

Electroweak precision tests and global fits

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1. Electroweak precision tests
2. Global electroweak fit
3. Precision physics with future e^+e^- colliders
4. Theoretical calculations

Experimental data:

- LEP/SLC: Z(W)-boson properties
- LHC/TeVatron: M_W , M_H , m_t , $\sin^2 \theta_{\text{eff}}^l$
- Other experiments:
 a_μ , PVES, G_μ , α_s , ...

Fit theory model to data:

- **SM** parameters
 M_Z , M_W , M_H , m_t , α_s , α^* , $m_{f \neq t}^{**}$
* fixed ** mostly fixed/negligible
- **BSM** models:
SM + new particle masses/couplings
- **SMEFT/HEFT**:
SM + Wilson coeff. of higher-dim. ops.



LEP
(CERN)



SLC
(SLAC)



LHC
(CERN)

Fermi constant (from μ decay):

$$G_\mu = \frac{\pi\alpha}{\sqrt{2}(1-M_W^2/M_Z^2)M_W^2}(1 + \Delta r)$$

→ prediction of M_W

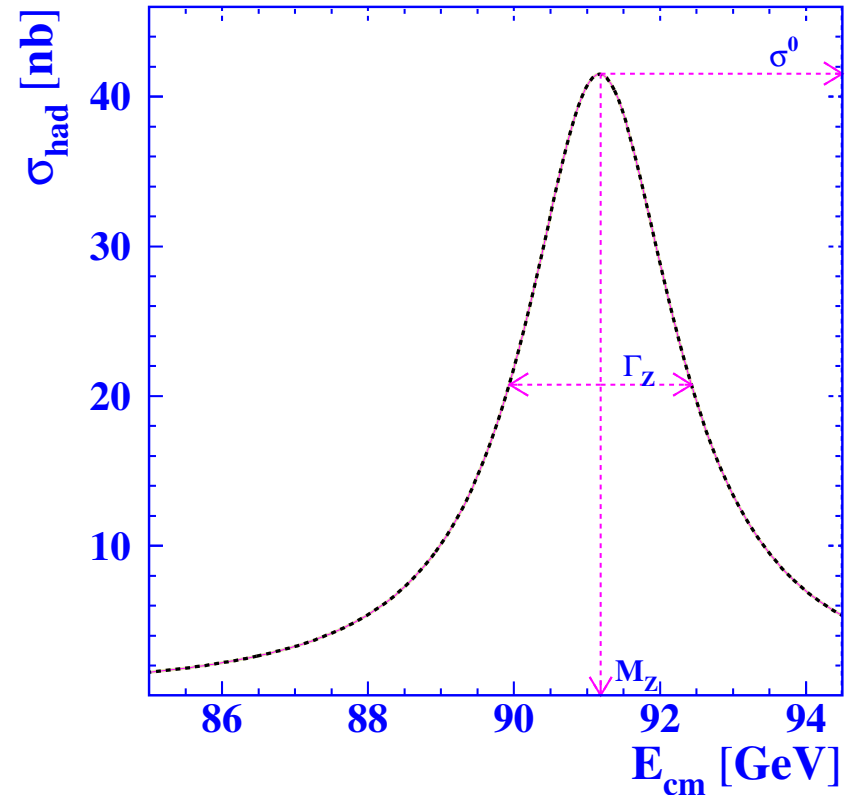
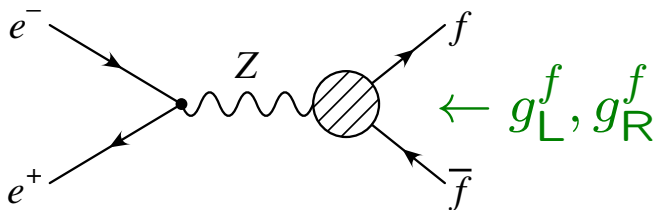
$e^+e^- \rightarrow f\bar{f}$ for $E_{\text{CM}} \sim M_Z$:

Width $\Gamma_Z = \sum_f \Gamma_{ff}$

Branching ratios $R_f = \Gamma_{ff}/\Gamma_Z$

$$\sigma^0 \approx \frac{12\pi\Gamma_{ee}\Gamma_{ff}}{(s-M_Z^2)^2 + M_Z^2\Gamma_Z^2} = \frac{12\pi}{M_Z^2} R_e R_f$$

$$\Gamma_{ff} = C [(g_L^f)^2 + (g_R^f)^2]$$



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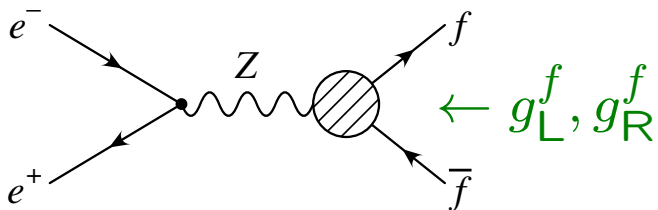
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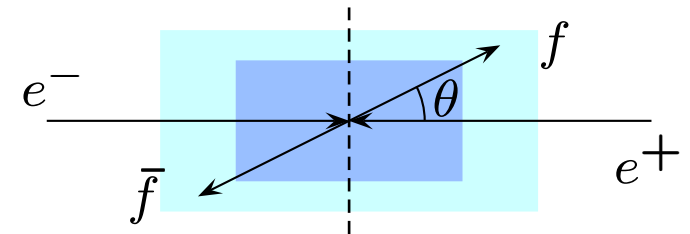


Asymmetries in $e^+e^- \rightarrow f\bar{f}$:

$$A_{\text{FB}} \equiv \frac{\int_{\theta > \frac{\pi}{2}} d\sigma - \int_{\theta < \frac{\pi}{2}} d\sigma}{\int_{\theta > \frac{\pi}{2}} d\sigma + \int_{\theta < \frac{\pi}{2}} d\sigma} = \frac{3}{4} A_e A_f$$

$$A_{\text{LR}} \equiv \frac{\sigma_{e_L} - \sigma_{e_R}}{\sigma_{e_L} + \sigma_{e_R}} = A_e$$

$$\langle \mathcal{P}_\tau \rangle = -A_\tau$$



$$A_f = \frac{2(1 - 4\sin^2 \theta_{\text{eff}}^f)}{1 + (1 - 4\sin^2 \theta_{\text{eff}}^f)^2}$$

$$\sin^2 \theta_{\text{eff}}^f = \frac{g_R^f}{2|Q_f|(g_R^f - g_L^f)}$$

- Deconvolution of initial-state QED radiation:

$$\sigma[e^+e^- \rightarrow f\bar{f}] = \mathcal{R}_{\text{ini}}(s, s') \otimes \sigma_{\text{hard}}(s')$$

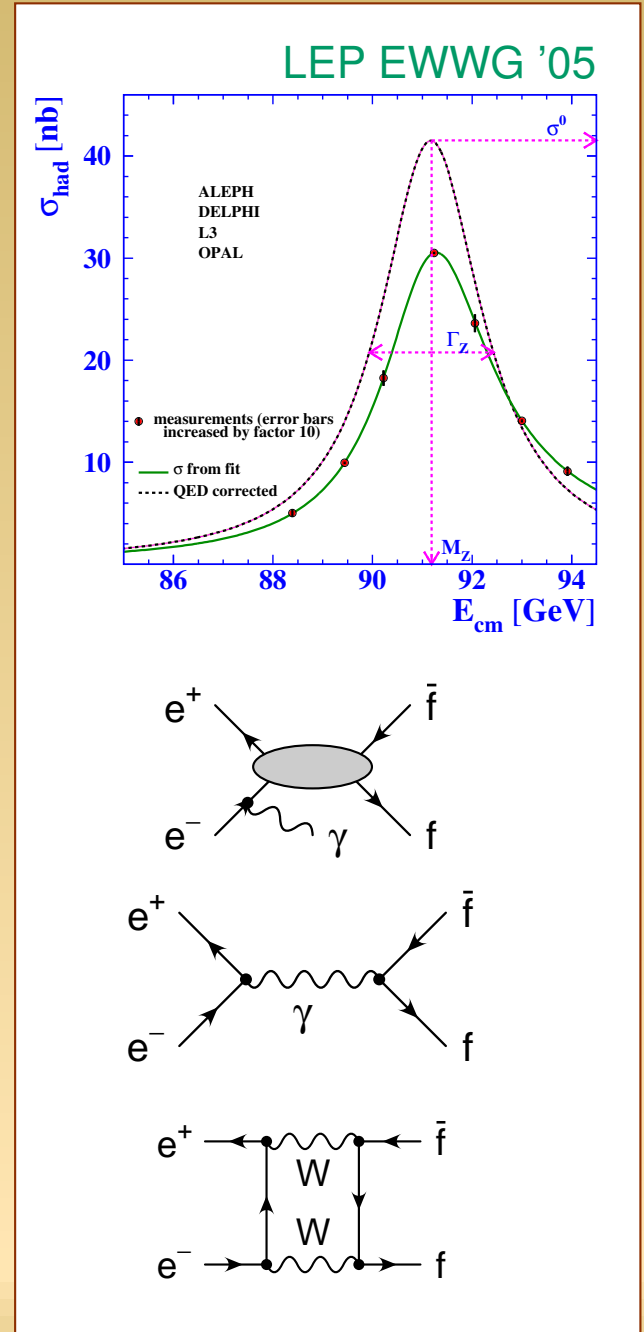
- Subtraction of γ -exchange, γ -Z interference, box contributions:

$$\sigma_{\text{hard}} = \sigma_Z + \sigma_\gamma + \sigma_{\gamma Z} + \sigma_{\text{box}}$$

- Z-pole contribution:

$$\sigma_Z = \frac{R}{(s - M_Z^2)^2 + M_Z^2 \Gamma_Z^2} + \sigma_{\text{non-res}}$$

- Final-state radiation, initial-final interference, etc.
 → Monte-Carlo programs, consistently matched to fixed-order calculations



- Deconvolution of initial-state QED radiation:

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- Subtraction of γ -exchange, γ -Z interference, box contributions:

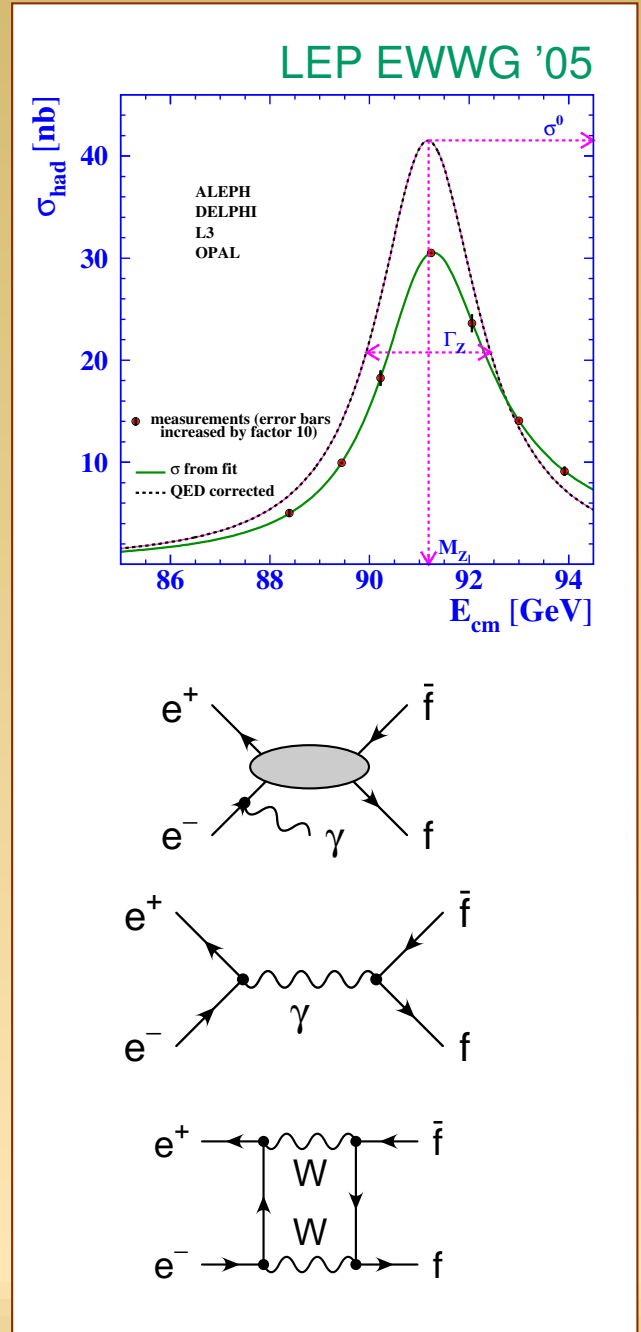
$$\sigma_{hard} = \sigma_Z + \underbrace{\sigma_\gamma + \sigma_{\gamma Z} + \sigma_{box}}_{\text{computed in SM (NLO)}}$$

- Z-pole contribution:

$$\sigma_Z = \frac{R}{(s - M_Z^2)^2 + M_Z^2 \Gamma_Z^2} + \sigma_{non-res}$$

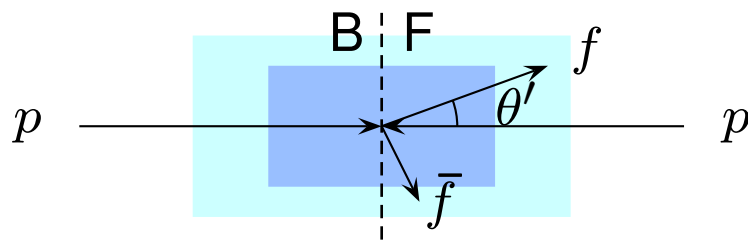
- Final-state radiation, initial-final interference, etc.
 → Monte-Carlo programs, consistently matched to fixed-order calculations

- possible BSM physics?

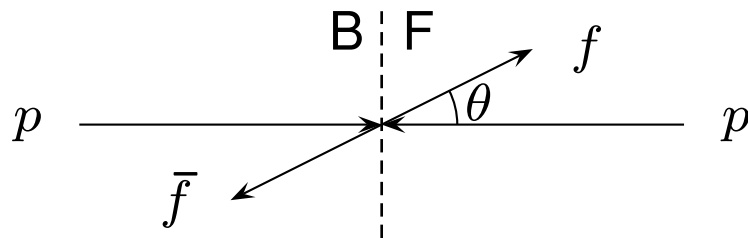


Forward-backward asymmetry:
“forward” defined through event
boost

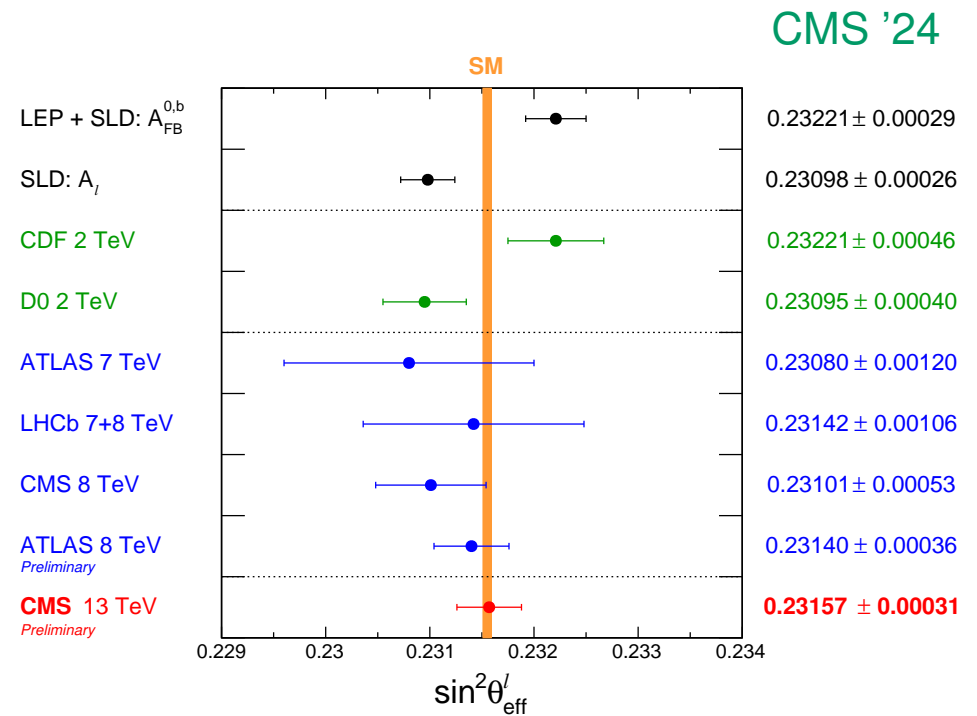
lab frame:



center-of-mass frame:



→ main systematics:
PDFs, QCD corrections



Averages:

PDG '24

$$\sin^2 \theta_{\text{eff}}^l = 0.23151 \pm 0.00016 \text{ (LEP+SLC)}$$

$$\sin^2 \theta_{\text{eff}}^l = 0.23146 \pm 0.00021 \text{ (LHC+TeV)}$$

$$\sin^2 \theta_{\text{eff}}^l = 0.23149 \pm 0.00013 \text{ (colliders)}$$

W mass:

from $pp \rightarrow W^\pm \rightarrow \ell^\pm \nu$,

using m_T and $p_{\ell,\perp}$ distributions

LEP combination
Phys. Rep. 532 (2013) 119

DO
PRL 108 (2012) 151804

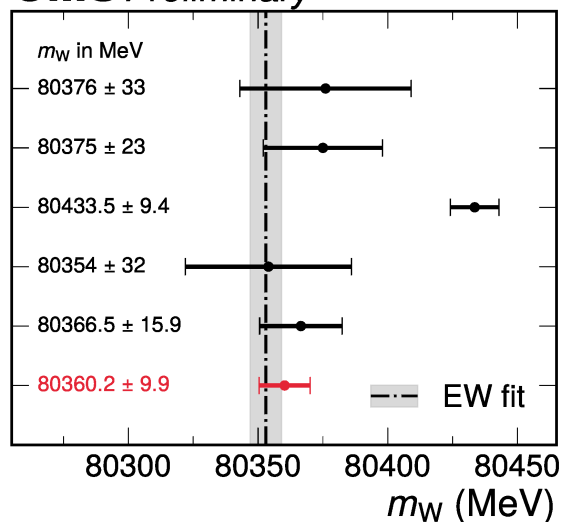
CDF
Science 376 (2022) 6589

LHCb
JHEP 01 (2022) 036

ATLAS
arxiv:2403.15085, subm. to EPJC

CMS
This Work

CMS Preliminary



Averages

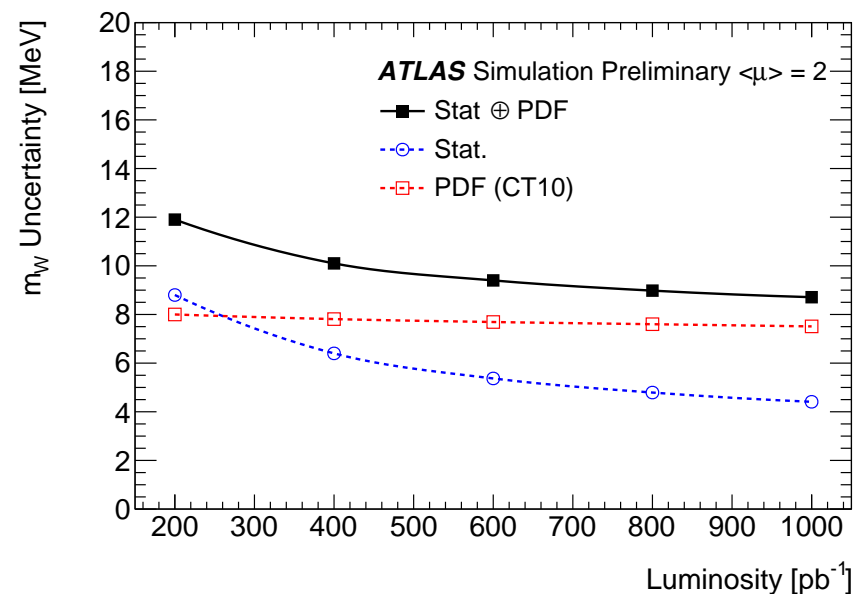
(approximate)

$$M_W = 80.360 \pm 0.008 \text{ GeV (w/o CDF II)}$$

$$M_W = 80.373 \pm 0.007 \text{ GeV (with CDF II)}$$

Ultimate precision at HL-LHC:

$$\delta M_W \sim 5\text{--}10 \text{ MeV}$$



→ talk by M. Velasco

- m_t : Most precise measurement at LHC: $\delta m_t \sim 0.3 \text{ GeV}$ PDG '24

Theoretical ambiguity in mass def.:

Hoang, Plätzer, Samitz '18

$$\begin{aligned} m_t^{\text{CB}}(Q_0) - m_t^{\text{pole}} &= -\frac{2}{3}\alpha_s(Q_0) Q_0 + \mathcal{O}(\alpha_s^2 Q_0) \\ &\approx 0.5 \pm 0.2_{\text{pert.}} \pm 0.2_{\text{np.}} \text{ GeV} \end{aligned}$$

- $m_{b,c}$: $\delta m_{b,c} \sim 8 \text{ MeV}$ (QCD sum rules) Erler, Masjuan, Spiesberger '16,22
- M_H : $M_H = 125.10 \pm 0.09 \text{ GeV}$ (LHC)

● α_s :

PDG '24

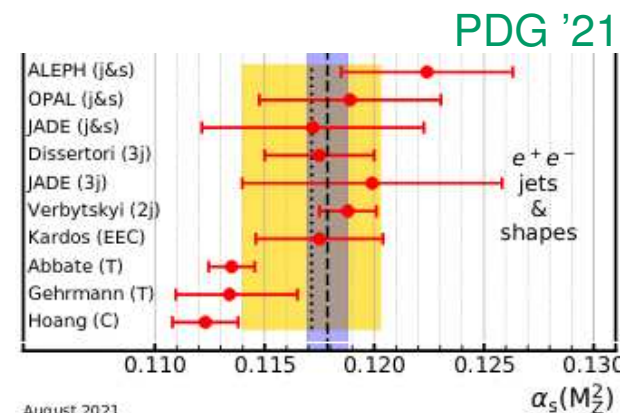
- Most precise determination using Lattice QCD:

$$\alpha_s = 0.1184 \pm 0.0008 \quad \text{FLAG '21}$$

- e^+e^- event shapes: $\alpha_s \sim 0.113 \dots 0.119$

→ Large non-perturbative power corrections

→ Systematic uncertainties?



