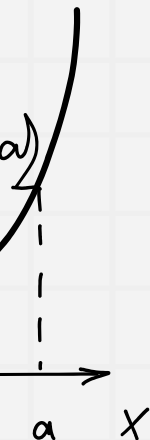


$$\int_0^1 dy \int_0^1 f(x) dx + \int_0^1 dy \int_{1/\sqrt{2}}^1 f(2x) dx =$$

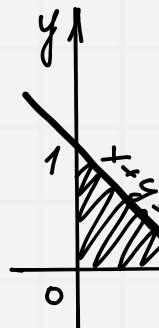
$$2\sqrt{y^2 - x^2}$$



SUSY TARASCA @ THE BEACH

Alfredo Raya

IFM-UMSNH



$$\begin{aligned} x &= 2y \\ z &= 1 + \sqrt{x} \\ z &= 4 + \sqrt{y} \end{aligned}$$

$$\int_0^1 dx \int_0^{1-x} x^2 z^{10(x+3y)} du =$$

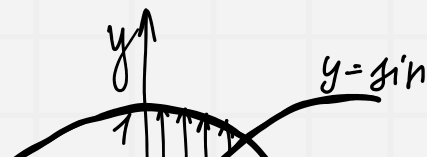


Table of contents

01

Fundamentals

From SUSY to SUSY-QM

02

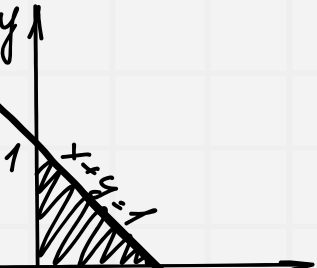
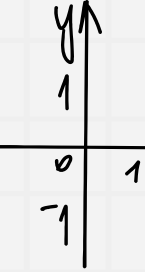
Dirac Theory

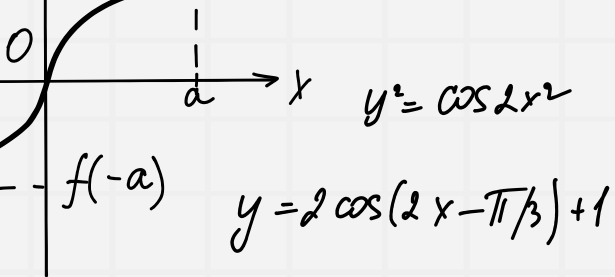
Examples with planar
fermions

03

Graphene

Ritus Propagator

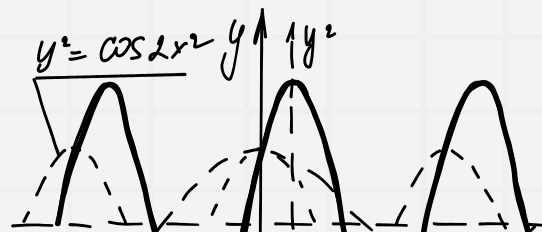
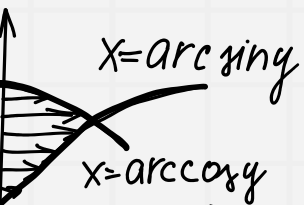




$$\int_0^1 dy \int_0^1 f dx + \int_{1/\sqrt{2}}^1 dy \int_0^1$$

SUSY!

At the interface of High Energy and Condensed Matter Physics



$$\int_0^x \int_0^y \int_0^{10(x+3y)} x^2 dz =$$

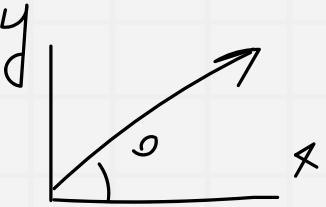
$$V: z=10(x+3y), x+y+z=1, x=0, y=0, z=0$$

$$\frac{3, x=5}{2+4y^2} \\ \frac{1}{2+4y^2}$$

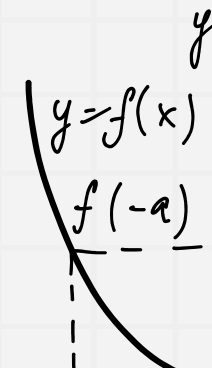
01

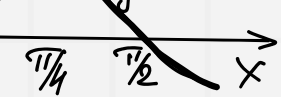
Fundamentals

SUSY and SUSY-QM

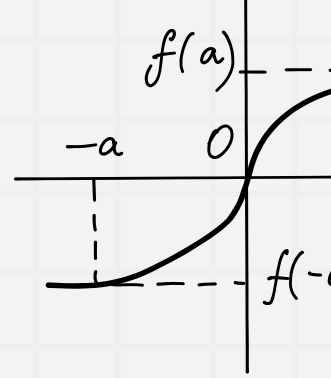


$$= \int_0^1$$





$$2\sqrt{y^2 - x^2}$$



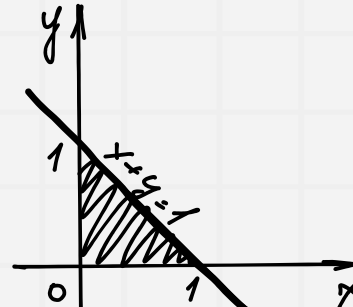
Standard Model of Particle Physics

Most successful theory. Many free parameters.

	MATTER FERMIONS			FORCE BOSONS	
QUARKS	u	c	t	g	Z
	d	s	b	γ	W
LEPTONS	e	μ	τ		H
	ν_e	ν_μ	ν_τ		

