

Using the IceCube data to constraint the $(3+2)$ neutrino scenario

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In this letter we are using the IceCube experiment to test, phenomenologically, the (3+2) sterile neutrino scenario. As far is known, the presence of sterile states with mass splitting $\Delta m^2 \sim 1\text{eV}^2$ distorts the angular and energy distributions of reconstructed muon events, observed by IceCube, through parametric and MSW resonances.

Since the distortions introduced by the sterile neutrino (3+2) scenario appear around of few TeV, where the Ice-Cube detection is highly optimized, here we have a great opportunity to prove scenarios with sterile states. Supporting by previous works, some considerations have be done in order to guarantee a conservative and robust bounds.

Then, by using one year data colleted by IceCube between 2011-2012 we constrain the (3+2) parameter space scenario

Over years, many experiments had proved that neutrinos oscillate

$$P_{\nu_{\alpha} \rightarrow \nu_{\beta}}(L, E) = |\langle \nu_{\beta} | \nu_{\alpha}(t) \rangle|^2 = \sum_{k,j} U_{\alpha k}^* U_{\beta k} U_{\alpha j} U_{\beta j}^* \exp \left(-i \frac{\Delta m_{kj}^2 L}{2E} \right)$$

were $|\nu_{\alpha}\rangle = \sum_k U_{\alpha k}^* |\nu_k\rangle$ where a flavor neutrino state is a linear combination with mass eigenstates

In this standard framework there are 3 active neutrinos with

3 masses $\Delta m_{21}^2, \Delta m_{31}^2, \Delta m_{32}^2$

3 mixing angles $\theta_{12}, \theta_{13}, \theta_{23}$

1 CP violating phase δ_{CP}

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with $|\nu_\alpha(t)\rangle = \sum U_{\alpha k}^*$

With the U's been PMNS matrix

$$\begin{bmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{bmatrix}$$

is a linear states

In this scenario

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$$\begin{bmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta_{CP}} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta_{CP}} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta_{CP}} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta_{CP}} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta_{CP}} & c_{23}c_{13} \end{bmatrix}$$

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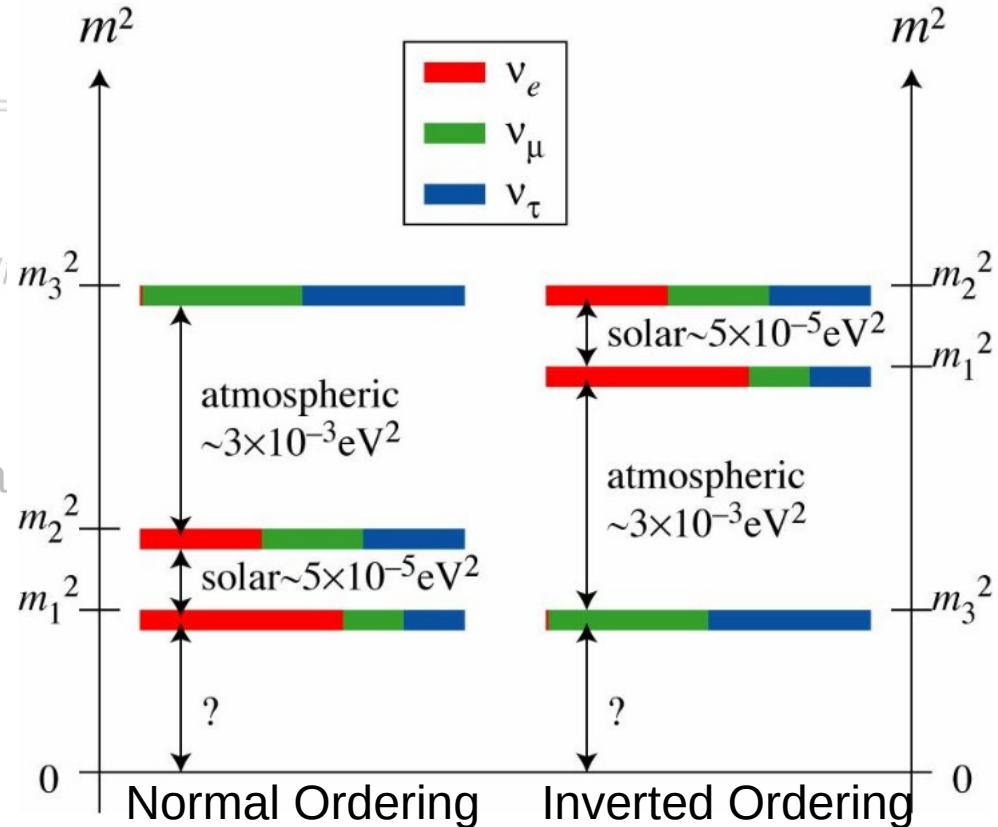
$$\text{with } |\nu_\alpha(t)\rangle = \sum_k U_{\alpha k}^* e^{-iE_k t} |\nu_k\rangle$$

In this standard framework there are 3

3 masses Δm_{21}^2 , Δm_{31}^2 , Δm_{32}^2

3 mixing angles θ_{12} , θ_{13} , θ_{23}

1 CP violating phase δ_{CP}



The three active neutrino framework can (almost?) fit, by global analysis, the oscillation parameters, but STILL, there are some anomalies that have been detected in:

The LSND anomaly

The Reactor neutrino anomaly

The Gallium anomaly

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The LSND anomaly [Phys. Rev. D 64 \(2001\) 112007 A. Aguilar et al. \(LSND Collaboration\)](#)

Research for $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$ which results suggest an oscillation with $\Delta m^2 > 0.4 \text{eV}^2$

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The Reactor neutrino anomaly PHYSICAL REVIEW C 83, 054615 (2011) Th. A. Mueller et al

After evaluating/compare the antineutrino spectral they found that have to be corrected by ~3% in normalization. This could mean $\bar{\nu}_e \rightarrow \nu_s$ if $\Delta m^2 \sim \Delta m_{LSND}^2$

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The Gallium anomaly Phys. Rev. C 80 (2009) 015807 SAGE Collaboration and Phys. Lett. B 685 (2010) 47 Kaether et al

Originating from the observed deficit in calibration tests of solar neutrino detectors, SAGE and GALLEX/GNO, by electron-capture source

The three active neutrino framework can (almost?) fit, by global analysis, the oscillation parameters, but STILL, there are some anomalies that have been detected in:

Experiment	Type	Channel	Significance
LSND	DAR	$\bar{\nu}_\mu \rightarrow \bar{\nu}_e$ CC	3.8σ
MiniBooNE	SBL accelerator	$\nu_\mu \rightarrow \nu_e$ CC	3.4σ
MiniBooNE	SBL accelerator	$\bar{\nu}_\mu \rightarrow \bar{\nu}_e$ CC	2.8σ
Gallex/SAGE	Radioactive	ν_e dissap	2.8σ
Reactors	Beta-decay	$\bar{\nu}_e$ dissap	3.0σ

K. N. Abazajian et al. "Light Sterile Neutrinos: A Whitepaper" arxiv:1204.5379

Each of them, individually, are not significant enough to confirm something, but They could show a directions to something else....

...this could be, as the most common interpretation as a glimpse of "sterile" neutrinos

If sterile neutrino exist, with mass-square differences $\sim \mathcal{O}(\text{eV}^2)$ then Earth matter induces a very interesting features

Parametric Resonance

Phys. Lett. B 440 (1998) 319 and
 Nucl. Phys. B 524 (1998) 505
 Liu, Smirnov, Mikheyev



Change the energy and zenith distributions of Atmospheric netrutrinos $\gtrsim 100 \text{ GeV}$

