



Measurement of the Higgs boson decay in $Z(\ell\ell)H, H \rightarrow WW$ at the FCC-ee at $\sqrt{s} = 240 \text{ GeV}$

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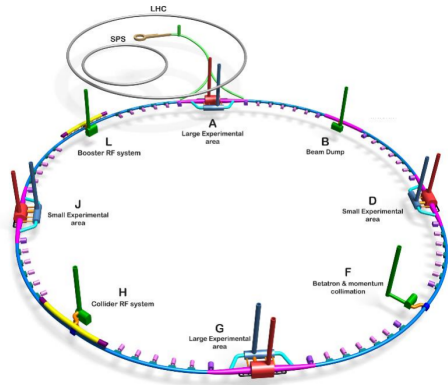
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Outline

- 1 FCC-ee: a Higgs factory
- 2 Relevance of $e^+e^- \rightarrow ZH$
- 3 Analysis of $e^+e^- \rightarrow Z(\ell\ell)H, H \rightarrow W(qq)W(qq)$
- 4 Results
- 5 Conclusions

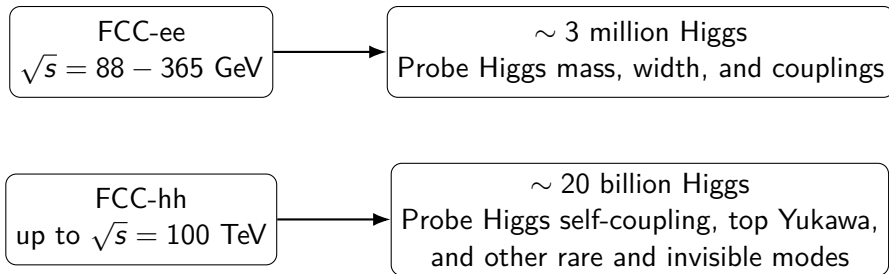


FCC opportunities*

- Map **Higgs and EW gauge boson properties** with unprecedented precision, probing interactions linked to the early universe ($10^{-12} - 10^{-10}$ s after the Big Bang).
- Improve knowledge of Standard Model phenomena (EW, QCD, flavour, Higgs, top) through **high-precision measurements**.
- Boost sensitivity to **rare low-energy processes**, including light, weakly coupled particles (e.g., neutrinos, axion-like particles), and decisively test dark matter candidates.
- Extend direct discovery reach for **new particles** by at least an order of magnitude at high energies.

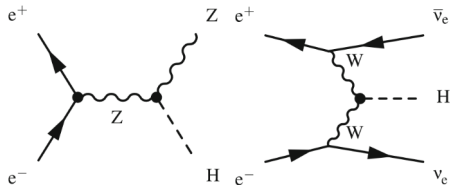
*See more information at the **Future Circular Collider Feasibility Study Report: Volume 1, Physics, Experiments, Detectors**. DOI: [10.17181/CERN.9DKX.TDH9](https://doi.org/10.17181/CERN.9DKX.TDH9).

FCC-ee and FCC-hh



Higgs production from e^+e^- colliders at $\sqrt{s} = 240$ GeV

In e^+e^- collisions at $\sqrt{s} = 240$ GeV, Higgs are mainly produced through two modes: "Higgsstrahlung" and WW fusion process (VBF):



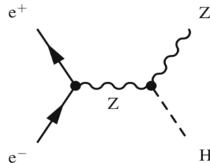
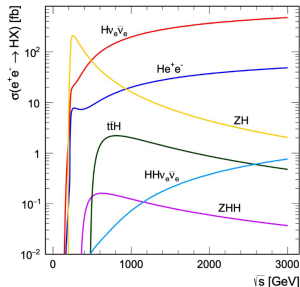
Higgs production modes for e^+e^- colliders at $\sqrt{s} = 240$ GeV *.

At the FCC-ee: significantly larger number of ZH events with a much smaller background will be observed compared with previous observations by ATLAS and CMS.

*Eysermans, J., & Selvaggi, M. (2025). DOI: 10.17181/9gzts-hsm97

Higgsstrahlung

Higgs model-independent studies can be carried out in Higgsstrahlung events: center-of-mass is known, information on the recoiling Higgs is accessible by the associated Z boson *.



