





New Physics Prospects in Higgs Couplings

RADPyC 2025

A. I. Hernández-Juárez, R Gaitán alan.hernandezjua@alumno.buap.mx May 22, 2025

The $H \to Z\gamma$ decay

PHYSICAL REVIEW LETTERS 132, 021803 (2024)

Editors' Suggestion

Featured in Physics

Evidence for the Higgs Boson Decay to a Z Boson and a Photon at the LHC

G. Aad et al.* (ATLAS and CMS Collaborations)

(Received 8 September 2023; accepted 27 November 2023; published 11 January 2024)

The first evidence for the Higgs boson decay to a Z boson and a photon is presented, with a statistical significance of 3.4 standard deviations. The result is derived from a combined analysis of the searches performed by the ATLAS and CMS Collaborations with proton-proton collision datasets collected at the CERN Large Hadron Collider (LHC) from 2015 to 2018. These correspond to integrated luminosities of around 140 fb⁻¹ for each experiment, at a center-of-mass energy of 13 TeV. The measured signal yield is 2.2 ± 0.7 times the standard model prediction, and agrees with the theoretical expectation within 1.9 standard deviations.

DOI: 10.1103/PhysRevLett.132.021803

The $H \rightarrow Z\gamma$ decay

The signal strength is defined as

$$\mu_i^{Z\gamma} = \frac{\sigma_i \mathscr{B}^{Z\gamma}}{(\sigma_i) \mathsf{SM}^{(\mathscr{B}^{Z\gamma})} \mathsf{SM}} = 2.2 \pm 0.7,$$

- σ_i is the cross-section of the Higgs production
- $\mathscr{B}^{Z\gamma}$ is the branching ratio of the $H \to Z\gamma$ decay

The $H \to Z\gamma$ decay



- σ_i is the cross-section of the Higgs production
- $\mathscr{B}^{Z\gamma}$ is the branching ratio of the $H \to Z\gamma$ decay

The $H \rightarrow Z\gamma$ decay

The signal strength is defined as

$$\mu_{i}^{Z\gamma} = \frac{\sigma_{i} \mathscr{B}^{Z\gamma}}{(\sigma_{i}) \mathsf{SM}^{(\mathscr{B}^{Z\gamma})} \mathsf{SM}} = 2.2 \pm 0.7,$$
Possibilities of new physics!

- σ_i is the cross-section of the Higgs production
- $\mathscr{B}^{Z\gamma}$ is the branching ratio of the $H \to Z\gamma$ decay

The $H \rightarrow Z\gamma$ decay

The signal strength is defined as

$$\mu_i^{Z\gamma} = \frac{\sigma_i \mathscr{B}^{Z\gamma}}{(\sigma_i) \mathsf{SM}^{(\mathscr{B}^{Z\gamma})} \mathsf{SM}} = 2.2 \pm 0.7,$$

Where is the new physics?

•
$$\mathscr{B}^{Z\gamma}$$
 is the branching ratio of the $H \to Z\gamma$ decay

The $H \rightarrow Z\gamma$ decay

The Higgs production is well measured and agrees with the SM production through a top quark loop in gluon fusion.



Therefore, the new physics can only arise from the $H \rightarrow Z\gamma$ decay:

$$\mu_i^{Z\gamma} = \frac{\mathscr{B}^{Z\gamma}}{\mathscr{B}^{Z\gamma}_{SM}} = 2.2 \pm 0.7,$$

A lot of models trying to explain this excess (THDM, MSSM, left-right models, new particles,...)

The $H \rightarrow Z\gamma$ decay

The Higgs production is well measured and agrees with the SM production through a top quark loop in gluon fusion.



Therefore, the new physics can only arise from the $H \rightarrow Z\gamma$ decay:



The $H \to Z\gamma$ decay

The general form of the vertex function $\Gamma^{\mu\nu}_{Z\gamma H}$ is given as follows

$$\Gamma^{\mu\nu}_{Z\gamma H} = h_1^{Z\gamma} g^{\mu\nu} + \frac{1}{m_Z^2} \Big\{ h_2^{Z\gamma} p_1^{\nu} p_2^{\mu} + h_3^{Z\gamma} \epsilon^{\mu\nu\alpha\beta} p_{1\alpha} p_{2\beta} \Big\},$$

