



Universidad
de Huelva

Covariant Extension of Generalized Parton Distributions using Artificial Neural Networks

Pietro Dall'Olio

Collaborators: J. M. Morgado, J. Rodríguez Quintero, C. Mezrag, P. Sznajder

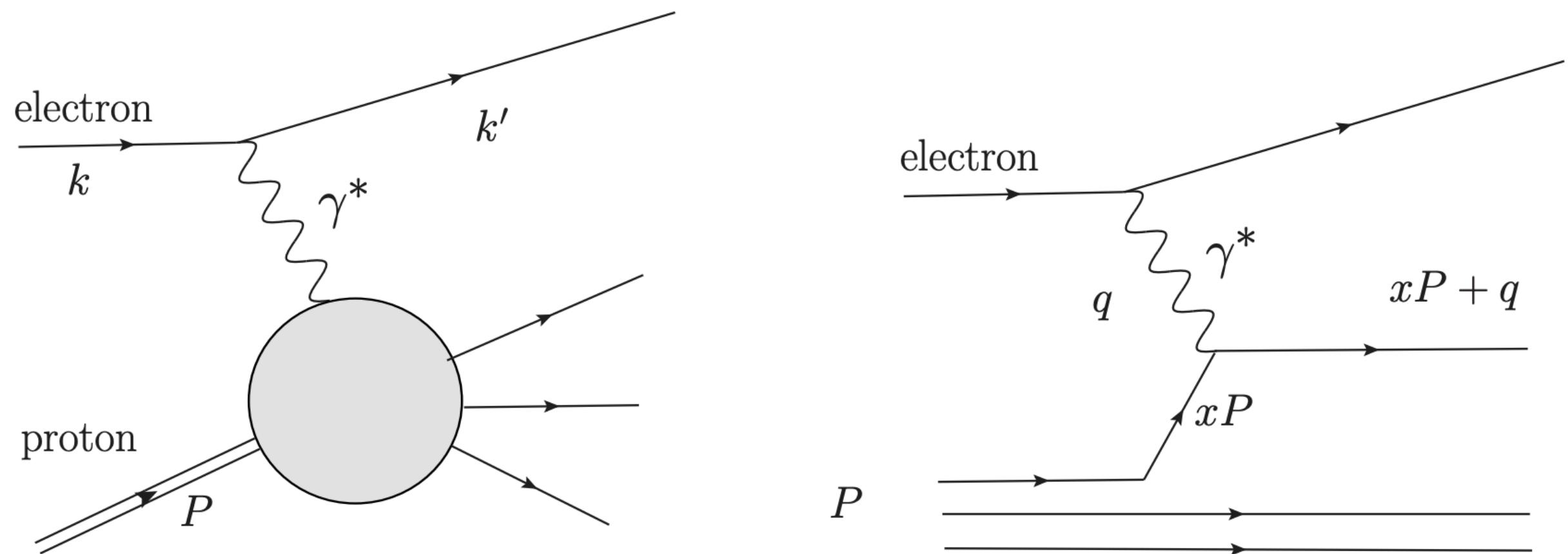
Non-Perturbative Physics: Tools and Applications 06 Sept 2023

Outline

- (Generalized) Parton Distribution Functions
- Covariant Extension via Radon Transform (RT) inversion
- RT inversion using Artificial Neural Networks
- Results on analytical models
- Conclusions and outlook

Parton Distribution Function (PDF)

Inclusive processes (DIS)

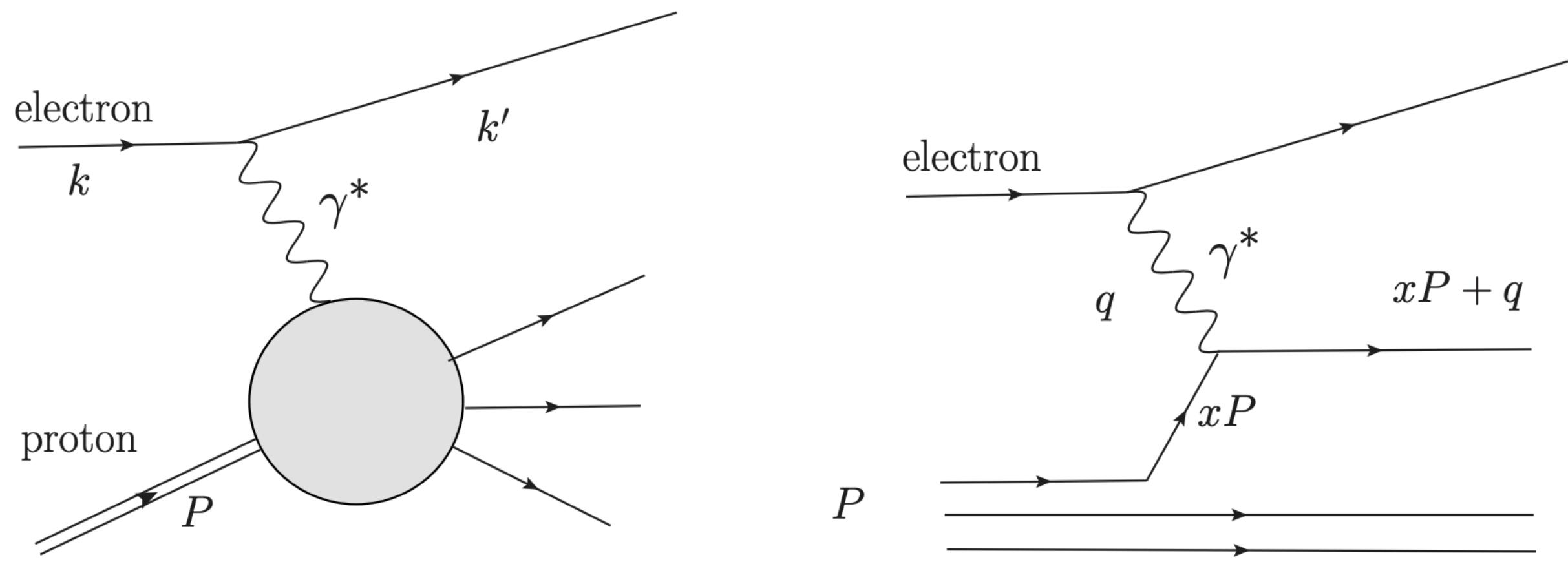


$$Q^2 = -q^2$$

$$\frac{d^2\sigma}{d\Omega dE} = \frac{\alpha^2}{Q^4} \frac{E'}{E} l_{\mu\nu} W^{\mu\nu}$$

Parton Distribution Function (PDF)

Inclusive processes (DIS)



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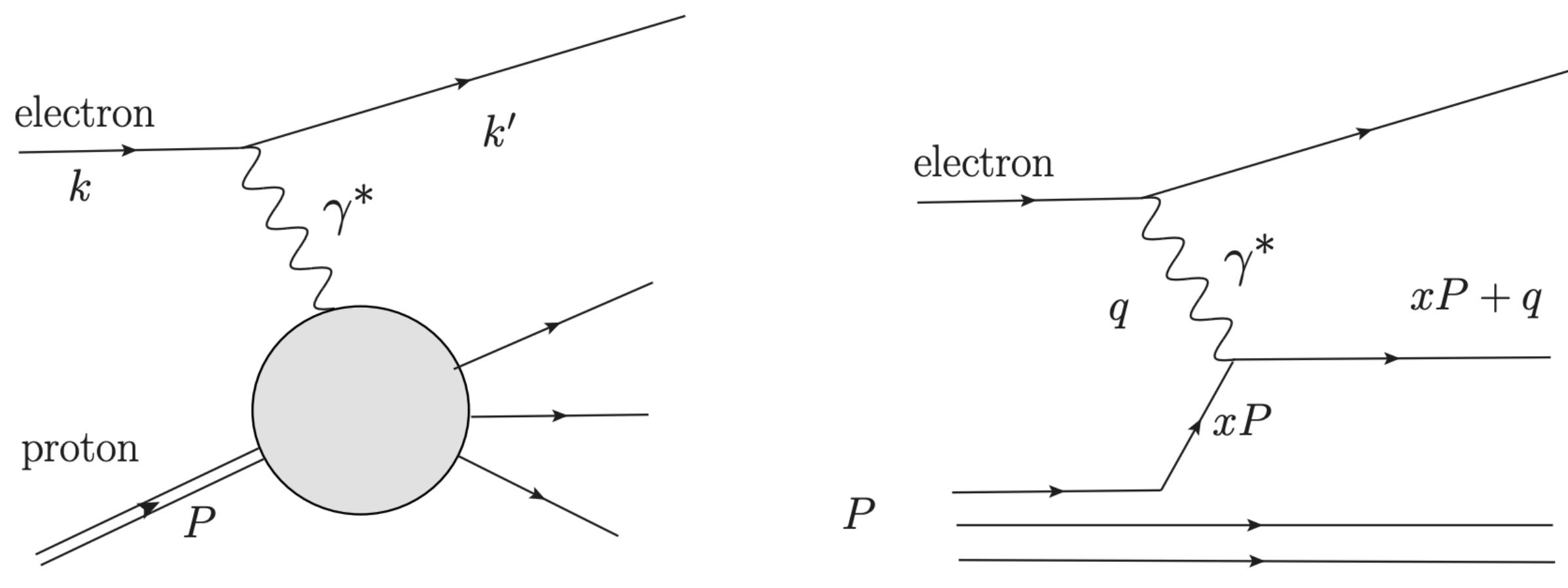
$$\frac{d^2\sigma}{d\Omega dE} = \frac{\alpha^2}{Q^4} \frac{E'}{E} l_{\mu\nu} W^{\mu\nu}$$

leptonic tensor

hadronic tensor

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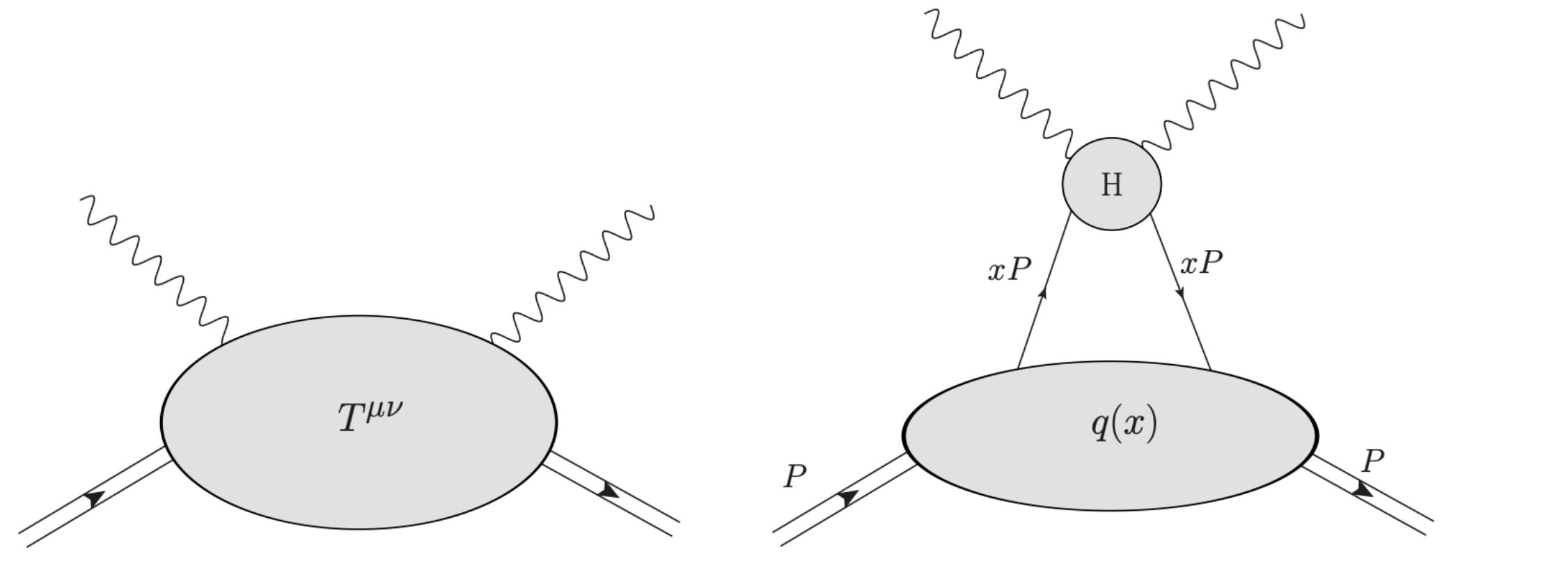
leptonic tensor

hadronic tensor

Optical theorem: $W^{\mu\nu} = \Im T^{\mu\nu}$ Compton tensor

$$T^{\mu\nu}(q, P) = \frac{i}{4\pi} \int d^4x e^{iq \cdot x} \sum_{\sigma} \langle P, \sigma | T J^{\mu}(x) J^{\nu}(0) | P, \sigma \rangle$$

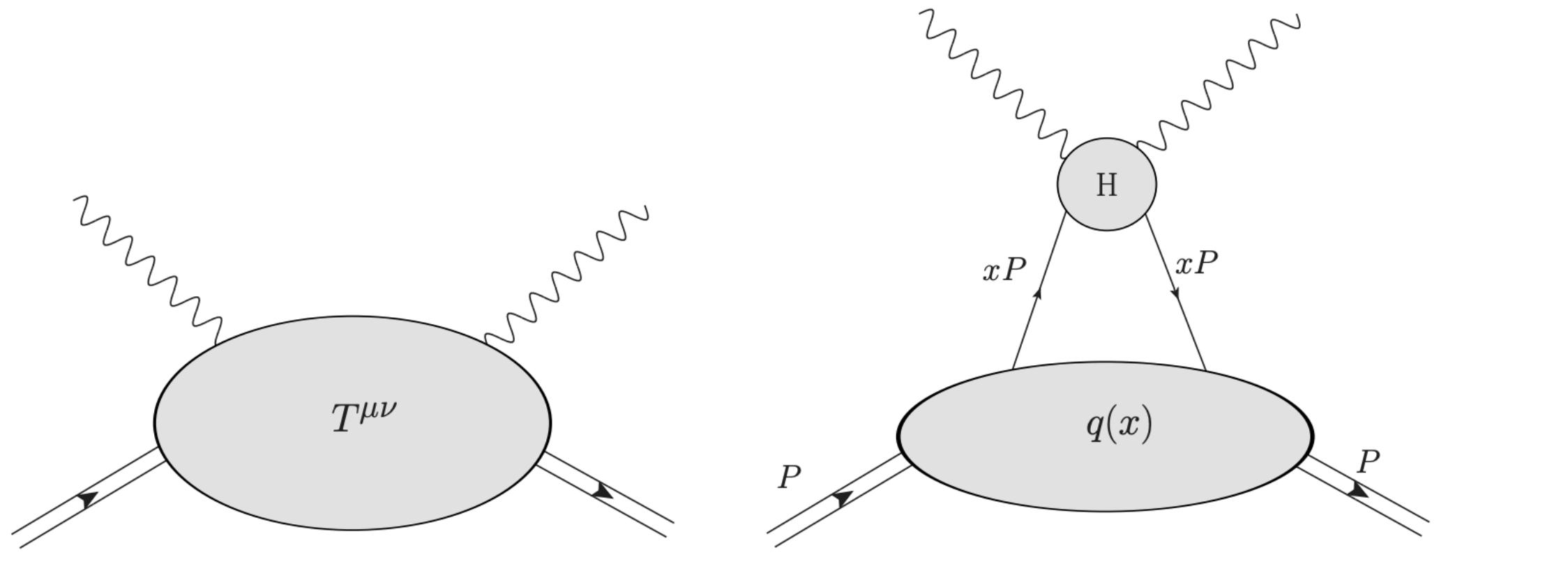
$$W^{\mu\nu}(q, P) = -\textcolor{teal}{W}_1 \left(g^{\mu\nu} - \frac{q^\mu q^\nu}{q^2} \right) + \textcolor{teal}{W}_2 \frac{1}{M^2} \left(P^\mu - \frac{P \cdot q}{q^2} q^\mu \right) \left(P^\nu - \frac{P \cdot q}{q^2} q^\nu \right)$$



Factorization

$$Q^2 \rightarrow \infty , \quad x_B \equiv \frac{Q^2}{2P \cdot q}$$

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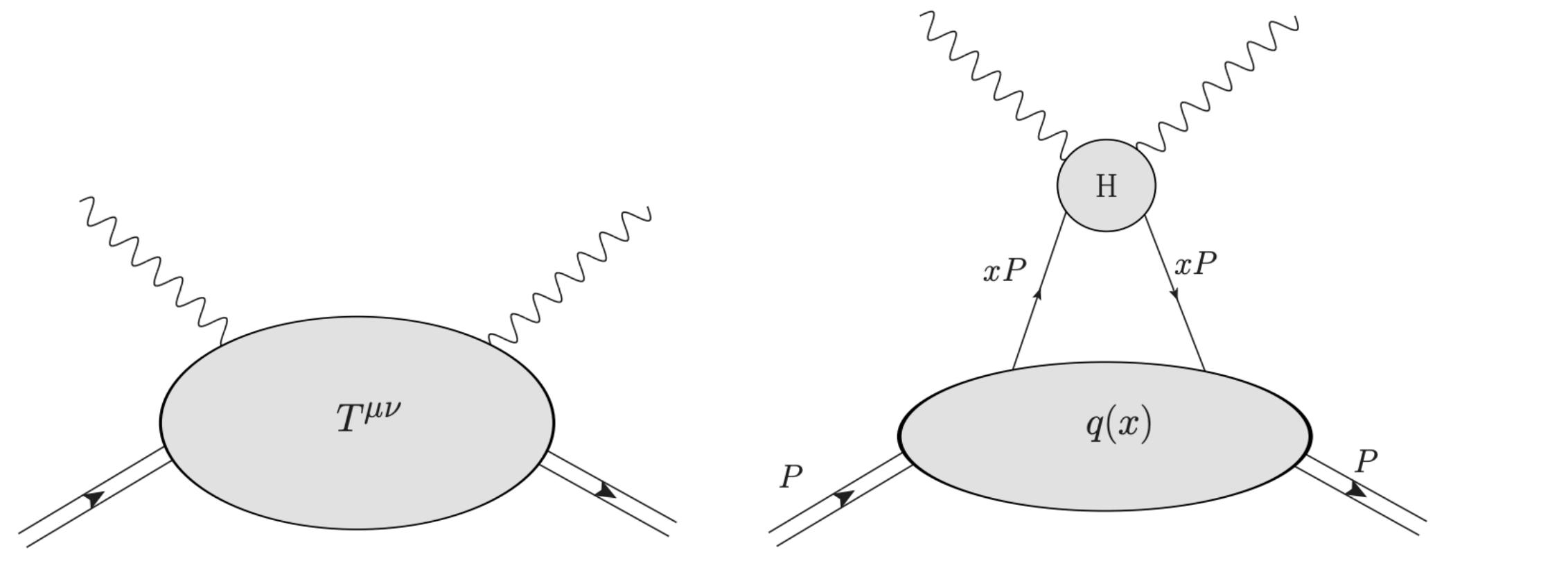


Factorization

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$$\textcolor{teal}{W}_k(x_B, Q^2) = \sum_{p=q,g} \int_0^1 dx C_k^{(i)}(x, x_B, Q^2, \mu^2) \textcolor{red}{q}_p(x, \mu^2) + O\left(1/Q^2\right)$$

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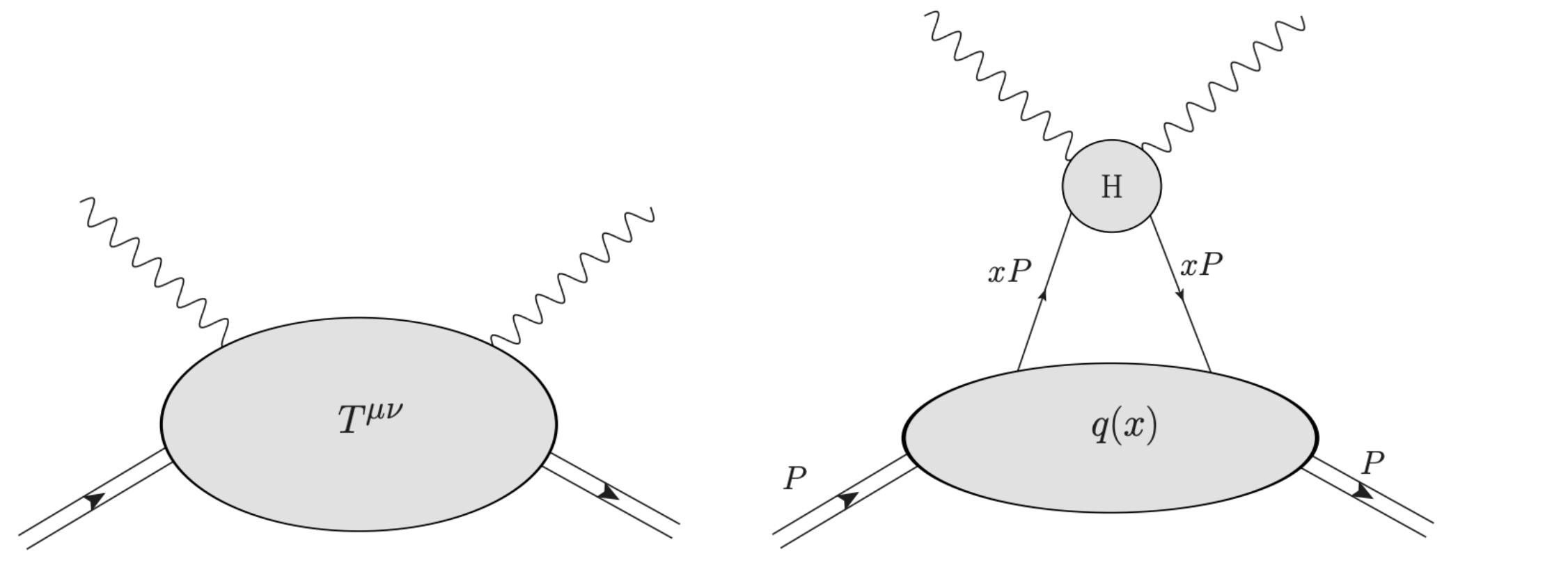
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Probability density of finding a parton with longitudinal momentum fraction x

