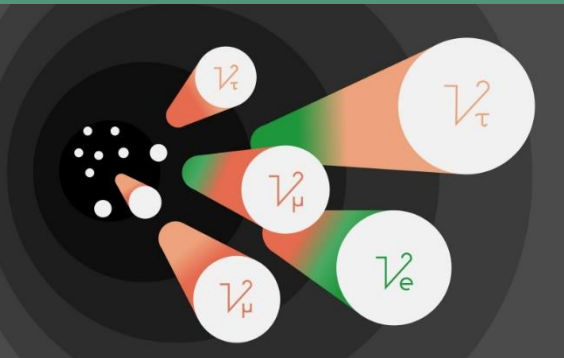


Cosmological limits to Non-Standard Interactions of the neutrino



Abdel Pérez-Lorezana
Cinvestav-IPN

ICN, JANUARY 2023



In coll. with J.A. Venzor, J. De Santiago and G. Garcia-Arroyo.



Neutrinos have mass and mix



BSM



New Fermions (N)

or

New Scalars (flavor symmetries)

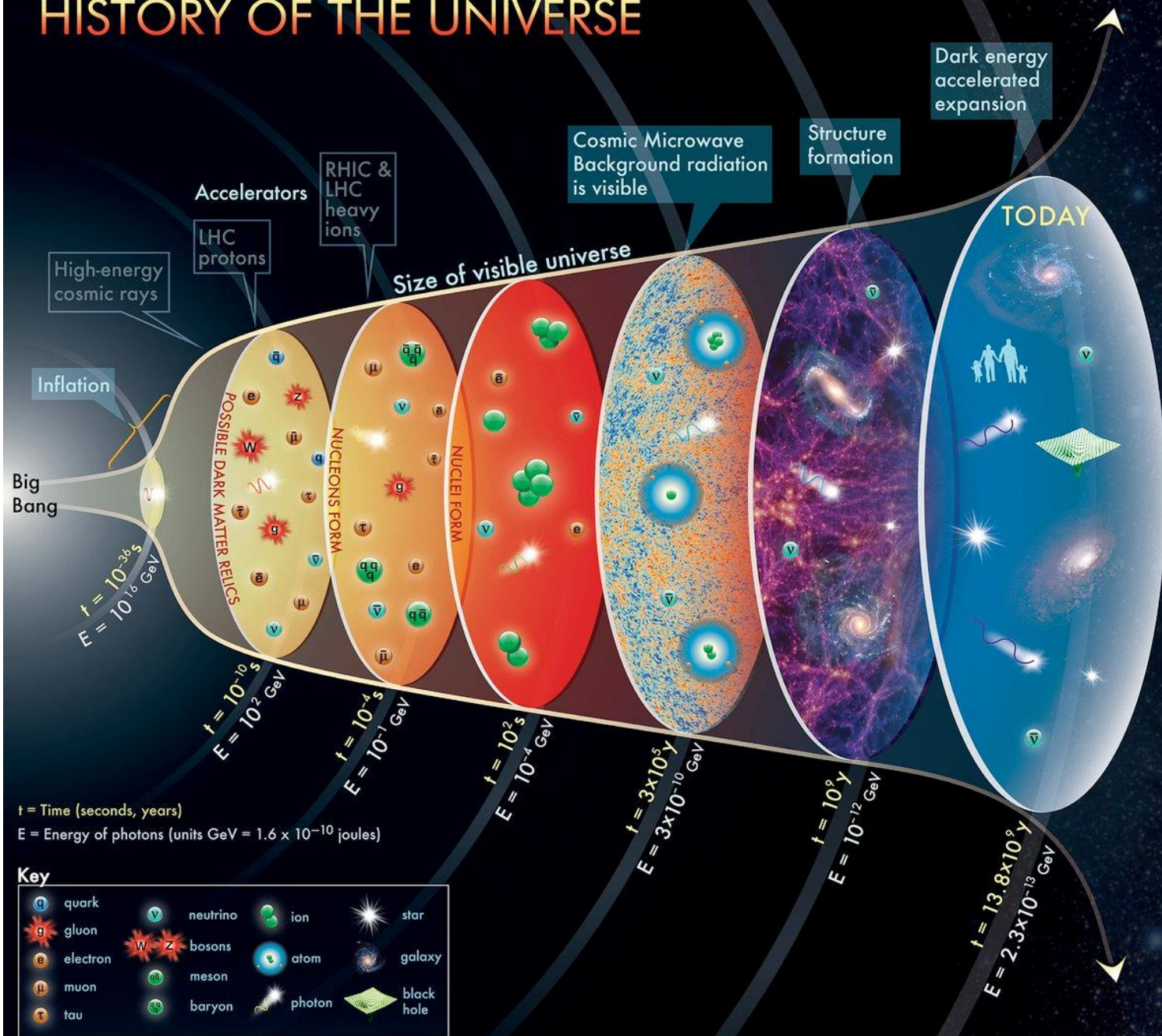


New particles come with new phenomenology



Early Universe is a good playground to look for it

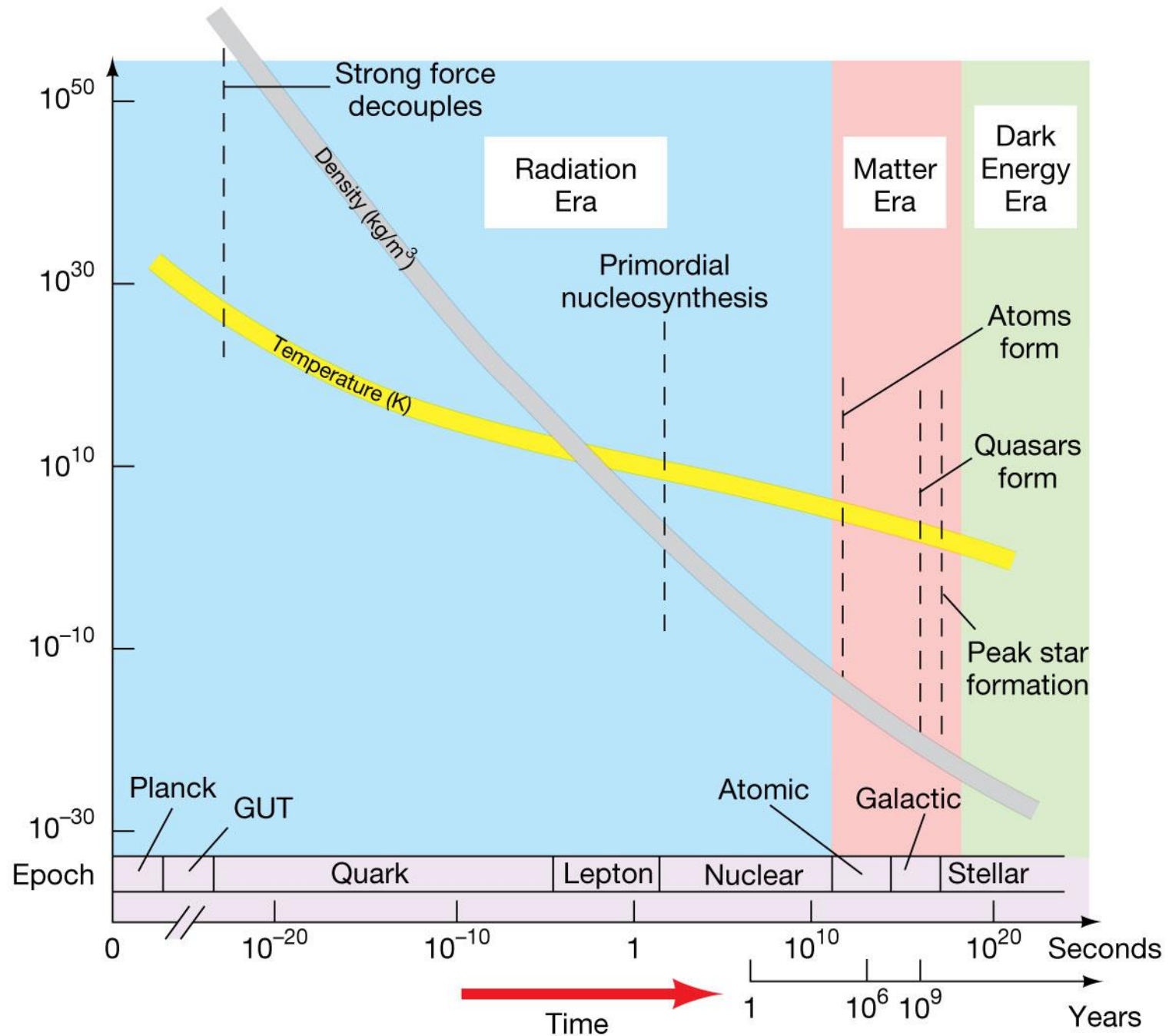
HISTORY OF THE UNIVERSE

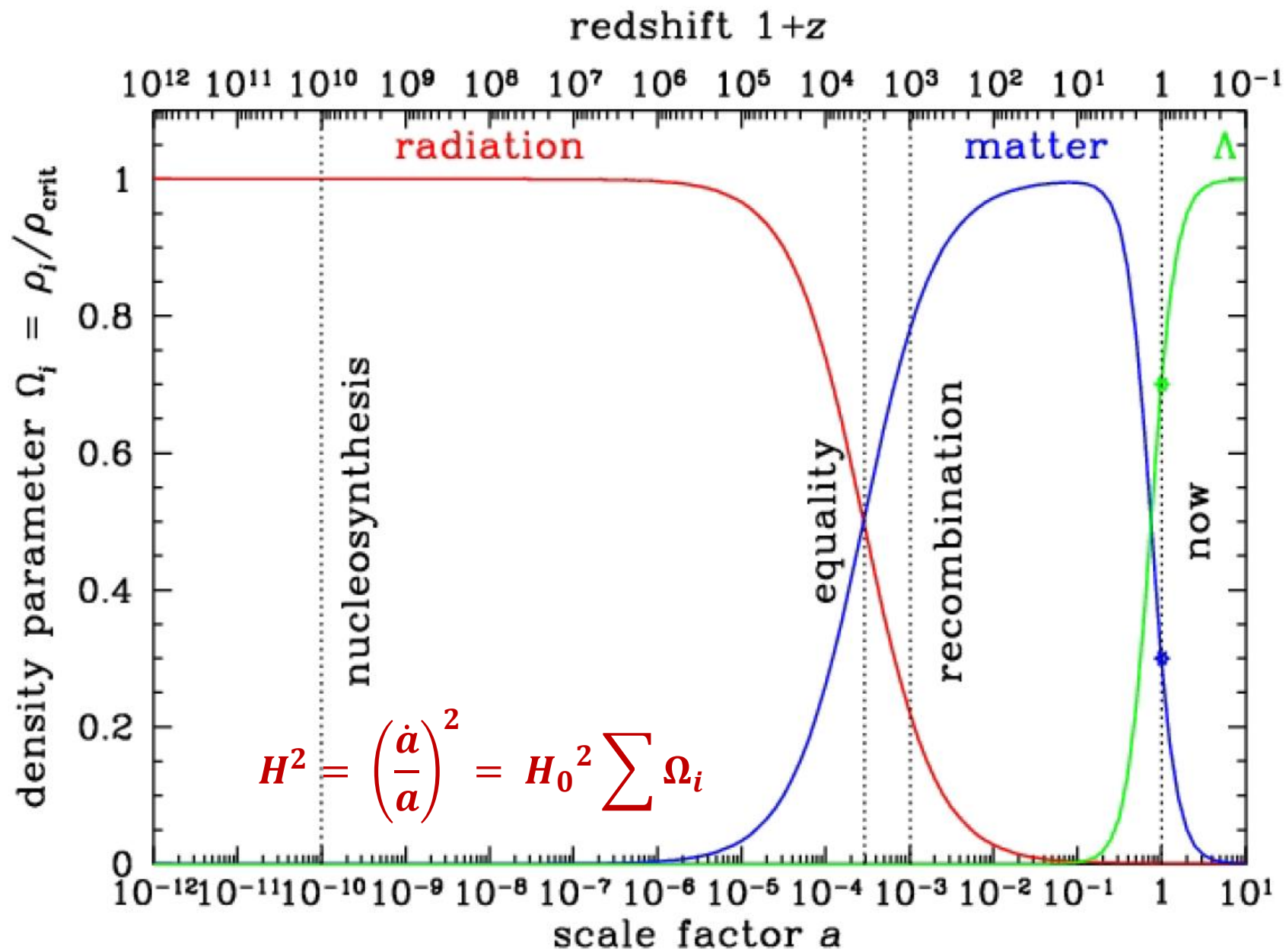


The concept for the above figure originated in a 1986 paper by Michael Turner.

Particle Data Group, LBNL © 2015

Supported by DOE





Relic Neutrino Background



At decoupling $e^+e^- \leftrightarrow \nu\bar{\nu}$ and $e\nu \leftrightarrow e\nu$: $T_\nu = T_\gamma$

Decoupling occurs when: $H \approx \Gamma = \langle \sigma v \rangle \eta$

Relic Neutrino Background



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$$T \approx \mathcal{O}(1) \left(\frac{M_W}{M_P} \right)^{1/3} \frac{M_W}{\alpha_w^{2/3}} \approx 1.5 \text{ MeV} \quad \text{A neutrino background}$$

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At $T \sim 0.5 \text{ MeV}$ ($z \sim 10^9$) : there is no energy to sustain e^+e^- pair production

Entropy balance: $S_{e^\pm} + S_\gamma$ and S_ν are diluted at the same rate

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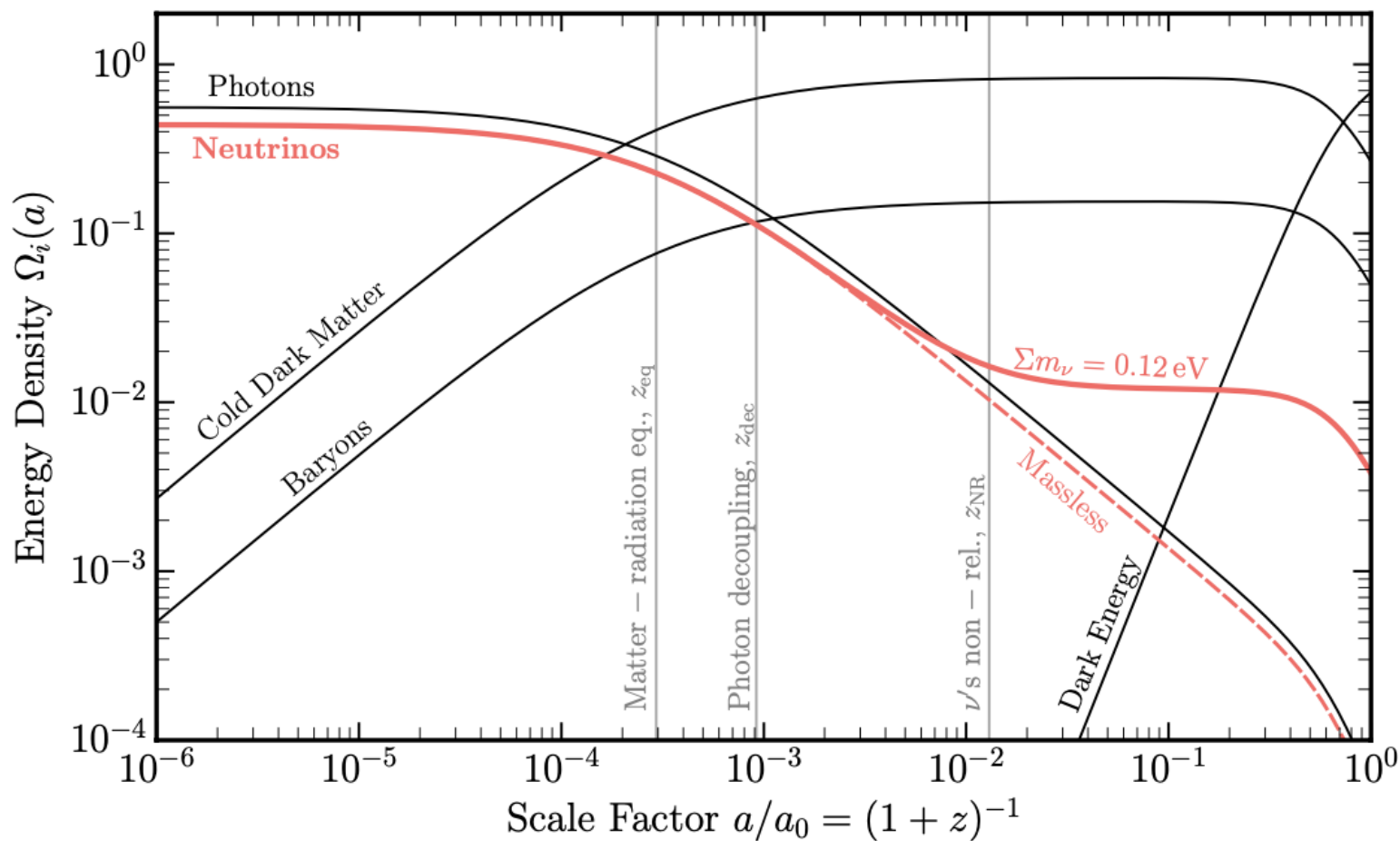
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Entropy balance: $S_{e^\pm} + S_\gamma$ and S_ν are diluted at the same rate

$$\text{Thus, after pair annihilation: } \left(\frac{T_\gamma}{T_\nu} \right) = \left(\frac{11}{4} \right)^{1/3} \quad \therefore T_\nu = 1.95 \text{ K} \\ (1.68 \times 10^{-4} \text{ eV})$$

