## Can the QCD axion feed a DE component? [Can the QCD axion play a role in accelerating the expansion?]



### [arXiv:2405.00090] with: K. Müürsepp (NICPB, Tallinn) & C. Smarra (SISSA, Trieste)

### XV Latin American Symposium on High Energy Physics November 4 - 8, 2024, Cinvestav, Mexico City

# Enrico Mardi





### The axion defining interaction

 $\mathscr{L}_{a} = \frac{\alpha_{s}}{8\pi} \left( \frac{a(x)}{F} + \bar{\theta} \right) G\tilde{G} + \mathscr{L} \left( \partial_{\mu} a(x), \psi, \varphi, A_{\mu} \right) + \left[ \delta \mathscr{L}_{\text{eff}}(a(x), \ldots) \right]$ 

 $a \rightarrow a + \text{const.}$ 

 $a \rightarrow a + const$  invariant

Absent or suppressed  $\Lambda_{\text{eff}} \sim m_P \& d \ge 10$ 

### The axion defining interaction

$$\mathscr{L}_{a} = \frac{\alpha_{s}}{8\pi} \left( \frac{a(x)}{F} + \bar{\theta} \right) G\tilde{G} + \mathscr{L} \left( \underbrace{a(x)}_{A \to a + \text{const.}} - \widehat{\theta} \right) G\tilde{G} + \mathscr{L} \left( \underbrace{a(x)}_{A \to a + \text{const.}} - \underbrace{a(x)}_{a \to a + \text{const.}} \right)$$

- $\theta$  is removed via a shift of the axion field  $a \to a \theta F$
- 3. The  $a \ GG$  interaction generates a mass term:

$$F^2 m_a^2 = i \int d^4 x \left\langle \frac{\alpha_s}{8\pi} G \tilde{G}(x) \frac{\alpha_s}{8\pi} G \tilde{G}(0) \right\rangle$$

 $\partial_{\mu}a(x), \psi, \varphi, A_{\mu} + [\delta \mathscr{L}_{eff}(a(x), \ldots)]$ 

 $\rightarrow a + const$  invariant

Absent or suppressed  $\Lambda_{\rm eff} \sim m_P \& d \ge 10$ 

2. Minimum of the vacuum energy occurs for  $\langle a(x) \rangle \rightarrow 0$ : solves strong CP problem

))  $\equiv \chi \leftarrow$  "Topological susceptibility"



In a hot plasma, at T >> T<sub>c</sub>, free color charges screen the correlator:  $\chi = 0$ 

- At T < T<sub>c</sub> color charges are confined in SU(3) singlets, no screening:  $\chi = (160 \text{ MeV})^4$

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$$m_a^2 = m_a^2(T)$$

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$$\mathcal{X} = \mathcal{X}(\mathsf{T}) \Longrightarrow$$

What is the T dependance?

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 $m_a^2(T) \sim T^{-n} [n \sim n(T)]$ 

In a hot plasma, at T >> T<sub>c</sub>, free color charges screen the correlator:  $\chi = 0$ 

What is the T dependance?

DIGA (lowest order):  $n = \beta_0$  -

IILM (more appropriate for [Interacting inst. liquid model: Shella

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$$m_a^2 = m_a^2(T)$$

$$m_a^2(T) \sim T^{-n} [n \sim n(T)]$$

$$-n_{f} - 4 = \frac{11}{3}N + \frac{1}{3}n_{f} - 4 \quad n = 8 \text{ (QCD)}$$

$$T \sim T_{osc}): \qquad n \sim 6.68$$
ard & Wanz, 2010]



#### **Effective mass, lattice calculations**

#### Lattice QCD: we can compute axion mass

$$m_a^2 f_a^2 = \chi(T)$$



At high T (no mesons) we can analytically compute potential (DIGA)  $V(\theta) = -\chi(T)\cos\theta$ 

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$$V(\theta) = -\chi(T)\cos\theta$$

10<sup>-1</sup>

 $10^{-2}$ 

 $10^{-3}$ 

Take away message: Even in canonical QCD [SU(3),  $n_f=3$ ]

is a reasonable T dependence, at least in some transient regime

Lattice QCD 2+1+1 [Borsanyi] Lattice QCD 2+1 [Bonati] Lattice QCD (DWF) 2+1 [Buchoff] (points) DIGA (T>>Tc) [Borsanyi] IILM [Wantz]