At $\sqrt{s} = 240$ GeV, the ZH rate is enhanced, sensitive to the eeZ and the ZZH vertices, along with the couplings to the Higgs. Images from: Sally Dawson, [ZH at one loop in SMEFT](#)

*Eysermans, J., Bernardi, G., Ang, L. (2025). DOI: 10.17181/jfb44-s0d81

Higgs couplings

From the Higgsstrahlung and WW fusion processes we can extract the different couplings of various Higgs decay modes ($H \rightarrow X\bar{X}$) using the following relations*:

$$\sigma_{ZH} \times Br(H \rightarrow X\bar{X}) \propto \frac{g_{HZZ}^2 g_{HXX}^2}{\Gamma_H} \quad (1)$$

$$\sigma_{H\nu_e\nu_e} \times Br(H \rightarrow X\bar{X}) \propto \frac{g_{HWW}^2 g_{HXX}^2}{\Gamma_H}$$

By analyzing the $H \rightarrow ZZ$ decay, we can extract the g_{HZZ} coupling from the first equation, which can then be used to obtain all other Higgs couplings for other decay channels.

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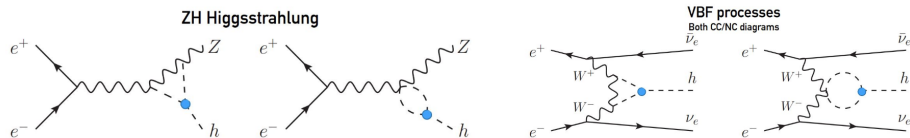
Other opportunities

Apart from determining the Higgs couplings with great precision, studying the ZH process will allow to

- Constraint the Higgs total width, by measuring the $H \rightarrow ZZ^*$ decay and the relation:

$$\Gamma_H \propto \frac{\sigma(e^+e^- \rightarrow ZH, H \rightarrow ZZ)^2}{\sigma(e^+e^- \rightarrow ZH)}$$

- Probe the Higgs self-coupling through NLO deviations*:



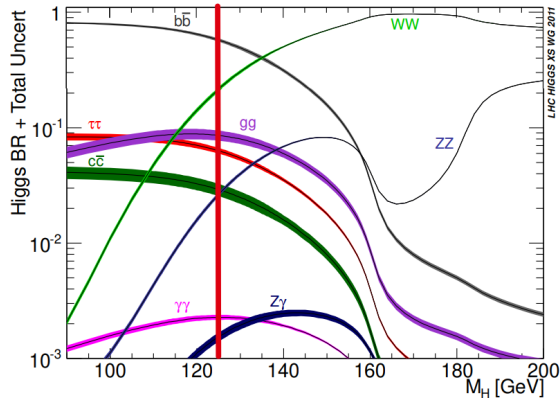
*Eysermans, J., Selvaggi, M. (2025) DOI:10.17181/9gzts-hsm97

$Z(\ell\ell)$ and $H \rightarrow WW$

- **Leptonic decays** of the Z boson $\rightarrow$ greater precision and efficient selection of the ZH events^a.
- **Mass recoil** against the lepton pair in the final state (m_{recoil}) $\rightarrow$ distinguish the ZH events from the main backgrounds ($WW, ZZ, Z\gamma$):

$$m_{recoil}^2 = (\sqrt{s} - E_{\ell\ell})^2 - p_{\ell\ell}^2 = s - 2E_{\ell\ell}\sqrt{s} + m_{\ell\ell}^2 \quad (2)$$

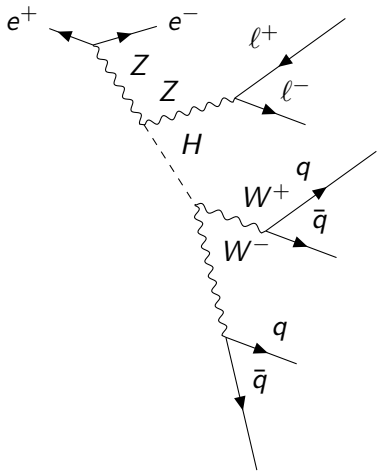
- Our study will focus on the $H \rightarrow W^+W^-$ decay, particularly the **hadronic channel** ($\Gamma_{W,hadrons}/\Gamma_W \sim 67\%$).



