

UNIVERSIDAD DE GUANAJUATO

DEPARTAMENTO DE FISICA

M. Sabido

Nonconmmutativity and Λ

Work in progress in collaboration S. Pérez-Payán, E. Mena
and C. Yee



XIII Mexican Workshop on
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PLAN OF THE TALK

INTRODUCTION, MOTIVATION

NC-GRAVITY

NC-COSMOLOGY

NC AND LAMBDA

FINAL REMARKS



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INTRODUCTION



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Noncommutativity of space time is an old idea (Snyder Phys. Rev. 71 (1947))



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There are physical systems that can easily be described by assuming that coordinates do not commute.



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An immense amount of work has been done in noncommutative gauge theory.



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Unfortunately on the gravity side of things, the story has been more complex.



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NC- GRAVITY



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NC- GRAVITY

It was shown from string theory that noncommutative space-time arises in the low energy limit, of open strings in the presence of NS B field. Using different regularization methods a commutative or a noncommutative gauge theory arises.



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Consider a theory invariant under the action of a Lie group G , with gauge group A_μ and a mater field Φ , that transforms in the adjoint representation



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$$\delta_\lambda A_\mu = \partial_\mu \lambda + i [\lambda, A_\mu], \quad \delta_\lambda \Phi = i [\lambda, \Phi],$$



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For the noncomutative theory, the gauge transformations are:



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For the noncommutative theory, the gauge transformations are:

$$\delta_{\hat{\lambda}} \hat{A}_\mu = \partial_\mu \hat{\lambda} + i [\hat{\lambda} *; \hat{A}_\mu], \quad \delta_{\hat{\lambda}} \hat{\Phi} = i [\hat{\lambda} *; \hat{\Phi}].$$



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NC- GRAVITY

Because the result should be independent of the regularization scheme both theories should be equivalent the a map between them must exist. This is the Seiberg-Witten map, which lets us write the noncommutative gauge fields from the commutative ones.



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$$\widehat{\lambda}(\lambda, A) = \lambda + \frac{1}{4}\theta^{\mu\nu} \{\partial_\mu \lambda, A_\nu\} + \mathcal{O}(\theta^2),$$



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From these equations we can write the noncommutative lagrangian and the noncommutative theory.



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NC- GRAVITY

$$I = \int d^4x e_a^\mu e_b^\nu R_{ab}^{\mu\nu}$$



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NC- GRAVITY

$$I = \int d^4x e_a^\mu e_b^\nu R_{ab}^{\mu\nu}$$

GR



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GR

$$\widehat{I}_{\theta^2} = \frac{1}{2^4} \theta^{\gamma\delta} \theta^{\tau\xi} \int dx^4 \left\{ 4e \left[4R_\delta{}^\rho (R_{\rho\tau}{}^{ab} R_{\gamma\xi ab} - \omega_\tau{}^{ab} \partial_\xi R_{\rho\gamma ab}) + \omega_\gamma{}^{\rho\sigma} \partial_\tau \omega_\delta{}^{ab} \partial_\xi R_{\rho\sigma ab} \right] \right.$$



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MOTIVATION



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MOTIVATION

We need an alternative



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MOTIVATION

We need an alternative

Minisuperspace Models

+



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MOTIVATION

We need an alternative

Minisuperspace Models

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Canonical Quantization

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Noncommutative Quantum Mechanics



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Noncommutative Quantum Cosmology

H. Compeán, O. Obregon and C. Ramirez PRL 88 (2002)



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MOTIVATION

In quantum mechanics we have

$$[x, P] = i\hbar \quad \rightarrow \quad \Delta x \Delta P > \frac{\hbar}{2}$$



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Being gravity the theory of space time, quantization gravity could be interpreted as a quantization of space time



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What happens if space time does not commute?



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MOTIVATION

In analogy with quantum mechanics



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MOTIVATION

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$$[x, y] = i\theta \quad \rightarrow \quad \Delta x \Delta y > \frac{\theta}{2}$$



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MOTIVATION

In analogy with quantum mechanics

$$[x, y] = i\theta \quad \rightarrow \quad \Delta x \Delta y > \frac{\theta}{2}$$

We have a minimal scale we can probe



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This deformation can be encoded in Moyal product of functions



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$$f(x, y) * g(x, y) = f(x, y) e^{\left(i\frac{\theta}{2} \left(\overleftarrow{\partial}_x \overrightarrow{\partial}_y - \overleftarrow{\partial}_y \overrightarrow{\partial}_x\right)\right)} g(x, y).$$



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which yields the noncommutative Schrodinger equation

