

XIX Mexican Workshop on Particles and Fields

October 21, 2025

Sub-GeV leptophobic B-boson exchange in $\Phi \rightarrow \eta \pi^0 \gamma$ decays

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Barcelona Institute of
Science and Technology



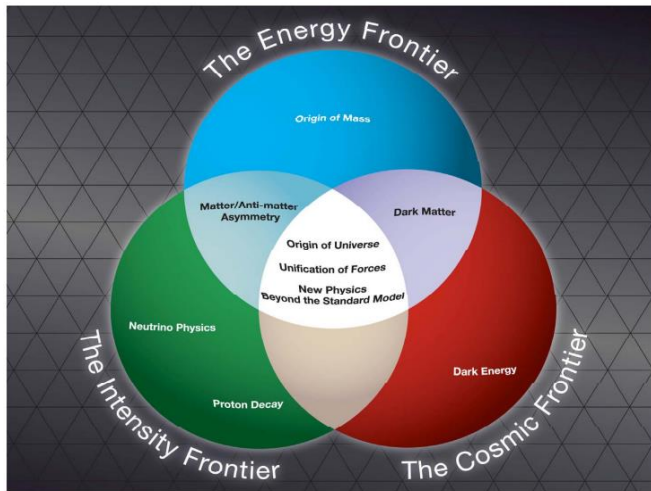
Generalitat de Catalunya
Departament de Recerca
i Universitats

Exploring Weakly Coupled New Physics

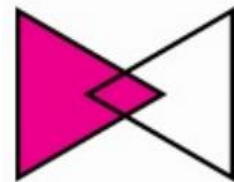
- **Energy Frontier:** Direct searches

- **Intensity Frontier:** High-luminosity, rare process detection

2008 Frontiers of Particle Physics



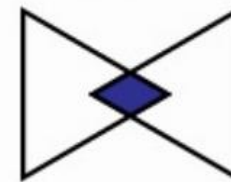
2023 Science Drivers



Decipher
the
Quantum
Realm

Elucidate the Mysteries
of Neutrinos

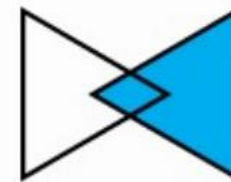
Reveal the Secrets of
the Higgs Boson



Explore
New
Paradigms
in Physics

Search for Direct Evidence
of New Particles

Pursue Quantum Imprints
of New Phenomena



Illuminate
the
Invisible
Universe

Determine the Nature
of Dark Matter

Understand What Drives
Cosmic Evolution

Report of the P5 Panel

- **Key Research Themes:**

- **Hidden Sector Particles:** axions, axion-like, dark photons, heavy neutral leptons
- **Dark Matter Candidates:** WIMPs, dark sector particles

Searches of a leptophobic B boson

- What if a **new force couples** predominantly to **quarks** over leptons?

- **Simplest model:** A new gauge boson **B** that couples to the baryon number

$$\mathcal{L}_{\text{int}} = \left(\frac{1}{3} g_B + \varepsilon Q_q e \right) \bar{q} \gamma^\mu q B_\mu - \varepsilon e \bar{\ell} \gamma^\mu \ell B_\mu \quad \Rightarrow \quad \varepsilon \sim e g_B / (4\pi)^2 \quad \alpha_B = \frac{g_B^2}{4\pi}$$

- It **preserves** the low-energy symmetries of QCD: **C, P, SU(3)_F**
- **B** is a **singlet** under isospin
- $I^G(J^{PC}) = 0^-(1^{--}) \Rightarrow$ **B is ω meson like!** $m_\pi \lesssim m_B \lesssim \text{GeV}$

	Dark-photon like	m_B [MeV]	Novel signatures	
Decay $\rightarrow$	$B \rightarrow e^+ e^-$	$B \rightarrow \pi^0 \gamma$	$B \rightarrow \pi^+ \pi^- \pi^0$	$B \rightarrow \eta \gamma$
Production $\downarrow$	$m_B \sim 1 - 140 \text{ MeV}$	140 – 620 MeV	620 – 1000 MeV	
$\pi^0 \rightarrow B \gamma$	$\pi^0 \rightarrow e^+ e^- \gamma$	...	...	...
$\eta \rightarrow B \gamma$	$\eta \rightarrow e^+ e^- \gamma$	$\eta \rightarrow \pi^0 \gamma \gamma$	...	...
$\eta' \rightarrow B \gamma$	$\eta' \rightarrow e^+ e^- \gamma$	$\eta' \rightarrow \pi^0 \gamma \gamma$	$\eta' \rightarrow \pi^+ \pi^- \pi^0 \gamma$	$\eta' \rightarrow \eta \gamma \gamma$
$\omega \rightarrow \eta B$	$\omega \rightarrow \eta e^+ e^-$	$\omega \rightarrow \eta \pi^0 \gamma$	...	...
$\phi \rightarrow \eta B$	$\phi \rightarrow \eta e^+ e^-$	$\phi \rightarrow \eta \pi^0 \gamma$	...	

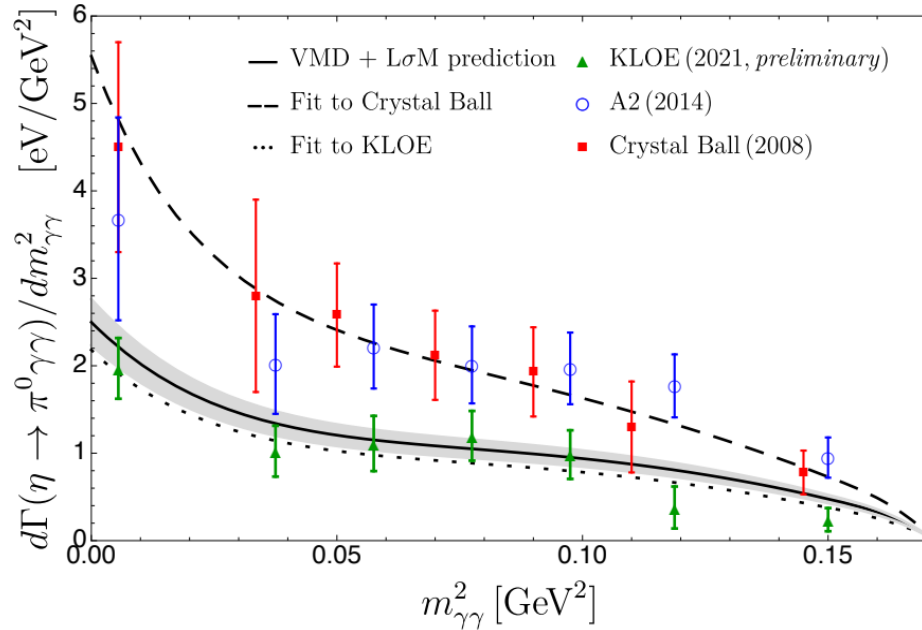
□ Escribano, González Solís, Royo.
Phys.Rev.D 106, 114007 (2022)

■ Escribano, Miranda. To appear.