- Hadronic τ decays: $\alpha_s = 0.1173 \pm 0.0017$

→ Non-perturbative uncertainties in OPE and from duality violation

Pich '14; Boito et al. '15,18

- Hadron colliders: jj , W/Z , $t\bar{t}$, DIS

→ Most precise determination from $pp \rightarrow Z + X$ at 8 TeV:

$$\alpha_s = 0.1183 \pm 0.0009$$

ATLAS '23

- α_s :

- Electroweak precision ($R_\ell = \Gamma_Z^{\text{had}} / \Gamma_Z^\ell$):

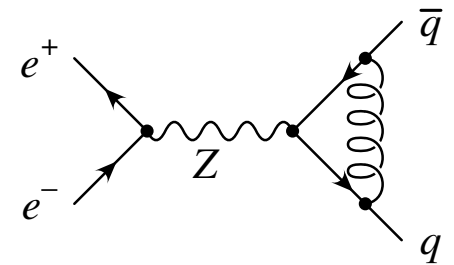
$$\alpha_s = 0.122 \pm 0.003$$

PDG '24

→ Negligible non-perturbative QCD effects

Theory input: $N^4\text{LO QCD corr.} + \text{NNLO EW}$

Caviat: R_ℓ could be affected by new physics

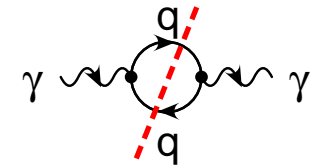


- $\Delta\alpha \equiv 1 - \frac{\alpha(0)}{\alpha(M_Z)} \approx 0.059 = 0.0315_{\text{lept}} + 0.0276_{\text{had}}$

a) $\Delta\alpha_{\text{had}}$ from $e^+e^- \rightarrow \text{had.}$ using dispersion relation

→ Current precision $\sim 10^{-4}$

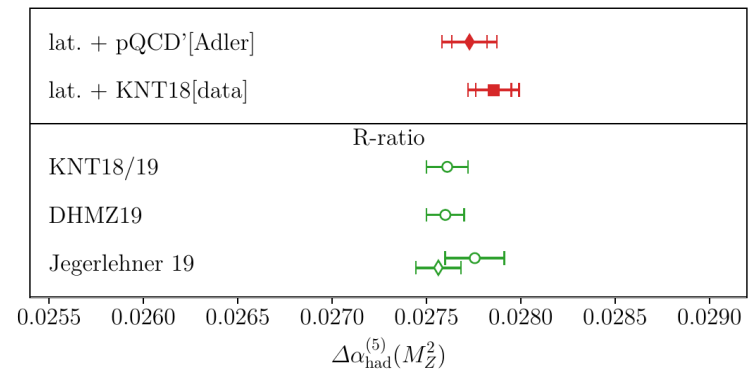
Davier et al. '19; Jegerlehner '19; Keshavarzi, Nomura, Teubner '19



b) $\Delta\alpha_{\text{had}}$ from Lattice QCD
(challenging but much progress)

Burger et al. '15

Cè et al. '22



Future improvements for methods (a) and (b):

- More precise exp./lattice data
- Full 4-loop pQCD for R-ratio / Adler function (for $|Q^2| \gg \Lambda_{\text{QCD}}$)
- More precise inputs for m_b, m_c, α_s

→ $\delta(\Delta\alpha_{\text{had}}) \lesssim 5 \times 10^{-5}$ likely achievable

Jegerlehner '19

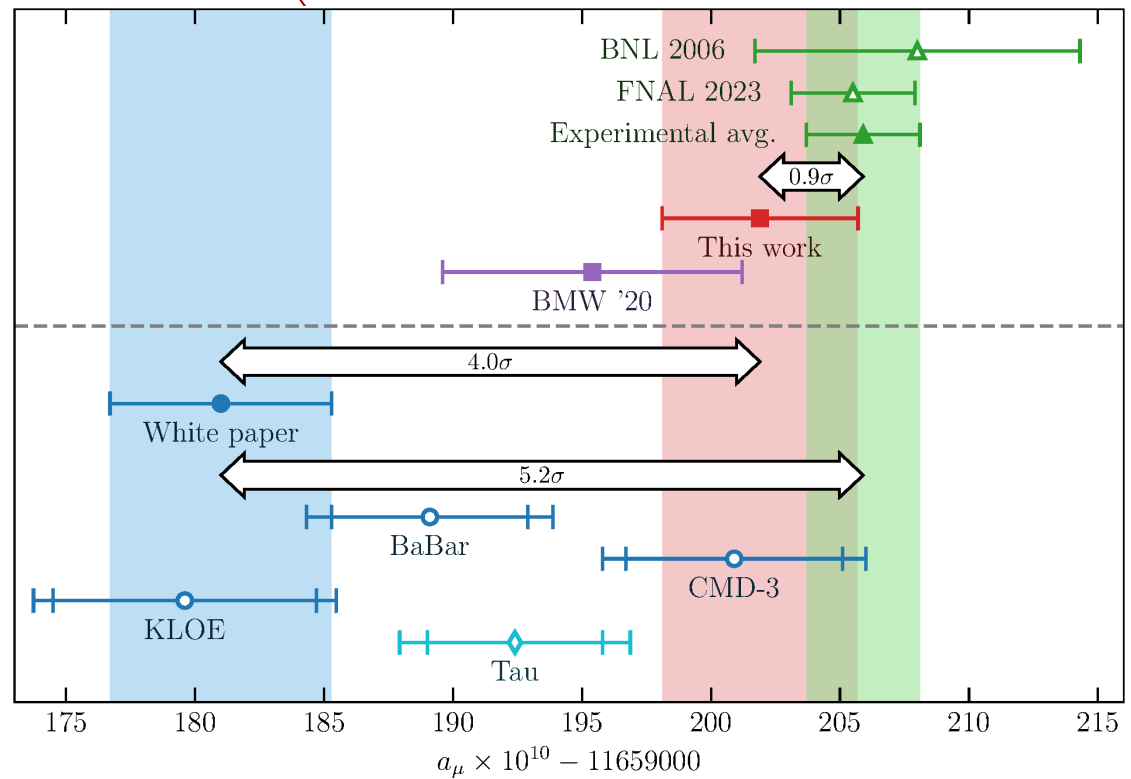
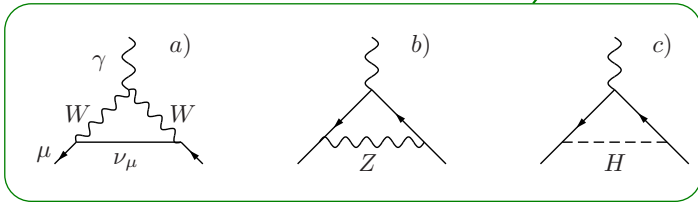
$a_\mu = \frac{g_\mu - 2}{2}$: minor correction from EW physics, but high precision

$a_\mu^{\text{exp}} = (1165920.59 \pm 0.22) \times 10^{-9}$

Muon g-2 '23

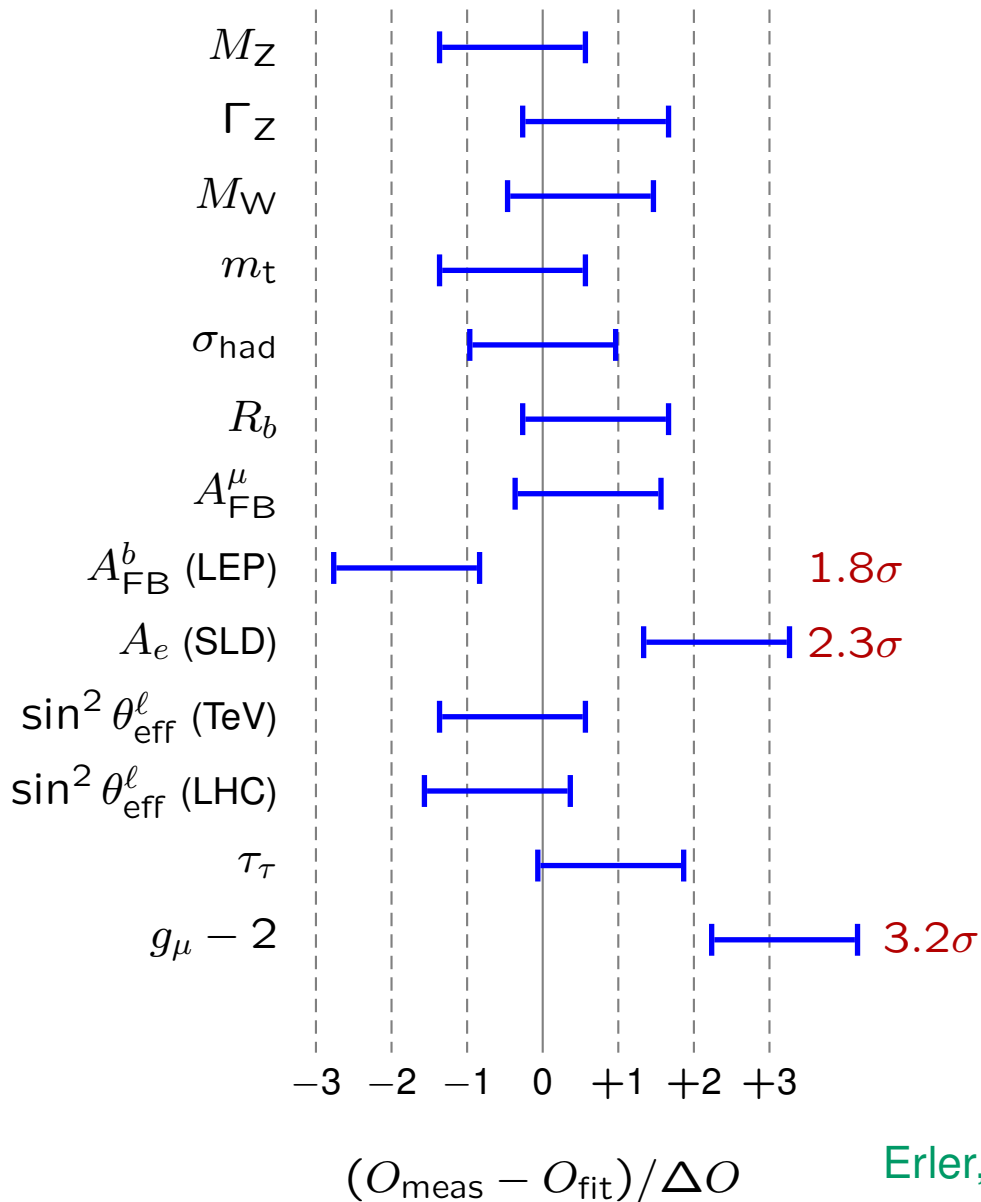
$a_\mu^{\text{th}} = (\underbrace{1165847.19}_{\text{QED}} + \underbrace{1.54}_{\text{EW}} + \underbrace{[65.58 \pm 0.25]}_{\text{HVP}} + \underbrace{[1.11 \pm 0.10]}_{\text{LBL}}) \times 10^{-9}$

