

Cosmology as a neutrino lab

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6 de junio 2024 RADPyC

Outline

- State-of-the-art of standard neutrinos.
- Neutrino NSI in cosmology
- Neutrino self-interactions with light and not so light mediators

Neutrinos

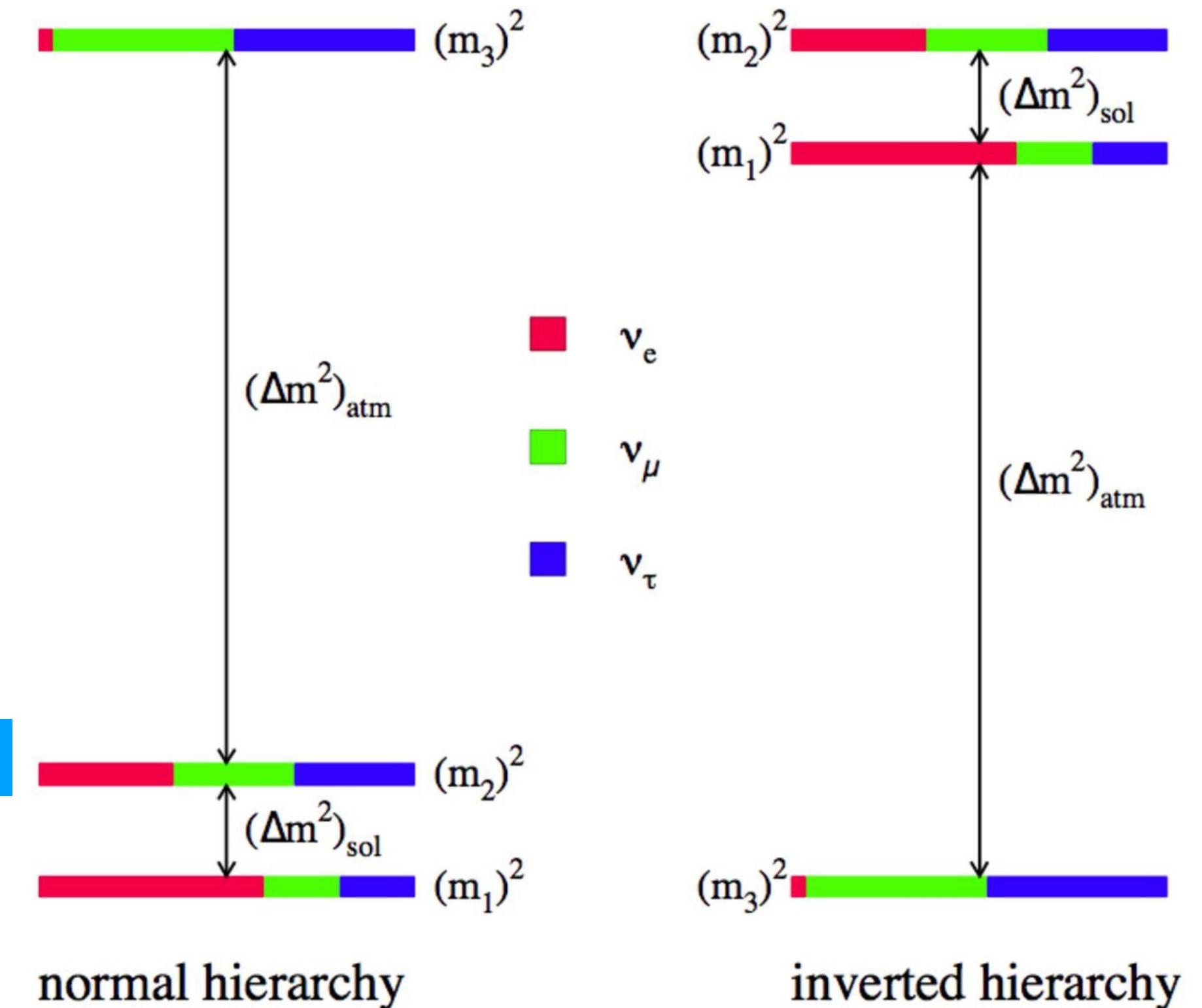
$$\Delta m_{\text{sol}}^2 = \Delta m_{21}^2 = m_2^2 - m_1^2 = 7.42^{+0.21}_{-0.20} \times 10^{-5} \text{ eV}^2$$

$$\text{NH} \longrightarrow \Delta m_{\text{ATM}}^2 = \Delta m_{31}^2 = m_3^2 - m_1^2 = 2.517^{+0.026}_{-0.028} \times 10^{-3} \text{ eV}^2$$

$$\text{IH} \longrightarrow \Delta m_{\text{ATM}}^2 = \Delta m_{23}^2 = m_2^2 - m_3^2 = 2.498^{+0.028}_{-0.028} \times 10^{-3} \text{ eV}^2$$



2020 Esteban+ JHEP



The standard model does not explain neutrino mass

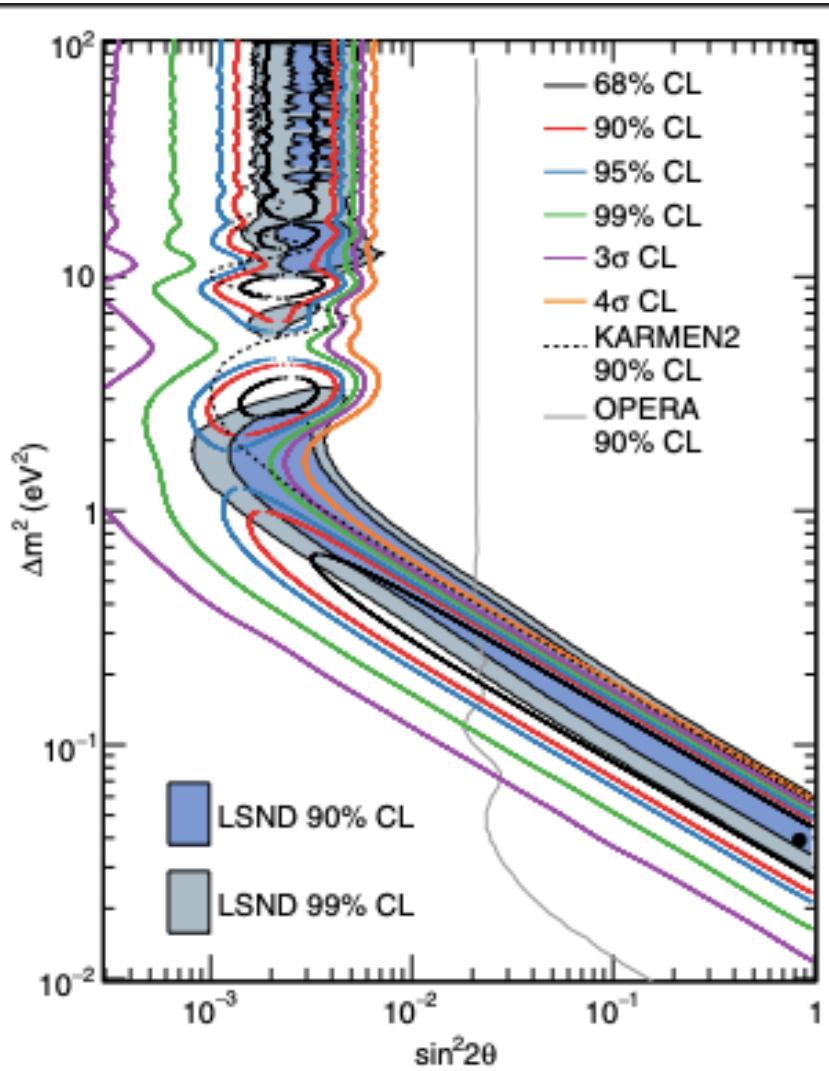
We know that neutrinos have mass. We don't have a direct measurement of them

★ The most popular neutrino mass models require a neutrino NSI

Anomalies

LSND \rightarrow MiniBoone $\nu_\mu \rightarrow \nu_e$

Two decade anomaly with a S/N of 6 sigmas:



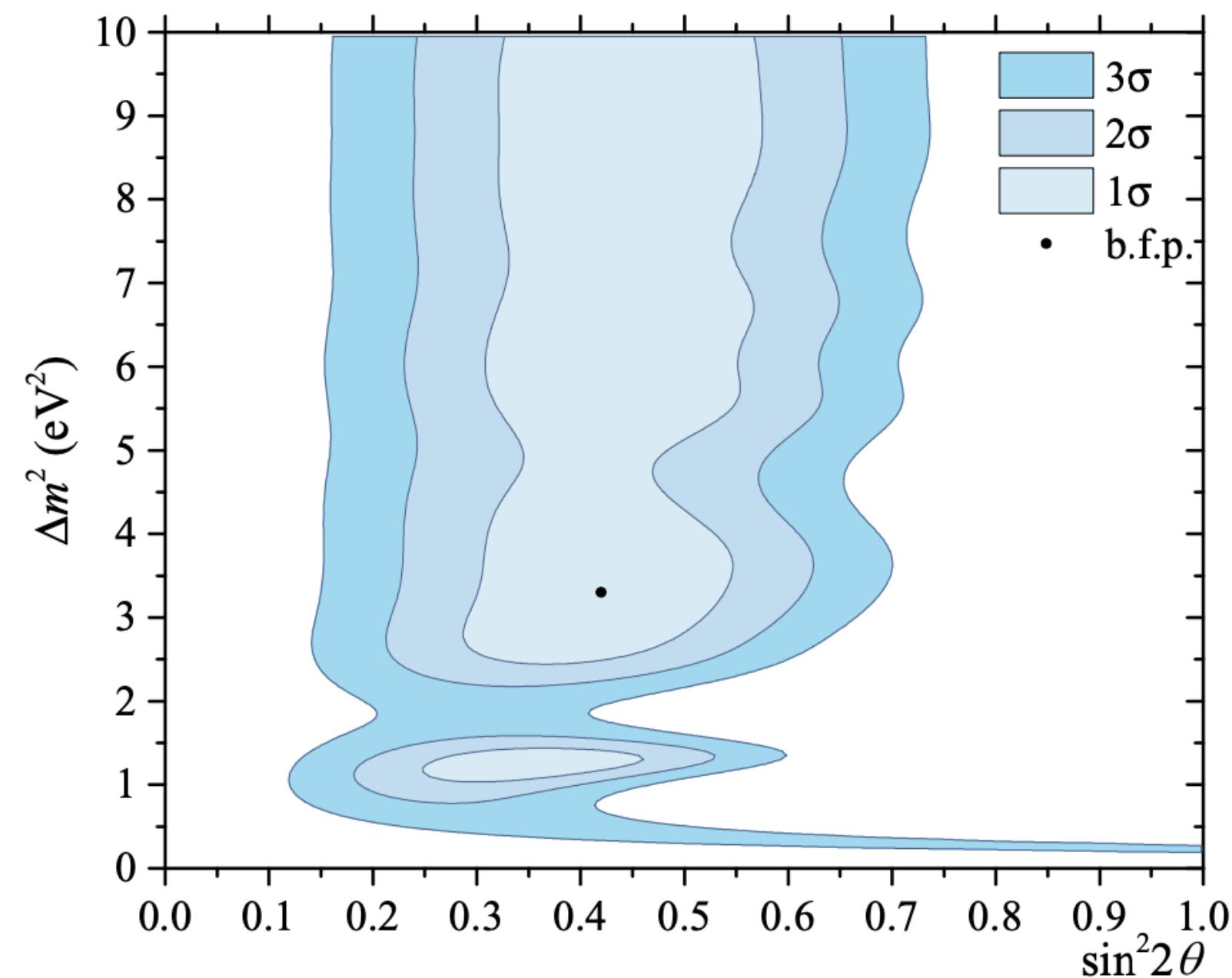
Excess of ν_e

2018 MiniBooNE
PRL

Wasn't confirmed by MicroBoone

2022
MicroBooNE PRL

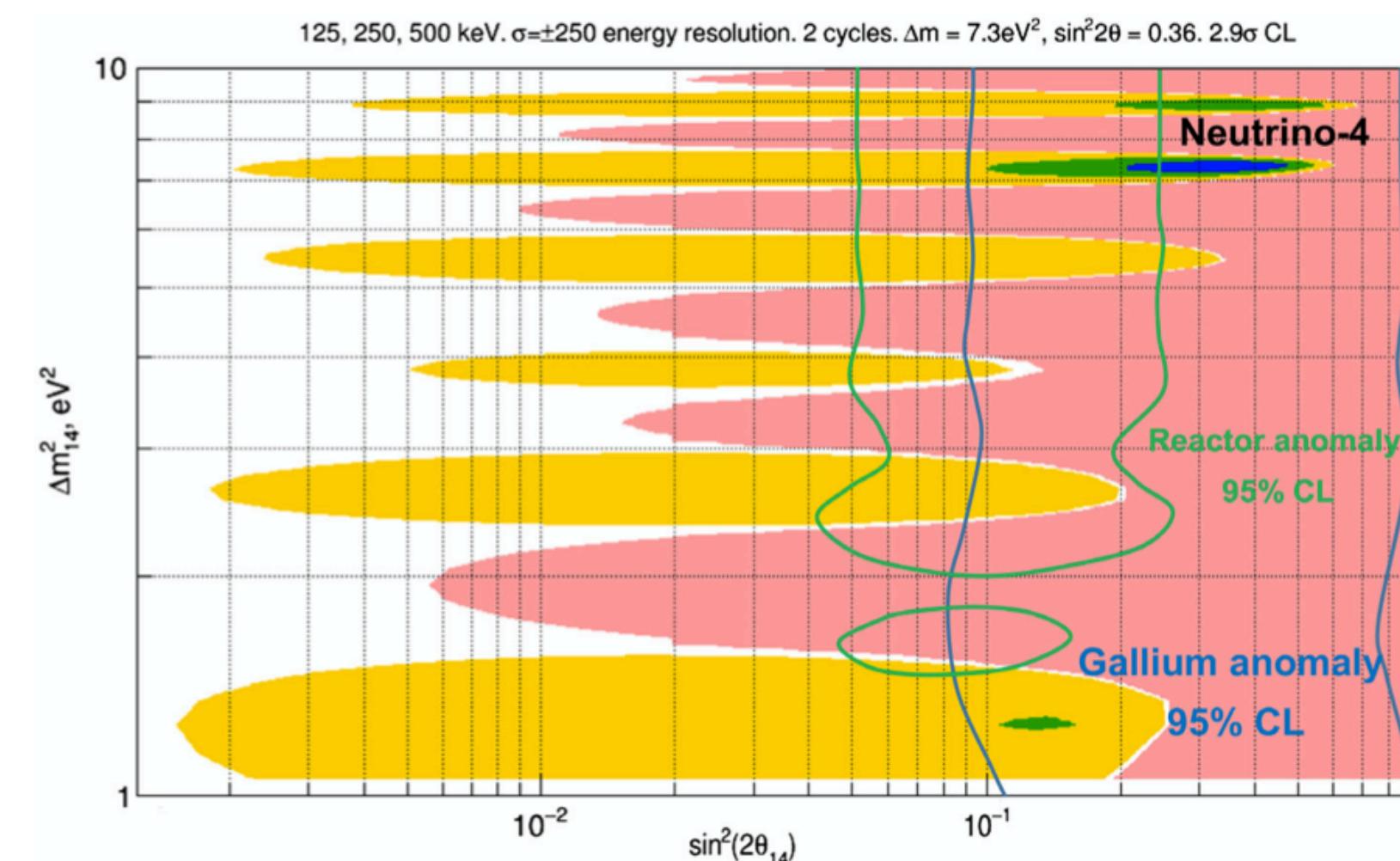
Galium:



SAGA/Gallex

They observe less events than expected in Ga-71 detectors

2022 BEST PRC

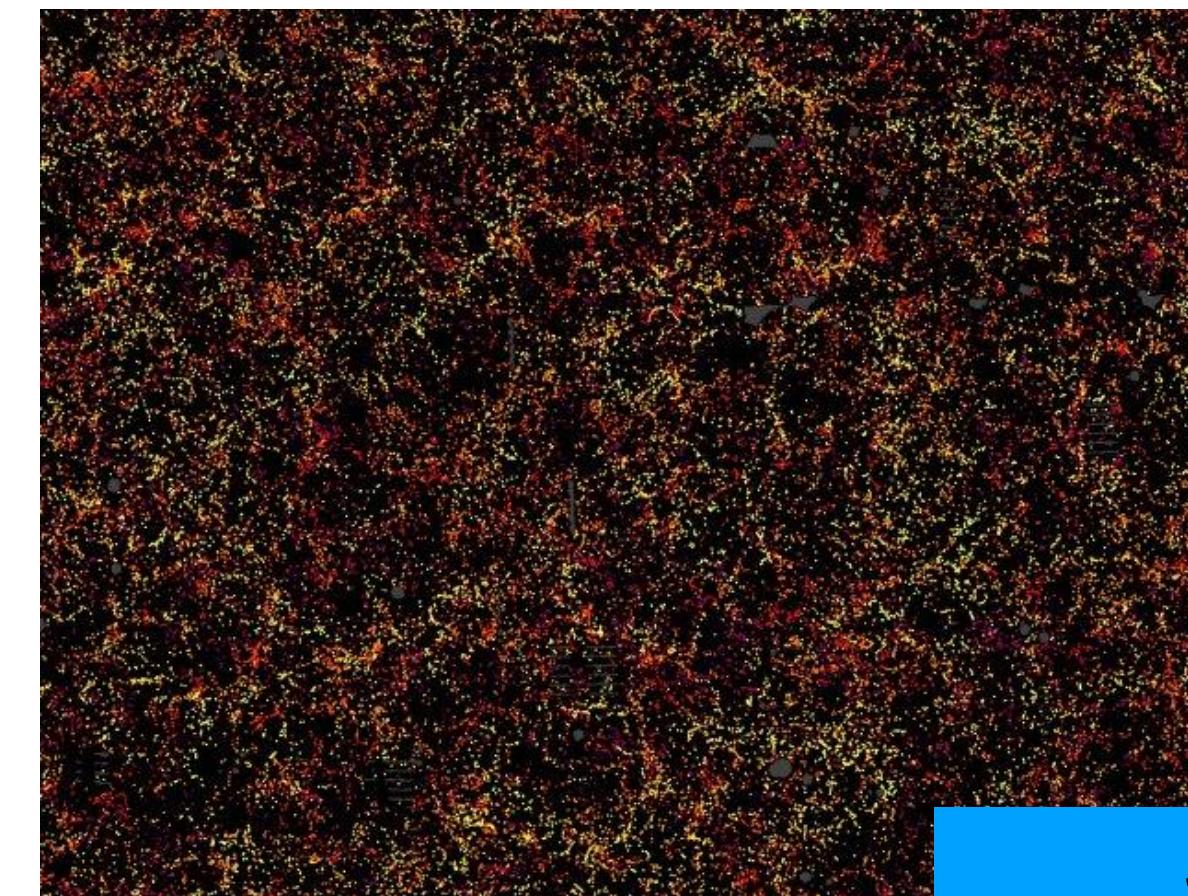
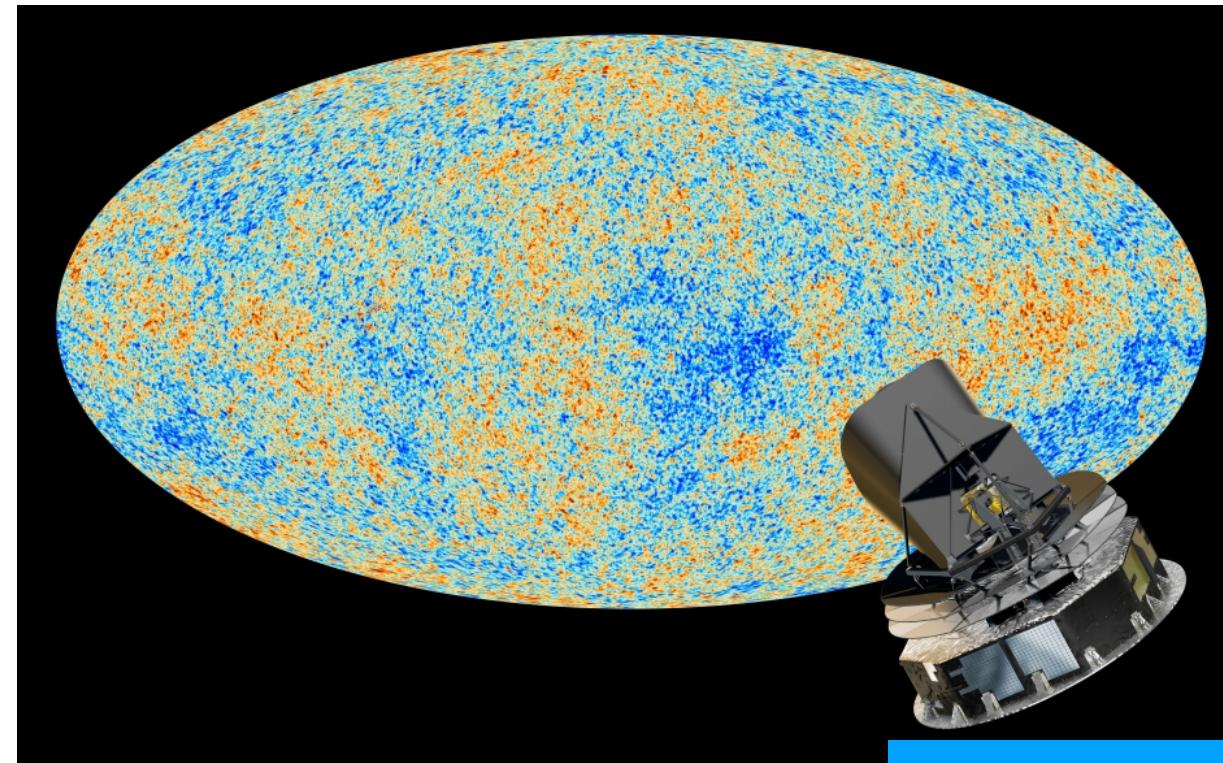
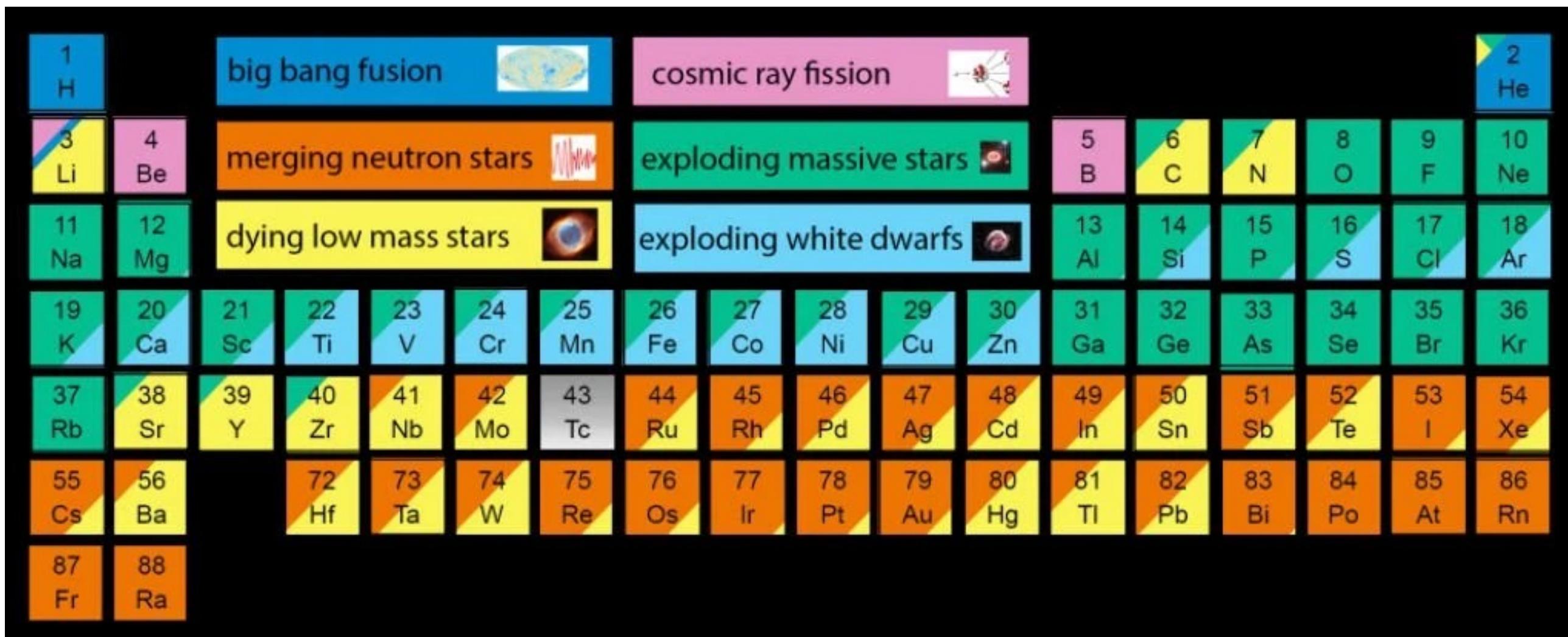


Deficit of anti(ν_e)

2021 Neutrino-4
PRD

Neutrinos are a main character in cosmology

BBN, supernovas



Basically, they destroy
structure

Neutrino cosmology

Theoretical computation

$$N_{\text{eff}} = 3.044$$

We need two params

$$N_{\text{eff}}$$

They behave as radiation in
the early Universe

(Around 40% of the total rad)

$$\sum m_\nu = m_1 + m_2 + m_3$$

Lately, they behave as
matter

$$N_{\text{eff}} = 2.99 \pm 0.17 \text{ (68 \% C.L.)}$$

$$\sum m_\nu < 0.12 \text{ (95 \% C.L.)}$$

2021 Planck collab. A&A

When the rad/matter transition occurs?

It depends on the
neutrino mass

Neutrino cosmology

$$N_{\text{eff}} = 2.99 \pm 0.17 \text{ (68 \% C.L.)}$$

$$\sum m_\nu < 0.12 \text{ (95 \% C.L.)}$$

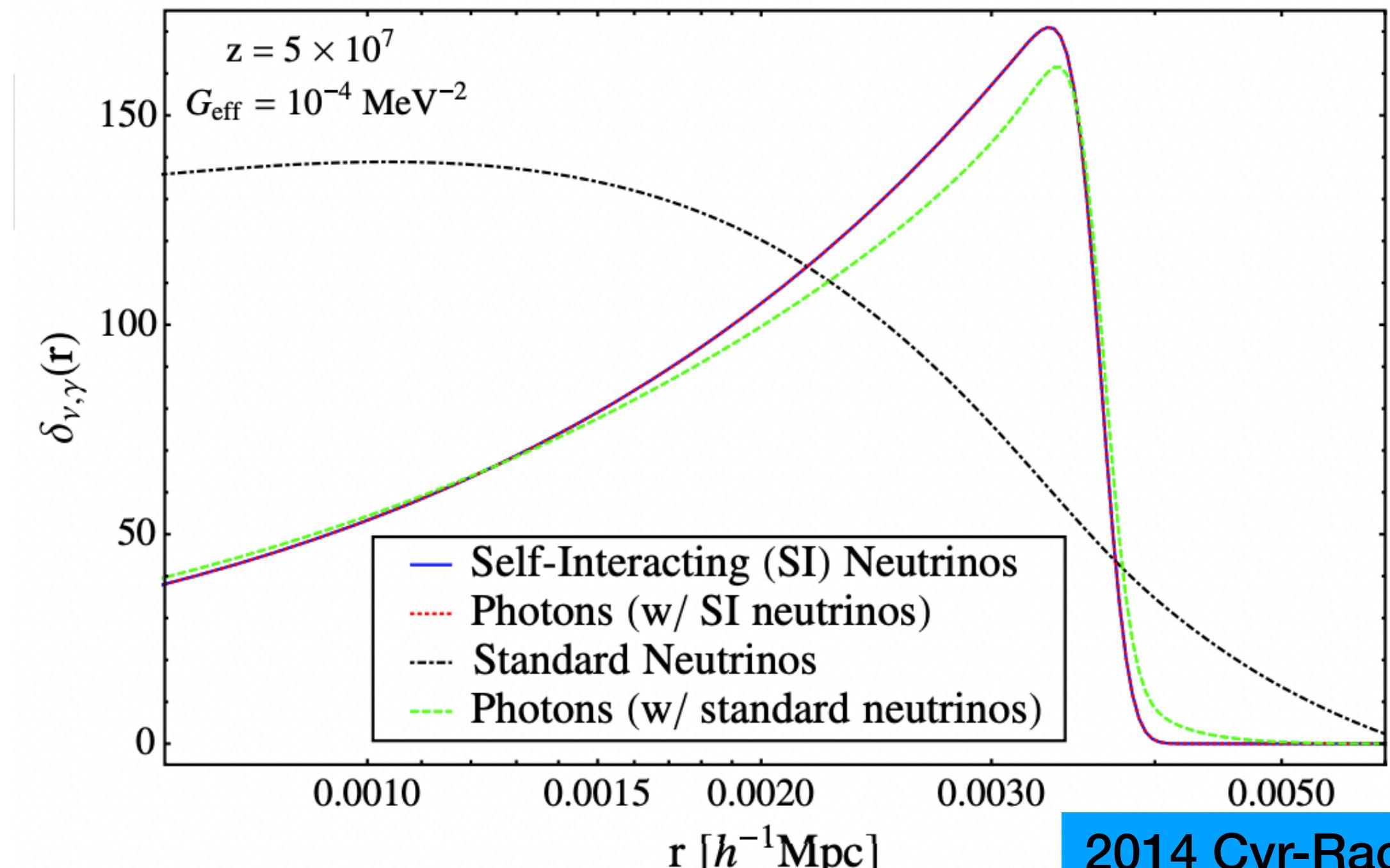
2021 Planck collab. A&A

★ In general, they suppress the matter and temperature spectra

★ They also induce a phase in the acoustic peaks that cannot be mimicked by other cosmo params



This is due to the fact that they propagate faster than sound speed on the plasma



2014 Cyr-Racine+
PRD

2004 Bashinsky & Seljak PRD

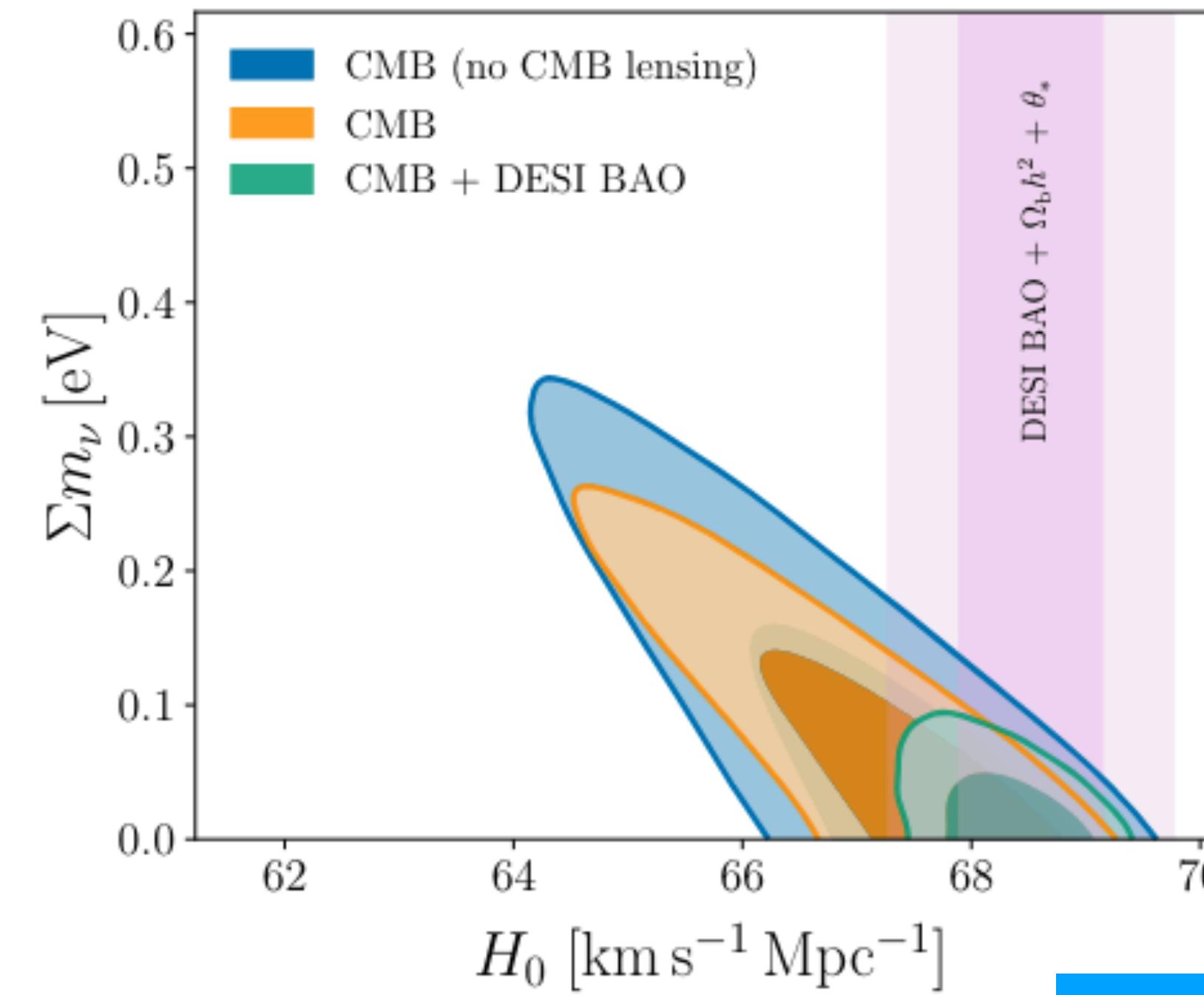
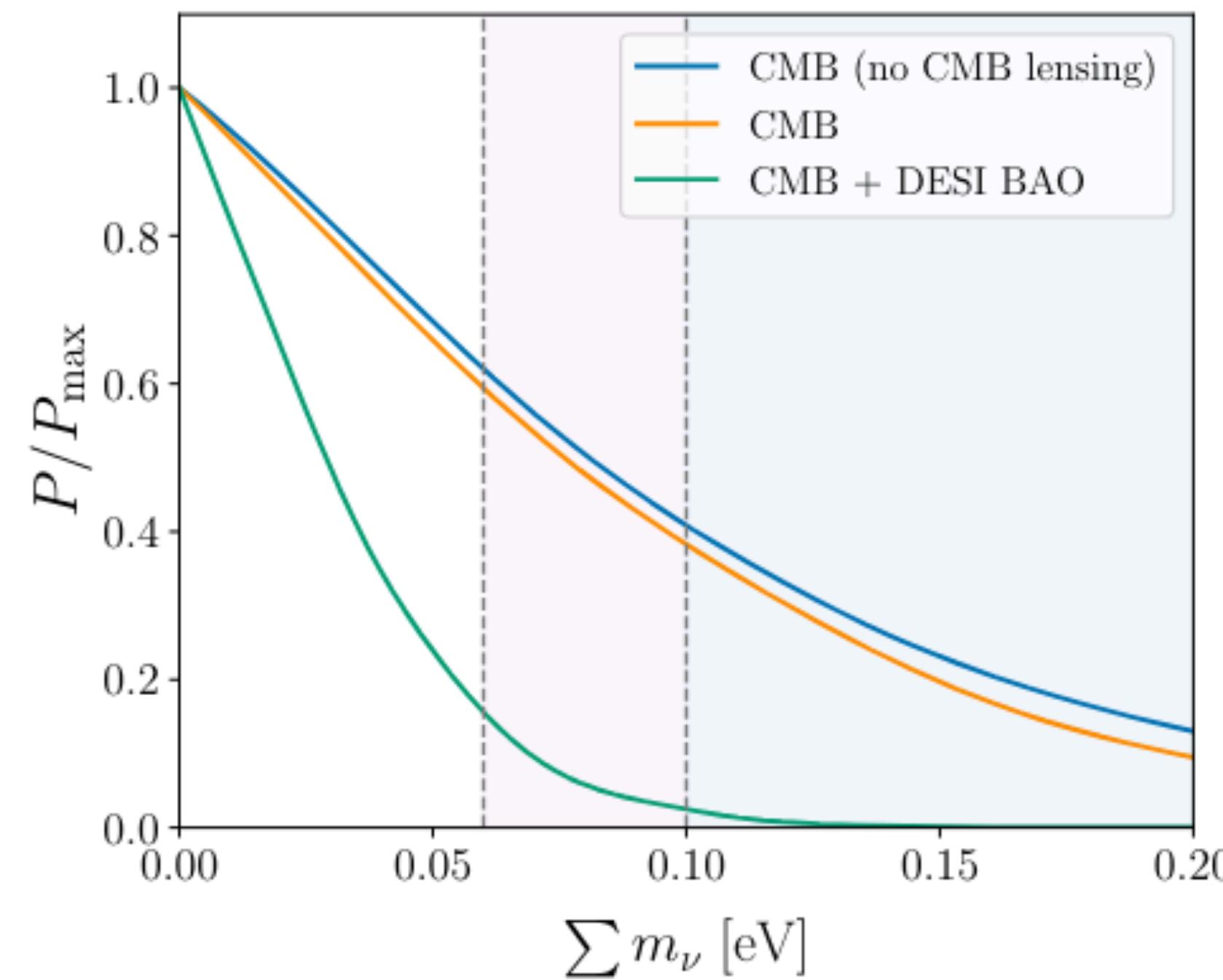
2015 Follin+ PRL

2019 Baumann+ Nat. Astrom.



Neutrino free-streaming

Neutrino cosmology - 2024



2024 DESI collab. 2404.03002

$\sum m_\nu < 0.072$ (95 % C.L.)

