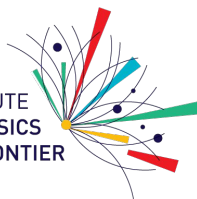


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MILLENNIUM INSTITUTE
FOR SUBATOMIC PHYSICS
AT HIGH-ENERGY FRONTIER
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FACULTAD DE FÍSICA

Long-lived particles

Giovanna Cottin

Pontificia Universidad Católica de Chile
& SAPHIR Millennium Institute

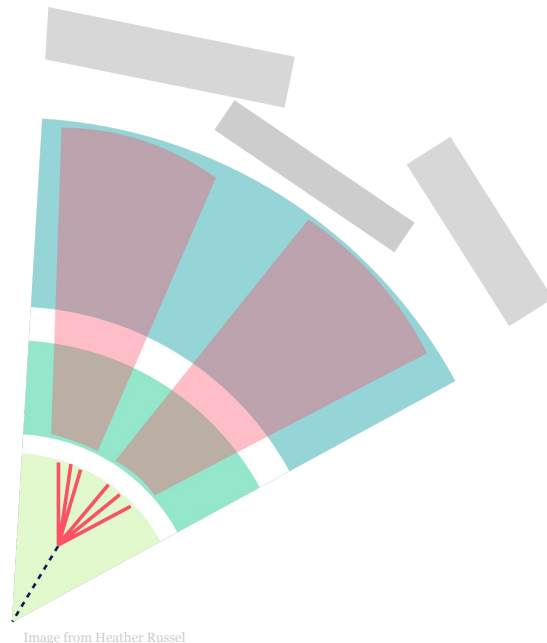


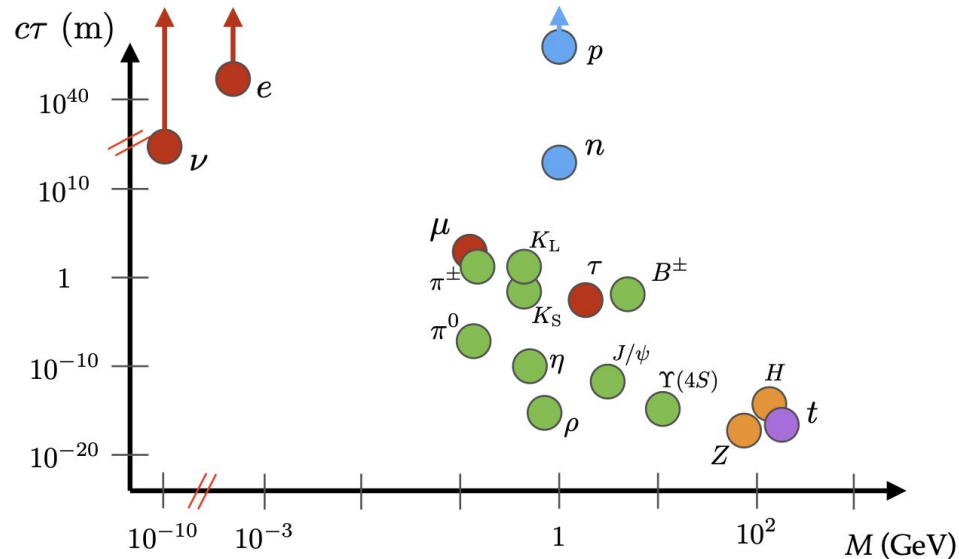
Image from Heather Russel

*XV Latin American Symposium on High Energy Physics (SILAFAE)
Mexico City, Cinvestav, November 4th 2024*

- What is a long-lived particle?
- Why look for long-lived particles?
- What's the long-lived particle phenomenology @ LHC?
- Why should we keep looking for long-lived particles?

● What is a long-lived particle?

Long-lived particles (LLP) can travel macroscopic distances before decaying inside a particle detector. Our world is full of them



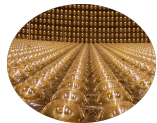
Theory motivation: their presence comes from conserved symmetries, small couplings, heavy mediators/hierarchy of mass scales, small phase space. Why a BSM sector shouldn't have some of these?

LLP = “**BSM particle** that dies (gives up all its energy or decays to SM) somewhere in the detector acceptance” J. Beacham

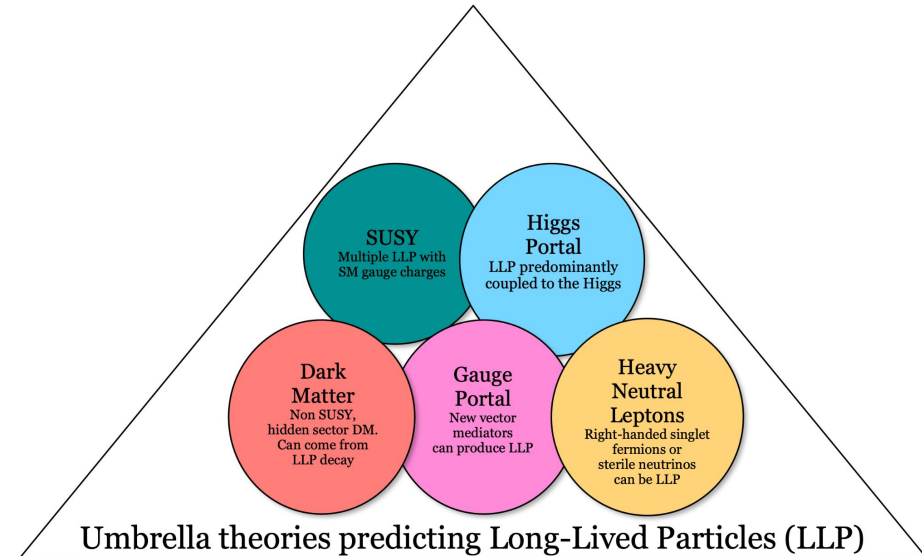
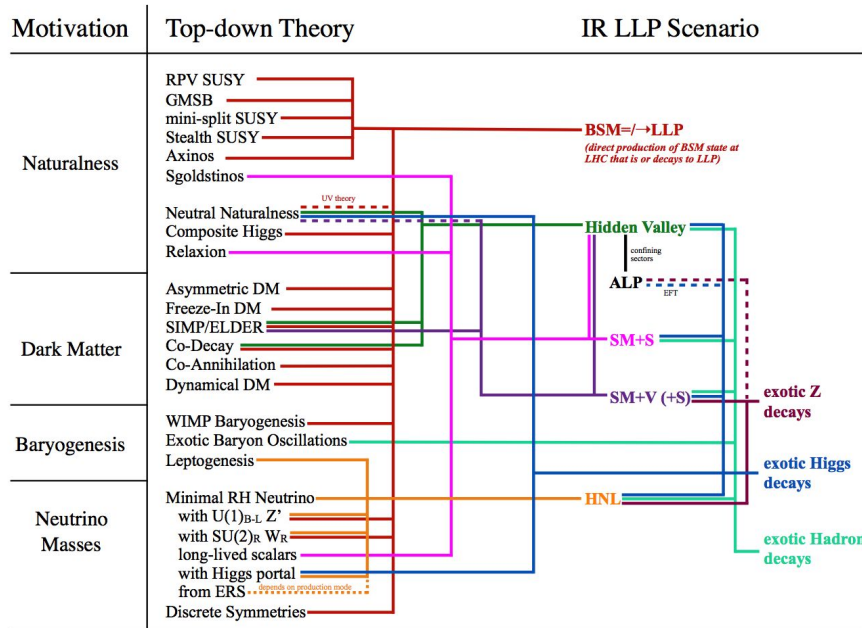
$$c\tau \sim \Gamma^{-1} \gtrsim 0.001 \text{ [mm]}$$

- What is a long-lived particle?
- Why look for long-lived particles?
- What's the long-lived particle phenomenology @ LHC?
- Why should we keep looking for long-lived particles?

LLPs are predicted in many BSM models



baryons antibaryons



No signs of new physics @ LHC

New physics may be so *feebly* coupled to our Standard Model that their signatures may have been overlooked or miss identified by LHC searches not dedicated to LLPs

LLP

$$c\tau \sim \Gamma^{-1} \gtrsim 0.001 \text{ [mm]}$$

$$\Gamma \sim \underbrace{c^2}_{\text{Feebly (small) couplings}} \underbrace{\left(\frac{\Delta m}{\Lambda}\right)^n}_{\text{Large mass hierarchies/ heavy mediators}} \underbrace{\Delta m}_{\text{Small mass difference or "compressed spectra"}}$$

Three reasons why

Three reasons why is hard

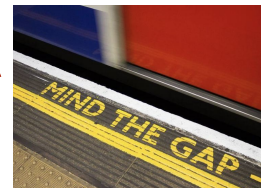
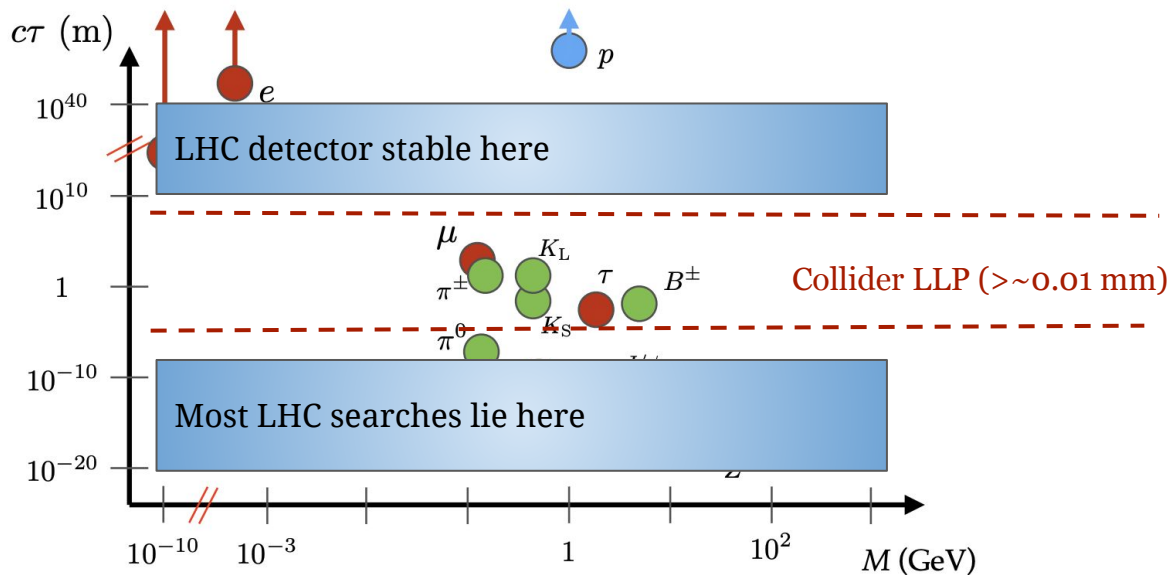
Low rates

Large energies
(LHC inaccessible)

Low efficiency (soft
particles/limited object
reconstruction)

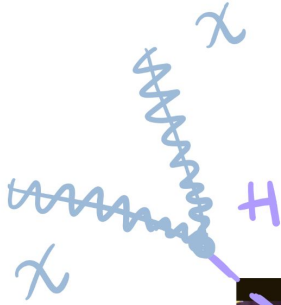
No signs of new physics @ LHC

Are we looking deep in all regions of parameter space?

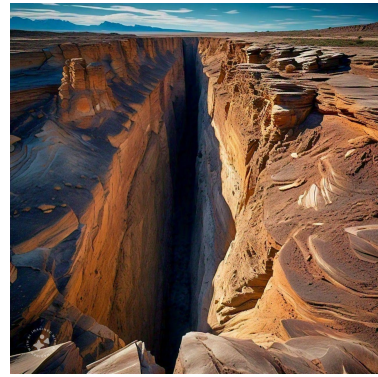
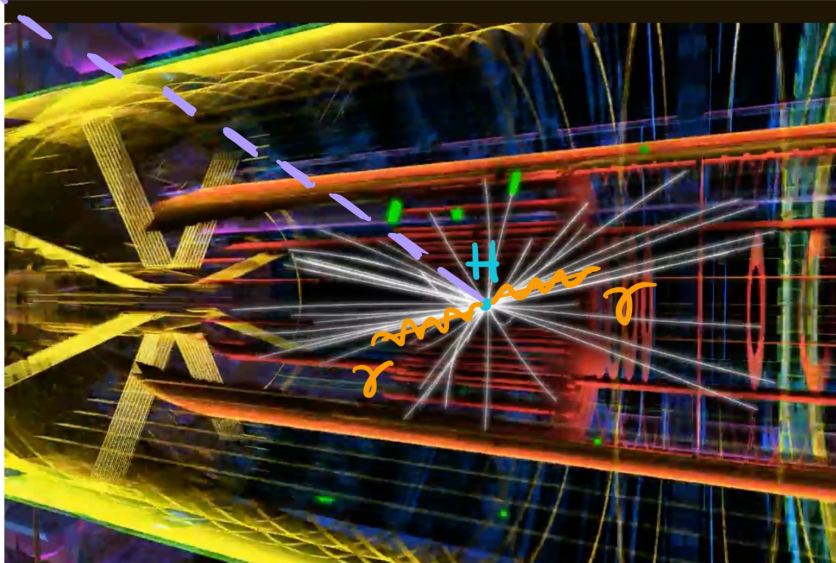


17/201 ATLAS LLP searches (13 TeV results using the full 2015 and 2016 data-set)

No signs of new physics @ LHC



Perhaps new physics is escaping completely !