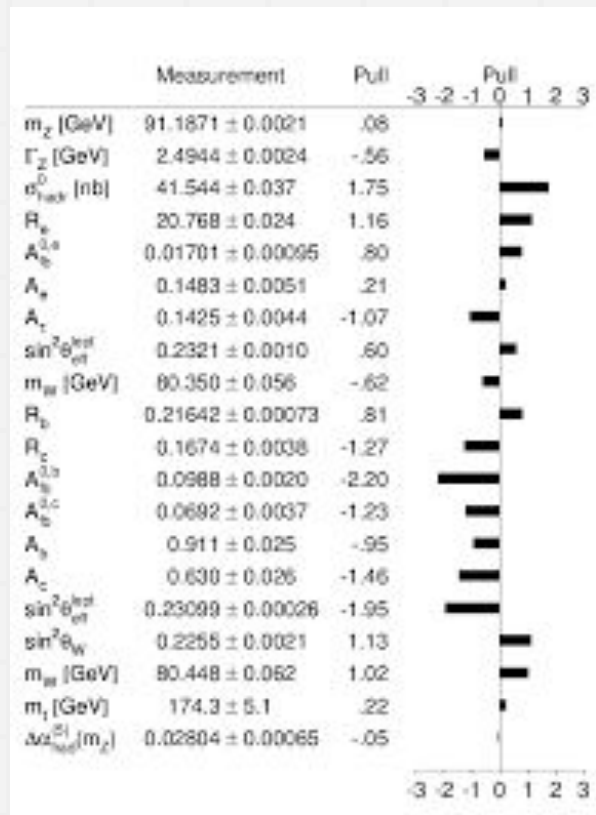
$$3, x=5$$

$$\frac{2+4y^2}{x^2+4y^2}$$



fdy

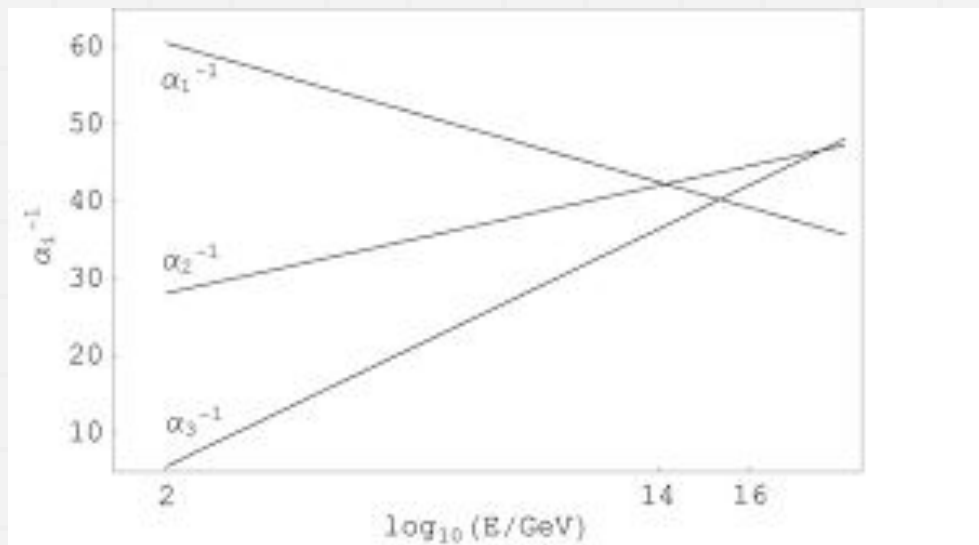
Great success with many free parameters



$$5 = 2$$

fdy

Grand Unification



Unification of the fundamental interactions

- Almost the same strength
- Pairwise unification
- Can there be exact unification?