(L. Mezincescu ArXiv:hep-th\0007046)



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$$\left[-\nabla^2 + V(x, y)\right] * \Psi(x, y) = E\Psi(x, y)$$



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NONCOMMUTATIVE COSMOLOGY



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NONCOMMUTATIVE COSMOLOGY

We may recover the original commutation relations with the variables



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NONCOMMUTATIVE COSMOLOGY

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$$x \rightarrow x + \frac{1}{2}\theta P_y, \quad y \rightarrow y - \frac{1}{2}\theta P_x$$

so the effects are encoded only on the potential



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NONCOMMUTATIVE COSMOLOGY

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$$x \rightarrow x + \frac{1}{2}\theta P_y, \quad y \rightarrow y - \frac{1}{2}\theta P_x$$

so the effects are encoded only on the potential

$$V(x, y) * \Psi(x, y) = V\left(x + \frac{\theta}{2}P_y, y - \frac{\theta}{2}P_x\right),$$



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$$\left[-\nabla^2 + V\left(x + \frac{\theta}{2}P_y, y - \frac{\theta}{2}P_x\right)\right] \Psi(x, y) = E\Psi(x, y).$$



PRELIMINARIES: NC CLASSICAL MECHANICS



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PRELIMINARIES: NC CLASSICAL MECHANICS

We still ask our selfs for the classical model from which this quantum mechanics arises.

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$$\{x_i, x_j\} = \theta_{ij}, \quad \{P_{x_i}, P_{x_j}\} = 0 \quad \{x_i, P_{x_j}\} = \delta_{ij}.$$



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This can be constructed from Hamiltonian Manifolds [\(W. Guzman, M.S., J. Socorro PLB 2011\)](#)



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PRELIMINARIES: NC CLASSICAL MECHANICS



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In local coordinates x^μ we have $X = X^\mu \partial_\mu$. We say is locally Hamiltonian if



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PRELIMINARIES: NC CLASSICAL MECHANICS



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PRELIMINARIES: NC CLASSICAL MECHANICS

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The canonical 2-form



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PRELIMINARIES: NC CLASSICAL MECHANICS



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PRELIMINARIES: NC CLASSICAL MECHANICS

A noncommutative classical mechanics, is a Hamiltonian manifold, for which the exact 2-form is given by a non degenerate ω_{nc} .



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$$\frac{dq^i}{dt} = \gamma^{ij} \frac{\partial H}{\partial p^j} + \theta^{ij} \frac{\partial H}{\partial q^j},$$
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NONCOMMUTATIVE COSMOLOGY



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NONCOMMUTATIVE COSMOLOGY

If we consider that the noncommutative field can be written as a series expansion on the noncommutative parameter, furthermore this expansion can be calculated from the full noncommutative theory.



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This is an effective noncommutativity between the minisuperspace fields, and is similar to the modification used in noncommutative quantum

mechanics. H. Compean, O. Obregón and C. Ramírez PRL 88 (2002), W. Guzmán, M.S, J. Socorro Phys. Rev D (2007), W. Guzmán, M.S, J. Socorro, Submitted PLB(2010))



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NONCOMMUTATIVE COSMOLOGY

One of a noncommutative mini-superspace point to the possibility of late time acceleration of the universe.

W. Guzmán, M.S. and J.Socorro PRD (2007).



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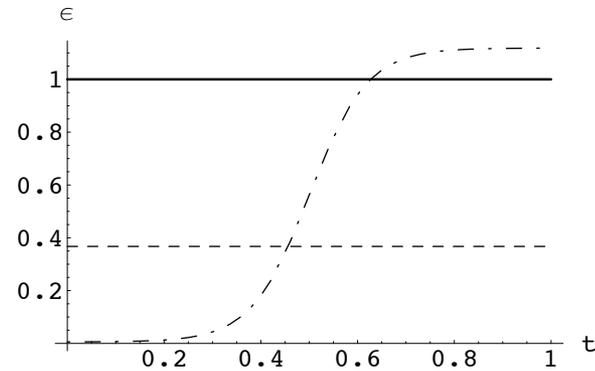
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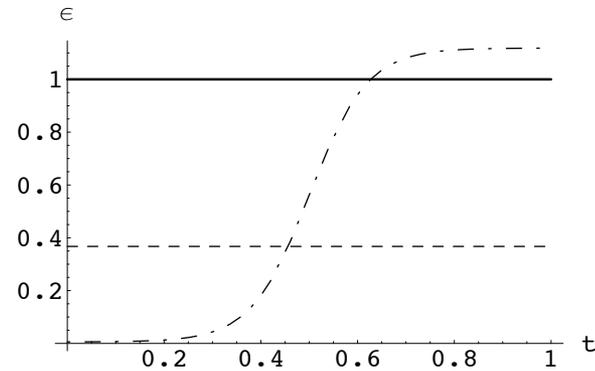
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$$\epsilon = \frac{\lambda^2}{2} \left[1 - \theta\mu \left(1 - \frac{\sqrt{6}}{\lambda} \tanh [3\mu(\lambda^2 - 6)t] \right) \right],$$



