

Investigating the soft and hard limits in transverse momentum spectra and exploring isotropic events in pp collisions

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Collective effects in proton-proton collisions?

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• Collective effects in proton-proton collisions has been a highly active area of research over the past decade: production of **hot QCD matter in small systems**?

	Physics Letters B 765 (2017) 193-220			Physics Letters B 789 (2019) 444-471	
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Evidence for collectivity The CMS Collaboration* CERN, Switzerland	y in pp collisions at the LHC	CrossMark	Correlated long-range p+Pb and low-multipli The ATLAS Collaboration*	mixed-harmonic fluctuations measured in <i>pp</i> , icity Pb+Pb collisions with the ATLAS detector	Check for updates
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Article history: Receive 20 June 2016 Receive 20 June 2016 Accepted 5 December 2016 Available online 13 December 2016 Editor: M. Doser Editor: M. Doser Reywords: CMS Physics Heavy ion Ridge Correlation pp	Measurements of two- and multi-particle angular correlations in pp collisions at $\sqrt{s} = 5$, 7, and 13 TeV are presented as a function of charged-particle multiplicity. The data, corresponding to integrated furminosities of 1.0pc ⁻¹ (5 TeV), 6 2pc ⁻¹ (7 TeV), and 0.7 pc ⁻¹ (13 TeV), were collected using the CMS detector at the LHC. The second-order (v_2) and third-order (v_2) azimuthal anisotropy harmonics of unidentified charged particles, as well as v_3 of x_3^0 and A/Λ particles, are extracted from long-range two- particle correlations as functions of particle multiplicity and transverse momentum. For high-multiplicity pp events, a mass ordering is observed for the v_2 values of charged hadrons (mostly pions), k_3^0 , and A/Λ , with lighter particle species exhibiting a stronger azimuthal anisotropy signal below $p_1 \approx 2$ CeV/c. For 13 TeV data, the v_2 signals are also extracted from four- and six-particle correlations for the first time in pp collisions, with comparable magnitude to those from two-particle correlations for the observed ions in high-multiplicity portices observations are similar to those seen in pPb and PMPb collisions. © 2016 The Author. Published by Elsevier B.V. This is an open access article under the CC BY license (http://creativecommons.org/licenses/tb/4.0). Funded by SCOAP ³ .		Article history: Correlations of two flow harmonics v, and v_w, via three- and four-particle cumulants are measured in 13 TeV peripheral Pb-Pb Collisions with the ATLAS detector at the LHC. The goal is to understand the multi-particle nature? Received in revised form 7 Octore 2018 Accepted 13 November 2018 Available online 2 January 2019 Editor: D.F. Geesaman Editor: D.F. Geesaman Editor: D.F. Geesaman Initiation of the the three of the long-range collective phenomenon in these sequences of the page on flow background from digit production present in the standard cumulant method is suppressed using a method of subsevent cumulants involving two, three and four subsevents expanding. The results show a negative correlation strems, the question of the cumulants involving two, three and four subsevents expanding and collision systems. The relation strems and over the full multiplicity range. However, the correlation strems and over the full multiplicity range. However, the unulants with the (V2) from a two-particle correlation strems momentum. Tange and collision systems. The relation strems and over the full multiplicity range. However, the subsevent methods provide turge correlation strems momentum. These results based on the subsevent methods provide turge viewer multiplicity and transverse momentum. These results based on the subsevent methods provide turge viewers de similar in the three collision system. The relation strems and policy and transverse momentum. These results based on the subsevent methods provide turge viewers described on the subsevent methods provide turge viemes described on the subsevent methods provi		



nature physics PUBLISHED ONLINE: 24 APRIL 2017 [DOI: 10.1038/APRH25411] OPEN Enhanced production of multi-strange hadrons in

high-multiplicity proton-proton collisions

ALICE Collaboration[†]

At sufficiently high temperature and energy density, nuclear matter undergoes a transition to a phase in which quarks and gluons are not confined: the quark-gluon plasma (QGP)1. Such an exotic state of strongly interacting quantum chromodynamics matter is produced in the laboratory in heavy nuclei high-energy collisions, where an enhanced production of strange hadrons is observed²⁻⁶. Strangeness enhancement, originally proposed as a signature of QGP formation in nuclear collisions⁷, is more pronounced for multi-strange baryons. Several effects typical of heavy-ion phenomenology have been observed in high-multiplicity proton-proton (pp) collisions^{8,9} but the enhanced production of multi-strange particles has not been reported so far. Here we present the first observation of strangeness enhancement in high-multiplicity proton-proton collisions. We find that the integrated yields of strange and multi-strange particles, relative to pions, increases significantly with the event charged-particle multiplicity. The measurements are in remarkable agreement with the p-Pb collision results^{10,11}, indicating that the phenomenon is related to the final system created in the collision. In high-multiplicity events strangeness production reaches values similar to those observed in Pb-Pb collisions, where a QGP is formed.

