

# Particle Physics & String Theory: Type II/F-Theory Perspective

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# Outline:

Focus on particle physics & D-branes:

I. Type II (w/ D-branes at small string coupling)

→ Standard Model & GUT's

Recent developments: Non-perturbative effects (D-instantons)

→ new hierarchy for couplings

II. F-theory (string theory w/ D-branes at finite coupling)

→ primarily (local) SU(5) GUT's

→ instantons

c.f. J. Heckman's talk

time permitting

III. Conclusions/outlook

## I. Particle Physics implications (Type II):

Recent focus: landscape of realistic D-brane quivers w/ D-instantons

### String Consistent MSSM Quivers w/ realistic fermion textures

M.C., J. Halverson, R. Richter; & P. Langacker '09-'10

Most MSSM Quivers-string inconsistent

**What are the simplest extensions?**

With no additional nodes:

**landscape analysis all MSSM quivers & additional matter**

(compatible with string constraints) → stringy inputs on exotic matter

M.C., J. Halverson & P. Langacker, 1108.5187

With additional  $U(1)$ 's or  $U(N)$  node:

**implications for SUSY breaking, dark matter,  $Z'$**

M.C., J. Halverson & H. Piragua, to appear

## II. D-instantons – formal developments

No time

Recent focus on F-theory (finite string coupling)

→ multi pronged approach:

(a) string junctions

(b) anomaly inflow

(c) via dualities (F-theory/Heterotic; F-theory/M-theory)

**Goals:** Zero modes and superpotential structure (Pfaffians): →

**Pfaffians via heterotic duality:**

Geometric interpretation of zero loci (including  $E_8$  symmetric point)

M.C., I. Garcia-Etxebarria & J. Halverson, 1107.2388

Inclusion of fluxes & direct F-theory results

M.C., R. Donagi, J. Halverson & J. Marsano, to appear

**Superpotential via  $N=2$   $D=3$  M-theory**

M.C., T. Grimm, J. Halverson & D. Klevers, to appear

# Perturbative String Theories → (finite) theory of quantum gravity

Green&Schwarz'84

## Phenomenologically promising

Recent (MSSM): Bouchard, M.C., Donagi'05

....Anderson, Gray, Lukas'09-'12

Lebedev, Nilles, Raby, Ramos, Ratz, Vaudrevange, Wingerter'07-'10

Heterotic  $E_8 \times E_8$  string

Type IIA superstring  
(closed)

Type IIB superstring  
(closed)

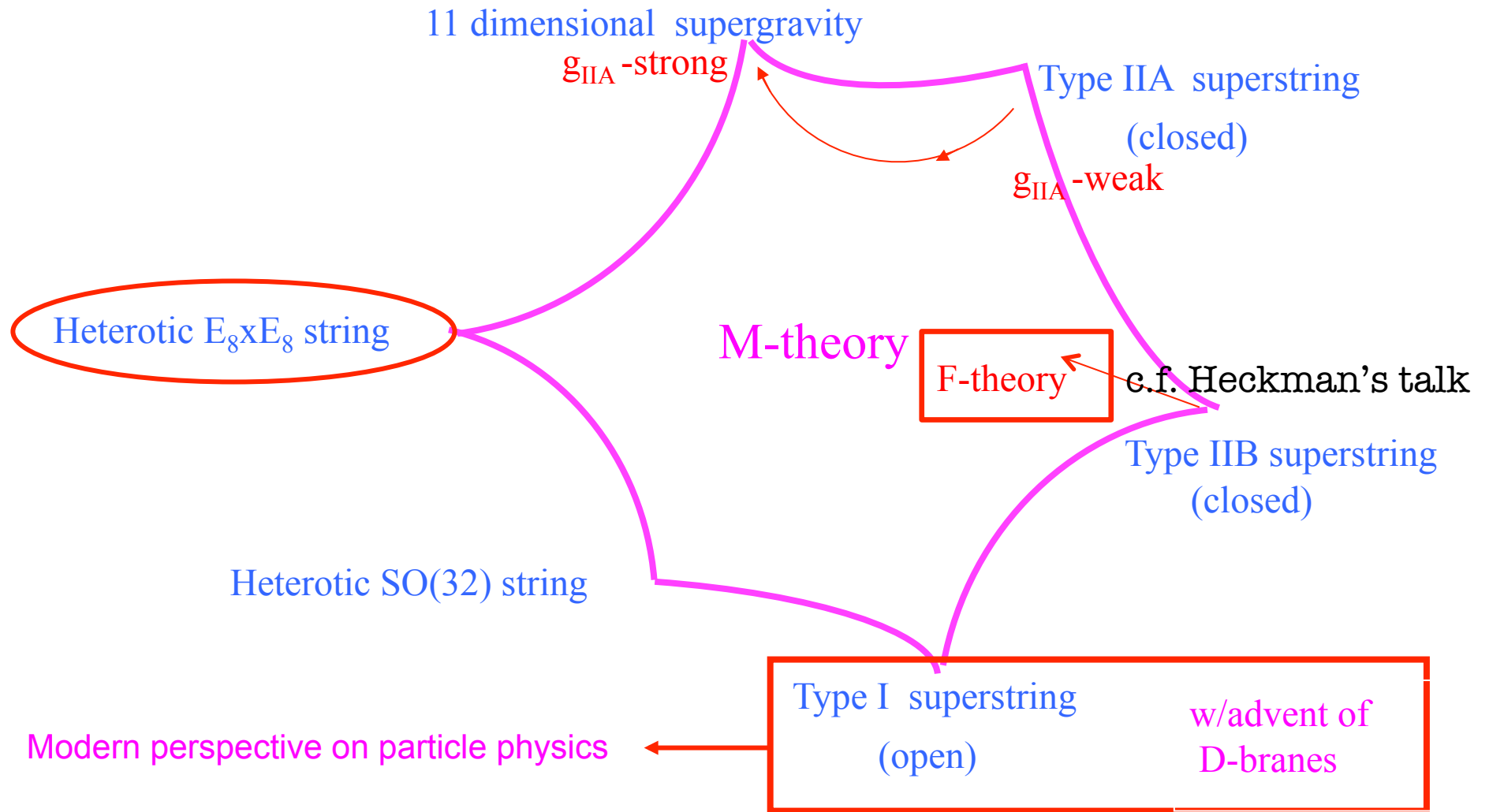
Heterotic  $SO(32)$  string

Type I superstring  
(open)

# Perturbative String Theories → (finite) theory of quantum gravity

Hull&Townsend'94  
Witten'95

## Non-perturbative Unification



Different String Theories related to each other by Weak-Strong Coupling **DUALITY**