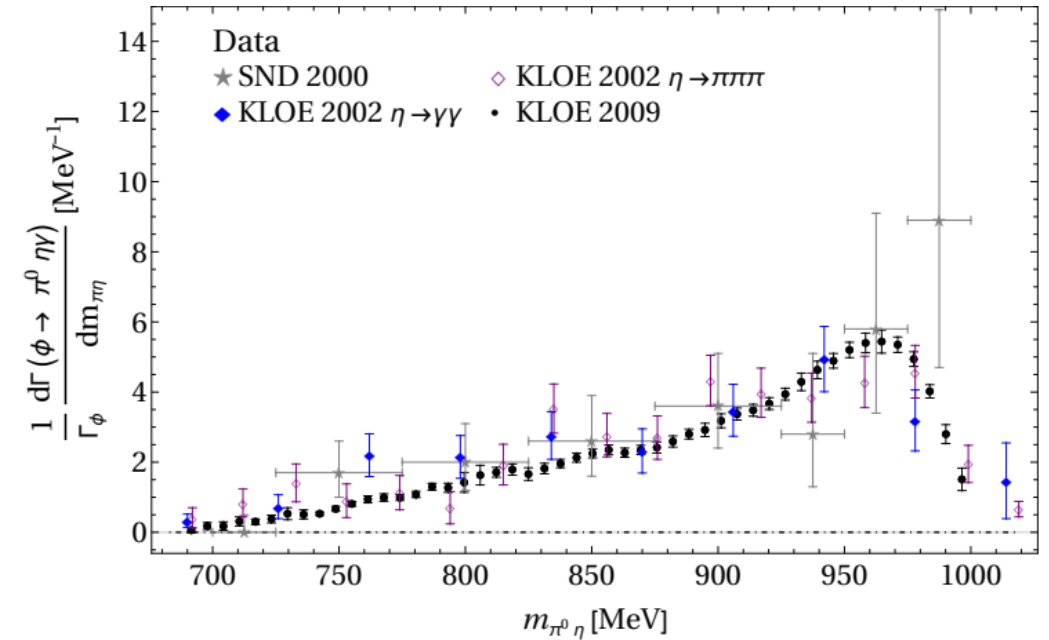
S. Tulin. Phys.Rev.D 89 (2014) 11, 114008

Motivation

- Tensions between measurements of the $\gamma\gamma$ spectrum for the $\eta \rightarrow \pi^0\gamma\gamma$ decay.
- Higher statistics for the $\eta\pi^0$ invariant mass distribution for the $\phi \rightarrow \pi^0\eta\gamma$ decays from KLOE.
- Better agreement between different datasets from SND00, KLOE02 and KLOE09.



R. Escribano et al. Phys.Rev.D 106 (2022) 11, 114007



Previous estimates

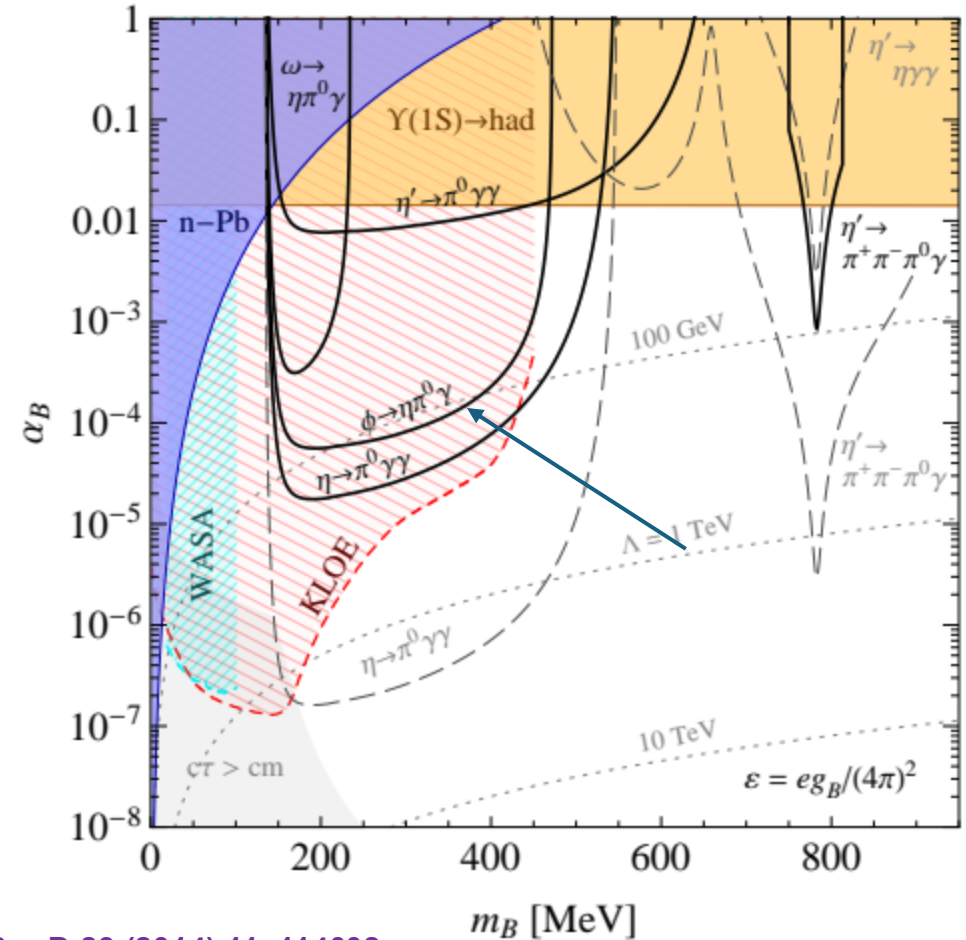
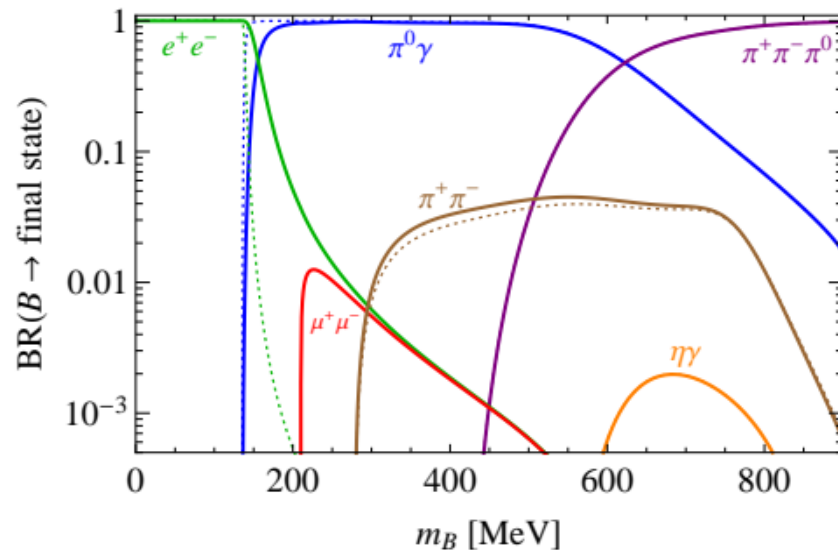
- Assuming the **Narrow Width Approximation (NWA)**:

