

Reconciling Supersymmetry and Thermal Leptogenesis

Jörn Kersten

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Based on Jasper Hasenkamp, JK, [arXiv:1008.1740](https://arxiv.org/abs/1008.1740) [hep-ph]

Outline

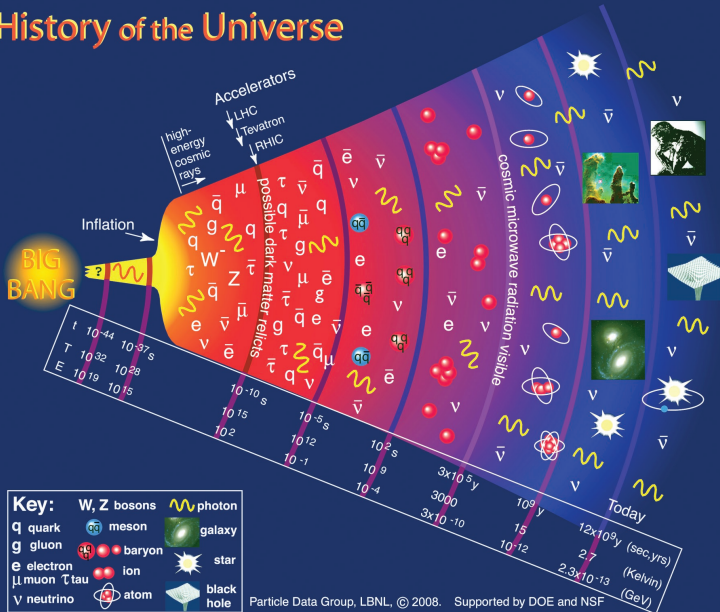
- 1 The Gravitino Problem
- 2 Entropy Production
- 3 Candidates for Entropy Producers

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History of the Universe



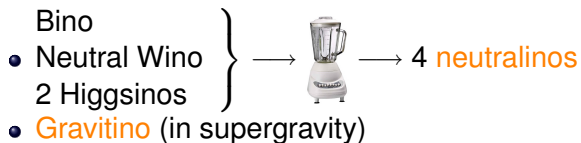
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- Symmetry between fermions and bosons
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different **spin**, other properties equal

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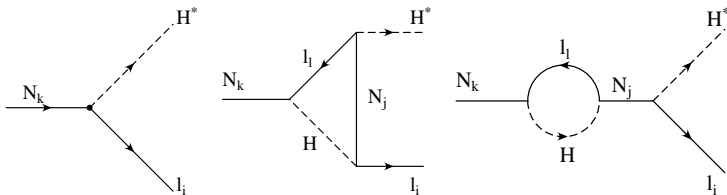
Supersymmetry

- Symmetry between fermions and bosons
- **Superpartner** for each Standard Model particle: different **spin**, other properties equal
 - Bino
 - Neutral Wino
 - 2 Higgsinos }  → 4 **neutralinos**
- **Gravitino** (in supergravity)
- Lightest superpartner (**LSP**) stable \leadsto **dark matter** candidate



Leptogenesis

- Gauge singlet neutrinos N
- Large Majorana masses $M_R \gtrsim 10^9 \text{ GeV}$
- Related to **light** neutrino masses: **see-saw mechanism**
- **C, CP violation** in decays



$$|\epsilon| = \frac{\Gamma(N \rightarrow \ell H) - \Gamma(N \rightarrow \bar{\ell} \bar{H})}{\Gamma(N \rightarrow \ell H) + \Gamma(N \rightarrow \bar{\ell} \bar{H})} < \frac{3}{16\pi} \frac{M_R \sqrt{\Delta m_{\text{atm}}^2}}{v^2}$$

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- Sphalerons (non-perturbative processes)
 \leadsto **baryon asymmetry** $\eta_B = \frac{n_B}{n_\gamma} \propto |\epsilon| < M_R \cdot \dots$
- Observed $\eta_B \sim 6 \cdot 10^{-10} \leadsto M_R \gtrsim 2 \cdot 10^9 \text{ GeV}$