Fixed Neff, LCDM

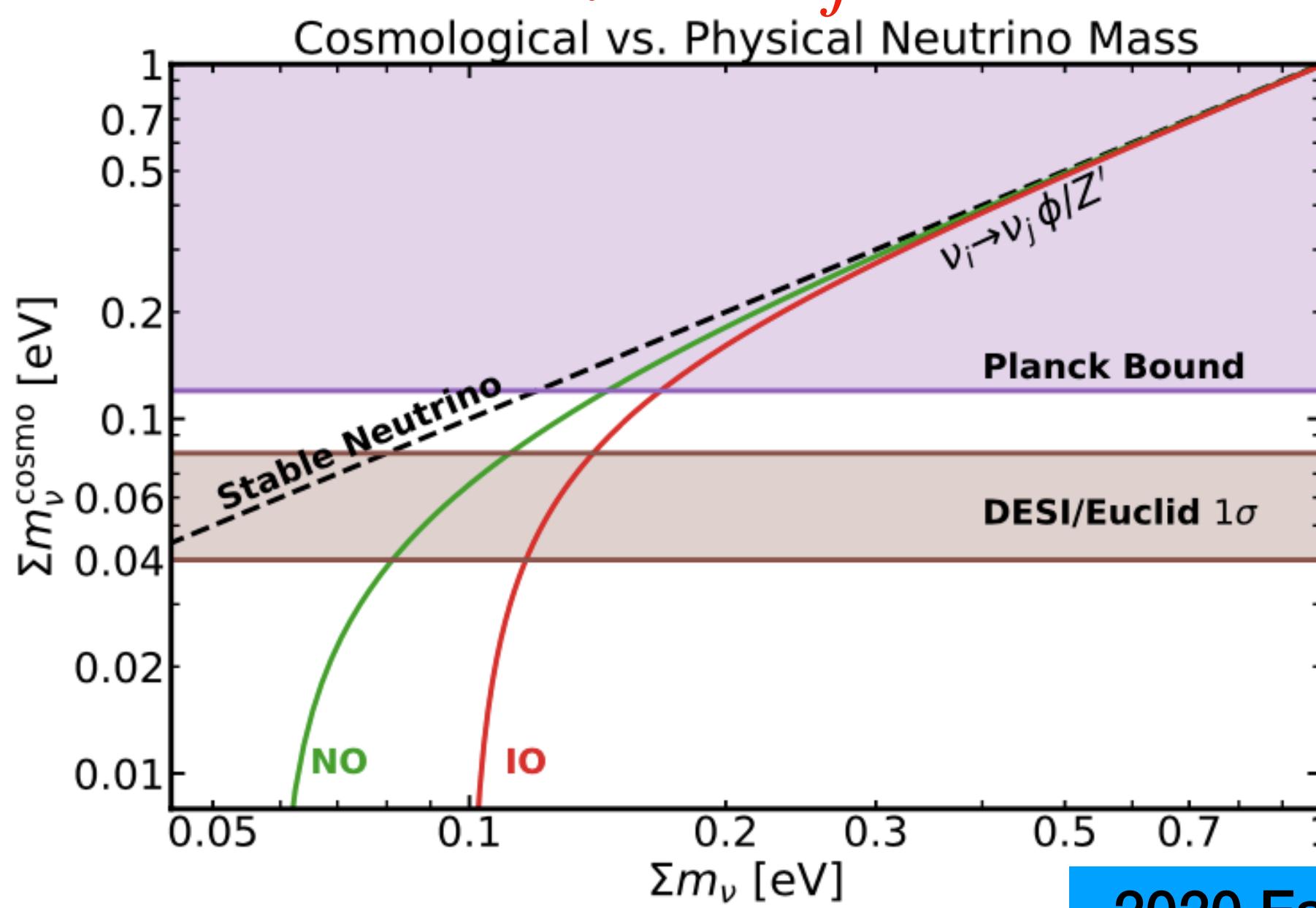
Neutrinos NSI (order zero effects)

..... a very incomplete list

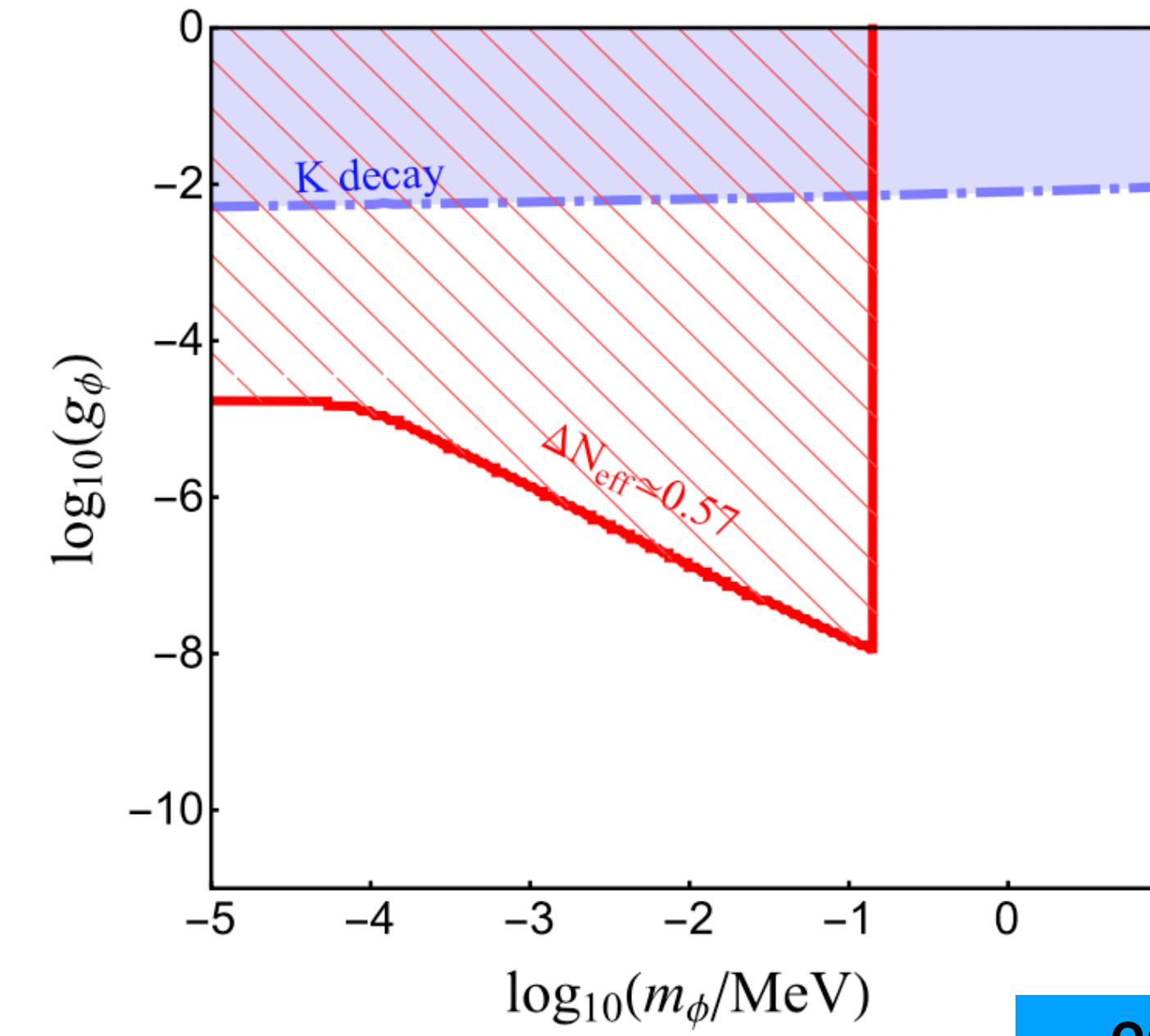
NSI can explain why cosmology
can have strong mass bounds

