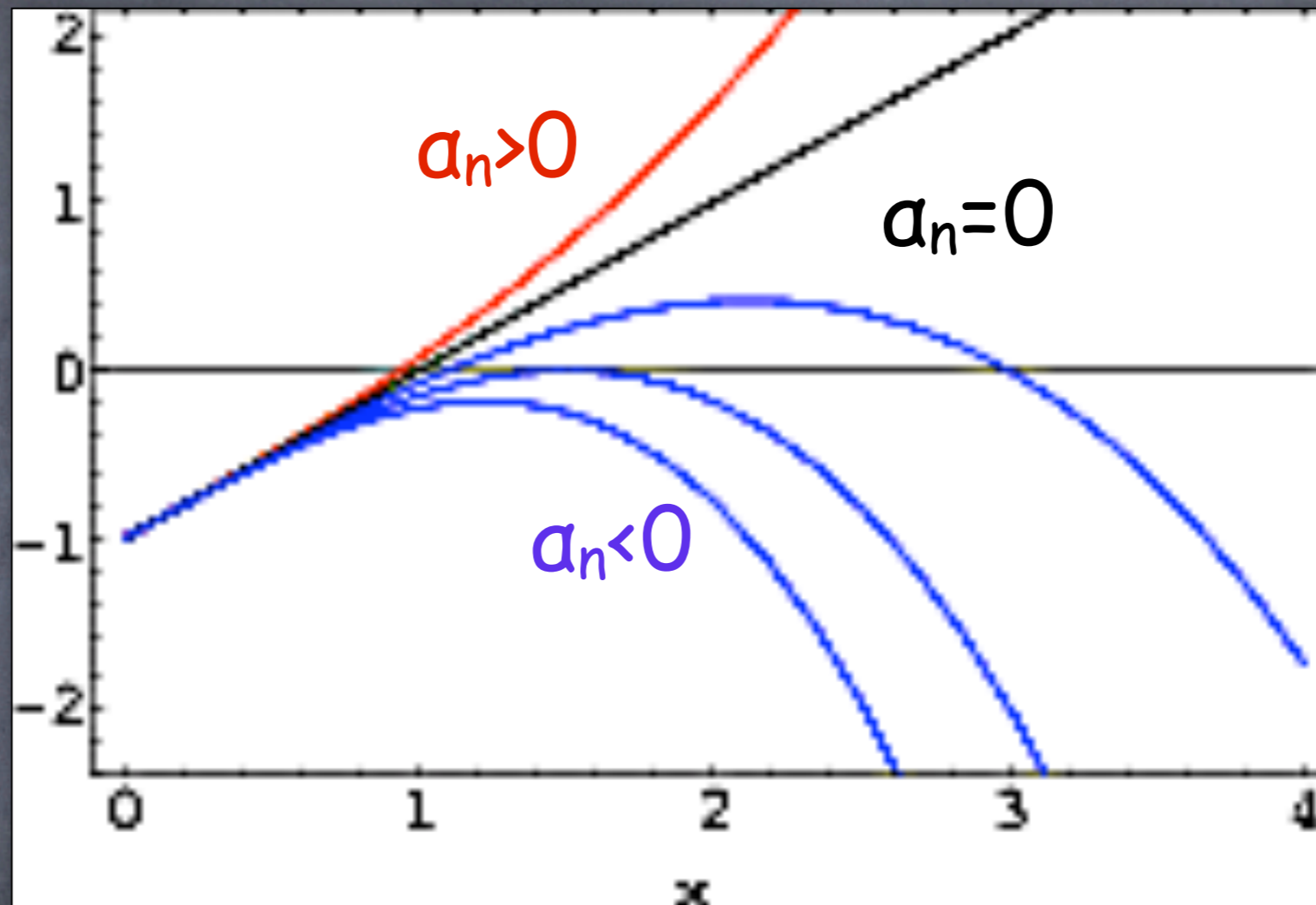
$$\alpha_n = \frac{\xi_n - (-1)^n \eta_n^\mp y^{n+1} - \eta_n^\pm (1-y)^{n+1}}{2^{2(n+2)} y^{n+1} (1-y)^{n+1}} \frac{m_e^{2(n+1)}}{k_b^{n+2} M_{Pl}^n}.$$

All combinations of $\xi_n, \eta_n^+, \eta_n^-$ can occur, depending on the partial wave of the pair, governed by total angular momentum conservation. All partial waves are allowed away from the thresholds.

The condition for photon decay is

$$\alpha_n x^{n+2} - 1 \geq 0$$

There are at least two real solutions $0 \leq x_n^l \leq x_n^r$ for pair production (lower and upper thresholds)



Galaverni, Sigl, Phys. Rev. Lett. 100 (2008) 021102.

For photon decay there is at most one positive real threshold.

Minimize/maximize these wrt. y

A given combination $\xi_n, \eta_n^+, \eta_n^-$ is ruled out if, for $10^{19} \text{ eV} < \omega < 10^{20} \text{ eV}$, at least one photon polarization state is stable against decay and does not pair produce for any helicity configuration of the final pair.

In the absence of LIV in pairs for $n=1$, this yields:

$$\xi_1 \leq 2.4 \times 10^{-15}$$

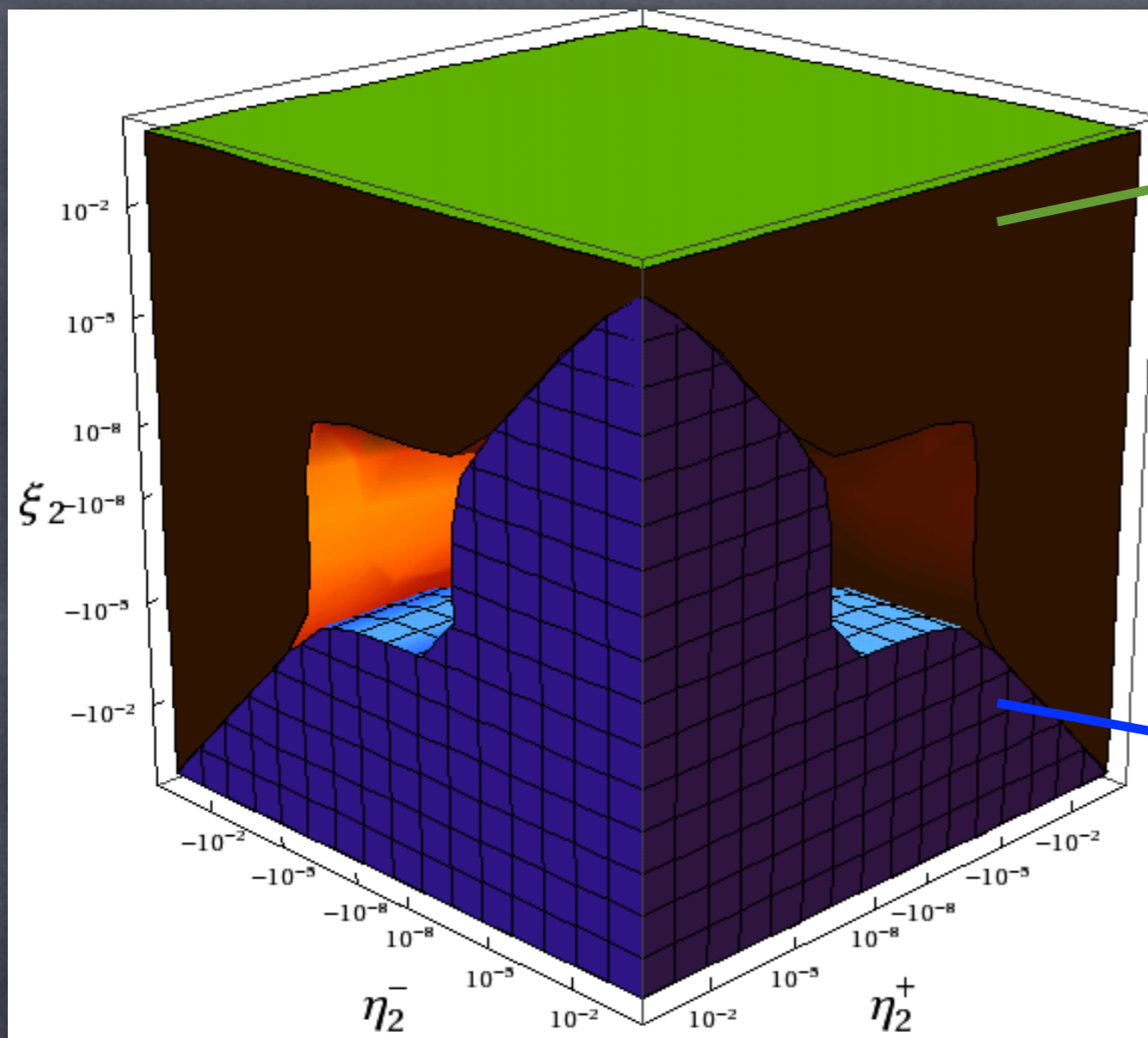
and for $n=2$:

$$\xi_2 \geq -2.4 \times 10^{-7}$$

If a UHE photon were detected, any LIV parameter combination would be ruled out for which photon decay is allowed for both photon polarizations for at least one helicity configuration of the final pair.