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- **Thermal** leptogenesis: N produced thermally at $T > M_R$

$$T_R \gtrsim 2 \cdot 10^9 \text{ GeV}$$

Gravitino Production

- Thermal production at high temperature

$$\Omega_{3/2}^{\text{tp}} h^2 \simeq 0.11 \left(\frac{T_R}{2 \cdot 10^9 \text{ GeV}} \right) \left(\frac{M_{\tilde{g}}}{10^3 \text{ GeV}} \right)^2 \left(\frac{67 \text{ GeV}}{m_{3/2}} \right)$$

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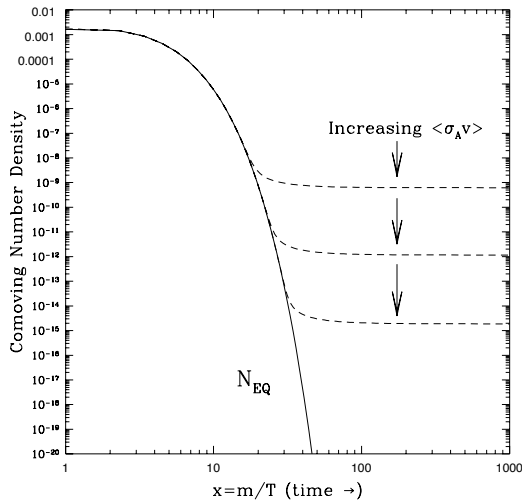
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- Observed dark matter abundance: $\Omega_{\text{DM}} h^2 \simeq 0.11$

→ **Compatible** with thermal leptogenesis:

- Gravitino **LSP** with mass $\gtrsim 60 \text{ GeV}$
- Heavier non-LSP gravitino

WIMP Freeze-Out



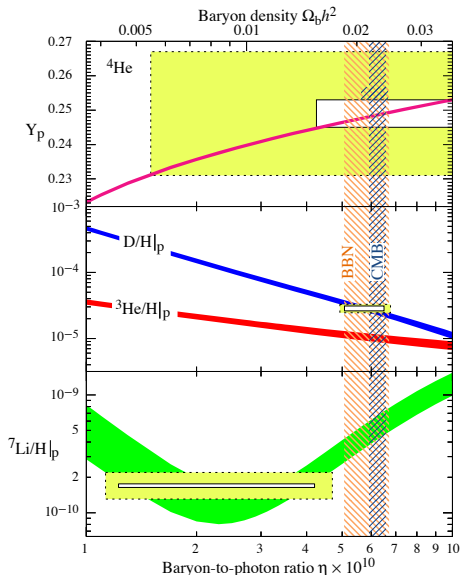
- Weakly interacting stable particle χ
- Thermal equilibrium:
 $N_\chi \propto e^{-T/m_\chi}$
- Annihilation rate
 $\Gamma(\chi\bar{\chi} \rightarrow xy) < H$
 \leadsto freeze-out:
 $N_\chi = \text{const.}$
 \leadsto relic density Ω_χ determined
- $T_{\text{fo}} \sim \frac{m_\chi}{25}$

Big Bang Nucleosynthesis

- $T \sim 1 \text{ MeV}$ or $t \sim 1 \text{ s}$:
freeze-out of $n \leftrightarrow p$
 $\leadsto n/p$ ratio fixed
- $T \sim 0.1 \text{ MeV}$: $p + n \rightarrow \text{D}$
- Afterwards formation of
 ${}^3\text{He}$, ${}^4\text{He}$, ${}^7\text{Li}$
- Abundances depend on
baryon density (Ω_B or η_B)
- Agree with observations for
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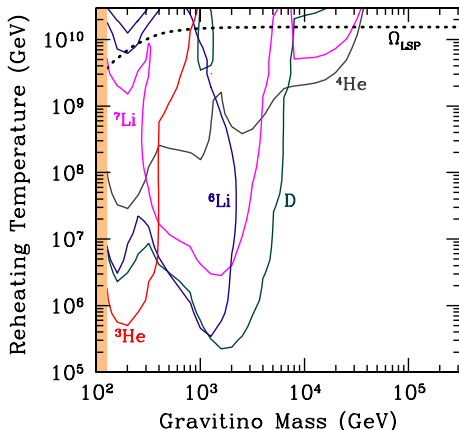


Gravitino Problem

- Gravitino interacts via **gravity**
 - ↪ extremely weakly
 - ↪ **lifetime** $\sim 10^{-2}$ s ... years
- Energetic decay products destroy nuclei produced in **Big Bang Nucleosynthesis**
- Distortions of the **Cosmic Microwave Background**
(less constraining)