SUSY comes to rescue



$$5 = 2$$

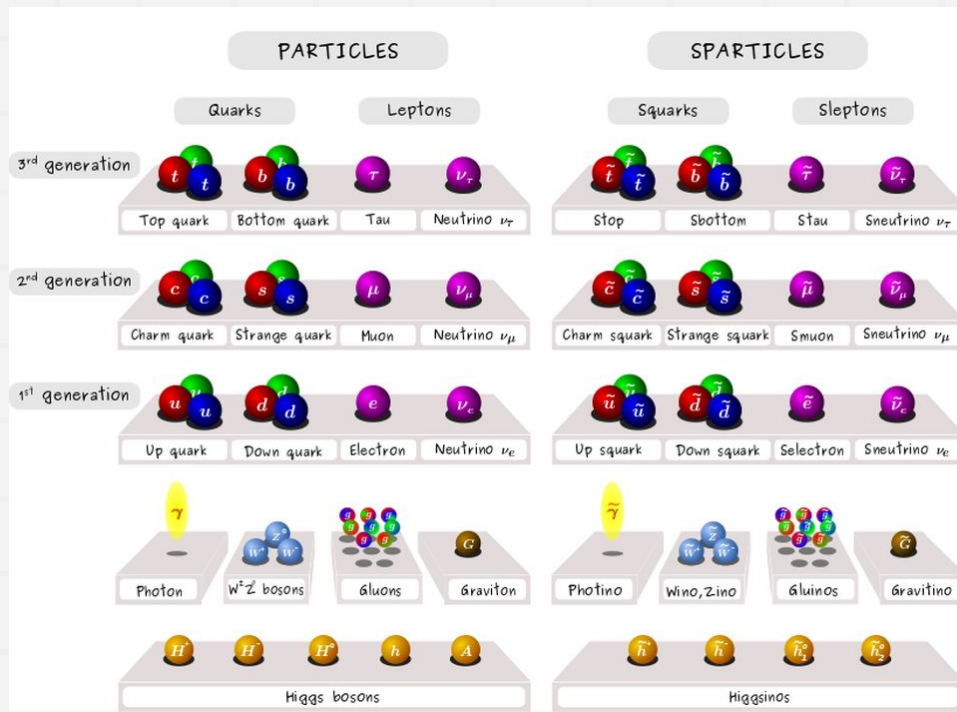
fdy

Minimal Supersymmetric Standard Model

$$Q_\alpha^A |B\rangle = |F\rangle$$

$$Q_\alpha^A |F\rangle = |B\rangle$$

NoetherSym



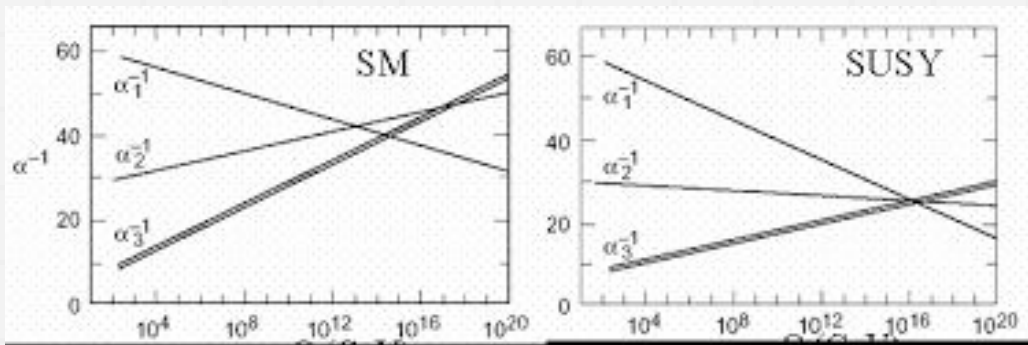
$$5 = 2$$

fdy

Minimal Supersymmetric Standard Model

$$Q_\alpha^A |B\rangle = |F\rangle$$
$$Q_\alpha^A |F\rangle = |B\rangle$$

NoetherSym



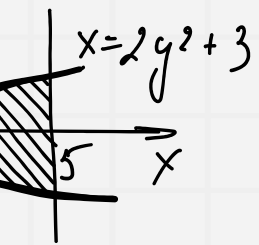
$$5 = 2$$

$$x^2 dz =$$

$$= \int_0^1 dx$$

SUSY is broken

$$dy dz =$$



No sparticle has been observed

- How is SUSY broken?
- SUSY-QM
- Beyond the MSSM

From high-energy physics to other realizations



$$V: z = 10(x + 3y),$$

$$x = 0, y = 0, z =$$

$$x^2 dz =$$

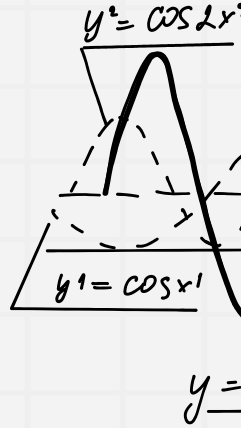
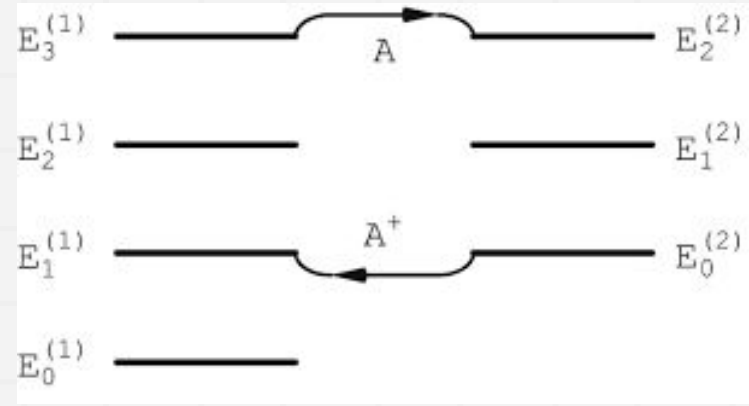
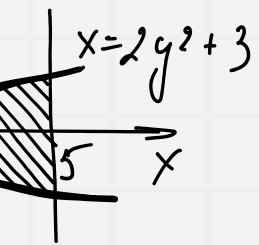
$$= \int_0^1 dx$$

SUSY-QM

$$dy dz =$$

(0+1)-QFT

- 2 Schrödinger Hamiltonians
- Quasi-isospectrality



$$V: z = 10(x + 3y), x = 0, y = 0, z =$$

$$\int_0^1 \int_0^{1-x} \int_0^{10(x+3y)} x^2 dz dy =$$

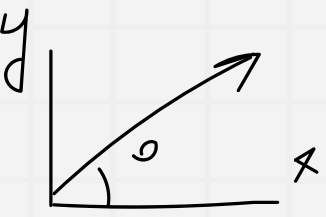
$$V: z=10(x+3y), x+y+z=1, x=0, y=0, z=0$$

$$\frac{3, x=5}{2+4y^2} \\ \frac{2+4y^2}{2+4y^2}$$

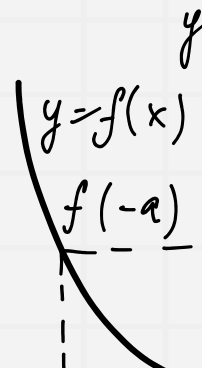
02

Dirac Theory

Planar Fermions



$$= \int_0^1$$

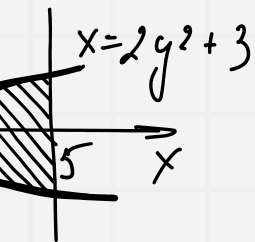


$$\int x^2 dz =$$

$$= \int_0^1 dx$$

Relativistic Wave Equation

$$dydz =$$



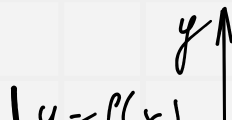
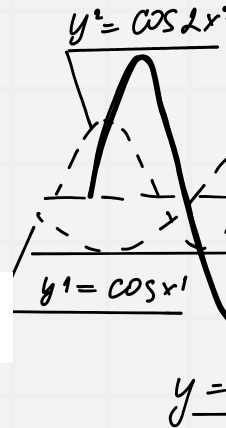
Trying to unify QM and SR, the Klein Gordon equation emerged as the first attempt

- Negative Energy States
- Probabilistic Interpretation of the wave function

Need a better idea!

$$(\partial_\mu \partial^\mu + m^2)\phi = 0$$

$$\rho \sim \phi^* \partial_t \phi - \phi \partial_t \phi^*$$



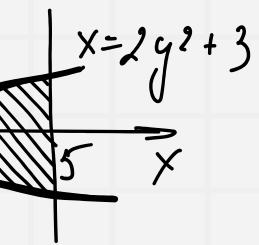
$$V: z = 10(x + 3y), \\ x = 0, y = 0, z =$$

$$x^2 dz =$$

$$= \int_0^1 dx$$

Relativistic Wave Equation

$$dydz =$$



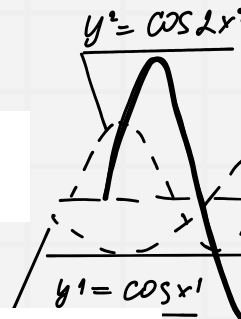
Dirac:

- Need a Hamiltonian linear in space derivatives
- Its squared is equivalent to the KG equation
- Has a neat probabilistic interpretation
- No negative energy states

$$H_D = \vec{\alpha} \cdot \vec{p} + \beta m$$

α_i, β matrices

$$\psi = \begin{pmatrix} \psi_1 \\ \vdots \\ \psi_n \end{pmatrix}, \quad |\psi|^2 \geq 0$$



Oops!

