



Photon production in high-energy collisions expectations, lessons and questions

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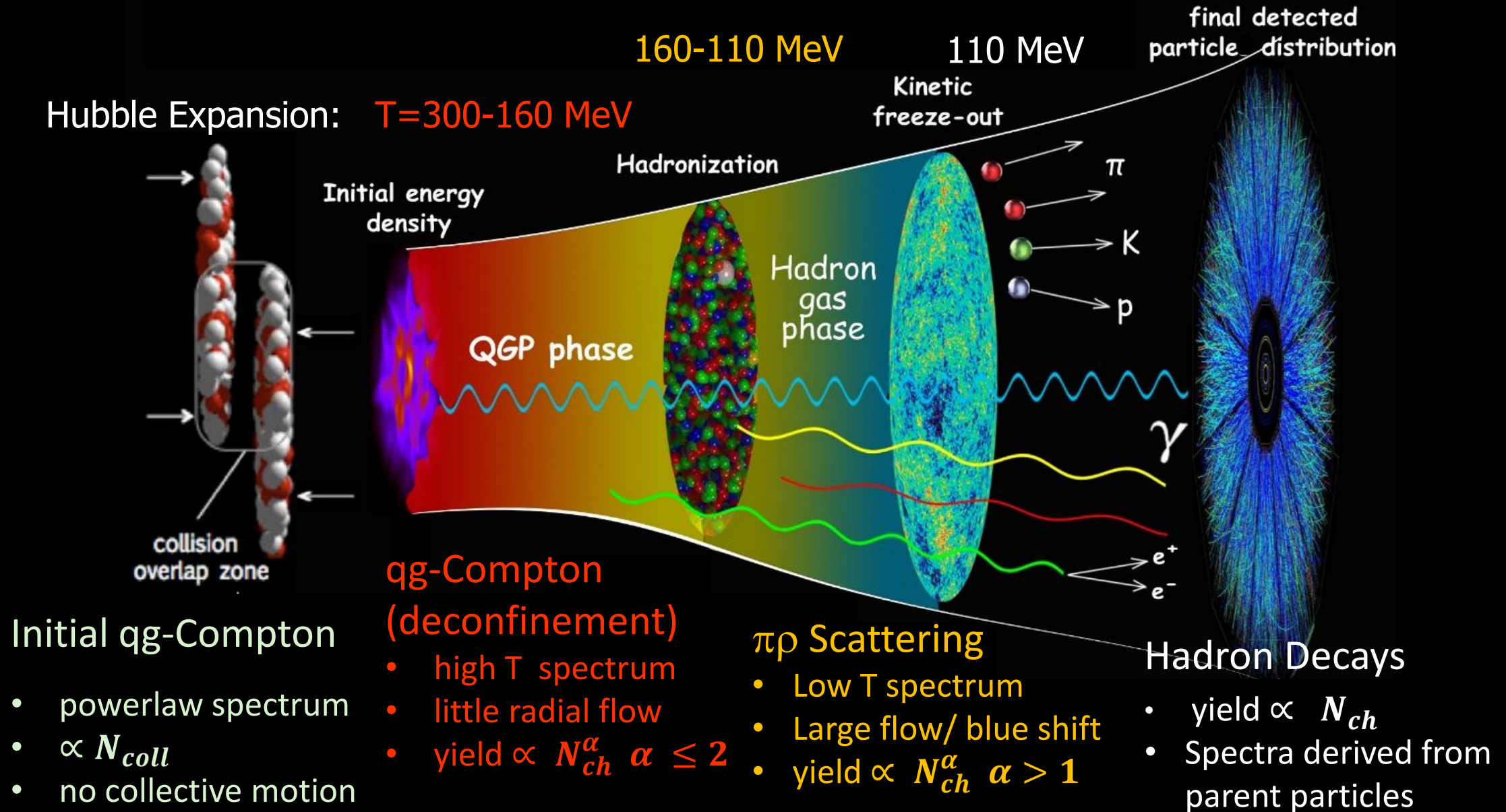




The big picture



Electromagnetic Radiation in A+A Collisions:



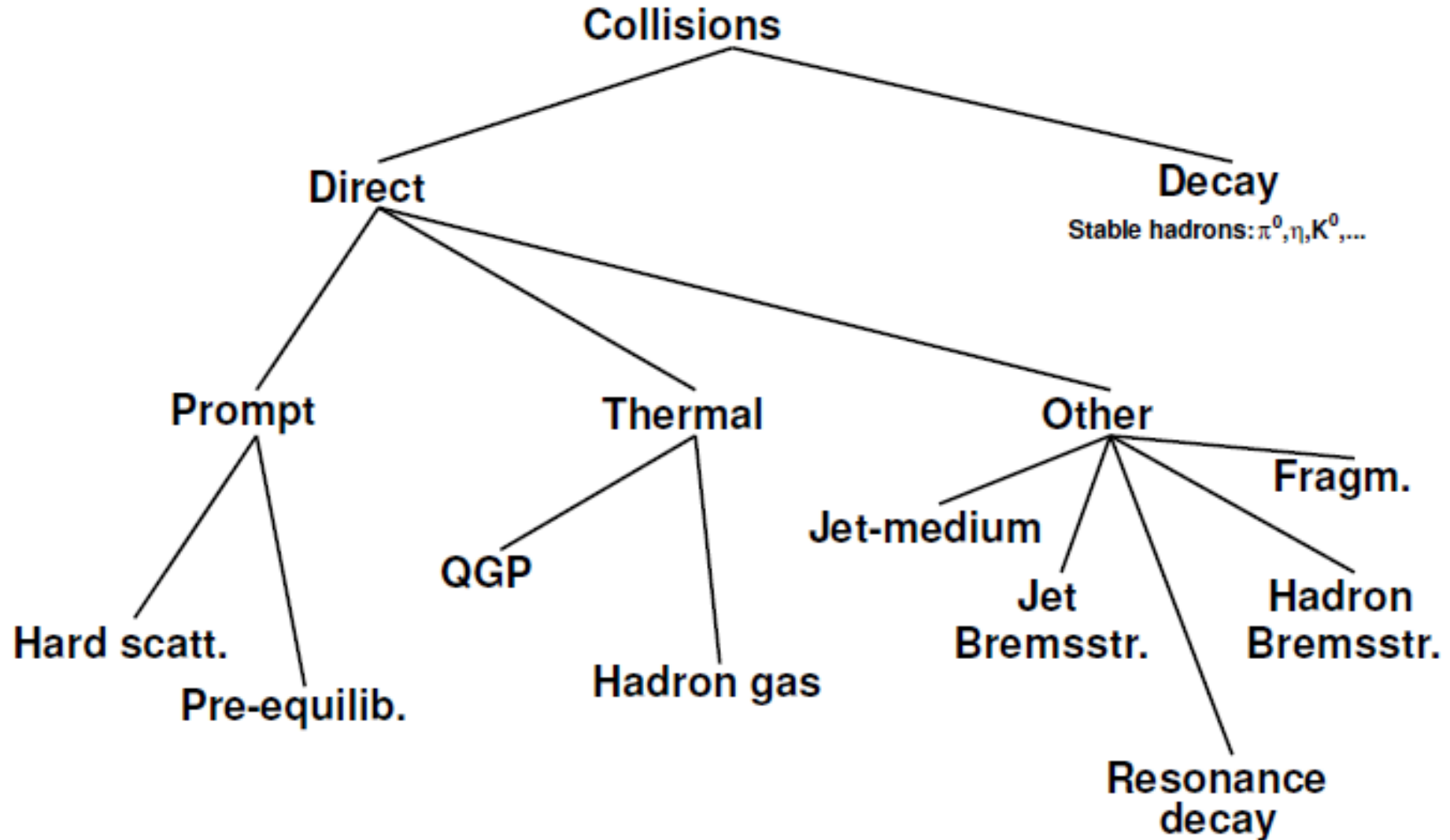


The terminology Historians of the collision

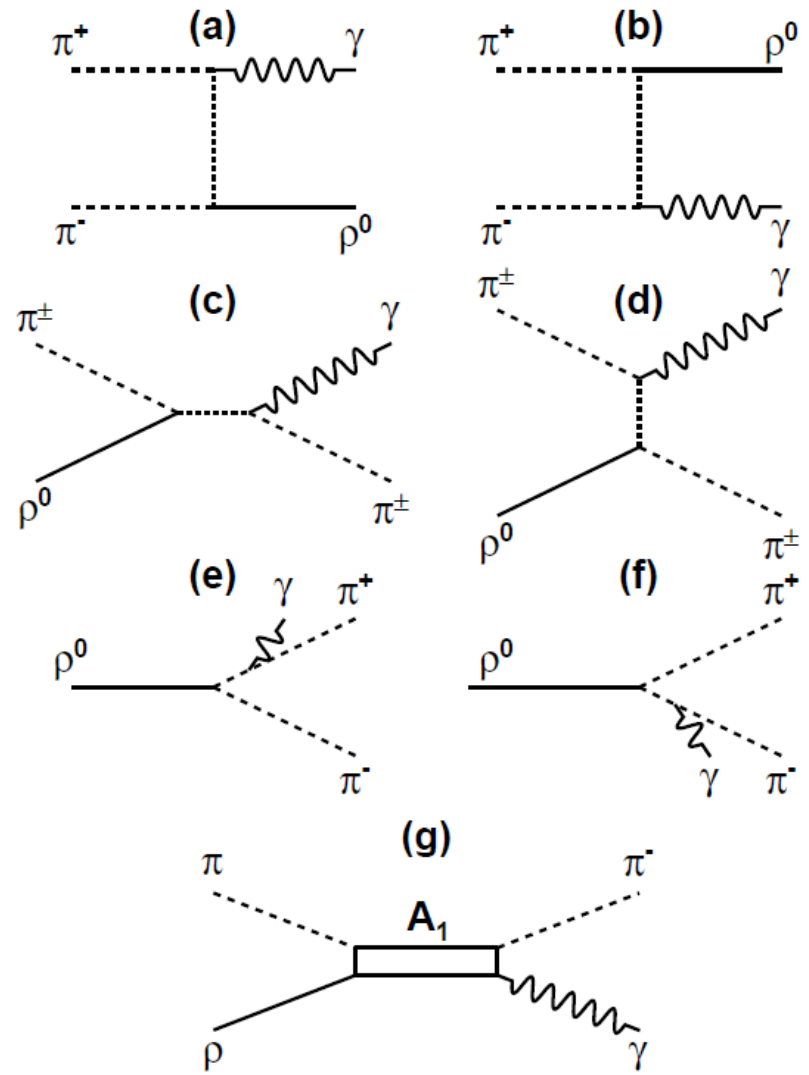
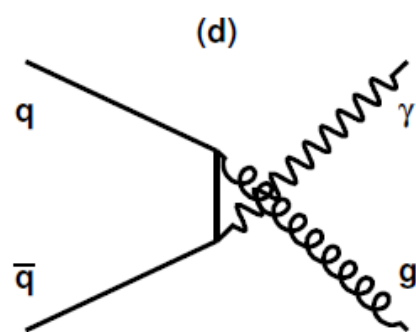
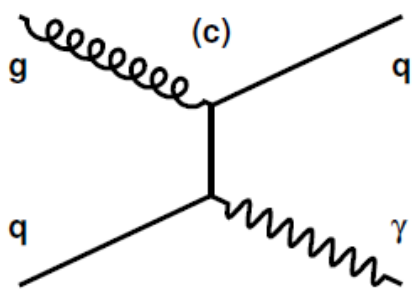
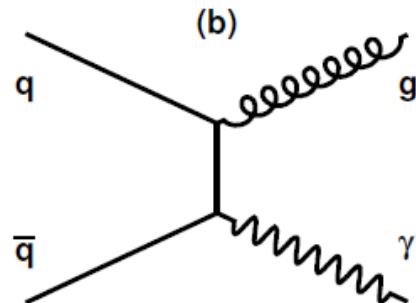
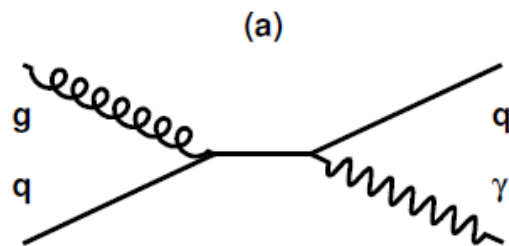




Terminology – reflecting different sources



The basic partonic and hadronic sources



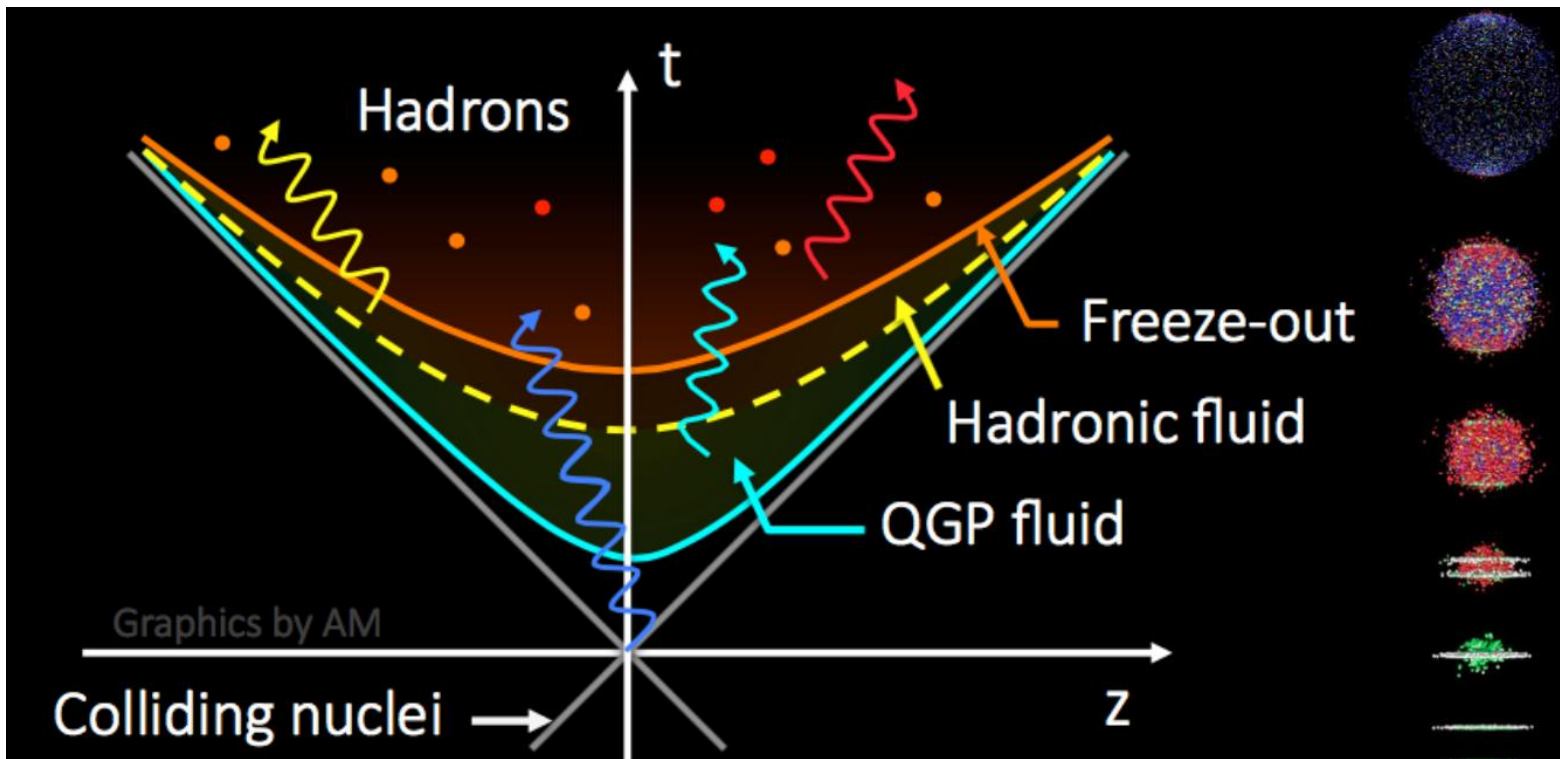
The blessing – and curse – of photons



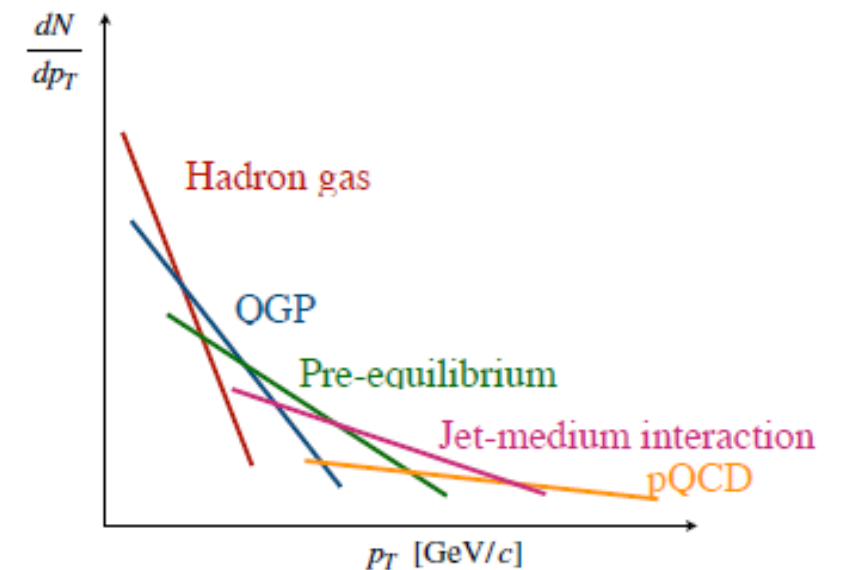
Penetrating probes, created at all stages

Huge background (FS decays, $\sim 10\%$ S/B)

Hard to differentiate between sources



Photons are “color blind” probe of Quark Gluon Plasma



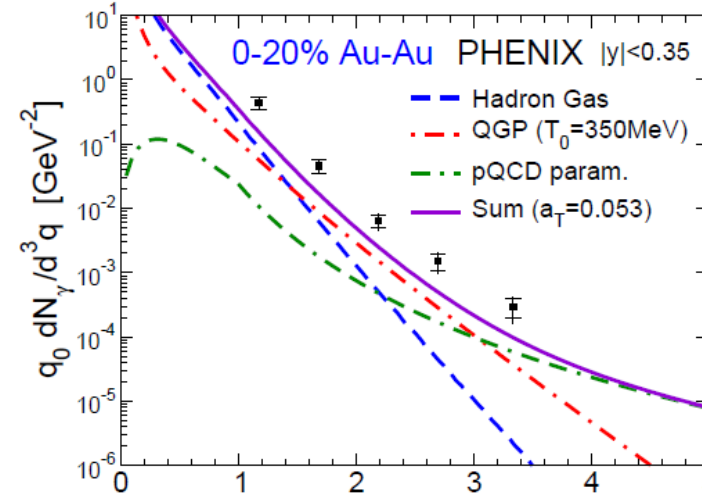
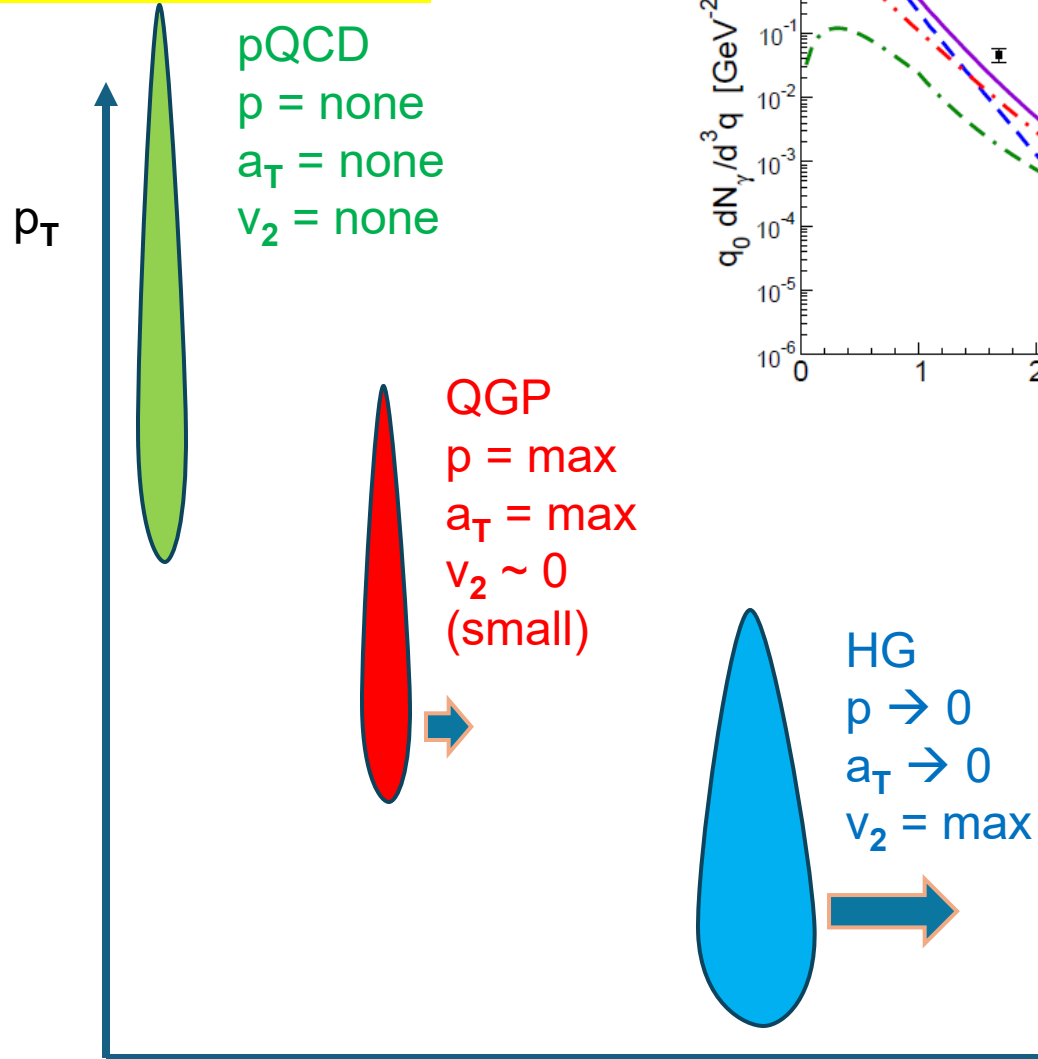
“Historians” of the collision – but convoluted



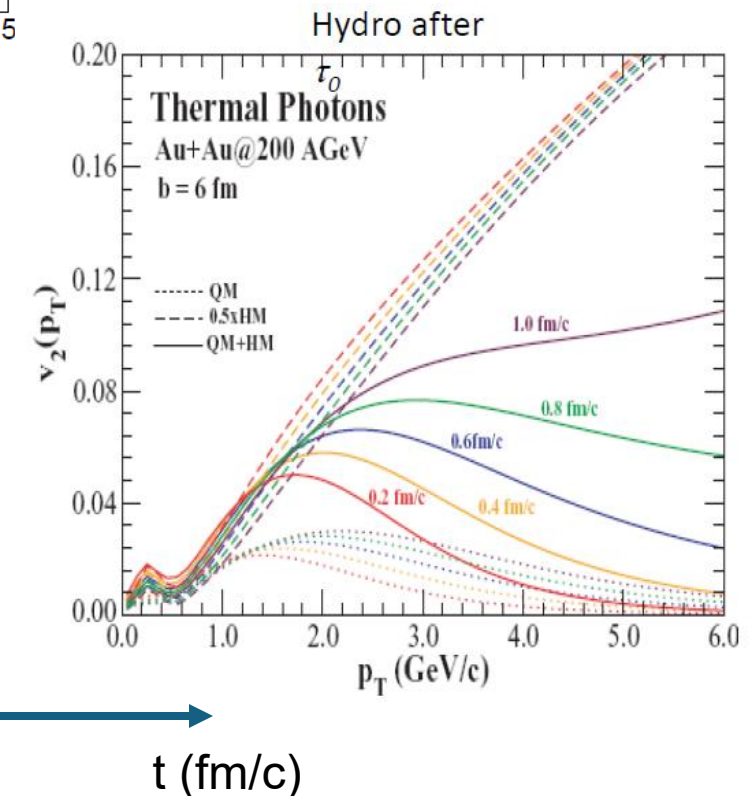
Dominant photon sources (simplified)



Prelude to the “puzzle”



In principle one can try to deconvolute the individual contributions starting from the highest p_T



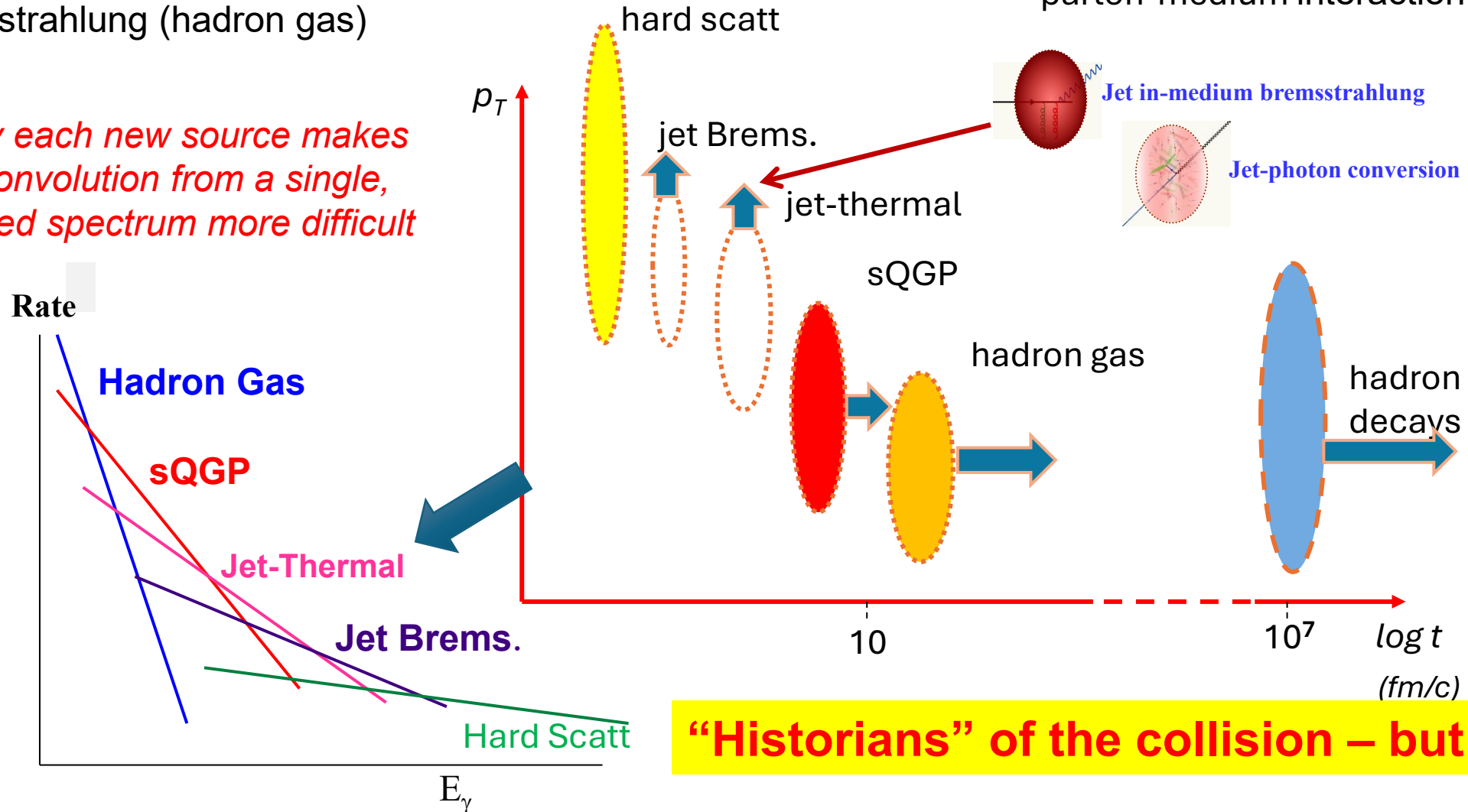
More sources...

Jet fragmentation,
 Jet-thermal interaction
 (jet-photon conversion)
 Initial magnetic field
 Bremsstrahlung (hadron gas)
 ...???

See e.g., Turbide, Gale, Jeon and Moore, PRC 72, 014906 (2005)

parton-medium interaction

Obviously each new source makes the deconvolution from a single, integrated spectrum more difficult



“Historians” of the collision – but convoluted





The reference – or is it? (p+p)

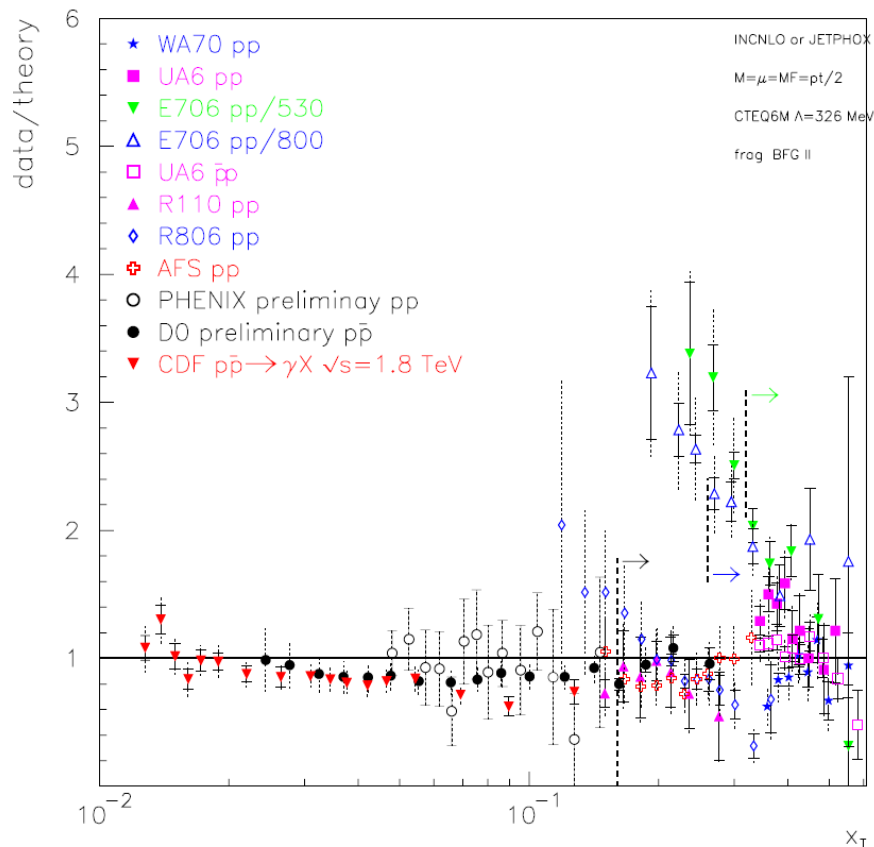


Direct photons in hadron (p+p) collisions, 22 – 7000 GeV

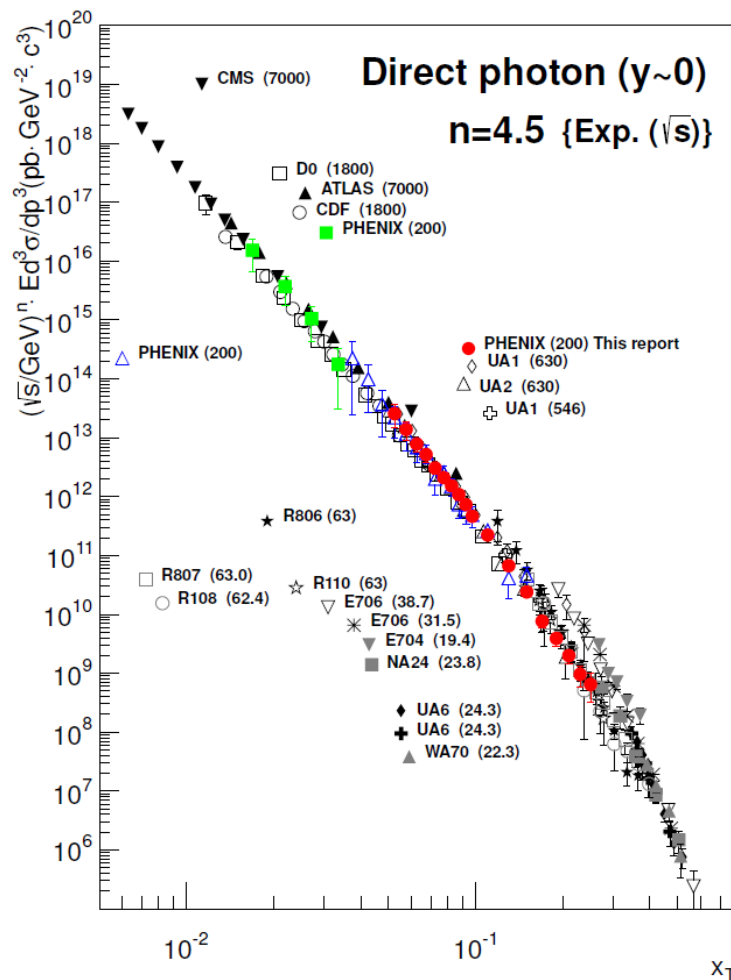


High p_T , scaling with x_T

PRD 73, 094007 (2006)



PRD 86, 072008 (2012)



$$E \frac{d^3 \sigma}{dp^3} = \frac{d^3 \sigma}{p_T dp_T dy d\phi} = \frac{1}{p_T^{n_{\text{eff}}(x_T, \sqrt{s})}} F\left(\frac{p_T}{\sqrt{s}}\right) = \frac{1}{(\sqrt{s})^{n_{\text{eff}}(x_T, \sqrt{s})}} G(x_T),$$

Over 18 orders of magnitude

Without evolution of α_s and the structure and fragm. functions $n=4$

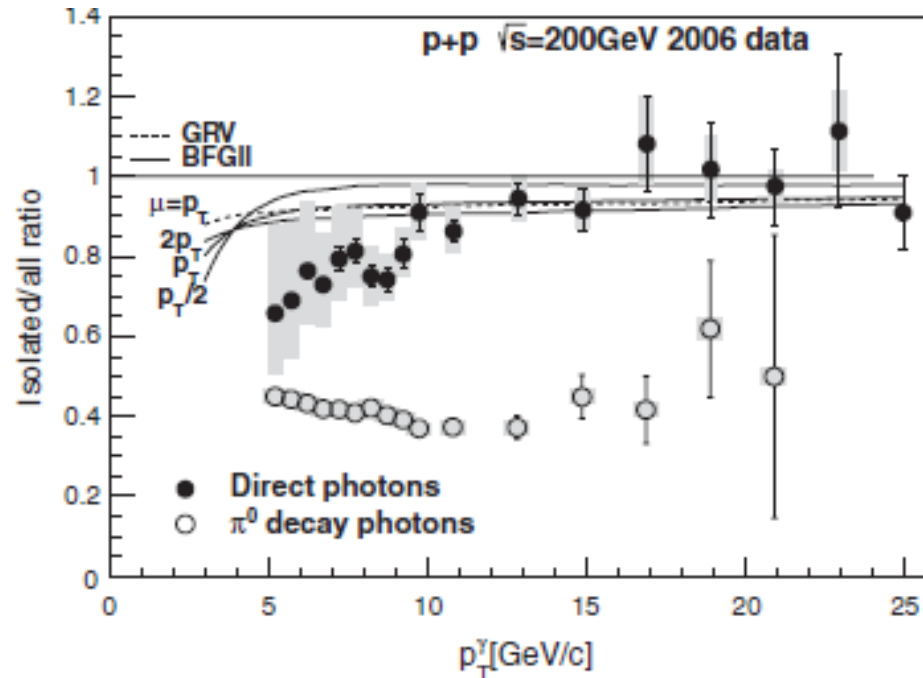
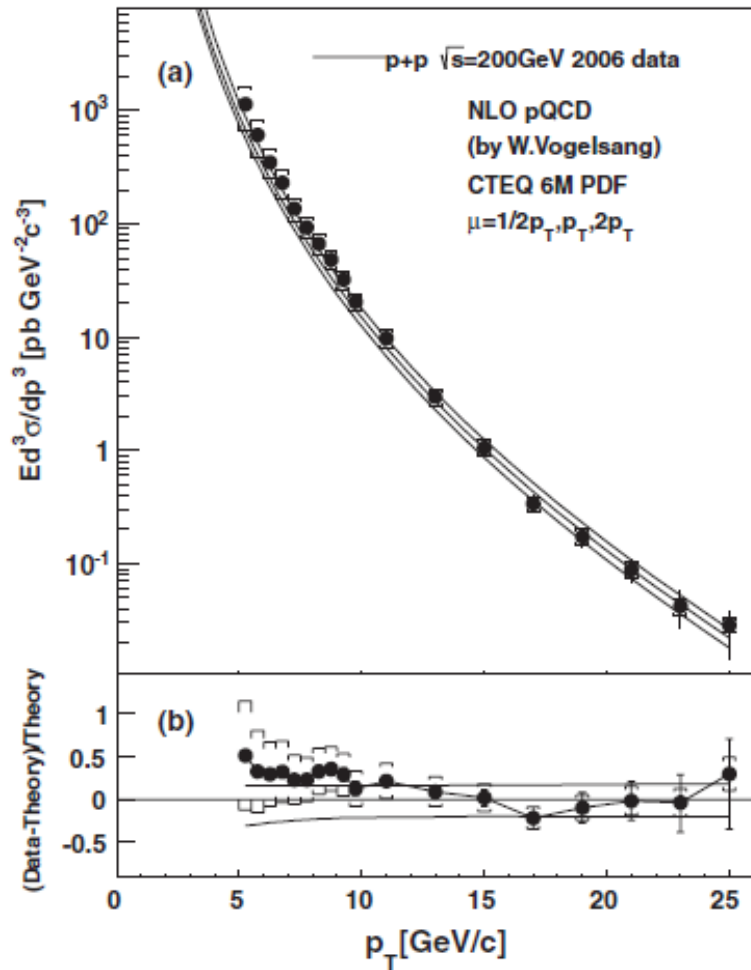
High p_T reasonably well understood



Direct photons in hadron (200 GeV p+p) collisions, 2012

Photons vs NLO pQCD

PRD 86, 072008 (2012)



Fraction of isolated direct photons

Fragmentation very small above 10 GeV/c

Possible excess below 10 GeV/c?

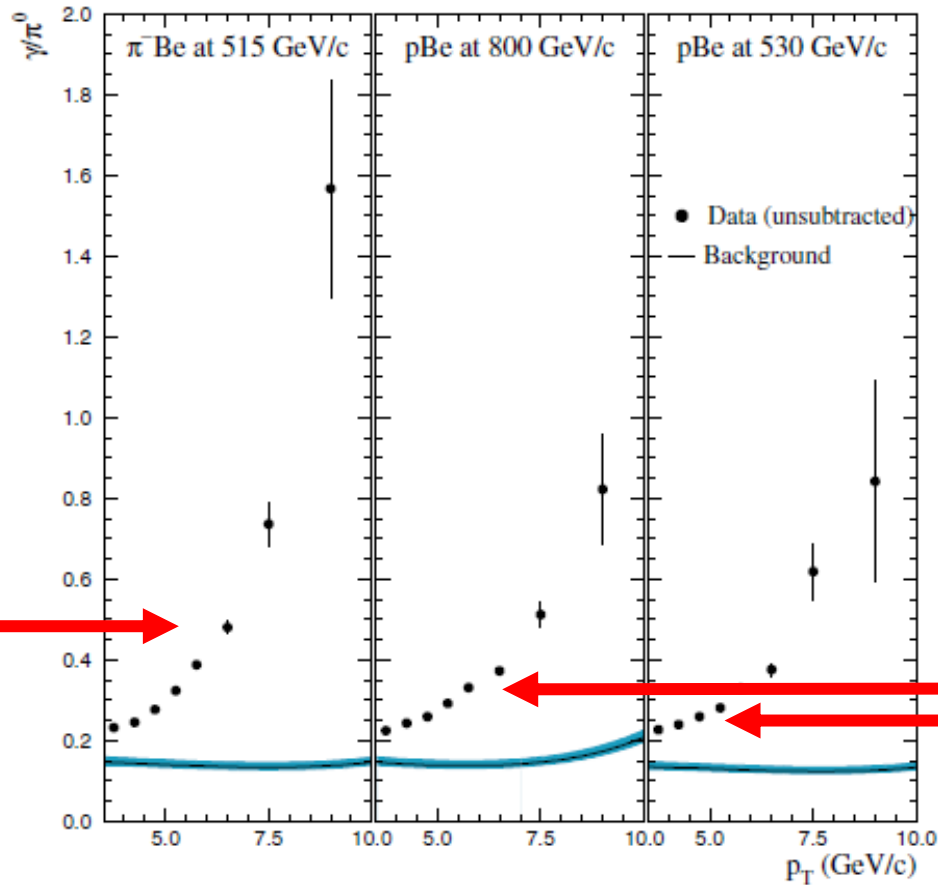


Interlude: Compton scattering dominates over annihilation



E706, fixed target, p and π^- beam

PRD 70, 092009 (2004)



Only sea antiquarks in proton

→ small probability for high p_T photon from annihilation

Valence antiquark in p^-

→ high p_T photon production increased

Not quite direct proof (γ/π^0), but strong

Valence antiquark present

No valence antiquark

We knew this, but good to see experimentally



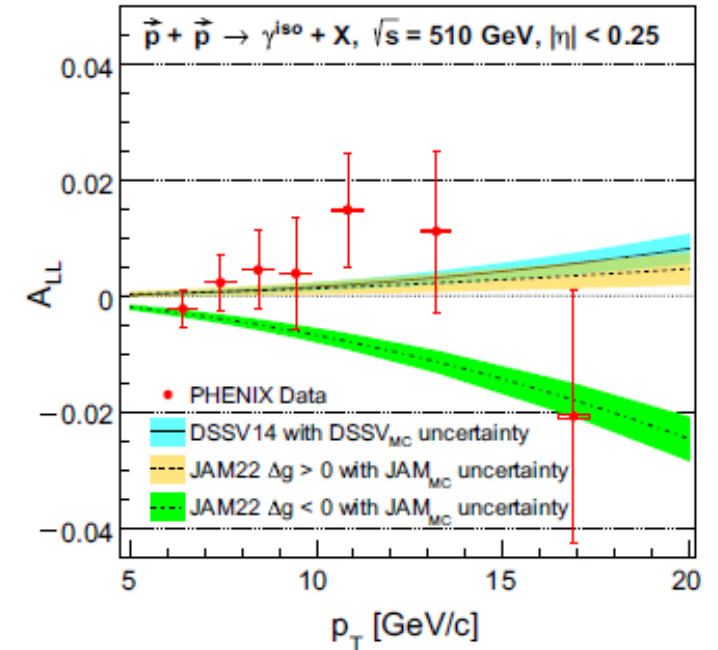
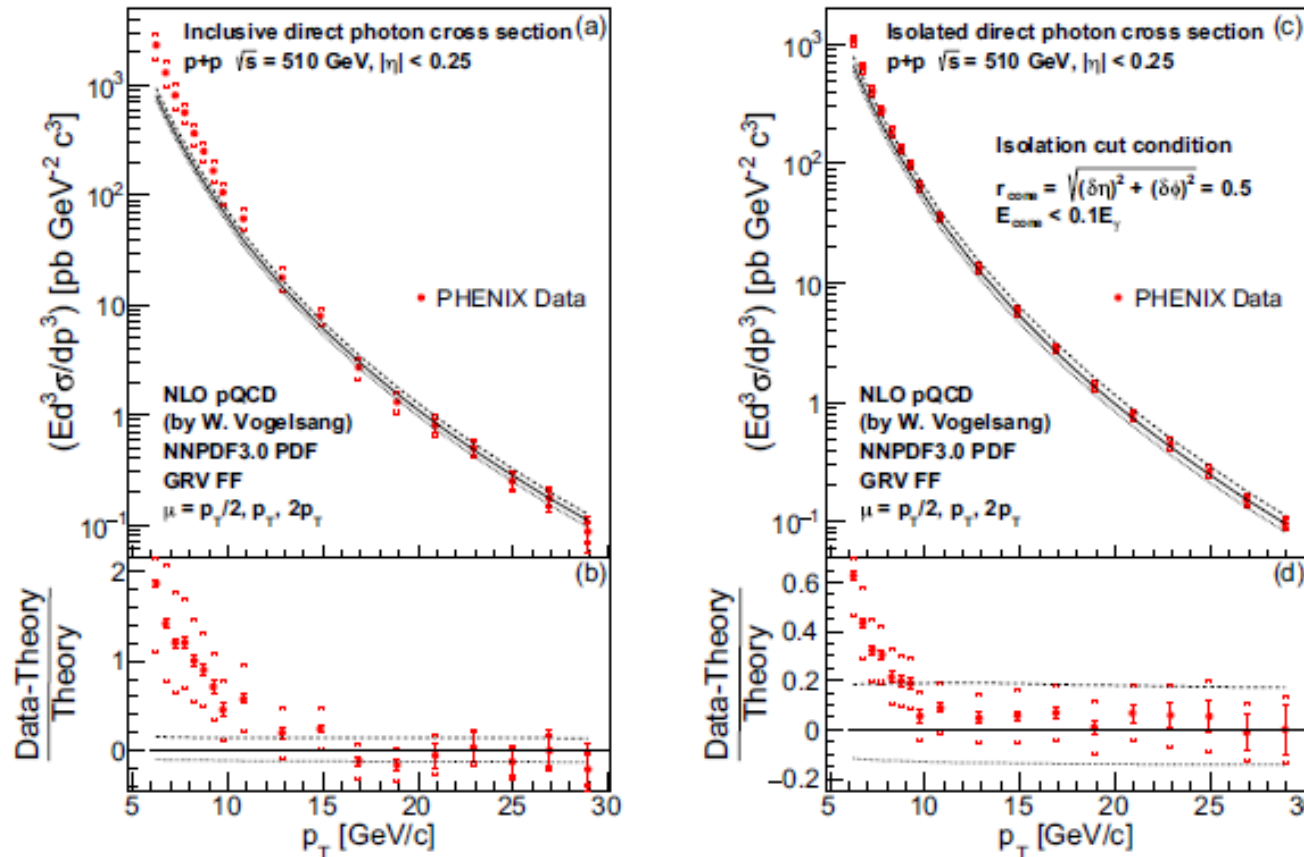
Direct photons in polarized 510 GeV p+p collisions



Polarized gluon contribution to proton spin?