$$\nu + \nu \rightarrow \phi + \phi$$

$$\nu_i \rightarrow \nu_j + \phi$$



The mediator can contribute to N_{eff}



$$\nu + \nu \leftrightarrow \phi + \phi$$

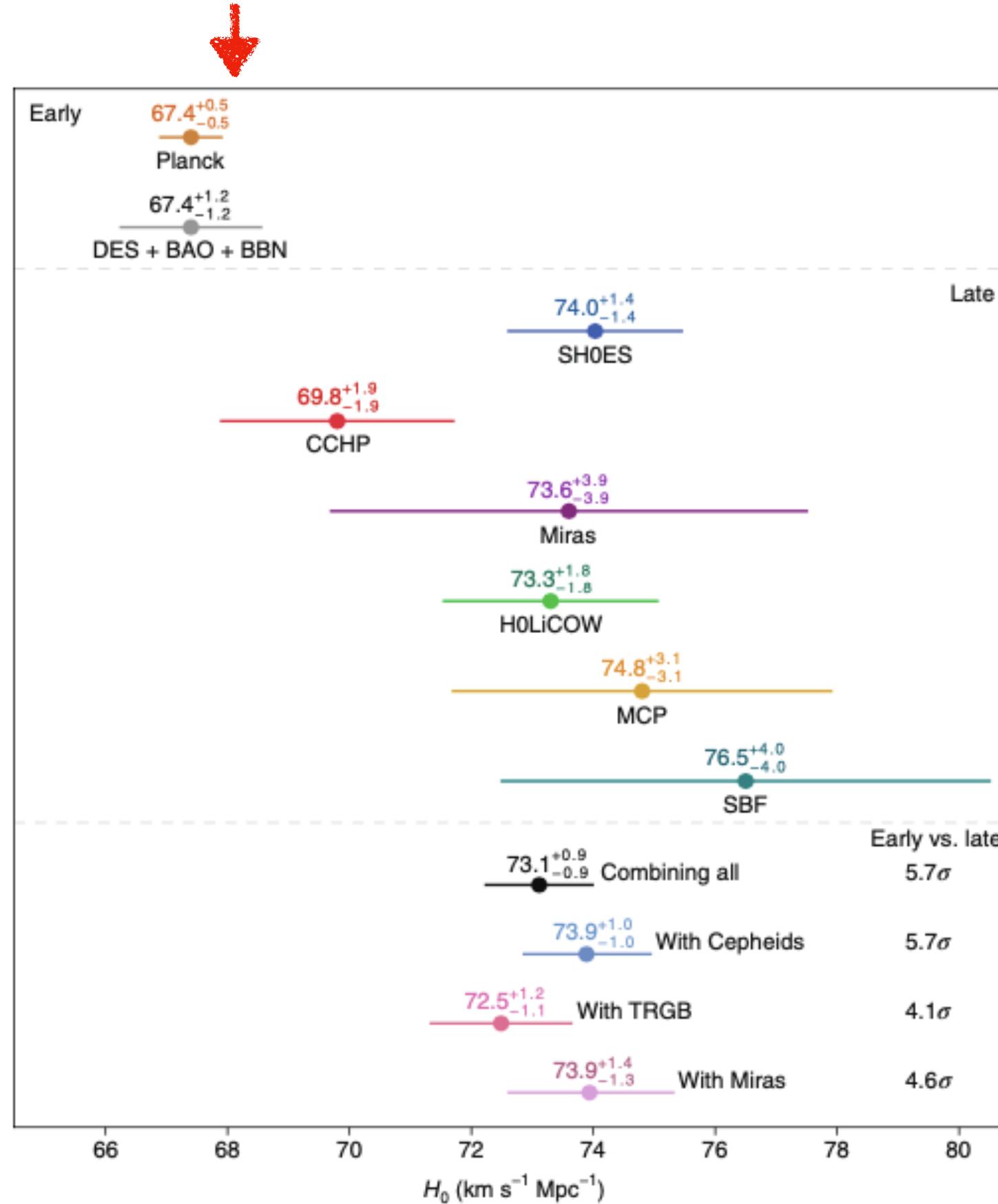
$$\nu + \phi \rightarrow \nu + \phi$$

Even if there is not
thermal equilibrium

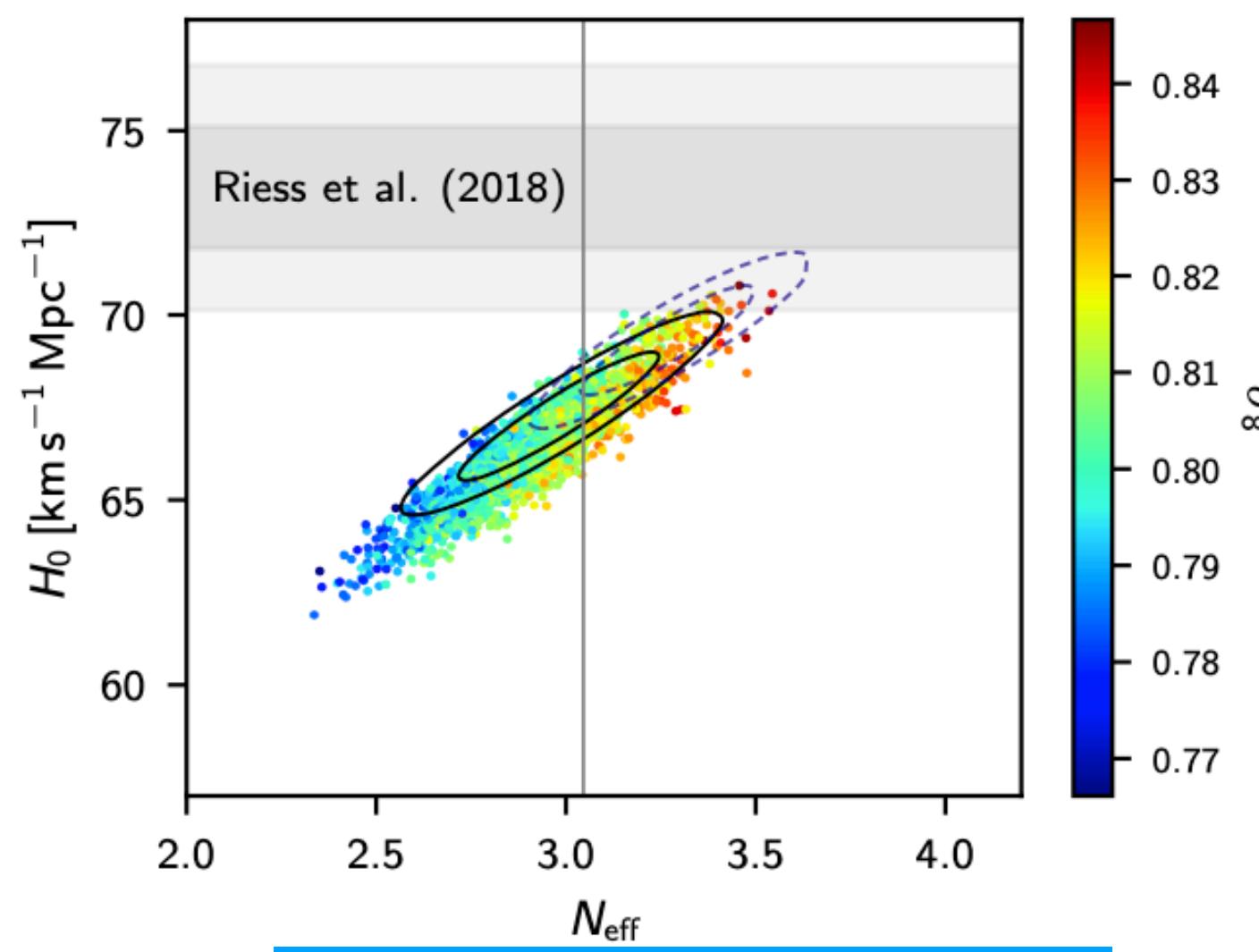
2004 Beacom+ PRL

Neutrino self-interactions... and the H0 tension

Similar results with ACT, SPT, BICEP



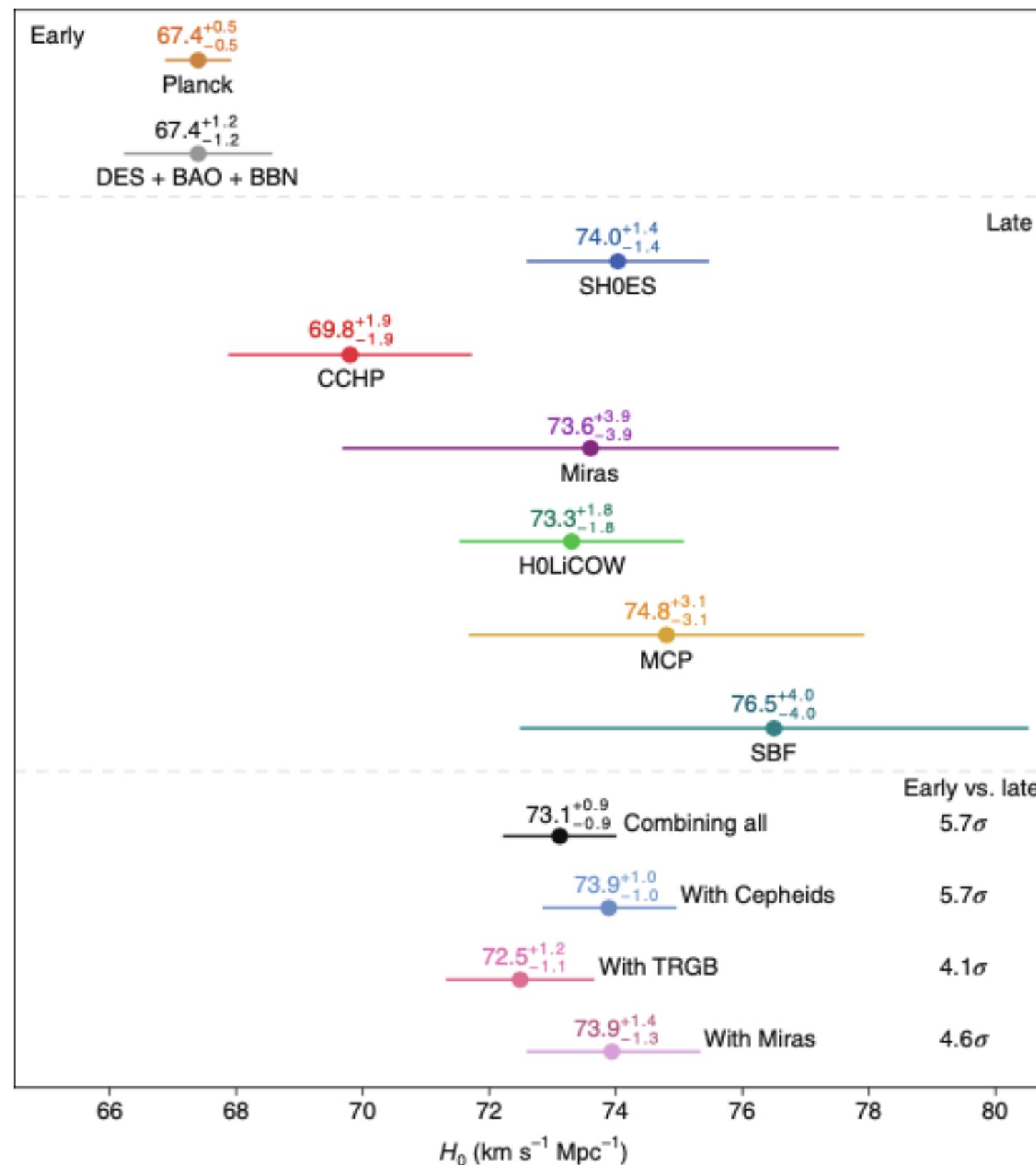
Positive correlation between N_{eff} H_0



2019 Verde+ Nature

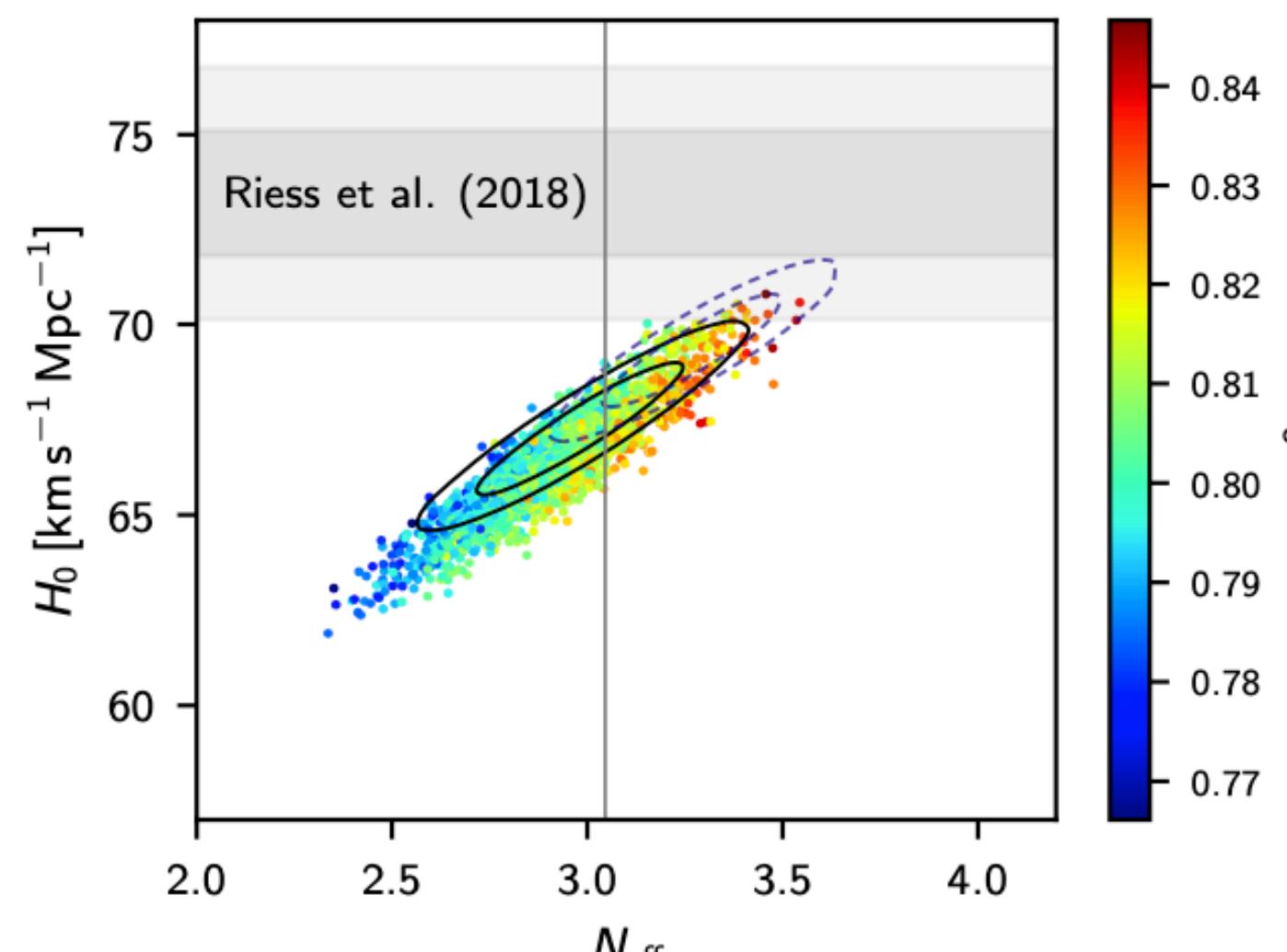
Neutrino self-interactions... and the H0 tension

Similar results with ACT, SPT, BICEP



2019 Verde+ Nature

Positive correlation between N_{eff} H_0



$$H_0 = 73.04 \pm 1.04 \text{ km s}^{-1}\text{Mpc}^{-1}$$

2112.04510

Self-interacting neutrinos lose their free-streaming

Params adjust to have N_{eff} bigger
than 3.04

Neutrino NSI imply a larger H_0

Perturbations (cosmological)

$\nu\nu \rightarrow \nu\nu$

Relaxation Time
Approximation (RTA)

$$\frac{1}{f_0} C(f) = - a\Gamma_{\text{scatt}} \Psi$$

$$\Gamma_{\text{scatt}} = \langle \sigma_0 v \rangle n_\nu$$

Density fluctuation

$$\dot{\Psi}_0 = - \frac{q}{\epsilon} k \Psi_1 + \frac{1}{6} \dot{h} \frac{d \ln f_0}{d \ln q}$$

For $l=0$ y $l=1$ null change mean
Neutrino number and momenta is conserved

Peculiar velocity

$$\dot{\Psi}_1 = \frac{qk}{3\epsilon} (\Psi_0 - 2\Psi_2)$$

Anisotropic stresses

$$\dot{\Psi}_2 = \frac{qk}{5\epsilon} (2\Psi_1 - 3\Psi_3) - \left(\frac{1}{15} \dot{h} + \frac{2}{5} \dot{\eta} \right) \frac{d \ln f_0}{d \ln q} - a\Gamma_{\text{scatt}} \Psi_2$$

Self-interactions suppress
shear y and higher
multipoles

$$\dot{\Psi}_l = \frac{qk}{(2l+1)\epsilon} [l\Psi_{l-1} - (l+1)\Psi_{l+1}] - a\Gamma_{\text{scatt}} \Psi_l \quad l \geq 3$$

Neutrino self-interactions

.... The mediator mass dictates the dynamics

$$\Gamma_{\text{scatt}} = \rightarrow$$

$$0.183 g_{\text{eff}}^4 T_\nu$$

Para

$$m_\phi < 10^{-3} \text{ eV}$$

Light mediators

$$\Gamma_{\text{res}}(g_\nu^2, m_\phi; T_\nu)$$

Para

$$10^{-3} \text{ eV} < m_\phi < 10^3 \text{ eV}$$

Resonances

$$\alpha_l G_H^2 T_\nu^5$$

Para

$$m_\phi > 10^3 \text{ eV}$$

Heavy mediator

$$G_H = \frac{|g_\nu|^2}{m_\phi^2}$$

$$g_{\text{eff}} = \xi^{1/4} g_\nu$$

G_H can be up to 9 order of magnitude
larger than G_F

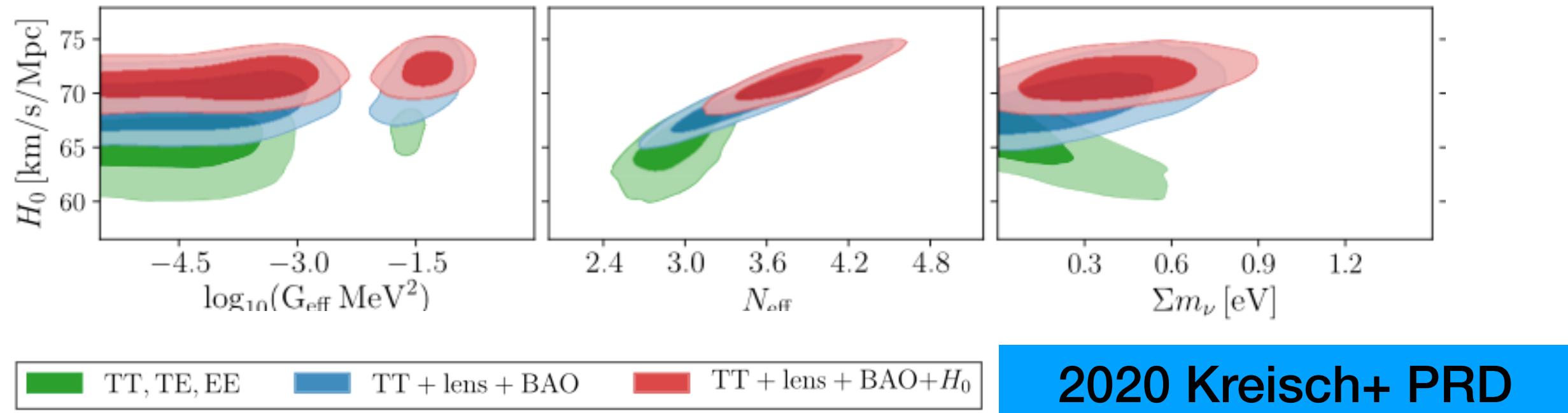
★. It all depends on WHEN



Heavy mediator affect the early Universe, while the light mediator is important lately

Self-interactions... up to 2022

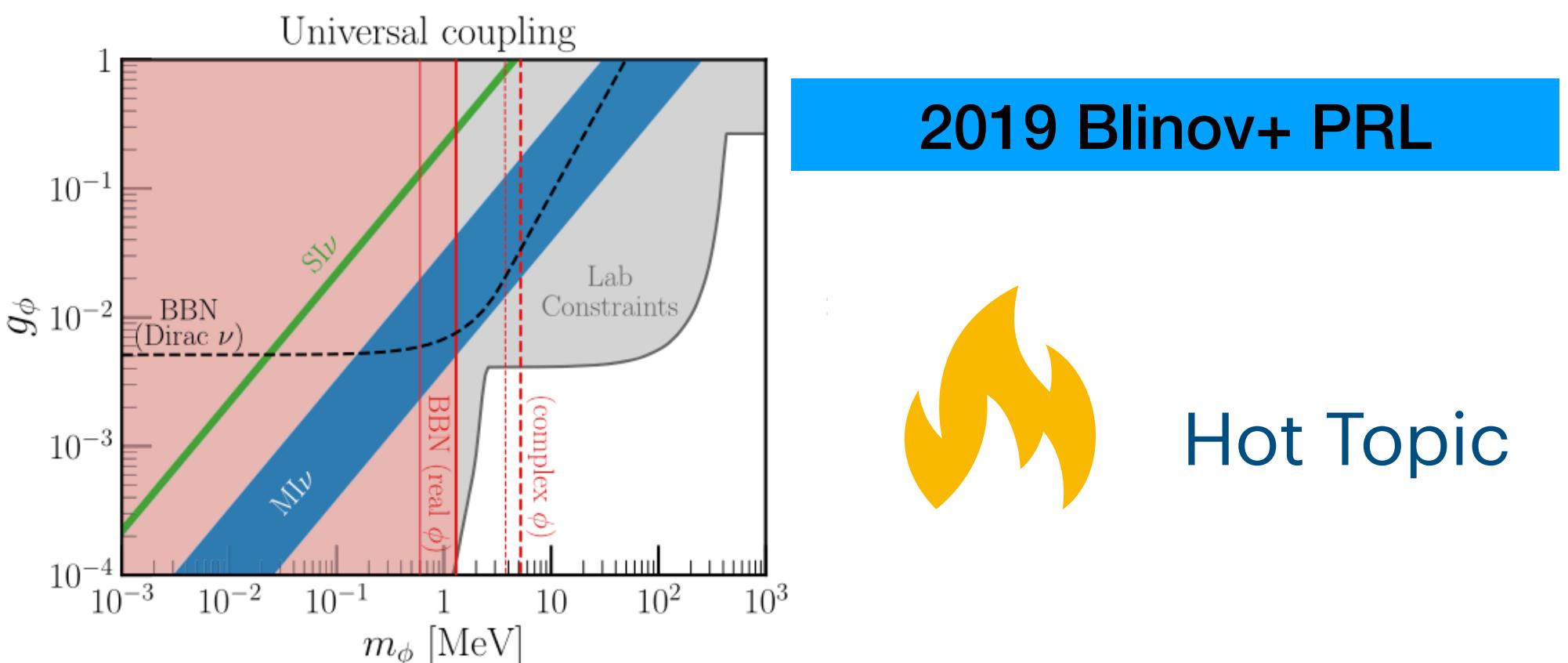
Heavy mediator



The tension can be solved (partially)

2021 Choudhury+ JCAP

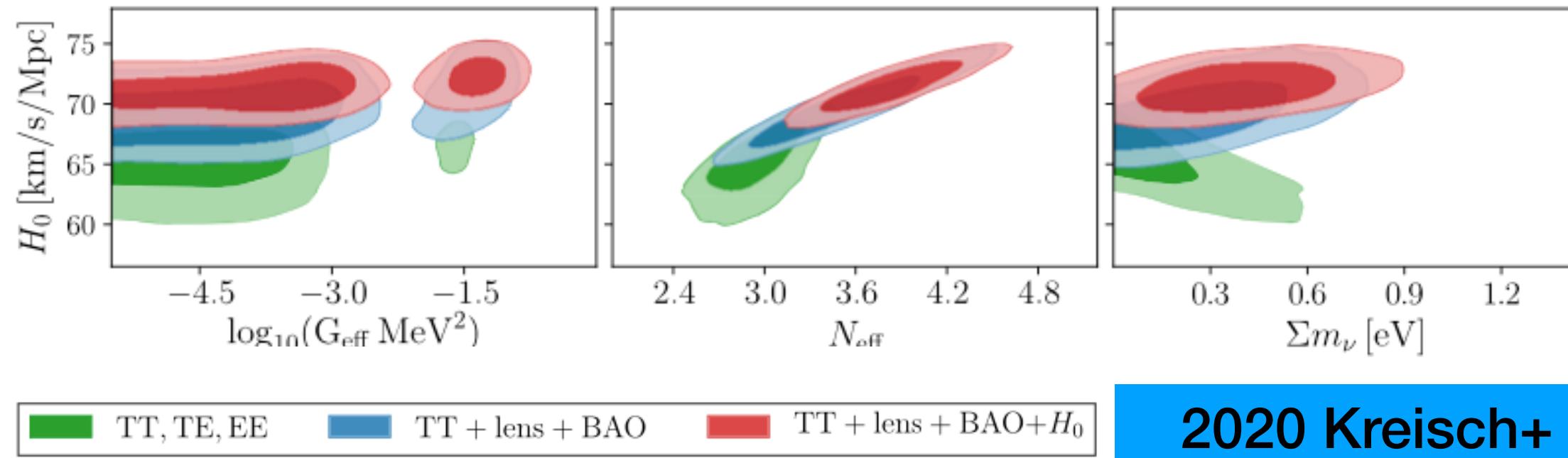
The GH values are so large to not to be seen on experiments, astrophysics and BBN