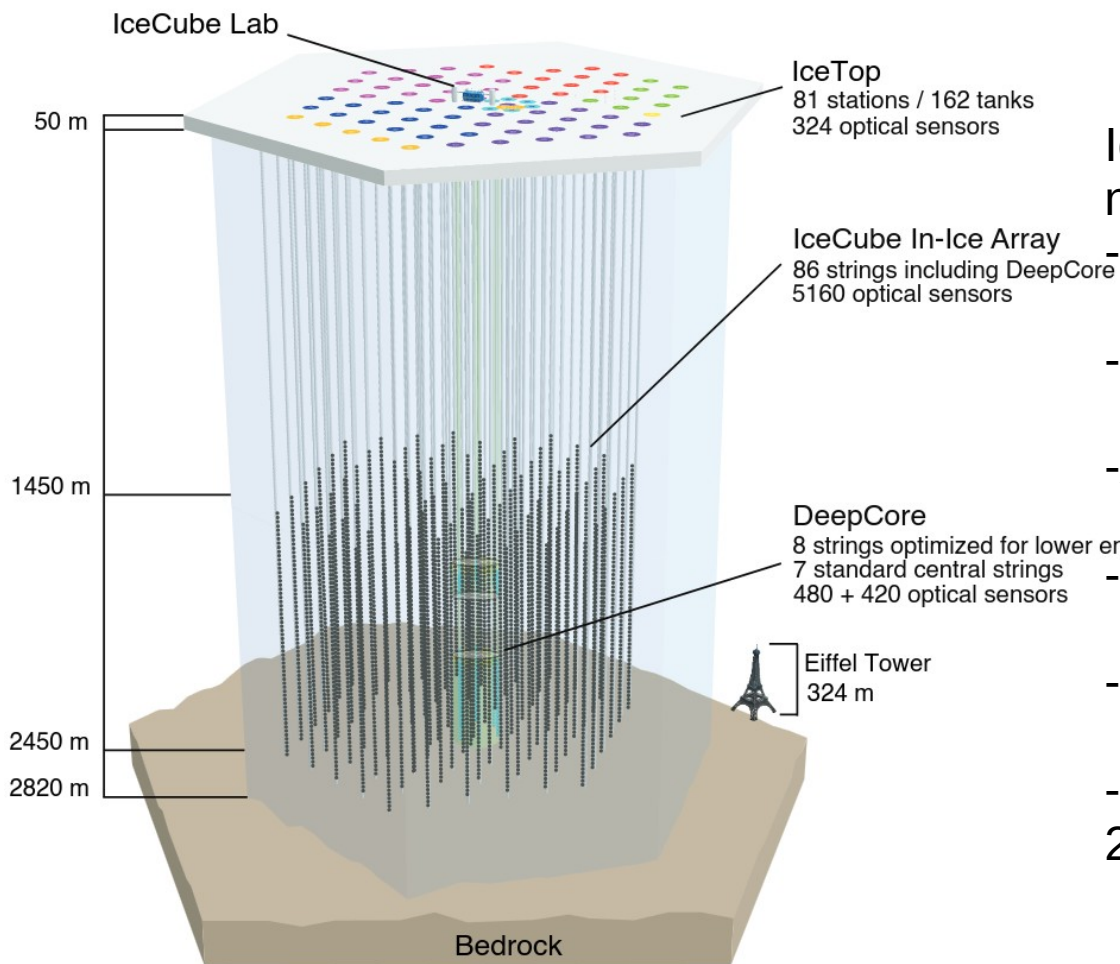
MSW Resonance

Phys. Lett. B 562 (2003) 279 and
 Hep-ph/0102166
 Nunokawa, Peres, Funchal, Yasuda

THEN



IceCube is sensitive to these scanarios



IceCube is gigaTon ice-Cherenkov neutrino detector:

-at South Pole

-can observe TeV neutrinos

-Atmospheric and astrophysical

-Energies above 100Gev and ~PeV

-5160 DOM in 86 strings

-over a cubic kilometer 1450m to 2450m

The evolution sterile neutrino equation is given by:

$$i \frac{d}{dr} \nu_f = \left[U \frac{M^2}{2E_\nu} U^\dagger + A \right] \nu_f$$

Where U is:

$$\begin{array}{c}
 \nu_e \\
 \nu_\mu \\
 \nu_\tau
 \end{array}
 \begin{array}{c}
 \nu_1 \quad \nu_2 \quad \nu_3 \\
 \left(\begin{array}{ccc}
 \boxed{} & \boxed{} & \boxed{} \\
 \boxed{} & \boxed{} & \boxed{} \\
 \boxed{} & \boxed{} & \boxed{}
 \end{array} \right)
 \end{array}$$

For the three active neutrinos scenario

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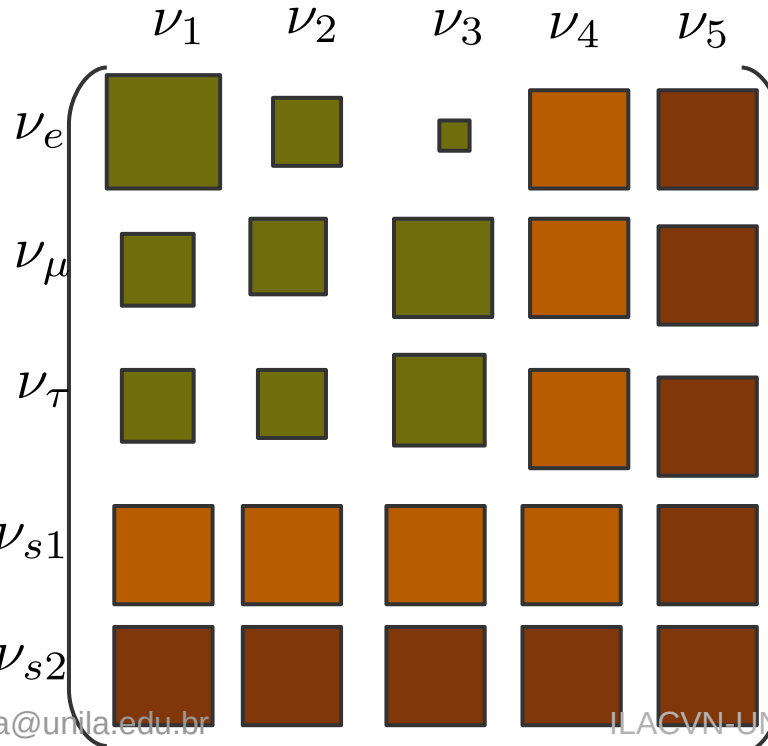
Where U is:

	ν_1	ν_2	ν_3	ν_4
ν_e				
ν_μ				
ν_τ				
ν_s				

For the 3+1 sterile neutrino scenario

The evolution sterile neutrino equation is given by:

$$i \frac{d}{dr} \nu_f = \left[U \frac{M^2}{2E_\nu} U^\dagger + A \right] \nu_f$$



For the 3+2 sterile neutrino scenario

The evolution sterile neutrino equation is given by:

$$i \frac{d}{dr} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \\ \nu_4 \\ \nu_5 \end{pmatrix} = \left[U \frac{M^2}{2E_\nu} U^\dagger + A \right] \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \\ \nu_4 \\ \nu_5 \end{pmatrix}$$

Where U is

$$U = R^{45}(\theta_{45}) R_{\delta_{35}}^{35}(\theta_{35}) R_{\delta_{25}}^{25}(\theta_{25}) R_{\delta_{15}}^{15}(\theta_{15}) R^{34}(\theta_{34}) R_{\delta_{24}}^{24}(\theta_{24}) R_{\delta_{14}}^{14}(\theta_{14}) R^{23}(\theta_{23}) R_{\delta_{13}}^{13}(\theta_{13}) R^{12}(\theta_{12})$$

R^{ij} Is the rotation matrix in the i-j plane with the angle θ_{ij}

$R_{\delta_{ij}}^{ij}$ Is the rotation matrix which contain a CP-Violating angle δ_{ij}

The evolution sterile neutrino equation is given by:
For this 3+2 sterile neutrino scenario them:

$$M^2 = \begin{pmatrix} 0 & 0 & 0 & 0 & 0 \\ 0 & \Delta m_{21}^2 & 0 & 0 & 0 \\ 0 & 0 & \Delta m_{31}^2 & 0 & 0 \\ 0 & 0 & 0 & \Delta m_{41}^2 & 0 \\ 0 & 0 & 0 & 0 & \Delta m_{51}^2 \end{pmatrix} \quad A = \begin{pmatrix} A_{CC} & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & A_{NC} & 0 \\ 0 & 0 & 0 & 0 & A_{NC} \end{pmatrix}$$

with $\Delta m_{ij}^2 = m_j^2 - m_i^2$

$$A_{CC} = 2\sqrt{2}G_F E_\nu N_e(r)$$

$$A_{NC} = \sqrt{2}G_F E_\nu N_n(r)$$

In this scenario there are **20** parameters, 10 mixing angles, 6 CP-Violating phases and 4 mass-square differences....but we dont need to use all them..