For $n = 1$, all parameters of absolute value $< 10^{-14}$ ruled out

For $n = 2$, if absolute value of both the photon and one of the electron parameters is $< 10^{-6}$, the second electron parameter can be arbitrarily large even once a UHE photon is seen.



UHE photons are detected

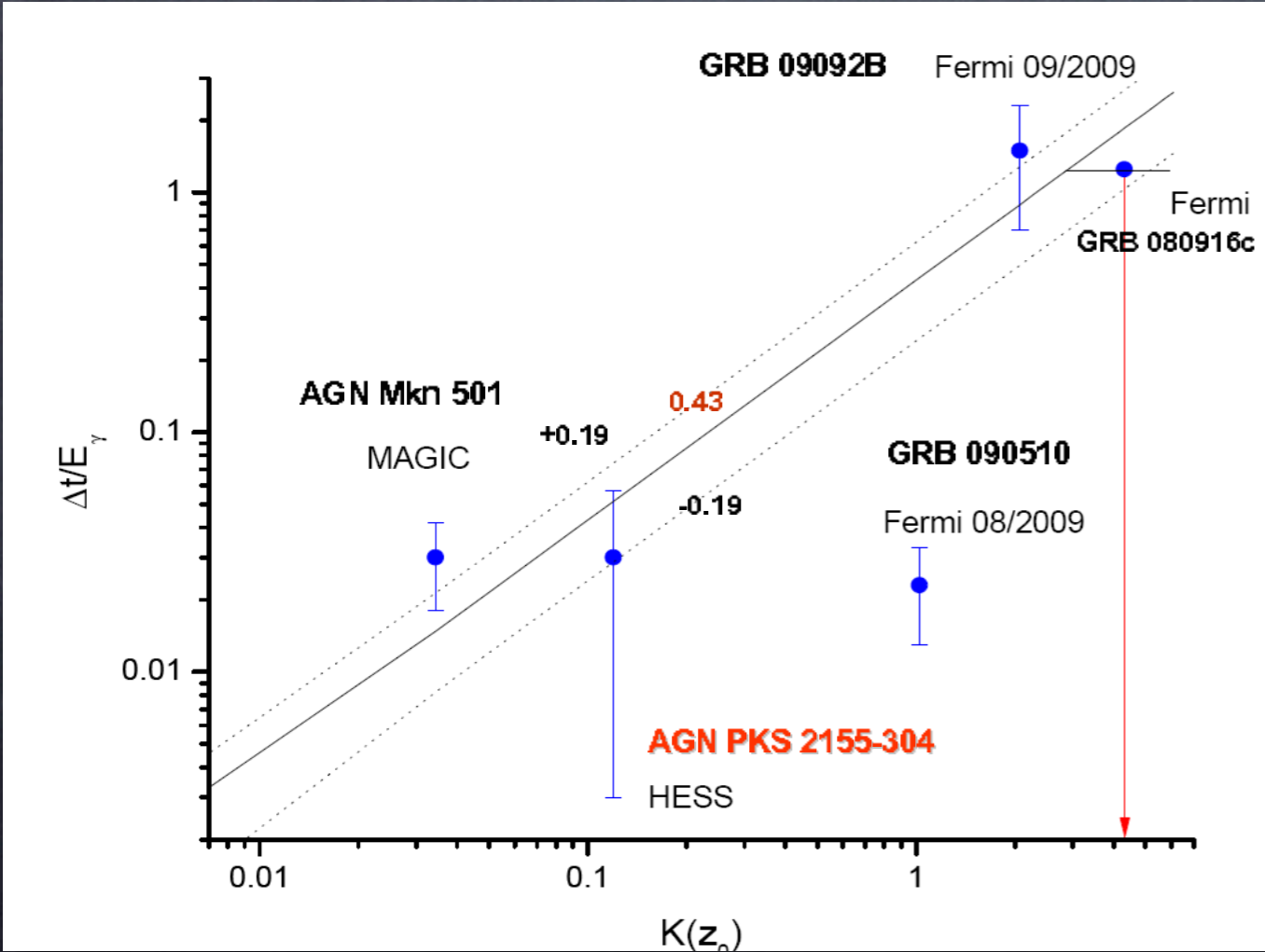
UHE photon absorption takes place

Such strong limits suggest that Lorentz invariance violations are completely absent !

The modified dispersion relation also leads to energy dependent group velocity $V = \partial E / \partial p$ and thus to an energy-dependent time delay over a distance d :

$$\Delta t = -\xi d \frac{E}{M_{\text{Pl}}} \simeq -\xi \left(\frac{d}{100 \text{ Mpc}} \right) \left(\frac{E}{\text{TeV}} \right) \text{ sec}$$

for linearly suppressed terms. GRB observations in TeV γ -rays can therefore probe quantum gravity and may explain that higher energy photons tend to arrive later (Ellis et al.).

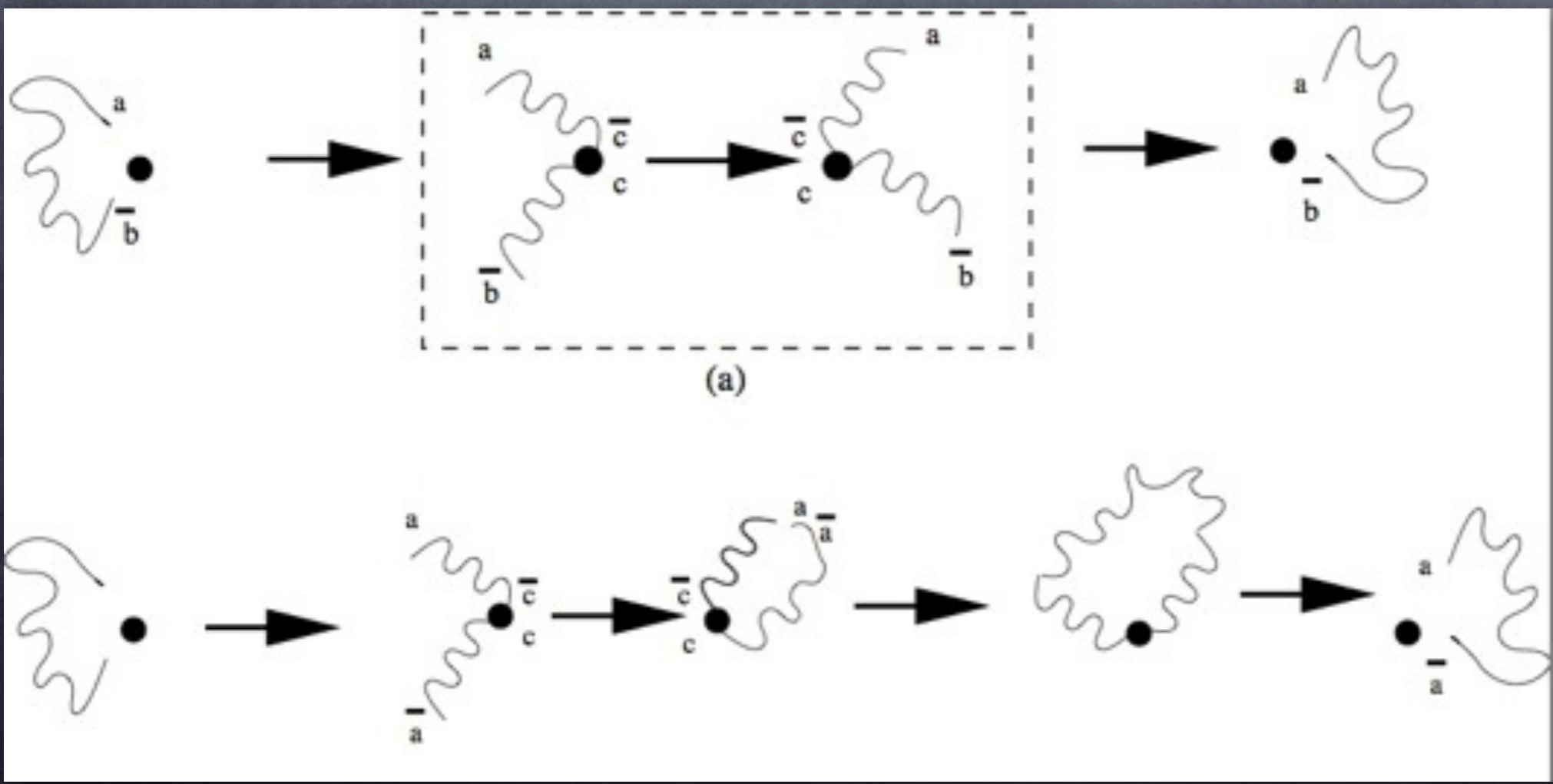


But the UHE photon limits are inconsistent with interpretations of time delays of high energy gamma-rays from GRBs within quantum gravity scenarios based on effective field theory