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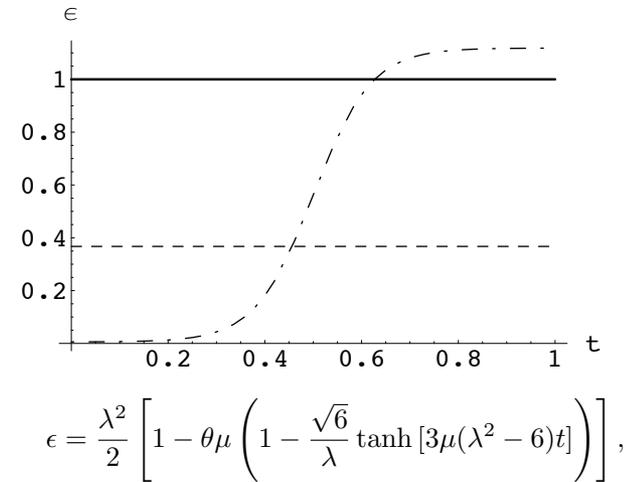
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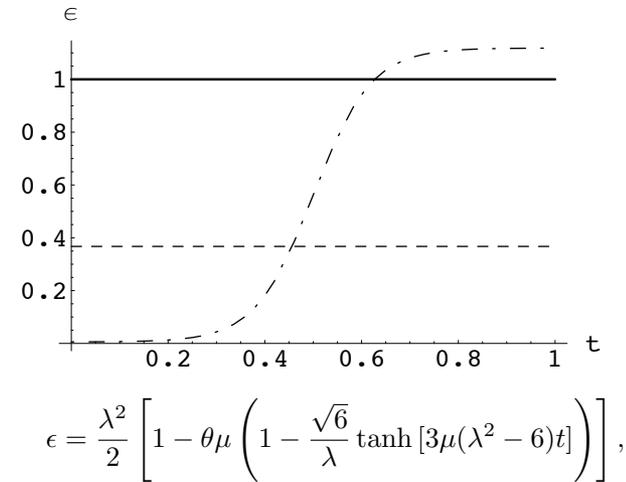
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Finally a dynamical systems analysis gave more evidence of this fact. I. Quiros and O. Obregón, PRD (2011).



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NONCOMMUTATIVE COSMOLOGY



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NONCOMMUTATIVE COSMOLOGY

Let us consider the action for a flat FRW universe a scalar field and cosmological constant as the matter content of the the model.



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NONCOMMUTATIVE COSMOLOGY

Let us consider the action for a flat FRW universe a scalar field and cosmological constant as the matter content of the the model.

$$S = \int dt \left\{ -\frac{3a\dot{a}^2}{\kappa^2 N} + a^3 \left(\frac{\dot{\phi}^2}{2N} + N\Lambda \right) \right\},$$



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$$S = \int \left[\frac{1}{2N} (\dot{x}^2 - \dot{y}^2) - N \frac{\omega^2}{2} (x^2 - y^2) \right] dt.$$



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NONCOMMUTATIVE COSMOLOGY



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NONCOMMUTATIVE COSMOLOGY

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Where ω^2 is proportional to Λ . For the noncommutative case, we impose noncommutativity among the minisuperspace coordinates and canonical momentum,



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$$\begin{aligned} \hat{x} &= x + \frac{\theta}{2} p_y, & \hat{y} &= y - \frac{\theta}{2} p_x, \\ \hat{p}_x &= p_x - \frac{\beta}{2} y, & \hat{p}_y &= p_y + \frac{\beta}{2} x, \end{aligned}$$



NONCOMMUTATIVE COSMOLOGY



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NONCOMMUTATIVE COSMOLOGY

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NONCOMMUTATIVE COSMOLOGY

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Where we defined

$$\omega_1^2 = \frac{4(\beta - \omega^2\theta)}{4 - \omega^2\theta^2}, \quad \omega_2^2 = \frac{4(\omega^2 - \beta^2/4)}{4 - \omega^2\theta^2},$$