Parameter	Reason for elimination
Standard oscillations	For $E_\nu \geq 100\text{GeV}$ ν_e mixing strongly suppressed and L_{osc} for ν_μ and ν_τ is bigger than earth
θ_{45}	Can be rotated away by redefinition of the sterile state
$\theta_{14} \theta_{15} \delta_{14} \delta_{15}$	Do not impact survival probabilities ν_μ and $\bar{\nu}_\mu$
$\theta_{34} \theta_{35} \delta_{24} \delta_{25} \delta_{35}$	If fixed to zero allow conservative constraints

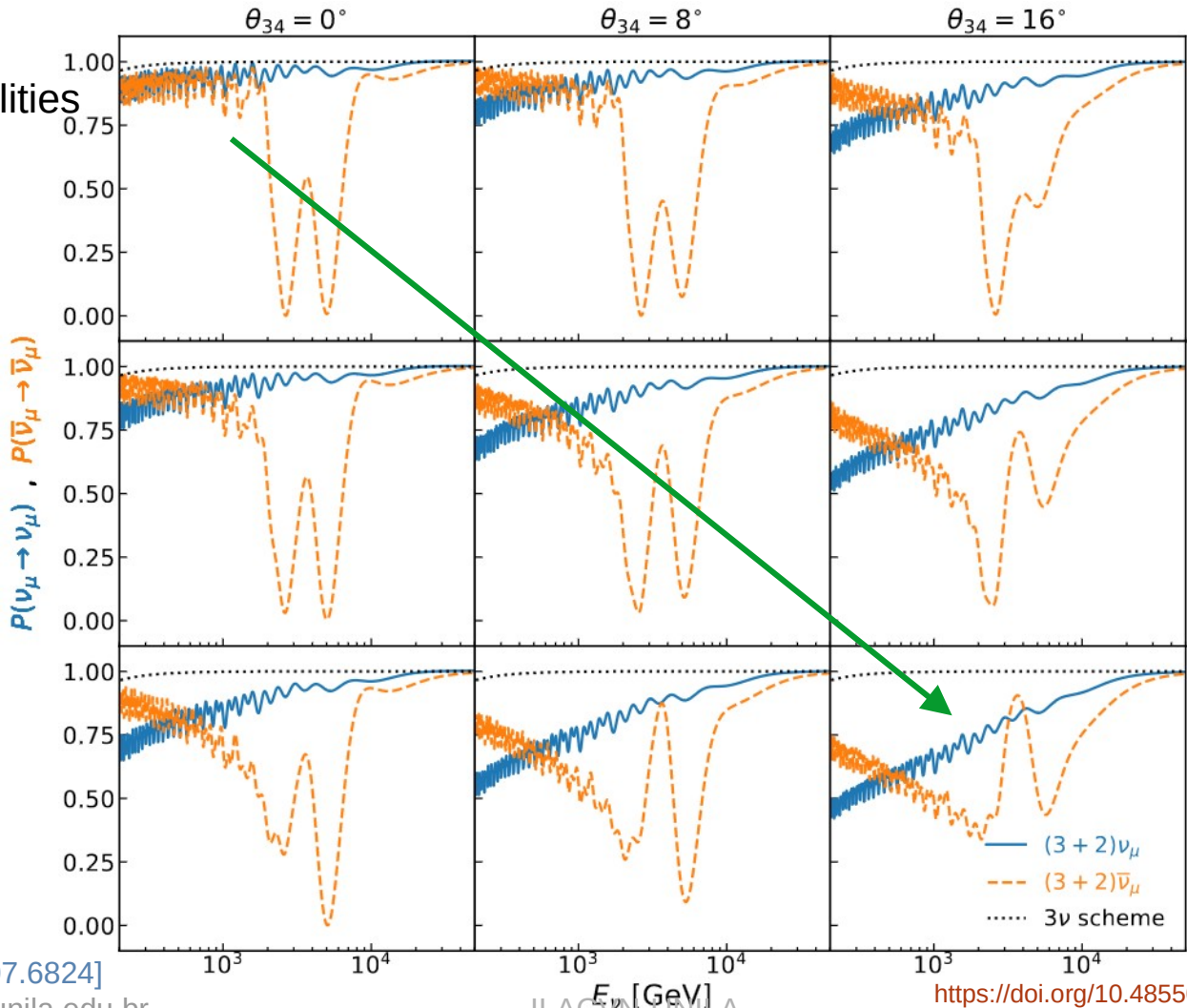
Survival probabilities

$$\cos \theta_z = -1$$

$$\theta_{24} = \theta_{25} = 8^\circ$$

$$\Delta m_{41}^2 = 1 \text{ eV}$$

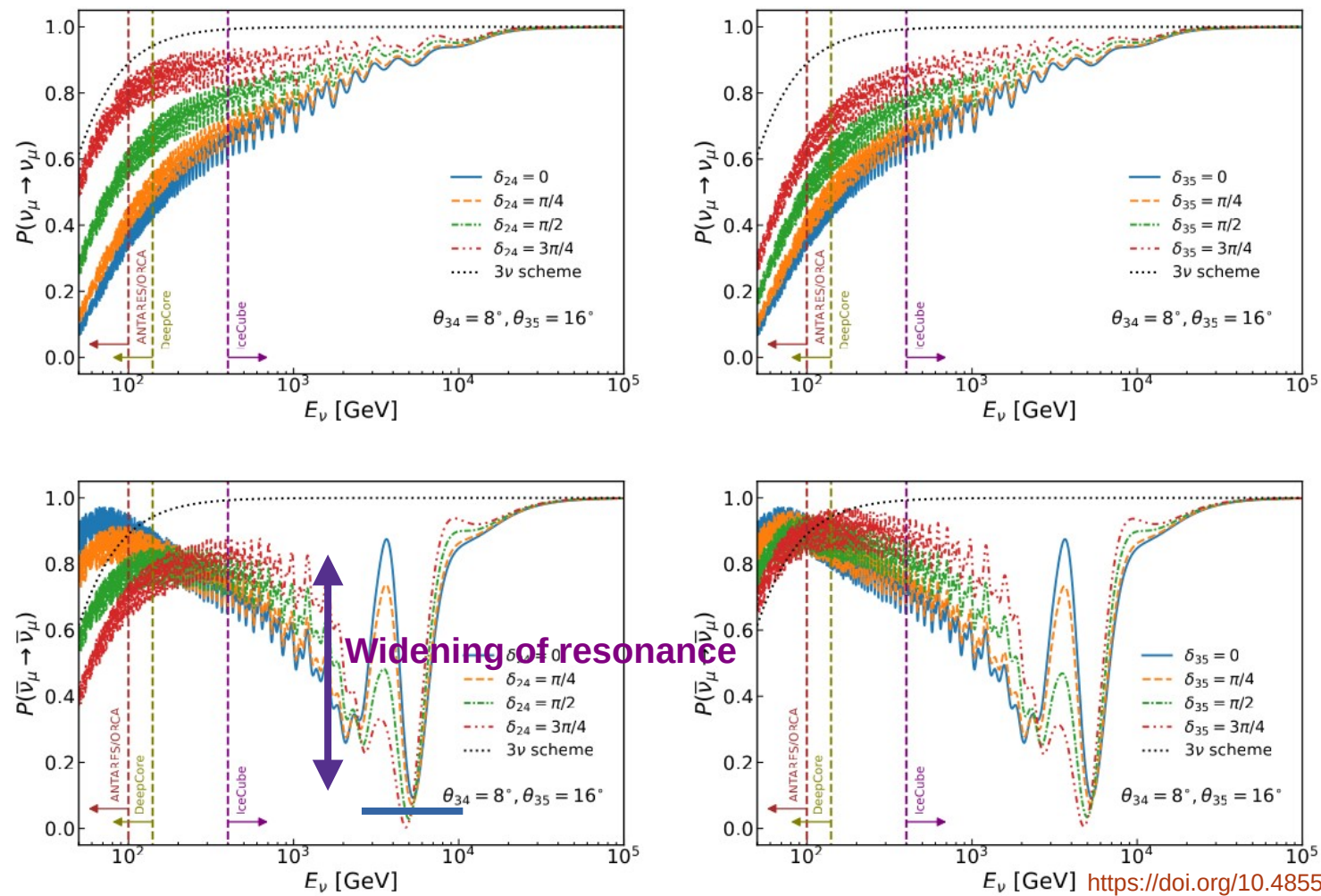
$$\Delta m_{51}^2 = 2 \text{ eV}$$



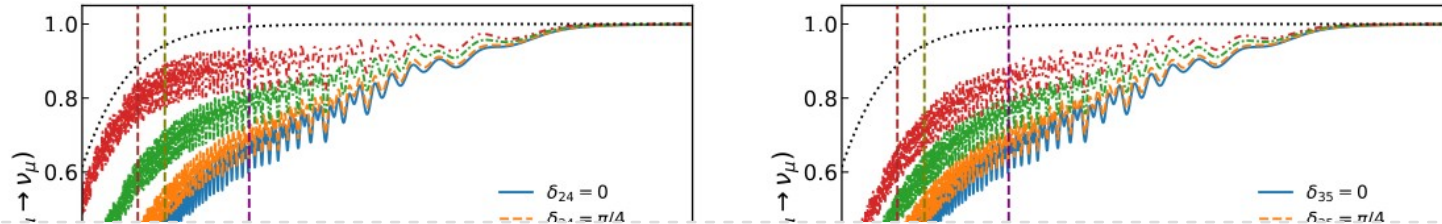
The orange curves become shallower

Is strengthening the limits on the (3+2) scenario

Survival probabilities



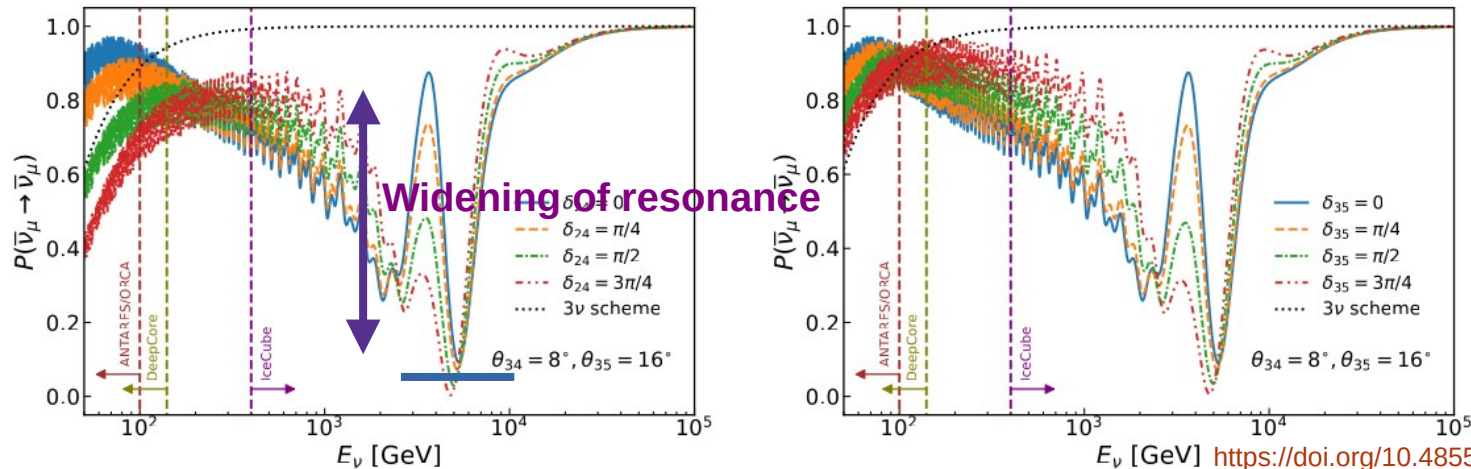
Survival probabilities



We can conclude that in order to probe the 3+2 neutrino scenario using the atmospheric neutrino data from IceCube is just necessarily to manage:

$$\theta_{24}, \theta_{25}, \Delta m_{41}^2, \Delta m_{51}^2$$

.....the rest go to zero...



Analysis We compute the survival probabilities of atmospheric ν_μ and $\bar{\nu}_\mu$ numerically, as function of E_ν and $\cos \theta_z$ and scanning over the parameter space of $(\theta_{24}, \theta_{25}, \Delta m_{41}^2, \Delta m_{51}^2)$.

$$\chi^2(\vec{\theta}) = -2 \ln \mathcal{L}(\vec{\theta}) =$$

$$\min_{\vec{\xi}, \{d\}} \left(2 \sum_{j,k} \left[N_{jk}^{\text{exp}}(\vec{\theta}, \vec{\xi}, d) - N_{jk}^{\text{obs}} + N_{jk}^{\text{obs}} \ln \frac{N_{jk}^{\text{obs}}}{N_{jk}^{\text{exp}}(\vec{\theta}, \vec{\xi}, d)} \right] + \sum_{\eta} \frac{(\xi_{\eta} - \Xi_{\eta})^2}{\sigma_{\eta}^2} \right)$$

