

Generalized susceptibilities near the QCD critical point

Rodrigo García Formentí Mendieta
Eleazar Cuautle Flores

Instituto de Ciencias Nucleares
UNAM



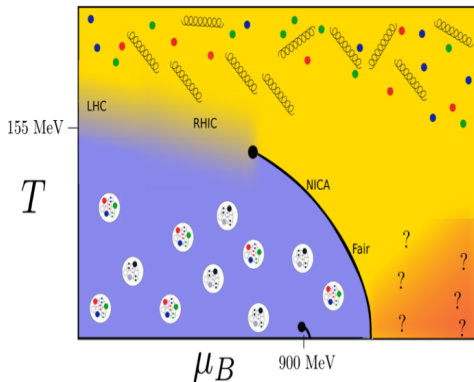
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Outline

- 1 Motivation: QCD phase diagram
- 2 Generalized susceptibilities
- 3 Study with Hadron Resonance Gas

Motivation: QCD phase diagram

- Calculations indicate a possible critical point in the QCD phase diagram.
- The critical point can be studied from the derivatives of the partition function.
- Experimentally, heavy-ion collisions are used.



Generalized susceptibilities

Assuming the system is described by a grand canonical ensemble, with partition function

$$Z(T, \mu_q, V). \quad (1)$$

The generalized susceptibility of order n associated with the conserved quantity q is the n -th derivative of the pressure with respect to μ_q/T :

$$\chi_q^{(n)} = \frac{1}{VT^3} \frac{\partial^n \log Z}{\partial (\mu_q/T)^n} = \frac{\partial^n p(T, \mu_q)}{\partial (\mu_q/T)^n}. \quad (2)$$

These quantities are sensitive to the critical point. They related to the correlation length ξ ¹

$$\chi_q^{(2)} \sim \xi^2, \quad \chi_q^{(3)} \sim \xi^{4,5}, \quad \chi_q^{(4)} \sim \xi^7. \quad (3)$$

In the thermodynamic limit, ξ diverges at the critical point².

¹M.A. Stephanov, Phys.Rev.Lett.102:032301,2009

²Erithiof Karsch et al. Phys Lett B695:136-142 2011

How do we measure susceptibilities?

Susceptibilities are accessible experimentally via event-by-event fluctuations of conserved charges. Let ΔN be the net charge in a given event and $\delta N = \Delta N - \langle \Delta N \rangle$ its deviation. The statistical cumulants are defined as:

$$C_1 = \langle \Delta N \rangle, \quad C_2 = \langle (\delta N)^2 \rangle, \quad C_3 = \langle (\delta N)^3 \rangle, \quad (4)$$

$$C_4 = \langle (\delta N)^4 \rangle - 3\langle (\delta N)^2 \rangle^2.$$

Cumulants are related to the generalized susceptibilities³

$$C_n = VT^3 \chi_q^{(n)} \quad (5)$$

³M.A. Stephanov, Phys.Rev.Lett.102:032301,2009

Phase diagram using the Hadron Resonance Gas

Hadron Resonance Gas model


The system is modeled as a non-interacting gas of hadrons and resonances, described by the partition function:

$$\log Z^{\text{HRG}}(T, V, \mu) = \sum_{i \in \text{PDG}} \pm \frac{V g_i}{2\pi^2} \int_0^\infty dp p^2 \log(1 \pm e^{-(E_i - \mu_j)/T}) \quad (6)$$

with $+$ or $-$ for a boson or a fermion. With chemical potential

$$\mu_i = B_i \mu_B + Q_i \mu_Q + S_i \mu_S \quad (7)$$

The ideal HRG model lacks a critical point. However, it can be phenomenologically introduced by incorporating interactions in the gas ⁴.