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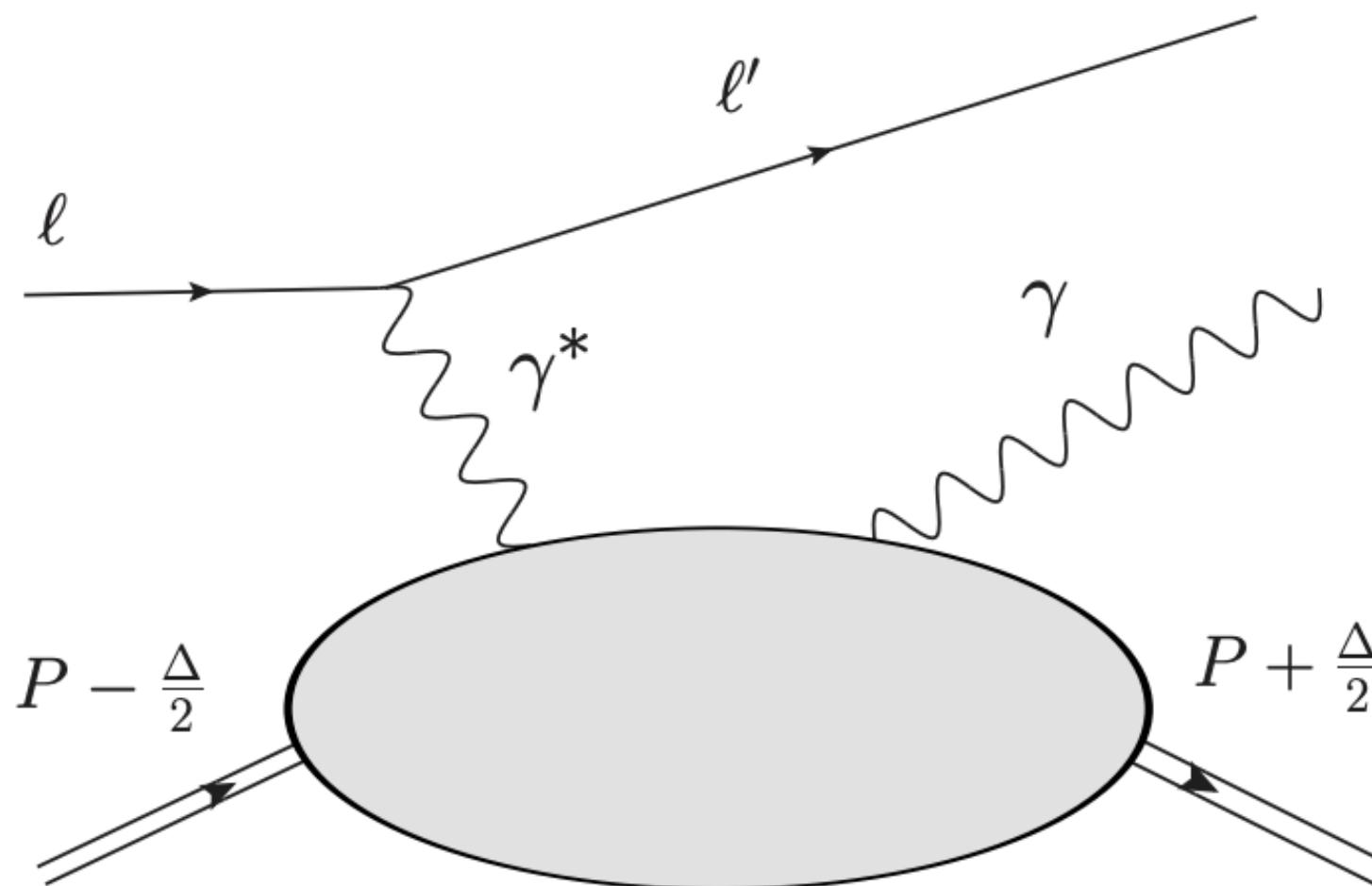
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No Wilson line in light-cone gauge $A^+ = 0$

Probability density of finding a parton with longitudinal momentum fraction x

Generalized Parton Distribution (GPD)

Exclusive processes (DVCS)



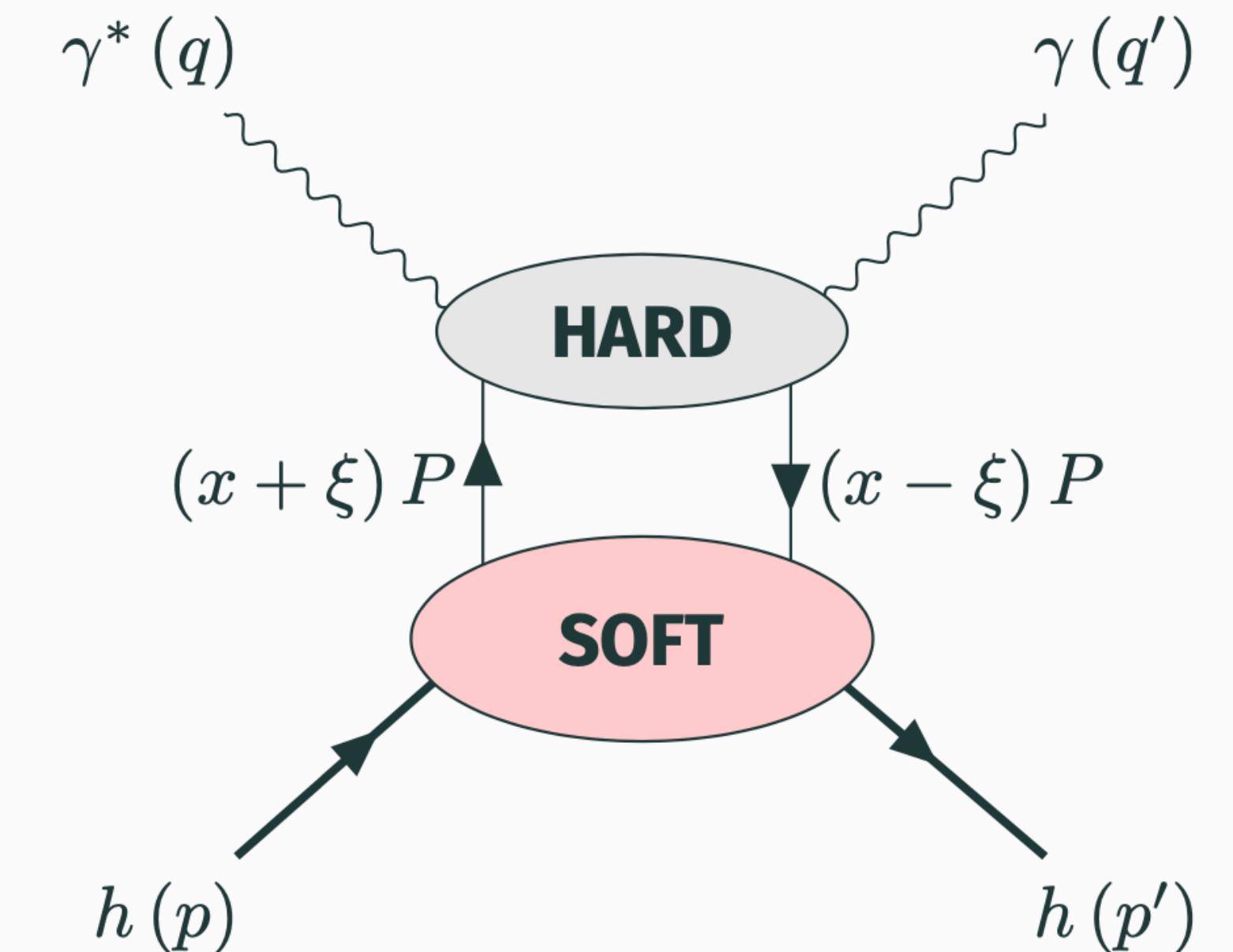
$$Q^2 \rightarrow \infty$$

➡

$$P = \frac{p + p'}{2}, \quad \Delta = p' - p$$

$$\xi = \frac{(p' - p)^+}{2P^+} = -\frac{\Delta^+}{2P^+}$$

$$t = \Delta^2$$



Factorization of the amplitude

$$\mathcal{M}(\xi, t; Q^2) = \sum_{p=q,g} \int_{-1}^1 \frac{dx}{\xi} K^p \left(\frac{x}{\xi}, \frac{Q^2}{\mu_F^2}, \alpha_s(\mu_F) \right) \mathcal{H}^p(x, \xi, t; \mu_F) + O(1/Q^2)$$

↓

hard / perturbative soft / non perturbative GPD

Ex: Chiral even twist-2 operators

Nucleon quark GPD

$$\begin{aligned}\mathcal{H}^q(x, \xi, t) &= \frac{1}{2} \int dz^- e^{ixP^+z^-} \langle P + \Delta/2 | \bar{\psi}^q(-z/2) \gamma^+ \psi^q(z/2) | P - \Delta/2 \rangle \Big|_{z^+=z_\perp=0} \\ &= \frac{1}{2P^+} \left[\mathbf{H}^q(x, \xi, t) \bar{u}(p') \gamma^+ u(p) + \mathbf{E}^q(x, \xi, t) \bar{u}(p') \frac{i\sigma^{+\mu} \Delta_\mu}{2M} u(p) \right]\end{aligned}$$

Nucleon gluon GPD

$$\begin{aligned}\mathcal{H}^g(x, \xi, t) &= \frac{1}{2} \int dz^- e^{ixP^+z^-} \langle P + \Delta/2 | G^{\mu+}(-z/2) G_\mu^+(z/2) | P - \Delta/2 \rangle \Big|_{z^+=z_\perp=0} \\ &= \frac{1}{2P^+} \left[\mathbf{H}^g(x, \xi, t) \bar{u}(p') \gamma^+ u(p) + \mathbf{E}^g(x, \xi, t) \bar{u}(p') \frac{i\sigma^{+\mu} \Delta_\mu}{2M} u(p) \right]\end{aligned}$$

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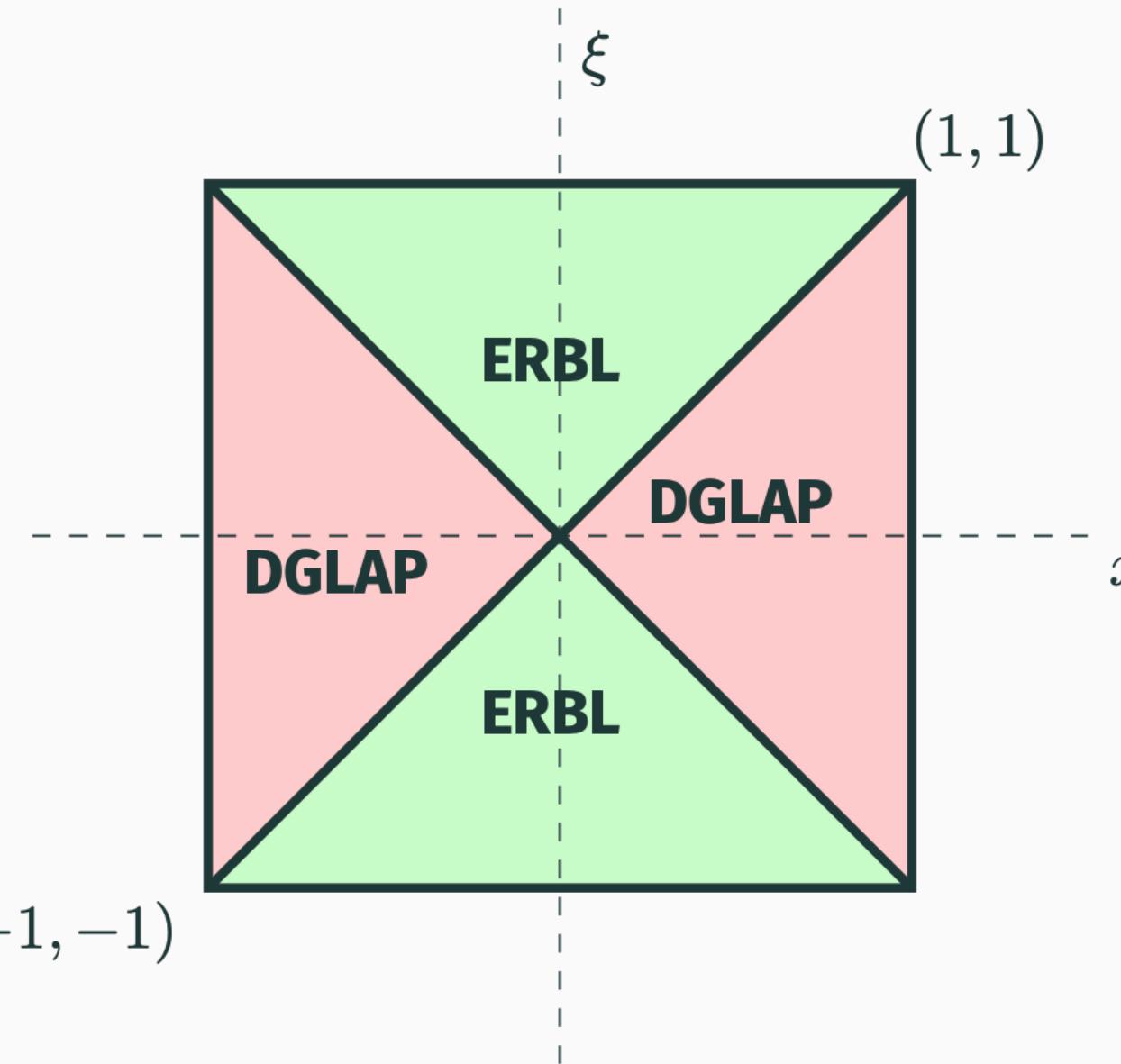
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quark GPD of spin zero hadron

$$H_\pi^q(x, \xi, t) = \frac{1}{2} \int \frac{dz^-}{2\pi} e^{ixP^+z^-} \langle \pi(P + \Delta/2) | \bar{\psi}^q(-z/2) \gamma^+ \psi^q(z/2) | \pi(P - \Delta/2) \rangle \Big|_{z^+=0, z^\perp=0}$$

GPD properties

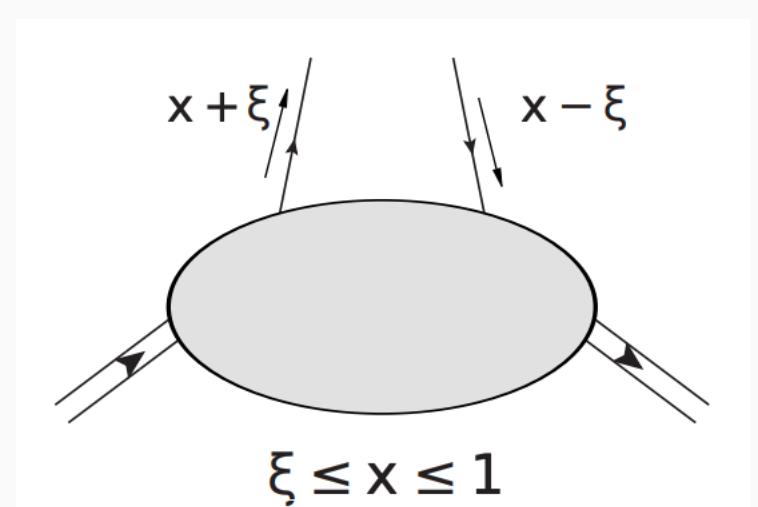
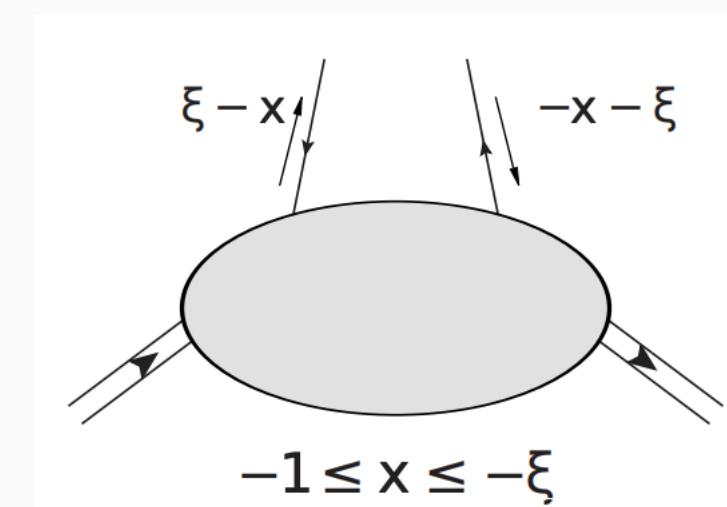
- **Support** $(x, \xi) \in [-1,1] \otimes [-1,1]$



DGLAP

$$|x| > |\xi|$$

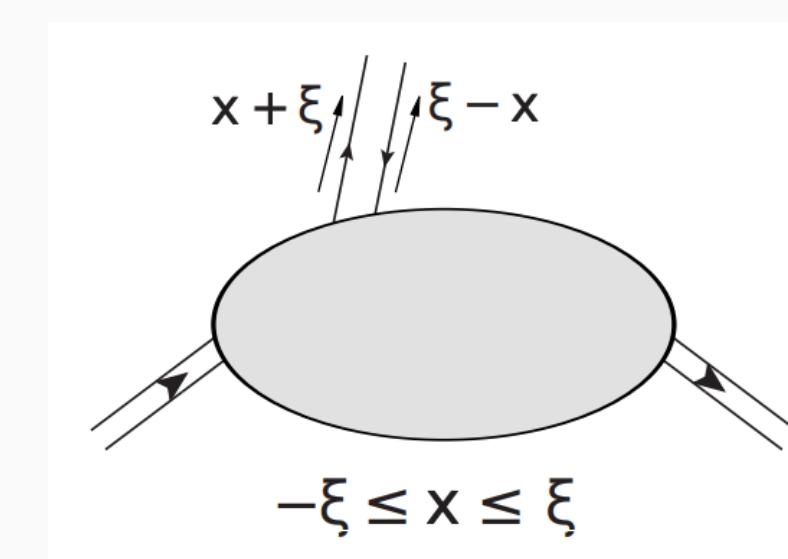
Emission/absorption quark ($x > 0$) or antiquark ($x < 0$)



ERBL

$$|x| < |\xi|$$

Emission of quark/antiquark pair



- **ξ -parity**

$$H^q(x, -\xi, t) = H^q(x, \xi, t) \quad \text{Time inversion symmetry}$$

- **Polynomiality**

$$\mathcal{A}_m(\xi, t) = \int_{-1}^1 dx x^m H^q(x, \xi, t) = \sum_{\substack{k=0 \\ k \text{ even}}}^{m+1} C_{k,m}(t) \xi^k \quad \text{Lorentz symmetry}$$

- **Positivity**

$$|H^q(x, \xi, t=0)| \leq \sqrt{q \left(\frac{x+\xi}{1+\xi} \right) q \left(\frac{x-\xi}{1-\xi} \right)}, \quad |x| > |\xi| \quad \text{Hilbert space norm}$$

- **Form factors**

$$\int_{-1}^1 dx H^q(x, \xi, t) = F^q(t), \quad \int_{-1}^1 dx x H^q(x, \xi, t) = A^q(t) + \xi^2 C^q(t)$$

- Forward limit

$$H^q(x,0,0) = q(x)\theta(x) - \bar{q}(-x)\theta(-x)$$

- Hadron 3D tomography

$$\rho^q(x, b_\perp) = \int \frac{d^2 \Delta_\perp}{(2\pi)^2} e^{-ib_\perp \cdot \Delta_\perp} H^q(x, 0, -\Delta_\perp^2)$$

Probability density of finding a parton with longitudinal momentum fraction x and position b_\perp in the transverse plane.

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It is hard to build GPDs from first principles that satisfy both **positivity** and **polynomiality**

Truncation in Overlap representation of LCWF is consistent only in **DGLAP** region

$$H^q(x, \xi, t) \Big|_{|x| > |\xi|} = \sum_{N,\beta} \sqrt{1 - \xi^2}^{2-N} \int [d\bar{x}]_N [d^2 \bar{\mathbf{k}}_\perp]_N \delta(x - \bar{x}_i) \psi_{N,\beta}^*(x_i^{out}, \mathbf{k}_{i\perp}^{out}) \psi_{N,\beta}(x_i^{in}, \mathbf{k}_{i\perp}^{in})$$

Positivity is built in, but for **polynomiality** GPD must be extended to **ERBL** region

Covariant extension via Radon transform inversion

Is it possible to extend the GPD from DGLAP to ERBL preserving polynomiality?

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$$\langle P + \Delta/2 | \bar{\psi}(-z/2) \gamma_\mu \psi(z/2) | P - \Delta/2 \rangle \Big|_{z^2=0} = \int d\beta d\alpha e^{-i\beta(P \cdot z) + i\alpha \frac{\Delta \cdot z}{2}} \left[2P_\mu f(\beta, \alpha, t) - \Delta_\mu g(\beta, \alpha, t) \right]$$

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Radon transform

- Integral over line parametrized by x, ξ
- It guarantees polynomiality

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Double Distributions (DD)

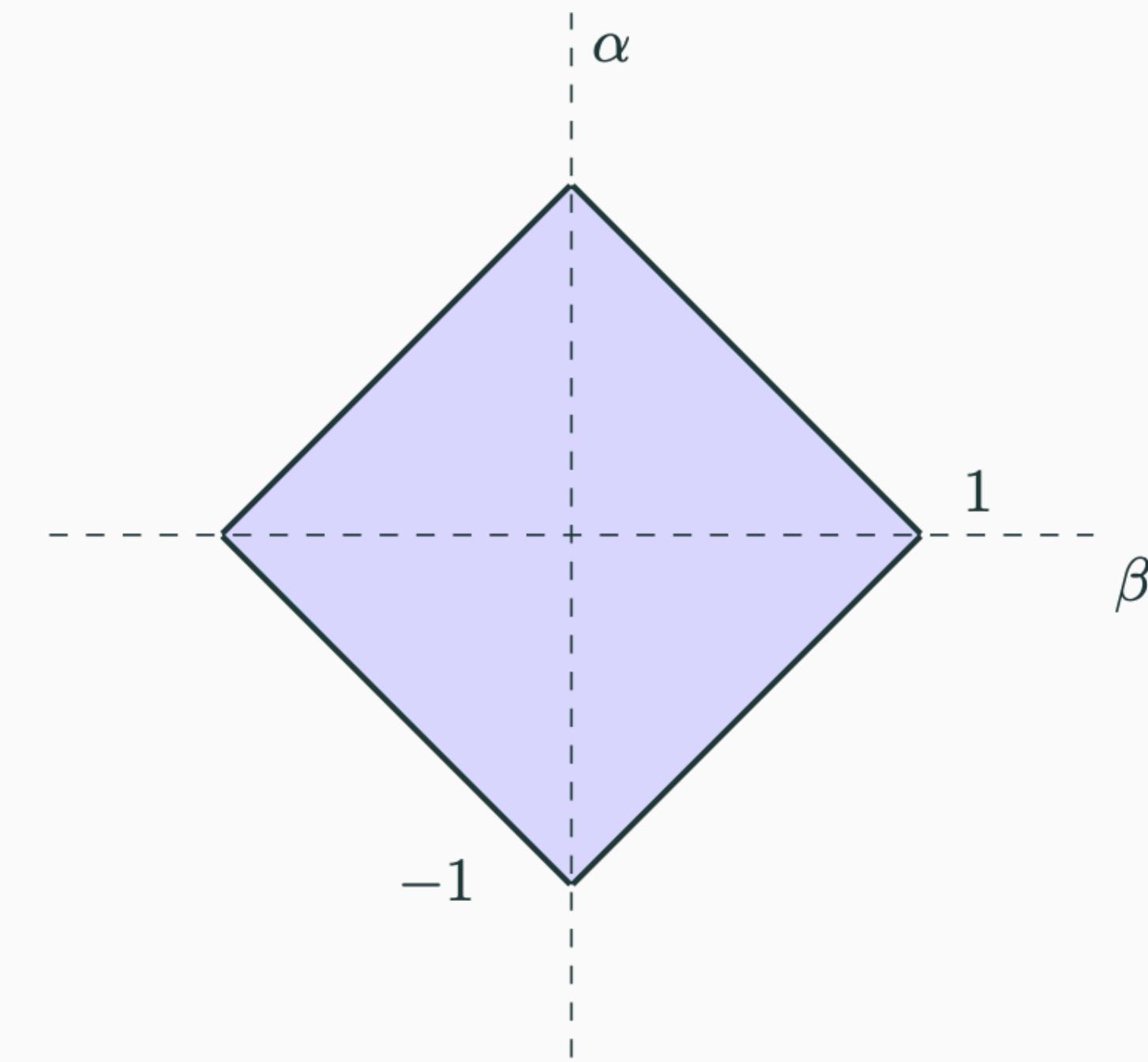
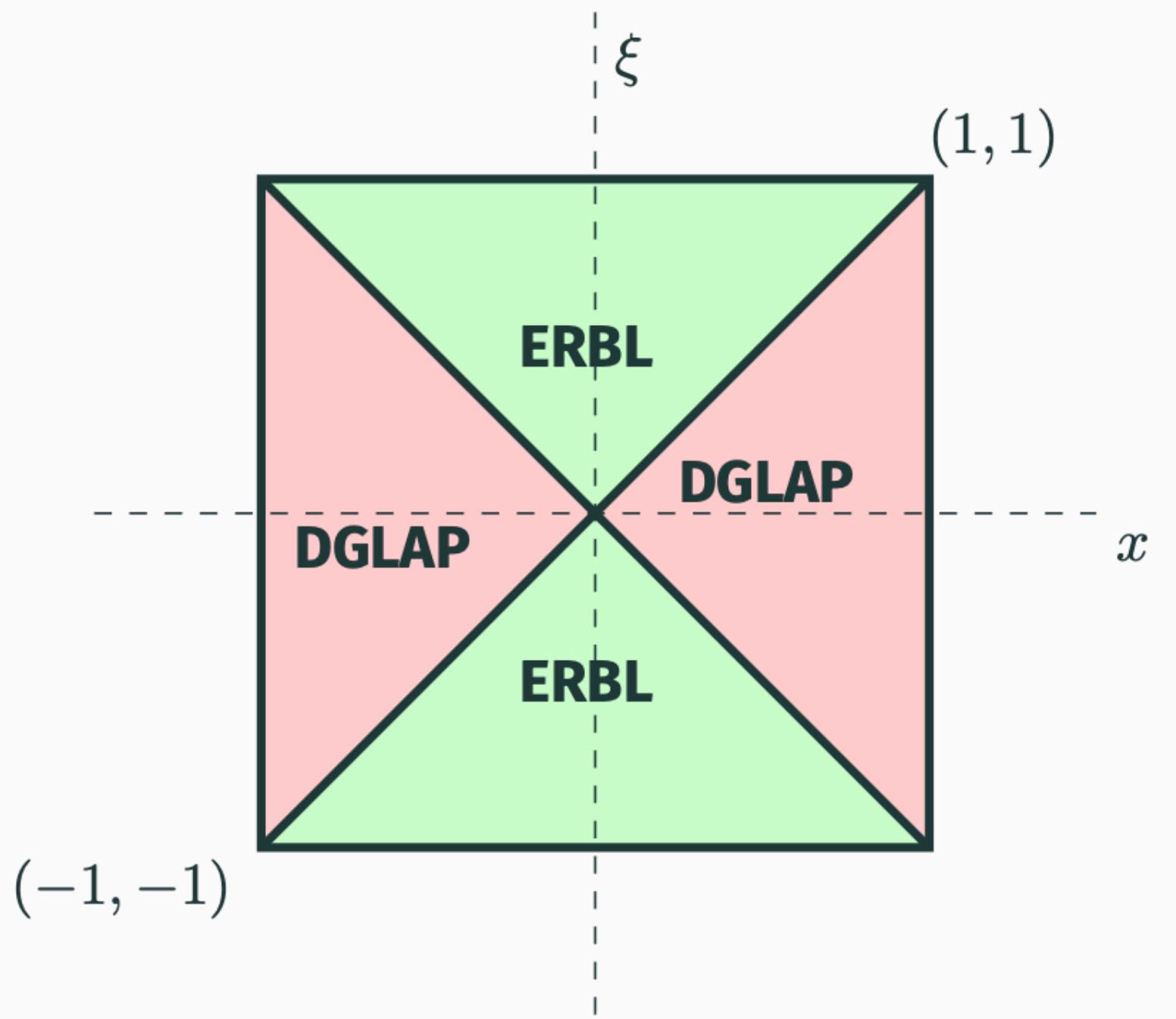
Not uniquely fixed

