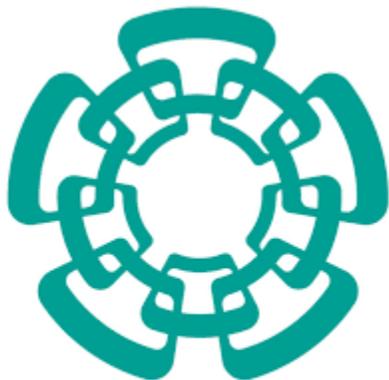
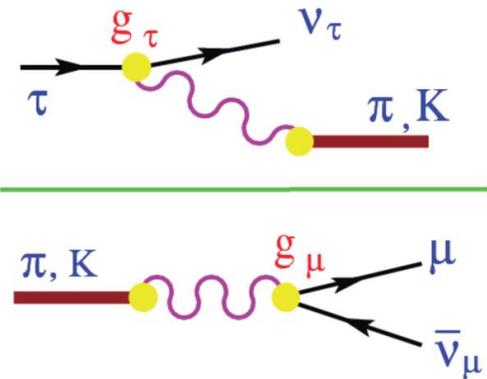


## Radiative corrections in semileptonic tau decays and reliable new physics tests



Cinvestav

Pablo Roig  
Cinvestav (Mexico)



## Radiative corrections in **one-meson** tau decays and reliable new physics tests



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In collaboration with:

M.A. Arroyo-Ureña (BUAP & Cinvestav, Mexico)

G. Hernández-Tomé (IF-UNAM, Mexico)

G. López-Castro (Cinvestav, Mexico)

I. Rosell (UCH-CEU, Valencia, Spain)

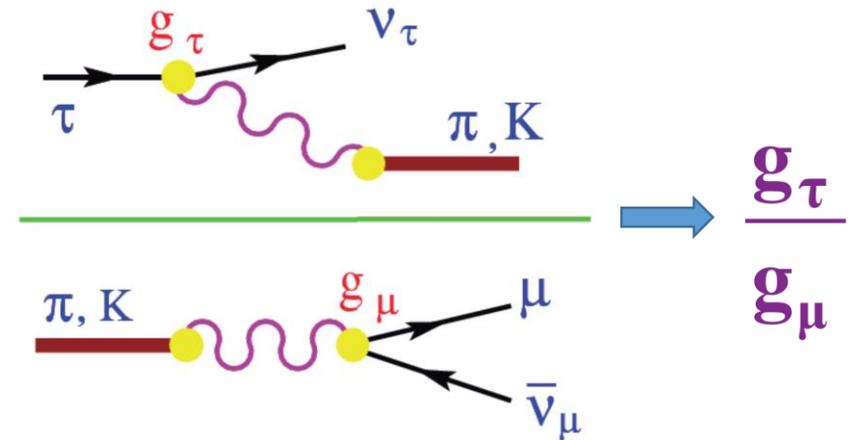
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[JHEP 02 \(2022\) 173 \[arXiv:2112.01859\]](#)

[PRD 104 \(2021\) 9, L091502 \[arXiv:2107.04603\]](#)

# OUTLINE

- 1) Motivation
- 2)  $P \rightarrow \mu \nu_\mu [\gamma]$  ( $P=\pi, K$ )
- 3)  $\tau \rightarrow P \nu_\tau [\gamma]$  ( $P=\pi, K$ )
- 4) Calculation of  $R_{\tau/P} \equiv \frac{\Gamma(\tau \rightarrow P \nu_\tau [\gamma])}{\Gamma(P \rightarrow \mu \nu_\mu [\gamma])}$
- 5) Results
- 6) Applications
- 7) Conclusions



# 1. Motivation

- ✓ **Lepton Universality (LU)** as a basic tenet of the Standard Model (SM).
  - ✓ A few **anomalies** observed in semileptonic B meson decays\*. (See talks by Irina and E. Rojas)
  - ✓ Lower energy observables currently provide the most precise test of LU\*\*.
- ✓ We aim to test **muon-tau lepton universality** through the ratio (P = π, K)\*\*\*:

$$R_{\tau/P} \equiv \frac{\Gamma(\tau \rightarrow P\nu_\tau[\gamma])}{\Gamma(P \rightarrow \mu\nu_\mu[\gamma])} = \left| \frac{g_\tau}{g_\mu} \right|_P^2 R_{\tau/P}^{(0)} (1 + \delta R_{\tau/P})$$

- ✓  $g_\tau = g_\mu$  according to LU.