“A huge gap” by Meta AI

Many LLP efforts worldwide !

LHC-LLP WG: <https://lpc.web.cern.ch/lhc-llp-wg>

LLP Community Workshops: <https://indico.cern.ch/event/1381368/>

LLP Community White Paper: [arXiv:1903.04497](https://arxiv.org/abs/1903.04497)

LLP repo: <https://github.com/llprecasting/recastingCodes>

RAMP seminars, including LLP searches: <https://indico.cern.ch/category/14155/>

CERNCOURIER

Reporting on international
high-energy physics

Physics ▾ Technology ▾ Community ▾ In focus Magazine
SEARCHES FOR NEW PHYSICS | MEETING REPORT

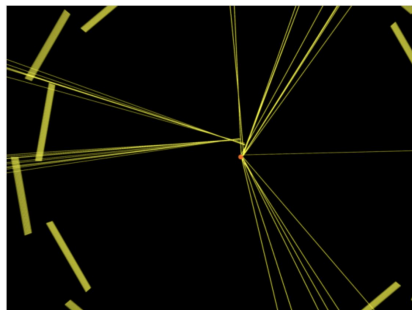
Long-lived particles gather interest

21 July 2021

SEARCHES FOR NEW PHYSICS | MEETING REPORT

A long-lived paradigm shift

27 November 2020

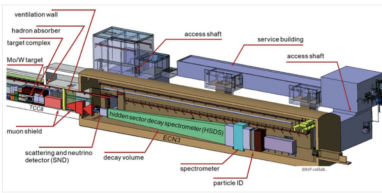


Displaced vertex A simulated CMS collision where a long-lived particle travels a short distance before it decays. Credit: CMS-PH-0-2020-002-5

SEARCHES FOR NEW PHYSICS | NEWS

SHIP to chart hidden sector

3 May 2024

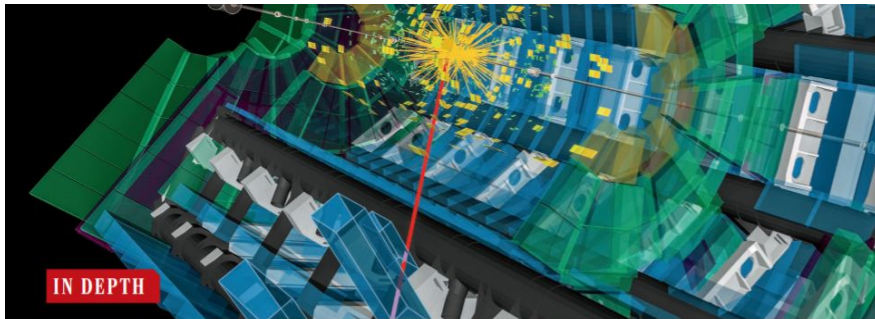
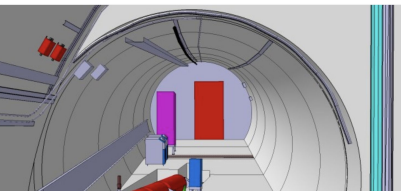


Full speed ahead Layout of the SHIP experiment, with the target on the left and the experiment in the ECN3 hall. Credit: SHIP collab.

FASE: CERN approves new experiment to look for long-lived, exotic particles

The experiment, which will complement existing searches for dark matter at the LHC, will be operational in 2021

7 MARCH, 2019 | By Cristina Agrigoroese



In a simulated event, the track of a decay particle called a muon (red), displaced slightly from the center of particle collisions, could be a sign of new physics.

PARTICLE PHYSICS

A hunt for long-lived particles ramps up

The Large Hadron Collider could be making new particles that are hiding in plain sight

By Adrian Cho

Are new particles materializing right under physicists' noses and going unnoticed? The world's great atom smasher, the Large Hadron Collider (LHC), could be making long-lived particles that slip through its detectors, some researchers say. Next week, they will gather at the LHC's home, CERN, the European particle physics laboratory near Geneva, Switzerland, to discuss how to capture them. They argue the LHC's next run should emphasize such searches, and some are calling for new detectors that could sniff out the fugitive particles.

It's a push born of anxiety. In 2012, experiments at the \$5 billion LHC discovered

simple strategy to look for new particles: Smash together protons or electrons at ever-higher energies to produce heavy new particles and watch them decay instantly into lighter, familiar particles within the huge, barrel-shaped detectors. That's how CMS and its rival detector, A Toroidal LHC Apparatus (ATLAS), spotted the Higgs, which in a trillionth of a nanosecond can decay into, among other things, a pair of photons or two "jets" of lighter particles.

Long-lived particles, however, would zip through part or all of the detector before decaying. That idea is more than a shot in the dark, says Giovanna Cottin, a theorist at National Taiwan University in Taipei. "Almost all the frameworks for beyond-the-standard-model physics predict

of subsystems—trackers that trace charged particles, calorimeters that measure particle energies, and chambers that detect penetrating and particularly handy particles called muons—all arrayed around a central point where the accelerator's proton beams collide. Particles that fly even a few millimeters before decaying would leave unusual signatures: kinked or offset tracks, or jets that emerge gradually instead of all at once.

Standard data analysis often assumes such oddities are mistakes and junk, notes Tova Holmes, an ATLAS member from the University of Chicago in Illinois who is searching for the displaced tracks of decays from long-lived supersymmetric particles. "It's a bit of a challenge because the way we've designed things, and the software



LLP2024: Fourteenth workshop of the Long-Lived Particle Community

Adding your recasting code

This is an open repository and if you have developed a code for recasting a LLP analysis, we encourage you to include it here. Please contact llp-recasting@googlegroups.com and we will provide you with the necessary information for including your code.

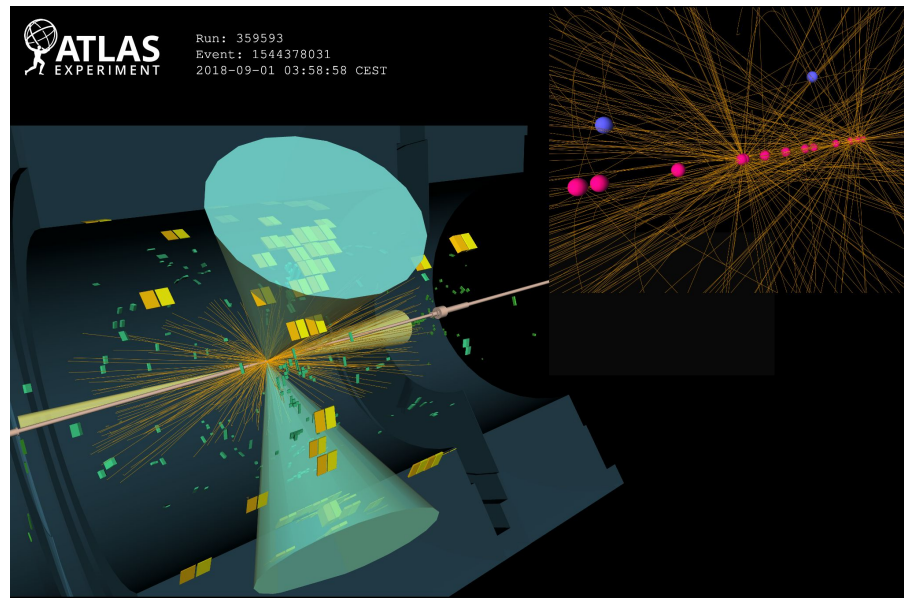
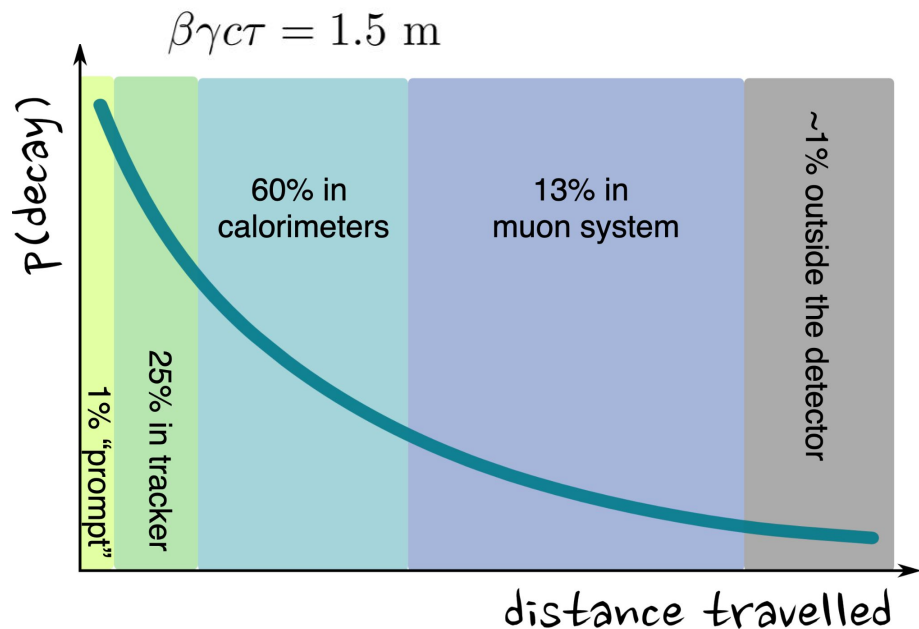
Repository Structure

The repository folder structure is organized according to the type of LLP signature and the corresponding analysis and authors:

- Displaced Vertices
 - 13 TeV ATLAS Displaced Vertex plus MET by ALESSA
 - 13 TeV ATLAS Displaced Vertex plus MET by GCottin
 - 8 TeV ATLAS Displaced Vertex plus jets by GCottin
- Heavy Stable Charged Particles
 - 8 TeV CMS HSCP
 - 13 TeV ATLAS HSCP
- Disappearing Tracks
- Displaced Jets

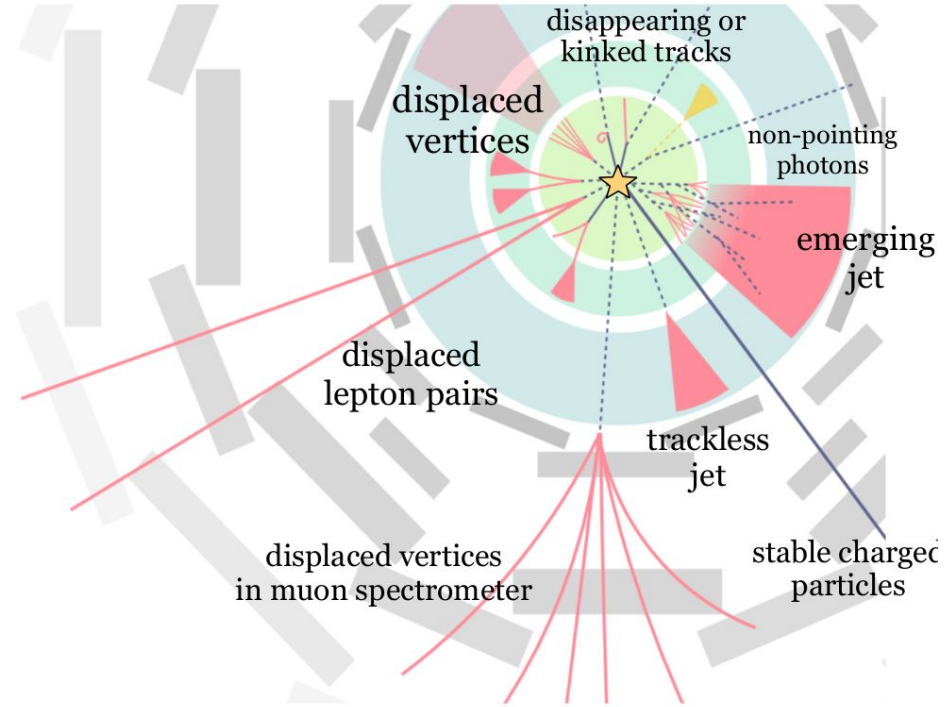
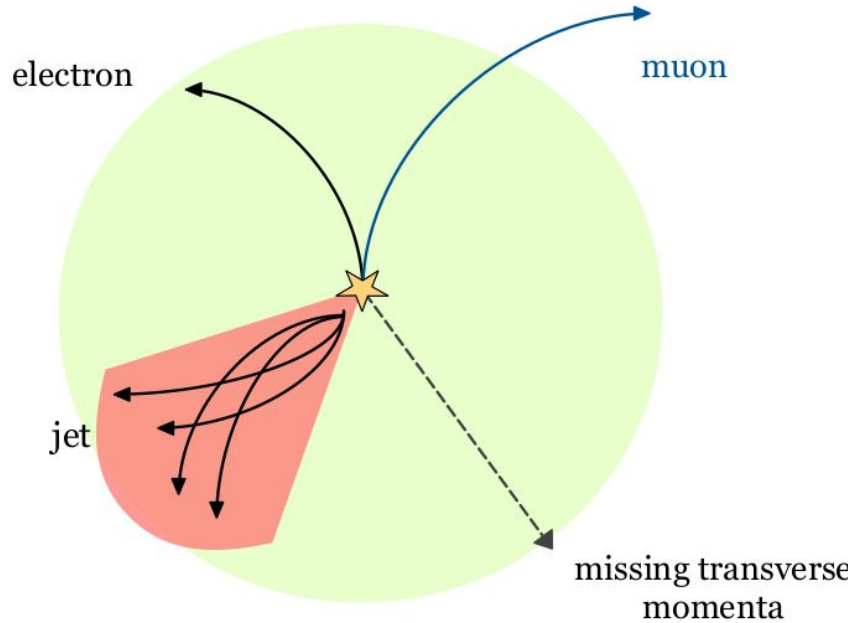
- What is a long-lived particle?
- Why look for long-lived particles?
- What's the long-lived particle phenomenology @ LHC?
- Why should we keep looking for long-lived particles?