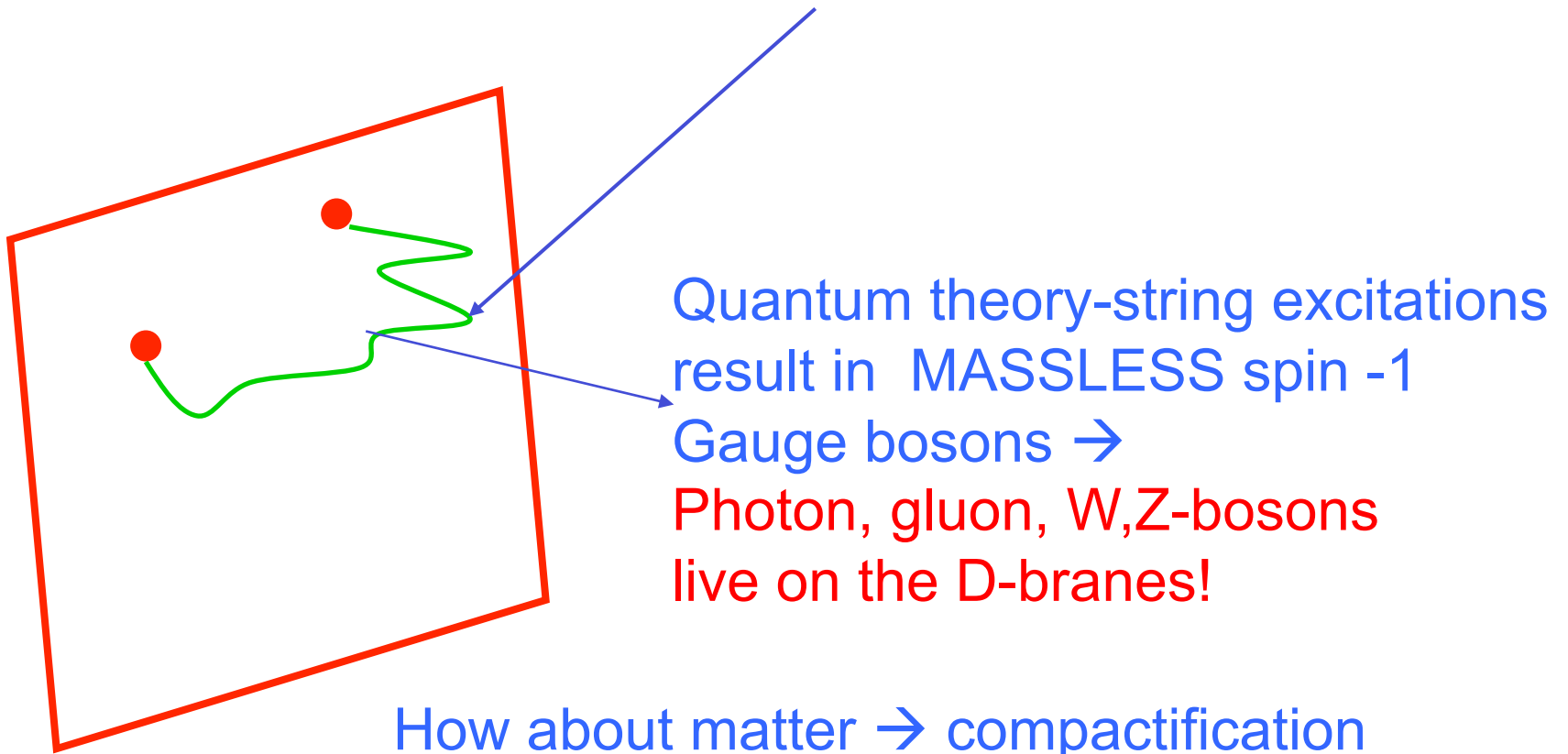
# D-branes & Particle Physics →

Beautiful relation to particles & forces of nature - geometric

Open strings w/ charges at the ends

Ends ``attached`` to boundary **Dp – branes** Polchinski'95

Extends in  $p+1$  dim. world-volume

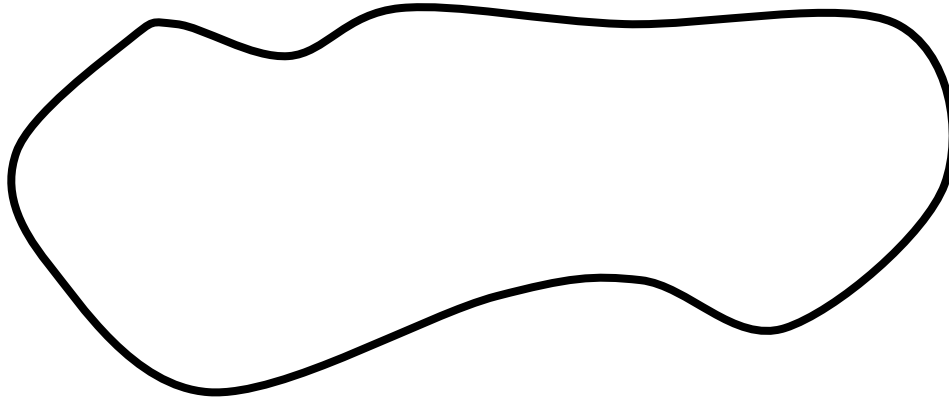


# Compactification

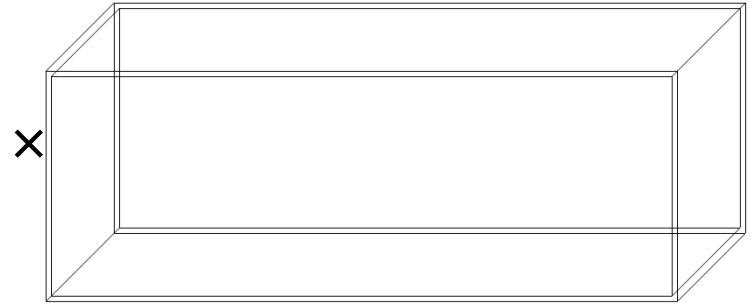
$D=9+1$   $\longrightarrow$   $D=3+1$



$X_6$  (Calabi-Yau)