Hot Topic

Self-interactions... up to 2022

Heavy mediator

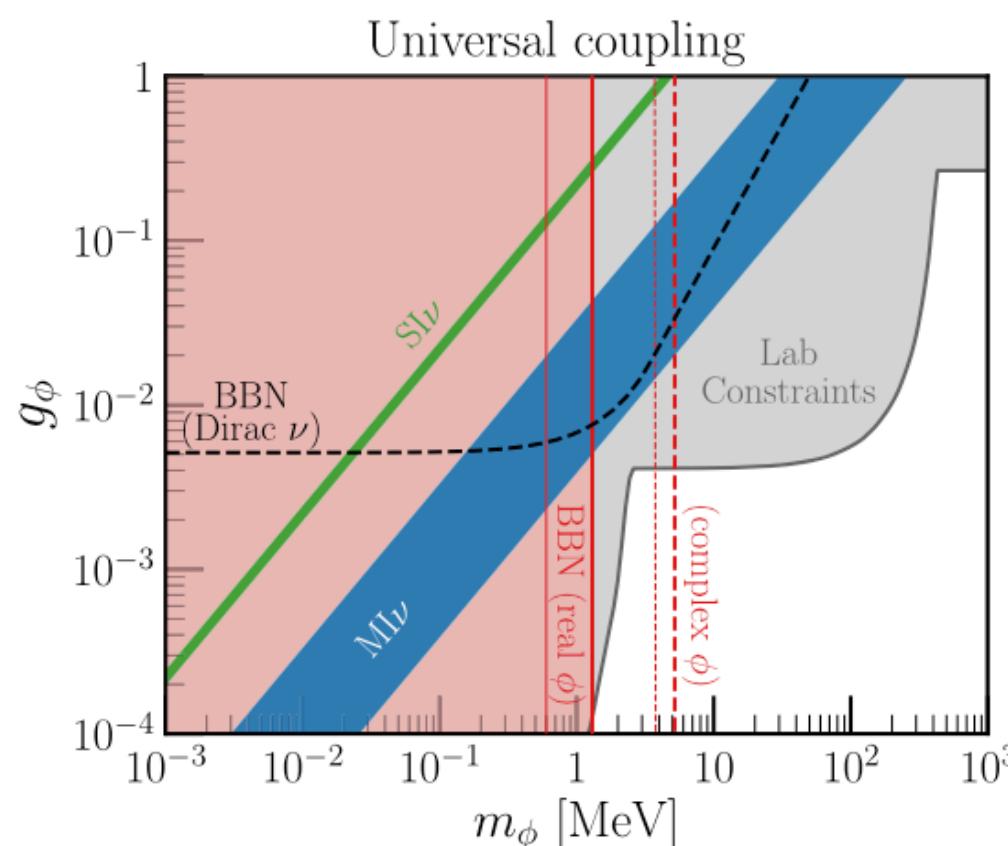


2020 Kreisch+ PRD

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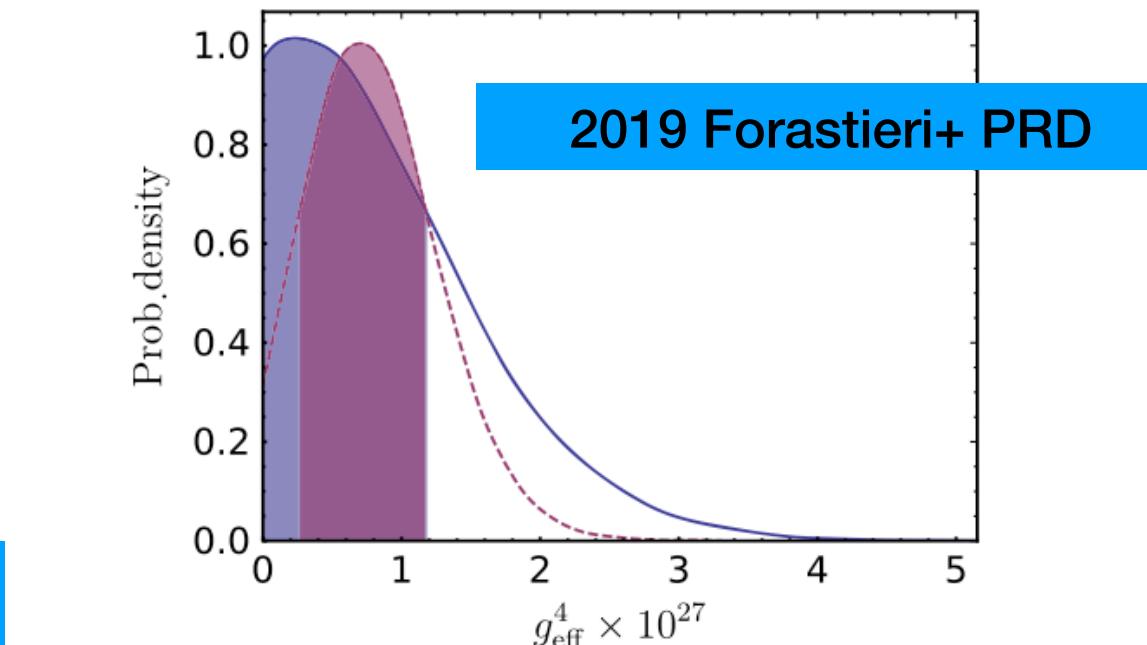
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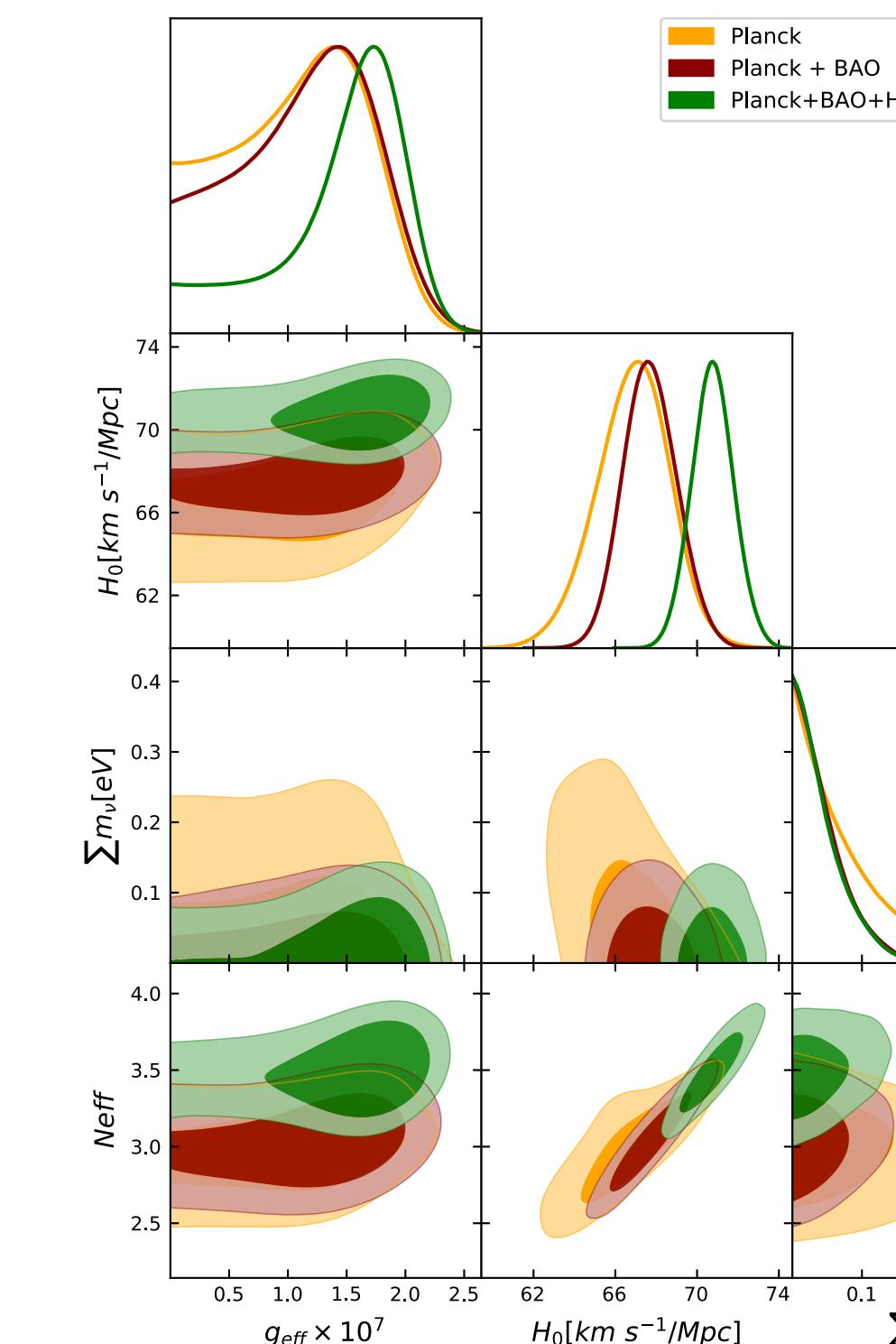
2019 Blinov+ PRL

Hot Topic

Light mediator

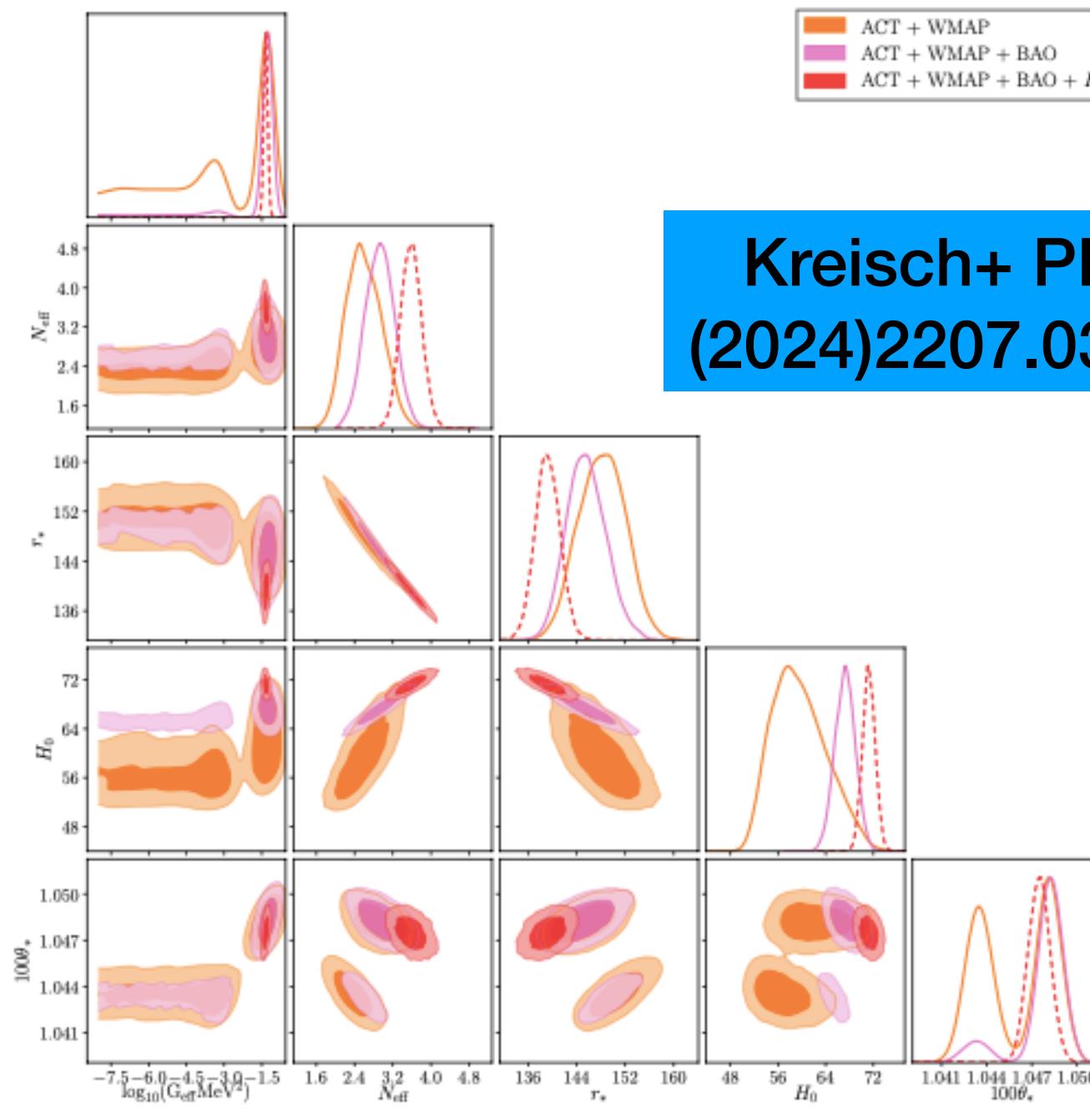


$g_{\text{eff}} < 2 \times 10^{-7}$ (95 % CL)

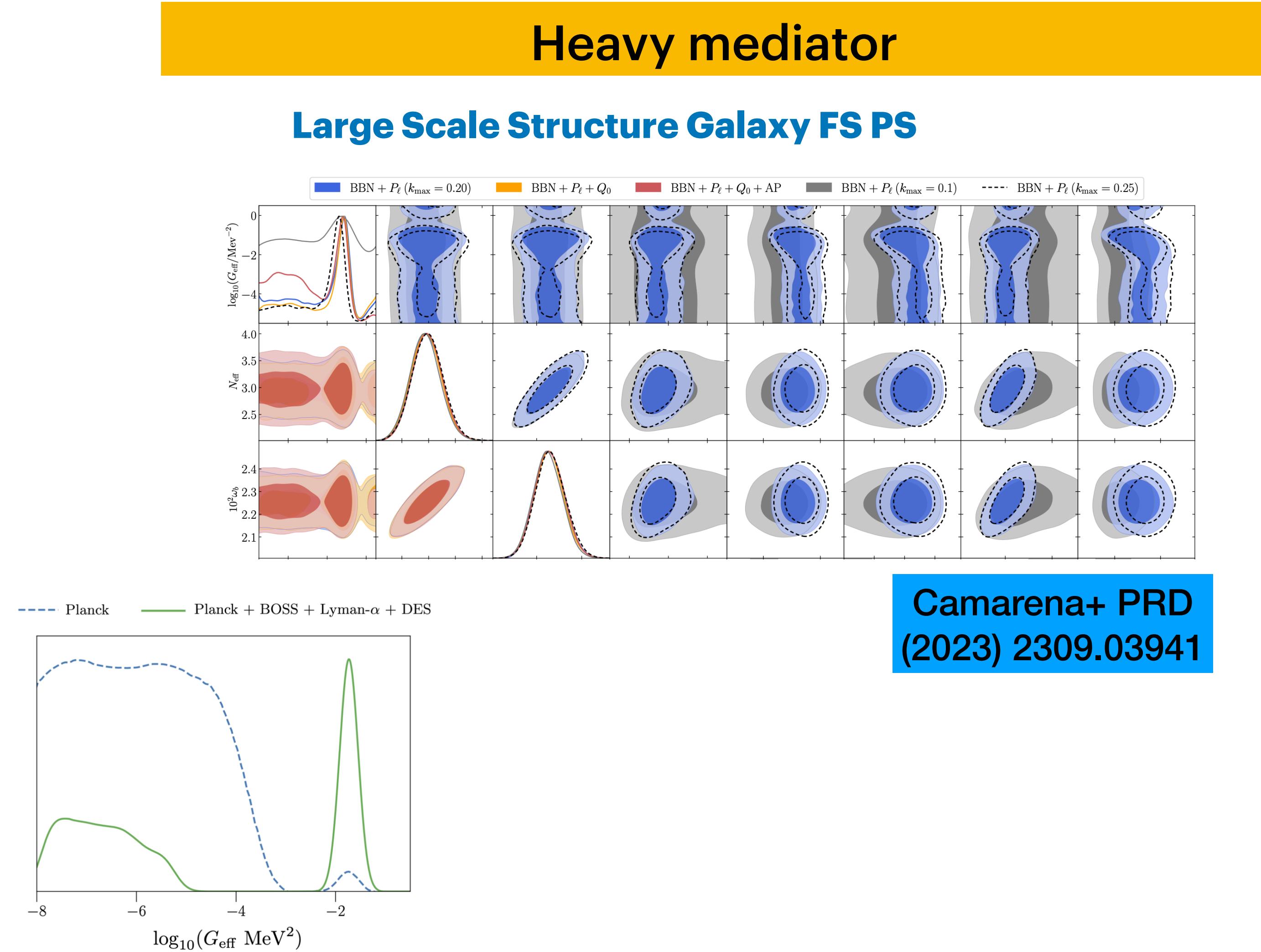


2022 Venzor+ PRD

Update...



