

# D-brane Inflation in String Theory and its Cosmological Footprints

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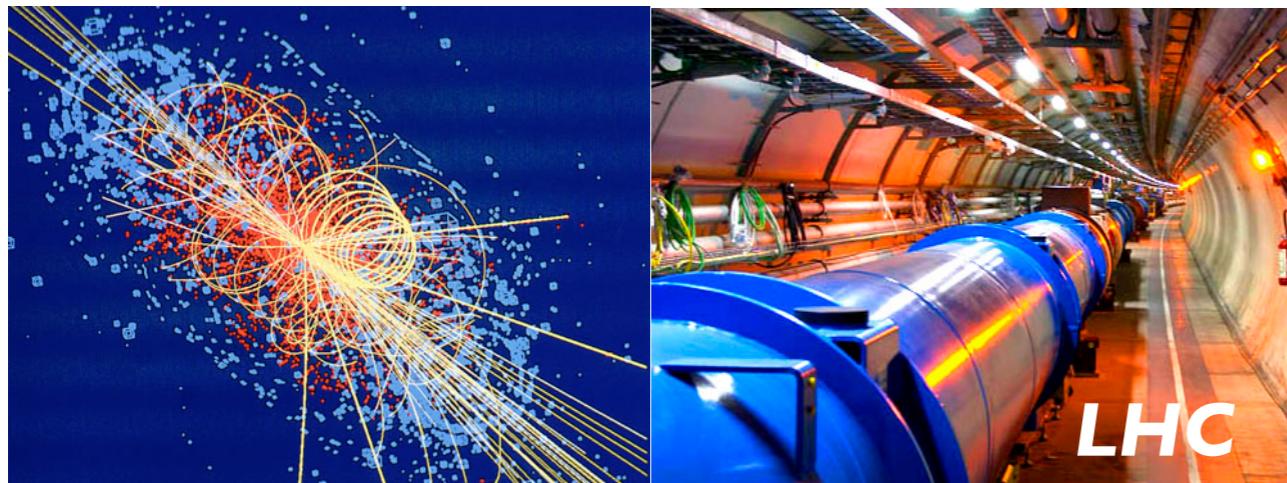
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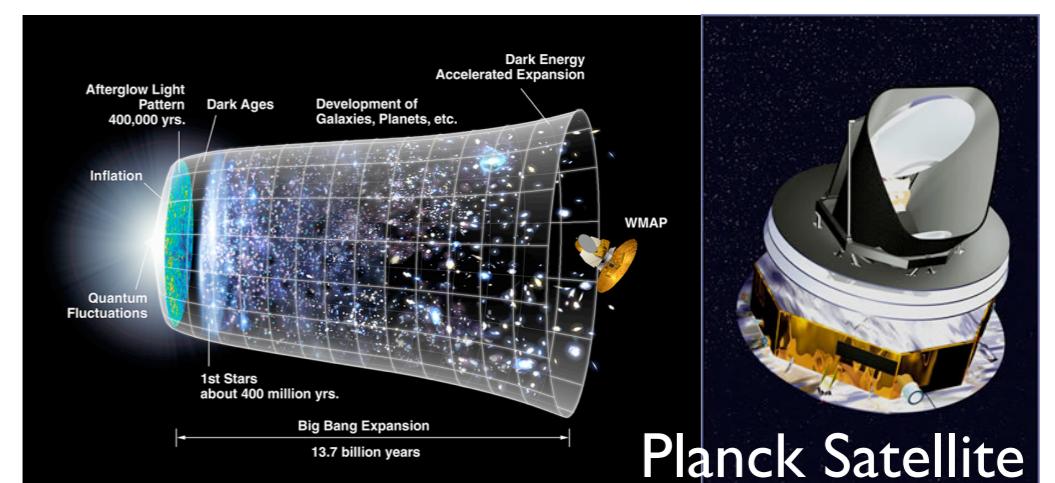
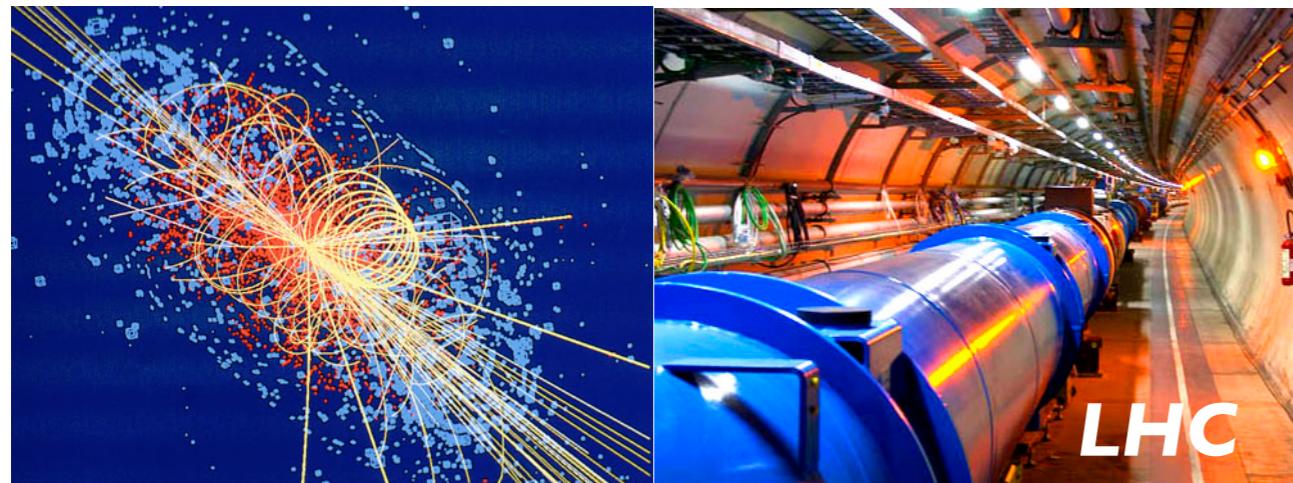
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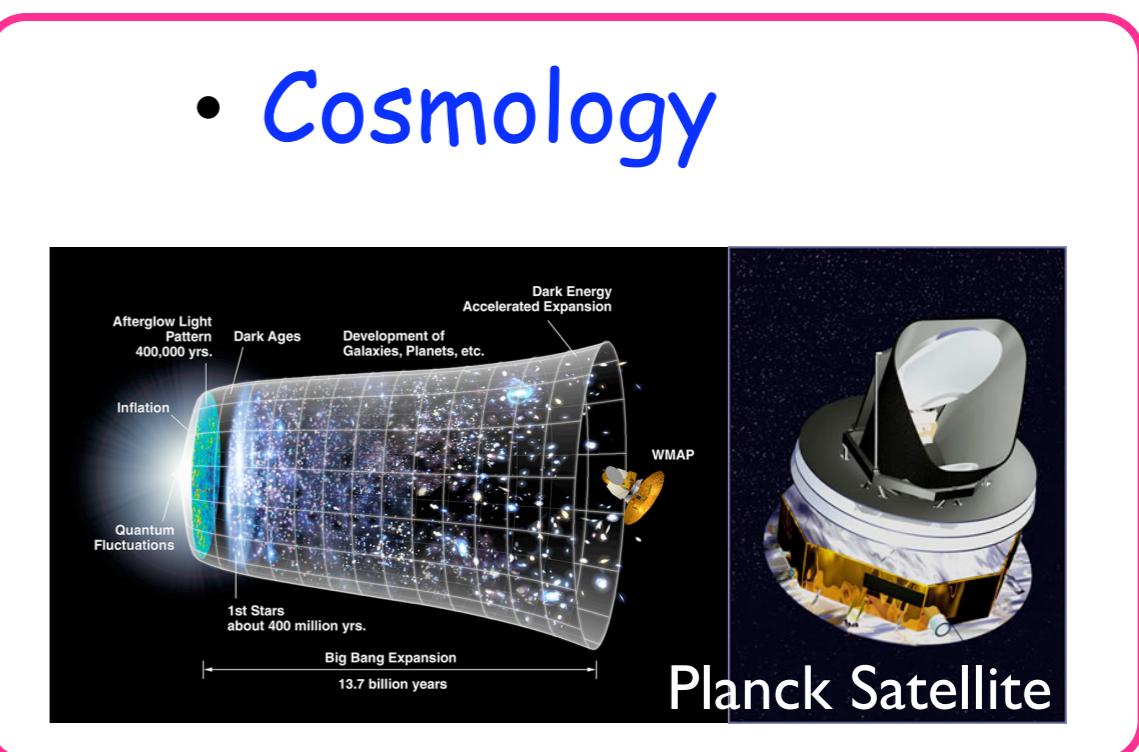
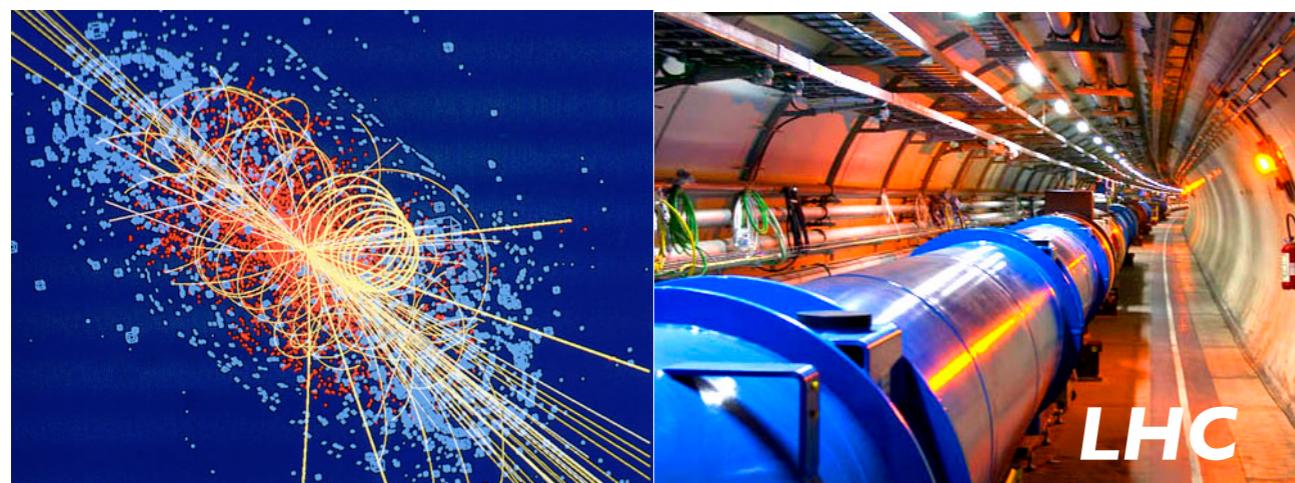
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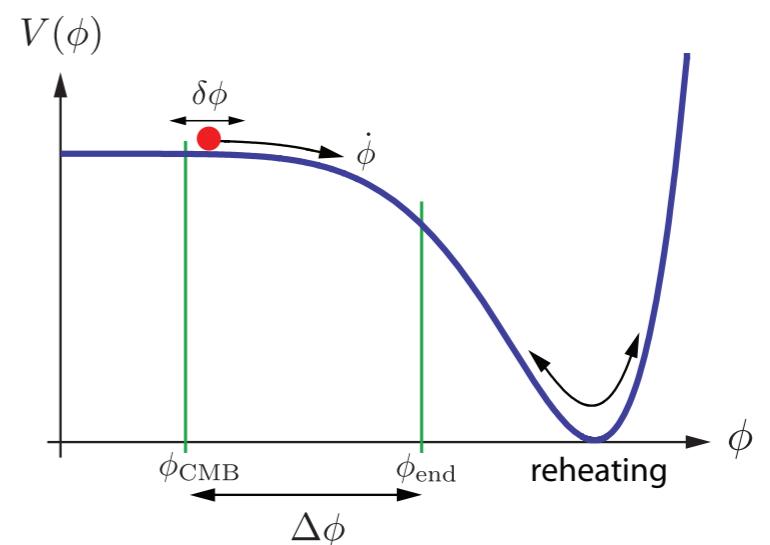
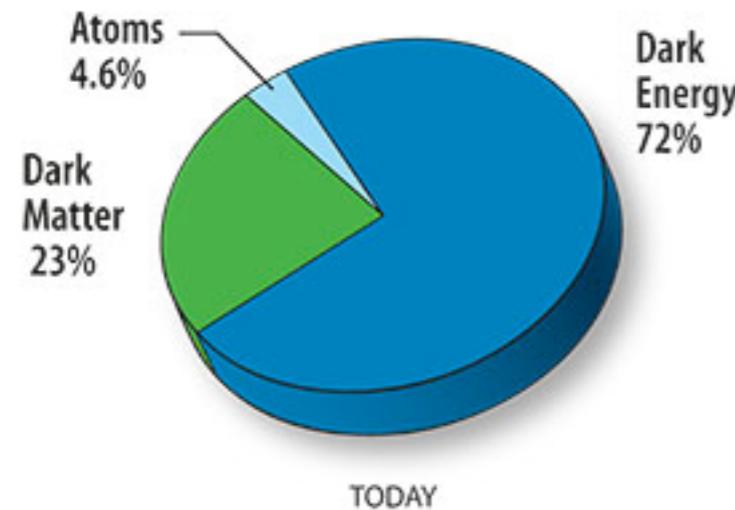