we use a public data of IceCube collected over a period of 343.7 days between 2011-2012.

Analysis

$$\chi^2(\vec{\theta}) = -2 \ln \mathcal{L}(\vec{\theta}) =$$

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$$N_{jk}^{\text{exp}} = \sum_i \left[\eta_{ijk} \phi_{ik}^{\nu, \text{atm}} \langle P(\nu_{\mu} \rightarrow \nu_{\mu}) \rangle_{ik} + \bar{\eta}_{ijk} \phi_{ik}^{\bar{\nu}, \text{atm}} \langle P(\bar{\nu}_{\mu} \rightarrow \bar{\nu}_{\mu}) \rangle_{ik} \right]$$

Response array tensor

$$ijk = E_{\nu}, E_{\mu}, \cos \theta_z$$

- incoming neutrino energy
- muon energy
- direction E_{ν}, E_{μ}
- optical efficiency

averaged survival probability

Analysis

$$\chi^2(\vec{\theta}) = -2 \ln \mathcal{L}(\vec{\theta}) =$$

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the average atmospheric ν_{μ} flux in the (i, k)th bin originating from the decay of pions and kaons

Analisis

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the average atmospheric ν_{μ} flux in the (i, k)th bin originating from the decay of pions and kaons

$$\phi_{ik}^{\nu, \text{atm}} = N_0 \left[\phi_{ik}^{\nu, \pi} + (1.1 - R_{\pi/K}) \phi_{ik}^{\nu, K} \right] \left(\frac{E_{\nu, i}}{E_0} \right)^{-\gamma} \quad \phi_{ik}^{\bar{\nu}, \text{atm}} = N' N_0 \left[\phi_{ik}^{\bar{\nu}, \pi} + (1.1 - R_{\pi/K}) \phi_{ik}^{\bar{\nu}, K} \right] \left(\frac{E_{\nu, i}}{E_0} \right)^{-\gamma}$$

Provided by IceCube

Results

After performing the analysis with the $x(\theta)$ was found the allowed region of the 3+2 parameters space by IceCube.

The best fit point is

$$\Delta m_{41}^2 = \Delta m_{51}^2 = 16 \text{ eV}^2$$
$$\sin^2 2\theta_{24} = 0.3 \quad \sin^2 2\theta_{25} = 0.23$$

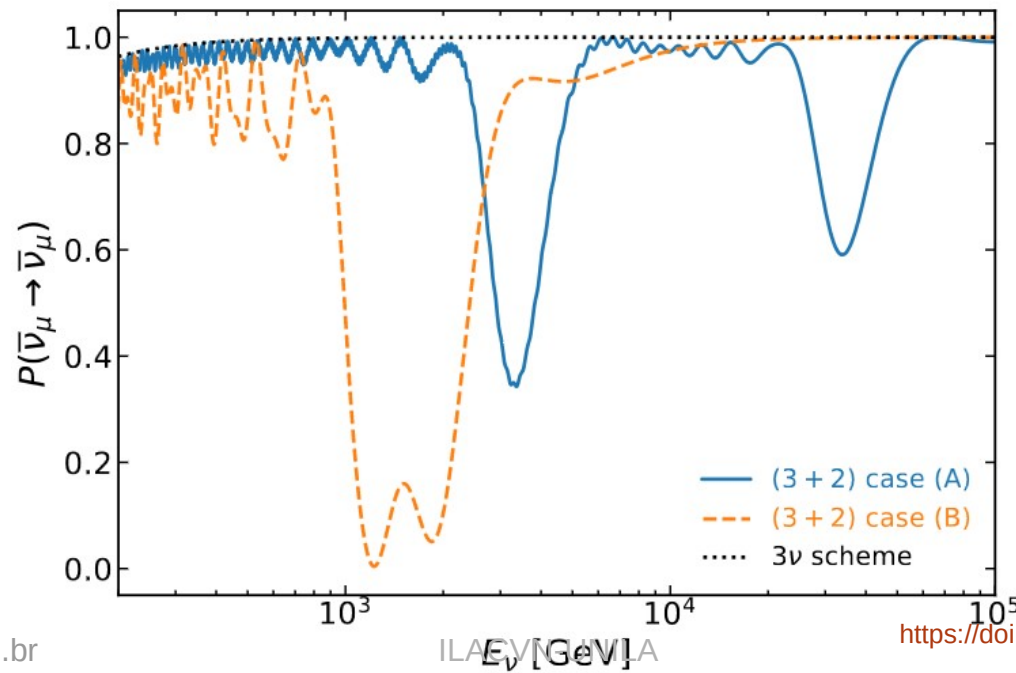
Which means that the IceCube data do not show preference to (3+2) scenario over the standard one