From: *Handbook of LHC Higgs Cross Sections*,
DOI: 10.5170/CERN-2012-002

^aEysermans, J., Bernardi, G., Ang, L. (2025). DOI: 10.17181/jfb44-s0d81

Final State of the Studied Process



We study the process

$$e^+e^- \rightarrow ZH \rightarrow Z(\ell^+\ell^-)H(W^+W^-) \rightarrow Z(\ell^+\ell^-)W(qq)W(qq).$$

For the final state we look for events with:

- 1 Two opposite-charged electrons/muons.
- 2 At least one isolated, with $p > 20$ GeV.
- 3 Four jets.

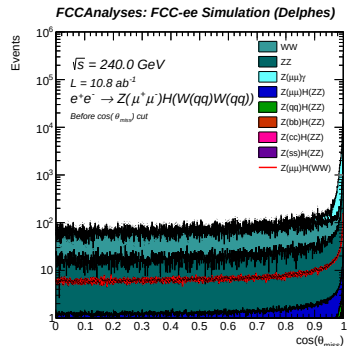
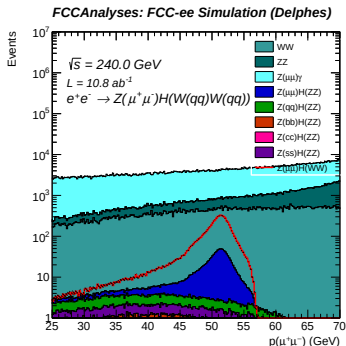
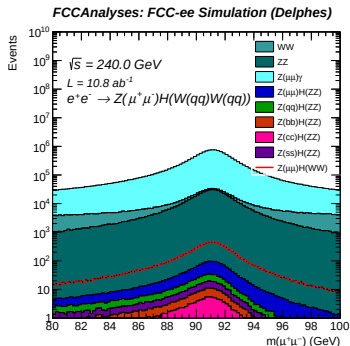
Signal and Background Samples

Monte Carlo event samples generated with WHIZARD and PYTHIA8, passed through a parametric simulation of the IDEA detector using DELPHES ([winter2023 campaign](#)).

Process	Simulation Label	$\sigma[fb]$	Events Generated
Signal:			
$e^+e^- \rightarrow Z(\mu\mu)H, H \rightarrow WW$	wzp6_ee_mumuH_HWW_ecm240	1.456	400,000
$e^+e^- \rightarrow Z(ee)H, H \rightarrow WW$	wzp6_ee_eeH_HWW_ecm240	1.541	400,000
Background:			
$e^+e^- \rightarrow ZZ$	p8_ee_ZZ_ecm240	1358.99	11,300,000
$e^+e^- \rightarrow WW$	p8_ee_WW_ecm240	16438.5	74,728,800
$e^+e^- \rightarrow Z(\mu\mu)\gamma$	wzp6_ee_mumu_ecm240	5288	10,700,000
$e^+e^- \rightarrow ee$	wzp6_ee_ee_Mee_30_150_ecm240	8305	85,400,000
$e^+e^- \rightarrow Z(\mu\mu)H, H \rightarrow ZZ$	wzp6_ee_mumuH_HZZ_ecm240	0.1786	400,000
$e^+e^- \rightarrow Z(ee)H, H \rightarrow ZZ$	wzp6_ee_eeH_HZZ_ecm240	0.1891	400,000
$e^+e^- \rightarrow Z(qq)H, H \rightarrow ZZ$	wzp6_ee_qqH_HZZ_ecm240	1.409	1,200,000
$e^+e^- \rightarrow Z(bb)H, H \rightarrow ZZ$	wzp6_ee_bbH_HZZ_ecm240	0.7915	1,000,000
$e^+e^- \rightarrow Z(cc)H, H \rightarrow ZZ$	wzp6_ee_ccH_HZZ_ecm240	0.6164	1,200,000
$e^+e^- \rightarrow Z(ss)H, H \rightarrow ZZ$	wzp6_ee_ssH_HZZ_ecm240	0.7912	600,000

