



Pions and **nucleons** at counter**point**

Khépani Raya Montaña



Universidad
de Huelva

Khépani Raya · Adnan Bashir · Daniele Binosi ·
Craig D. Roberts · José Rodríguez-Quintero

Pseudoscalar Mesons and Emergent Mass

<https://doi.org/10.1007/s00601-024-01924-2>

MWPF 2025. León, México

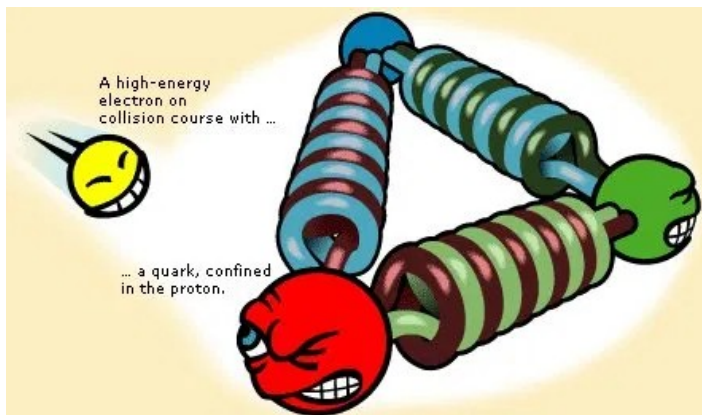
October 20-24, 2025

QCD: Emergent Phenomena

- **QCD** is characterized by two **emergent** phenomena:
confinement and dynamical generation of mass (**DGM**).



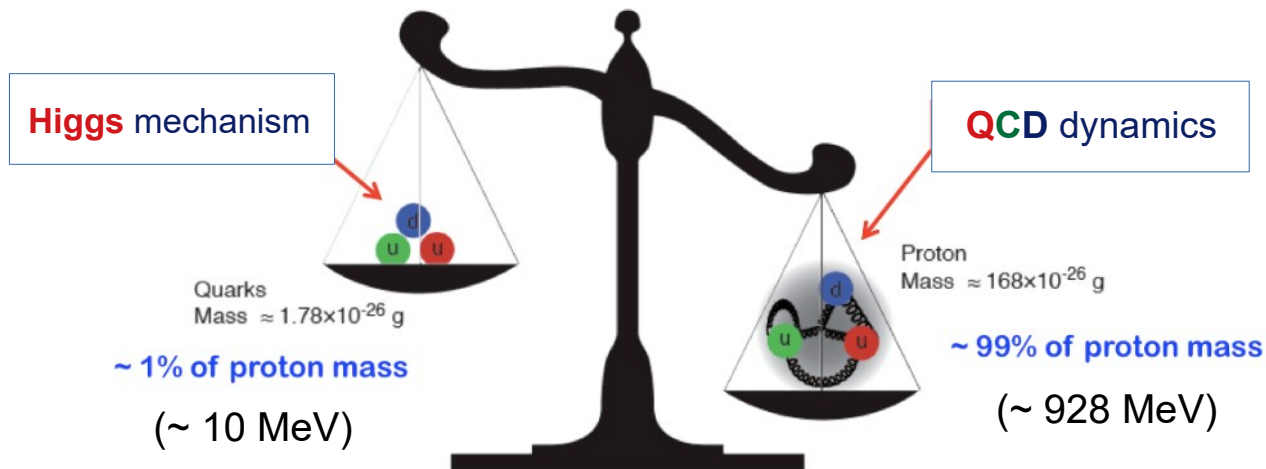
- ♦ Quarks and gluons not *isolated* in nature.
- ➔ Formation of colorless bound states: “**Hadrons**”
- ➔ **1-fm scale** size of hadrons?



$$\begin{aligned}\mathcal{L}_{\text{QCD}} &= \sum_{j=u,d,s,\dots} \bar{q}_j [\gamma_\mu D_\mu + m_j] q_j + \frac{1}{4} G_{\mu\nu}^a G_{\mu\nu}^a, \\ D_\mu &= \partial_\mu + ig \frac{1}{2} \lambda^a A_\mu^a, \\ G_{\mu\nu}^a &= \partial_\mu A_\nu^a - \partial_\nu A_\mu^a - gf^{abc} A_\mu^b A_\nu^c,\end{aligned}$$



- ♦ Emergence of hadron masses (**EHM**) from QCD **dynamics**



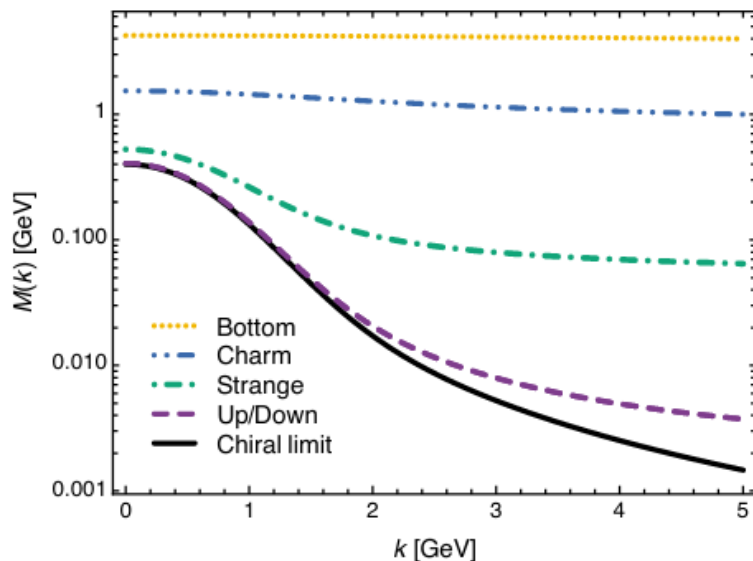
QCD: Emergent Phenomena

- **QCD** is characterized by two **emergent** phenomena:
confinement and dynamical generation of mass (**DGM**).

Can we trace them down to fundamental d.o.f?

Constituent masses

Dynamical Chiral Symmetry Breaking



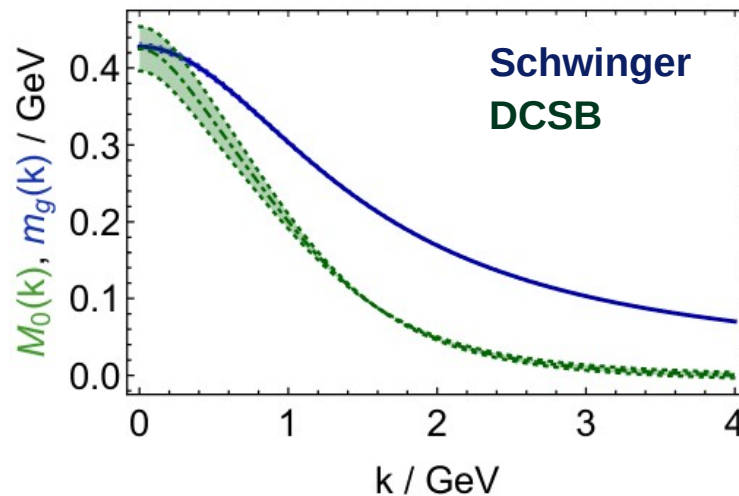
$$S_f^{-1}(p) = Z_f^{-1}(p^2)(i\gamma \cdot p + \mathbf{M}_f(p^2))$$

“Higgs” masses



$$\begin{aligned}\mathcal{L}_{\text{QCD}} &= \sum_{j=u,d,s,\dots} \bar{q}_j [\gamma_\mu D_\mu + m_j] q_j + \frac{1}{4} G_{\mu\nu}^a G_{\mu\nu}^a, \\ D_\mu &= \partial_\mu + ig \frac{1}{2} \lambda^a A_\mu^a, \\ G_{\mu\nu}^a &= \partial_\mu A_\nu^a - \partial_\nu A_\mu^a - gf^{abc} A_\mu^b A_\nu^c,\end{aligned}$$

- Emergence of hadron masses (**EHM**) from QCD **dynamics**

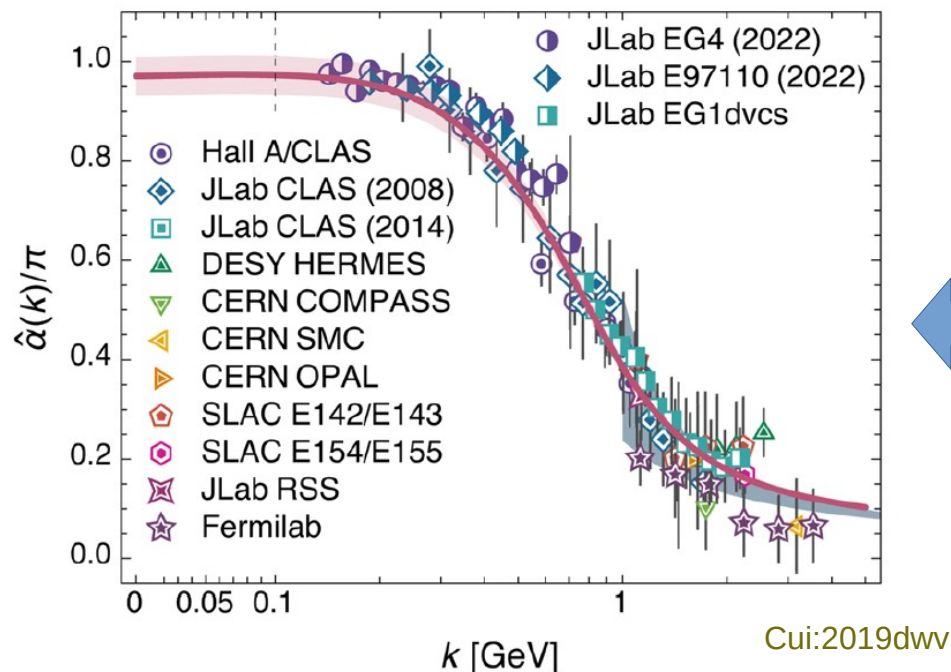


Gluon and quark *running masses*

QCD: Emergent Phenomena

- **QCD** is characterized by two **emergent** phenomena:
confinement and dynamical generation of mass (**DGM**).

Can we trace them down to fundamental d.o.f?

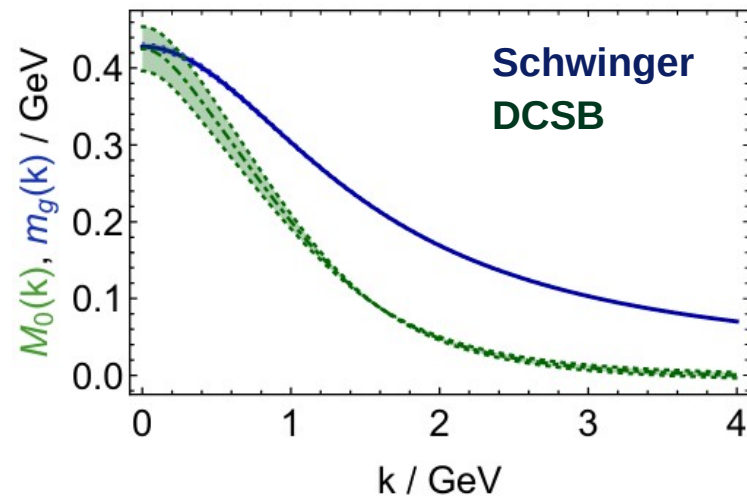


$$\mathcal{L}_{\text{QCD}} = \sum_{j=u,d,s,\dots} \bar{q}_j [\gamma_\mu D_\mu + m_j] q_j + \frac{1}{4} G_{\mu\nu}^a G_{\mu\nu}^a,$$

$$D_\mu = \partial_\mu + ig \frac{1}{2} \lambda^a A_\mu^a,$$

$$G_{\mu\nu}^a = \partial_\mu A_\nu^a - \partial_\nu A_\mu^a - gf^{abc} A_\mu^b A_\nu^c,$$

- ♦ Emergence of hadron masses (**EHM**) from QCD **dynamics**



Gluon and quark *running masses*

Mass Budgets

$$M_{u/d} \approx 0.3 \text{ GeV}$$

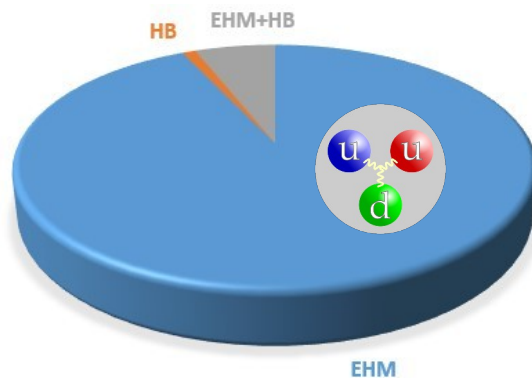
- What is the origin of **EHM**?
... its connection with e.g. **confinement** and **DCSB**?