→ talk by A. El Khadra



Discrepancies in evaluation of HVP contribution

BMW '24



Very good agreement:

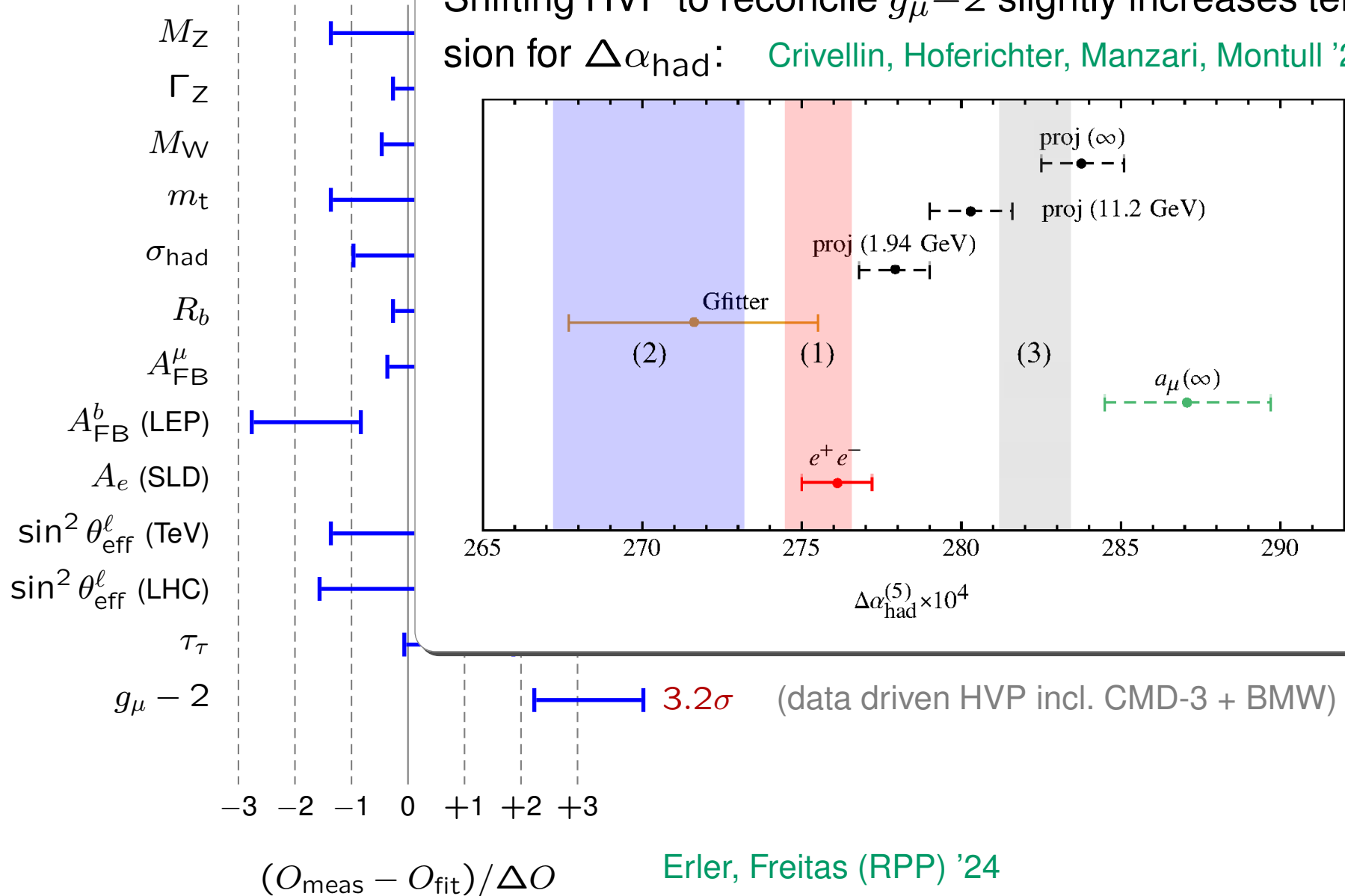
$$\chi^2/\text{d.o.f.} = 49.5/47 \quad (p = 37\%)$$

[without M_W from CDF II]

Most quantities measured with
1%–0.1% precision

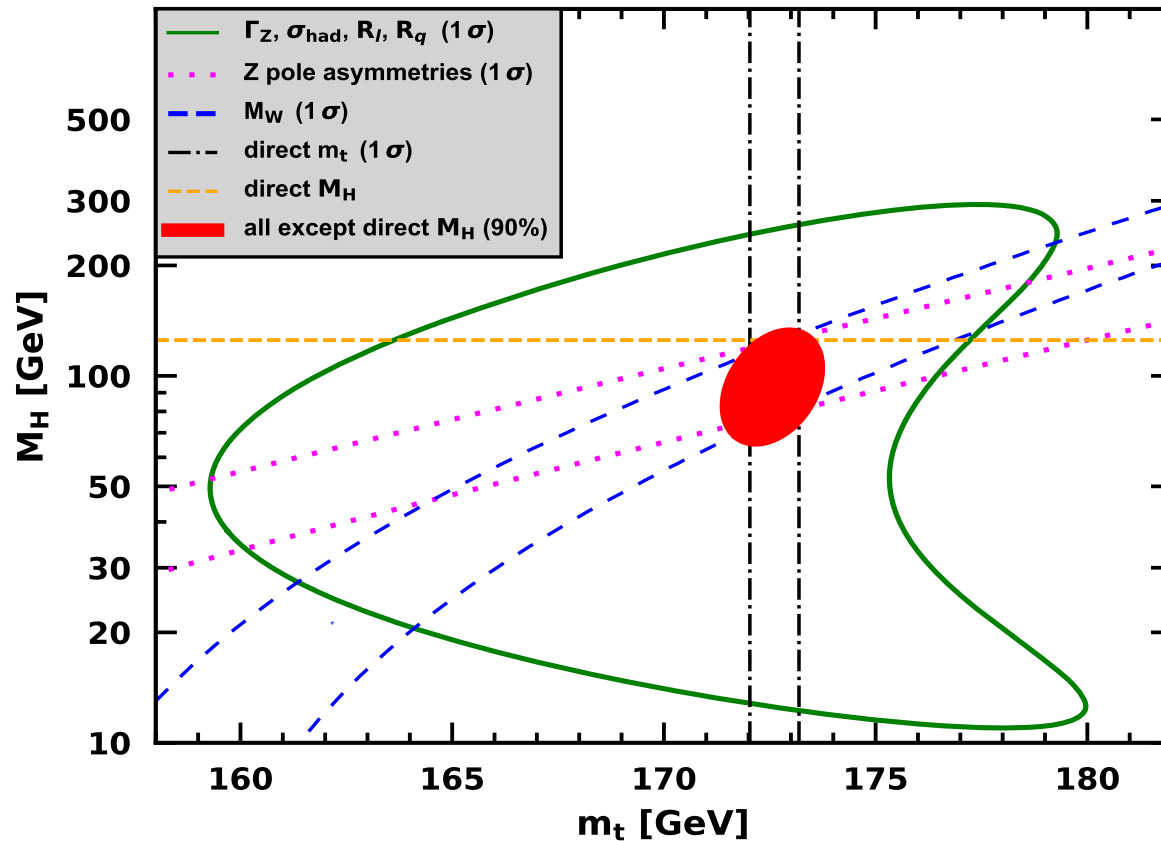
(data driven HVP incl. CMD-3 + BMW)

Shifting HVP to reconcile $g_\mu - 2$ slightly increases tension for $\Delta\alpha_{\text{had}}$: [Crivellin, Hoferichter, Manzari, Montull '20](#)



Erlar, Freitas (RPP) '24

Erler '24



Direct measurements:

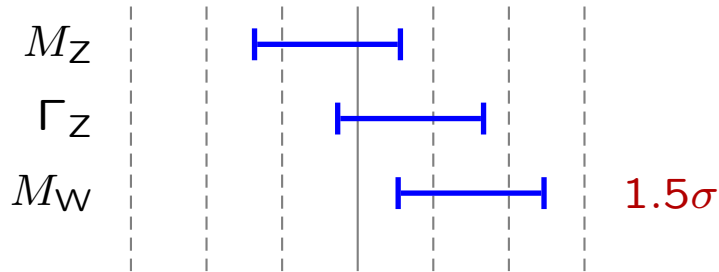
$$M_H = 125.10 \pm 0.09 \text{ GeV}$$

$$m_t = 172.61 \pm 0.58 \text{ GeV}$$

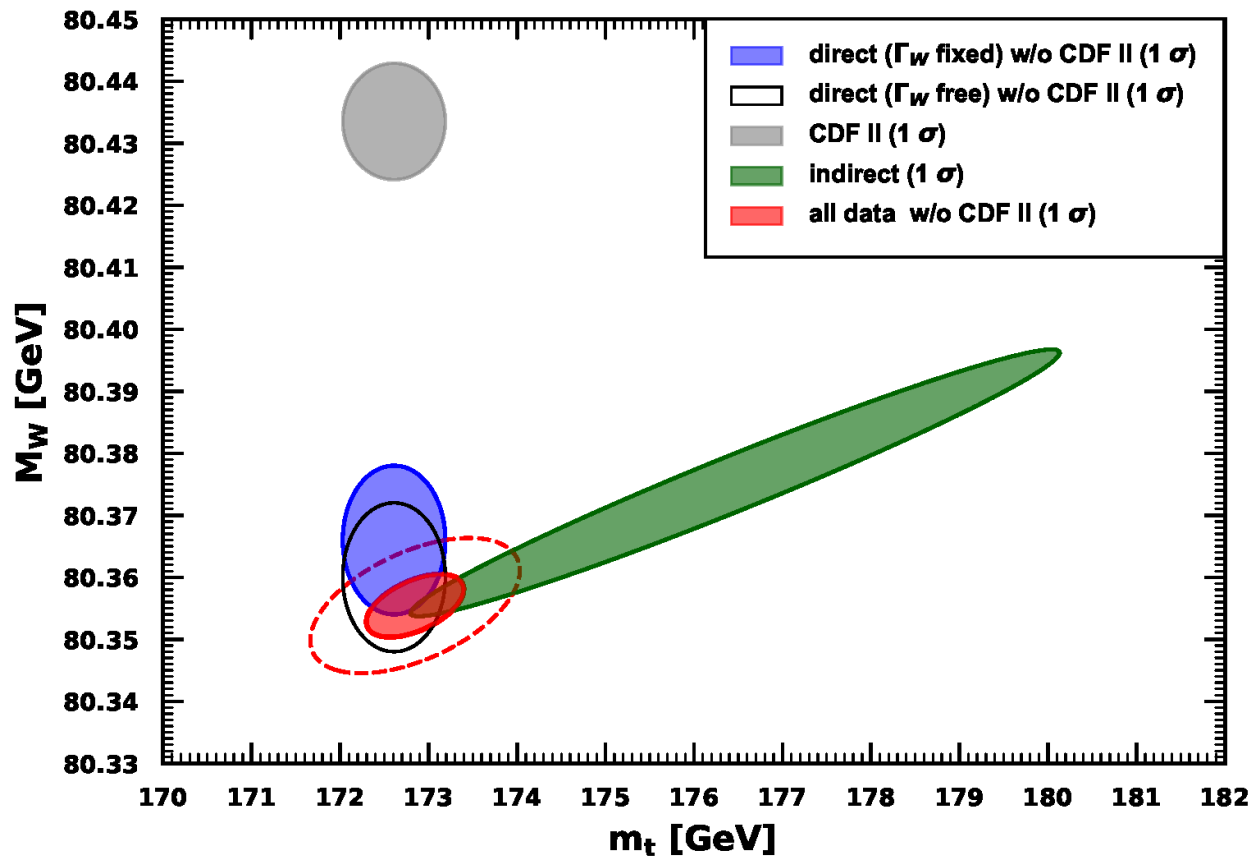
Indirect prediction:

$$M_H = 97^{+18}_{-16} \text{ GeV}$$

$$m_t = 175.2 \pm 1.8 \text{ GeV}$$



Including M_W from CDF II:
 Marginal agreement:
 $\chi^2/\text{d.o.f.} = 70.1/48$ ($p = 2\%$)