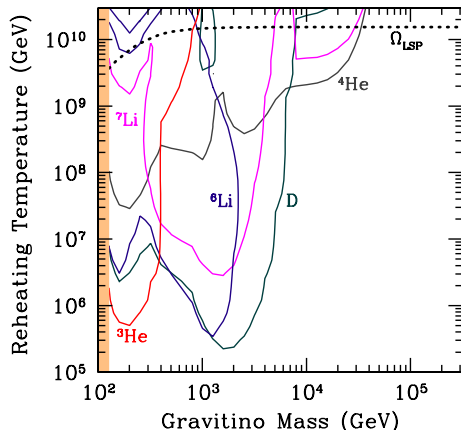
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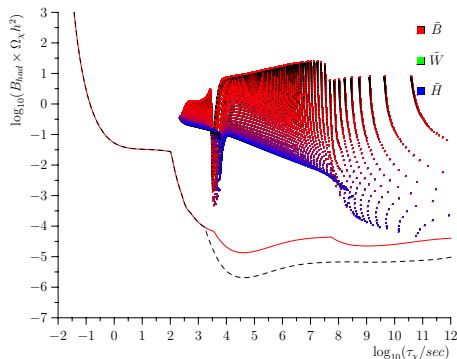
$\leadsto T_R \lesssim 10^7$ GeV or $m_{3/2} \gg 1$ TeV

\leadsto **Conflict** with thermal **leptogenesis**, or **unnatural spectrum**

Big Bang Nucleosynthesis with Gravitino LSP

- Gravitino LSP: Next-to-LSP (**NLSP**) long-lived
- BBN bounds depend on kind of NLSP
- Assume Ω_{NLSP} to be given by thermal relic density
- **Neutralino** ruled out unless very heavy

$m_{3/2} = 100 \text{ GeV}$ with: $M_2 = 2200$, $M_3 = 2200$, $\tan \beta = 10$, $\text{sign}(\mu) = 1$.



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- **Neutralino** ruled out unless very heavy
- **Stau** decays can be ok, but **bound states** with nuclei change BBN reaction rates \leadsto overproduction of ${}^6\text{Li}$
- **Sneutrino** mostly harmless

\leadsto Gravitino problem remains as **NLSP decay problem**

Solutions

- Abandon SUSY
- Abandon thermal leptogenesis
- Fine-tune to exploit loopholes
- Very heavy gravitino
- Gravitino LSP + harmless NLSP
 - New interactions \leadsto faster decay
 - Very light gravitino \leadsto faster decay, $\Omega_{3/2} \not\propto T_R$
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 - Abundance smaller than thermal relic abundance
- Arbitrary combinations

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NLSP Dilution by Entropy Production

- BBN bounds depend on $\Omega_{\text{NLSP}} \propto \frac{N}{S}$
 S = comoving **entropy** density
- Increase of entropy after freeze-out: $S \rightarrow S \Delta$
 \leadsto **dilution** of NLSP density: $\Omega_{\text{NLSP}} \rightarrow \frac{\Omega_{\text{NLSP}}}{\Delta}$
 \leadsto reduction of impact on BBN

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- Entropy from decay of non-relativistic particle ϕ

$$\frac{\rho_{\phi}}{\rho_{\text{rad}}} \propto \frac{R^{-3}}{R^{-4}} = R$$

- $\leadsto \phi$ **dominates** energy density at some time $t_{=}$, temperature $T_{=}$
- Candidates: later

Constraints

- Radiation domination at NLSP freeze-out:

$$T_{\text{=}} < T_{\text{fo}} \sim \frac{m_{\text{NLSP}}}{25}$$

\leadsto standard calculation of Ω_{NLSP} applies (ϕ can be ignored)

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$$T_{\text{dec}} = T(\tau_{\phi}) \gtrsim (0.7 \dots 4) \text{ MeV}$$

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→ Maximal dilution factor:

$$\Delta \simeq 0.75 \frac{T_{\text{=}}}{T_{\text{dec}}} \lesssim 750 \left(\frac{m_{\text{NLSP}}}{100 \text{ GeV}} \right) \sim 10^3$$

Other Effects of Entropy

😊 $\Omega_{\text{NLSP}} \rightarrow \frac{\Omega_{\text{NLSP}}}{\Delta}$

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Remember $\eta_{\text{B}} \propto M_{\text{R}}$ and $T_{\text{R}} \gtrsim M_{\text{R}}$