$Z(I^+I^-)H$ event selection: kinematic cuts

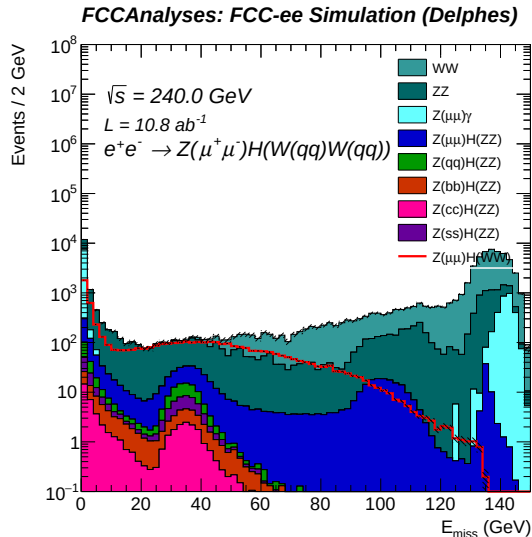
- 1 2 op. ch. muons, $p > 20$ GeV, at least one isolated
- 2 Veto on ℓ 's with $5 < p < 20$ GeV
- 3 $m_{\ell\ell} \in [86, 96]$ GeV, $p_{\ell\ell} \in [30, 57]$ GeV
- 4 $|\cos \theta_{\text{miss}}| < 0.98$



Jet Clustering and Di-Jet Pair Selection

- Restrict the $W \rightarrow \ell\nu$ decays by cutting events with missing energy $E_{miss} > 10$ GeV.
- Create a collection of particles removing both muons and electrons with $p > 20$ GeV.
- Exclusive $N = 4$ “Durham” algorithm is used for jet clustering, which is a sequential spherical k_t clustering method^a.

^aS. Catani et al Physics Letters B, 269(3):432–438, 1991



Jet Clustering and Di-Jet Pair Selection

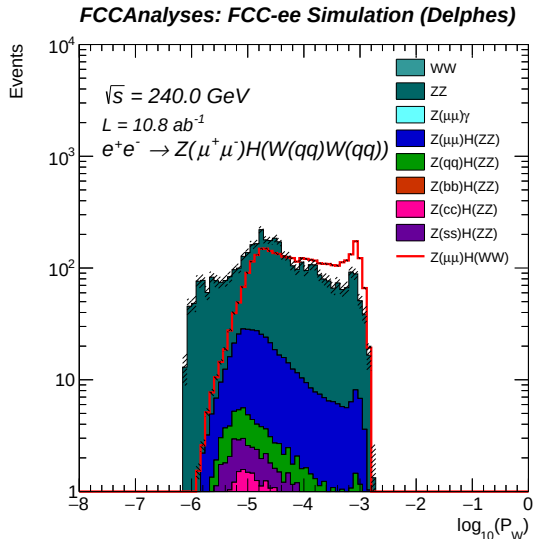
We choose the best dijet pair to reconstruct the two W bosons by maximizing^a :

$$P_{W}^{jj} = \frac{M_W^2 \Gamma_W^2}{\left(m_{jj}^2 - M_W^2\right)^2 + M_W^2 \Gamma_W^2}. \quad (3)$$

- $jj[0]$: Maximizing pair ("on-shell").
- $jj[1]$: The remaining pair ("off-shell").

Same procedure but with M_Z and Γ_Z instead: use the resulting dijet-pair invariant mass ($m_{jj[0]}^Z$) to further reduce the ZZ background.