 $\gamma_{\nu}(p_2)$



The $H \to Z\gamma$ decay

The $H\to Z\gamma$ decay can be parametrized by the vertex function $\Gamma^{\mu\nu}_{Z\gamma H}$, $Z_\mu(p_1)$



The general form of the vertex function $\Gamma^{\mu\nu}_{Z\gamma H}$ is given as follows

$$\Gamma_{Z\gamma H}^{\mu\nu} = h_1^{Z\gamma} g^{\mu\nu} + \frac{1}{m_Z^2} \Big\{ h_2^{Z\gamma} p_1^{\nu} p_2^{\mu} + h_3^{Z\gamma} e^{\mu\nu\alpha\beta} p_{1\alpha} p_{2\beta} \Big\},$$

$$h_2^{Z\gamma} = \frac{2 m_Z^2}{m_Z^2 - m_H^2} h_1^{Z\gamma}. \longrightarrow \begin{bmatrix} \text{Complex} \\ \sim 10^{-1} \text{ in the SM} \end{bmatrix} \begin{bmatrix} \text{CP-violating and zero} \\ \text{in the SM} \end{bmatrix}$$

The
$$H \rightarrow Z\gamma$$
 decay

$$\begin{split} \Gamma(H \to Z\gamma) &= g^2 \frac{m_H^2 - m_Z^2}{32 \ \pi m_H^3 m_Z^4} \Big(4 \, |\, h_1^{Z\gamma}|^2 \, m_Z^4 + |\, h_3^{Z\gamma}|^2 \, \big(m_H^2 - m_Z^2 \big)^2 \Big) \\ &= \Gamma^{\text{SM}}(H \to Z\gamma) + \delta \Gamma(H \to Z\gamma), \end{split}$$

The general form of the vertex function $\Gamma^{\mu\nu}_{Z\gamma H}$ is given as follows

The $H \rightarrow Z\gamma$ decay

The signal strength $\mu^{Z\gamma}$ can be expressed as follows



The $H \rightarrow Z\gamma$ decay

Effective Lagrangian that induces FCNC of the Higgs and Z boson:

$$\mathscr{L} = \frac{g}{c_W} \bar{f}_i \left(g_V^{ij} - g_A^{ij} \gamma^5 \right) f_j Z^\mu + \frac{g}{2m_W} \bar{f}_i \left(g_S^{ij} + g_P^{ij} \gamma^5 \right) f_j H,$$

 g_V^{ij} , g_A^{ij} , g_S^{ij} and g_P^{ij} complex constants

A possible new physics contribution:



The $H \rightarrow Z\gamma$ decay $h_{3}^{Z\gamma} = \frac{g \ Q \ e \ m_{Z}^{2} \ N_{c}}{4\pi^{2} c_{W} m_{W}} \left\{ m_{j} \mathbf{C}_{0} \left(0, m_{H}^{2}, m_{Z}^{2}, m_{j}^{2}, m_{j}^{2}, m_{i}^{2} \right) \left[-\operatorname{Im} \left\{ g_{A}^{ij} \left(g_{S}^{ij} \right)^{*} \right\} + \operatorname{Im} \left\{ g_{V}^{ij} \left(g_{P}^{ij} \right)^{*} \right\} \right] \right\}$ $+m_i \mathbf{C}_0 \left(0, m_H^2, m_Z^2, m_i^2, m_i^2, m_j^2\right) \left[\operatorname{Im}\left\{g_A^{ij}\left(g_S^{ij}\right)^*\right\} + \operatorname{Im}\left\{g_V^{ij}\left(g_P^{ij}\right)^*\right\}\right]\right\}$ Calculated for the first time also A possible new physics contribution: f_j It does not contribute Η f_i to the $H \rightarrow \gamma \gamma$ decay f_i

The $H \rightarrow Z\gamma$ decay

$$h_{3}^{Z\gamma} = \frac{g \ @ \ e \ m_{Z}^{2} \ N_{c}}{4\pi^{2}c_{W}m_{W}} \left\{ m_{j}\mathbf{C}_{0} \left(0, m_{H}^{2}, m_{Z}^{2}, m_{j}^{2}, m_{j}^{2}, m_{i}^{2} \right) \left[-\ln\left\{ g_{A}^{ij} \left(g_{S}^{ij} \right)^{*} \right\} + \ln\left\{ g_{V}^{ij} \left(g_{P}^{ij} \right)^{*} \right\} \right] \right\} + m_{i}\mathbf{C}_{0} \left(0, m_{H}^{2}, m_{Z}^{2}, m_{i}^{2}, m_{i}^{2}, m_{i}^{2} \right) \left[\ln\left\{ g_{A}^{ij} \left(g_{S}^{ij} \right)^{*} \right\} + \ln\left\{ g_{V}^{ij} \left(g_{P}^{ij} \right)^{*} \right\} \right] \right\} \right]$$