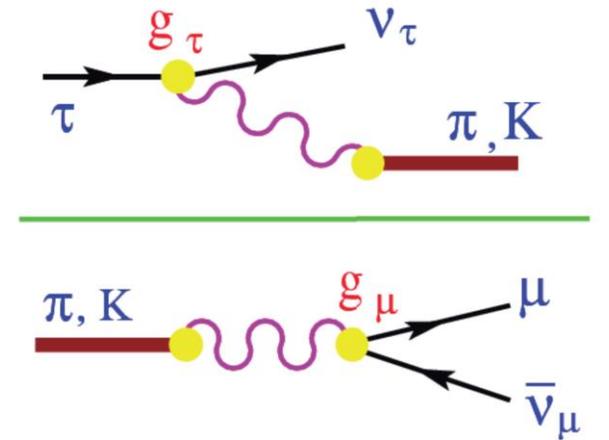
- ✓  $R_{\tau/P}^{(0)}$  is the LO result  $R_{\tau/P}^{(0)} = \frac{1}{2} \frac{M_\tau^3}{m_\mu^2 m_P} \frac{(1 - m_P^2/M_\tau^2)^2}{(1 - m_\mu^2/m_P^2)^2}$ .

- ✓  $\delta R_{\tau/P}$  encodes the **radiative corrections**.

- ✓  $\delta R_{\tau/P}$  was calculated by **Decker & Finkemeier (DF'95)** ^ :

- ✓  $\delta R_{\tau/\pi} = (0.16 \pm 0.14)\%$  and  $\delta R_{\tau/K} = (0.90 \pm 0.22)\%$ .

- ✓ Important **phenomenological and theoretical reasons** to address the analysis again.



\* Albrecht et al.'21  
 \*\* Bryman et al.'21

\*\*\* Marciano & Sirlin'93  
 ^ Decker & Finkemeier'95

# 1. Motivation

✓ Phenomenological disagreement in LU tests:

✓ Using  $\frac{\Gamma(\tau \rightarrow P\nu_\tau[\gamma])}{\Gamma(P \rightarrow \mu\nu_\mu[\gamma])}$  and DF'95\*, HFLAV\*\* reported:

✓  $|g_\tau/g_\mu|_\pi = 0.9958 \pm 0.0026$  (at  $1.6\sigma$  of LU)

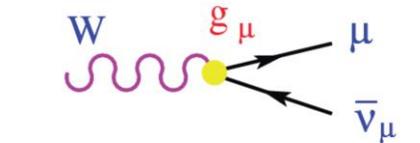
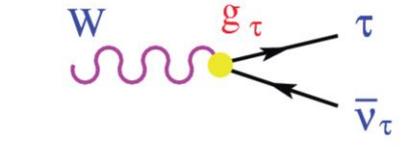
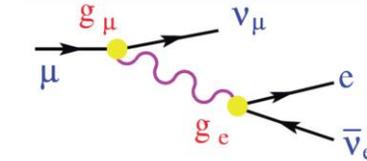
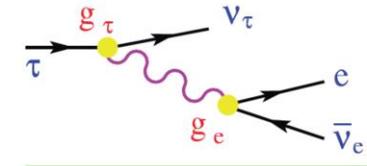
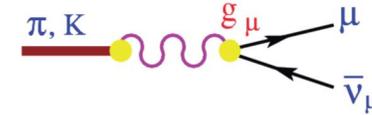
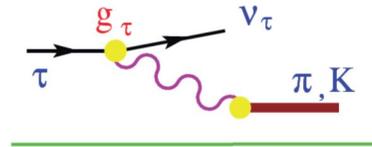
✓  $|g_\tau/g_\mu|_K = 0.9879 \pm 0.0063$  (at  $1.9\sigma$  of LU)

✓ Using  $\frac{\Gamma(\tau \rightarrow e\bar{\nu}_e\nu_\tau[\gamma])}{\Gamma(\mu \rightarrow e\bar{\nu}_e\nu_\mu[\gamma])}$ , HFLAV\*\* reported:

✓  $|g_\tau/g_\mu| = 1.0010 \pm 0.0014$  (at  $0.7\sigma$  of LU)

✓ Using  $\frac{\Gamma(W \rightarrow \tau\nu_\tau)}{\Gamma(W \rightarrow \mu\nu_\mu)}$ , CMS and ATLAS\*\*\* and reported:

✓  $|g_\tau/g_\mu| = 0.995 \pm 0.006$  (at  $0.8\sigma$  of LU)



\* Decker & Finkemeier'95

\*\* HFLAV'21

\*\*\* CMS'21, ATLAS'21

# 1. Motivation

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## ✓ Theoretical issues within DF'95\*:

- ✓ **Hadronic form factors** are different for **real-** and **virtual-photon** corrections, do not satisfy the correct QCD short-distance behavior, violate **unitarity**, **analyticity** and the **chiral limit** at leading non-trivial orders.
- ✓ A **cutoff** to regulate the loop integrals (separating **long-** and **short-distance** corrections)
- ✓ **Unrealistic uncertainties** (purely  $O(e^2p^2)$  ChPT size).

