SEMINARIO DE FÍSICA DE ALTAS ENERGÍAS

Instituto de Ciencias Nucleares UNAM



UNAM Seminar June 20, 2024

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An introduction to Heavy Ion Physics and CLASH

P. Christiansen (Lund University)



or "Why are heavy-ion physicists the Slytherins of particle physics!"





Outline

- A brief introduction to QCD and the Quark Gluon Plasma (QGP)
- The good: quarkonium
- The bad (?): the perfect liquid
- The ugly (?): QGP in small systems
- The CLASH
- The crazy (?): perfect QCD



Outline

- A brief introduction to QCD and the Quark Gluon Plasma (QGP)
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- The CLASH
- The crazy (?): perfect QCD
 - No time! Read it on arXiv:2301.13467 🙂



elt by:

Electricity, magnetism and chemistry

are all the results of electro-magnetic

force

quarks and charged leptons

Felt by:

Some forms of radio-activity are the

result of the weak force

quarks and leptons

although

gravitons

have so far not been

discovered

Feit by:

quarks

The explosive release of nuclear

energy is the result of the strong force

DACS. BETER CROWTHER

Felt by:

All the weight we experience is the

result of the gravitational force

all particles with mass

QED vs QCD

QED: superposition principle

(P. Christiansen, Lund)

An intro to Heavy Ion Physics and CLASH



QCD: color fields interact (form flux tube at long distance)





The strong interaction: Quantum Chromo Dynamics (QCD)

3 strong charges (red, green, blue) Particles in nature are color neutral Quarks are "confined"







The first challenge: the two limits of QCD





Solution (?): use lattice QCD

First question we will ask: What happens to the heavy-quark potential when we heat up the system

Lattice QCD results for the heavy quark potential (free energy): arXiv:0710.0498



The long-range force will be screened in the plasma. This should lead to deconfinement!

Deconfinement at high energy densities



Tc ~ 160 MeV (~2.000.000.000 K)

At high energy densities a new form of matter exists: Quark-Gluon Plasma, where quarks&gluons are deconfined I will only talk about the high temperature transition which is what we probe in heavy-ion collisions

QGP – the phase of the universe a few micro[®] seconds after The Big Bang







The second challenge: we cannot easily directly observe the QGP





The good: quarkonium

J/ψ: 1S, (ψ': 2S)



Υ: 1S, 2S, 3S





LHC has delivered



CMS has proven to be a marvelous detector for bottomonium

ALICE can complement by going to lower \boldsymbol{p}_{T} for charmonium



Suppression of heavy quarkonia can work as a thermometer

Note: 6.5<p_<30 GeV for J/ ψ and ψ (2s)



Suppression qualitatively depends on binding energy as predicted A lot of devils in the details that I skip.



J/ψ in the QGP at high energies



Requires high charm denisty



(P. Christiansen, Lund)

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J/ψ recombination



One needs to include recombination/regeneration to explain the data.



The bad (?): the perfect liquid



Lattice QCD calculation of the energy density





Several things pointed to that the QGP should behave more like a gas:

- The energy density found in lattice QCD calculations
- The screening of the long-range force
- The running of the coupling constant
- Ideas inspired by bag models



Radial flow



Aligned



- Flow in general plays a very important role in heavy-ion collisions.
- We believe that flow in the partonic phase is imprinted on the final state hadrons at freeze out.



Flow velocity $\beta_r \rightarrow mass$ dependent boost:

 $p_T \sim \gamma \beta_r m$ (for particle initially at rest)





Shear viscosity



The shear force is given as $F=\eta Av/d$ The shear vicosity-to-entropy density ratio, η/s , is a unitless quantity for characterizing fluids.



Elliptic flow and triangular flow is almost ideal!



- Huge flow at intermediate p_T:
 - 2 times more particles in plane than out Nearly ideal fluid
- Significant higher order flow caused by fluctuations also described by nearly ideal hydro + initial state



The QGP fluid compared to other fluids





Recent comparison



Because of the very low η /s (shear viscosity-to-entropy density) we think the QGP is a perfect liquid!