$$E = m_0, m_0, -m_0, -m_0$$

$$V: z = 10(x + 3y), x = 0, y = 0, z =$$



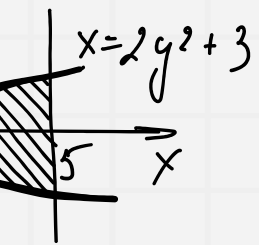
$$|u = p(x)|$$

$$\int x^2 dz =$$

$$= \int_0^1 dx$$

Relativistic Wave Equation

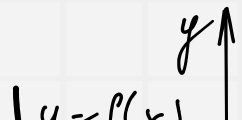
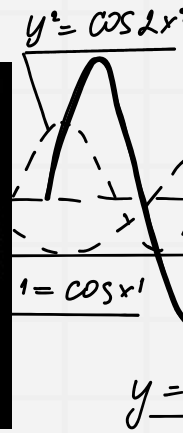
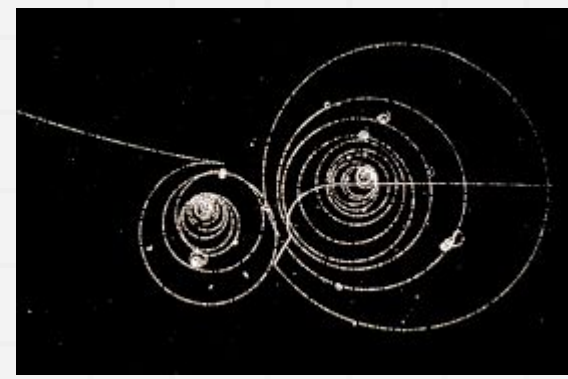
$$dy dz =$$



Dirac:

- QM+SR = antimatter
- More intelligent equation than its creator

$$\psi_{\uparrow\downarrow}^{\gt}, \quad \psi_{\uparrow\downarrow}^{\lt}$$



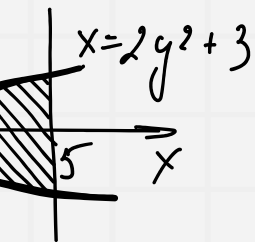
$$V: z = 10(x + 3y), \\ x = 0, y = 0, z =$$

$$x^2 dz =$$

$$= \int_0^1 dx$$

Relativistic Wave Equation in 2D

$$dy dz =$$

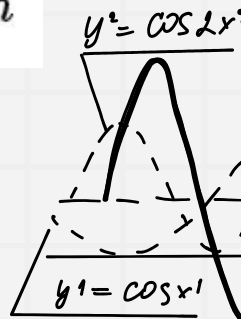


Dirac:

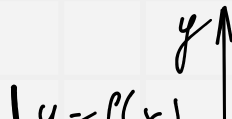
- 2x2 Dirac Matrices
- 2 Irreducible inequivalent representations
- 1 Spin orientation in each
- Mass breaks parity
- No chiral symmetry

$$\gamma_5 = \pm iI$$

$$H_D = \sigma_x p_x + \sigma_y p_y + \sigma_3 m$$



$$y =$$



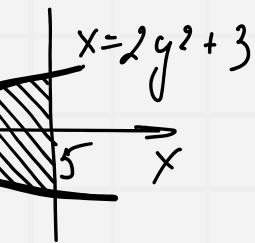
$$V: z = 10(x + 3y), \\ x = 0, y = 0, z =$$

$$x^2 dz =$$

$$= \int_0^1 dx$$

Relativistic Wave Equation in 2D

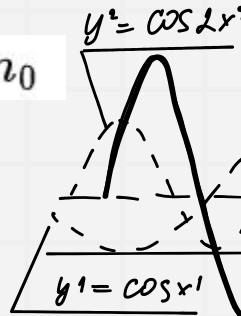
$$dy dz =$$



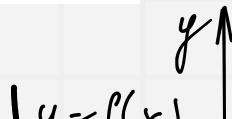
Dirac:

$$H_D = \alpha_x p_x + \alpha_y p_y + \beta m + \alpha_x \alpha_y m_0$$

- 4x4 Dirac Matrices
- 1 reducible representation
- 2 Spin orientations, 2 positive to negative energy solution
- Dirac mass preserves parity
- 2 chiral transformations, $\gamma_3 \gamma_5$
- New mass term $m_0 \frac{1}{2} \gamma_3 \gamma_5$



$$y =$$



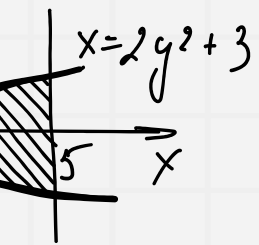
$$V: z = 10(x + 3y), \\ x = 0, y = 0, z =$$

$$\int_{x^2+y^2}^1 x^2 dz =$$

$$= \int_0^1 dx$$

Relativistic Wave Equation in 2D

$$dydz =$$



Graphinos

- 4x4 Dirac Matrices
- Uniform magnetic field
- Dirac and Haldane masses

SHernández Ortiz, G Murguía, AR
JPCM24(2012)015304

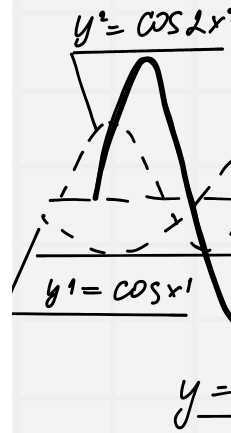
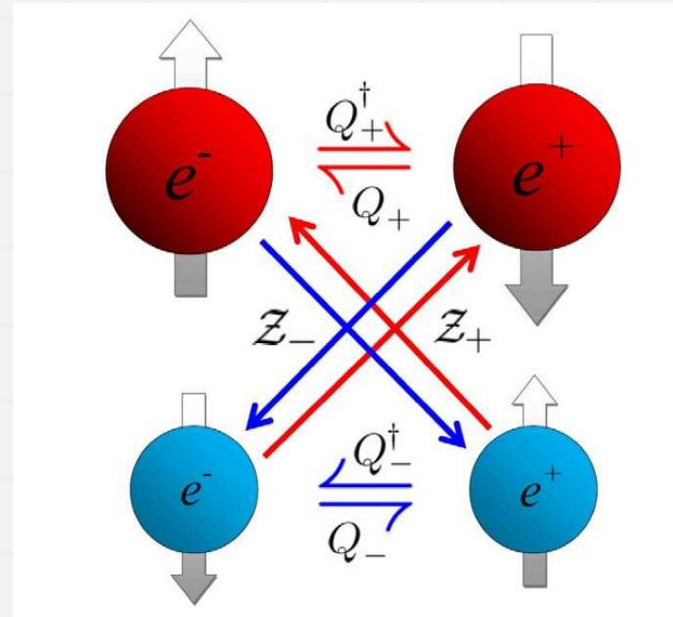


