



# Lorentz violation in electromagnetic moments of fermions

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# Field theories

Two main ingredients:

Dynamic variables



Symmetries



SM fields

- EW gauge bosons
- QCD gauge bosons
- Dirac spinors
- Higgs scalar

SM symmetries

- Gauge symmetry
- Lorentz symmetry
- $\Rightarrow$  CPT symmetry
- Other global symmetries

# Beyond the SM

Electrodynamics:  $\mathcal{L}_{U(1)_e} = \frac{1}{4}(F_{\mu\nu})^2 + \bar{\psi}_i(i\not{D} - m_j)\psi_j + \text{electroweak stuff}$



EW symmetry breaking

Standard Model:

$$\begin{aligned}\mathcal{L}_{\text{SM}} = & -\frac{1}{4}(F_{\mu\nu})^2 + i\bar{\psi}\not{D}\psi \\ & -\bar{\psi}_i Y_{ij}\psi_j + \text{H.c.} \\ & + |D_\mu\phi|^2 - V(\phi)\end{aligned}$$

Not perfect, but  
our best theory  
so far!!!

+ LV stuff, perhaps?



- Fundamental description:
- Planck Scale
  - Dynamic variables and symmetries
  - Lorentz violation?  $\Rightarrow$  Deviations from SM observables

# Observer VS particle

The gravitational linear potential:\*

$$L = \frac{m\dot{\mathbf{r}}^2}{2} - m\mathbf{g} \cdot \mathbf{r}$$

Observer LT

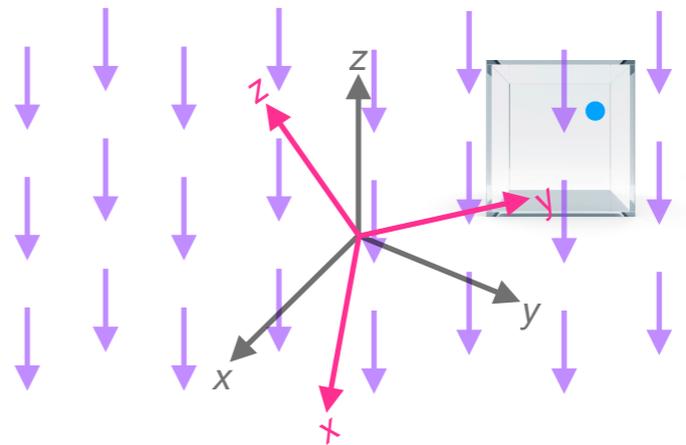
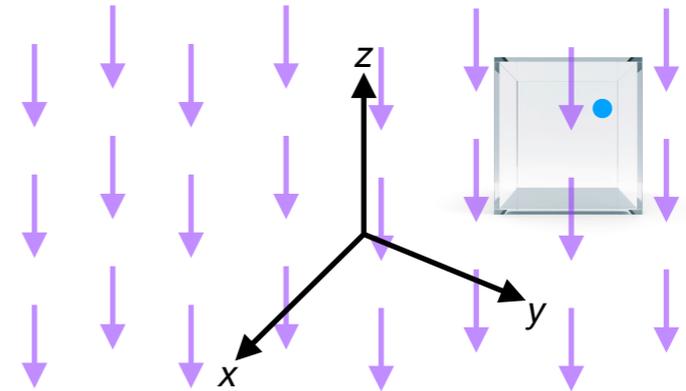
$$L \xrightarrow{\text{obs.}} L' = \frac{m\dot{\mathbf{r}}^2}{2} - m\mathbf{g} \cdot \mathbf{r} = L$$