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NONCOMMUTATIVE COSMOLOGY

The resulting noncommutative Hamiltonian for this model is

$$\hat{H} = \frac{1}{2} \{ \hat{p}_x^2 - \hat{p}_y^2 + \omega_1^2 (\hat{x}\hat{p}_y + \hat{y}\hat{p}_x) + \omega_2^2 (\hat{x}^2 - \hat{y}^2) \}.$$

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The equations of motion are

$$\begin{aligned} \dot{\hat{x}} &= \{ \hat{x}, \hat{H} \} = \frac{1}{2} [2\hat{p}_x + \omega_1^2 \hat{y}], & \dot{\hat{y}} &= \{ \hat{y}, \hat{H} \} = \frac{1}{2} [-2\hat{p}_y + \omega_1^2 \hat{x}], \\ \dot{\hat{p}}_x &= \{ \hat{p}_x, \hat{H} \} = \frac{1}{2} [-\omega_1^2 \hat{p}_y - 2\omega_2^2 \hat{x}], & \dot{\hat{p}}_y &= \{ \hat{p}_y, \hat{H} \} = \frac{1}{2} [-\omega_1^2 \hat{p}_x + 2\omega_2^2 \hat{y}], \end{aligned}$$



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NONCOMMUTATIVE COSMOLOGY

To solve the equations of motion we propose the change of variables



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NONCOMMUTATIVE COSMOLOGY

To solve the equations of motion we propose the change of variables $\eta = \hat{x} + \hat{y}$, $\zeta = \hat{y} - \hat{x}$



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NONCOMMUTATIVE COSMOLOGY

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NONCOMMUTATIVE COSMOLOGY

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An arrive to a simplified set of equations

$$\ddot{\eta} - \omega_1^2 \dot{\eta} + \frac{1}{4}(4\omega_2^2 + \omega_1^4)\eta = 0,$$

$$\ddot{\zeta} + \omega_1^2 \dot{\zeta} + \frac{1}{4}(4\omega_2^2 + \omega_1^4)\zeta = 0,$$



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This equations are very easy to solve for any value of Λ , θ and β .



NONCOMMUTATIVE COSMOLOGY

Let analyze the case when the cosmological constant is zero, as well as the parameter θ . In this case the solution is very simple, in particular the scale factor is

$$a^3(t) = V_0 \cosh^2 \left(\frac{1}{4} t \beta \right),$$

Comparing the to models we arrive to the relationship

$$\Lambda \sim \beta^2.$$

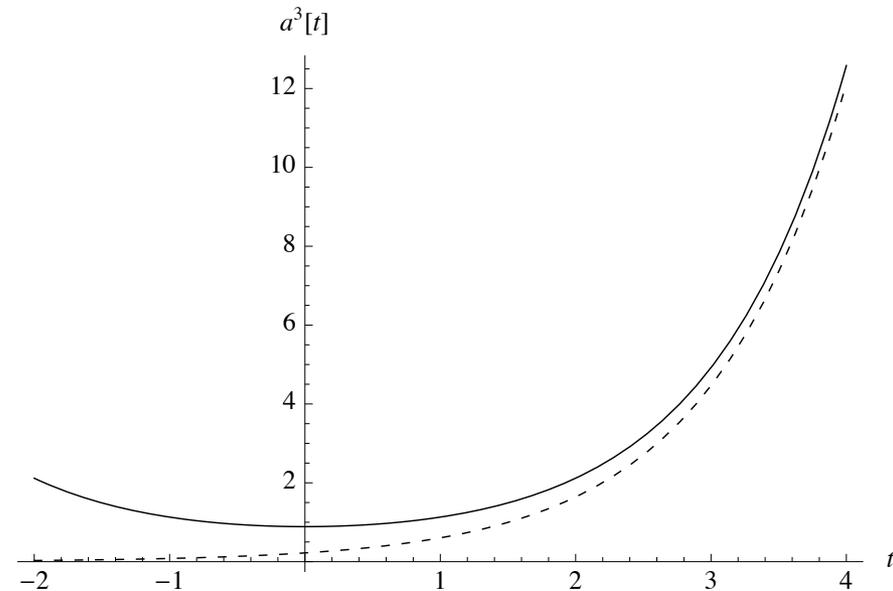


Figure 1: Dynamics of the phase space deformed model for the values $X_0 = Y_0 = 1, \delta_2 = \delta_1 = 0, \omega = 0$ and $\beta = 1$. The solid line corresponds to the volume of the universe, calculated with the noncommutative model. The dotted line corresponds to the volume of the de Sitter spacetime. For large values of t the behavior is the same.



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NONCOMMUTATIVITY AND LAMBDA

With this result lets find a model from which to derive Λ . We start with a (4+1) gravity.