Radon transform

- Integral over line parametrized by x, ξ
- It guarantees polynomiality

$$H^q(x, \xi) = \mathcal{R}[h] = \int_{\Omega} d\beta d\alpha \delta(x - \beta - \alpha\xi) h(\beta, \alpha) + \theta(|\xi| - |x|) \text{ D-terms}$$

Support: $\Omega = \{(\beta, \alpha) \mid |\beta| + |\alpha| \leq 1\} = \Omega^+ \theta(\beta) + \Omega^- \theta(-\beta)$

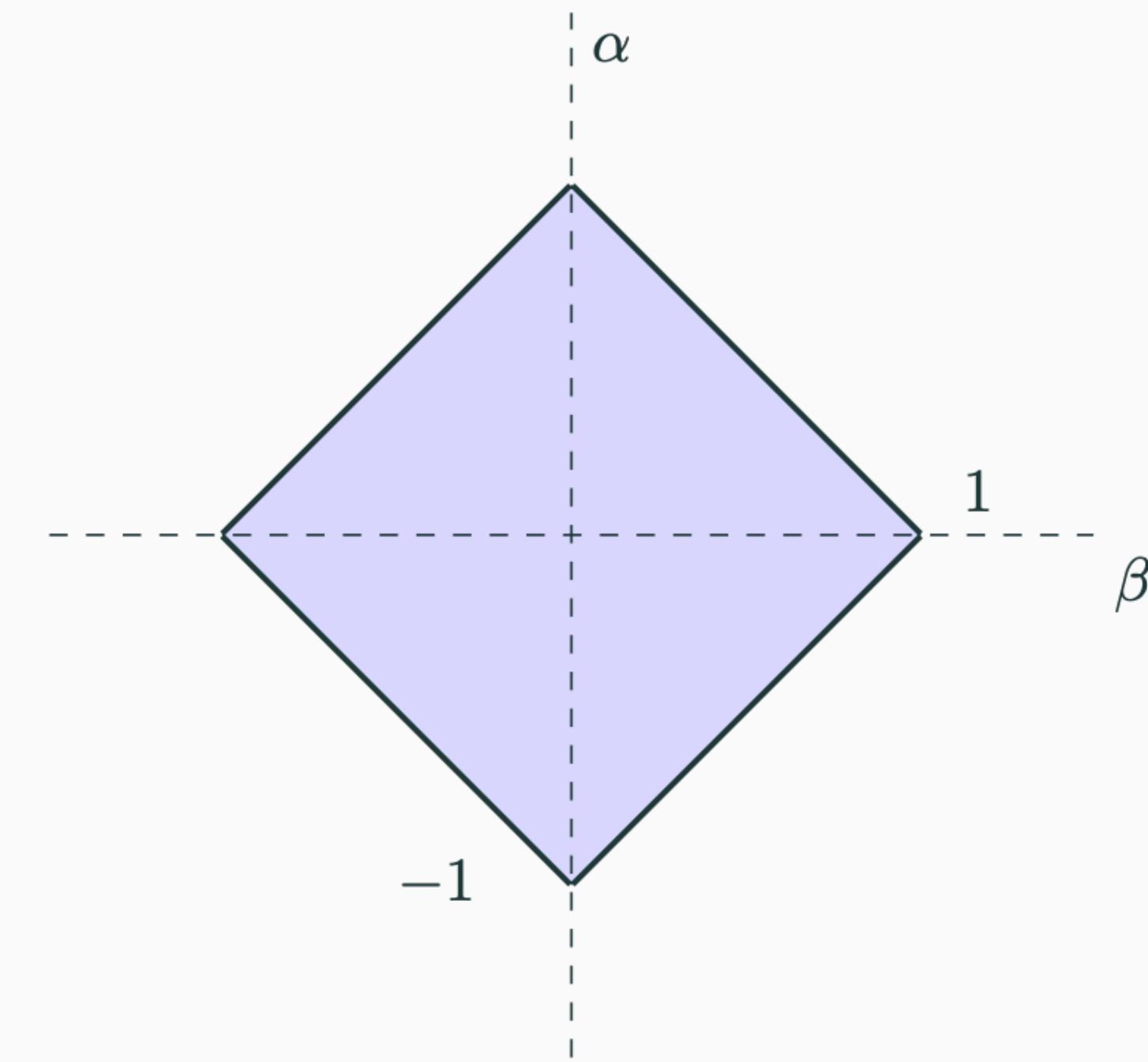
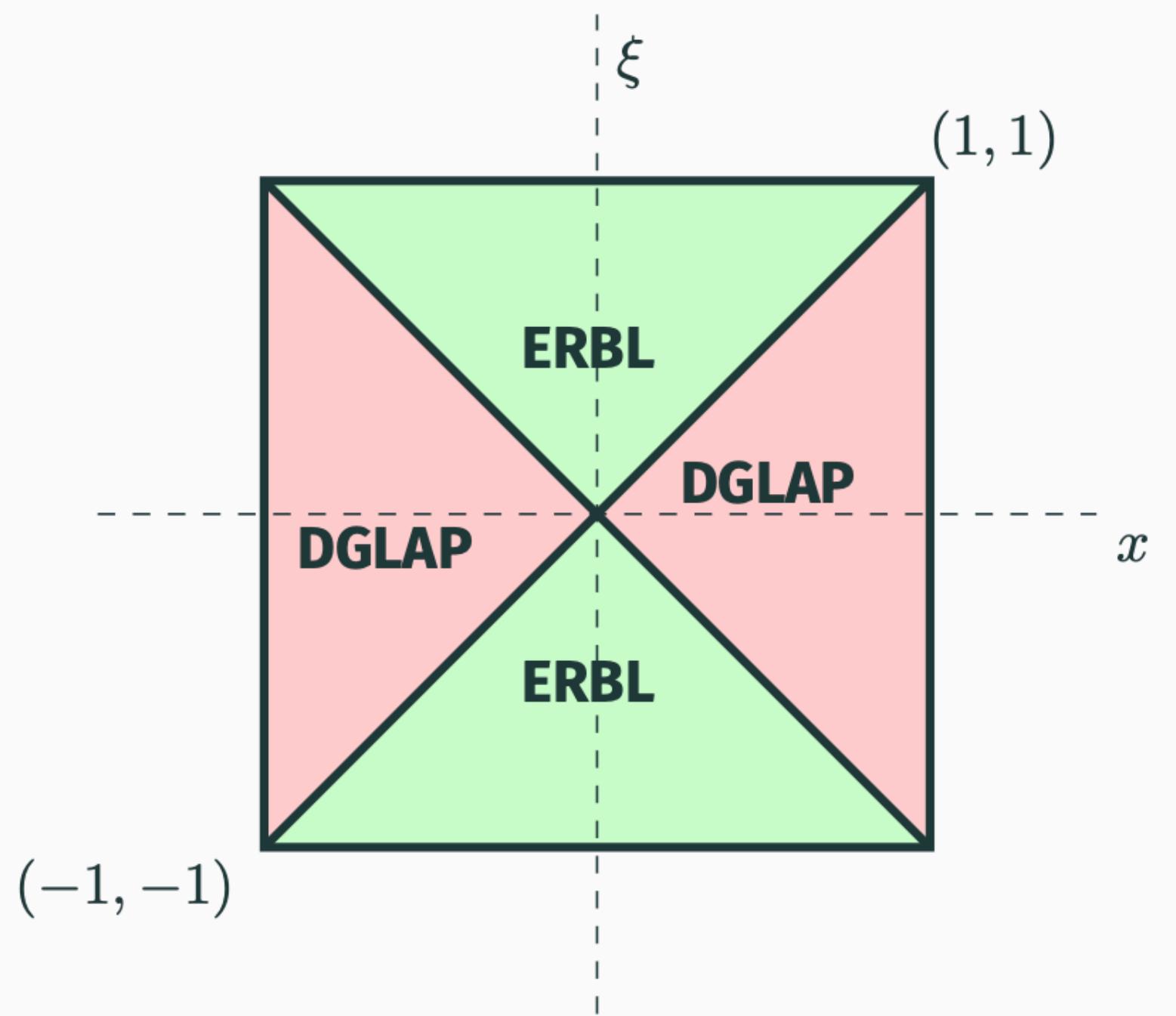


line $\alpha = -\frac{\beta}{\xi} + \frac{x}{\xi}$

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- Cannot be fixed by DGLAP data.
- Its DD has support only in $\beta = 0$
- Contribute only to $C_{m,m+1}(t)$

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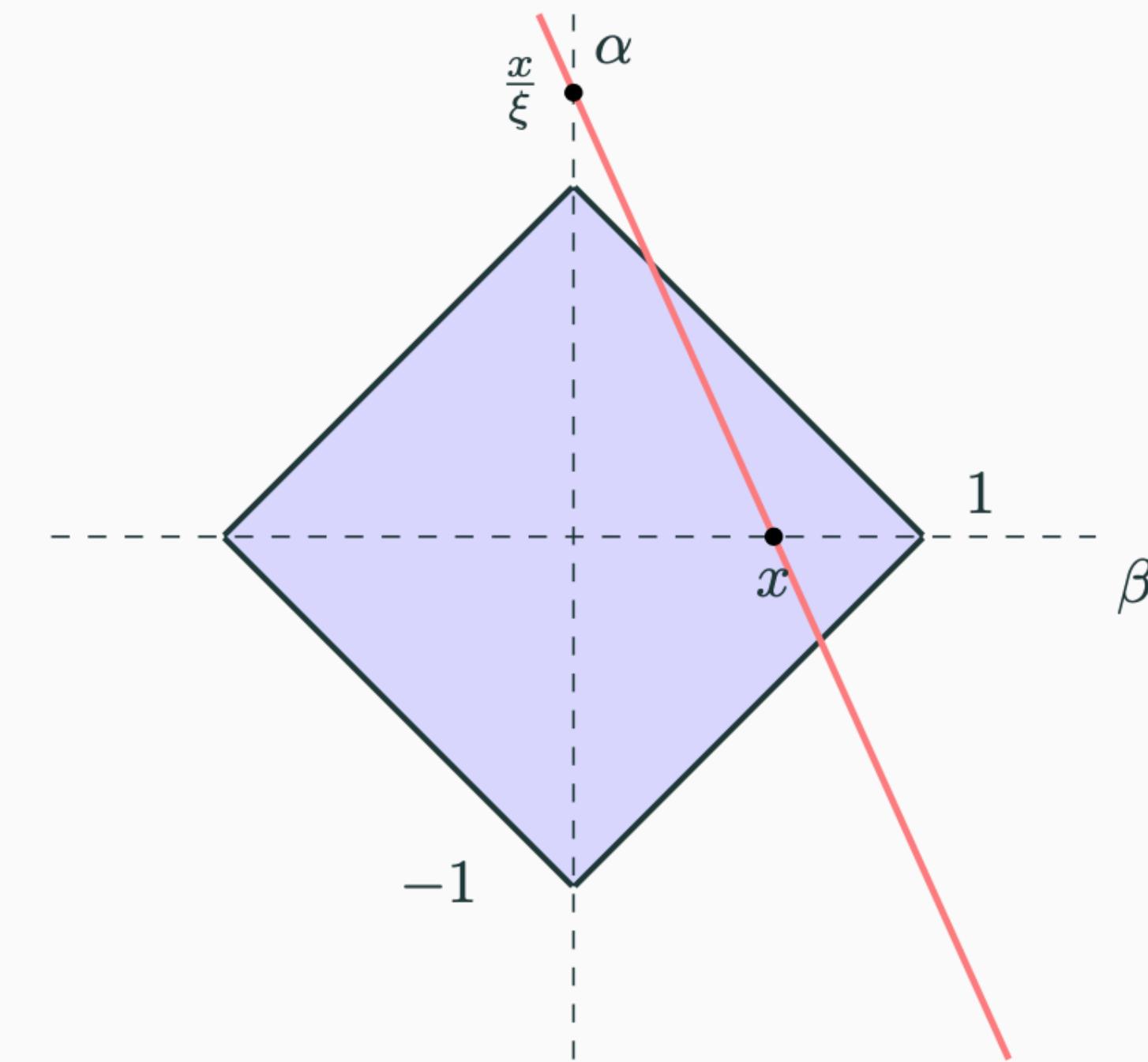
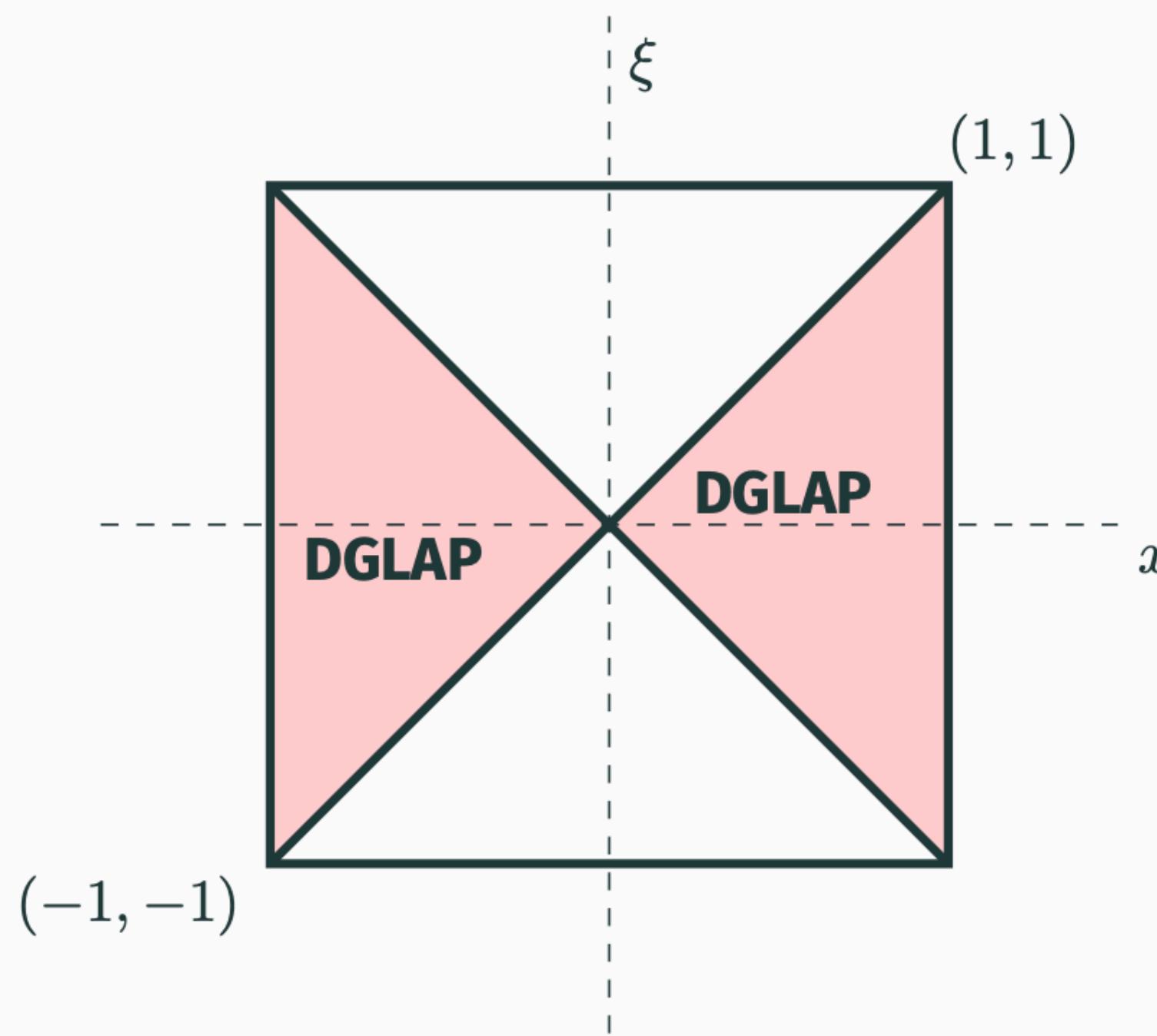


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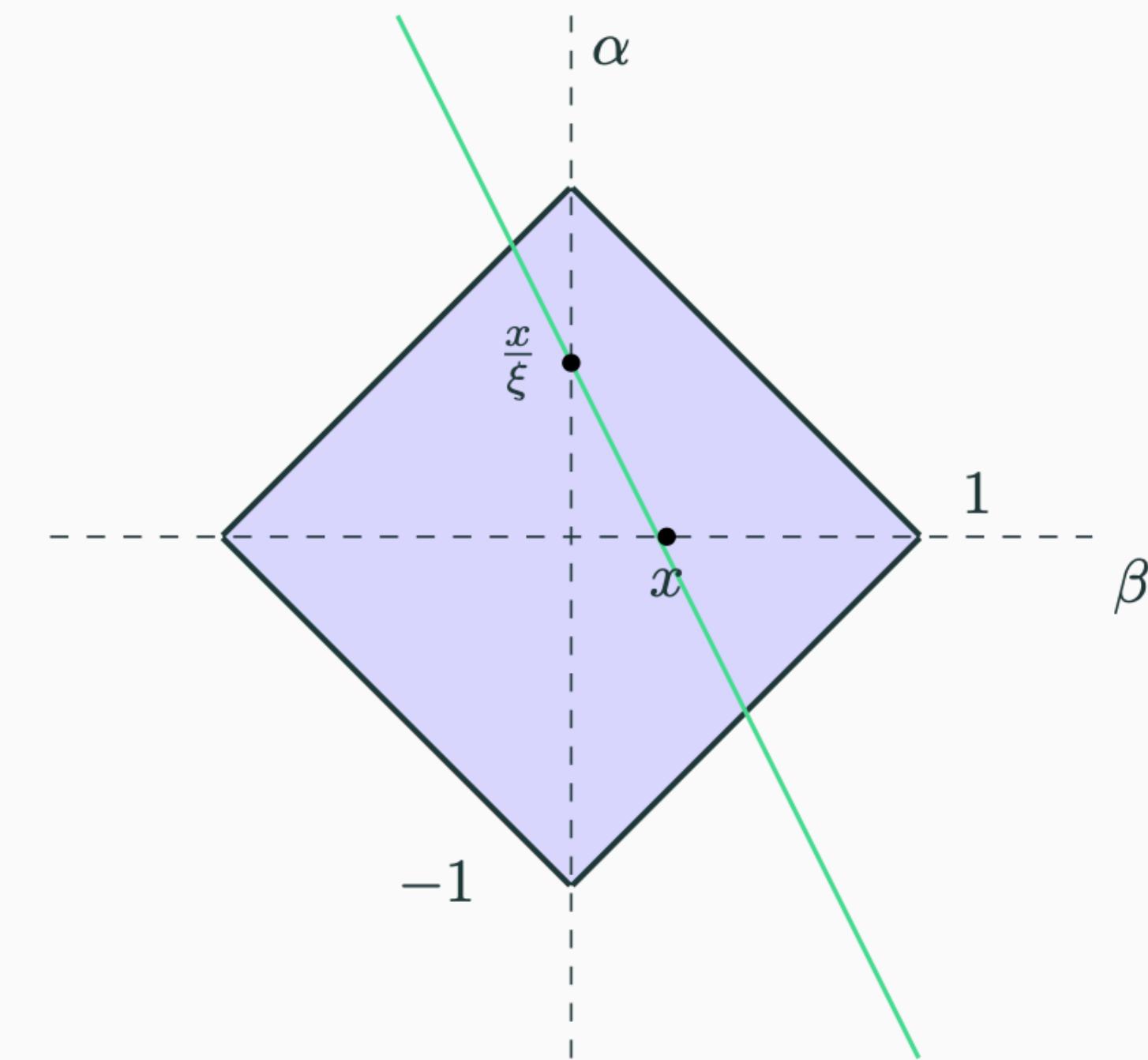
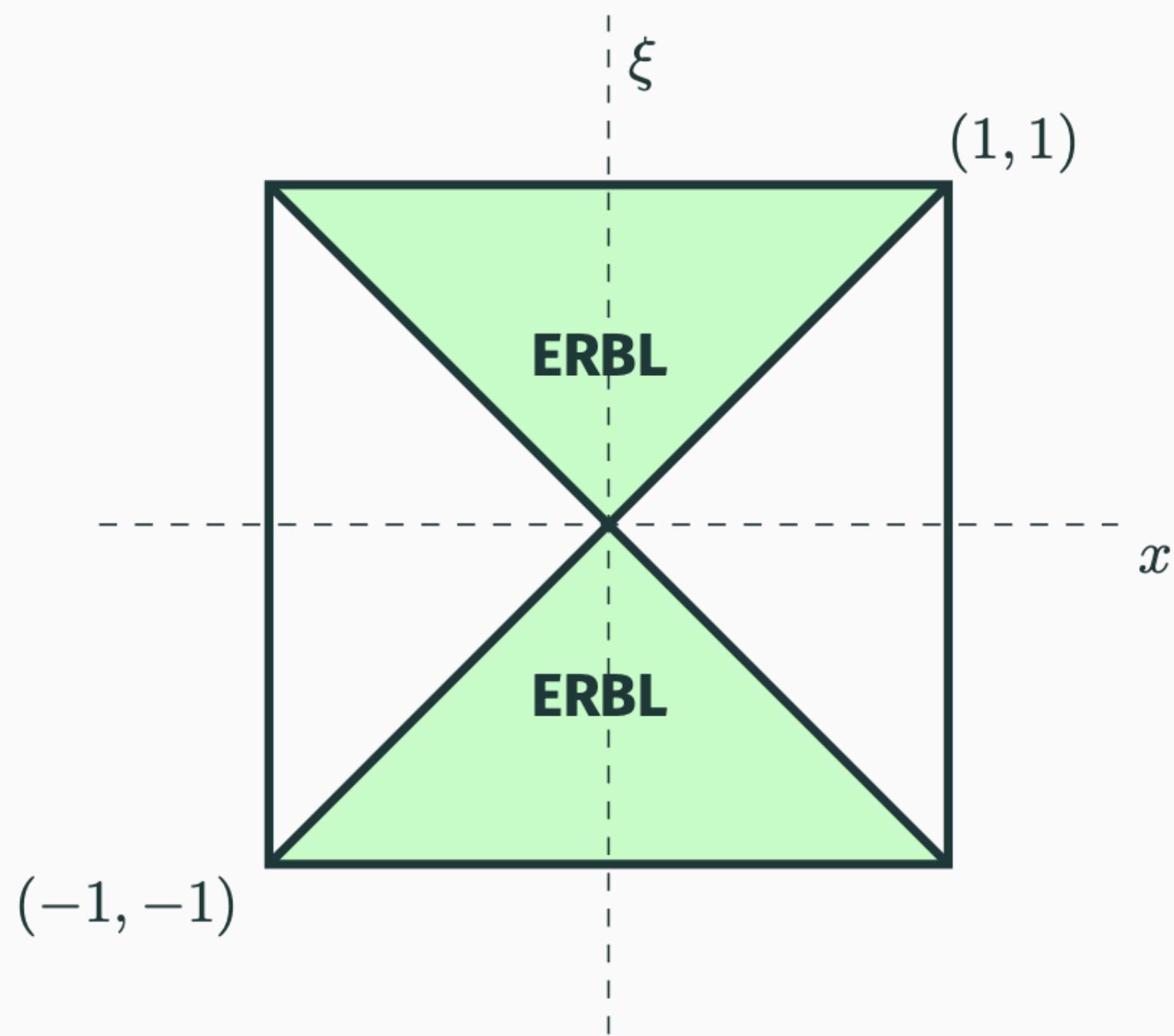