➤ **Most** of the **mass** in the visible universe is contained within **nucleons**

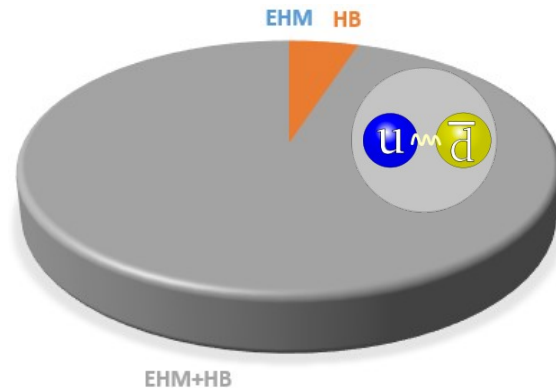
→ Their role in the universe is **crucial**

➤ The **Pion** is the most essential **NG boson** of **DCSB**.

→ It's mere **existence** is connected with **mass** generation in the **SM**



Proton mass budget



Pion meson mass budget

$$m_p = 0.938 \text{ GeV} \approx 2M_u + M_d$$
$$m_\pi = 0.14 \text{ GeV} \neq M_u + M_d$$

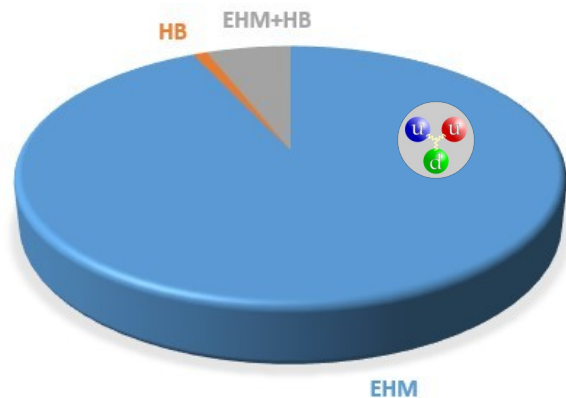
With the same building blocks, Nature makes one **heavy** and one **too light**. What's happening?

Mass Budgets

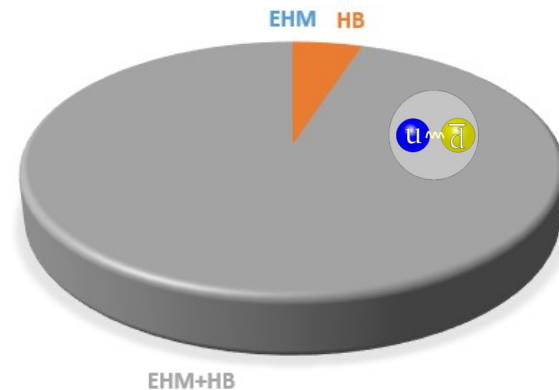
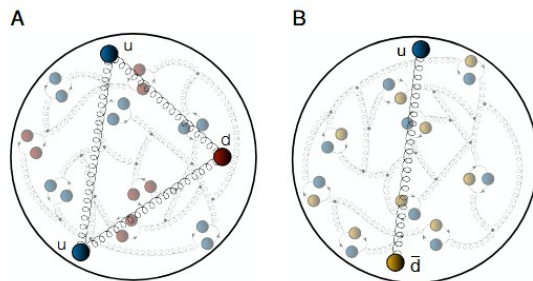
$$m_s/m_u \sim 20$$

$$f_K/f_\pi \sim M_s/M_u \sim 1.2$$

- The same **mechanisms** that make the **proton massive**, makes the **pion 'massless'**



Proton mass budget



Pion mass budget

$$m_p = 0.938 \text{ GeV} \approx 2M_u + M_d$$

$$m_\pi = 0.14 \text{ GeV} \neq M_u + M_d$$

- This **dichotomy** needs to be understood.

EIC, EicC, JLab, Amber...



Valence-quark distribution amplitudes (**DAs**)

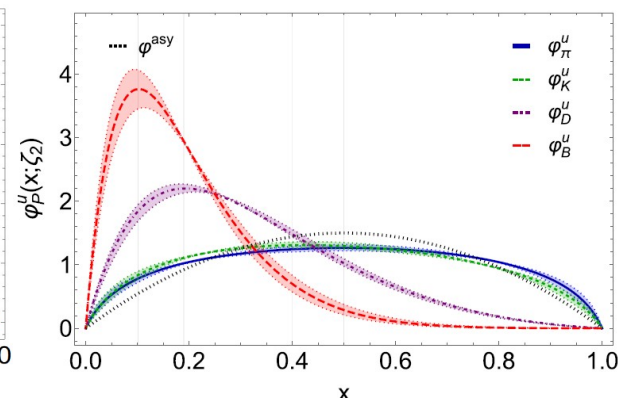
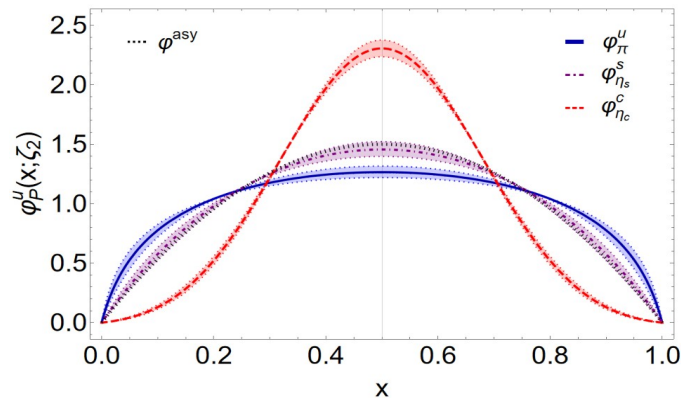
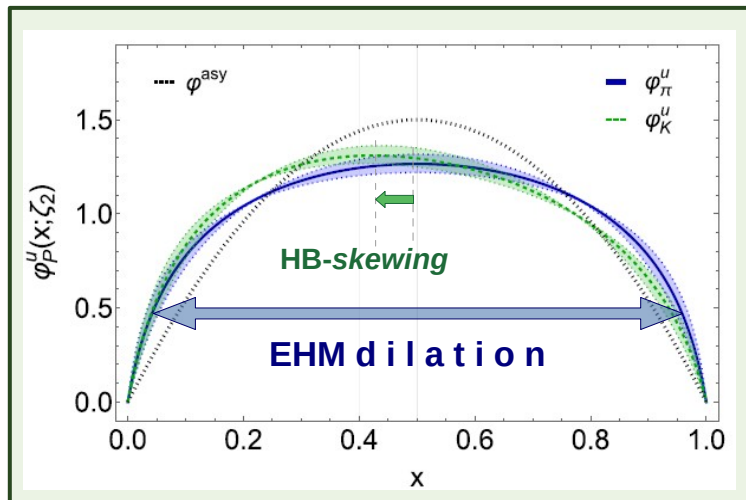
$$f_M \phi_M^q(x) = \text{tr} \int_{dk} \delta_n^x(k_M) \gamma_5 \gamma \cdot n \chi_M(k_-, P)$$

Light-front momentum fraction

Written in terms of **BSWF**

- **1-dimensional** projection of the **light-front wavefunction**.
- Clear **probe of EHM**, related with hard **exclusive processes**, etc.

Pseudoscalars' DAs



$$f_M \phi_M^q(x) = \text{tr} \int_{dk} \delta_n^x(k_M) \gamma_5 \gamma \cdot n \chi_M(k_-, P)$$

- **Broad** in the **light-sector**, **narrow** in the **heavy-sector**.
- Mild **skewness** for the **Kaon**, striking for **heavy-lights**.

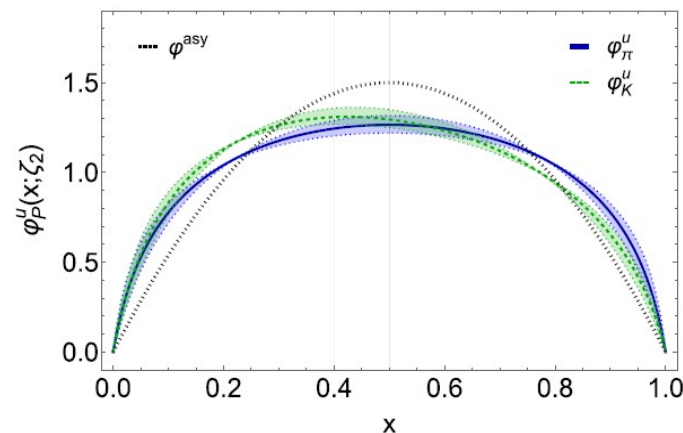
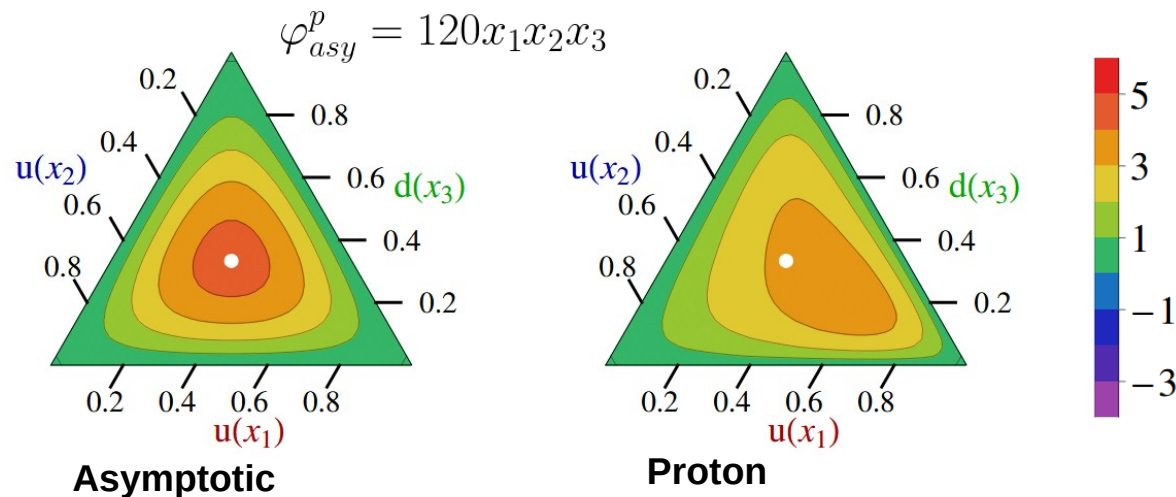
Ssbar in the middle

Proton DAs

$$(M_u = M_d)$$

- ✓ As the **pion**, the **proton DA** is **broader** than its asymptotic counterpart:

Mezrag:2017znp



- ✓ Nonetheless, even assuming **isospin symmetry**, the proton exhibits a **non-symmetric DA**.
 - This signals the **formation** of non-trivial quark-quark correlations (dynamical **diquarks**)

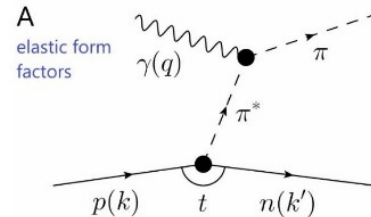
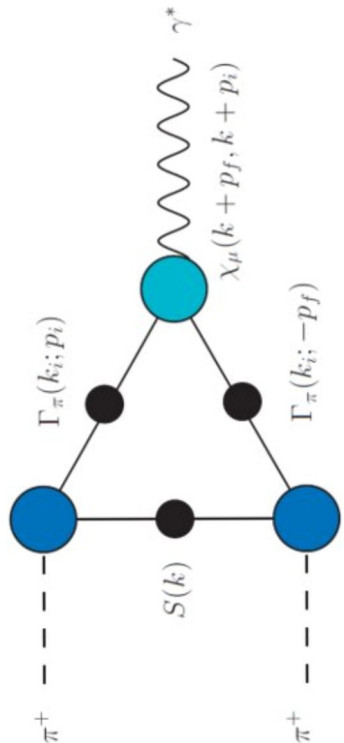
$$\left(\frac{d\sigma}{d\Omega}\right) = \left(\frac{d\sigma}{d\Omega}\right)_{\text{point-like}} \times |F_\pi(q^2)|^2$$

Electromagnetic Elastic Form Factors (**EFFs**)

$$P_\mu F_{\mathbf{P}}^q(Q^2) = \text{tr}_{\text{CD}} \int_{dk}^\Lambda \chi_\mu^q(k + p_o, k + p_i) \Gamma_{\mathbf{P}}(k_i; p_i) S_h(k) \Gamma_{\mathbf{P}}(k_o; -p_o)$$

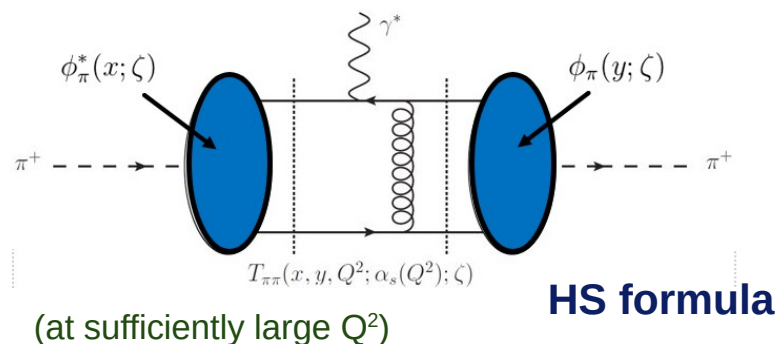
All can be written in terms of **propagators** and **vertices**

- Gives information on **momentum/charge** distribution.
- **Pion EFF** highly relevant for contemporary physics.