Neutrino cosmology parameters



$$H^2(a) = H_0^2 \frac{\rho(a)}{\rho_{crit}} \quad \rho_{crit} \equiv \frac{3H_0^2}{8\pi G} \quad \Omega_r h^2 = \left(1 + \frac{7}{8} \left(\frac{4}{11}\right)^{4/3} N_{\text{eff}}\right) \Omega_\gamma h^2$$

Free streaming and structure formation



Over dense regions collapse

while

neutrinos fly away

Collapse time scale:

$$\Delta t_{\text{collapse}} \equiv (4 \pi G \rho a)^{-1/2}$$

Escape time scale:

$$\Delta t_{\text{escape}} \equiv \frac{\lambda}{v_{\text{thermal}}}$$

Evolution of primordial (matter) perturbations:

$$\ddot{\delta} + (\textit{Pressure} - \textit{gravity})\delta = 0$$

Neutrino free streaming scale:

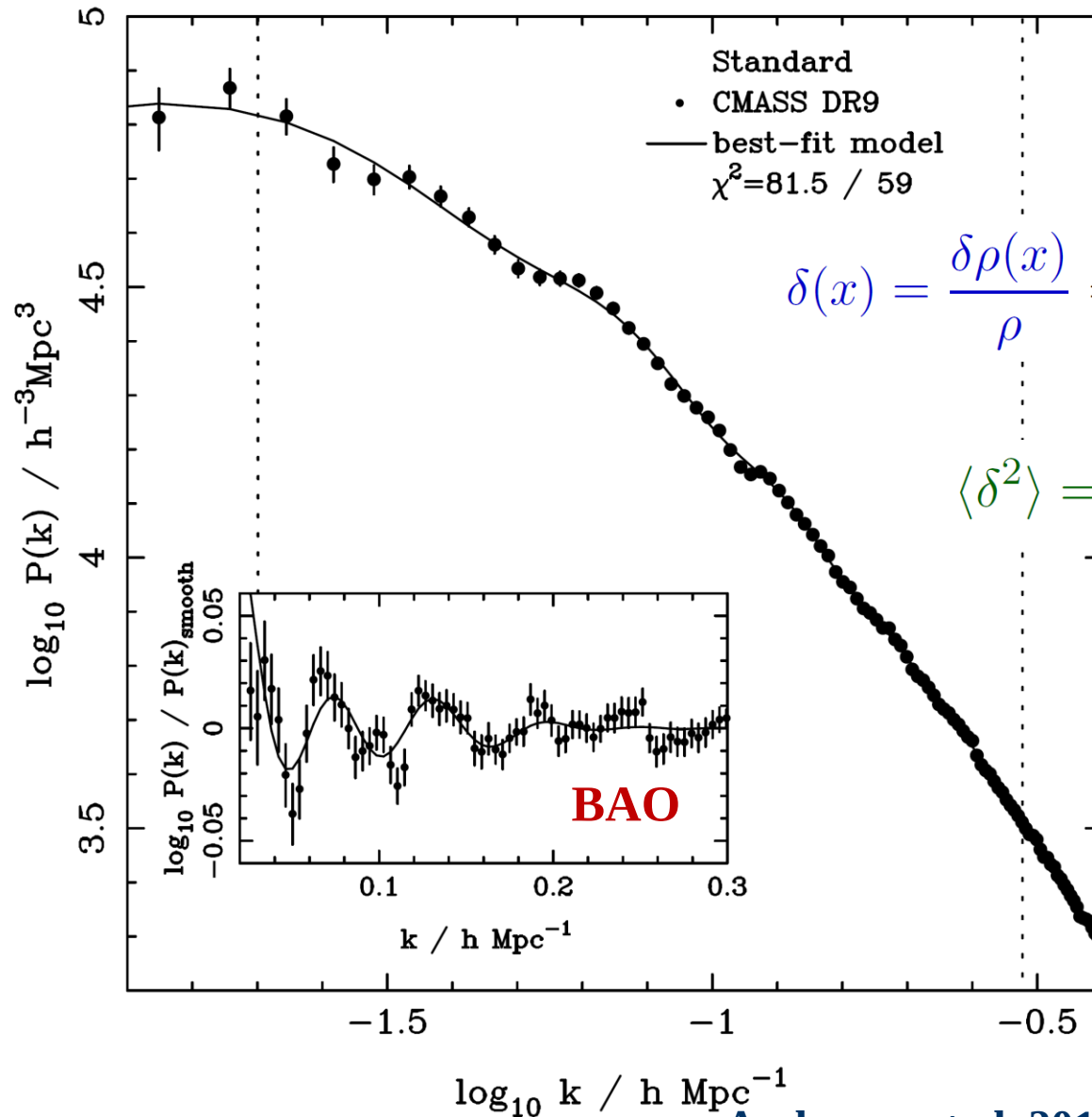
$$\lambda_{\text{FS}} \equiv \sqrt{\frac{8 \pi^2 c_v^2}{3 \Omega_m H^2}} \simeq 4.2 \sqrt{\frac{1+z}{\Omega_{m,0}}} \left(\frac{\text{eV}}{m_v} \right) h^{-1} \text{ Mpc}$$
$$k_{\text{FS}} \equiv \frac{2 \pi}{\lambda_{\text{FS}}}$$

$k \gg k_{\text{FS}}$: δ_v vanishes
metric pert. are reduced

- Shifts CMB power spectra & damps its amplitude
- Structures don't form

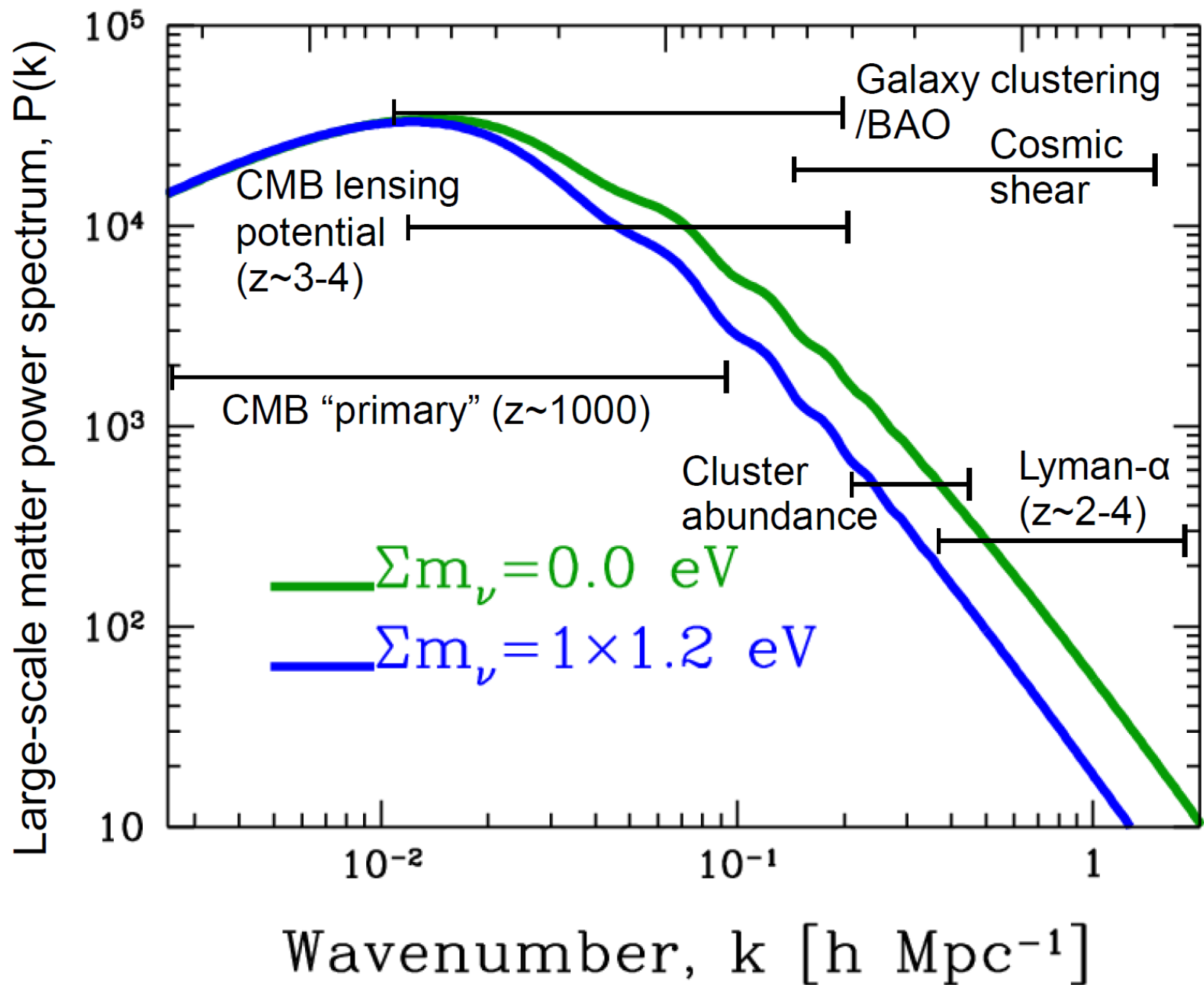


Matter power spectrum

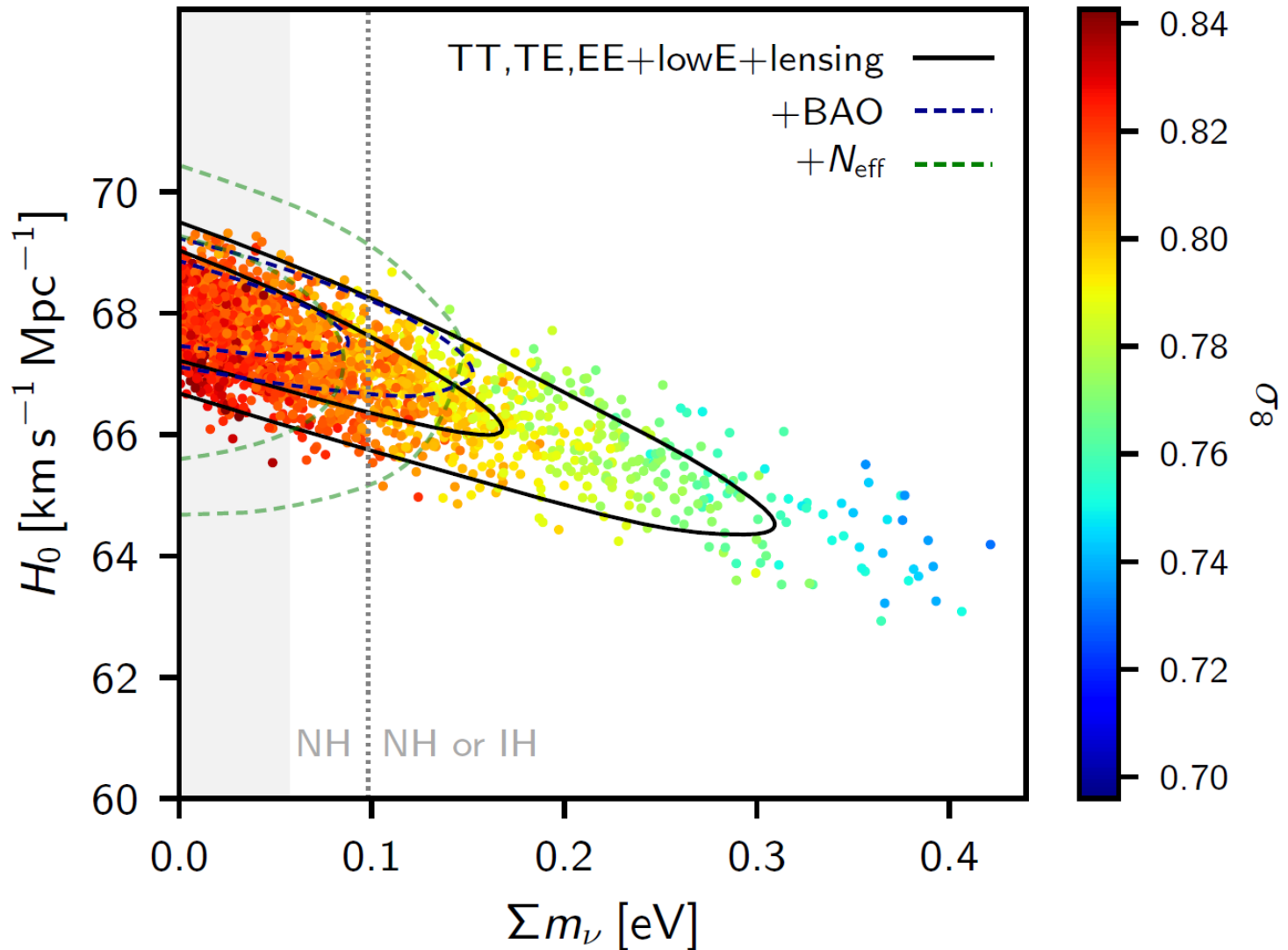


$$\delta(x) = \frac{\delta\rho(x)}{\rho} = \sum \delta_k e^{-i\vec{k}\cdot\vec{x}}$$

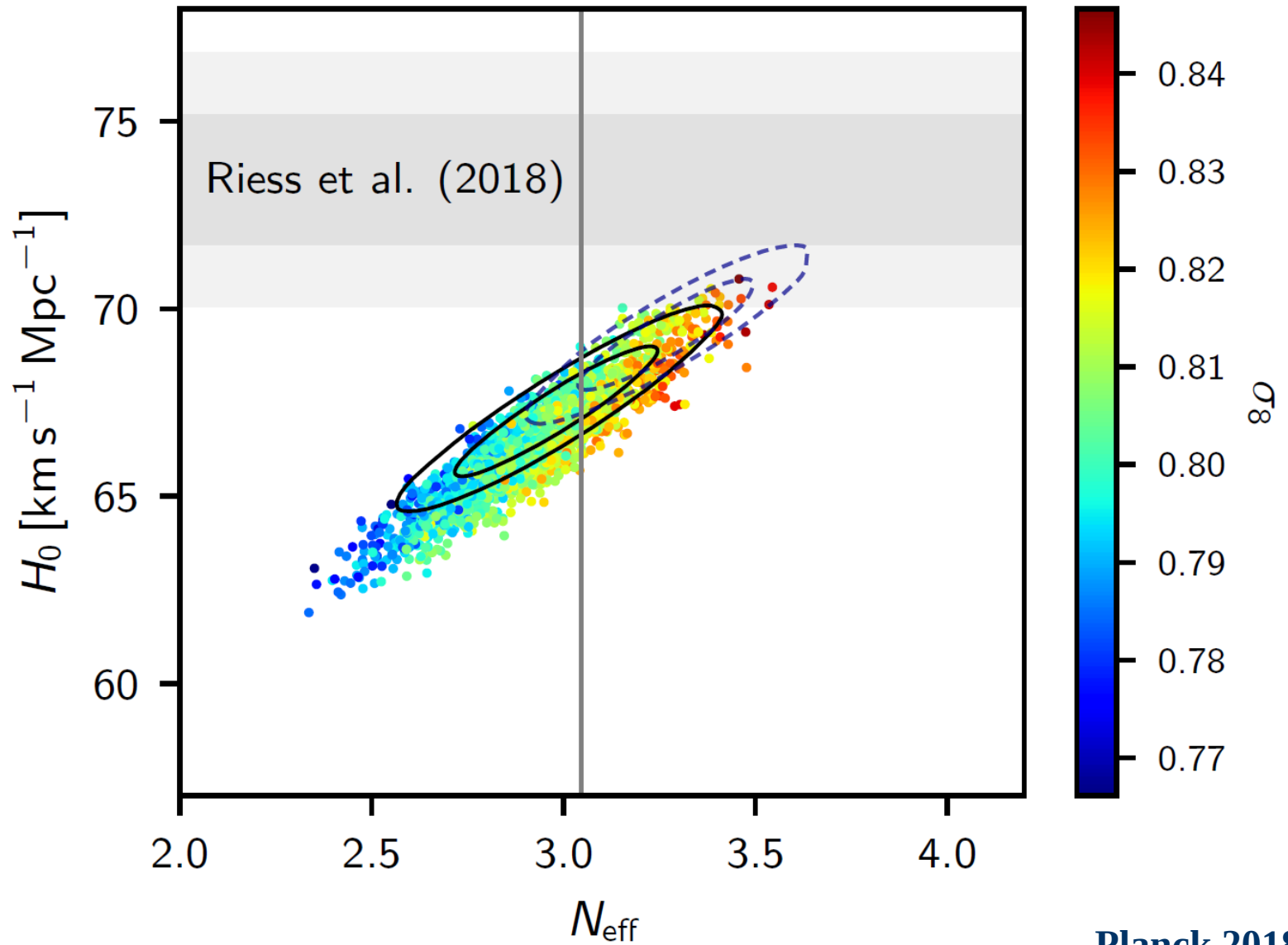
$$\langle \delta^2 \rangle = \sum P(k)$$



Neutrino mass limits



Planck 2018: $\Sigma m_\nu < 0.12$ eV (95 %, *Planck* TT,TE,EE+lowE+lensing+BAO).