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Covariant Extension



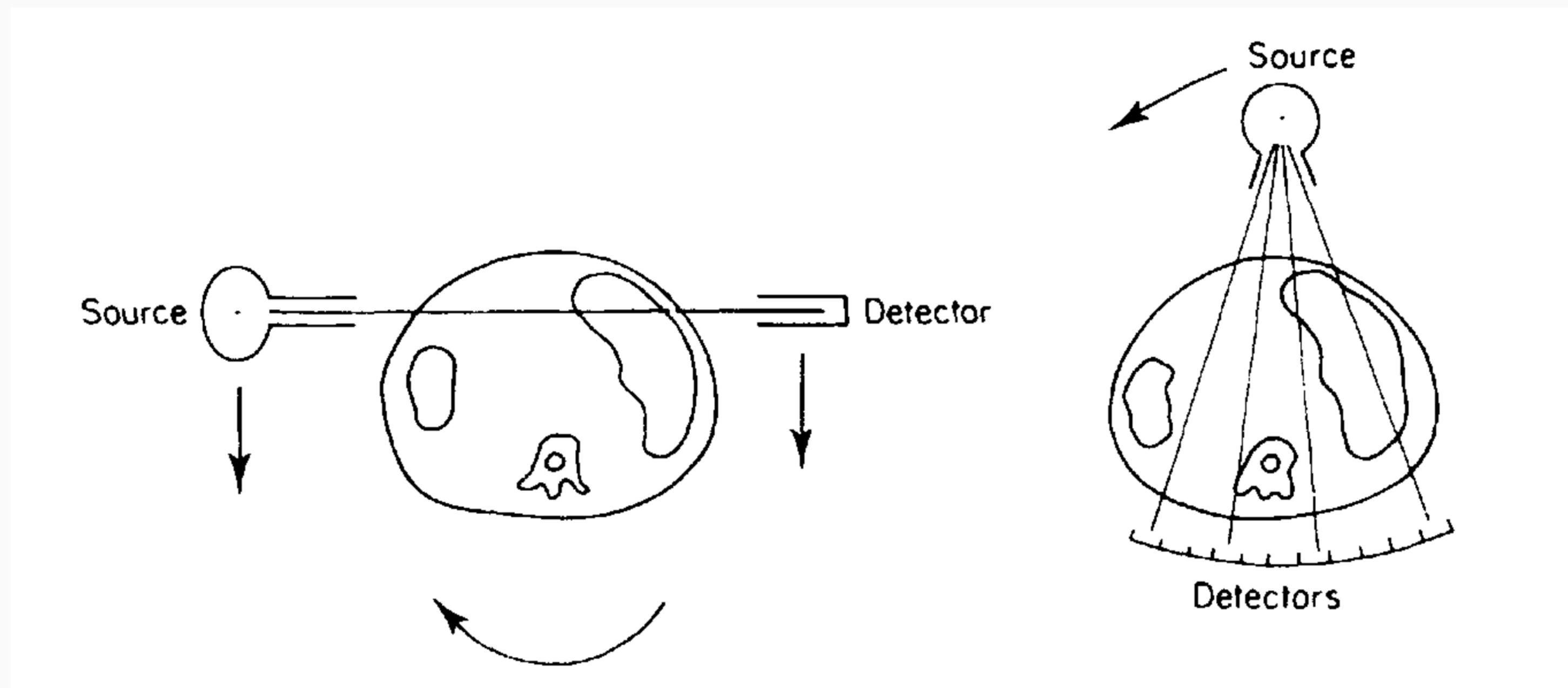
Radon transform inversion

Is it possible to invert the Radon transform knowing the GPD only in DGLAP?

Covariant Extension \longleftrightarrow Radon transform inversion

Is it possible to invert the Radon transform knowing the GPD only in DGLAP?

Answer from Computerized Tomography



Parallel scanning

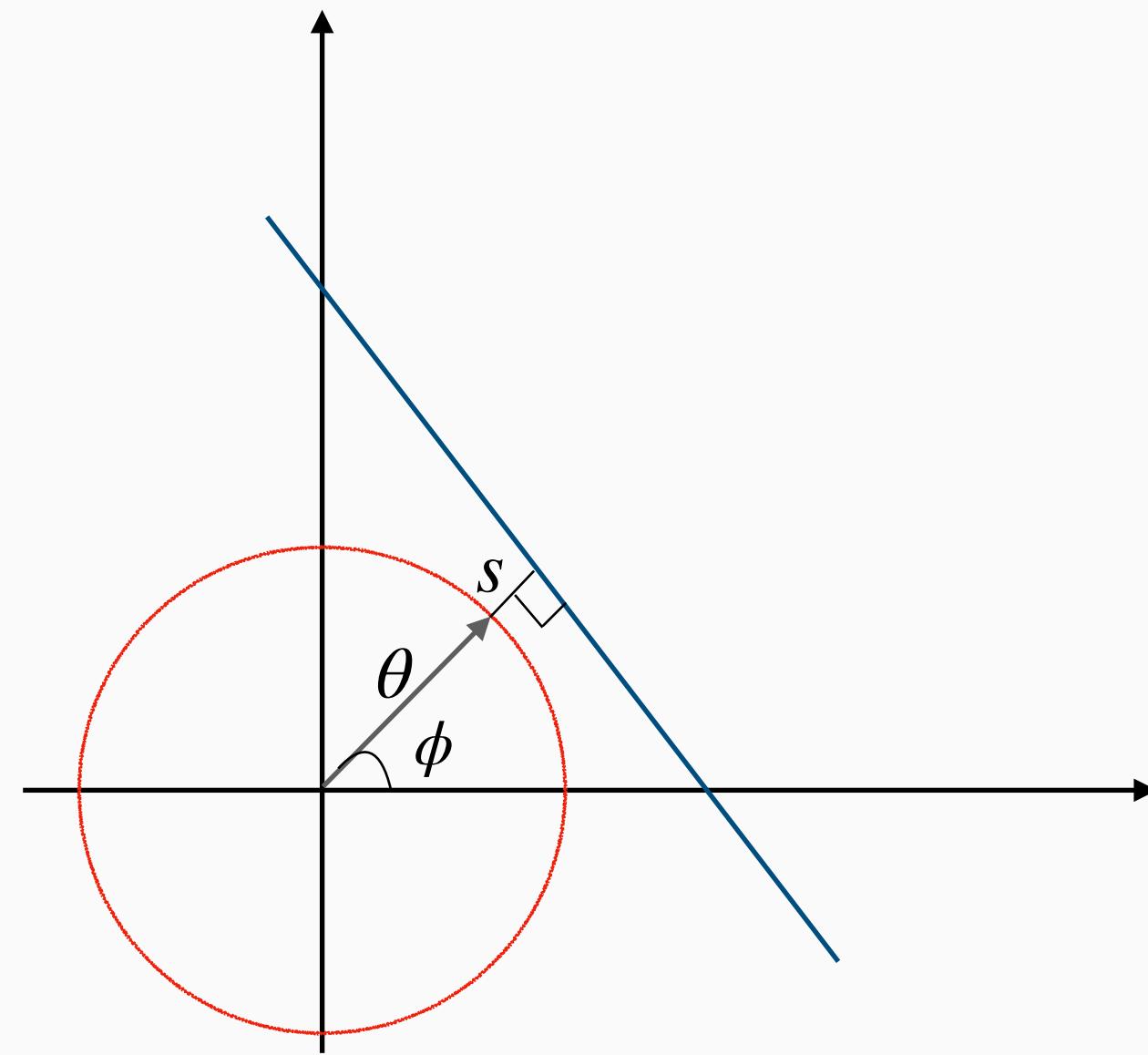
Fan-beam scanning

$$\frac{\Delta I}{I} = f(x) dx$$

$$\frac{I_{out}}{I_{in}} [L] = \exp \left\{ - \int_L f(x) dx \right\}$$

$$\mathcal{R}f(\theta, s) \equiv \int_{z \cdot \theta = s} dz f(z) \quad f \in \mathcal{S}(\mathbb{R}^n), \quad \theta \in S^{n-1}, \quad s \in \mathbb{R}$$

GPD: $z \rightarrow (\beta, \alpha), \quad \theta \rightarrow (\cos(\phi), \sin(\phi)), \quad x = \frac{s}{\cos(\phi)}, \quad \xi = \tan(\phi)$



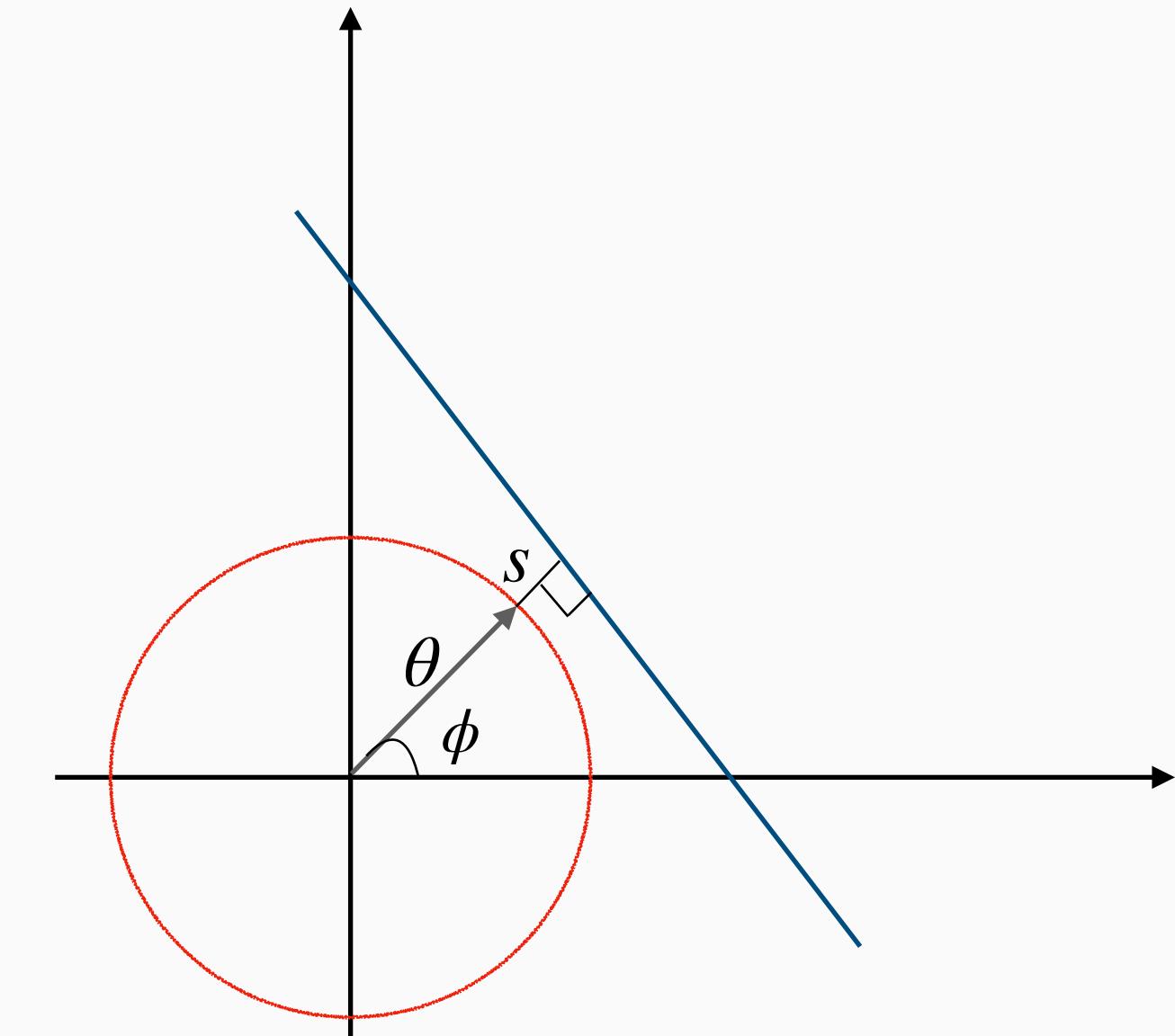
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GPD: $z \rightarrow (\beta, \alpha), \quad \theta \rightarrow (\cos(\phi), \sin(\phi)), \quad x = \frac{s}{\cos(\phi)}, \quad \xi = \tan(\phi)$

Uniqueness theorems:

[F. Natterer, *The Mathematics of Computerized Tomography*]

Given $f \in \mathcal{S}(\mathbb{R}^n)$, if $\mathcal{R}f(\theta, s) = 0 \quad \forall$ hyperplane $z \cdot \theta = s$ that does not intersect a compact convex set $K \subset \mathbb{R}^n \Rightarrow f(z) = 0$ outside K

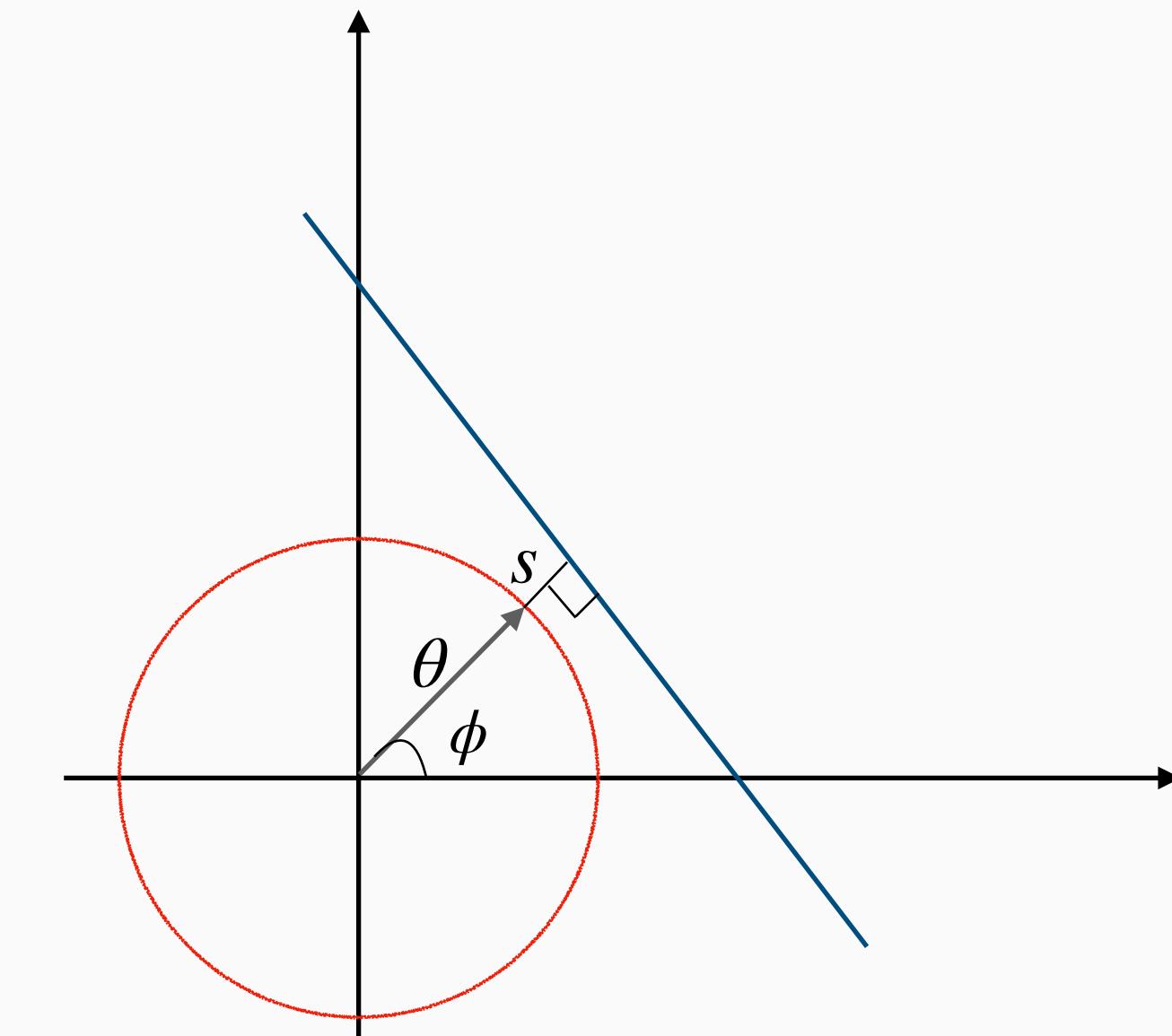


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→ DD is uniquely fixed by GPD in DGLAP (except for $\beta = 0$)

Given f in \mathbb{R}^2 compactly-supported and locally summable, $(\theta_0, s_0) \in S^1 \times \mathbb{R}$, U_0 open neighborhood of θ_0 :

$$\mathcal{R}f(\theta, s) = 0 \quad \forall s > s_0, \theta \in U_0 \Rightarrow f(z) = 0 \quad \text{if } z \cdot \theta_0 > s_0$$

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→ DD is uniquely fixed (except for $\beta = 0$) by knowing GPD in:

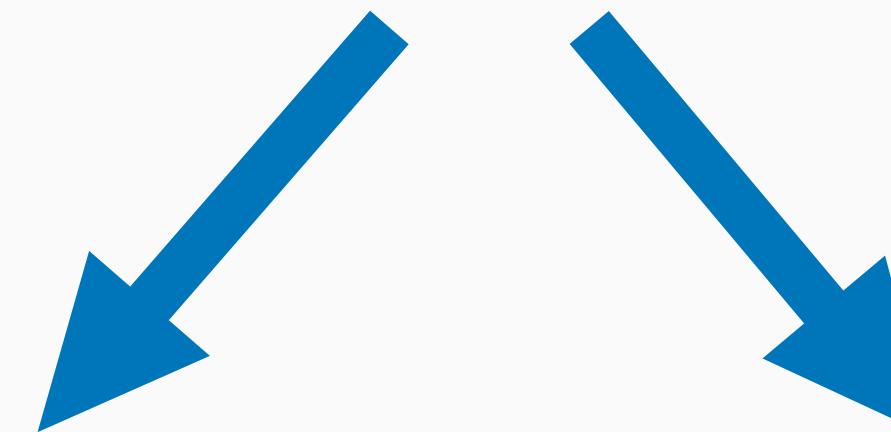
$$x \in [-1, 1], \quad \xi \in [0, \lambda x], \quad 0 < \lambda \leq 1 \quad (\text{H. Moutarde})$$

In experiments $\xi \lesssim 0.2$

Numerical Inversion of Radon Transform

Artificial Neural Networks

H. Dutrieux et al. *Eur.Phys.J.C.* 82 (2022)



Finite Elements Methods

N. Chouika et al. *Eur.Phys.J.C.* 77 (2017)

J.M. Morgado et al. *Phys.Rev.D* 105 (2022)

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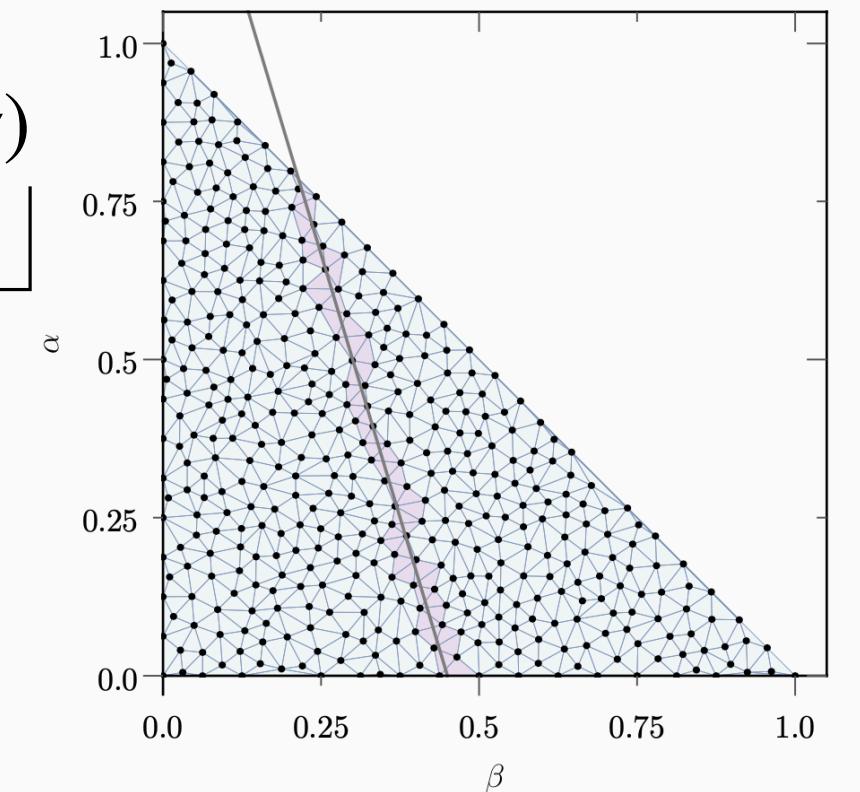
J.M. Morgado et al. *Phys.Rev.D* 105 (2022)

$$H_i \equiv H(x_i, \xi_i) = \sum_j h_j \int_{\Omega^+} d\beta d\alpha \delta(x_i - \beta - \alpha \xi_i) v_j(\beta, \alpha)$$