PRL 130, 251901 (2023)

A_{LL} – longitudinal double spin asymmetry



Disfavors negative Δg

Similar discrepancy below 10 GeV/c, as in 200 GeV collisions – even for isolated photons

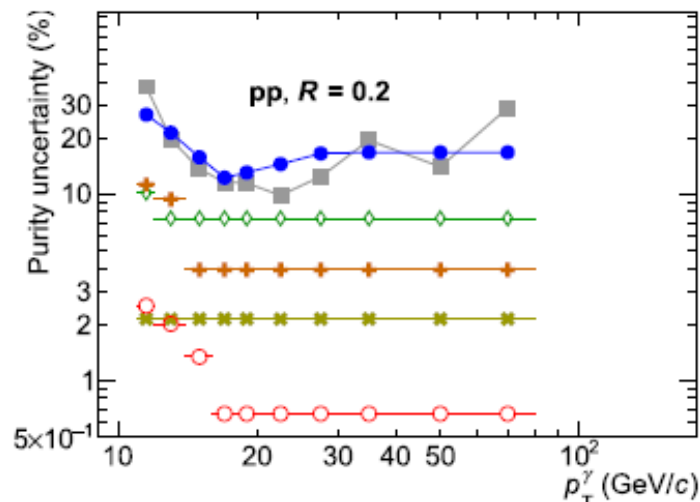
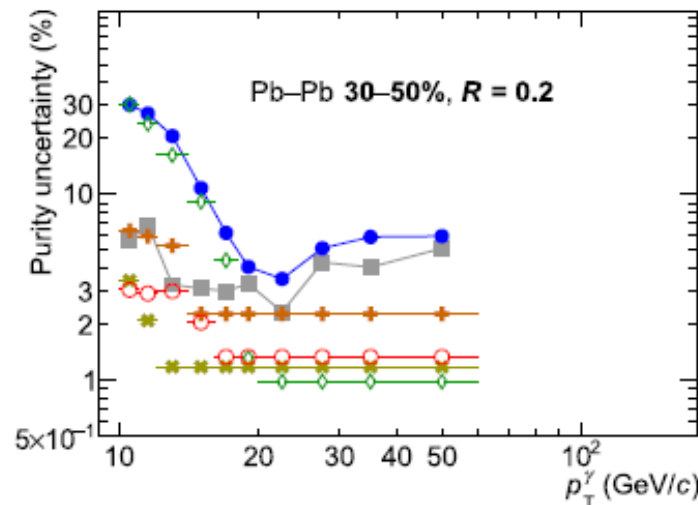
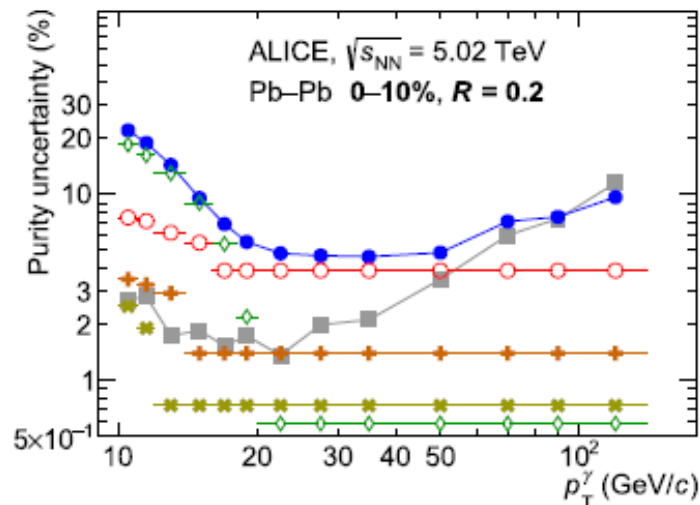
Small, positive contribution from gluon



Technical – isolated photon uncertainties (ALICE)

Can you still isolate in PbPb?

Eur. Phys. J. C (2025) 85:553



- Total, including fit.
- Statistical
- ◇ Isolation probability
- * Bkg. $\sigma^2_{\text{long}, 5 \times 5}$
- + Bkg. $p_T^{\text{iso, ch}}$
- MC signal amount

Isolation cut purity – uncertainties
→ this is the overwhelming term
in the total uncertainty on
cross-section

Isolation probability
→ ratio of shape distributions
of iso and non-iso constant?

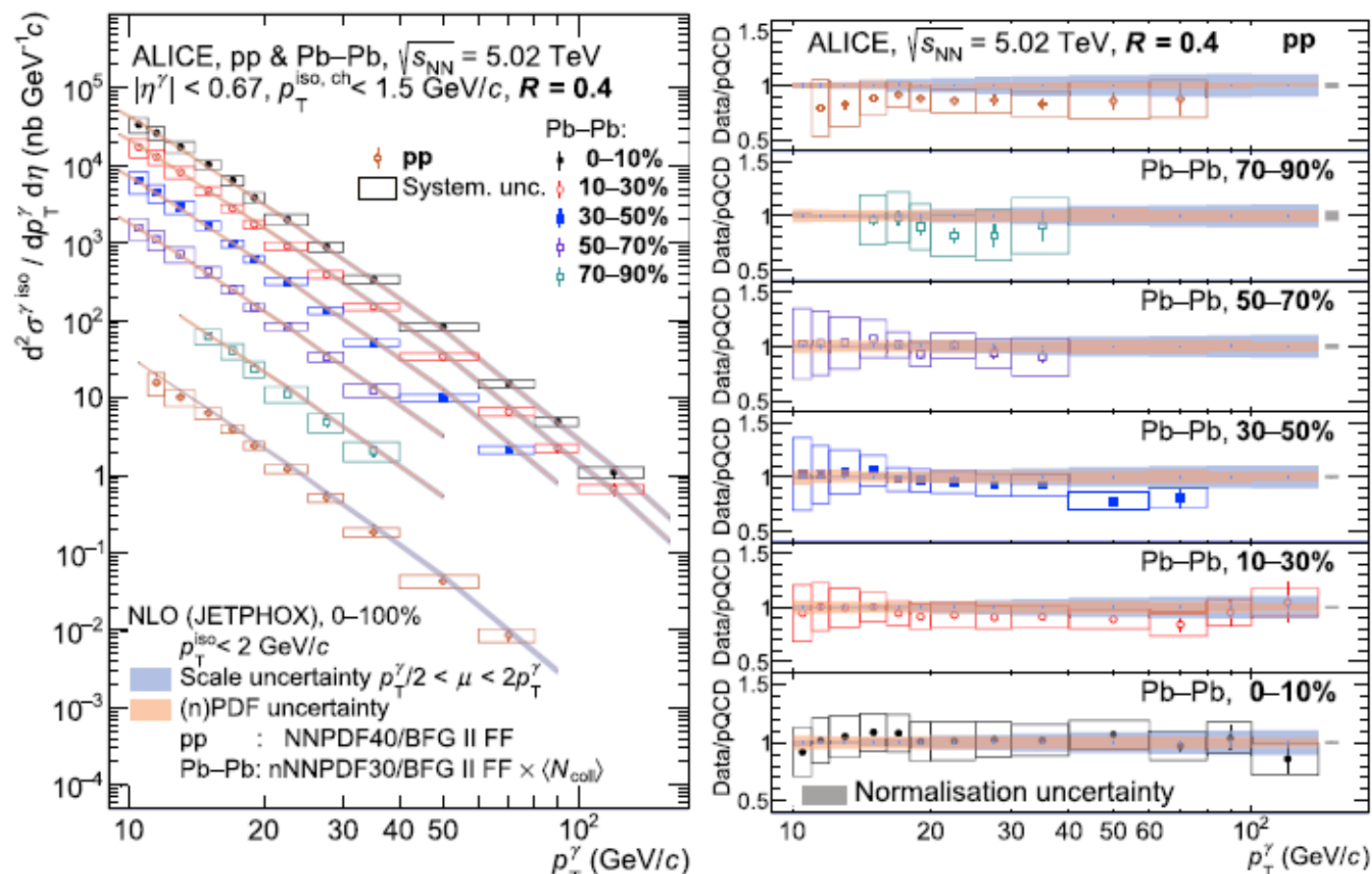
MC signal amount
→ varying signal (γ -jet) by 20%
w.r.t. bgd (jet-jet)

Reverse ordering in PbPb vs pp



Is the high p_T region understood?

Eur. Phys. J. C (2025) 85:553

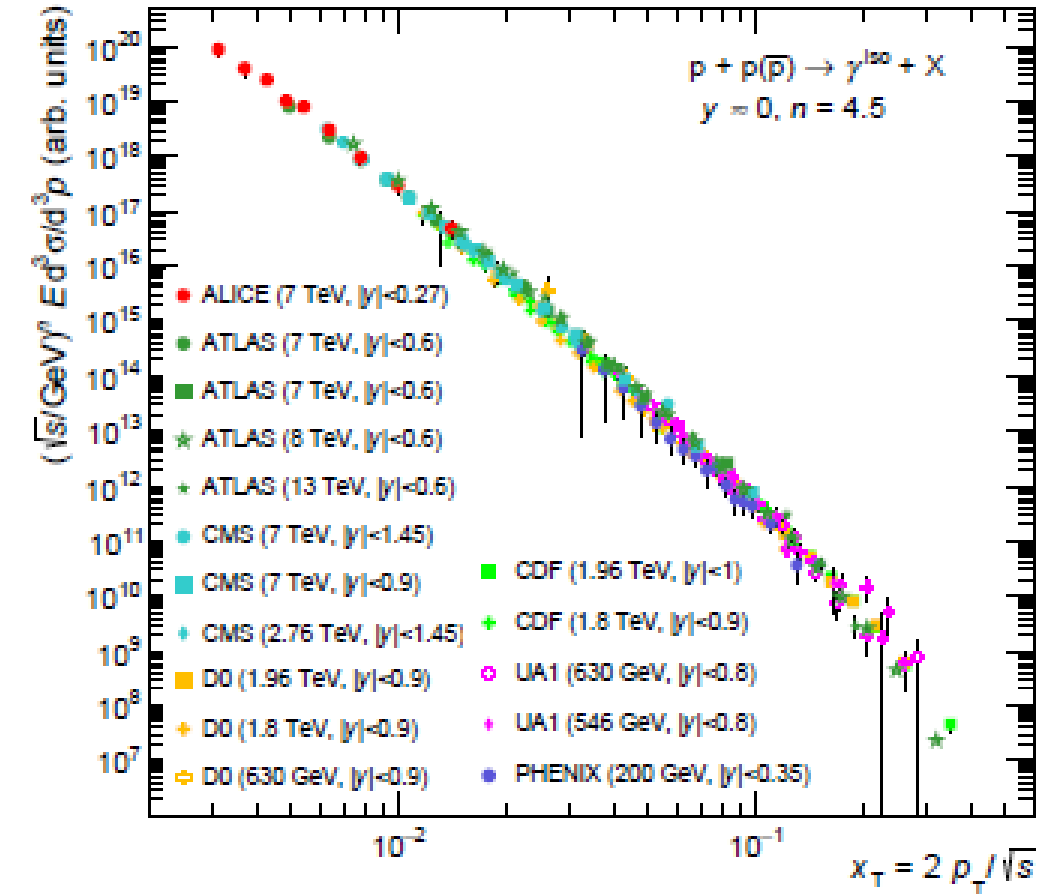
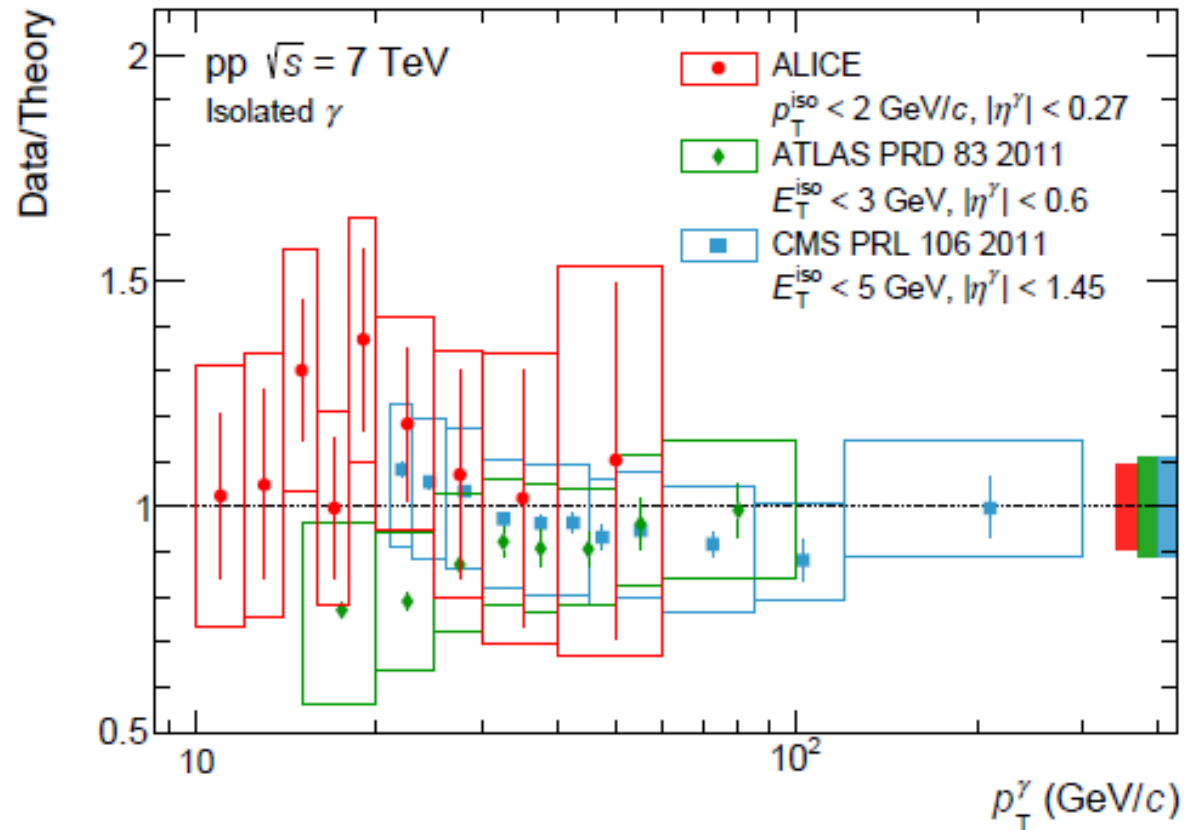


All consistent with JETPHOX pQCD



Is the high p_T region understood?

Eur. Phys. J. C (2019) 79:896



All consistent with JETPHOX pQCD

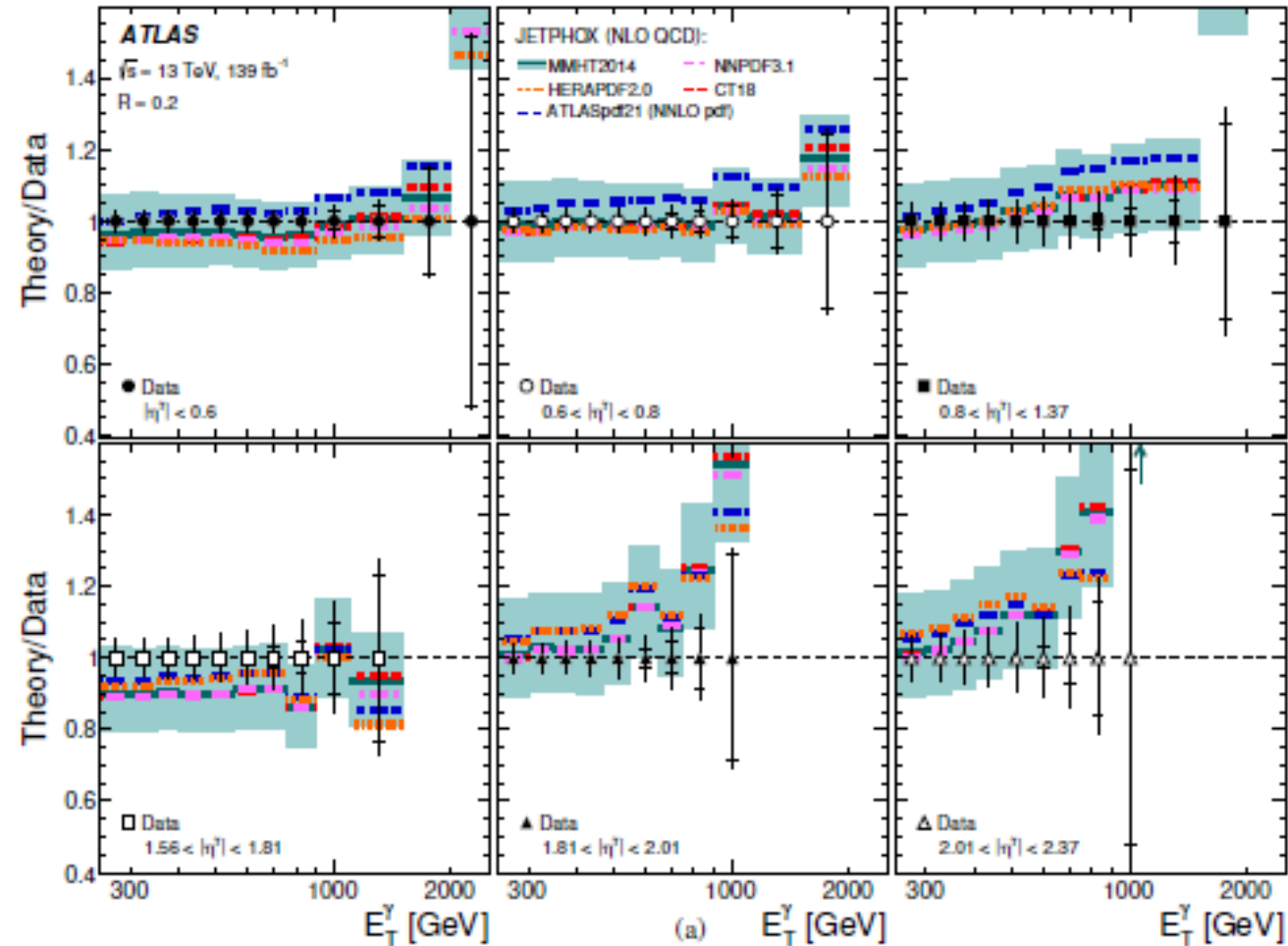


High p_T isolated photons,, pp 13 TeV (ATLAS)



Rapidity dependence?

JHEP07 (2023) 086



Not well described by JETPHOX pQCD



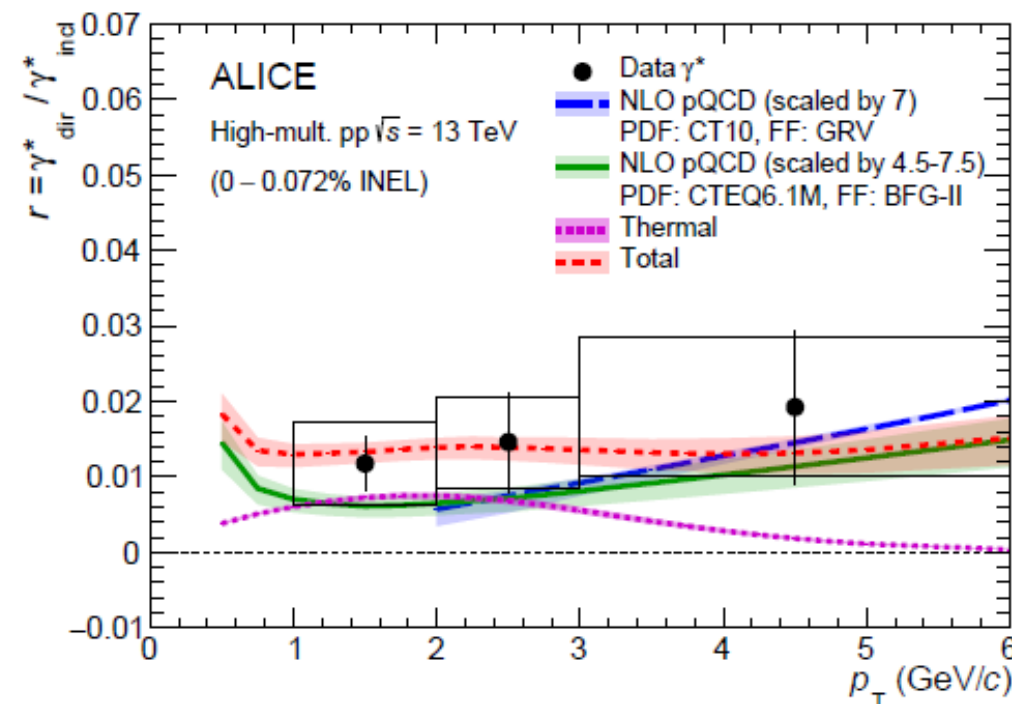
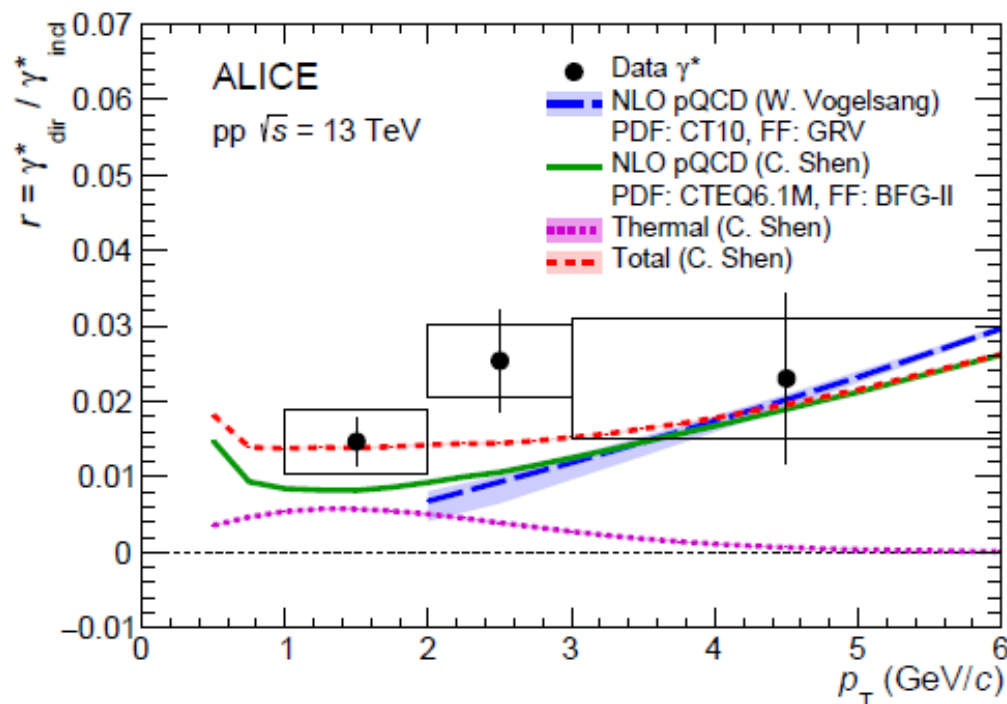
QGP even in high multiplicity pp?

Something like QGP formed already in the highest multiplicity pp events, too?

Internal conversion $\rightarrow \gamma^*$

PLB 868 (2025) 139645

$$r = \gamma_{\text{dir}}^* / \gamma_{\text{had}}^*$$



No sign of additional “thermal” radiation





Experimental techniques



Photon measurements in PHENIX



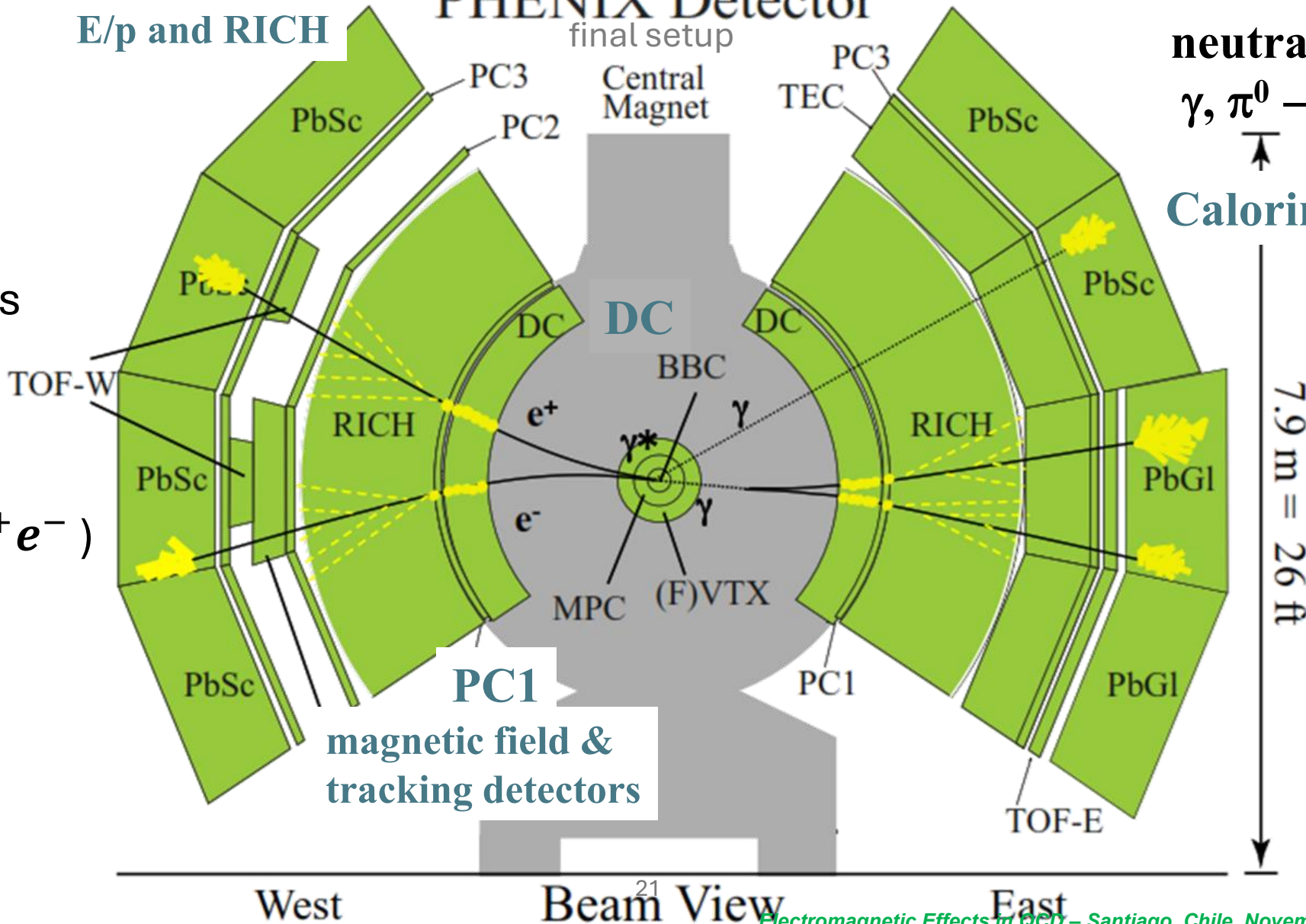
e^+e^- identification
E/p and RICH

PHENIX Detector
final setup

Photons,
neutral pion
 $\gamma, \pi^0 \rightarrow \gamma\gamma$
Calorimeter

Virtual Photons
 $\gamma^* \rightarrow e^+e^-$

Photons
 $\lim_{m_{ee} \rightarrow 0} (\gamma^* \rightarrow e^+e^-)$
 $\gamma \rightarrow e^+e^-$



Basic techniques of photon measurements: calorimetry

Real photon captured in an electromagnetic calorimeter

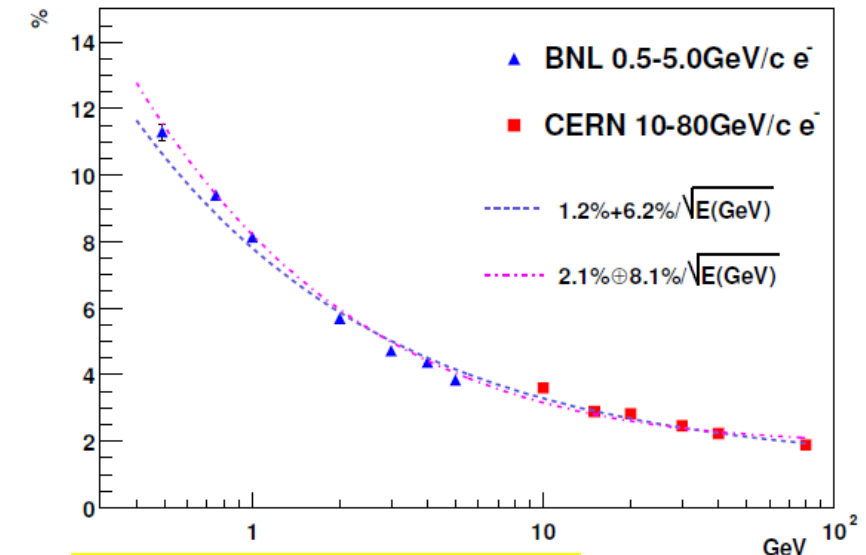
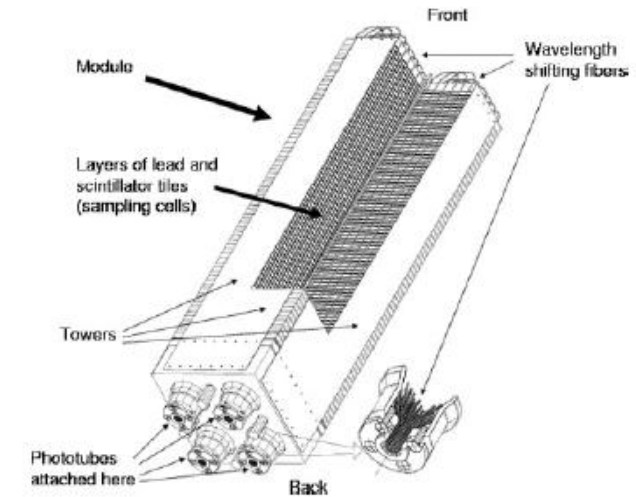
Electromagnetic calorimeter

- relatively cheap, large acceptance possible (sampling)
- resolution improves with energy (ideal at high p_T)
- “thin” for hadrons
 - high energy cluster (almost) always electromagnetic
- deposited energy counts, not particle momentum!**
- measurement can be almost standalone

But

- issues with resolution and hadron background at low p_T
- issues with merged π^0 decay photons at very high p_T

These can be mitigated: shower maximum detectors (SMD), small Moliere-radius crystal calorimeters, digital calorimeters
→ but at substantial cost increase

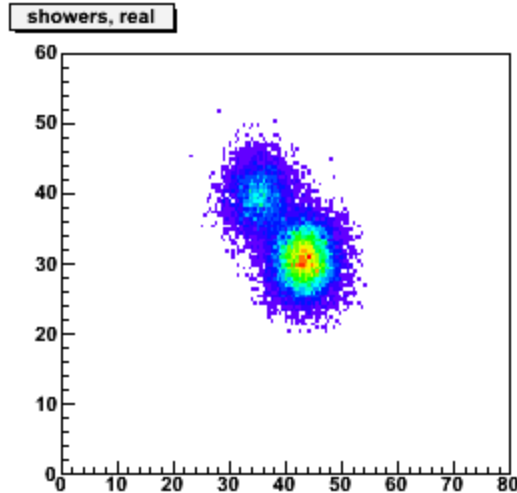


Best at high p_T

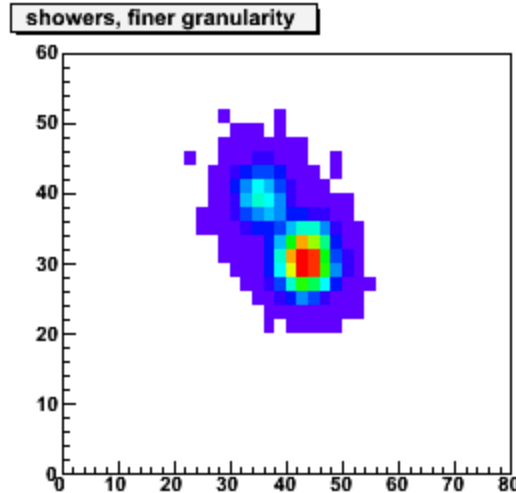
Close-by showers – real and as seen



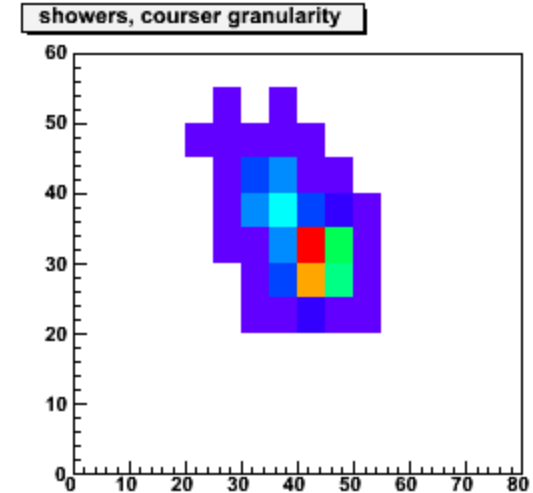
Nature meets detector...



Very high granularity:
well separated,
good measurement
for both energy and
position



Medium granularity:
well separated,
decent measurement
for energy, some
shift in position,
pair p_T still correct



Low granularity:
poor separation.
**questionable energy
sharing**, position
shift,
pair p_T still ~correct

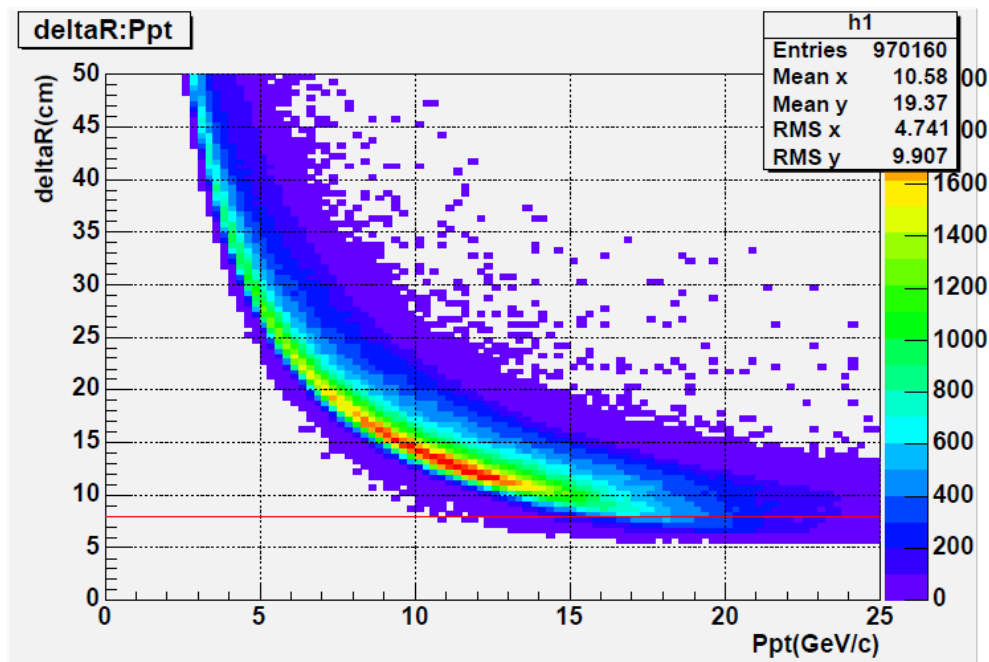
... and loses sometimes



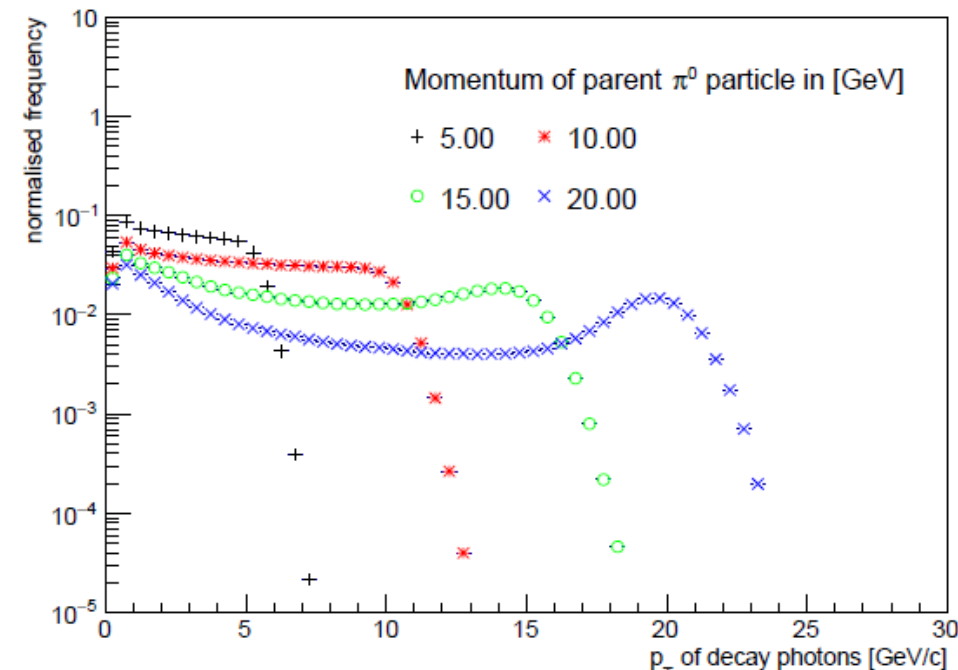
Zooming in on the π^0 problem (at high p_T)



Do you think π^0 decays to two ordinary photons?
WRONG! It decays to two correlated photons



PHENIX – distance between two clusters
photons from a π^0 decay



The distribution of measured **decay photons**
should be **flat** up to the kinematic limit.
Instead, you see depletion at medium, **enhancement**
at high p_T . Unless you fully understand this, your
direct photons at high p_T will be strongly overestimated!

Possible clue to E706?