(plot without new CMS result)

Erlar '24

Adding BSM oblique parameters:

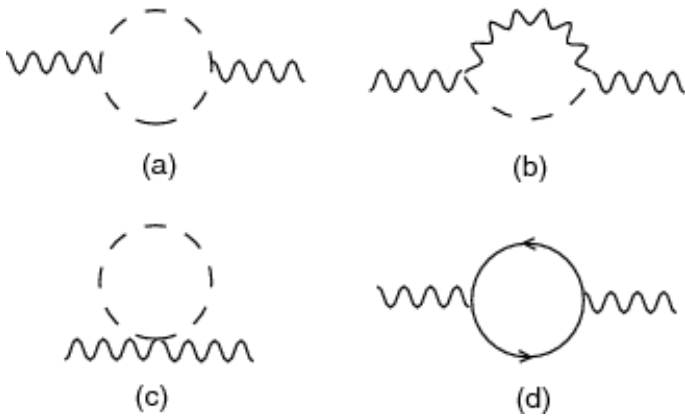
$$\alpha T = \frac{\Sigma_{WW}(0)}{M_W} - \frac{\Sigma_{ZZ}(0)}{M_Z}$$

→ $T = 0.08 \pm 0.02$

Including M_W from CDF II:

Good agreement:

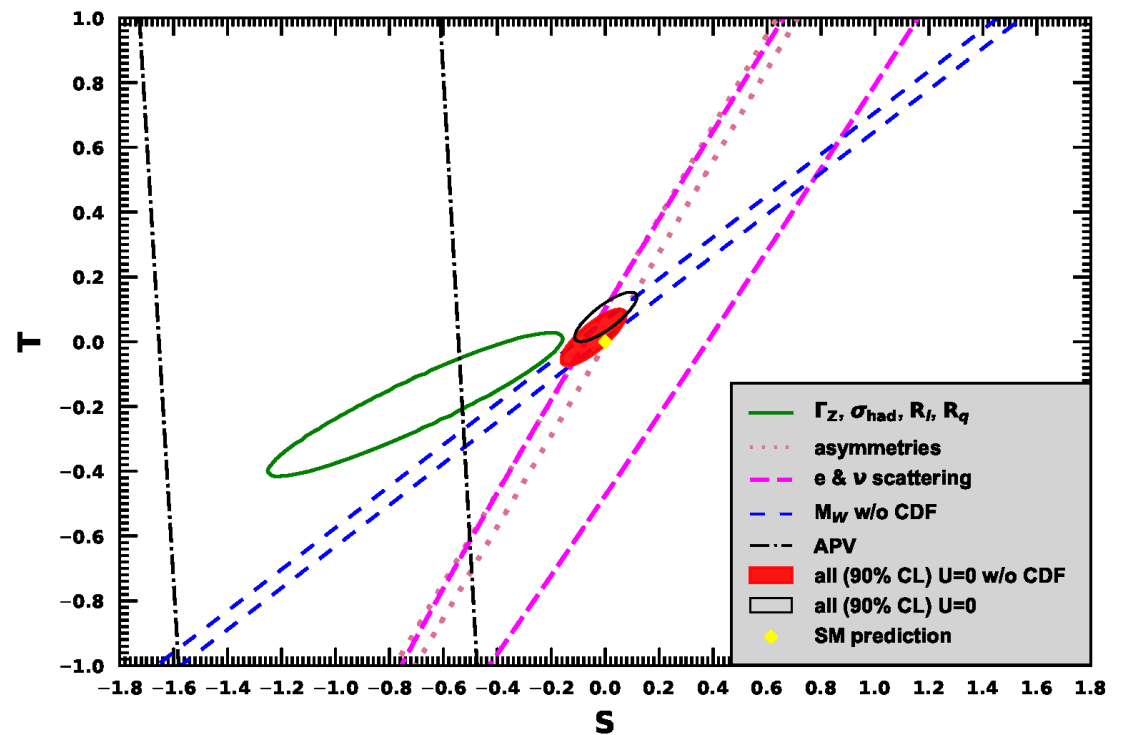
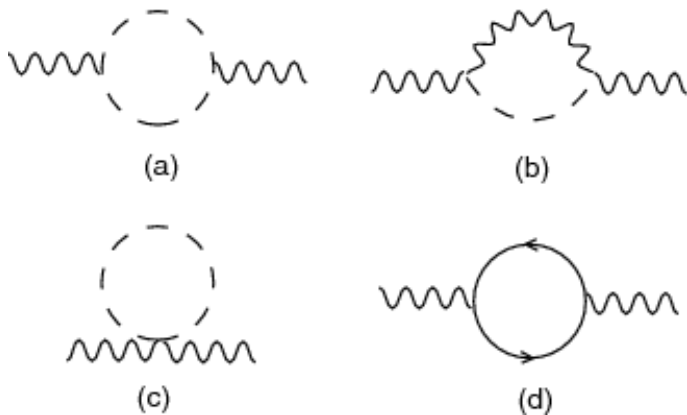
$$\chi^2/\text{d.o.f.} = 58.5/47 \quad (p = 12\%)$$



Adding BSM oblique parameters:

$$\alpha T = \frac{\Sigma_{WW}(0)}{M_W} - \frac{\Sigma_{ZZ}(0)}{M_Z}$$

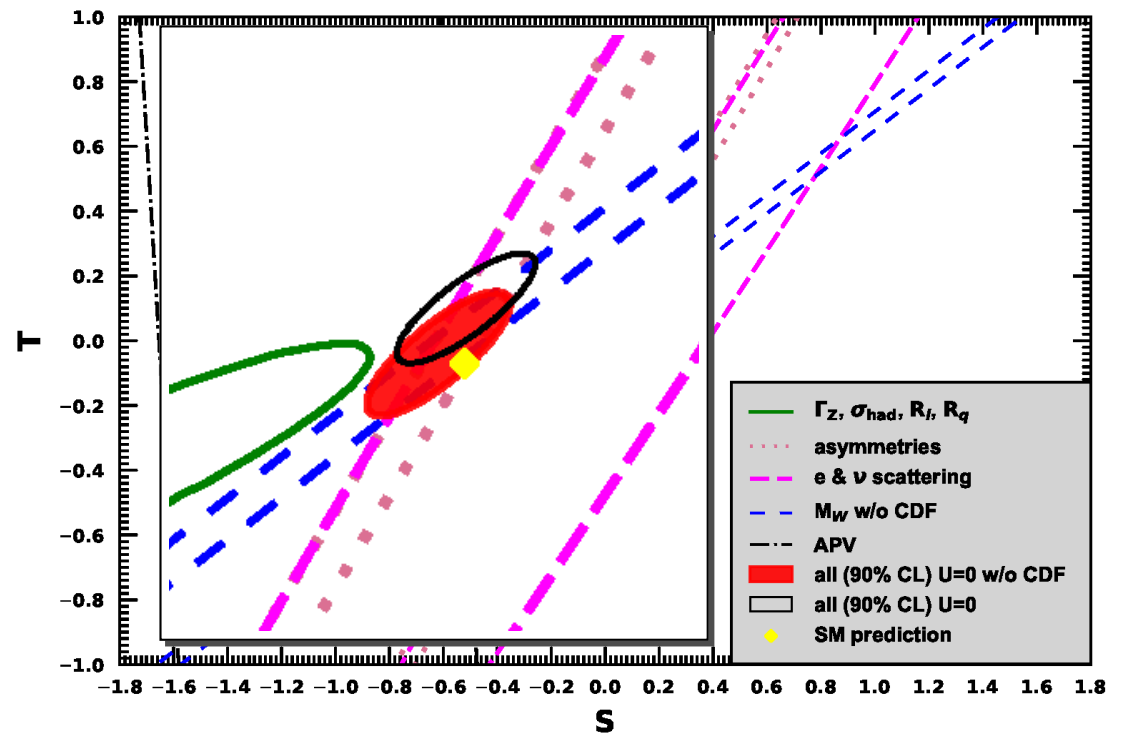
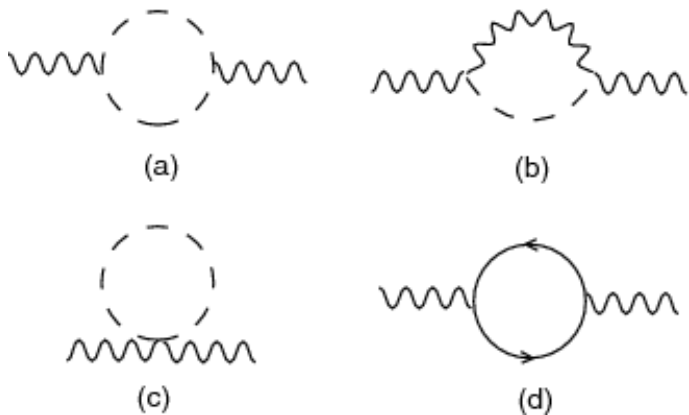
$$\frac{\alpha}{4s^2c^2} S = \frac{\Sigma_{ZZ}(M_Z^2) - \Sigma_{ZZ}(0)}{M_Z} + \frac{s^2 - c^2}{sc} \frac{\Sigma_{Z\gamma}(M_Z^2)}{M_Z} - \frac{\Sigma_{\gamma\gamma}(M_Z^2)}{M_Z}$$



Adding BSM oblique parameters:

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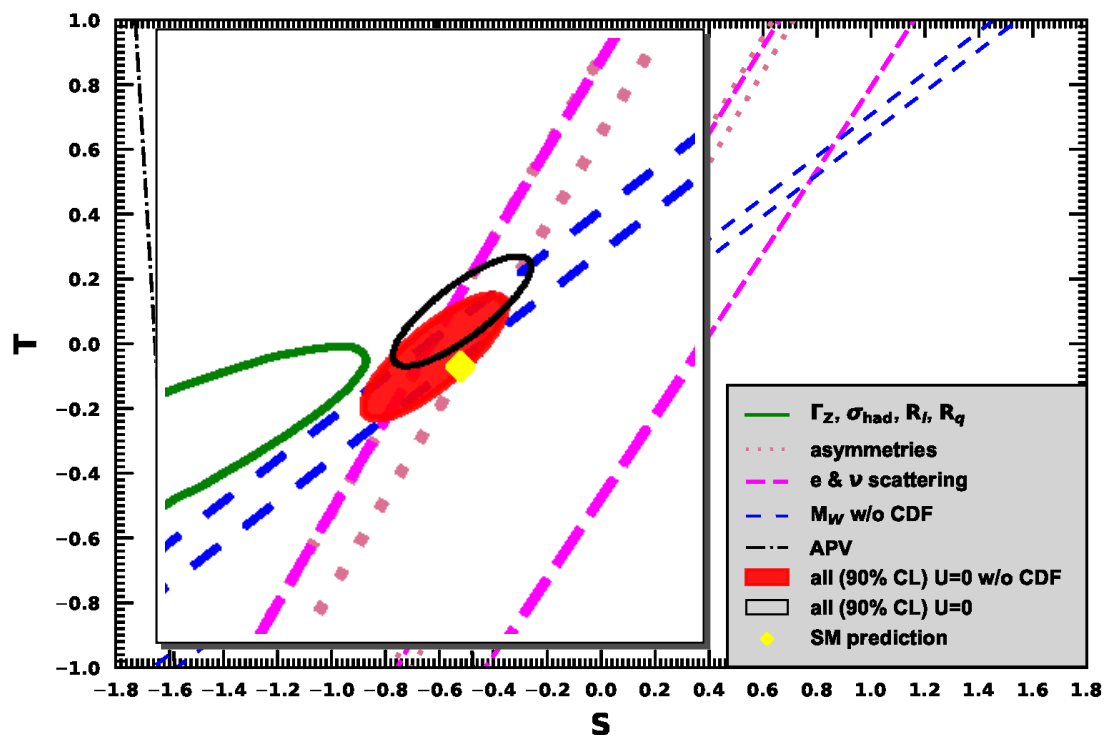
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Dim-6 SMEFT studies confirm that T parameter effectively absorbs M_W shift

de Blas, Pierini, Reina, Silvestrini '22
Bagnaschi et al. '22; Balkin et al. '22



Objective: Comparison of measurements for pseudo-obs. (M_W , $\sin^2 \theta_{\text{eff}}^\ell$, ...) with SM theory predictions

	δM_W [MeV]	$\delta \sin^2 \theta_{\text{eff}}^\ell$ [10^{-5}]
now	± 8	± 13
1-loop	± 450	± 1000
2-/3-loop QCD	± 70	± 45
ferm. 2-loop EW	± 50	± 90
bos. 2-loop EW	± 2	± 1
leading 3-loop	± 5	± 25

computed from G_μ

Experimental precision sensitive to 2-/3-loop effects

Marciano, Sirlin '80

Djouadi et al. '88

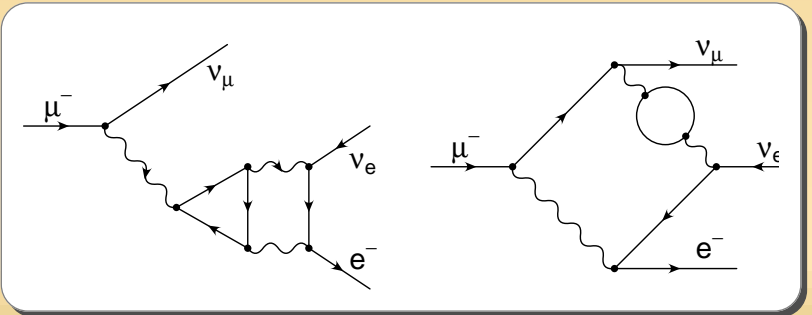
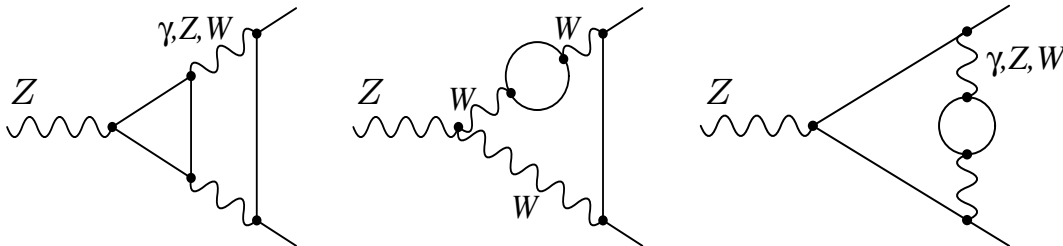
Chetyrkin, Kühn, Steinhauser '95

Freitas et al. '00 Awramik, Czakon '03

Awramik, Czakon, Freitas, Weiglein '04

Awramik, Czakon, Freitas '06

Faisst, Kühn, Seidensticker, Veretin '03



- Many seminal works on 1-loop and leading 2-loop corrections

Veltman, Passarino, Sirlin, Marciano, Bardin, Hollik, Riemann, Degrassi, Kniehl, ...

- Full 2-loop results for M_W , Z -pole observables

Freitas, Hollik, Walter, Weiglein '00

Hollik, Meier, Uccirati '05,07

Awramik, Czakon '02

Awramik, Czakon, Freitas, Kniehl '08

Onishchenko, Veretin '02

Freitas '14

Awramik, Czakon, Freitas, Weiglein '04

Dubovyk, Freitas, Gluza, Riemann, Usovitsch '16,18

Awramik, Czakon, Freitas '06

- Approximate 3- and 4-loop results (enhanced by Y_t and/or N_f)

Chetyrkin, Kühn, Steinhauser '95

Chetyrkin et al. '06

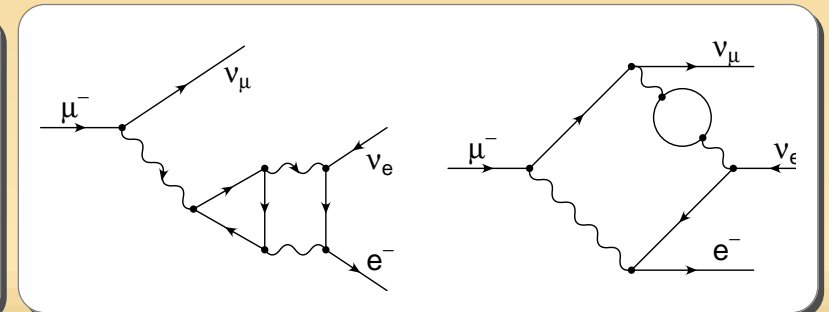
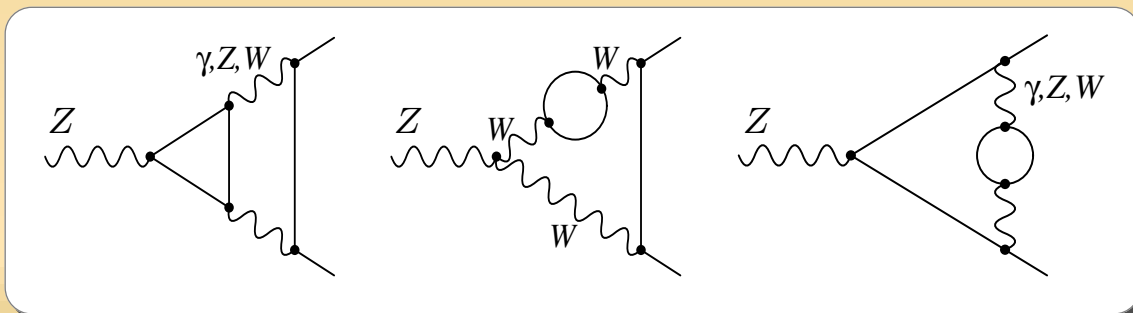
Faisst, Kühn, Seidensticker, Veretin '03

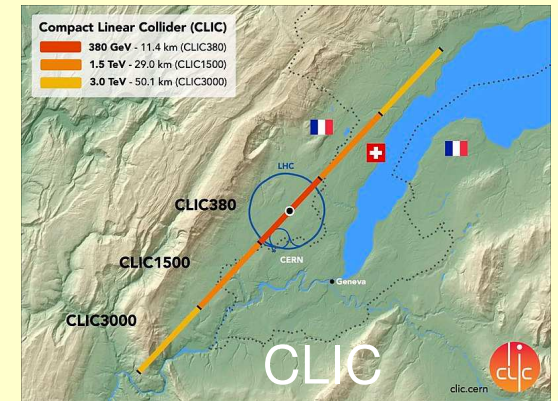
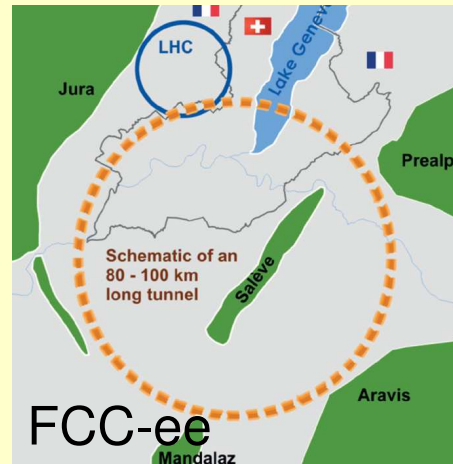
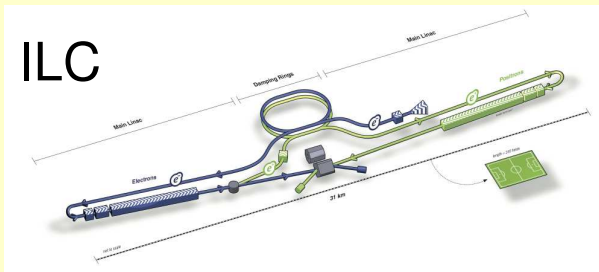
Boughezal, Czakon '06