^aPiotr Niezurawski et al JHEP11(2002)034

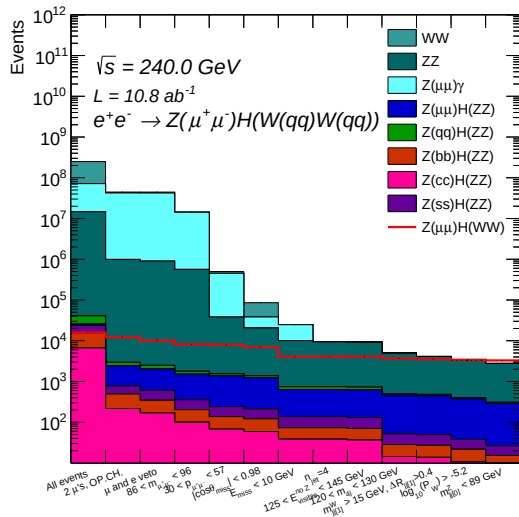


Final Selection Cuts

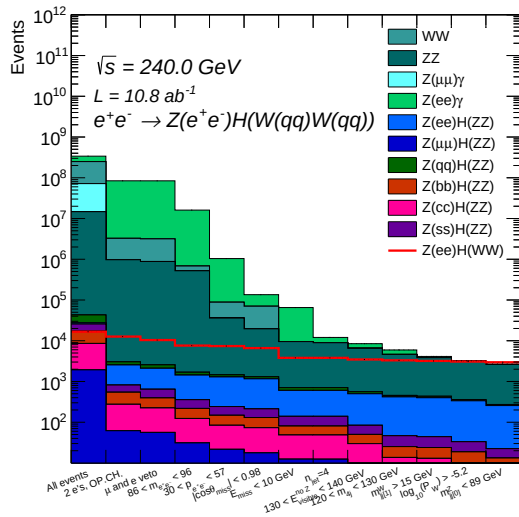
- Cut 1: Events with only two opposite charge μ' s.
- Cut 2: Veto on μ' s and e' s with momentum within $[5, 20]$ GeV.
- Cut 3: $86 < m_{\mu\mu} < 89$ GeV.
- Cut 4: $30 < p_{\mu\mu} < 57$ GeV.
- Cut 5: $|\cos\theta_{miss}| < 0.98$.
- Cut 6: Missing energy $E_{miss} < 10$ GeV.
- Cut 7: $n_{jet} = 4$.
- Cut 8: $125 < E_{visible}^{no Z} < 145$ GeV, where $E_{visible}^{no Z}$ is the visible energy of the particles excluding the two Z muons.
- Cut 9: $120 < m_{4j} < 130$ GeV, where m_{4j} is the four-jet invariant mass.
- Cut 10: $m_{jj[1]} > 15$ GeV and $\Delta R_{jj[1]} > 0.4$, where $m_{jj[1]}$ is the invariant mass of the off-shell-W jet pair, and $\Delta R_{jj[1]}$ is the corresponding ΔR .
- Cut 11: $\log_{10} P_W > -5.2$, where P_W is the quantity described previously.
- Cut 12: $m_{jj[0]}^Z < 89$ GeV, where $m_{jj[0]}^Z$ is the invariant mass of the on-shell dijet pair selected to align to the mass of the Z boson.

$Z(l^+l^-)H, H \rightarrow W(qq)W(qq)$ **Cut Flow**

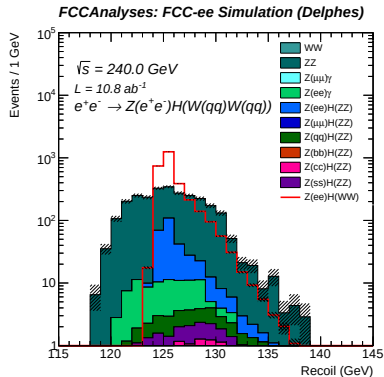
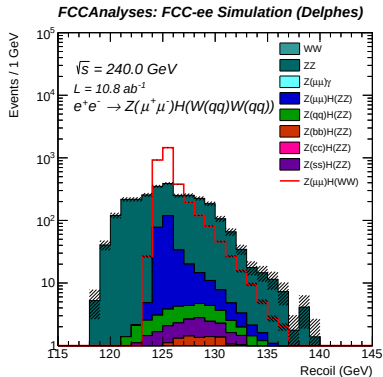
FCCAnalyses: FCC-ee Simulation (Delphes)



FCCAnalyses: FCC-ee Simulation (Delphes)



Final Results: Fit to recoil mass after selections



Uncertainty after best-fit
 (with log-normal
 uncertainty of 3% for
 $H \rightarrow ZZ$ background and
 1% for the rest):

- $\mu^+\mu^-$ channel: 2.15.
- e^+e^- channel: 2.30.

The uncertainty for the combined contributions is 1.60.

Results in Technical Note (DOI:10.17181/x48j4-3kg84) part of the Future Circular Collider Feasibility Study Report, to be presented to the European Strategy for Particle Physics (ESPP).