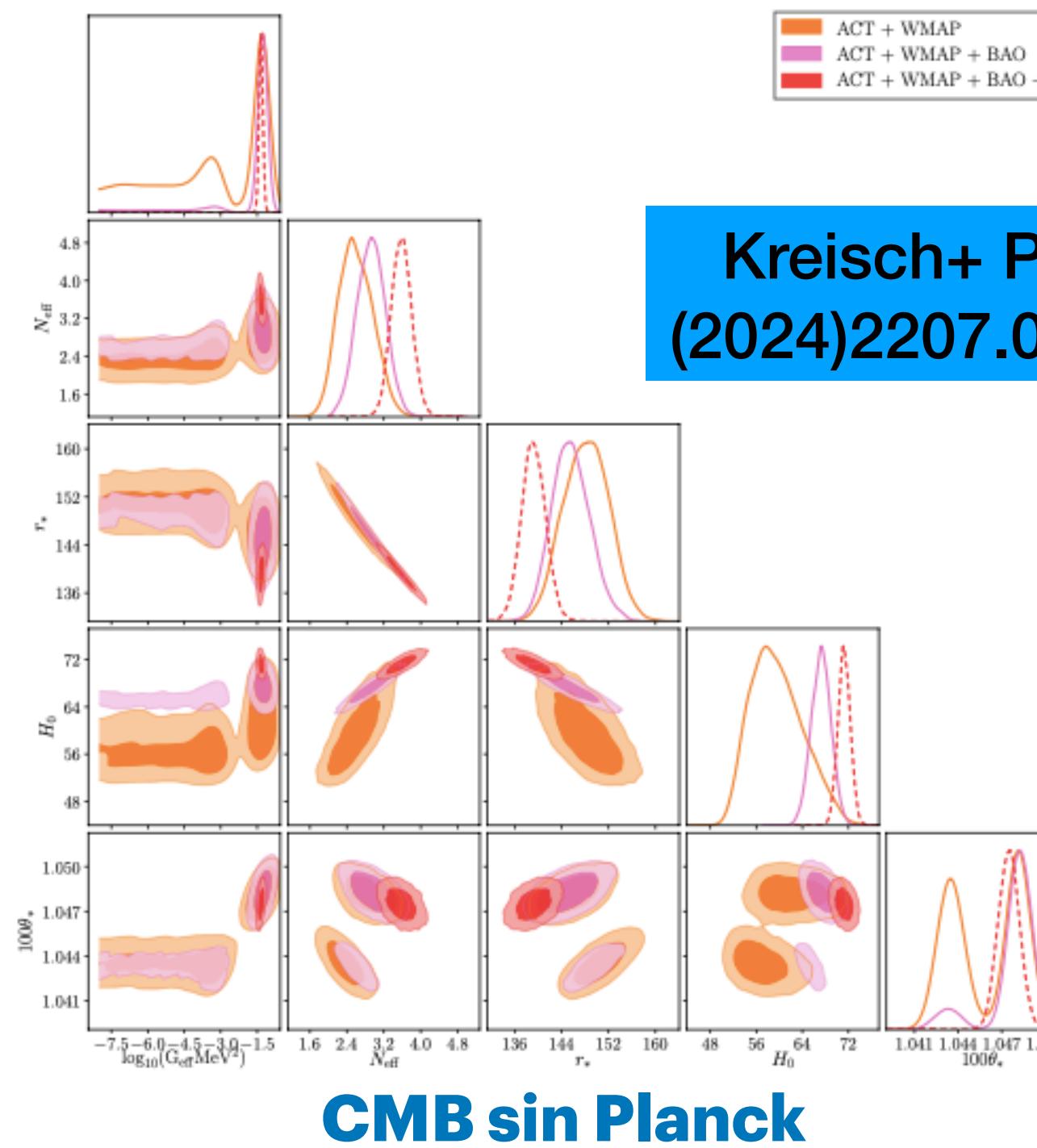
Kreisch+ PRD
(2024)2207.03164



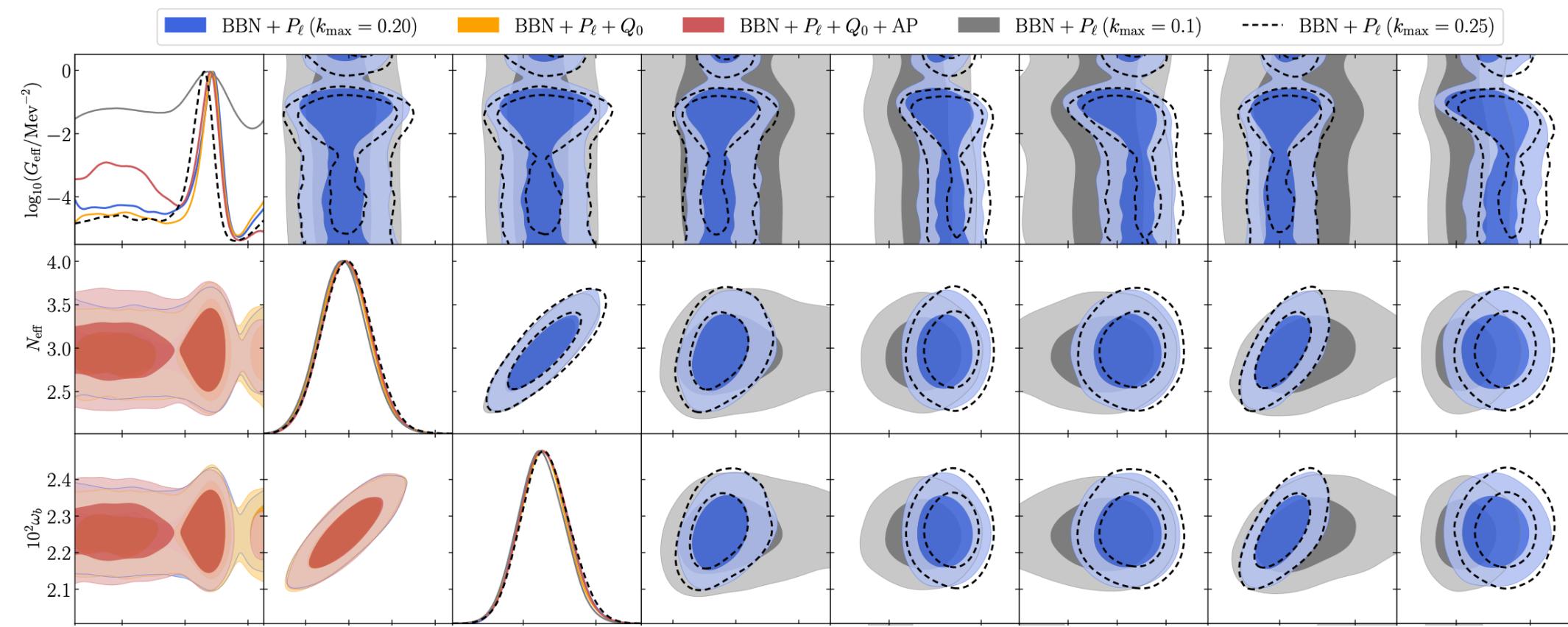
Update...

Heavy mediator

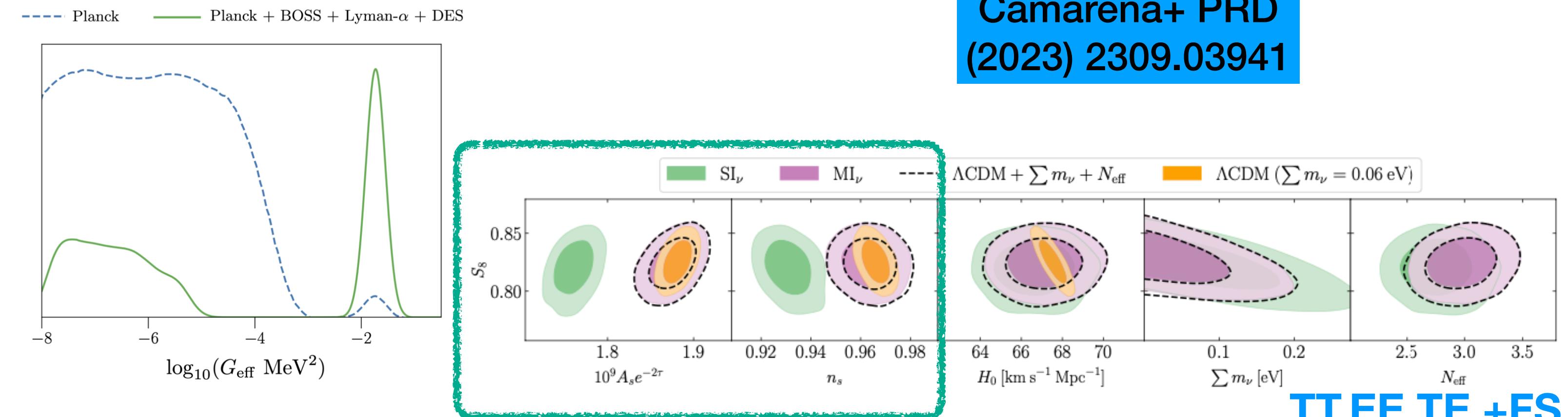
Large Scale Structure Galaxy FS PS



Kreisch+ PRD
(2024)2207.03164



Camarena+ PRD
(2023) 2309.03941

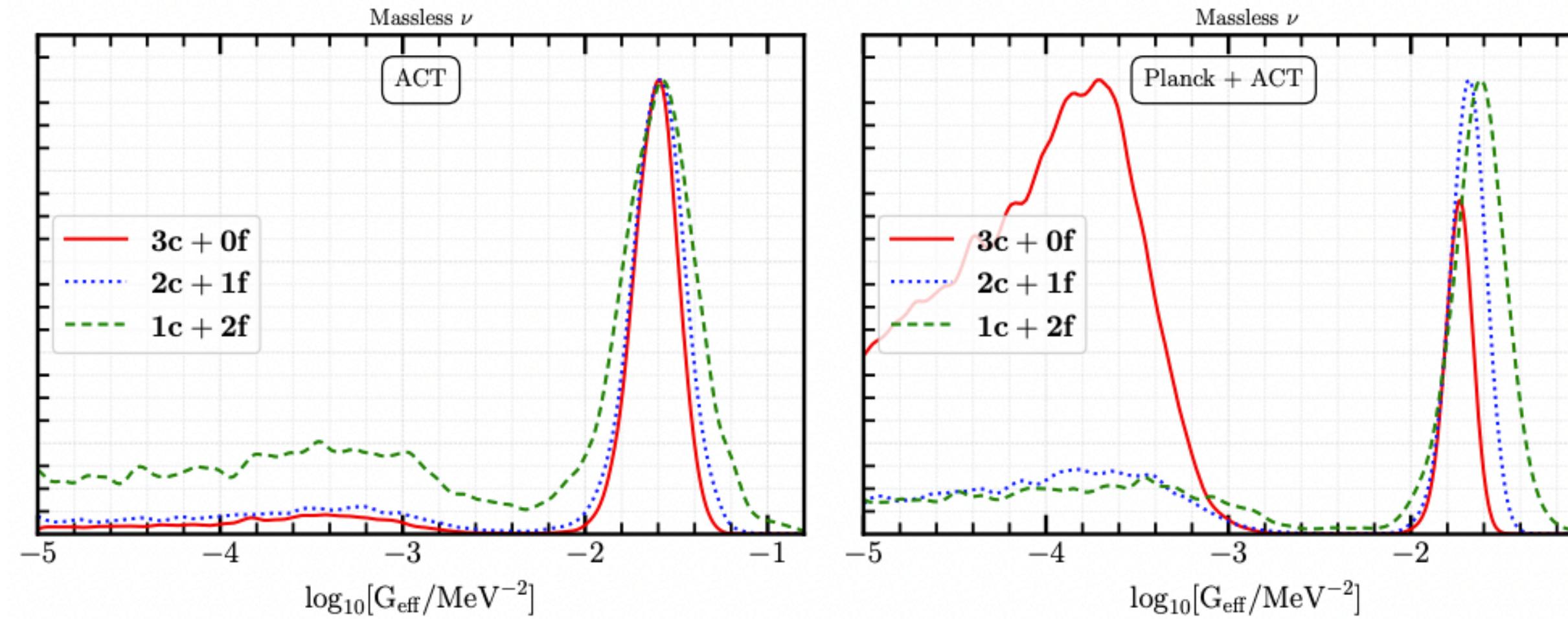


He+ PRD (2024)
2309.03956

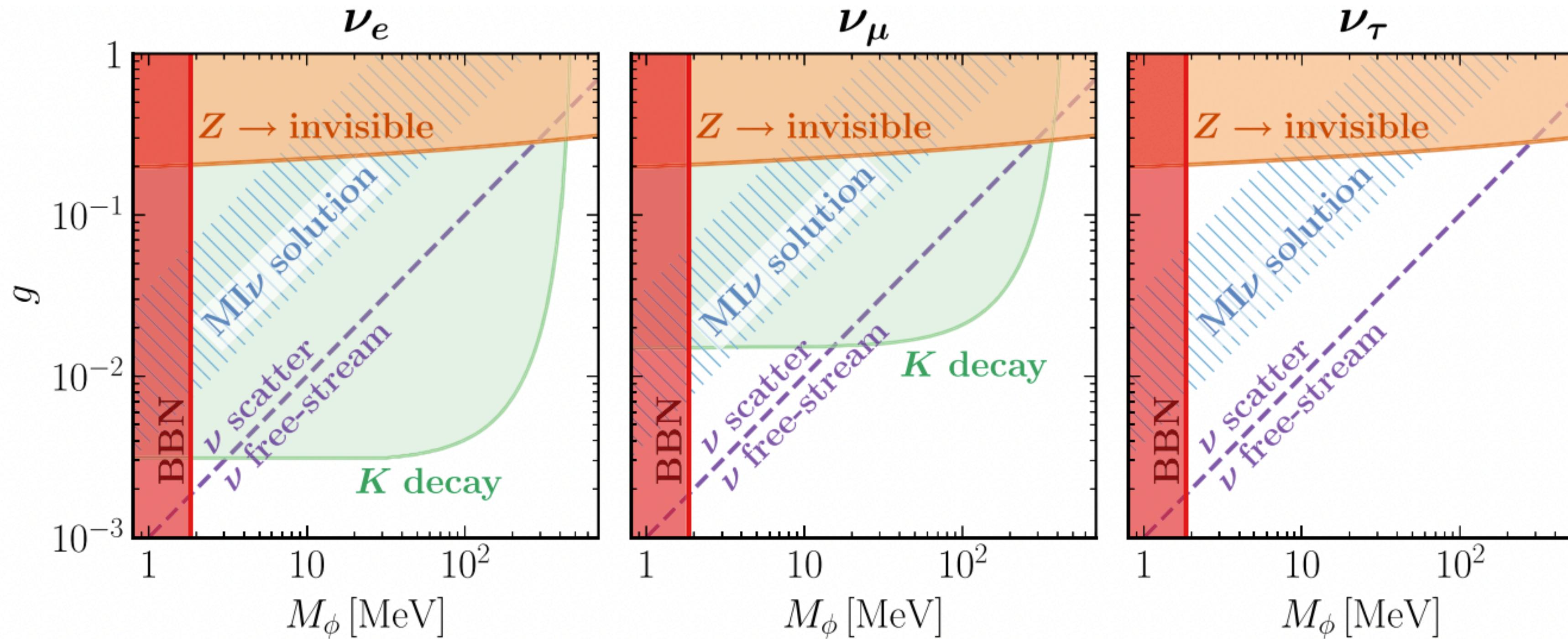
Camarena+
2403.05496

Update...