Big Bang Nucleosynthesis



Nucleosynthesis era:

$$100 \text{ MeV} \gtrsim T \gtrsim 0.05 \text{ MeV}$$

$$t \sim 3 - 5 \text{ min}$$

$$10^8 < z < 10^{12}$$

**But at $T \sim 1 - 2 \text{ MeV}$
($t \sim 0.2 \text{ sec}$)**

neutrinos decouple

$$e + p \leftrightarrow \nu + n$$

$$e + n \leftrightarrow \bar{\nu} + p$$

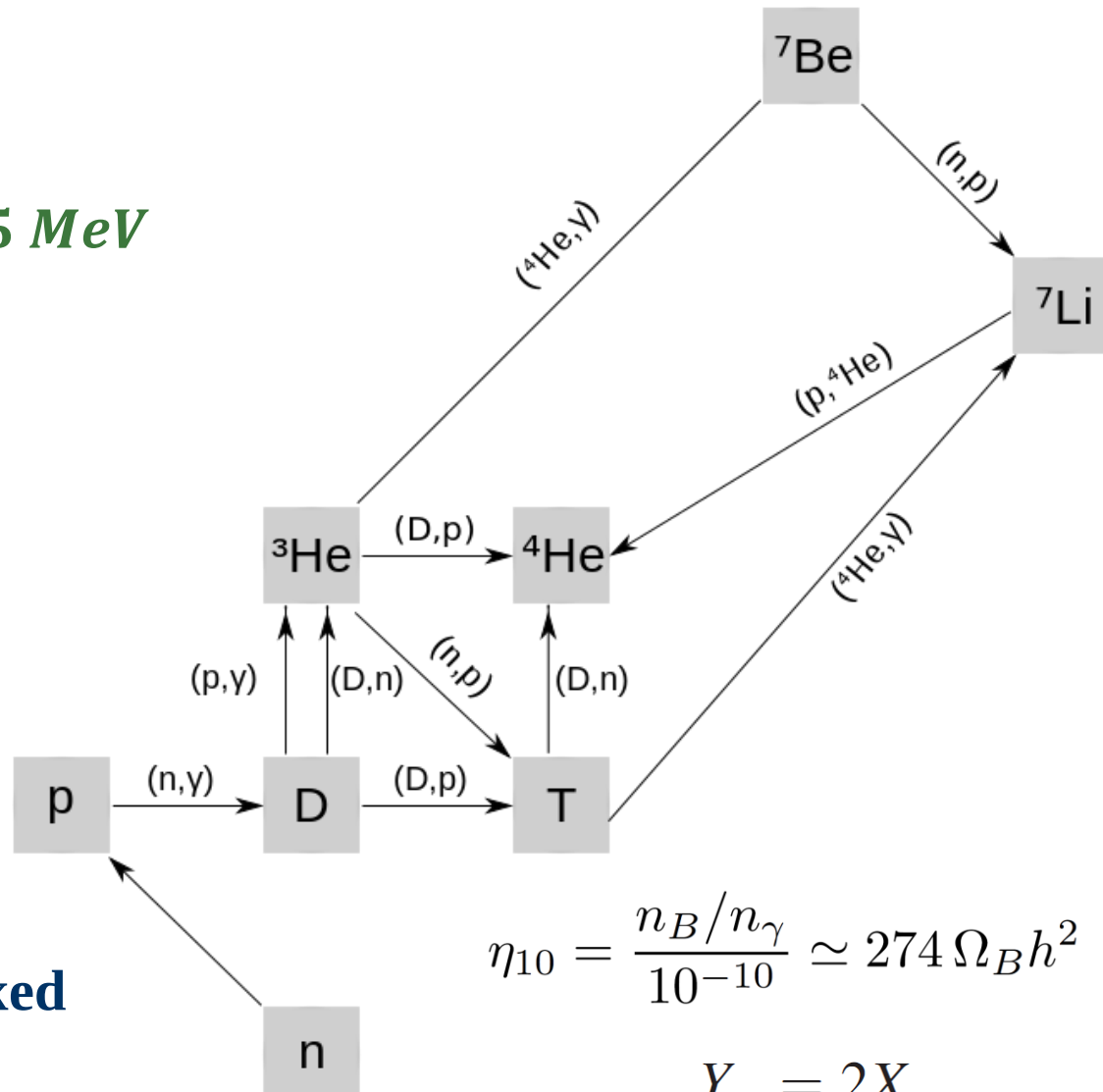
$$Y_p \approx \frac{2(n/p)}{1 + (n/p)} \text{ gets fixed}$$

$$\frac{dX_n}{dt} = \lambda_{np} [(1 - X_n)e^{-\Delta m/T} - X_n]$$

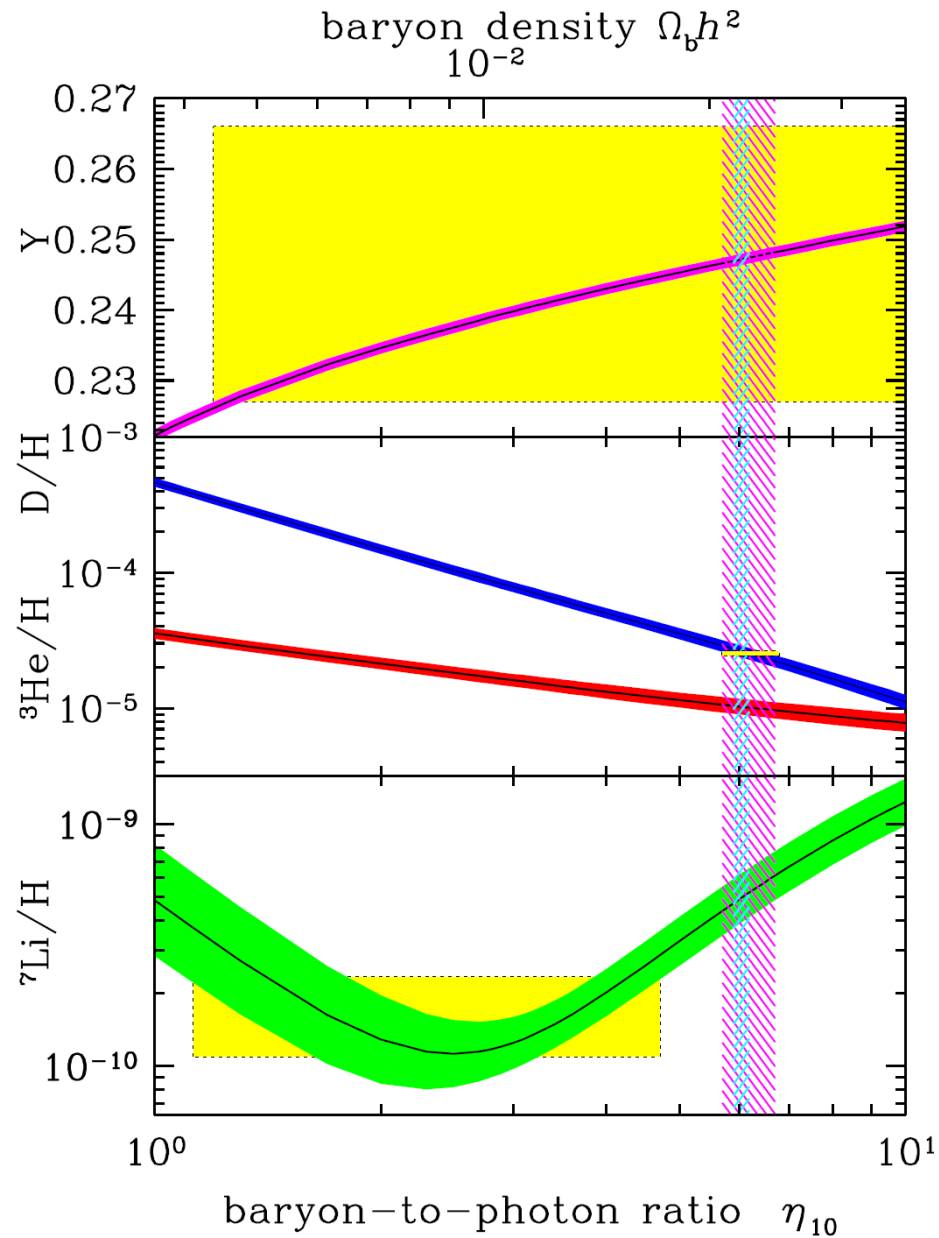
$$\eta_{10} = \frac{n_B/n_\gamma}{10^{-10}} \simeq 274 \Omega_B h^2$$

$$Y_p = 2X_n$$

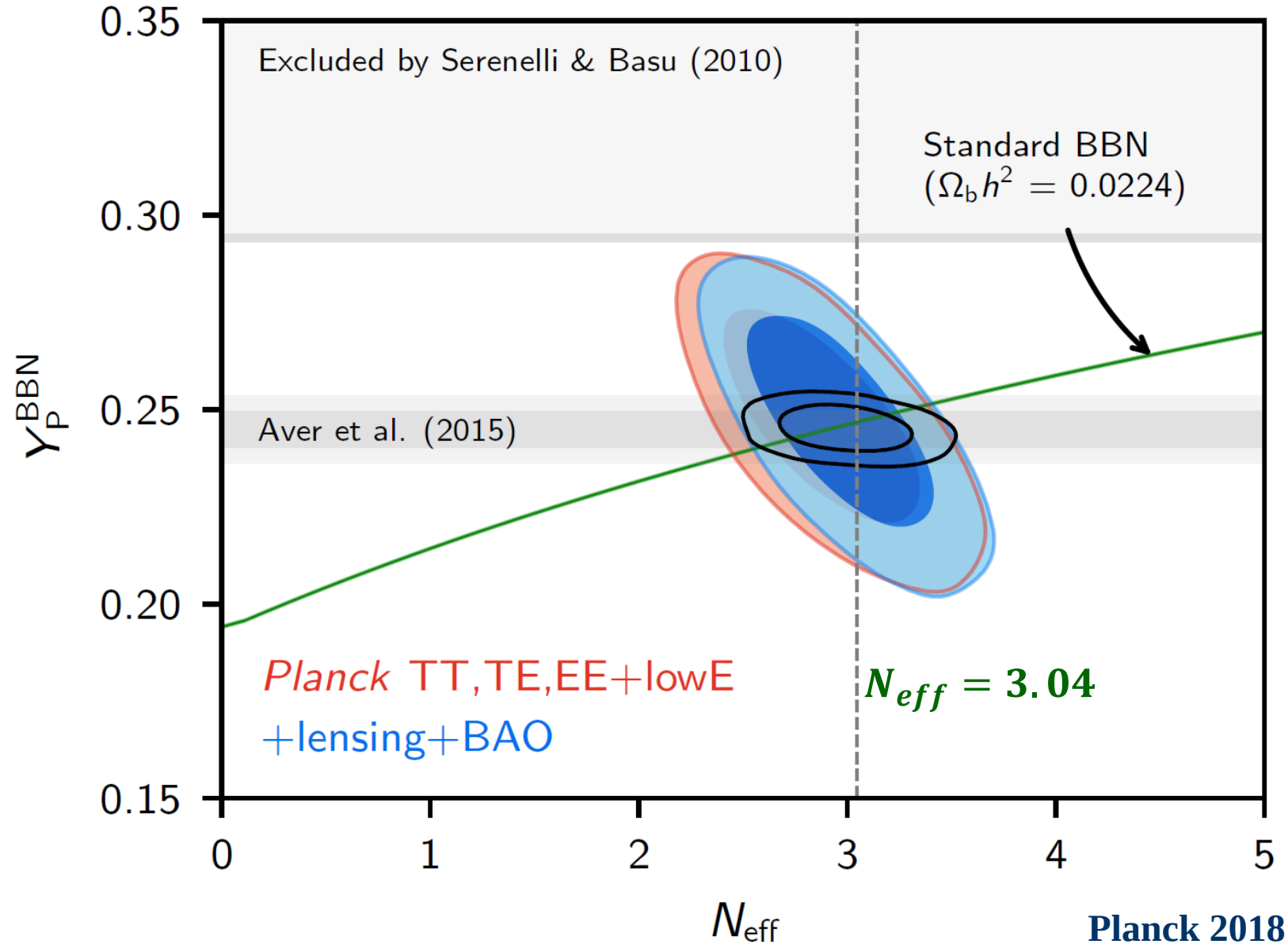
$$\lambda_{np} = n_\nu^{(0)} \langle \sigma v \rangle$$



Nucleosynthesis outputs

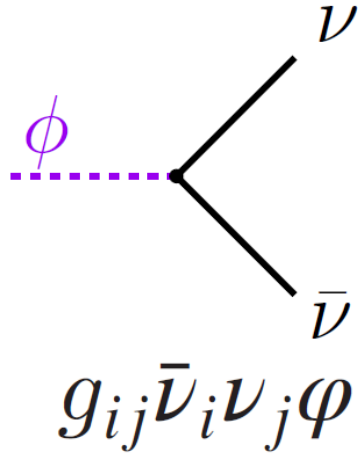


N_{eff}

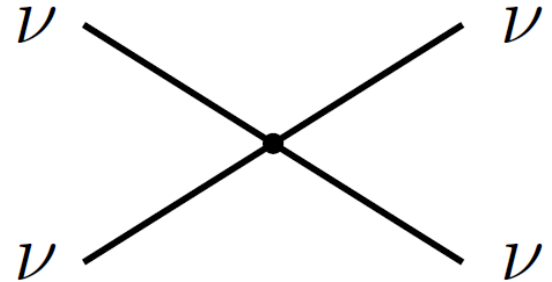




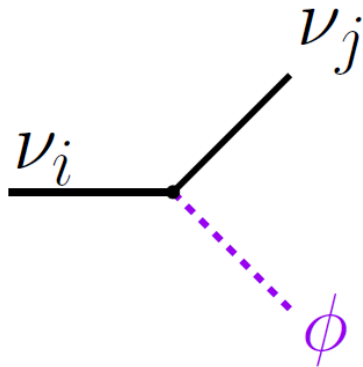
Neutrinophilic bosons



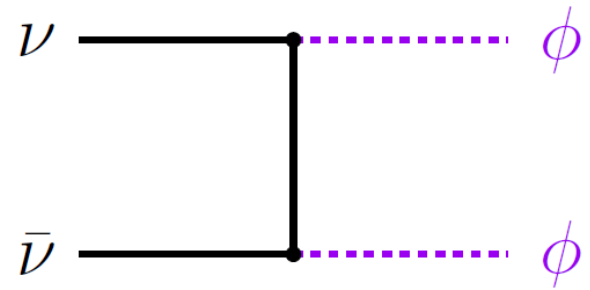
Neutrino scatterings



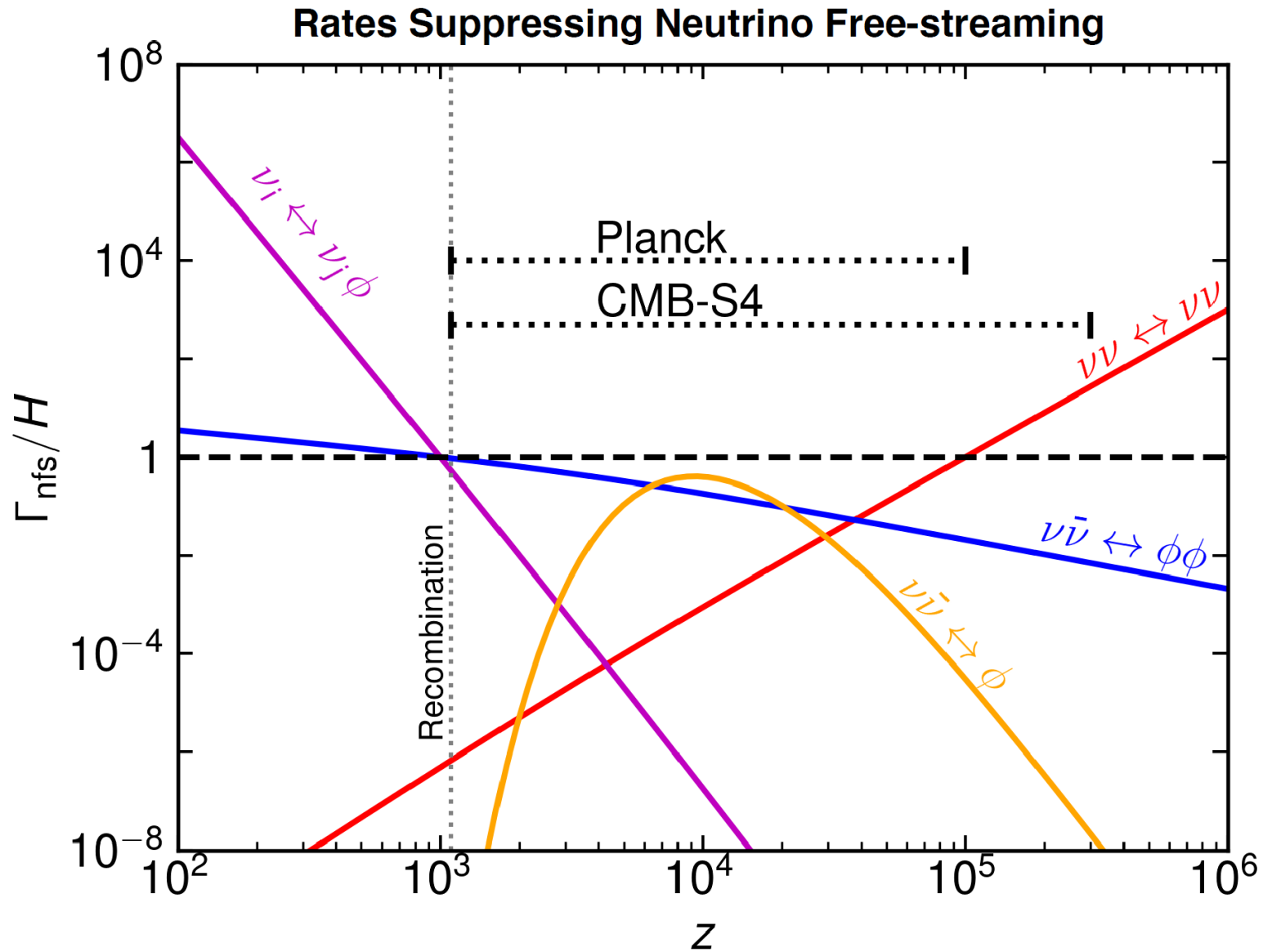
Neutrino decays



Neutrino annihilations



Neutrino - scalar NSI





$$\sigma_\nu \rightarrow \delta G_{\mu\nu} \rightarrow \delta T_{\mu\nu}|_\gamma \rightarrow \delta T_\gamma$$



$$\sigma_\nu \rightarrow \delta G_{\mu\nu} \rightarrow \delta T_{\mu\nu}|_\gamma \rightarrow \delta T_\gamma$$