Figure 2. Action of the operators Z_{\pm} in the reducible representation of graphinos. These operators reveal an exact SUSY-QM structure in the massless limit and hardly break such a structure if a Haldane mass term is included.



$$|u = p(x)|$$

$$\bar{x} = 0, y = 0, z = 10(x+3y)$$

$$\int_{1-x}^{10(x+3y)} \int x^2 dz =$$

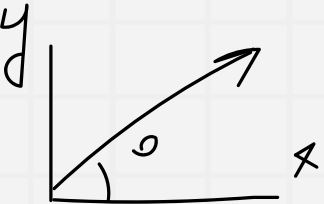
$$V: z=10(x+3y), x+y=1, x=0, y=0, z=0$$

$$\frac{3, x=5}{2+4y^2} \\ \frac{2+4y^2}{2+4y^2}$$

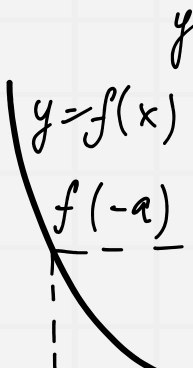
03

Graphene

The Age of Quantum Materials



$$= \int_0^1$$



$$\int x^2 dy dz =$$

$$\int dy dz =$$

$$x = 2y^2 + 3$$



Graphene

Graphene

- 1 atom thick layer
- Genuine 2D crystal
- Remarkable properties
- Low energy, massless Dirac eq.

$$H_g = v_F \vec{\sigma} \cdot \vec{p}$$

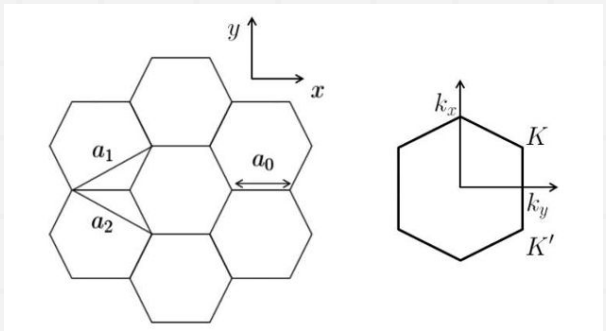
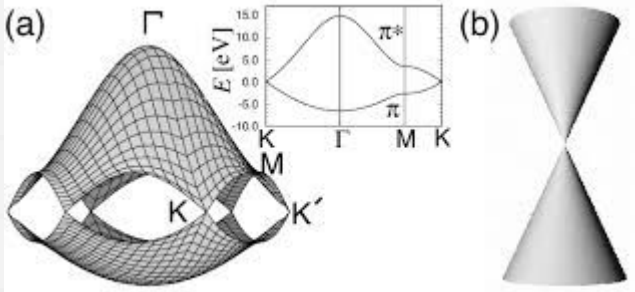


Figure 1. Crystallographic structure of graphene (left) and the Dirac points in the reciprocal k -space (right). The lattice parameters are $a_0 = 1.42 \text{ \AA}$, $\mathbf{a}_1 = a_0\sqrt{3}(1/2, \sqrt{3}/2)$ and $\mathbf{a}_2 = a_0\sqrt{3}(-1/2, \sqrt{3}/2)$.



$$y^2 = \cos 2x$$

$$y^2 = \cos^2 x$$

$$y =$$

$$V: z = 10(x + 3y), x = 0, y = 0, z =$$

$$u = p(x)$$

$$\int_0^{10(x+3y)} dy =$$

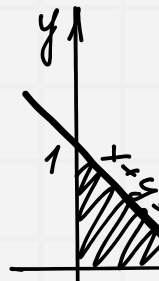
$$\sqrt{2-x^2}$$

Graphene



Propagator

Ritus Method

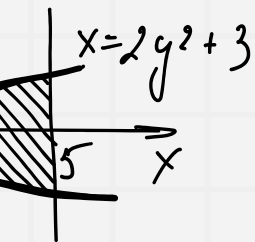


$$x^2 dz =$$

$$= \int_0^1 dx$$

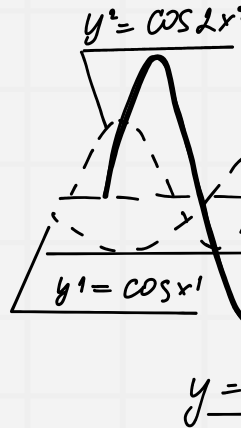
Propagator

$$dy dz =$$



Propagator

- Non-diagonal in momentum space
- Spectral representation
- Schwinger Method
- Ritus method
- Very few field configurations known



$$|u = p(x)|$$

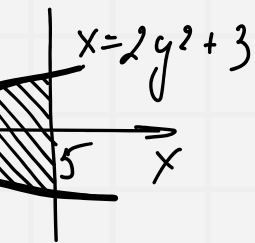
$$V: z = 10(x + 3y), \\ x = 0, y = 0, z =$$

$$x^2 dz =$$

$$= \int_0^1 dx$$

Ritus Method

$$dy dz =$$



Ritus Method

- Eigenfunctions of $(\gamma \Pi)^2$

$$(\gamma \cdot \Pi)^2 \mathbb{E}_p = p^2 \mathbb{E}_p$$

- Ritus functions

$$\int d^d z \bar{\mathbb{E}}_{p'}(z) \mathbb{E}_p(z) = \mathbb{I} \delta(p - p'),$$

$$\int d^d p \mathbb{E}_p(z') \bar{\mathbb{E}}_p(z) = \mathbb{I} \delta(z - z'),$$

- Propagator

$$S(z, z') = \int d^d p \mathbb{E}_p(z) \frac{1}{\gamma \cdot \bar{p} - m} \bar{\mathbb{E}}_{p'}(z')$$



$$|u = p(x)|$$



$$y =$$

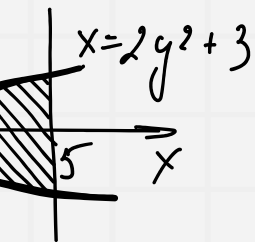
$$V: z = 10(x + 3y), \\ x = 0, y = 0, z =$$

$$x^2 dz =$$

$$= \int_0^1 dx$$

Ritus Method

$$dy dz =$$



Ritus Method

- Perpendicular magnetic field (Landau-like gauge)

$$[-\partial_x^2 + (p_2 + eW_0(x))^2 - e\sigma W_0'(x)] F_{k,p_2,\sigma} = k F_{k,p_2,\sigma}$$