Including flavor physics



Das&Ghosh JCAP
2303.08843



2021 Esteban+
PRD

Neutrino self-interactions

.... The mediator mass dictates the dynamics

$$\Gamma_{\text{scatt}} = \rightarrow$$

$$0.183 g_{\text{eff}}^4 T_\nu \quad \text{Para} \quad m_\varphi < 10^{-3} \text{ eV}$$

$$\Gamma_{\text{res}}(g_\nu^2, m_\varphi; T_\nu) \quad \text{Para} \quad 10^{-3} \text{ eV} < m_\varphi < 10^3 \text{ eV} \quad \leftarrow$$

Resonances

$$\alpha_l G_H^2 T_\nu^5 \quad \text{Para} \quad m_\varphi > 10^3 \text{ eV}$$

Resonant region

$$10^{-3} \text{ eV} < m_\phi < 10^3 \text{ eV}$$

$$\Gamma_{\text{scatt}} = \langle \sigma_0 v \rangle n_\nu$$

We use a Breit-Wigner
cross-section

$$\sigma_0(s) = \frac{g_\nu^4}{4\pi} \frac{s}{[s - m_\phi^2]^2 + \Gamma_\phi^2 m_\phi^2}$$

$$\Gamma_\phi = \frac{g_\nu^2 m_\phi}{4\pi} \quad s = E_{\text{CM}}^2$$

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$$\Gamma_\phi = \frac{g_\nu^2 m_\phi}{4\pi} \quad s = E_{\text{CM}}^2$$

When the width is pretty small

$$\lim_{\Gamma_\phi \rightarrow 0} \sigma_0(s) = \frac{4\pi^2 \Gamma_\phi s}{m_\phi^3} \lim_{\Gamma_\phi \rightarrow 0} \frac{\Gamma_\phi m_\phi}{[s - m_\phi^2]^2 + \Gamma_\phi^2 m_\phi^2} = \boxed{\frac{\pi g_\nu^2}{m_\phi^2} s \delta(s - m_\phi^2)}$$

The cross-section is a Dirac delta

Resonant region

$$10^{-3} \text{ eV} < m_\varphi < 10^3 \text{ eV}$$

$$\langle \sigma_0 v \rangle = \frac{1}{n_\nu^2} \int \frac{d^3 p_1}{(2\pi)^3} \int \frac{d^3 p_2}{(2\pi)^3} f_{FD}(p_1) f_{FD}(p_2) \sigma_0(s) v$$



$$\Gamma_{\text{scatt}} = \langle \sigma_0 v \rangle n_\nu$$

We have to take into account all the available energies

Resonant region

$$10^{-3} \text{ eV} < m_\varphi < 10^3 \text{ eV}$$

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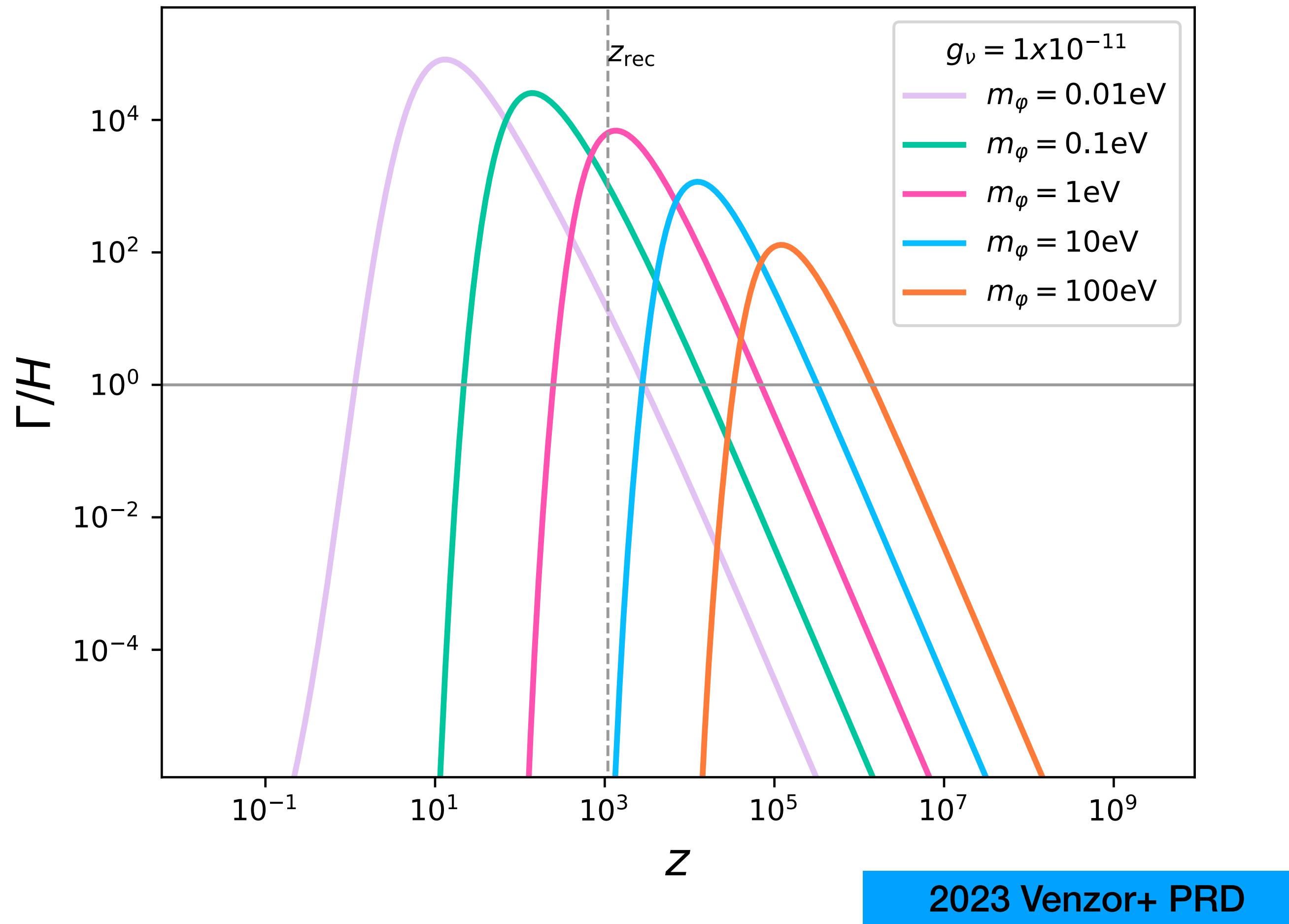
$$\Gamma_{\text{scatt}} = \langle \sigma_0 v \rangle n_\nu$$

We have to take into account all the available energies

$$\langle \sigma_0 v \rangle \rightarrow \langle \sigma_0 v_{\text{MOL}} \rangle = \frac{4\pi^2 T^2}{(2\pi)^6 n_\nu^2} \int_0^\infty \sigma_0(s) s F(s; T) ds$$

$$F(s; T) = \int_{\sqrt{s/T}}^\infty dx \frac{e^{-x}}{1 - e^{-x}} \left[\frac{\sqrt{x^2 - s/T^2}}{2} + \ln(G(x; s, T)) \right]$$

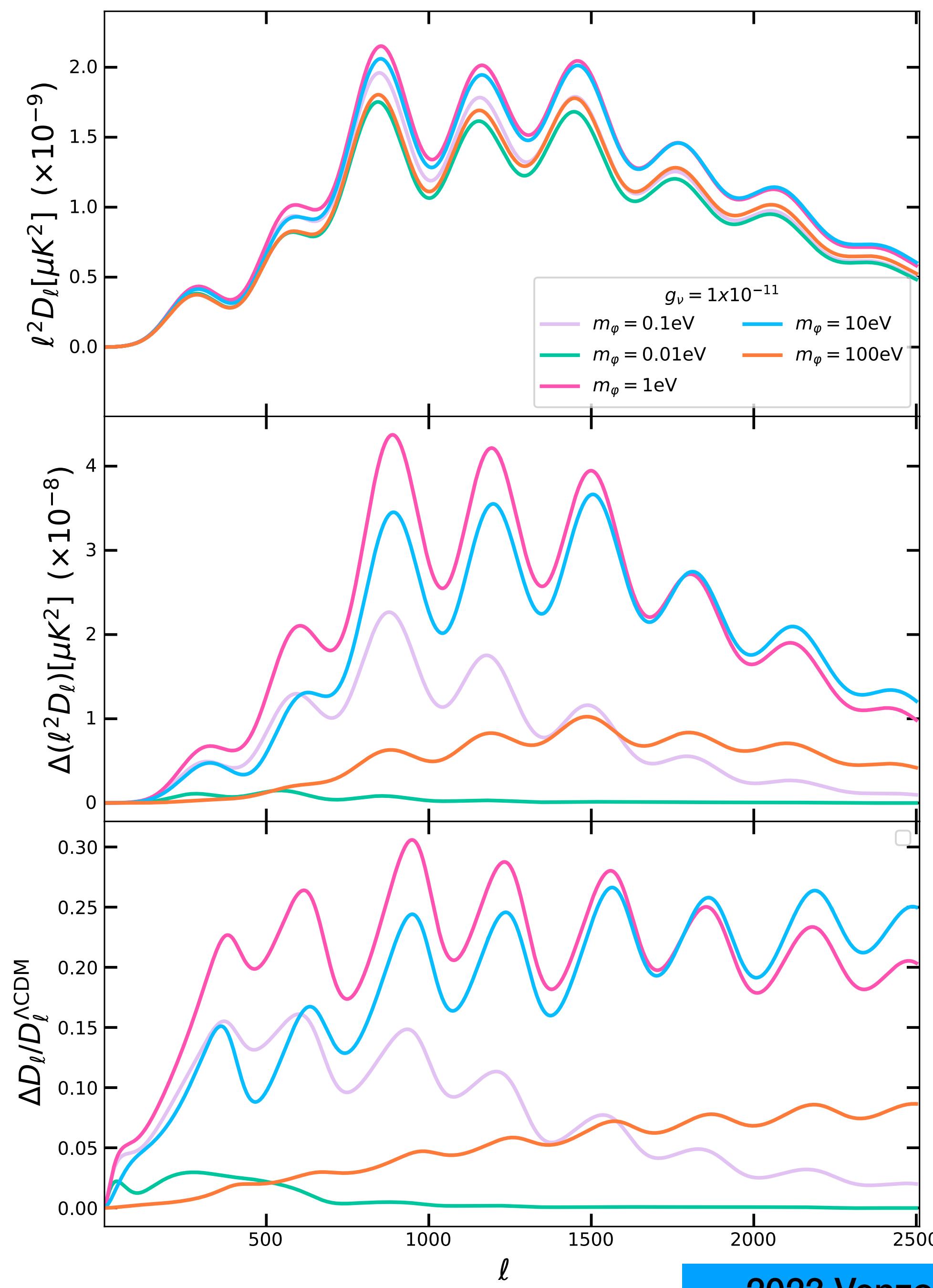
$$G(x; s, T) = \frac{1 + e^{-1/2[x + \sqrt{x^2 - s/T^2}]}}{1 + e^{-1/2[x - \sqrt{x^2 - s/T^2}]}}$$



$$\Gamma_{\text{res}} = \frac{g_\nu^2 \pi^5 m_\varphi^2}{24 \zeta(3) T_\nu} F(m_\varphi^2; T_\nu)$$

$$z_{\text{peak}} \sim 2.13 \times 10^3 \left(\frac{m_\varphi}{\text{eV}} \right)$$

We have to fix the mediator mass, the only new free parameter is the coupling



We can see an enhancement on the temperature spectral

1 and 10 eV have a stronger effect (with the same coupling)