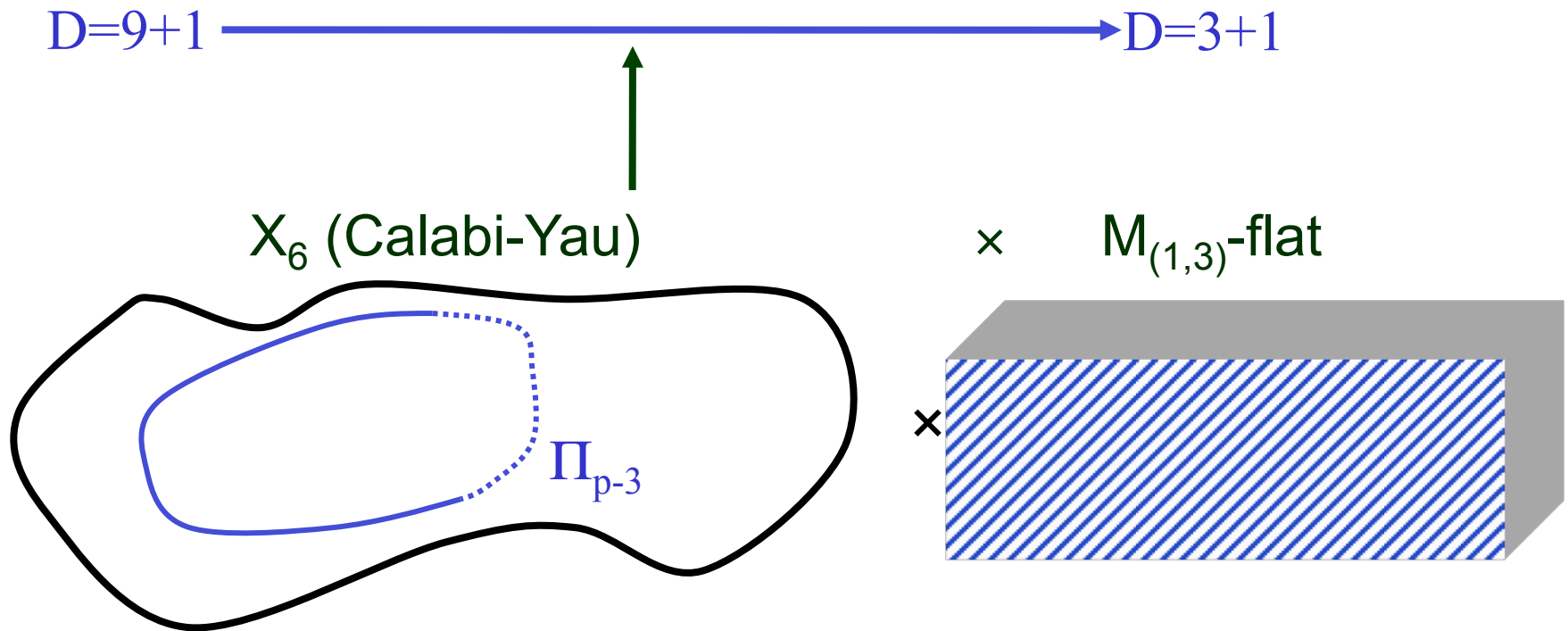


$\times$   $M_{(1,3)}$ -flat





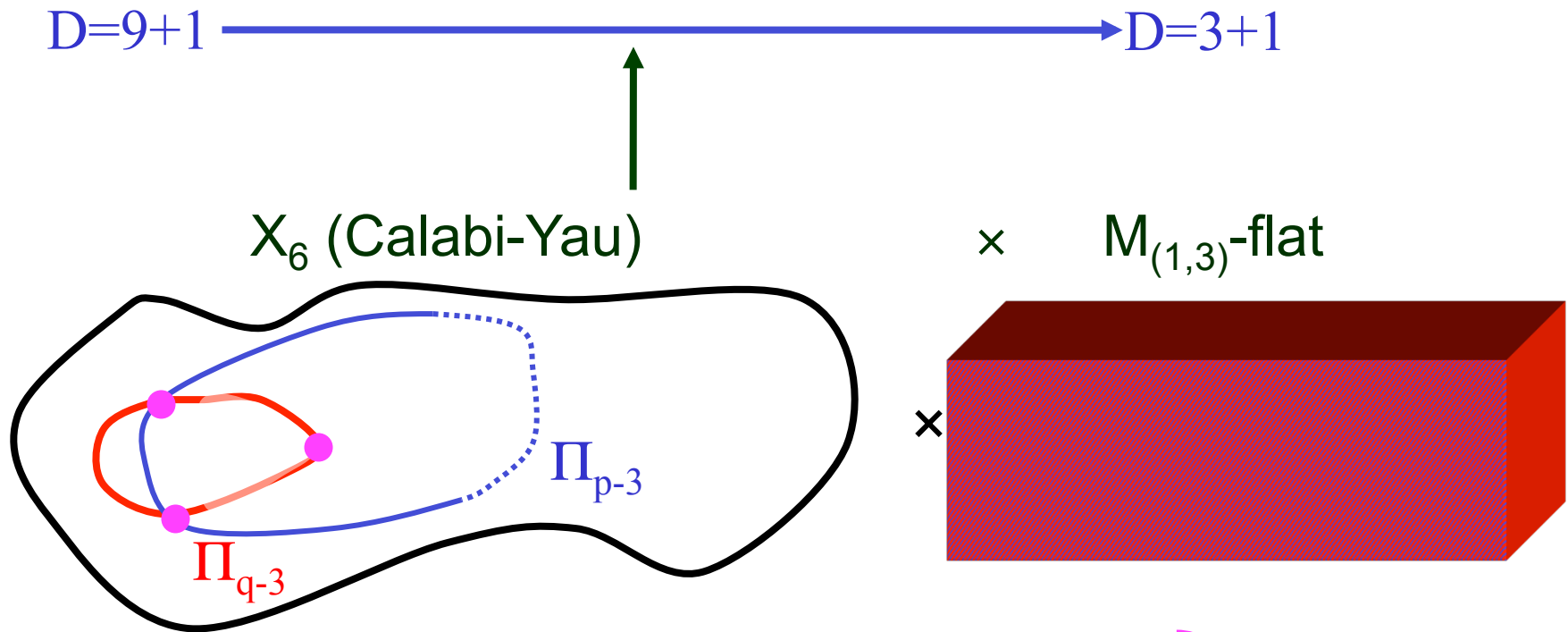
# Compactification



D p-branes – extend in  $p+1$  dimensions:

3+1-our world  $M_{(3,1)}$ ;  $(p-3)$ -wrap  $\Pi_{p-3}$  cycles of  $X_6$

# Compactification



D p-branes – extend in  $p+1$  dimensions:  
 3+1-our world  $M_{(3,1)}$ ;  $(p-3)$ -wrap  $\Pi_{p-3}$  cycles of  $X_6$

D q-branes – extend in  $q+1$  dimensions:  
 3+1-our world  $M_{(3,1)}$ ;  $(q-3)$ -wrap  $\Pi_{q-3}$  cycles of  $X_6$

$\Pi_{q-3} \cap \Pi_{p-3}$   
 Two sets of  
 D-branes  
 typically  
 intersect at  
 # of points

matter at each intersection  $\rightarrow$  family replication Geometric!

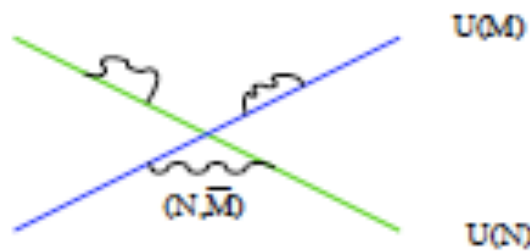
## Type II w/ D-branes →

fertile ground for particle physics model building

[Pedagogical review: TASI'10 lectures, M.C., Halverson arXiv:1101.2907]

Type IIA w/intersecting D-branes → key features of SM & SU(5) GUT spectrum: non-Abelian gauge symmetry, chirality & family replication

Geometric



Representation	Multiplicity
$(\bar{a}, b)$	$\pi_a \circ \pi_b$
$(\bar{a}, \bar{b})$	$\pi_a \circ \pi'_b$
$\begin{array}{ c } \hline \square_a \\ \hline \end{array}$	$\frac{1}{2} (\pi'_a \circ \pi_a + \pi_{O6} \circ \pi_a)$
$\begin{array}{ c } \hline \square_a \\ \hline \end{array}$	$\frac{1}{2} (\pi'_a \circ \pi_a - \pi_{O6} \circ \pi_a)$

In addition global consistency : Gauss's law for D-brane charge on internal space. → stringy



Large classes (order of 100's) of globally consistent supersymmetric, SM-like & GUT constructions; also couplings

[M.C. ,Shiu, Uranga'01]...