- Uniform Magnetic Field and exponentially decaying magnetic field
G. Murguía, A. Sánchez, E. Reyes, AR, Am. J. Phys. 78, 700-707 (2010)



$$y =$$



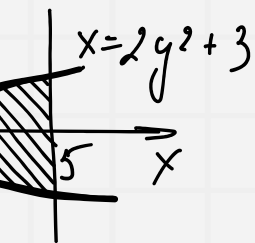
$$V: z = 10(x + 3y), \\ x = 0, y = 0, z =$$

$$x^2 dz =$$

$$= \int_0^1 dx$$

Ritus Method

$$dy dz =$$



Ritus Method $(\gamma \cdot \Pi)^2 \mathbb{E}_p(z) = p^2 \mathbb{E}_p(z)$

$$\mathbb{E}_p^{(\mathcal{A})}(z) = \begin{pmatrix} E_{p,+1}(z) & 0 \\ 0 & E_{p,-1}(z) \end{pmatrix}$$

$$\mathbb{E}_p^{(\mathcal{A})}(z) = E_{p,+1}(z)\Delta(+)+E_{p,-1}(z)\Delta(-)$$

$$\Delta(\pm) = \frac{\mathbb{1} \pm i\gamma^1 \gamma^2}{2}$$

$$|u = p(x)|$$



$$y =$$

$$V: z = 10(x+3y), \\ x=0, y=0, z=$$



$$x^2 dz =$$

$$= \int_0^1 dx$$

Ritus Method

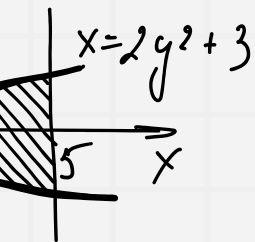
$$dy dz =$$

Ritus Method $(\gamma \cdot \Pi)^2 \mathbb{E}_p(z) = p^2 \mathbb{E}_p(z)$

$$E_{p,\sigma}(z) = N_\sigma e^{-i(p_0 t - p_2 y)} F_{k,p_2}^\sigma(x)$$

$$[\partial_x^2 - (-p_2 + eW(x))^2 + e\sigma W'(x) + k] F_{k,p_2}^\sigma(x) = 0,$$

$$V_\pm(x) = (-p_2 + eW(x))^2 \pm eW'(x)$$



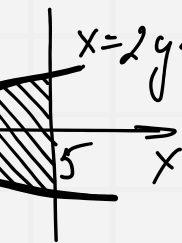
$$V: z = 10(x + 3y), \\ x = 0, y = 0, z =$$

$$x^2 dz =$$

$$= \int_0^1 dx$$

Ritus Method

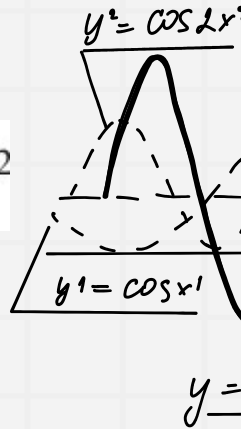
$$dy dz =$$



Ritus Method $(\gamma \cdot \Pi)^2 \mathbb{E}_p(z) = p^2 \mathbb{E}_p(z)$

$$B(x) = Be^{-\hat{\alpha}x} \Rightarrow W(x) = -\frac{B}{\hat{\alpha}} (e^{-\hat{\alpha}x} - 1)$$

$$k_n = \hat{p}_2^2 - (\hat{p}_2 - n\hat{\alpha})^2$$



$$E_{p,+1}(z) = \frac{\hat{\alpha}}{2\pi} \left(\frac{2n!(s-n)}{\Gamma(2s-n+1)} \right)^{1/2} e^{-ip_0t+ip_2y} e^{-\varrho/2} \varrho^{(s-n)} \times L_n^{2(s-n)}(\varrho),$$

$$E_{p,-1}(z) = \frac{\hat{\alpha}}{2\pi} \left(\frac{2(n-1)!(s-n)}{\Gamma(2s-n)} \right)^{1/2} e^{-ip_0t+ip_2y} \times e^{-\varrho/2} \varrho^{(s-n)} L_{n-1}^{2(s-n)}(\varrho).$$

$$\therefore z = 10(x+3y), \quad x=0, y=0, z=$$

(31)



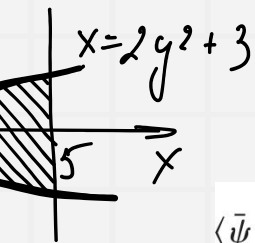
$$x^2 dz =$$

$$= \int_0^1 dx$$

Ritus Method

$$dy dz =$$

Ritus Method $(\gamma \cdot \Pi)^2 \mathbb{E}_p(z) = p^2 \mathbb{E}_p(z)$



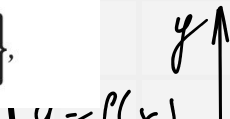
$$\langle \bar{\psi} \psi \rangle = \text{Tr}\{iS(z, z)\},$$

$$\langle \bar{\psi} \psi \rangle_{\mathcal{A}} = i \int d^3 p \frac{m}{\bar{p}^2 - m^2} [|E_{p,+1}(z)|^2 + |E_{p,-1}(z)|^2]$$



$$\begin{aligned} \langle \bar{\psi} \psi \rangle_{\mathcal{A}} = & \frac{m\hat{\alpha}^2}{2\pi} \left\{ \sum_{s=0}^{\infty} \frac{1}{|m|} \left(\frac{s}{\Gamma(2s+1)} \right) e^{-\varrho} \varrho^{2s} [L_0^{2s}(\varrho)]^2 \right. \\ & + \sum_{n=1}^{\infty} \sum_{s=n+1}^{\infty} \frac{e^{-\varrho} \varrho^{2(s-n)}}{\sqrt{\hat{\alpha}^2(2sn - n^2) + m^2}} \\ & \times \left[\left(\frac{n!(s-n)}{\Gamma(2s - (n-1))} \right) [L_n^{2(s-n)}(\varrho)]^2 \right. \\ & \left. \left. + \left(\frac{(n-1)!(s-n)}{\Gamma(2s-n)} \right) [L_{(n-1)}^{2(s-n)}(\varrho)]^2 \right] \right\}, \end{aligned}$$

$$\langle \bar{\psi} \psi \rangle_{\mathcal{A}}^{n=0, \sigma=1} = \frac{e}{4\pi} B e^{-\hat{\alpha}x} \text{sgn}(m),$$