\* Decker & Finkemeier'95

\*\* HFLAV'21

\*\*\* CMS'21, ATLAS'21

# 1. Motivation

## ✓ Phenomenological disagreement in LU tests:

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- ✓ Using  $\frac{\Gamma(W \rightarrow \tau\nu_\tau)}{\Gamma(W \rightarrow \mu\nu_\mu)}$ , CMS and ATLAS\*\*\* and reported:
  - ✓  $|g_\tau/g_\mu| = 0.995 \pm 0.006$  (at  $0.8\sigma$  of LU)

## ✓ By-products of the project:

- ✓ Radiative corrections in  $\Gamma(\tau \rightarrow P\nu_\tau[\gamma])$ .
- ✓ CKM unitarity test via  $\Gamma(\tau \rightarrow K\nu_\tau[\gamma])$  or via the ratio  $\Gamma(\tau \rightarrow K\nu_\tau[\gamma]) / \Gamma(\tau \rightarrow \pi\nu_\tau[\gamma])$ .
- ✓ Constraints on possible non-standard interactions in  $\Gamma(\tau \rightarrow P\nu_\tau[\gamma])^\wedge$ .

\* Decker & Finkemeier'95

\*\* HFLAV'21

\*\*\* CMS'21, ATLAS'21

## ✓ Theoretical issues within DF'95\*:

- ✓ Hadronic form factors are different for real- and virtual-photon corrections, do not satisfy the correct QCD short-distance behavior, violate unitarity, analyticity and the chiral limit at leading non-trivial orders.
- ✓ A cutoff to regulate the loop integrals (separating long- and short-distance corrections)
- ✓ Unrealistic uncertainties (purely  $O(e^2p^2)$  ChPT size).

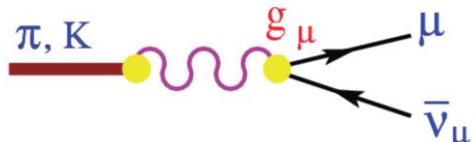
<sup>^</sup> Cirigliano et al.'10 '19, '21

<sup>^</sup> González-Alonso & Martín-Camalich '16

<sup>^</sup> González-Solis et al. '20

## 2. $P \rightarrow \mu \nu_\mu [\gamma]$ ( $P=\pi, K$ )

- ✓ Calculated unambiguously within the **Standard Model (Chiral Perturbation Theory, ChPT\*)**.
- ✓ Notation by **Marciano & Sirlin\*\*** and numbers by **Cirigliano & Rosell\*\*\*** ( $D=d,s$  for  $\pi, K$  and  $F_\pi \approx 92.2$  MeV):



$$\Gamma(P \rightarrow \mu \nu_\mu [\gamma]) = \underbrace{\frac{G_F^2 |V_{uD}|^2 F_P^2}{4\pi} m_P m_\mu^2 \left(1 - \frac{m_\mu^2}{m_P^2}\right)^2}_{\text{LO result}} \underbrace{S_{EW}}_{\substack{\text{short-distance} \\ \text{EW correction} \\ \approx 1.0232^{**}}} \underbrace{\left\{ 1 + \frac{\alpha}{\pi} F(m_\mu^2/m_P^2) \right\}}_{\substack{\text{structure independent (SI)} \\ \text{contributions (point-like} \\ \text{approximation)}^\wedge}} \times \\
 \left\{ 1 - \frac{\alpha}{\pi} \left[ \frac{3}{2} \log \frac{m_\rho}{m_P} + c_1^{(P)} + \frac{m_\mu^2}{m_\rho^2} \left( c_2^{(P)} \log \frac{m_\rho^2}{m_\mu^2} + c_3^{(P)} + c_4^{(P)} (m_\mu/m_P) \right) - \frac{m_P^2}{m_\rho^2} \tilde{c}_2^{(P)} \log \frac{m_\rho^2}{m_\mu^2} \right] \right\}$$

↑ structure-dependent (SD) contributions [coefficients reported in Cirigliano & IR'07]

- ✓ The only **model-dependence** is the determination of the **counterterms** in  $c_1^{(P)}$  and  $c_3^{(P)}$ :
  - ✓ **Large- $N_c$  expansion of QCD**: ChPT is enlarged by including the lightest multiplets of spin-one **resonances** such that the relevant Green functions are **well-behaved at high energies**<sup>†</sup>.