Results

	$\sin^2 2\theta_{24}$	$\sin^2 2\theta_{25}$	Δm_{41}^2 [eV ²]	Δm_{51}^2 [eV ²]	$\Delta\chi^2$
Case (A) from Ref. [35]	2.8×10^{-2}	1.5×10^{-2}	1.32	13.9	4.32
Case (B) from Ref. [36]	9.1×10^{-2}	6.8×10^{-2}	0.46	0.77	36.43

[35] A. Diaz, C.A. Argüelles, G.H. Collin, J.M. Conrad and M.H. Shaevitz - Phys. Rept. 884 (2020) 1

[36] D. Cianci, A. Furmanski, G. Karagiorgi and M. Ross-Lonergan - Phys. Rev. D 96 (2017) 055001

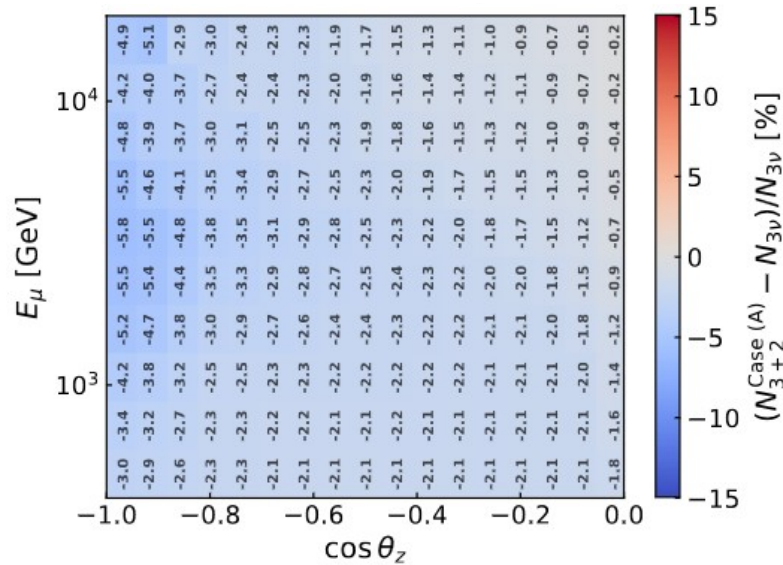


Results

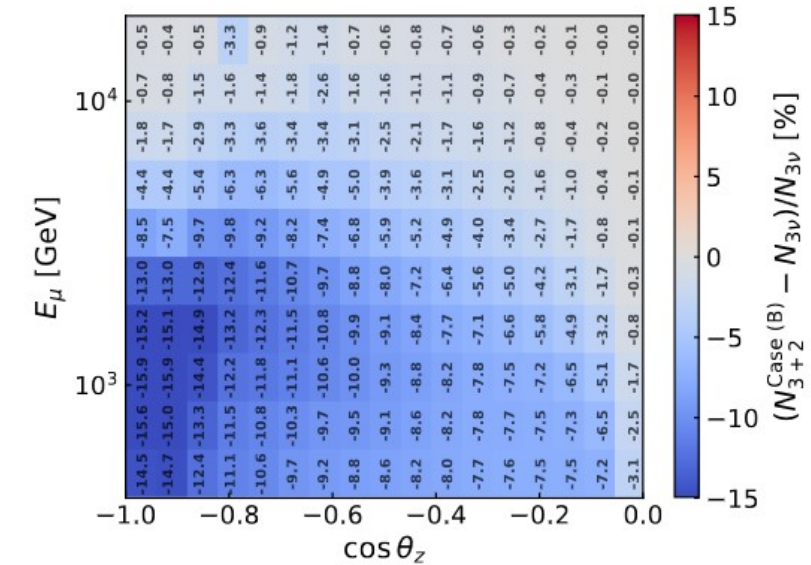
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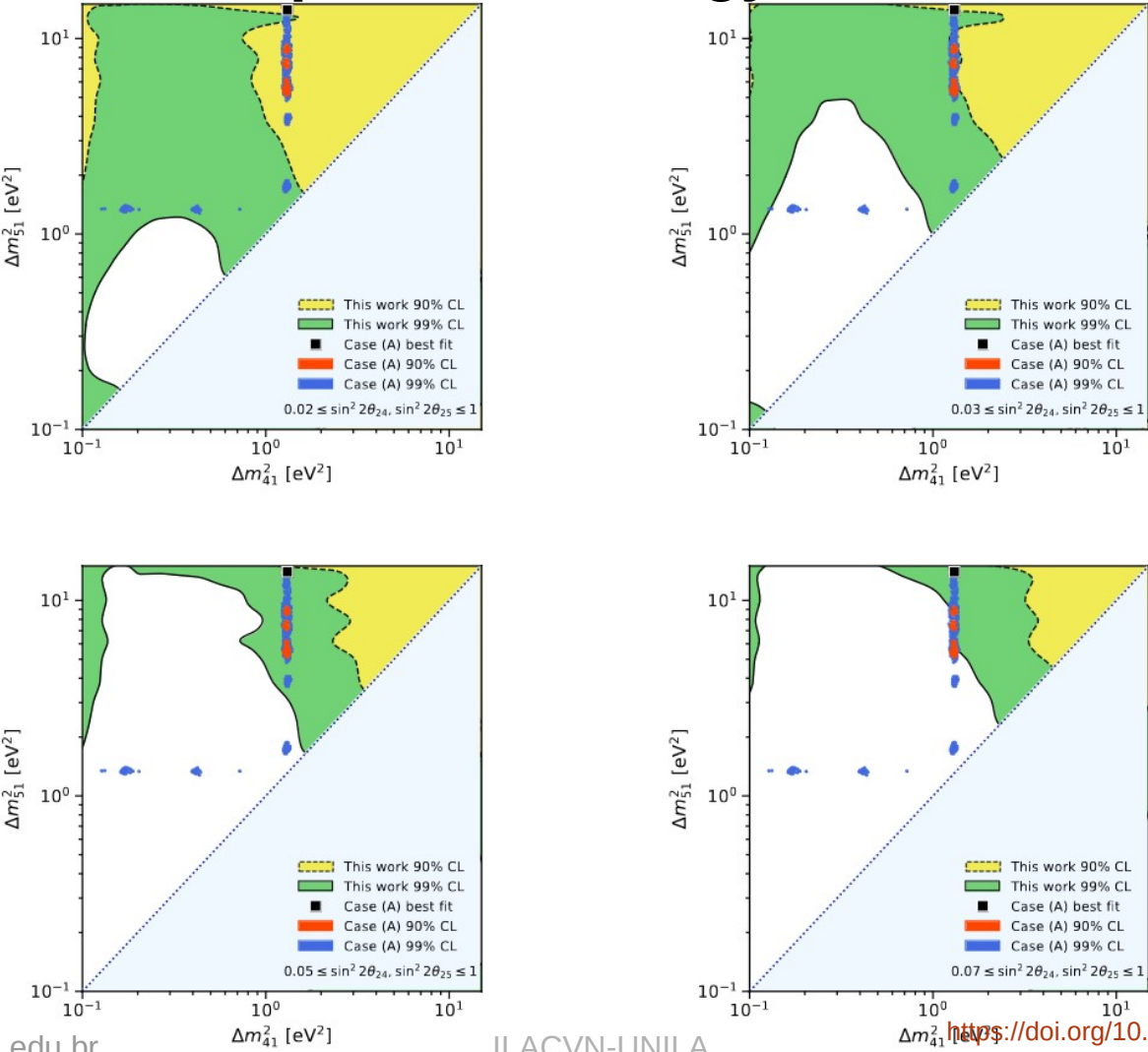


(a)