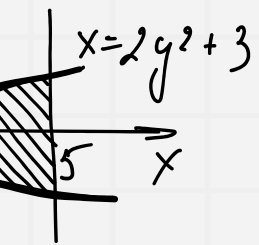
$$\begin{aligned} V: z &= 10(x+3y), \\ x &= 0, y = 0, z = \end{aligned}$$

$$x^2 dz =$$

$$= \int_0^1 dx$$

Ritus Method

$$dy dz =$$



Ritus Method

- Perpendicular magnetic field (Landau-like gauge)

$$[-\partial_x^2 + (p_2 + eW_0(x))^2 - e\sigma W_0'(x)] F_{k,p_2,\sigma} = k F_{k,p_2,\sigma}$$

- Second Order SUSY for these fields

Y. Concha, E. Díaz-Bautista, AR, Phys. Scr. 97, 095203 (2022)



$$y =$$



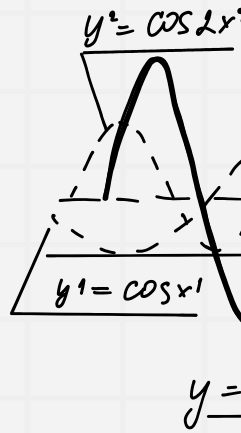
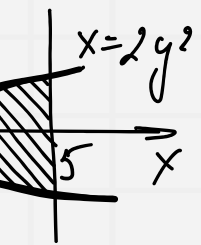
$$V: z = 10(x + 3y), \\ x = 0, y = 0, z =$$

$$x^2 dz =$$

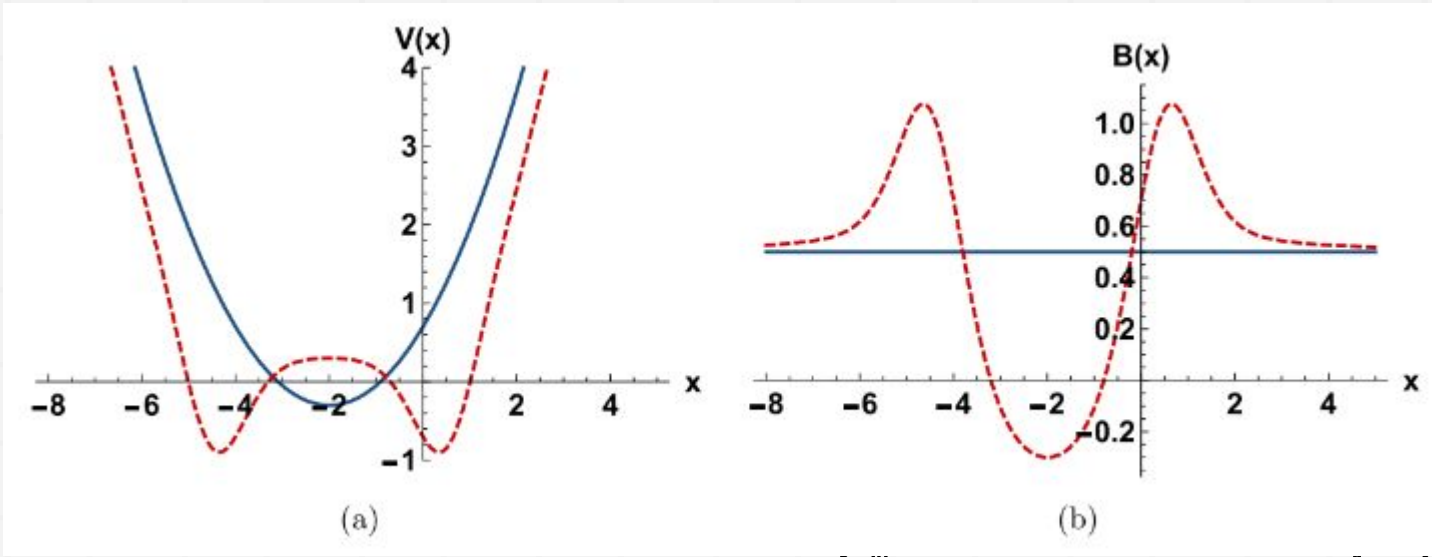
$$= \int_0^1 dx$$

Ritus Method

$$dy dz =$$



Ritus Method



Y. Concha, E. Díaz-Bautista, AR, Phys. Scr. 97, 095203 (2022)

$$z = 10(x + 3y),$$

$$x = 0, y = 0, z =$$

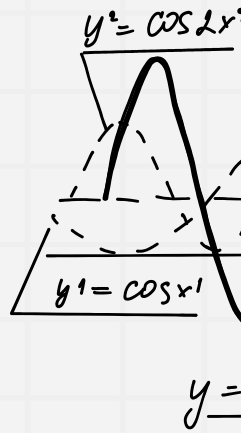
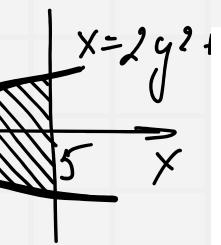