[M.C. Papadimitriou '03],  
[Cremades, Ibáñez, Marchesano'03]...

Recent developments:

New types of D-instantons: introduced to generate certain perturbatively absent couplings for charged sector matter

[Blumenhagen, M.C., Weigand, hep-th/0609191]

[Ibañez, Uranga, hep-th/0609213]

- charged matter coupling corrections

...

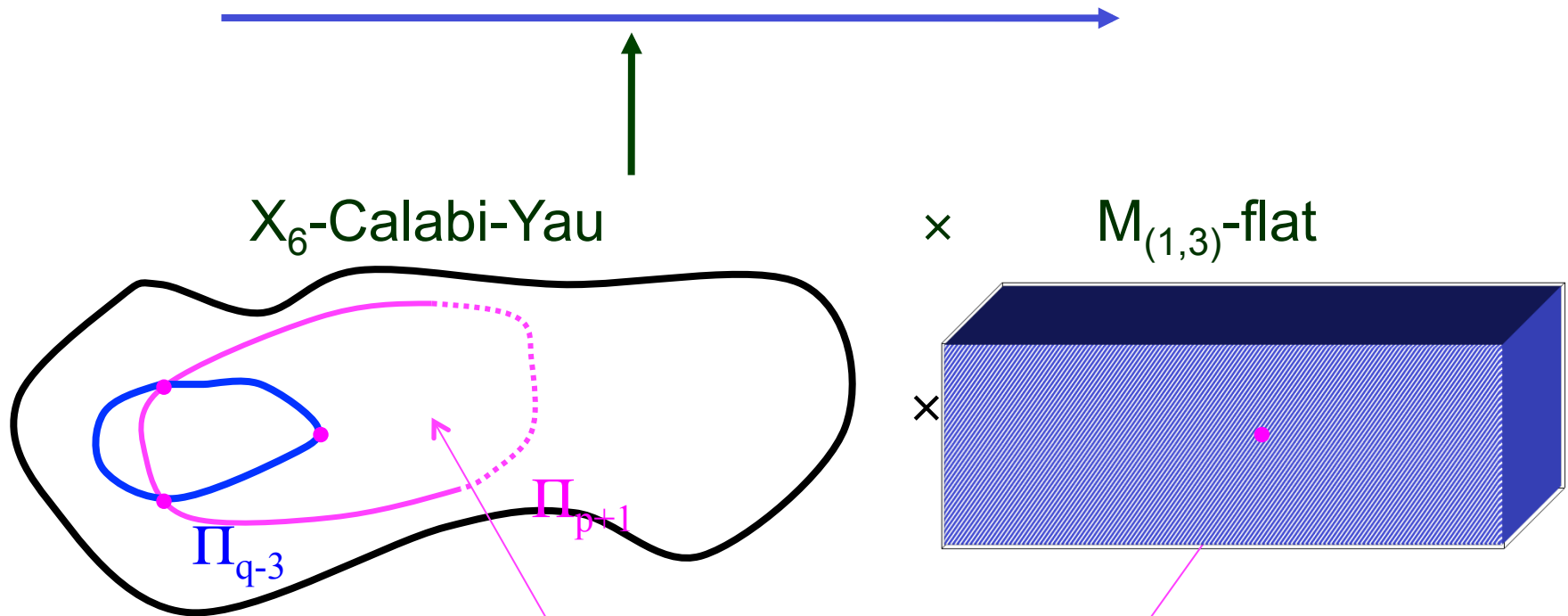
[Florea, Kachru, McGreevy, Saulina, 0610003]

- supersymmetry breaking

Review: [Blumenhagen, M.C., Kachru, Weigand, 0902.3251]

Encoded in non-perturbative violation of ``anomalous'' U(1)'s

# D-Instanton - Euclidean D-brane background



Wraps cycle  $\Pi_{p+1}$  cycles of  $X_6$  point-in 3+1 space-time

New geometric hierarchies for couplings:

$$\text{Re}(e^{-S_{E2}}) = e^{-\frac{2\pi}{\ell_s^3 g_s} \text{Vol}_{E2}} = e^{-\frac{2\pi}{\alpha_{\text{GUT}}} \frac{\text{Vol}_{E2}}{\text{Vol}_{D6}}} \quad \text{stringy!}$$

Instanton can intersect with physical D-brane (charged zero modes)

→ generate non-perturbative couplings of charged matter

## Specific Examples:

- i Majorana neutrino masses original papers, ...
- ii Nonpert. Dirac neutrino masses [M.C., Langacker, 0803.2876]

### iii 10 10 5 GUT couplings

[Blumenhagen, M.C., Lüst, Richter, Weigand, 0707.1871]

one-instanton effect  $\longrightarrow g_s \rightarrow 1$  (M-theory on  $G_2$ )

### iv Polonyi-type couplings $\longrightarrow$

[Aharony, Kachru, Silverstein 0708.0493] [MC, Weigand 0711.0209, 0807.3953]

[Heckman, Marsano, Saulina, Schafer-Nameki, Vafa 0808.1286] ...

Examples of such instanton induced hierarchical couplings primarily for local Type IIA toroidal orbifolds  $SU(5)$  GUT's



Challenge: global models  $\rightarrow$  Type I/IIB/F-theory (algebr. geom.)

- i. Type I GUT's on compact elliptically fibered Calabi-Yau  
First global chiral (four-family)  $SU(5)$  GUT's w/ D-instanton generated Polonyi & Majorana neutrino masses

[M.C., T.Weigand, 0711.0209, 0807.3953]

- ii. Global Type IIB GUT's :  $10^{10} 5_H$  non-perturbative coupling  
(two family)  $SU(5)$  GUT on CY as hypersurface in toric variety

[Blumenhagen, Grimm, Jurke, Weigand, 0811.2938]

- iii. Global F-theory lift [M.C., I. Garcia-Etxebarria, J. Halverson, 003.5337]

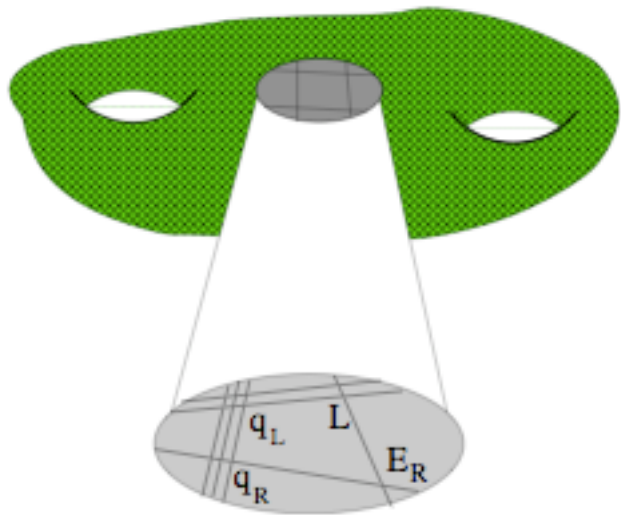
[Develop a code to calculate zero modes/spectrum in Type IIB and F-theory on toric varieties; code w/ new efficient technique  $\rightarrow$

[Blumenhagen, Jurke, Rahn & Roschy, 1003.5217] ]

Most examples w/  $O(1)$  instantons addressed  $SU(5)$  GUT's  
How about directly Standard Model?