↑
Lagrange Polynomial
↓

$$R_{ij}$$

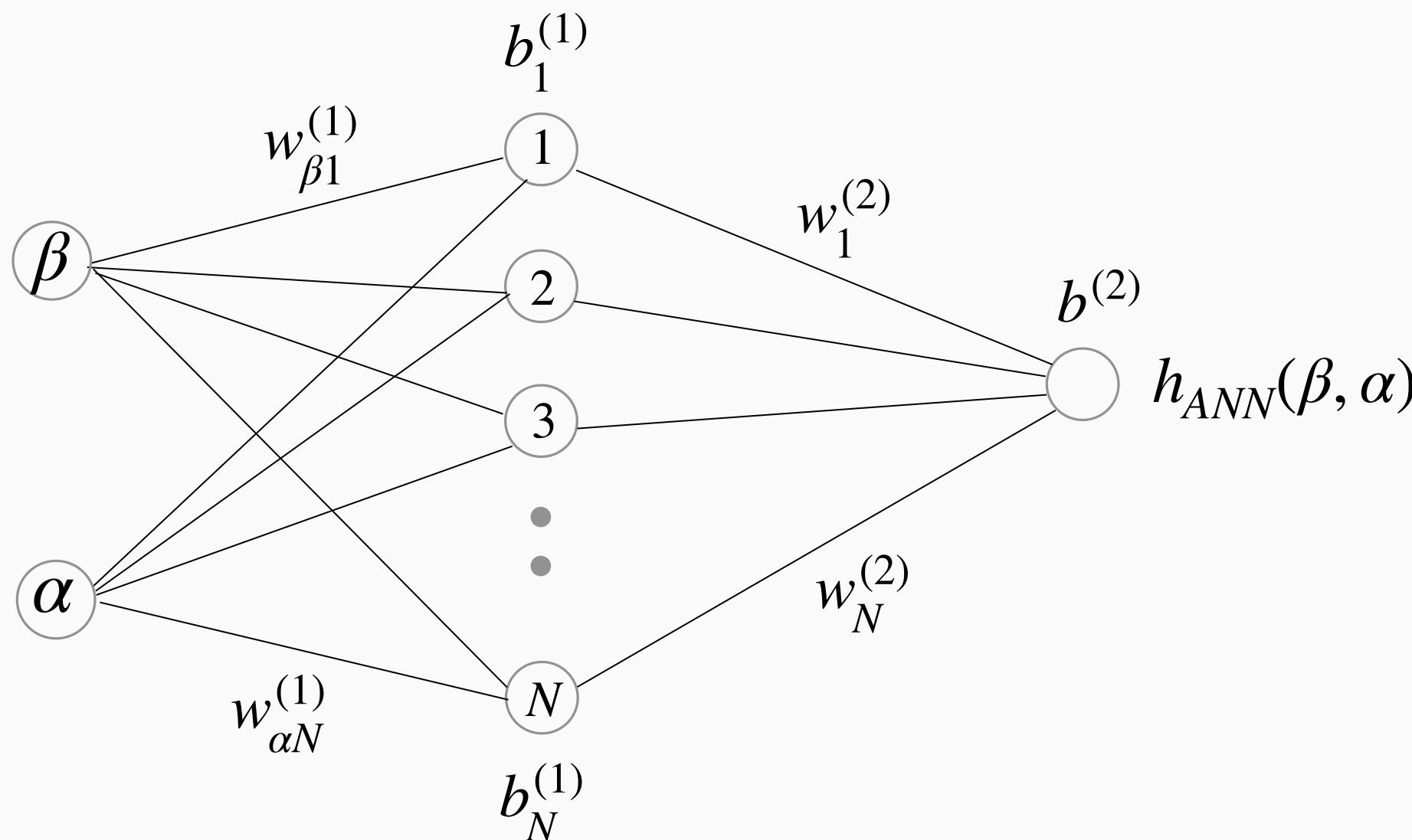
$$h = (R^T R)^{-1} R^T H$$



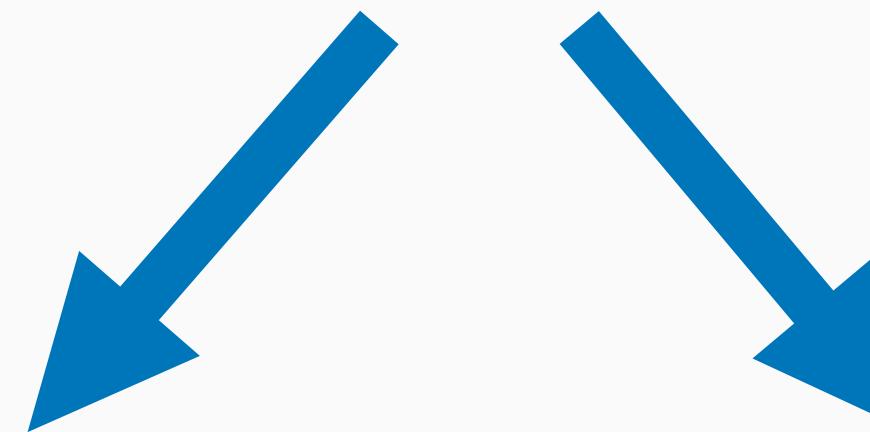
Numerical Inversion of Radon Transform

Artificial Neural Networks

H. Dutrieux et al. *Eur.Phys.J.C.* 82 (2022)



Unbiased parametrization of DD



Finite Elements Methods

N. Chouika et al. *Eur.Phys.J.C.* 77 (2017)

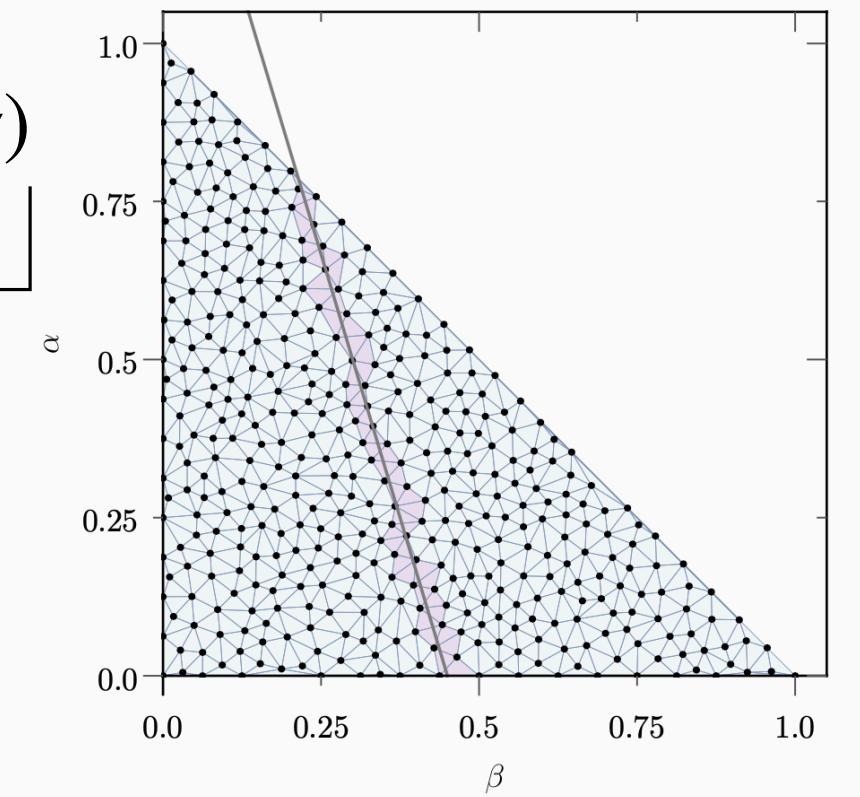
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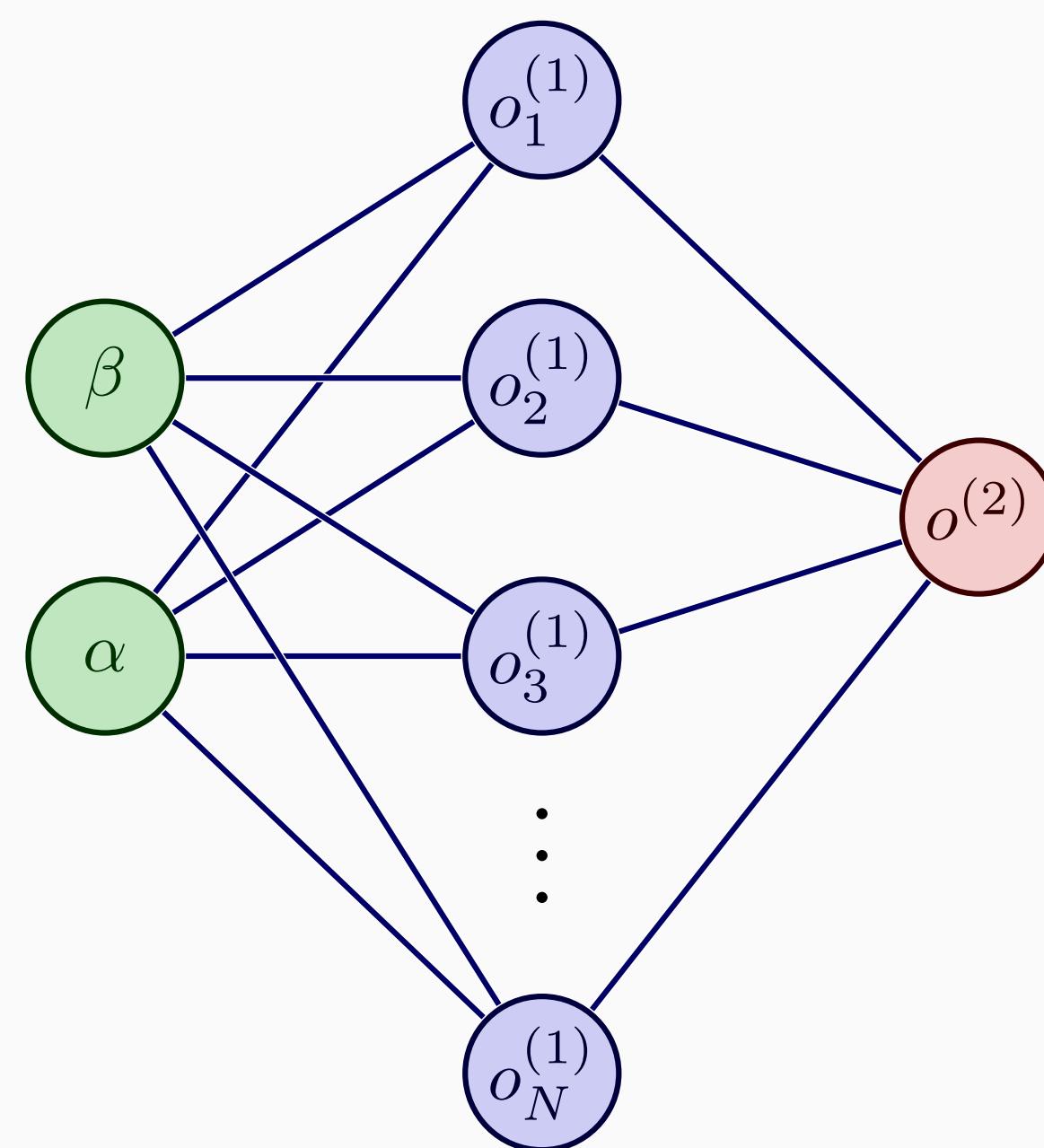
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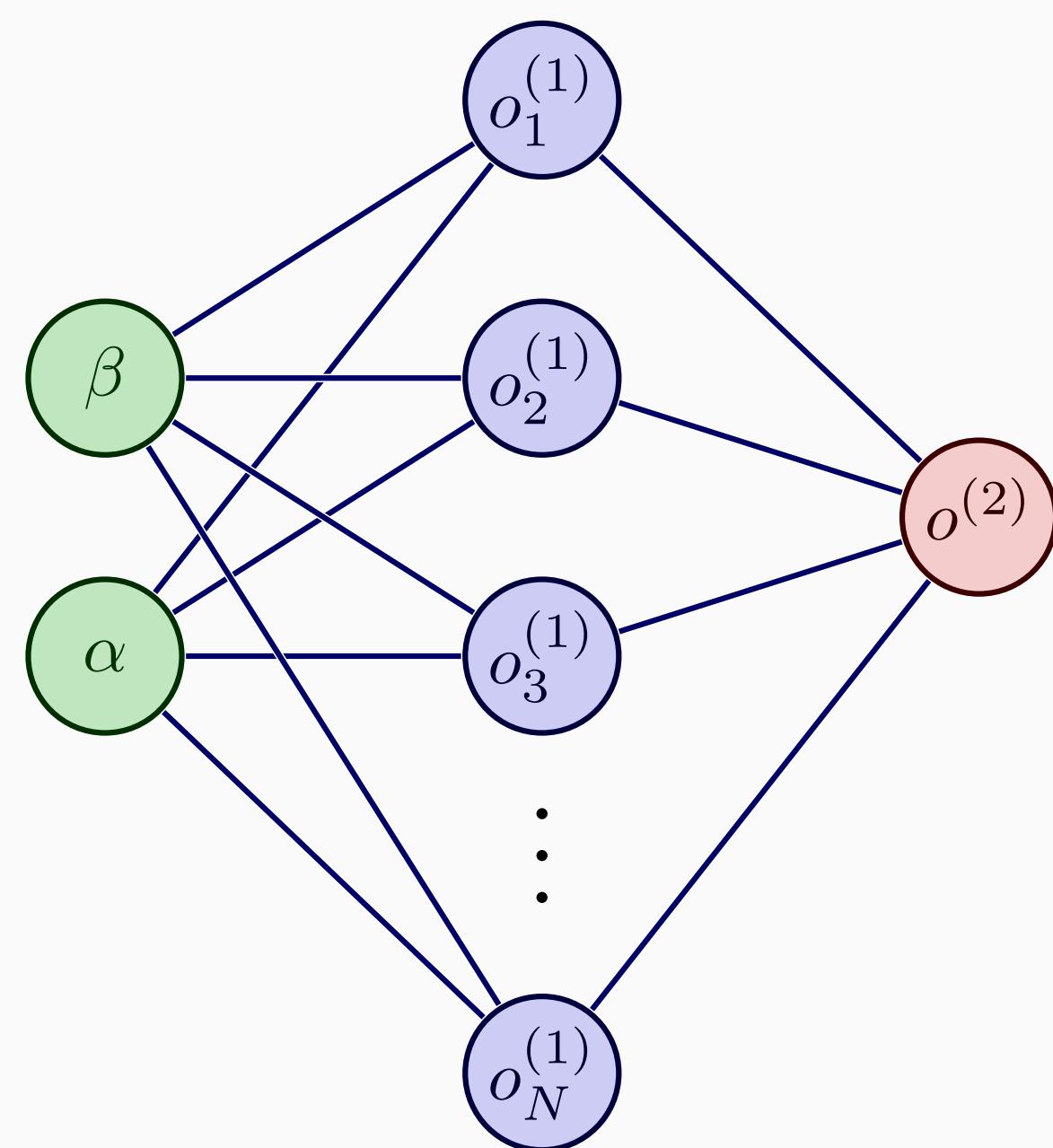
Parametrizing $h(\beta, \alpha)$ using Artificial Neural Networks (ANN)



$$\begin{aligned} &= h_{ANN}(\beta, \alpha) = \sum_{i=1}^N w_i^{(2)} o_i^{(1)} + b^{(2)} \\ &= \sum_{i=1}^N w_i^{(2)} \left[\sigma \left(w_{\beta i}^{(1)} \beta + w_{\alpha i}^{(1)} \alpha + b_i^{(1)} \right) + \sigma \left(w_{\beta i}^{(1)} \beta - w_{\alpha i}^{(1)} \alpha + b_i^{(1)} \right) \right] + b^{(2)} \end{aligned}$$

- Adam optimizer (gradient descent)
- Dropout regularization

Parametrizing $h(\beta, \alpha)$ using Artificial Neural Networks (ANN)



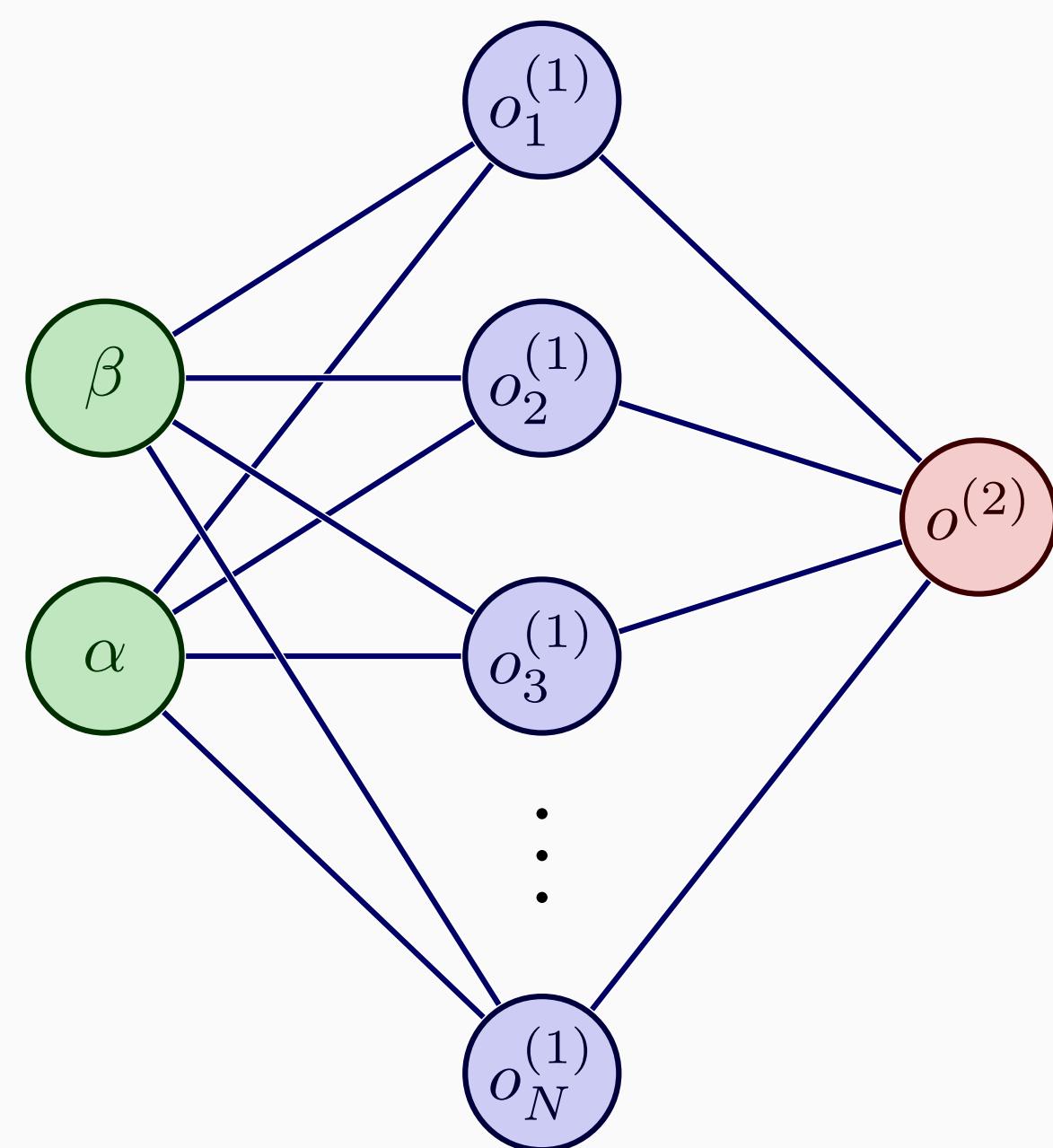
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↑

$$\sigma(x) = \frac{1}{1 + e^{-x}}$$

- Adam optimizer (gradient descent)
- Dropout regularization

Algorithm

- Initialize the ANN parameters (randomly).
- Given a sampling set of GPD values $H_i(x_i, \xi_i)$ in the DGLAP region, numerically evaluate the RT along each line (x_i, ξ_i) using $h_{ANN}(\beta, \alpha)$ as DD. $\mathcal{R}h_{ANN}(x_i, \xi_i) = \hat{H}(x_i, \xi_i)$.
- Update the ANN parameters using Adam optimization algorithm in order to minimize
$$MSE = \frac{1}{N_{sample}} \sum_{i=1}^{N_{sample}} \left(H_i - \hat{H}_i \right)^2$$
- Iterate until convergence.