LHC-LLP searches can target different lifetimes using different parts of the detectors. Detection usually requires special triggers and reconstruction



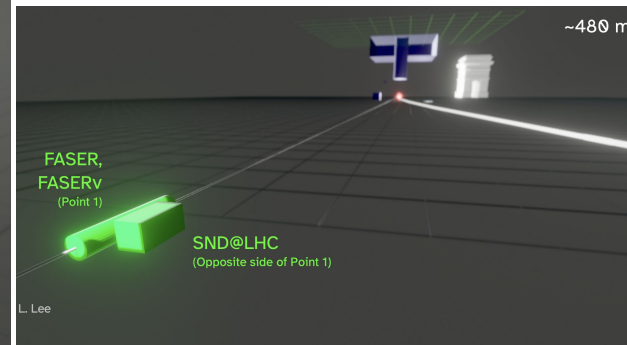
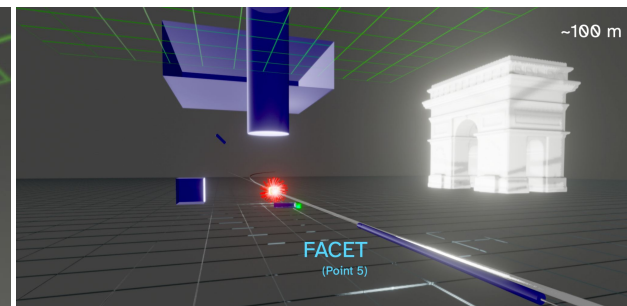
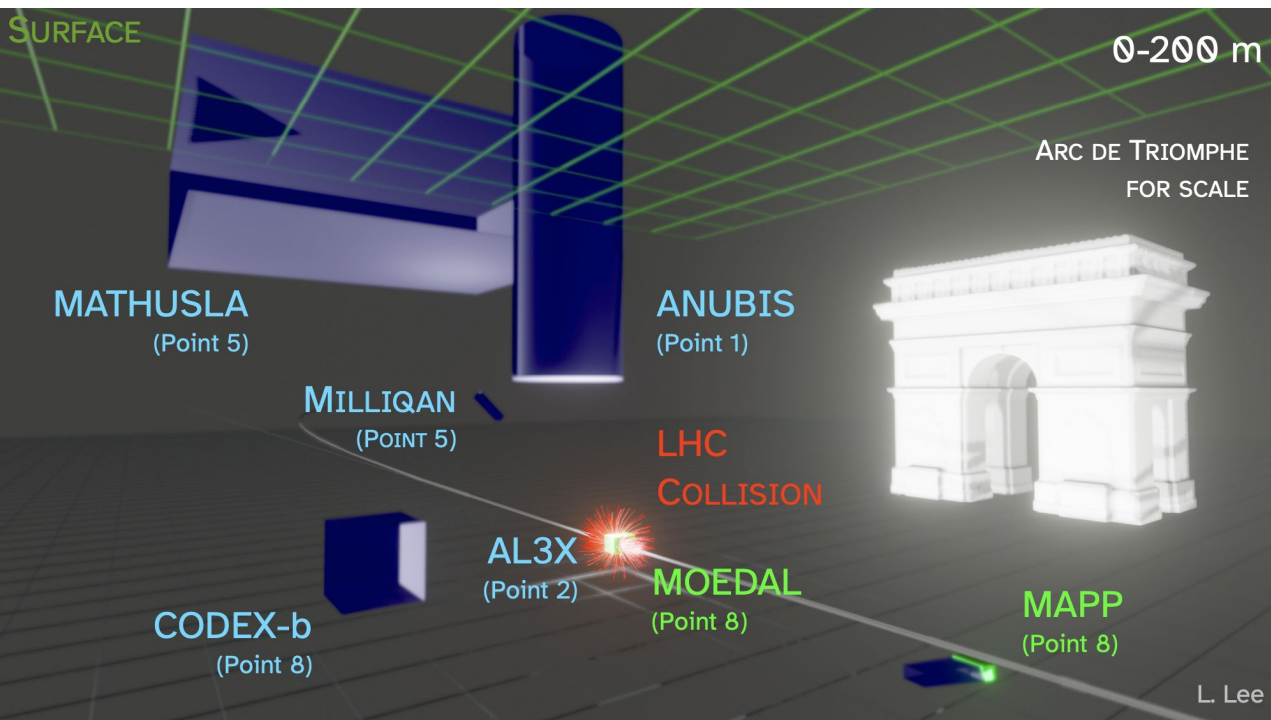
[ATLAS Event Display with Displaced Vertices](#)

Depending on where the LLP decays, it's mass and which quantum numbers it has, this will give rise to different exotic collider signatures

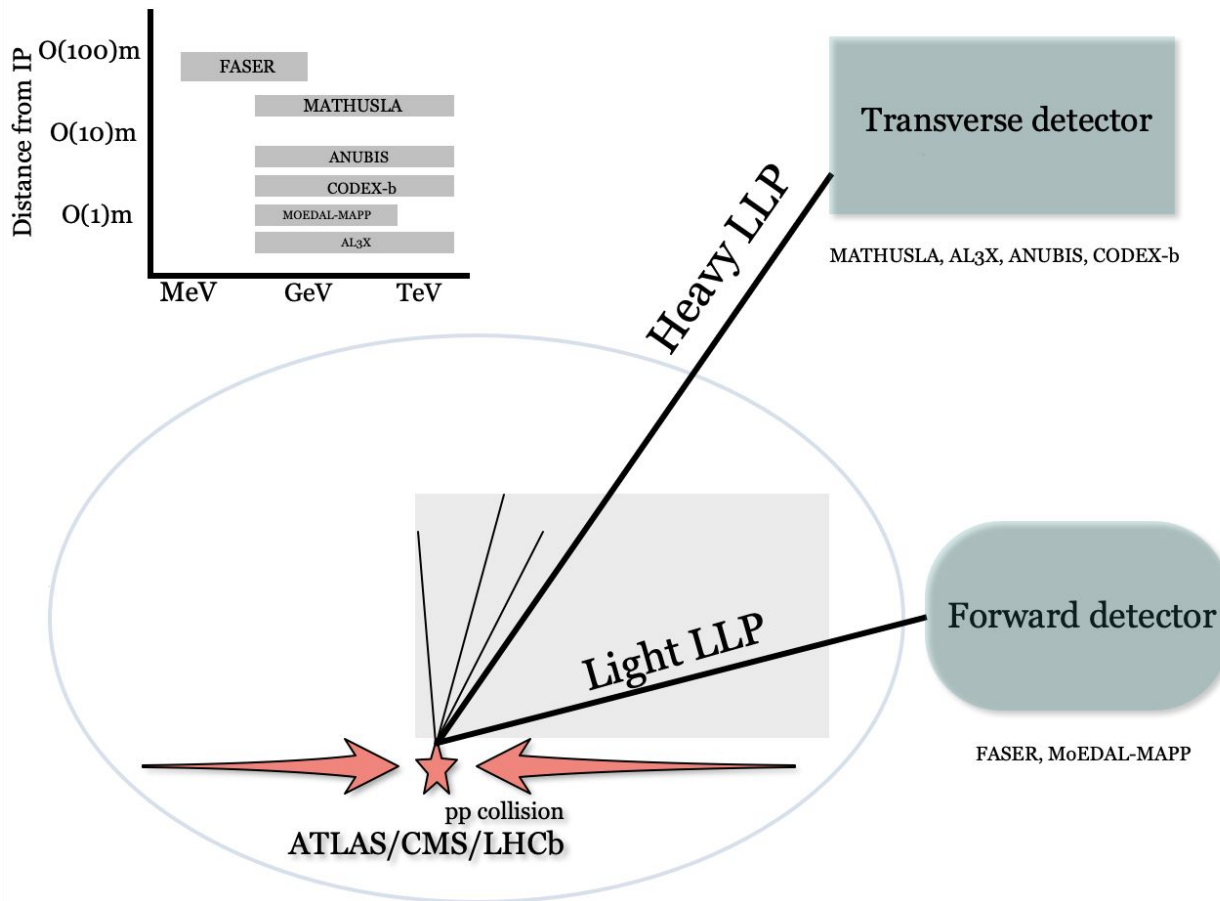


New approved/proposed LHC-LLP detector experiments

All these experiments have the potential to search for LLPs with different geometrical acceptances



Many LLP detectors needed to target complementary regions in LLP model parameter space



General purpose detectors are “complex” (large angular acceptance, many sub-detectors to exploit LLP signature)/ no LLP triggers/challenging triggers, large backgrounds, difficult-non-efficient reconstruction

Dedicated detectors are “simpler” (no trigger, targeted reconstruction, are cheaper), limited/zero bkg/smaller angular acceptance

LHC-LLP phenomenology in selected models

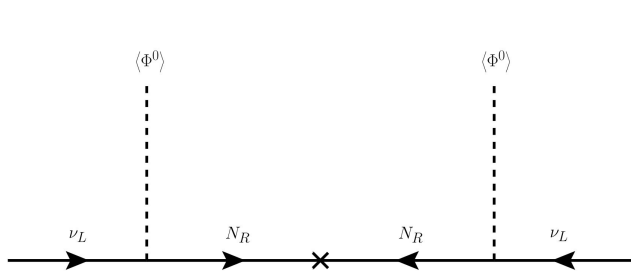
Theoretical motivation: Several models that explain **neutrino masses** and **dark matter** predict LLPs



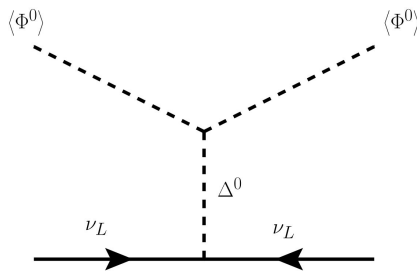
$$\mathcal{L}_5 = \frac{C_5^{\ell\ell'}}{\Lambda} [\Phi \cdot \bar{L}_\ell^c] [L_{\ell'} \cdot \Phi] + \text{H.c.}$$

Seesaw
P. Minkowski, [Phys. Lett. 67B \(1977\)](#)
R. N. Mohapatra and G. Senjanovic, [Phys. Rev. Lett. 44 \(1980\)](#)
J. Schechter and J. W. F. Valle, [Phys. Rev. D22, 2227 \(1980\)](#)
Diagrams from S. Centelles, R. Srivastava, J.W.F. Valle, [Phys. Lett. B 781 \(2018\)](#)

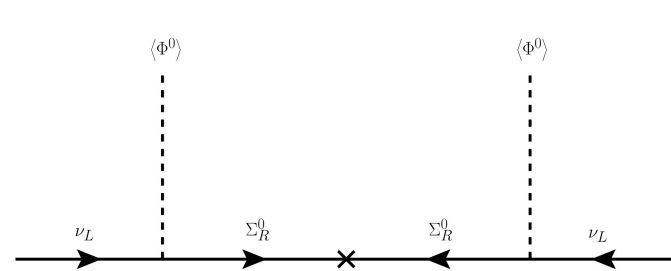
At tree level, the d=5 Weinberg operator can be generated in three ways (type-I,II and III seesaw)



$$N_R \equiv F(1, 1, 0)$$



$$\Delta^0 \equiv S(1, 3, 1)$$



$$\Sigma_R^0 \equiv F(1, 3, 0)$$

At one-loop, there are many more possibilities, where neutrino masses and dark matter (DM) might have a common origin

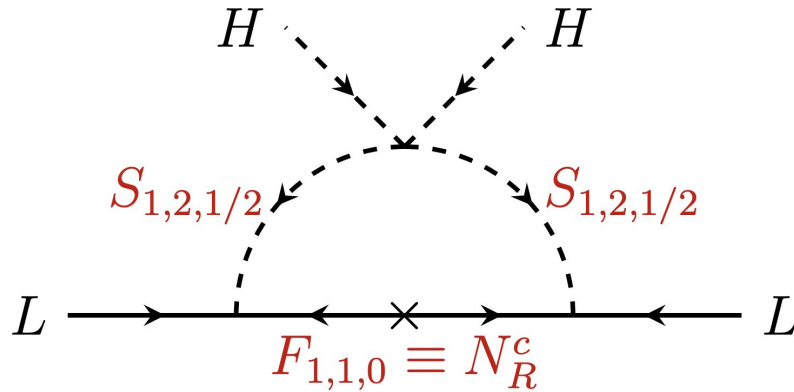
Scotogenic Model

E. Ma, [Phys.Rev.D 73 \(2006\) 077301](#)

LLP phenomenology in Scotogenic model in

A. G. Hessler, A. Ibarra, E. Molinaro, S. Vogl, [JHEP 01 \(2017\)](#)

I. M. Ávila, G. Cottin, M. A. Díaz, [J.Phys.G 49 \(2022\)](#)

SmodelS2 in G. Alguero et al, [JHEP 08 \(2022\) 068](#)