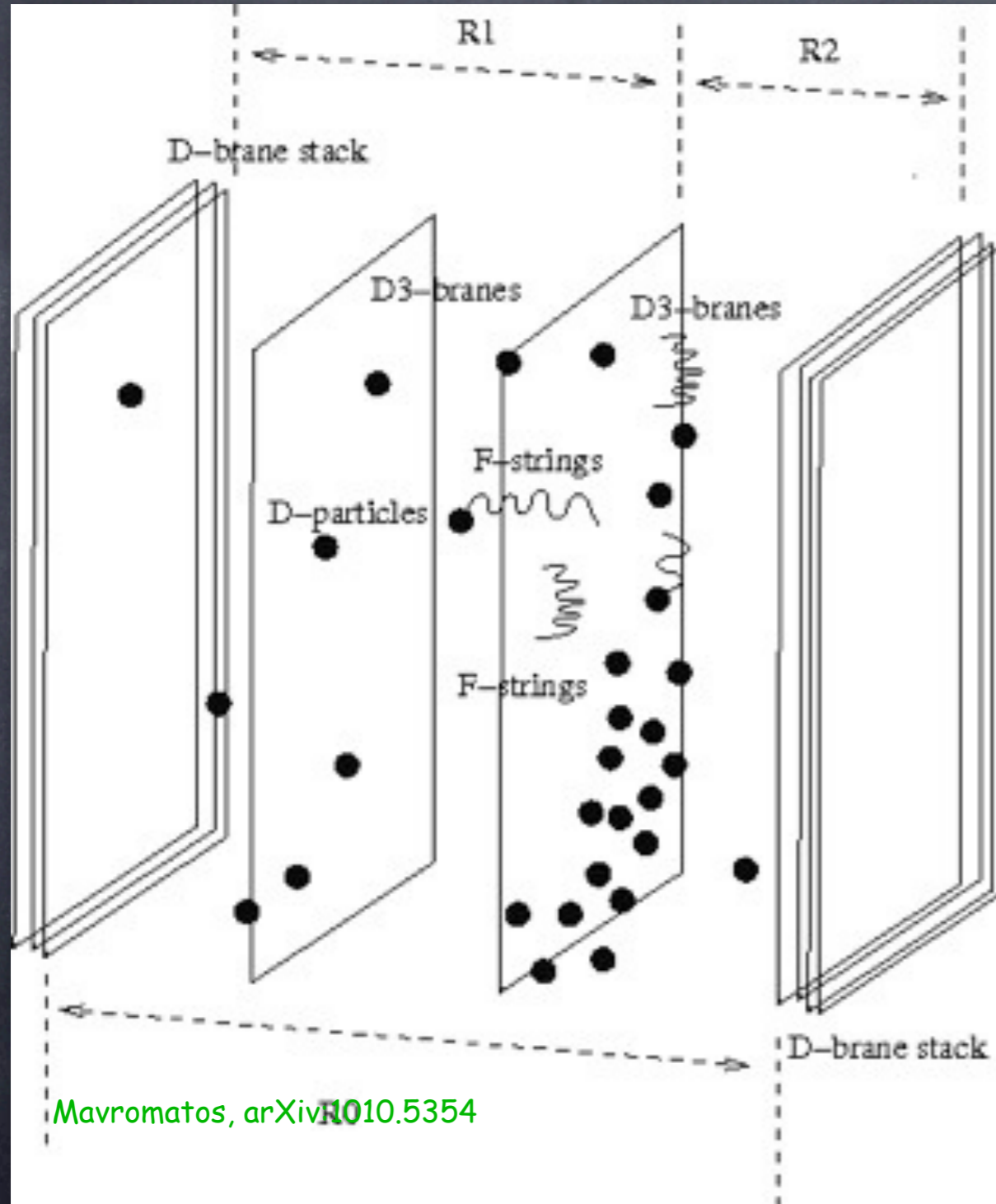
Maccione, Liberati, Sigl, PRL 105 (2010) 021101

Possible exception in space-time foam models,

Ellis, Mavromatos, Nanopoulos, arXiv:1004.4167

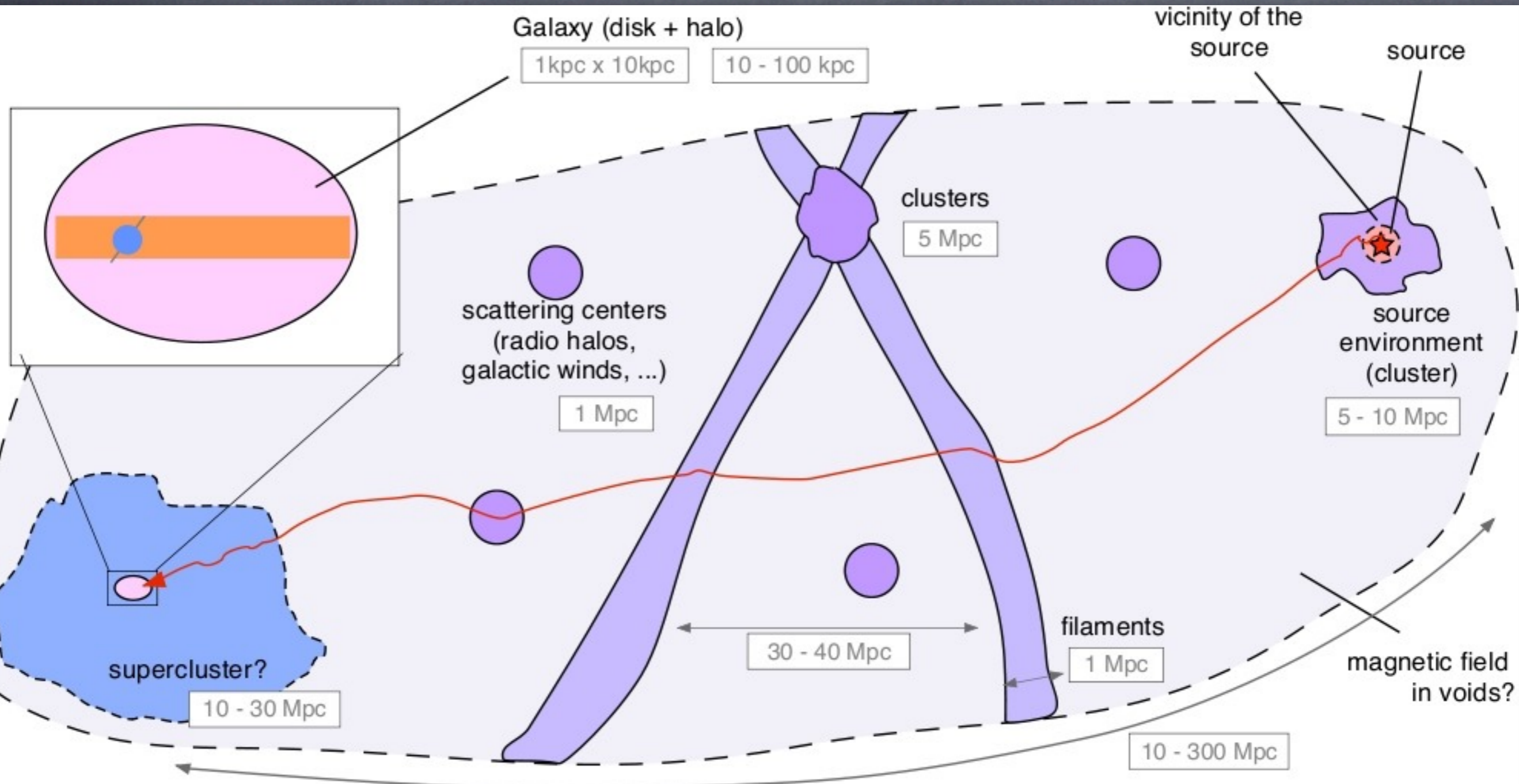


In space-time foam models there may be fluctuating terms in dispersion relation, thus no strict energy-momentum conservation. This could circumvent pair production limits, allowing to interpret time dispersion by quantum gravity effects