Invariant

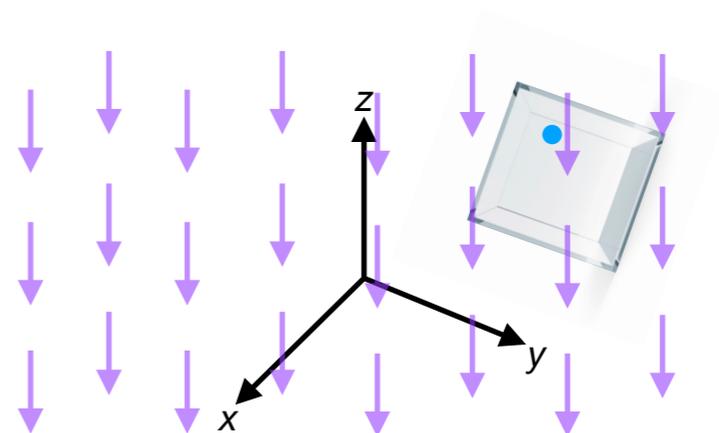
Particle LT

$$L \xrightarrow{\text{part.}} L' = \frac{m\dot{\mathbf{r}}^2}{2} - m g_k (R_{kj} r^j) \neq L$$

Non invariant



New frame



Same frame

\*Jay D. Tasson, What do we know about Lorentz invariance?, *Rep. Prog. Phys.* **77** 062901 (2014).

# The Lorentz-violating SME

Effective Lagrangian: SM symmetries  
 SM dynamic variables  
 Effects from high energies  $\Rightarrow$  Novel phenomena  
 Deviations in SM observables

$$\mathcal{L}_{\text{SME}} = \mathcal{L}_{\text{mSME}} + \mathcal{L}_{\text{nrSME}}$$

■ Renormalizable or minimal SME: mass units = 4

■ Non-renormalizable SME: mass units > 4

$$\mathcal{L}_{\text{mSME}} = \mathcal{L}_{\text{SM}} + \mathcal{L}_{\text{LV}}$$

■ The SM Lagrangian

■ Lorentz-violating terms:  $T_{\mu_1 \dots \mu_n} \mathcal{O}^{\mu_1 \dots \mu_n}$

■ Background fields

■ Field content

Transformations: OT:  $T_{\mu_1 \mu_2 \dots \mu_n} \mathcal{O}^{\mu_1 \mu_2 \dots \mu_n} \xrightarrow{\text{obs.}} T_{\mu_1 \mu_2 \dots \mu_n} \mathcal{O}^{\mu_1 \mu_2 \dots \mu_n}$  **Invariant!**

PT:  $T_{\mu_1 \dots \mu_n} \mathcal{O}^{\mu_1 \dots \mu_n} \xrightarrow{\text{part.}} T_{\mu_1 \dots \mu_n} \Lambda^{\mu_1}_{\nu_1} \dots \Lambda^{\mu_n}_{\nu_n} \mathcal{O}^{\nu_1 \dots \nu_n}$  **Non-invariant!**

# The Yukawa sector\*

Leptons

Quarks

Antisymmetric (spacetime)

$$\mathcal{L}_Y^{\text{SME}} = -\frac{1}{2} \underbrace{(H_L)_{\mu\nu}^{AB} \overline{L}_A \phi \sigma^{\mu\nu} R_B}_{\text{Leptons}} - \frac{1}{2} \underbrace{(H_U)_{\mu\nu}^{AB} \overline{Q}_A \tilde{\phi} \sigma^{\mu\nu} U_B}_{\text{Quarks}} - \frac{1}{2} \underbrace{(H_D)_{\mu\nu}^{AB} \overline{Q}_A \phi \sigma^{\mu\nu} D_B}_{\text{Antisymmetric (spacetime)}} + \text{H.c.}$$

SSB + unitary gauge  $\Rightarrow$

$$\mathcal{L}_Y^{\text{SME}} = -\frac{1}{2} (v + H) \sum_{f=l,u,d} \overline{f}_A \left[ (Y_f)_{\mu\nu}^{AB} P_L + (Y_f)_{\mu\nu}^{BA*} P_R \right] \sigma^{\mu\nu} f_B$$