equilibrium and can be described using a grand-canonical statistical model^{12,13}. In peripheral collisions, where the overlap of the colliding nuclei becomes very small, the relative yields of strange particles to pions decrease and tend toward those observed in pp collisions for which a statistical-mechanics approach can also be applied14,15 Extensions of a pure grand-canonical description of particle pro duction, such as statistical models implementing strangeness canonical suppression16 and core-corona superposition17.18 models, can effectively produce a suppression of strangeness production in small systems. However, the microscopic origin of enhanced strangeness production is not known, and the measurements presented in this Letter may contribute to its understanding. Several effects, such as azimuthal correlations and mass-dependent hardening of p_T distributions, which in nuclear collisions are typically attributed to the formation of a strongly interacting quark-gluon medium, have been observed in high-multiplicity pp and proton-nucleus collisions at the LHC^{8-11,19-25}. Yet, enhanced production of strange particles as a function of the charged-particle multiplicity density $(dN_{ch}/d\eta)$ has so far not been observed in pp collisions. The study of pp collisions at high multiplicity is thus of considerable interest as it opens the exciting possibility of a microscopic understanding of phenomena known from nuclear reactions.



The pT spectra

• Transverse momentum (pT) spectra of identified particles are fundamental physical observables in high-energy collisions

• The evolution of the pT distribution with multiplicity in pp collisions is remarkably similar to the evolution observed in larger colliding systems



- Can be interpreted as originating from the hydrodynamical radial expansion of the produced medium
- **Thermal description:** final state particles are stemming from a system in thermal and chemical equilibrium in the hadronic phase

• The hadronic pT spectra can be described by exponential distributions based on the **Boltzmann-Gibbs** statistics-based effective models :

$$f(E) \sim \exp\left(-\frac{E-\mu}{T}\right)$$

where μ is chemical potential and T is the associated temperature, or simplified to:

$$f(p_T) = A \exp\left(-\frac{p_T}{T}\right)$$

• The main assumption is that the hadron production is governed by physically different, non-thermal and non-equilibrium perturbative QCD processes at low and high pT values

• In order to connect these two parts in a meaningful way, an improvement in the statistical picture is the generalization of the Boltzmann-Gibbs theory using a particular generalization of the exponential and logarithm functions:

$$\exp_q(x) = [1 + (1 - q)x]^{1/(1 - q)}$$

• The distributions with this form are usually called **Tsallis-Pareto** distributions, and they connect the powerlaw tailed and exponential-like (i.e. Boltzmann-Gibbs) distributions in a smooth way, simplified to:

$$f(p_T) = A \left(1 + p_T \frac{q-1}{T}\right)^{-\frac{q}{q-1}}$$

The pT spectra fitting

• The hadronic pT spectra can also be studied by means of the Boltzmann-Gibbs blast-wave model in terms of a few collective variables, in particular temperature, longitudinal, and transverse flow:

$$\frac{dN}{p_T dp_T} \propto \int_0^R r \ dr \ m_T \ I_0 \left(\frac{p_T \sinh \rho}{T_{kin}}\right) K_1 \left(\frac{m_T \cosh \rho}{T_{kin}}\right)$$
where $\rho = \tanh^{-1}(\beta_T)$

But the temperature is related to the mean transverse momentum... and to the **fit range**



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ALICE data in different event multiplicity classes: I – highest, X – lowest class.





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• Let's suppose that the region where a Boltzmannian fit gives a satisfactory fit correspond to soft interactions

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- Then the pT-spectra may be roughly separated into three regions:
 - 1) low-pT part representing purely soft interaction
 - 2) intermediate part where the hard interactions get mixed with soft ones
 - 3) hard interactions beyond pT ~ 4 GeV



Universal crossing for all LHC energies

Crossing at the unity value at ~0.6 GeV: hint for soft limit?



Soft/hard limit?

• Two different observations lead to the same conclusion: the pT-spectra in a given multiplicity class stems from at least two different processes:

1. Predominantly soft region (low- $p_{\rm T}$): 0.15 GeV/ $c \le p_{\rm T} \le 0.6$ GeV/c, 2. Mixed (soft and hard) region (higher- $p_{\rm T}$): 0.6 GeV/ $c < p_{\rm T} \le 4$ GeV/c.