$$\text{Br}(\phi \rightarrow \pi^0 \eta \gamma) = \text{Br}(\phi \rightarrow B \eta) \times \text{Br}(B \rightarrow \pi^0 \gamma)$$

$$\text{Br}(B \rightarrow \pi^0 \gamma) = 1$$

- **QCD contribution off**

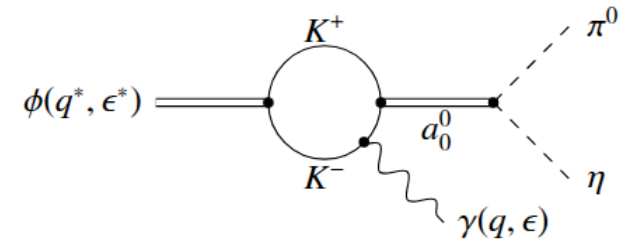
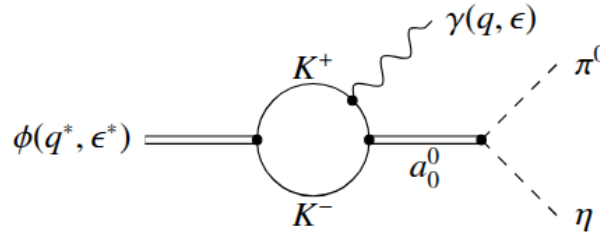
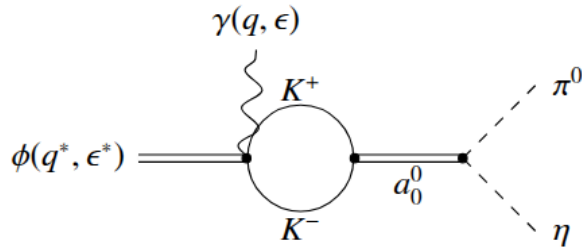
$$\text{Br}(\phi \rightarrow \pi^0 \eta \gamma) < \text{Br}_{\text{exp}} \text{ at } 2\sigma$$



S. Tulin. Phys.Rev.D 89 (2014) 11, 114008

Scalar meson exchange contributions

- In the **LσM** framework, the **$\phi \rightarrow \pi^0 \eta \gamma$ decay** proceed through the **exchange** of intermediate **scalar mesons**



$$\mathcal{A}_{\phi \rightarrow \pi^0 \eta \gamma}^{L\sigma M} = \frac{eg}{2\pi^2 m_{K^+}^2} \{a\} L(m_{\pi^0 \eta}^2) \mathcal{A}_{K^+ K^- \rightarrow \pi^0 \eta}^{L\sigma M}$$

$$\mathcal{A}_{K^+ K^- \rightarrow \pi^0 \eta}^{L\sigma M} = \frac{s - m_\eta^2}{D_{a_0}(s)} \frac{g_{a_0 K^+ K^-} g_{a_0 \pi^0 \eta}}{m_\eta^2 - m_{a_0}^2} + \frac{m_\eta^2 + m_\pi^2 - s}{4f_\pi f_K} (c\varphi_P - \sqrt{2}s\varphi_P)$$

$$g_{a_0 \pi^0 \eta} = \frac{m_\eta^2 - m_{a_0}^2}{f_\pi} c\varphi_P,$$

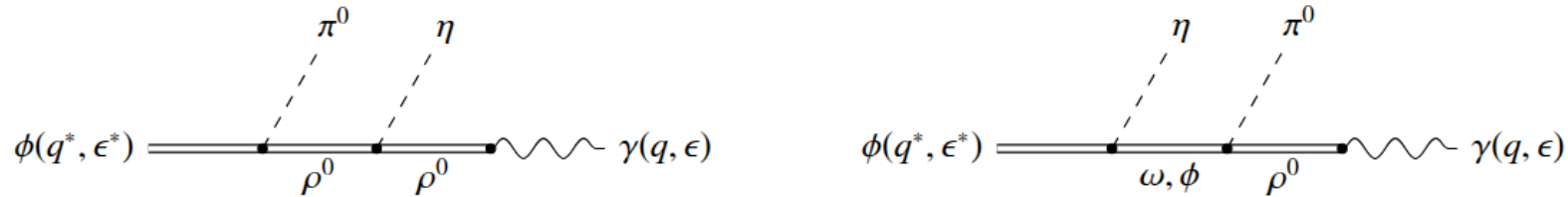
$$g_{a_0 K^+ K^-} = -g_{a_0 K^0 \bar{K}^0} = \frac{m_K^2 - m_{a_0}^2}{2f_K}$$

A. Bramon et al. Phys.Lett.B 494 (2000) 221-228

$$D(s) = s - m_R^2 + \text{Re}\Pi(s) - \text{Re}\Pi(m_R^2) + i\text{Im}\Pi(s)$$

Vector meson exchange contributions

- In the **VMD** framework, the $\phi \rightarrow \pi^0 \eta \gamma$ decay proceed through the **exchange** of intermediate **vector mesons**



$$\mathcal{A}_{\phi \rightarrow \pi^0 \eta \gamma}^{\text{VMD}} = \frac{\epsilon e G^2}{\sqrt{2} g} \left\{ c \varphi_P \left(\frac{P'^2 \{a\} + \{b(P')\}}{D_\rho(P'^2)} + \frac{P^2 \{a\} + \{b(P)\}}{D_\omega(P^2)} \right) + \sqrt{2} s \varphi_P \left(\frac{P^2 \{a\} + \{b(P)\}}{D_\omega(P^2)} - \frac{P^2 \{a\} + \{b(P)\}}{D_\phi(P^2)} \right) \right\}$$