Biunitary transformation:  $(Y_f)_{\mu\nu} = U_L^{f\dagger} (H_f)_{\mu\nu} U_R^f$

Mass eigenspinors

Higgs field,  $H$

\* D. Colladay and V. A. Kostelecky, *Lorentz-violating extension of the standard model*, Phys. Rev. D **58**, 116002 (1998).

$$\mathcal{L}_Y^{\text{SME}} = -\frac{1}{2}(v + H) \sum_{f=l,u,d} \bar{f}_A \left[ (Y_f)^{AB}_{\mu\nu} P_L + (Y_f)^{BA*}_{\mu\nu} P_R \right] \sigma^{\mu\nu} f_B$$

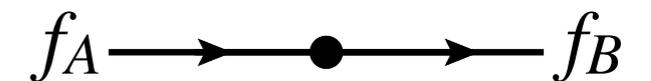
■ Parameters:

$$\begin{aligned} Y_{0i}^{AB} &= e_i^{AB}, \\ Y_{ij}^{AB} &= \epsilon_{ijk} b^{ABk}, \end{aligned} \quad \Rightarrow \quad \mathbf{e}^{AB}, \mathbf{b}^{AB}$$

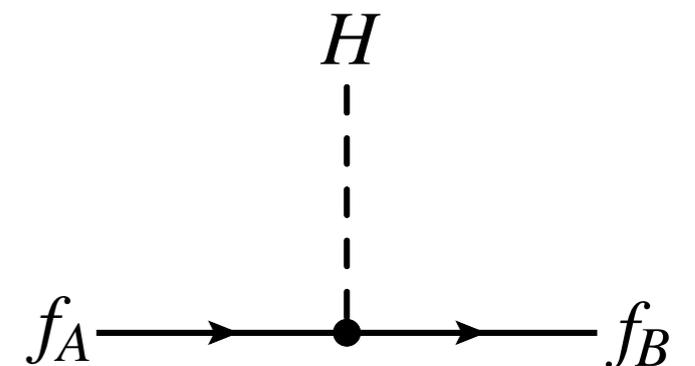
$\mathcal{L}_Y^{\text{SME}} \Rightarrow$

Vertices:

■ Two-point insertion



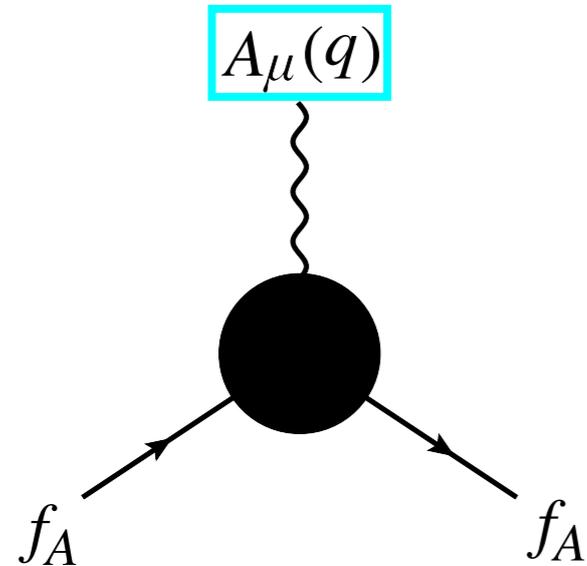
■ Three-point vertex



$\Rightarrow$

**Contributing Feynman diagrams!**

# Contributions to EMMs



$$= \bar{\mathcal{U}}_A \left( \underbrace{f_A^{\text{M}}(q^2)}_{\text{blue}} \sigma_{\mu\nu} q^\nu + \underbrace{f_A^{\text{E}}(q^2)}_{\text{blue}} \sigma_{\mu\nu} q^\nu \gamma_5 \right) \mathcal{U}_A \boxed{+ \dots}$$