Basic techniques of photon measurements: external conversion

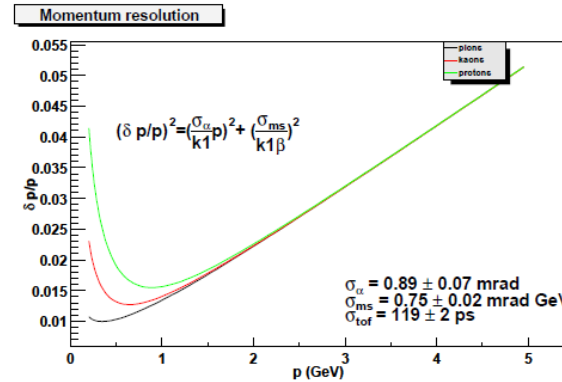


Wenqing Fan, SBU, PhD thesis

Real photon converts to e^+e^- on detector material

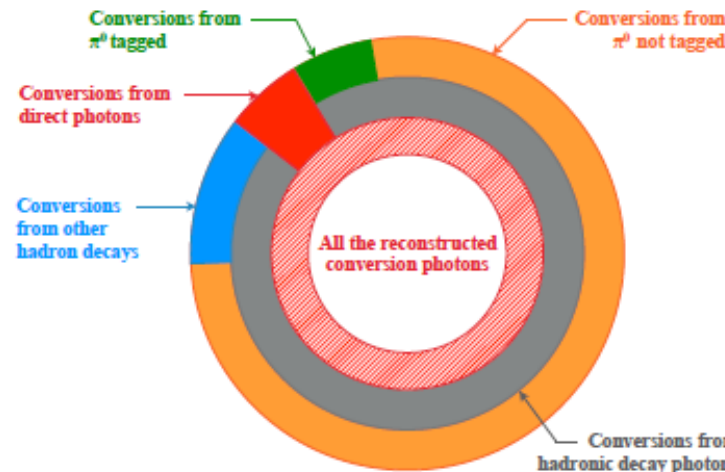
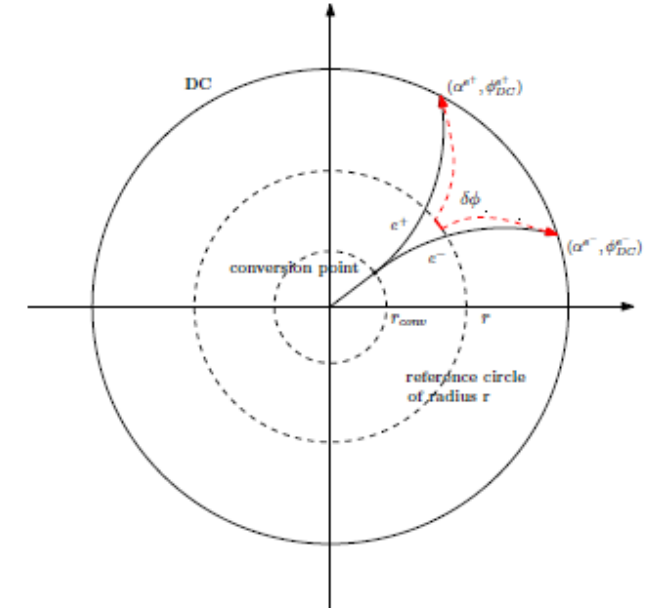
Pro:

- tracking: higher resolution
- viable at low p_T
- not sensitive to Dalitz (origin!)
- feasible in “crowded” detectors



Con

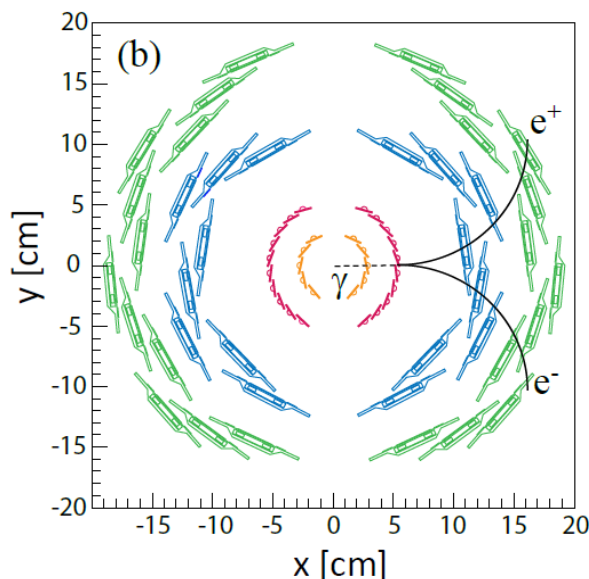
- electron ID difficult
- small acceptance*efficiency
- sensitive to material location
- cocktail (discuss later)



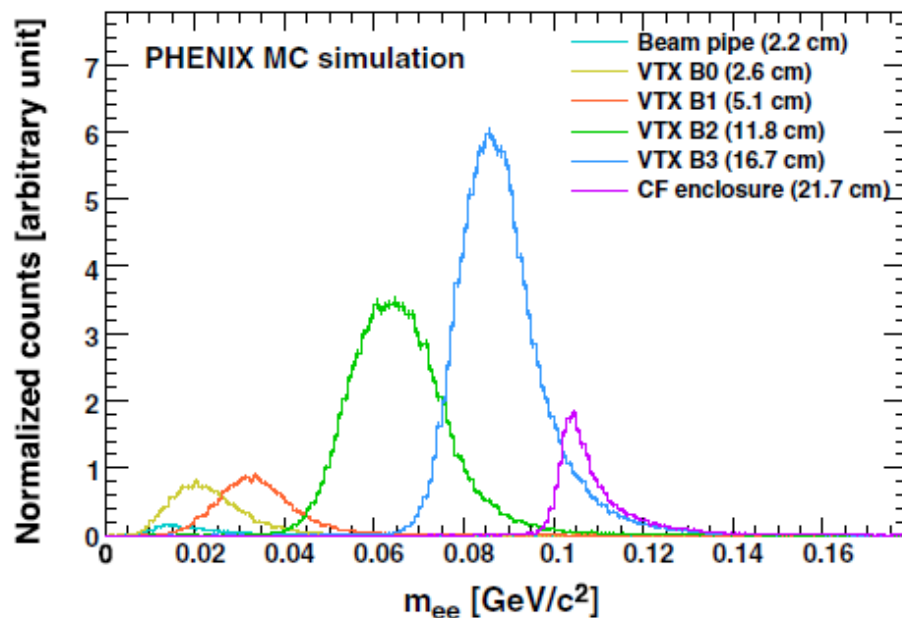
Method of choice at low p_T



Real photon converts to e^+e^- on detector material



VTX detector



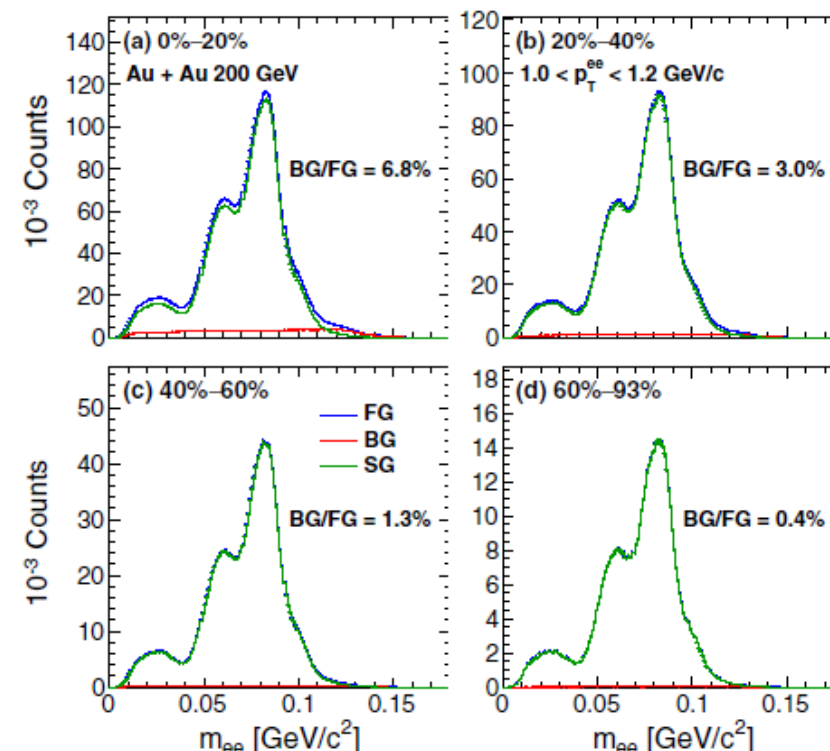
Simulation

Reconstructed conv. photon “mass”
(DC tracking assumes collision vtx!)

Photon purity: ~99%

But is it a direct photon?

$1.0 < p_T < 1.2$ GeV/c



Data

Reconstructed conv.
photon “mass”

Viable at low p_T – complements calorimetry



Tagging photons from hadron decay



Majority of photon are from hadron decays

On the face of it simple:

- pair up conversion photons with real photons in the calorimeter
- tag those that reconstruct a π^0

But:

- tremendous combinatorial at low p_T
- limited acceptance, efficiency
- ignores double conversions

Delicate balance:

- increasing material increases conversion (i.e. our signal)
- too much material promotes double conversions

$M_{ee\gamma}$

Conv + real
photon
invariant mass

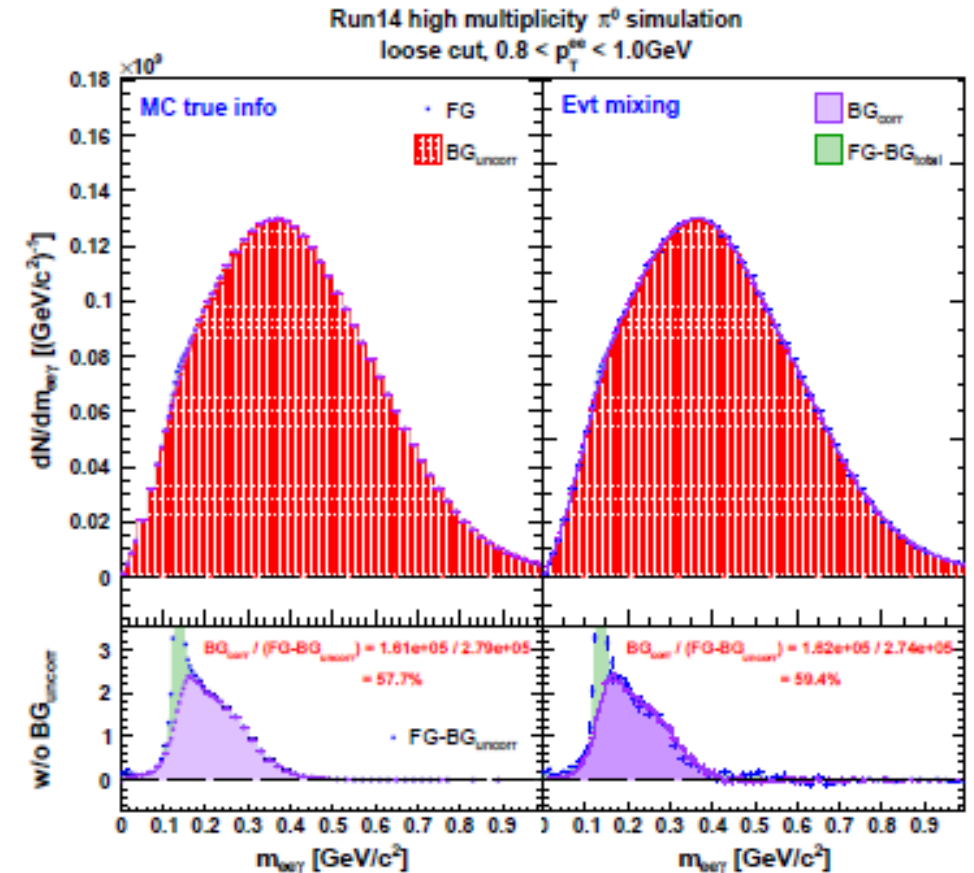


Figure 3.26: Invariant mass distributions of $e^+e^-\gamma$ pairs.

Large combinatorial



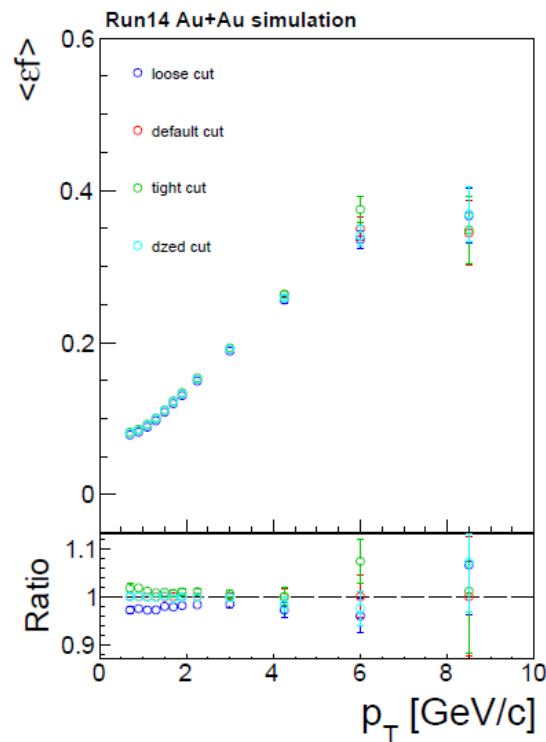


Conditional probability of tagging, R_γ

How often is the partner π^0 photon found in the calorimeter?

$$\left(\frac{N_\gamma^{\text{incl}}}{N_\gamma^{\pi^0, \text{tag}}} \right)_{\text{Data}} = \frac{\gamma^{\text{incl}} \cdot \cancel{p_{e\text{conv}}} \cdot \cancel{a_{e^+e^-}} \cdot \cancel{\epsilon_{e^+e^-}}}{\gamma^{\pi^0} \cdot \cancel{p_{e\text{conv}}} \cdot \cancel{a_{e^+e^-}} \cdot \cancel{\epsilon_{e^+e^-}} \cdot \langle \epsilon_\gamma f \rangle} = \frac{\gamma^{\text{incl}}}{\gamma^{\pi^0} \cdot \langle \epsilon_\gamma f \rangle}$$

Many systematics cancel



After accounting for other hadrons (cocktail, to be discussed later)

$$R_\gamma = \frac{N_\gamma^{\text{incl}}}{N_\gamma^{\text{hadr}}} = \frac{\langle \epsilon f \rangle \times \left(\frac{N_\gamma}{N_\gamma^{\pi^0 \text{tag}}} \right)^{\text{Data}}}{\left(\frac{N_\gamma^{\text{hadr}}}{N_\gamma^{\pi^0}} \right)^{\text{MC}}}$$

Powerful for R_γ in terms of uncertainties





Virtual photon (“internal conversion”) method

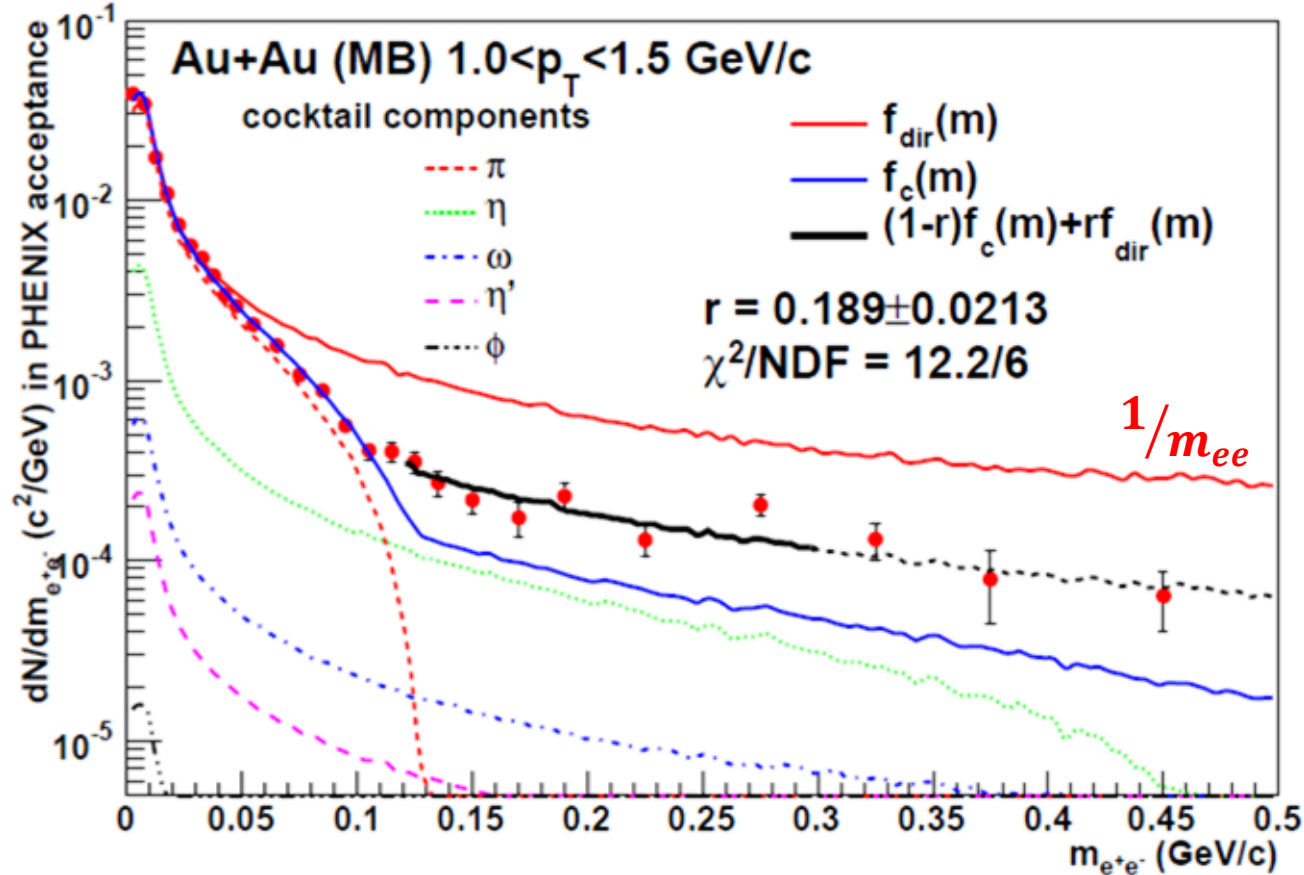
Kroll-Wada

$$\frac{d^2 n_{ee}}{dm_{ee}} = \frac{2\alpha}{3\pi} \frac{1}{m_{ee}} \sqrt{1 - \frac{4m_e^2}{m_{ee}^2}} \left(1 + \frac{2m_e^2}{m_{ee}^2}\right) S dn_\gamma$$

$$\gamma_{dir}: \quad S \sim 1 \text{ for } m_{ee} \ll p_T$$

$$\gamma_{had}: \quad S = \left(1 - \frac{m_{ee}^2}{M_{had}^2}\right)^3 |F(m_{ee}^2)|^2$$

PHENIX: Phys. Rev. Lett. 104 (2010) 132301



- Using virtual photons $\gamma^* \rightarrow e^+ e^-$:
 - any process that radiates γ will also radiate $\gamma^* \rightarrow e^+ e^-$
 - for $m_{ee} \ll p_T$ extrapolate γ^* to $m_{ee} = 0$
 - $m_{ee} > m_\pi$ cut improves S/B by factor 10
 - sys. uncertainty cancelation in ratio $\gamma_{dir}^* / \gamma_{incl}^*$

Works above 1 GeV/c
Hadronic “cocktail” important

Direct γ^* yield fitted in range 120 to 300 MeV
 Insensitive to π^0 yield





Direct photons in heavy ion collisions

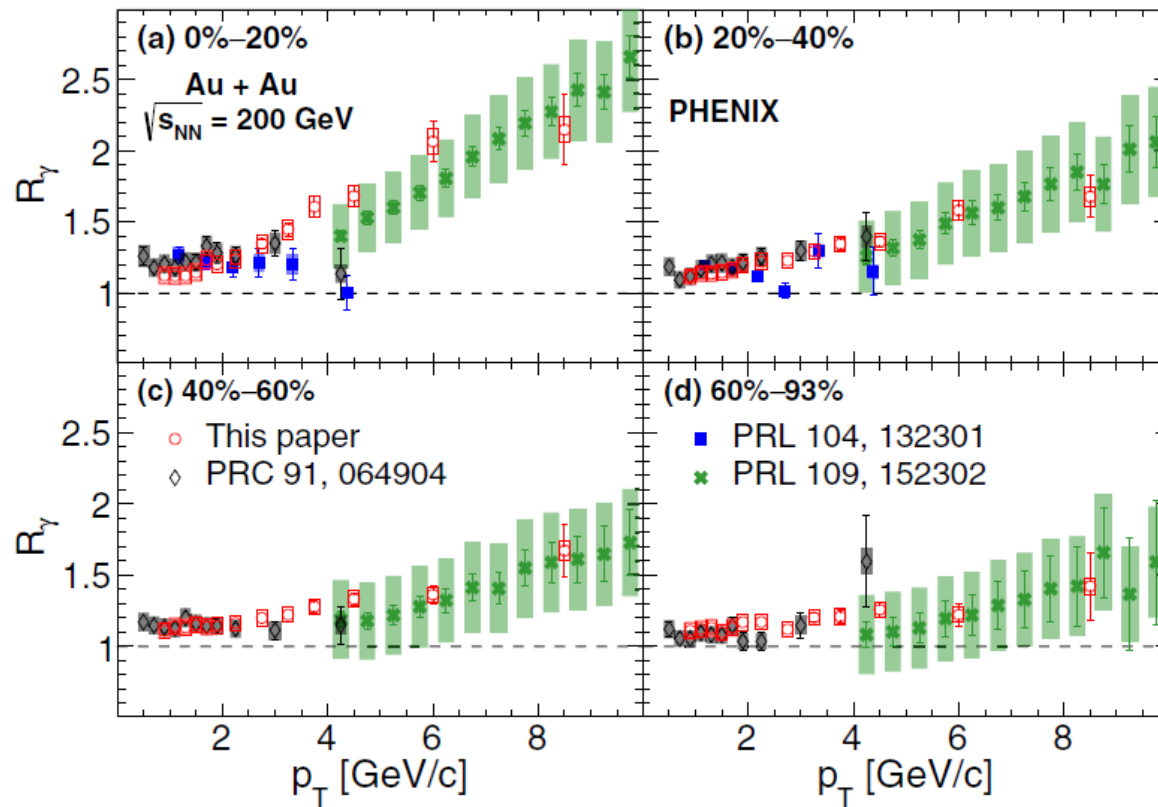


Direct photons with multiple methods (PHENIX, ALICE)

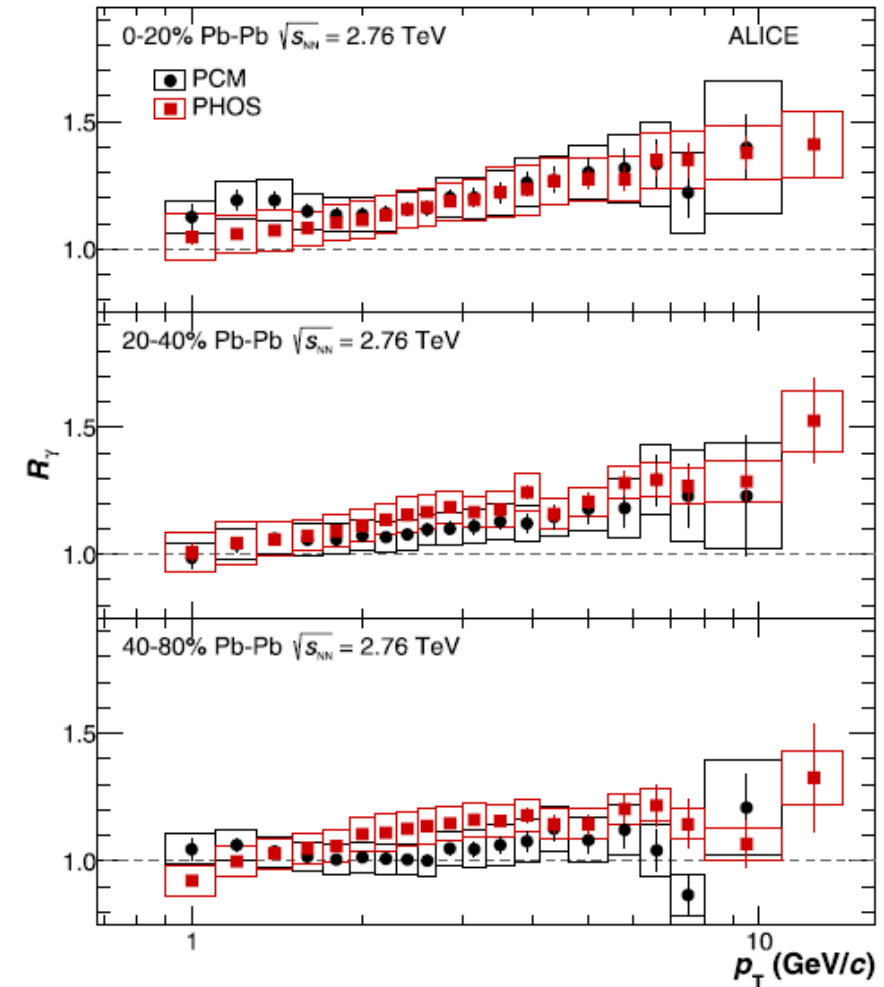


R_γ with different methods

PRC 109, 044912 (2024) (PHENIX)



PLB 754 (2016) 235-248 (ALICE)



$$\gamma^{\text{direct}} = (R_\gamma - 1)\gamma^{\text{hadron}}$$

Remember this, when interpreting spectra

Direct photons and non-prompt photons

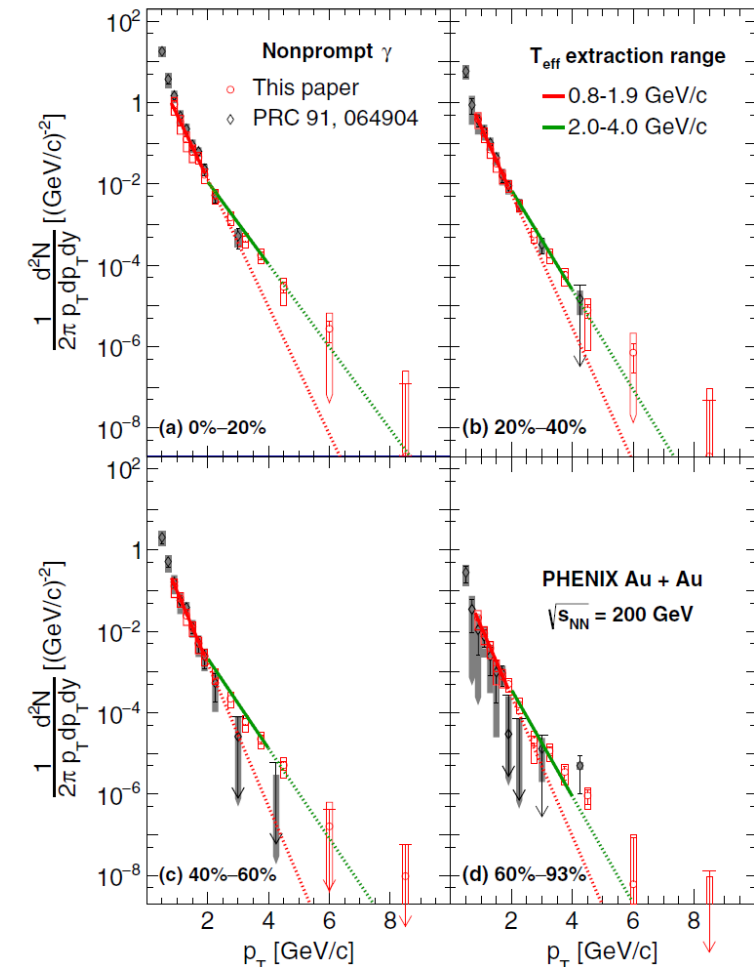
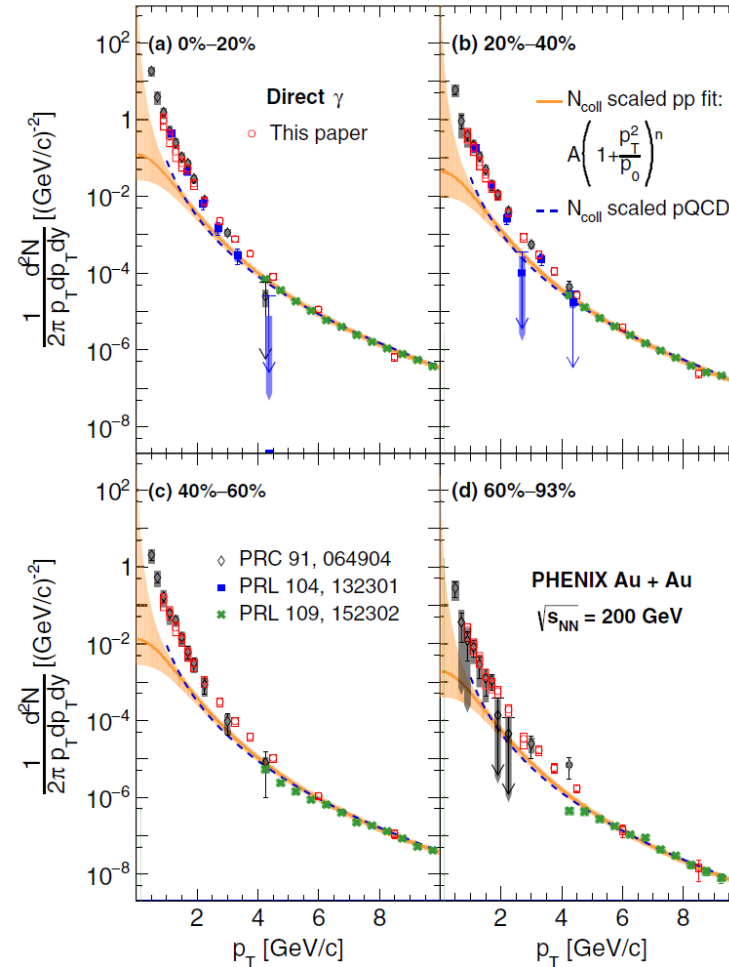
PRC 109, 044912 (2024)

Non-prompt: properly scaled
p+p yields subtracted

$N_{\text{coll}} \rightarrow$ number of binary
nucleon-nucleon collisions
(estimated with Glauber MC)

At lower p_T varying exponential
slopes (T_{eff})

Above 4 GeV/c hard scattering
dominates



Three methods, consistent results



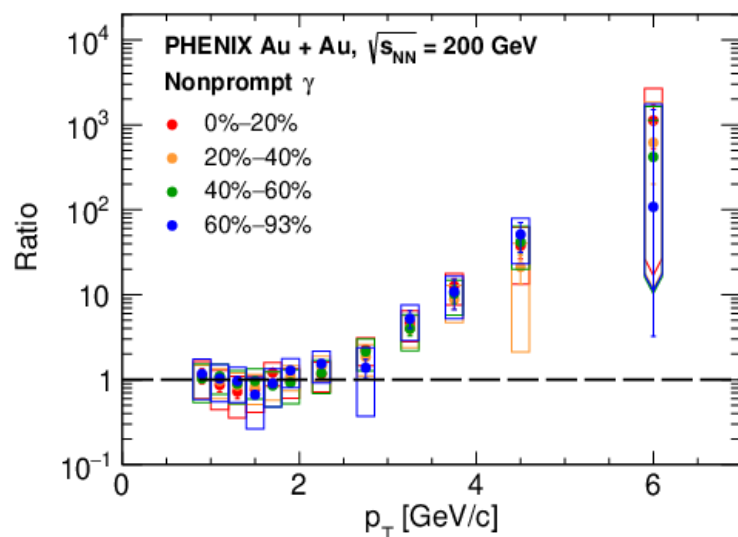
T_{eff} vs the p_T range and centrality



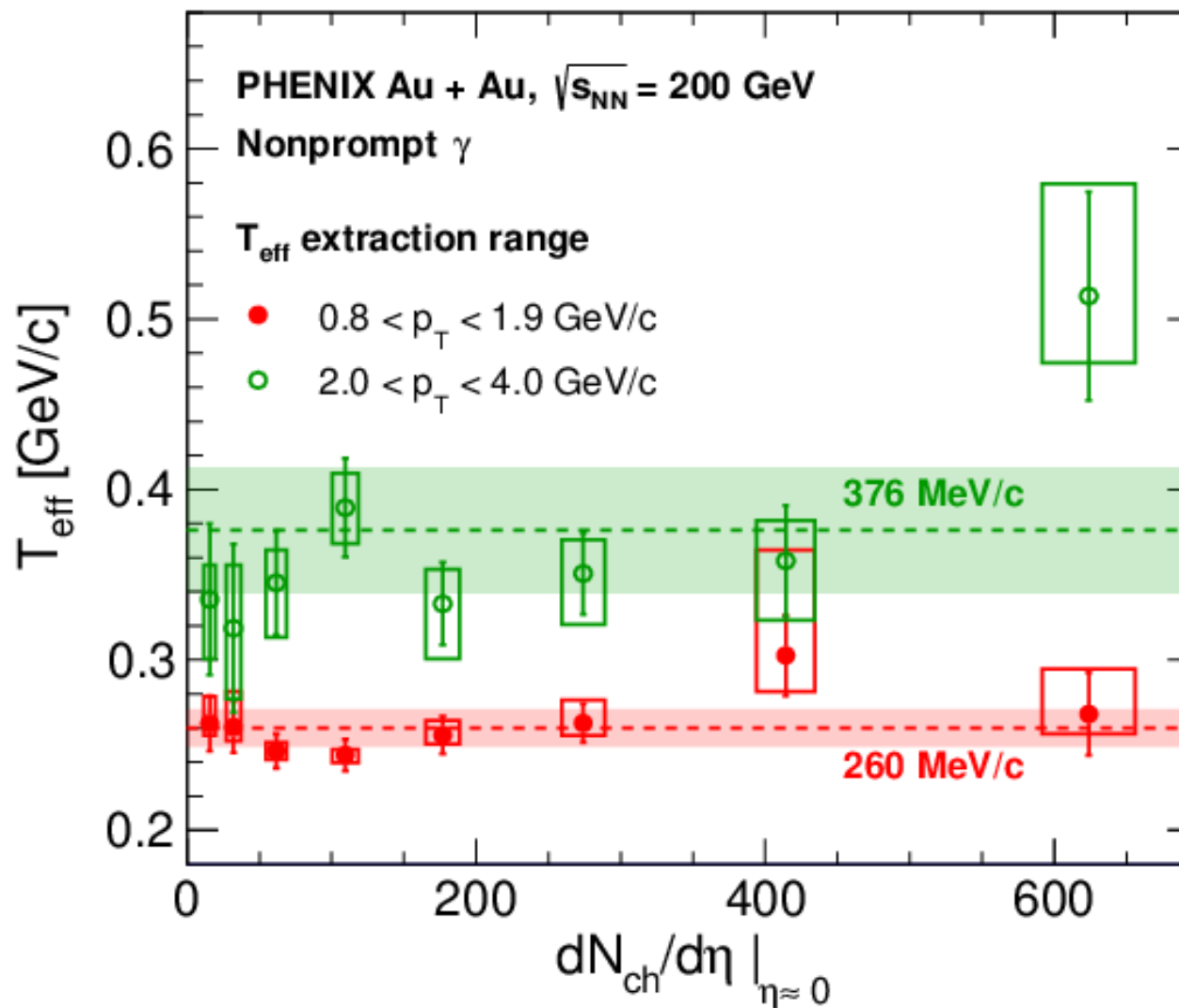
PRC 109, 044912 (2024)

No dependence on centrality

Strong dependence on p_T
Composition and relative weight of different sources always the same?



Ratio of the yield to a fixed exponential ($T_{\text{eff}} = 0.26$ GeV)



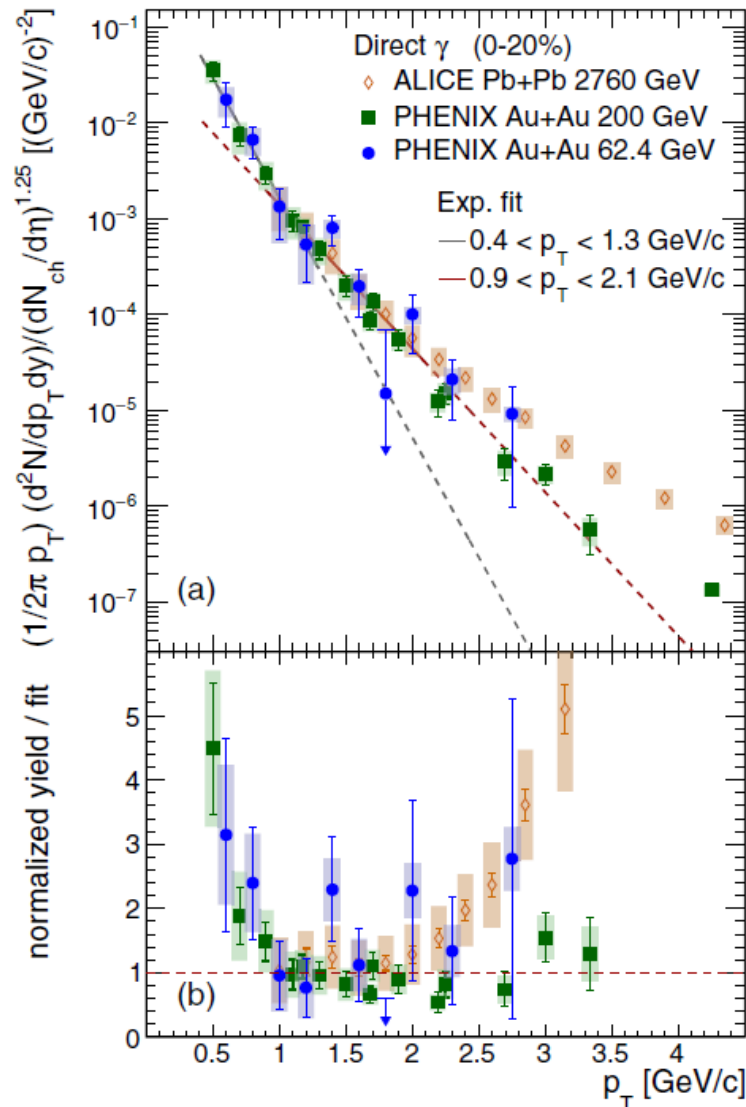
Yields: universal scenario?



Slope vs p_T

Dependence on collision energy

PRC 107, 024914 (2023)



Data vary from 62 to 2760 GeV!

Slopes vary with p_T

→ imprint of T_{eff} evolution?
(different times, different sources, different p_T)

Normalized yields very similar at low p_T

Possible interpretation:

no matter where the QGP started, the endgame
will always be the same
(no surprise, but good to see)

By freeze-out everything is similar

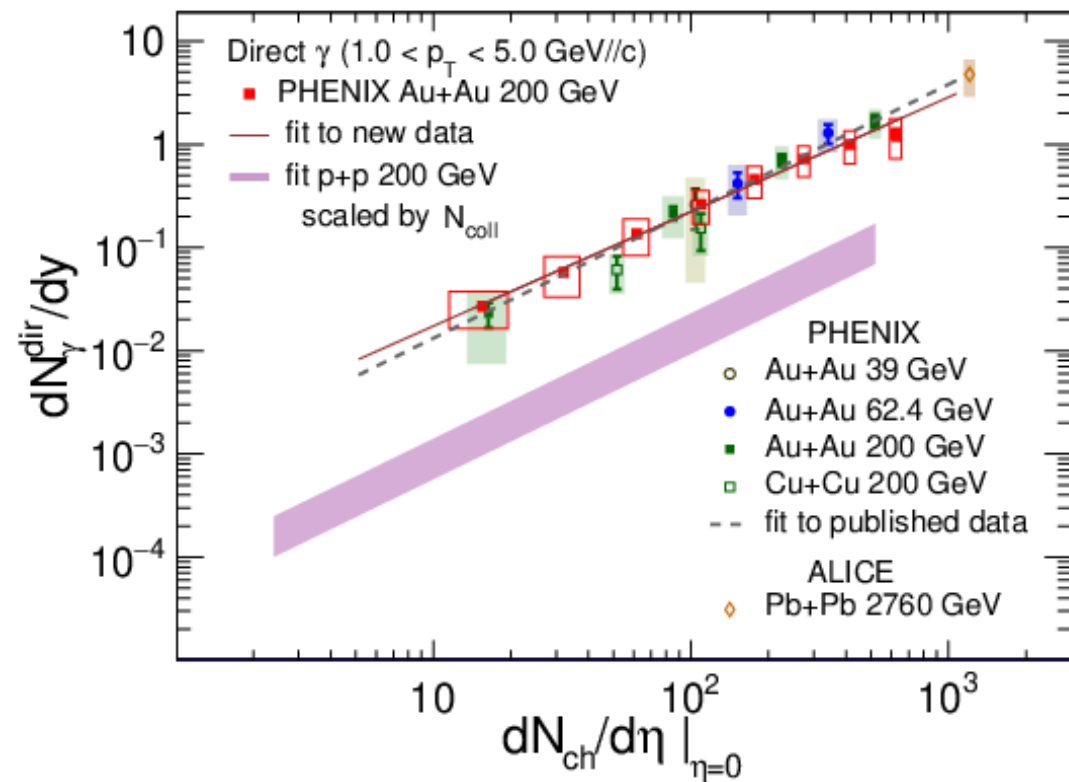
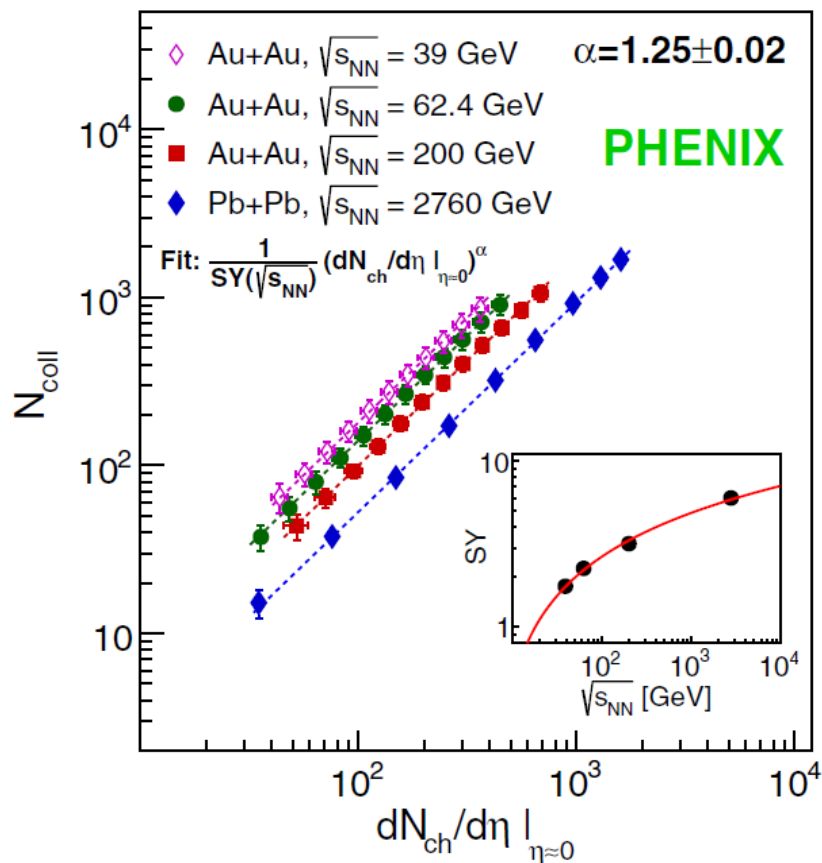


Scaling with N_{ch} (N_{coll})?

In larger systems yields scale with N_{ch} ?