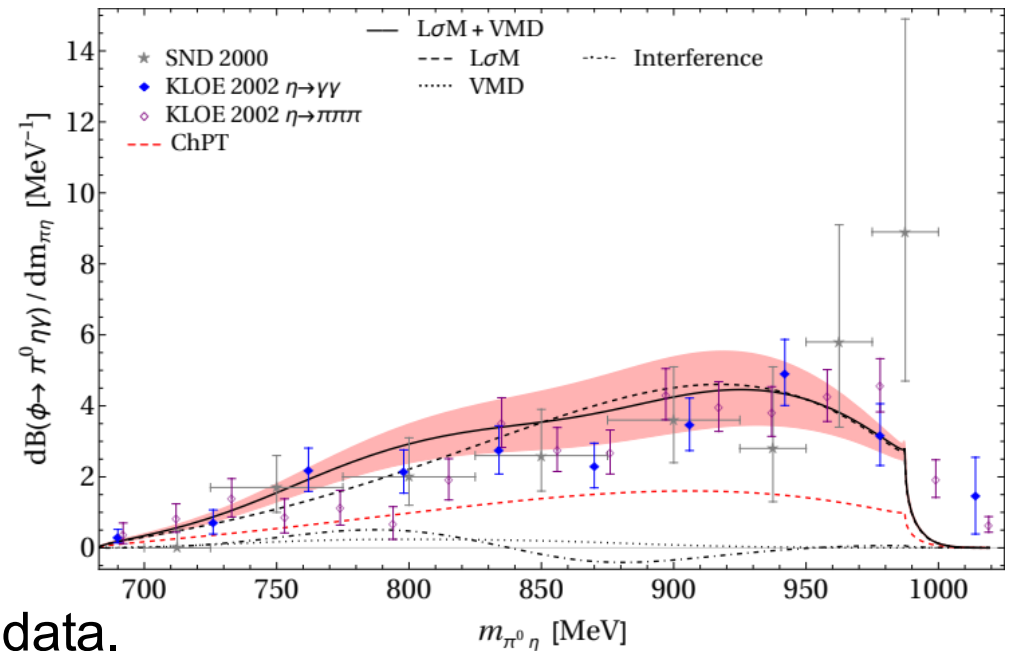
A. Bramon et al. Phys.Lett.B 283 (1992) 416-420

$$\{a\} = (\epsilon^* \epsilon)(q^* q) - (\epsilon^* q)(\epsilon q^*)$$

$$\begin{aligned} \{b(P)\} = & (\epsilon^* q)(\epsilon P)(q^* P) + (\epsilon q^*)(\epsilon^* P)(q P) \\ & - (\epsilon^* \epsilon)(q^* P)(q P) - (\epsilon^* P)(\epsilon P)(q^* q) \end{aligned}$$

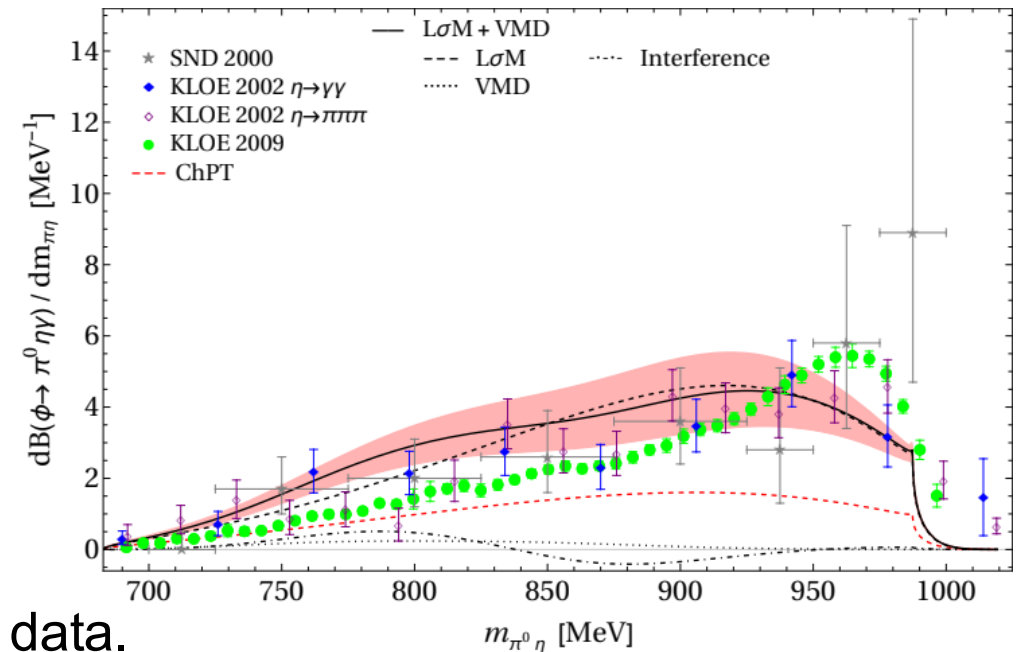
$\phi \rightarrow \pi^0 \eta \gamma$ predictions

- In **LσM**, the theoretical prediction is **$BR=8.78(1.76) \times 10^{-5}$**
 - **Scalar** contribution dominates:
 - a_0 : **95%**
 - **VMD** contribution:
 - ρ , ω , and ϕ : **1%**
 - **Interference**:
 - S-V: **4%**
- **Compatible** with the present available **data**:
 - **$BR=7.27(30) \times 10^{-5}$** **PDG 2024**
- Comparison with **experimental data**:
 - **Good agreement** with **SND'00** and **KLOE'02** data.



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 - **$BR=7.27(30) \times 10^{-5}$** **PDG 2024**
- Comparison with **experimental data**:
 - **Good agreement** with **SND'00** and **KLOE'02** data.
 - **Fails** to describe the **KLOE'09** data.