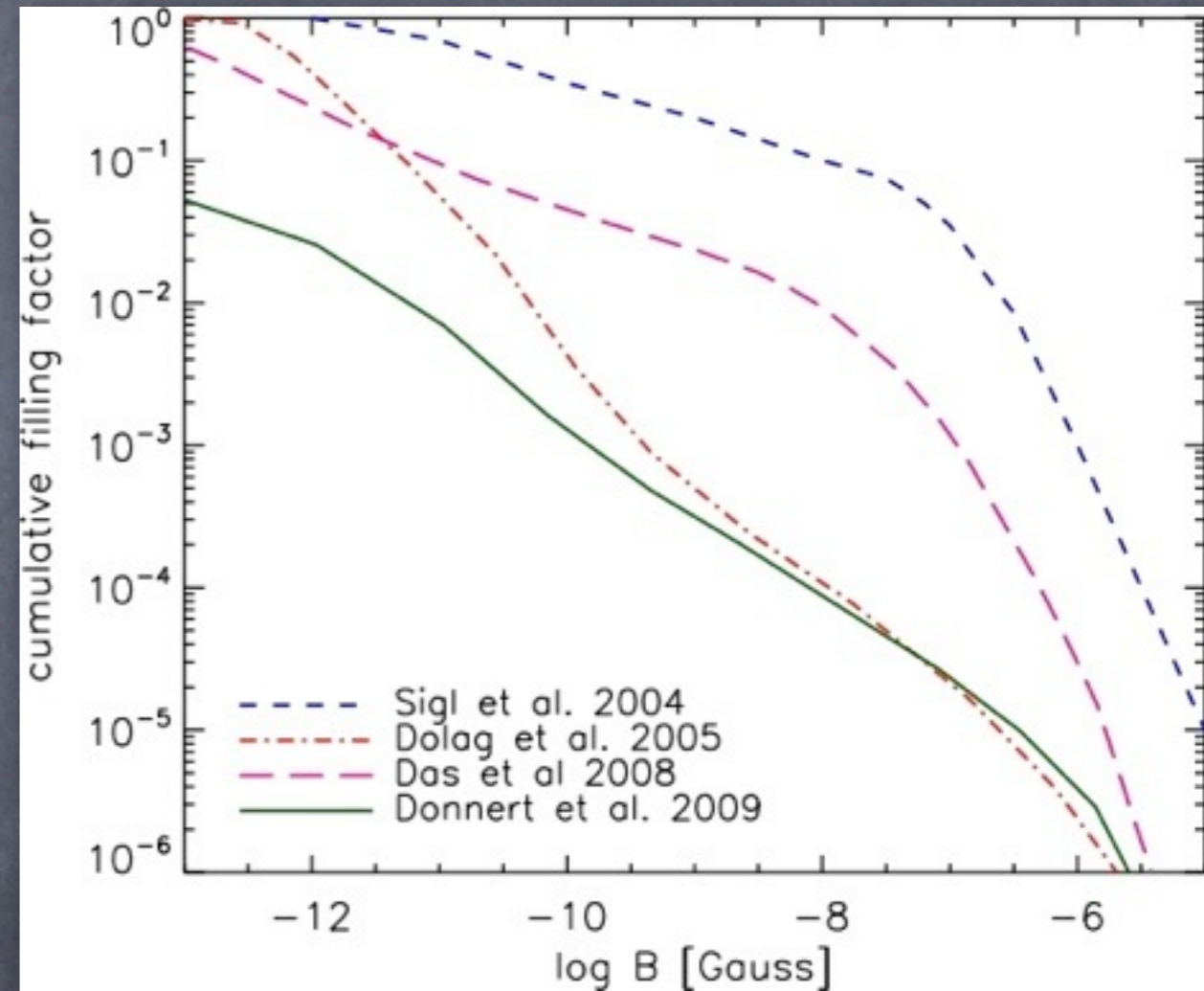
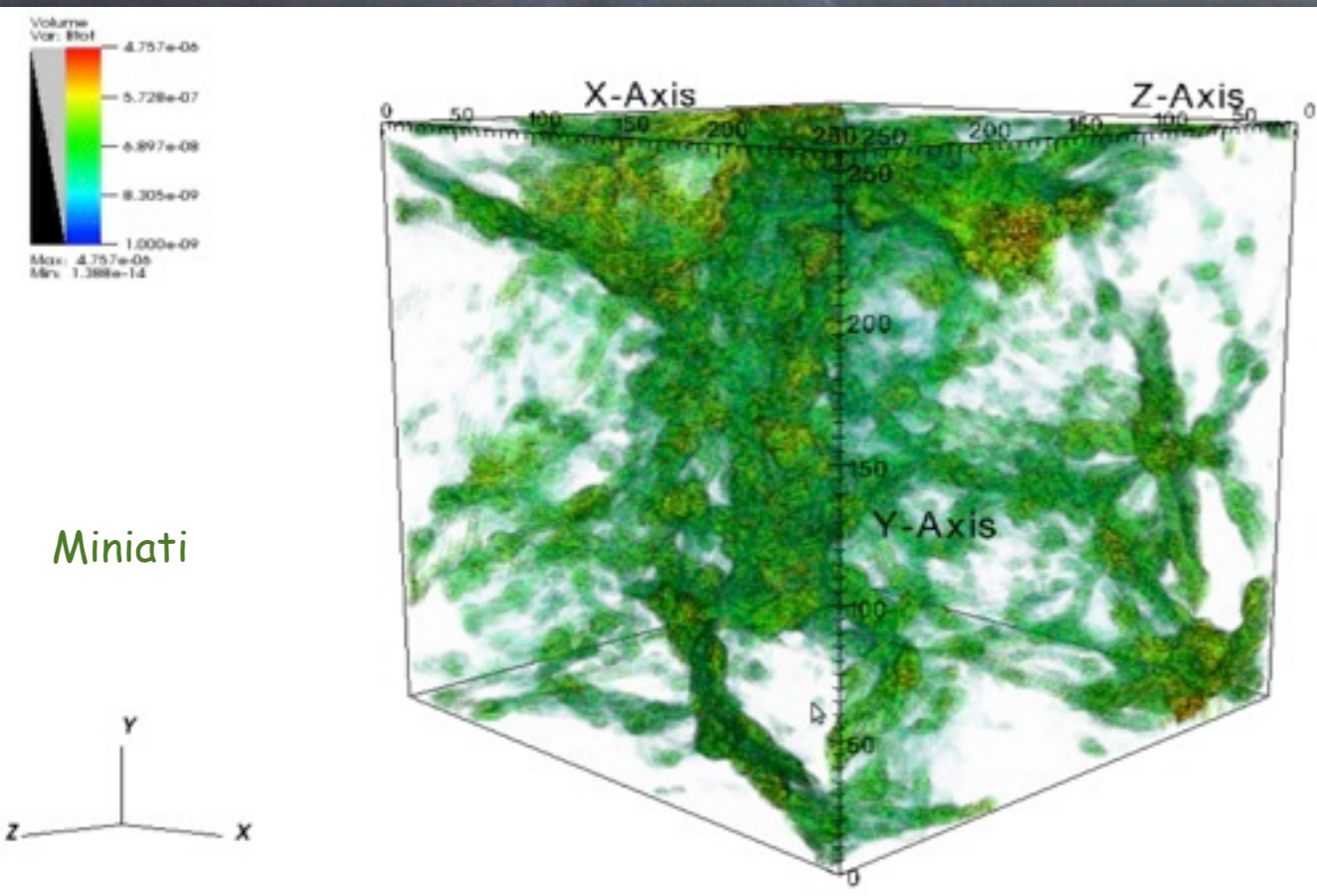


