Method and results of the analysis of data on vertical rigidities of cosmic rays cutoff in the geomagnetic field

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Abstract: The method of the analysis of data on vertical rigidities of cosmic rays cutoff is presented based on the particle trajectory calculations in the geomagnetic field described by the Tsyganenko-89 model. The essence of the method consists in the fact that both the experimental data and the calculation results are described in the form of change of their value relative to the values related to IGRF. The value of these relative changes depends quite certainly on the rigidity itself and on the level of geomagnetic field disturbance (the $K_p$-index). The proposed model describes well calculation data and agrees with the results of cutoff rigidity measurements carried out by satellites.

Introduction

The effective cutoff rigidity (ECR) is an important parameter, that characterizes the given point of the near-Earth space, and defines the fluxes at this point as both galactic cosmic rays, and solar energetic particles (SEP) generated during solar flares. These fluxes, along with trapped particles of the Earth’s radiation belt, determine the radiation situation onboard the spacecraft. In addition, the geomagnetic cutoff effect allows one to study SEP spectra and fluxes using the data from both neutron monitors (NM), and spacecraft situated inside the magnetosphere.

In practice, the vertical ECR is usually applied (EVCR), which represents a fine estimate averaged over the total ECR body angle [1].

The ECR calculations are based on numerical integration of the equations of motion of charged particles in the geomagnetic field described by any model. Till now, the ECRs were calculated by this technique both for the points of the global NM network, and for many typical orbits of spacecraft, the International Space Station for instance [2, 3]. Note that the ECR depends both on geomagnetic disturbance level, and on the local time [4]. The possibilities of direct ECR calculation permanently grow due to growing power of computers; however, some practical tasks require simpler and less labor-consuming ECR calculation techniques, such as interpolation ones. So, Smart et al have calculated a set of EVCR tables for various conditions and developed the tools for interpolation [5].

Proposed EVCR determination technique

The approach we propose is based on interpolation of a basic set of ECRs calculated with using a series of physical concepts and computational models. The changes of the EVCR value under an effect of magnetosphere disturbance (described by $K_p$-index) and local time $(T)$ using the Tsyganenko’s-89 geomagnetic model [6] are used in the given technique as corrections, whose values are described by the attenuation factor $\Delta$, determined by Nymmik [7] as:

$$\Delta(R_0, K_p, T) = \frac{R_0}{R(R_0, K_p, T)} - 1$$  \hspace{1cm} (1)$$

where $R_0$ is the initial rigidity calculated for the IGRF model. The reverse conclusion follows from this statement, namely: if the value $\Delta(R_0, K_p, T)$ is known for the given point, then the real EVCR can be determined by the same formula. The analysis of the results of particles trajectory calculations, carried out according to Tsyganenko’s mag-
Figure 1: The values of factors of attenuation of effective vertical rigidities in terms of initial rigidity IGRF and magnetosphere disturbance levels (from below upwards: $K_p = 0, 1, 3, 5$). Calculation data (points) and their approximations (lines).

The certain distinction lies in the fact, that, with the purpose of increasing the accuracy, in determining the penumbra parameters we have used both the more detailed scale of rigidities - $\Delta R = 0.003 \, GV$ (and $\Delta R = 0.001 \, GV$ in separate cases), and the algorithm of the fifth order of accuracy [8] for trajectory integration with accounting for the time dependence of the geomagnetic field.

Our calculations of EVCR for the epoch 2005 for the IGRF model are presented in Table 1. Due to calculation’s performance limits some values in this table with $R < 0.01 \, GV$ are not real EVCRs, but their’s upper possible values. The EVCR values from Tsyganenko’s model for calculating $\Delta_1(R_0, K_p)$ (see Eq. 3) were determined by averaging the data calculated for the set of local time moments $T = 1, 2, \ldots 24$ (hours).

Formulas (3) and (4) are applicable for determining EVCR at

$$R_0 \geq R_{crit} = 0.238 + 0.029 \cdot K_p \quad (5)$$

At smaller rigidities, the zero rigidities appear in the evening-night sector - during these periods the given point becomes accessible for penetrating of particles with rigidities smaller than the initial value of the rigidity range looked through, that is equal to $0.001 \, GV$.