Phase space distribution function

$$f(\mathbf{x}, \mathbf{q}, \tau) = f^0(q) [1 + \Theta(\mathbf{x}, q, \mathbf{n}, \tau)]$$



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Phase space distribution function

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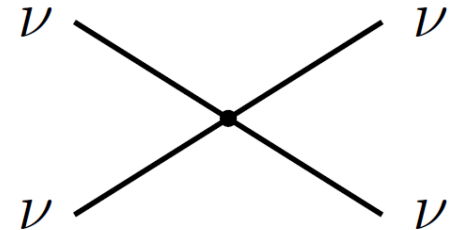
In synchronous gauge:

Boltzmann Eq. (in Fourier space) $\mu = \hat{k} \cdot \hat{n}$

$$f^0 \left[\frac{\partial \Theta}{\partial \tau} + ik \frac{q}{\epsilon} \mu \Theta + \frac{d \ln f_\alpha^0}{d \ln q} \left[\dot{\eta} - \frac{\dot{h} + 6\dot{\eta}}{2} \mu^2 \right] \right] = a \tilde{C}^{(1)}[f]$$



$$\nu_i(\mathbf{p}_1) + \nu_j(\mathbf{p}_2) \leftrightarrow \nu_k(\mathbf{p}_3) + \nu_l(\mathbf{p}_4).$$



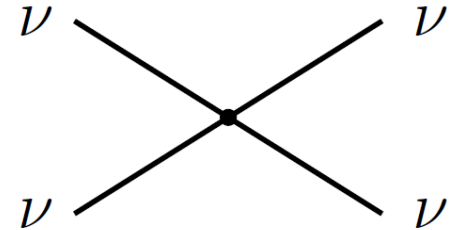
$$\tilde{C} = \frac{1}{2E} \int d\Pi_2 d\Pi_3 d\Pi_4 (2\pi)^4 \delta^4(P + P_2 - P_3 - P_4) \mathcal{M}^2 F(\mathbf{k}, \mathbf{p}, \mathbf{p}_2, \mathbf{p}_3, \mathbf{p}_4, \tau)$$

$$d\Pi = \frac{g_* d^3 p}{2E (2\pi)^3}$$

$$F(\mathbf{x}, \mathbf{p}_1, \mathbf{p}_2, \mathbf{p}_3, \mathbf{p}_4, \tau) = f(\mathbf{x}, \mathbf{p}_3, \tau) f(\mathbf{x}, \mathbf{p}_4, \tau) [1 - f(\mathbf{x}, \mathbf{p}_1, \tau)] [1 - f(\mathbf{x}, \mathbf{p}_2, \tau)] \\ - f(\mathbf{x}, \mathbf{p}_1, \tau) f(\mathbf{x}, \mathbf{p}_2, \tau) [1 - f(\mathbf{x}, \mathbf{p}_3, \tau)] [1 - f(\mathbf{x}, \mathbf{p}_4, \tau)]$$



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Legendre mode decomposition:

$$\Theta(\mathbf{k}, q, \mathbf{n}, \tau) = \sum_{\ell=0}^{\infty} (-i)^\ell (2\ell + 1) \theta_\ell(k, q, \tau) P_\ell(\mu)$$



$$\frac{\partial \theta_\ell}{\partial \tau} - \frac{kq}{(2\ell + 1)\epsilon} [\ell \theta_{\ell-1} - (\ell + 1)\theta_{\ell+1}] + \frac{d \ln f^0}{d \ln q} \left[\frac{\dot{h} + 6\dot{\eta}}{15} \delta_{\ell 2} - \frac{\dot{h}}{6} \delta_{\ell 0} \right] = \frac{a}{f^0} \tilde{C}_\ell$$

$$\tilde{C}^{(1)} = \frac{g_*^3 T}{128\pi^3 z} \frac{d \ln f^{(0)}(z)}{d \ln p} \sum_{\ell=0}^{\infty} (-i)^\ell (2\ell + 1) \vartheta_\ell P_\ell(\mu) [\mathcal{A}(z) + \mathcal{B}_\ell(z) - 2\mathcal{D}_\ell(z)]$$

$$z = \epsilon/T_0$$



$$f(x^i, q_j, \tau) = f_0(q) [1 + \Psi(x^i, q, \hat{n}_j, \tau)]$$

Relaxation time approximation $c/f_0 \approx a \Gamma \Psi$

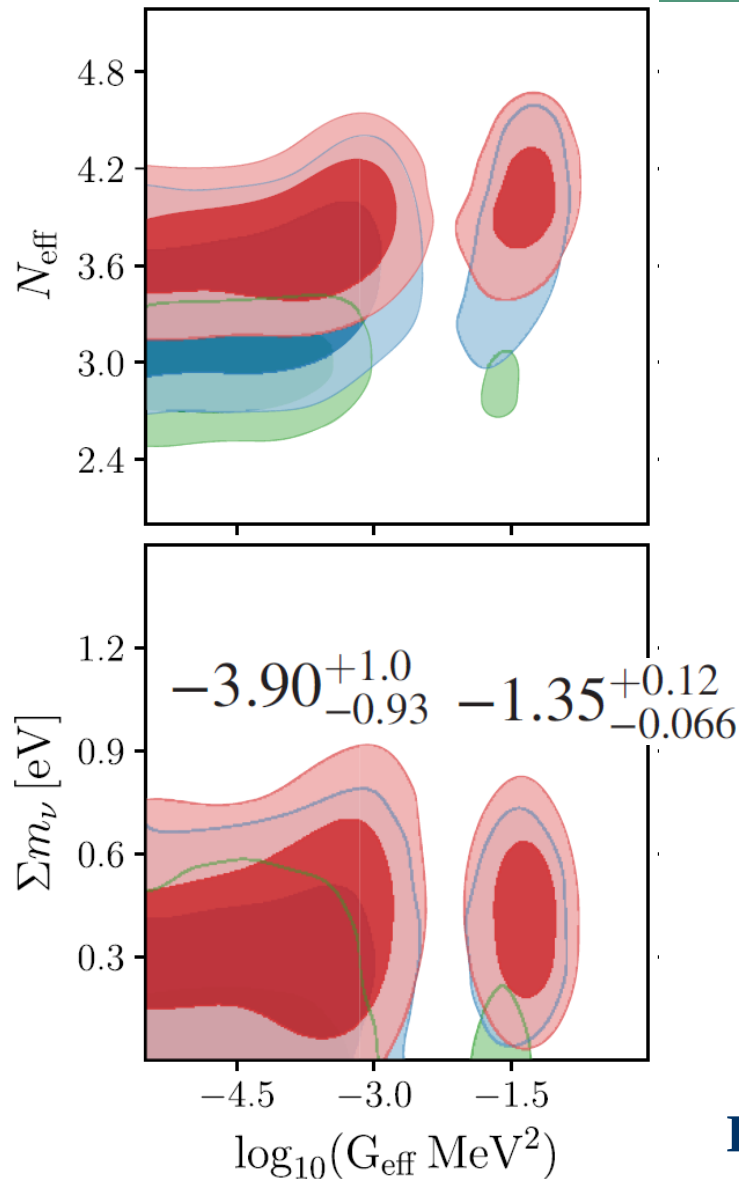
$$\dot{\Psi}_0 = -\frac{qk}{\epsilon} \Psi_1 + \frac{\dot{h}}{6} \frac{d \ln f_0}{d \ln q},$$

$$\dot{\Psi}_1 = \frac{qk}{3\epsilon} (\Psi_0 - 2\Psi_2),$$

$$\dot{\Psi}_2 = \frac{qk}{5\epsilon} (2\Psi_1 - 3\Psi_3) - \left(\frac{\dot{h}}{15} + \frac{2\dot{\eta}}{5} \right) \frac{d \ln f_0}{d \ln q} - a\Gamma\Psi_2,$$

$$\dot{\Psi}_{l \geq 3} = \frac{qk}{(2l+1)\epsilon} (l\Psi_{l-1} - (l+1)\Psi_{l+1}) - a\Gamma\Psi_l.$$

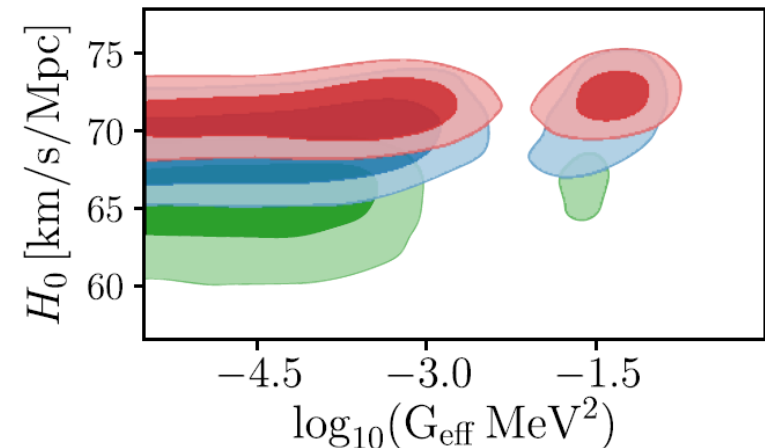
NSI bounds for heavy scalar mediator



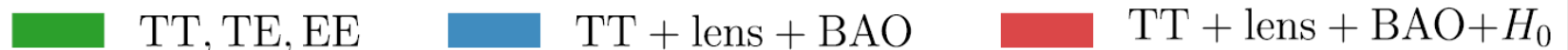
Heavy scalar mediator:

$$|\mathcal{M}|^2 = 2G_{\text{eff}}^2(s^2 + t^2 + u^2)$$

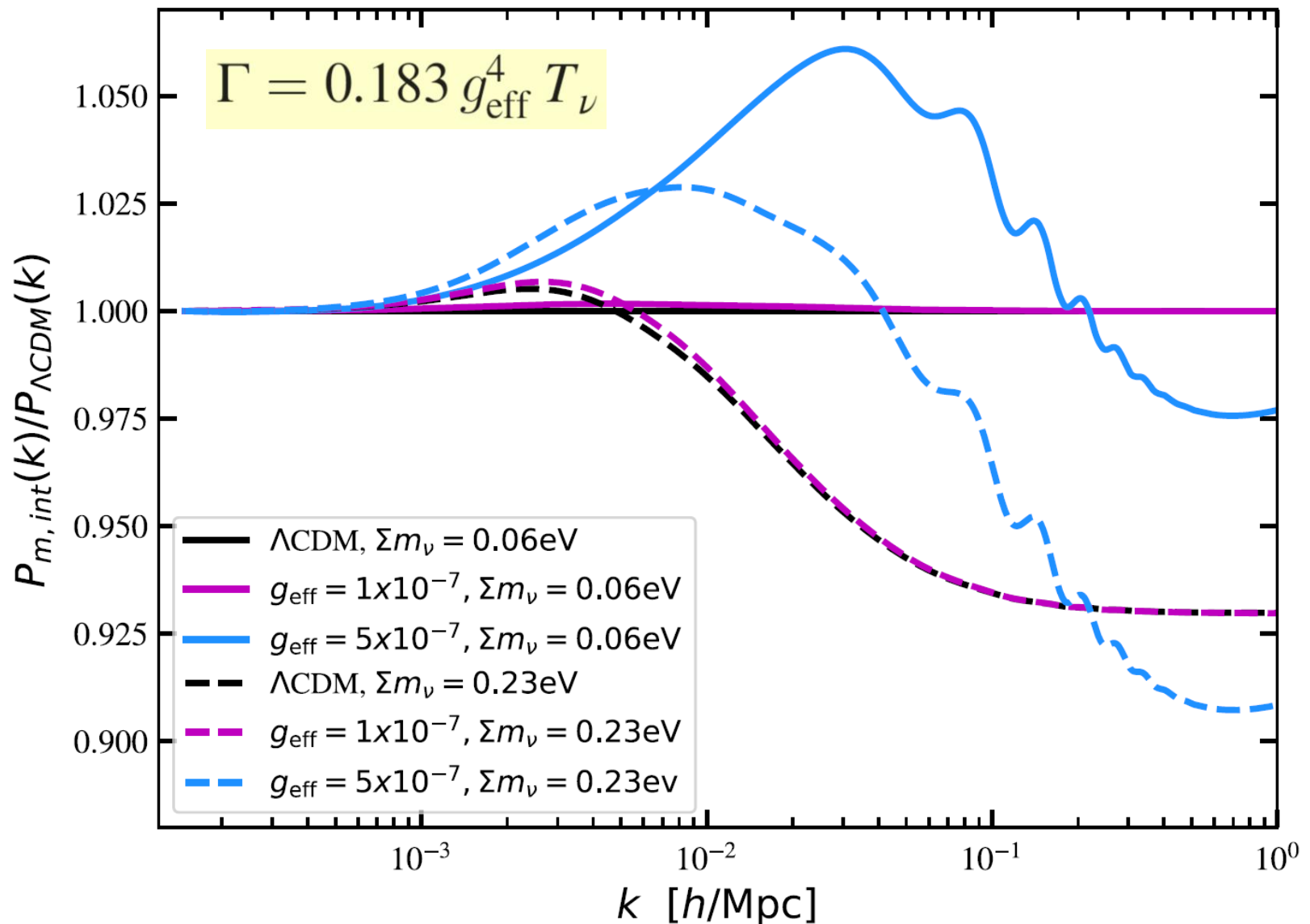
$$G_{\text{eff}} = |g_\nu|^2 / m_\phi^2$$