A. I. Hernández-Juárez, R. Gaitán and R. Martinez, H - Zy decay and CP violation, Phys. Rev. D 111, 015001 (2025), arXiv: 2405.03094 [hep-ph].

- We estimate that for FCNC of the top quark $h_3^{Z\gamma} \approx 10^{-5}$, too small to explain the $\mu^{Z\gamma}$ excess.
- Contributions from new quarks are also possible and close to the bounds on $h_3^{Z\gamma}$.

The HZZ vertex

nature physics

ARTICLES https://doi.org/10.1038/s41567-022-01682-0

Check for updates

OPEN

Measurement of the Higgs boson width and evidence of its off-shell contributions to ZZ production

The CMS Collaboration^{*⊠}







Evidence of off-shell Higgs boson production from ZZ leptonic decay channels and constraints on its total width with the ATLAS detector

The ATLAS Collaboration

$$\Gamma_H = 3.2^{+2.4}_{-1.7}$$
 MeV

The
$$H \rightarrow ZZ^*$$
 well measured at the LHC

The Higgs boson must to be off-shell to produce two on-shell Z bosons

The HZZ vertex

Anomalous couplings for the ZZH vertex can be induced



Similar for the $HW^{\pm}W^{\mp}$ case

 h_i^V in terms of the anomalous couplings

$$\Gamma^{ZZH}_{\mu\nu} = h_1^V g_{\mu\nu} + \frac{h_2^V}{m_Z^2} p_{1\nu} p_{2\mu} + \frac{h_3^V}{m_Z^2} \epsilon_{\mu\nu\alpha\beta} p_1^{\alpha} p_2^{\beta},$$















- Small effects of h_3^H in the process $gg \to H^* \to ZZ \to 4l$

¿Polarized process
$$gg \to H^* \to Z_{\lambda}Z_{\lambda} \to 4l$$
? $\lambda = R, L$ and 0

• Left-Right asymmetry:

$$\mathscr{A}_{LR}^{H} = \frac{\Gamma_{H^* \to Z_L Z_L} - \Gamma_{H^* \to Z_R Z_R}}{\Gamma_{H^* \to Z_L Z_L} + \Gamma_{H^* \to Z_R Z_R}}$$



• Small effects of h_3^H in the process $gg \to H^* \to ZZ \to 4l$

¿Polarized process $gg \to H^* \to Z_{\lambda}Z_{\lambda} \to 4l$? $\lambda = R, L$ and 0

• Left-Right asymmetry:

$$\mathscr{A}_{LR}^{H} \sim \operatorname{Re}\left[h_{1}^{H}\right]\operatorname{Im}\left[h_{3}^{H}\right] - \operatorname{Re}\left[h_{3}^{H}\right]\operatorname{Im}\left[h_{1}^{H}\right]$$



• Small effects of h_3^H in the process $gg \to H^* \to ZZ \to 4l$

¿Polarized process $gg \to H^* \to Z_{\lambda}Z_{\lambda} \to 4l$? $\lambda = R, L$ and 0

• Left-Right asymmetry:

$$\mathscr{A}_{LR}^{H} \sim \operatorname{Re}\left[h_{1}^{H}\right]\operatorname{Im}\left[h_{3}^{H}\right] - \operatorname{Re}\left[h_{3}^{H}\right]\operatorname{Im}\left[h_{1}^{H}\right]$$
$$h_{1}^{H} \text{ complex in the SM}$$

CP-violation

imaginary parts

Similar asymmetries are posible in the $Z^* \to Z_{\!\lambda} H$ and the HWW vertex



The observation of the \mathscr{A}^{H}_{LR} asymmetry would imply a new source of CP violation

Summary

- Effects of new physics are still possible in Higgs couplings.
- CP-violation can explain the reported excess in the $H \rightarrow Z\gamma$ decay.
- New sources of CP-violation in the HZZ and HWW.