$$x^2 dz =$$

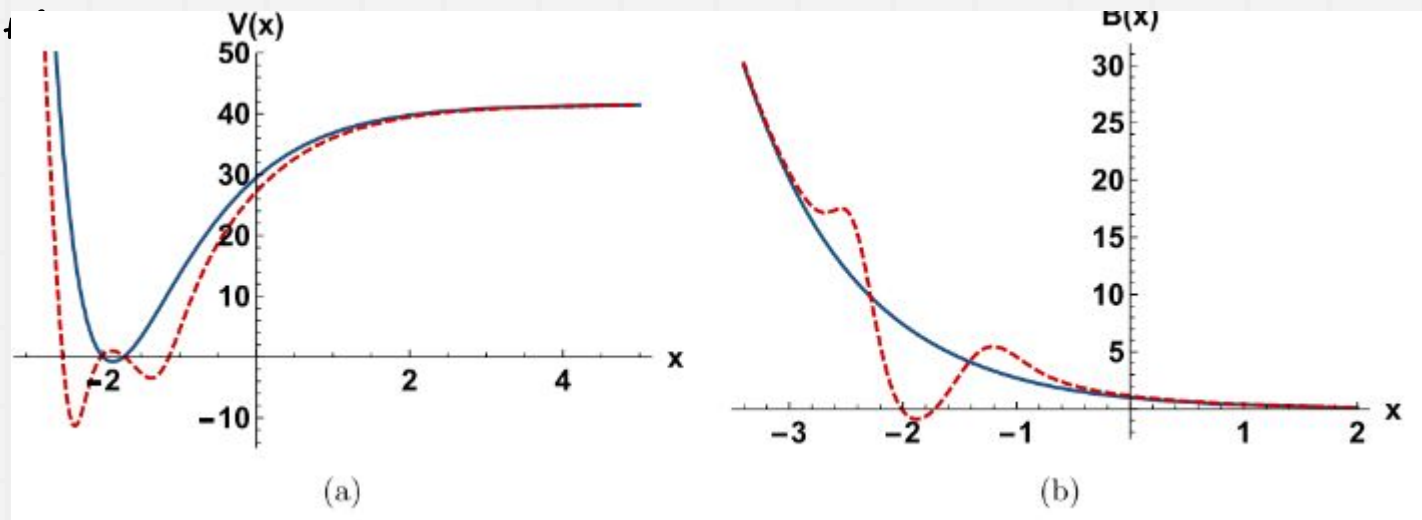
$$= \int_0^1 dx$$

Ritus Method

$$dy dz =$$



Ritus Method



Y. Concha, E. Díaz-Bautista, AR, Phys. Scr. 97, 095203 (2022)

$$z = 10(x + 3y),$$

$$x = 0, y = 0, z =$$

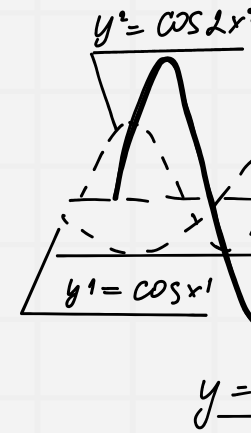
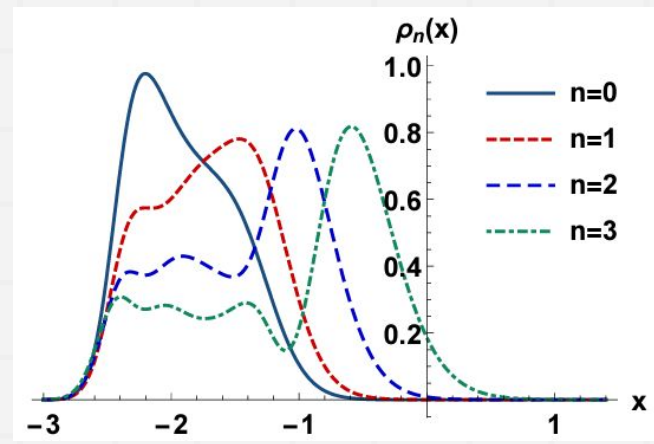
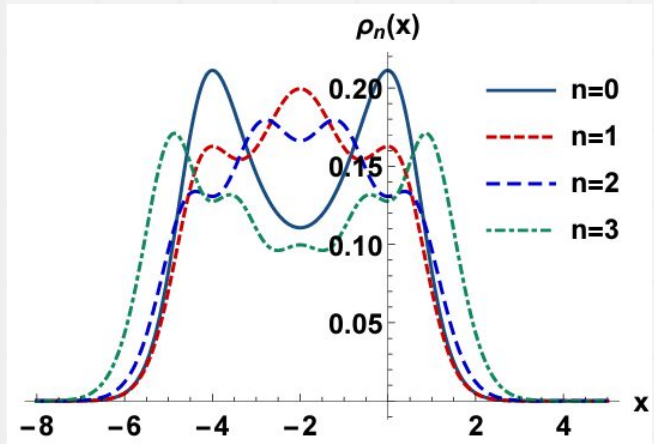
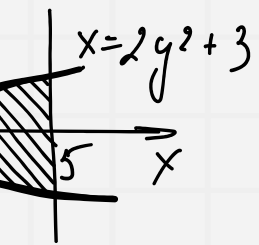


$$x^2 dz =$$

$$= \int_0^1 dx$$

Ritus Method

$$dy dz =$$



Y. Concha, E. Díaz-Bautista, AR, Phys. Scr. 97, 095203 (2022)



$$V: z = 10(x + 3y),$$

$$x = 0, y = 0, z =$$

$$\int_0^{1-x} \int_0^{10(x+3y)} \int_0^x x^2 dz =$$

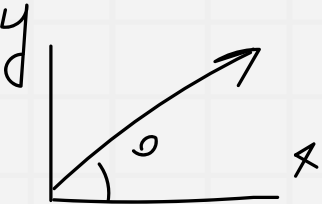
$$V: z=10(x+3y), x+y+z=1, x=0, y=0, z=0$$

$$\int_0^3 \frac{x^2+4y^2}{x^2+4y^2} dx = 5$$

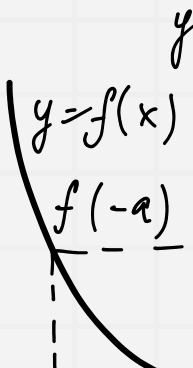
05

Final Remarks

Good to be in vacations while working!



$$= \int_0^1$$

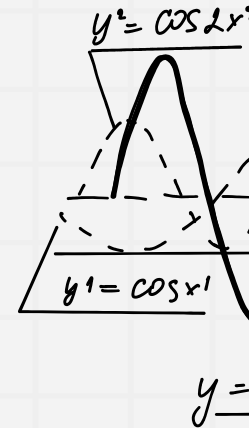
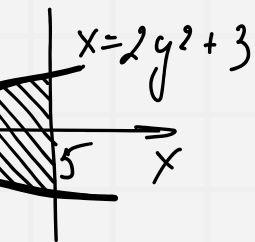


$$\int (x+y) x^2 dz =$$

$$= \int_0^1 dx$$

n-th Latin American Workshop on Electromagnetics Effects in QCD

$$dy dz =$$



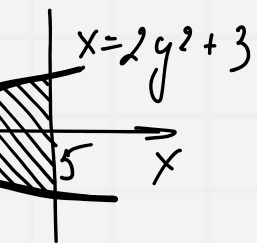
$$u = p(x)$$

$$V: z = 10(x + 3y), x = 0, y = 0, z =$$

$$x^2 dz =$$

$$= \int_0^1 dx$$

$$dy dz =$$



Thanks!



$$y =$$



$$|u = p(x)|$$

$$V: z = 10(x + 3y), \\ x = 0, y = 0, z =$$