Hedgehogs (erizos) at the LHC?









Introduction to "hedgehog-like" events

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Looking for the presence of events with a very
 extended structure of low momentum tracks filling in a uniform way the pseudorapidity-azimuth (η-φ) space.

- First dedicated analysis of highest transverse energy (E_T) events seen in the UA1 detector at the SppS collider at CERN in proton-antiproton collisions at $\sqrt{s} = 630$ GeV
- Several isotropic events with $E_T \sim 210$ GeV in UA1 observed (even tested for top quark production), no evidence for non-QCD mechanism for these events.
- Similar unusual events observed in p-pbar collisions at $\sqrt{s} = 1.8$ TeV by CDF's Run 1 detector with more than 60 charged particles and $E_T \sim 320$ GeV
- Called "hedgehog-like" events by C. Quigg
- Taken for granted that in these events with high ${\sf E}_{\sf T}$ perturbative aspects of QCD dominate the event properties: multi-jet events.



<u>UA1 Collaboration, Zeit. für Phys. C,</u> <u>V. 36, p. 33 (1987)</u>



Geometry of the final state: event shapes

• Event shape variables: instrumental in classifying the **geometrical and topological configurations** of the final-state particles produced in high-energy collisions at PETRA, ISR, SppS, SLD, LEP, HERA, Tevatron and the LHC.



Commonly used event shape variables

• Study the isotropy of the final-state energy distribution by defining the **linearized sphericity tensor of the event**, where the Greek indices denote the x, y, and z components of the momentum of the charged particle *i*. The eigenvalues must satisfy the normalization condition: $\lambda_1 \ge \lambda_2 \ge \lambda_3$ and $\lambda_1 + \lambda_2 + \lambda_3 = 1$.

$$S^{\alpha\beta} = \frac{\displaystyle\sum_{i} p^{\alpha}_{i} p^{\beta}_{i}}{\displaystyle\sum_{i} |\vec{p}_{i}|^{2}} \ , \label{eq:sample}$$

$$A = \frac{3}{2}\lambda_3; \quad S = \frac{3}{2}(\lambda_2 + \lambda_3).$$

• Aplanarity (A) serves as a measure of how planar an event is. A balanced pencil-like event corresponds to A = 0, and an isotropic event corresponds to A = 1/2.

• **Sphericity (S)** quanties the isotropy of an event, representing the degree to which energy and momentum are evenly distributed in all directions. S= 0 denotes a balanced dijet event, and S = 1 for an isotropic event.

Commonly used event shape variables

• **Centrality**: a measure of how much of the event is contained within the central part of the detector; ranges between 0 and 1, where a pencil-like has C = 0 and a centrally contained event corresponds to C = 1.

$$C = \frac{\sum_{i} p_{T,i}}{\sum_{i} E_{i}},$$

$$T = 1 - \max_{\hat{n}} \frac{\sum_{i} |\vec{p}_{\mathrm{T},i} \cdot \hat{n}|}{\sum_{i} |\vec{p}_{\mathrm{T},i}|},$$

• **Transverse thrust**: a widely used event shape ranging from 0 for a pencil-like topology to 1/3 for a circularly symmetric distribution of particles in the transverse plane.

• Transverse spherocity: infrared and colliner safe event shape that ranges from $S_0 = 0$ for events with back-to-back multijet final states to $S_0 = 1$ for isotropic event topologies.

$$S_0 = \frac{\pi^2}{4} \min_{\hat{n}_s} \left(\frac{\sum_i |\vec{p}_{T,i} \times \hat{n}_s|}{\sum_i p_{T_i}} \right)^2.$$

Characterisation of high-multiplicity events

• Attempts to characterise these high-multiplicity events: use of event shapes, i.e. using **transverse sphericity**: $2\lambda^{xy}$ $\sum_{n=1}^{\infty} 1 [n^2 + n + n]$

$$S_{\perp} = \frac{2\lambda_2^{xy}}{\lambda_1^{xy} + \lambda_2^{xy}} , \quad S^{xy} = \sum_i \frac{1}{|\vec{p}_{\mathrm{T},i}|^2} \begin{bmatrix} p_{x,i}^2 & p_{x,i} p_{y,i} \\ p_{x,i} p_{y,i} & p_{y,i}^2 \end{bmatrix}$$

- Both ALICE and ATLAS observed an under-estimation of isotropic events by MC generators at high charged multiplicity ($N_{ch} \ge 30$)
 - Suggest that a very active underlying event (UE) is needed by the MC event generators in order to explain these high-multiplicity events





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• ALICE measurement shows that $< p_T >$ as a function of N_{ch} in isotropic events was found to be smaller than that measured in jet-like events, and that for jet-like events, the $< p_T >$ is over-estimated by PYTHIA 6 and 8 models.