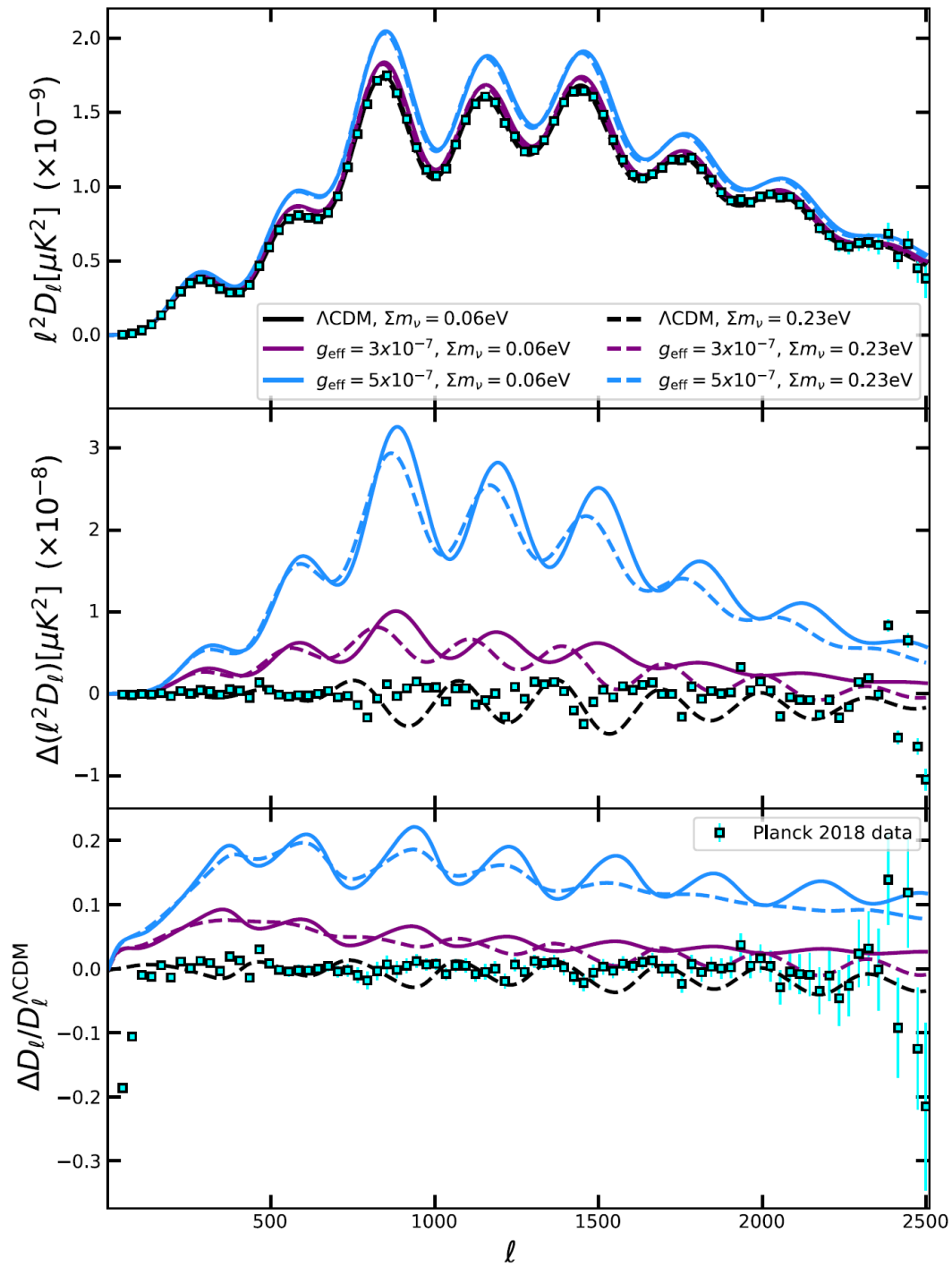


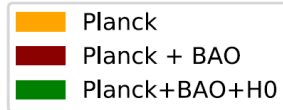
Kreisch, Cyr-Racine, Dore, PRD 2020



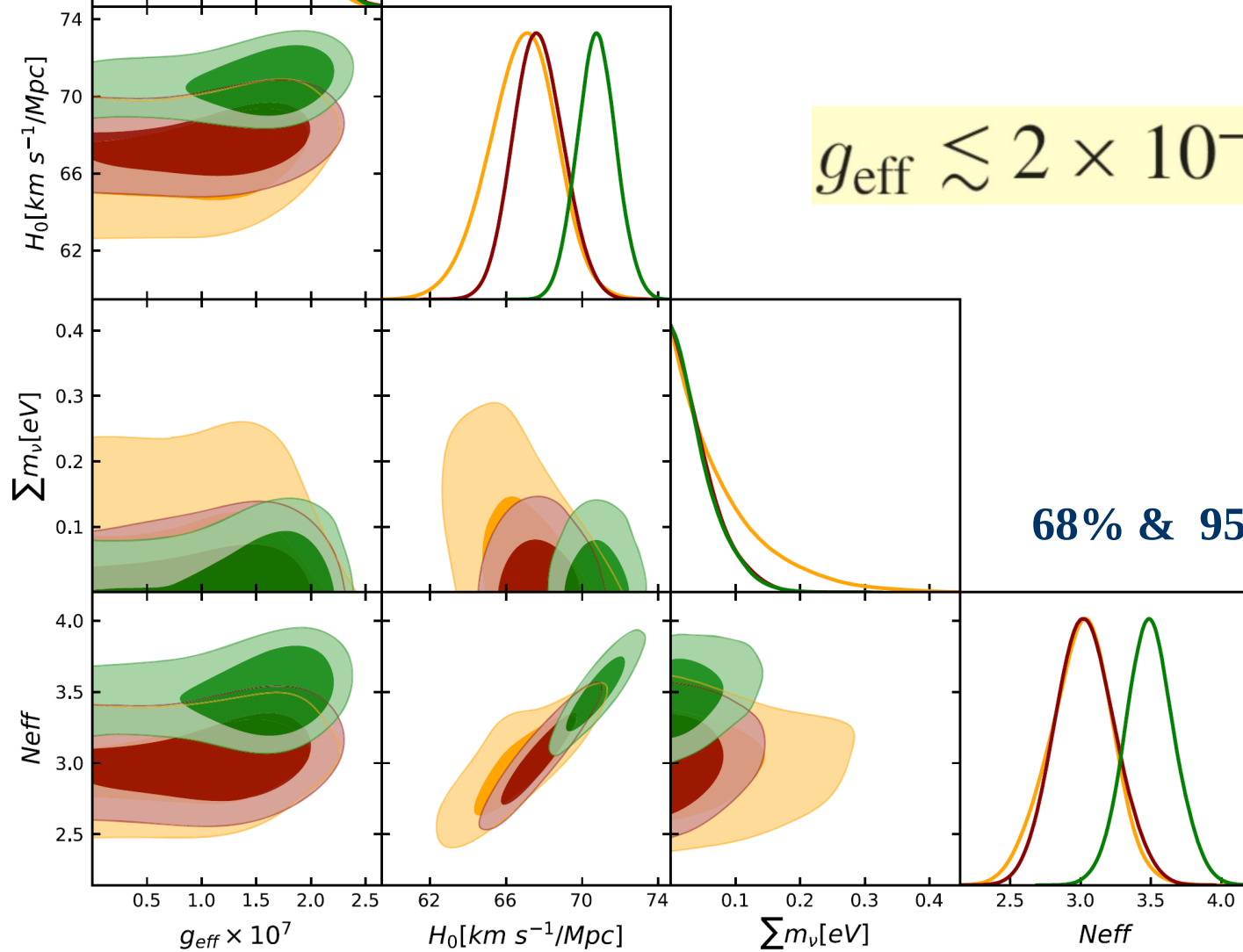
NSI bounds for light scalar mediator







**J. Venzor, G. García-Arrollo, APL, J. de Santiago,
PRD 105, 123539 (2022)**

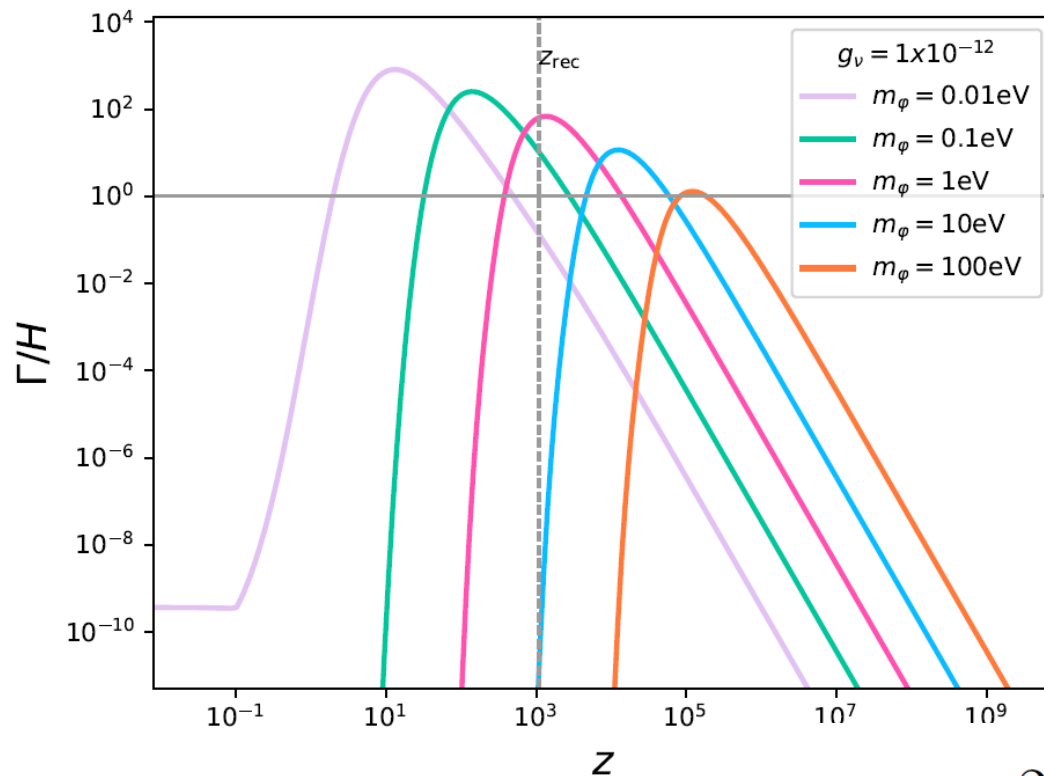




Resonant NSI effects

$$10^{-3} \text{eV} \ll m_\phi \ll 10^3 \text{eV}$$

$$\sigma(s) = \frac{g_v^4}{4\pi} \frac{s}{(s - m_\phi^2)^2 + \Gamma_\phi^2 m_\phi^2} \xrightarrow{\Gamma_\phi \rightarrow 0} \frac{g_v^2}{m_\phi^2} s \delta(s - m_\phi^2)$$



$$\Gamma_{\text{scatt}} = \langle \sigma v_{\text{MOL}} \rangle n_\nu,$$

$$\Gamma_{\text{scatt}} = \frac{g_v^2 \pi^5 m_\phi^2}{24 \zeta(3) T_\nu} F(m_\phi^2; T)$$

$$z_{\text{peak}} \sim 2.13 \times 10^3 \left(\frac{m_\phi}{\text{eV}} \right)$$

Preliminary!

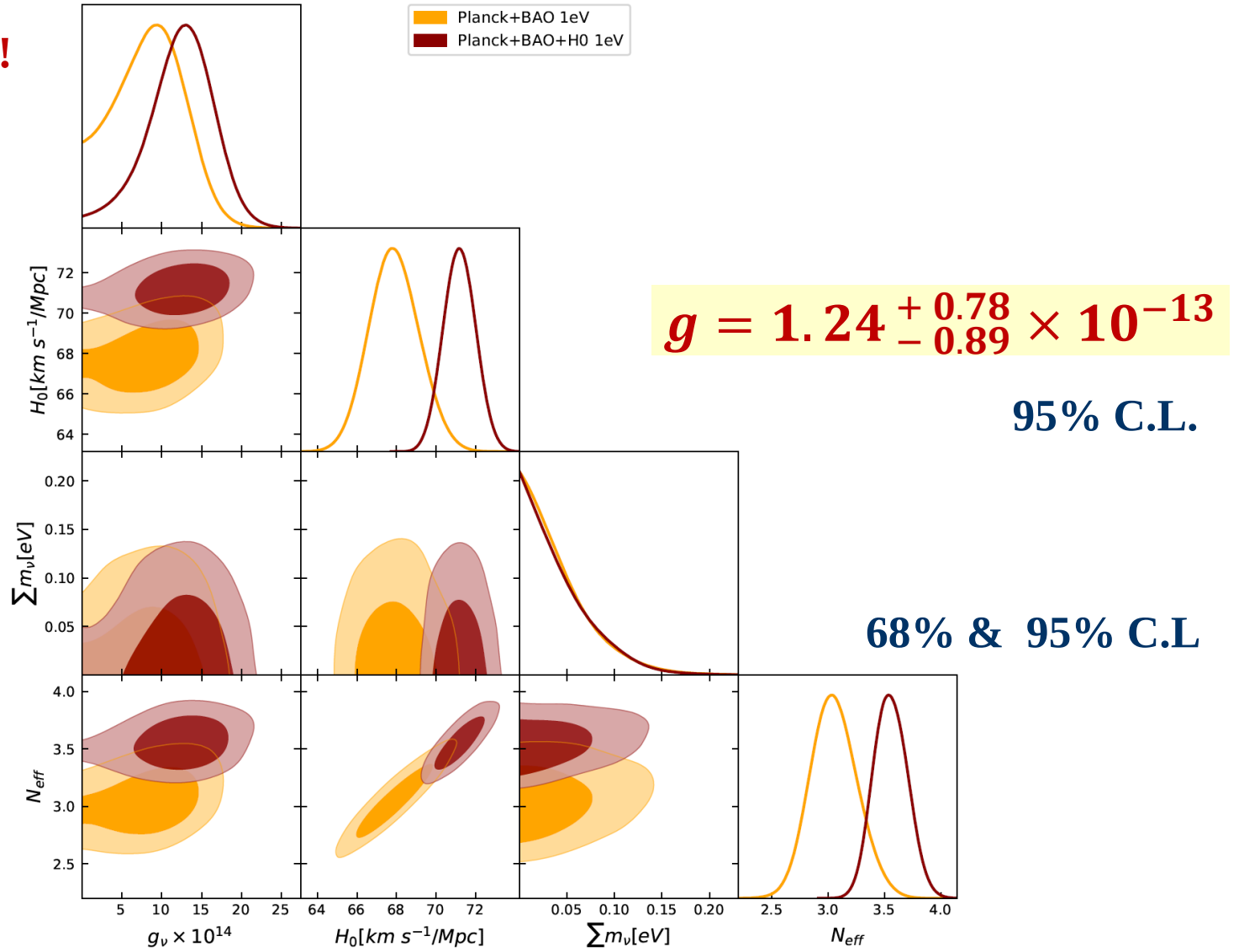


FIG. 11. Posterior probability for different sets of data and for $M_\phi = 1\text{eV}$

Preliminary!