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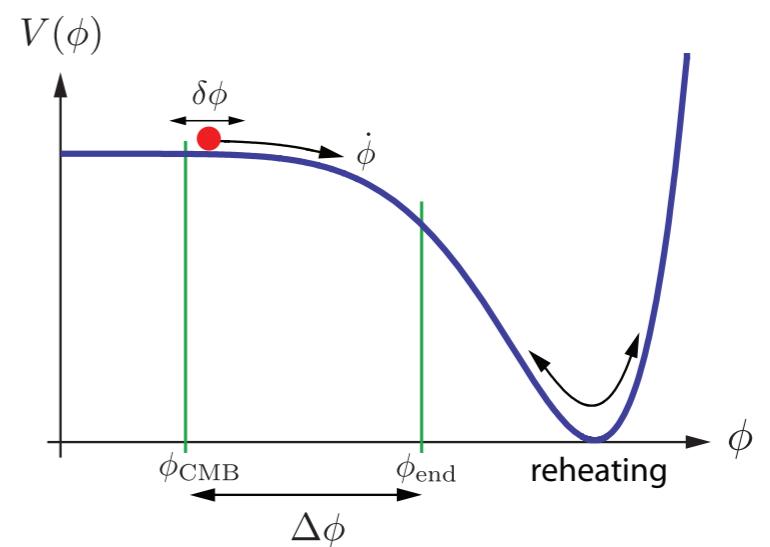
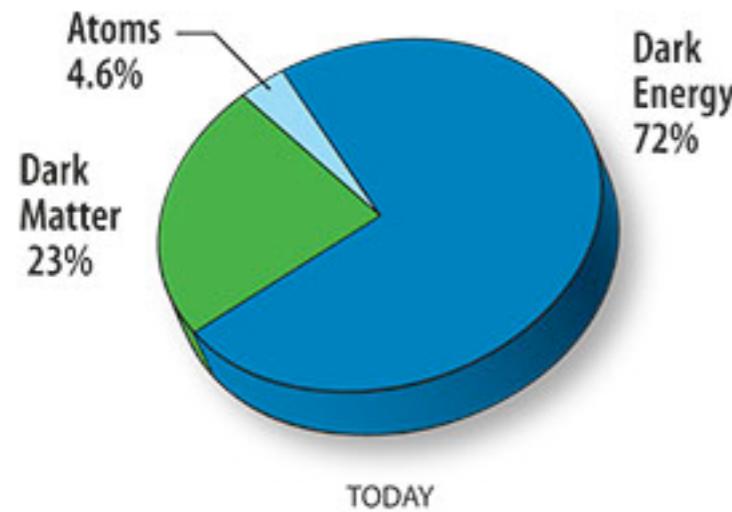


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- Current observations support  $\Lambda$ CDM\* model, complemented with inflationary paradigm, to generate the primordial perturbations that seed the large structures we observe



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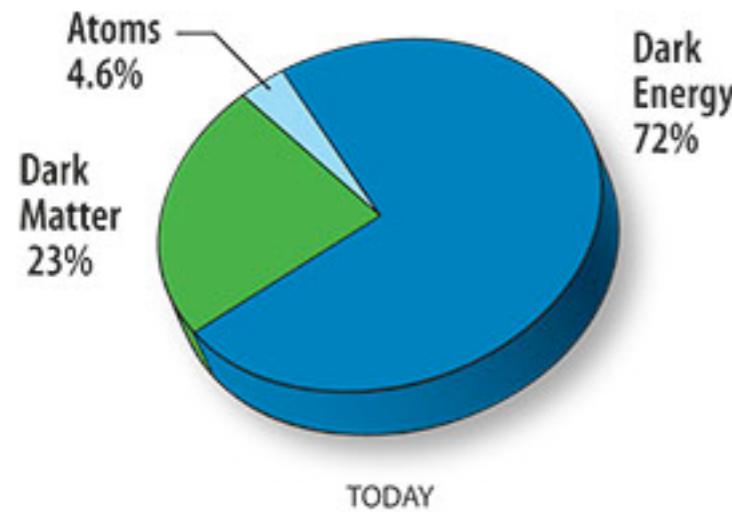
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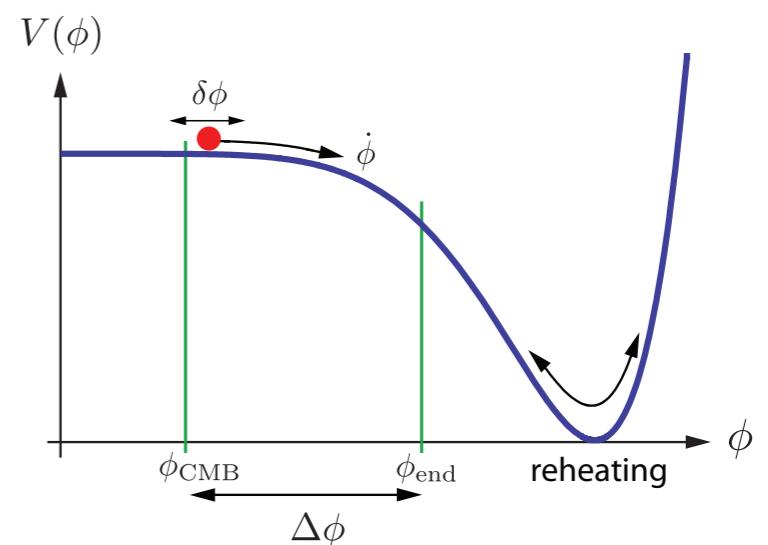
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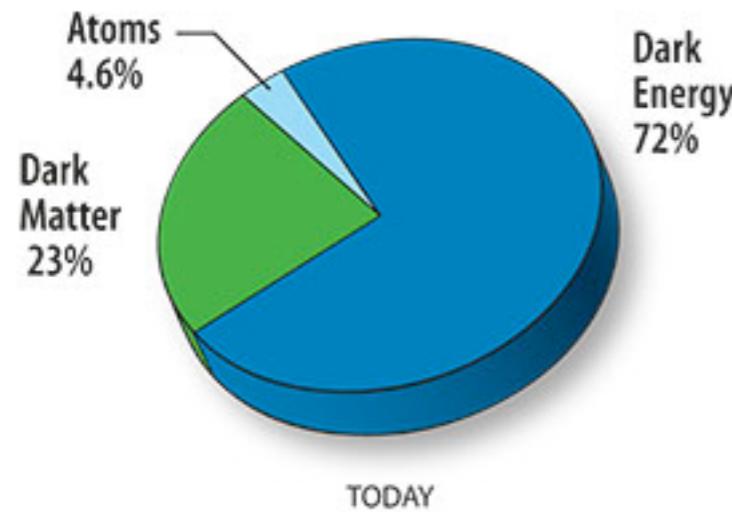
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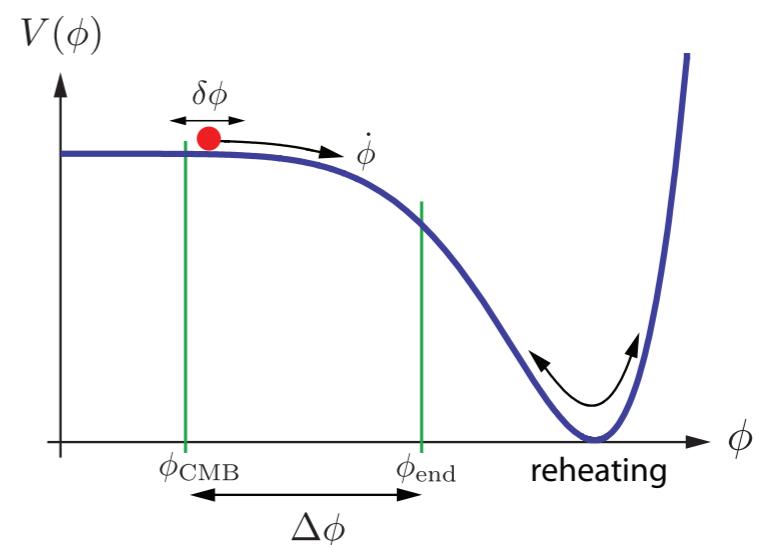
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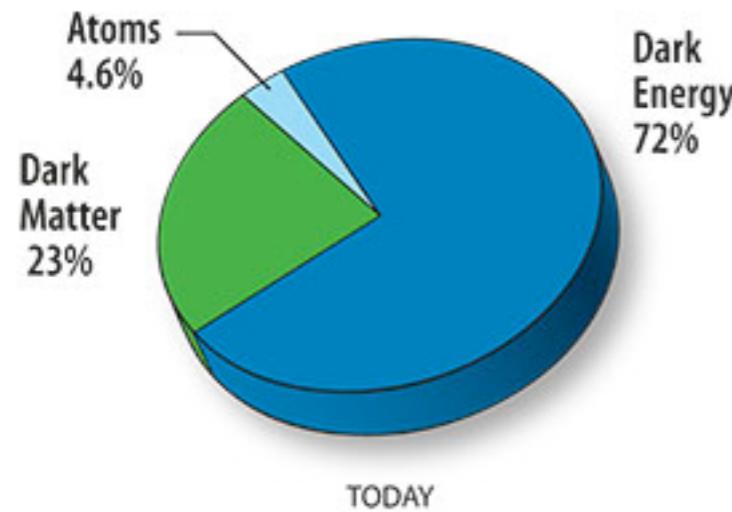
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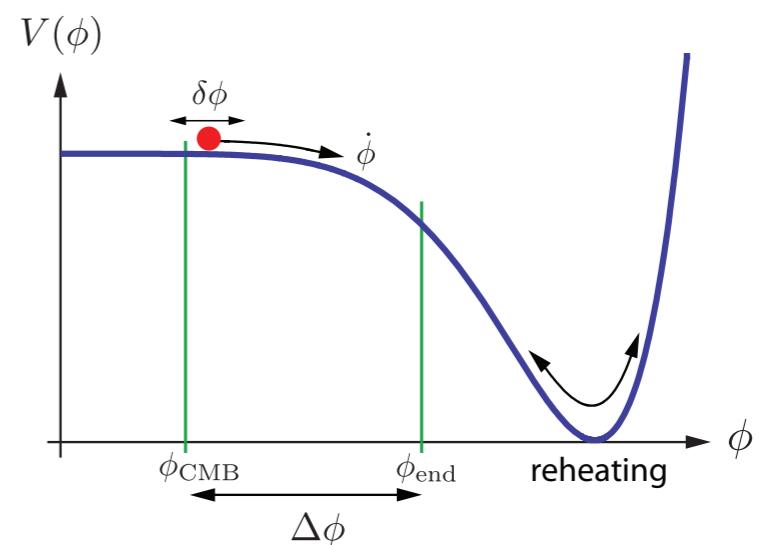
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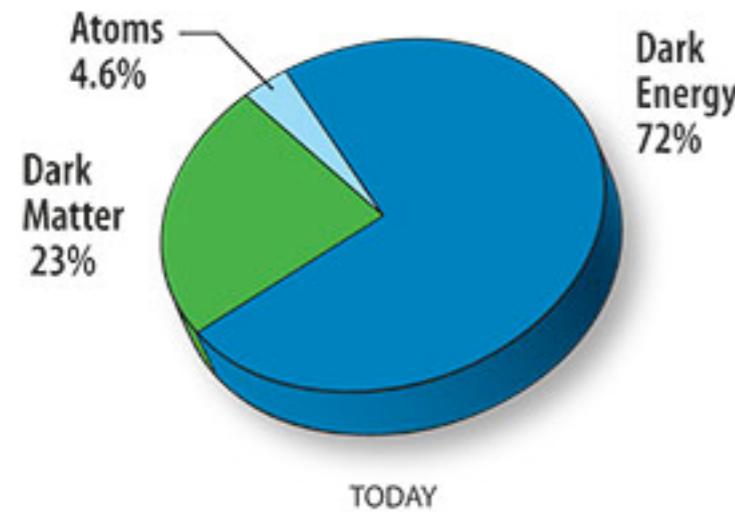
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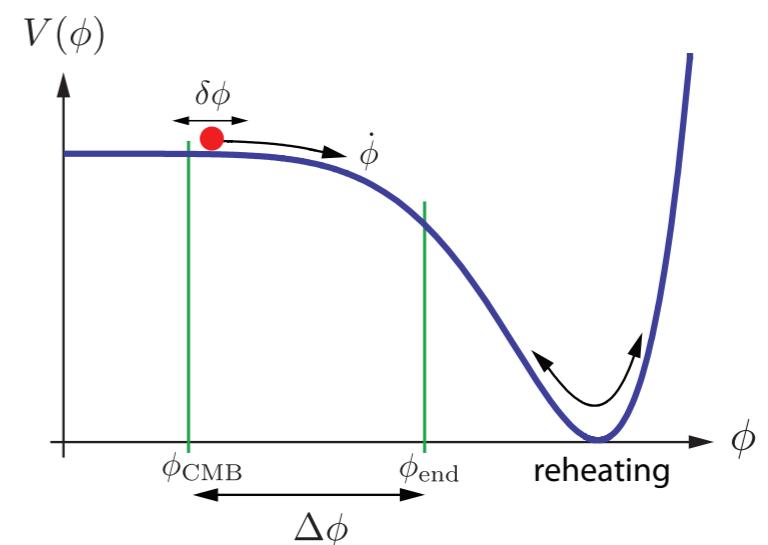
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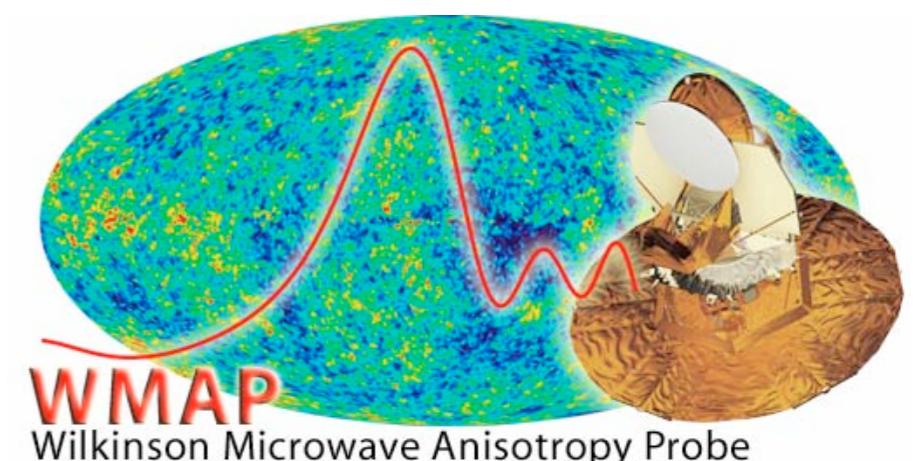
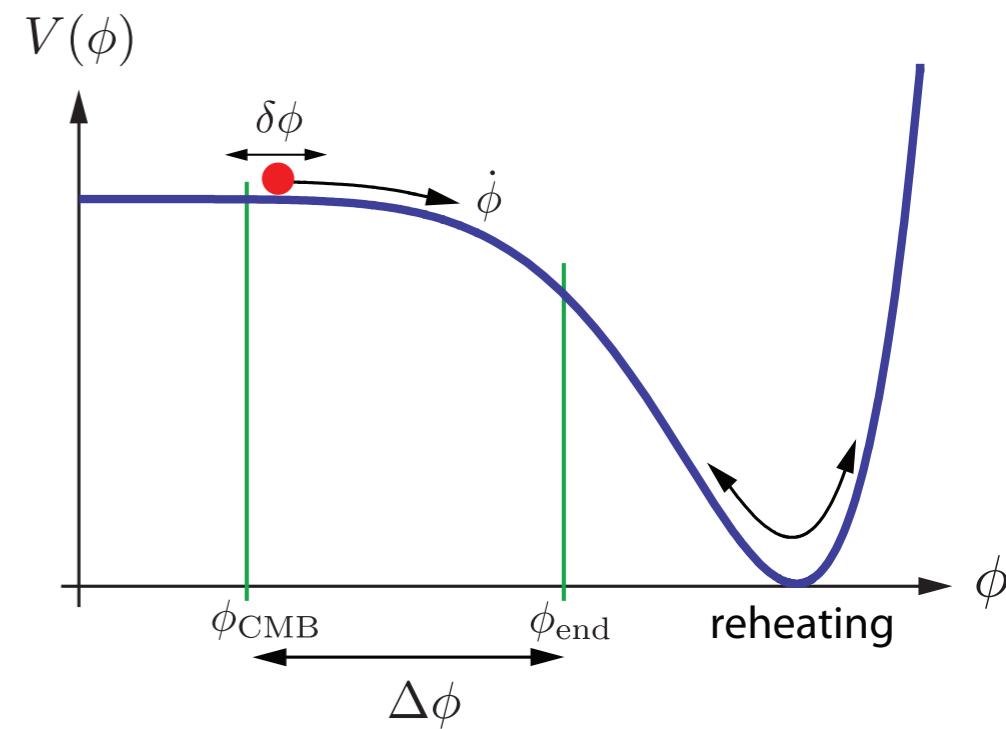
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## D-brane Inflation