Scalar-Singlet-Triplet (SST) Model

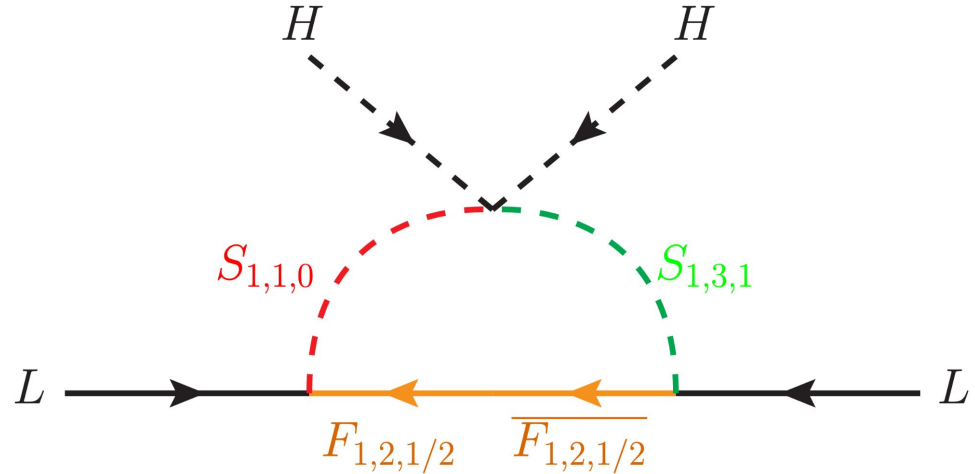
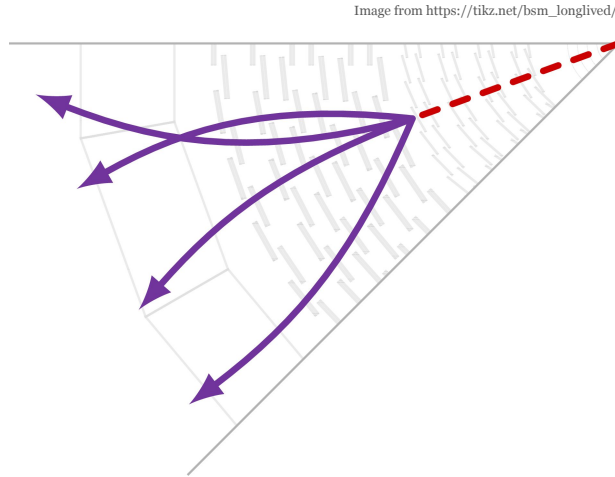
C. Arbeláez, G. Cottin, J. C. Helo, M. Hirsch, T. B. de Melo, [arXiv:2408.03364](https://arxiv.org/abs/2408.03364)[illegible]

Table 8: DM models with exotics (DM-E) which need a stabilizing symmetry to give an acceptable WIMP candidate.

LLP Phenomenology for Heavy Neutral Leptons (Neutral LLP)

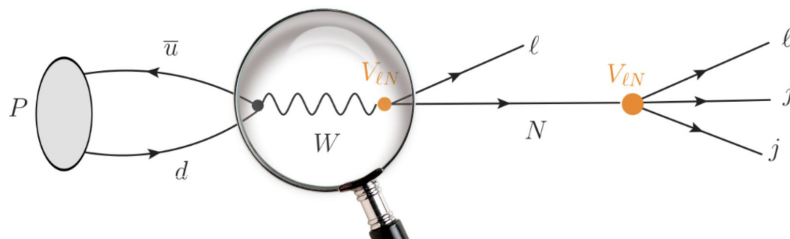


Minimal Heavy Neutral Lepton (HNL) model

- HNLs (fermionic singlets) motivated by seesaw models
- HNLs mix with SM neutrinos
- Can be automatically long-lived! $\Gamma \sim G_F^2 |V_{lN}|^2 m_N^5$

$$\mathcal{L}_{\min} = -\frac{g}{\sqrt{2}} V_{\ell N} \bar{\ell} \gamma^\mu P_L N W_\mu^\dagger - \frac{g}{2 \cos \theta_W} U_{\ell i}^* V_{\ell N} \bar{\nu}_i \gamma^\mu P_L N Z_\mu + \text{h.c.}$$

$$m_N < m_B$$

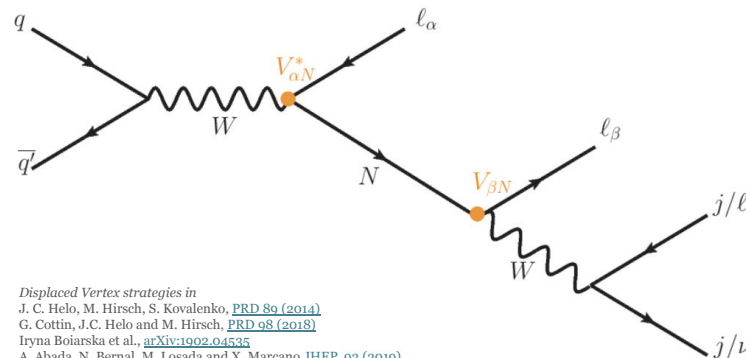


$\mathcal{O}(\text{GeV})$

K. Bondarenko, et al, [JHEP 11 \(2018\) 032](#)

Diagrams by R. Beltrán

$$m_B < m_N < m_W$$



Displaced Vertex strategies in

J. C. Helo, M. Hirsch, S. Kovalenko, [PRD 89 \(2014\)](#)

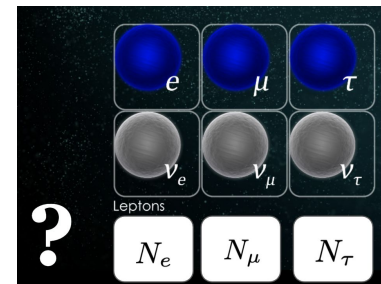
G. Cottin, J.C. Helo and M. Hirsch, [PRD 98 \(2018\)](#)

Iryna Boiarska et al., [arXiv:1902.04535](#)

A. Abada, N. Bernal, M. Losada and X. Marciano [JHEP 03 \(2019\)](#)

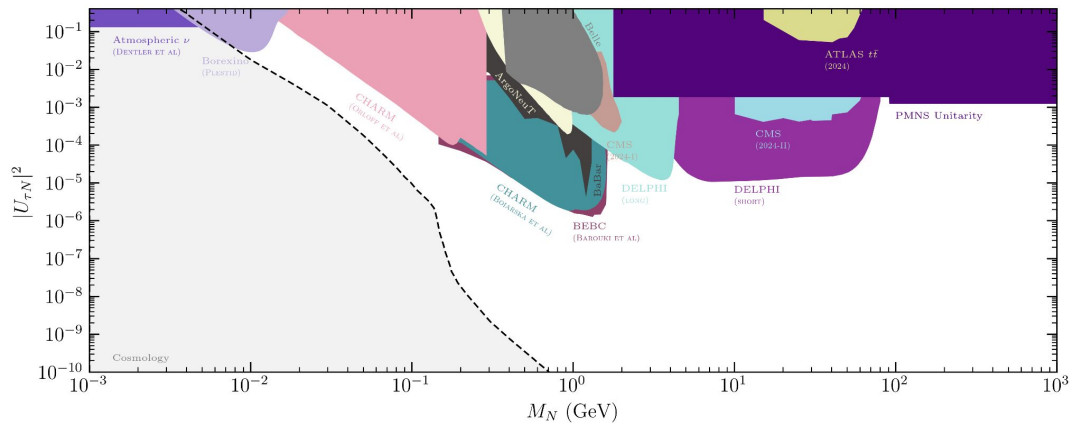
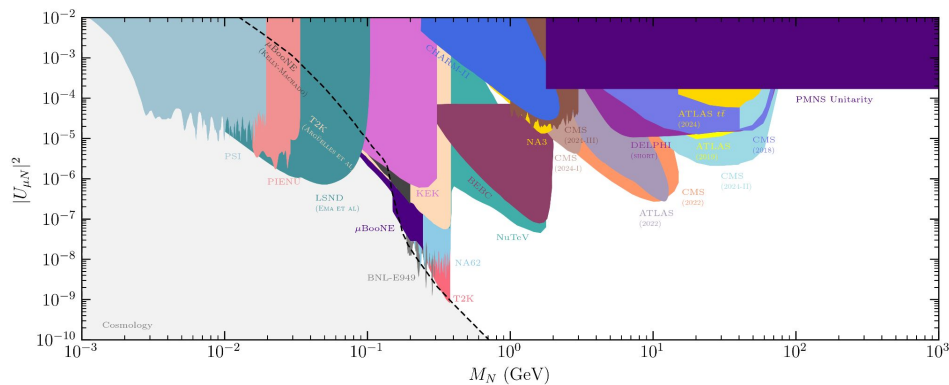
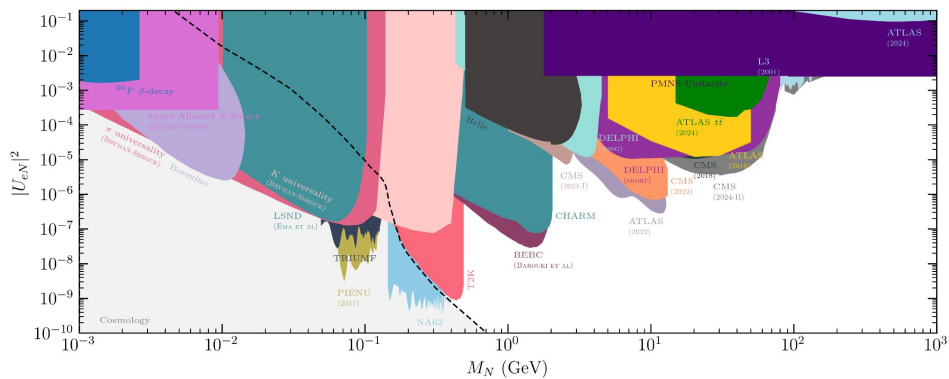
M. Drewes, J. Hajer, [JHEP 70 \(2020\)](#)

See reviews for HNL phenomenology
M. Drewes, Int.J.Mod.Phys.E 22 (2013) 1330019, [arXiv:1303.6912](#)
F. Deppisch, P. S. Bhupal Dev, Apostolos Pilaftsis, New J.Phys. 17 (2015) 7, 075019, [arXiv:1502.06541](#)
Y. Cai, T. Han, Tong Li, R. Ruiz, Front.in Phys. 6 (2018) 40, [arXiv:1711.02180](#)
The Present and Future Status of Heavy Neutral Leptons, Snowmass, [arXiv:2203.08039](#)

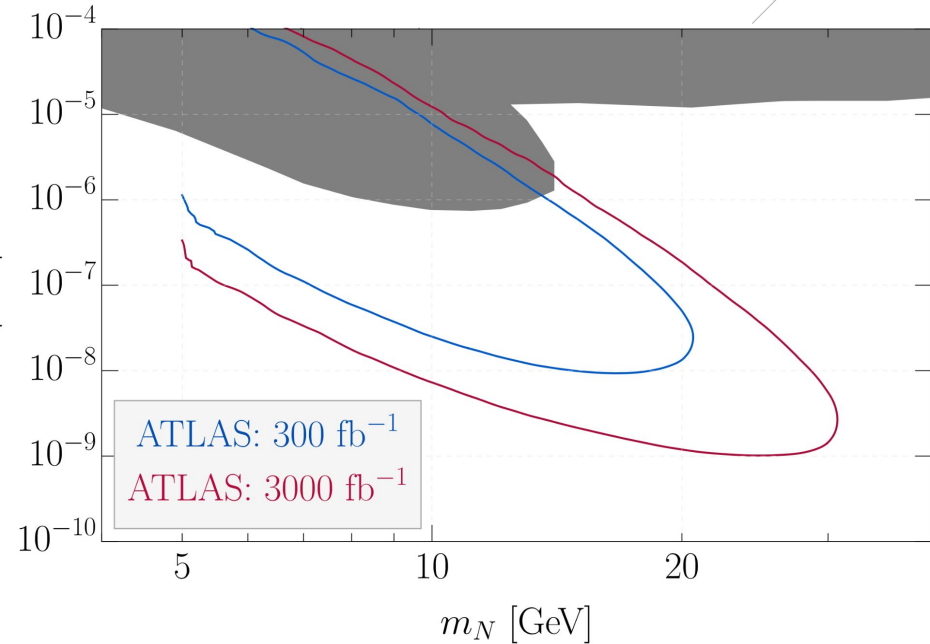
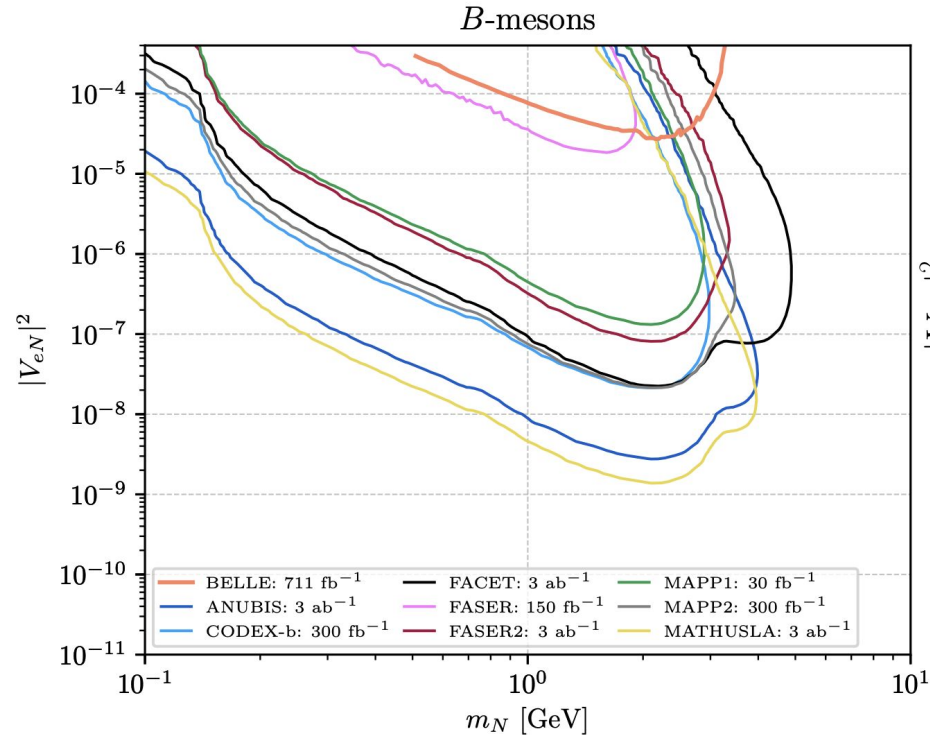
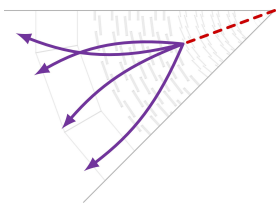