■ Off-shell photon

■ Lorentz invariant contributions

■ Several more terms with Lorentz violation

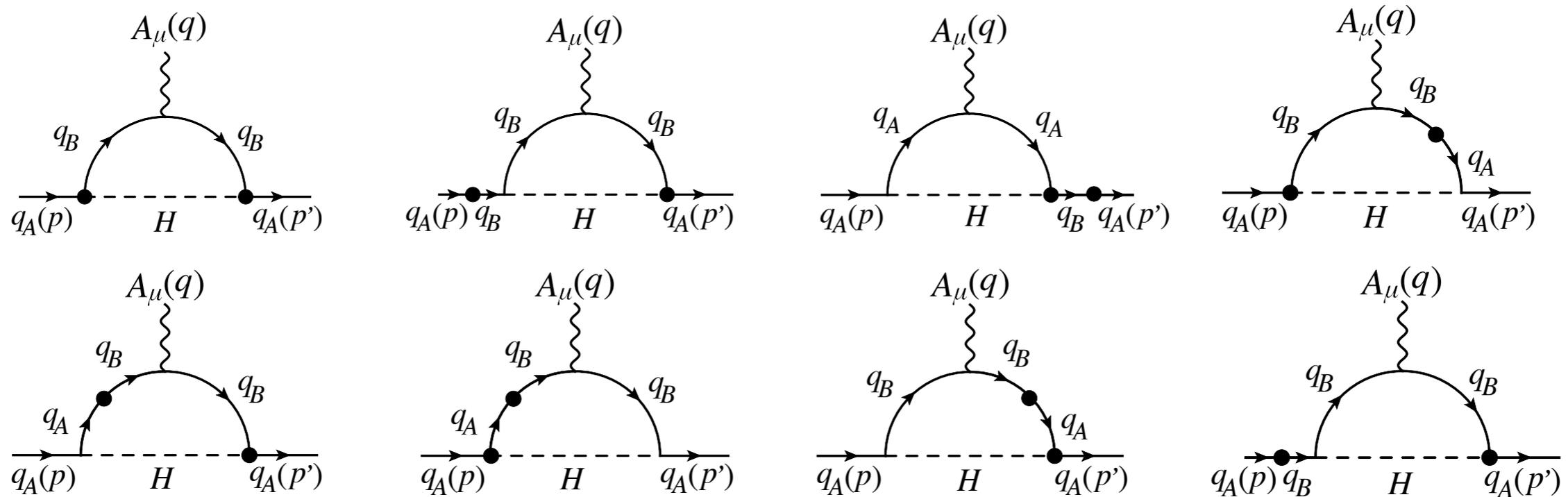
■ Electromagnetic moments:

$$\text{AMM} \quad a_A^{\text{SME}} = f_A^{\text{M}}(q^2 = 0)$$

$$\text{EDM} \quad d_A^{\text{SME}} = f_A^{\text{E}}(q^2 = 0)$$

# Contributing diagrams

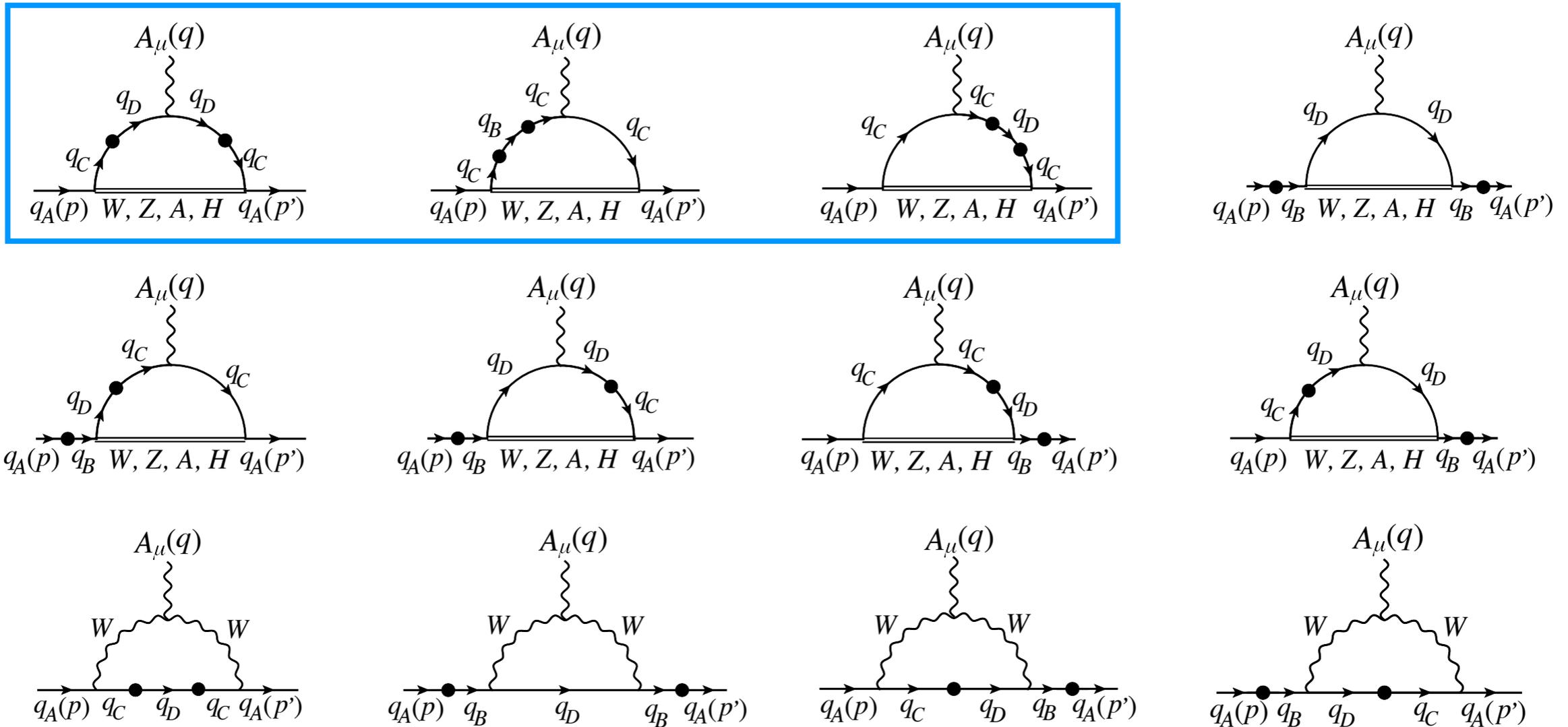
Three-point vertices and two-point insertions



Always a virtual Higgs-boson line

# Contributing diagrams

Only two-point insertions



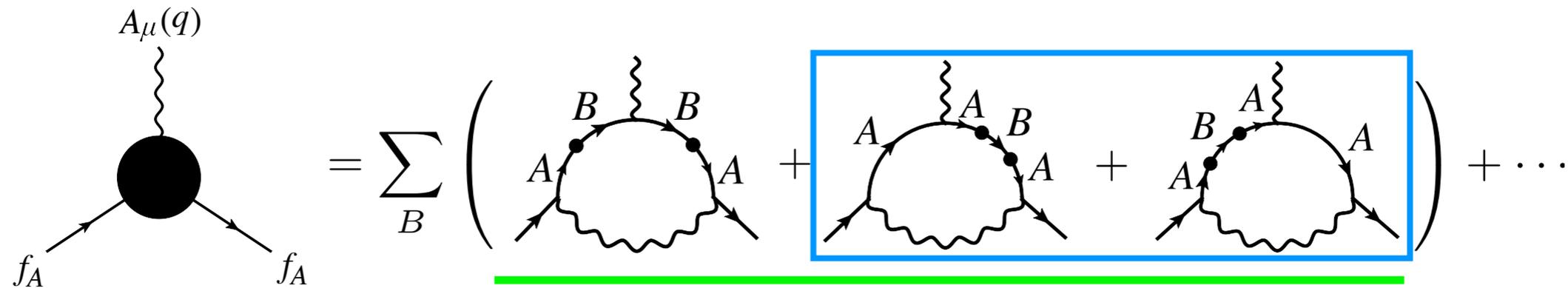
■ Dominant contributions from virtual-photon diagrams