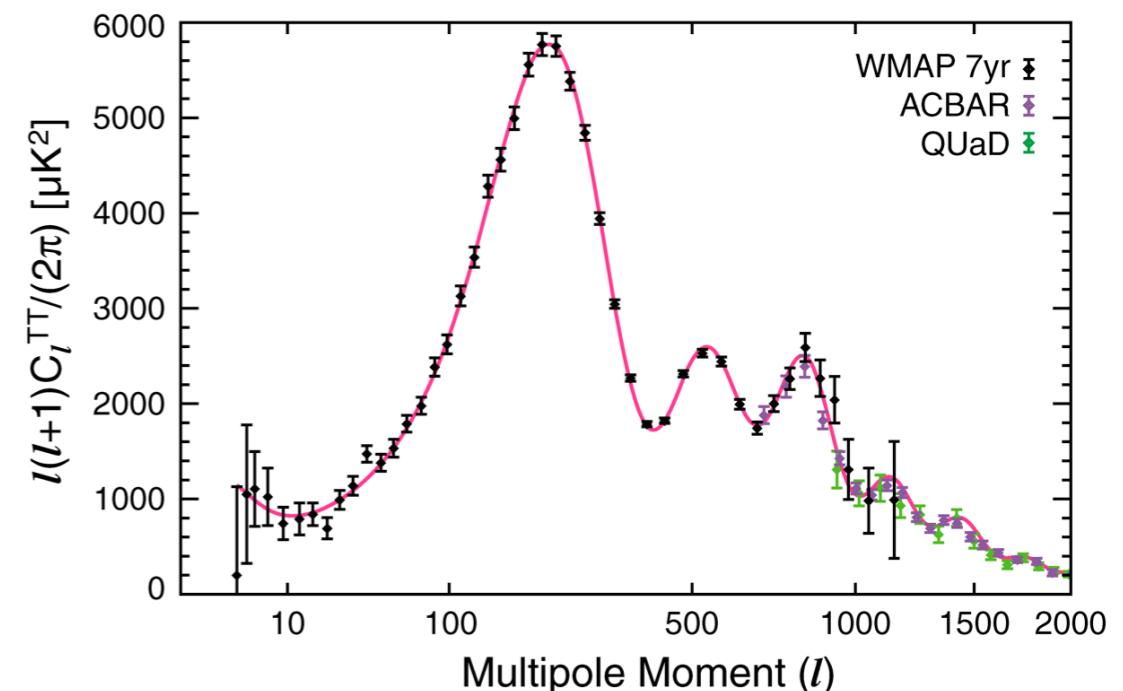
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# The Inflationary Paradigm

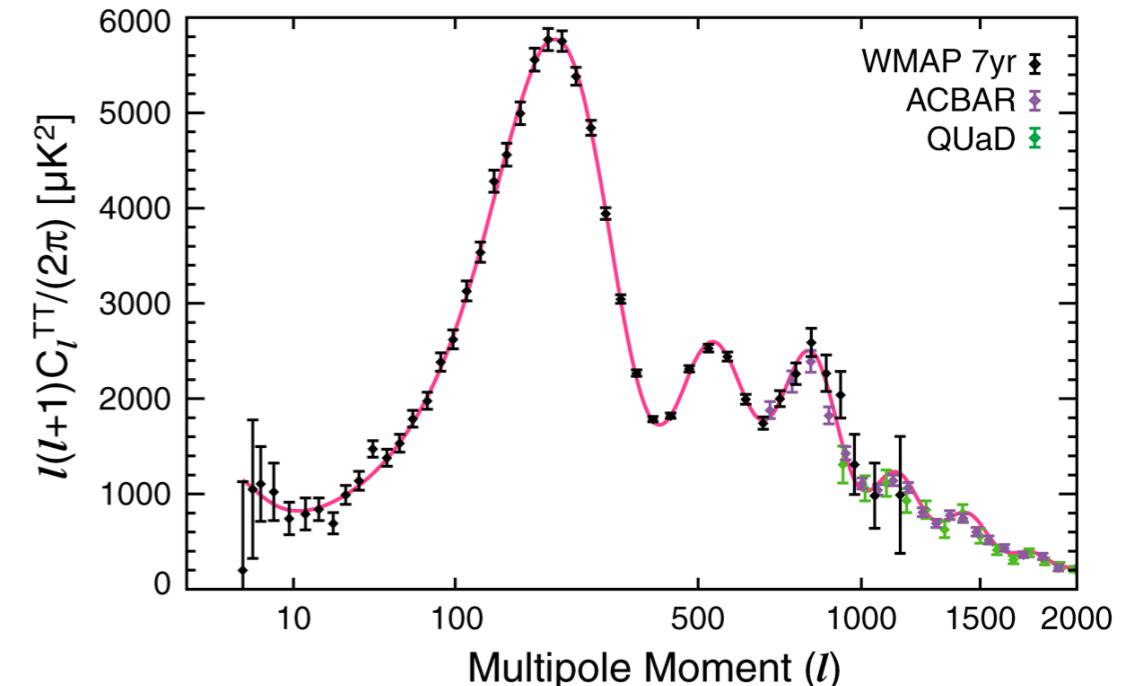
- A period of quasi exponential **accelerated expansion** in the very early Universe due to vacuum energy domination.
- Governed by the dynamics of a **scalar field with very flat potential**.
- Explains why the Universe is approximately **homogeneous and spatially flat**.
- Can explain the absence of **relic particles** and defects predicted by fundamental physics.
- Leading mechanism to explain observed inhomogeneities in the Universe: provides the **seeds for large scale structure formation**.  
Fluctuations of  $\delta\phi \rightarrow$  fluctuations of energy density  $\delta\rho \rightarrow$  temperature fluctuations  $\delta T$  in the CMB (Cosmic Microwave Background).



- Simplest inflationary model
  - single scalar field with very flat potential in a flat Universe
  - nearly scale invariant adiabatic Gaussian spectrum of scalar perturbations. In good agreement with observations of the CMB (WMAP)



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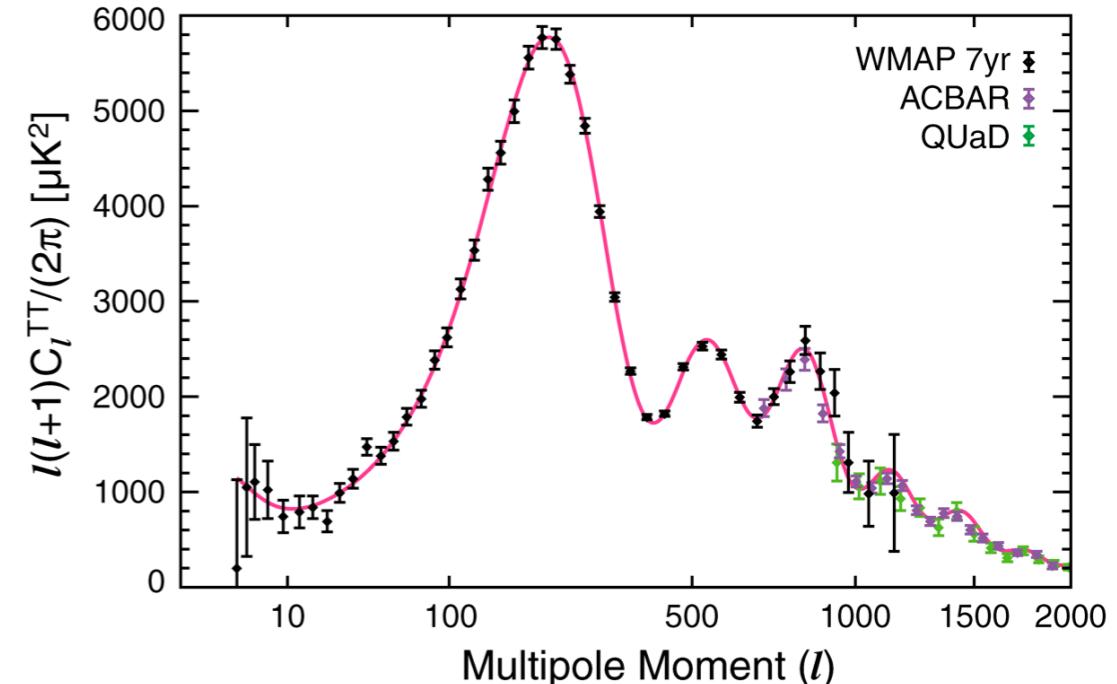