\* Weinberg'79

\* Gasser & Leutwyler'84 '85

\*\* Marciano & Sirlin'93

\*\*\* Cirigliano & IR'07

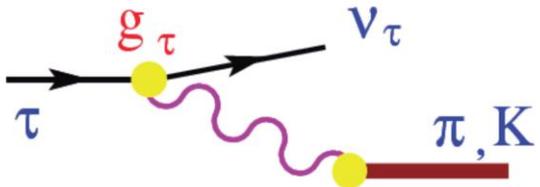
^ Kinoshita'59

† Ecker et al.'89

† Cirigliano et al.'06

### 3. $\tau \rightarrow P \nu_\tau [\gamma]$ ( $P=\pi, K$ )

- ✓ Calculated within an effective approach encoding the hadronization:
- ✓ Large- $N_c$  expansion of QCD: ChPT is enlarged by including the lightest multiplets of spin-one resonances such that the relevant Green functions are well-behaved at high energies\*.
- ✓ We follow a similar notation to  $P \rightarrow \mu \nu_\mu [\gamma]$  ( $D=d,s$  for  $\pi, K$  and  $F_\pi \approx 92.2$  MeV):



$$\Gamma(\tau \rightarrow P \nu_\tau [\gamma]) = \underbrace{\frac{G_F^2 |V_{uD}|^2 F_P^2}{8\pi} M_\tau^3 \left(1 - \frac{m_P^2}{M_\tau^2}\right)^2}_{\text{LO result}} \underbrace{S_{EW}}_{\substack{\text{short-distance} \\ \text{EW correction} \\ \approx 1.02321^{**}}} \underbrace{\left\{ 1 + \frac{\alpha}{\pi} G(m_P^2/M_\tau^2) \right\}}_{\text{structure independent (SI) contributions (point-like approximation)***}} \times \\
 \left\{ 1 - \frac{3\alpha}{2\pi} \log \frac{m_\rho}{M_\tau} + \delta_{\tau P}|_{rSD} + \delta_{\tau P}|_{vSD} \right\}$$

↑
↑

real-photon structure-dependent (rSD) contributions
virtual-photon structure-dependent (vSD) contributions

- ✓ Real-photon structure-dependent (rSD) contributions from Guo & Roig'10<sup>^</sup>.
- ✓ Virtual-photon structure-dependent (vSD) contributions not calculated in the literature.

\* Ecker et al.'89  
 \* Cirigliano et al.'06  
 \*\* Erler'02

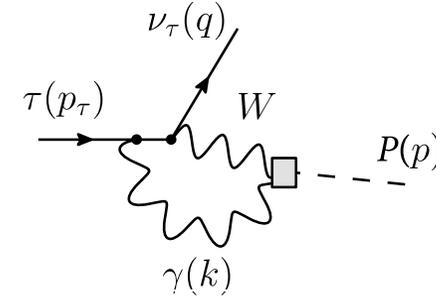
\*\*\* Kinoshita'59  
 ^ Guo & Roig'10

### 3. $\tau \rightarrow P \nu_\tau [\gamma]$ ( $P=\pi, K$ )

- ✓ Virtual-photon structure-dependent contribution (vSD):

$$i\mathcal{M}[\tau \rightarrow P \nu_\tau]_{\text{SD}} = G_F V_{uD} e^2 \int \frac{d^d k}{(2\pi)^d} \frac{\ell^{\mu\nu}}{k^2 [(p_\tau + k)^2 - M_\tau^2]} [i\epsilon_{\mu\nu\lambda\rho} k^\lambda p^\rho F_V^P(W^2, k^2) + F_A^P(W^2, k^2) \lambda_{1\mu\nu} + 2B(k^2) \lambda_{2\mu\nu}]$$

$$\begin{aligned} \ell^{\mu\nu} &= \bar{u}(q) \gamma^\mu (1 - \gamma_5) [(p_\tau + k) + M_\tau] \gamma^\nu u(p_\tau) \\ \lambda_{1\mu\nu} &= [(p + k)^2 + k^2 - m_P^2] g_{\mu\nu} - 2k_\mu p_\nu \\ \lambda_{2\mu\nu} &= k^2 g_{\mu\nu} - \frac{k^2 (p + k)_\mu p_\nu}{(p + k)^2 - m_P^2} \end{aligned}$$



- ✓ Form factors from Guo & Roig'10 and Guevara et al.'13,'21\*:

$$F_V^P(W^2, k^2) = \frac{-N_C M_V^4}{24\pi^2 F_P (k^2 - M_V^2)(W^2 - M_V^2)}$$

$$F_A^P(W^2, k^2) = \frac{F_P}{2} \frac{M_A^2 - 2M_V^2 - k^2}{(M_V^2 - k^2)(M_A^2 - W^2)}$$

$$B(k^2) = \frac{F_P}{M_V^2 - k^2}$$

- ✓ Well-behaved two- and three-point Green functions.

- ✓ Chiral and U(3) limits.

- ✓  $M_V$  and  $M_A$  vector- and axial-vector resonance mass:  $M_V=M_\rho$  and  $M_A=M_{a_1}$  ( $\pi$  case);  $M_V=M_{K^*}$  and  $M_A \approx M_{f_1}$  ( $K$  case).