### Particles with varying mass: Effective Equation of State



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implies an effective EoS:  $p_a = w \rho_a$  with w = -n/6

- Taking  $m_a^2(T) \sim T^{-n}$ , the conserv. law  $d(\rho_a a^3) = -p_a da^3$
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No. Not enough energy density:  $\rho_b \lesssim \Lambda_b^4 < T_0^4 \sim \rho_{rad} \ll \rho_{DE}$ 

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- If a dominates puniverse, => <u>acceleration</u> already for n > 2
- Could a PNGB b(x), coupled to a "dark" gauge group Gb that is undergoing a confining PT <u>now</u> ( $\Lambda_b < T_0$ ) produce Cosmo accel.?







### Take $G_a \times G_b$ , $G_a = SU(3)_{QCD}$ ; $G_b = SU(3)$ or SU(2); $\Lambda_a \gg \Lambda_b$





### $\mathscr{L}_{V} \sim \bar{\psi}_{I} \psi_{R} \Phi_{1} + \bar{\chi}_{I} \chi_{R} \Phi_{2} \rightarrow$ $\psi \sim (1,3), \ \chi \sim (3,3)$

Take  $G_a \times G_b$ ,  $G_a = SU(3)_{QCD}$ ;  $G_b = SU(3)$  or SU(2);  $\Lambda_a \gg \Lambda_b$ 

$$\rightarrow \quad \bar{\psi}_L \psi_R v_1 e^{i\frac{a_1}{v_1}} + \bar{\chi}_L \chi_R v_2 e^{i\frac{a_2}{v_2}}$$







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 $\mathscr{L}_V \sim \bar{\psi}_I \psi_R \Phi_1 + \bar{\chi}_I \chi_R \Phi_2$  $\psi \sim (1,3), \ \chi \sim (3,3)$ 

This generates the potenti  $V = \Lambda_a^4 \left[ 1 - \cos\left(\frac{\varphi_a}{r}\right) \right] + \Lambda_b^4 \left[ 1 - \frac{\varphi_a}{r} \right] + \Lambda_b^4 \left[ 1 - \frac{\varphi_a}$ 

$$G_{b} = SU(3) \text{ or } SU(2); \Lambda_{a} \gg \Lambda_{b}$$

$$\rightarrow \quad \bar{\psi}_L \psi_R v_1 e^{i\frac{a_1}{v_1}} + \bar{\chi}_L \chi_R v_2 e^{i\frac{a_2}{v_2}}$$

ial:  

$$F, F' \propto v_2, \ f \propto v_1$$

$$COS\left(\frac{\varphi_a}{F'} + \frac{\varphi_b}{f}\right)$$

$$\left(\begin{array}{c}\varphi_a\\\varphi_b\end{array}\right) = \begin{pmatrix}cos\beta & sin\beta\\-sin\beta & cos\beta\end{pmatrix}\begin{pmatrix}a_1\\a_2\end{pmatrix}$$









 $\ddot{A} + 3H\dot{A} +$  $A = \begin{pmatrix} \varphi_a \\ \varphi_b \end{pmatrix}; \quad \mathcal{M}^2 = m_a^2 \begin{pmatrix} 1 & \epsilon r(T) \\ \epsilon r(T) & r(T) \end{pmatrix}$ 

$$\mathcal{M}^{2}A = 0$$
  
$$(f); \qquad m_{a} = \frac{\Lambda_{a}^{2}}{F}, \quad r(T) = \frac{m_{b}^{2}(T)}{m_{a}}, \quad \epsilon = \frac{f}{F'}$$





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Assumption: at T=0  $m_b = \Lambda_b^2 / f > m_a$  [f<<F, i.e. v<sub>1</sub> << v<sub>2</sub>]



$$\ddot{A} + 3H\dot{A} + \mathscr{M}^{2}A = 0$$

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This implies a Level Crossing  $m_b(T_{LC}) = m_a$  (width  $\Gamma_{LC} \sim 3\epsilon$ ) where QCD axions  $\varphi_a$  can partially convert into b-axions  $\varphi_b$ 