Testing with analytical models

Nakanishi based model for pion

N.Chouika et al. *Phys.Lett.B*:780(2018)

$$H(x, \xi, t = 0) = \begin{cases} 30 \frac{(1-x)^2(x^2 - \xi^2)}{(1-\xi^2)^2}, & |x| > \xi \\ 15 \frac{(1-x)(\xi^2 - x^2)(x + 2x\xi + \xi^2)}{2\xi^3(1+\xi)^2}, & |x| < \xi \end{cases}$$

$$H(x, \xi, t = 0) = (1-x) \int_{\Omega^+} d\beta d\alpha \delta(x - \beta - \alpha\xi) h(\beta, \alpha)$$
$$h(\beta, \alpha) = \frac{15}{2} (1 - 3(\alpha^2 - \beta^2) - 2\beta)$$

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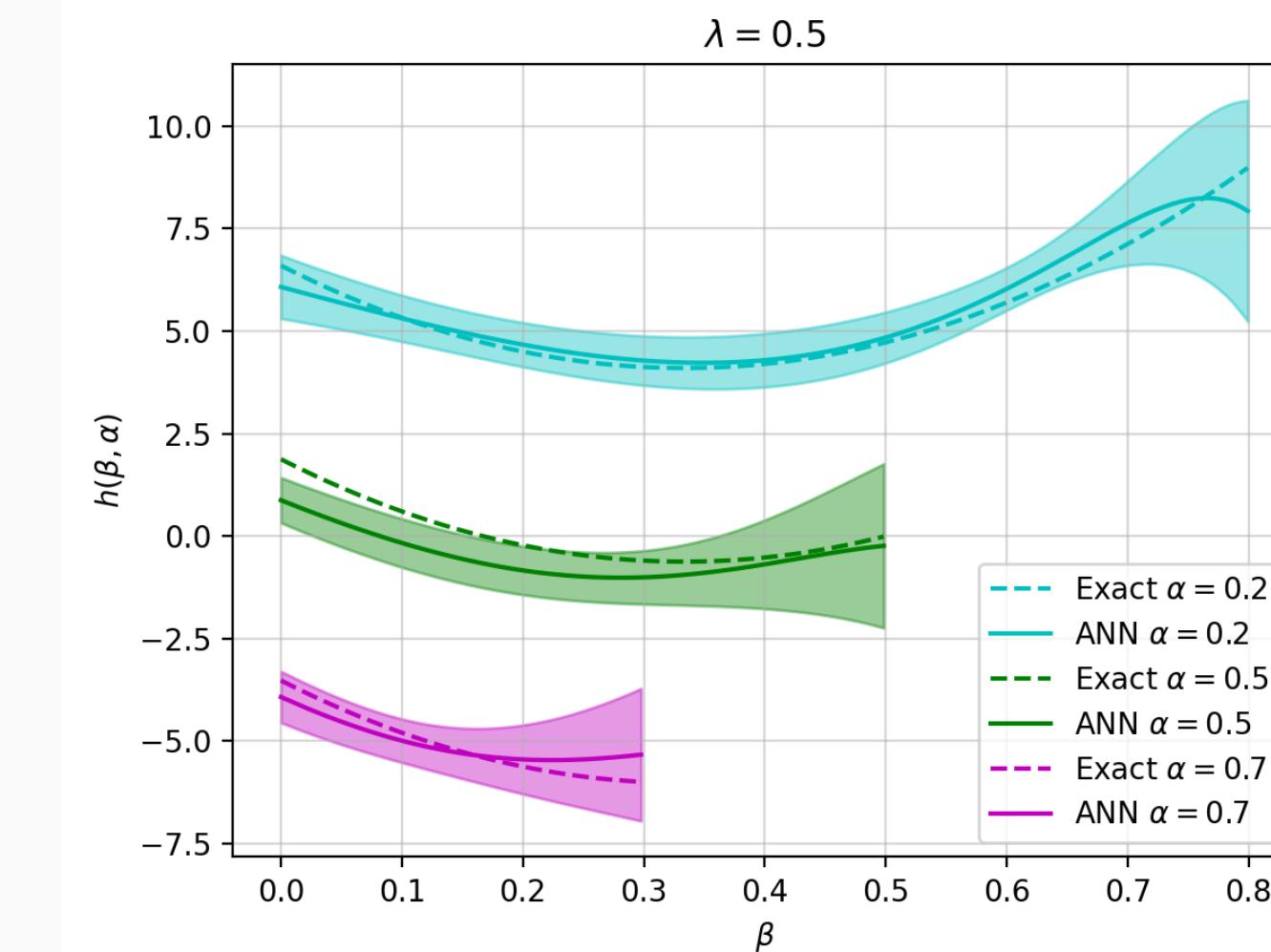
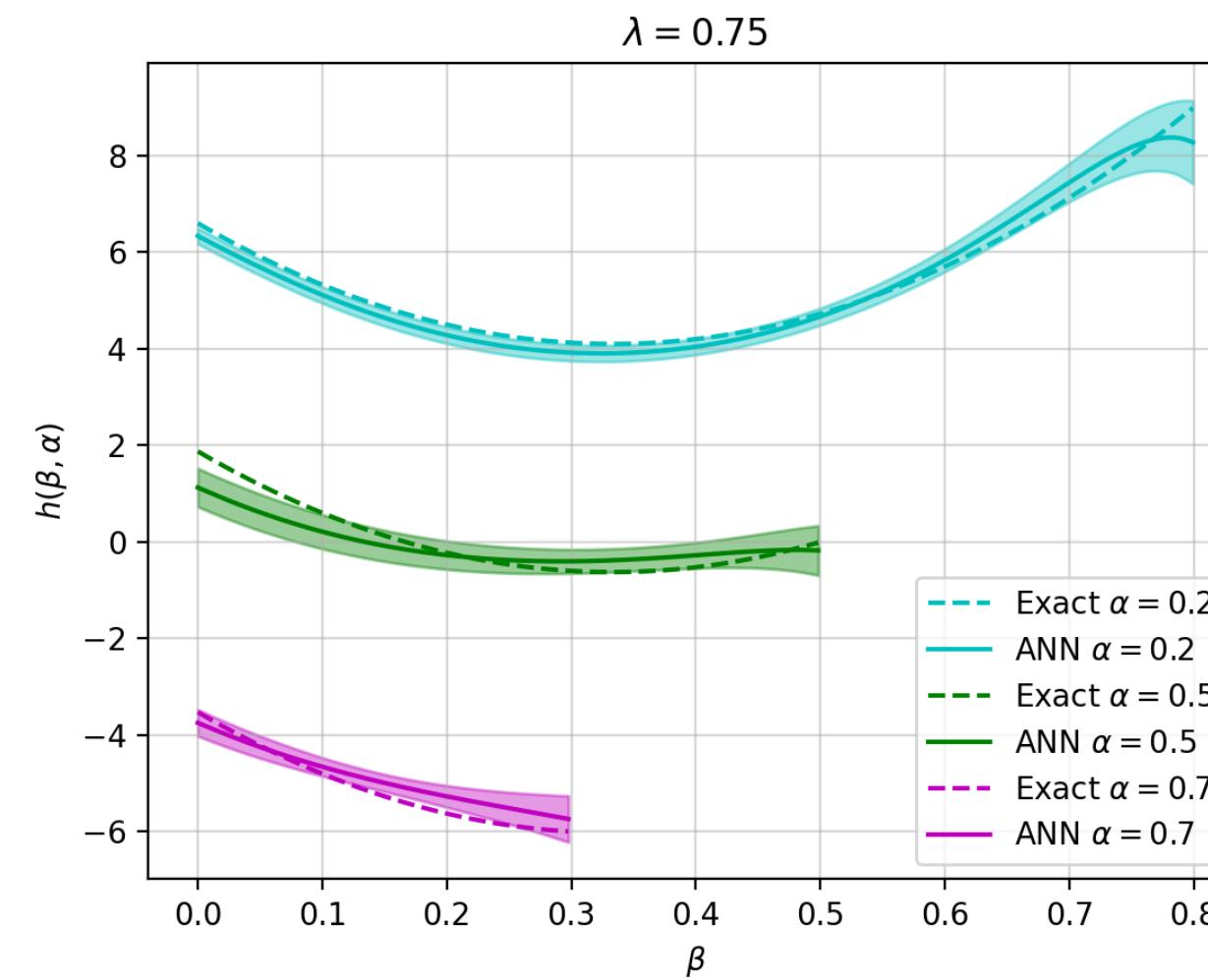
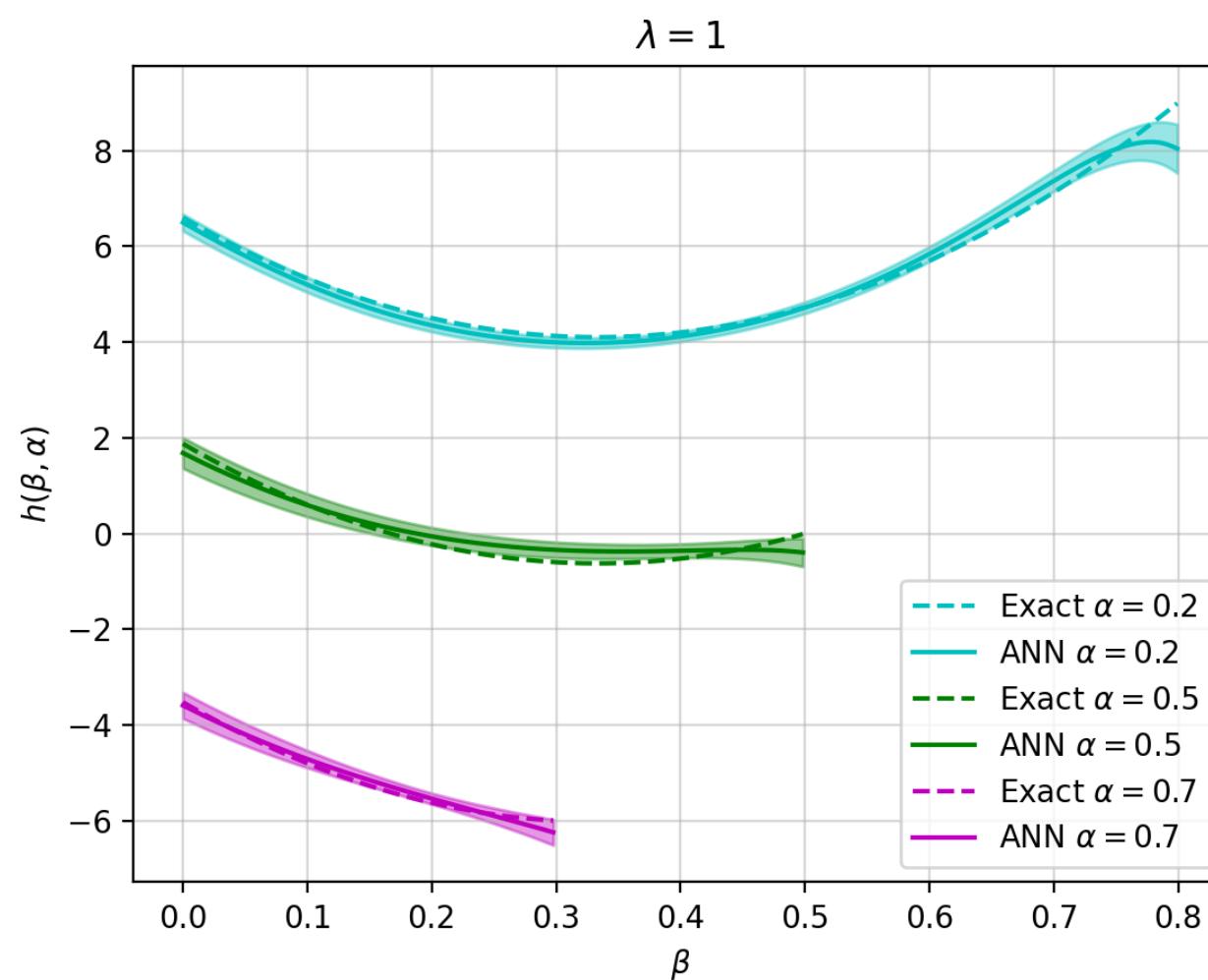
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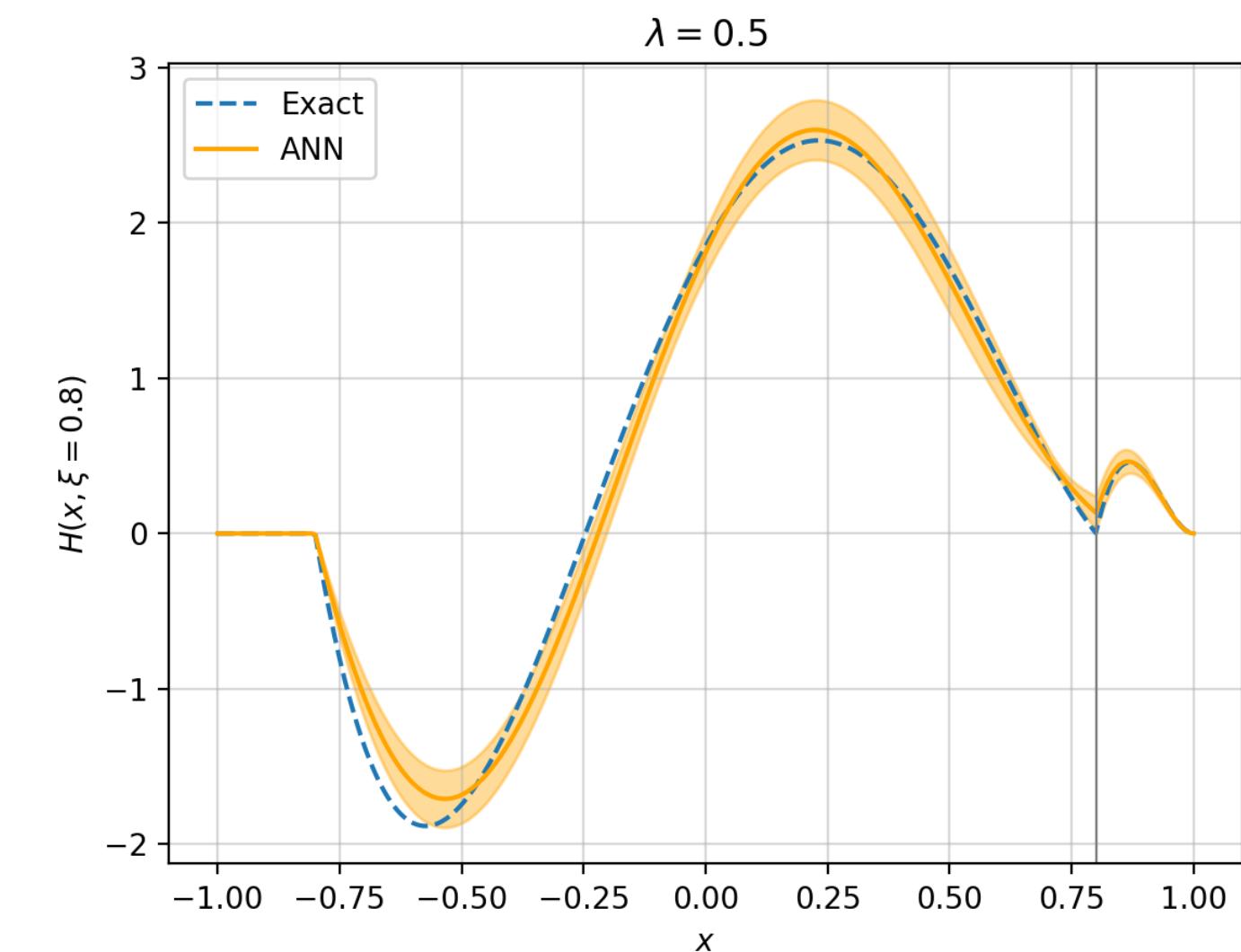
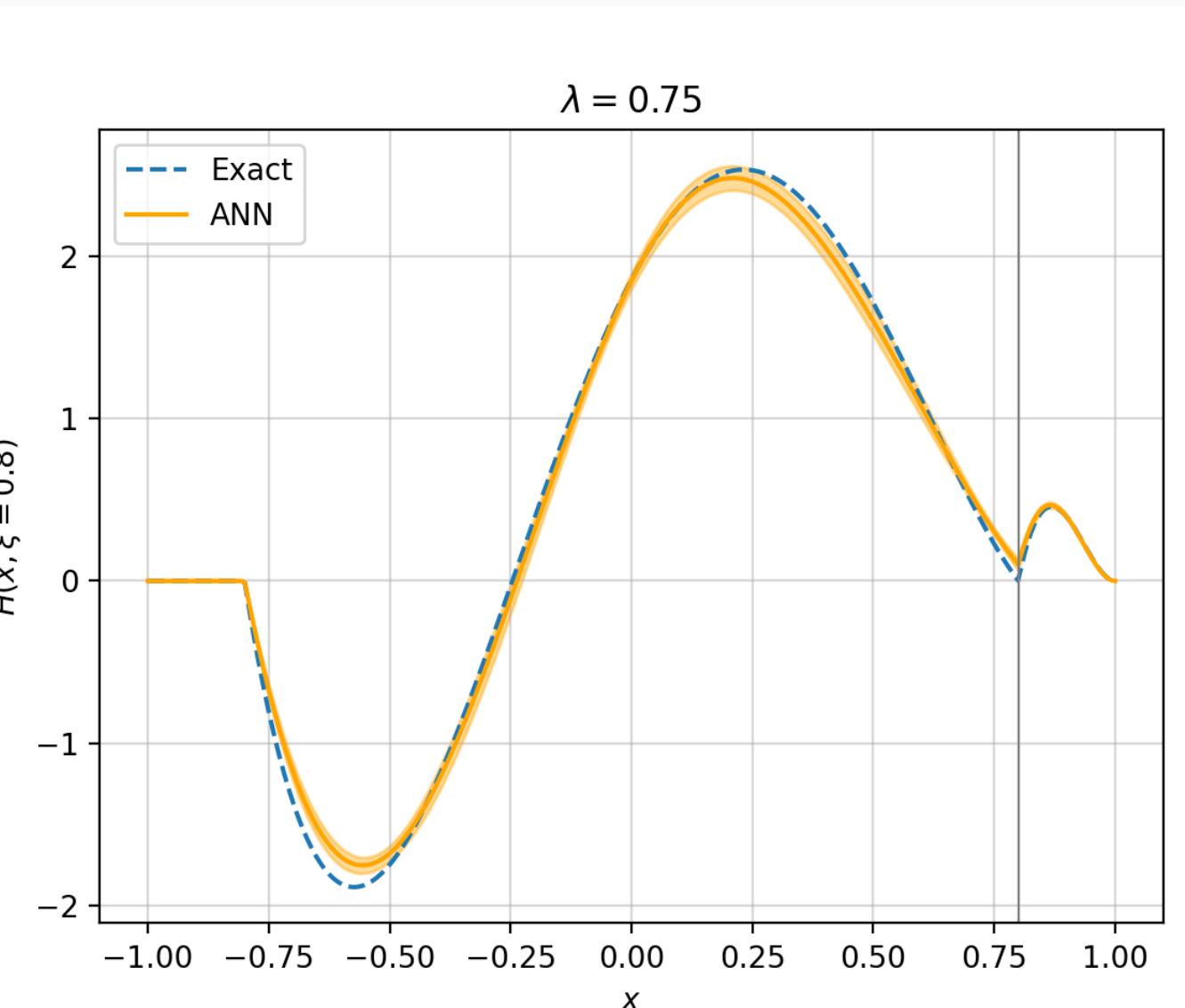
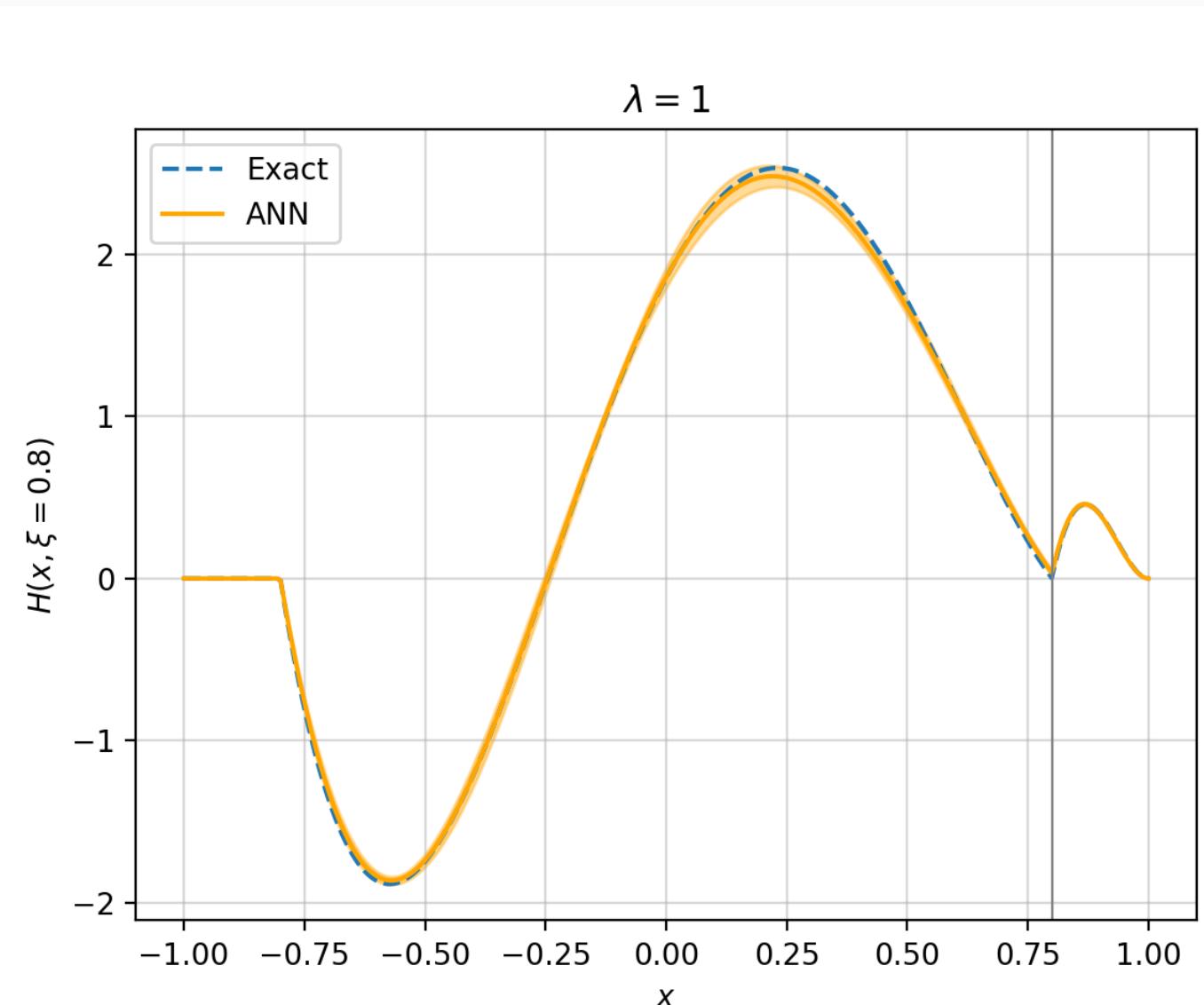
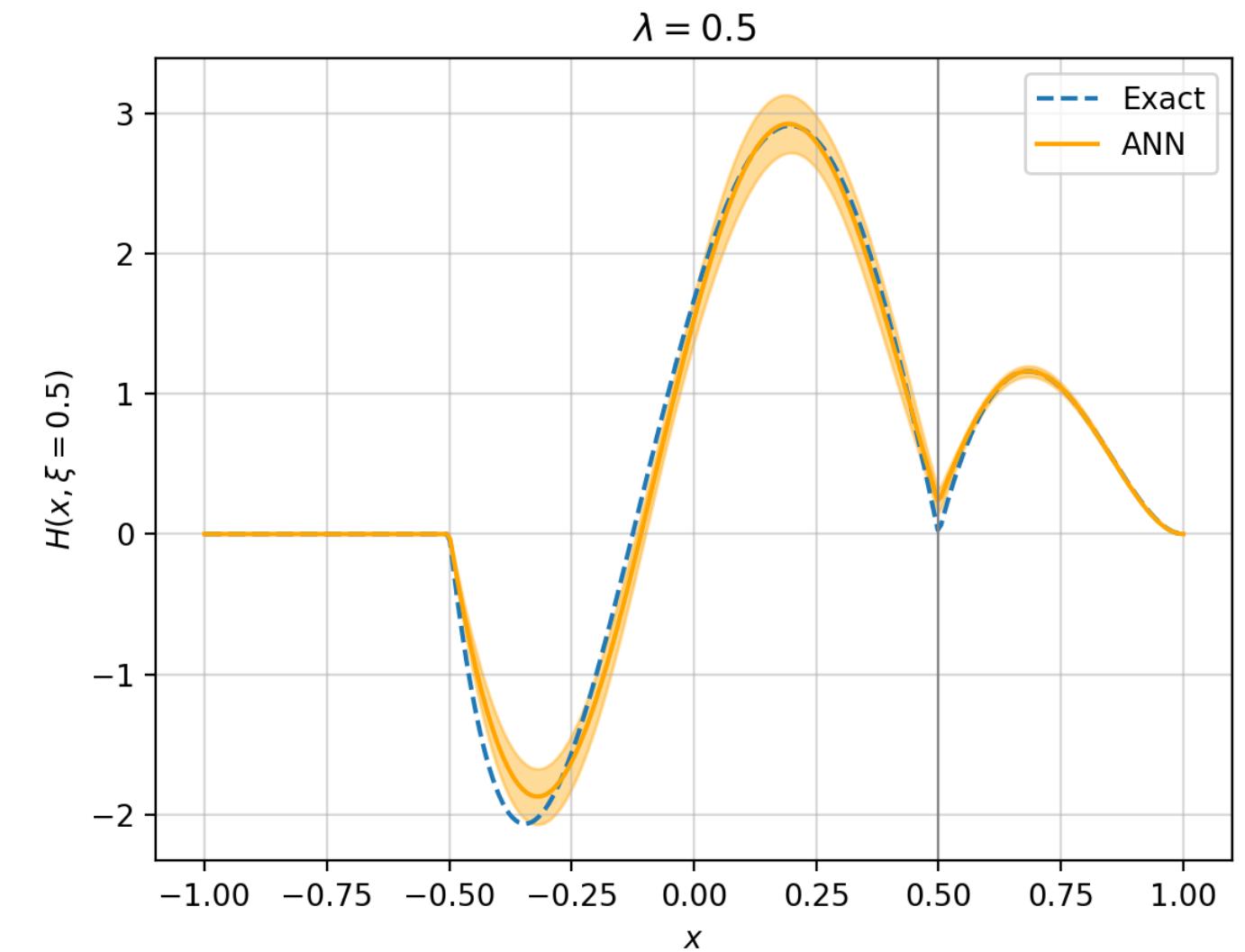
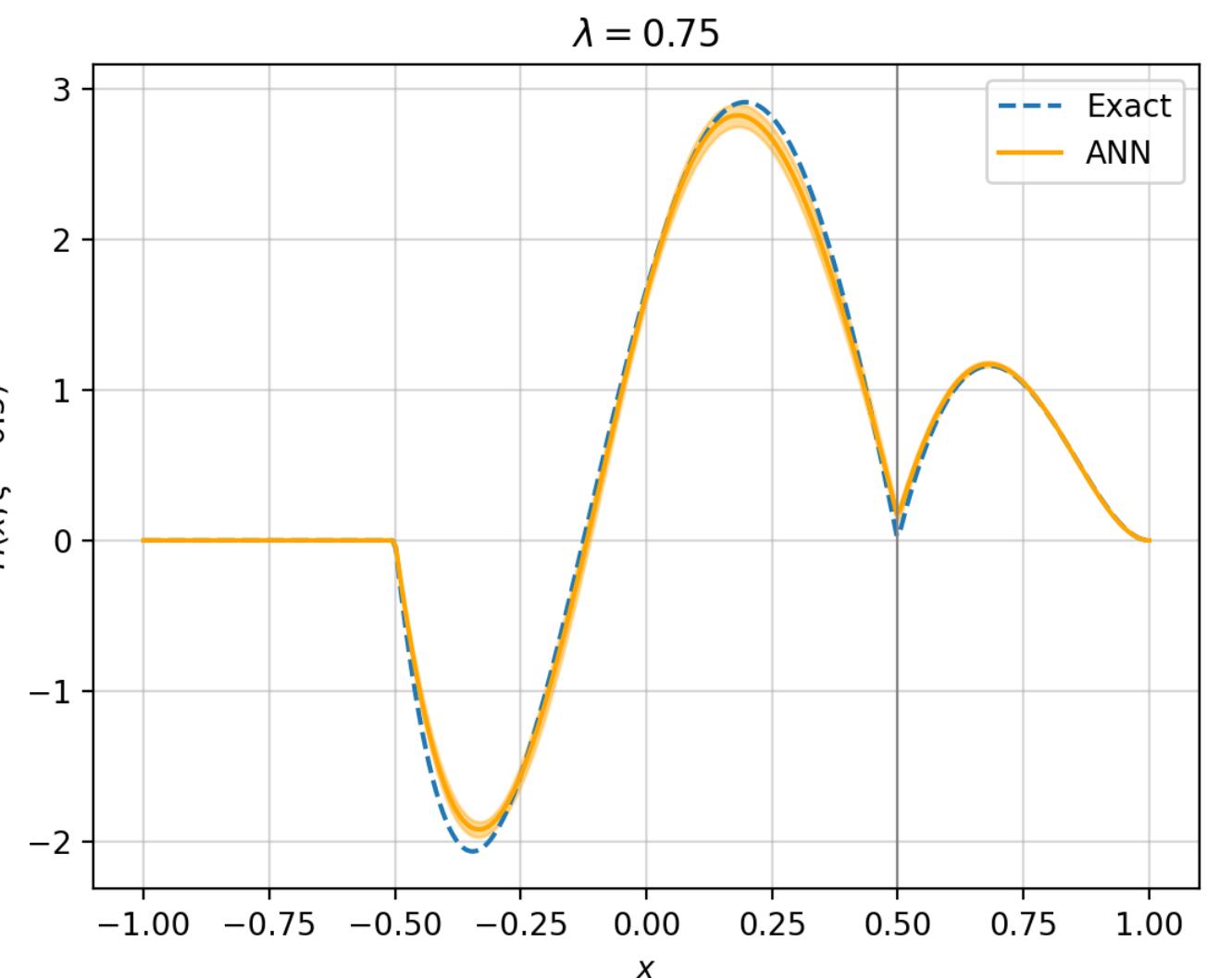
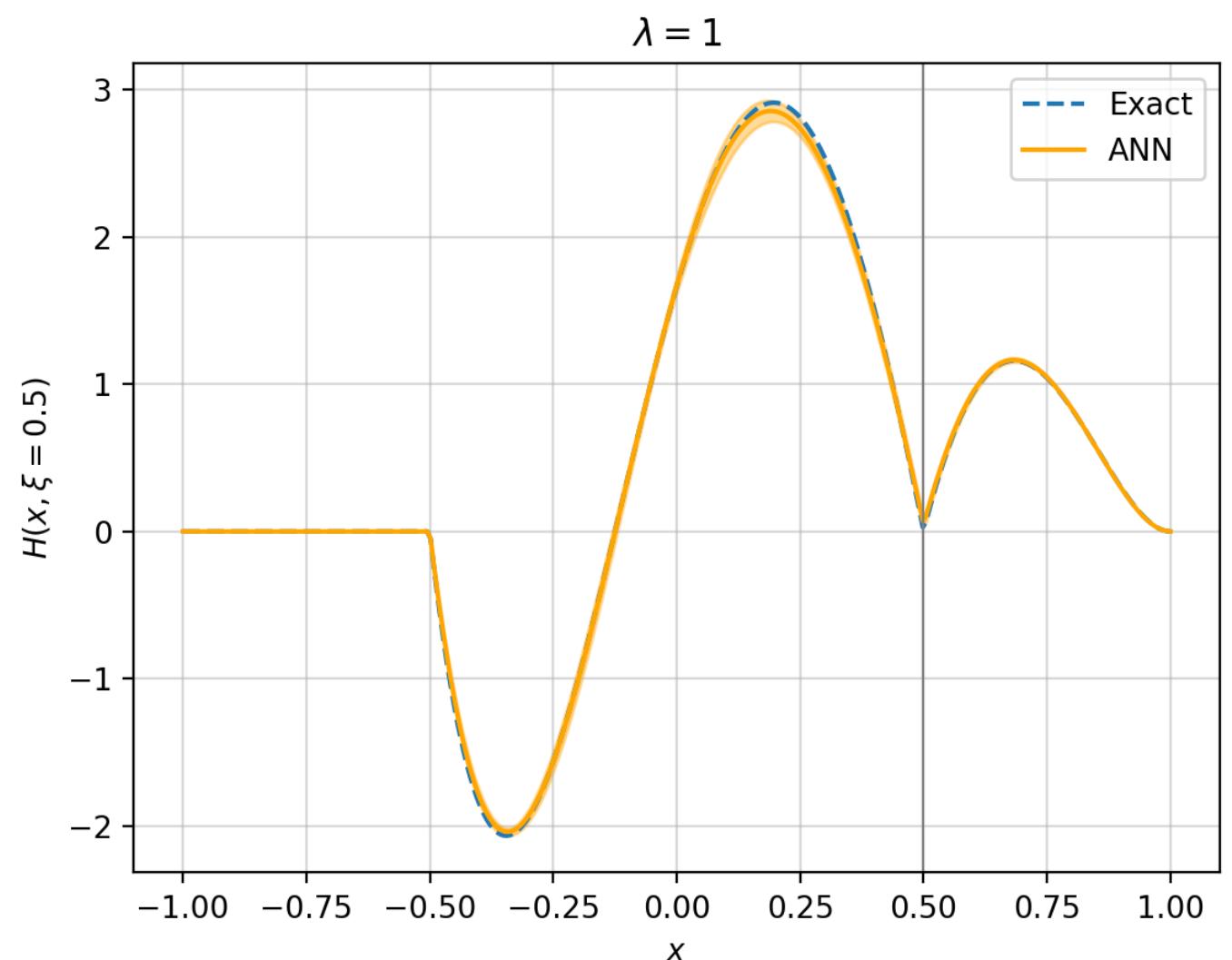
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$$N_{neurons} = 10^2, \quad N_{sample} = 10^4$$