Boughezal, Tausk, v. d. Bij '05

Chen, Freitas '20

Schröder, Steinhauser '05



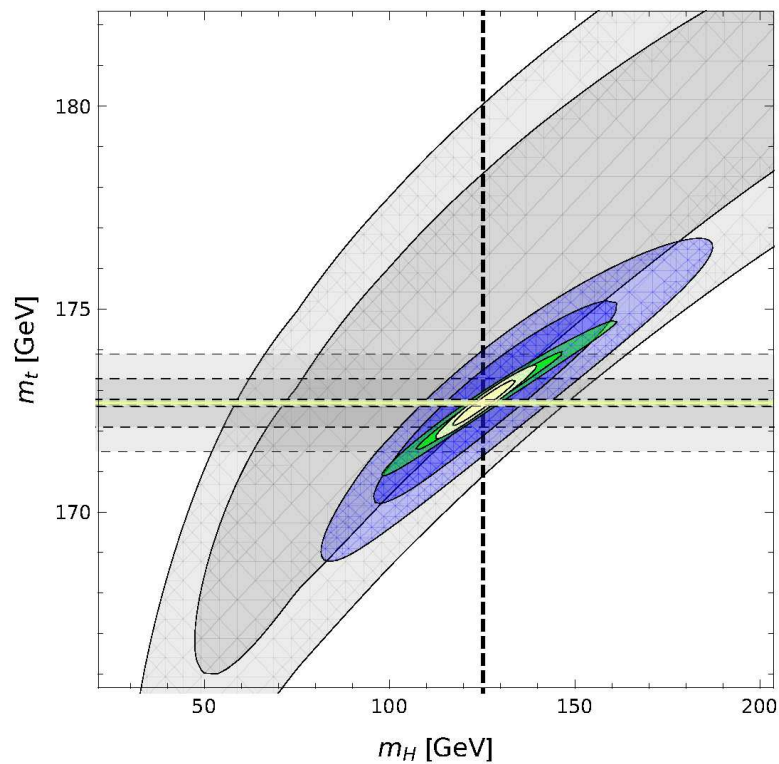


- circular colliders: high-lumi run at $\sqrt{s} \sim M_Z$
- linear colliders: radiative return $e^+e^- \rightarrow \gamma Z$

\sqrt{s}	M_Z	$2M_W$	240–250 GeV	350–380 GeV
ILC	100 fb^{-1}	500 fb^{-1}	2 ab^{-1}	200 fb^{-1} (10 pts.)
CLIC	—	—	—	1 ab^{-1}
FCC-ee	150 ab^{-1}	10 ab^{-1} (2 pts.)	5 ab^{-1}	1 ab^{-1} (8 pts.)
CEPC	100 ab^{-1}	6 ab^{-1} (3 pts.)	20 ab^{-1}	1 ab^{-1} ?

→ talks by J. Guimaraes da Costa, G. Bernardi

	Current exp.	ILC250	CEPC	FCC-ee
M_W [MeV]	8	2.4	0.5	0.4
Γ_Z [MeV]	2.3	1.5	0.025	0.025
$R_\ell = \Gamma_Z^{\text{had}}/\Gamma_Z^\ell$ [10^{-3}]	25	20	2	1
$R_b = \Gamma_Z^b/\Gamma_Z^{\text{had}}$ [10^{-5}]	66	23	4.3	6
$\sin^2 \theta_{\text{eff}}^\ell$ [10^{-5}]	13	2	0.3	0.4



Snowmass '21/22

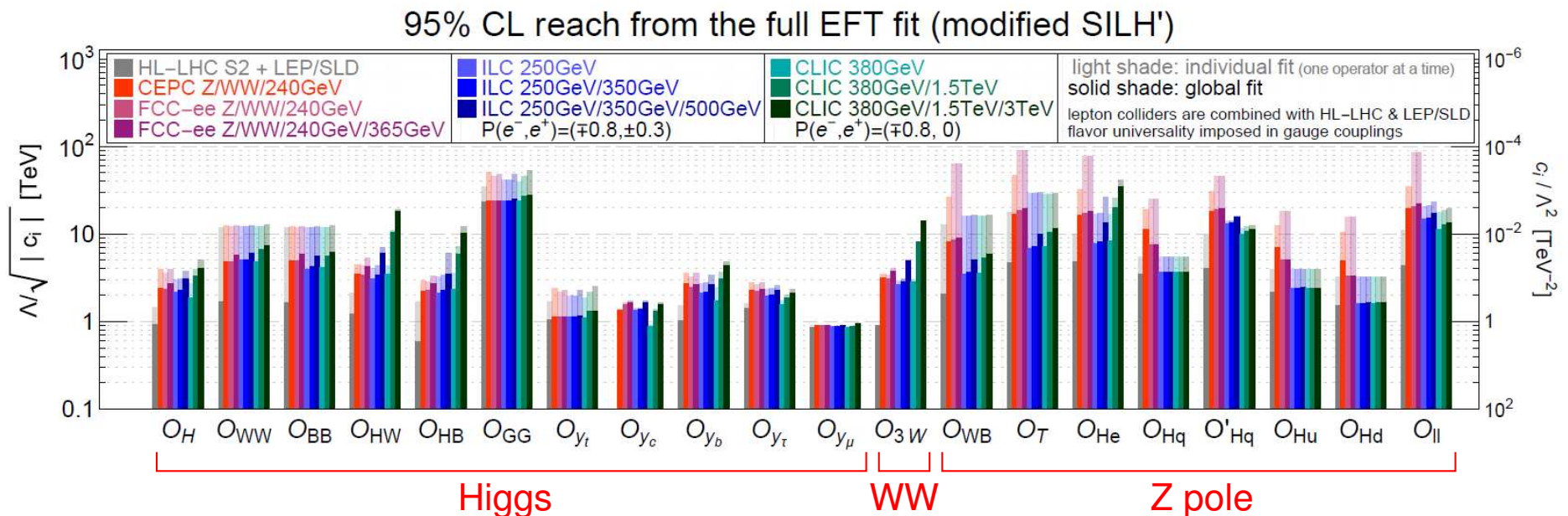
- Current
- ILC250 + ILC-GigaZ
- CEPC
- FCC-ee

- Extension of SM by **higher-dimensional operators**:

Wilson '69
Weinberg '79

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \frac{1}{\Lambda} \dots + \frac{1}{\Lambda^2} \sum_i c_i \mathcal{O}_i^{(6)} + \dots$$

- SMEFT dim-6 operators provide framework for comparing experiments



- To probe new physics, compare EWPOs with SM theory predictions
- Need to take theory error into account:

	Current exp.	Current th.	CEPC	FCC-ee
M_W [MeV]	11–12	4	0.5	0.4
Γ_Z [MeV]	2.3	0.4	0.025	0.025
$R_\ell = \Gamma_Z^{\text{had}} / \Gamma_Z^\ell$ [10^{-3}]	25	5	2	1
$R_b = \Gamma_Z^b / \Gamma_Z^{\text{had}}$ [10^{-5}]	66	10	4.3	6
$\sin^2 \theta_{\text{eff}}^\ell$ [10^{-5}]	13	4.5	0.3	0.4

- Theory error estimate is not well defined, ideally $\Delta_{\text{th}} \ll \Delta_{\text{exp}}$
- Common methods:
 - Count prefactors (α , N_c , N_f , ...)
 - Extrapolation of perturbative series
 - Renormalization scale dependence
 - Renormalization scheme dependence

■ Estimated impact of future higher-order calculations

Freitas et al. '19

	Current th.	Projected th. [†]	CEPC	FCC-ee
M_W [MeV]	4	1	0.5	0.4
Γ_Z [MeV]	0.4	0.15	0.025	0.025
$R_\ell = \Gamma_Z^{\text{had}} / \Gamma_Z^\ell$ [10^{-3}]	5	1.5	2	1
$R_b = \Gamma_Z^b / \Gamma_Z^{\text{had}}$ [10^{-5}]	10	5	4.3	6
$\sin^2 \theta_{\text{eff}}^\ell$ [10^{-5}]	4.5	1.5	0.3	0.4

[†] **Theory scenario:** $\mathcal{O}(\alpha\alpha_s^2)$, $\mathcal{O}(N_f\alpha^2\alpha_s)$, $\mathcal{O}(N_f^2\alpha^2\alpha_s)$, leading 4-loop
 ($N_f^n =$ at least n closed fermion loops)

Note: Estimates (based on extrapolation of perturb. series and prefactors)
 are unreliable and only provide a rough guess

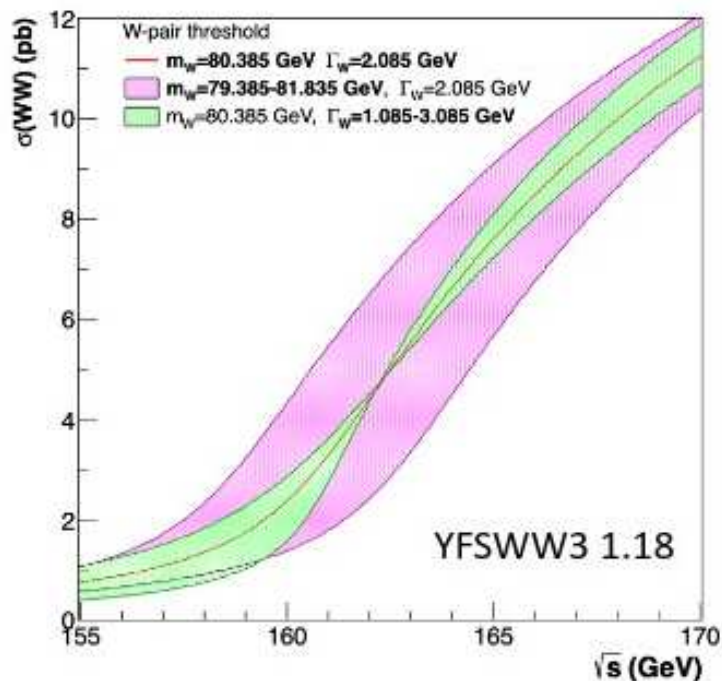
■ Also need NNLO corrections for subtracted “backgrounds”

■ More precise MC tools for multi-photon emission, hadronization, etc.

Jadach, Skrzypek '19

WW threshold : W mass and width

Scans of possible E_1 E_2 data taking energies and luminosity fractions f (at the E_2 point)



A - minimum of $\Delta\Gamma_W=0.91$ MeV with $\Delta m_W=0.55$ MeV
 taking data at $E_1=156.6$ GeV $E_2=162.4$ GeV $f=0.25$
 yields $\Delta m_W=0.47$ MeV (as single par)

B- minimum of $\Delta m_W=0.28$ MeV $\Delta\Gamma_W=3.3$ MeV with
 $E_1=155.5$ GeV $E_2=162.4$ GeV $f=0.95$
 yields $\Delta m_W=0.28$ MeV (as single par)

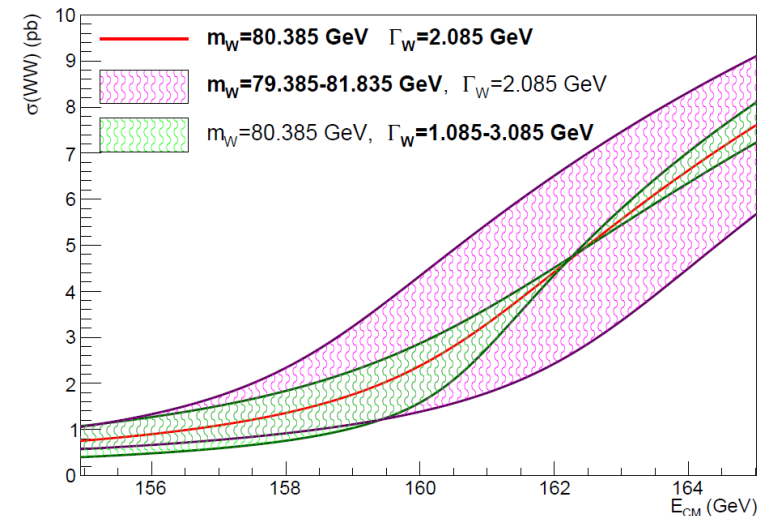
C- minimum of $\Delta\Gamma_W=0.96$ MeV + $\Delta m_W=0.41$ MeV with
 $E_1=157.5$ GeV $E_2=162.4$ GeV $f=0.45$
 yields and $\Delta m_W=0.37$ MeV (as single par)