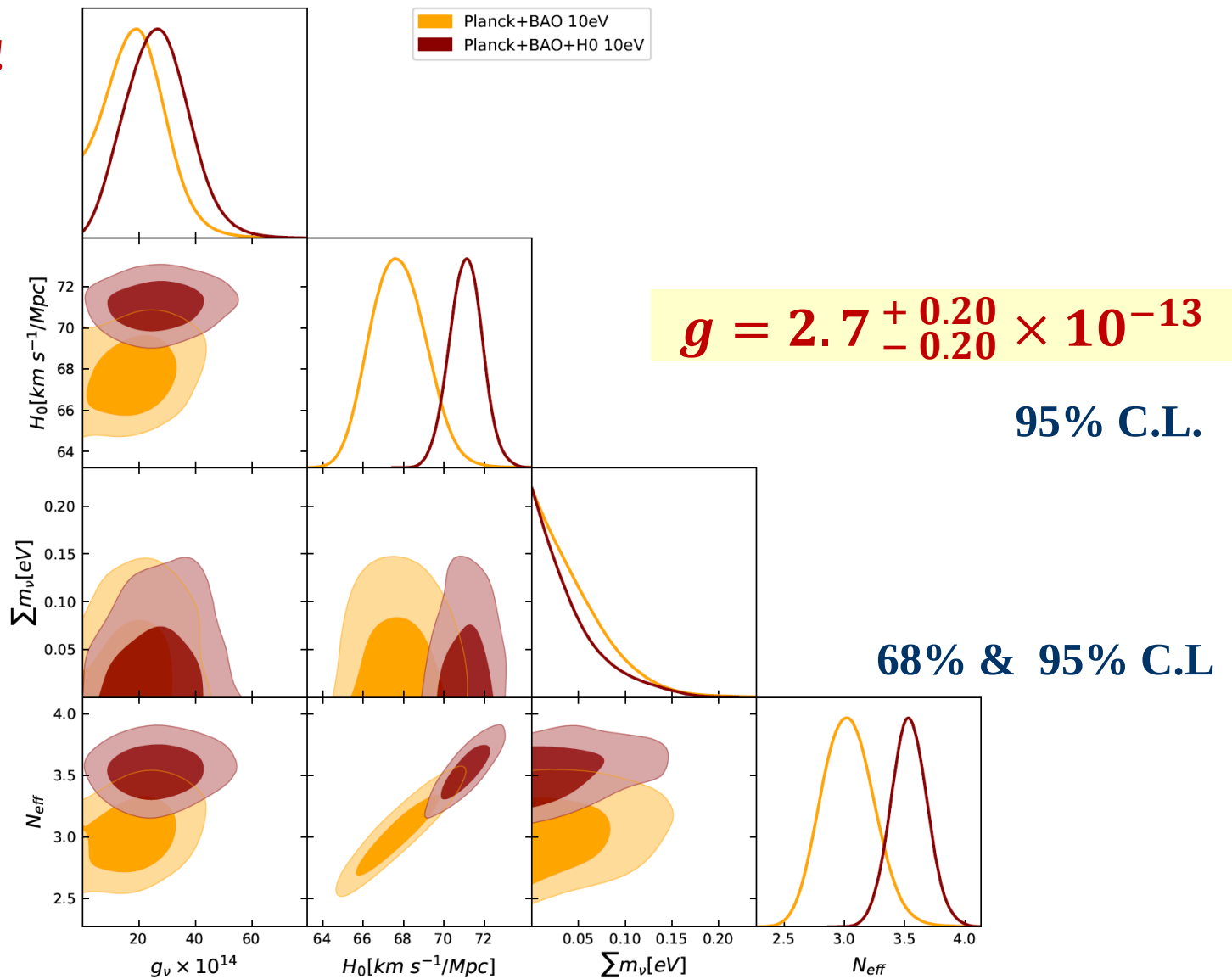
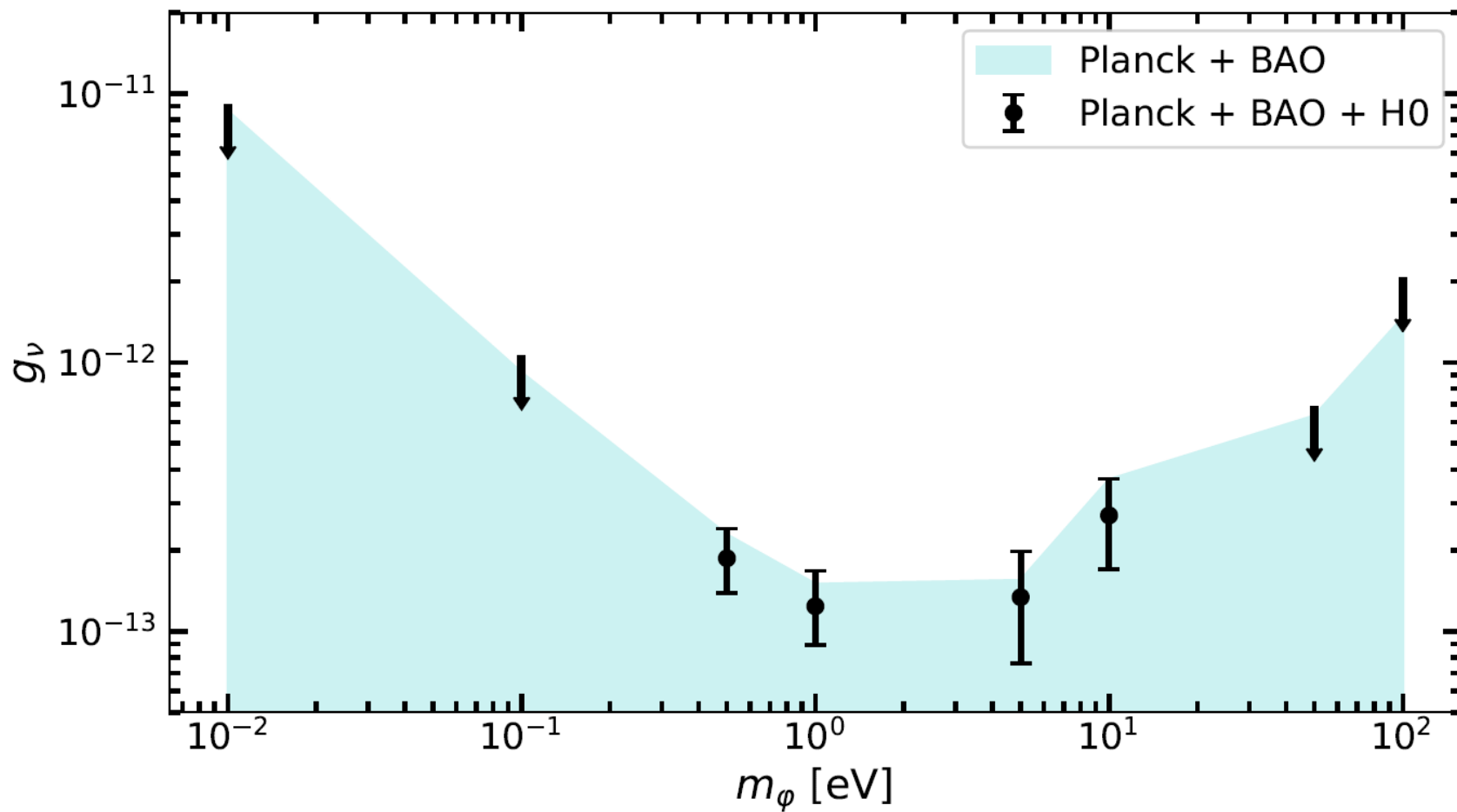


FIG. 12. Posterior probability for different sets of data and for $M_\phi = 10\text{eV}$

Resonant NSI bounds

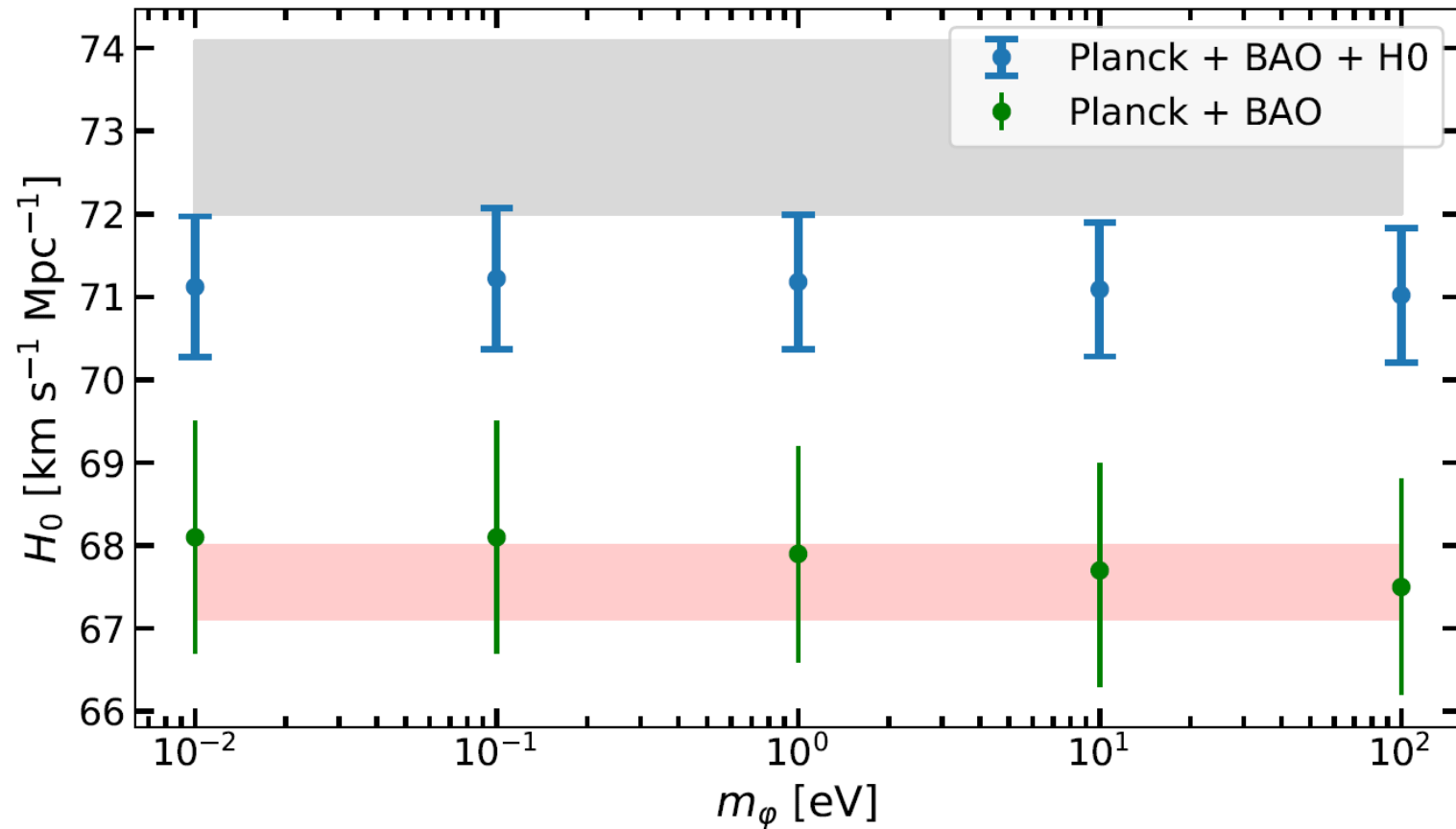


Preliminary!

Resonant NSI effect on H_0



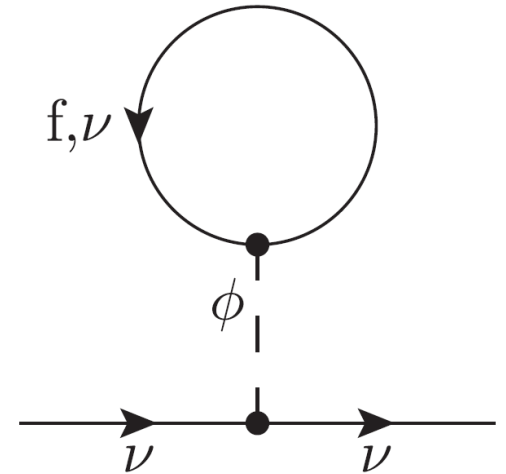
Preliminary!



Tension decreased down to about $2.8\sigma - 3.3\sigma$



$$\Delta m(m_f; T) = \frac{m_f}{\pi^2} \int_{m_f}^{\infty} dk \sqrt{k^2 - m_f^2} f(k)$$

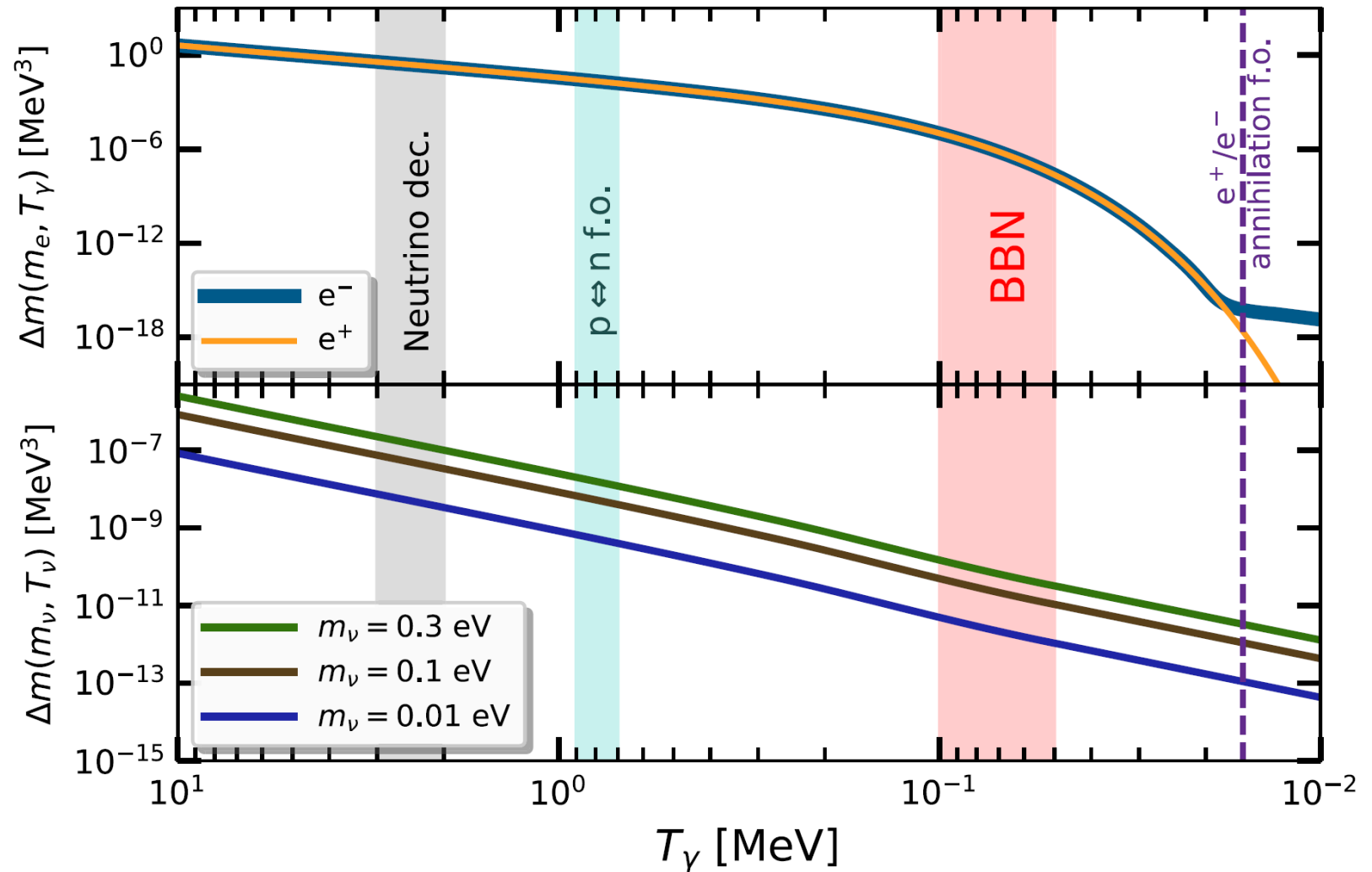


$$m_{\text{eff}} = m_{\nu} + 2G_{\text{eff}}\Delta m(m_e; T_{\gamma}) + 3G_S\Delta m(m_{\nu}; T_{\nu})$$

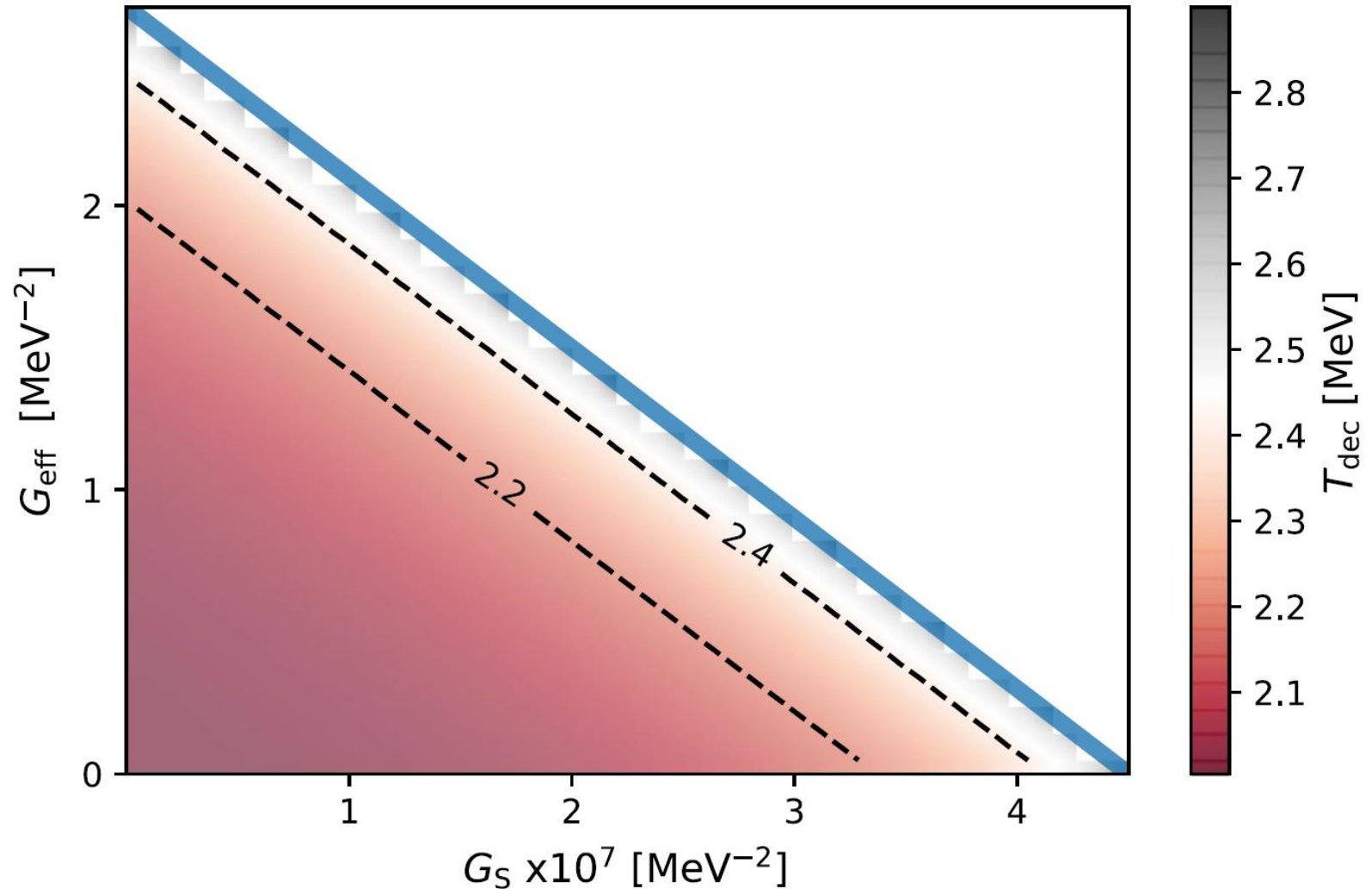
$$G_{\text{eff}} = \frac{g_f g_{\nu}}{m_{\phi}^2}$$