$$\xi \in [0, \lambda x]$$



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Goloskokov-Kroll model

S.V. Goloskokov, P. Kroll, Eur.Phys.J.C. 50 (2007)

$$H(x, \xi) = \int_{\Omega} d\beta d\alpha \delta(x - \beta - \alpha\xi) q(\beta) h_{GK}(\beta, \alpha),$$

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Valence distribution $q_{val}^u(\beta) = \beta^{-\delta}(1-\beta)^{2n+1} \sum_{j=0}^2 c_j \beta^{j/2}, \quad n = 1, \quad \delta = 0.48$

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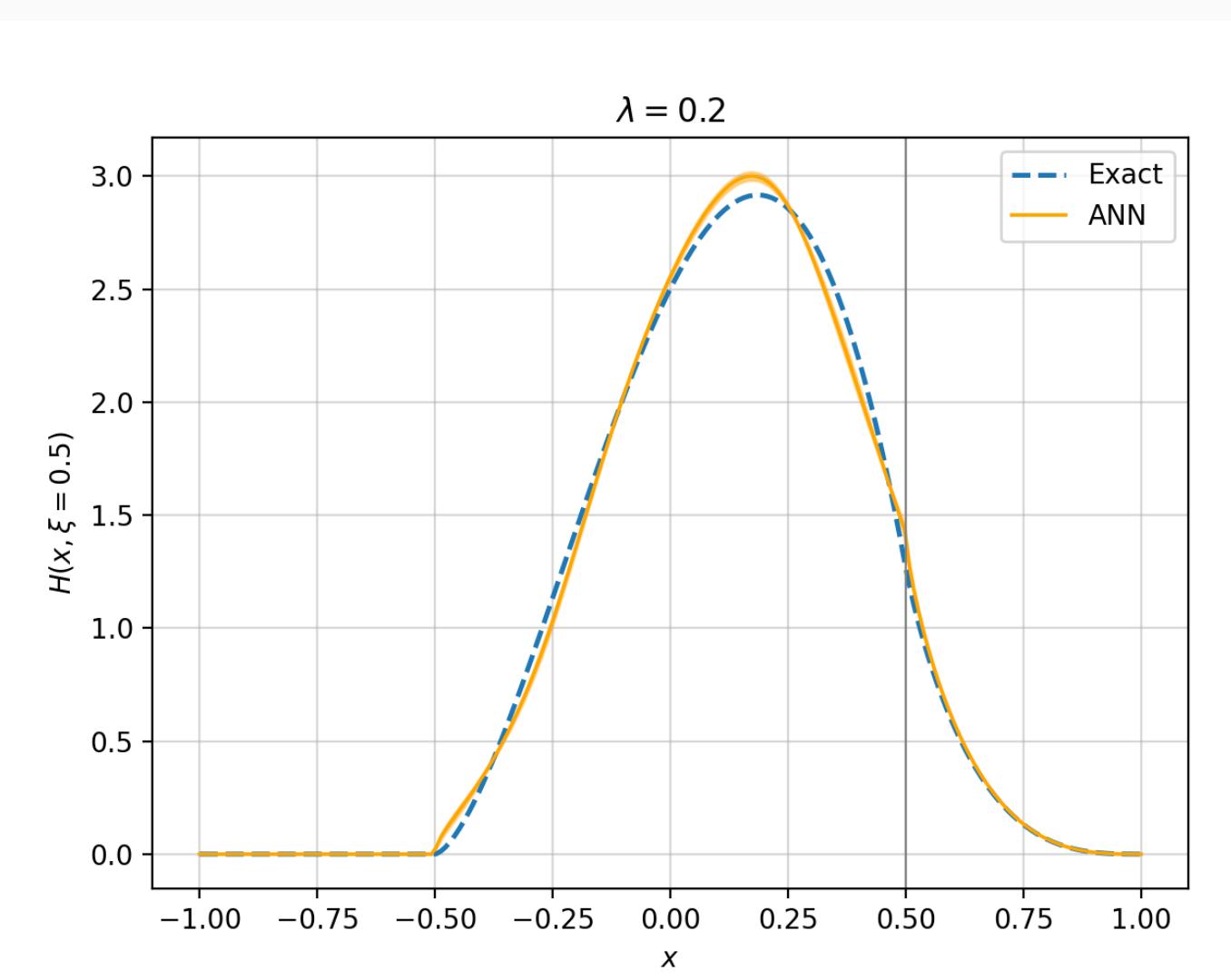
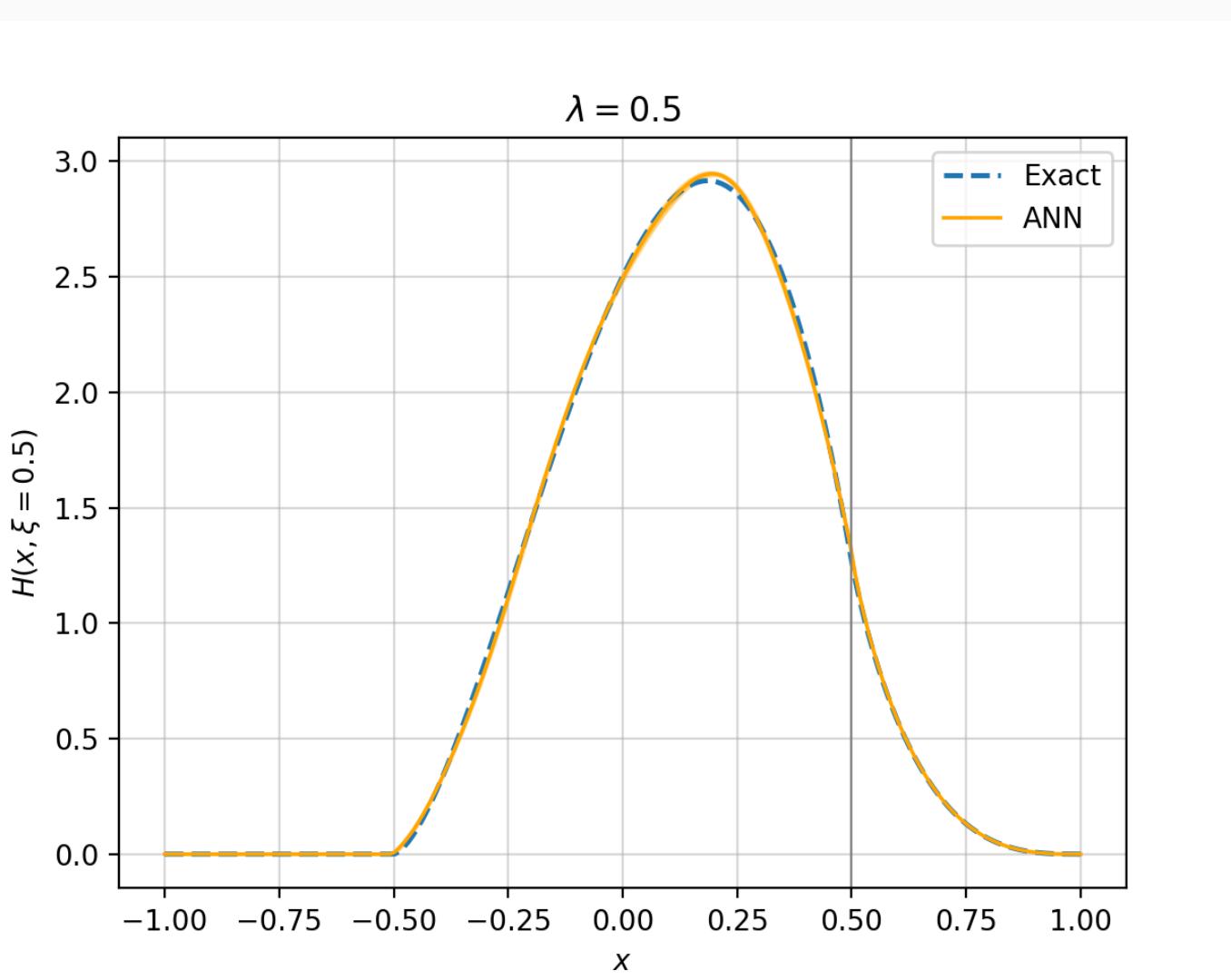
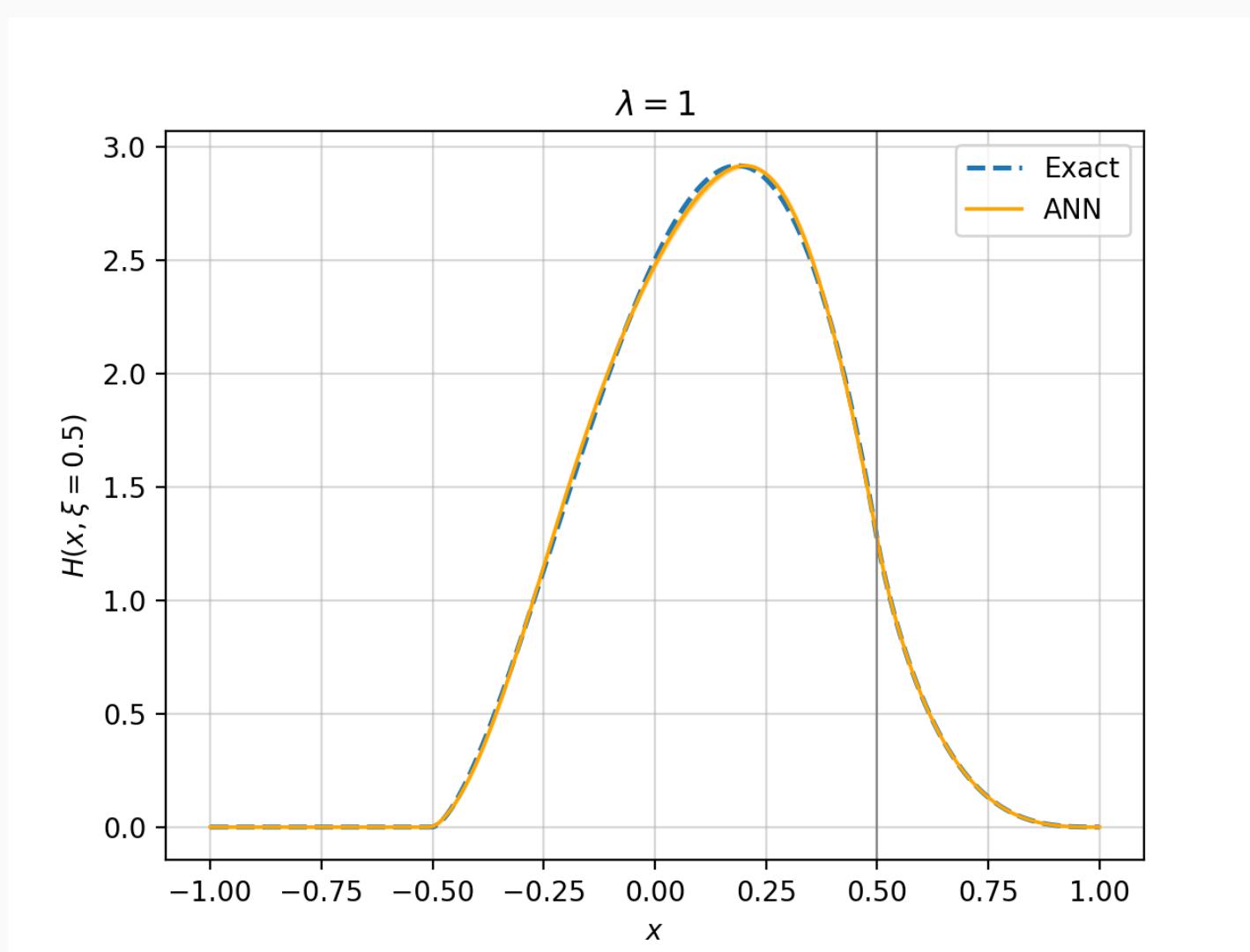
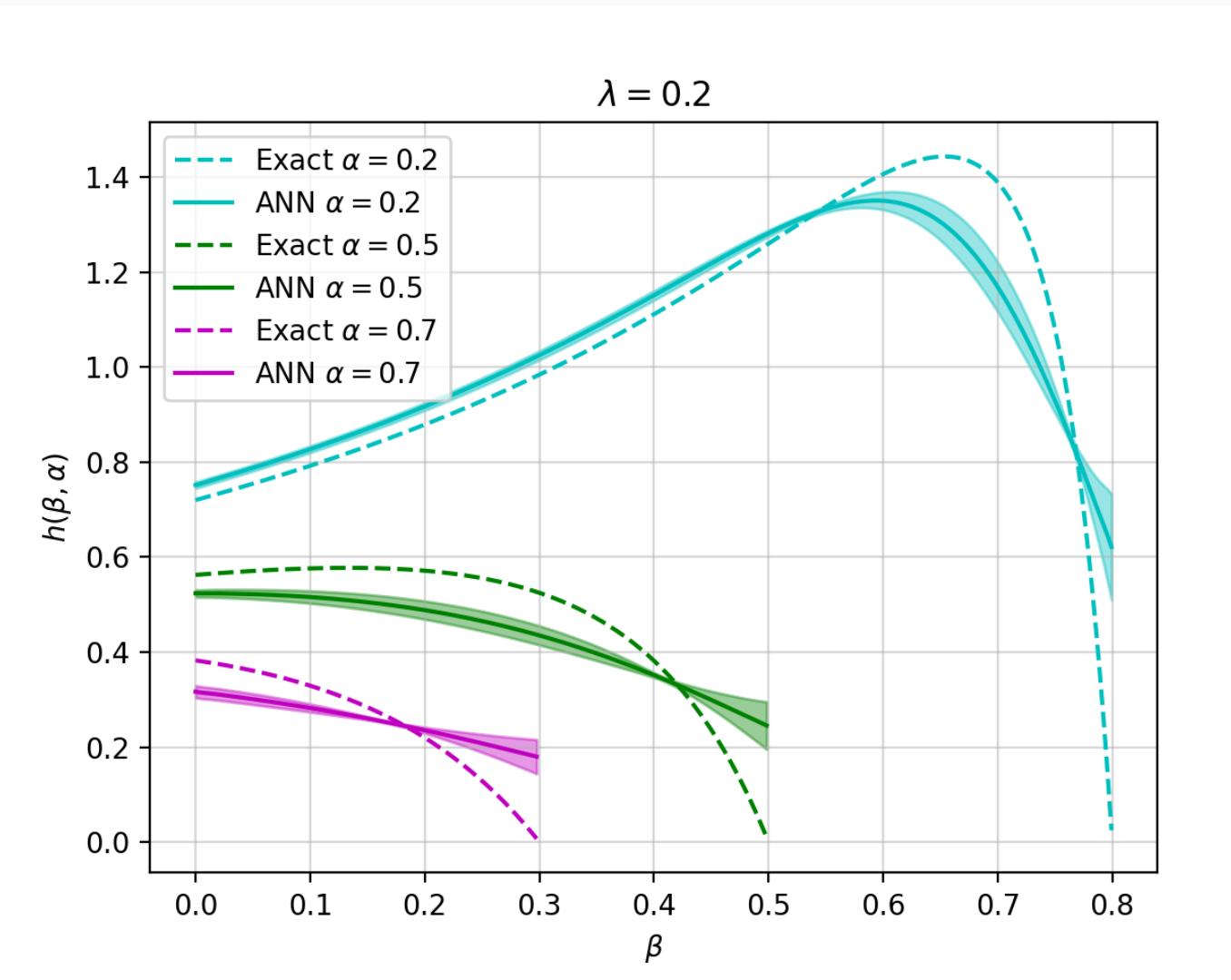
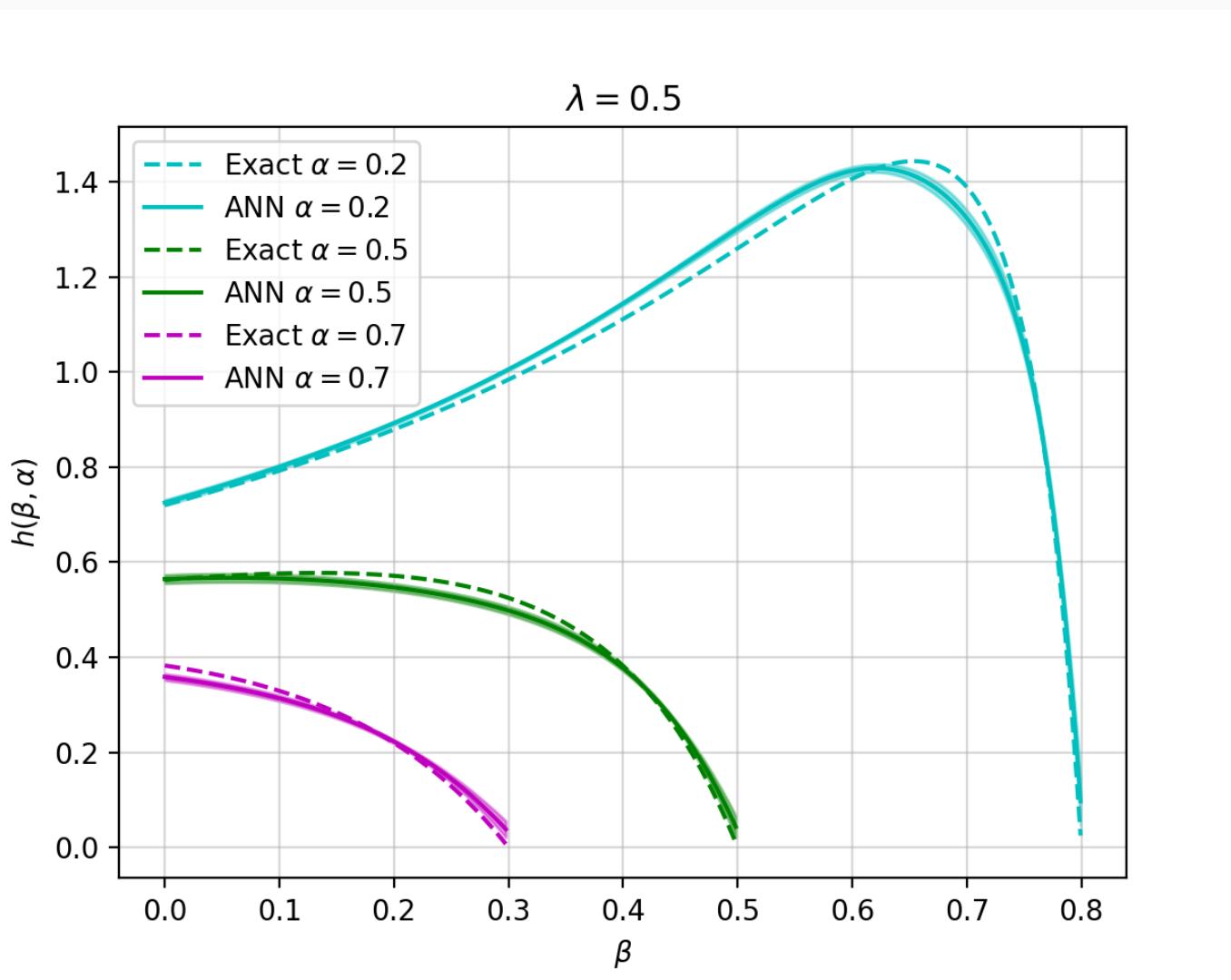
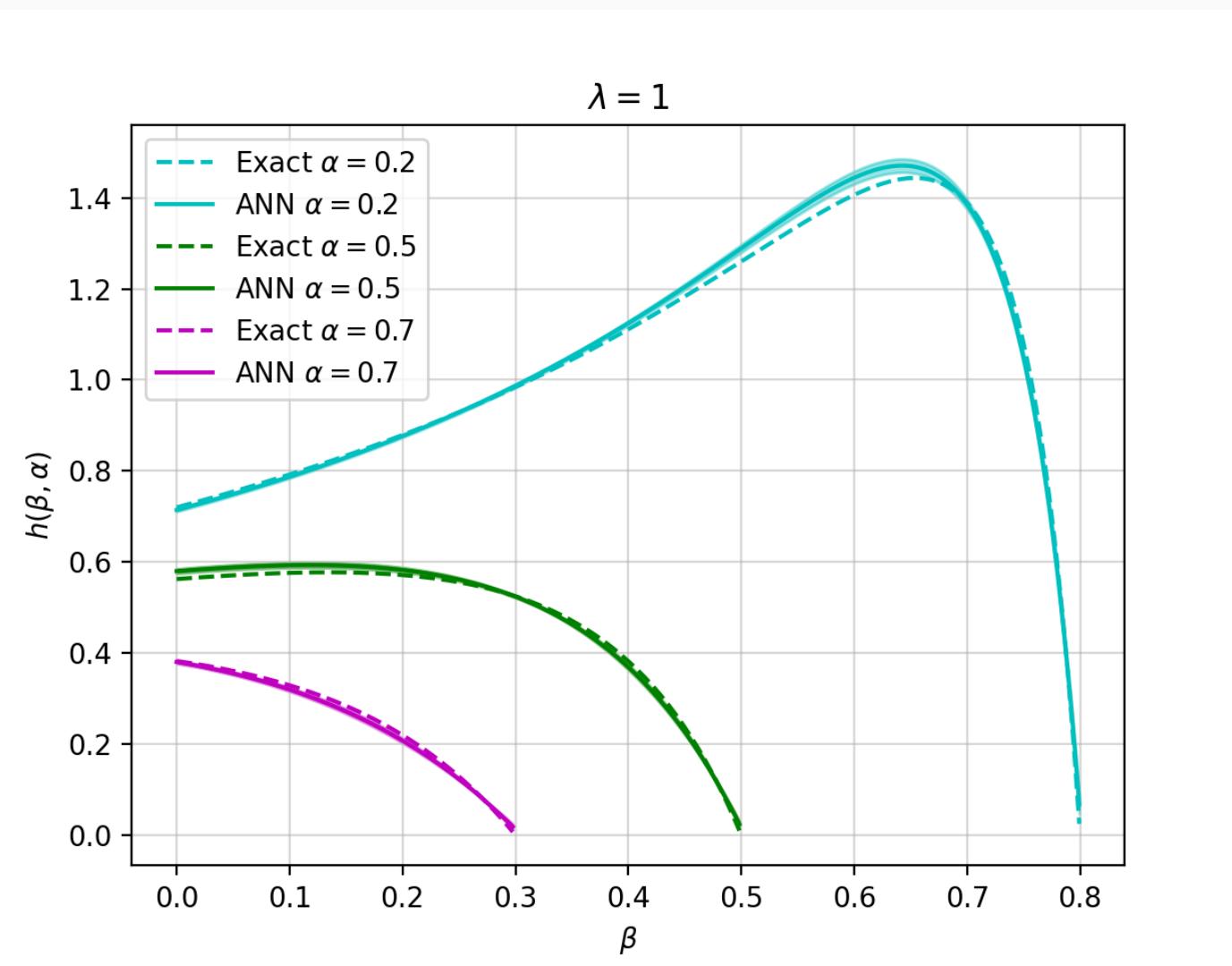