Mavromatos, arXiv:1010.5354

3-Dimensional Effects in Propagation



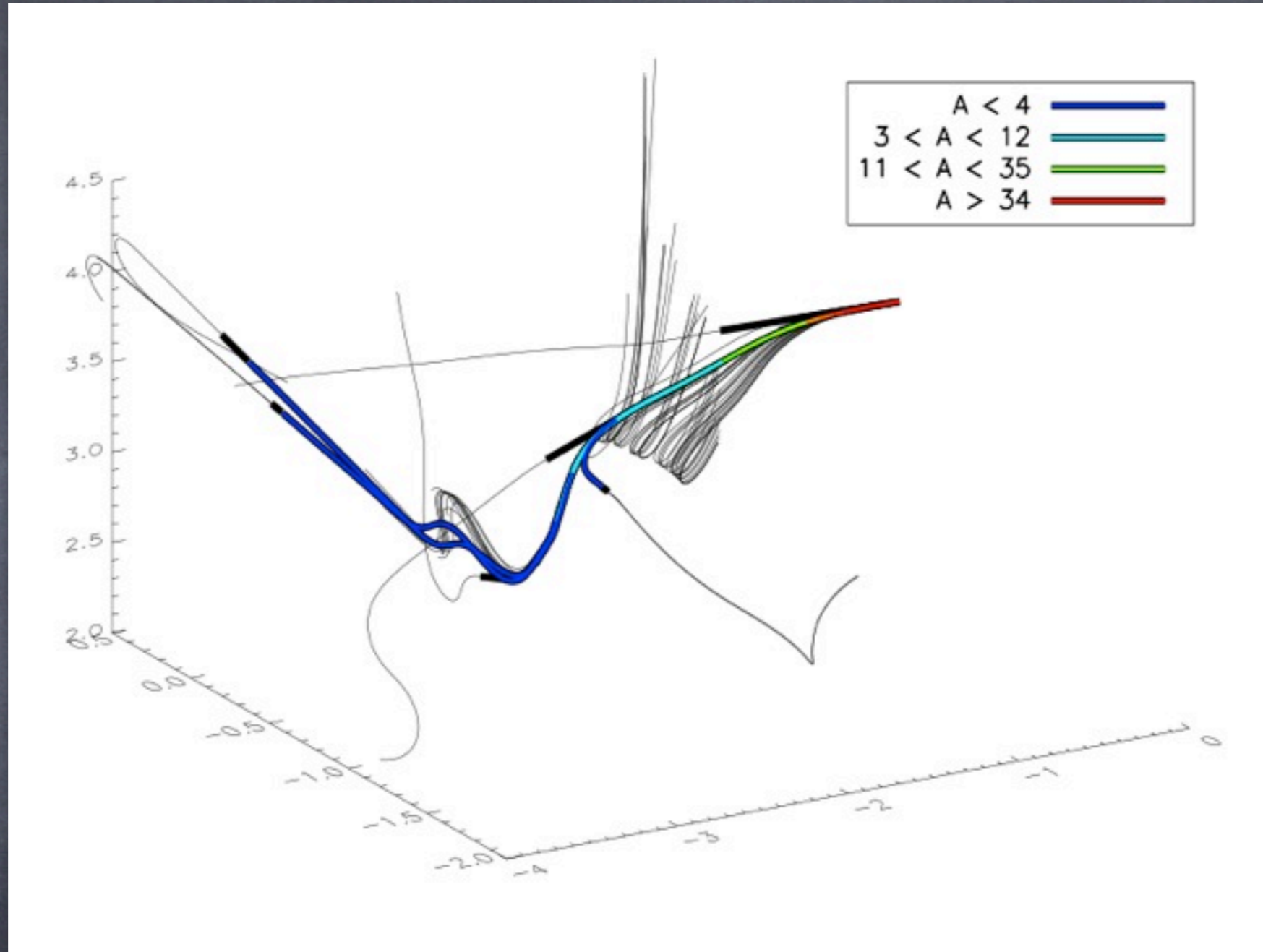
Structured Extragalactic Magnetic Fields



Kotera, Olinto, *Ann.Rev.Astron.Astrophys.* 49 (2011) 119

Filling factors of extragalactic magnetic fields are not well known and come out different in different large scale structure simulations

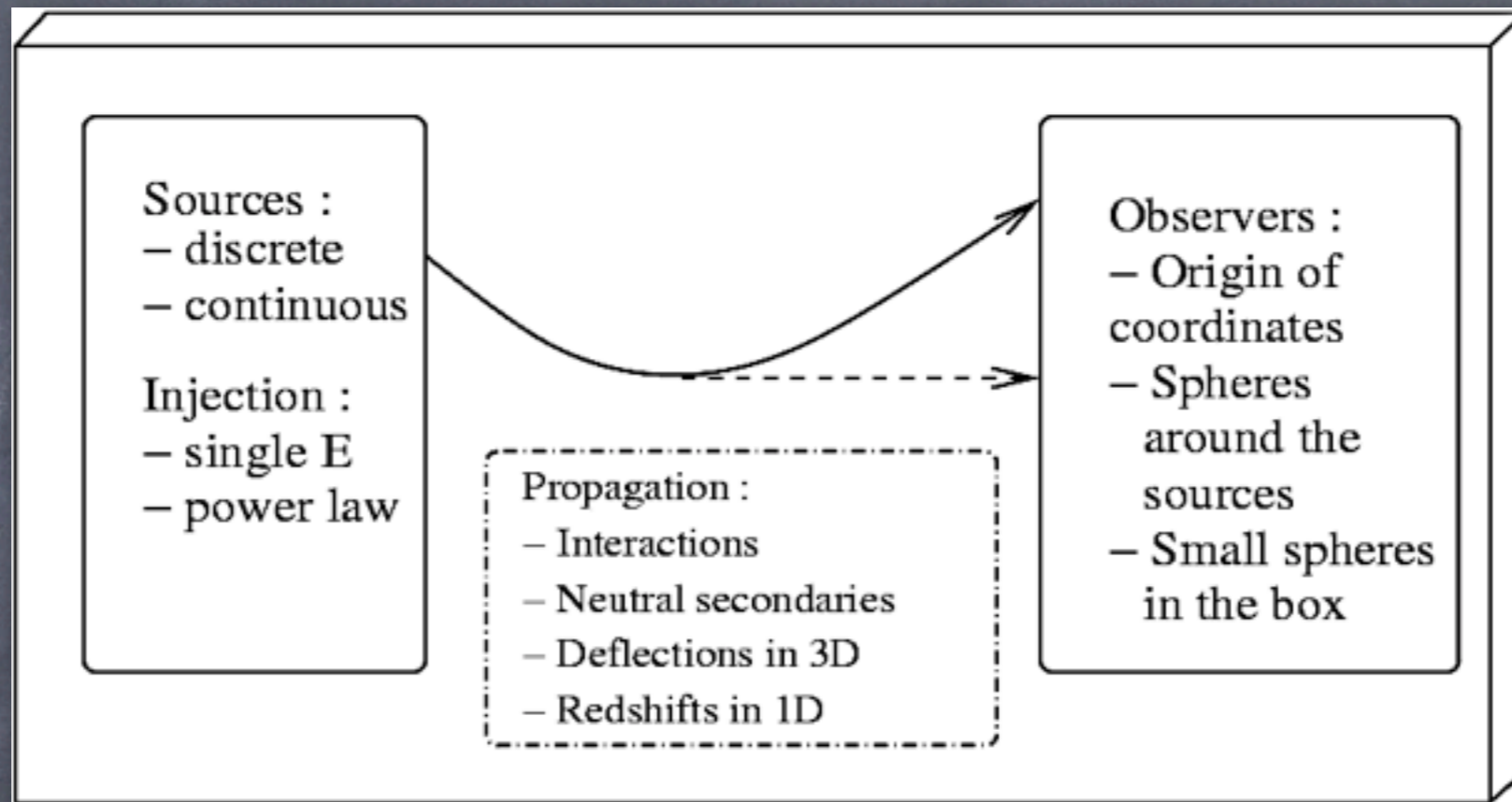
Extragalactic iron propagation produces nuclear cascades in structured magnetic fields:



Initial energy 1.2×10^{21} eV, magnetic field range 10^{-15} to 10^{-6} G. Color-coded is the mass number of secondary nuclei

CRPropa 2.0

CRPropa is a public code for UHE cosmic rays, neutrinos and γ -rays being extended to heavy nuclei and hadronic interactions



Eric Armengaud, Tristan Beau, Günter Sigl, Francesco Miniati,
*Astropart.Phys.*28 (2007) 463.

Version 1.4 at <http://apcauger.in2p3.fr/CRPropa/index.php>

Now including: Jörg Kulbartz, Luca Maccione,
Nils Nierstenhoefer, Karl-Heinz Kampert, Peter Schiffer, Arjen van Vliet

ask for beta version CRPropa 2.0 !

The main part of the code is written in C++ and calls some Fortran routines
(mainly SOPHIA for interactions photo-pion production of nucleons)
nuclear interactions based on TALYS