Assumption: at T=0  $m_b = \Lambda_b^2 / f > m_a$  [f<<F, i.e. v<sub>1</sub> << v<sub>2</sub>]







t<sub>LC</sub>

Adiabatic  $m_a (\epsilon t_{LC}) \gg 1$ Plot: [ $\epsilon t_{LC} m_a = 50$ ]





t<sub>LC</sub>

Adiabatic  $m_a (\epsilon t_{LC}) >> 1$ Plot: [ $\epsilon t_{LC} m_a = 50$ ]









Adiabatic  $m_a (\epsilon t_{LC}) >> 1$ Plot:  $[\epsilon t_{LC} m_a = 50]$ 

Diabatic











Adiabatic  $m_a$  ( $\epsilon$  t<sub>L</sub>c) >> 1 Plot: [ $\epsilon$ t<sub>L</sub>c m<sub>a</sub> =50]

Diabatic  $m_a$  ( $\epsilon$  t<sub>L</sub>c)  $\lesssim 1$ Plot: [ $\epsilon$ t<sub>L</sub>c m<sub>a</sub>=1]







Severe Constraining Conditions



 $f > T_{\rm LC} > T_{\rm DE} > T_0 > \Lambda_b$ 



Severe Constraining Conditions

$$m_b(T_{\rm LC}) \sim \frac{\Lambda_b^2}{f} \left(\frac{T_b}{T_{\rm LC}}\right)^3 = m_a = \frac{\Lambda_a^2}{F}$$

 $f > T_{\rm LC} > T_{\rm DE} > T_0 > \Lambda_b$ 

Which imply a pre-inflation scenario  $F \gtrsim 10^{14} \,\text{GeV}, \ [m_a \lesssim 6 \cdot 10^{-8} \,\text{eV}], \ \theta_a \lesssim 6\%$ 



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$$F \gtrsim 10^{14} \,\text{GeV}, \quad [m_{a} \lesssim 6 \cdot 10^{-8} \,\text{eV}], \quad \theta_{a} \lesssim 6 \,\%$$
And a non-adiabatic level crossing
$$\epsilon \sim 10^{-25} \left(\frac{\Lambda_{b}}{10^{-4} \,\text{eV}} \frac{160 \,\text{MeV}}{\Lambda_{a}}\right)^{2}$$

 $t_{\rm LC} = 10^9 \,\mathrm{yr}, \ [z_{\rm LC} \sim 5] \quad \Rightarrow \quad m_a t_{\rm LC} \lesssim 10^{25}$ 

![](_page_30_Figure_4.jpeg)

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![](_page_31_Figure_4.jpeg)

### Theoretical Cosmology, first half of XX century: Two confirmed predictions

![](_page_32_Picture_2.jpeg)

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#### The Universe is expanding: Friedmann (1922), Lemaitre (1927)

### Observational confirmation Hubble (1929)

![](_page_33_Picture_4.jpeg)

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2. The expansion is accelerating: Bondi & Gold (1948); Hoyle (1948)

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Observational confirmation Riess (1998) et al.; Perlmutter et al. (1999)

![](_page_34_Picture_6.jpeg)

![](_page_34_Picture_7.jpeg)

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1. Follows from Cosmological Principle: Universe homogeneous and isotropic on large scales 2. Follows from <u>Perfect</u> Cosmological Principle: Universe unchanging in time on large scales

![](_page_35_Picture_9.jpeg)

![](_page_35_Picture_10.jpeg)

![](_page_35_Picture_11.jpeg)

### Theoretical Cosmology, first half of XX century: Two confirmed predictions

- 1. The Universe is expanding: Friedmann (1922), Lemaitre (1927)
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Perfect Cosmological Principle (Bondi & Gold, 1948): Cosmological principle extended by assuming the Universe to be homogeneous in space and in time (i.e. stationary).

### Observational confirmation Hubble (1929)

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1. Follows from Cosmological Principle: Universe homogeneous and isotropic on large scales 2. Follows from <u>Perfect</u> Cosmological Principle: Universe unchanging in time on large scales

![](_page_36_Picture_10.jpeg)

"Present observations indicate that the universe is expanding. This suggests that the mean density in the past has been greater than it is now. If we are now to make any statement regarding the behaviour of such a denser universe [...] then we have to know the physical laws and constant applicable in a denser universe. But we have no determination for those."