Experimental Landscape

E. Fernandez-Martinez, et. al, [2304.06772](#), <https://github.com/mhostert/Heavy-Neutrino-Limits>



Sensitivity prospects in the Minimal Heavy Neutral Lepton (HNL) model with displaced activity



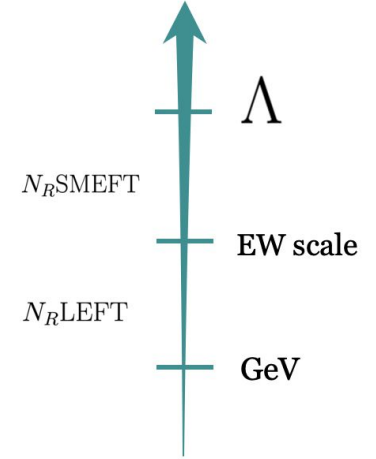
Beyond the minimal HNL model

- We can study HNLs systematically in an EFT approach, with NRO which are suppressed by a new physics scale
- HNL production and/or decay can be dominated by the operator

$$\mathcal{L}_{N_R\text{SMEFT}} = \mathcal{L}_{\text{SM}+N_R} + \sum_{d \geq 5} \frac{1}{\Lambda^{d-4}} \sum_i c_i^{(d)} \mathcal{O}_i^{(d)}$$

Example of $d=6$
four-fermion
operators with
pairs of HNL

Name	Structure
\mathcal{O}_{dN}	$(\bar{d}_R \gamma^\mu d_R) (\bar{N}_R \gamma_\mu N_R)$
\mathcal{O}_{uN}	$(\bar{u}_R \gamma^\mu u_R) (\bar{N}_R \gamma_\mu N_R)$
\mathcal{O}_{QN}	$(\bar{Q} \gamma^\mu Q) (\bar{N}_R \gamma_\mu N_R)$
\mathcal{O}_{eN}	$(\bar{e}_R \gamma^\mu e_R) (\bar{N}_R \gamma_\mu N_R)$
\mathcal{O}_{NN}	$(\bar{N}_R \gamma_\mu N_R) (\bar{N}_R \gamma_\mu N_R)$
\mathcal{O}_{LN}	$(\bar{L} \gamma^\mu L) (\bar{N}_R \gamma_\mu N_R)$



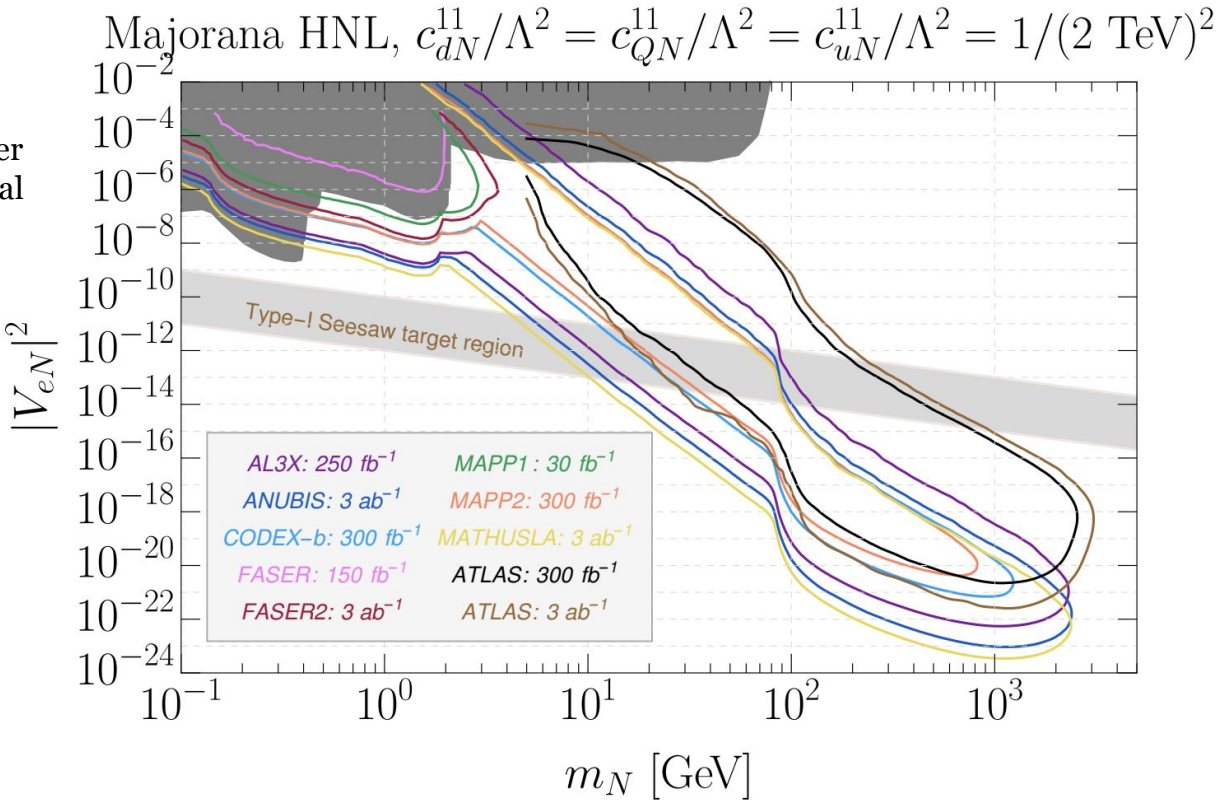
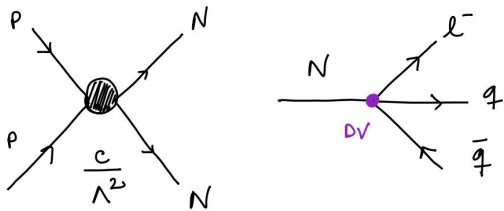
First developed in
F. del Aguila, S. Bar-Shalom, A. Soni, J. Wudka, [0806.0876](#) (Phys.Lett.B670, 2008)
A. Aparici, K. Kim, A. Santamaria, J. Wudka, [0904.3244](#) (Phys.Rev.D80, 2009)
Basis for $d \leq 9$ in
H.-L. Li, Z. Ren, M.-L. Xiao, J.-H. Yu, Y.-H. Zheng, [2105.09329](#)

HNLs in EFT with LLPs at the LHC
(d=5) A. Caputo, et al, [JHEP 06 \(2017\)](#)
(NRLEFT, singleN) Jordy de Vries, et al, [JHEP 03 \(2021\)](#)
(NRSMEFT, pair N) G. Cottin, et al, [JHEP 01 \(2021\)](#)
(NRSMEFT, single N) R. Beltrán, et al, [JHEP 09 \(2022\)](#)
(NRLEFT, pair N) R. Beltrán, et al, [JHEP 01 \(2023\)](#)
(NRLEFT), E. Fernandez-Martinez, et. al, [JHEP09\(2023\)](#)

Sensitivity prospects beyond the Minimal HNL model with displaced activity

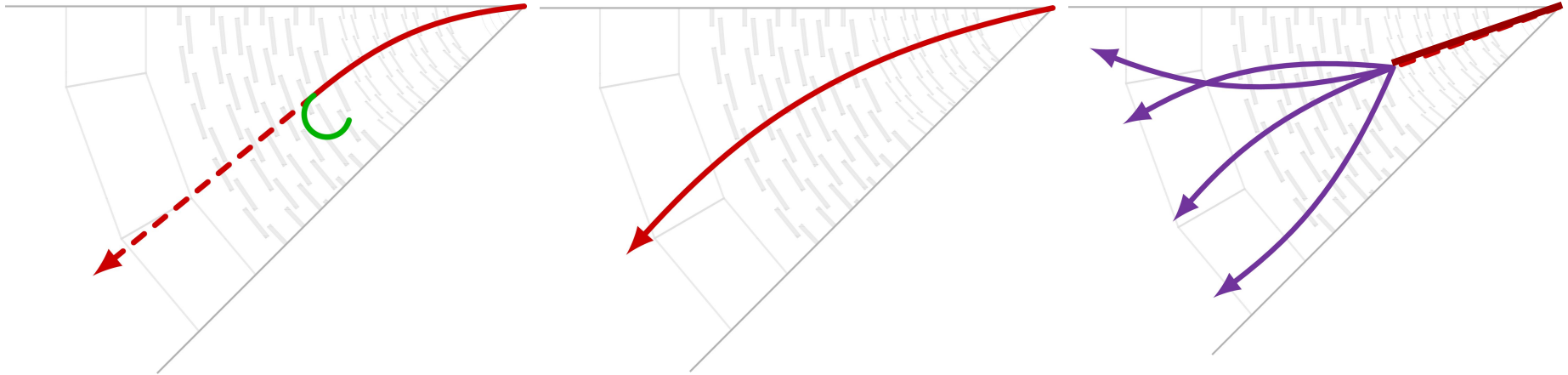
G. Cottin, J. C. Helo, M. Hirsch, A. Titov, Z. S. Wang, [JHEP 09 \(2021\) 039](#)

- Production dominated by the operator
- HNLs decay only via mixing
- DV strategy proposed for ATLAS tracker
- Probability of displaced decay in fiducial volume in far detectors



LLP Phenomenology for Scotogenic-like dark matter models (multi-charged LLPs)

Image from https://tikz.net/bsm_longlived/
modified by G.Cottin

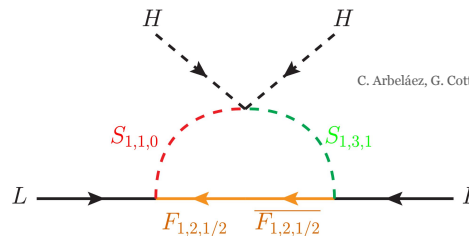
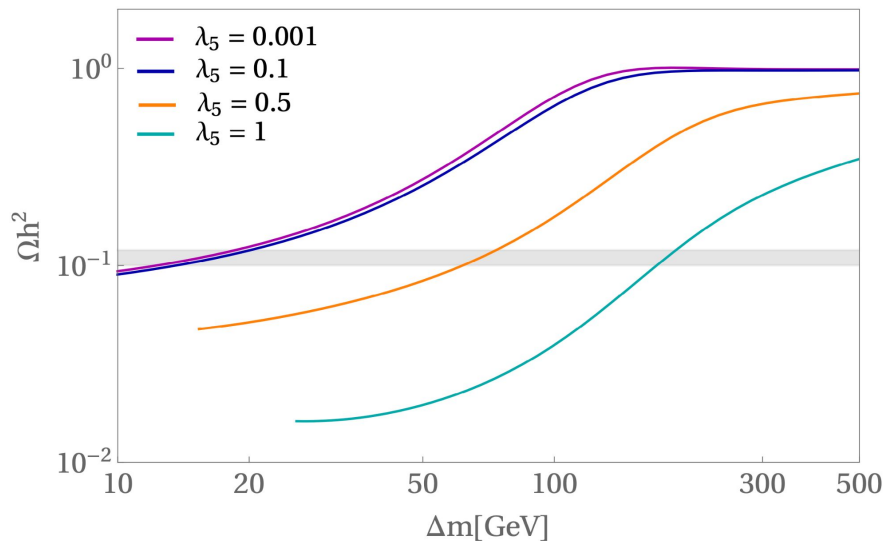