Adressed for local Madrid quiver [Ibanez,Richter, 0811.1583] 

Systematic Analysis of D-Instanton effects for MSSM's quivers  
(compatible with global constraints)



Landscape analysis of MSSM w/  
realistic fermion textures

[M.C., J. Halverson, R. Richter, 0905.3379;  
0909.4292; 0910.2239]

Stringy Weinberg operator neutrino masses  
(examples of low string scale)

[M.C., J. Halverson, P. Langacker, R. Richter, 1001.3148]

Singlet-extended MSSM landscape

[M.C. J. Halverson, P. Langacker, 1006.3341]

Bottom-up approach initiated [Aldazabal,Ibanez,Quevedo,Uranga'00]..

Related recent works: Specific 3-stack [Leontaris, 0903.3691]

Madrid quiver [Anastasopoulos, Kiritsis, Lionetto, 0905.3044]

$SU(5)$  GUT's [Kiritsis, Lennek, Schellekens, 0909.0271]...

MSSM at toric singularities: [Krippendorff, Dolan,Maharana,Quevedo,1002.1790, 1106.6039]



## Approach: Bottom-up quivers

Spectrum and couplings **geometric** 

efficient **classification of key physics**

[compatible w/ global constraints] **→stringy**, but without delving into specifics of globally defined string compactifications]

Quiver data: massless **spectrum &**

examination of **couplings**

[both **perturbative & non-perturbative-instantons**]

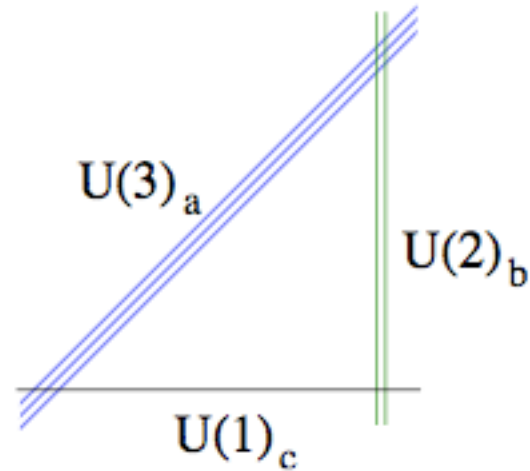


Probe “**quiver landscape**”

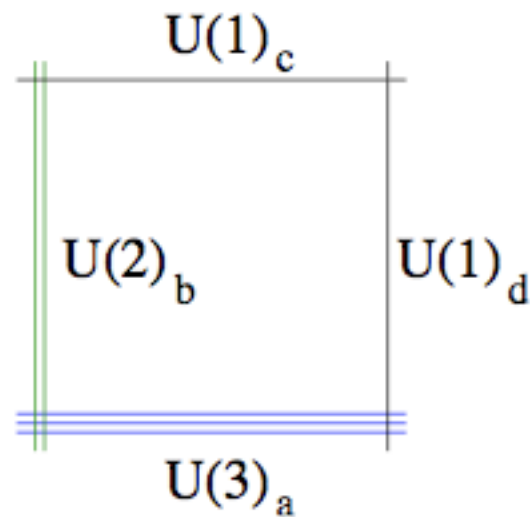
to identify **realistic quivers** in the landscape of string vacua

## Multi-stack MSSM quivers

Employ three-stack MSSM  $U(3)_a \times U(2)_b \times U(1)_c$



& four-stack MSSM  $U(3)_a \times U(2)_b \times U(1)_c \times U(1)_d$



five-stack....

# Four-stack set of MSSM models w/ 3 $N_R$ & potentially viable fermion textures

[M.C., J. Halverson, R. Richter, 0905.3379]

Solution #	$q_L$		$d_R$			$u_R$		$L$			$E_R$		$N_R$			$H_u$				$H_d$		
	$(a, b)$	$(\bar{a}, \bar{b})$	$(\bar{a}, c)$	$(\bar{a}, d)$	$\Gamma_a$	$(\bar{a}, \bar{c})$	$(\bar{a}, d)$	$(b, \bar{c})$	$(b, d)$	$(\bar{b}, d)$	$(c, d)$	$\Gamma_c$	$\Gamma_d$	$\Gamma_e$	$\Gamma_f$	$(c, d)$	$(\bar{c}, \bar{d})$	$(b, c)$	$(\bar{b}, c)$	$(b, d)$	$(\bar{b}, \bar{d})$	$(b, \bar{c})$
1	3	0	3	0	0	0	3	0	0	3	0	0	3	2	0	0	1	0	0	0	1	1
2	3	0	2	0	1	0	3	0	0	3	0	0	3	2	0	0	1	0	0	0	1	1
3	3	0	1	0	2	0	3	0	0	3	0	0	3	2	0	0	1	0	0	0	1	1
4	3	0	0	0	3	0	3	0	0	3	0	0	3	2	0	0	1	0	0	0	1	1
5	3	0	0	0	3	0	3	0	0	3	0	0	3	3	0	0	0	1	0	0	0	1
6	3	0	3	0	0	2	1	0	0	3	0	2	1	2	0	1	0	0	0	0	1	1
7	3	0	3	0	0	3	0	0	0	3	2	1	0	2	0	0	1	0	1	0	0	1
8	3	0	3	0	0	3	0	0	0	3	0	2	1	2	0	0	1	0	1	0	0	1
9	3	0	3	0	0	3	0	0	0	3	1	2	0	2	0	1	0	0	1	0	0	1
10	2	1	3	0	0	1	2	0	0	3	0	0	3	0	0	0	3	1	0	0	0	1
11	2	1	3	0	0	1	2	0	0	3	3	0	0	0	0	3	0	1	0	0	0	1
12	2	1	3	0	0	3	0	0	0	3	3	0	0	0	0	0	3	1	0	0	0	1
13	2	1	3	0	0	3	0	0	1	2	3	0	0	0	0	0	3	0	1	0	0	1
14	1	2	3	0	0	3	0	0	3	0	3	0	0	0	0	0	3	1	0	0	0	1
15	0	3	0	3	0	0	3	3	0	0	0	1	2	0	3	0	0	0	0	1	0	1
16	0	3	0	0	3	0	3	0	3	0	0	0	3	0	3	0	0	1	0	0	0	1
17	0	3	0	0	3	0	3	1	2	0	1	0	2	0	3	0	0	0	0	1	0	1
18	0	3	0	0	3	0	3	3	0	0	2	0	1	0	3	0	0	0	0	1	0	1
19	0	3	0	0	3	0	3	3	0	0	0	1	2	0	3	0	0	0	0	1	0	1
20	0	3	0	0	3	1	2	1	2	0	2	0	1	0	3	0	0	1	0	0	0	1
21	0	3	0	0	3	1	2	1	2	0	0	1	2	0	3	0	0	1	0	0	0	1
22	0	3	0	0	3	1	2	3	0	0	3	0	0	0	3	0	0	1	0	0	0	1
23	0	3	0	0	3	1	2	3	0	0	1	1	1	0	3	0	0	1	0	0	0	1
24	0	3	0	0	3	2	1	0	3	0	3	0	0	0	3	0	0	1	0	0	0	1
25	0	3	0	0	3	2	1	0	3	0	1	1	1	0	3	0	0	1	0	0	0	1
26	0	3	0	0	3	2	1	2	1	0	2	1	0	0	3	0	0	1	0	0	0	1
27	0	3	0	0	3	2	1	2	1	0	0	2	1	0	3	0	0	1	0	0	0	1
28	0	3	0	3	0	3	0	3	0	0	0	3	0	0	3	0	0	1	0	0	0	1
29	0	3	0	2	1	3	0	3	0	0	0	3	0	0	3	0	0	1	0	0	0	1
30	0	3	0	1	2	3	0	3	0	0	0	3	0	0	3	0	0	1	0	0	0	1
31	0	3	0	0	3	3	0	1	2	0	1	2	0	0	3	0	0	1	0	0	0	1