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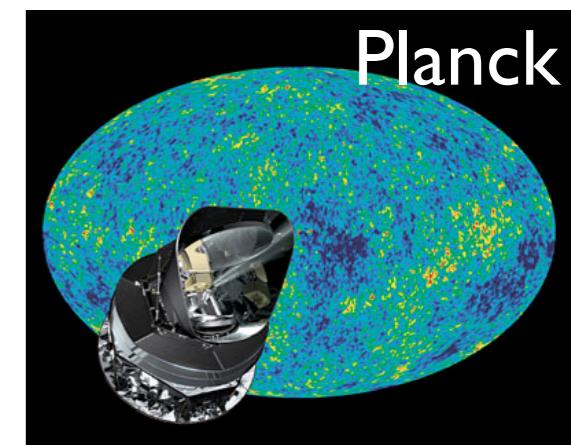
- \* Observable spectrum of (primordial) gravitational waves  $r \gtrsim 10^{-2}$

- \* Significant non-Gaussianity  $-10 < f_{NL} < 50$

- \* Scale dependence of the spectral index  $-0.084 < \frac{dn_s}{d \ln k} < 0.02$

- \* Isocurvature density perturbations

- \* Anisotropies in the spectrum



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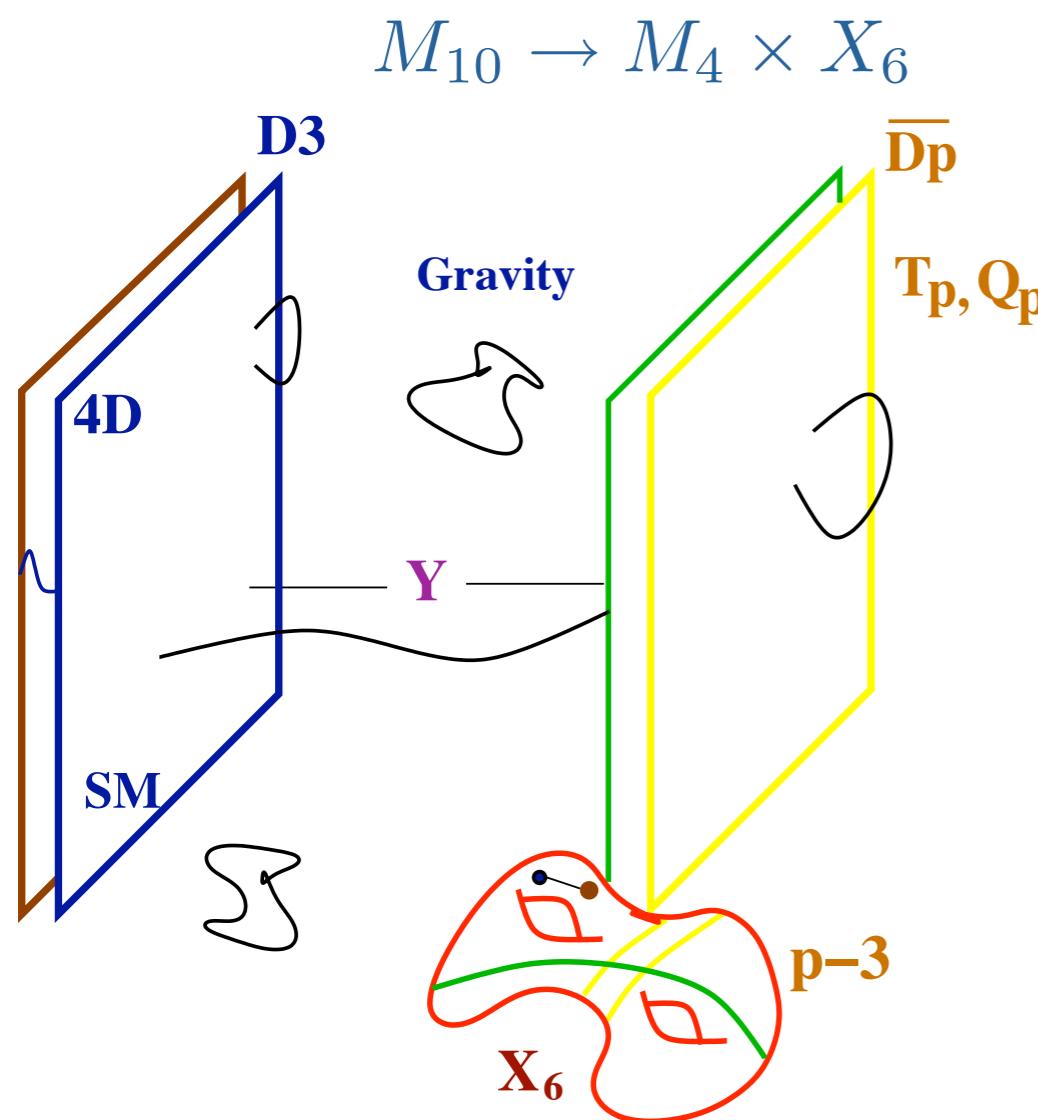
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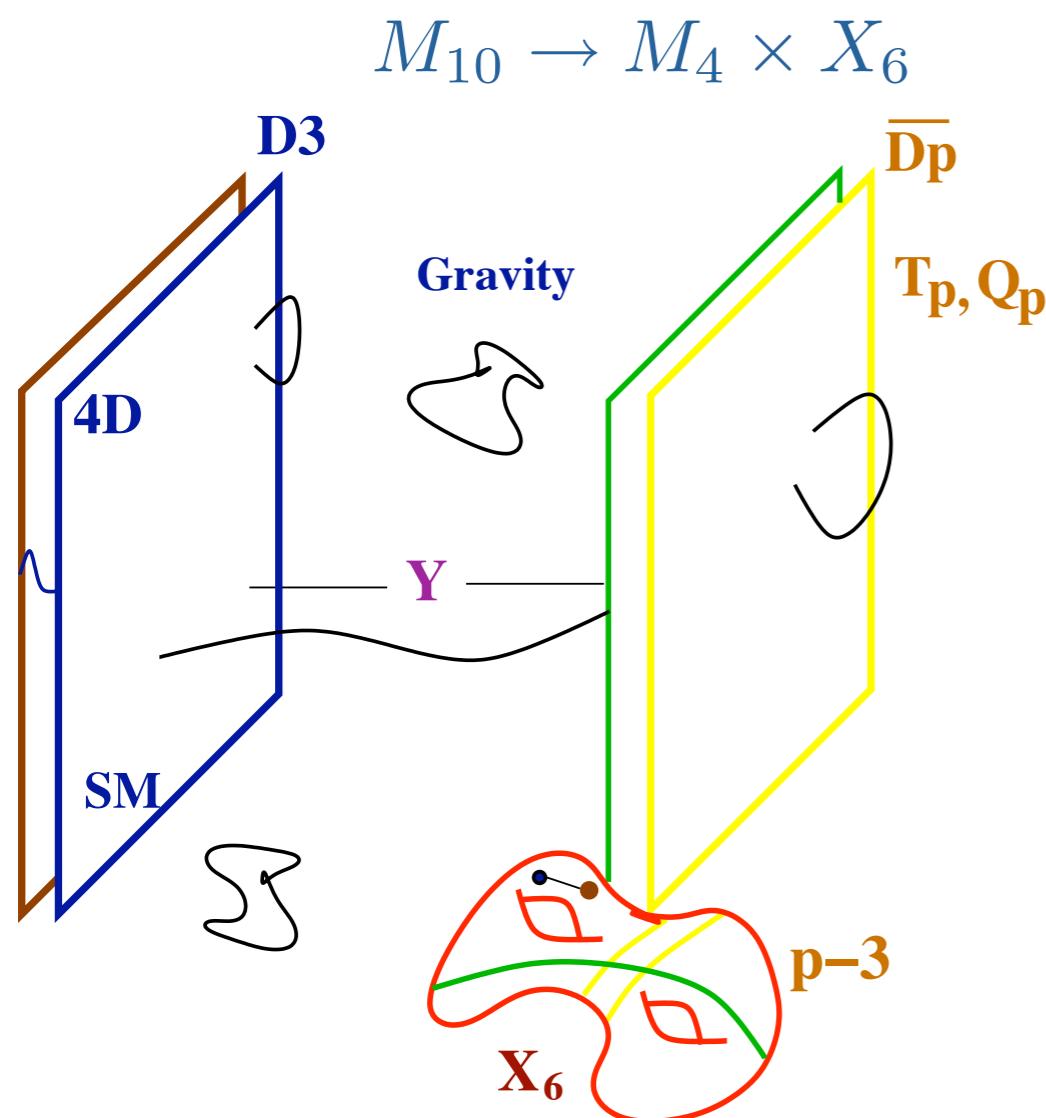