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NONCOMMUTATIVITY AND LAMBDA

With this result lets find a model from which to derive Λ . We start with a (4+1) gravity.

$$I = \int \sqrt{-g} (R - \Lambda) dt d^3x d\rho,$$



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$$L = \frac{1}{2} \left(a\phi\dot{a}^2 + a^2\dot{a}\phi - \kappa a\phi + \frac{1}{3}\Lambda a^3\phi \right)$$



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NONCOMMUTATIVITY AND LAMBDA

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The hamiltonian for the model is, after an appropriate change of variables

$$H = \frac{1}{2} \left[(\dot{x}^2 + \omega^2 x^2) - (\dot{y}^2 + \omega^2 y^2) \right]$$



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NONCOMMUTATIVITY AND LAMBDA

Following the same steps as before we calculate the noncommutative hamiltonian



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NONCOMMUTATIVITY AND LAMBDA

Following the same steps as before we calculate the noncommutative hamiltonian

$$H = \left\{ \left(\hat{p}_u - \frac{2(\beta - \theta\omega^2)}{4 - \omega^2\theta^2} \hat{v} \right)^2 - \left(\hat{p}_v + \frac{2(\beta - \theta\omega^2)}{4 - \omega^2\theta^2} \hat{u} \right)^2 + \left(\frac{4(\beta - \theta\omega^2)^2}{(4 - \omega^2\theta^2)^2} + \frac{4(\omega^2 - \beta^2/4)}{4 - \omega^2\theta^2} \right) \hat{u}^2 - \left(\frac{4(\beta - \theta\omega^2)^2}{(4 - \omega^2\theta^2)^2} + \frac{4(\omega^2 - \beta^2/4)}{4 - \omega^2\theta^2} \right) \hat{v}^2 \right\}$$



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This can be written in simpler form if we define a 2 dimensional vector potential



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NONCOMMUTATIVITY AND LAMBDA

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This can be written in simpler form if we define a 2 dimensional vector potential

$$H = \left\{ \left[(\hat{p}_u - A_{\hat{u}})^2 + \omega'^2 \hat{u}^2 \right] - \left[(\hat{p}_v - A_{\hat{v}})^2 + \omega'^2 \hat{v}^2 \right] \right\}$$



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$$\omega'^2 \equiv \frac{4(\beta - \theta\omega^2)^2}{(4 - \omega^2\theta^2)^2} + \frac{4(\omega^2 - \beta^2/4)}{4 - \omega^2\theta^2}$$

$$A_{\hat{u}} \equiv \frac{-2(\beta - \theta\omega^2)}{4 - \omega^2\theta^2} \hat{v}, \quad A_{\hat{v}} \equiv \frac{2(\beta - \theta\omega^2)}{4 - \omega^2\theta^2} \hat{u}$$



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We can calculate the effective magnetic field



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NONCOMMUTATIVITY AND LAMBDA

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We can calculate the effective magnetic field

$$B = \frac{4(\beta - \omega^2\theta)}{4 - \omega^2\theta^2}$$



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NONCOMMUTATIVITY AND LAMBDA

We can now find the relationship between the commutative and noncommutative cosmological constants



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NONCOMMUTATIVITY AND LAMBDA

We can now find the relationship between the commutative and noncommutative cosmological constants

$$\tilde{\Lambda}_{eff} = \Lambda_{eff} + \frac{3\beta^2}{8}$$



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NONCOMMUTATIVITY AND LAMBDA

We can now find the relationship between the commutative and noncommutative cosmological constants

$$\tilde{\Lambda}_{eff} = \Lambda_{eff} + \frac{3\beta^2}{8}$$

So even if in the original theory $\Lambda_{eff}=0$, we find that the noncommutative universe has an effective Λ



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FINAL REMARKS



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FINAL REMARKS

- Using the ideas of noncommutative quantum mechanics and effective noncommutative in the minisuperspace is used for noncommutative quantum cosmology.



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FINAL REMARKS

- Using the ideas of noncommutative quantum mechanics and effective noncommutative in the minisuperspace is used for noncommutative quantum cosmology.
- There is a possibility that Λ is related to the noncommutative parameter. This was found in a simple cosmological model.



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FINAL REMARKS

- Using the ideas of noncommutative quantum mechanics and effective noncommutative in the minisuperspace is used for noncommutative quantum cosmology.
- There is a possibility that Λ is related to the noncommutative parameter. This was found in a simple cosmological model.
- A simple 5 dimensional model was proposed, where the origin of Λ is related to the noncommutativity between the compact dimension and the radius of the noncompact dimension.



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