"Physical laws cannot be assumed to be independent on the structure of the Universe. Conversely, the structure of the Universe depends upon the physical laws." Then there may be a stable, self–perpetuating state with constant physical laws.

![](_page_37_Figure_3.jpeg)

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<u>Steady State Universe (SSU)</u>: to counterbalance dilution from the expansion, matter is constantly created at the rate of 1 H atom (or 1 neutron)/cm<sup>3</sup>/10<sup>12</sup> yrs.

![](_page_38_Figure_5.jpeg)

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Pmatter = const. => accelerated expansion (i.e. and effective EoS: w = -1)

![](_page_40_Figure_6.jpeg)

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Our construction also yields  $\rho_b = const$ . but is not "steady state". Standard cosmological history unaltered until LC at  $z \sim 2 - 10$ . "matter creation" => "mass generation from phase transition in a dark plasma"

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**P**matter = const. => accelerated expansion (i.e. and effective EoS: w = -1)

![](_page_41_Figure_9.jpeg)

![](_page_41_Picture_10.jpeg)

### Conclusions

![](_page_42_Picture_2.jpeg)

![](_page_43_Picture_0.jpeg)

- A coupled 2 axions system can generate DE from DM, and explain both phenomena

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![](_page_43_Picture_4.jpeg)

![](_page_43_Picture_5.jpeg)

- It is consistent with different evolving EoS Quintessence (w > -1),  $\Lambda$  (w = -1), Phantom (w < -1), Quintom [ $w(t) < -1 \rightarrow w(t) > -1$ ]

- A coupled 2 axions system can generate DE from DM, and explain both phenomena

![](_page_44_Picture_6.jpeg)

![](_page_44_Figure_7.jpeg)

- It is consistent with different evolving EoS
- It can shed light on the "why now ?" puzzle

- A coupled 2 axions system can generate DE from DM, and explain both phenomena

Quintessence (w > -1),  $\Lambda(w = -1)$ , Phantom (w < -1), Quintom  $[w(t) < -1 \rightarrow w(t) > -1]$ 

![](_page_45_Picture_9.jpeg)

![](_page_45_Figure_10.jpeg)

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- If the QCD axion constitutes the DM, there is not much freedom for model building. Only viable for pre-inflationary axion scenarios.

![](_page_46_Picture_6.jpeg)

![](_page_46_Figure_7.jpeg)

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![](_page_47_Picture_5.jpeg)

![](_page_47_Picture_7.jpeg)

![](_page_47_Figure_8.jpeg)

![](_page_47_Picture_12.jpeg)

# Can the QCD axion feed a DE component?

with: K. Müürsepp & C. Smarra [arXiv:2405.0009

![](_page_49_Picture_3.jpeg)

# Enrico Nardi

![](_page_49_Picture_6.jpeg)

Laboratori Nazionali di Frascati

### CATCH22+2 DIAS, Dublin, May 1-5 2024

![](_page_49_Picture_9.jpeg)

# Can the QCD axion feed a DE component? arXiv:2405.00090 Enrico Nardi

with: Kristjan. Müürsepp (HEPC-NICPB, Tallinn) & Clemente Smarra (SISSA & INFN, Trieste)

![](_page_50_Picture_2.jpeg)

![](_page_50_Picture_3.jpeg)

![](_page_50_Picture_4.jpeg)

lstituto Nazionale di Fisica Nucleare

Laboratori Nazionali di Frascati

![](_page_50_Picture_7.jpeg)

# Can the QCD axion feed a DE component? arXiv:2405.00090 Enrico Mardi

![](_page_51_Picture_2.jpeg)

![](_page_51_Picture_3.jpeg)

#### with: Kristjan. Müürsepp (HEPC-NICPB, Tallinn) & Clemente Smarra (SISSA, Trieste)

![](_page_51_Picture_5.jpeg)

![](_page_51_Picture_6.jpeg)

di Fisica Nucleare

![](_page_51_Picture_8.jpeg)