Scalar-Singlet-Triplet (SST) Model

- Neutrino masses at 1-loop
- DM candidate is a combination of singlet+triplet
- Predicts a doubly-charged long-lived particle (in the co-annihilation region)

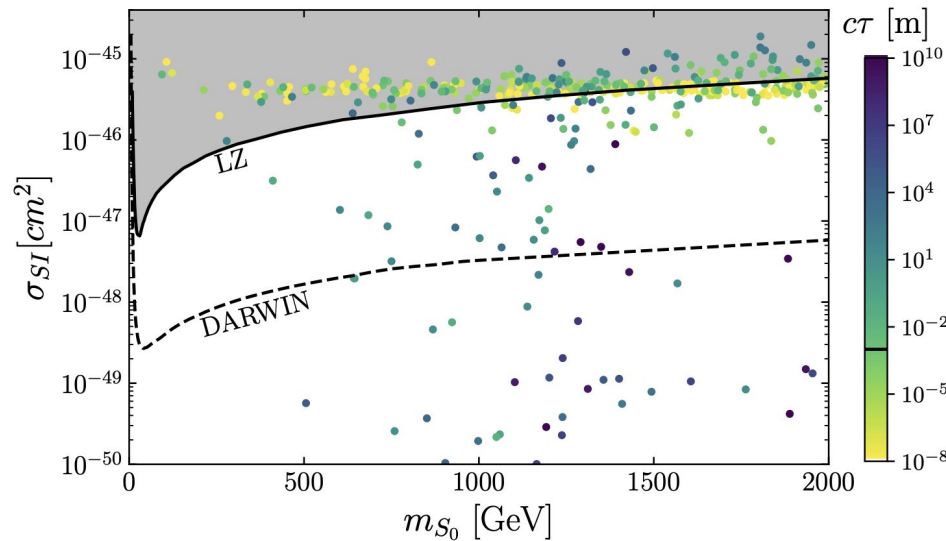
Relic density favours points with small mass splittings between scalar singlet and triplet

$$\lambda_5 H H S_1 S_3^\dagger \quad \Delta m \equiv m_{t^{++}} - m_{S_0}$$

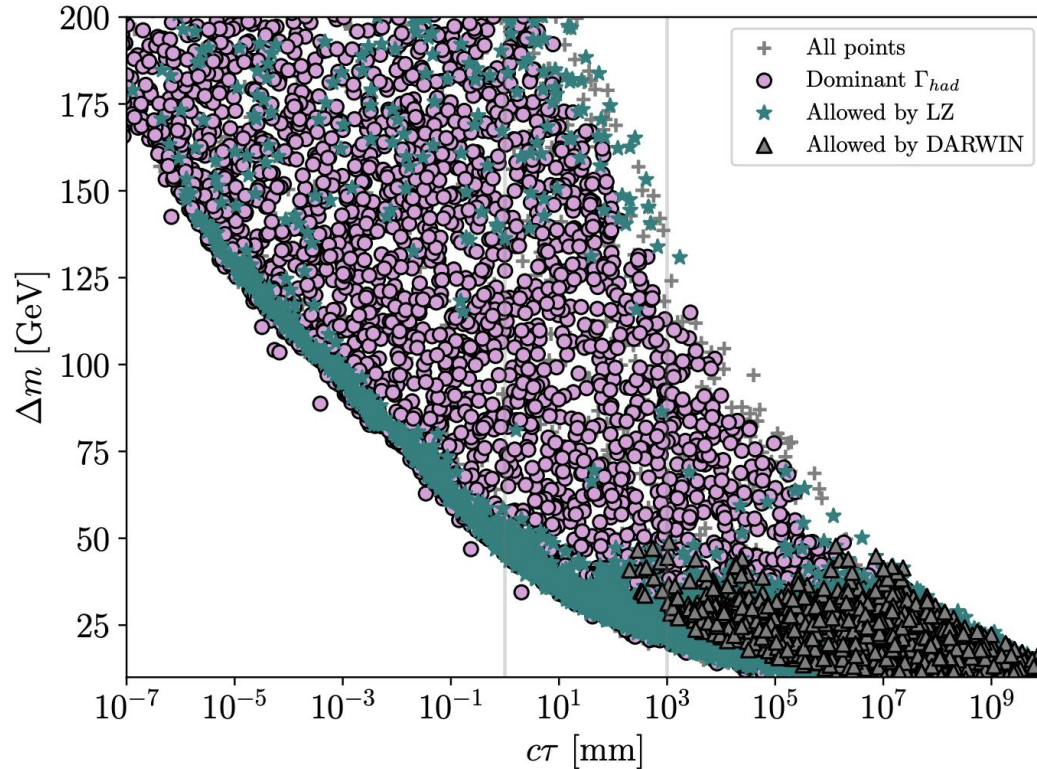


Coannihilations in
K. Griest and D. Seckel, [Phys. Rev. D 43 \(1991\)](#)
M. Klasen, C. E. Yaguna, et al, [JCAP 04, 044 \(2013\)](#)
1-loop neutrino mass models with dark matter in
C. Arbeláez, et al, [JHEP 08 \(2022\) 023](#)
LLP phenomenology for SST model in
C. Arbeláez, G. Cottin, J. C. Helo, M. Hirsch, T. B. de Melo, [arXiv:2408.03364](#)

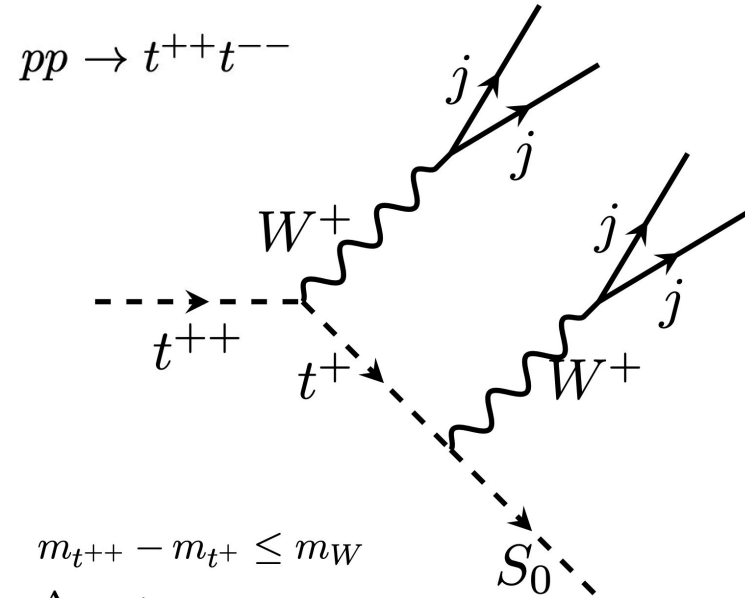
Points in this region that pass future direct detection bounds have small decay width for t^{++}



The presence of the scalar triplet will give rise to signals (potentially very different) from the scotogenic model

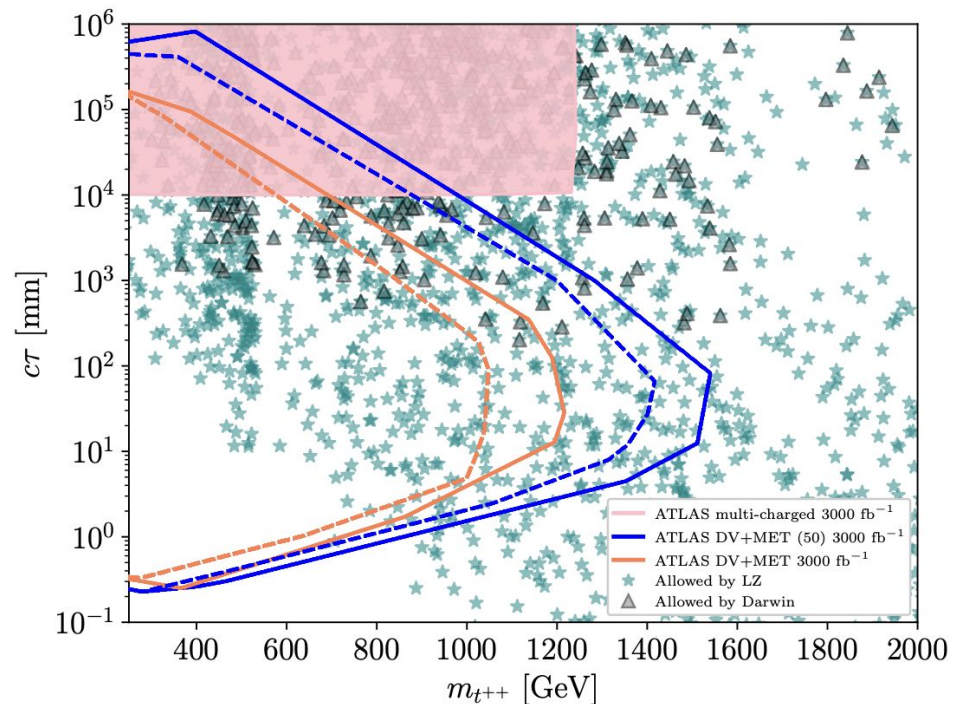
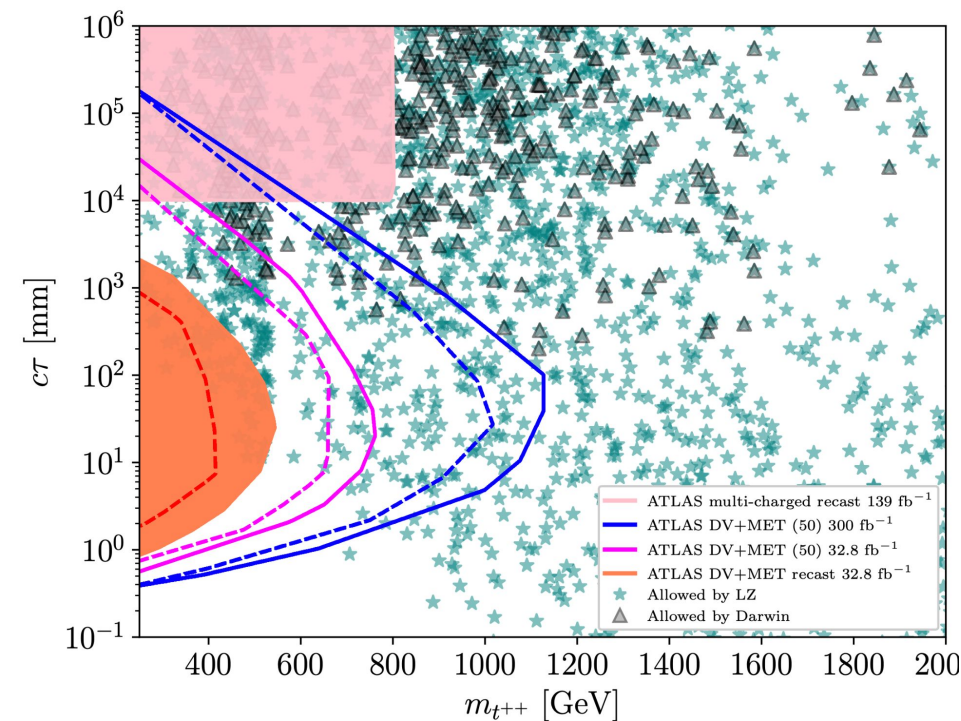
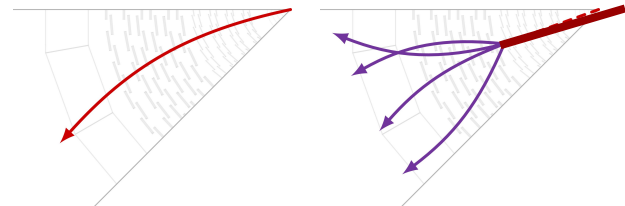


$$S_3 = \begin{pmatrix} t^+/\sqrt{2} & t^{++} \\ t^0 & -t^+/\sqrt{2} \end{pmatrix}$$



Sensitivity prospects for SST model with displaced activity

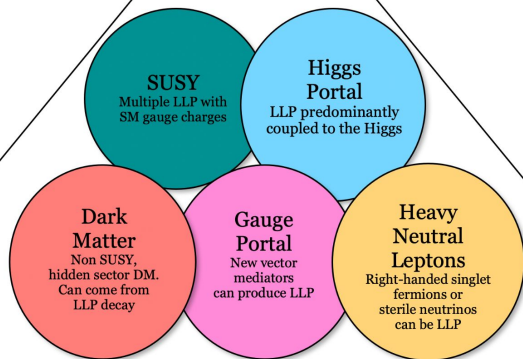
LLP phenomenology in C. Arbeláez, G. Cottin, J. C. Helo, M. Hirsch, T. B. de Melo, [arXiv:2408.03364](https://arxiv.org/abs/2408.03364)



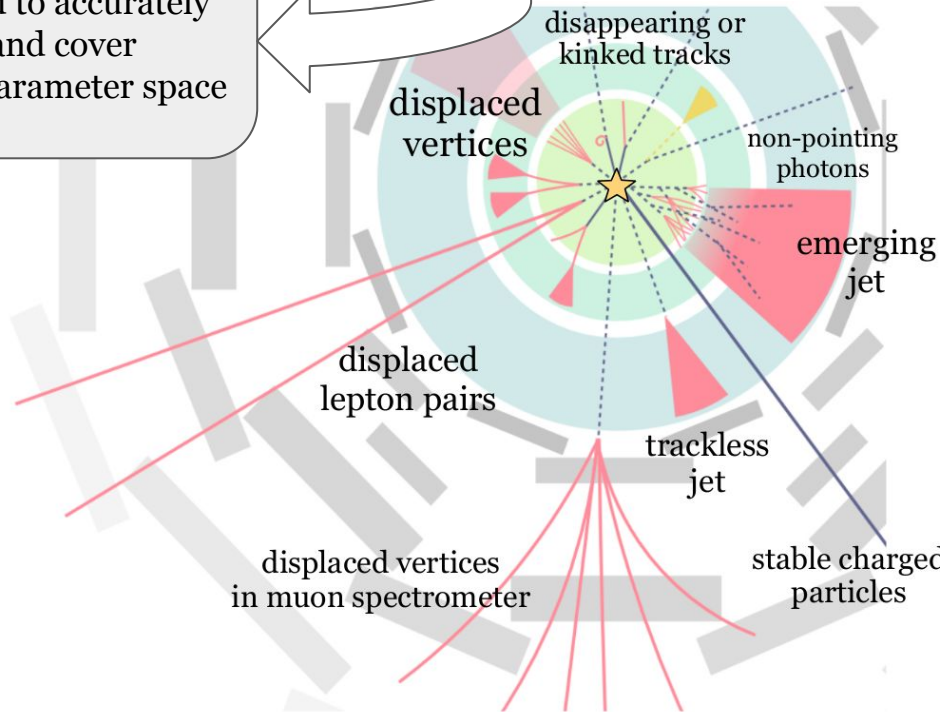
- What is a long-lived particle?
- Why look for long-lived particles?
- What's the long-lived particle phenomenology @ LHC?
- Why should we keep looking for long-lived particles?