\* Guo & Roig'10

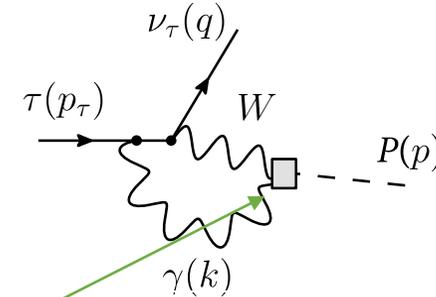
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$$\begin{aligned} \ell^{\mu\nu} &= \bar{u}(q) \gamma^\mu (1 - \gamma_5) [(p_\tau + k) + M_\tau] \gamma^\nu u(p_\tau) \\ \lambda_{1\mu\nu} &= [(p + k)^2 + k^2 - m_P^2] g_{\mu\nu} - 2k_\mu p_\nu \\ \lambda_{2\mu\nu} &= k^2 g_{\mu\nu} - \frac{k^2 (p + k)_\mu p_\nu}{(p + k)^2 - m_P^2} \end{aligned}$$



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- ✓ Well-behaved two- and three-point Green functions.
- ✓ Chiral and U(3) limits.
- ✓  $M_V$  and  $M_A$  vector- and axial-vector resonance mass:  $M_V=M_\rho$  and  $M_A=M_{a_1}$  ( $\pi$  case);  $M_V=M_{K^*}$  and  $M_A \approx M_{f_1}$  (K case).

\* Guo & Roig'10

\* Guevara et al.'13,'21

## 4. Calculation of $R_{\tau/P} = R_{\tau/P}^{(0)} (1 + \delta R_{\tau/P}) = R_{\tau/P}^{(0)} (1 + \delta_{\tau P} - \delta_{P\mu})$

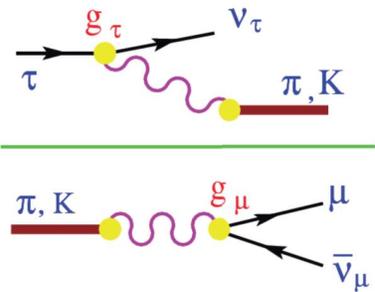
### 1. Structure-independent contribution (point-like approximation): SI.

✓ We confirm the results by DF'95\*.

$$\delta R_{\tau/P}|_{SI} = \frac{\alpha}{2\pi} \left\{ \frac{3}{2} \log \frac{M_\tau^2 m_P^2}{m_\mu^4} + \frac{3}{2} + g \left( \frac{m_P^2}{M_\tau^2} \right) - f \left( \frac{m_\mu^2}{m_P^2} \right) \right\}$$

$$f(x) = 2 \left( \frac{1+x}{1-x} \log x - 2 \right) \log(1-x) - \frac{x(8-5x)}{2(1-x)^2} \log x + 4 \frac{1+x}{1-x} \text{Li}_2(x) - \frac{x}{1-x} \left( \frac{3}{2} + \frac{4}{3} \pi^2 \right)$$

$$g(x) = 2 \left( \frac{1+x}{1-x} \log x - 2 \right) \log(1-x) - \frac{x(2-5x)}{2(1-x)^2} \log x + 4 \frac{1+x}{1-x} \text{Li}_2(x) + \frac{x}{1-x} \left( \frac{3}{2} - \frac{4}{3} \pi^2 \right)$$



$$\delta R_{\tau/\pi}|_{SI} = 1.05\% \text{ and } \delta R_{\tau/K}|_{SI} = 1.67\%$$

Real-photon structure-dependent contribution: rSD.

- ✓  $\delta_{P\mu}|_{rSD}$  from Cirigliano & IR'07\*\*:  $\delta_{\pi\mu}|_{rSD} = -1.3 \cdot 10^{-8}$  and  $\delta_{K\mu}|_{rSD} = -1.7 \cdot 10^{-5}$ .
- ✓  $\delta_{\tau P}|_{rSD}$  from Guo & Roig'10\*\*\*:  $\delta_{\tau\pi}|_{rSD} = 0.15\%$  and  $\delta_{\tau K}|_{rSD} = (0.18 \pm 0.05)\%$ .

$$\delta R_{\tau/\pi}|_{rSD} = 0.15\% \text{ and } \delta R_{\tau/K}|_{rSD} = (0.18 \pm 0.15)\%$$

\* Decker & Finkemeier'95

\*\* Cirigliano & Rosell'07

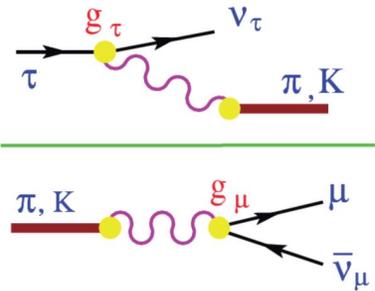
\*\*\* Guo & Roig'10

#### 4. Calculation of $R_{\tau/P} = R_{\tau/P}^{(0)} (1 + \delta R_{\tau/P}) = R_{\tau/P}^{(0)} (1 + \delta_{\tau P} - \delta_{P\mu})$