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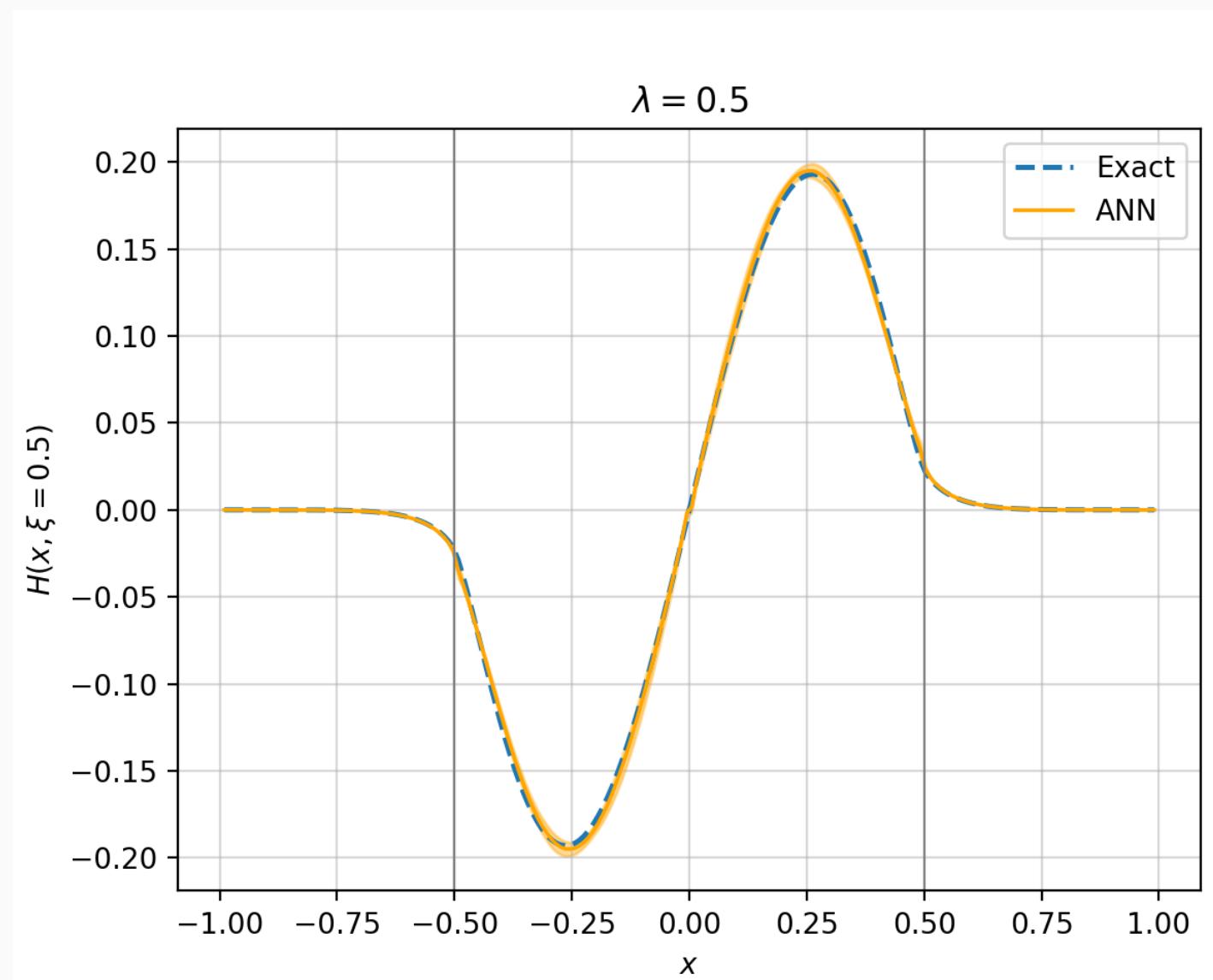
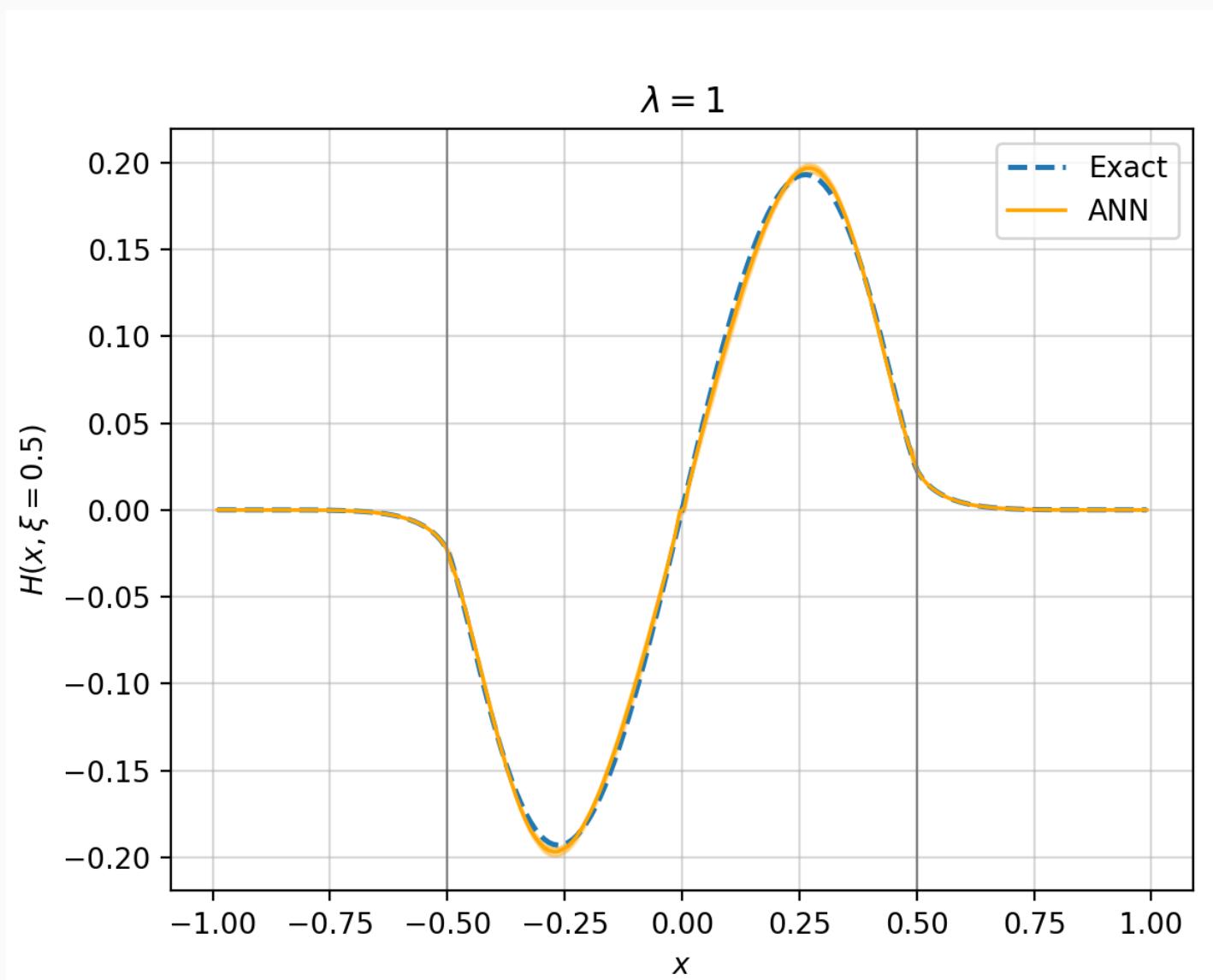
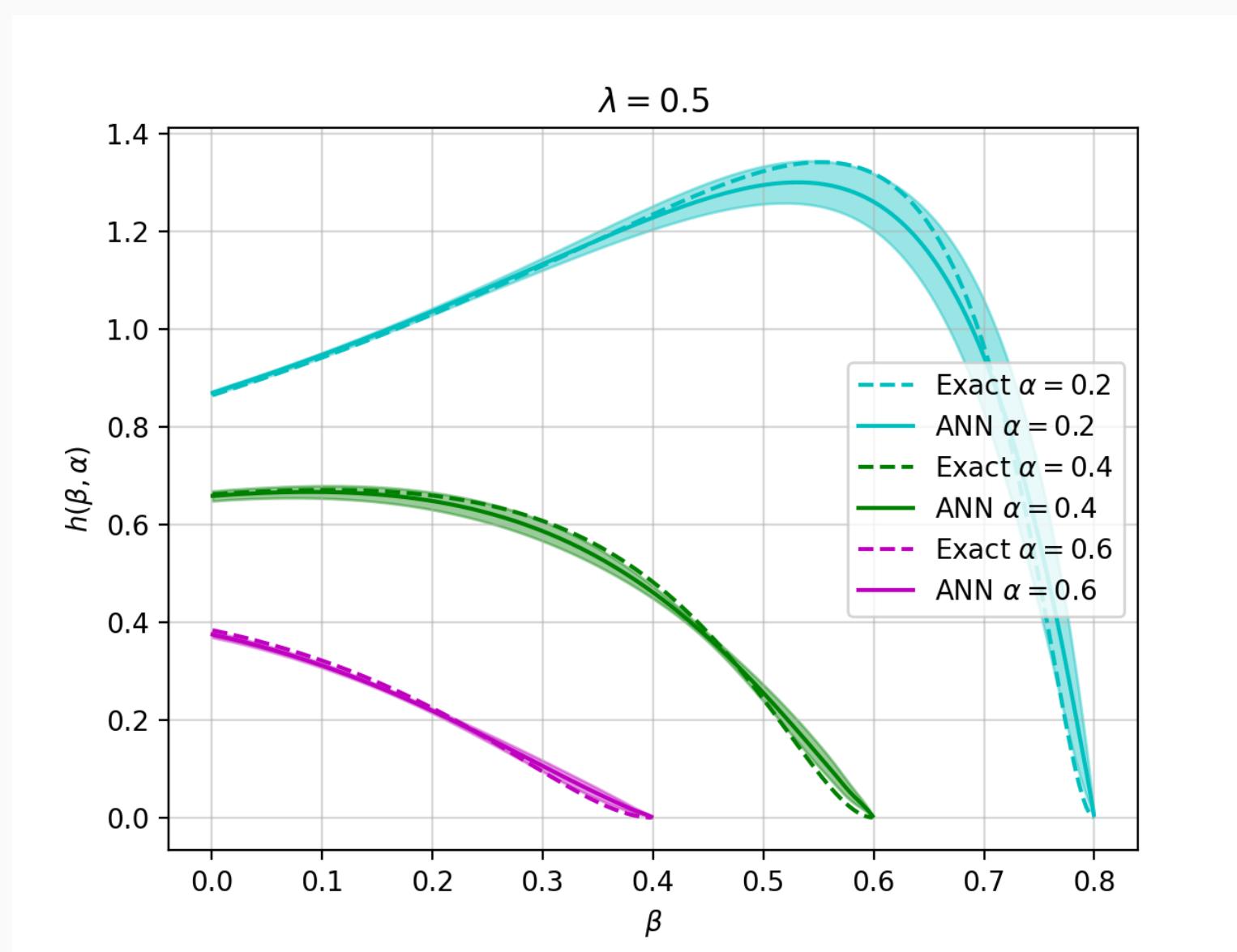
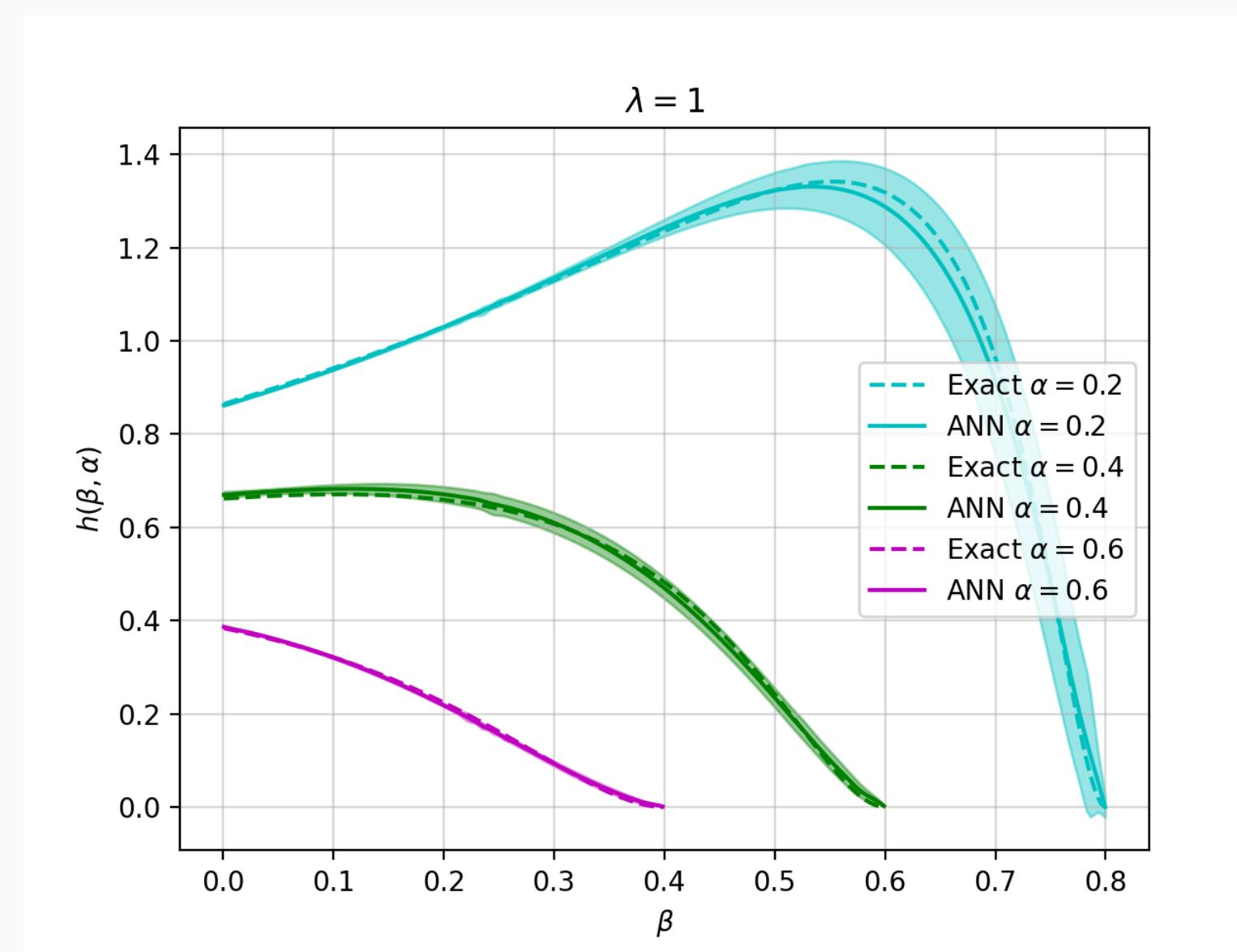
$N_{neurons} = 25, \quad N_{sample} = 5 \times 10^3$

$$h_{GK}(\beta, \alpha) \simeq \frac{h_{ANN}(\beta, \alpha)}{\int_{-1+|\beta|}^{1-|\beta|} d\alpha h_{ANN}(\beta, \alpha)}$$

Valence DD/GPD



Sea DD/GPD



Conclusions and Outlook

- GPD can be extended from DGLAP (and proper subsets) to ERBL inverting its RT
- ANNs are a good tool for inverting RT
- Testing the method on experimental results (EIC)
- Applying ANNs to PDF reconstruction by a sequence of Mellin moments
(with Khépani Raya and PP)