Solutions for hypercharge  $U(1)_Y = \frac{1}{6}U(1)_a + \frac{1}{2}U(1)_c - \frac{1}{2}U(1)_d$  - Madrid embedding

# Concrete 5-stack model (benchmark)

(w/ three mass scales in top, bottom in charged lepton sector)

Sector	Matter Fields	Transformation	Multiplicity	Hypercharge
$ab$	$q_L$	$(a, \bar{b})$	1	$\frac{1}{6}$
$ab'$	$q_L$	$(a, b)$	2	$\frac{1}{6}$
$ac'$	$u_R$	$(\bar{a}, \bar{c})$	2	$-\frac{2}{3}$
$ad'$	$u_R$	$(\bar{a}, \bar{d})$	1	$-\frac{2}{3}$
$aa'$	$d_R$	$\begin{array}{ c } \hline \square \\ \hline \end{array}_a$	3	$\frac{1}{3}$
$bc'$	$H_u$	$(b, c)$	1	$\frac{1}{2}$
$bd'$	$L$	$(\bar{b}, \bar{d})$	3	$-\frac{1}{2}$
$be'$	$H_d$	$(\bar{b}, \bar{e})$	1	$\frac{1}{2}$
$ce'$	$E_R$	$(c, e)$	2	1
$ce$	$N_R$	$(\bar{c}, e)$	1	0
$dd'$	$E_R$	$\begin{array}{ c c } \hline \square & \square \\ \hline \end{array}_d$	1	1
$de$	$N_R$	$(\bar{d}, e)$	2	0

Allows for full (inter- & intra-) family mass hierarchy via “factorization of Yukawa matrices” due to vector-pairs of zero fermion modes-stringy (technical, no time)

## Recent: String constraints & matter beyond the MSSM

[M.C., J. Halverson, P. Langacker, 1108.5387]

### I. Classify ALL possible MSSM quivers (three & four stacks)

irrespective of global conditions → most quivers inconsistent

What is additional matter to be compatible w/ global constraints?

→ stringy inputs on exotic matter

**3-stack analysis:** global conditions ( $T_{a,b,c}=0$ ) constraining, e.g., MSSM w/

$$U(1)_Y = \frac{1}{6}U(1)_a + \frac{1}{2}U(1)_c \quad T_a = 0 \quad T_b = \pm 2n \quad T_c = 0 \text{ mod } 3 \quad \text{with } n \in \{0, \dots, 7\},$$

w/ preferred additions: quasi-chiral Higgs pairs, MSSM singlets

hypercharge-less SU(2) triplets, &

various quark anti-quark pairs, all w/ integer el. ch.;

one (massless) Z' quiver

**4-stack analysis:** richer structure

sizable number of quivers w/ Z', including leptophobic (tuned);

additional structures: possible  $SH_{\underline{u}}H_{\underline{d}}$ ;  $\nu$ -masses;

exotics w/ fractional el. ch. ...

## Possible additions

Transformation	$T_a$	$T_b$	$T_c$	$M_a$	$M_b$	$M_c$
$\boxplus_a \quad (6, 1)_{\frac{1}{3}}$	7	0	0	$-\frac{1}{2}$	0	0
$\boxminus_a \quad (\bar{6}, 1)_{-\frac{1}{3}}$	-7	0	0	$\frac{1}{2}$	0	0
$\boxplus_a \quad (3, 1)_{\frac{1}{3}}$	-1	0	0	$-\frac{1}{2}$	0	0
$\boxminus_a \quad (3, 1)_{-\frac{1}{3}}$	1	0	0	$\frac{1}{2}$	0	0
$\boxplus_b \quad (1, 3)_0$	0	6	0	0	0	0
$\boxminus_b \quad (1, 3)_0$	0	-6	0	0	0	0
$\boxplus_b \quad (1, 1)_0$	0	-2	0	0	0	0
$\boxminus_b \quad (1, 1)_0$	0	2	0	0	0	0
$(\bar{b}, c) \quad (1, 2)_{\frac{1}{2}}$	0	-1	2	0	$-\frac{1}{2}$	$-\frac{1}{3}$
$(b, \bar{c}) \quad (1, 2)_{-\frac{1}{2}}$	0	1	-2	0	$\frac{1}{2}$	$\frac{1}{3}$
$(b, c) \quad (1, 2)_{\frac{1}{2}}$	0	1	2	0	$-\frac{1}{2}$	$-\frac{1}{3}$
$(\bar{b}, \bar{c}) \quad (1, 2)_{-\frac{1}{2}}$	0	-1	-2	0	$\frac{1}{2}$	$\frac{1}{3}$
$\boxplus_c \quad (1, 1)_1$	0	0	5	0	0	$-\frac{4}{3}$
$\boxminus_c \quad (1, 1)_{-1}$	0	0	-5	0	0	$\frac{4}{3}$
$(a, \bar{b}) \quad (3, 2)_{\frac{1}{6}}$	2	-3	0	0	$-\frac{1}{2}$	0
$(\bar{a}, b) \quad (\bar{3}, 2)_{-\frac{1}{6}}$	-2	3	0	0	$\frac{1}{2}$	0
$(a, b) \quad (3, 2)_{\frac{1}{6}}$	2	3	0	0	$-\frac{1}{2}$	0
$(\bar{a}, \bar{b}) \quad (\bar{3}, 2)_{-\frac{1}{6}}$	-2	-3	0	0	$\frac{1}{2}$	0
$(a, \bar{c}) \quad (3, 1)_{-\frac{1}{3}}$	1	0	-3	$\frac{1}{2}$	0	0
$(\bar{a}, c) \quad (\bar{3}, 1)_{\frac{1}{3}}$	-1	0	3	$-\frac{1}{2}$	0	0
$(a, c) \quad (3, 1)_{\frac{2}{3}}$	1	0	3	$-\frac{1}{2}$	0	-1
$(\bar{a}, \bar{c}) \quad (\bar{3}, 1)_{-\frac{2}{3}}$	-1	0	-3	$\frac{1}{2}$	0	1