# The calculation

The diagram shows a vertex where an incoming fermion  $f_A$  and an outgoing fermion  $f_A$  meet at a black circle. A wavy line labeled  $A_\mu(q)$  is attached to the top of the circle. This is equated to a sum over particles  $B$  of three diagrams in parentheses, followed by an ellipsis. Each diagram shows a fermion loop with an internal fermion line and a wavy photon line. The photon line couples to the incoming fermion, the outgoing fermion, or the internal fermion line.

- Full contribution:
  - Diagonal electromagnetic form factors
  - Unitary gauge  $\longrightarrow$  No pseudo-Goldstones
  - Passarino-Veltman tensor-reduction method
  - Mathematica, FeynCalc, and Package-X
  - Ward identity
  - AMMs  $a_A^{\text{SME}}$  and EDMs  $d_A^{\text{SME}}$
  - UV-finite electromagnetic moments

# The calculation



■ Dominant contributions

■ IR divergences remain  $\longrightarrow$  To be canceled from cross section

The contributions:

$$a_A^{\text{SME}} = \sum_B \left[ \tilde{a}_{AB} (|\text{Re } \mathbf{e}^{AB}|^2 + |\text{Re } \mathbf{b}^{AB}|^2) + \hat{a}_{AB} (|\text{Im } \mathbf{e}^{AB}|^2 + |\text{Im } \mathbf{b}^{AB}|^2) \right]$$

$$d_A^{\text{SME}} = \sum_B \tilde{d}_{AB} (|\text{Re } \mathbf{e}^{AB}| |\text{Im } \mathbf{b}^{AB}| + |\text{Re } \mathbf{b}^{AB}| |\text{Im } \mathbf{e}^{AB}|)$$

# IR divergences

$$a_A^{\text{SME}} = \frac{e^3 v^2}{2\pi^2 m_A^2} \left( \Delta_{\text{IR}} + \log \frac{\mu^2}{m_A^2} \right) (|\text{Re } \mathbf{e}^{AA}|^2 + |\text{Re } \mathbf{b}^{AA}|^2) + \text{IR finite}$$

$$d_A^{\text{SME}} = \frac{e^3 v^2}{4\pi^2 m_A^3} \left( \Delta_{\text{IR}} + \log \frac{\mu^2}{m_A^2} \right) |\text{Im } \mathbf{e}^{AA}| |\text{Re } \mathbf{b}^{AA}| + \text{IR finite}$$

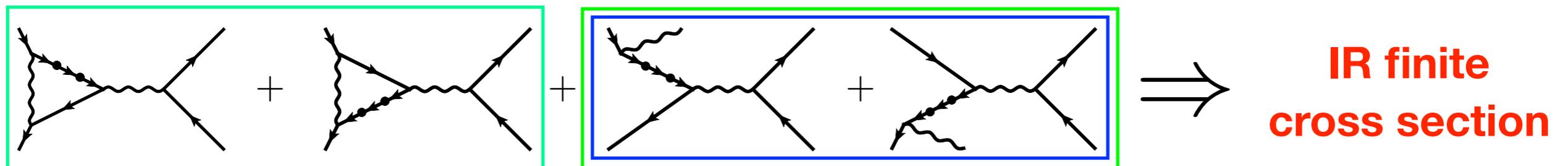
■ IR divergences

■ Contributions are not observables

Soft final-state photon

IR divergent

Forthcoming work: cancelation of IR divergences at cross-section level



■ Soft final-state photon

■ IR divergent

**IR finite  
cross section**

# SM quarks

Proton magnetic moment:

$$a_p = \frac{4}{3} a_u - \frac{1}{3} a_d$$

■ Up quark AMM  
■ Down quark AMM

Current measurement\*:

$$\mu_p = 2.7928473446(8) \mu_N$$

■ Error  $\Rightarrow$  Bounds on SME

Most restrictive bounds  $\sim 10^{-11}$

Access to 2nd and 3rd quark families

LVP	Bounds
$ \text{Re}\{\mathbf{e}^{uu}, \mathbf{b}^{uu}\} $	$6.828 \times 10^{-10}$
$ \text{Im}\{\mathbf{e}^{uu}, \mathbf{b}^{uu}\} $	$1.526 \times 10^{-9}$
$ \text{Re}\{\mathbf{e}^{uc}, \mathbf{b}^{uc}\} $	$1.921 \times 10^{-8}$
$ \text{Im}\{\mathbf{e}^{uc}, \mathbf{b}^{uc}\} $	$1.925 \times 10^{-8}$
$ \text{Re}\{\mathbf{e}^{ut}, \mathbf{b}^{ut}\} $	$7.156 \times 10^{-11}$
$ \text{Im}\{\mathbf{e}^{ut}, \mathbf{b}^{ut}\} $	$7.156 \times 10^{-11}$
$ \text{Re}\{\mathbf{e}^{dd}, \mathbf{b}^{dd}\} $	$6.828 \times 10^{-9}$
$ \text{Im}\{\mathbf{e}^{dd}, \mathbf{b}^{dd}\} $	$1.526 \times 10^{-8}$
$ \text{Re}\{\mathbf{e}^{ds}, \mathbf{b}^{ds}\} $	$3.267 \times 10^{-8}$
$ \text{Im}\{\mathbf{e}^{ds}, \mathbf{b}^{ds}\} $	$3.441 \times 10^{-8}$
$ \text{Re}\{\mathbf{e}^{db}, \mathbf{b}^{db}\} $	$2.194 \times 10^{-7}$
$ \text{Im}\{\mathbf{e}^{db}, \mathbf{b}^{db}\} $	$2.220 \times 10^{-7}$

\*G. Schneider *et al.*, *Double-trap measurement of the proton magnetic moment at 0.3 parts per billion precision*, Science **358**, 1081 (2017).

# SM quarks

## Neutron EDM:

$$d_n = \frac{4}{3} d_d - \frac{1}{3} d_u$$

■ Down quark EDM

■ Up quark EDM

## Current bound\*:

$$|d_n| < 1.8 \times 10^{-26} e \cdot \text{cm}$$

⇒ Bounds on SME

■ Most restrictive bound  $\sim 10^{-12}$

■ Access to 2nd and 3rd quark families

LVP	Bounds
$ e^{uu}, b^{uu} $	$4.308 \times 10^{-12}$
$ e^{uc}, b^{uc} $	$9.401 \times 10^{-11}$
$ e^{ut}, b^{ut} $	$1.096 \times 10^{-9}$
$ e^{dd}, b^{dd} $	$1.703 \times 10^{-11}$
$ e^{ds}, b^{ds} $	$6.449 \times 10^{-11}$
$ e^{db}, b^{db} $	$4.264 \times 10^{-10}$

\*P. J. Mohr, D. B. Newell, and B. N. Taylor, *CODATA recommended values of the fundamental physical constants: 2014*, Rev. Mod. Phys. **88**, 035009 (2016).