$\Delta m_W, \Delta\Gamma_W$: error on W mass and width from fitting both
 Δm_W : error on W mass from fitting only m_W

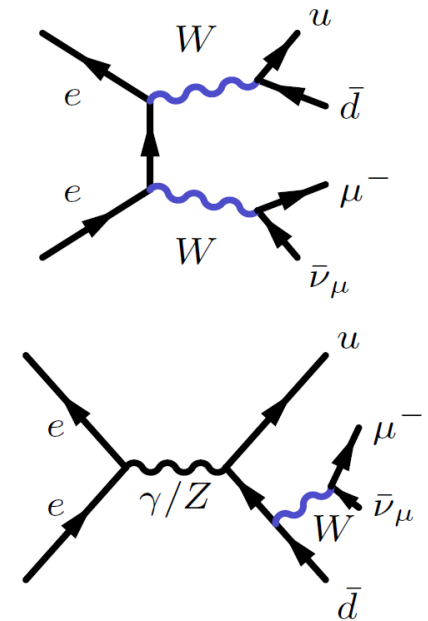
- a) Corrections near threshold enhanced by $1/\beta$ and $\ln \beta$

$$\beta \sim \sqrt{1 - 4 \frac{M_W^2 - iM_W \Gamma_W}{s}} \sim \sqrt{\Gamma_W / M_W}$$

- b) Non-resonant contributions are important



- Full $\mathcal{O}(\alpha)$ calculation of $e^+e^- \rightarrow 4f$
Denner, Dittmaier, Roth, Wieders '05
- EFT expansion in $\alpha \sim \Gamma_W / M_W \sim \beta^2$
Beneke, Falgari, Schwinn, Signer, Zanderighi '07
 - NLO corrections with NNLO Coulomb correction ($\propto 1/\beta^n$): $\delta_{\text{th}} M_W \sim 3 \text{ MeV}$
Actis, Beneke, Falgari, Schwinn '08
 - Adding NNLO corrections to $ee \rightarrow WW$ and $W \rightarrow f\bar{f}$ and NNLO ISR: $\delta_{\text{th}} M_W \lesssim 0.6 \text{ MeV}$



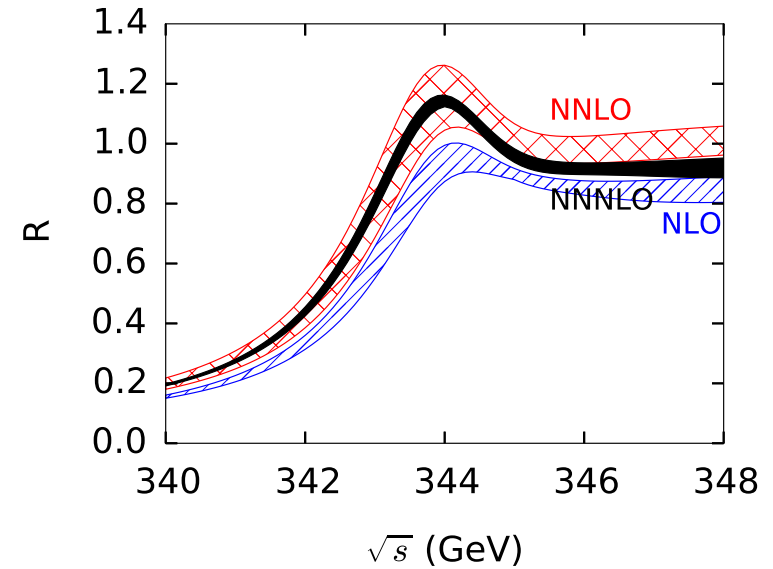
From $e^+e^- \rightarrow t\bar{t}$ at $\sqrt{s} \sim 350$ GeV:

Impact of theory modelling:

$$\delta m_t^{\overline{\text{MS}}} = [\]_{\text{exp}}$$

- $\oplus [50 \text{ MeV}]_{\text{QCD}}$
- $\oplus [10 \text{ MeV}]_{\text{mass def.}}$
- $\oplus [70 \text{ MeV}]_{\alpha_s}$

> 100 MeV



Beneke et al. '15

From $e^+e^- \rightarrow t\bar{t}$ at $\sqrt{s} \sim 350$ GeV:

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- $\oplus [10 \text{ MeV}]_{\text{mass def.}}$
- $\oplus [70 \text{ MeV}]_{\alpha_s}$

> 100 MeV

future improvements:

- [20 MeV]_{exp} (FCC-ee, CEPC)
- $\oplus [30 \text{ MeV}]_{\text{QCD}}$ (h.o. resumm., N⁴LO?)
- $\oplus [10 \text{ MeV}]_{\text{mass def.}}$
- $\oplus [15 \text{ MeV}]_{\alpha_s}$ ($\delta\alpha_s \lesssim 0.0002$)

$\lesssim 50 \text{ MeV}$

- α_s :

- Electroweak precision ($R_\ell = \Gamma_Z^{\text{had}} / \Gamma_Z^\ell$):

$$\alpha_s = 0.122 \pm 0.003$$

PDG '18

→ Negligible non-perturbative QCD effects

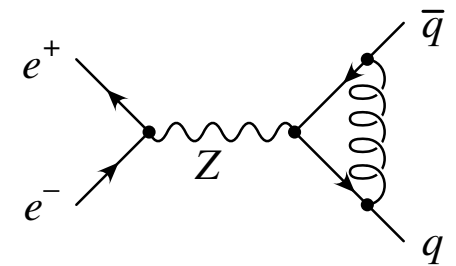
$$\text{FCC-ee: } \delta R_\ell \sim 0.001$$

$$\Rightarrow \delta \alpha_s < 0.0001$$

Theory input: $N^3\text{LO EW corr.} + \text{leading } N^4\text{LO}$

to keep $\delta_{\text{th}} R_\ell \lesssim \delta_{\text{exp}} R_\ell$

Caviat: R_ℓ could be affected by new physics



- α_s :

- Electroweak precision ($R_\ell = \Gamma_Z^{\text{had}} / \Gamma_Z^\ell$):

$$\alpha_s = 0.122 \pm 0.003$$

PDG '18

→ Negligible non-perturbative QCD effects

$$\text{FCC-ee: } \delta R_\ell \sim 0.001$$

$$\Rightarrow \delta \alpha_s < 0.0001$$

Caviat: R_ℓ could be affected by new physics

- $R = \frac{\sigma[ee \rightarrow \text{had.}]}{\sigma[ee \rightarrow \mu\mu]}$ at lower \sqrt{s}

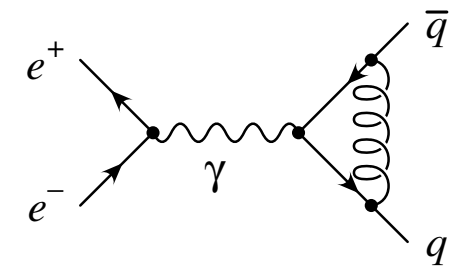
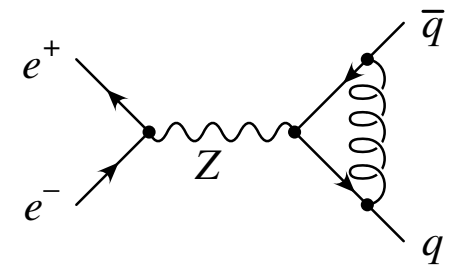
$$\text{e.g. CLEO } (\sqrt{s} \sim 9 \text{ GeV}): \alpha_s = 0.110 \pm 0.015$$

Kühn, Steinhauser, Teubner '07

→ dominated by s -channel photon, less room for new physics

→ QCD still perturbative

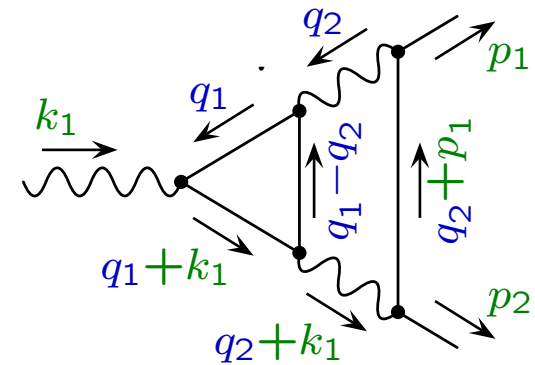
$$\text{naive scaling to } 50 \text{ ab}^{-1} \text{ (BELLE-II): } \delta \alpha_s \sim 0.0001$$



Experimental precision requires inclusion of **multi-loop corrections** in theory

Integrals over loop momenta:

$$\int d^4 q_1 d^4 q_2 f(q_1, q_2, p_1, k_1, \dots, m_1, m_2, \dots)$$



Challenges:

1. $\mathcal{O}(1000)$ – $\mathcal{O}(10000)$ integrals
2. Individual integrals can be divergent (drop out for physical results)
→ Regularization, renormalization
3. Multi-dimensional integrations, depending on multiple mass/momentum scales

General approaches:

- Analytical
- Numerical
- Approximations (expansions), specialized techniques, ...

Analytical techniques:

- Reduction to master integrals (MIs), reduced number of ints. by 10-100
 - computationally intensive
 - public programs `Reduze`, `FIRE`, `LiteRed`, `KIRA`, `FiniteFlow`, ...
von Manteuffel, Studerus '12; Smirnov '13,14; Lee '13; Maierhoefer, Usovitsch, Uwer '17; Peraro '19
- Not fully understood function space of MIs
(Goncharov polylogs, iterated elliptic integrals, hypergeometric functions, ...)
- Work best for problems with few (no) masses

Numerical techniques:

- Multi-dim. numerical integrations:
 - in momentum space: $4L$ dimensions ($L = \#$ of loops)
 - in Feynman par. space: $P - 1$ dimensions ($P = \#$ of propagators)
 - slowly converging, limited precision
- Numerical instabilities, in particular for diagrams with physical cuts
- Works best for problems with many masses

New techniques, e.g.:

- Numerical reduction to MIs, numerical MIs via differential equations (DEs)
Mandal, Zhao '18, Czakon, Niggetiedt '20
- DEs with respect to auxiliary parameter, $\frac{1}{k_i^2 - m_i^2 + i\epsilon}$
Liu, Ma, Wang '17
Liu, Ma '18,21,22
- Series solutions of DEs
Moriello '19, Hidding '20
- Dispersion relations + Feynman parameters
Song, Freitas '21, 22

- **Electroweak precision tests** have played an important role in testing the Standard Model
- Today they probe physics beyond the Standard Model at **TeV scale**
- More data from **LHC** and **future e^+e^- colliders** will push the reach into the **multi-TeV** regime
- **Electroweak fits** rely on detailed theory calculations for QED effects, backgrounds, SM predictions, etc.
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