PRC109, 044912 (2024)

PRL 123, 022301 (2019)



Yes, but α (slope) controversial

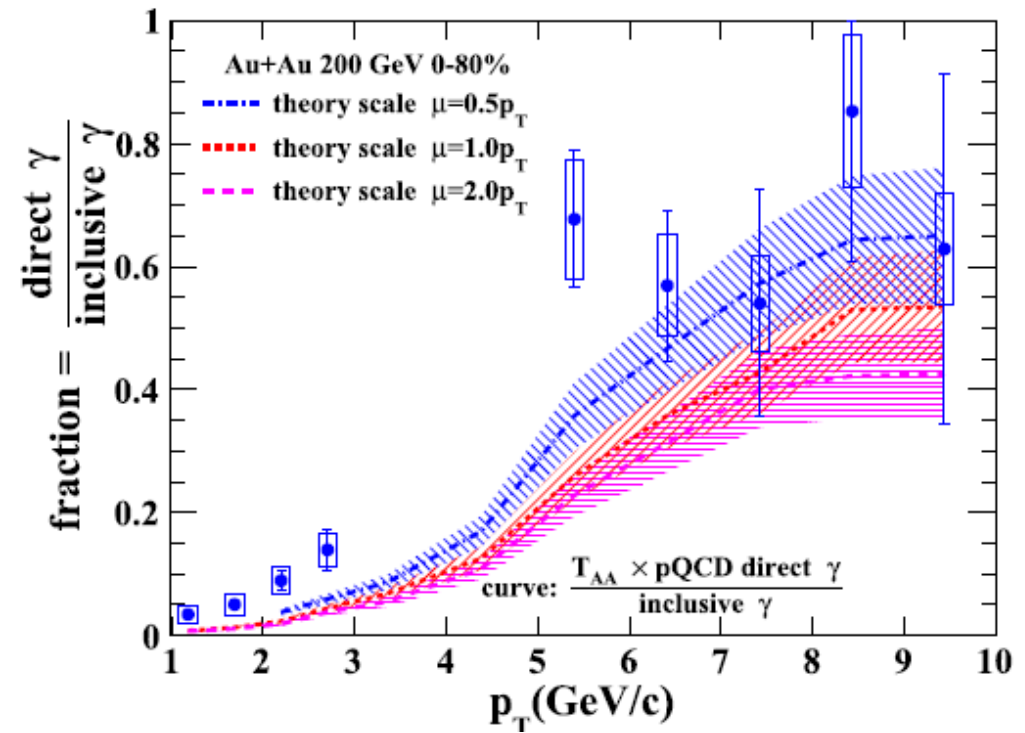
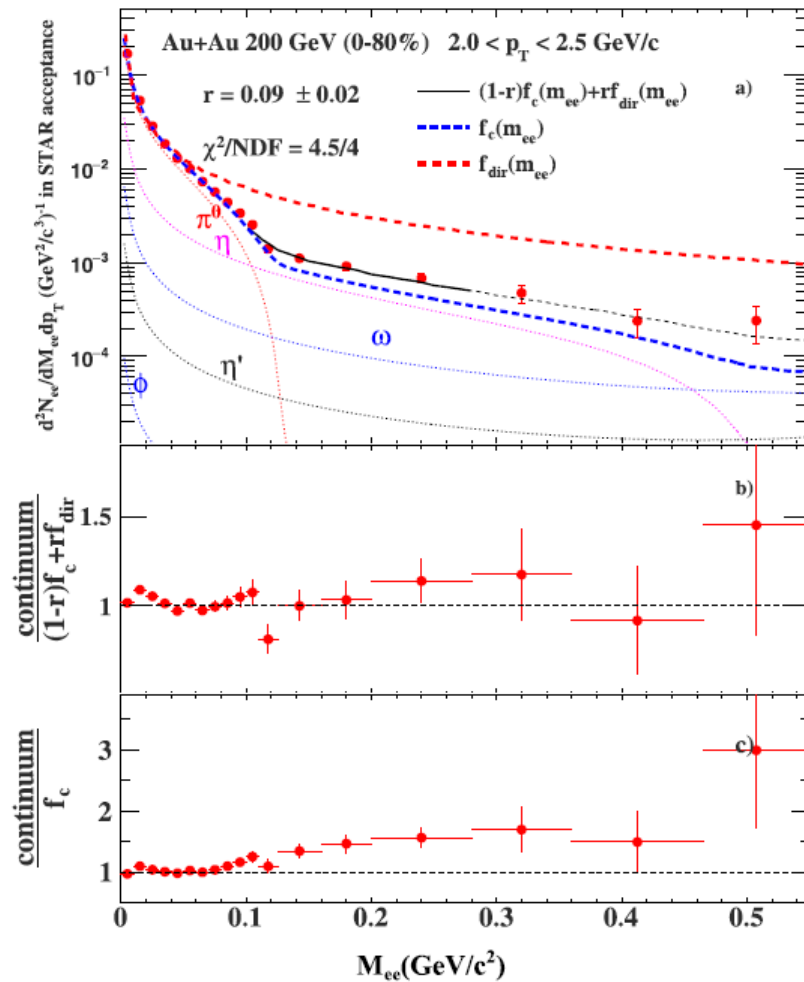


Two methods: internal conversion / calorimeter

PLB 770 (2017) 451-458

(same as the first PHENIX measurement)

fraction $\rightarrow r = (R_\gamma - 1) / R_\gamma$

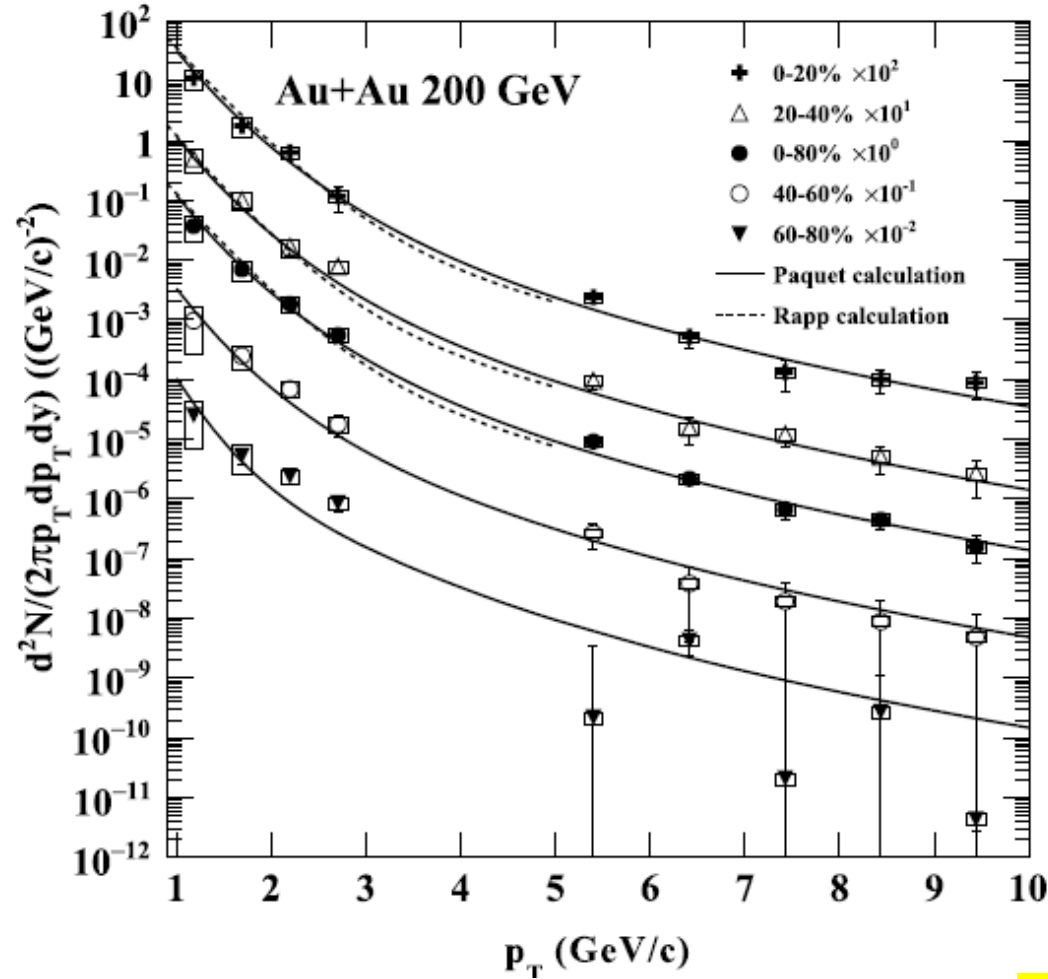


“Tension” with PHENIX, will discuss later

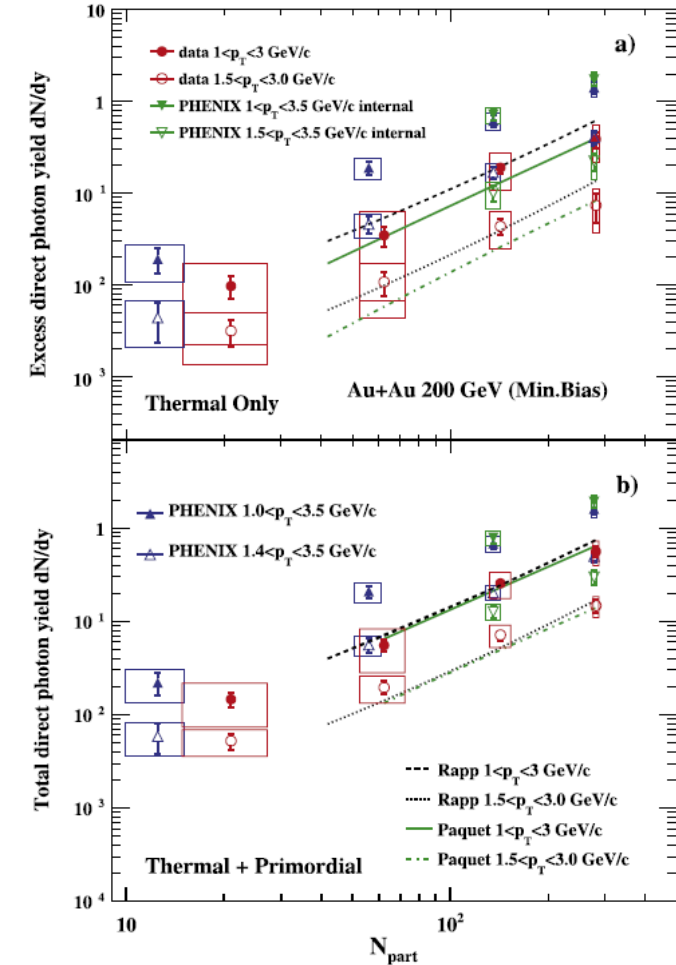


STAR 200 GeV Au+Au

Two methods: internal conversion / calorimeter



PLB 770 (2017) 451-458



Integrated yields disagree with PHENIX





ALICE 2.76 TeV PbPb -- updated

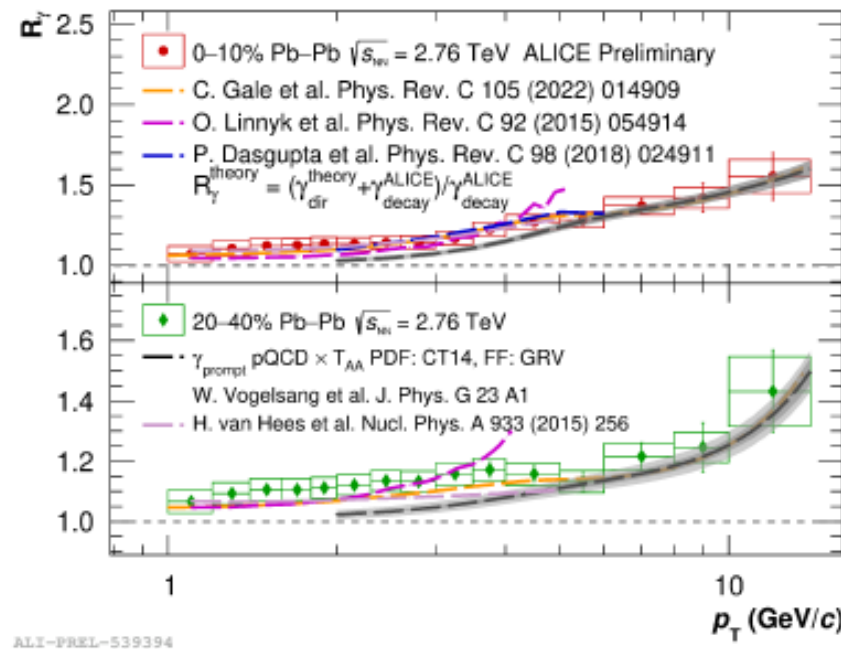
Improved significance

PoS (HardProbes2023) 061

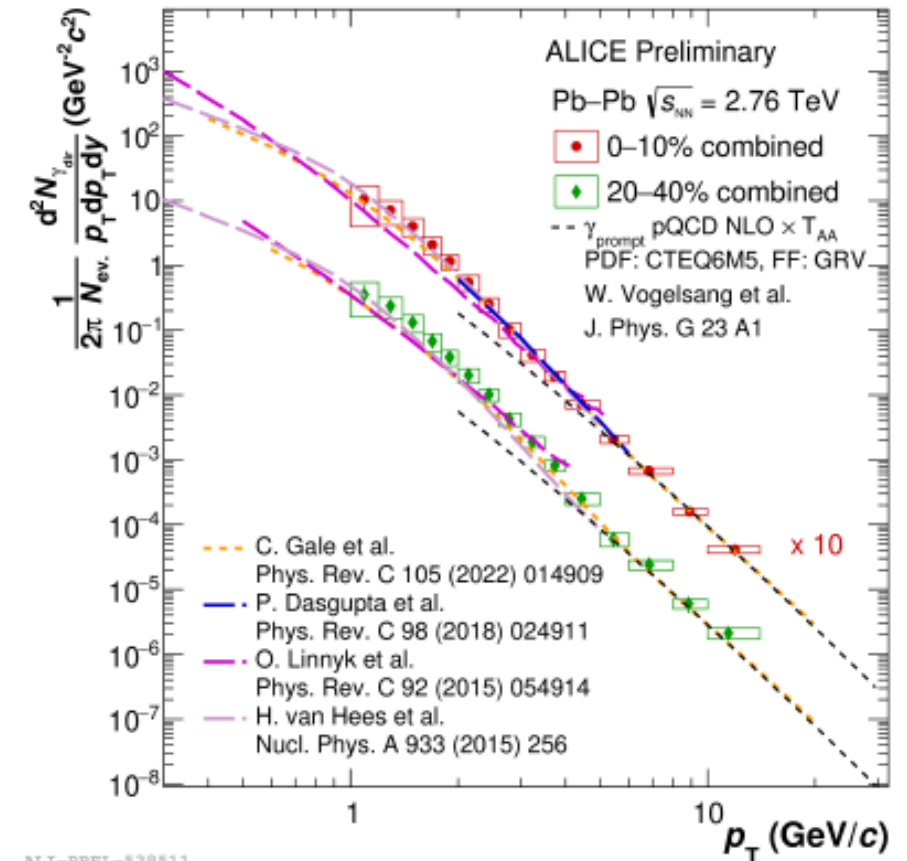
First publication: PLB 754 (2016) 235-248

Direct photon signal
down to 1 GeV/c
but quite weak

$T_{\text{eff}} \sim 340$ MeV for
both centralities
(1.1-2.1 GeV/c region)



ALI-PREL-539394



ALI-PREL-538511

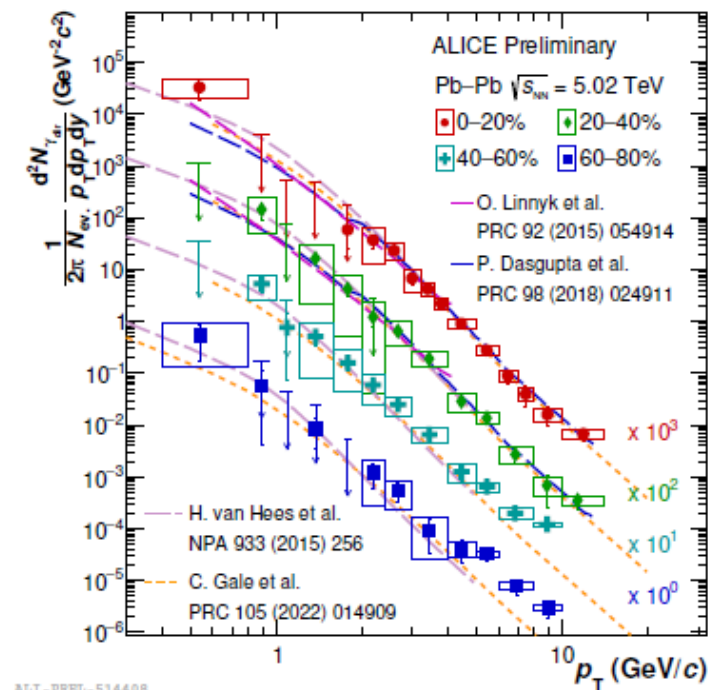
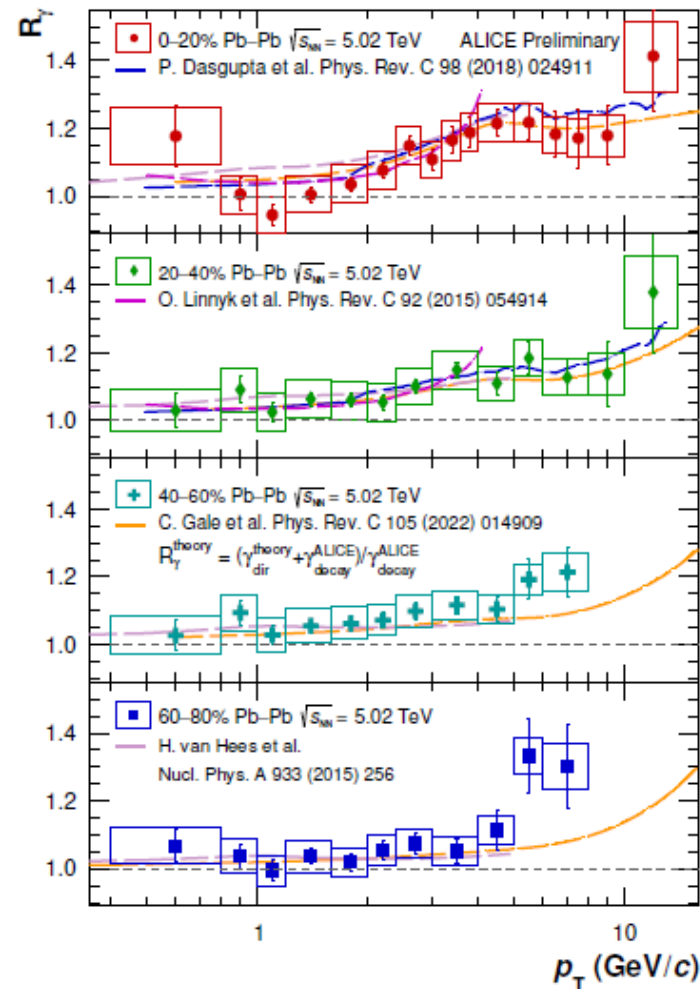
“Thermal” radiation not exceeding PHENIX





Increase energy

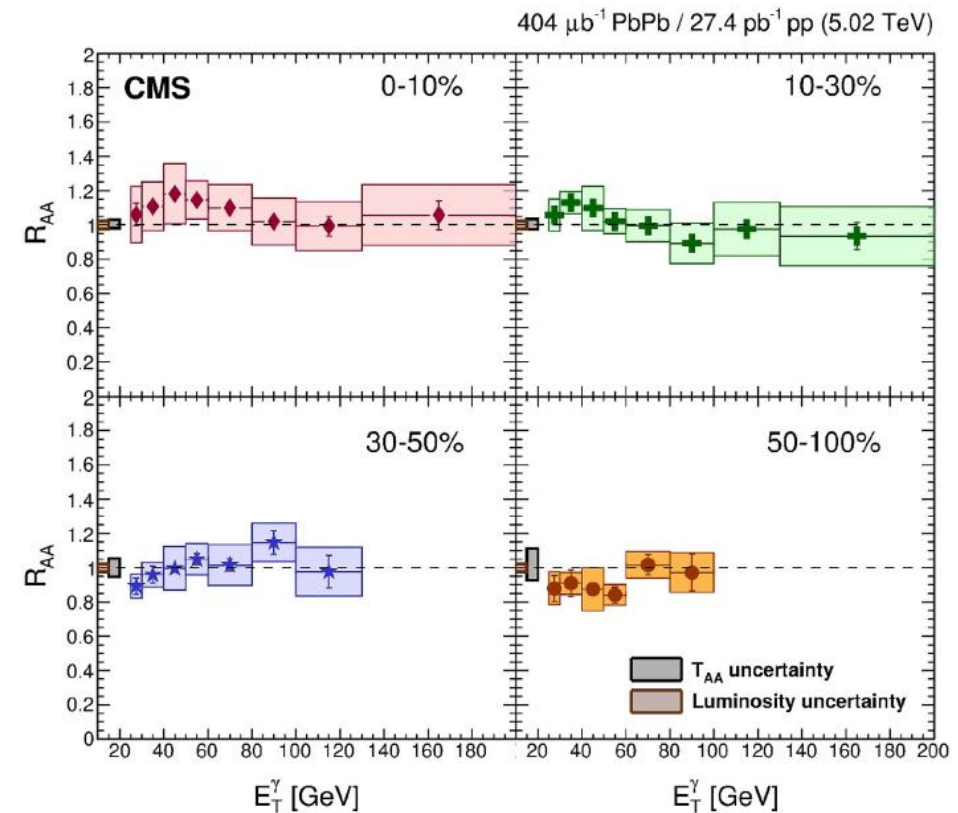
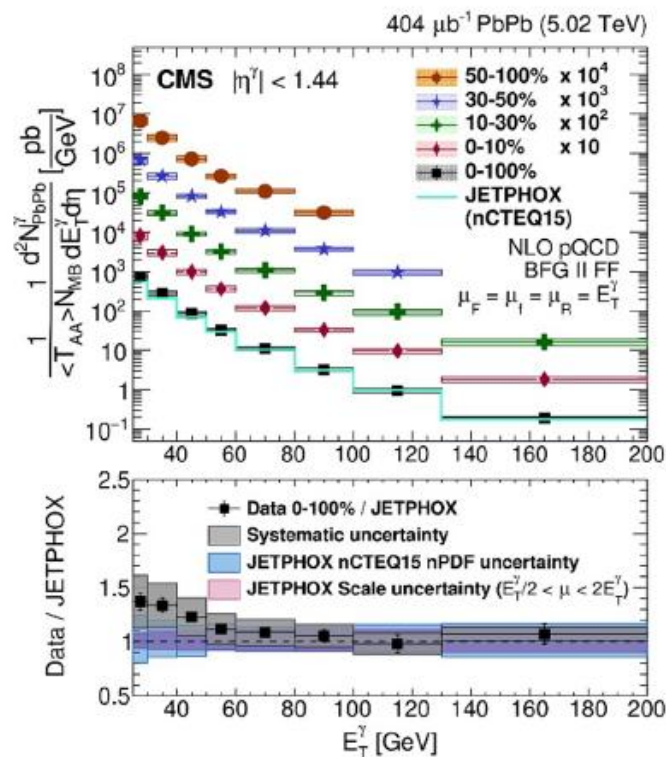
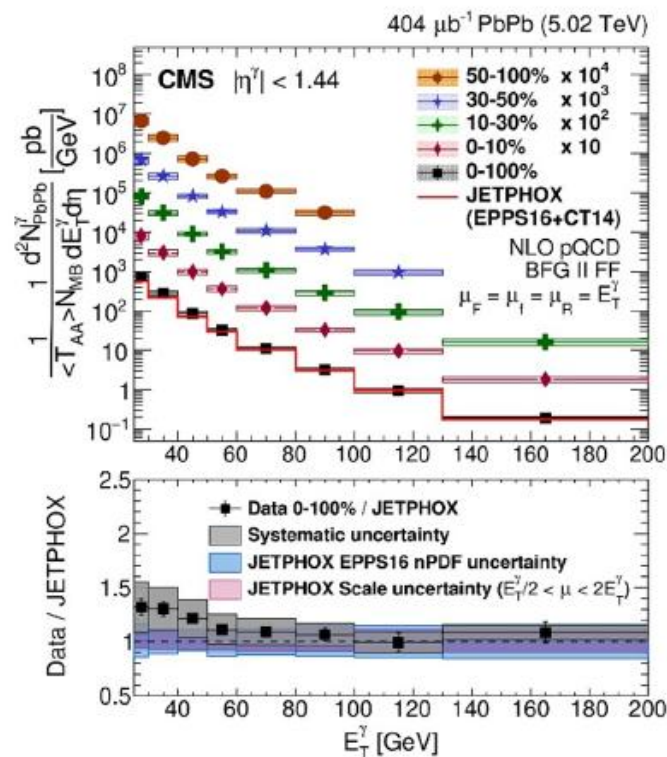
No significant
direct photon signal
at low p_T



Thermal signal fading out???



Does high p_T still look as expected?



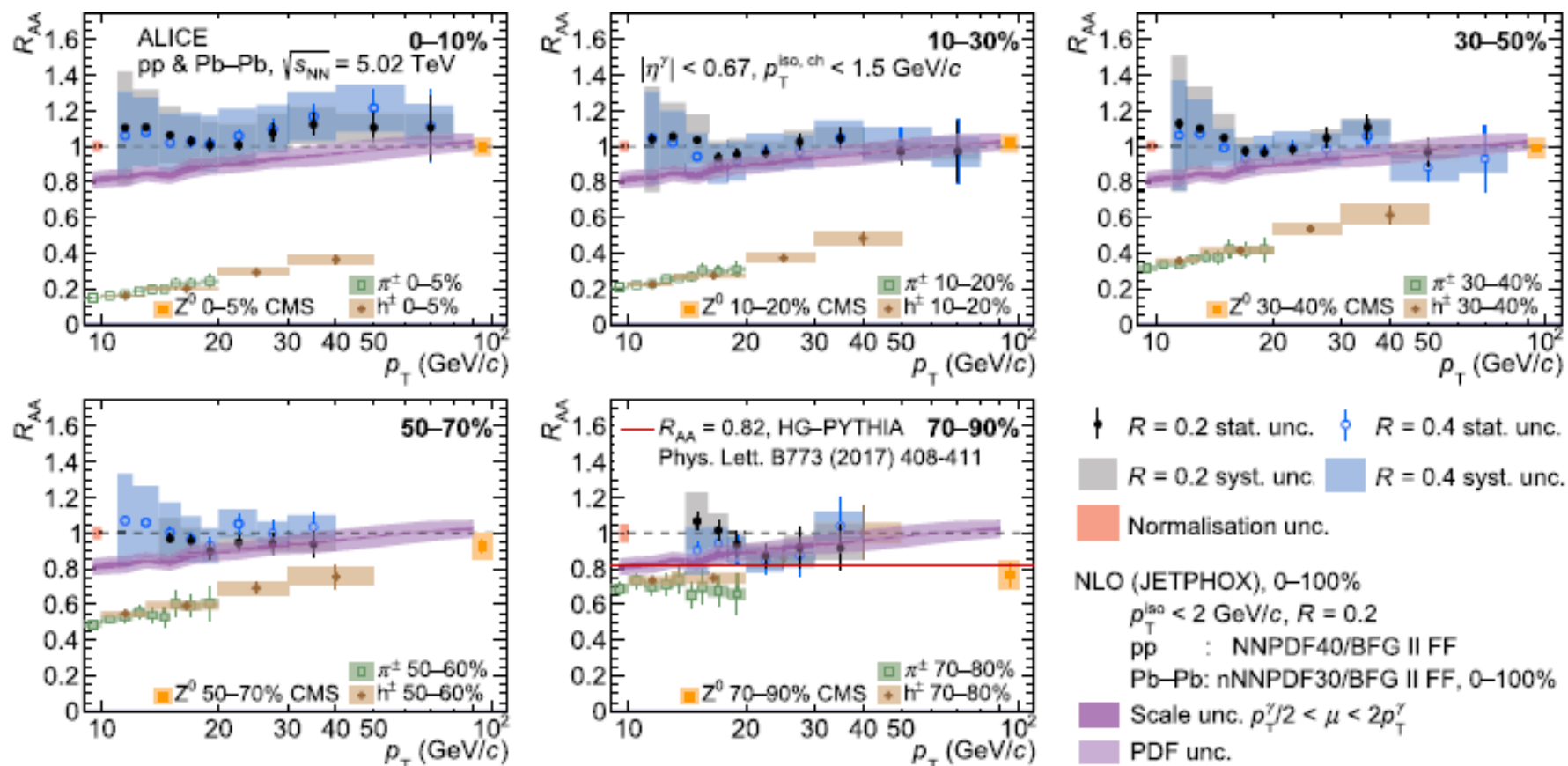
Good news: no nuclear modification at high p_T



R_{AA} of isolated photons, ALICE PbPb 5.02 TeV

Any nuclear modification?

Eur. Phys. J. C (2025) 85:553



$R_{AA} = 1 \rightarrow$ isolated photons scale with N_{coll}





The centrality, R_{AA} and photon saga





What's the big deal about $R_{AA} = 1$?

Connecting impact parameter to observables

PRC 90, 034902 (2014)

Ann.Rev.Nucl.Part.Sci.57:205-243,2007

3.1 Methodology

N_{ch}
charged mult
(large η gap)

N_{part}
participating
(wounded) nucleons

N_{coll}
binary (NN)
collisions

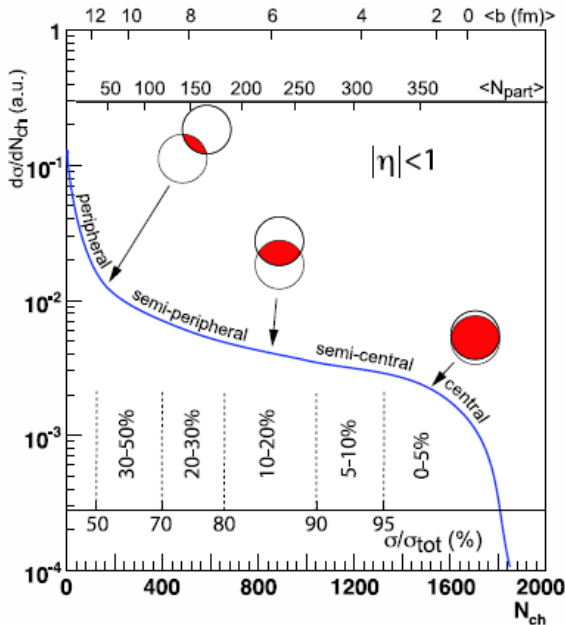
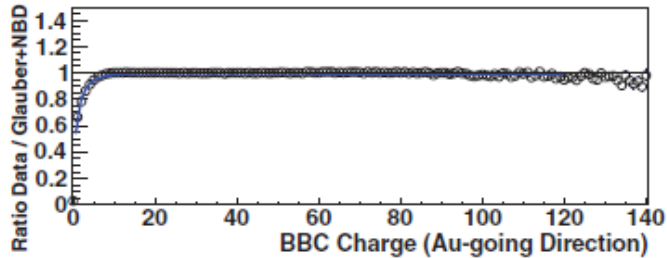
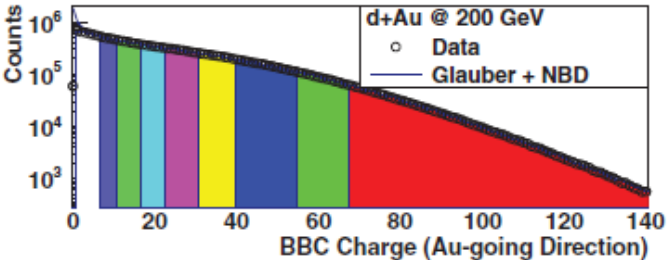
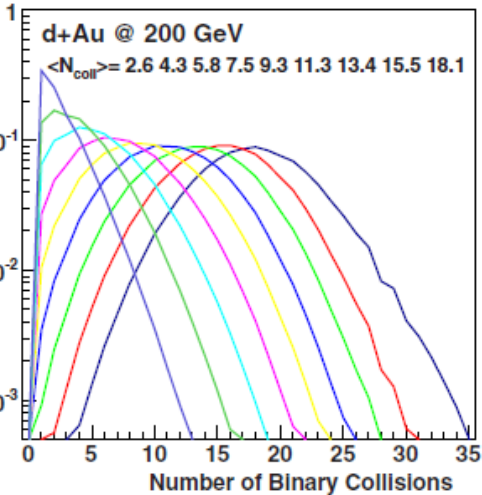


Figure 8: A cartoon example of the correlation of the final state observable N_{ch} with Glauber calculated quantities (b , N_{part}). The plotted distribution and various values are illustrative and not actual measurements (T. Ullrich, private communication).



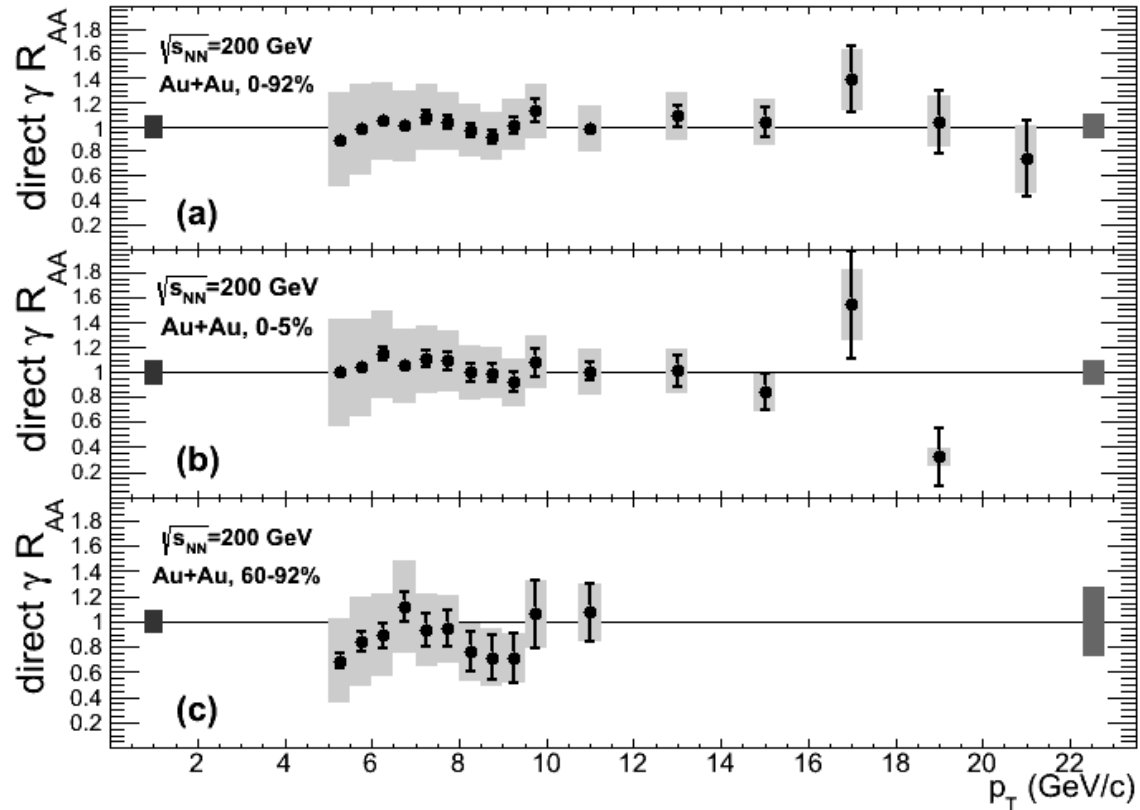
“In heavy ion collisions, we manipulate the fact that the majority of the initial state nucleon-nucleon collisions will be analogous to minimum bias p+p collisions...”

Glauber MC works in A+A – experimental proof



Hard e.m. probes (mostly) immune to FS effects

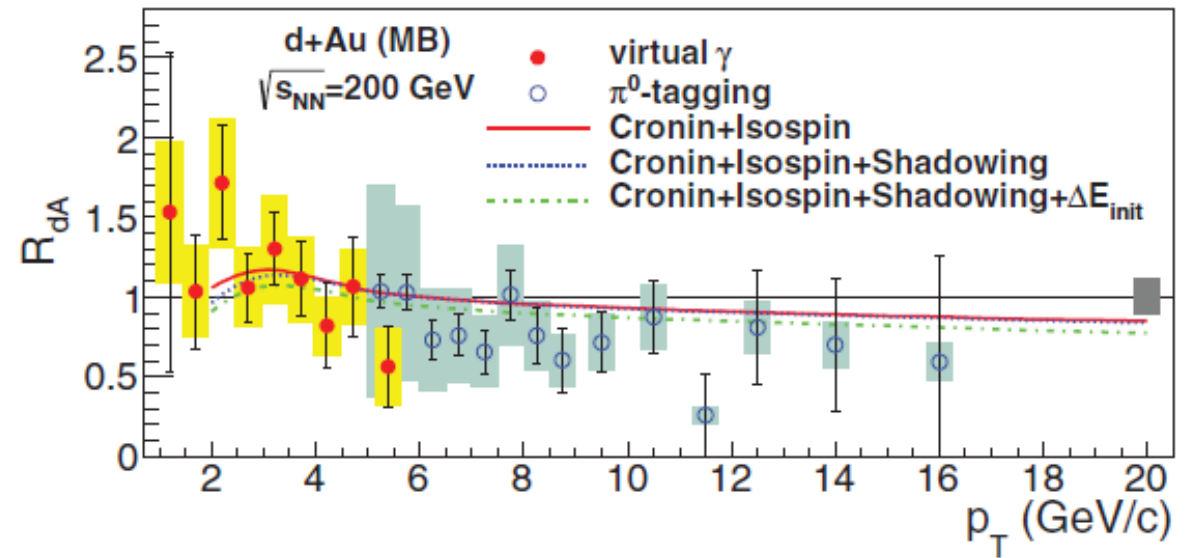
Au+Au



PHENIX, PRL 109, 152302 (2012)

d+Au

PHYSICAL REVIEW C 87, 054907 (2013)



Large uncertainties, but hints of isospin effect

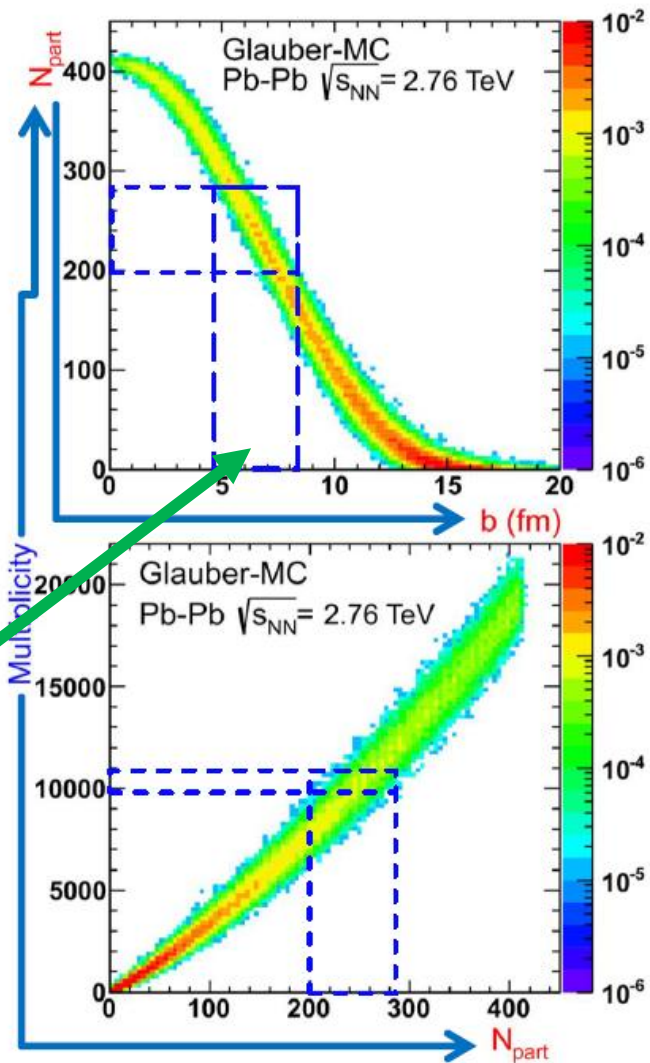
Some caveats (like isospin effect)



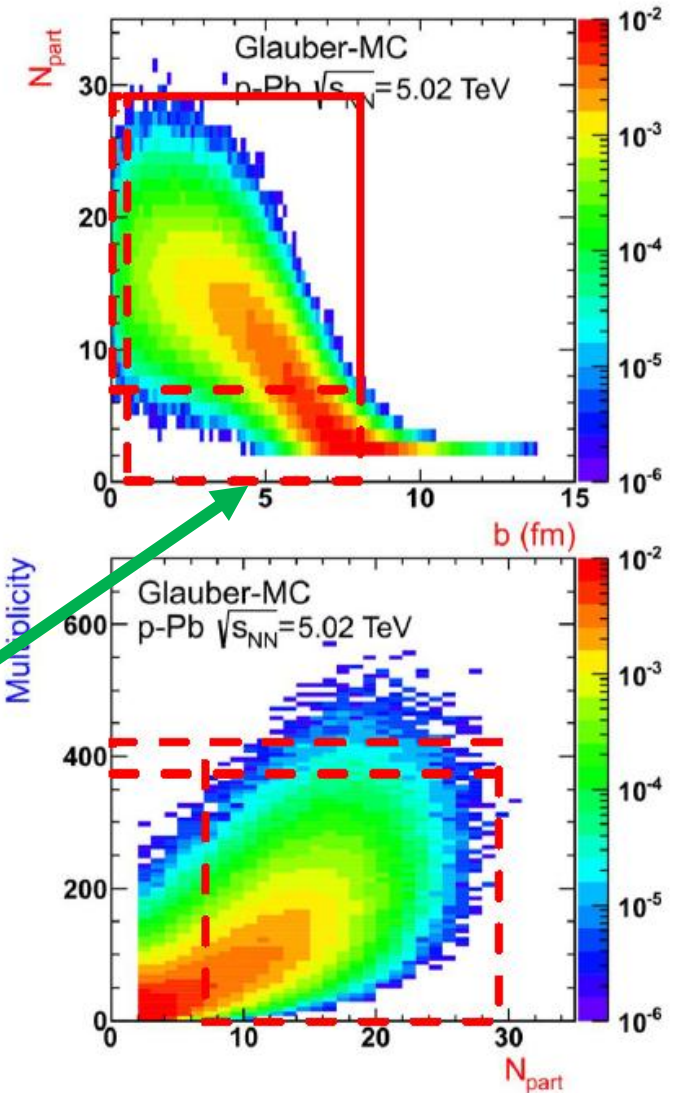
(Un)ambiguity of centrality determination – PbPb, pPb



Connecting b (th) $\rightarrow N_{\text{part}}$ (th) $\rightarrow N_{\text{ch}}$ (exp)



N_{ch} , N_{part} , b correlation
tight in PbPb,
very loose in pPb



Ambiguity in
small systems

Also: centrality from bulk observables – but where?

Signal: $\eta \sim 0$, “bulk”: $|\eta| \gg 0$

ALICE PRC 91 (2015) 064905

CL1 $\rightarrow |\eta| < 0.9$

V0A $\rightarrow 2.8 < \eta < 5.1$ (Pb-going side)

V0C $\rightarrow -3.7 < \eta < -1.7$ (p-going side)

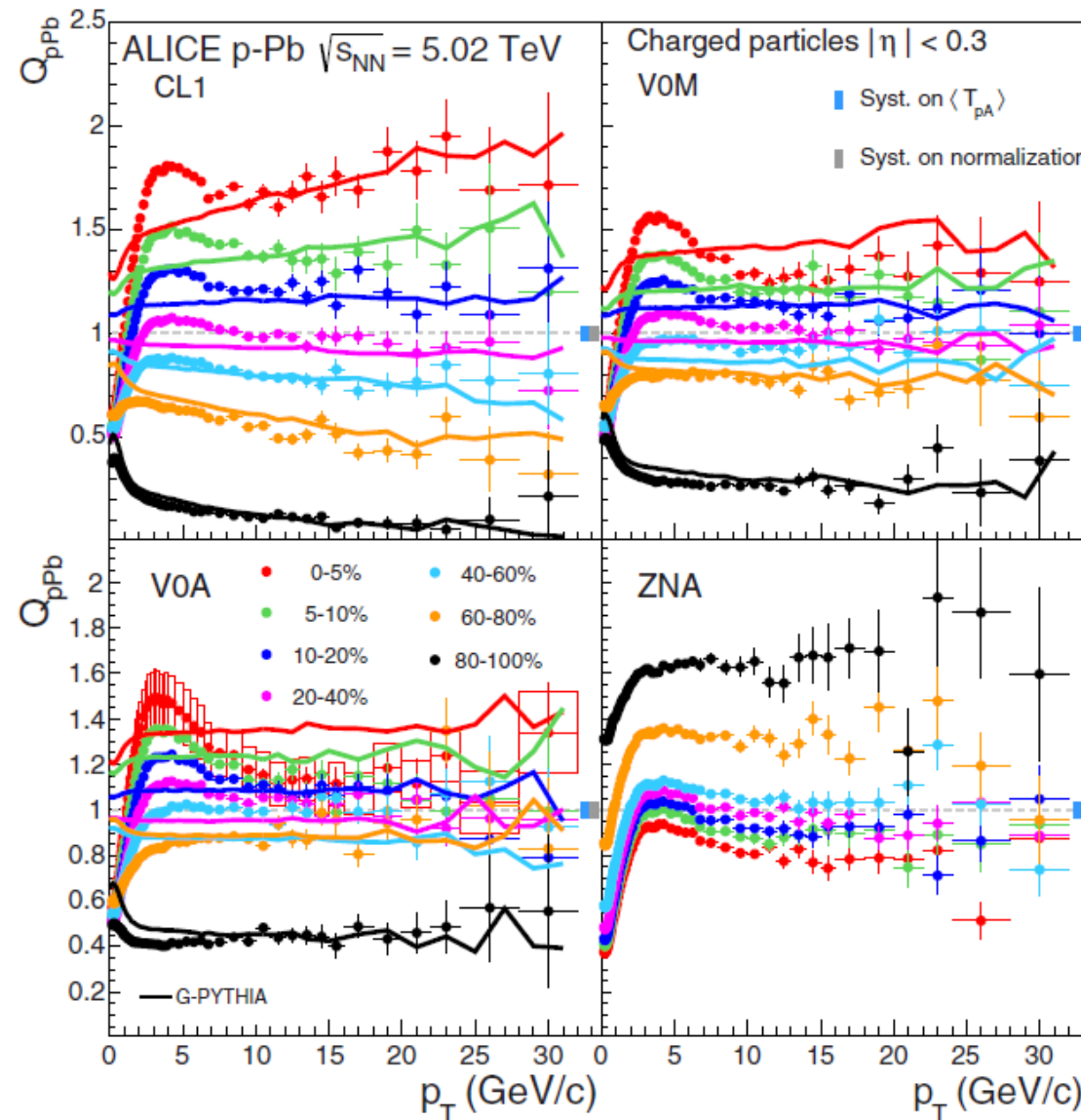
V0M \rightarrow V0A + V0C

ZNA \rightarrow ZDC on Pb-going side

Watch strong auto-correlation
in CL1 central, jet veto bias in
peripherals

Smaller fluctuations in V0A,
mostly around unity, except
vastly displaced peripheral due to
multiplicity bias (?)

Reverse ordering for ZNA, as
expected (as expected???)



How about energy conservation?

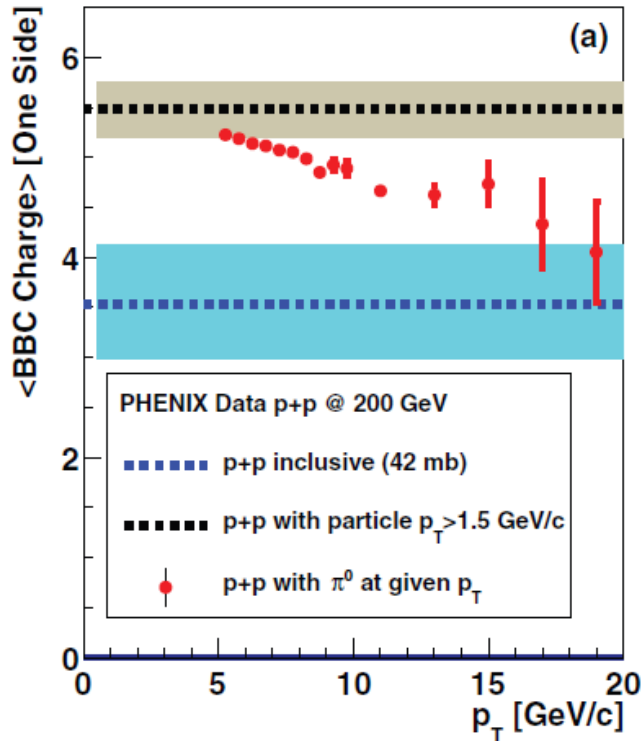


Color transparency or energy conservation in R_{xA} ?