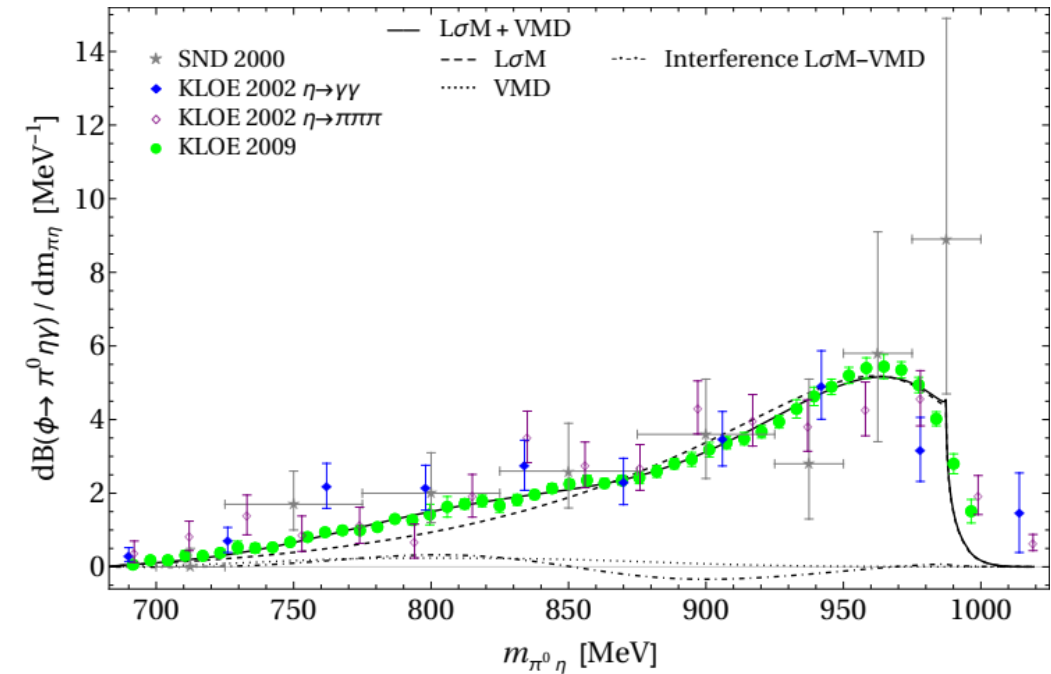


$\phi \rightarrow \pi^0 \eta \gamma$ predictions

- **Advantages** of our model:
 - Incorporation of **scalar resonances** in an explicit way.
 - Includes the **constraints** required by **chiral symmetry**.
- We relax the **theoretical framework**:
 - **Three parameters:** m_{a_0} , $g_{a_0 \pi^0 \eta}$ and $g_{a_0 K^+ K^-}$
- We perform a **fit** to the **experimental data**:

	Fit 1	Fit 2
$\chi^2/d.o.f$	69.3/46 = 1.5	68.0/45 = 1.5
m_{a_0}	$986.6^{+0.7}_{-1.2}$	$986.6^{+0.7}_{-1.1}$
$g_{a_0 K \bar{K}}$	$\pm(2838.2^{+80.7}_{-77.8})$	$\pm(2844.6^{+81.9}_{-79.0})$
$g_{a_0 \pi \eta}$	$\pm(3692.0^{+56.5}_{-55.1})$	$\pm(3691.3^{+57.1}_{-55.7})$
δ_{+0}	$0^\dagger$	0.12 ± 0.10

Vector-Scalar interference



This approach improves the agreement between theory and experiment.

SM contribution to the branching ratio

Contribution	$L\sigma M$	Fit 1	Fit 2
Scalar	$8.29^{+1.78}_{-1.74}$	$6.52^{(24)}_{(20)}{}^{(23)}_{(22)}(1)$	$6.54^{(22)}_{(20)}(23)(1)$
Vector		$0.36(0)$	
Interference	$0.13(2)$	$0.02(0)(0)(0)$	$0.01(0)(0)(0)(1)$
Total	$8.78^{+1.78}_{-1.74}$	$6.90^{(33)}_{(30)}$	$6.91^{(32)}_{(31)}$
PDG [15]		7.27 ± 0.30	

Good agreement between theory and experiment.

Contributions to $\mathcal{B}(\phi \rightarrow \pi^0 \eta \gamma)$ in units of 10^{-5} .

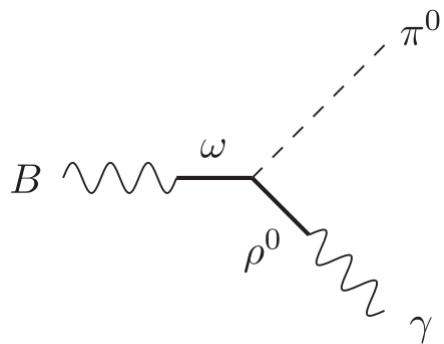
Calculation of hadronic processes

- Using the **hidden local symmetry** (HLS) for **VMD**

$$\mathcal{L}_{VVP} = \frac{G}{\sqrt{2}} \varepsilon^{\mu\nu\alpha\beta} \text{tr} [\partial_\mu V_\nu \partial_\alpha V_\beta P] \quad G = \frac{3g^2}{4\pi^2 f_\pi}$$

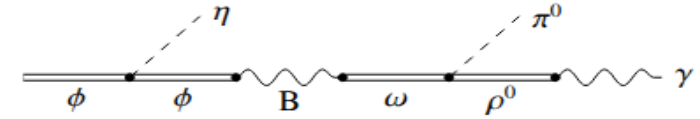
- P** is the pseudoscalar meson nonet
- V** is the vector meson nonet (**gauge boson** of a hidden **U(3)_V symmetry**)
- In conventional **VMD**: $A^\mu \rightarrow B^\mu, \quad e \rightarrow g_B, \quad Q = \frac{1}{3}$

$$\mathcal{L}_{V\gamma} = -4egf_\pi^2 A^\mu \text{tr} [QV_\mu] \longrightarrow \mathcal{L}_{VB} = -4\frac{1}{3}g_Bgf_\pi^2 B^\mu \text{tr} [V_\mu]$$



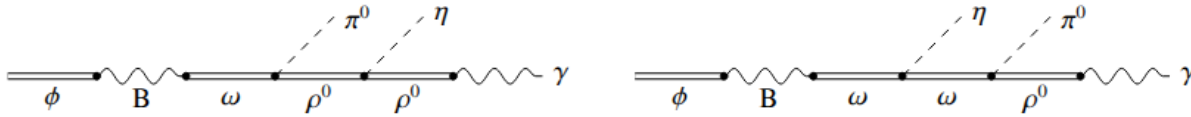
$$\Gamma_{B \rightarrow \pi^0 \gamma} = \frac{\alpha_{\text{em}} \alpha_B m_B^3}{96\pi^3 f_\pi^2} \left(1 - \frac{m_\pi^2}{m_B^2}\right)^3 |F_\omega(m_B^2)|^2$$