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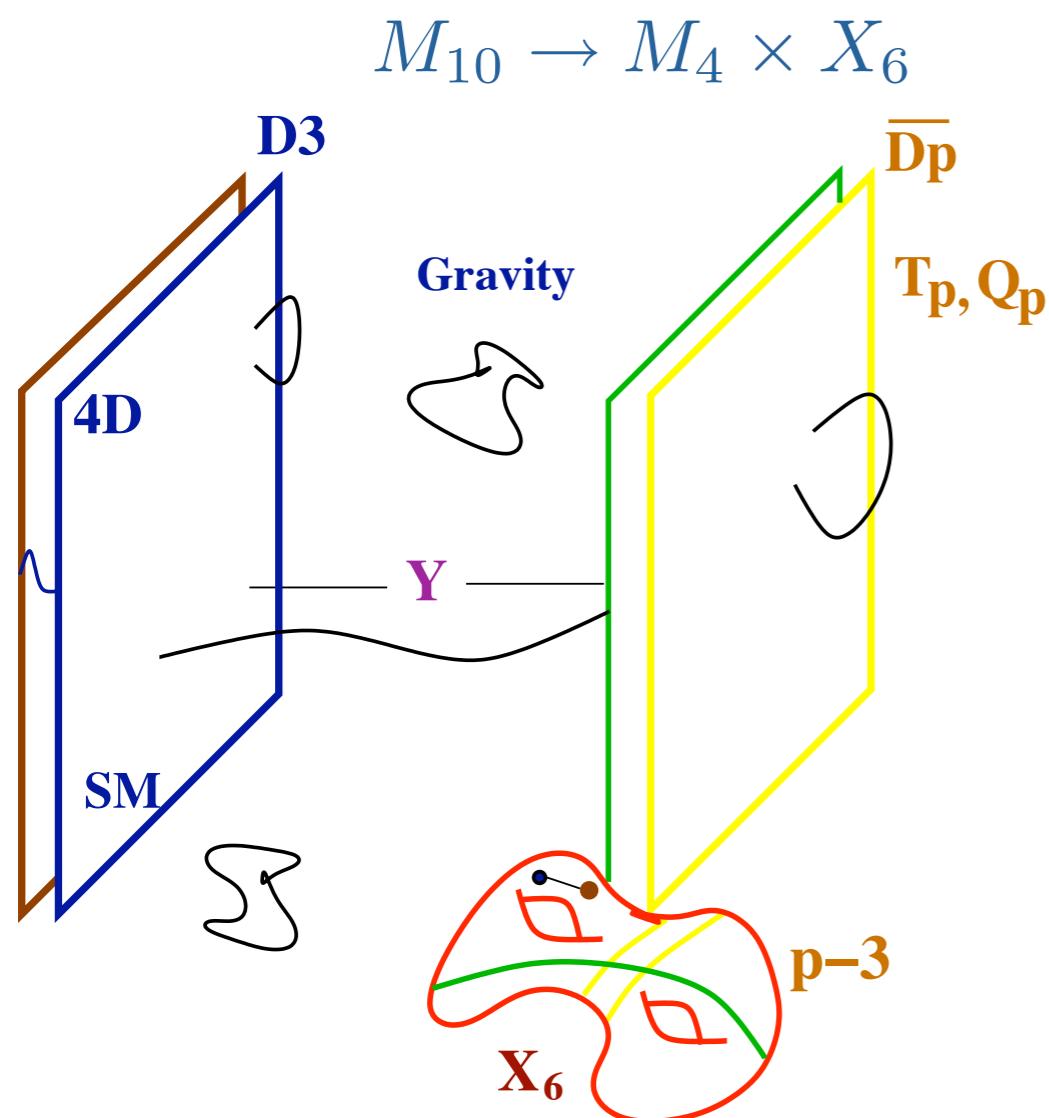
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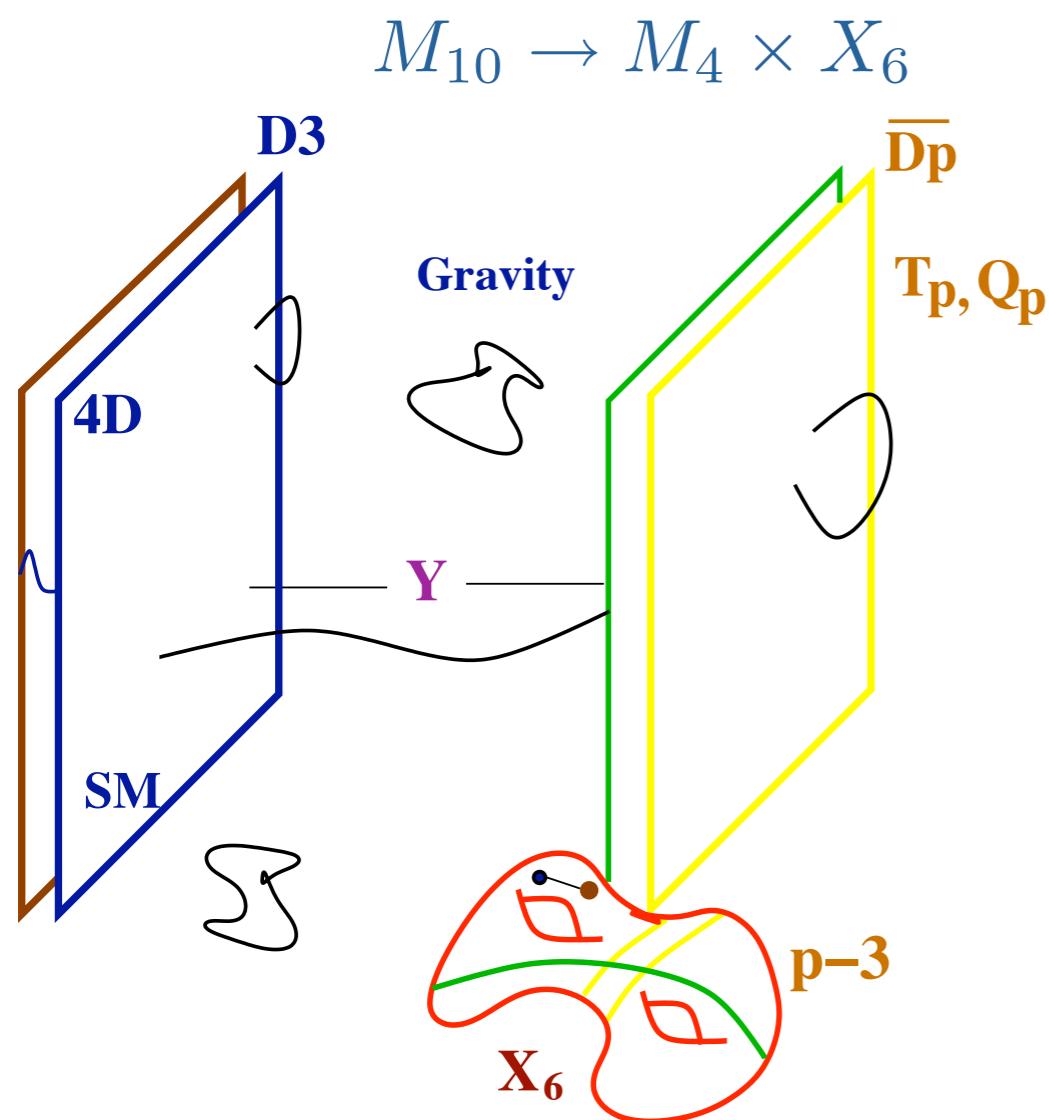
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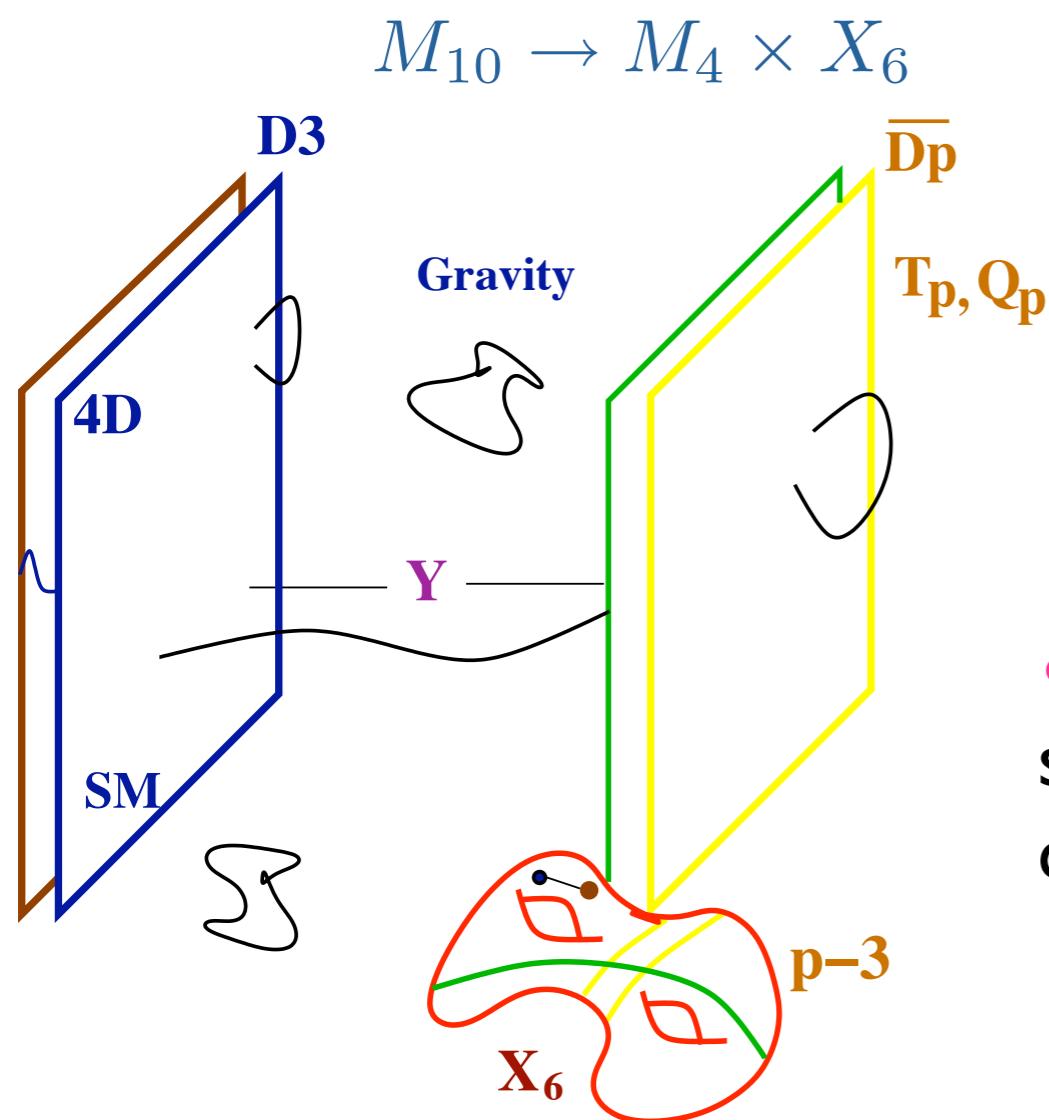
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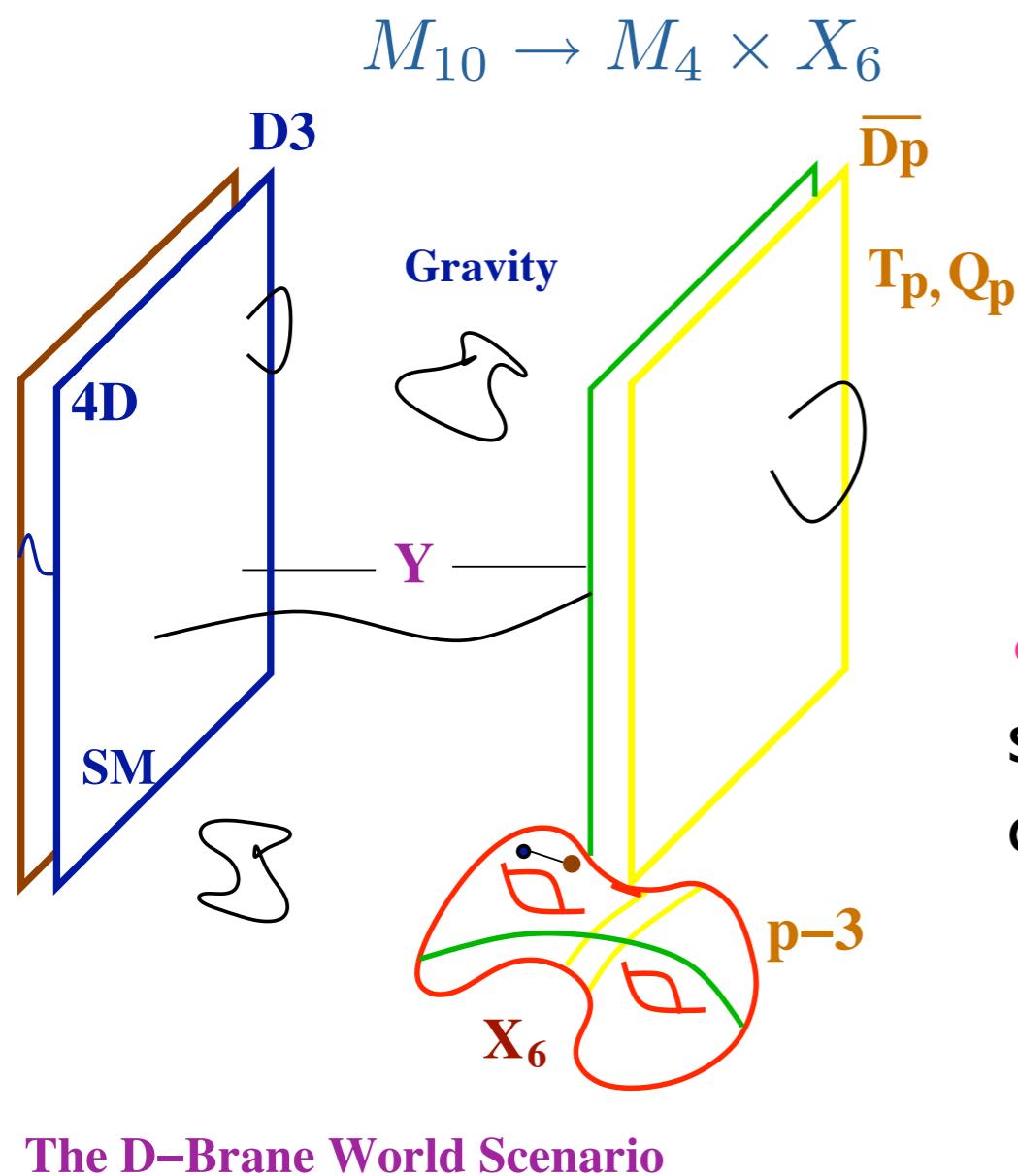