#### 3. Virtual-photon structure-dependent contribution: vSD.

- ✓  $\delta_{P\mu}|_{\text{vSD}}$  from Cirigliano & IR'07\*:  $\delta_{\pi\mu}|_{\text{vSD}} = (0.54 \pm 0.12)\%$  and  $\delta_{K\mu}|_{\text{vSD}} = (0.43 \pm 0.12)\%$ .
- ✓  $\delta_{\tau P}|_{\text{vSD}}$ , **new calculation**:  $\delta_{\tau\pi}|_{\text{vSD}} = (-0.48 \pm 0.56)\%$  and  $\delta_{\tau K}|_{\text{vSD}} = (-0.45 \pm 0.57)\%$ .

$$\delta R_{\tau/\pi}|_{\text{vSD}} = (-1.02 \pm 0.57)\% \text{ and } \delta R_{\tau/K}|_{\text{vSD}} = (-0.88 \pm 0.58)\%$$



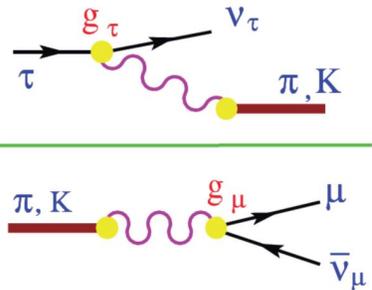
\* Cirigliano & IR'07

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$$\delta R_{\tau/\pi}|_{\text{vSD}} = (-1.02 \pm 0.57)\% \text{ and } \delta R_{\tau/K}|_{\text{vSD}} = (-0.88 \pm 0.58)\%$$



- ✓ **Uncertainties dominated by  $\delta_{\tau P}|_{\text{vSD}}$ :**

- ✓ **P decays** within ChPT [counterterms can be determined by **matching** ChPT with the resonance effective approach at higher energies], whereas **tau decays** within **resonance effective approach** [no **matching** to determine the counterterms].
- ✓ Estimation of the **model-dependence** by comparing our results with a less general scenario where **only well-behaved two-point Green functions** and a **reduced resonance Lagrangian** is used:  $\pm 0.22\%$  and  $\pm 0.24\%$  for the pion and the kaon case.
- ✓ Estimation of the **counterterms** by considering the **running between 0.5 and 1.0 GeV**:  $\pm 0.52\%$  (similar procedure in Marciano & Sirlin'93). **Conservative estimate**, since vSD counterterms affecting in **P decays** imply similar corrections to our estimation of the vSD counterterms in **tau decays**.

\* Cirigliano & Rosell'07

## 5. Results

Contribution	$\delta R_{\tau/\pi}$	$\delta R_{\tau/K}$	Ref.
SI	+1.05%	+1.67%	*
rSD	+0.15%	$+(0.18 \pm 0.05)\%$	**
vSD	$-(1.02 \pm 0.57)\%$	$-(0.88 \pm 0.58)\%$	new
<b>Total</b>	<b><math>+(0.18 \pm 0.57)\%</math></b>	<b><math>+(0.97 \pm 0.58)\%</math></b>	new

Errors are not reported if they are lower than 0.01%.

Central values agree remarkably with DF'95, merely a coincidence:  $\delta R_{\tau/\pi} = (0.16 \pm 0.14)\%$  and  $\delta R_{\tau/K} = (0.90 \pm 0.22)\%$ , but in that work:

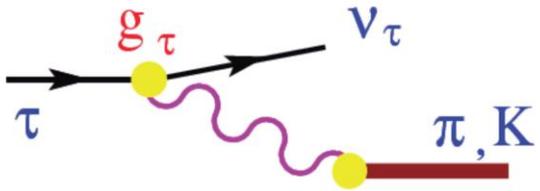
- ✓ **problematic hadronization**: form factors are different for real- and virtual-photon corrections, do not satisfy the correct QCD short-distance behavior, violate unitarity, analyticity and the chiral limit at leading non-trivial orders.
- ✓ a **cutoff** to regulate the loop integrals, splitting unphysically long- and short-distance regimes.
- ✓ **unrealistic uncertainties** (purely  $O(e^2 p^2)$  ChPT size).