Formulas (1-5) determine the EVCR value for any point of the near-Earth spacecraft’s orbit as a function of coordinates, magnetosphere disturbance level and local time value. The use of these formulas allows one to avoid the necessity of performing resource-consuming trajectory calculations of charged particles in the magnetosphere. The validity of this model is confirmed by the results of our numerical calculation of trajectories of charged particles in the geomagnetic field described by the superposition of the IGRF and Tsyganenko-89 models. One of difficulties of such calculations is the verification of numerical integration results. In addition to standard techniques of checking the accuracy control, we have compared our calculations at fixed points with the Smart’s et al data [2, 3]. So, we convinced in well coincidence of $R > 0.5 \, GV$. 

The mean value of the attenuation factor, which does not depend on the local time ($T$):

$$\Delta_1 = 0.174 \cdot R_0^{-0.7} \cdot \exp(0.304 \cdot K_p) \quad (3)$$

where $\gamma = 1.7$ at $R_0 \geq 1 \, GV$ and $\gamma = 1.7 \cdot R_0^{-0.1}$ at $R_0 < 1 \, GV$. Figure 1 presents the calculated values of the correction quantity $\Delta_1$, which are well described by formula (3). The second multiplier describes the time dependence of EVCR on the mean value of rigidity, $K_p$-index, and on the local time $T$:

$$\Delta_2(R_0, K_p, T) = [1 + A \cdot B \cdot C] \quad (4)$$

where $A = (0.025 + 0.0344 \cdot K_p)$, $B = R_0^{(2.42 - 0.075 \cdot K_p)}$ and $C = \sin(\frac{2\pi}{24}(T + \Delta T))$

and where $\Delta T$ for the range of rigidities $R = 0.2 \pm 0.6 \, GV$ has the value of 2 hours approximately.
Table 1: Basic EVCR table calculated for IGRF epoch 2005.

<table>
<thead>
<tr>
<th>λ, o</th>
<th>0</th>
<th>10</th>
<th>20</th>
<th>30</th>
<th>40</th>
<th>50</th>
<th>60</th>
<th>70</th>
<th>80</th>
<th>90</th>
<th>100</th>
</tr>
</thead>
<tbody>
<tr>
<td>Geographic East Longitude, o</td>
<td>1.778</td>
<td>2.018</td>
<td>2.165</td>
<td>2.357</td>
<td>2.711</td>
<td>2.339</td>
<td>1.460</td>
<td>0.713</td>
<td>0.389</td>
<td>0.560</td>
<td>1.166</td>
</tr>
</tbody>
</table>

Calculation results and comparison with experiments

As noted above, the model is correct for EVCR values, which are greater than some quantities, whose particular values depend on \( K_p \). We have estimated this critical value rigidity \( R_{crit} \) (see Eq.5). We have compared the calculated values of quantity \( \Delta \) with EVCR values obtained in various experiments, determined from observations of the boundary of penetration of SEP (protons and alpha-particles). Figure 2 presents the model approximations of attenuation factors for the magneto-quiet conditions, characterized by values \( K_p = 0 \) and \( K_p = 1 \). A small arrow stands at the place of critical rigidities \( R_{crit} \). The set of experimental data is well described by model functions, as a rule.

Conclusion

This model represents the first approximation to the version of the express-calculation technique called for providing a possibility of fast (online) determination of fluxes of particles penetrating to the Earth satellite and orbital stations. In addition, the model is designed for fast determination of the function of particles penetration to projected Earth satellite orbits. This model does not cover the situations, where, along with the magnetosphere disturbance described by the \( K_p \)-index, the ring current, described by the \( Dst \)-index, is simultaneously strengthened in the magnetosphere. In our opinion, such extreme situations occur rather rarely and can be neglected in the first approximation. Besides, some parts of the model need updating.
Figure 2: Model dependencies of the time averaged attenuation factor on the RIGRF value, for \(K_p = 0\) and \(K_p = 1\), and the \(\Delta_1\) values calculated from the results of experiments on determining the boundaries of penetration of particles of various rigidity to the near-Earth satellite orbits. The sources of information\[9, 10, 11, 12, 13\] are deciphered in the Figure by the first author’s name.

Acknowledgment

This work is a tribute to the memory of Prof. S.N. Kuznetsov, whose advices helped us greatly.

References