Elastic Form Factors

- In the **large- Q^2** regime, **QCD** connects the **EFF** and **DA**:



- At **leading-order**:

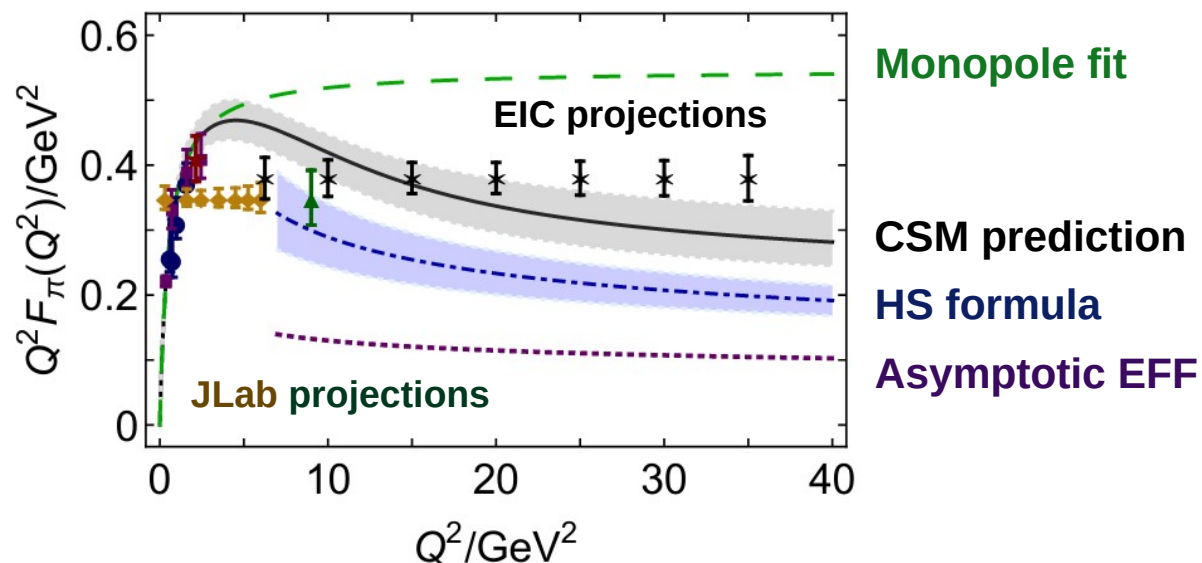
$$Q^2 F_P(Q^2) \stackrel{Q^2 \gg m_p^2}{\approx} 16\pi\alpha_s(Q^2) f_P^2 w_P^2(Q^2)$$

$$w_P = \frac{1}{3} \int_0^1 dx \frac{1}{x} \varphi_P(x; Q^2) \leftarrow \text{PDA} \rightarrow$$

→ The asymptotic behavior is weighted by f_P , a measure of **EHM**.

→ **Factorization/scaling violations** are proof of the validity of **QCD itself**.

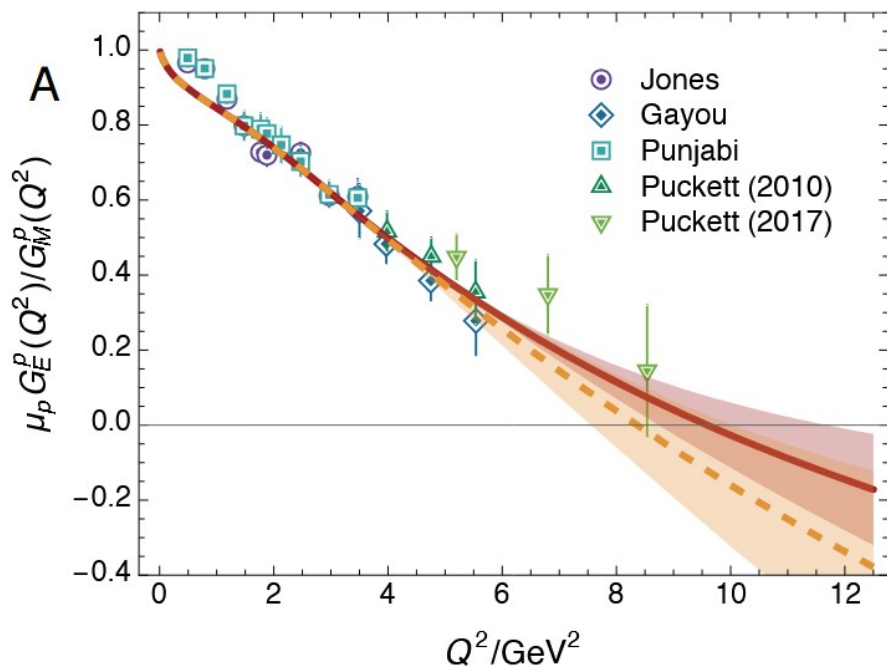
B



Elastic Form Factors

Cheng:2025yij

- For the **proton**, different theoretical and phenomenological analysis suggest a **zero-crossing** in the **GE/GM** ratio



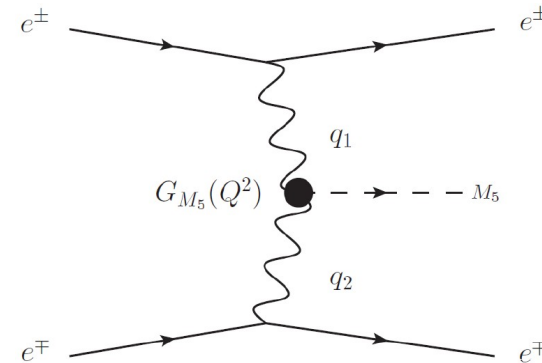
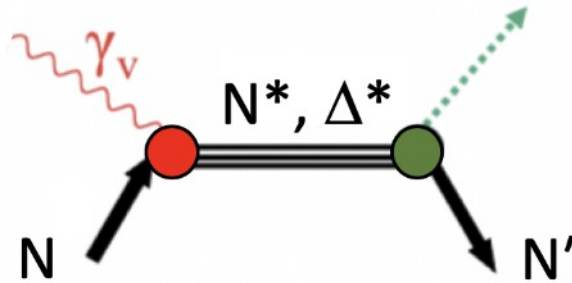
- This reflection on a **destructive interference** between the individual **valence-quark** contributions.
- In fact connected with non-trivial effects coming from orbital **angular momentum**.
- The **location** of the zero is highly sensitive to the **diquark** content of the **nucleon**.

- Modern analyses favor a **non-negligible** component of **axial-vector** diquark:

~ 25 – 35 %

- The rest comes from the expected **scalar** diquark.

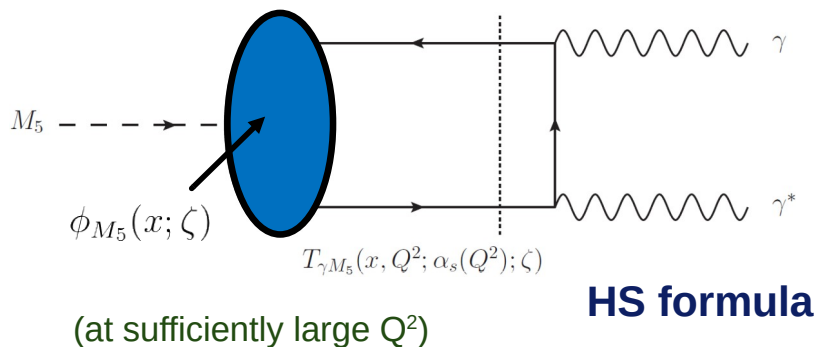
Transition Form Factors (TFFs)



- Gives information on **momentum/charge** distribution.
- Highly relevant for contemporary physics.

Two-photon TFFs

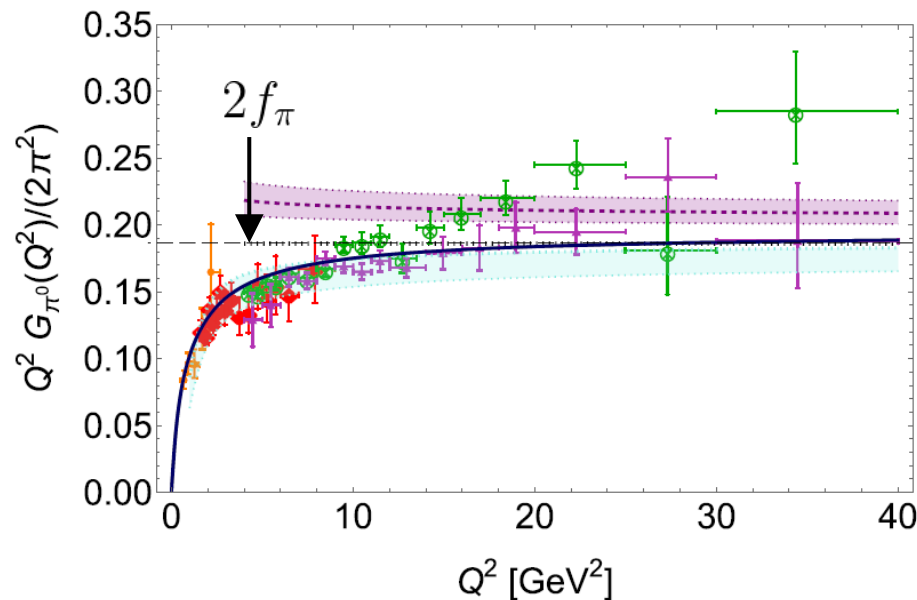
- In the **large- Q^2** regime, **QCD** connects the **TFF** and **DA**:



- At **leading-order**:

$$Q^2 G_{\mathbf{P}}^q(Q^2) \stackrel{Q^2 \gg m_p^2}{\approx} 12\pi^2 f_{\mathbf{P}}^q e_q^2 w_q(Q^2)$$

$$w_{\mathbf{P}} = \frac{1}{3} \int_0^1 dx \frac{1}{x} \varphi_{\mathbf{P}}(x; Q^2) \quad \leftarrow \text{PDA} \rightarrow$$

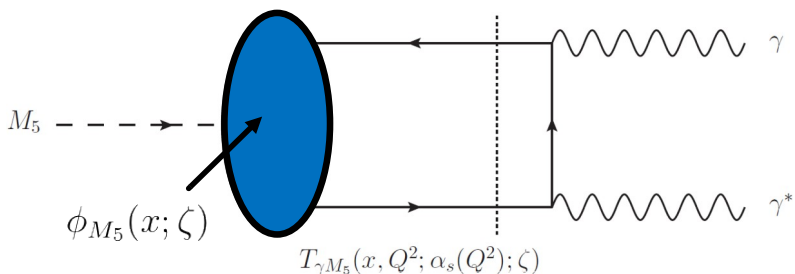


HS formula
CSM prediction
Asymptotic TFF

- The asymptotic behavior is weighted by $f_{\mathbf{P}}$, a measure of **EHM**.
- Factorization/scaling violations** are proof of the validity of **QCD itself**.