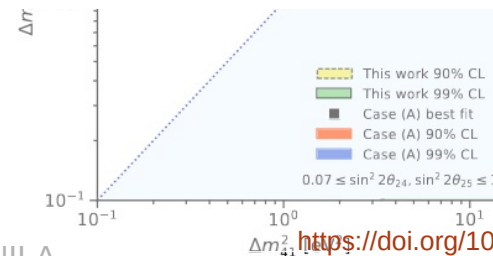
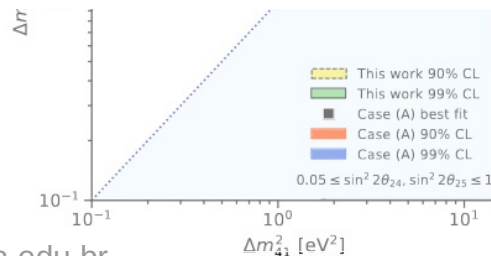
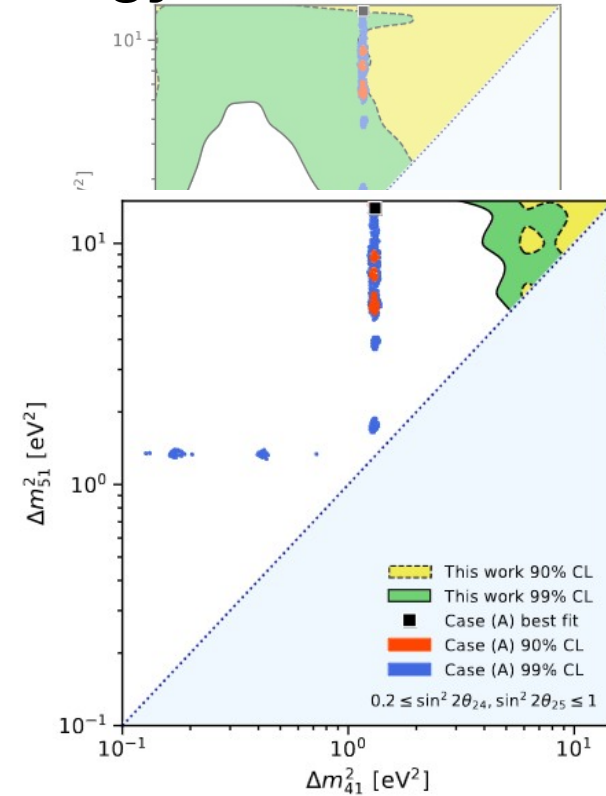
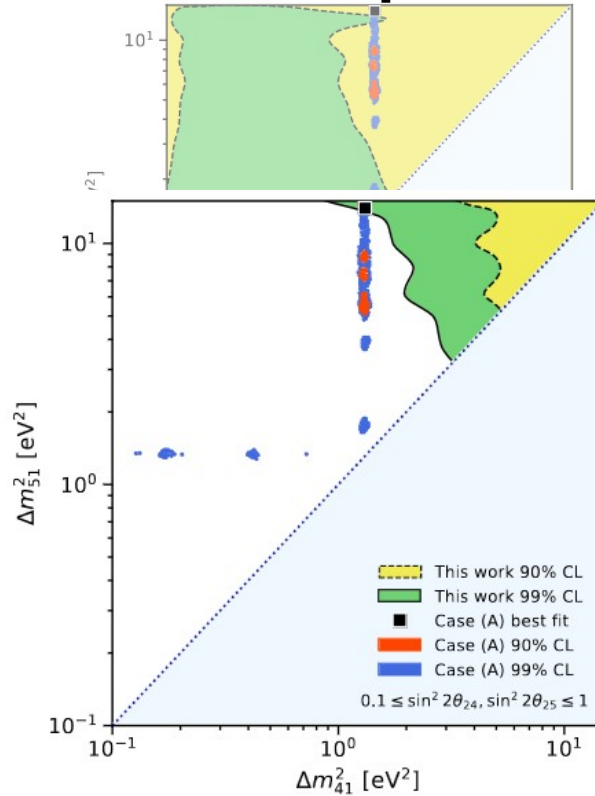


(b)

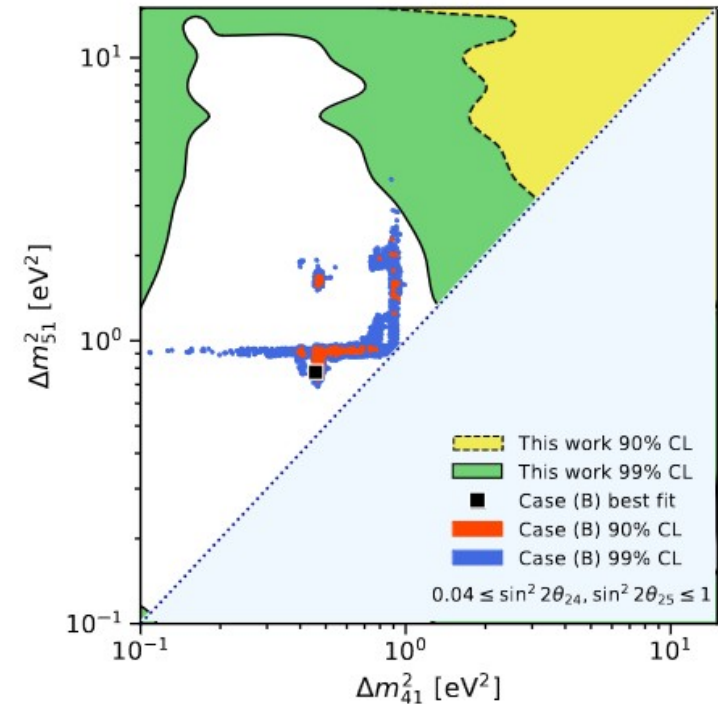
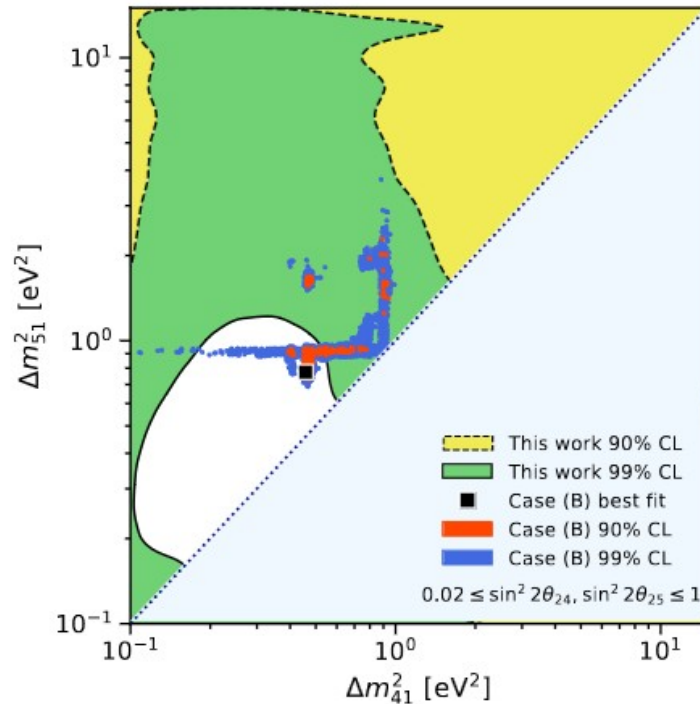
Results



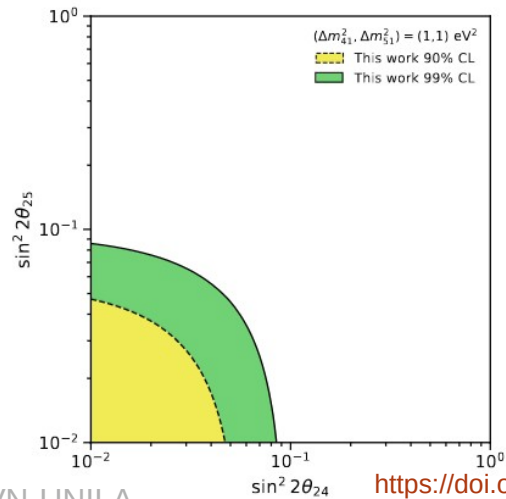
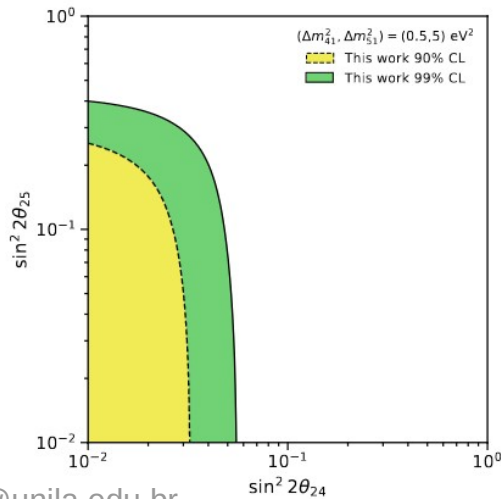
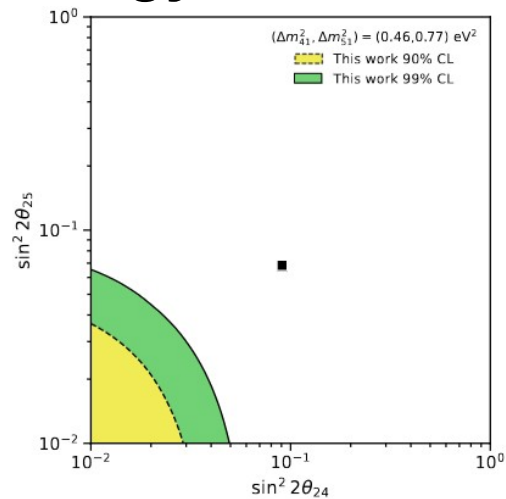
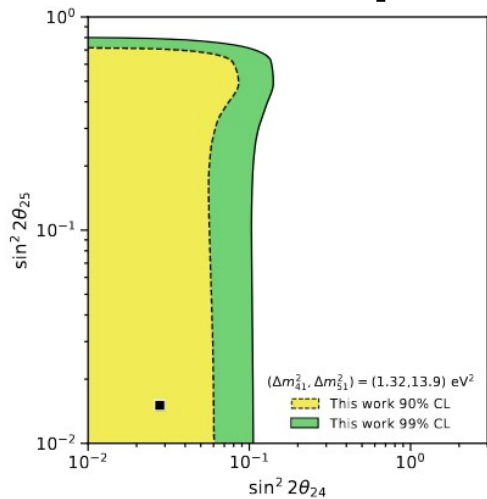
Results