We went looking for a gas, but we found the • strongest interacting liquid known to mankind



Third challenge

- We have no precise idea of what a Quark-Gluon Plasma should look like
- That is IMO why we are the Slytherins of the LHC experiments: we do not have a solid theoretical foundation
 - New paradigm (?) for why we study the QGP
 - Circumvent traditional problem that while quarks and gluons are the fundamental degrees of freedom we observe hadrons to directly study the QCD dynamics of quarks and gluons

We are a mix of explorers... ...and gold diggers!



(P. Christiansen, Lund)

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Perfect liquid – the purest gold!

Perfect liquid expansion is almost reversible → Almost no entropy production!? \rightarrow We can "photograph" the initial overlap







The ugly (?): QGP in small systems



(P. Christiansen, Lund)

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Now we keep exploring and digging!







(P. Christiansen, Lund

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The effect of system size: Macroscopic effects in small systems?



Hot nuclear matter

Cold nuclear matter

QCD baseline





The rise of the double ridge





 Double ridge structure reminiscent of azimuthal flow in Pb-Pb collisions



Ridges in all systems



The perfect liquid is produced in all systems suggesting that small QCD systems produce "macroscopic" matter



Particle ratios in p-Pb and Pb-Pb show similar radial flow features



• Characteristic evolution of p/π and Λ/K_{S}^{0} with multiplicity is reminiscent of Pb-Pb where it is believed to be due to radial flow



The Mexican angle



- Visited UNAM 1 month in 2011 (EPLANET)
 - Ongoing collaboration since then
 - Common workshops: QCD challenges from pp to AA collisions, Taxco (2016), Puebla (2017), Lund (2019), Padova (2023), Muenster (2024)



(P. Christiansen, Lund)

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The "flow peak" in pp



Realized that Color Reconnection (CR) in PYTHIA gives rise to flow like boosts Antonio Ortiz Velasquez, Peter Christiansen, Eleazar Cuautle Flores, Ivonne Maldonado Cervantes, Guy Paić, PRL 111, 042001 (2013). For details, see T. Sjöstrand, arXiv:1310.8073.

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CR can be a microscopic model of flow

 \rightarrow Renewed interest in CR

Alternative to hydrodynamics



pp phenomenologists' favorite figure from ICHEP 2016



Need new physics mechanisms to describe increase!



Integrated particle ratios



DIPSY Color rope model: C. Bierlich, G. Gustafson, L. Lönnblad, A. Tarasov (Jefferson Lab), JHEP 1503 (2015) 148

The CLASH: Macroscopic (top-down) vs microscopic (bottom up) models

- Stat. thermal model
 - Canonical
 - Grand-canonical
- Hydrodynamics
 - Radial flow
 - Azimuthal anisotropic

- Tunneling of $q\overline{q}$ -pairs
 - Strings
 - Ropes
- String interactions
 - Color reconnection
 - Shoving





CLASH experimental angle: "Event Engineering"



Figure 1: Schematic of the structure of a $pp \rightarrow t\bar{t}$ event, as modelled by Pythia.

- Discovery \rightarrow Control/Isolation
- Question: can we control strangeness enhancement?
 E.g., switch on and off strangeness enhancement for a fixed multiplicity



Two ideas tested in CLASH

- No time to show: Relative Transverse Activity (R_{T})
 - <u>PhD thesis</u>: Omar Vazquez Rueda (UNAM → Lund → University of Houston)
 - ALICE, JHEP 06 (2023) 027 (π, K, p)
 - <u>PhD thesis</u>: Oliver Matonoha
 - To be published (K_s^0 , ϕ , Λ , Ξ)
- Transverse Spherocity (S₀)
 - Extension (N_{ch} \rightarrow Particle identification) of ideas and work proposed by Antonio Ortiz (see later)
 - PhD thesis: Adrian Nassirpour
 - ALICE, JHEP 05 (2024) 184





Transverse Spherocity S₀



Define the unweighted transverse spherocity: $S_{O}^{p_{T}=1} = \frac{\pi^{2}}{4} \min_{\hat{n}} \left(\frac{\sum_{tracks} |\hat{p}_{T} \times \hat{n}|}{N_{tracks}} \right)^{2}$

- Most other ALICE results were for the p_{T} -weighted S_{O}
 - We need this change because we study shortlived and neutral particles
 - Will call it S_0 in the following