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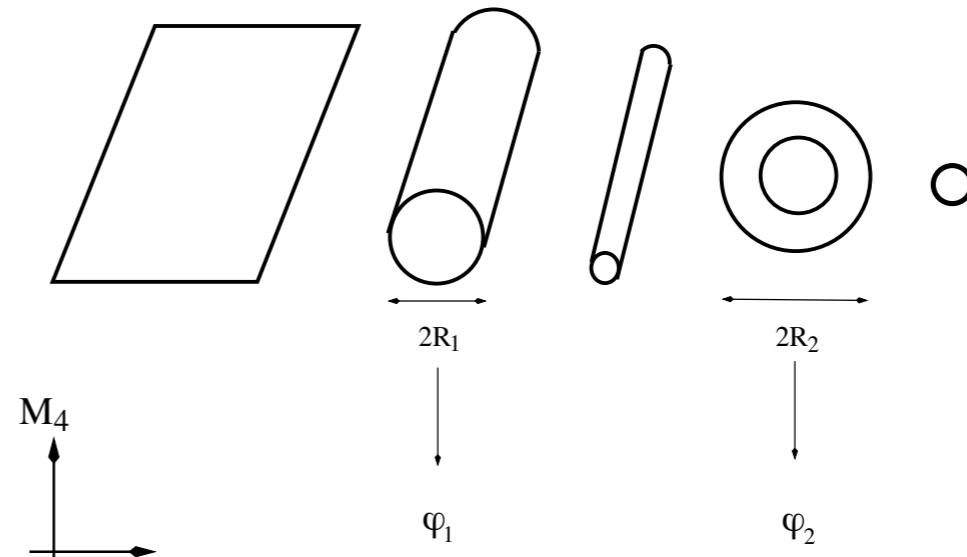
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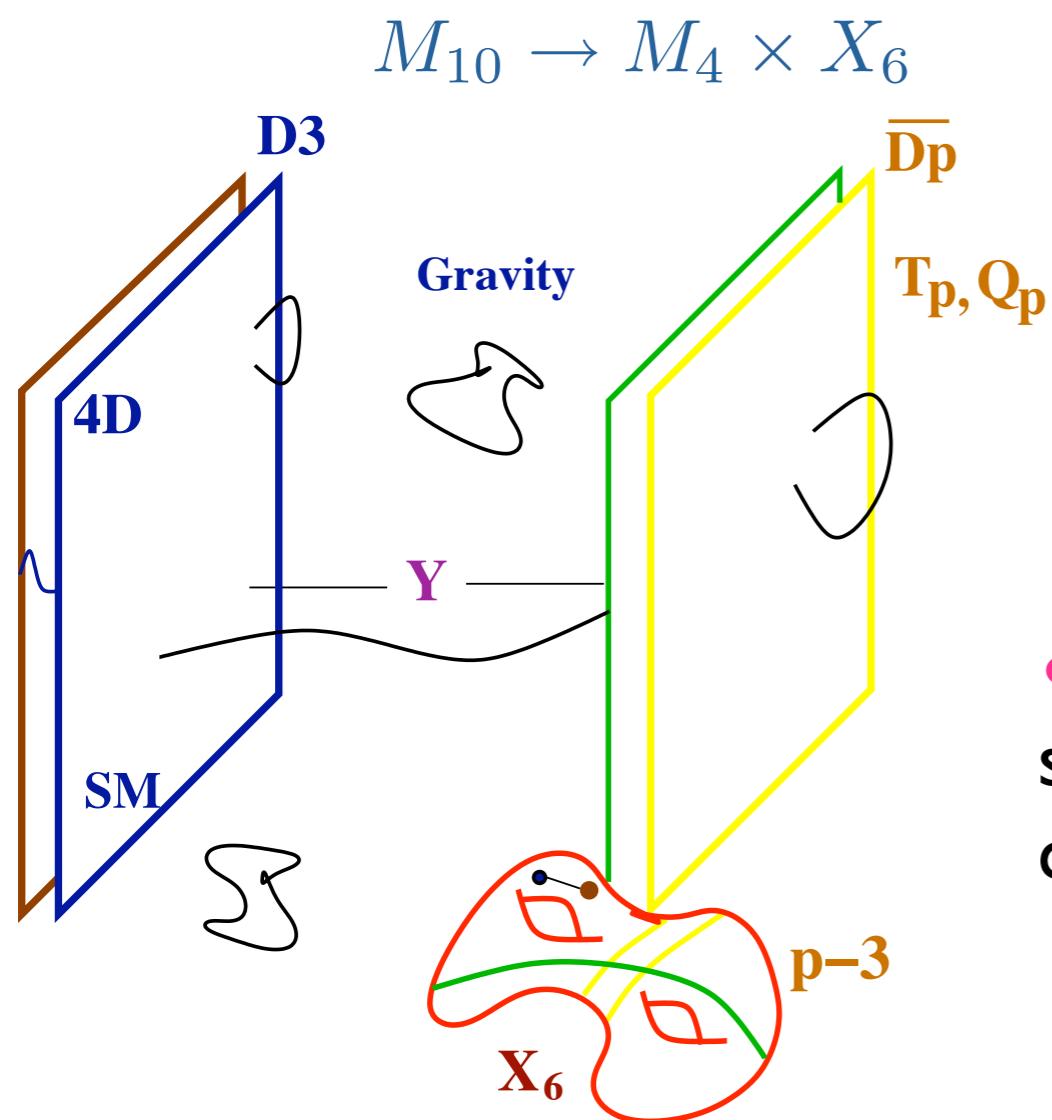


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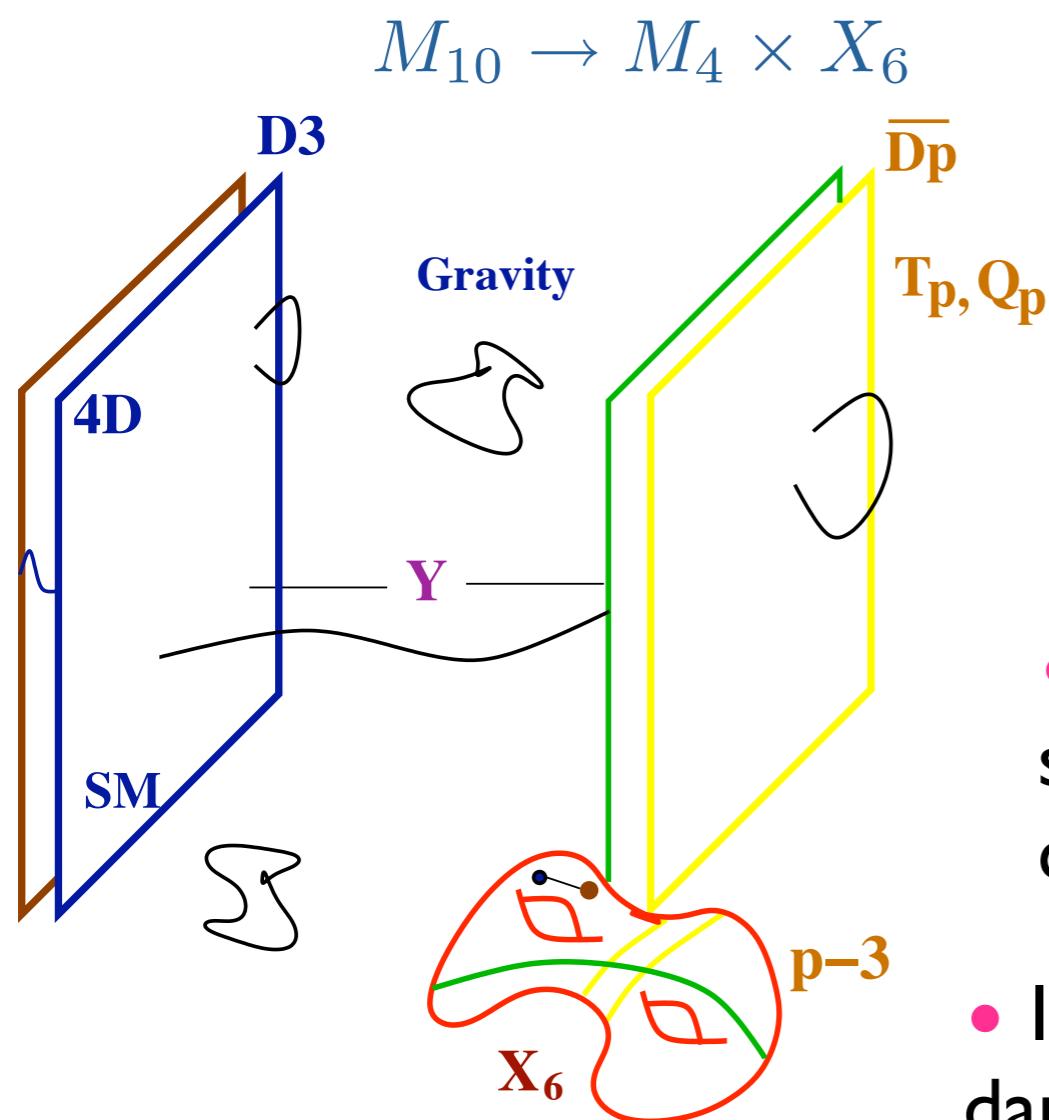


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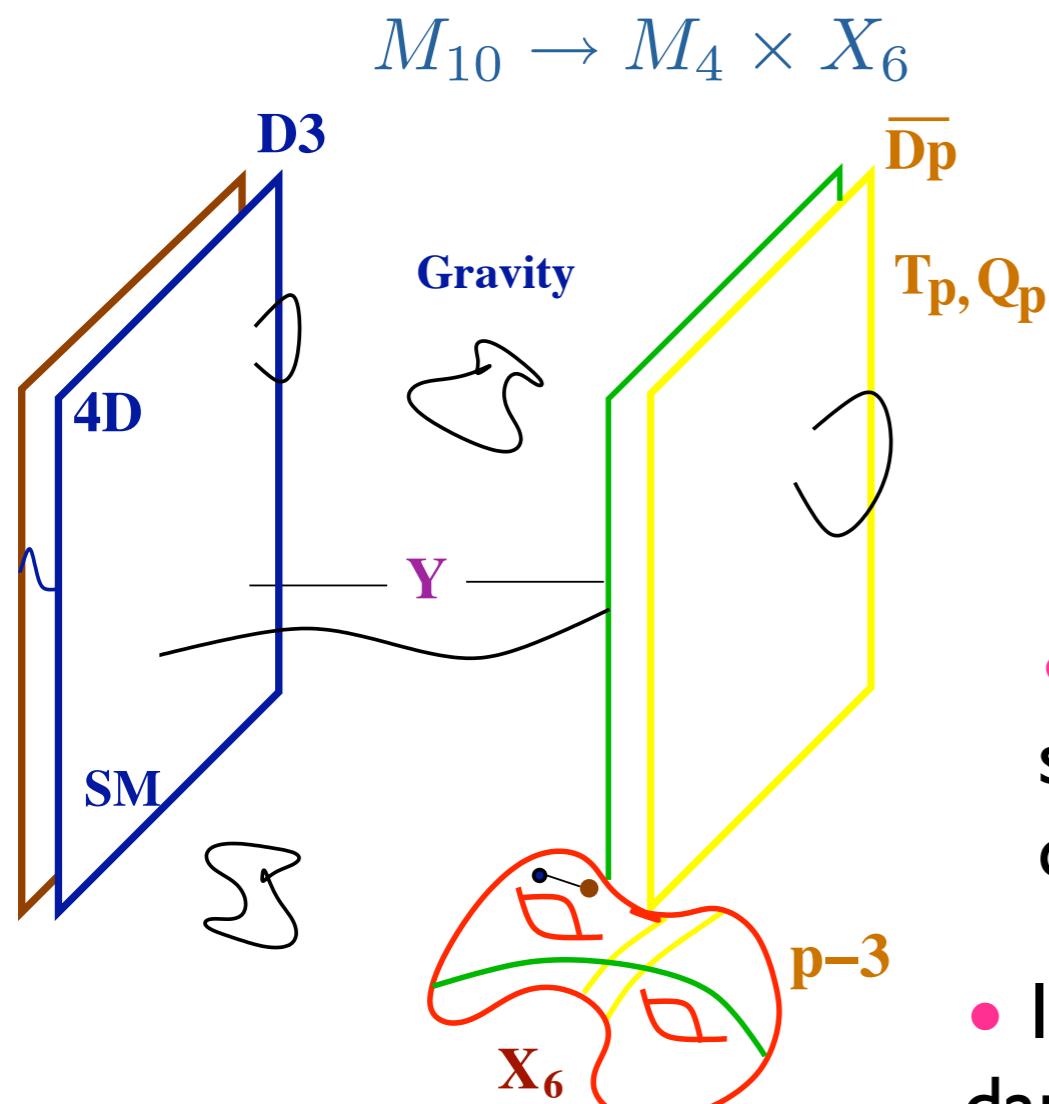