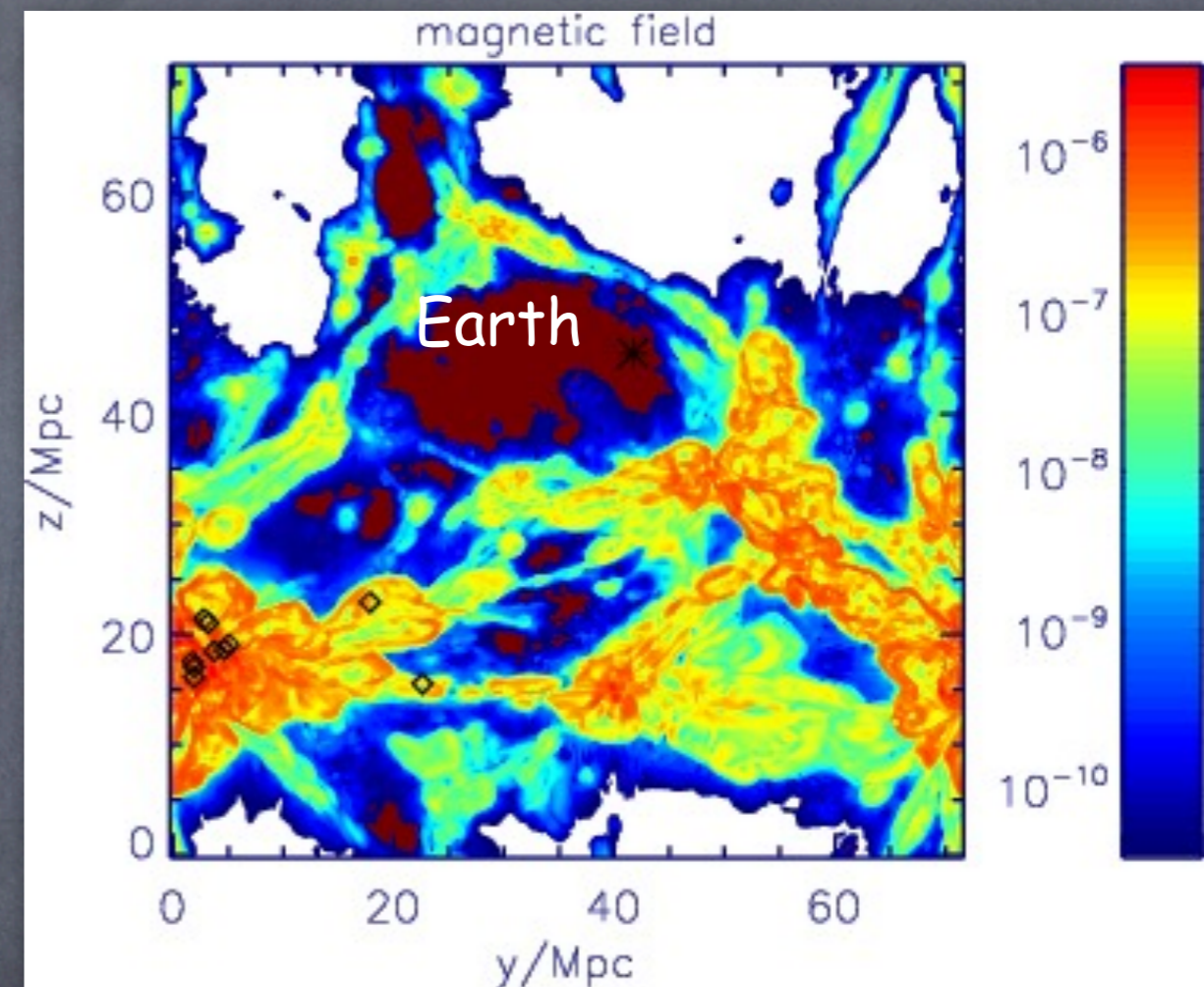
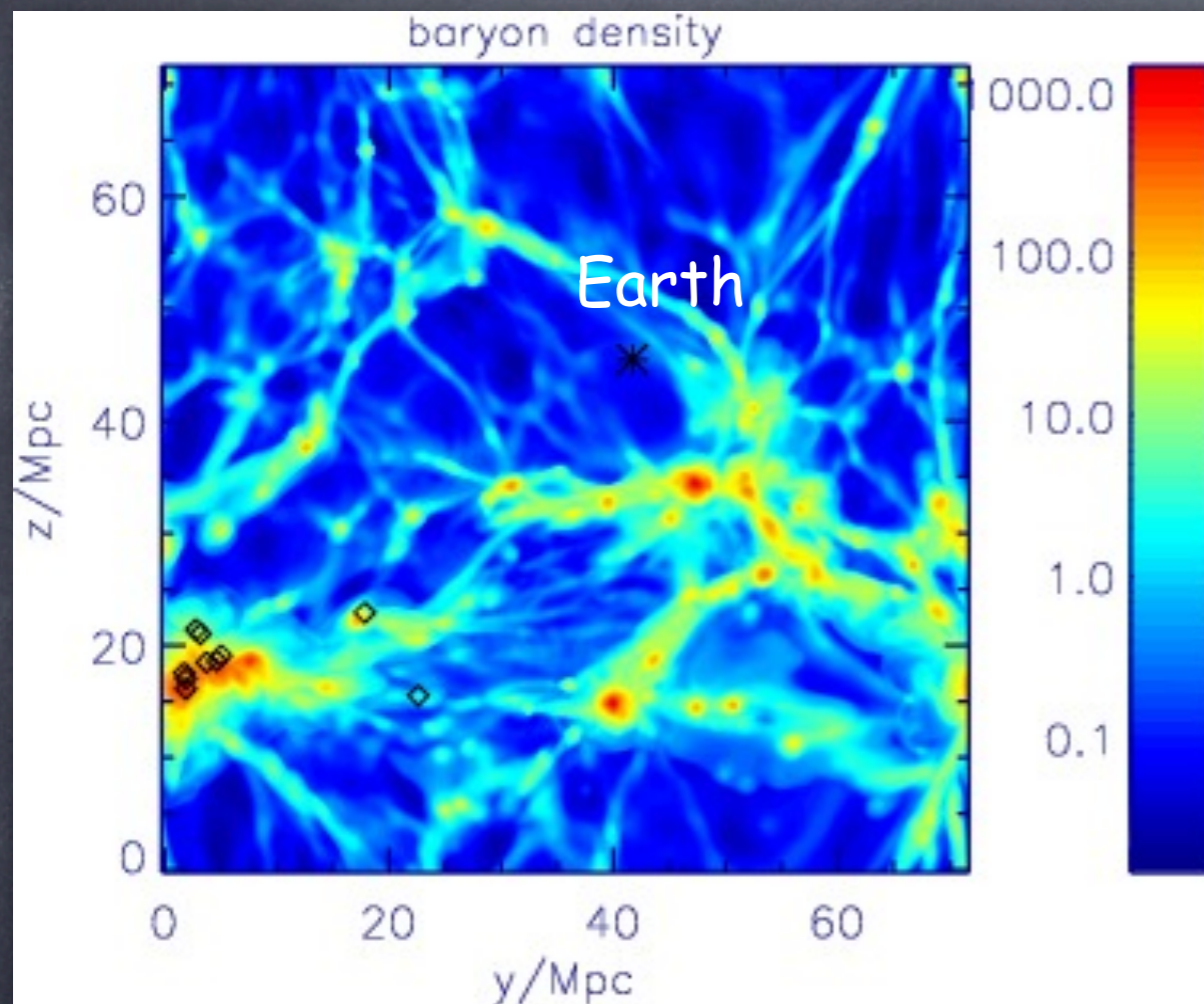
Electromagnetic cascades are treated by solving one-dimensional transport
equations

The set-up (source distributions, environment, magnetic fields, low energy
photon backgrounds, injection spectrum, arbitrary composition at fixed energy per
nucleon, which interactions/secondaries to take into account)
can be provided with xml files.

Output can be in form of whole trajectories or events; possible output formats are
ASCII, FITS or ROOT.