Two-photon TFFs

- In the **large- Q^2** regime, **QCD** connects the **TFF** and **DA**:



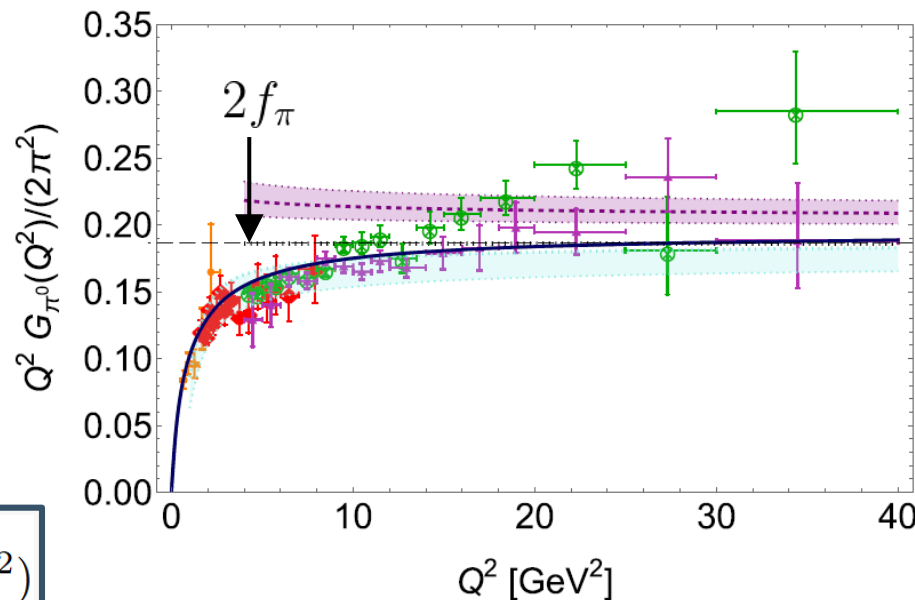
HS formula

(at sufficiently large Q^2)

$$Q^2 G_{\mathbf{P}}^q(Q^2) \stackrel{Q^2 \gg m_p^2}{\approx} 12\pi^2 f_{\mathbf{P}}^q e_q^2 w_q(Q^2)$$

- In the **opposing** end, the **chiral anomaly** entails:

$$2f_{\mathbf{P}}^0 G_{\mathbf{P}^0}^0(Q^2 = 0) = 1$$

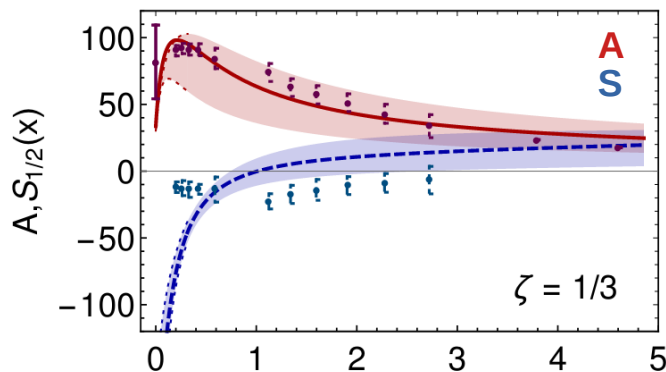
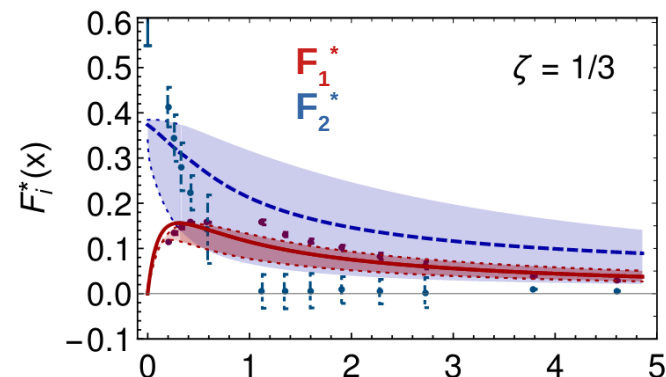


HS formula
CSM prediction
Asymptotic TFF

- Thus **EHM** (via **DCSB**) sets the infrared scale as well.
- Any deviations from this result are a measure of **EHM+HB** interplay.

Example: $N \rightarrow N(1535)$

- Transition form factors and **helicity amplitudes**:



x

$$x = Q^2/\bar{m}^2, \bar{m} = (m_+ + m_-)/2$$

Raya:2021pyr

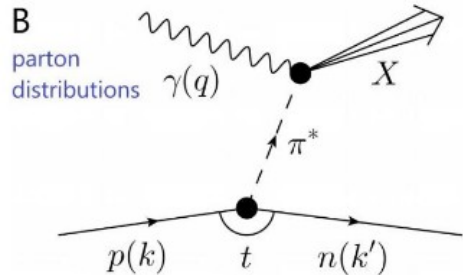
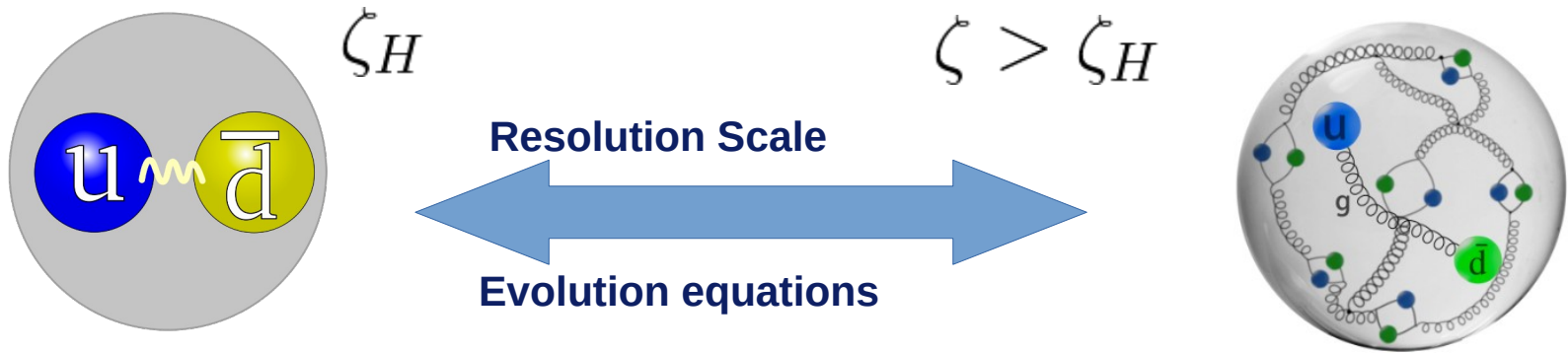
- In the absence of **DCSB**, the **nucleon** and its **parity partner** would be structurally **alike**. This is not the case...
- Their **mass** and internal **structure** is in fact quite **different**. We need to elucidate how and why.
 - Something not trivial is happening internally.

- Nucleon **TFFs** provide stringent **constraints** to the **wavefunctions**.

If one varies $g_{DB} \rightarrow g_{DB}(1 \pm 0.5)$, then $m_{N(1535)}$
 $= (1.67, 1.82)$ GeV and

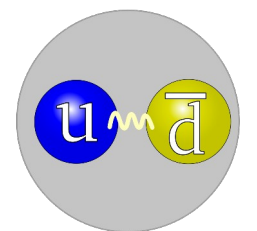
$N(1535)\frac{1}{2}^-$	s	a_1^1	a_2^1	p	v_1	v_2
$g_{DB} 1.5$	0.76	0.27	0.18	0.49	0.12	0.08
$g_{DB} 1.0$	0.66	0.20	0.14	0.68	0.11	0.09
$g_{DB} 0.5$	0.35	0.04	0.00	0.92	-0.05	0.18

Distribution functions (DFs)



- Yields e.g. information on **momentum** distribution.
- Evolution disentangles **valence**, **sea** and **gluon** contributions.

π -K DFs: **hadronic scale**



$(M_u = M_d)$ ζ_H : hadronic scale

- Fully-dressed **valence quarks** (quasiparticles)

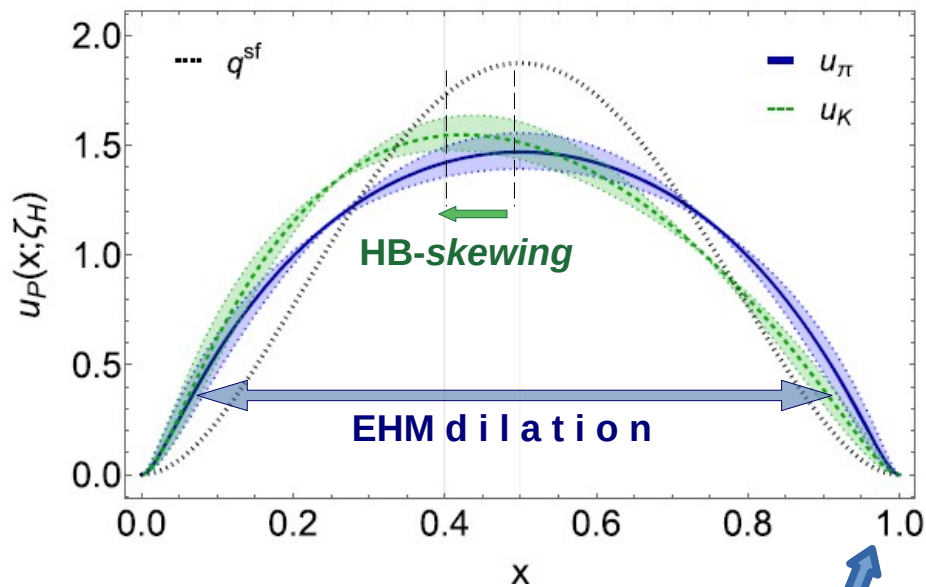
➤ At this scale, **all properties** of the hadron are contained within their valence quarks.

➔ Equally massive quarks symmetric distributions and equitable distribution of momentum fraction:

$$\langle x \rangle_\pi^u = 0.5, \quad u_\pi(x; \zeta_H) = u_\pi(1 - x; \zeta_H)$$

➤ The **kaon** distributions are only-shifted by a **few-percentage**.

$$\langle x \rangle_K^u = 0.48, \quad \langle x \rangle_{\bar{K}}^{\bar{s}} = 0.52$$



Endpoint **smoothness** is a reflection of the underlying interaction

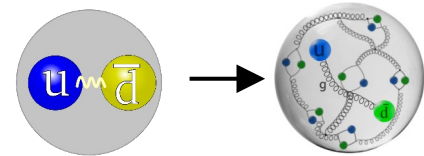
$$1/(k^2)^\beta \rightarrow (1 - x)^{2\beta}$$

Farrar:1975yb

Berger:1979du

Holt:2010vj

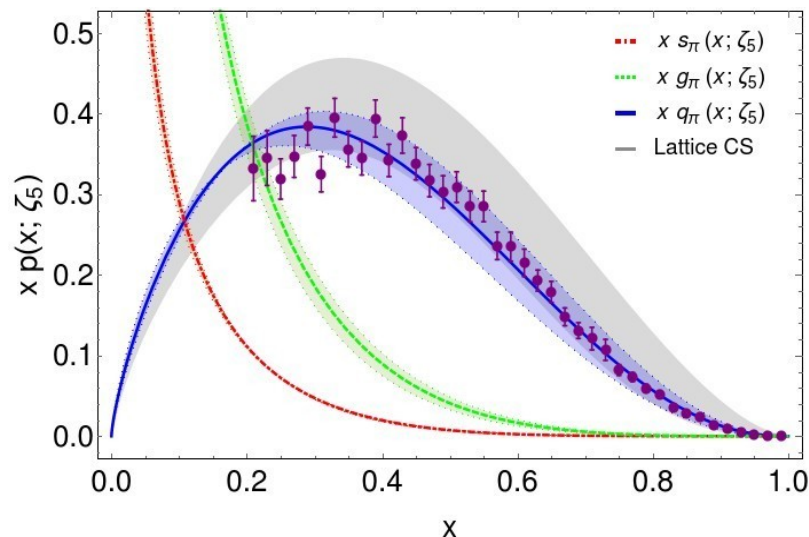
Pion DFs: Lattice & Experiment



- At **5.2 GeV**, the experimental scale, our predictions matches that from Aicher *et al.*

Aicher:2010cb

$$\langle x_{\text{gluon}} \rangle = 0.45(1), \quad \langle x_{\text{sea}} \rangle = 0.14(2)$$



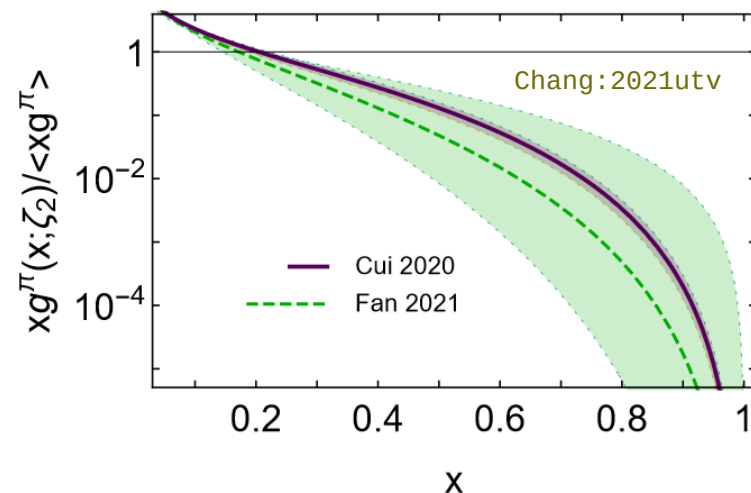
- An agreement with novel **lattice** “Cross Section” results is also obtained.