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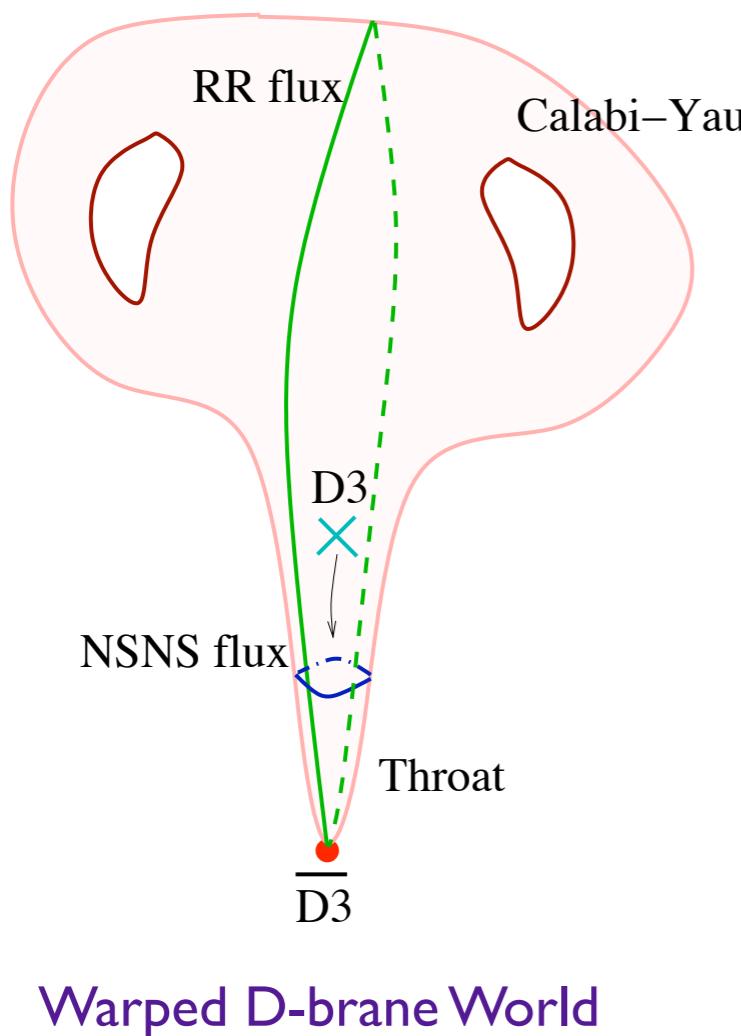
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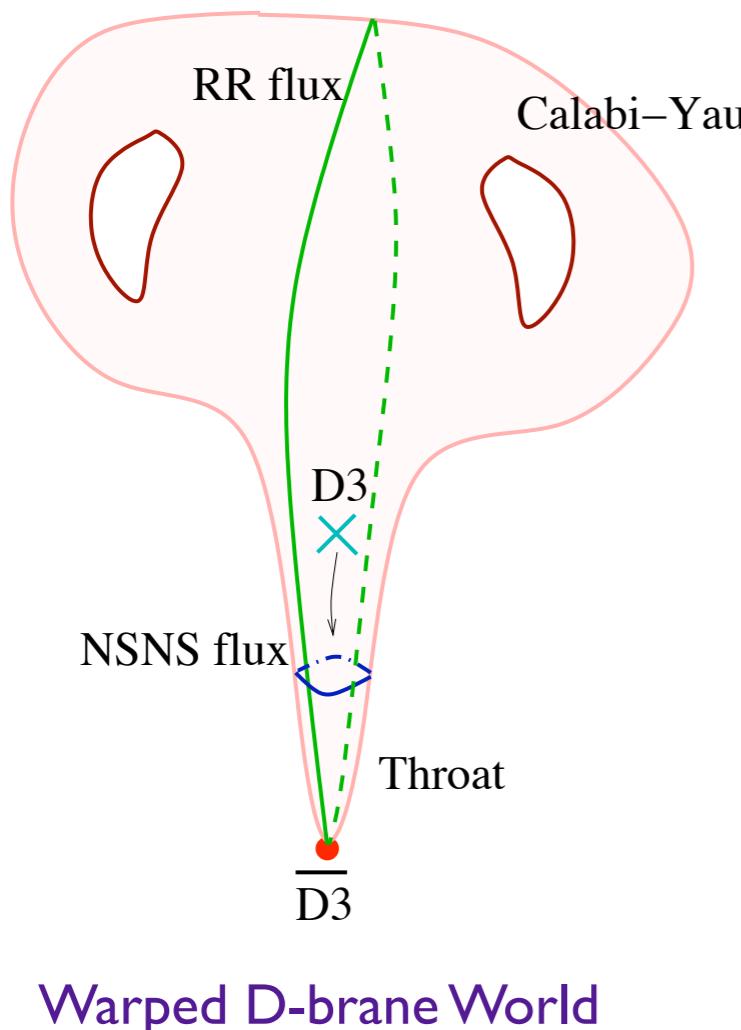
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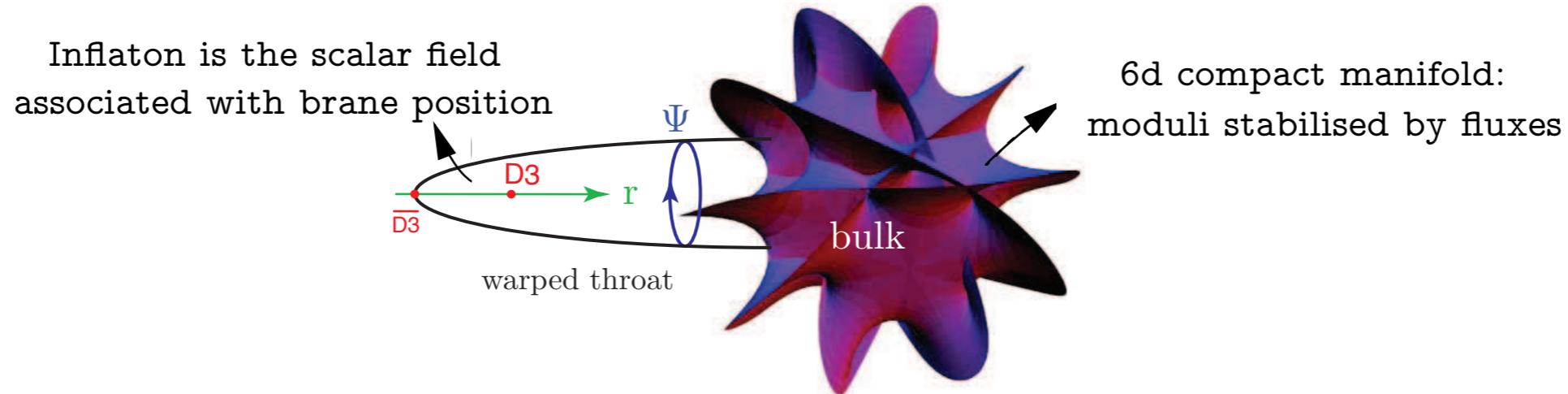
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Application to cosmology: D-brane Inflation

# D-Brane Inflation

- A D3-brane moves through a warped space. Internal fluxes help stabilise dangerous light scalar fields. [Kachru et al. '03].



- D-brane **position(s)** play role of inflaton(s)  $\Rightarrow$  Geometrical interpretation of inflaton. Scalar potential generated due to anti-D3-brane at end of throat.
- Two different scenarios (generically multifield inflation):

- \* **Slow roll inflation:** flat potential required. **Standard predictions for single field case**  $n_s \sim 0.96$   $r \ll 1$   $|f_{NL}| \ll 1$ . **Multifield dynamics** also arises. [Kachru et al. '03].

$$S \supset \int d^4x \sqrt{-g_4} \left[ \frac{M_p^2}{2} R + \partial\phi^i \partial\phi^j g_{ij} + V(\phi^i) + \dots \right]$$

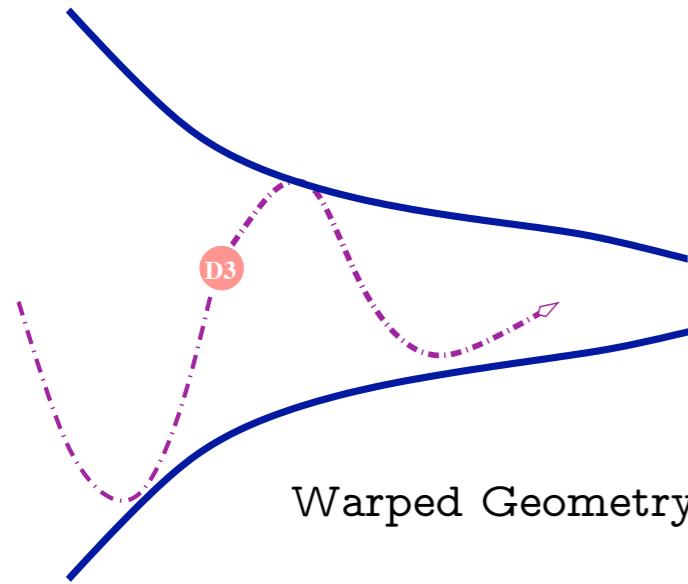
- \* **DBI inflation:** no flat potential is required. Inflaton Lagrangean characterised by non-canonical kinetic terms. [Silverstein-Tong '04]

$$S \supset \int d^4x \sqrt{-g_4} \left[ \frac{M_p^2}{2} R + P(X^{ij}, \phi^i) + \dots \right]$$

# DBI Inflation and its signatures

- D-brane motion in internal 6D space described by Dirac-Born-Infeld + Wess-Zumino action  $\Rightarrow$  multifield cosmological evolution

[Easson-Gregory-Mota-Tasinato-Zavala '07]



$$S_4 = \int d^4x \sqrt{-g} \left\{ \frac{M_{Pl}^2}{2} R - \left[ h^{-1} \sqrt{1 - h \dot{\phi}^n \dot{\phi}^m g_{mn}} - h^{-1} + V(\phi^l) \right] \right\}$$

$$M_{Pl}^2 = V_6 / \kappa_{10}^2 \quad \kappa_{10}^2 = (2\pi)^7 g_s^2 \alpha'^4$$

- Non-standard kinetic terms give rise to speed limit for D-brane motion in compact space  $h^{-1} > \dot{\phi}^2$ 
  - Combined with strong warping  $\Rightarrow$  accelerating trajectories without requiring a flat potential (cf. slow roll)

- DBI cosmological signatures:

- Enough inflation requires brane to move close to speed of light:  $c_S^2 \ll 1$   
 where  $c_S^2 = (1 - h\dot{\phi}^2)$

- Small sound velocity implies large departures from Gaussian spectrum.  
 Quantified by non-negligible bispectrum  $\langle \zeta \zeta \zeta \rangle \sim f_{NL} \langle \zeta \zeta \rangle^2$

$$f_{NL} = -\frac{1}{3} \left( \frac{1}{c_S^2} - 1 \right)$$

$$-114 < f_{NL} < 166 \quad (68\% cl \ 1\sigma) \Rightarrow$$

$$c_S^2 \gtrsim 3 \times 10^{-3}$$

- In Wilson line case also observable gravitational wave spectrum can be generated  $0.16 \lesssim r \lesssim 0.24$

[Avgoustidis-Zavala '09]

# Cosmic superstring production at end of D-brane inflation

- A key side effect of D-brane inflation is the formation of **cosmological defects** from due to  $Dp/\bar{D}p$  annihilation. [Tye et al. '02]

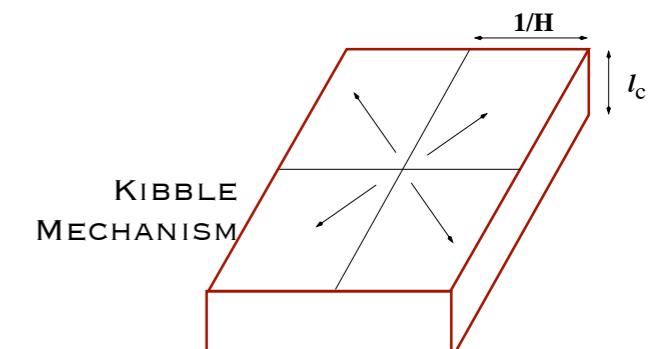
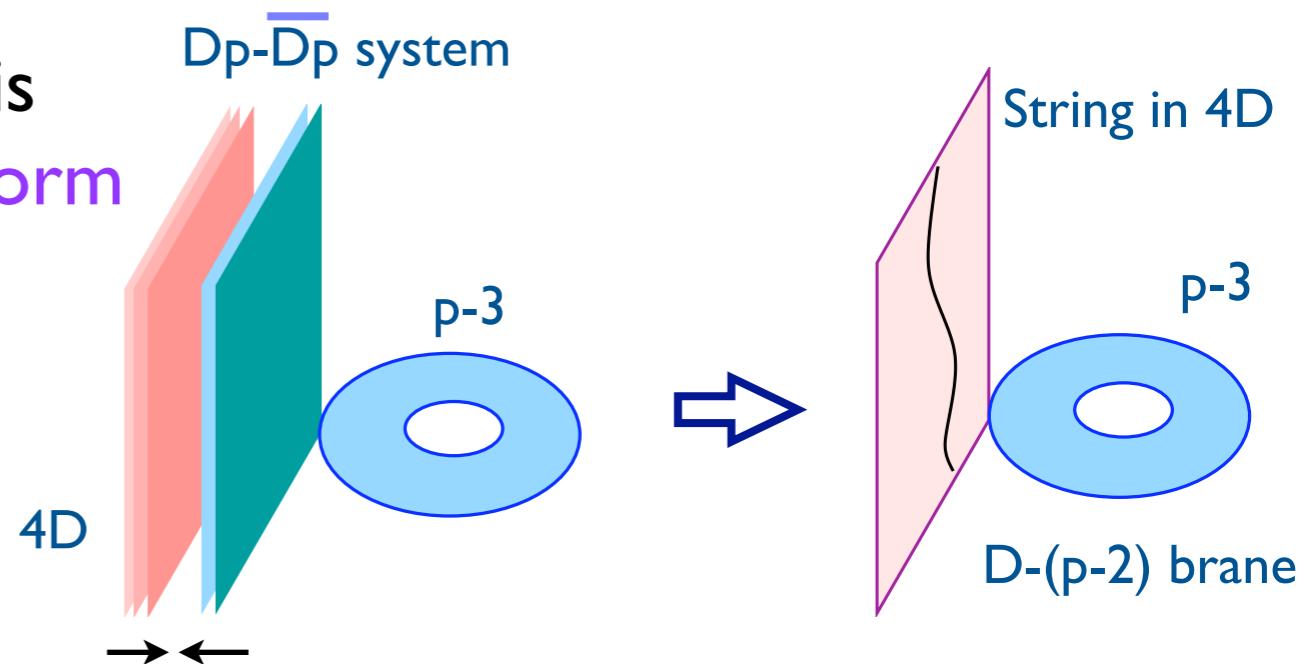
- When  $Dp/\bar{D}p$  collide  $D(p-2)$ -branes form.