↪ To keep observed η_{B} :

$$M_{\text{R}} \rightarrow M_{\text{R}}\Delta \text{ and } T_{\text{R}} \rightarrow T_{\text{R}}\Delta$$

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Without Δ	With Δ
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T_R	$T_R \Delta$
$\Omega_{3/2}$	$\Omega_{3/2}$
Ω_{NLSP}	$\frac{\Omega_{\text{NLSP}}}{\Delta}$

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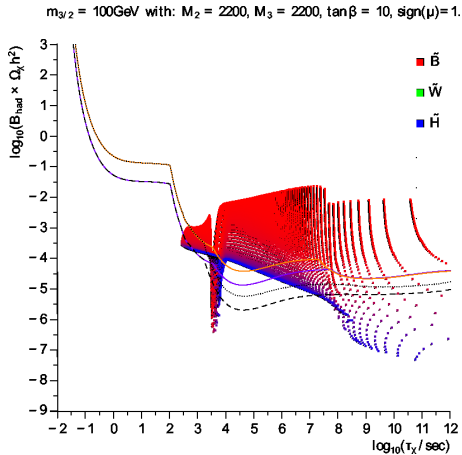
Strong **washout** of η_B for $M_R \gtrsim 10^{13}$ GeV → slower increase

→ observed η_B can only be reached for $\Delta \lesssim 10^3 \dots 10^4$

Without Δ	With Δ
η_B	η_B
T_R	$T_R \Delta$
$\Omega_{3/2}$	$\Omega_{3/2}$
Ω_{NLSP}	$\frac{\Omega_{\text{NLSP}}}{\Delta}$

Neutralino NLSP with Entropy Production

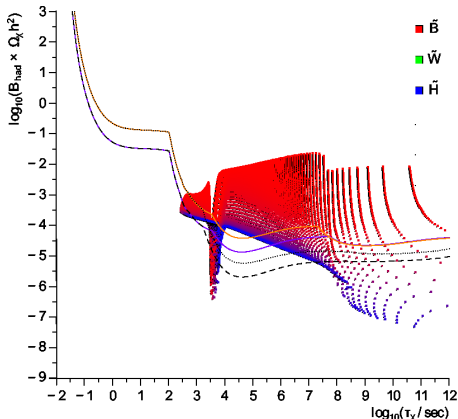
- Gravitino LSP,
 $m_{3/2} = 100 \text{ GeV}$
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- Light neutralinos allowed for significant higgsino or wino content
- Pure binos remain excluded

→ Thermal leptogenesis possible

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Generic or necessary for long-lived particles even **without** demanding entropy production

Entropy from Saxion Decays

- Strong CP problem \leadsto Peccei-Quinn mechanism \leadsto axion
- SUSY: axion **supermultiplet** (axion, **saxion** ϕ , axino)
- Interactions suppressed by characteristic scale $f_a \gtrsim 10^9$ GeV

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- Dominant decay $\phi \rightarrow gg \leadsto$ dilution factor:

$$\Delta \gtrsim 55 \left(\frac{f_a}{10^{12} \text{ GeV}} \right)^{\frac{2}{3}} \ll 10^3 \text{ ☹️}$$

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- Further problem with **axino**

Non-Thermally Produced Saxion

- Saxion field displaced from potential minimum during inflation
- Oscillations around minimum \leadsto non-relativistic particles
- Production and decay decoupled \leadsto consistent scenario

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- Example with maximal dilution factor:

$$\Delta \sim 10^3$$

$$\text{Saxion mass} \sim 10 \text{ GeV}$$

$$\text{Axino mass} \sim 1 \text{ TeV}$$

$$f_a \sim 10^{10} \text{ GeV}$$

$$\text{Initial amplitude} \sim 10^4 f_a$$

$$m_{\text{NLSP}} \simeq 200 \text{ GeV}$$

$$m_{3/2} \simeq 100 \text{ GeV}$$

Conclusions

- Gravitino problem in SUSY scenarios with thermal leptogenesis
- Solution: gravitino LSP, dilution of NLSP by entropy
- Neutralino NLSP with large higgsino or wino component ok
- Constraints on entropy-producing particle
- Thermally produced particles fail
- Saxion produced in oscillations works