# SM charged leptons

## Electron AMM\*

$$\Delta a_e = -1.06(082) \times 10^{-12}$$

## Muon AMM\*\*

$$\Delta a_\mu = 249(87) \times 10^{-11}$$

## Tau AMM $\Rightarrow$ Poor bounds

## Most restrictive bounds $\sim 10^{-11}$

LVP	Bounds
$ \text{Re}\{\mathbf{e}^{ee}, \mathbf{b}^{ee}\} $	$1.619 \times 10^{-11}$
$ \text{Im}\{\mathbf{e}^{ee}, \mathbf{b}^{ee}\} $	$3.620 \times 10^{-11}$
$ \text{Re}\{\mathbf{e}^{e\mu}, \mathbf{b}^{e\mu}\} $	$2.596 \times 10^{-10}$
$ \text{Im}\{\mathbf{e}^{e\mu}, \mathbf{b}^{e\mu}\} $	$2.609 \times 10^{-10}$
$ \text{Re}\{\mathbf{e}^{e\tau}, \mathbf{b}^{e\tau}\} $	$1.067 \times 10^{-9}$
$ \text{Im}\{\mathbf{e}^{e\tau}, \mathbf{b}^{e\tau}\} $	$1.067 \times 10^{-9}$
$ \text{Re}\{\mathbf{e}^{\mu e}, \mathbf{b}^{\mu e}\} $	$1.436 \times 10^{-7}$
$ \text{Im}\{\mathbf{e}^{\mu e}, \mathbf{b}^{\mu e}\} $	$1.451 \times 10^{-7}$
$ \text{Re}\{\mathbf{e}^{\mu\mu}, \mathbf{b}^{\mu\mu}\} $	$1.622 \times 10^{-7}$
$ \text{Im}\{\mathbf{e}^{\mu\mu}, \mathbf{b}^{\mu\mu}\} $	$3.628 \times 10^{-7}$
$ \text{Re}\{\mathbf{e}^{\mu\tau}, \mathbf{b}^{\mu\tau}\} $	$7.243 \times 10^{-7}$
$ \text{Im}\{\mathbf{e}^{\mu\tau}, \mathbf{b}^{\mu\tau}\} $	$7.686 \times 10^{-7}$

\*T. Aoyama, M. Hayakawa, T. Kinoshita, and M. Nio, *Tenth-Order QED Contribution to the Electron  $g-2$  and an Improved Value of the Fine Structure Constant*, Phys. Rev. Lett. **109**, 111807 (2012).

\*\*T. Aoyama, M. Hayakawa, T. Kinoshita, and M. Nio, *Complete Tenth-Order QED Contribution to the Muon  $g-2$* , Phys. Rev. Lett. **109**, 111808 (2012).

# SM charged leptons

## ■ Electron EDM\*

$$|d_e| < 8.7 \times 10^{-29} e \cdot \text{cm}$$

## ■ Muon EDM\*\*

$$|d_\mu| < 1.8 \times 10^{-19} e \cdot \text{cm}$$

■ Tau EDM  $\Rightarrow$  Poor bounds

■ Most restrictive bounds  $\sim 10^{-14}$

LVP	Bounds
$ \mathbf{e}^{ee}, \mathbf{b}^{ee} $	$5.282 \times 10^{-14}$
$ \mathbf{e}^{e\mu}, \mathbf{b}^{e\mu} $	$6.577 \times 10^{-13}$
$ \mathbf{e}^{e\tau}, \mathbf{b}^{e\tau} $	$2.697 \times 10^{-12}$
$ \mathbf{e}^{\mu e}, \mathbf{b}^{\mu e} $	$5.516 \times 10^{-3}$
$ \mathbf{e}^{\mu\mu}, \mathbf{b}^{\mu\mu} $	$7.143 \times 10^{-6}$
$ \mathbf{e}^{\mu\tau}, \mathbf{b}^{\mu\tau} $	$2.530 \times 10^{-5}$

J. Baron *et al.* (ACME Collaboration), *Order of magnitude smaller limit on the electric dipole moment of the electron*, Science **343**, 269 (2014).

G. W. Bennett *et al.*, *Improved limit on the muon electric dipole moment*, Phys. Rev. D **80**, 052008 (2009).

**Thank you!!!**