$\phi \rightarrow \pi^0 \eta \gamma$ decays: B boson calculation



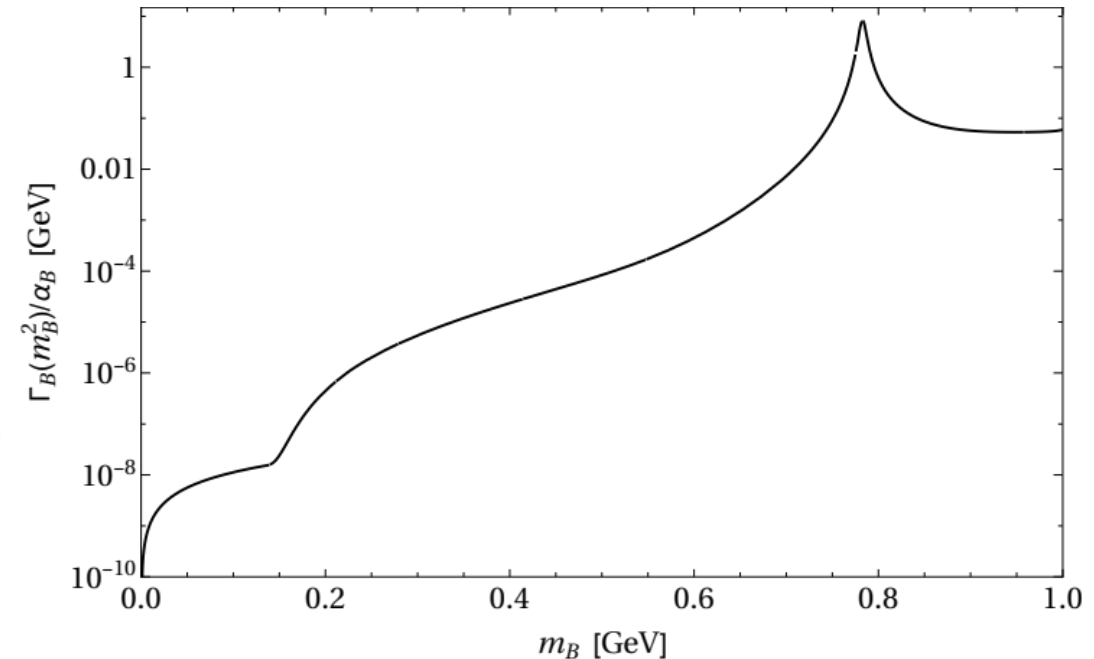
- **Three diagrams** corresponding to the exchange of a **B boson**:

$$\mathcal{A}_{\phi \rightarrow \pi^0 \eta \gamma}^{\text{B-boson}} = \frac{e g g_B^2 m_\rho^4}{16 \pi^4 f_\pi^2} \left\{ -\sqrt{2} s \varphi_P \frac{P^2 \{a\} + \{b(P)\}}{D_\phi [P^2] D_B [P^2] D_\omega [P^2]} + \frac{c \varphi_P}{D_B [m_\phi^2] D_\omega [m_\phi^2]} \left(\frac{P'^2 \{a\} + \{b(P')\}}{D_\rho [P'^2]} + \frac{P^2 \{a\} + \{b(P)\}}{D_\omega [P^2]} \right) \right\}$$



- **B boson width:**

$$\begin{aligned} \Gamma_B(s) = & \frac{\gamma_{B \rightarrow \pi^0 \gamma}(s)}{\gamma_{B \rightarrow \pi^0 \gamma}(m_B^2)} \Gamma_{B \rightarrow \pi^0 \gamma} \theta(s - m_\pi^2) + \frac{\gamma_{B \rightarrow e^+ e^-}(s)}{\gamma_{B \rightarrow e^+ e^-}(m_B^2)} \Gamma_{B \rightarrow e^+ e^-} \theta(s - 4m_e^2) \\ & + \frac{\gamma_{B \rightarrow \mu^+ \mu^-}(s)}{\gamma_{B \rightarrow \mu^+ \mu^-}(m_B^2)} \Gamma_{B \rightarrow \mu^+ \mu^-} \theta(s - 4m_\mu^2) + \frac{\gamma_{B \rightarrow \pi^+ \pi^-}(s)}{\gamma_{B \rightarrow \pi^+ \pi^-}(m_B^2)} \Gamma_{B \rightarrow \pi^+ \pi^-} \theta(s - 4m_\pi^2) \\ & + \frac{\gamma_{B \rightarrow \pi^+ \pi^- \pi^0}(s)}{\gamma_{B \rightarrow \pi^+ \pi^- \pi^0}(m_B^2)} \Gamma_{B \rightarrow \pi^+ \pi^- \pi^0} \theta(s - 9m_\pi^2), \end{aligned}$$

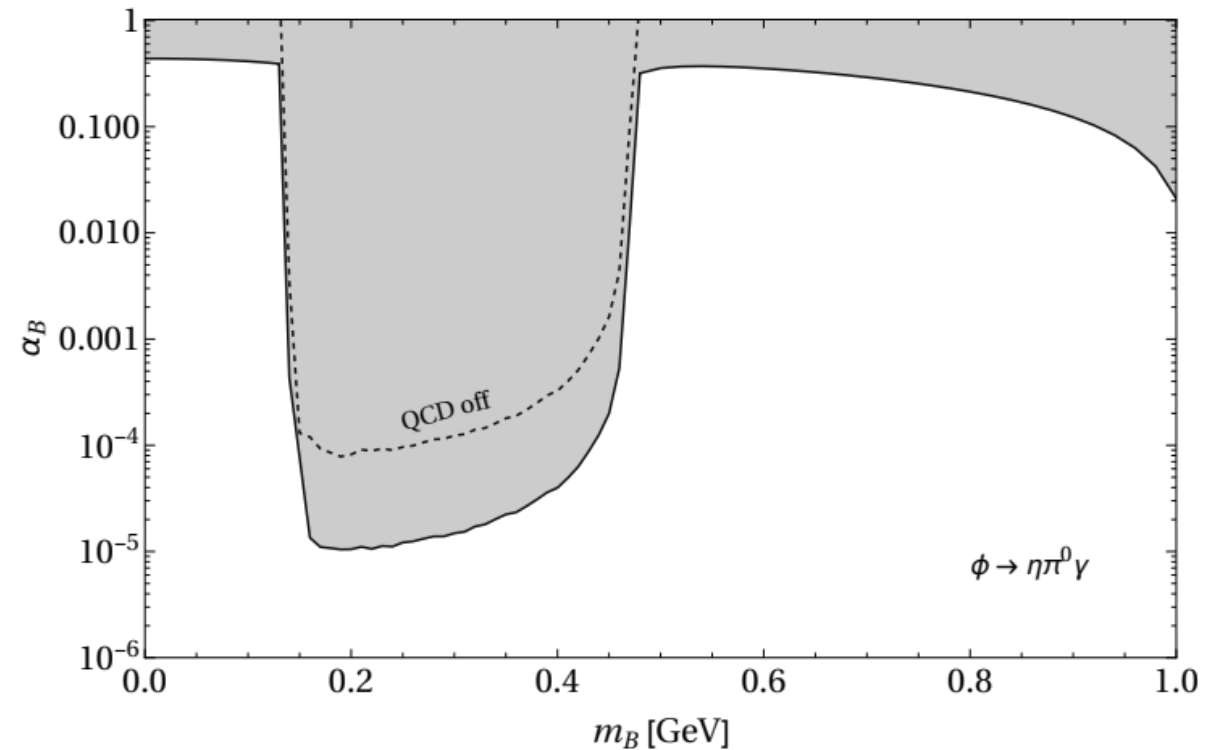


New limits on α_B and m_B

- **Not** assuming the **NWA**
- **QCD** contribution **on**
- $\text{Br}(\phi \rightarrow \pi^0 \eta \gamma) < \text{Br}_{\text{exp}}$ at 2σ