● 105

3-node quivers ( $\leq 5$  additions)

Multiplicity	Matter Additions				
4	$\boxplus_b, (1, 3)_0$	$\boxplus_b, (1, 3)_0$	$\boxplus_b, (1, 1)_0$	$(a, \bar{b}), (3, 2)_{\frac{1}{6}}$	$(\bar{a}, \bar{b}), (3, 2)_{-\frac{1}{6}}$
4	$\boxplus_b, (1, 3)_0$	$\boxplus_b, (1, 1)_0$			
4	$\boxplus_b, (1, 3)_0$	$\boxplus_b, (1, 1)_0$			
4	$\boxplus_b, (1, 3)_0$	$\boxplus_b, (1, 1)_0$	$\boxplus_b, (1, 1)_0$	$(b, \bar{c}), (1, 2)_{-\frac{1}{2}}$	$(b, c), (1, 2)_{\frac{1}{2}}$
4	$\boxplus_b, (1, 3)_0$	$\boxplus_b, (1, 1)_0$	$\boxplus_b, (1, 1)_0$	$(b, \bar{c}), (1, 2)_{-\frac{1}{2}}$	$(b, c), (1, 2)_{\frac{1}{2}}$
4	$\boxplus_b, (1, 3)_0$	$\boxplus_b, (1, 1)_0$	$\boxplus_b, (1, 1)_0$	$(a, \bar{b}), (3, 2)_{\frac{1}{6}}$	$(\bar{a}, \bar{b}), (3, 2)_{-\frac{1}{6}}$
4	$\boxplus_b, (1, 1)_0$	$\boxplus_b, (1, 1)_0$			
4	$\boxplus_b, (1, 1)_0$	$(b, \bar{c}), (1, 2)_{-\frac{1}{2}}$	$(b, c), (1, 2)_{\frac{1}{2}}$		
4	$(b, \bar{c}), (1, 2)_{-\frac{1}{2}}$	$(b, \bar{c}), (1, 2)_{-\frac{1}{2}}$	$(b, c), (1, 2)_{\frac{1}{2}}$	$(b, c), (1, 2)_{\frac{1}{2}}$	
4	$(a, \bar{b}), (3, 2)_{\frac{1}{6}}$	$\boxplus_a, (3, 1)_{\frac{1}{3}}$	$(b, \bar{c}), (1, 2)_{-\frac{1}{2}}$	$(\bar{a}, \bar{c}), (3, 1)_{-\frac{2}{3}}$	$\boxplus_c, (1, 1)_1$
4	$\boxplus_b, (1, 3)_0$	$\boxplus_b, (1, 1)_0$	$\boxplus_b, (1, 1)_0$	$\boxplus_b, (1, 1)_0$	$\boxplus_b, (1, 1)_0$
4	$\boxplus_b, (1, 3)_0$	$\boxplus_b, (1, 1)_0$	$\boxplus_b, (1, 1)_0$	$\boxplus_b, (1, 1)_0$	$\boxplus_b, (1, 1)_0$
4	$\boxplus_b, (1, 3)_0$	$\boxplus_b, (1, 1)_0$	$\boxplus_b, (1, 1)_0$		
4	$\boxplus_b, (1, 3)_0$	$\boxplus_b, (1, 1)_0$	$(b, \bar{c}), (1, 2)_{-\frac{1}{2}}$	$(b, c), (1, 2)_{\frac{1}{2}}$	
4	$\boxplus_b, (1, 3)_0$	$(b, \bar{c}), (1, 2)_{-\frac{1}{2}}$	$(b, \bar{c}), (1, 2)_{-\frac{1}{2}}$	$(b, c), (1, 2)_{\frac{1}{2}}$	$(b, c), (1, 2)_{\frac{1}{2}}$
4	$\boxplus_b, (1, 1)_0$				
4	$\boxplus_b, (1, 1)_0$	$\boxplus_b, (1, 1)_0$	$(b, \bar{c}), (1, 2)_{-\frac{1}{2}}$	$(b, c), (1, 2)_{\frac{1}{2}}$	
4	$\boxplus_b, (1, 3)_0$	$\boxplus_b, (1, 3)_0$	$\boxplus_b, (1, 1)_0$	$\boxplus_b, (1, 1)_0$	
4	$\boxplus_b, (1, 3)_0$	$\boxplus_b, (1, 3)_0$	$\boxplus_b, (1, 1)_0$	$(b, \bar{c}), (1, 2)_{-\frac{1}{2}}$	$(b, c), (1, 2)_{\frac{1}{2}}$
4	$\boxplus_b, (1, 1)_0$	$\boxplus_b, (1, 1)_0$	$\boxplus_b, (1, 1)_0$	$\boxplus_b, (1, 1)_0$	