Sufian:2019bol

- At **2 GeV**, the **valence DF** shows agreement with lattice moments:

ζ_2	$\langle x \rangle_u^{\pi}$	$\langle x^2 \rangle_u^{\pi}$	$\langle x^3 \rangle_u^{\pi}$
Ref. [34]	0.24(2)	0.09(3)	0.053(15)
Ref. [35]	0.27(1)	0.13(1)	0.074(10)
Ref. [36]	0.21(1)	0.16(3)	
Herein	0.24(2)	0.098(10)	0.049(07)

- The **Gluon DF** profiles matches **lattice** expectations:

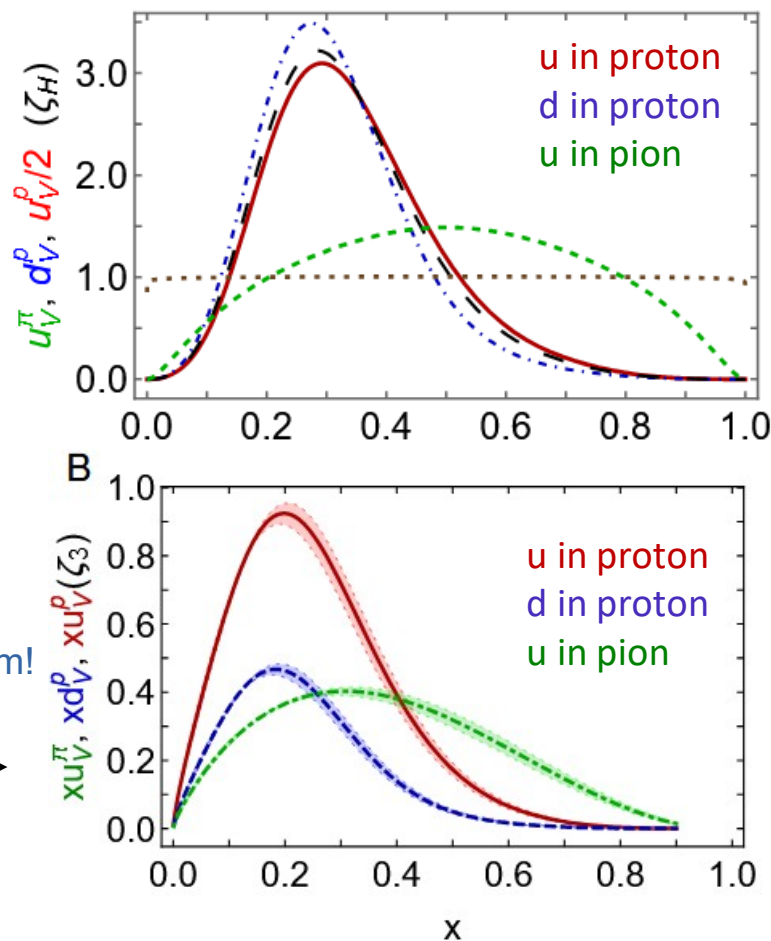


Pion vs Proton

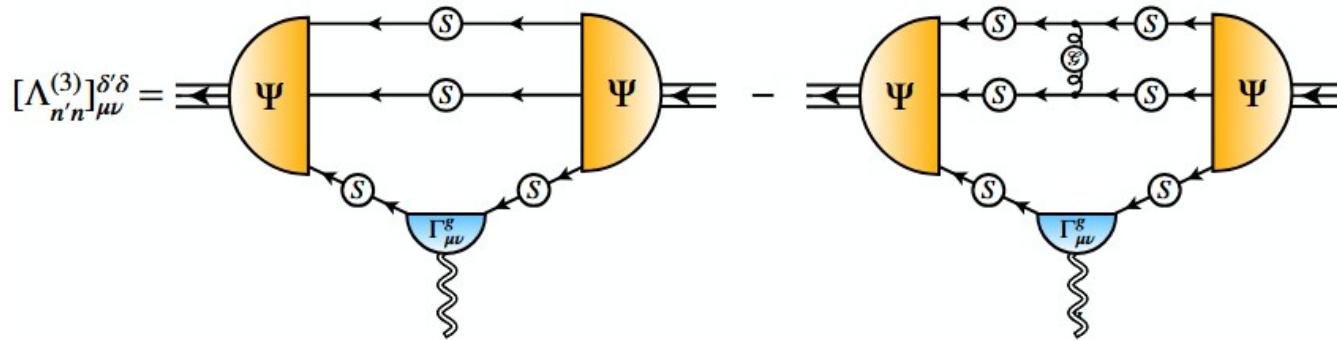
Y. Lu *et al.*

Phys.Lett.B 830 (2022) 137130

- The (nearly) massless **pion DFs** differs vastly from the massive **proton**. For instance:
 - ✓ **Counting rules** entail large- x behaviors $(1-x)^2$ and $(1-x)^3$ for the **pion** and **proton**, respectively.
 - ✓ The **momentum fractions** at ζ_H : $(M_u = M_d)$
 $\langle x \rangle_{u_p}^{\zeta_H} = 0.687$, $\langle x \rangle_{d_p}^{\zeta_H} = 0.313$, $\langle x \rangle_{u_\pi}^{\zeta_H} = 0.5$
 $\Rightarrow u_V(x) \neq 2d_V(x)$ EHM induced diquark correlations inside the proton:
 ➤ No equitable distribution of momentum!
 - ✓ Differences are **preserved** after evolution. \longrightarrow



Gravitational Form Factors (**GFFs**)



- Gives information on **mass/pressure** distribution.
- A first step towards **nucleon GFFs**

Gravitational form factors

- The expectation value of the energy-momentum tensor (**EMT**) in the **pseudoscalar meson** defines the **gravitational form factors**:

$$\Lambda_{\mu\nu}^{\mathbf{P}}(K, Q) = 2K_{\mu}K_{\nu}A^{\mathbf{P}}(Q^2) + \frac{1}{2}[Q_{\mu}Q_{\nu} - Q^2\delta_{\mu\nu}]D^{\mathbf{P}}(Q^2) + 2m_{\mathbf{P}}^2\delta_{\mu\nu}\bar{c}^{\mathbf{P}}(Q^2)$$

- Where **symmetry principles** entail:

$$A^{\mathbf{P}}(0) = 1$$

Momentum
conservation

$$D^{\mathbf{P}}(0) \stackrel{m_{\mathbf{P}}=0}{=} -1$$

Soft-pion
theorem

$$\bar{c}^{\mathbf{P}}(Q^2) = 0$$

EMT
conservation

- The deviation from **D(0)=-1**, is a manifestation of the **interplay** between **Higgs and QCD** mass generation mechanisms.

Gravitational form factors

- The expectation value of the energy-momentum tensor (**EMT**) in the **nucleon** defines the **gravitational form factors**:

$$m_N \Lambda_{\mu\nu}^N(K, Q) = -\Lambda_+(p_f) [K_\mu K_\nu A^N(Q^2) + i K_{\{\mu} \sigma_{\nu\}\rho} Q_\rho J^P(Q^2) + \frac{1}{4} (Q_\mu Q_\nu - Q^2 \delta_{\mu\nu}) D^N] \Lambda_+(p_i)$$

- Where **symmetry principles** entail:

$$A^P(0) = 1$$

Momentum
conservation

$$J^N(0) = 1/2$$

Spin sum-rule

$$\bar{c}^P(Q^2) = 0$$

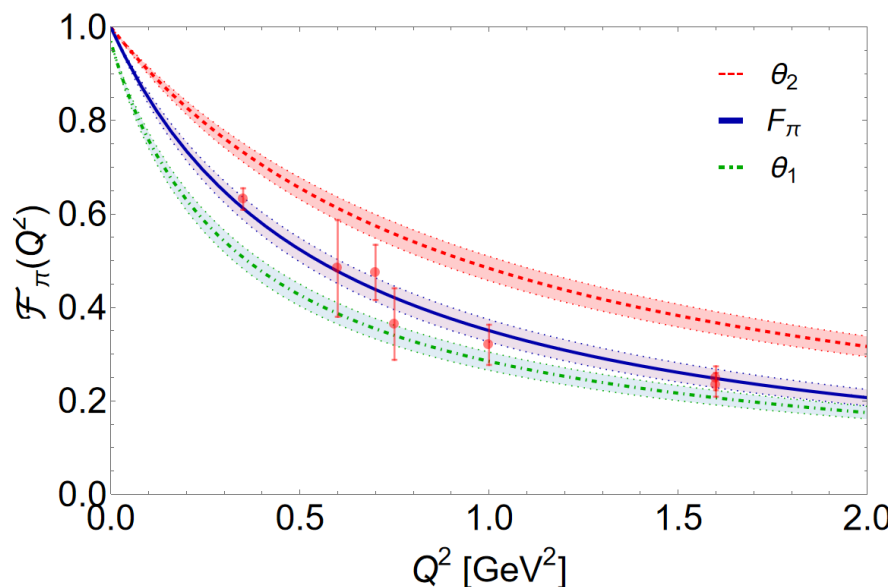
EMT
conservation

- There is **no constraint on $D(0)$** , the so called **D-term**, which is often referred to as **“The Last Unknown Global Property of the Nucleon”**.

Pion GFFs

$$r_{\theta_1}/r_F \approx r_F/r_{\theta_2} \approx 0.75$$

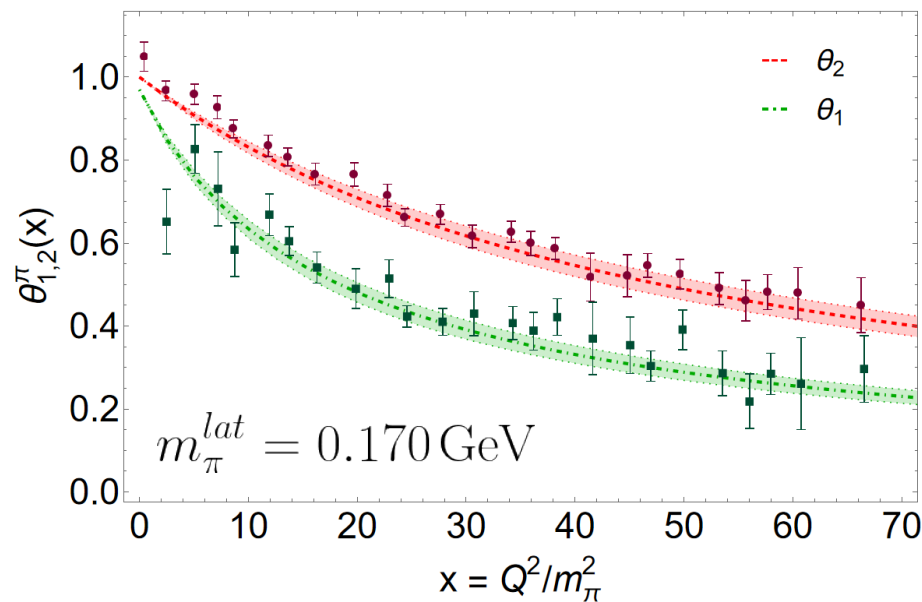
- For π (**K**) it is found: $r_{\theta_1}^{\text{P}} = 0.81 \text{ fm} (0.63 \text{ fm}) > r_F^{\text{P}} = 0.64 \text{ fm} (0.58 \text{ fm}) > r_{\theta_2}^{\text{P}} = 0.47 \text{ fm} (0.40 \text{ fm})$.
(mechanical)
(charge)
(mass)