\* Decker & Finkemeier'95

\*\* Cirigliano & Rosell'07

\*\* Guo & Roig'10

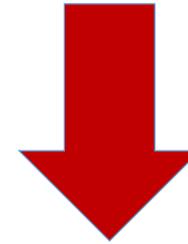
## 6. Application I: Radiative corrections in $\Gamma(\tau \rightarrow P\nu_\tau[\gamma])$



$$\Gamma(\tau \rightarrow P\nu_\tau[\gamma]) = \frac{G_F^2 |V_{uD}|^2 F_P^2}{8\pi} M_\tau^3 \left(1 - \frac{m_P^2}{M_\tau^2}\right)^2 S_{EW} (1 + \delta_{\tau P})$$

short-distance  
EW correction  
 $\approx 1.0201^*$

✓  $\delta_{\tau P}$  includes **SI** and **SD radiative** corrections.

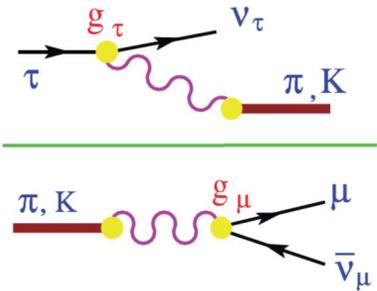


$$\delta_{\tau P} = \frac{\alpha}{2\pi} \left( g \left( \frac{m_P^2}{M_\tau^2} \right) + \frac{19}{4} - \frac{2\pi^2}{3} - 3 \log \frac{m_\rho}{M_\tau} \right) + \delta_{\tau P}|_{rSD} + \delta_{\tau P}|_{vSD} = \begin{cases} \delta_{\tau\pi} = (-0.24 \pm 0.56)\% \\ \delta_{\tau K} = (-0.15 \pm 0.57)\% \end{cases}$$

\* Erler'02

## 6. Application II: lepton universality test

$$R_{\tau/P} \equiv \frac{\Gamma(\tau \rightarrow P\nu_\tau[\gamma])}{\Gamma(P \rightarrow \mu\nu_\mu[\gamma])} = \left| \frac{g_\tau}{g_\mu} \right|_P^2 \frac{1}{2} \frac{M_\tau^3}{m_\mu^2 m_P} \frac{(1 - m_P^2/M_\tau^2)^2}{(1 - m_\mu^2/m_P^2)^2} (1 + \delta R_{\tau/P})$$



$$\left| \frac{g_\tau}{g_\mu} \right|_\pi = 0.9964 \pm 0.0028_{\text{th}} \pm 0.0025_{\text{exp}} = 0.9964 \pm 0.0038$$

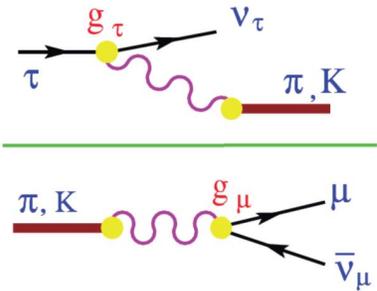
$$\left| \frac{g_\tau}{g_\mu} \right|_K = 0.9857 \pm 0.0028_{\text{th}} \pm 0.0072_{\text{exp}} = 0.9857 \pm 0.0078$$

## 6. Application II: lepton universality test

$$R_{\tau/P} \equiv \frac{\Gamma(\tau \rightarrow P\nu_\tau[\gamma])}{\Gamma(P \rightarrow \mu\nu_\mu[\gamma])} = \left| \frac{g_\tau}{g_\mu} \right|_P^2 \frac{1}{2} \frac{M_\tau^3}{m_\mu^2 m_P} \frac{(1 - m_P^2/M_\tau^2)^2}{(1 - m_\mu^2/m_P^2)^2} (1 + \delta R_{\tau/P})$$

PDG

$$\begin{aligned} \delta R_{\tau/\pi} &= (0.18 \pm 0.57)\% \\ \delta R_{\tau/K} &= (0.97 \pm 0.58)\% \end{aligned}$$



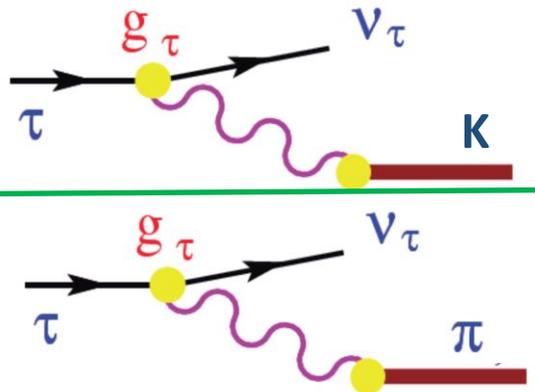
$$\begin{aligned} \left| \frac{g_\tau}{g_\mu} \right|_\pi &= 0.9964 \pm 0.0028_{\text{th}} \pm 0.0025_{\text{exp}} = 0.9964 \pm 0.0038 \\ \left| \frac{g_\tau}{g_\mu} \right|_K &= 0.9857 \pm 0.0028_{\text{th}} \pm 0.0072_{\text{exp}} = 0.9857 \pm 0.0078 \end{aligned}$$