- Topologically, a variety of defects may be produced at end of inflation: domain walls, cosmic strings, monopoles.

- Cosmologically, only strings can form via Kibble mechanism in 4D: D-strings, F-strings, bound states.

- Dimensionless tension depends on type of string formed. A spectrum of tensions for cosmic superstrings arises (cf. field theoretic strings):

$$10^{-12} \lesssim G\mu \lesssim 10^{-6}$$

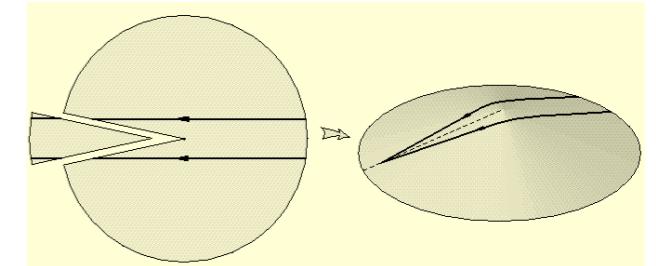


# String theory in the Sky?

Cosmic strings can be detected via gravity, rather than particle physics:

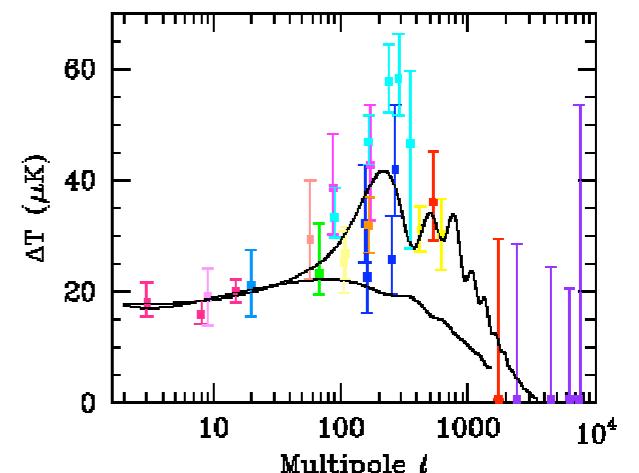
- Gravitational lensing. Requires appropriate alignment of source behind string.

Some candidates, no detection.



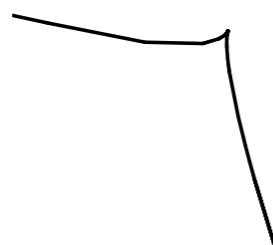
- Gravitational perturbations

CMB:  $G\mu < 10^{-7}$    Pulsars:  $G\mu < 10^{-6}$



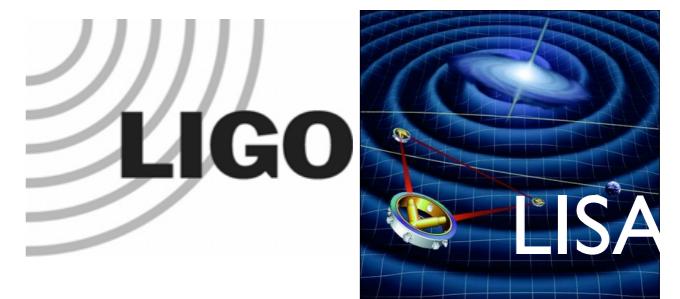
- **Gravitational waves.** Distortions of spacetime.

Gravitational wave bursts are produced by **cusps**, points at which the string reaches the speed of light. Possible detection at LISA, LIGO.



Cups formation in a loop

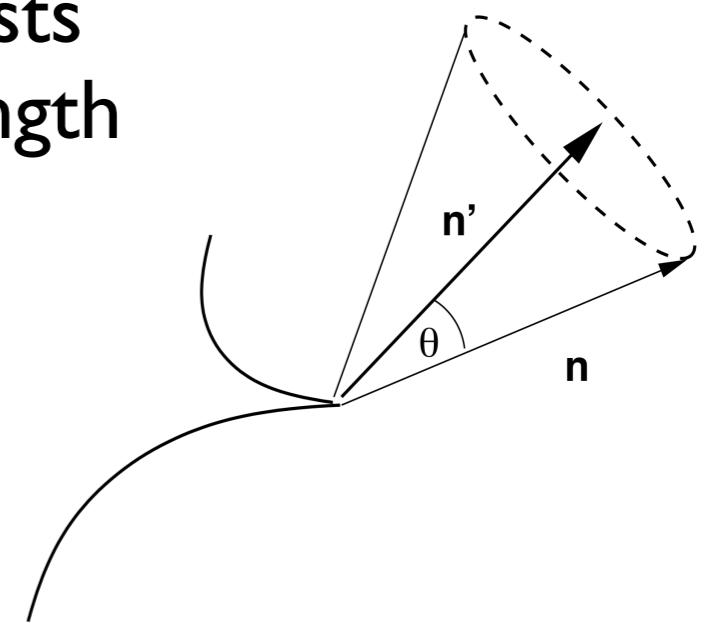
Act as a strong source of gravitational radiation



# Gravitational Wave Signal (4D)

[Damour-Vilenkin (DV) '00-'05]

- Cusp beams out a strong pulse of gravitational radiation in a cone centered on the cusp.
- The form and amplitude of cusp gravity wave bursts can be computed as a function of mass per unit length of the string.
- The resulting cusp wave form has a power law behaviour in frequency.
- If cosmic string arise from D-brane inflation, it moves in  $(4+n)$  dimensions



what is the effect of extra dimensions on the cusp wave form?

# Effects of extra dimensions in gravitational wave bursts (GWB)

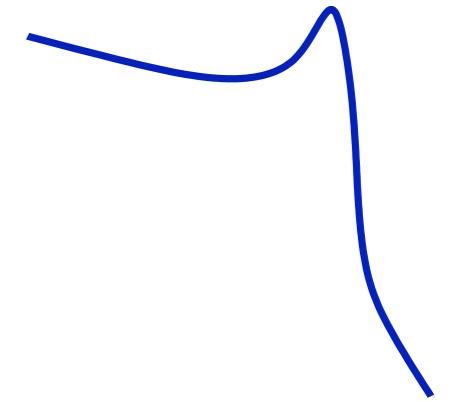
- \* Extra dimensions reduce the probability of intercommutation  
 $P \lesssim 1 \Rightarrow$  improve signal [DV '05]
- \* But also the kinematics of strings is different:  
[O'Callaghan-Chadburn-Geshnizjani-Gregory-Zavala '10]

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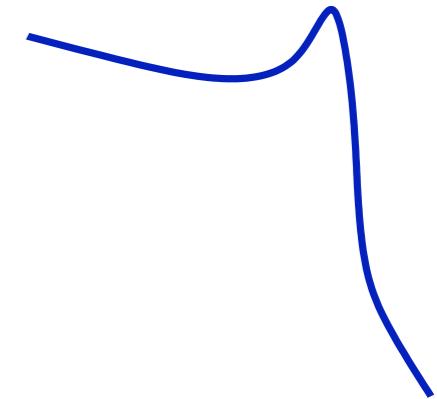
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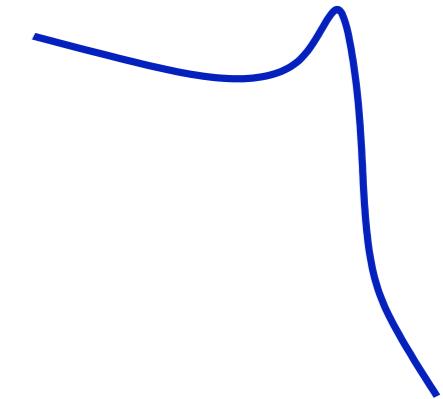
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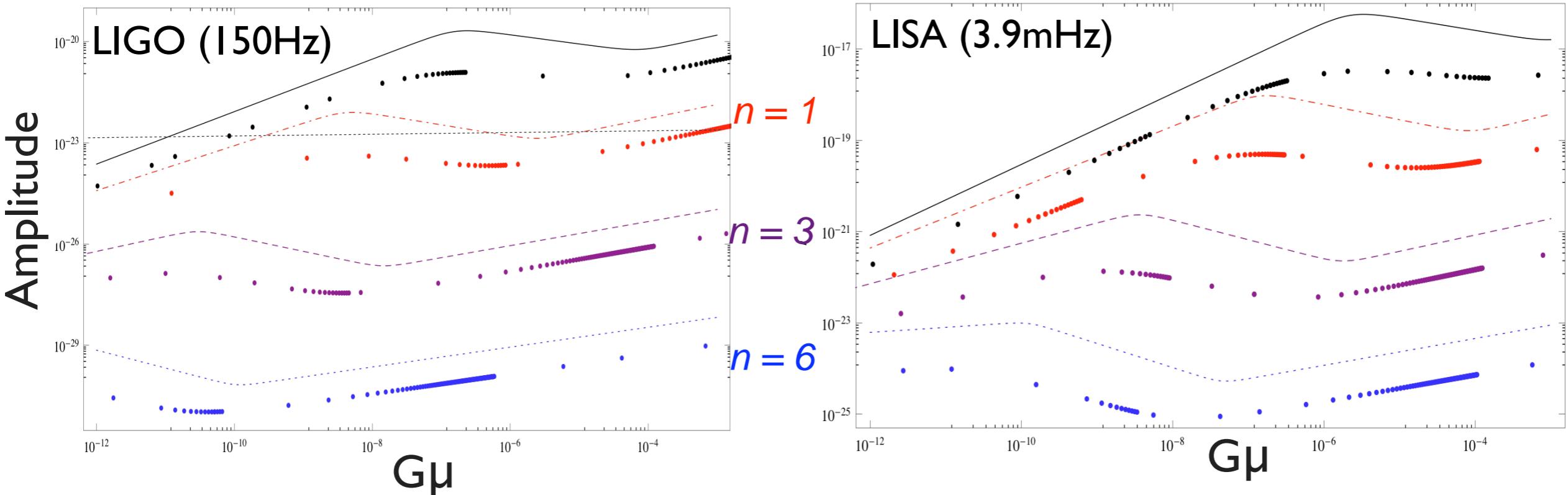
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# Effects of extra dimensions in gravitational wave bursts (GWB)

- ★ Extra dimensions reduce the probability of intercommutation  
 $P \lesssim 1 \Rightarrow$  improve signal [DV '05]
- ★ But also the **kinematics of strings** is different:  
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  - ★ Round off effect due to motion in extra dimension
  - ★ Decrease cusp formation probability and thus gravitational wave background as compared to 4D
  - ★ Amplitude depends on number of extra dimensions



# Summary

- \* String theory is very powerful guiding tool to investigate physics beyond standard models of Cosmology ( $\Lambda$ CDM) and particle physics.
- \* D-brane Inflation Models.
  - Distinctive features from generic models with more than one field, such as non-Gaussianities, isocurvature perturbations, detectable tensor perturbations.
  - End of inflation. Can give rise to detectable cosmic superstrings.
- \* Other stringy inflationary models:
  - Closed string modulus plays role of the inflaton: modular inflation.
  - Slow roll inflation with one or more fields.
  - No D-branes involved  no cosmic superstrings!