$$\text{Br}(\phi \rightarrow \pi^0 \eta \gamma)_{\text{exp}}^{\text{PDG}} = (7.27 \pm 0.30) \times 10^{-5}$$

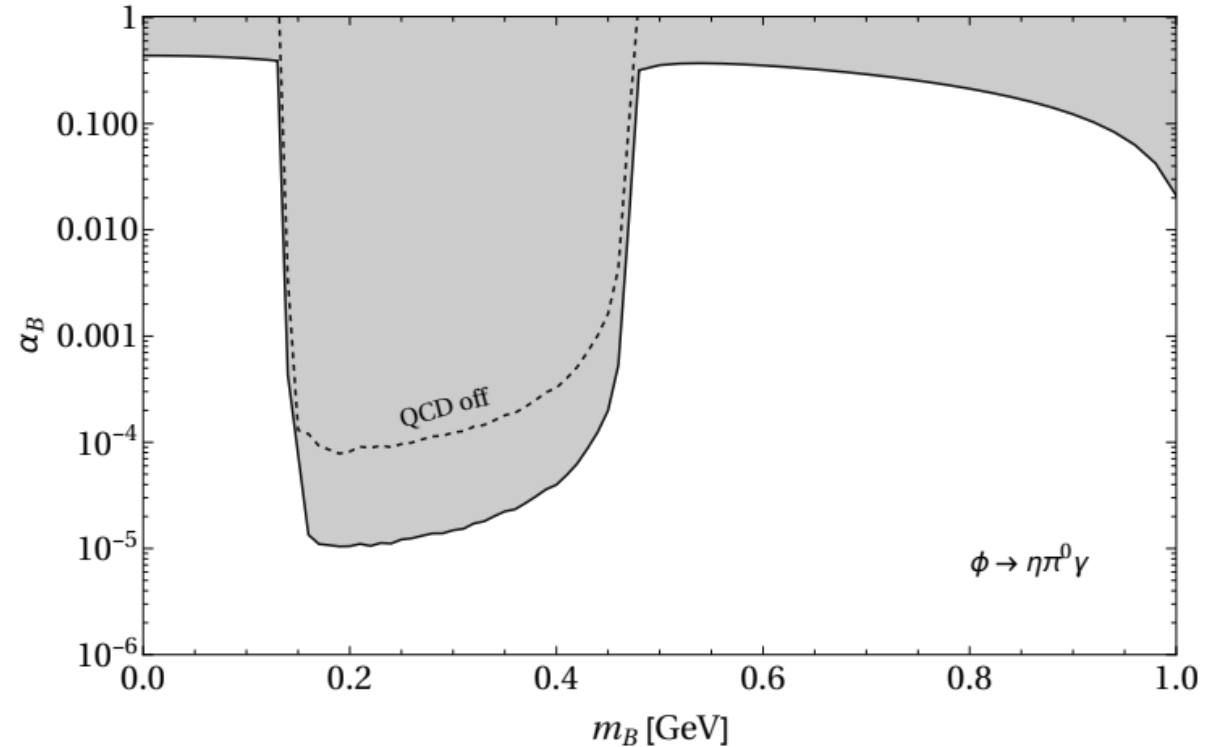
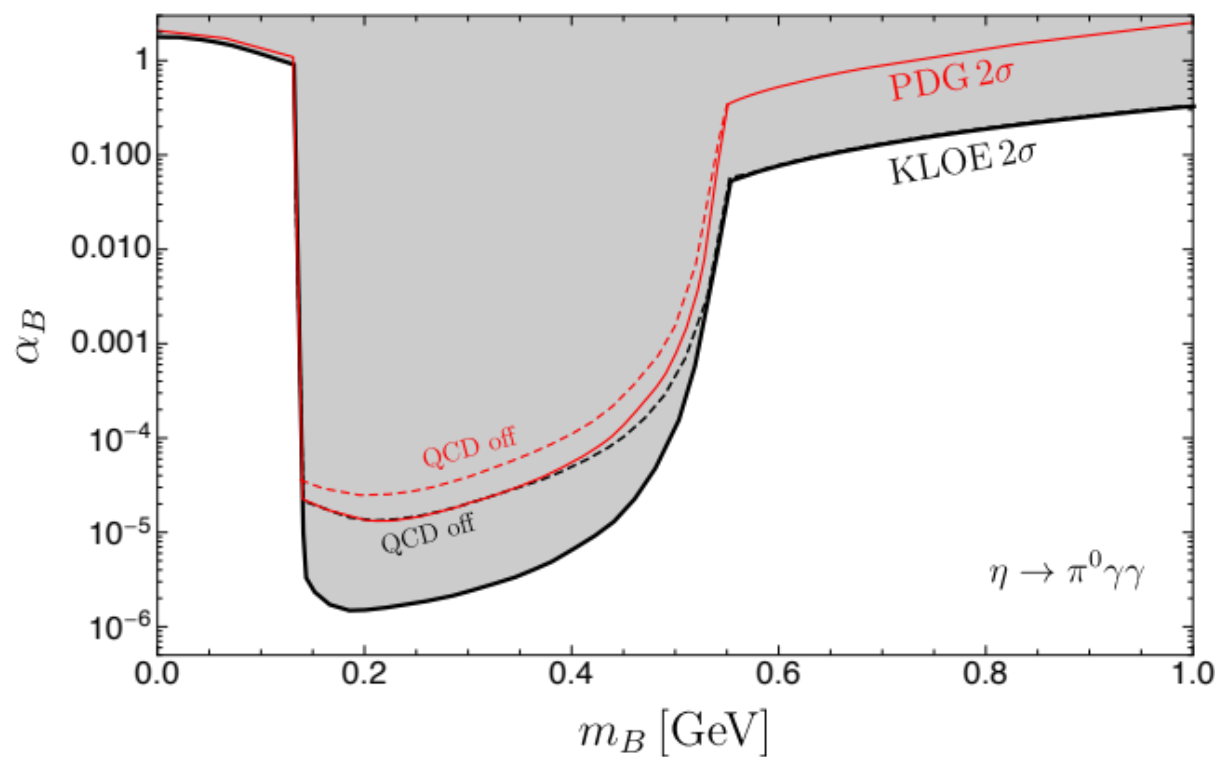
S. Navas et al. [PDG], Phys. Rev. D 110, 03001 (2024)



- **Limits** on the **B boson** parameters **improved by an order of magnitude** compared to previous analysis.