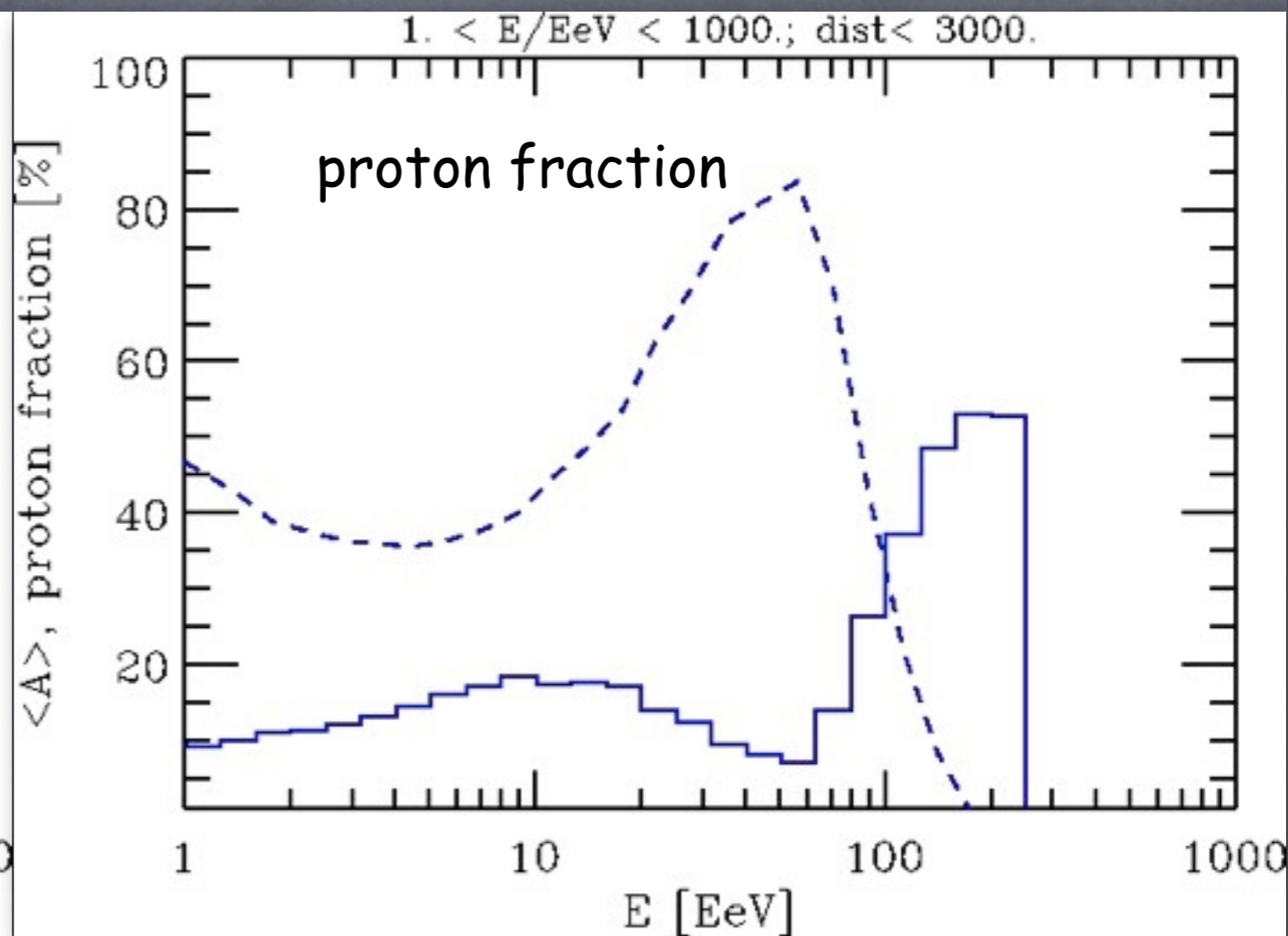
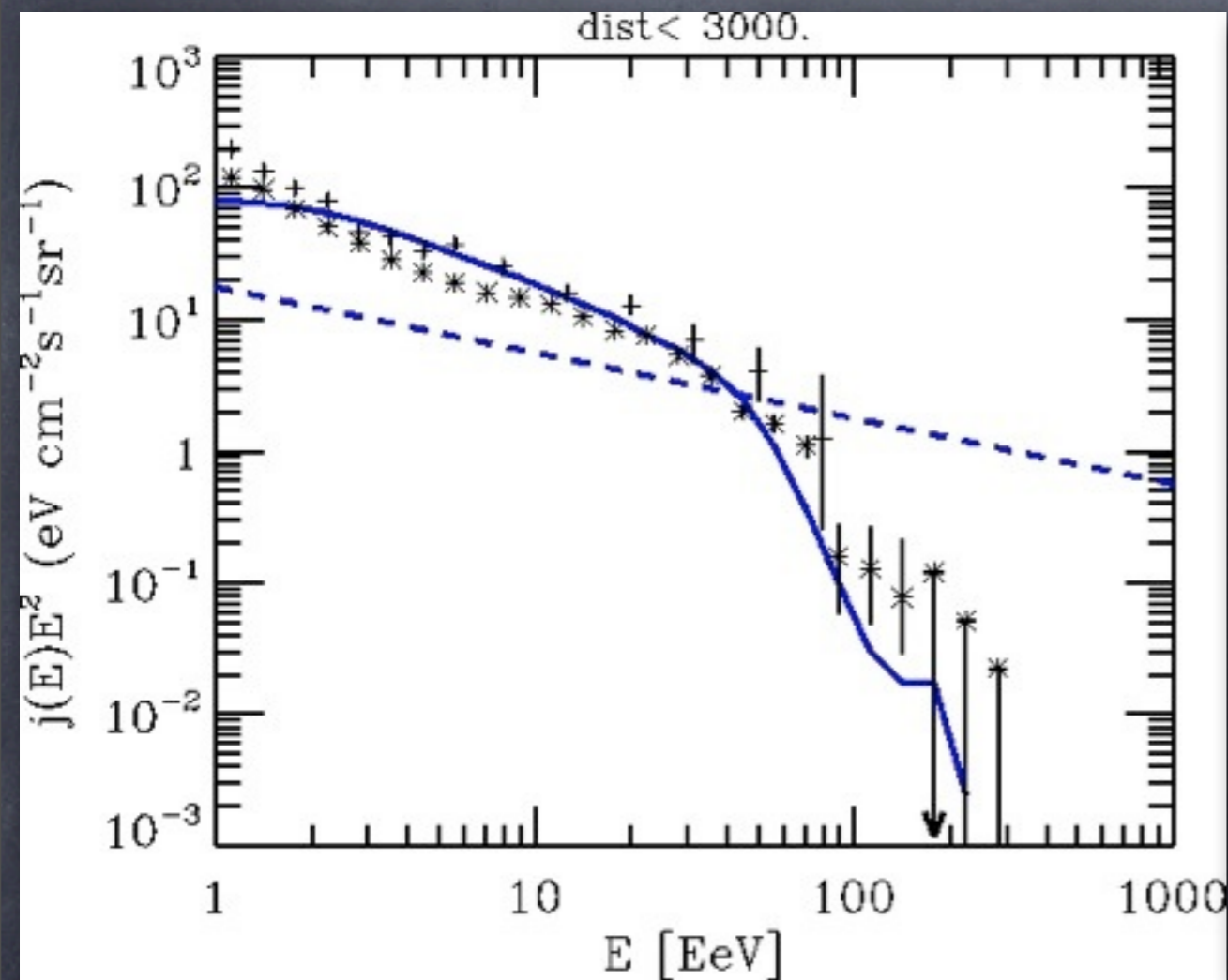
Presented are two examples for 1D and 3D simulations