Conclusions

- We have evaluated the **sensitivity of the measurement** of the Higgs boson branching ratio to the hadronic decay mode of the W^+W^- final state at the Future Circular Collider.
- The signal was extracted via a **binned likelihood fit** to the recoil mass distribution, utilizing the Combine tool. The fit was applied across the full range of the recoil mass, with a bin width of 1 GeV.
- The uncertainties in the measured branching ratio for the $H \rightarrow W^+W^-$ hadronic decay mode were found to be 2.15 and 2.30 for the $Z(\mu\mu)$ and $Z(ee)$ channels, respectively, with a **combined uncertainty of 1.60**.
- These results highlight the potential of the Future Circular Collider to provide a **precise determination of the Higgs branching ratio** to a pair of W bosons, with minimal uncertainty.

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THANK YOU!

Back-up Slides

Multivariate Analysis and BDT training

The classification of signal events in the $Z(\ell\ell)H, H \rightarrow WW$ process at FCC-ee is performed using a Multivariable Analysis (MVA) approach based on Boosted Decision Trees (BDT), implemented via the XGBoost library.

Features:

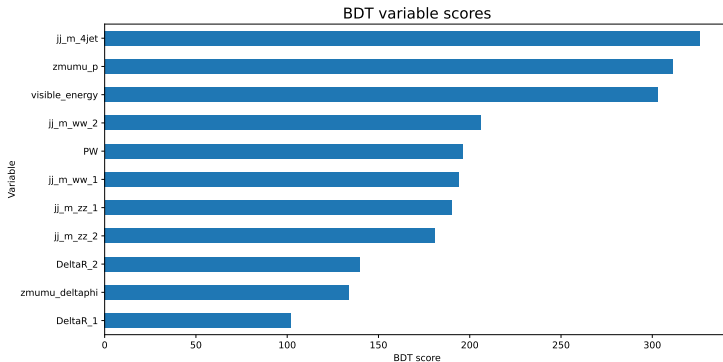
- **Four-jet invariant mass:** jj_m_4jet ,
- **On-shell-W dijet invariant mass:** $jj_m_ww_1$,
- **Off-shell-W dijet invariant mass:** $jj_m_ww_2$,
- **Probability of dijet pairings to correspond to the W boson:** PW ,
- **On-shell-Z dijet invariant mass:** $jj_m_zz_1$,
- **Off-shell-Z dijet invariant mass:** $jj_m_zz_2$
- **$Z_{\mu\mu}$ kinematics:** momentum ($zmumu_p$) and $\Delta\phi$ ($zmumu_deltaphi$).
- **Visible energy (excluding Z muons):** $visible_energy$.
- **On-shell-W and Off-shell-W dijet pairs separation:** ΔR_1 ($DeltaR_1$), ΔR_2 ($DeltaR_2$), respectively.

Hyperparameters:

- **objective:** binary:logistic
- **eval_metric:** auc
- **eta:** 0.1
- **max_depth:** 5
- **subsample:** 0.5
- **colsample_bytree:** 0.5
- **n_estimators:** 350
- **learning_rate:** 0.20
- **gamma:** 3
- **min_child_weight:** 10
- **early_stopping_rounds:** 25

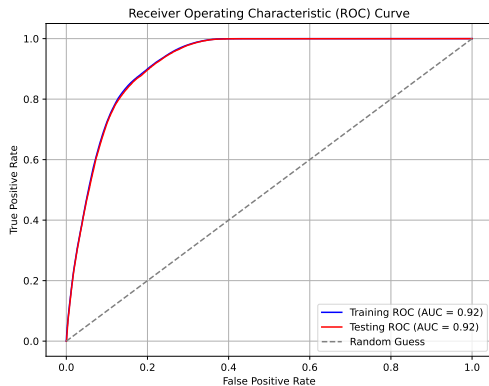
BDT variable scores

Scores for each of the features chosen for the classifier in the BDT training. A higher score suggests that the variable is more valuable for classification purposes.



ROC Curve

Receiver Operating Characteristic (ROC) curve. In our case, the value of the Area Under the Curve (AUC) indicates that the ability of the model to distinguish between signal and background is acceptable (above 0.9).



BDT Score Distribution

Distribution of the scores produced by the BDT classifier for the signal (blue) and background (red).