Multiplicity	Matter Additions				
4	$\overline{\square}b, (1, 3)_0$	$\overline{\square}b, (1, 3)_0$	$\overline{\square}b, (1, 3)_0$	$\overline{\square}b, (1, 1)_0$	$\overline{\square}b, (1, 1)_0$
4	$\overline{\square}b, (1, 3)_0$	$\overline{\square}b, (1, 3)_0$	$\square b, (1, 1)_0$		
1	$\square a, (\overline{3}, 1)_{\frac{1}{3}}$	$\square b, (1, 3)_0$	$\square b, (1, 1)_0$	$(a, \overline{c}), (3, 1)_{-\frac{1}{3}}$	
1	$\overline{\square}a, (3, 1)_{-\frac{1}{3}}$	$\square b, (1, 3)_0$	$\square b, (1, 1)_0$	$(\overline{a}, c), (\overline{3}, 1)_{\frac{1}{3}}$	
1	$\square a, (3, 1)_{\frac{1}{3}}$	$\overline{\square}b, (1, 3)_0$	$\square b, (1, 1)_0$	$(a, \overline{c}), (3, 1)_{-\frac{1}{3}}$	
1	$\overline{\square}a, (3, 1)_{-\frac{1}{3}}$	$\overline{\square}b, (1, 3)_0$	$\square b, (1, 1)_0$	$(\overline{a}, c), (\overline{3}, 1)_{\frac{1}{3}}$	
1	$\square a, (3, 1)_{\frac{1}{3}}$	$\overline{\square}b, (1, 1)_0$	$\overline{\square}b, (1, 1)_0$	$(a, \overline{c}), (3, 1)_{-\frac{1}{3}}$	
1	$\overline{\square}a, (3, 1)_{-\frac{1}{3}}$	$\overline{\square}b, (1, 1)_0$	$\overline{\square}b, (1, 1)_0$	$(\overline{a}, c), (\overline{3}, 1)_{\frac{1}{3}}$	
1	$\square a, (3, 1)_{\frac{1}{3}}$	$\overline{\square}b, (1, 1)_0$	$(b, \overline{c}), (1, 2)_{-\frac{1}{2}}$	$(b, c), (1, 2)_{\frac{1}{2}}$	$(a, \overline{c}), (3, 1)_{-\frac{1}{3}}$
1	$\overline{\square}a, (3, 1)_{-\frac{1}{3}}$	$\overline{\square}b, (1, 1)_0$	$(b, \overline{c}), (1, 2)_{-\frac{1}{2}}$	$(b, c), (1, 2)_{\frac{1}{2}}$	$(\overline{a}, c), (\overline{3}, 1)_{\frac{1}{3}}$
1	$(a, b), (3, 2)_{\frac{1}{6}}$	$(b, \overline{c}), (1, 2)_{-\frac{1}{2}}$	$(\overline{a}, c), (\overline{3}, 1)_{\frac{1}{3}}$	$(\overline{a}, \overline{c}), (\overline{3}, 1)_{-\frac{2}{3}}$	$\square c, (1, 1)_1$
1	$\square a, (3, 1)_{\frac{1}{3}}$	$\overline{\square}b, (1, 3)_0$	$\overline{\square}b, (1, 1)_0$	$\overline{\square}b, (1, 1)_0$	$(a, \overline{c}), (3, 1)_{-\frac{1}{3}}$
1	$\overline{\square}a, (3, 1)_{-\frac{1}{3}}$	$\overline{\square}b, (1, 3)_0$	$\overline{\square}b, (1, 1)_0$	$\overline{\square}b, (1, 1)_0$	$(\overline{a}, c), (\overline{3}, 1)_{\frac{1}{3}}$
1	$\square a, (3, 1)_{\frac{1}{3}}$	$\square a, (\overline{3}, 1)_{\frac{1}{3}}$	$\square b, (1, 1)_0$	$(a, \overline{c}), (3, 1)_{-\frac{1}{3}}$	$(a, \overline{c}), (3, 1)_{-\frac{1}{3}}$
1	$\overline{\square}a, (3, 1)_{-\frac{1}{3}}$	$\overline{\square}a, (3, 1)_{-\frac{1}{3}}$	$\square b, (1, 1)_0$	$(\overline{a}, c), (\overline{3}, 1)_{\frac{1}{3}}$	$(\overline{a}, c), (\overline{3}, 1)_{\frac{1}{3}}$
1	$\square a, (3, 1)_{\frac{1}{3}}$	$\square b, (1, 1)_0$	$(a, \overline{c}), (3, 1)_{-\frac{1}{3}}$		
1	$\overline{\square}a, (3, 1)_{-\frac{1}{3}}$	$\square b, (1, 1)_0$	$(\overline{a}, c), (\overline{3}, 1)_{\frac{1}{3}}$		
1	$\square a, (3, 1)_{\frac{1}{3}}$	$\overline{\square}b, (1, 3)_0$	$\overline{\square}b, (1, 3)_0$	$\square b, (1, 1)_0$	$(a, \overline{c}), (3, 1)_{-\frac{1}{3}}$
1	$\overline{\square}a, (3, 1)_{-\frac{1}{3}}$	$\overline{\square}b, (1, 3)_0$	$\overline{\square}b, (1, 3)_0$	$\square b, (1, 1)_0$	$(\overline{a}, c), (\overline{3}, 1)_{\frac{1}{3}}$



## II. Implications for MSSM's with additional nodes

Up-to two additional U(1)'s or one U(N)

Systematic search (implement global consistency conditions)

i) SM singlets by far the most common fields

& (light anomalous) U(1)'-monochromatic gamma ray line dark matter scenario

à la Dudas, Mambrini, Pokorski, Romagnoni 1205.1520

w/ coupling to SM automatically forbidden by anomalous U(1)

decay to Z  $\gamma$  possible (via Z' - B<sub>Y</sub> - B<sub>Y</sub> "Chern-Simons" vertex)

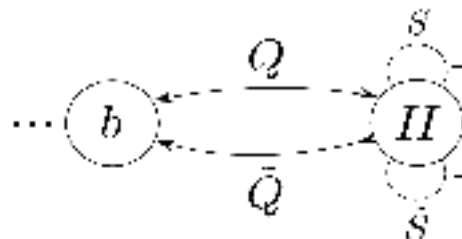
ii) E.g., Stringy dynamical SUSY breaking scenario (à la Fayet)

Aharony, Kachru, Silverstein '07

Quasi-chiral matter  $\rightarrow$  messenger masses instanton suppressed;

lifetime of metastable vacuum set by instanton superpotential

$$W \supset \lambda_{(0,0)} M_* S \tilde{S} + \lambda_{(-2,0)} M_* Q \tilde{Q} + \lambda_{(-2,0)} \frac{1}{M_*} S \tilde{S} Q \tilde{Q}$$



### III. Axiguons w/ quiver embedding $[SU(3)_{CL} \times SU(3)_{CR}]$

(playground to explain Tevatron  $t\bar{t}$  anomaly)

Accounting for the correct sign of the asymmetry requires a “mirror” generations, e.g.,

$$Q_{1,2} \sim (3, 1, 2)_{\frac{1}{6}} \quad Q_3 \sim (1, 3, 2)_{\frac{1}{6}}$$

**Quiver constraint: no tri-fundamental Higgs**  $\rightarrow$   
 need effective Yukawas (e.g.,  $A, C$ ) & w/ mirror generations  
 giving some perturbative Yukawas, (e.g.,  $P_u$ )

$$Y_{QH_u u^c} = \begin{pmatrix} A & A & P_u \\ A & A & P_u \\ P_u & P_u & C \end{pmatrix} \quad A = \frac{M_Q}{\lambda_Q v_\phi} \quad C = \frac{M_u}{\lambda_u v_\phi}$$

$\tilde{Q} H_u u_2^c, \lambda_Q \Phi \tilde{Q} \tilde{Q}^c, \text{ and } M_Q \tilde{Q}^c Q_2$

anomalous  $U(1)$ 's and instantons can help with **scales of  $A, C$**   
 (still major difficulties with Yukawa couplings)

Foresee further progress:

- a) DEVELOPMENT of TECHNIQUES! → generalize constructions to general Calabi Yau spaces (advanced algebraic geometry techniques);  
Fluxes in F-theory
- b) Quantitatively improve realistic model constructions, including further progress on non-perturbative effects

# Conclusions/Outlook

“Glimpses” of particle physics from String Theory  
Focus on D-branes → Type II (and F-theory)

- a) Progress: development of techniques for constructions  
sizable number of semi-realistic models
- b) “The devil is in the details!”-typically not fully realistic  
typically exotic matter → but viewed as a string prediction
- c) Systematic searches within local D-quivers w/SSM  
(compatible with string constraints)  
→ probe for SM exotics; nature of hidden; SUSY breaking...

Hopefully, LHC & Astrophysics experiments → inputs  
toward realistic particle physics from string theory  
w/ efforts presented here playing a role