Discrete Sources in nearby large scale structure

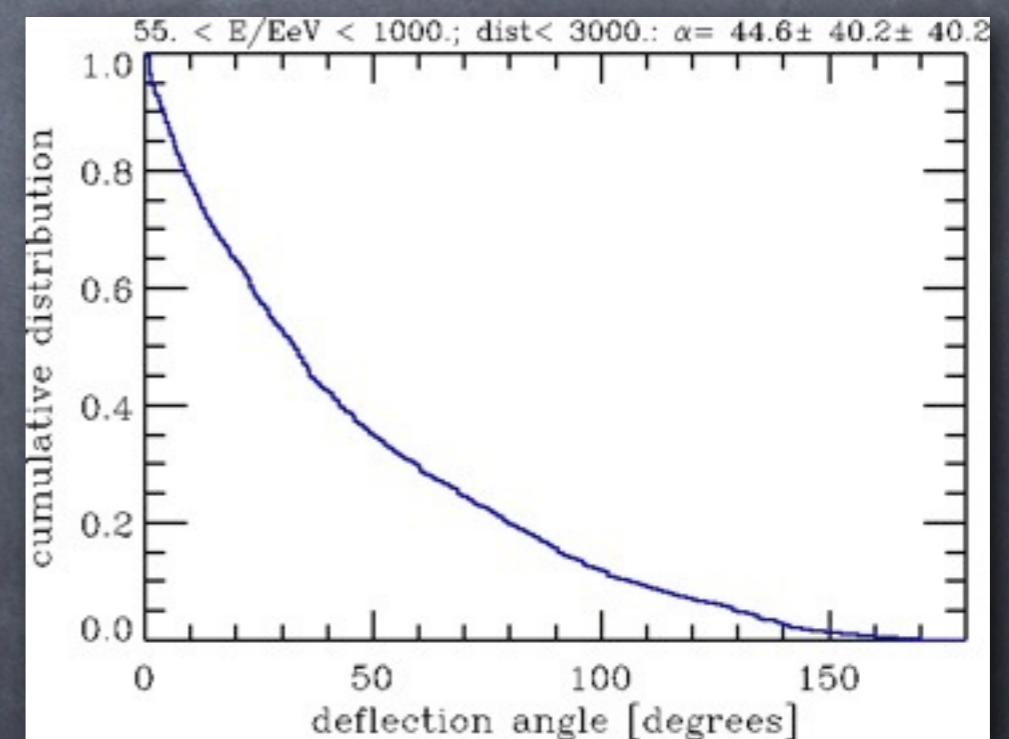
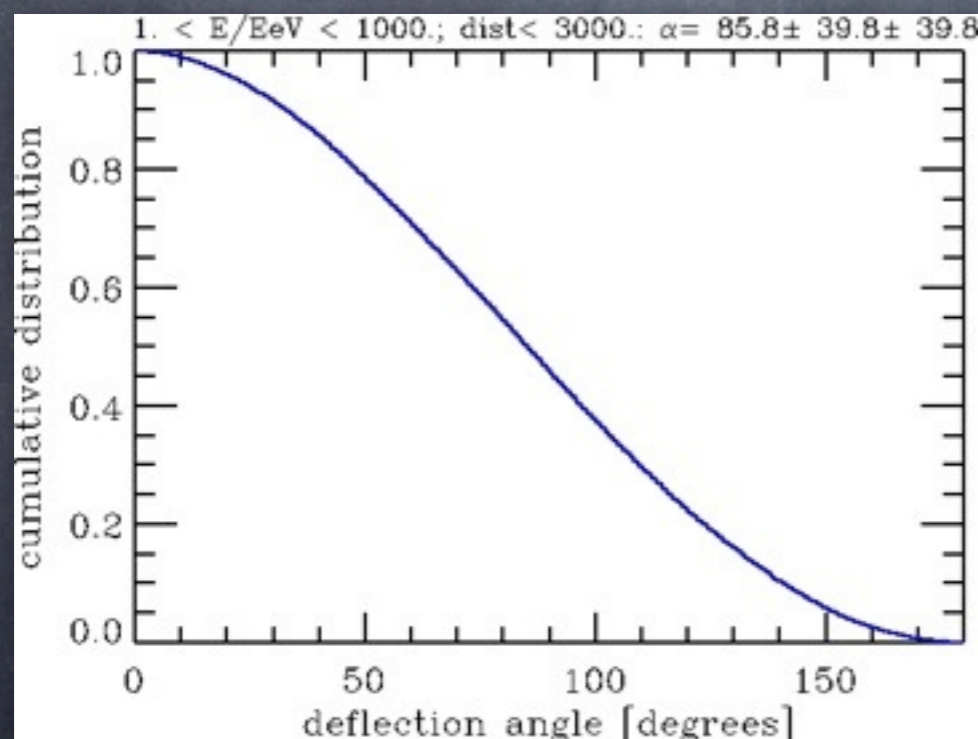
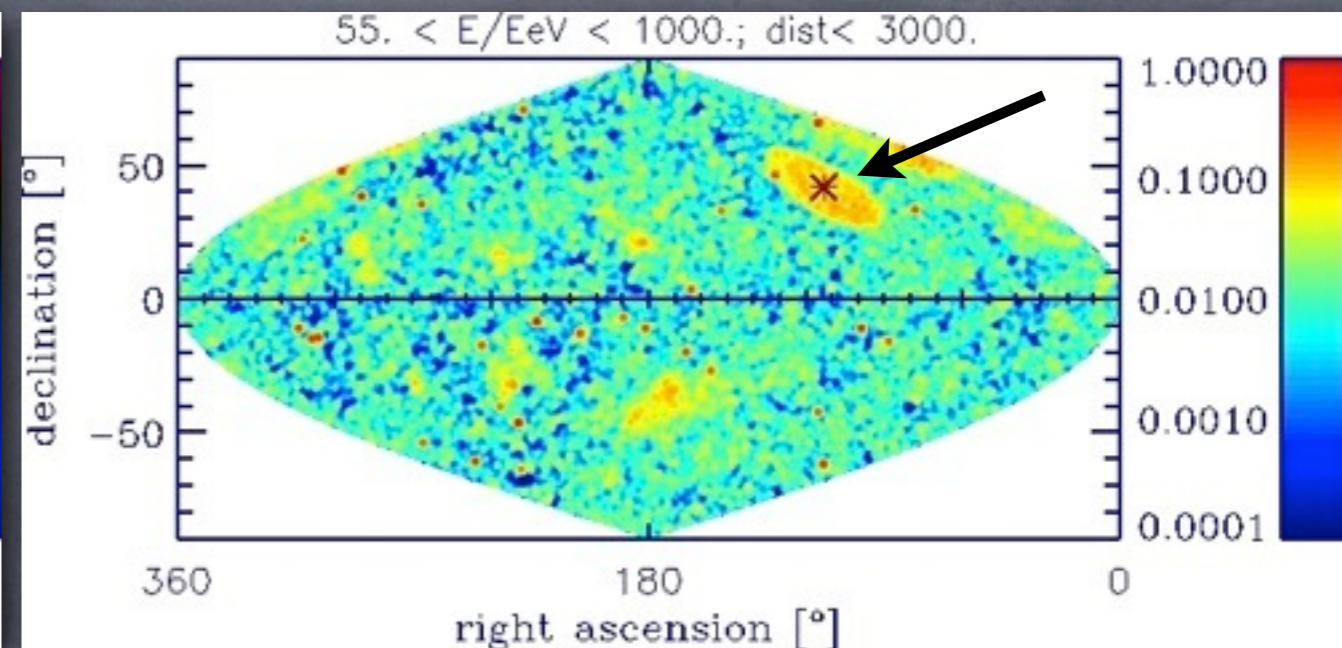
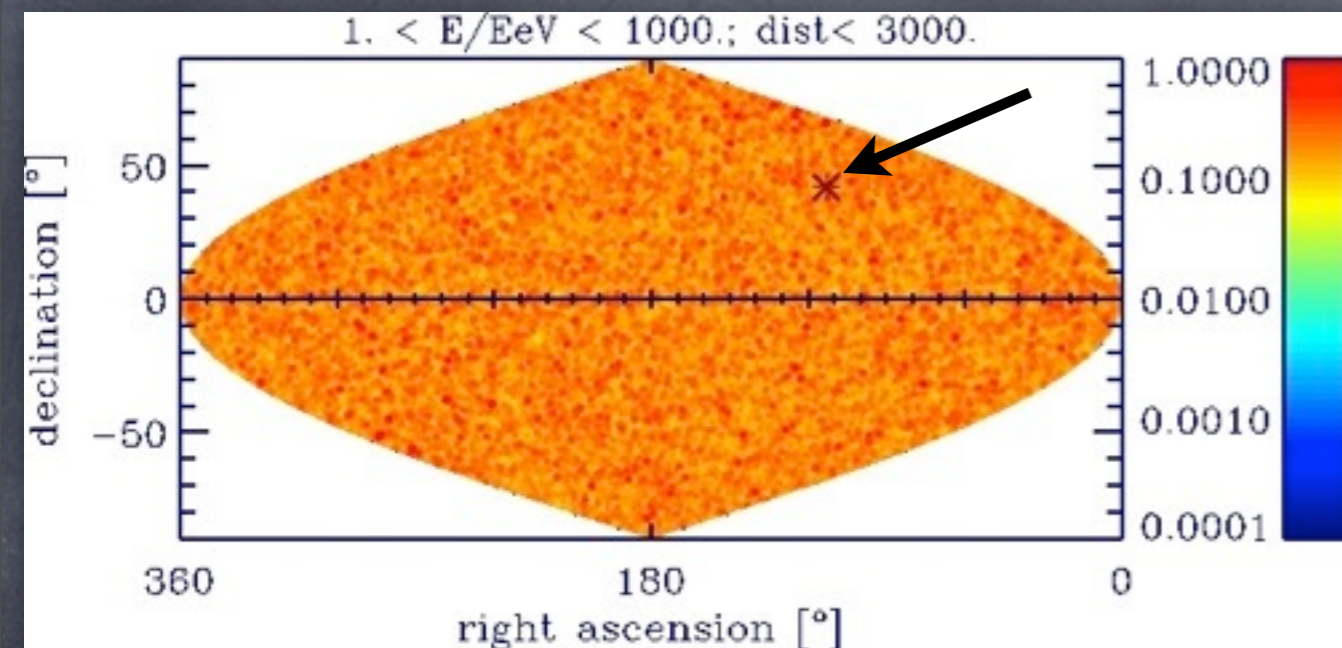


10 sources per $(75 \text{ Mpc})^3$ box, concentrated in a galaxy cluster at $\approx 30 \text{ Mpc}$, injecting $E^{-2.5}$ spectra up to $200xZ \text{ EeV}$ with $10 \times$ galactic abundance;
+ one source @ 4 Mpc of 0.002 relative strength injecting E^{-2} spectrum up to $10xZ \text{ EeV}$.

Results: Spectra and Composition



Results: Sky Distributions and Anisotropies



above 1 EeV

above 55 EeV

It is surprisingly difficult to construct simple scenarios with structured sources and magnetic fields that reproduce all observations: spectra, energy dependent composition and anisotropy; to explain them separately is quite easy

Conclusions

- 1.) It is surprisingly difficult to construct simple scenarios with structured sources and magnetic fields that reproduce all observations: spectra, energy dependent composition and anisotropy; to explain them separately is quite easy
- 2.) The observed X_{\max} distribution of air showers provides potential constraints on hadronic interaction models: Some models are in tension even when "optimizing" unknown mass composition; however, systematic uncertainties are still high.

Conclusions

3.) Both diffuse cosmogenic neutrino and photon fluxes mostly depend on chemical composition, maximal acceleration energy and redshift evolution of sources

4.) Multi-messenger modeling sources including gamma-rays and neutrinos start to constrain the source and acceleration mechanisms

5.) Highest Energy Cosmic Rays, Gamma-rays, and Neutrinos give the strongest constraints on violations of Lorentz symmetry => terms suppressed to first and second order in the Planck mass would have to be unnaturally small