- The **Kaon**, albeit more **compressed**, exhibit similar patterns.
 $\bar{r}_K/\bar{r}_\pi = 0.85(6)$

- Agreement with **Lattice QCD** is also obtained:

Hackett:2023nkr

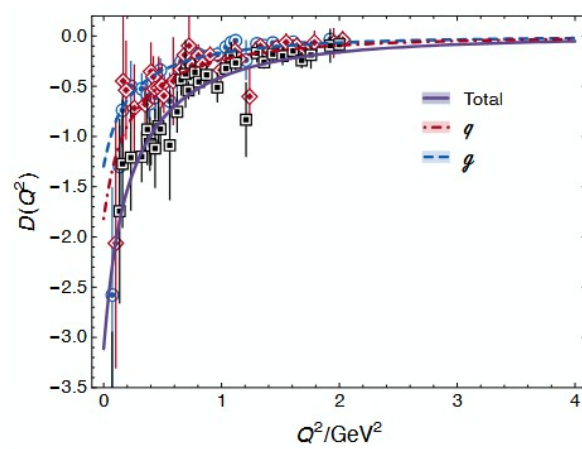
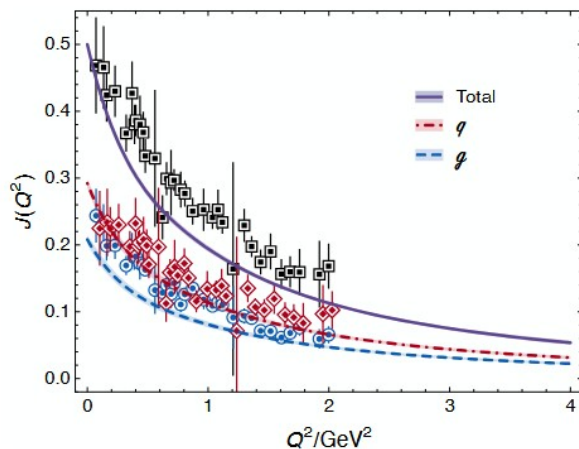
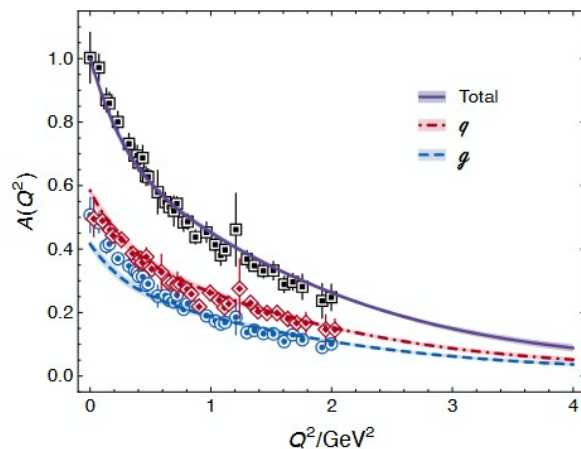


- Our **prediction** is compatible with more recent **lattice** expectations:

Hackett:2023rif

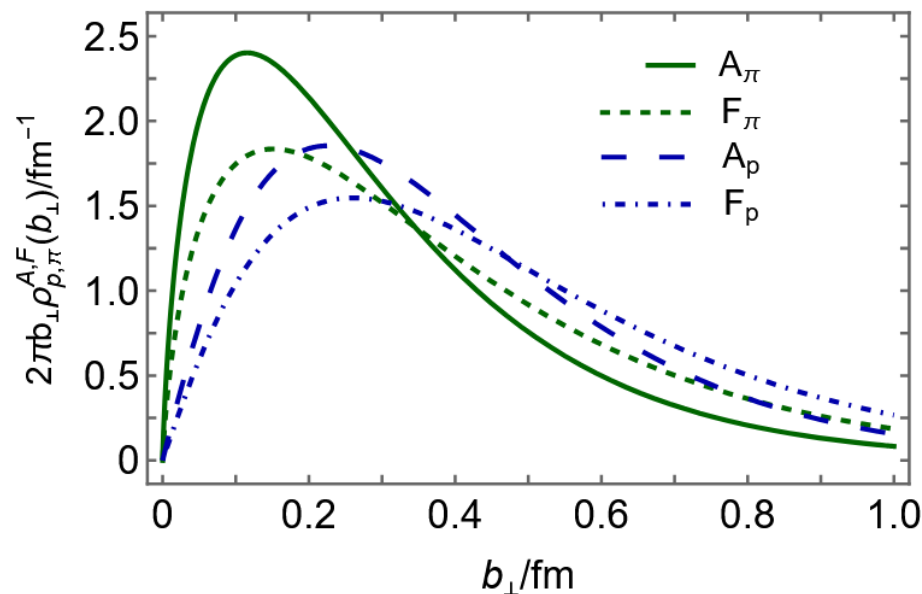
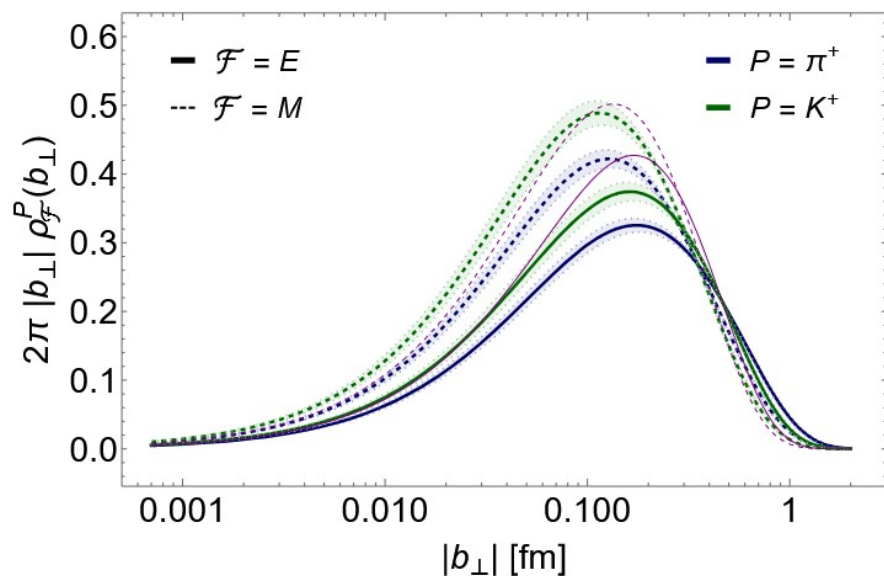
herein	$A(0)$	$J(0)$	$-D(0)$
Total	1.00	0.50	$3.114(10)_{\pm}$
q	$0.584(13)_{\pm}$	$0.292(06)_{\pm}$	$1.820(43)_{\pm}$
g	$0.416(13)_{\mp}$	$0.208(06)_{\mp}$	$1.294(33)_{\mp}$
IQCD	$A(0)$	$J(0)$	$-D(0)$
Total	1.011(37)	0.506(25)	3.87(97)
q	0.510(25)	0.251(21)	1.30(49)
g	0.501(27)	0.255(13)	2.57(84)

- The **symmetry-preserving** treatment guarantees the **mass and spin** sum rules
- At **2 GeV**, our prediction features slightly smaller gluon contribution
- As with the **π -K**, **charge** effects span over a larger domain that **mass** effects: $r_{\text{mass}} = 0.81(5)r_{\text{ch}}$



Spatial Distributions

- For π (**K**) it is found: $r_{\theta_1}^P = 0.81 \text{ fm} (0.63 \text{ fm}) > r_F^P = 0.64 \text{ fm} (0.58 \text{ fm}) > r_{\theta_2}^P = 0.47 \text{ fm} (0.40 \text{ fm})$.
(mechanical)
(charge)
(mass)



- The **Kaon**, albeit more **compressed**, exhibit similar patterns.
 $\bar{r}_K / \bar{r}_\pi = 0.85(6)$

- The **proton** is spatially more **extended**.
 $\bar{r}_p / \bar{r}_\pi \approx 1.25$

π -K Pressures

$$p_K^u(r) = \frac{1}{6\pi^2 r} \int_0^\infty d\Delta \frac{\Delta}{2E(\Delta)} \sin(\Delta r) [\Delta^2 \theta_1^{K_u}(\Delta^2)],$$

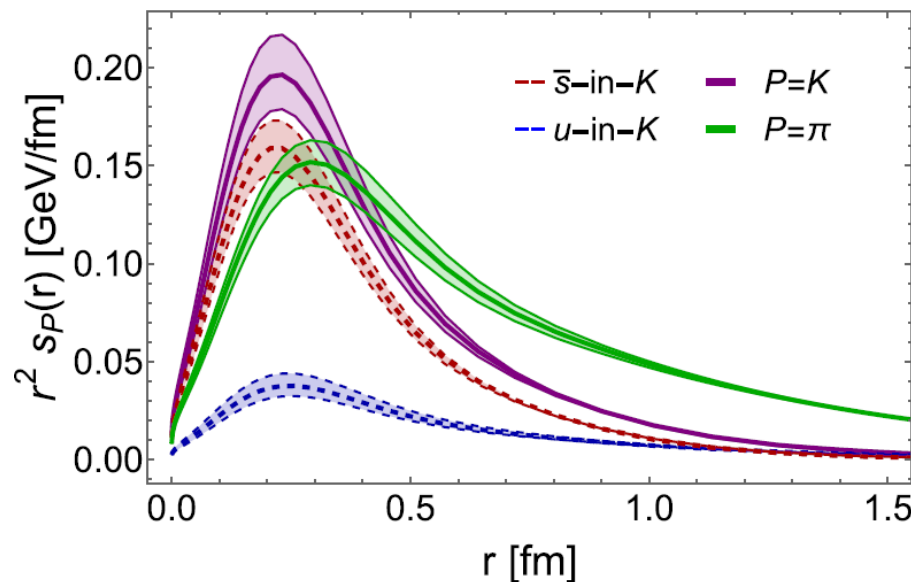
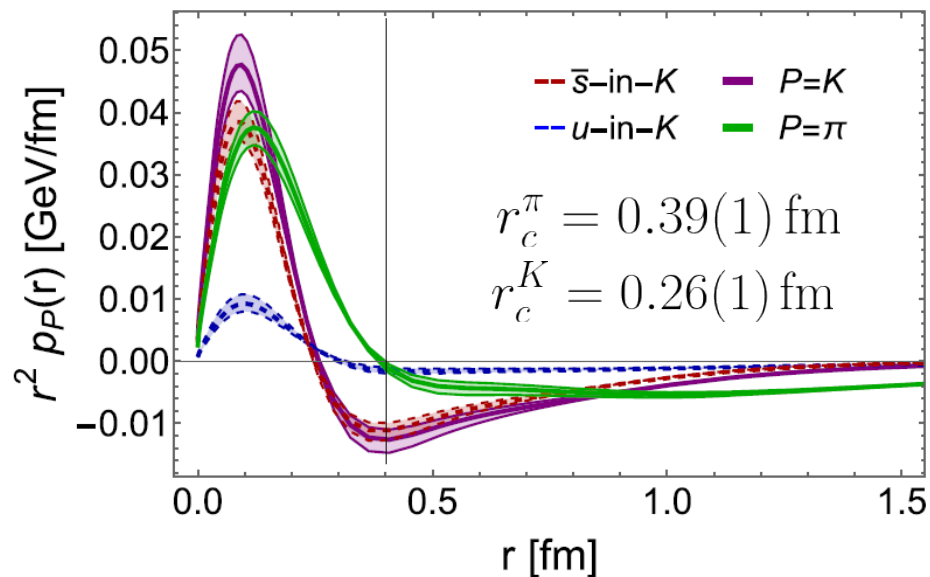
$$s_K^u(r) = \frac{3}{8\pi^2} \int_0^\infty d\Delta \frac{\Delta^2}{2E(\Delta)} j_2(\Delta r) [\Delta^2 \theta_1^{K_u}(\Delta^2)],$$

“Pressure” Quark attraction/repulsion

CONFINEMENT

“Shear”