Large mediator mass

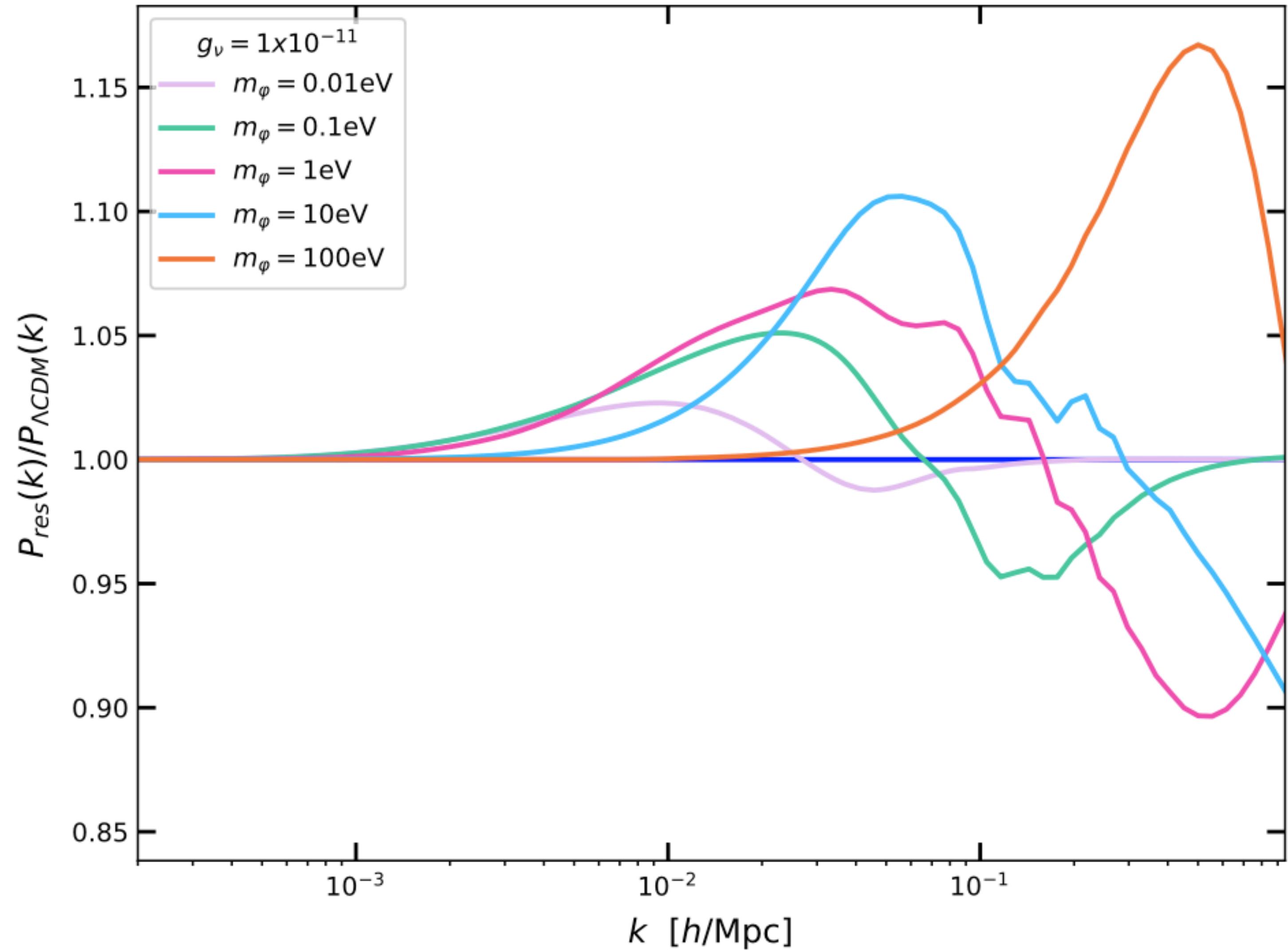
Small mass



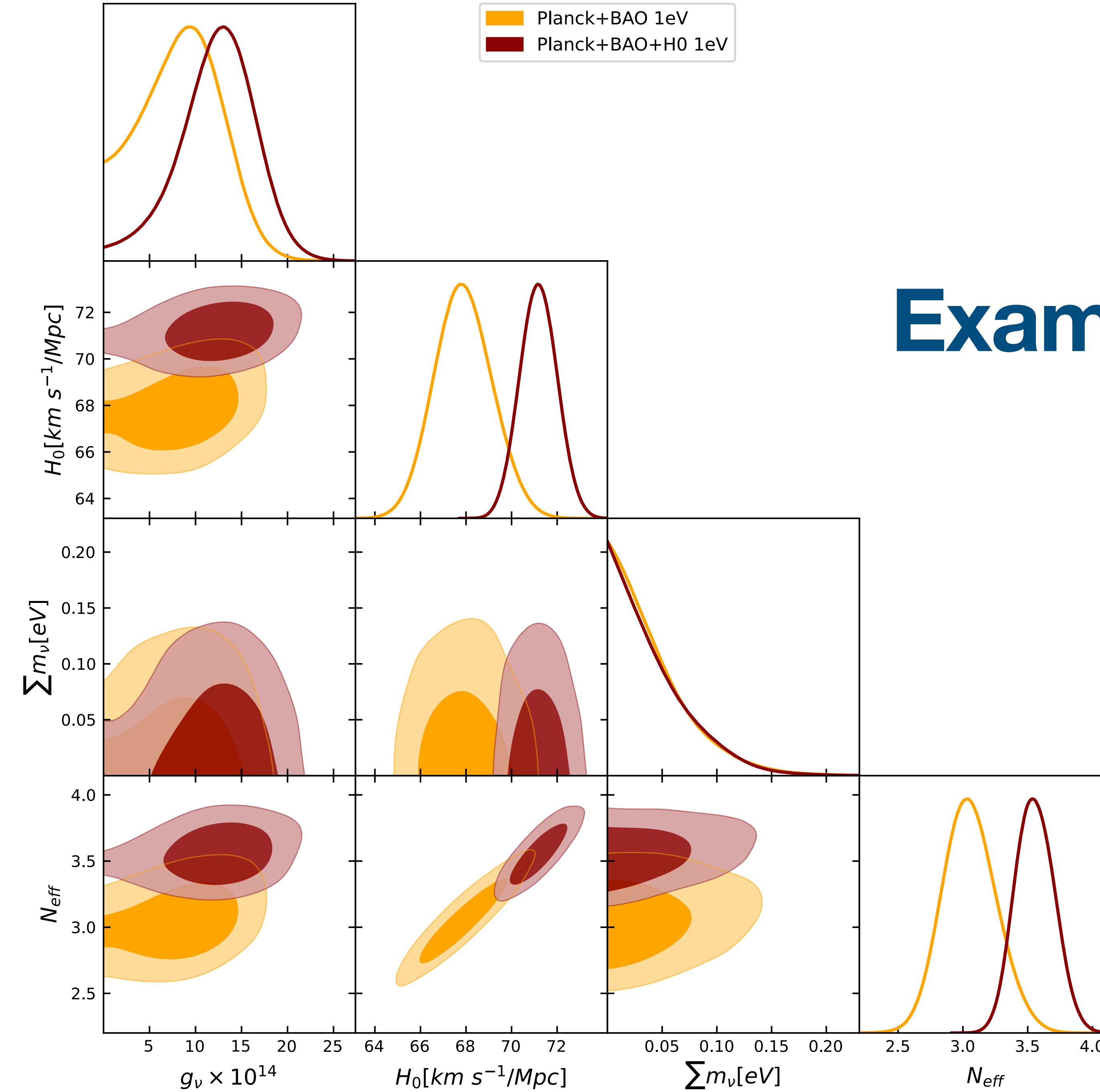
Small scales



Large scales

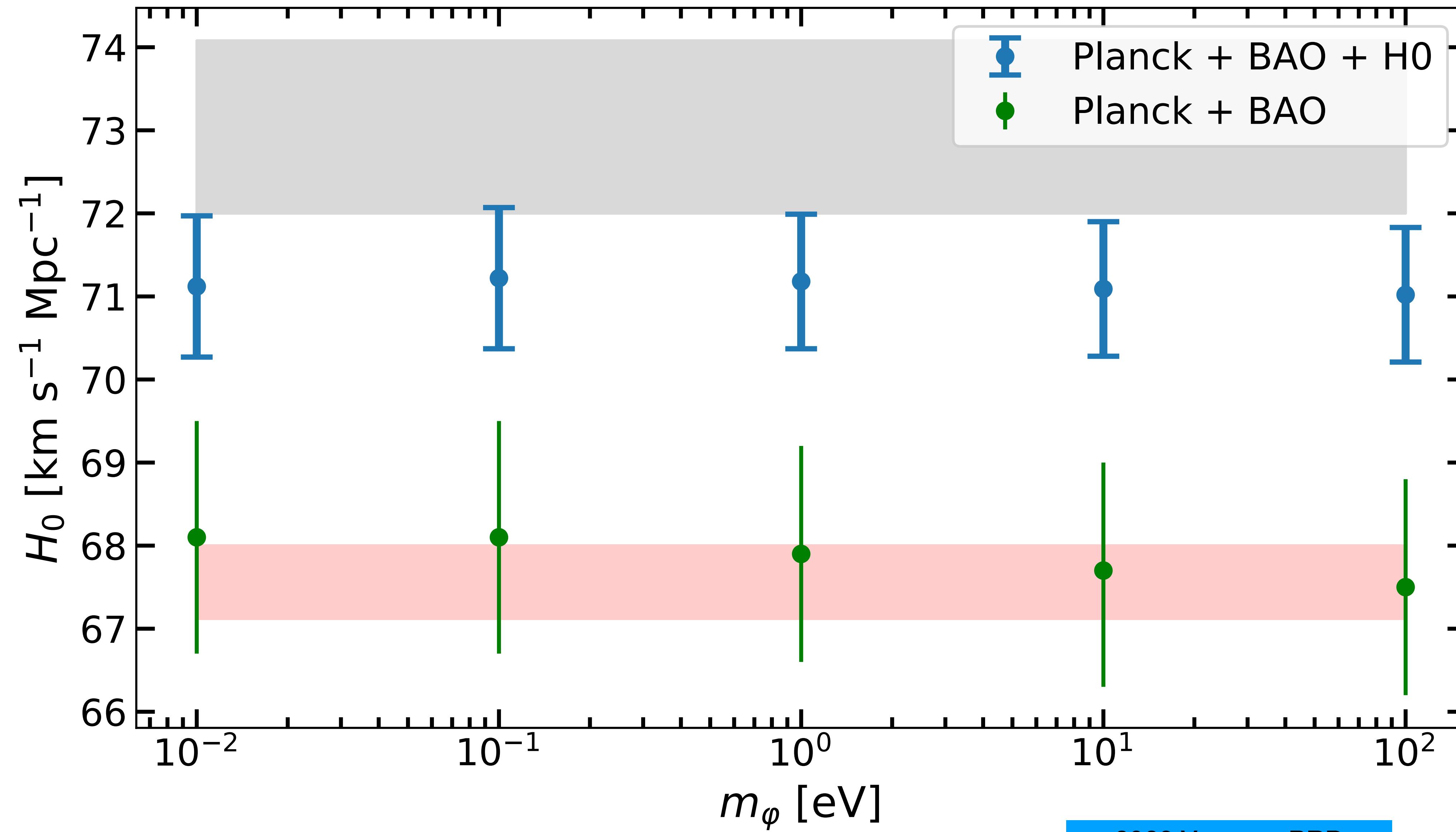


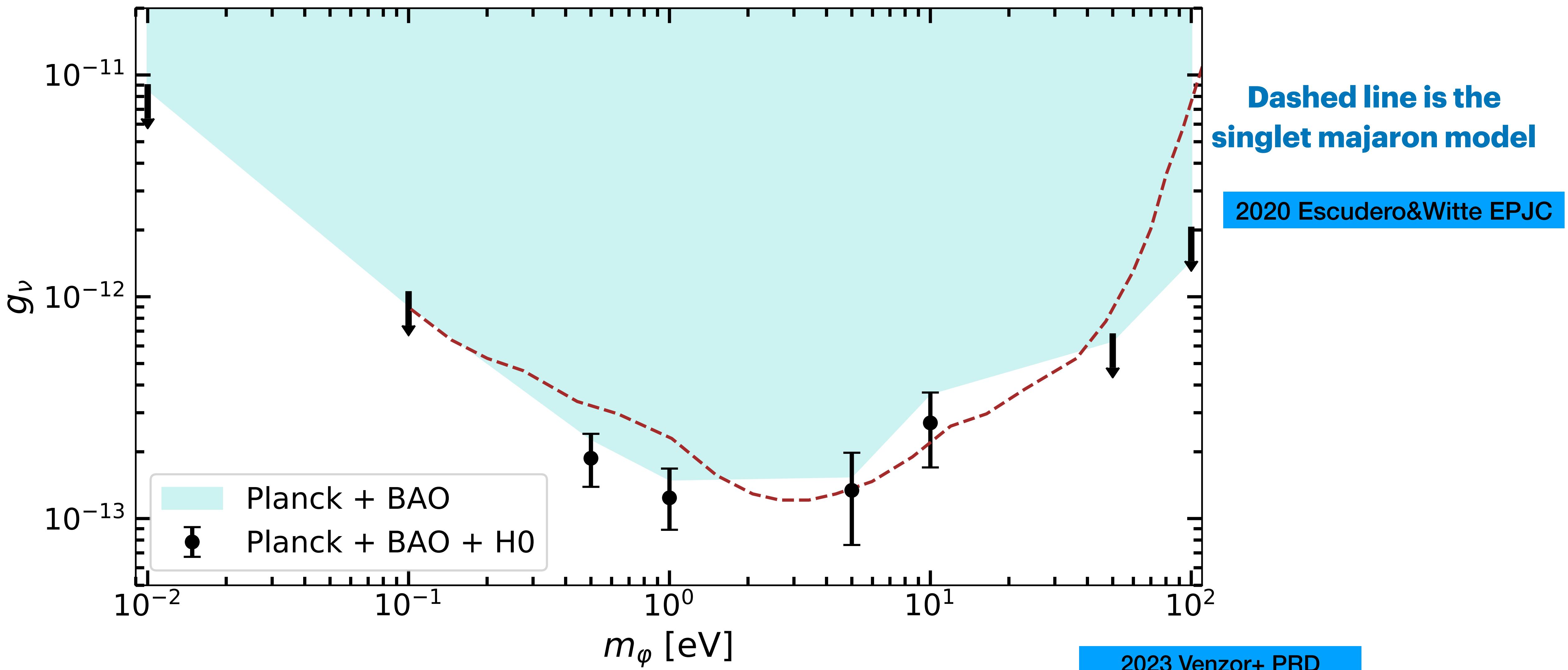
Larger masses enter to the horizon first



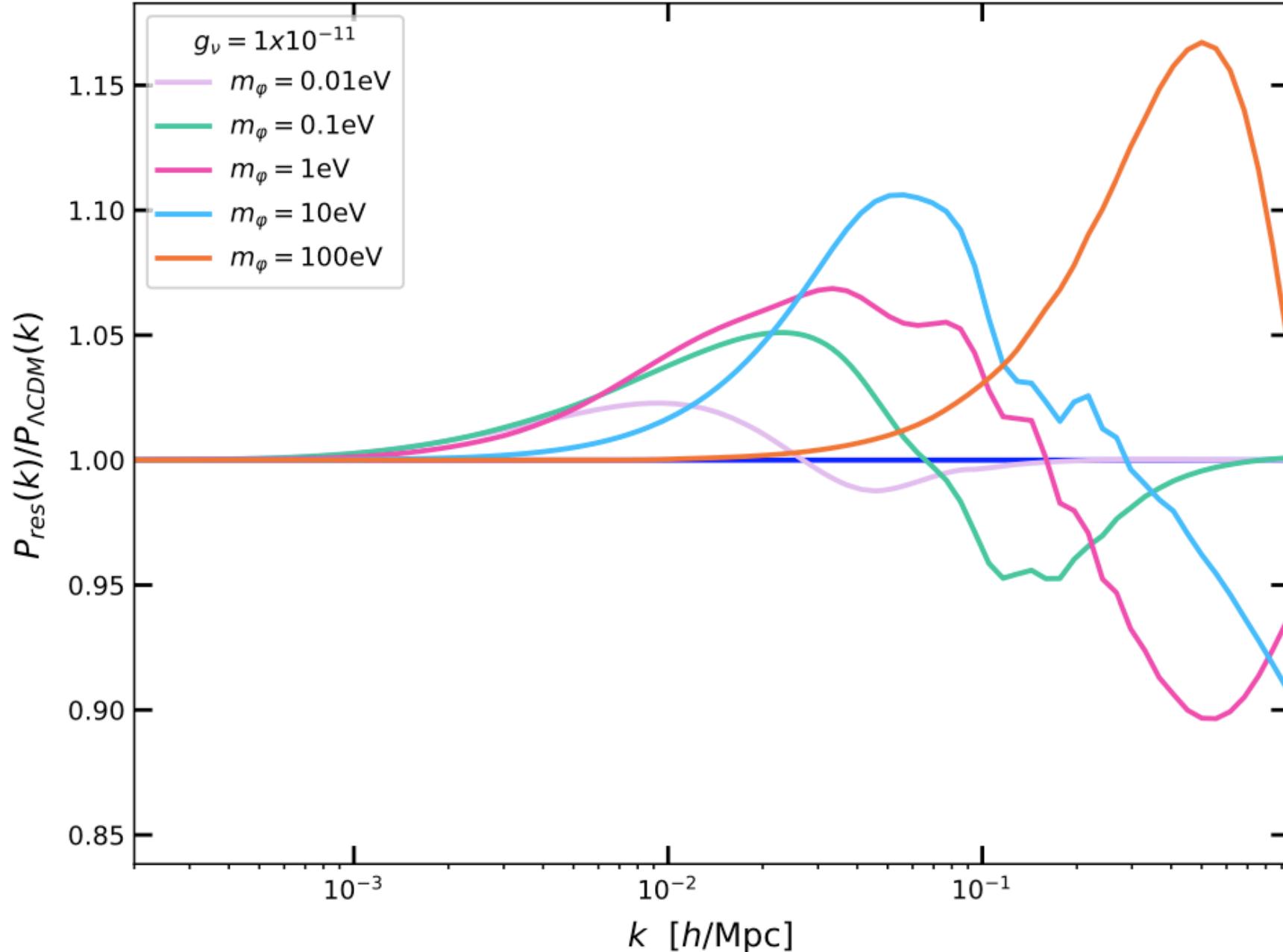
Example: mass=1eV

2023 Venzor+ PRD



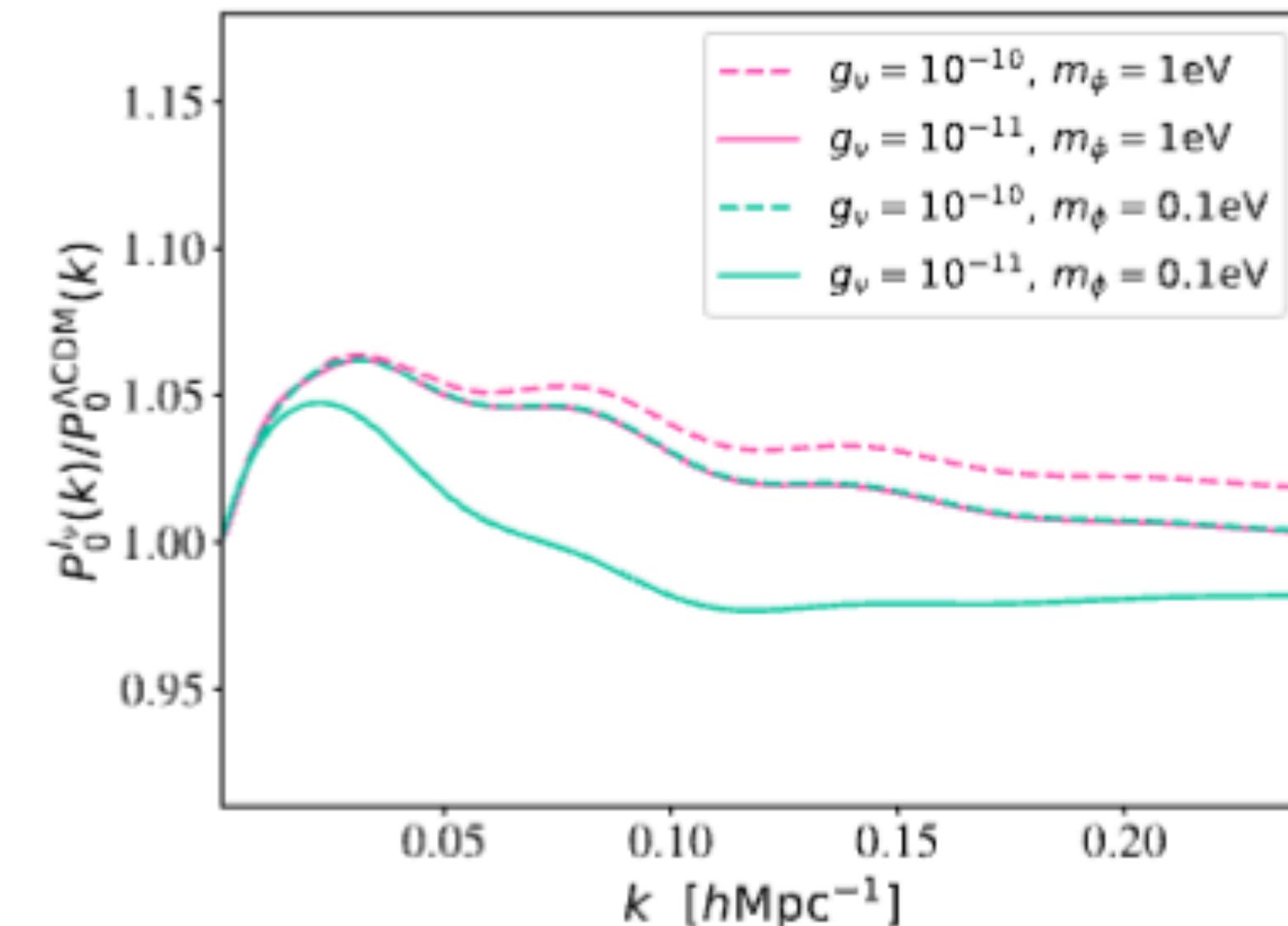


Full Shape Galaxy spectrum



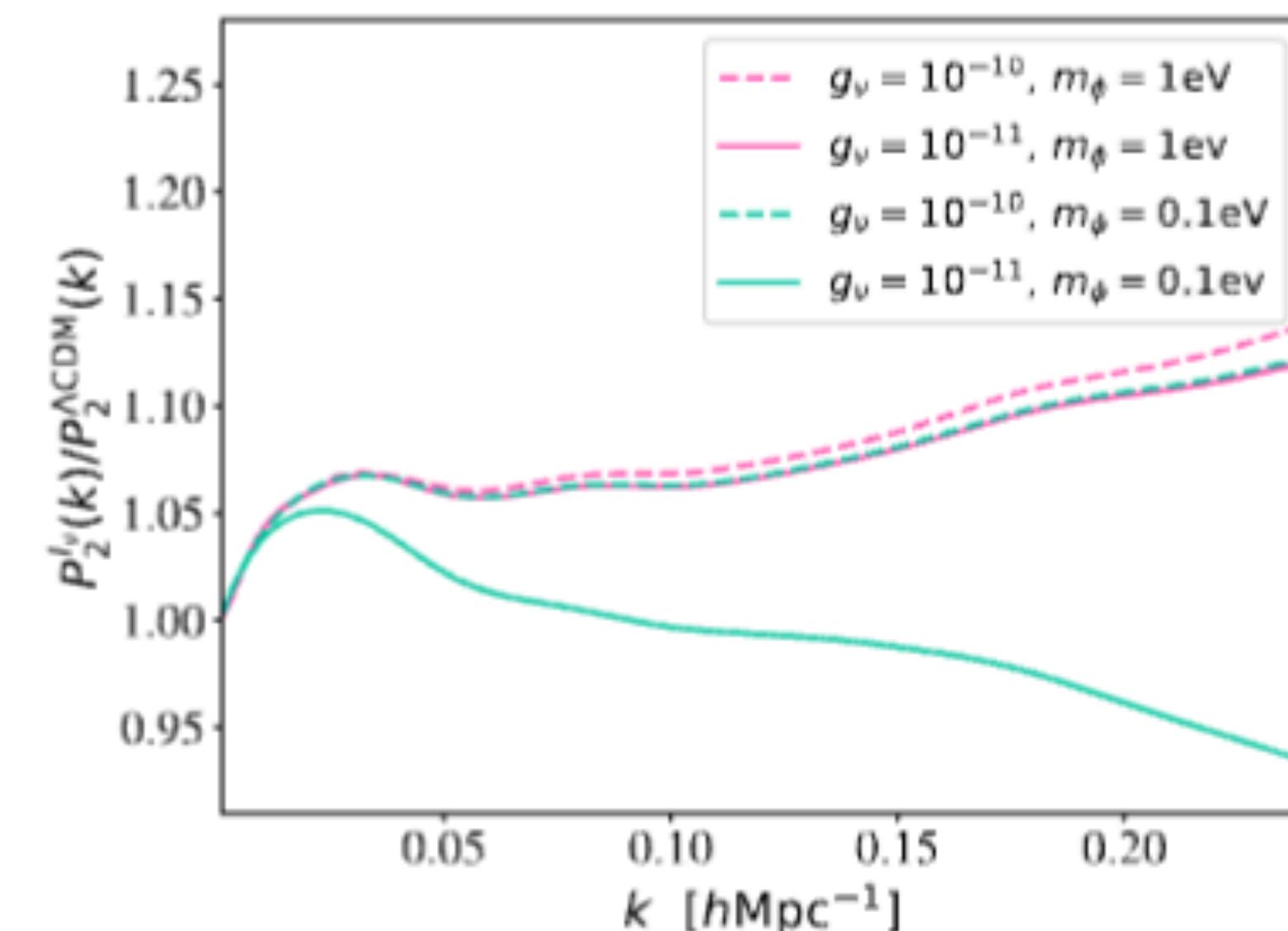
2023 Venzor+ PRD

Larger masses enter to the horizon first



Non linear effects on the Matter Power Spectrum

Recently developed
Zaldarriaga, Simonovic,
Ivanov, Cox,
Scoccimarro



Preliminary

Perspectives

- All neutrino self-interaction cases require extra radiation
- In the heavy mediator case the bimodality of the coupling posterior has not disappeared.
- For the light and resonant cases we need to include non-elastic processes into the analysis.

iThanks!