⁴V. Vovhenko et al., Comput. Phys. Commun. 244, 295 (2019) 

HRG, Quantum Van der Waals


The simplest way to include interactions and a critical point is by generalizing the van der Waals gas:

$$p(T, \mu) = \sum_i p_i^{\text{ideal}}(T, \mu_i^*) - \sum_{ij} a_{ij} n_i n_j, \quad (8)$$

with

$$\mu_i^* + \sum_j b_{ij} p_j^* - \sum_j (a_{ij} + a_{ji}) n_j = \mu_i. \quad (9)$$

Using experimental data, the a_{ij} and b_{ij} can be determined. THERMAL-FIST package⁵, to calculate the susceptibilities.

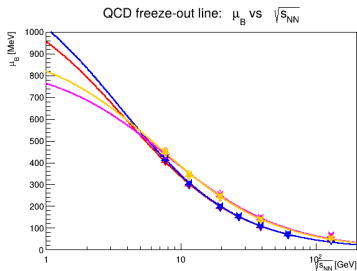
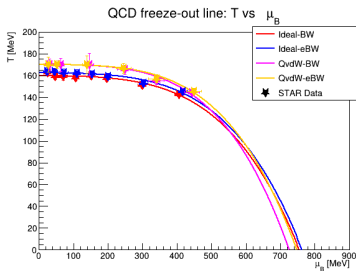
⁵V. Vovhenko et al., Comput. Phys. Commun. 244, 295 (2019) 

Freeze-out Line

Experimentally, measurements are restricted to the chemical freeze-out line. Using STAR data, and the parametrization⁶:

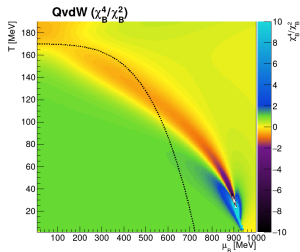
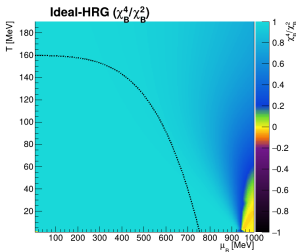
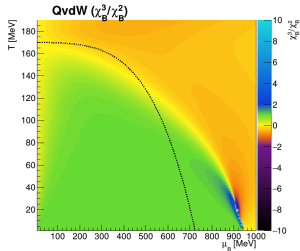
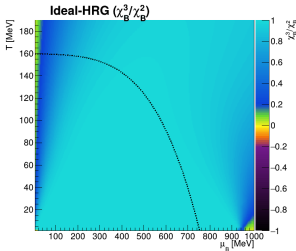
$$\mu_B(\sqrt{s_{NN}}) = \frac{a}{1 + b\sqrt{s_{NN}}}, \quad T(\mu_B) = c - d\mu_B^2 - e\mu_B^4, \quad (10)$$

Model	a [GeV]	b [GeV ⁻¹]	c[GeV ⁻³]	d[GeV]	e [GeV ⁻¹]
I-HRG	0.160 ± 0.001	0.035 ± 0.04	0.430 ± 0.28	1.201 ± 0.20	0.25 ± 0.05
QvdW	0.170 ± 0.001	0.030 ± 0.04	0.559 ± 0.20	0.858 ± 0.05	0.12 ± 0.01

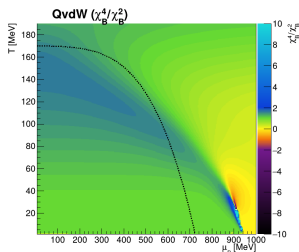
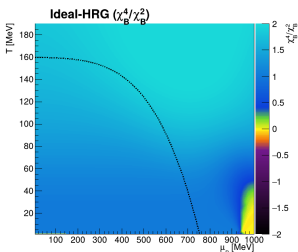
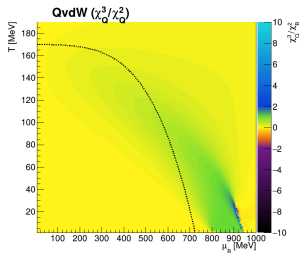
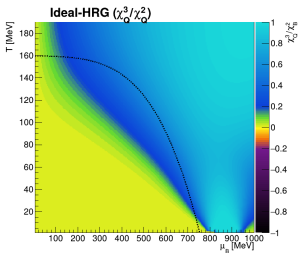