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• Recently, a new event shape parameter, flattenicity, was proposed [<u>A. Ortiz</u>, <u>G. Paic, Rev. Mex. Fis. Suppl. 3 (2022) 4, 040911</u>] that allows one to identify and characterise high-multiplicity events with a quasi-isotropic distribution in a wide pseudorapidity range in proton-proton collisions.

• MC event generators are able to model "hedgehog" events, which opens the possibility to study their properties and find a potential way to experimentally trigger these events.

• The idea: find out how uniform the p_T of tracks is distributed in a given event!



Calculating flattenicity

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• Build 10 x 12 grid in $(\eta - \phi)$ space:



Calculating flattenicity

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• Build 10 x 12 grid in $(\eta - \phi)$ space:



- Event-by-event, the relative standard deviation of the $p_{\rm T}^{\rm cell}$ distribution is obtained.
- Events with isotropic distribution of particles ("hedgehogs") are expected to have $\rho = 1$.

$$\rho = 1 - \frac{\sqrt{\sum_{i=1}^{i=120} \left(p_{\rm T}^{\rm cell,i} - \langle p_{\rm T}^{\rm cells} \rangle \right)^2 / \mathcal{N}_{\rm cell}^2}}{\langle p_{\rm T}^{\rm cells} \rangle},$$

Benchmark hedgehog-like structures



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 $\rho = 0.98$



p_T [GeV]

Flattenicity using different MC generators

• Two sets of inelastic events, including non-diffractive and diffractive components, were simulated using **Pythia 8.309**: ATLAS A14 tune and CMS CUETP8M1 tune.

- Herwig 7.2.0 with SoftTune based on the MMHT2014 LO PDF
- Recently released EPOS 4.0.0 framework using authors tune

• First, **revise the performance** of these MC models in reproducing the 13 TeV pp collision data.



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- Herwig 7.2.0 with SoftTune based on the MMHT2014 LO PDF
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- Then, compare transverse spherocity and flattenicity for the 4 MC settings



• To demonstrate flattenicity's value in identifying isotropic events and its complementary information about the global shape of an event, we study its correlation with other event shape variables.

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• Overall, there is clear indication that while flattenicity is related to other event shape variables, it **provides distinct and complementary** info of the isotropy of events



• Classify the **top 10%** of events in the distributions of flattenicity, transverse spherocity, and transverse thrust for each multiplicity class as **isotropic events**.

 Our observation is that for the highest charged-particle multiplicity class, only 3.1% of events meet all the three criteria: clearly indicates that flattenicity selects a different subset of events as compared to widely used S₀ and T.



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Conclusions

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• We notice that, traditionally, collision data observables are presented as **mean values** of more differential distributions, such as mean charged-particle multiplicity (Nch), mean of transverse momentum (pT), mean values of anisotropic flow, etc.

• The agreement or disagreement of the means of a model with the data may not really point to the details of the result, given that different models may have the same means as the data but with different underlying assumptions.

• We observe that the mean transverse momentum $\langle pT \rangle$ is highly sensitive to the selected pT range, and it is not sensitive within the soft region (low-pT) region studied of 0.15 GeV $\leq pT \leq 0.6$ GeV.

• Universal crossing point of the ratios to the inclusive spectrum at $pT \approx 0.6$ GeV.

• We are in presence of two different QCD processes: when studying **collective effects**, it is important to notice that when extending the pT range above the 0.6 GeV, one risks important contributions of hard QCD nature.

- Hedgehog events have never been seriously studied in pp collisions at the LHC.
- Flattenicity the new event structure parameter allows one to identify the hedgehog events and observe the evolution of events from jetty to hedgehog type.
- pT spectra as a function of flattenicity already studied in ALICE, see Paola Vargas talk
- When compared to other event shape variables widely used in the literature, such as transverse spherocity, we found that flattenicity is able to **identify** a subset of isotropic events with **hedgehog-like structures**.
- More experimental measurements using flattenicity are needed! Look for "hedgehog" events in data!





Muchas gracias por su atención!