$$G_S = \frac{g_{\nu}^2}{m_{\phi}^2}$$

Effective mass from ν NSI

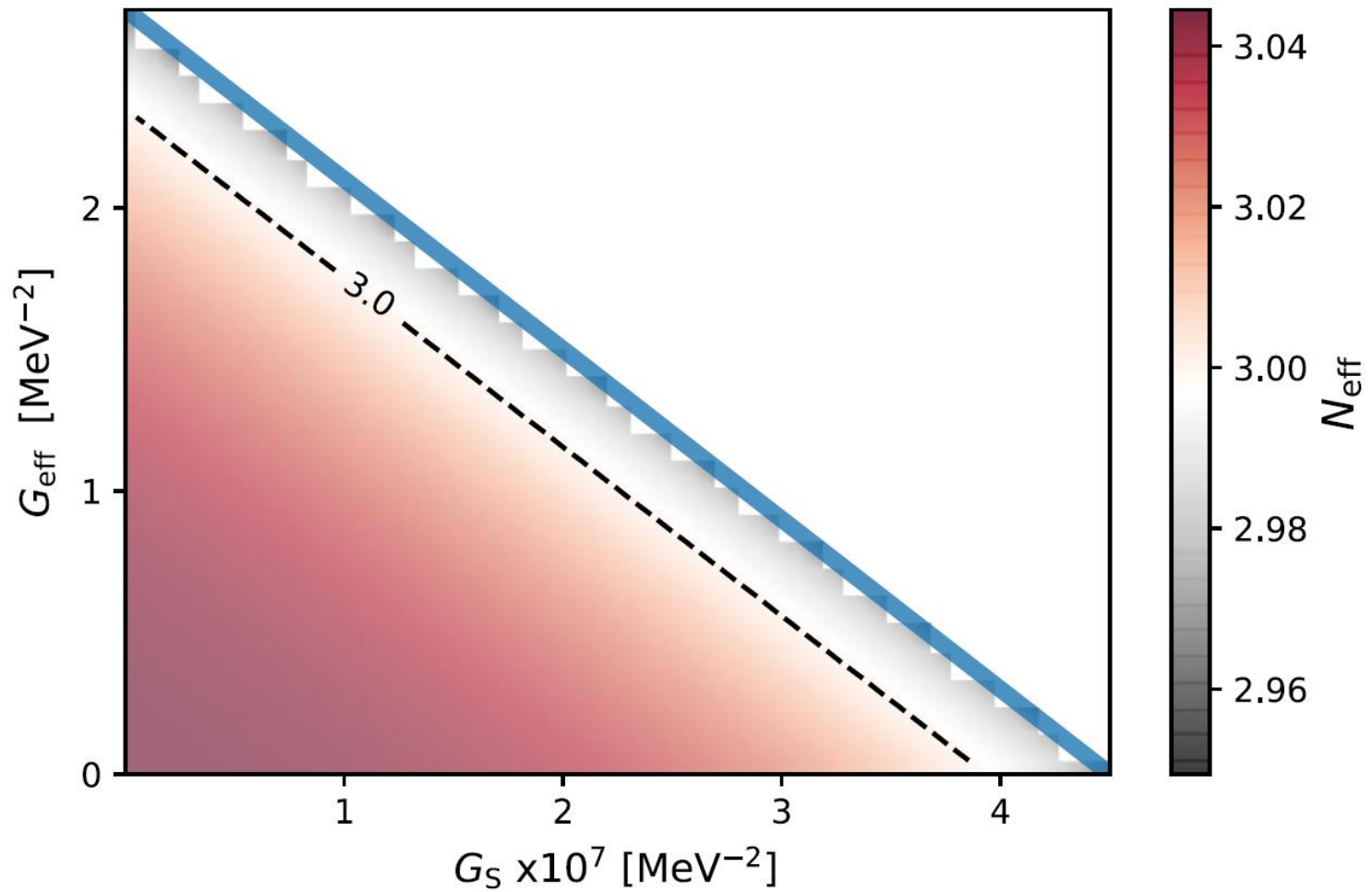


NSI effective mass effects on T_{dec}



J. Venzor, APL, J. de Santiago, PRD 103 (2020)

Light scalar NSI effects on N_{eff}

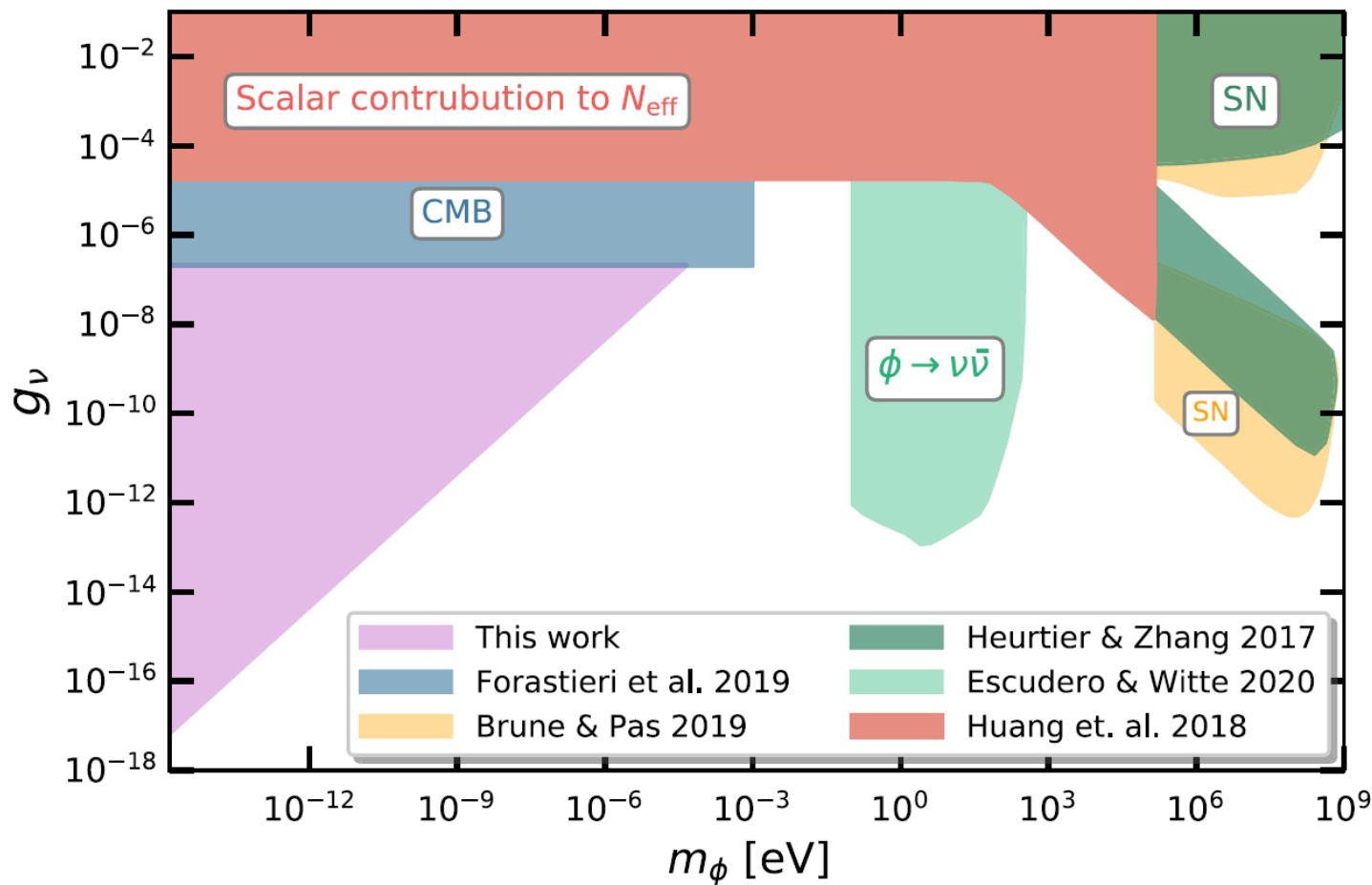


J. Venzor, APL, J. de Santiago, PRD 103 (2020)

BBN bounds on light scalar NSI



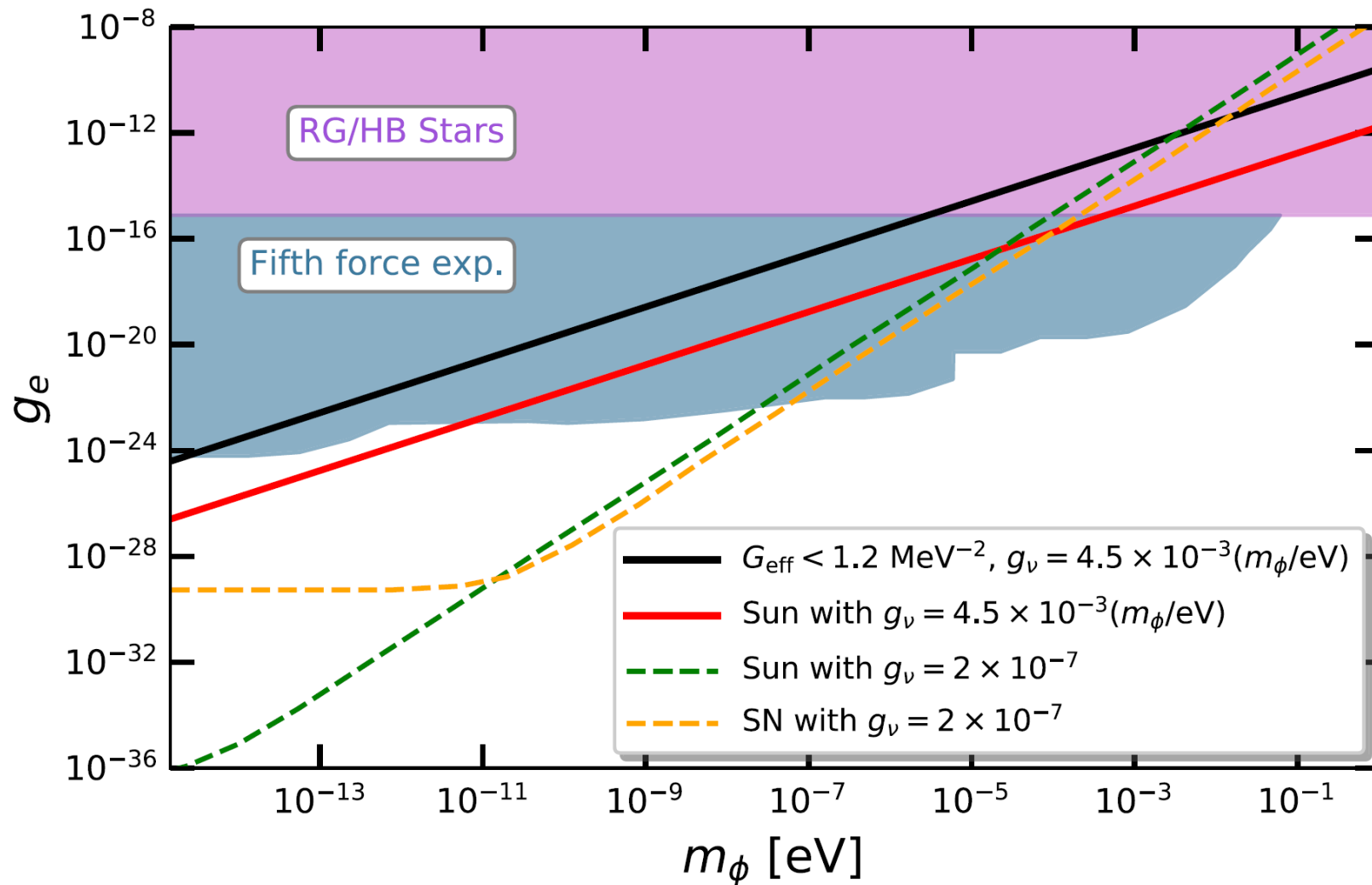
J. Venzor, APL, J. de Santiago, PRD 103 (2020)



$$G_{\text{eff}} < 1.2 \text{ MeV}^{-2} \quad (68\% \text{ C.L.})$$

$$g_\nu < 4.5 \times 10^{-3} \left(\frac{m_\phi}{\text{eV}} \right) \quad (68\% \text{ C.L.}).$$

BBN bounds on light scalar NSI



$$G_S < 2.0 \times 10^7 \text{ MeV}^{-2} \quad (68\% \text{ C.L.})$$

$$g_e < 2.7 \times 10^{-10} \left(\frac{m_\phi}{\text{eV}} \right)$$



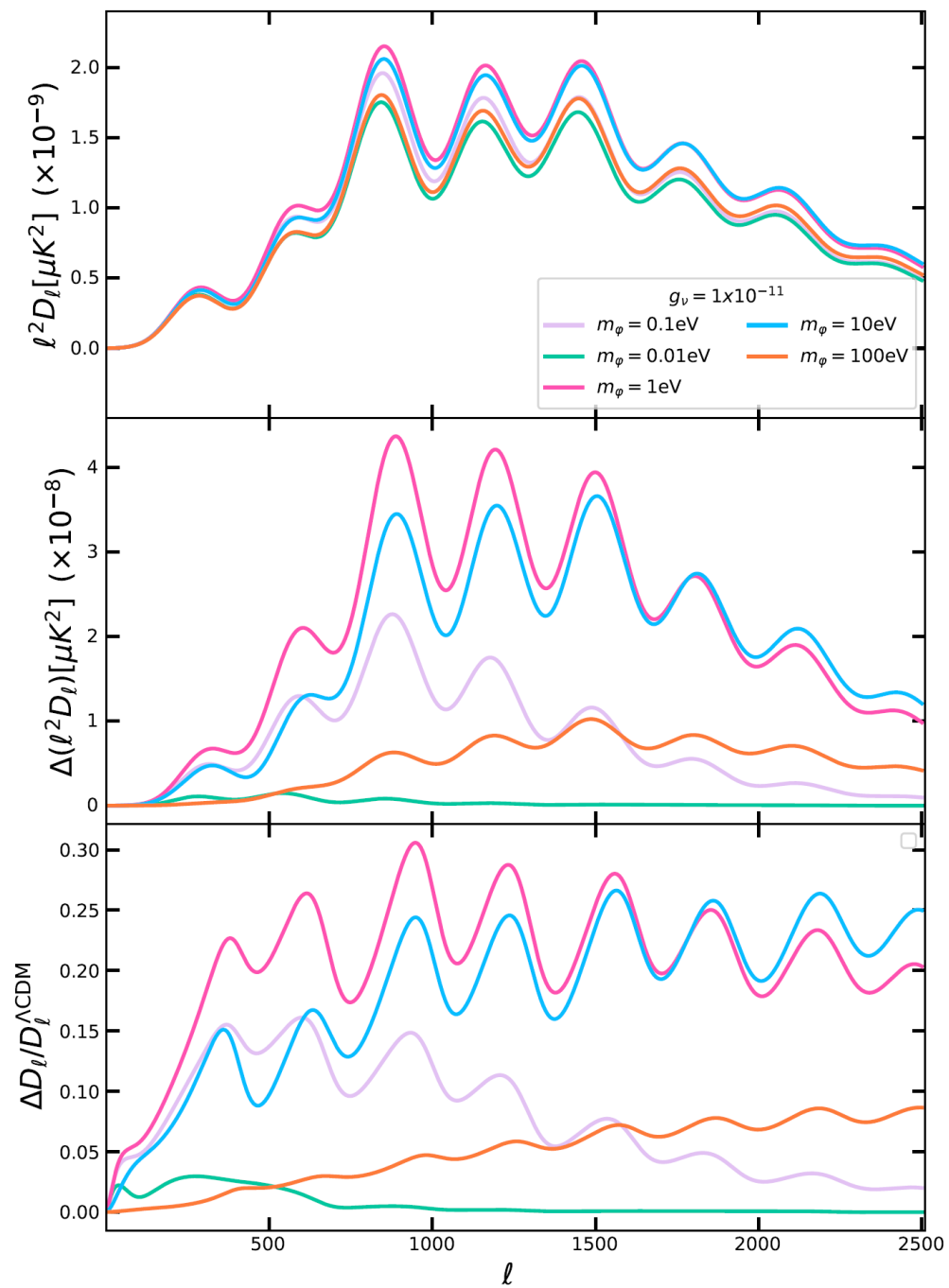
- Neutrinos play an important role along thermal history of the Universe.
- Cosmological data (CMB, Matter spectrum, BBN) are sensible to neutrino physics through