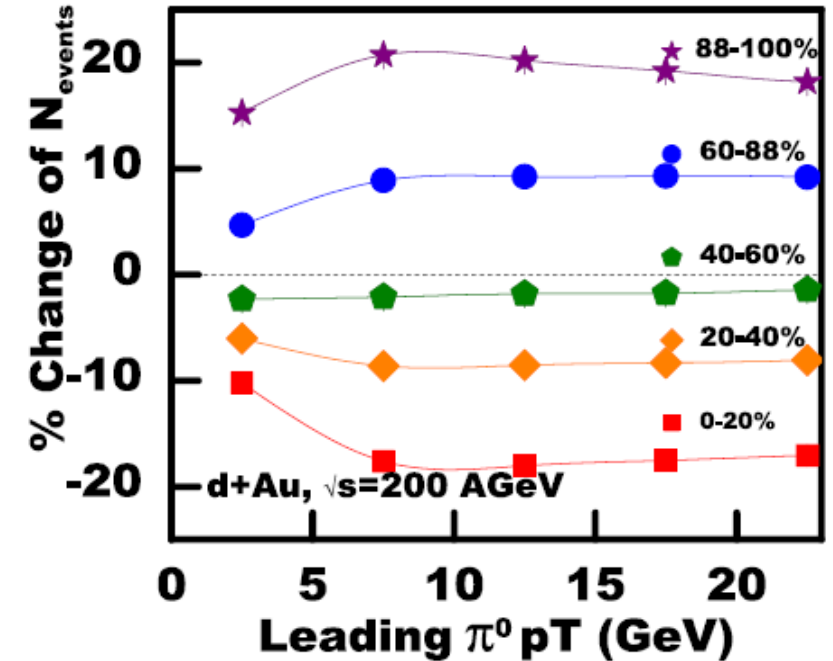
Max p_T at $\eta \sim 0$, N_{ch} at $|\eta| \gg 0$ in p+p

Kordell, Majumder PRC 97, 054907 (2018)

“...the puzzling enhancement in peripheral events ... as well as the suppression seen in central events... are possibly due to *mis*-binning of central and semicentral events, containing a jet, as peripheral events... due to suppression of soft particle production away from the jet, caused by the depletion of energy available in a nucleon of the deuteron in d-Au or proton in p-Pb after the production of a hard jet...”



Initial enhancement of N_{ch} forward with mid-rapidity p_T , then depletion (p+p data)



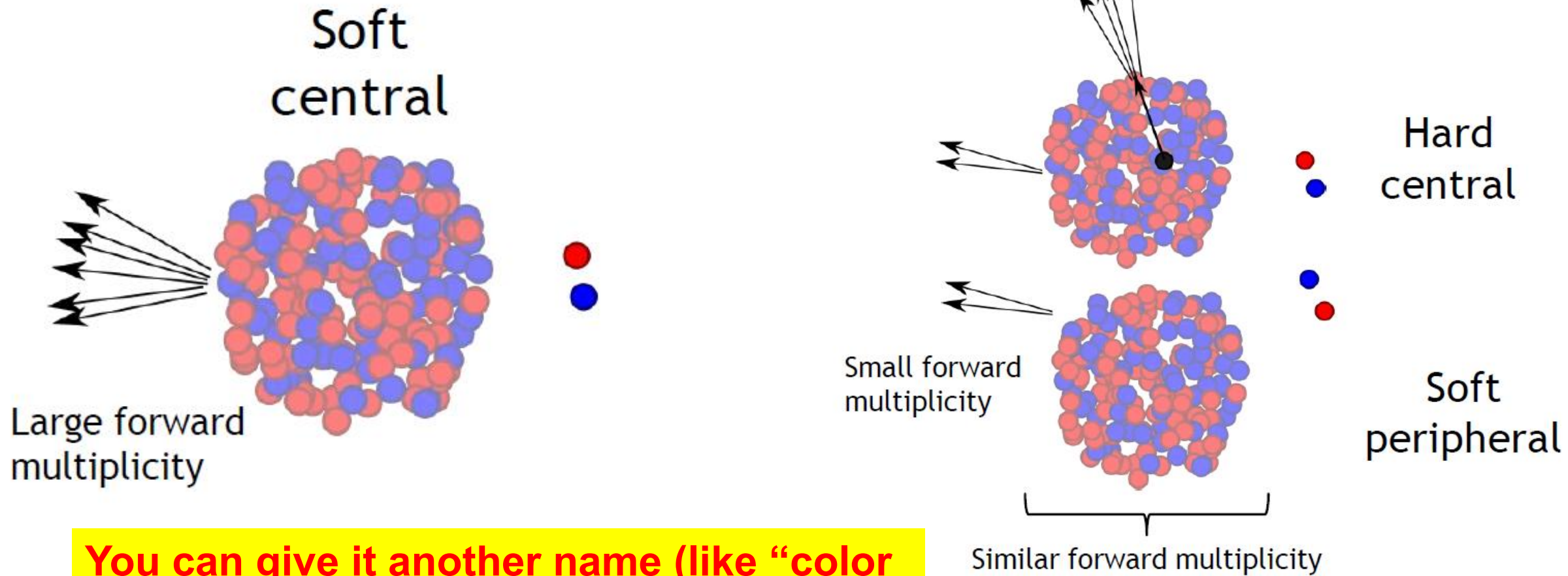
Anticorrelation at sufficiently high p_T

PRC 90, 034902 (2014)

Very large p_T at $\eta = 0$ – mis-binning of centrality?

The same (few) projectile nucleons have to produce both hard scattering and “bulk” forward

This is essentially an energy conservation argument

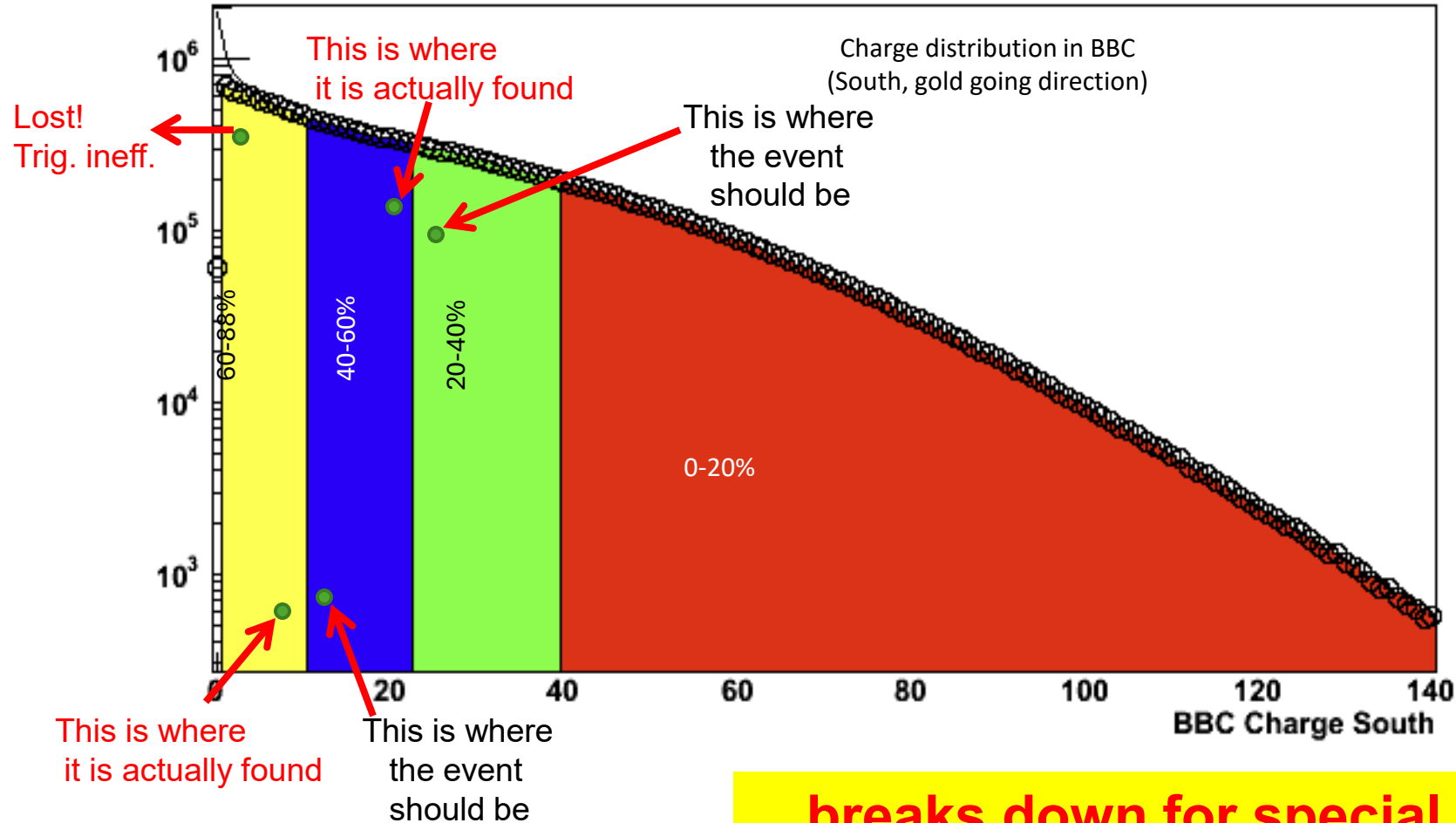


You can give it another name (like “color fluctuation”) but it is still the same thing



Illustration: shift between multiplicity classes

Mapping N_{ch} to b using average events...



...breaks down for special (high p_T) events

If (experimental) centrality is determined with fixed (forward) multiplicity thresholds, irrespective of what happened at $\eta \sim 0$, events may end up in the wrong centrality class – and attributed an incorrect $\langle N_{coll} \rangle$



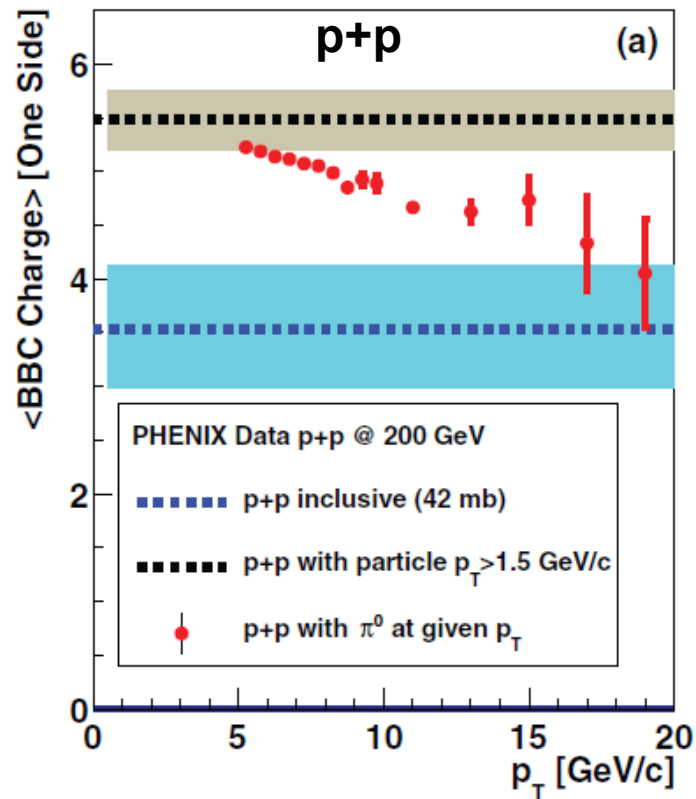
Does bulk observable-based centrality fix N_{coll} once and for all?



One more time...

“In heavy ion collisions, we manipulate the fact that the majority of the initial state nucleon-nucleon collisions will be analogous to minimum bias p+p collisions...”

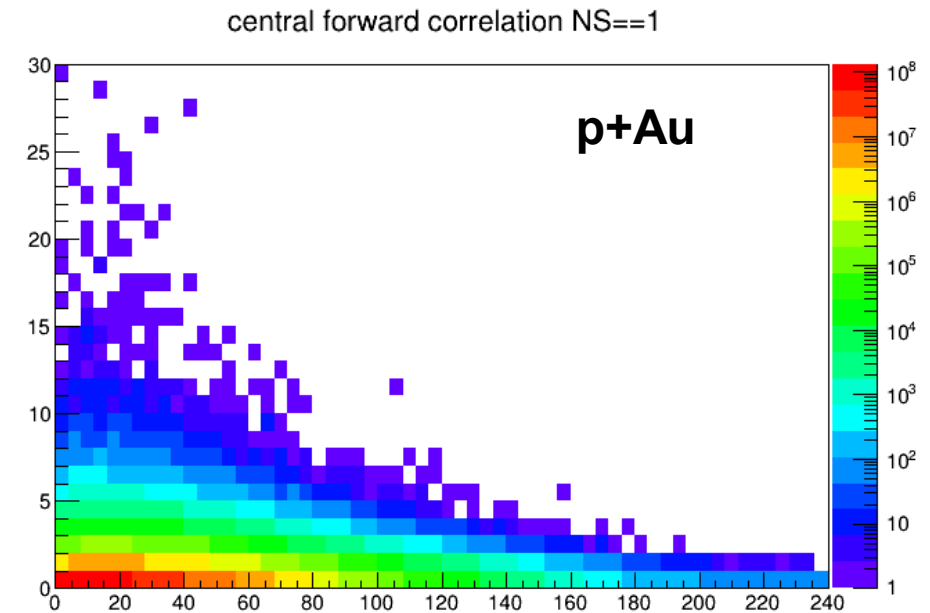
No, it is biased, and the bias changes as a function of the hardest scattering seen at mid-rapidity!



Highest p_T
seen at $\eta = 0$

Anticorrelation!

**In A+A the story
is very different!**



Charge seen at $-3.9 < \eta < -3.0$



A way out: actually measure N_{coll}



You still categorize events with N_{ch} , but override Glauber MC

Is it possible? Yes, at least you can get close, and at the very least get rid of fake final state effects in R_{xA} .

Remember, photons don't care about FS → mostly true, at high p_T most of them are from initial hard scattering and have 200+ *fm* mean free path in QGP (e.g. *Rept.Prog.Phys.* 83 (2020) 4, 046301)

For an arbitrary “centrality” classification just take the ratio of the direct photon and hadron spectra

→ pure centrality bias (even if p_T -dependent) will affect both similarly

→ if the ratios change with centrality, there's a genuine final state effect on hadrons

Same idea, different realization: you can **get N_{coll} experimentally** from the $Y^g(\text{AB}, p_T)/Y^g(\text{pp}, p_T)$ direct photon yield ratios

$$N_{\text{coll}}^{\text{EXP}} = \frac{\left(\frac{d^2 N_\gamma}{dp_T d\eta} \right)_{\text{AB}}}{\left(\frac{d^2 N_\gamma}{dp_T d\eta} \right)_{\text{pp}}} = \frac{Y_{\text{AB}}^\gamma}{Y_{\text{pp}}^\gamma}$$

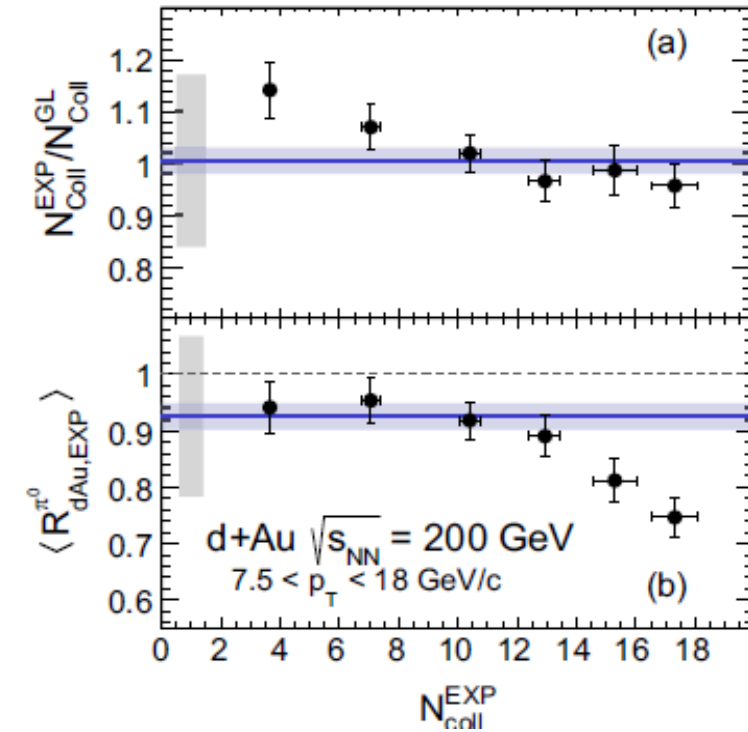
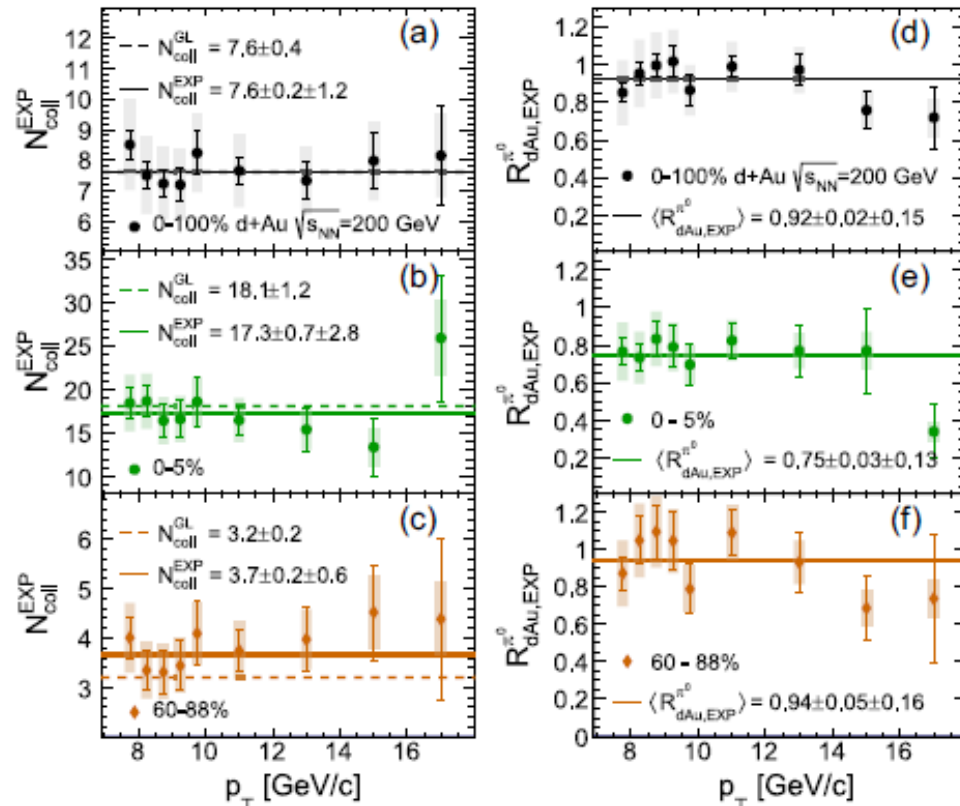
Some caveats...



Nuclear modification of π^0 in dAu

Earlier unphysical enhancement in peripherals gone

PRL 134, 022302 (2025)



Ultimate test: pAu, dAu, $^3\text{HeAu}$
“excitation” function (in the works)





PHENIX / STAR discrepancy

Main culprit: cocktail?



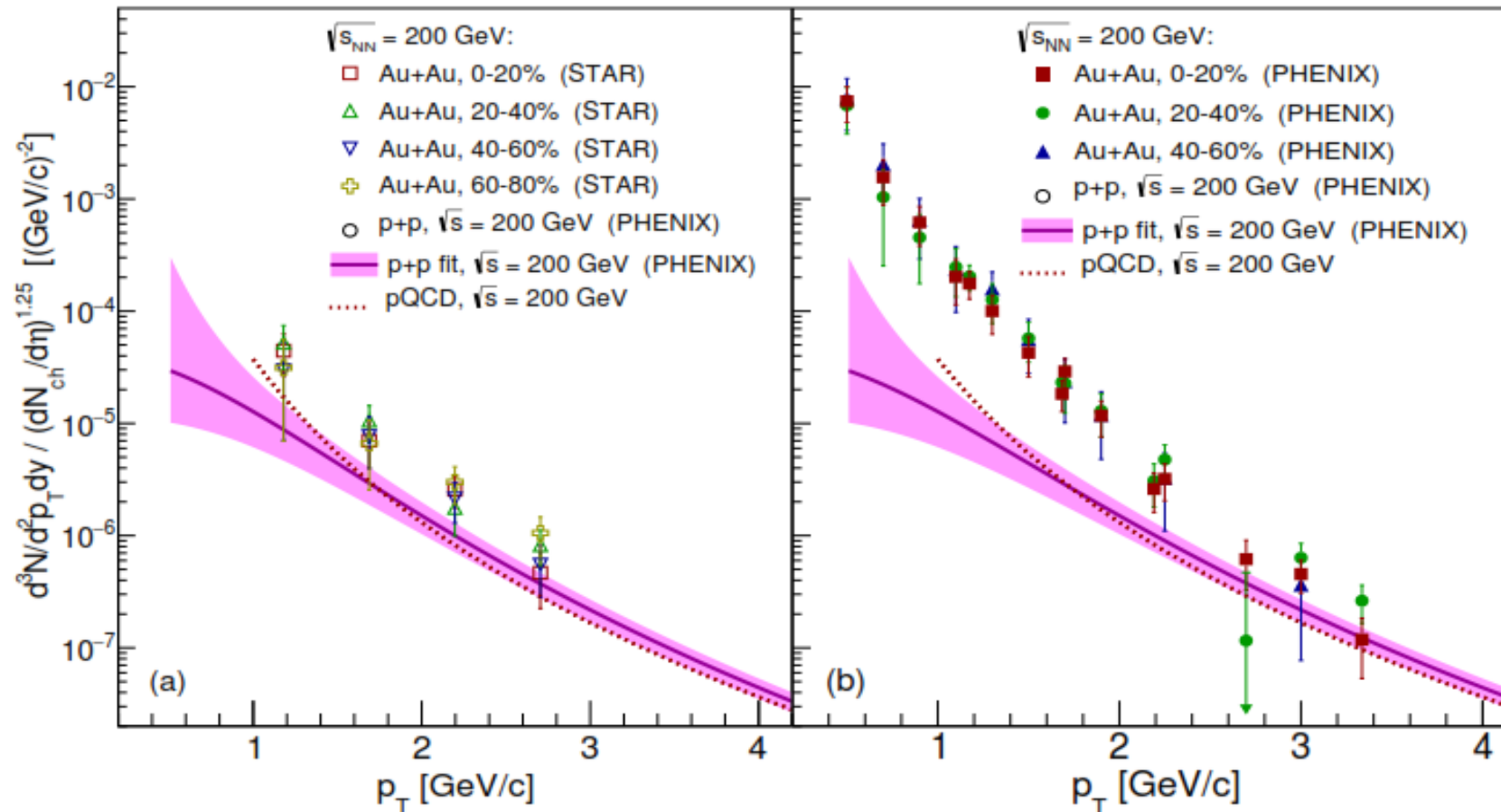
PHENIX / STAR discrepancy (at low p_T)



A sore point since 2016

PHENIX $\gamma \rightarrow e^+e^-$: Phys. Rev. C91 (2015) 064904
 γ^* : Phys. Rev. Lett. 104 (2010) 132301

STAR: Phys. Lett. B770 (2017) 451



Renewed efforts since QM2025

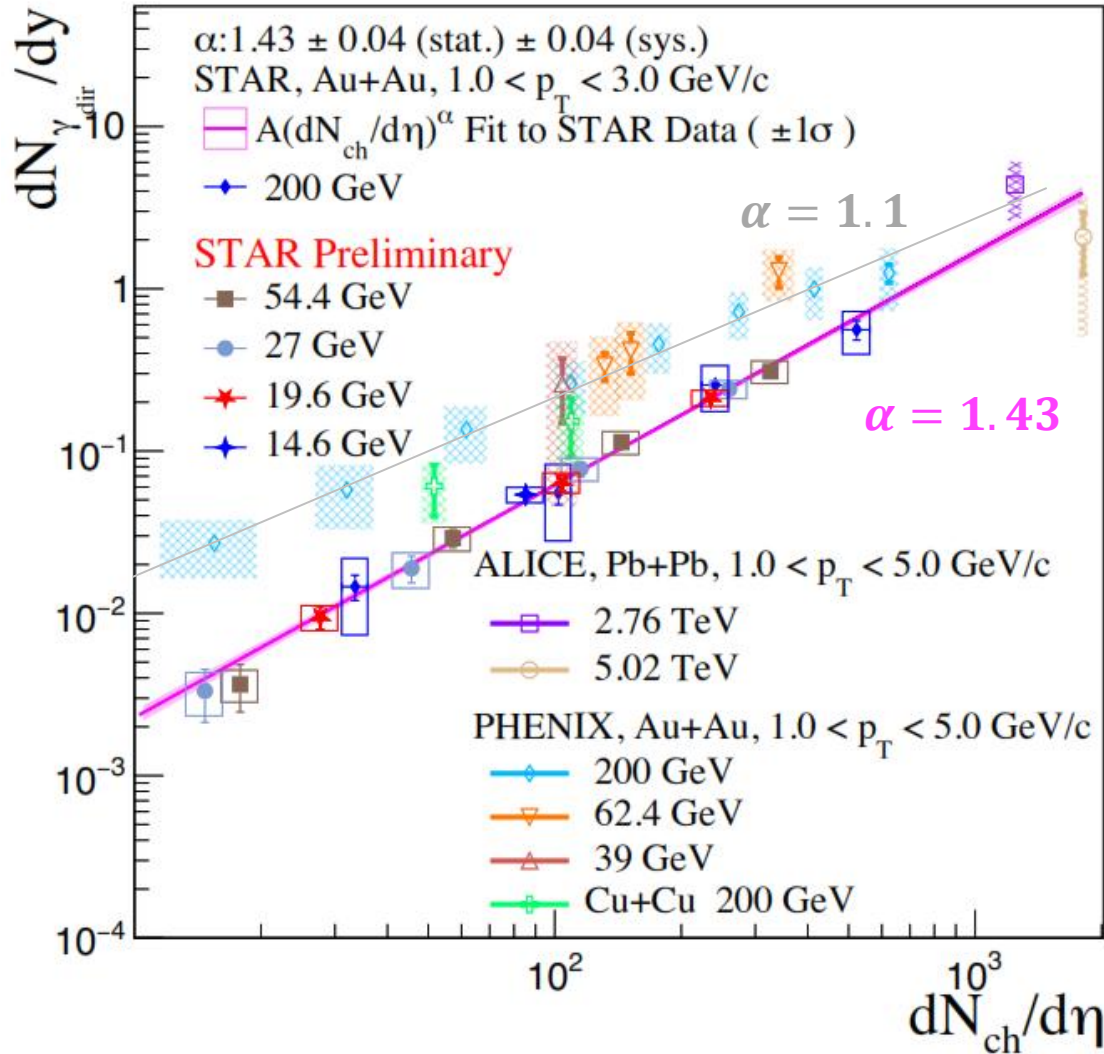


Credit: Axel Drees, SBU

Integrated yields vs N_{ch} (and slopes)



Different evolution with centrality



STAR at QM2025

● STAR results:

- Virtual photon analysis $\gamma^* \rightarrow e^+e^-$ for Au+Au at 14.6, 19.6, 27, 54.4, 200 GeV
- Self consistent analyses & results across energies
- Direct photon yield scales with $N_{\gamma}^{dir} \sim \left(\frac{dN_{ch}}{d\eta}\right)^{\alpha}$ with $\alpha \sim 1.43$

● Compared to PHENIX results:

- STAR yield a factor of 3-5 below PHENIX
- $\alpha \sim 1.43$ versus $\alpha \sim 1.1$ for PHENIX

Credit: Axel Drees, SBU

Very different physics message!



Compare in a particular p_T range



Au+Au minimum bias data

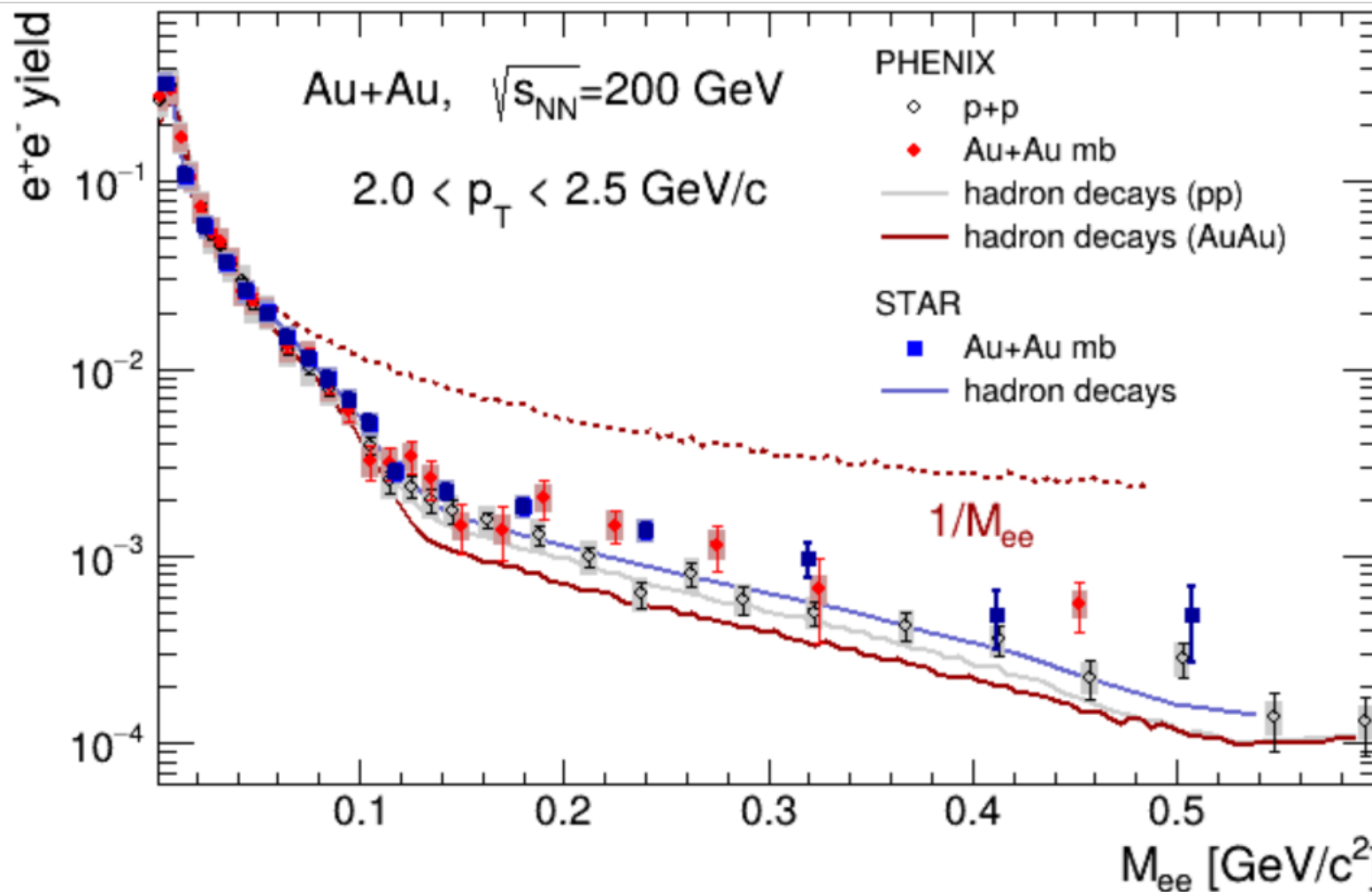
Credit: Axel Drees, SBU

Same p_T selection

All data is normalized
for $m < 30$ MeV

STAR and PHENIX data
agree well in π^0 and η
Dalitz region

η/π^0 at low p_T higher in STAR



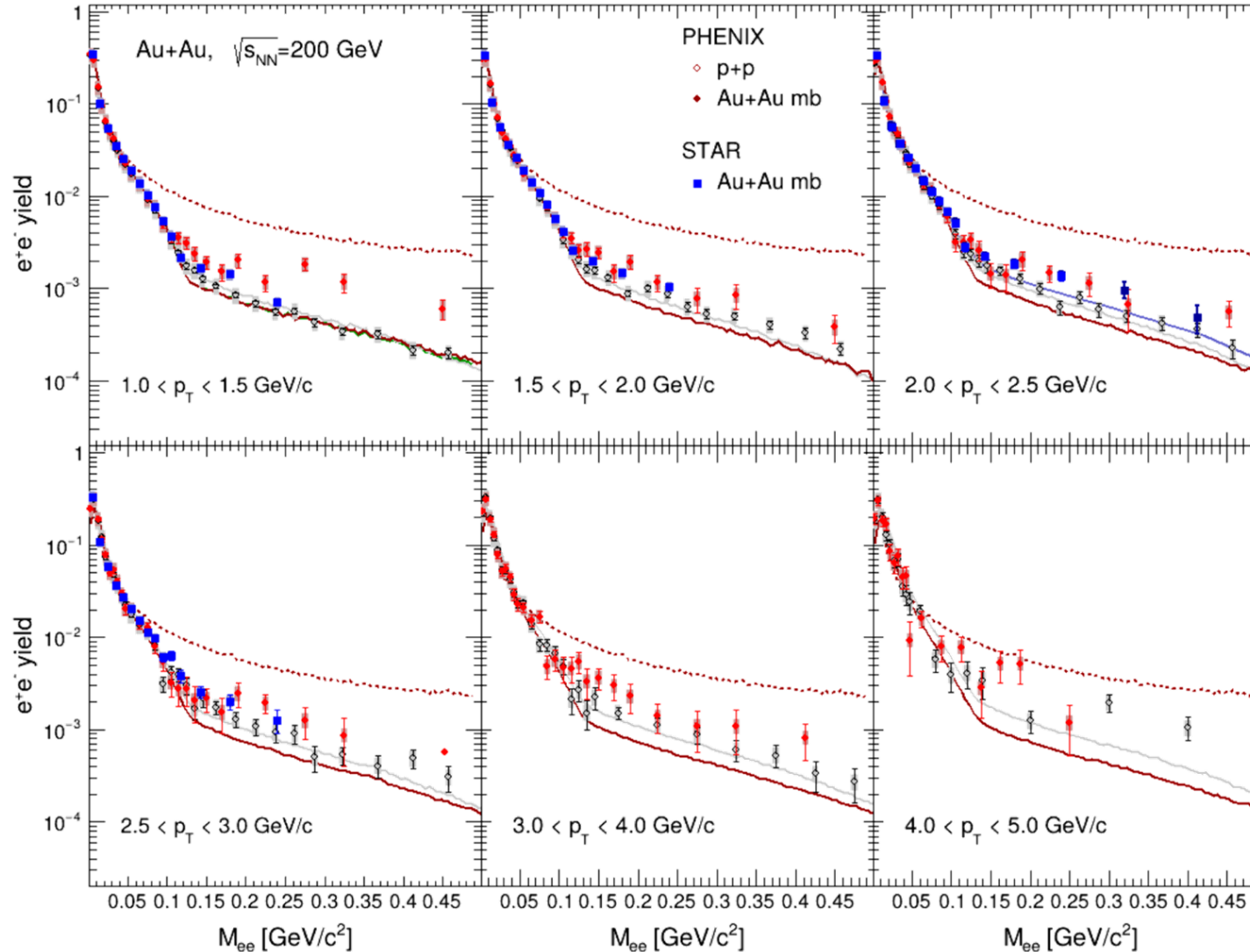
Cocktails do not agree!



Evolution with p_T

How does it evolve with p_T ?

Credit: Axel Drees, SBU



Data seem to be consistent for $p_T > 2$ GeV/c

Difference at lower p_T

Converges



Collaborate with STAR

Exchange AuAu virtual gamma data (not all in public domain or HEP data)

Provide standalone software that reproduces experimental data

Independently analyse data from both experiments with same cocktail

Compare & discuss

- On the PHENIX side

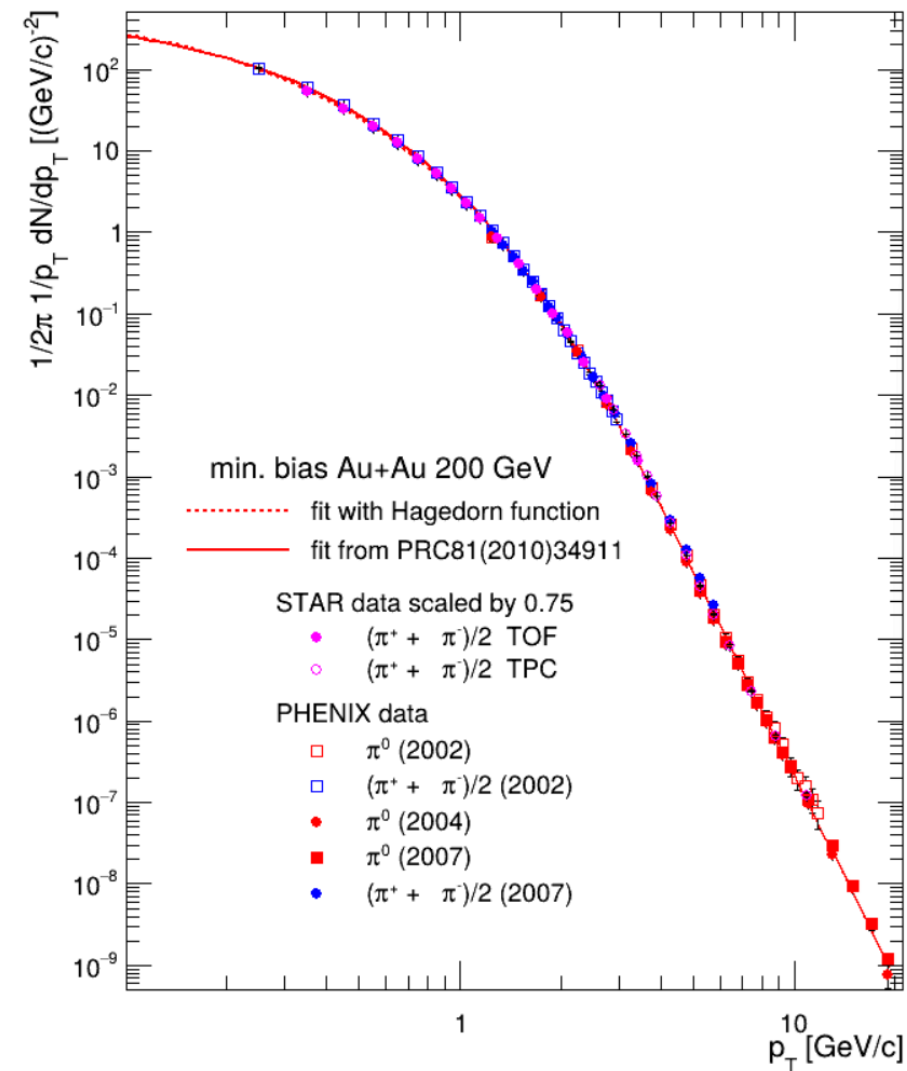
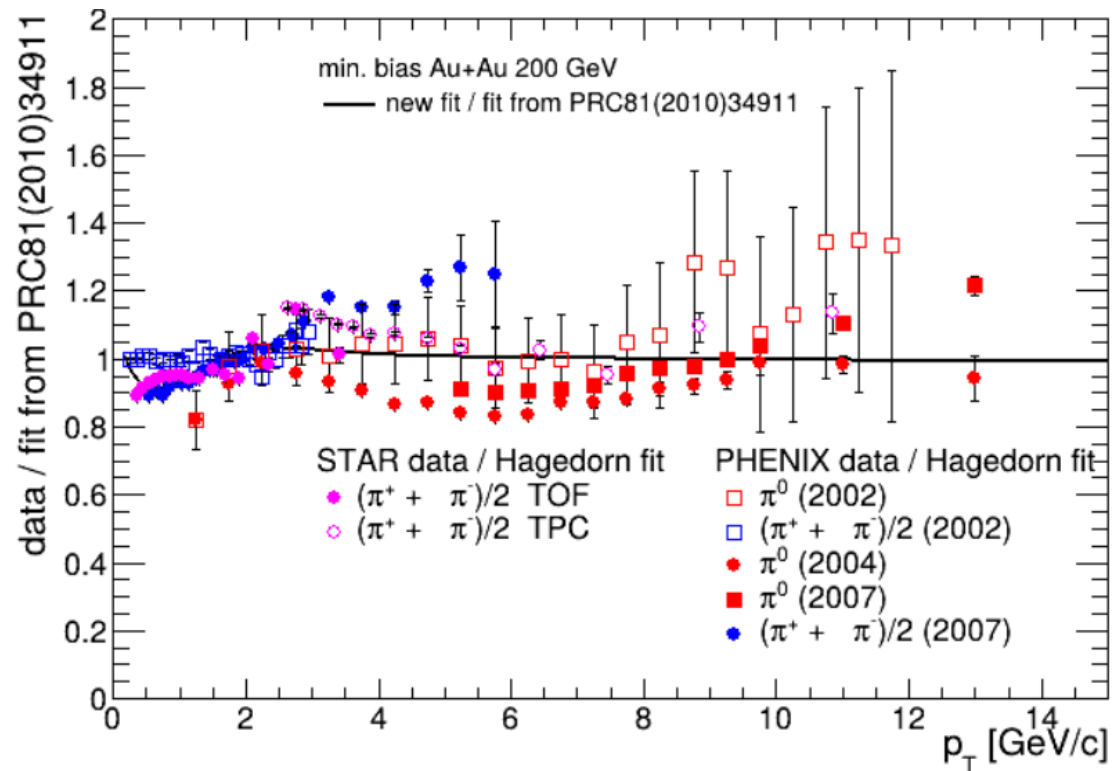
- Collect data in sharable form
- Develop standalone PHENIX simulator
- Create new cocktail based on today's knowledge

Ongoing work with STAR



Consistency of π^0

Compare to charged pions, STAR

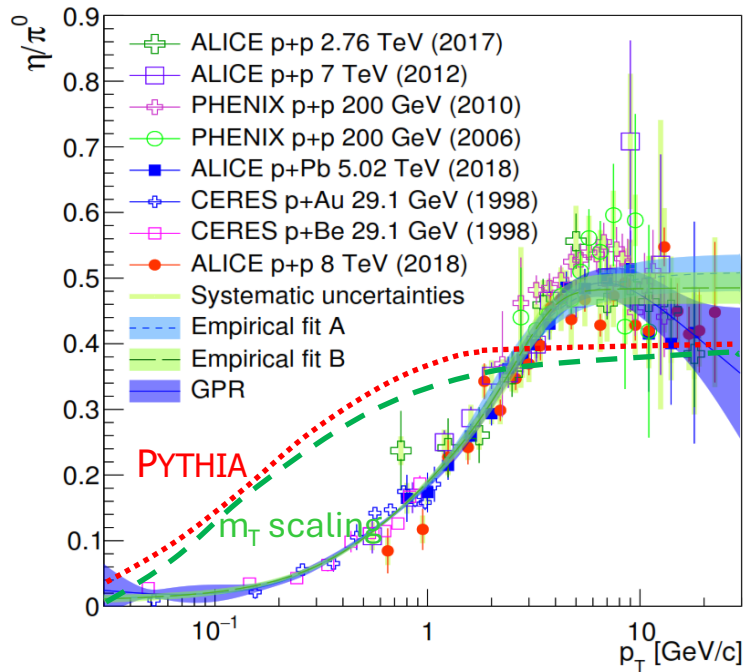


Reasonable agreement



The second highest contributor: η

Recent study of world data



η/π^0 PRC 104, 054902 (2021)

η/π^0 vs collision energy

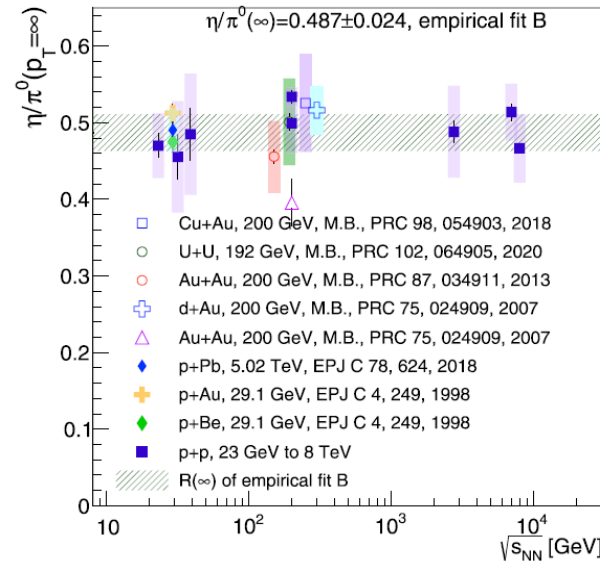


FIG. 4. Values of $R^\infty = \eta/\pi^0(p_T \rightarrow \infty)$ as a function of $\sqrt{s_{NN}}$ for the minimum bias $p+p$, $p+A$, and $A+B$ data sets. Statistical errors are shown as bars, and systematic uncertainties are shown as bands. Also shown is a band representing 0.487 ± 0.024 , the result of the empirical fit B to the combined $p+p$ and $p+A$ data. Note that the $A+B$ data at 200 GeV are offset in $\sqrt{s_{NN}}$ to avoid overlap of data sets.