The effect of S_o selection for different multiplicity estimators

Forward estimator Different region than where we measure S₀ Shown for top 10%. (typically used in ALICE to avoid autocorrelations)



- Physics we can address with S₀ depends on where we select the multiplicity
- The following results are all done with the mid-rapidity estimator
 - This ensures that multiplicity is almost constant so that we mainly select harder or softer events



Results top 1% multiplicity and top 1% S₀ (0.01% of events)

- Large differences between jetty and isotropic ratios ✓
- Events without S₀ selection are similar to isotropic
 - QGP-like effects dominates
 - Perfect liquid?
 - Hard physics is outlier
- Jet-like events
 - Radial-flow "peaks" are reduced
 - Strangeness is significantly reduced at high $p_{\rm T}$





Results top 1% multiplicity and top 10% S_0 (0.1% of events)



ALICE, JHEP 05 (2024) 184

- For top 10% we also have resonances (ϕ and K^{*0})
 - Require more statistics due to event mixing background
- Vs top 1%: effects are reduced but trends are the same



Strangeness enhancement vs S₀ (top 1% multiplicity)



- We can control the strangeness enhancement with $S_{\rm O}$ 🗸
 - The effect is bigger for Ξ (S=2) than for Λ (S=1)
- Pythia ropes can describe the enhancement qualitatively



Strangeness enhancement vs S_o (top 1% multiplicity)



- EPOS LHC captures the trend
 - The QGP core is reduced in jetty events
- HERWIG has opposite trend?! (next slide)



Why Herwig is wrong



- Herwig produces a baryon enhancement by allowing 3 mesons close in phase space to form a baryon-antibaryon pair
 - But this will be more likely to happen in pencil-like events!
 - What about quark coalescence models?



Strangeness enhancement vs S_o (top 10% multiplicity)



- ϕ ($\approx s\bar{s}$) and Ξ (ssd) follows different trends
- Data and models agree
 - Surprising for Pythia where ϕ is produced via 2 $s\bar{s}$ breakings
 - Suggests that the effect is mostly due to junctions
- How can we differentiate between EPOS and Pythia Ropes?

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Answer: look at the how the strange quarks are balanced

QGP:

 Ξ (Xi) baryon

 \overline{d}

S

S

We naively expect that in a QGP the quarks will be deconfined and so eventually the quark pairs will drift apart in phase space.

Lund string: Most quarks and antiquarks are produced together during hadronization.

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The easiest case: Ξ balanced by antiproton

QGP:

S

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We expect that the balancing occurs on a statistical basis so this can happen.

The easiest case: Ξ balanced by antiproton

Normal Lund string and ropes: Ξ almost never balanced by antiproton but instead typically by antistrange baryons and even anti- Ξ !

Idea from CLASH workshop write up: J. Adolfsson et al, Eur. Phys. J. A 56 (2020) 11, 288, "QCD challenges from pp to A–A collisions"





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The easiest case: Ξ balanced by antiproton

Junction:

 Ξ balanced more by kaons and less by antistrange baryons. Broader correlations in rapidity.

Idea from CLASH workshop write up: J. Adolfsson et al, Eur. Phys. J. A 56 (2020) 11, 288, "QCD challenges from pp to A–A collisions"

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Microscopic balance of Ξ by antiprotons: MB results



- EPOS (QGP) model: no structure due to extreme assumption of grand-canonical ensemble
- Pythia8 Monash: fails since this almost never happens
- Pythia8 Junctions: describes well the data

Microscopic balance of Ξ by antiprotons: low mult results



- Pythia8 Junctions: fails to describe the data since in the low multiplicity limit it must agree with Monash (no CR)
- But why does nature prefer such a complicated process where strangeness is balanced by two mesons?



Part of the work of Jonatan Adolfsson's PhD Thesis



 He studied many more combinations, see arXiv:2308.16706

Future outlook

- The study of the microscopic balance is something we want to keep pursuing
 - Xi (ssd) \rightarrow Omega (sss)
 - 2 particles \rightarrow
 - 3 particles
 - $-ssbar \rightarrow ccbar$

Thank You!