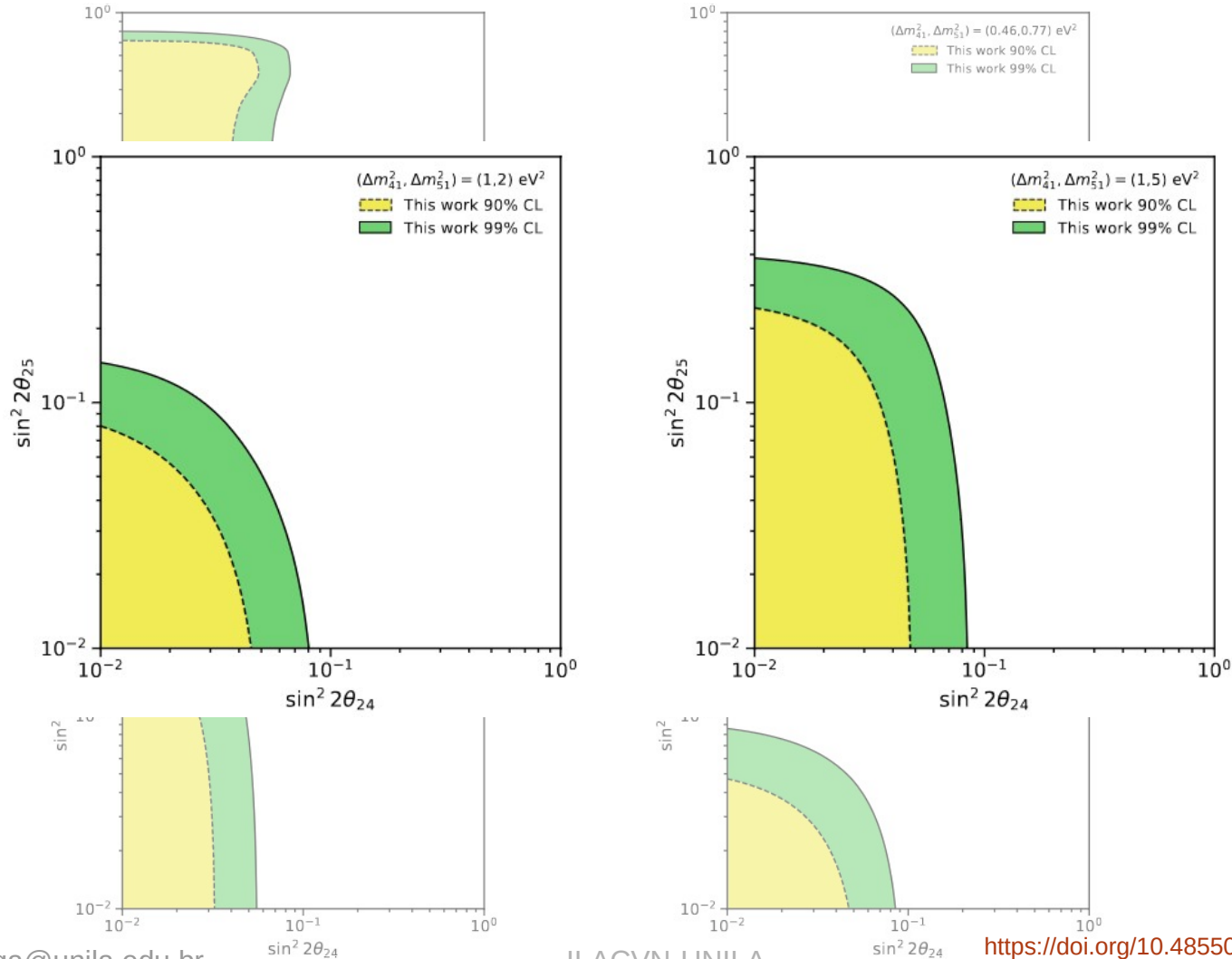
Results



Results



Results

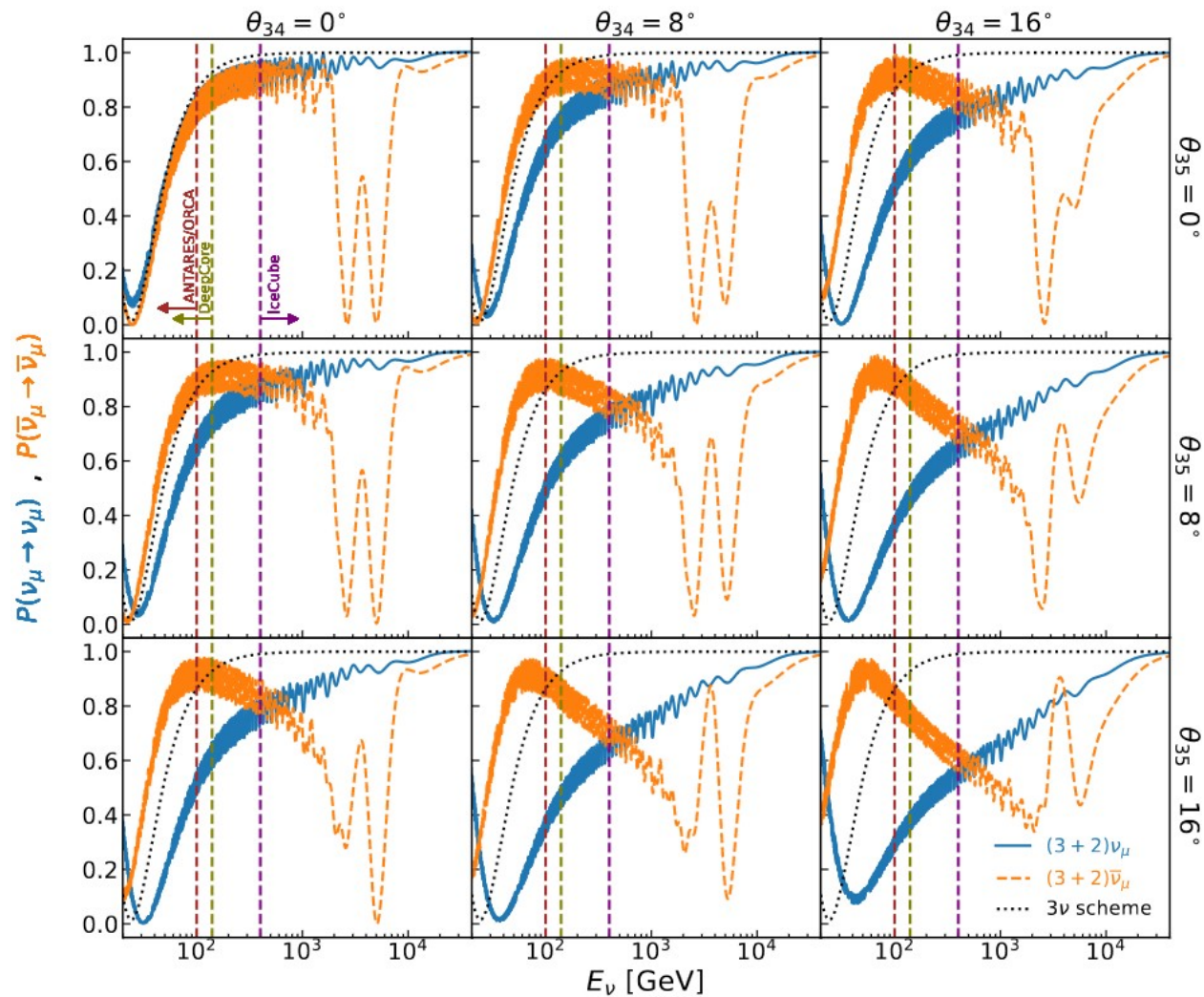


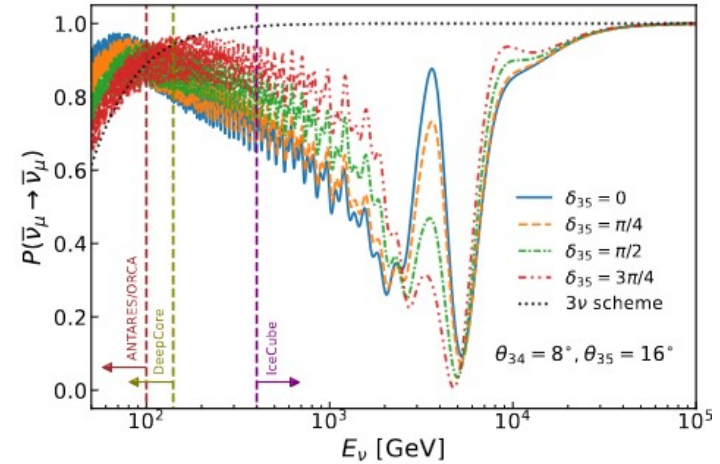
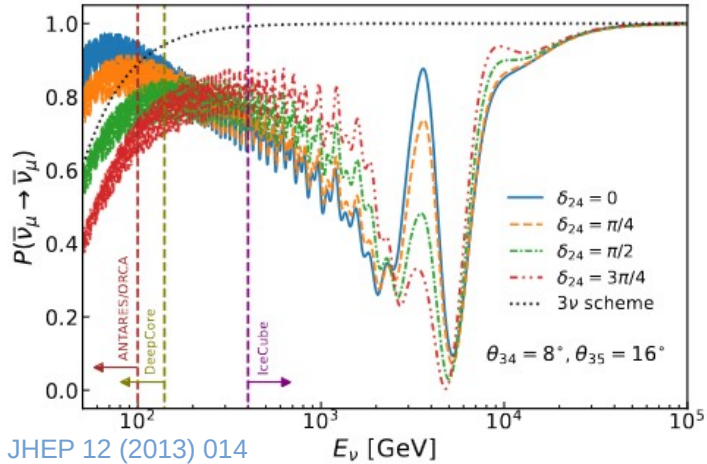
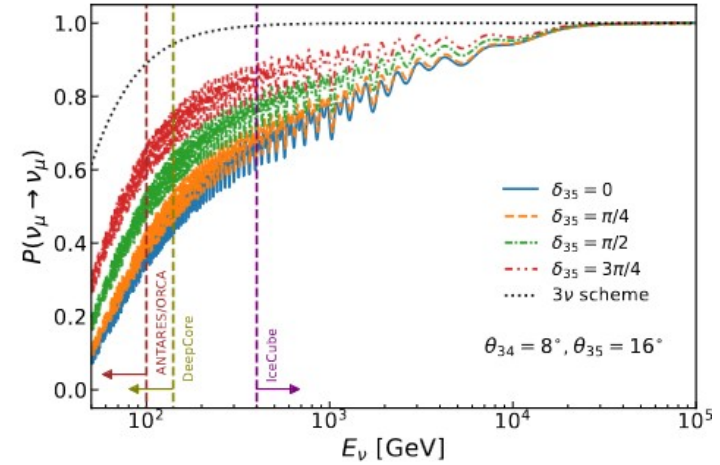
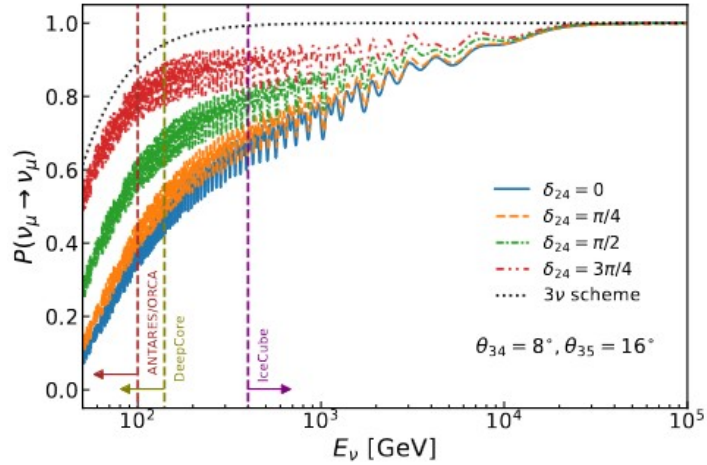
Conclusions

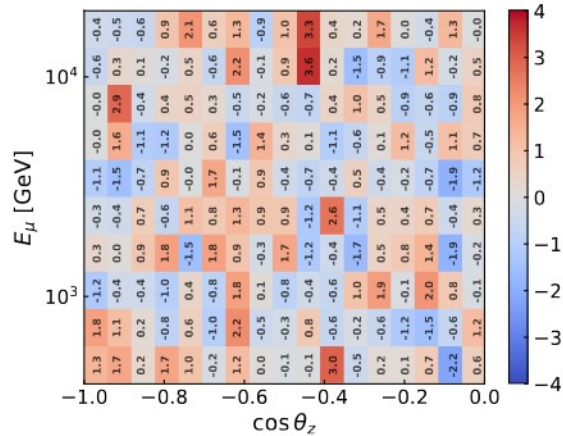
- Anomalies in the SBL, reactor and Solar experiments could indicate the presence of sterile(s) neutrinos,
- In principle the number of sterile neutrinos is free,
- To IceCube, more sterile states means more resonances that also means stronger constraints,
- Only 4 parameter are needed to get a conservative constraint,