⁶F. Reyes and E. Cuautle, eprint arXiv:2606.14018

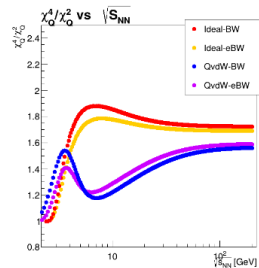
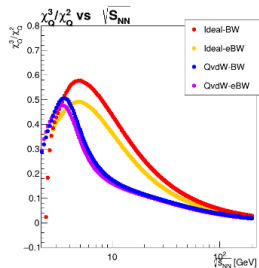
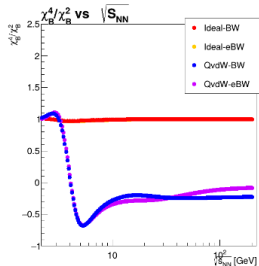
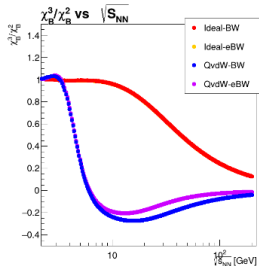
Ideal vs QvdW Susceptibilities (Baryon number) $(T_c, \mu_c) = (19.5 \text{ MeV}, 15 \text{ MeV})$



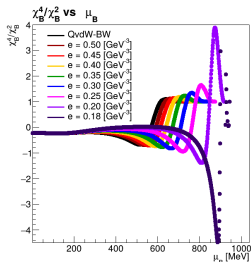
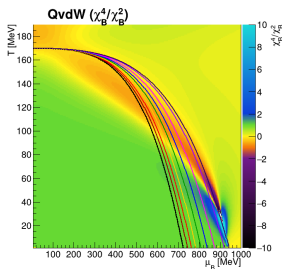
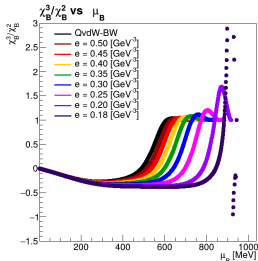
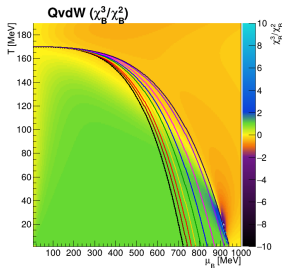
Ideal vs QvdW Susceptibilities (Electric Charge)



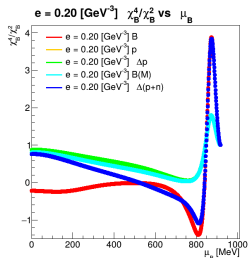
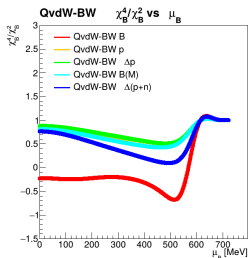
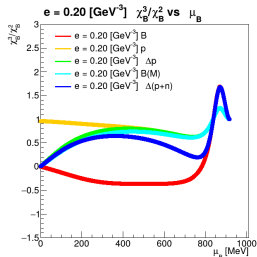
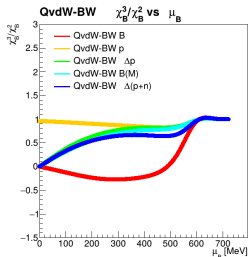
Susceptibilities in the freeze-out line



Susceptibilities in the freeze-out line



Susceptibilities of baryon number proxy



Summary

- Generalized susceptibilities are important observables in the search for the critical point.
- The HRG model allows estimating the phase diagram and comparing scenarios with and without a critical point against experimental data.
- Susceptibility ratios strongly depend on the proximity of the chemical freeze-out line to the critical point.
- The susceptibility depends on the proxy for baryon number. The neutrons measurements are very important to reach higher better behavior of susceptibilities and the critical end point.