S. Tulin. Phys.Rev.D 89 (2014) 11, 114008

New limits on α_B and m_B



- Limits on the **B boson** parameters improved by an order of magnitude,
- But weaker than those from $\eta \rightarrow \pi^0 \gamma \gamma$ decays.

R. Escribano et al. Phys.Rev.D 106 (2022) 11, 114007

B boson contributions to the branching ratio

Contribution	$\mathcal{B}(\phi \rightarrow \pi^0 \eta \gamma) \times 10^5$	
LσM	6.52	
VMD	0.36	
Interference SV	+0.02	
Contribution	Set 1	Set 2
<i>B</i> -boson	0.08	$2.8 \cdot 10^{-4}$
Interference SB	$+5.8 \cdot 10^{-7}$	$-7.0 \cdot 10^{-4}$
Interference VB	$+1.8 \cdot 10^{-6}$	+0.02



Results from Fit 1

$$\alpha_B = 10^{-6}, m_B = 250 \text{ MeV}$$

$$\alpha_B = 10^{-2}, m_B = 540 \text{ MeV.}$$

B-boson signature strongly suppressed.

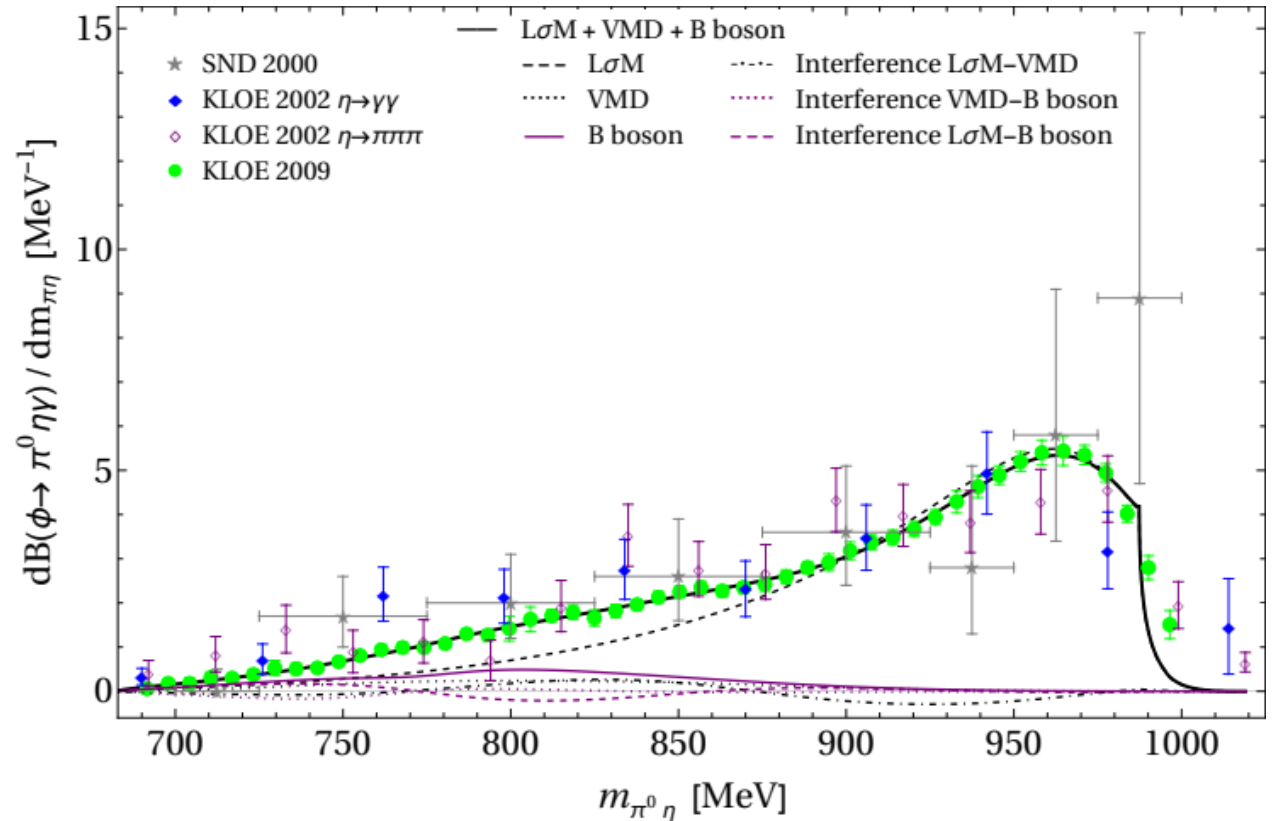
Fits to $\phi \rightarrow \pi^0 \eta \gamma$ decays

- Much **more statistics** from **KLOE**

$\chi^2/d.o.f$	$58.2/42 = 1.4$
m_{a_0}	980.14 ± 2.68
$g_{a_0 K \bar{K}}$	$\pm(2401.3 \pm 146.6)$
$g_{a_0 \pi \eta}$	$\pm(3190.1 \pm 174.8)$
δ_{+0}	$(48.1 \pm 29.7)^\circ$
m_B	475.4 ± 4.4
α_B	0.35 ± 0.07
δ_{0B}	$(165.8 \pm 13.3)^\circ$

- Signatures outside may be visible

$$m_{\pi^0} \lesssim m_B \lesssim m_\phi - m_\eta$$



Summary

- **Scalar** contribution **dominates**, **vector** contributions **subdominate**.
- **LσM fails** to accurately describe the data from **KLOE**.
- Allowing the parameters to **vary freely**, the SM description is **improved**.
- **Sensitivity to B boson** has been **analyzed**.
- **Limits on α_B and m_B** have been **strengthened** by **an order of magnitude**.
- The **$\phi \rightarrow \pi^0 \eta \gamma$ decays** provide **weaker constraints** on the B-boson parameters compared to **$\eta \rightarrow \pi^0 \gamma \gamma$** below the mass of the η meson.