- We performed a binned Poisson likelihood analysis with various systematics and statistical uncertainties,
- Best fit $\Delta m_{41}^2 = \Delta m_{51}^2 = 16eV^2, \sin^2(2\theta_{24}) = 0.3, \sin^2(2\theta_{25}) = 0.23$
- The bounds were compared with the allowed regions of the SBL (A) and (B),
- Case (A) is compatible with IceCube data; Case (B) seems in tension,

Thanks
Gracias
Obrigado
Mercy
Grazie
Danke
Arigato
Xièxiè

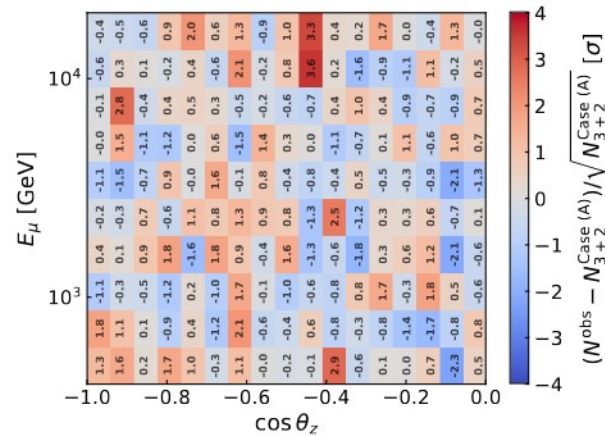
Backup



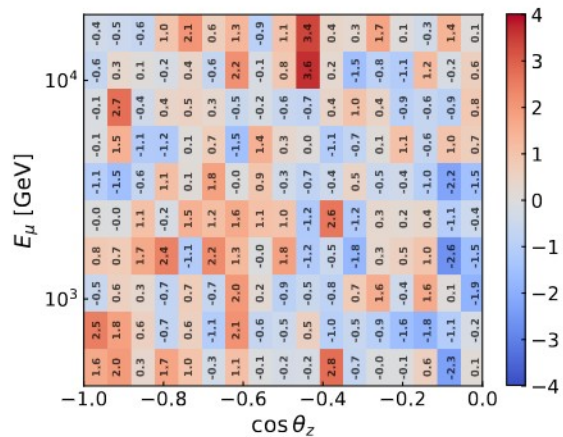




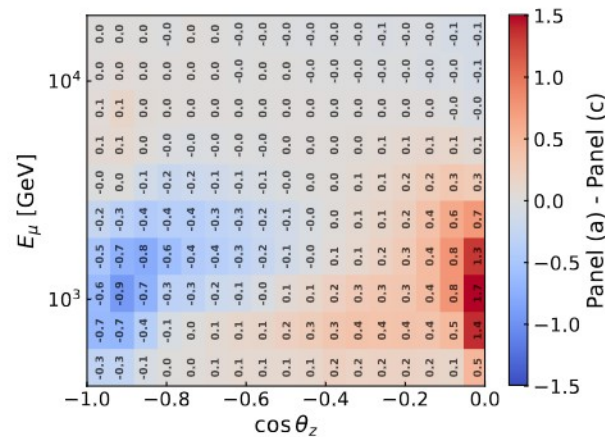
(a)



(b)



(c)



(d)