η/π^0 vs N_{ch} (centrality)

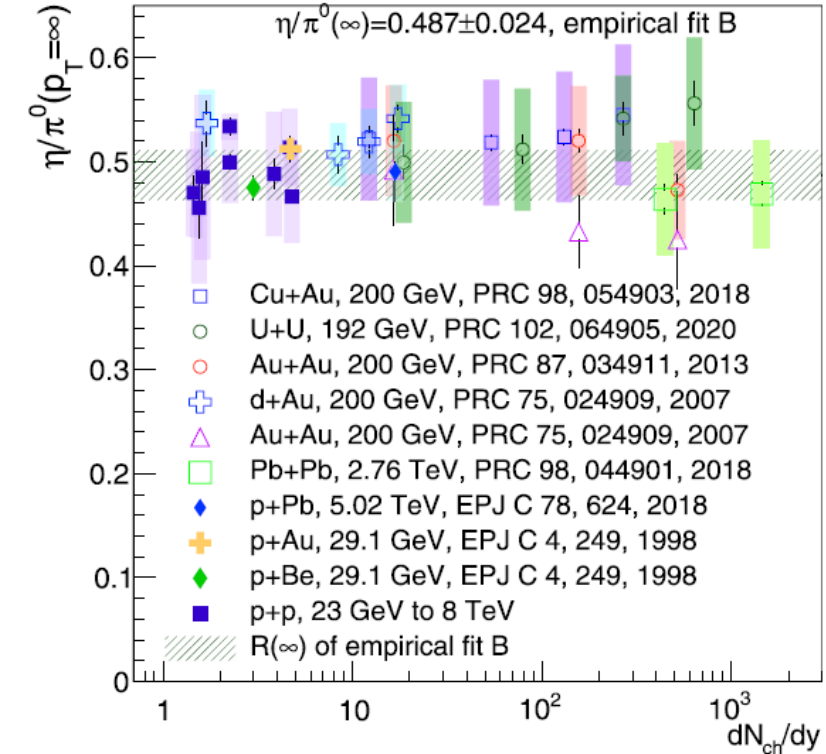
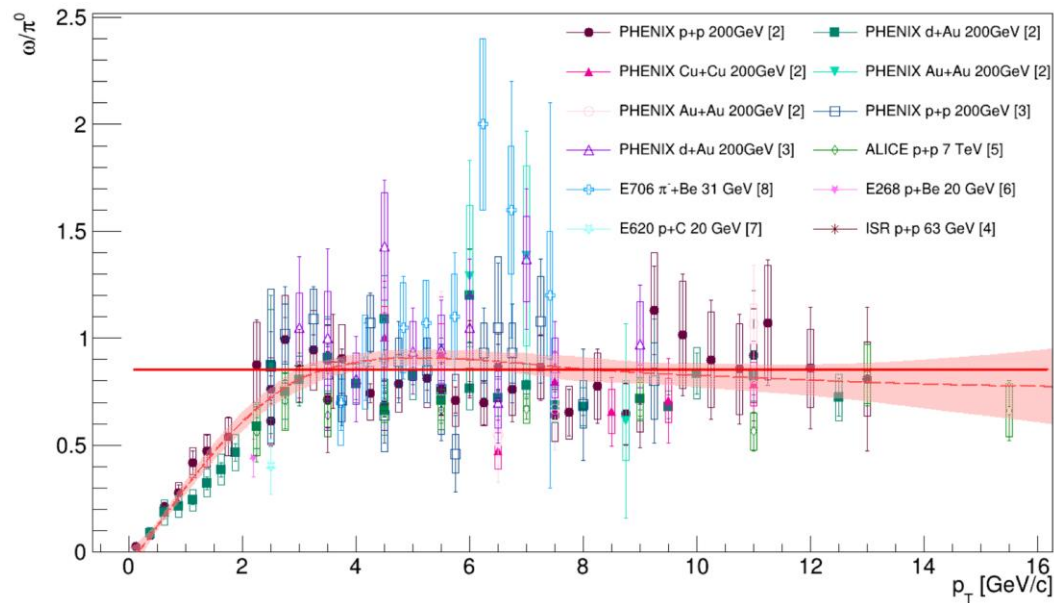


FIG. 5. Values of $R^\infty = \eta/\pi^0(p_T \rightarrow \infty)$ as a function of $dN_{ch}/d\eta$. The presentation is identical to Fig. 4; however, for $A+B$ collisions results from different centrality classes are shown rather than results for the minimum bias sample.

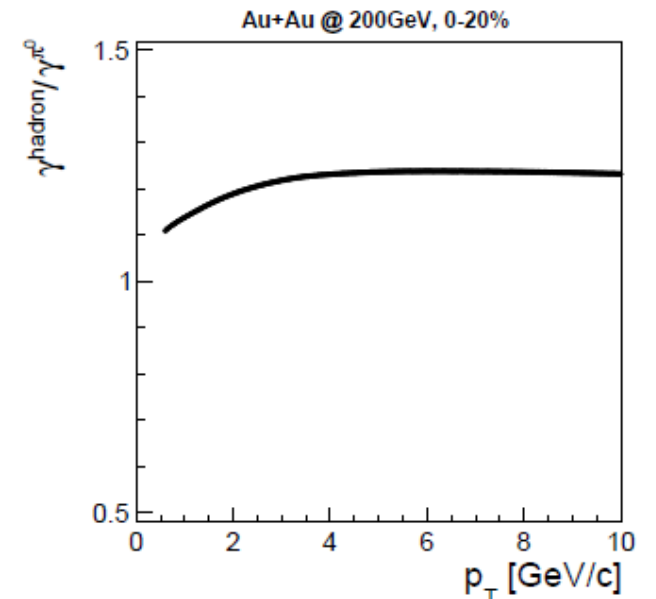
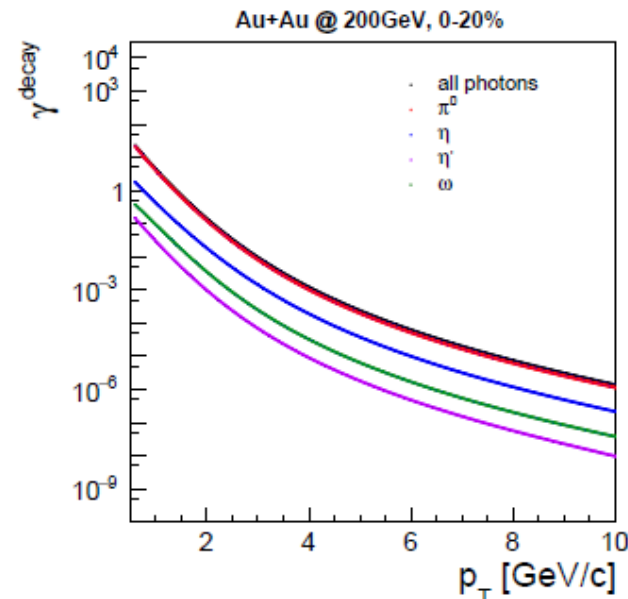
Is it truly constant at high p_T ?

Other contributors – ω , η' ...

Small contribution, but we need a precision measurement



Axel Drees, Konstantin Bauer,
SBU 2023 (unpublished)



Wenqing Fan, SBU, PhD thesis

Few measurements available
Changes with collision energy





“Direct photon puzzle”



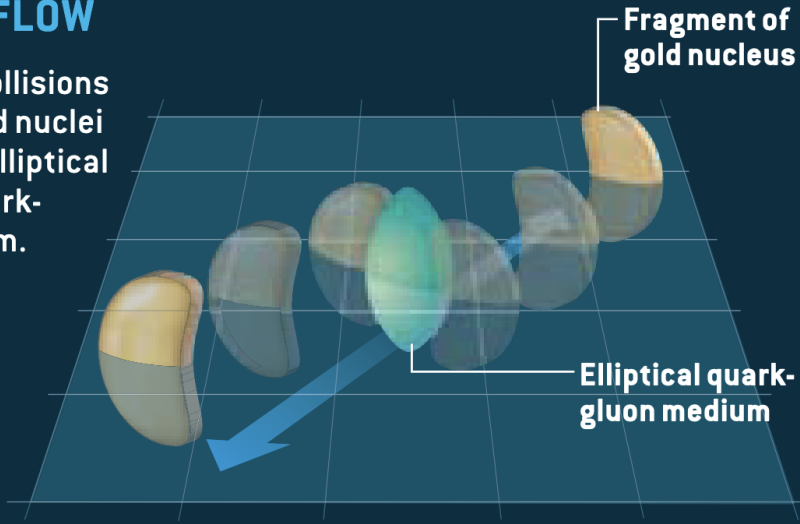
Azimuthal anisotropy in photon emission



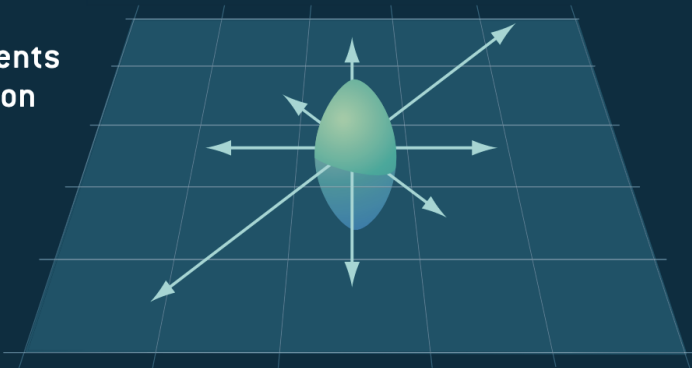
I hate the word “flow”... it implies a particular dynamics

ELLIPTIC FLOW

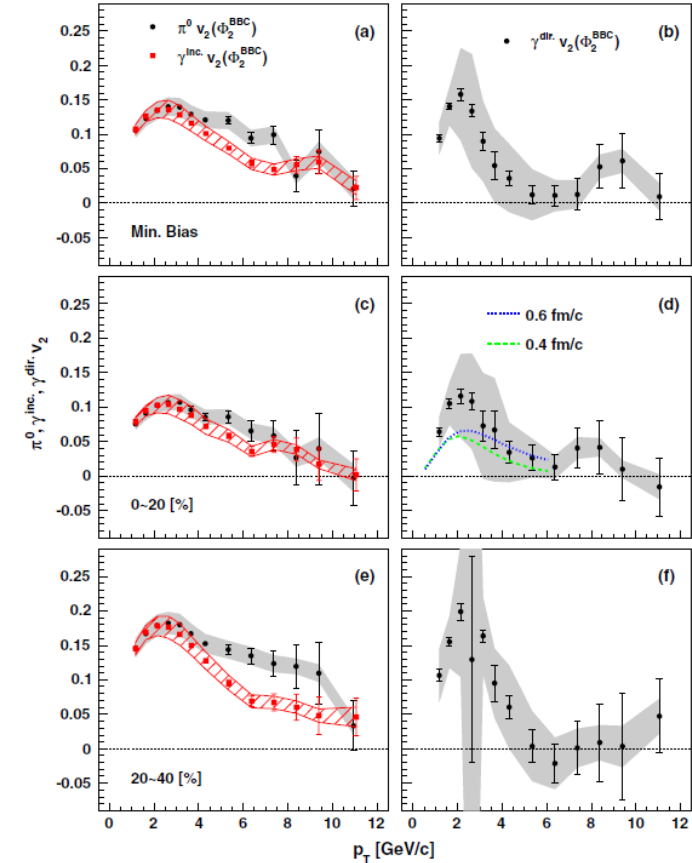
Off-center collisions between gold nuclei produce an elliptical region of quark-gluon medium.



The pressure gradients in the elliptical region cause it to explode outward, mostly in the plane of the collision (arrows).



PRL 109, 122302 (2012)



At low p_T π^0 and hadron v_2 similar!



What is it?

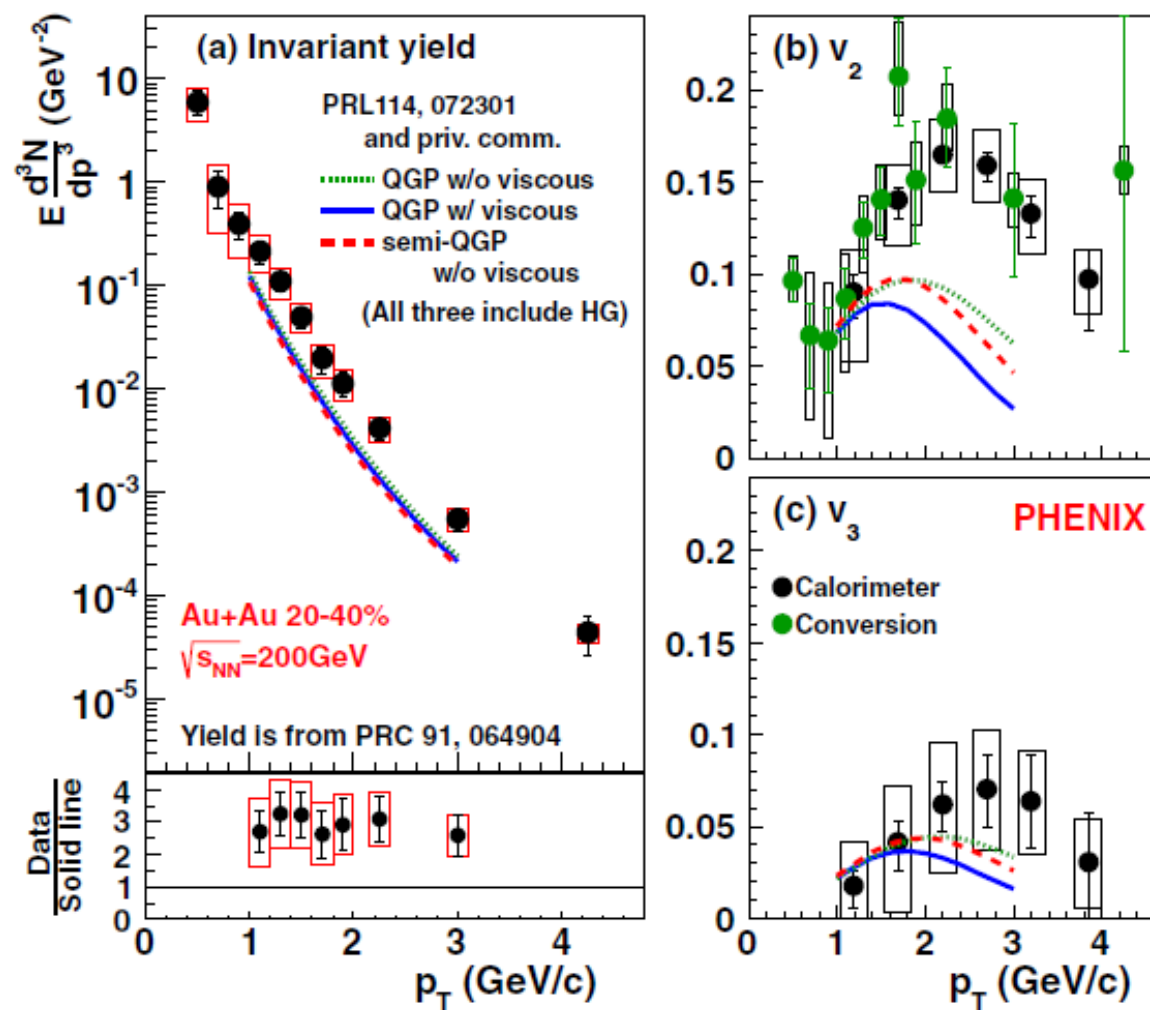
Long-time **paradigm**:
mostly **thermal radiation**
from QGP and HG

Large yield: early emission, high T

Large azimuthal anisotropies:
late emission, low T

Models could not explain
large yield and v_2 simultaneously

PHENIX PRC 94, 064901 (2016)



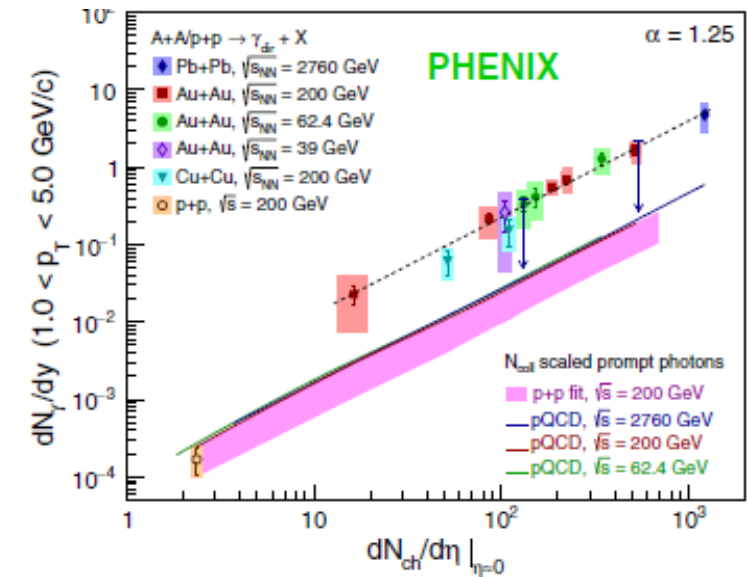
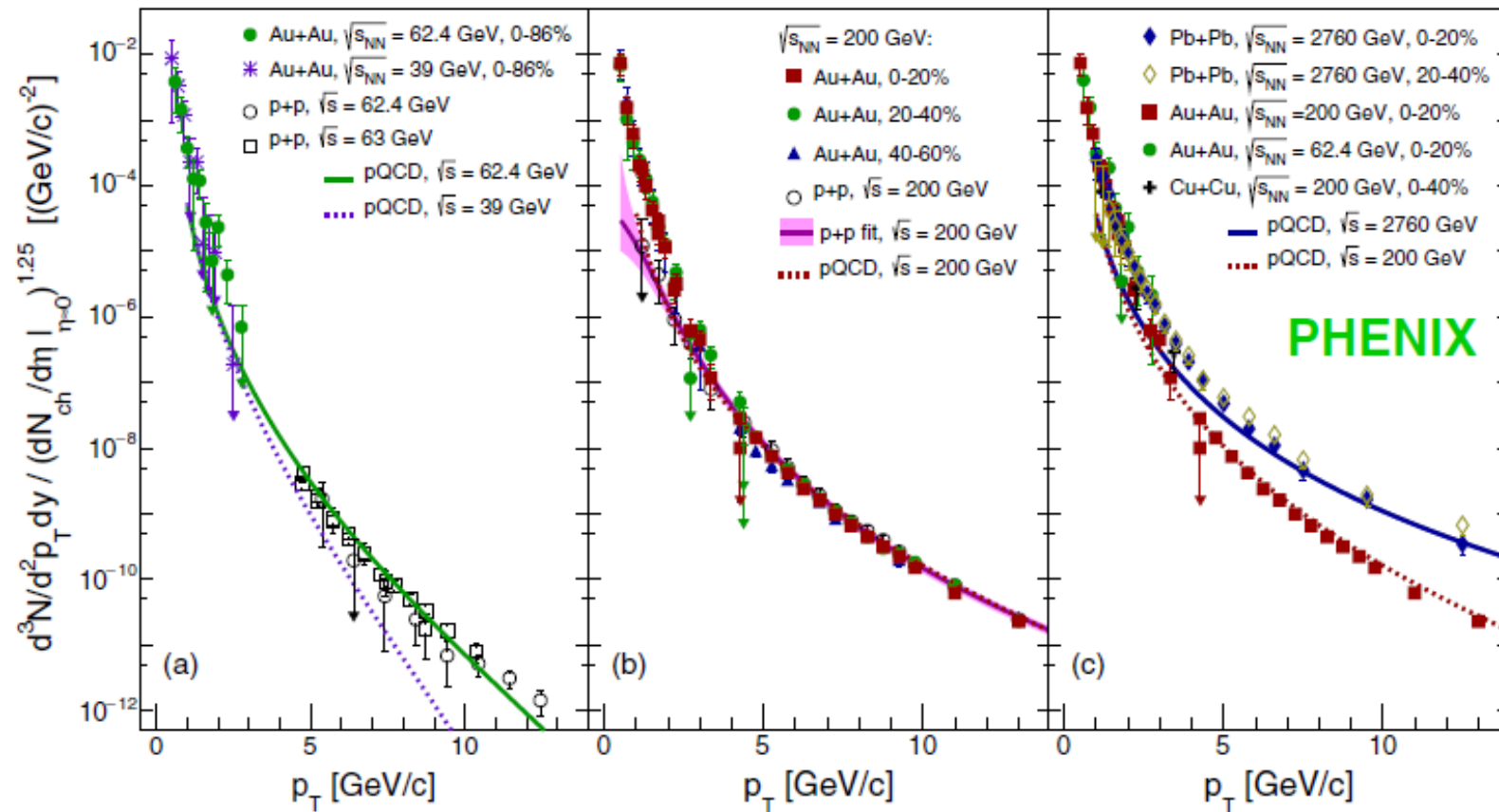
Paradigm shift needed?



Scaling: part of the puzzle?

Universality over system size, collision energy?

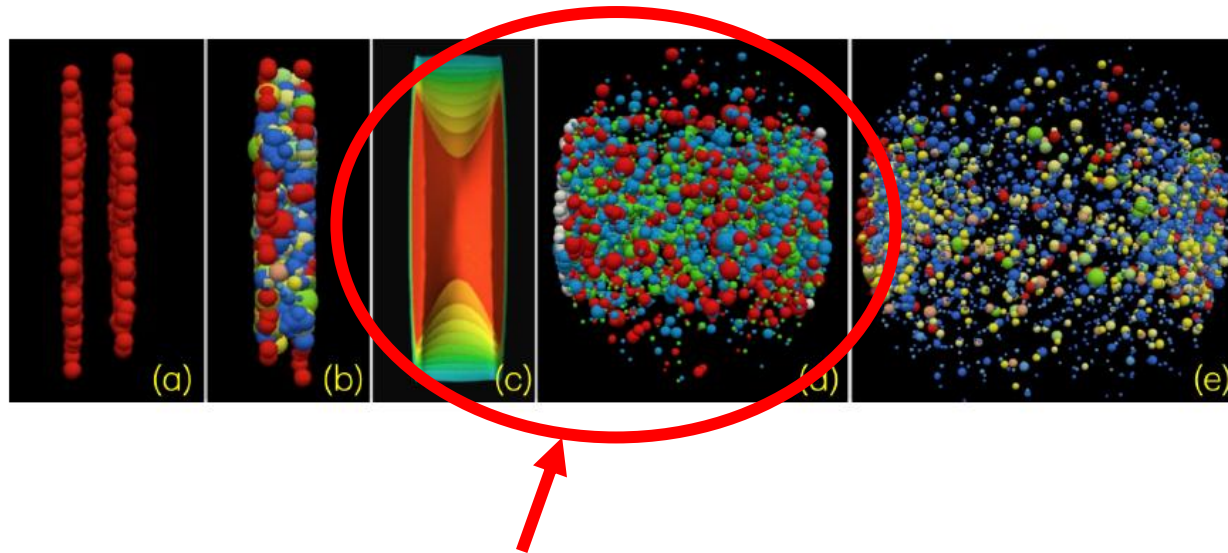
$$\frac{dN_\gamma}{dy} = \int_{p_{T,\min}}^{p_{T,\max}} \frac{dN_\gamma^{\text{dir}}}{dp_T dy} dp_T = A \times \left(\frac{dN_{\text{ch}}}{d\eta} \right)^\alpha$$



Surprisingly small power ($\alpha=1.25$)

Is there a “puzzle” at all?

Thermal/hydro paradigm: yield \rightarrow “early”, flow \rightarrow “late”



Logical – and **tractable!** –
assumption
but why would thermal source
be super-dominant?
No a priori reason
“Unconventional” sources?

Concentrating on this region and as a (locally)
thermalized hydro system is too narrow-minded?

Is the puzzle just due to some tunnel-vision?

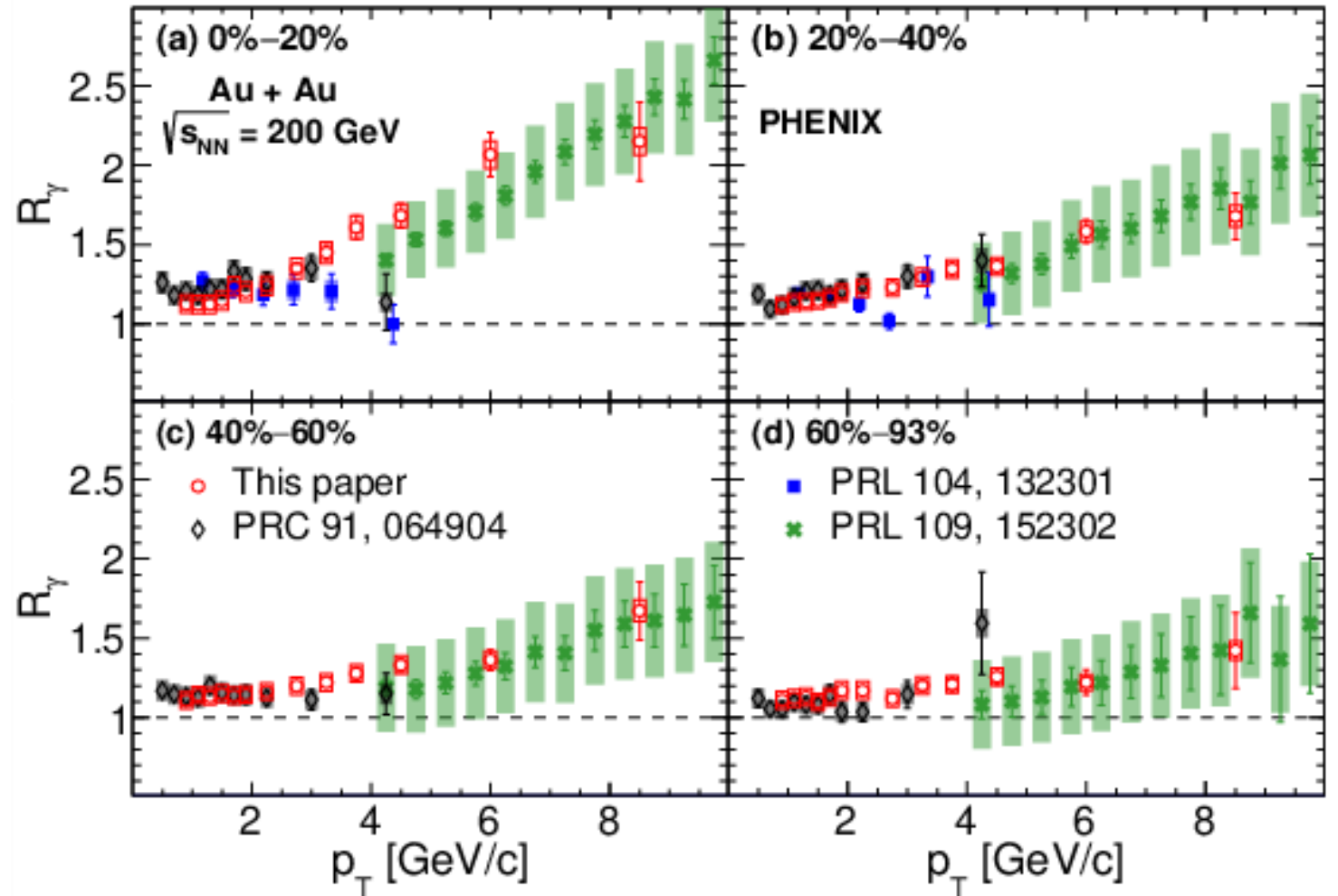
Gale arXiv:2502.13938



Four different analysis techniques...

$$R_\gamma = \frac{\gamma^{inclusive}}{\gamma^{decay}}$$

Clear excess in the
“thermal” region,
depending on
centrality



PRC 109, 044912 (2024)

...giving consistent results

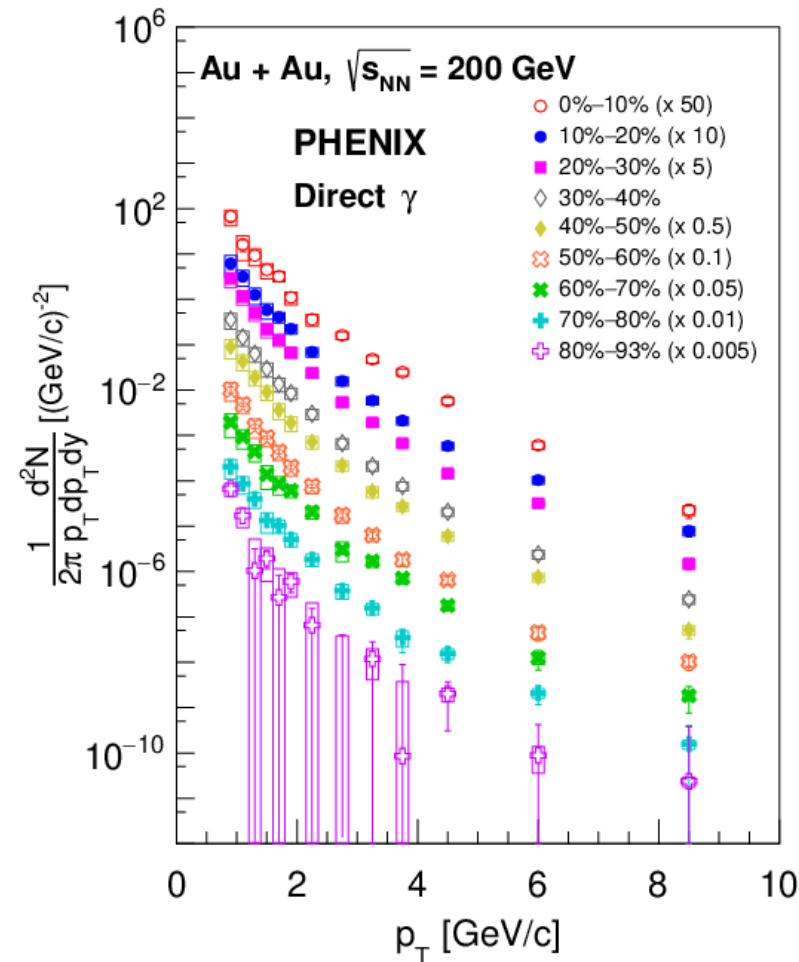
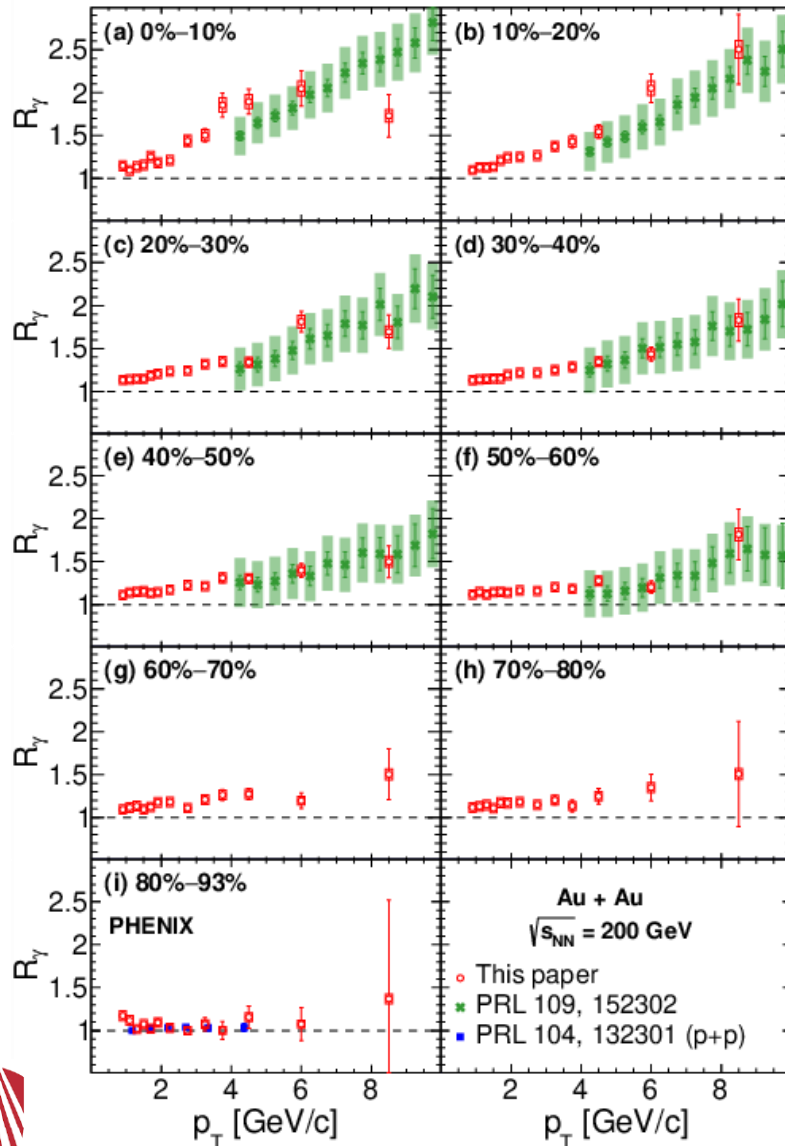


External conversion in VTX

A closer look at R_γ and direct γ

R_γ – and its uncertainty – is crucial in the analysis of azimuthal asymmetries

PRC 109, 044912 (2024)



Limited range
→ high precision
(Conversion)

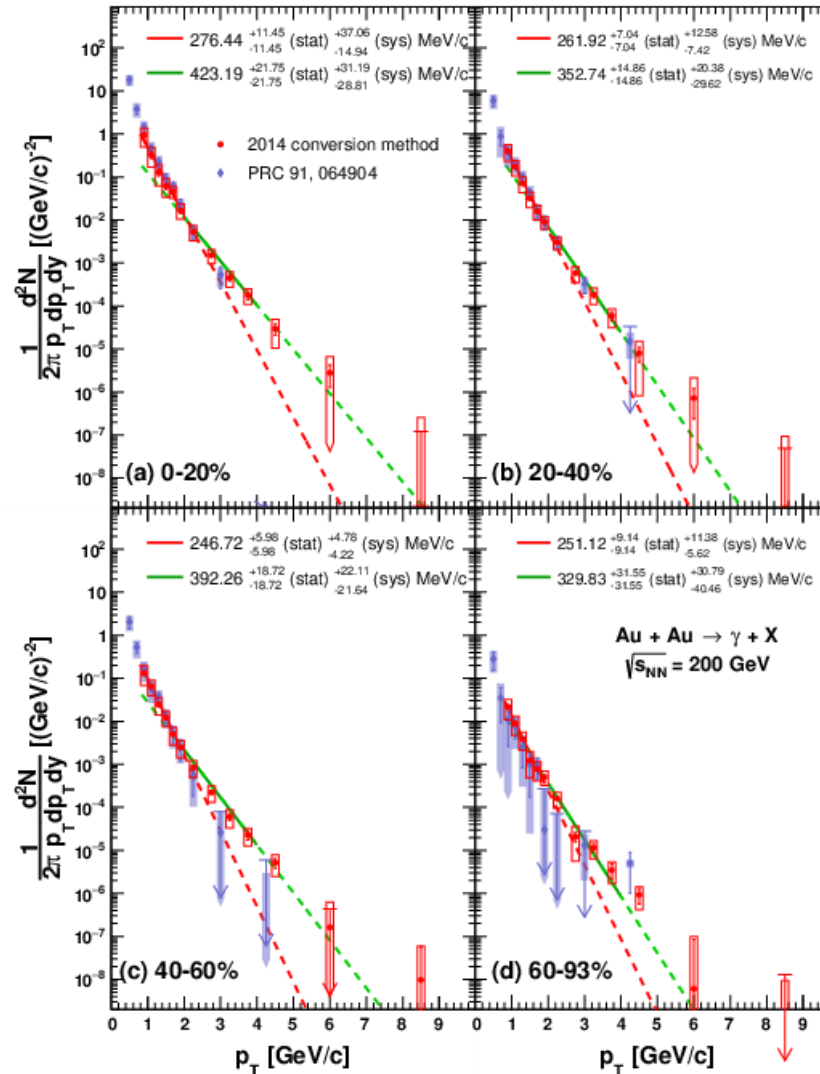
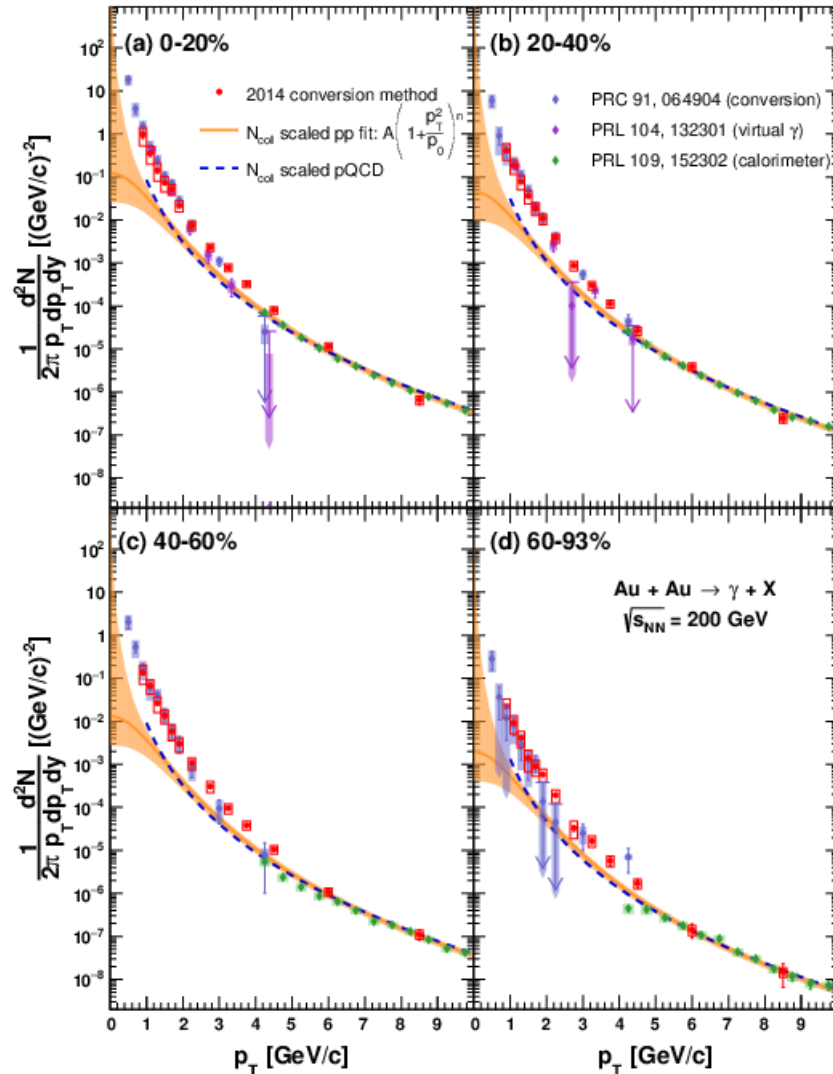
Consistent, solid overlap with calorimeter data

Direct and non-prompt photons



PRC 109, 044912 (2024)

Non-prompt: subtract p+p (extrapolated)



Non-prompt:
not a single
exponential
("Thermal" is
frowned upon ☹)

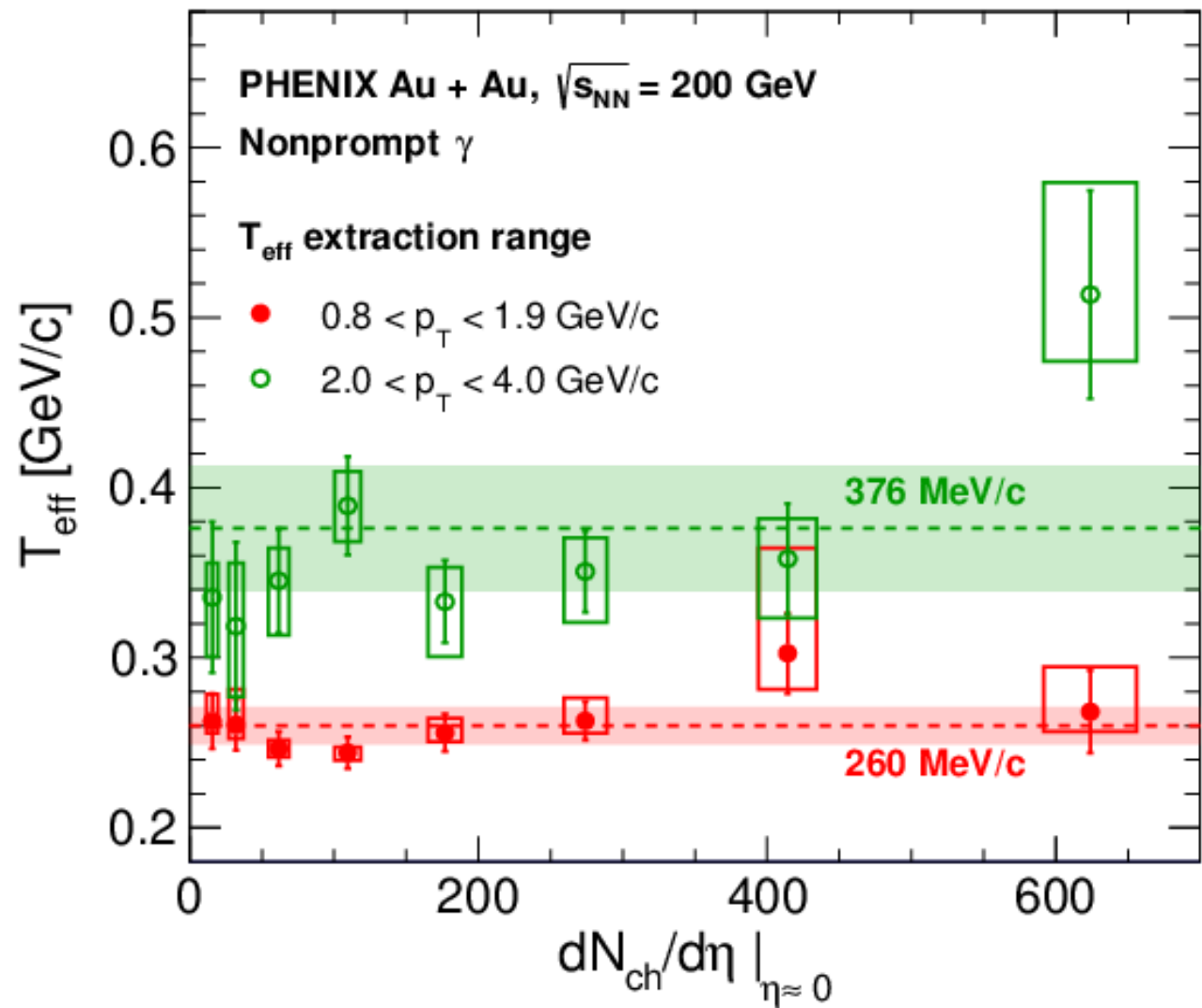
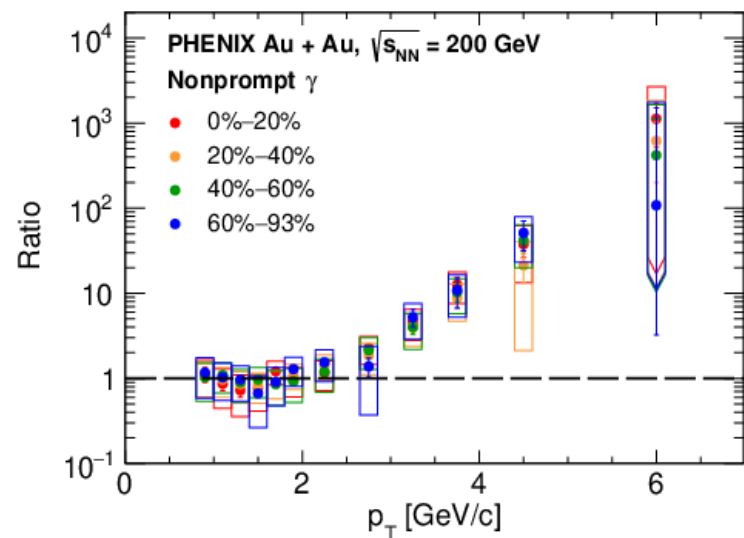
T_{eff} increases
with p_T



T_{eff} vs the p_T range and centrality

No dependence on centrality

Strong dependence on p_T
Composition and relative weight of different sources always the same?