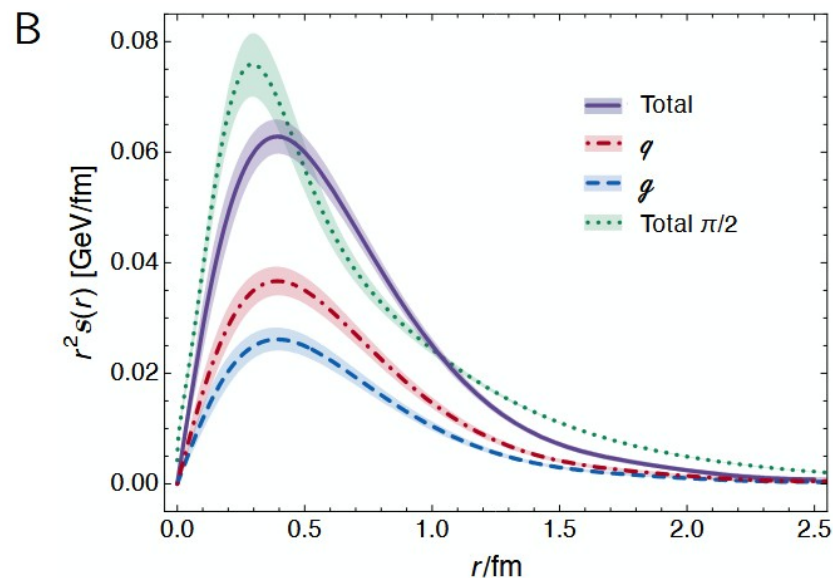
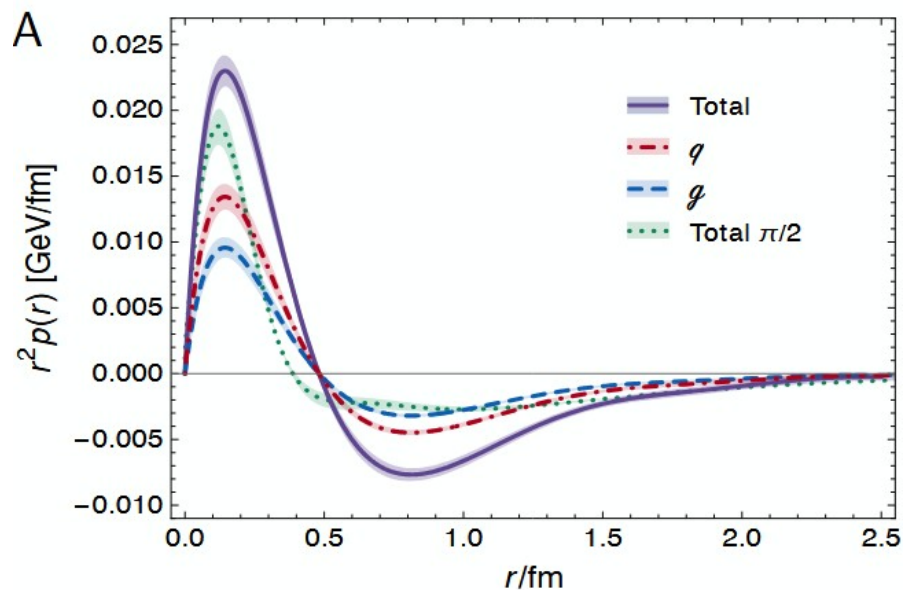
Deformation QCD forces



Nucleon Pressures

Z-Q Yao *et al.*,
Eur.Phys.J.A 61 (2025) 5, 92

- Nucleon pressures follow the same patterns as in the **pion**.



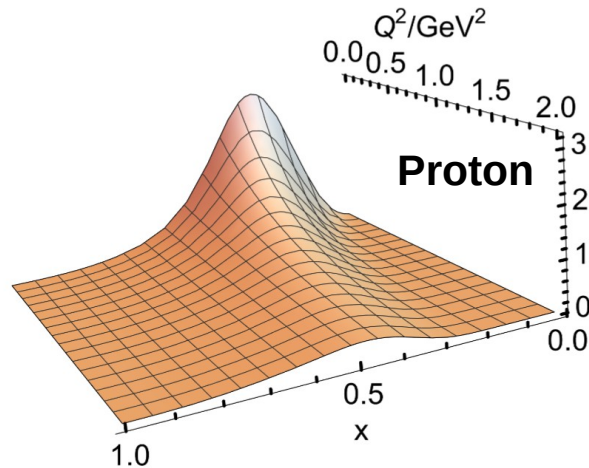
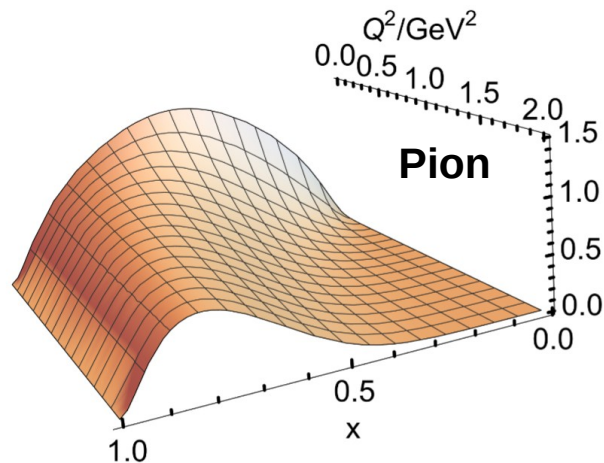
- Although only about **half** as large as those in the **pion**, the **nucleon pressures** are comparable to those found in the core of **neutron stars**.

$$\sim 10^{35} \text{ Pa}$$

Pion and Proton GPDs

Pion and Proton GPDs

- Pion and Proton GPDs, at the hadronic scale, inherit the **dilation/compression** patterns of the DF.



- In the **isospin symmetric** limit, the **pion GPD** is **symmetric** in the forward limit.

→ **Not** the case for the **proton**...

$$\langle x \rangle_{u_p}^{\zeta_H} = 0.687, \quad \langle x \rangle_{d_p}^{\zeta_H} = 0.313, \quad \langle x \rangle_{u_\pi}^{\zeta_H} = 0.5$$

- GPDs encode many **aspects** of the hadron internal **structure**.

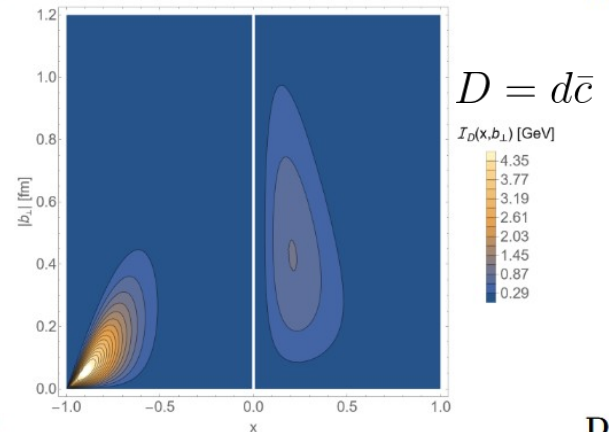
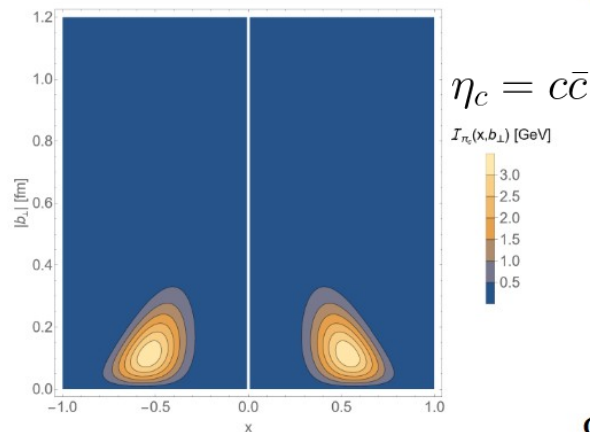
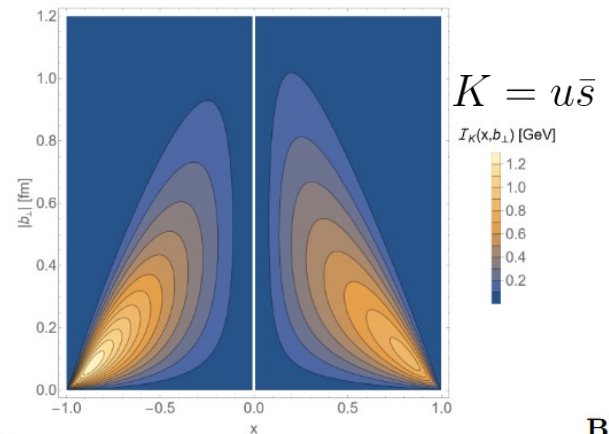
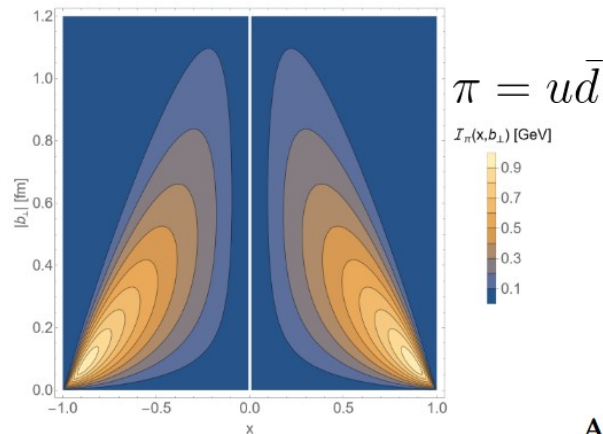
$$q_P(x) = H_P^q(x, 0, 0) \quad F_P^q = \int_{-1}^1 dx H_P^q(x, \xi, \Delta^2)$$

$$\theta_2^{Pq}(\Delta^2) + \xi^2 \theta_1^{Pq}(\Delta^2) = \int_{-1}^1 dx x H_P^q(x, \xi, \Delta^2)$$

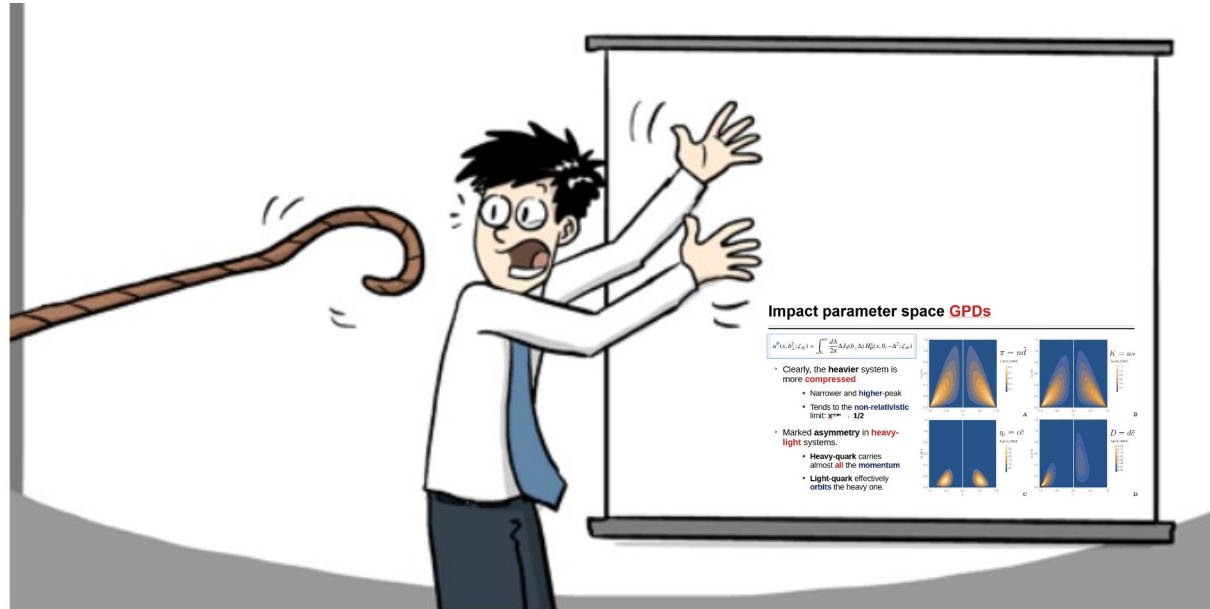
Impact parameter space **GPDs**

$$u^P(x, b_\perp^2; \zeta_{\mathcal{H}}) = \int_0^\infty \frac{d\Delta}{2\pi} \Delta J_0(b_\perp \Delta) H_P^u(x, 0, -\Delta^2; \zeta_{\mathcal{H}})$$

- Clearly, the **heavier** system is more **compressed**
 - Narrower and **higher**-peak
 - Tends to the **non-relativistic** limit: $x^{\max} \rightarrow 1/2$
- Marked **asymmetry** in **heavy-light** systems.
 - Heavy-quark** carries almost **all** the **momentum**
 - Light-quark** effectively **orbits** the heavy one.