I presented examples motivated by **neutrino masses** and **dark matter**

Many BSM models predict LLPs with exotic and challenging signatures. Phenomenology needed to accurately characterize them and cover (unconstrained) model parameter space

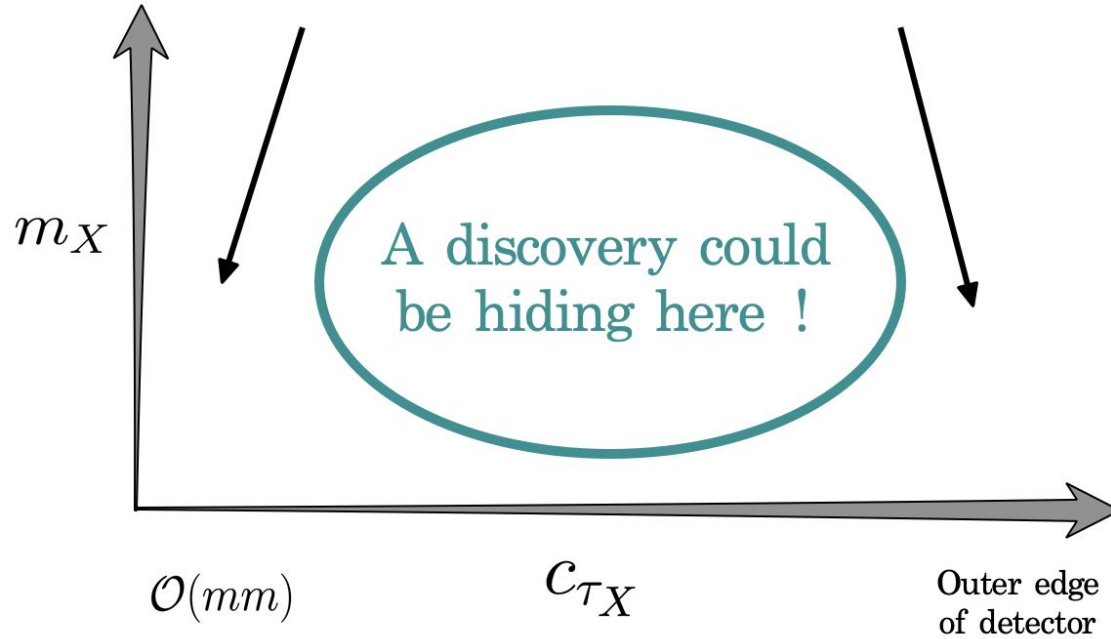


Umbrella theories predicting Long-Lived Particles (LLP)

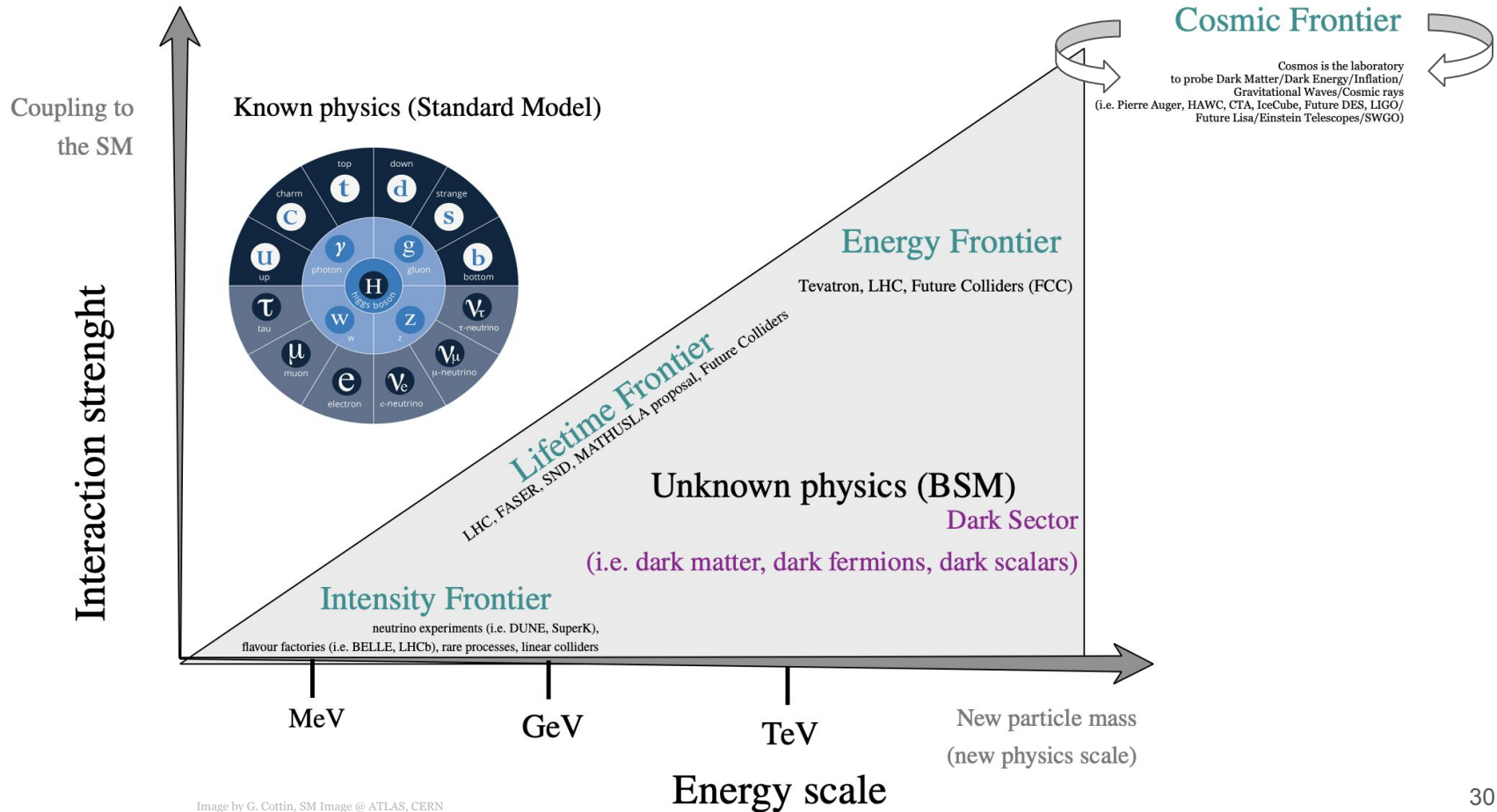


New physics "X" at colliders

most LHC searches here



The search for new physics: a tale of cross frontiers



A long-lived paradigm shift

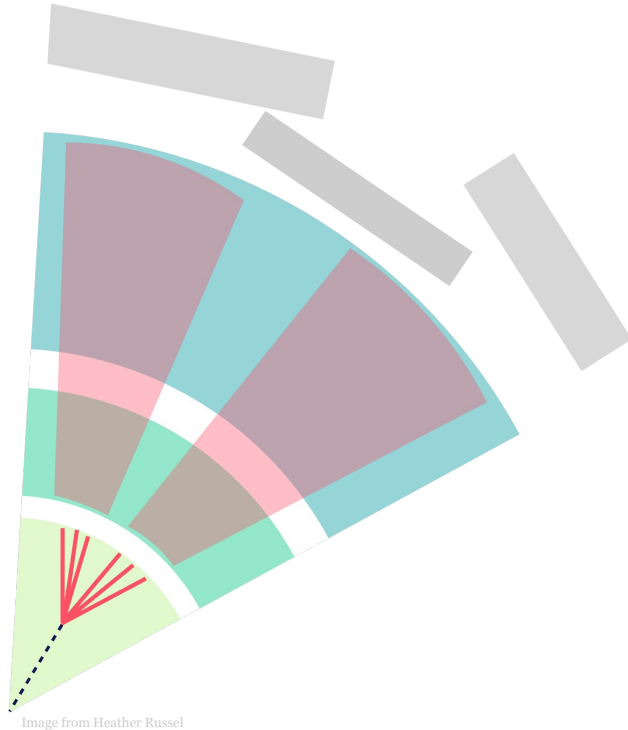
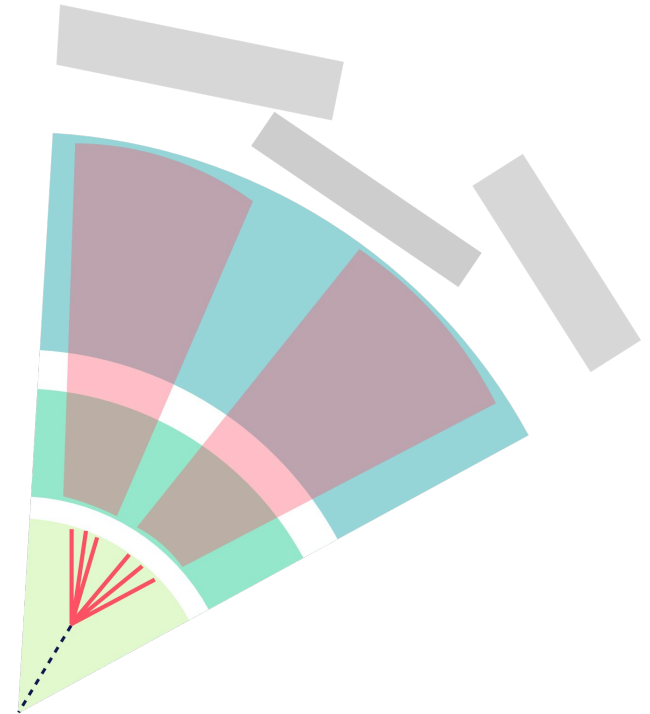
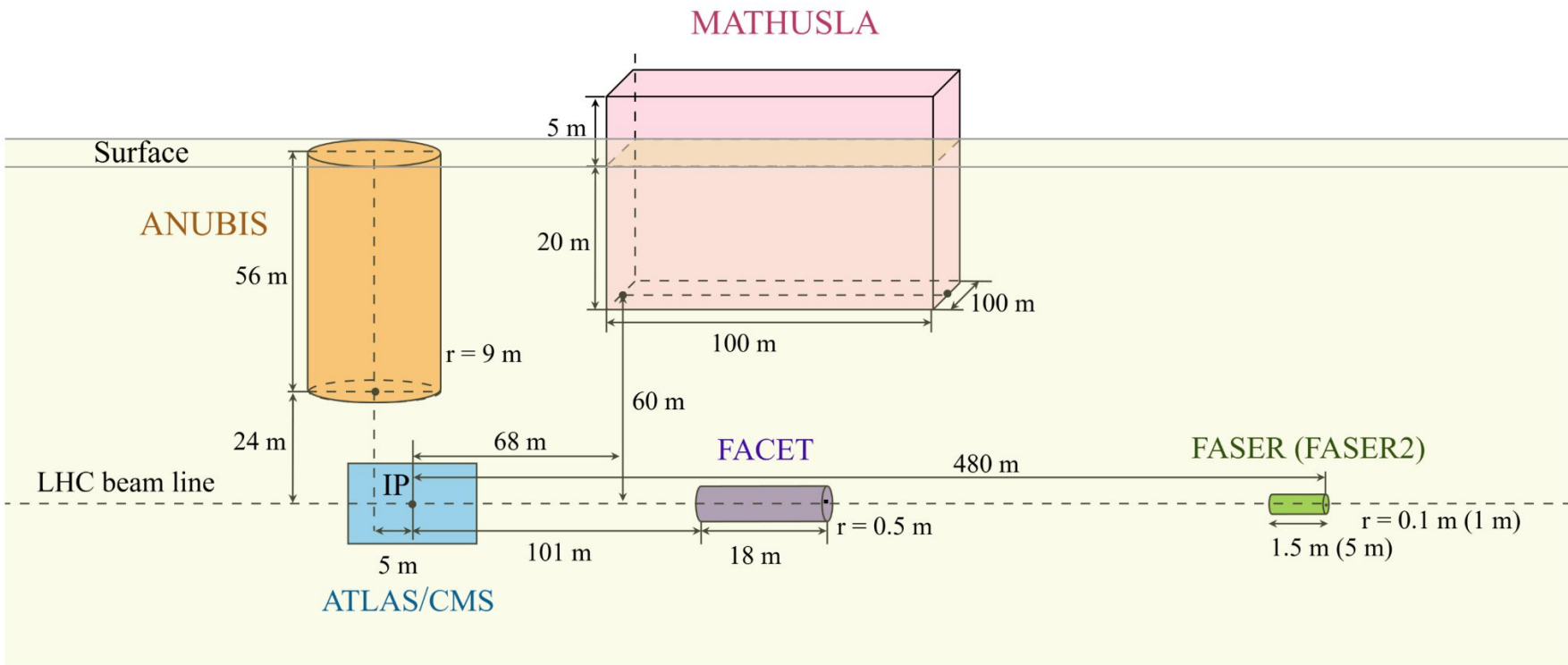