PRC 109, 044912 (2024)

Yields: universal scenario?



Easing a tension on α being too small?



PRC109, 044912 (2024)

Issue: photon yield vs N_{ch}

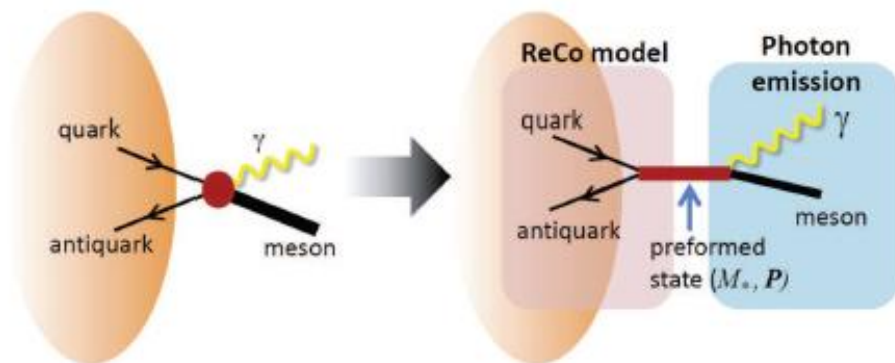
$$dN_{\gamma}/dy \sim (dN_{ch}/d\eta)^{\alpha}$$

α from QGP ~ 1.85

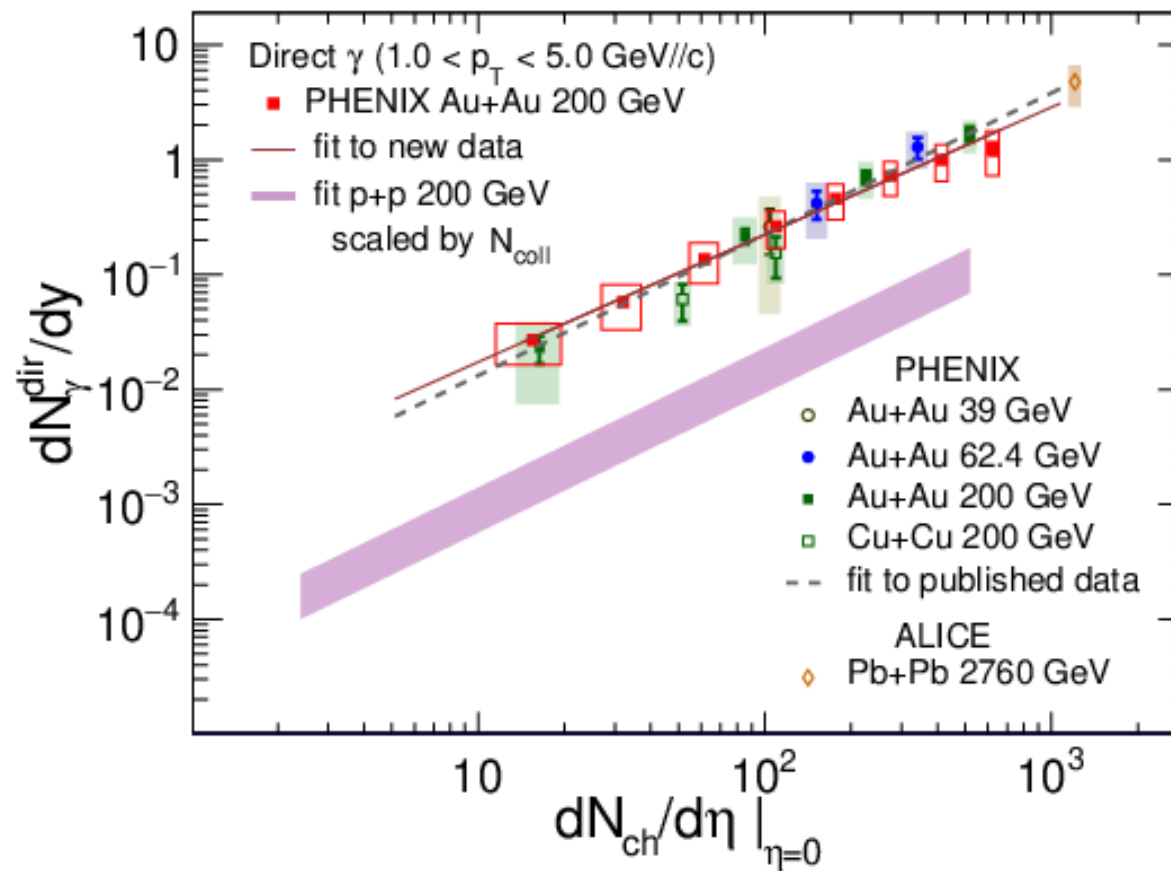
α from HG ~ 1.25

α **measured** ~ 1.11

But radiative hadronization: $N_{\gamma} \sim N_{ch}$



PRC 106, 034906 (2022)

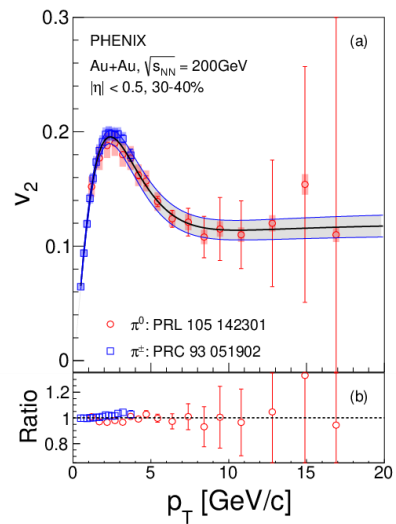


May help to explain small α observed

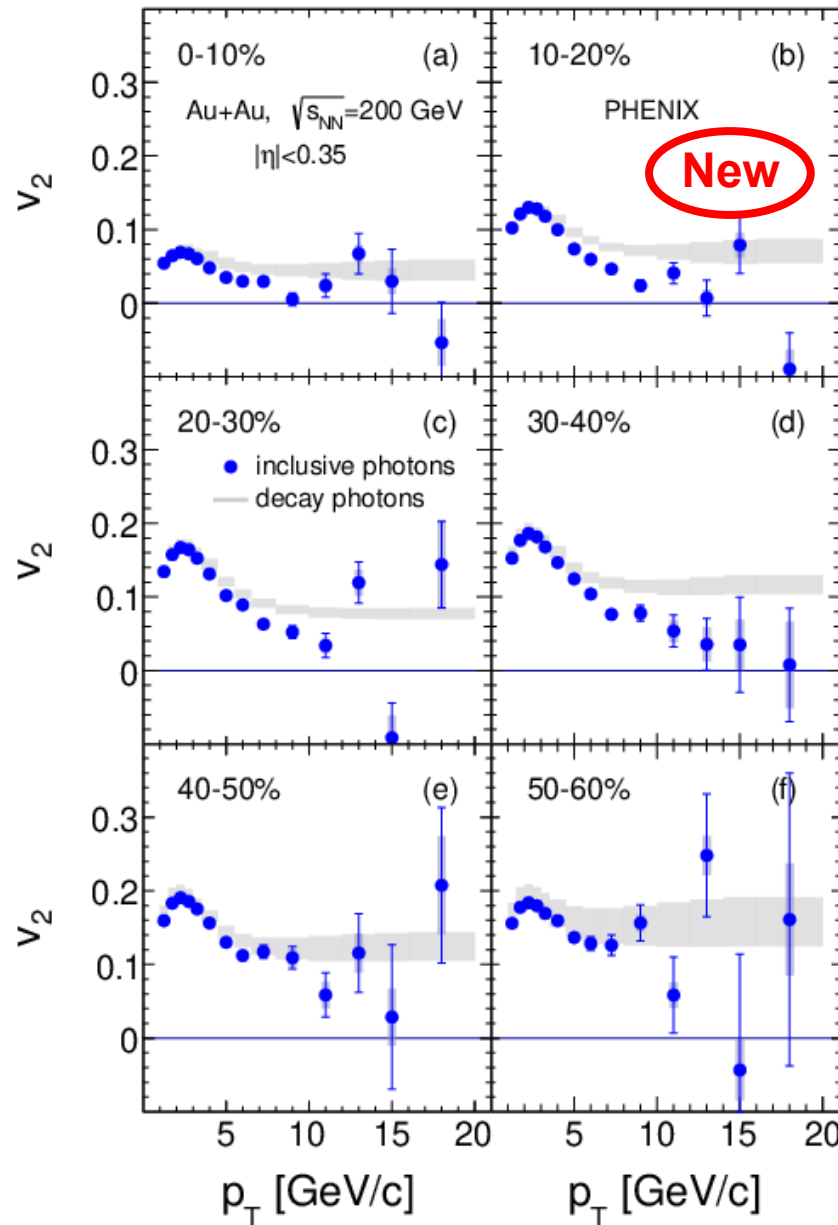


Direct photon v_2 – starts with inclusive and decay v_2

**Inclusive γ
Decay γ
(cocktail)**



Pion v_2 and fit
(used for decay v_2)



arXiv:2504.02955

At low p_T **inclusive (blue)**
and **decay (gray band)**
 v_2 almost identical

At high p_T strong separation
(prompt photon $v_2 \sim 0$)

$$v_2^{dir} = \frac{R_\gamma v_2^{inc} - v_2^{dec}}{R_\gamma - 1}$$

v_2 hyper-sensitive to R_γ

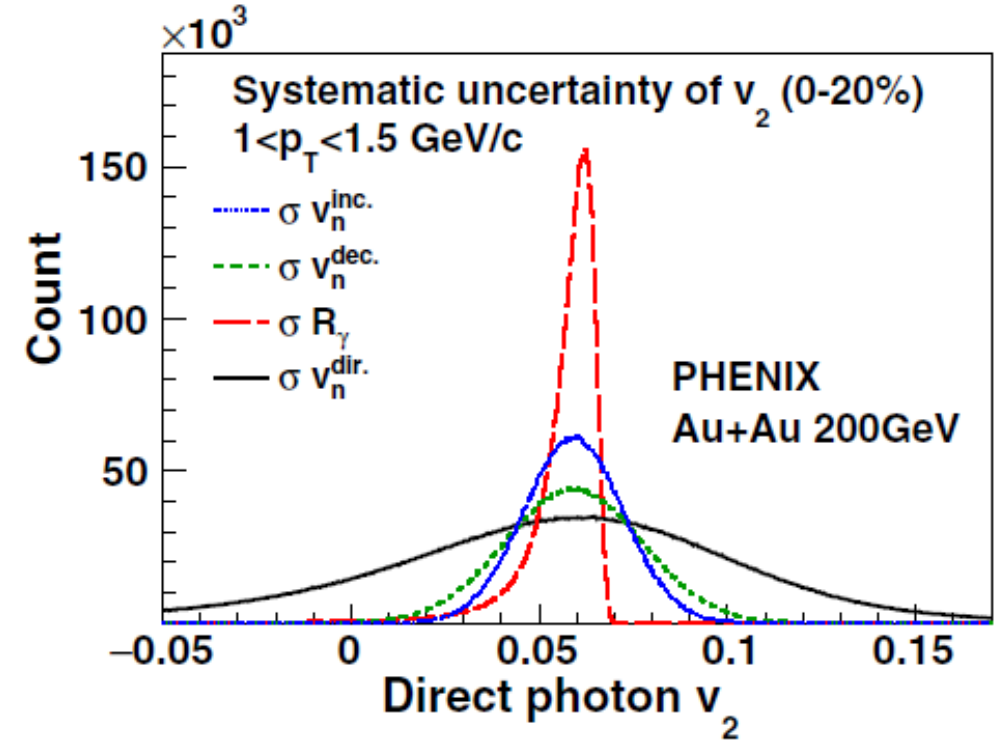
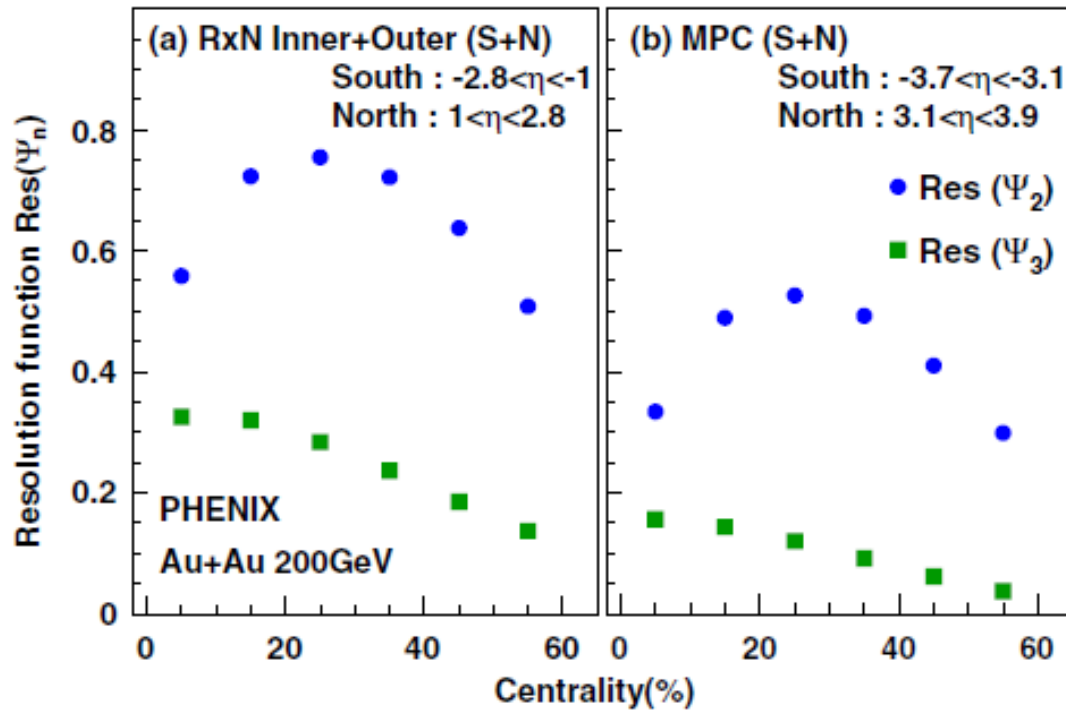
Reaction plane resolution, uncertainties on v_2



PRC 94, 064901 (2016)

Main sources of uncertainties

$$v_n = v'_n / \text{Res}(\Psi_n).$$



$$v_2^{\text{dir}} = \frac{R_\gamma v_2^{\text{inc}} - v_2^{\text{dec}}}{R_\gamma - 1}$$

v_2 hyper-sensitive to R_γ

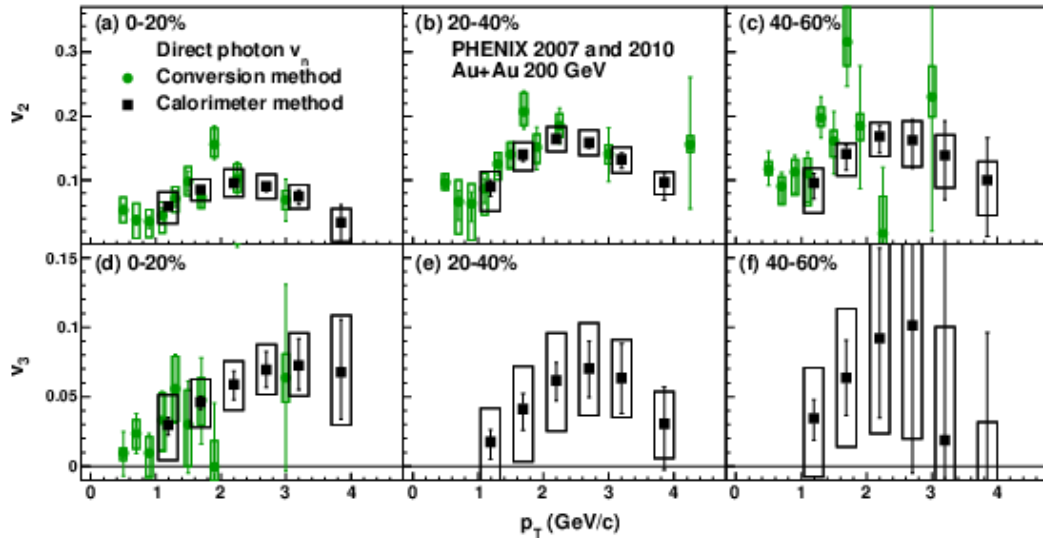




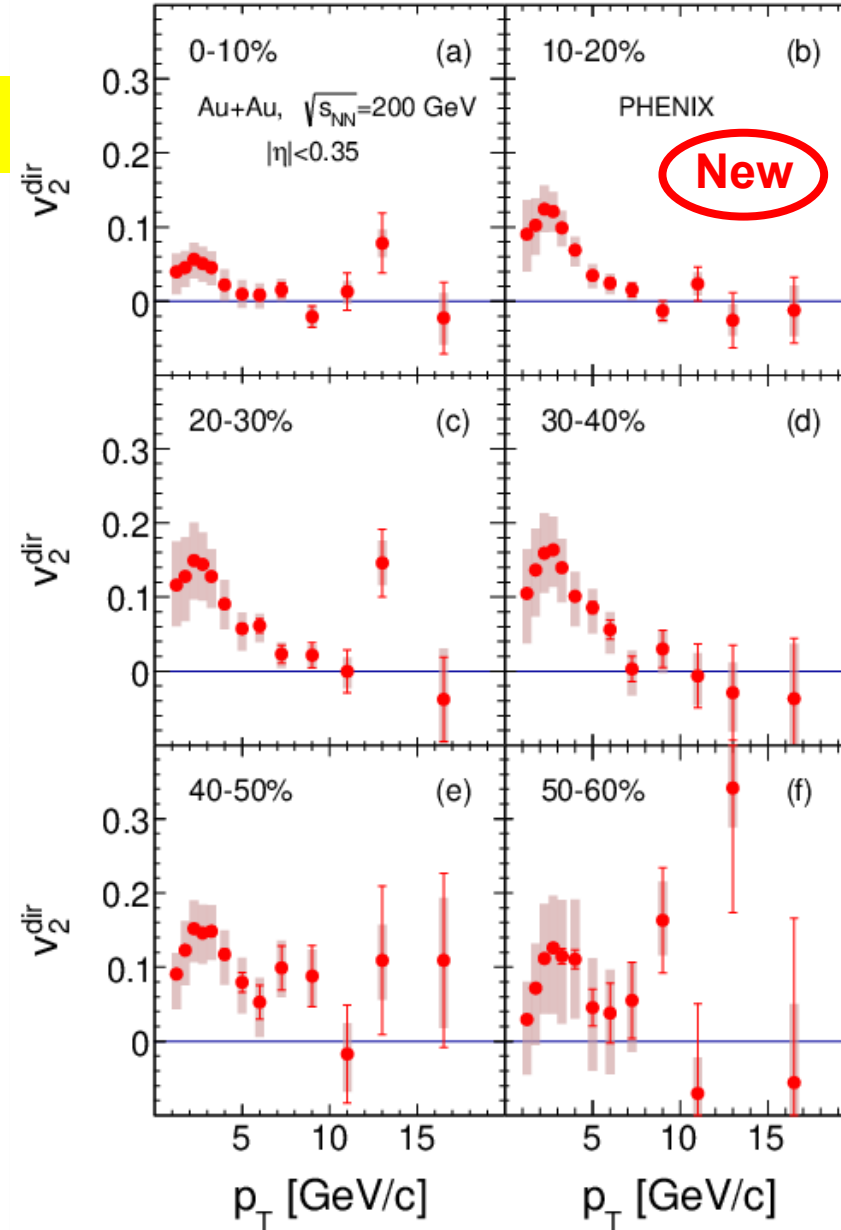
Direct photon v_2 – the new result

Issue: magnitude, p_T dependence

Previous: PRC 94, 064901 (2016)



Conversion, calorimeter
Limited p_T range
Significant v_2



arXiv:2504.02955

10% centrality

Wide p_T range

Calorimeter

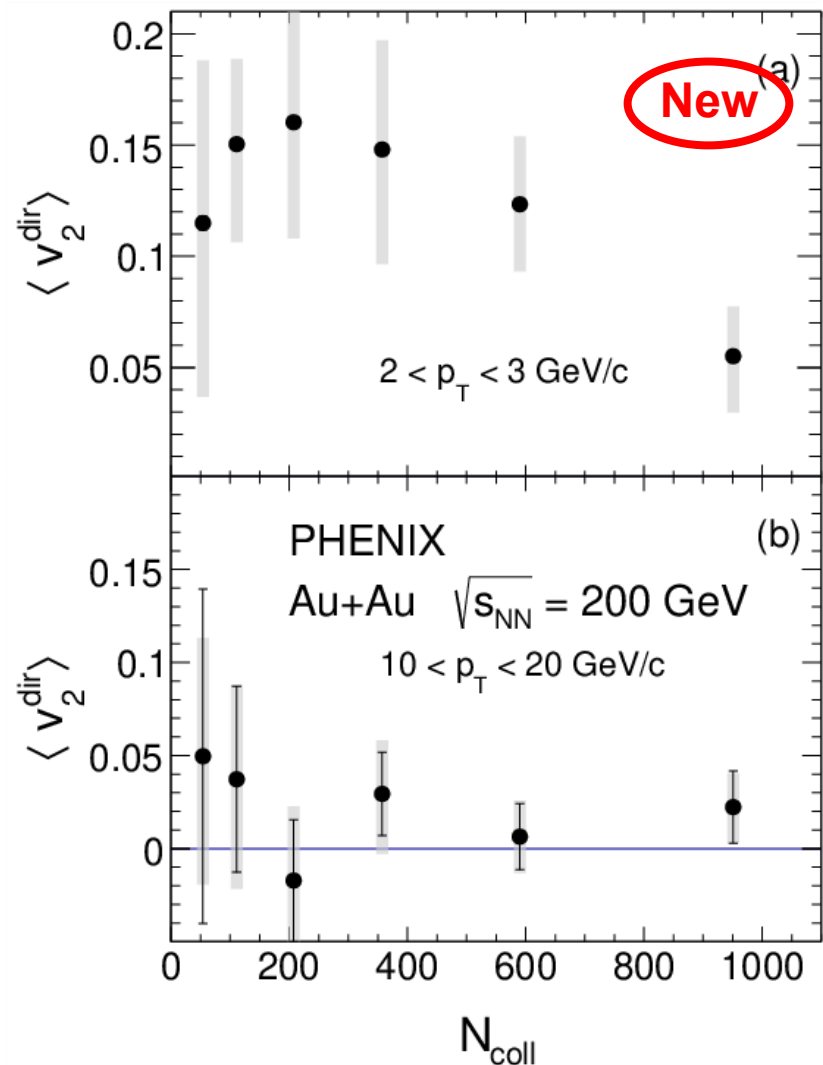
Large at low p_T
vanishes at
high p_T



Average v_2 vs p_T and centrality

arXiv:2504.02955

Low and high p_T regions



Low p_T (non-prompt)
 \rightarrow large v_2 , centrality dependence

High p_T (prompt)
 $\rightarrow v_2 \sim 0$, no centrality dependence

Hard scattering – $v_2 \sim 0$

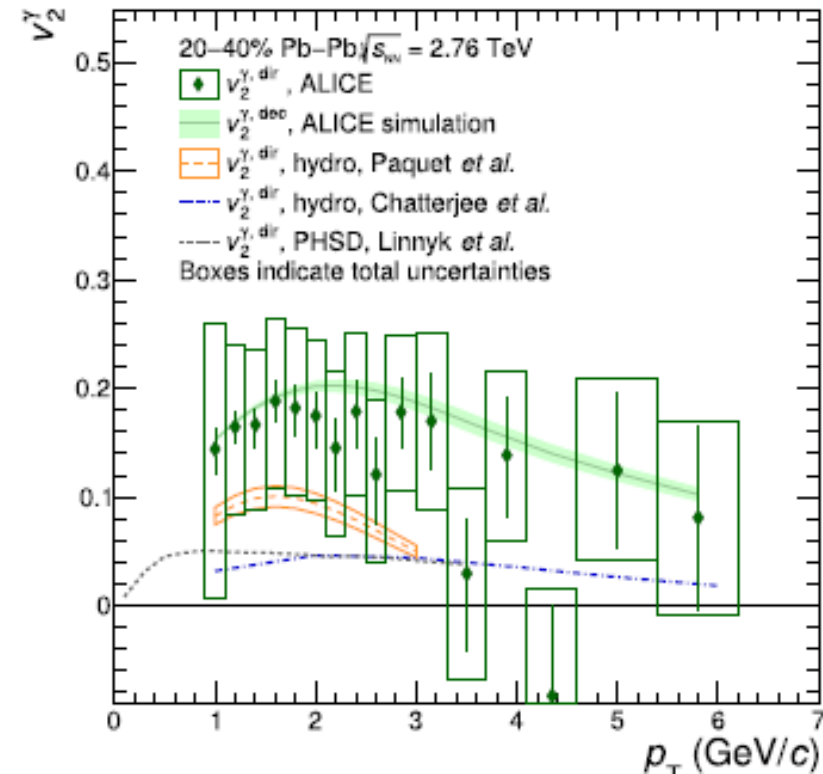
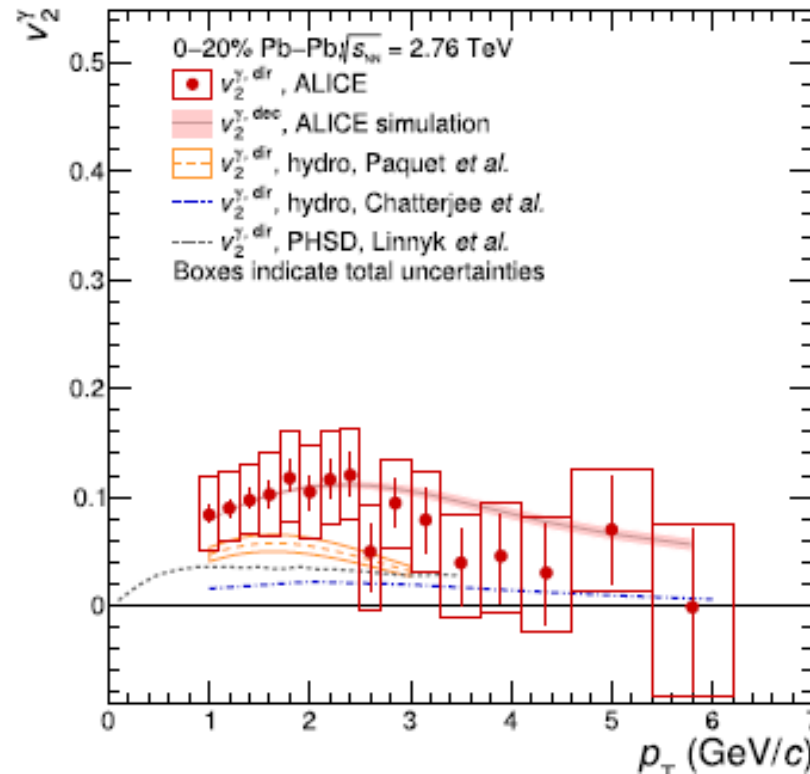




Direct photon v_2 in 2.76 PbPb (ALICE)

PLB 789 (2019) 308-322

Low and high p_T regions



Similar as at PHENIX



Precision on R_γ crucial

PLB 789 (2019) 308-322

$$v_2^{dir} = \frac{R_\gamma v_2^{inc} - v_2^{dec}}{R_\gamma - 1}$$

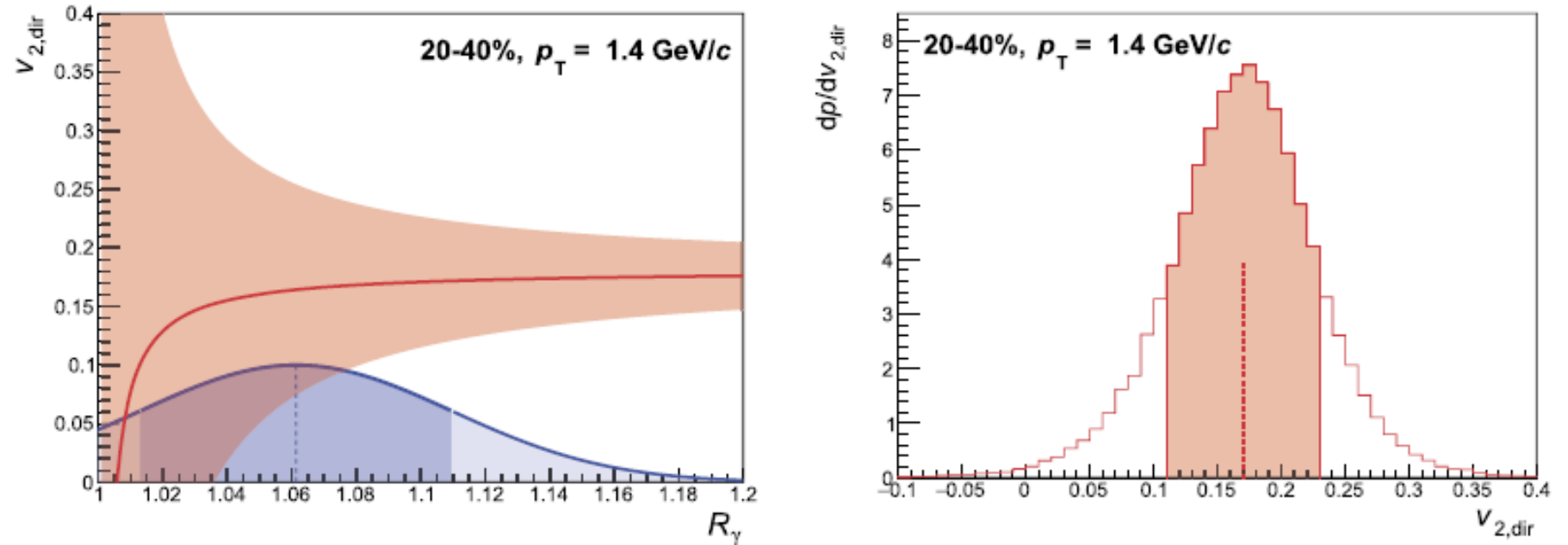


Fig. 4. Left: Central value (solid red line) and uncertainty of the direct-photon v_2 for a selected p_T interval. The upper and lower edges of the red shaded area correspond to the total uncertainty of $v_{2,dir}^{Y,dir}$ as obtained from linear Gaussian propagation of the uncertainties $\sigma(v_2^{Y,inc})$ and $\sigma(v_2^{Y,dec})$. The Gaussian (with arbitrary normalization) reflects the measured value of R_γ in this p_T interval (blue dashed line) and its $\pm 1\sigma$ uncertainty (dark-blue shaded interval). Right: Posterior distribution of the true value of $v_{2,dir}^{Y,dir}$ for the same interval in the Bayesian approach. Note that the distribution has a non-Gaussian shape, implying that the $\pm 2\sigma$ interval typically corresponds to a probability of less than 95.45% as would be the case for a Gaussian.

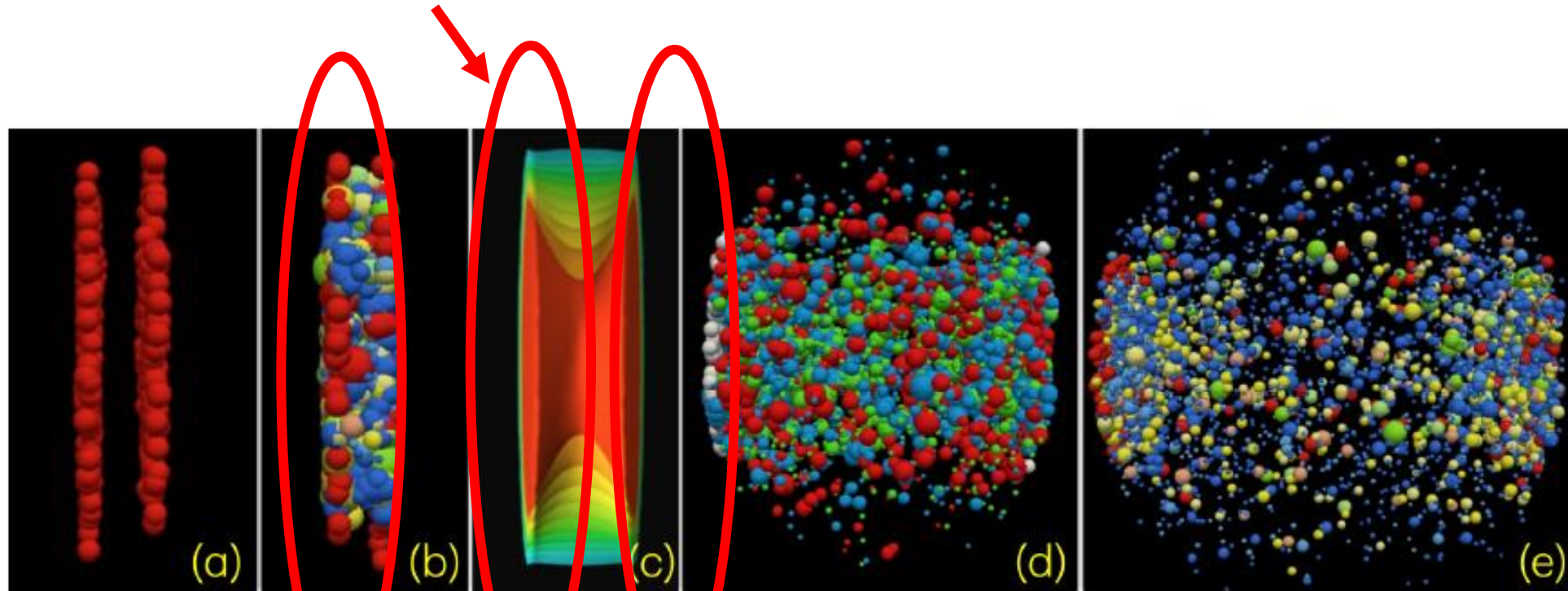
Asymmetric, and v_2 can even go negative



New photon sources considered

Weak magnetic γ from QGP

Add “unconventional” sources



Pre-equilibrium

Radiative hadronization

Stepping past the thermal paradigm



Comparisons of photon v_2 to recent models

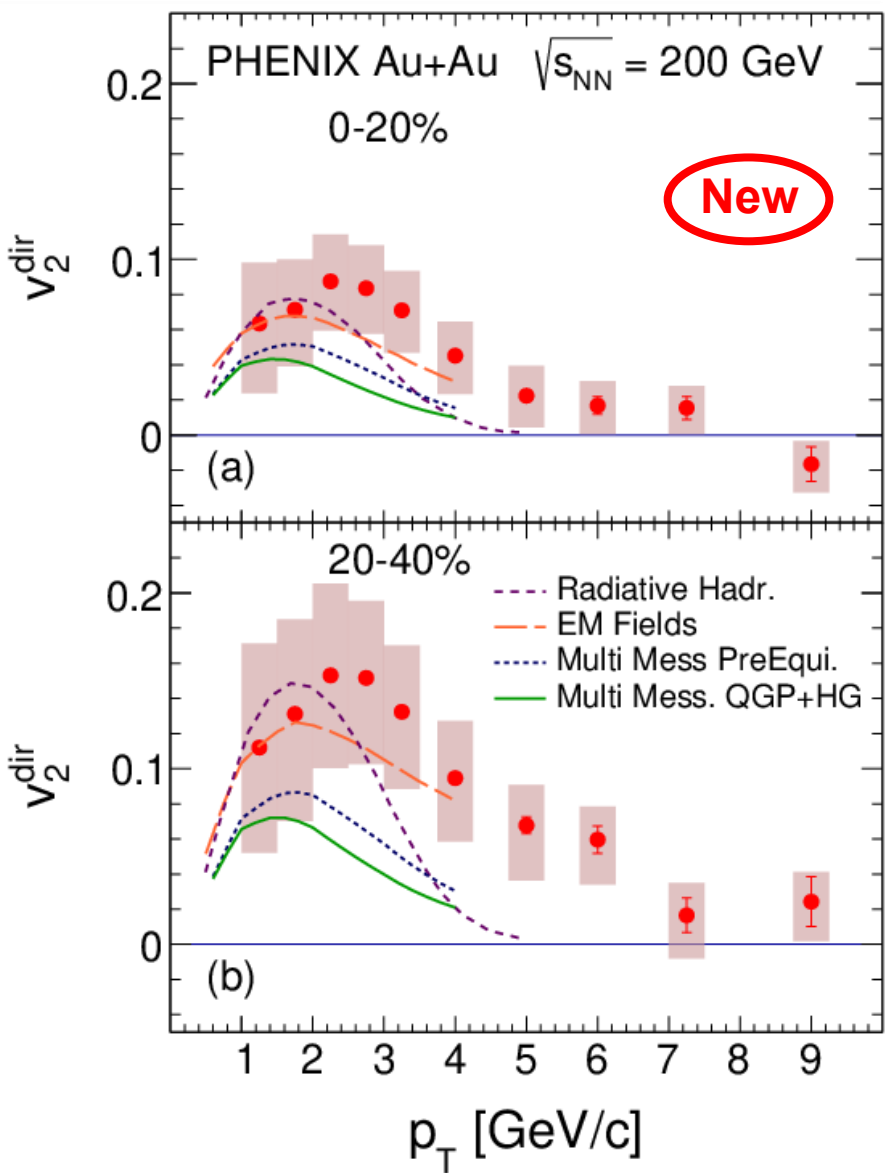


Data from arXiv:2504.02955

Different approaches

- Multi-messenger PRC 105, 014909 (2022)
- Radiative hadronization PRC 106, 034906 (2022)
- Magnetic emission Nucl. Phys. Rev 41, 1 (2024)

source	multim.	rad. hadr.	magnetic
prompt	X	X	X
magnetic			X
pre-eq	X		X
QGP th	X	X	X
HG th	X	X	X
rad. hadr.		X	



Could (should???) all those sources be combined?



A few recent models to resolve the puzzle



Strongly coupled plasma with constant magnetic field



An early attempt

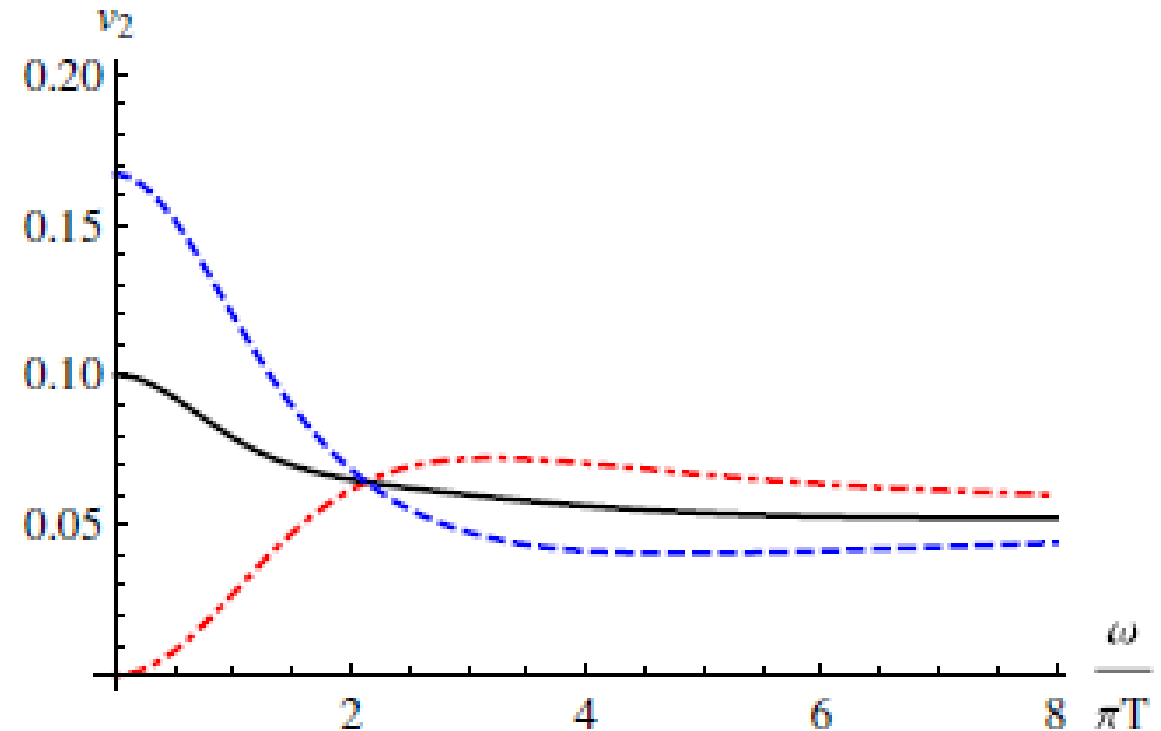
Mueller PRD 89, 026013 (2014)

Generated solely by magnetic field
in the strong coupling scenario

Photon v_2 does not vanish at low p_T

Two polarizations: in-plane (blue)
and out-of-plane (red)

Only a fraction of all v_2



Upper bound only



Non-prompt photon yields

Tuchin PRC 91, 014902 (2015)

PHENIX AuAu data

Authors conclusions:

- a significant fraction of photon excess in the region $k_{\perp} = 1\text{--}3$ GeV can be attributed to the synchrotron radiation
- this source alone would predict $v_2 = 4/7$ and $v_4 = 1/10$ independent of photon momentum and centrality
- odd terms (v_3, \dots) should vanish
- can contribute a significant fraction of non-prompt photons at 2-3 GeV/c
- temperatures should be below $T = 400$ MeV (since then synchrotron photons would account for all observed photons)

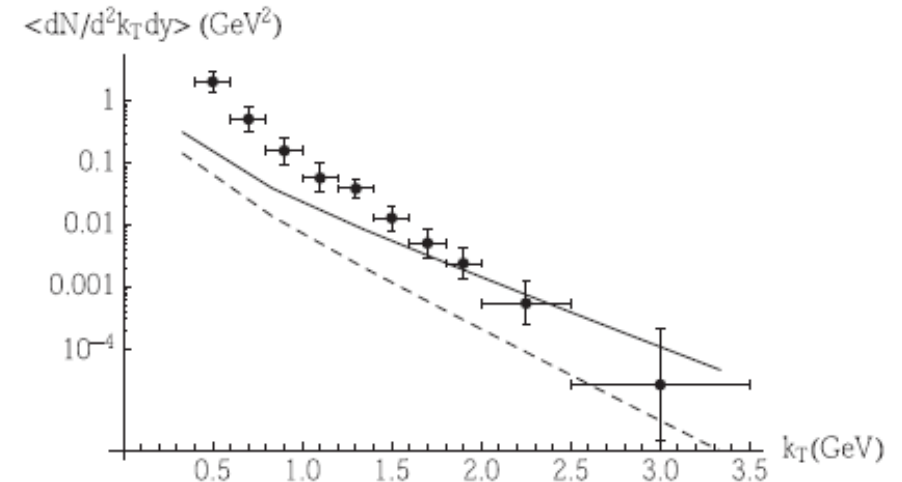


FIG. 2. Spectrum of synchrotron photons averaged over the azimuthal angle versus photon transverse momentum k_{\perp} at rapidity $y = 0$ and centrality 40%–60% ($b = 10.2$ fm [27]). Solid line: $T = 400$ MeV; dashed line: $T = 200$ MeV. Data are from [1]; they represent the direct photon k_T spectra after subtraction of the Ncoll scaled $p + p$ contribution (Fig. 8 there).