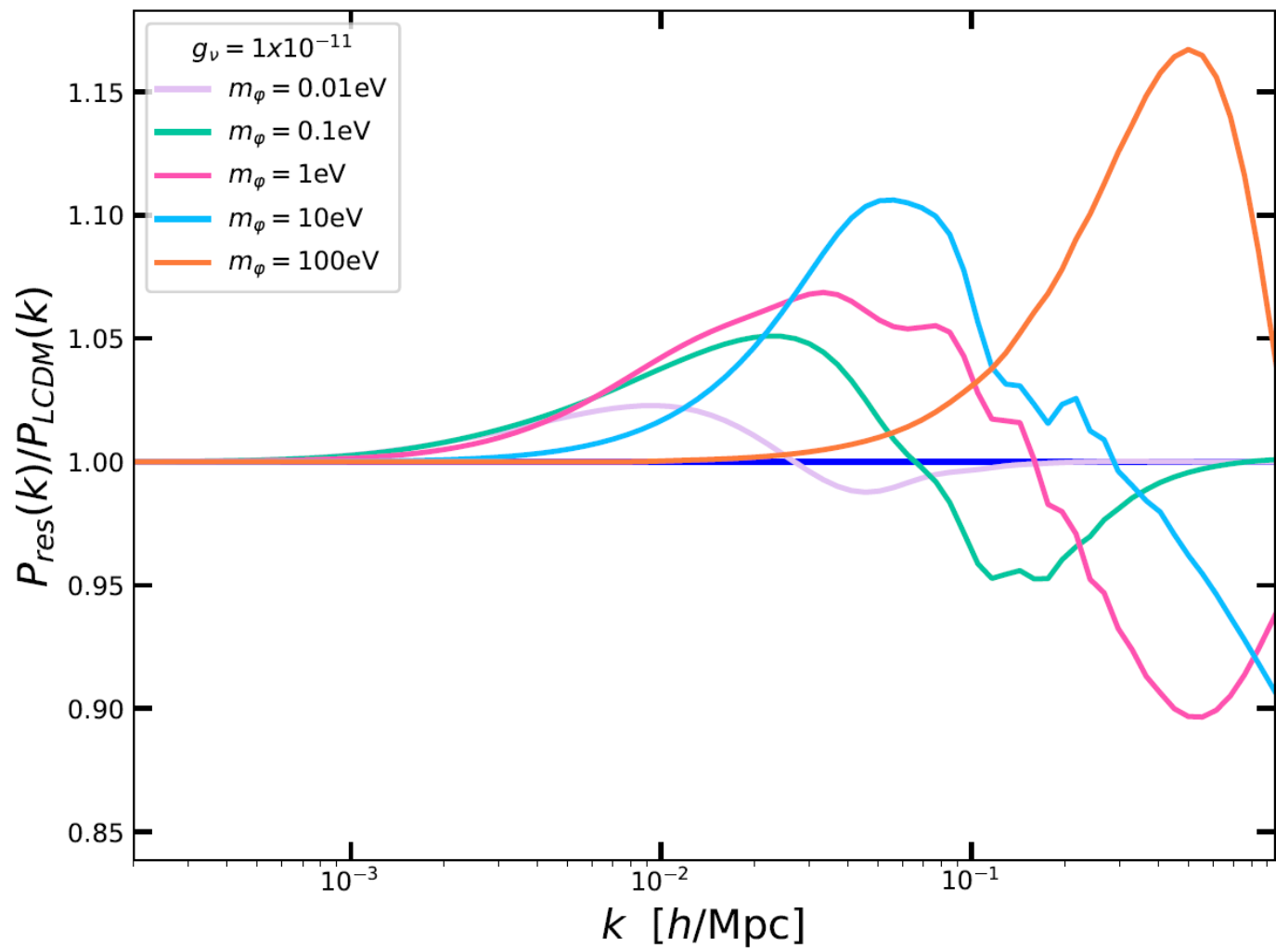
$$T_{dec}, \quad \sum m_\nu, \quad N_{eff}, \quad (k_{FS})$$

- NSI may alter such parameters, providing a way to explore for observational bounds to the new couplings
- We have explored NSI effects on CMB-BAO-H0 data and BBN for light scalar mediators, and around resonance

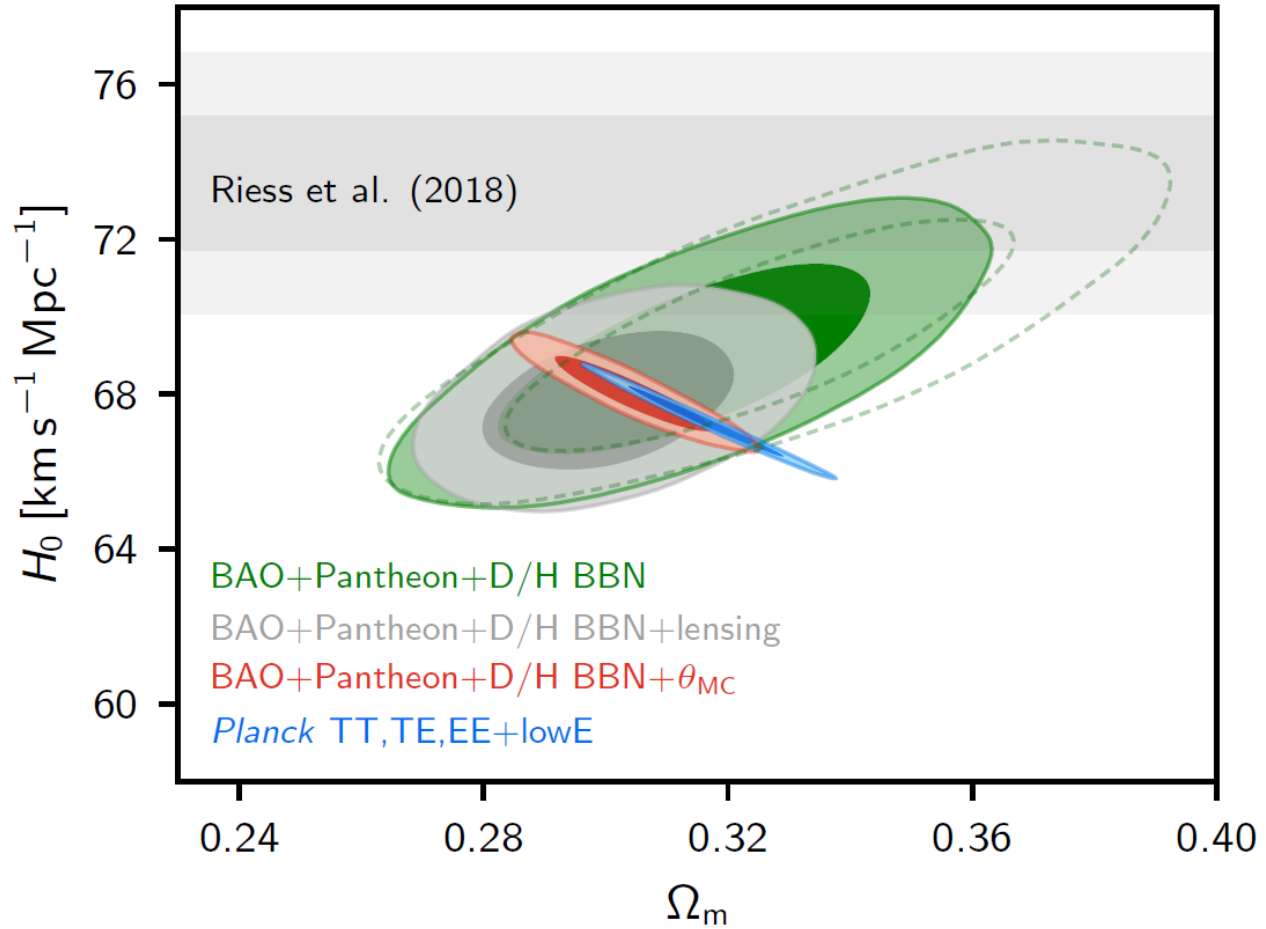
More work underway...

TT-PS





H_0 Tension



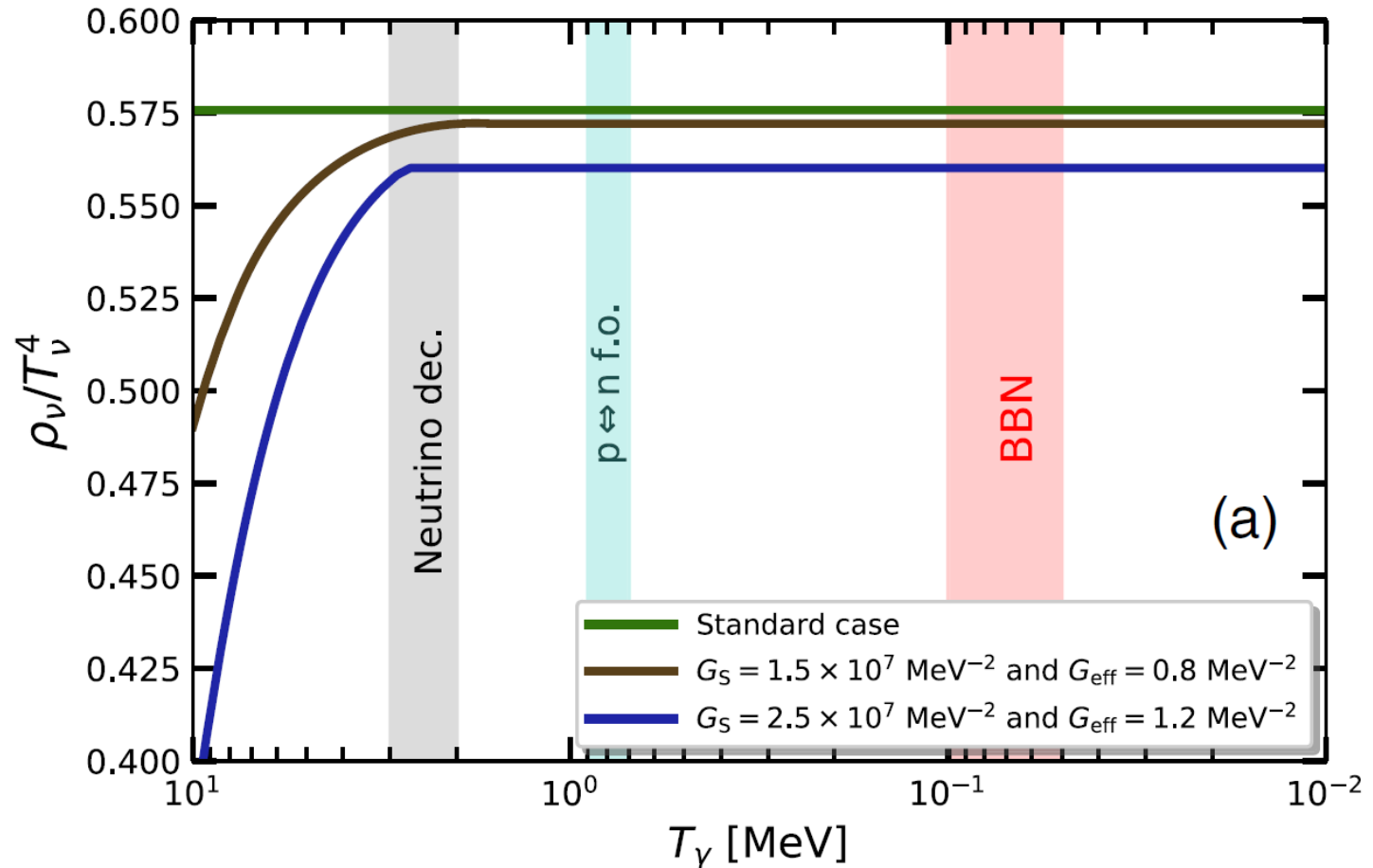
Resonant NSI interaction limits



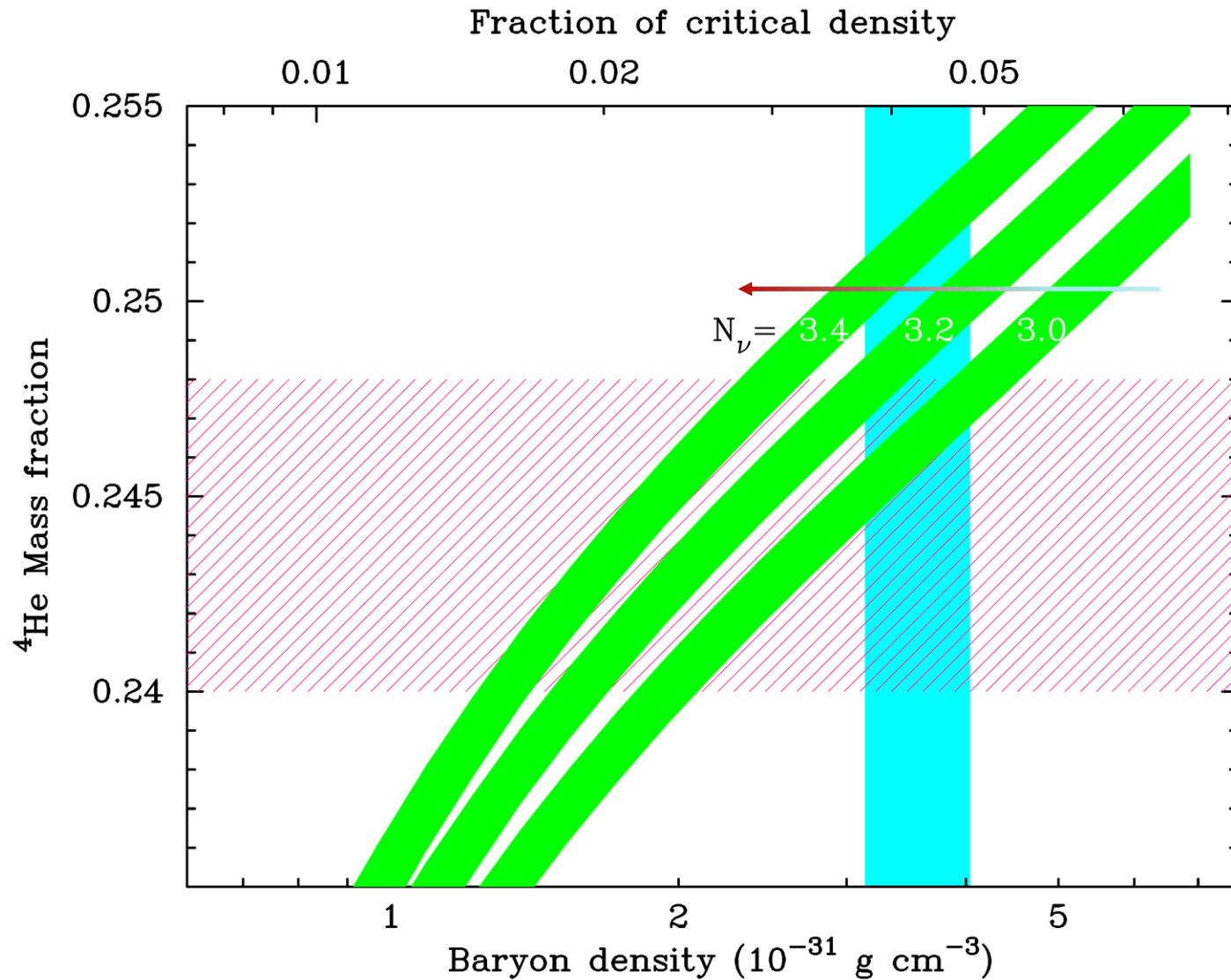
Data: Planck +	BAO	BAO+H0	BAO	BAO+H0	BAO	BAO+H0	BAO	BAO+H0	BAO	BAO+H0
M_ϕ [eV]	10^{-2}		10^{-1}		1		10		100	
H_0 [km s ⁻¹ /Mpc]	$68.1^{+2.8}_{-2.8}$	$71.1^{+1.8}_{-1.8}$	$68.1^{+2.7}_{-2.7}$	$71.2^{+1.6}_{-1.7}$	$67.9^{+2.5}_{-2.3}$	$71.2^{+1.6}_{-1.6}$	$67.7^{+2.6}_{-2.6}$	$71.1^{+1.6}_{-1.6}$	$67.5^{+2.6}_{-2.7}$	$71.0^{+1.7}_{-1.7}$
N_{eff}	$3.12^{+0.47}_{-0.44}$	$3.58^{+0.31}_{-0.29}$	$3.12^{+0.43}_{-0.43}$	$3.59^{+0.31}_{-0.31}$	$3.05^{+0.42}_{-0.39}$	$3.56^{+0.30}_{-0.27}$	$3.03^{+0.43}_{-0.39}$	$3.54^{+0.30}_{-0.30}$	$3.02^{+0.42}_{-0.41}$	$3.54^{+0.29}_{-0.28}$
$\sum m_\nu$ [eV]	< 0.110	< 0.119	< 0.114	< 0.123	< 0.112	< 0.110	< 0.120	< 0.117	< 0.110	< 0.0915
$g_\nu \times 10^{14}$	< 861	< 909	< 91.7	< 106	< 15.0	$12.4^{+7.8}_{-8.9}$	< 36.7	27^{+20}_{-20}	< 147	< 207

TABLE I. Observational limits at 95% confidence for different models with varying N_{eff} and m_ν . We see that for 1 and 10 M_ϕ the interaction is non zero at more than 2σ .

NSI effective mass effects on neutrino density



N_{eff} vs BBN



Burles, Nollett & Turner 1999