Image from Heather Russel

Backup



New approved/proposed transverse far detector experiments complement reach



Details of probability of decay formulas in fiducial volumes

See A C++ program for estimating detector sensitivities to long-lived particles, F. Domingo, J. Günther, J. S. Kim, Z.S. Wang, Eur.Phys.J.C 84 (2024), [arXiv:2308.07371](https://arxiv.org/abs/2308.07371)

Jordy de Vries, H. K. Dreiner, J. Y. Günther, Z. S. Wang, G. Zhou, [2010.07305](https://arxiv.org/abs/2010.07305) (JHEP 03 (2021))

For FASER, MATHUSLA and CODEX-b see D. Dercks, J. de Vries, H. K. Dreiner, Z. S. Wang, [1810.03617](https://arxiv.org/abs/1810.03617) (Phys. Rev. D 99, 055039 (2019)) and earlier in

J.C. Helo, M. Hirsch, Z. S. Wang, [1803.02212](https://arxiv.org/abs/1803.02212) (JHEP 07 (2018))

For ANUBIS see M. Hirsch, Z. S. Wang, [2001.04750](https://arxiv.org/abs/2001.04750) (Phys. Rev. D 101, 055034 (2020))

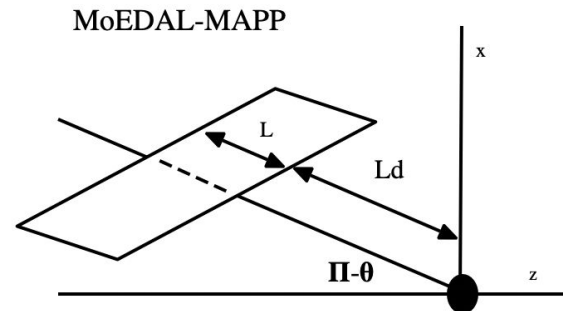
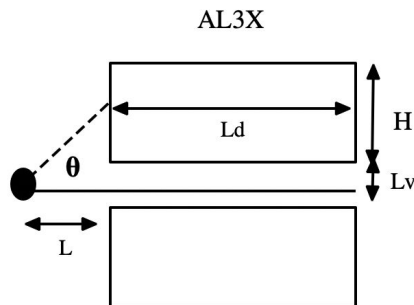
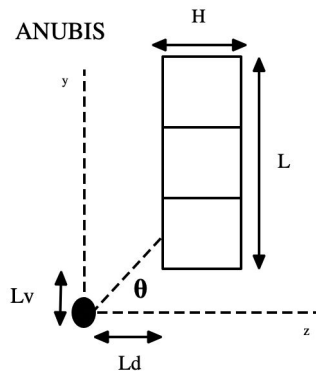
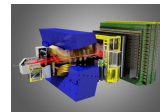
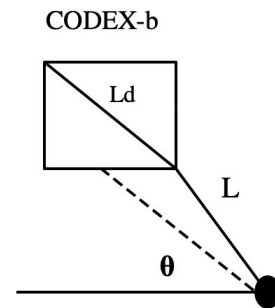
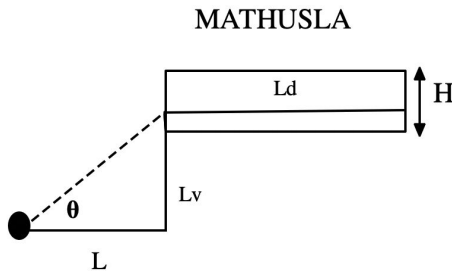
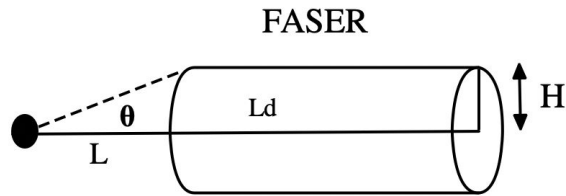
For AL3X see D. Dercks, H.K. Dreiner, M. Hirsch, Z. S. Wang, [1811.01995](https://arxiv.org/abs/1811.01995) (Phys.Rev. D 99, 055020 (2019))

For MAPP see H. K. Dreiner, J. Y. Günther, Z. S. Wang, [2008.07539](https://arxiv.org/abs/2008.07539) (Phys. Rev. D 103, 075013 (2021))

Diagram from C.-T. Lu, X. Wang, X. Wei, Y. Wu, [arXiv:2410.19561](https://arxiv.org/abs/2410.19561)

Geometry for LLP detectors

Decay probability of each simulated HNL takes into account the far detector geometry (L, L_d, L_v, H, θ) and their kinematics



Details of probability of decay formulas in fiducial volumes

See Jordy de Vries, H. K. Dreiner, J. Y. Günther, Z. S. Wang, G. Zhou, [2010.07305](#) (JHEP 03 (2021))

For FASER, MATHUSLA and CODEX-b see D. Dercks, J. de Vries, H. K. Dreiner, Z. S. Wang, [1810.03617](#) (Phys. Rev. D 99, 055039 (2019)) and earlier in

J.C. Helo, M. Hirsch, Z. S. Wang, [1803.02212](#) (JHEP 07 (2018))

For ANUBIS see M. Hirsch, Z. S. Wang, [2001.04750](#) (Phys. Rev. D 101, 055034 (2020))

For AL3X see D. Dercks, H. K. Dreiner, M. Hirsch, Z. S. Wang, [1811.01995](#) (Phys. Rev. D 99, 055020 (2019))

For MAPP see H. K. Dreiner, J. Y. Günther, Z. S. Wang, [2008.07539](#) (Phys. Rev. D 103, 075013 (2021))

LLP far detectors

FASER

Cylinder with $r = 10$ cm and $\ell = 1.5$ m
 $c\tau \sim 480$ m

Boosted cross section

$4\pi/10^3$ coverage

CoDEX-b

Box of 10 m \times 10 m \times 10 m
 $c\tau \sim 25$ m

$4\pi/10^2$ coverage

Lower luminosity

ANUBIS

Cylinder with $r = 9$ m and $h = 56$ m
 $c\tau \sim \text{few } 10$ m

$4\pi/50$ coverage

MATHUSLA

Box of 100 m \times 100 m \times 25 m
 $c\tau \sim \mathcal{O}(100)$ m

$4\pi/25$ coverage

AL3X: A Laboratory for Long-Lived eXotics
@ALICE

Cylinder with 0.85 m $< r < 5$ m and $\ell = 12$ m
 $c\tau \sim 10$ m

MoEDAL-MAPP: MoEDAL's Apparatus for Penetrating Particles
(MoEDAL: Monopole and Exotics Detector at the LHC)
@LHCb

MAPP1: ~ 130 m³

MAPP2: ~ 430 m³

$c\tau \sim 50$ m

“HNL optimized” displaced strategies/proposals in the Minimal HNL Model that build up on experimental LLP searches not necessarily designed for HNLs (apologies if I missed your work!)

Displaced vertex searches

- HNL decays leptonically and/or semileptonically
- prompt lepton trigger from W decay at LHC

Inner Trackers (IT)

- high-mass and displaced track multiplicity DVs in inner tracker to suppress hadronic bkg.

J. C. Helo, M. Hirsch, S. Kovalenko, [PRD 89 \(2014\)](#), G. Cottin, J.C. Helo and M. Hirsch, [PRD 98 \(2018\)](#), Iryna Boiarska et al., [arXiv:1902.04535](#), A. Abada, N. Bernal, M. Losada and X. Marcano [JHEP 03 \(2019\)](#), M. Drewes, J. Hajer, [JHEP 70 \(2020\)](#)

Muon Systems (MS)

- can access larger displacements (i.e lower HNL masses)
- purely muonic vertices give better sensitivity at low HNL masses (not necessarily limited by invariant mass cut of DV)

K. Bondarenko et al., [Phys. Rev. D 100 \(2019\)](#)
M. Drewes, J. Hajer, [JHEP 70 \(2020\)](#)

Future colliders

- can access lower mixings due to large elec-pos production at the Z pole

S. Antusch, O. Fischer, A. Hammad, [JHEP 03 \(2020\)](#)
J. Alimena, et. al, [Front. Phys. 10:967881 \(2022\)](#)
S Ajmal, et al., [arXiv:2410.03615](#)

Displaced lepton searches

- HNL decays leptonically
- prompt lepton trigger
- At least one displaced track from HNL is a non-isolated lepton
- Can do explicit vertex reconstruction of displaced di-muon/di-electron/emu vertices in ID or MS
- No explicit vertex requirement leads to additional sources of SM bkg.

J. Liu, Z. Liu, Lian-Tao Wang, Xiao-Ping Wang, [JHEP 07\(2019\)](#)
M. Drewes, J. Hajer, [JHEP 70 \(2020\)](#)
LHC HNL displaced searches for displaced dilepton DV: ATLAS [1905.09787](#), [2204.11988](#),
CMS [2201.05578](#), LHCb [2011.05263](#)

Displaced lepton-jets

- Prompt lepton trigger from W decay
- Reconstruct boosted HNL decay as a single displaced $l\bar{j}$ object (with more than one leptonic displaced track concentrated within a cone)

E. Izaguirre and B. Shuve, [Phys. Rev. D 91 \(2015\)](#)
S. Dube, D. Gadkari, and A. M. Thalapillil, [Phys. Rev. D 96 \(2017\)](#)

Displaced shower

- MET/displaced triggers
- Reconstruction of cluster of hits @CMS

G.Cottin, J.C. Helo, M. Hirsch, C. Peña, C. Wang, S. Xie, [JHEP02\(2023\)011](#)
CMS [PAS-EXO-22-017](#)

SST Model

Fields	$SU(3)_C$	$SU(2)_L$	$U(1)_Y$	\mathbb{Z}_2
(Real) S_1	1	1	0	—
S_3	1	3	1	—
F (\bar{F})	1	2	1/2 (−1/2)	—

Table 1. Quantum numbers for the particle content of the model. While neutrino masses could be explained already with one copy of F/\bar{F} , in our numerical study we use three families of F/\bar{F} to reflect the number of generations in the SM.

The triplet scalar can be written in the usual way:

$$S_3 = \begin{pmatrix} t^+/\sqrt{2} & t^{++} \\ t^0 & -t^+/\sqrt{2} \end{pmatrix}. \quad (2.1)$$

The Lagrangian of the model adds the following terms to the SM Lagrangian:

$$\begin{aligned} \mathcal{L} = & - (M_F F \bar{F} + M_{S_1}^2 |S_1|^2 + M_{S_3}^2 |S_3|^2) - (h_F L F S_1 + h_{\bar{F}} L \bar{F} S_3 + \text{h.c.}) \\ & + \lambda_2 |H|^2 |S_1|^2 + \left(\lambda_{3a} |H|^2 |S_3|^2 + \lambda_{3b} \epsilon_{l_1 l_3} \epsilon_{l_2 l_4} \epsilon_{l_3 b l_4 b} H_{l_1} H_{l_2}^* (S_3)_{l_3 l_{3b}} (S_3^*)_{l_4 l_{4b}} \right) \\ & + \lambda_4 (|S_1|)^4 - (\lambda_5 H H S_1 S_3^\dagger + \text{h.c.}) \end{aligned} \quad (2.2)$$

SST Model

$$(m_{s_1^0, s_2^0})^2 = \frac{1}{4} \left(2(M_{S_1}^2 + M_{S_3}^2) - (\lambda_2 + \lambda_{3_a})v^2 \right. \\ \left. \mp \sqrt{(2(M_{S_1}^2 - M_{S_3}^2) + v^2(\lambda_{3_a} - \lambda_2))^2 + 8\lambda_5^2 v^4} \right),$$

$$m_{t^{++}}^2 = M_{S_3}^2 - \frac{1}{2} (\lambda_{3_a} + \lambda_{3_b}) v^2, \\ m_{t^+}^2 = M_{S_3}^2 - \frac{1}{4} (2\lambda_{3_a} + \lambda_{3_b}) v^2.$$