Early attempt, very incomplete

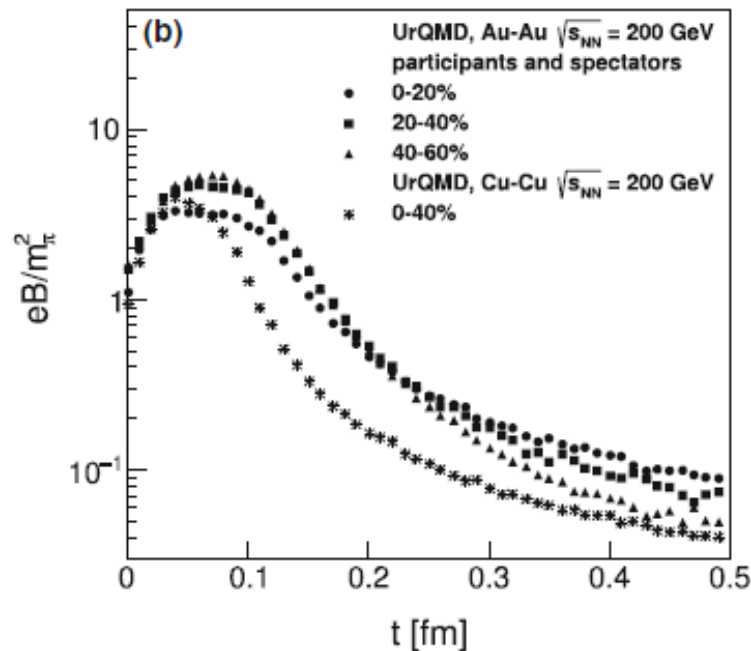
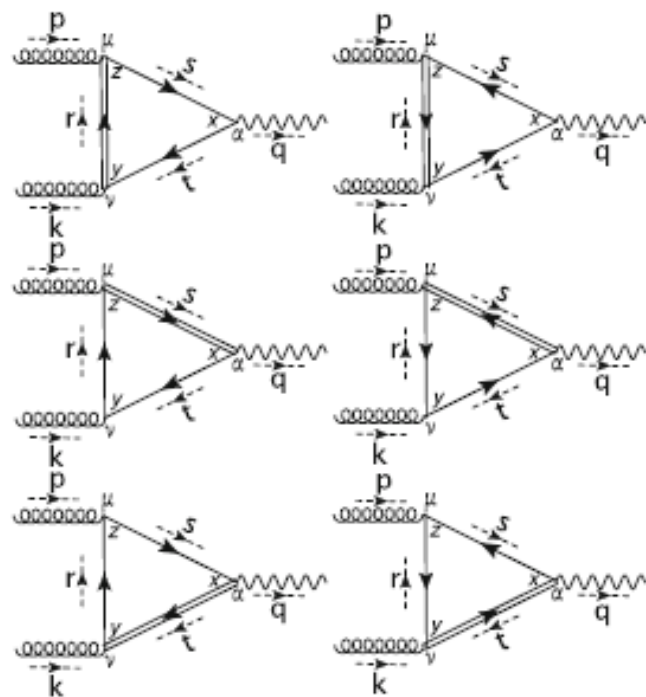


Magnetic field induced gluon fusion and splitting



Pre-hydro source

Ayala+ Eur. Phys. J. A (2020) 56:53



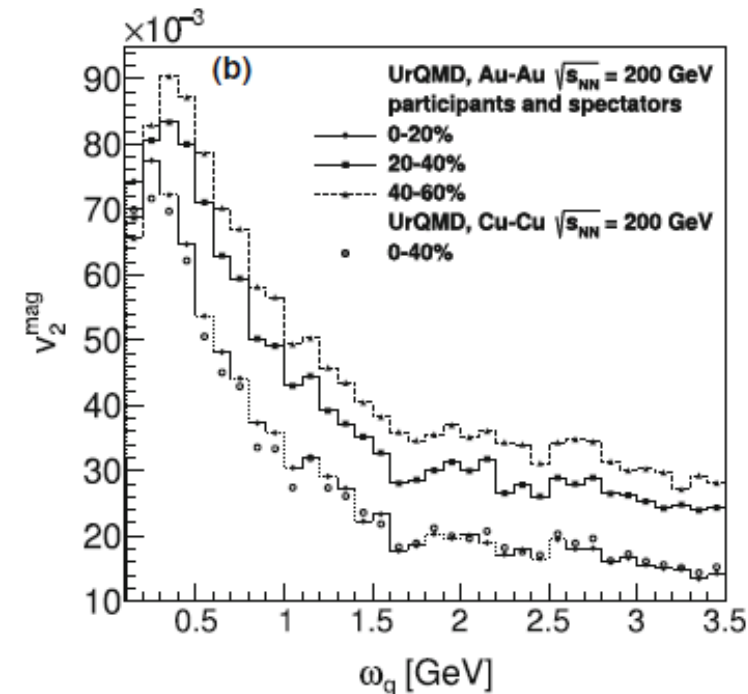
Gluon fusion / splitting

Issue: short lifetime of B

Thermal effects (hydro) start
after magnetic pulse

Similar magnitude as Mueller

Non-zero flow at low p_T !



“Magnetic flow” only –
needs to be weighted

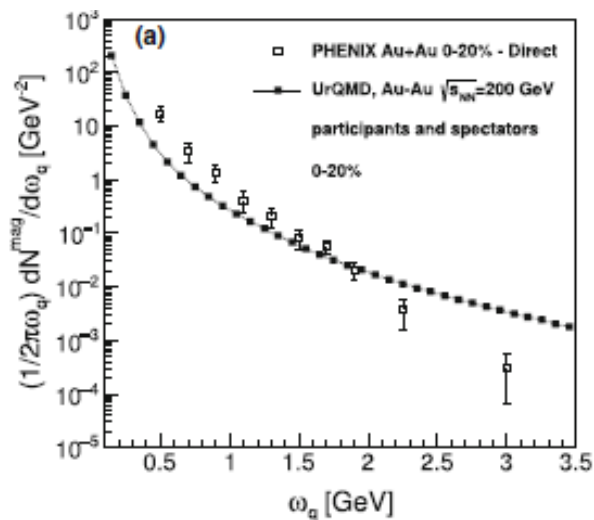
Short lifetime, $v_2(p_T=0) > 0$



Participants, spectators

Ayala+ Eur. Phys. J. A (2020) 56:53

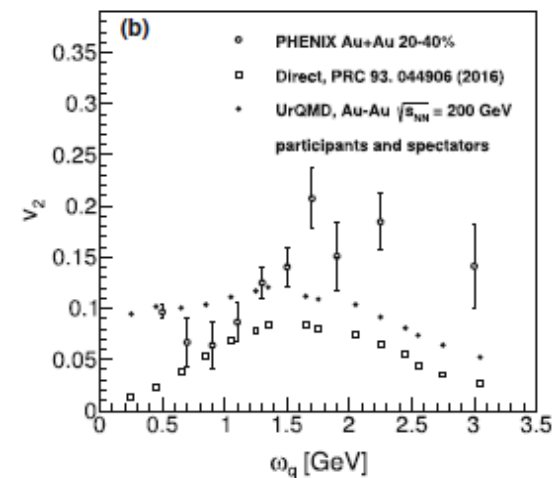
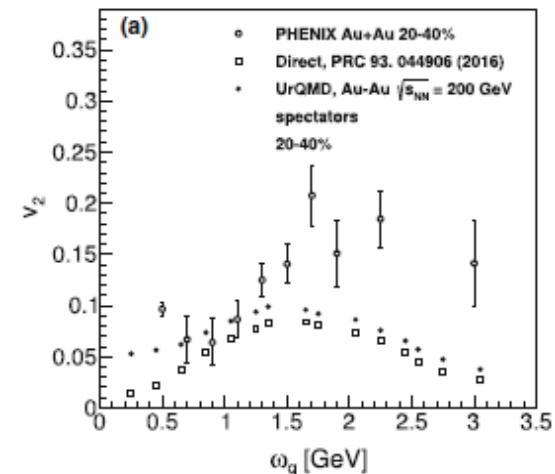
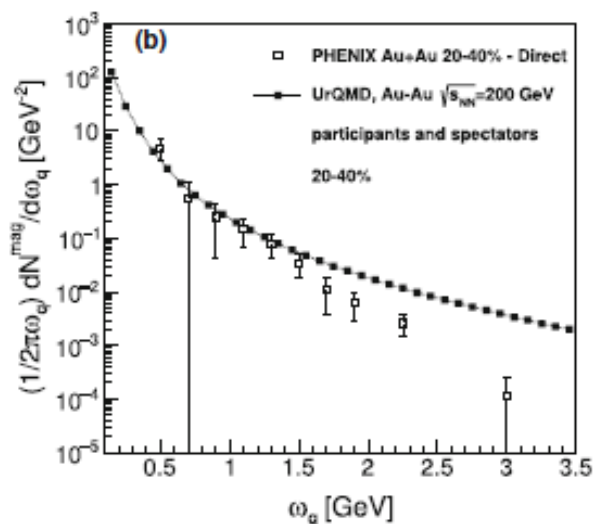
Yields and v_2 have to be explained simultaneously



Moving in the right direction
but significant shape difference

Overpredict yield, underpredict
 v_2 above 1.5 GeV/c

Enhancement at $p_T = 0$
(what is hard to measure)

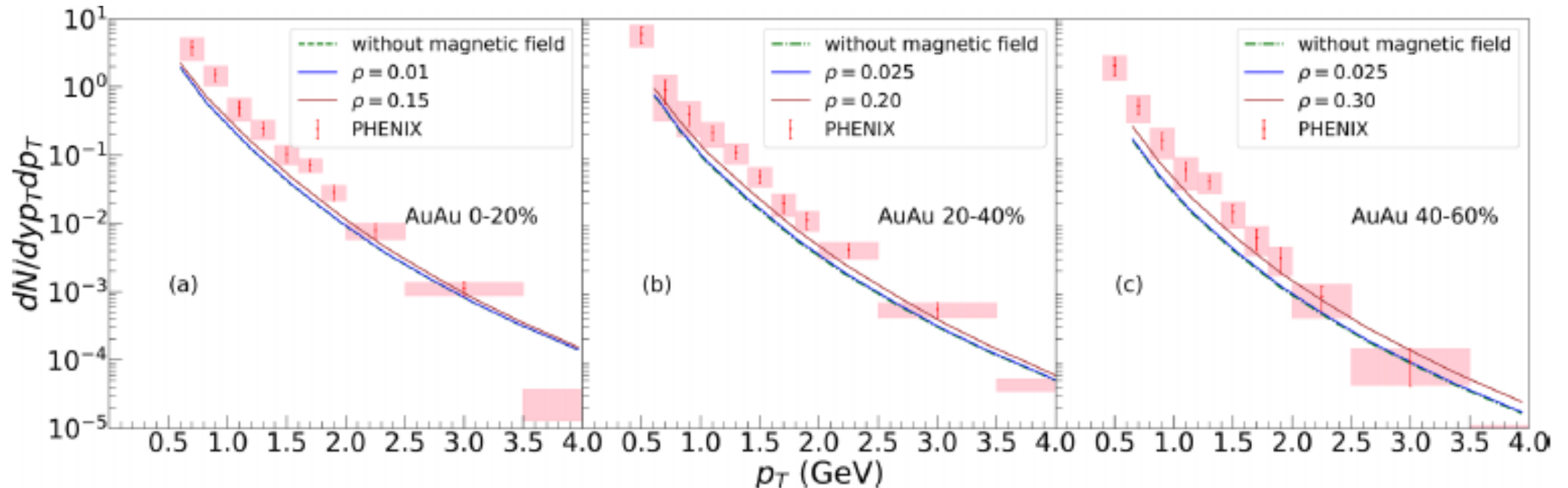


Challenge to experimentalists: get $v_2(p_T=0)$

“Weak” field in QGP, hydro, RHIC

Induced radiation, e-b-e viscous hydro

Sun, Yan PRC 109, 034917 (2024)



Effects of the weak magnetic field during QGP evolution
Interplay of magnetic field and longitudinal dynamics of medium

PHENIX data, PRC 91, 064904 (2015)

$$\rho \equiv \frac{\sigma_{\text{el}}}{T} \frac{\overline{eB_y}}{m_\pi^2},$$

Small correction to yield



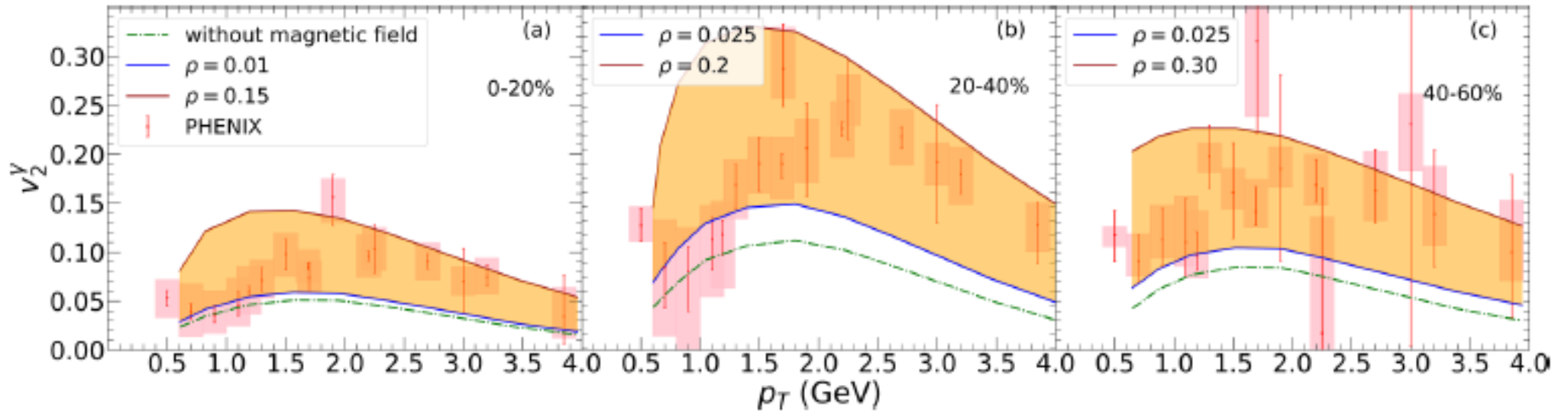


“Weak” field in QGP, hydro, RHIC

Induced radiation, e-b-e viscous hydro

Sun, Yan PRC 109, 034917 (2024)

$$\rho \equiv \frac{\sigma_{\text{el}}}{T} \frac{e\overline{B}_y}{m_\pi^2},$$



v_3 emerges (novel)
Yield increase marginal
flow increase large

PHENIX data PRC 94, 064901 (2016)

Get ρ (field strength) from flow

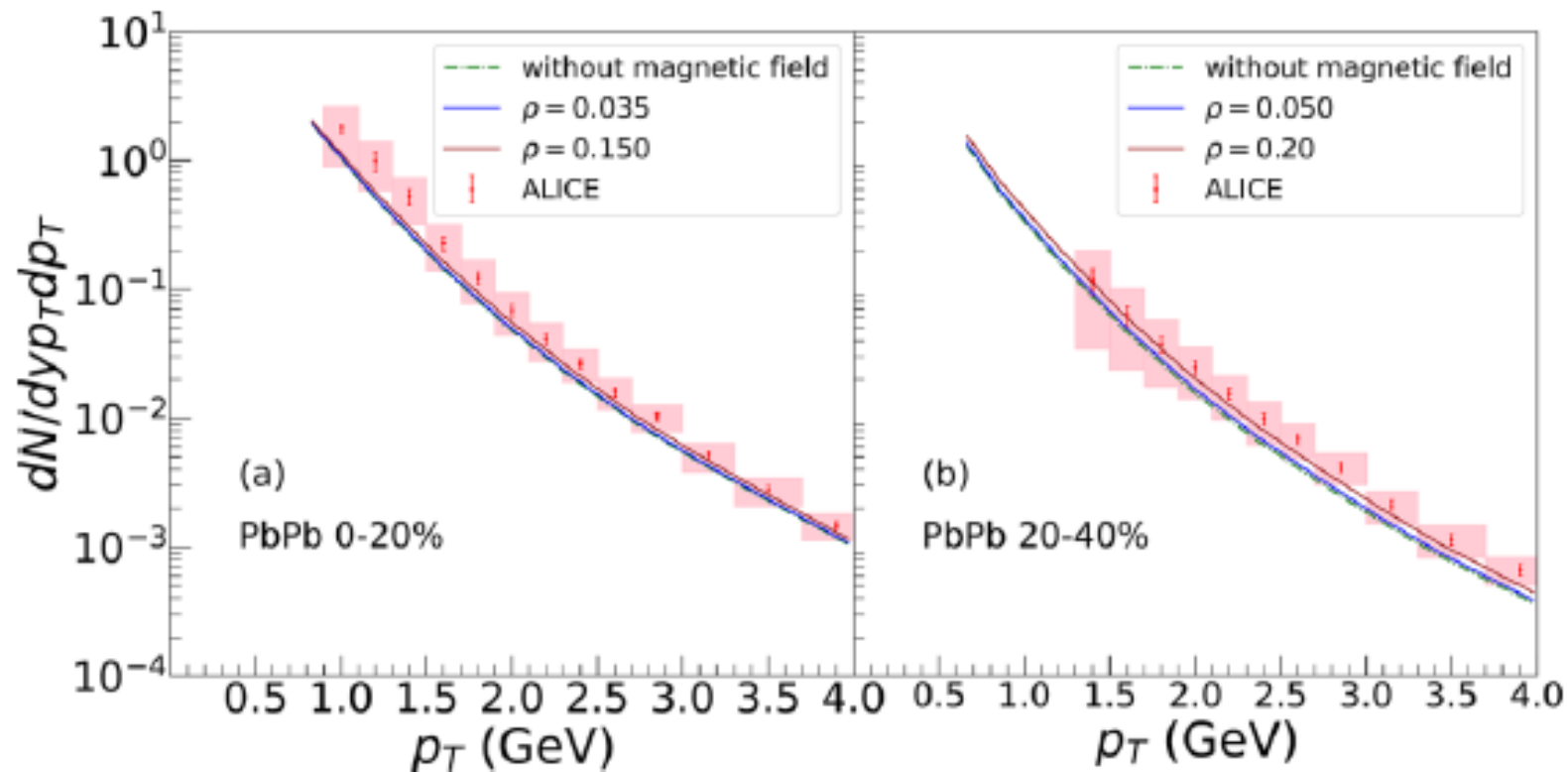
Field $0.1m_\pi^2 \sim 10^{16}$ G



“Weak” field in QGP, hydro, LHC

Induced radiation, e-b-e viscous hydro

Sun, Yan PRC 109, 034917 (2024)



ALICE, PbPb, 2.76 TeV (data PLB 789, 308 (2019))

Small correction to yield



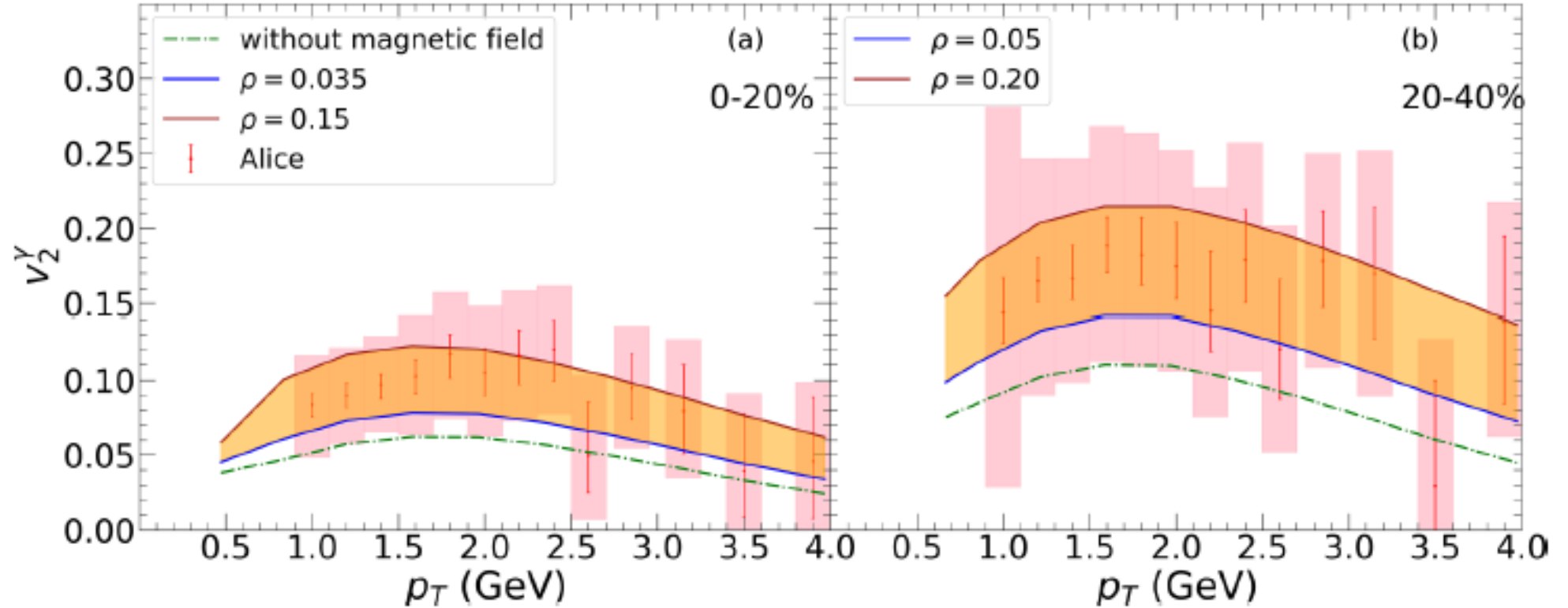


“Weak” field in QGP, hydro, LHC

Induced radiation, e-b-e viscous hydro

Sun, Yan PRC 109, 034917 (2024)

$$\rho \equiv \frac{\sigma_{\text{el}}}{T} \frac{e \overline{B}_y}{m_\pi^2},$$



ALICE, PbPb, 2.76 TeV (data PLB 789, 308 (2019))

ρ (field strength) somewhat higher than at RHIC

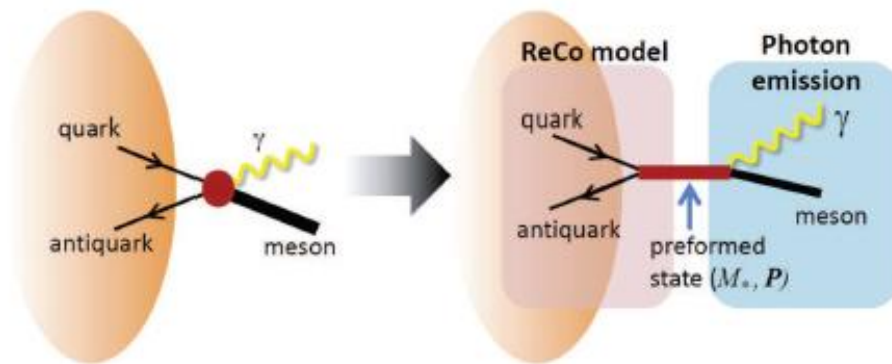


Radiative hadronization – RHIC

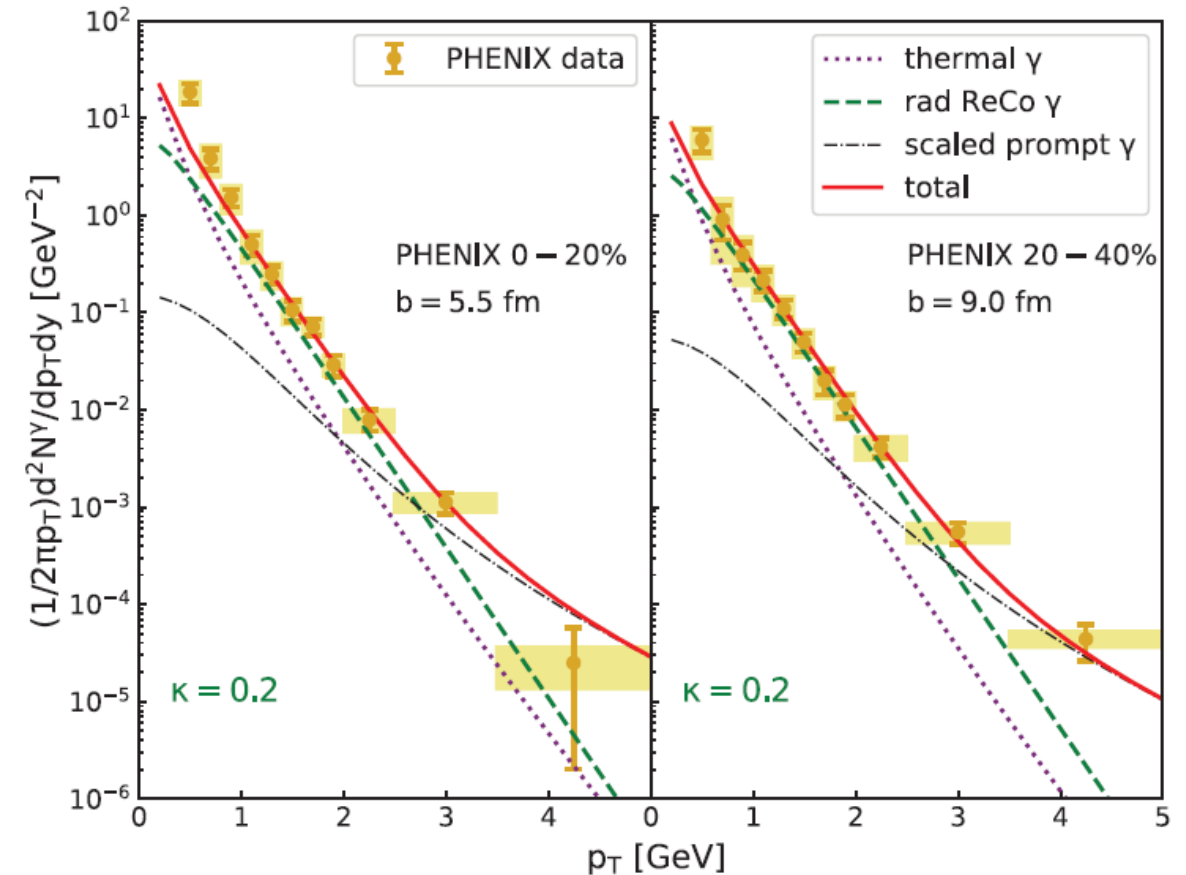
Circumvents the “old paradigm”

Old paradigm: yield early, high T , flow late
 Here: new source **at the time of hadronization**
 Recombination off-shell, then come on-shell
 by radiating a photon
 Will inherit v_2 at the time of phase transition

κ chosen so that thermal, prompt and radiative photons describe the data



Fujii et al, PRC 106, 034906 (2022)



PHENIX, PRC 91, 064904 (2015)

κ is adjusted to describe data

Radiative hadronization – RHIC

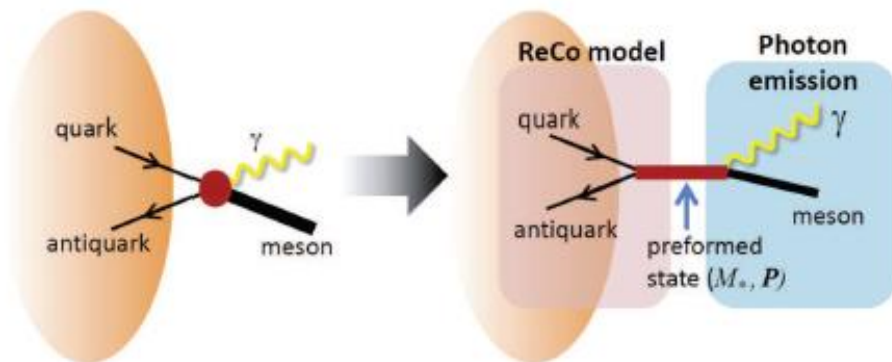
What about v_2 ?

Old paradigm: yield early, high T, flow late

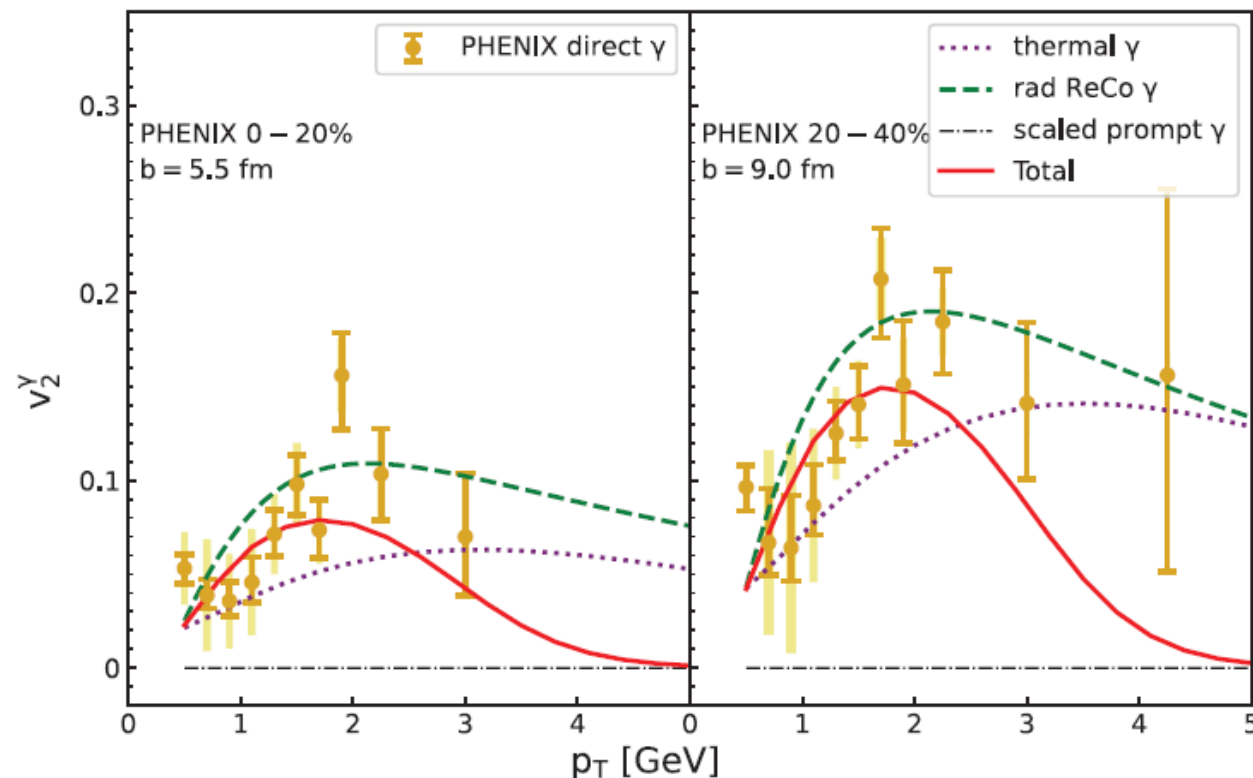
Here: new source **at the time of hadronization**

Recombination off-shell, then come on-shell by radiating a photon

Will inherit v_2 at the time of phase transition



Fujii et al, PRC 106, 034906 (2022)



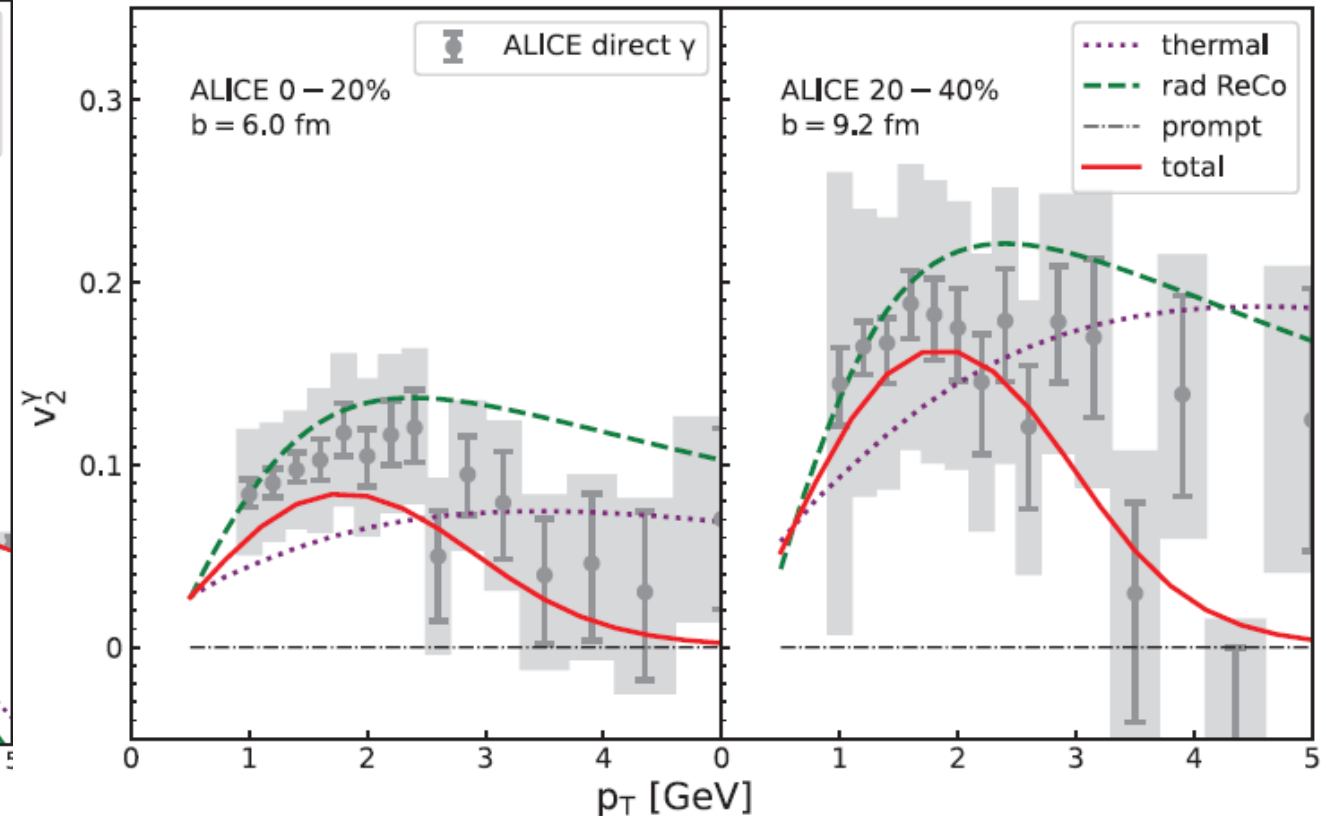
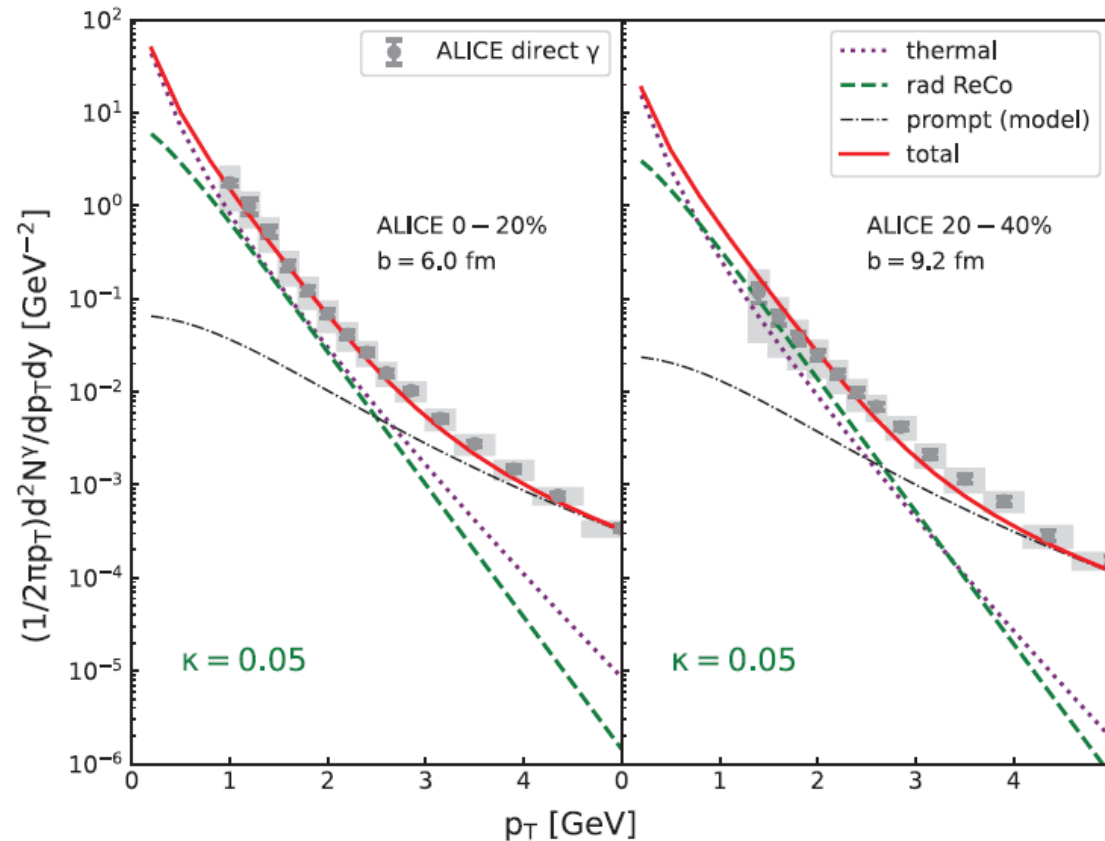
PHENIX, PRC 94, 064901 (2016)

Starts to fail above 2 GeV/c

Radiative hadronization – LHC

Does it work at LHC? (should!)

Fujii et al, PRC 106, 034906 (2022)



Remember: the transition itself at RHIC and LHC appeared to be similar (low p_T slopes, scaling)

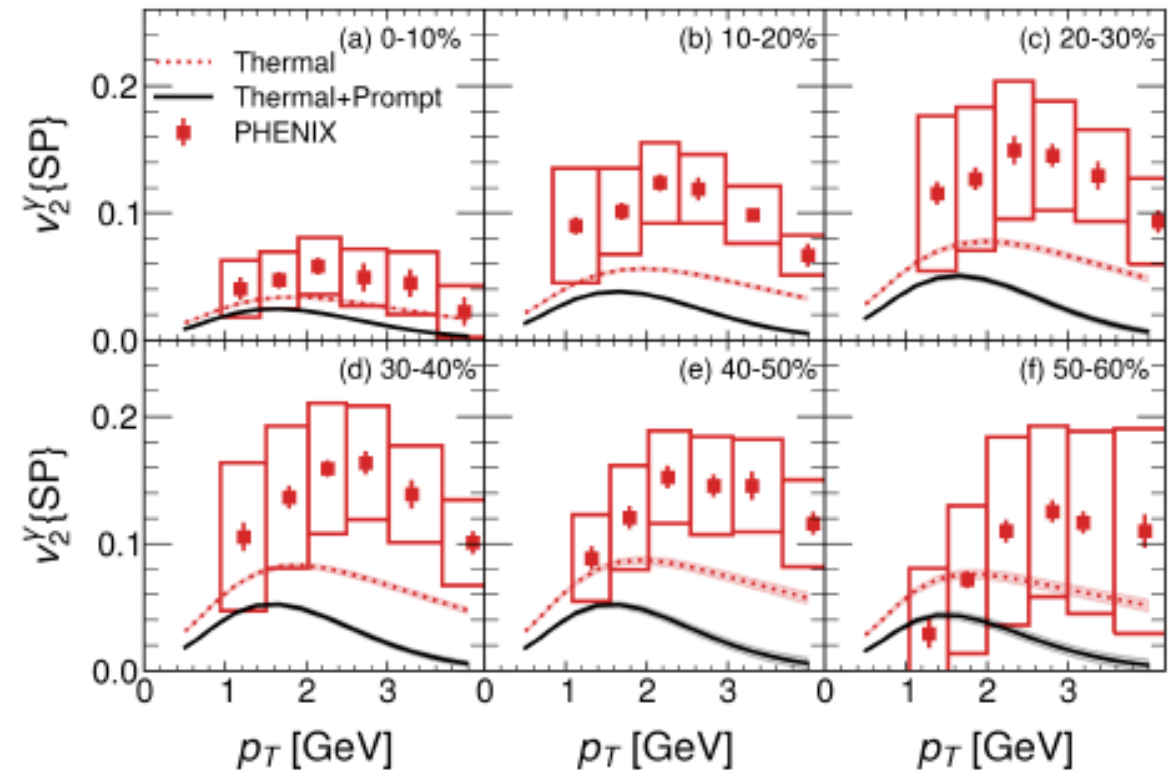
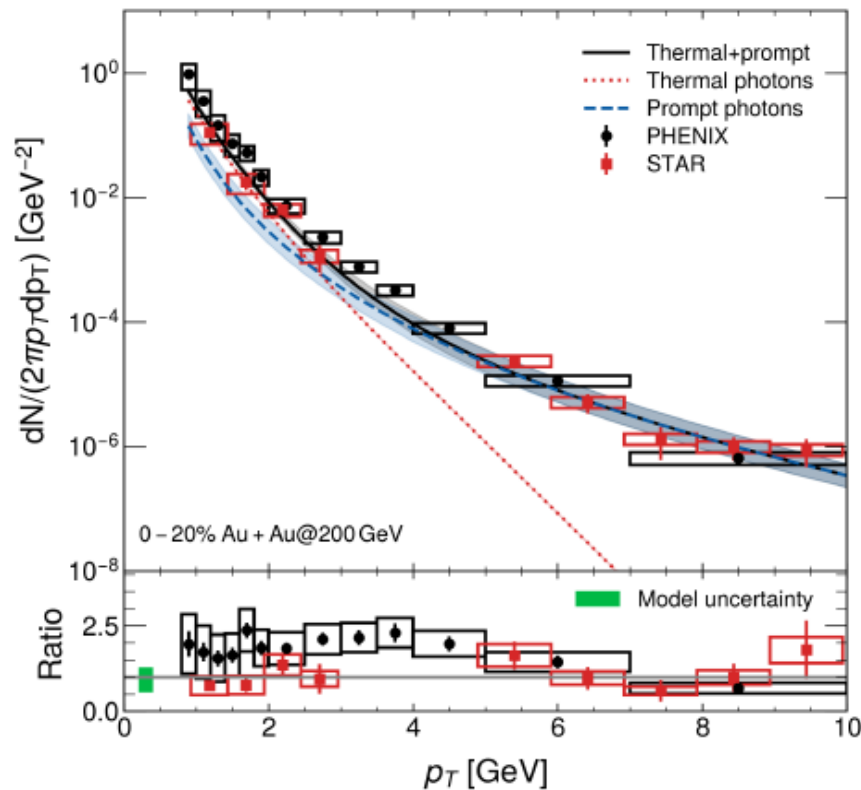
κ very different at RHIC and LHC



Multistage (multi-messenger) model

Pre-equilibrium, MUSIC, UrQMD afterburner

Gale et al. arXiv 2511.08773



Add rad. hadr., magnetic?





Where does this leave us?



Summary

Photon at high p_T well understood

At low p_T the direct photon puzzle still not resolved

Interesting and plausible new sources suggested

Experimental uncertainties still too large

We never had an honest-to-God direct photon experiment,
built without compromises to other physics (and we are paying the price)

We should have one!

Further argument in GD *Rept.Prog.Phys.* 83 (2020) 4, 046301





Backup