Final Highlights



Final Highlights

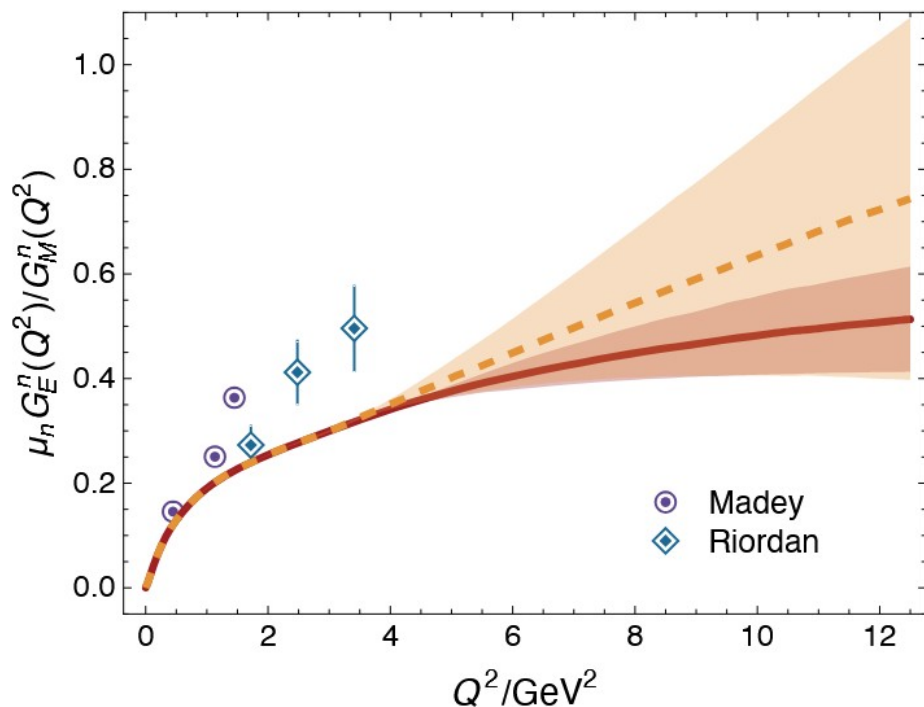
- The **emergent** phenomena in **QCD** produces unique outcomes:
 - **Confinement**, dynamical **mass** generation, and a peculiar effective **coupling**.
 - These orchestrate the formation of **hadrons** and their properties, and are responsible for almost all of the **mass** of the **VM**.
- **Nucleons** and **pseudoscalar** mesons take **center** stage in elucidating these aspects
 - **Other** hadrons provide a **complementary** picture.
- **Excited and Exotic** states push the limits of our understanding of how the **strong force** binds matter together.
- **Experimental** facilities around the globe are set to examine these aspects at an **unprecedented** level.
 - Convenient.. as **theory** has evolved to the point where all sorts of quantities are **within reach**.



Elastic Form Factors

Yao:2024uej

- For the **proton**, different theoretical and phenomenological analysis suggest a **zero-crossing** in the **GE/GM** ratio



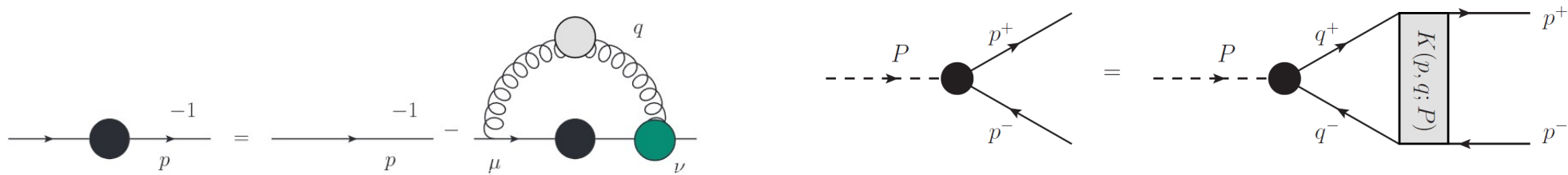
- This reflection on a **destructive interference** between the individual **valence-quark** contributions.
- In fact connected with non-trivial effects coming from orbital **angular momentum**.
- The **location** of the zero is highly sensitive to the **diquark** content of the **nucleon**.

- Modern analyses favor a **non-negligible** component of **axial-vector** diquark:

~ 25 – 35 %

the rest comes from the expected **scalar** diquark.

Continuum Schwinger Methods (CSM)



Dyson-Schwinger Equations

- Equations of motion of a **quantum field theory**
- Relate Green functions with higher-order Green functions
 - ➔ • **Infinite** tower of coupled equations.
 - × Systematic **truncation** required
- ✓ **No assumptions** on the **coupling** for their derivation.
 - ➔ ✓ Capture both **perturbative** and **non-perturbative** facets of **QCD**
- ✓ **Not limited** to a certain domain of current **quark masses**
- ✓ Maintain a **traceable connection** to QCD.

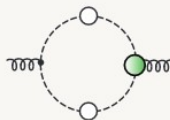
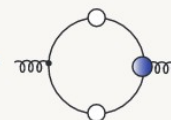
Example DSEs

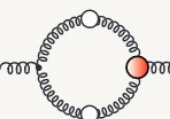

Quark propagator:

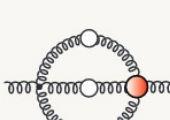

$$\text{---}\bigcirc\text{---}^{-1} = \text{---}^{-1} + \text{---}\bigcirc\text{---}$$

Gluon propagator:

$$\text{~~~~~}^{-1} = \text{~~~~~}^{-1} +$$

+  + 

+  + 

+  + 

Baryons: Faddeev equation

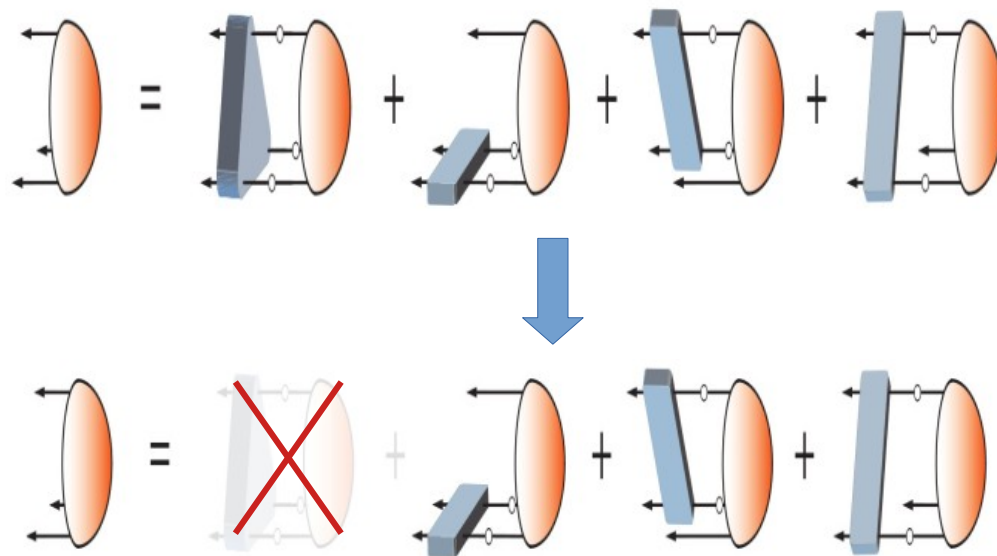
- Strong evidence anticipates the formation of **dynamical** quark-quark correlations (**diquarks**) within **baryons**, for instance:
 - The **primary three-body** force **binding** the quarks within the baryon vanishes when projected onto the color singlet channel.

Eichmann:2016yit

i.e. a 3-gluon vertex attached to each quark once (and only once)

- The dominant 3-gluon contribution is the one attaching twice to a quark
- This produces a strengthening of quark-quark interactions

Barabanov:2020jvn

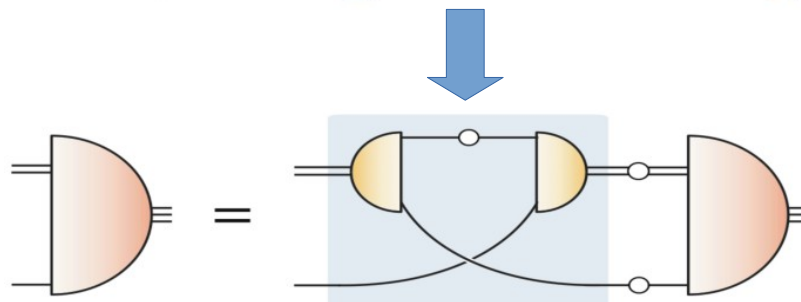
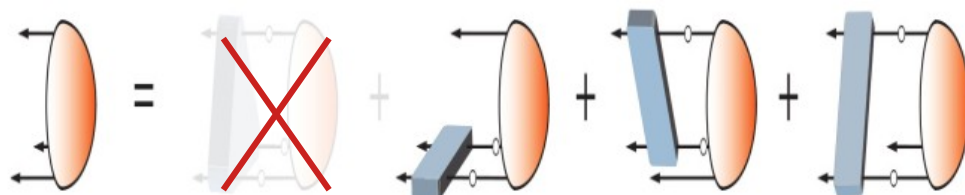


Baryons: Faddeev equation

- Strong evidence anticipates the formation of **dynamical** quark-quark correlations (**diquarks**) within **baryons**, for instance:

Barabanov:2020jvn

- The **primary three-body** force **binding** the quarks within the baryon vanishes when projected onto the color singlet channel.
- The **attractive** nature of **quark-antiquark** correlations in a color-singlet meson, is also **attractive** for $\bar{3}_c$ **quark-quark** correlations within a color singlet baryon.



Dynamical Quark-diquark picture

Non-pointlike **diquarks**:

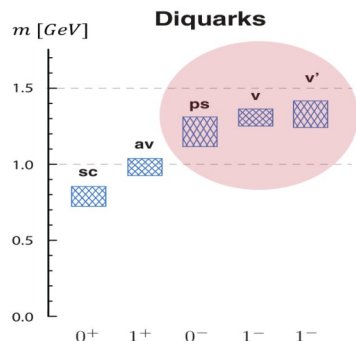
- Color anti-triplet
- Fully interacting
- Origins related to **EHM** phenomena

Baryons: Quark-diquark picture

Barabanov:2020jvn

→ The **attractive** nature of **quark-antiquark** correlations in a color-singlet meson, is also **attractive** for $\bar{3}_c$ **quark-quark** correlations within a color singlet baryon.

- Due to charge conjugation properties, a **J^P diquark** partners with an analogous **J^P meson**.
- We can thus establish a connection between the **meson** and **diquark** Bethe-Salpeter equations:



$$\Gamma_{q\bar{q}}(p; P) = - \int \frac{d^4 q}{(2\pi)^4} g^2 D_{\mu\nu}(p-q) \frac{\lambda^a}{2} \gamma_\mu S(q+P) \Gamma_{q\bar{q}}(q; P) S(q) \frac{\lambda^a}{2} \gamma_\nu$$

$$\Gamma_{qq}(p; P) C^\dagger = - \frac{1}{2} \int \frac{d^4 q}{(2\pi)^4} g^2 D_{\mu\nu}(p-q) \frac{\lambda^a}{2} \gamma_\mu S(q+P) \Gamma_{qq}(q; P) C^\dagger S(q) \frac{\lambda^a}{2} \gamma_\nu$$

Less tightly 'bound'

Non-pointlike diquarks:

- Color anti-triplet
- Fully interacting
- Origins related to **EHM** phenomena

- Computed '**masses**' should be interpreted as correlation **lengths**:

$$m_{[ud]_{0+}} = 0.7 - 0.8 \text{ GeV}, \quad m_{\{uu\}_{1+}} = 0.9 - 1.1 \text{ GeV}$$

- Stressing the fact that the **diquarks** have a **finite** size:

$$r_{[ud]_{0+}} \gtrsim r_\pi, \quad r_{\{uu\}_{1+}} \gtrsim r_\rho$$