- ✓  $\pi$  case: at  $0.9\sigma$  of LU vs.  $1.6\sigma$  of LU in HFLAV'21\* using DF'95\*\*
- ✓  $K$  case: at  $1.8\sigma$  of LU vs.  $1.9\sigma$  of LU in HFLAV'21\* using DF'95\*\*

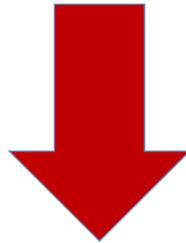
\* HFLAV'21

\*\* Decker & Finkemeier'95

## 6. Application III: CKM unitarity test in the ratio $\Gamma(\tau \rightarrow K\nu_\tau[\gamma]) / \Gamma(\tau \rightarrow \pi\nu_\tau[\gamma])$



$$\frac{\Gamma(\tau \rightarrow K\nu_\tau[\gamma])}{\Gamma(\tau \rightarrow \pi\nu_\tau[\gamma])} = \frac{|V_{us}|^2 F_K^2 (1 - m_K^2/M_\tau^2)^2}{|V_{ud}|^2 F_\pi^2 (1 - m_\pi^2/M_\tau^2)^2} (1 + \delta)$$



$$\left| \frac{V_{us}}{V_{ud}} \right| = 0.2288 \pm 0.0010_{\text{th}} \pm 0.0017_{\text{exp}} = 0.2288 \pm 0.0020$$

## 6. Application III: CKM unitarity test in the ratio $\Gamma(\tau \rightarrow K\nu_\tau[\gamma]) / \Gamma(\tau \rightarrow \pi\nu_\tau[\gamma])$

FLAG'20\*:  
 $F_K/F_\pi = 1.1932 \pm 0.0019$

$\delta = \delta_{\tau K} - \delta_{\tau\pi} = +(0.10 \pm 0.80)\%$

PDG

$$\frac{\Gamma(\tau \rightarrow K\nu_\tau[\gamma])}{\Gamma(\tau \rightarrow \pi\nu_\tau[\gamma])} = \frac{|V_{us}|^2 F_K^2 (1 - m_K^2/M_\tau^2)^2}{|V_{ud}|^2 F_\pi^2 (1 - m_\pi^2/M_\tau^2)^2} (1 + \delta)$$

$$\left| \frac{V_{us}}{V_{ud}} \right| = 0.2288 \pm 0.0010_{\text{th}} \pm 0.0017_{\text{exp}} = 0.2288 \pm 0.0020$$

- ✓ 2.1 $\sigma$  away from CKM unitarity, considering  $|V_{ud}| = 0.97373 \pm 0.00031$ \*\*.
- ✓ To be compared with  $|V_{us}/V_{ud}| = 0.2291 \pm 0.0009$ \*\*\*, obtained with kaon semileptonic decays. Our error does not reach this level due to lack of statistics in  $\tau$  decays.

\* FLAG'20  
 \*\* Hardy & Towner'20  
 \*\*\* Seng et al.'21

## 6. Application III: CKM unitarity test in the ratio $\Gamma(\tau \rightarrow K\nu_\tau[\gamma]) / \Gamma(\tau \rightarrow \pi\nu_\tau[\gamma])$

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PDG

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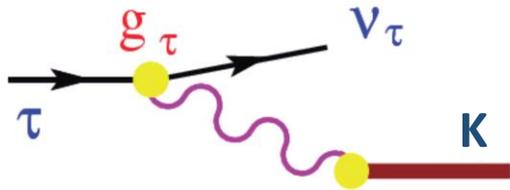
Conservative estimation of the errors in  $\delta$ , since we have directly propagated the uncertainties of  $\delta_{\tau K}$  and  $\delta_{\tau\pi}$ .

$$\left| \frac{V_{us}}{V_{ud}} \right| = 0.2288 \pm 0.0010_{\text{th}} \pm 0.0017_{\text{exp}} = 0.2288 \pm 0.0020$$

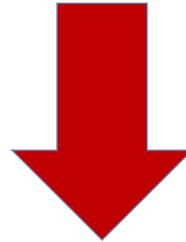
- ✓ 2.1 $\sigma$  away from CKM unitarity, considering  $|V_{ud}| = 0.97373 \pm 0.00031$ \*\*.
- ✓ To be compared with  $|V_{us}/V_{ud}| = 0.2291 \pm 0.0009$ \*\*\*, obtained with kaon semileptonic decays. Our error does not reach this level due to lack of statistics in  $\tau$  decays.

\* FLAG'20  
 \*\* Hardy & Towner'20  
 \*\*\* Seng et al.'21

## 6. Application IV: CKM unitarity test in $\Gamma(\tau \rightarrow K\nu_\tau[\gamma])$

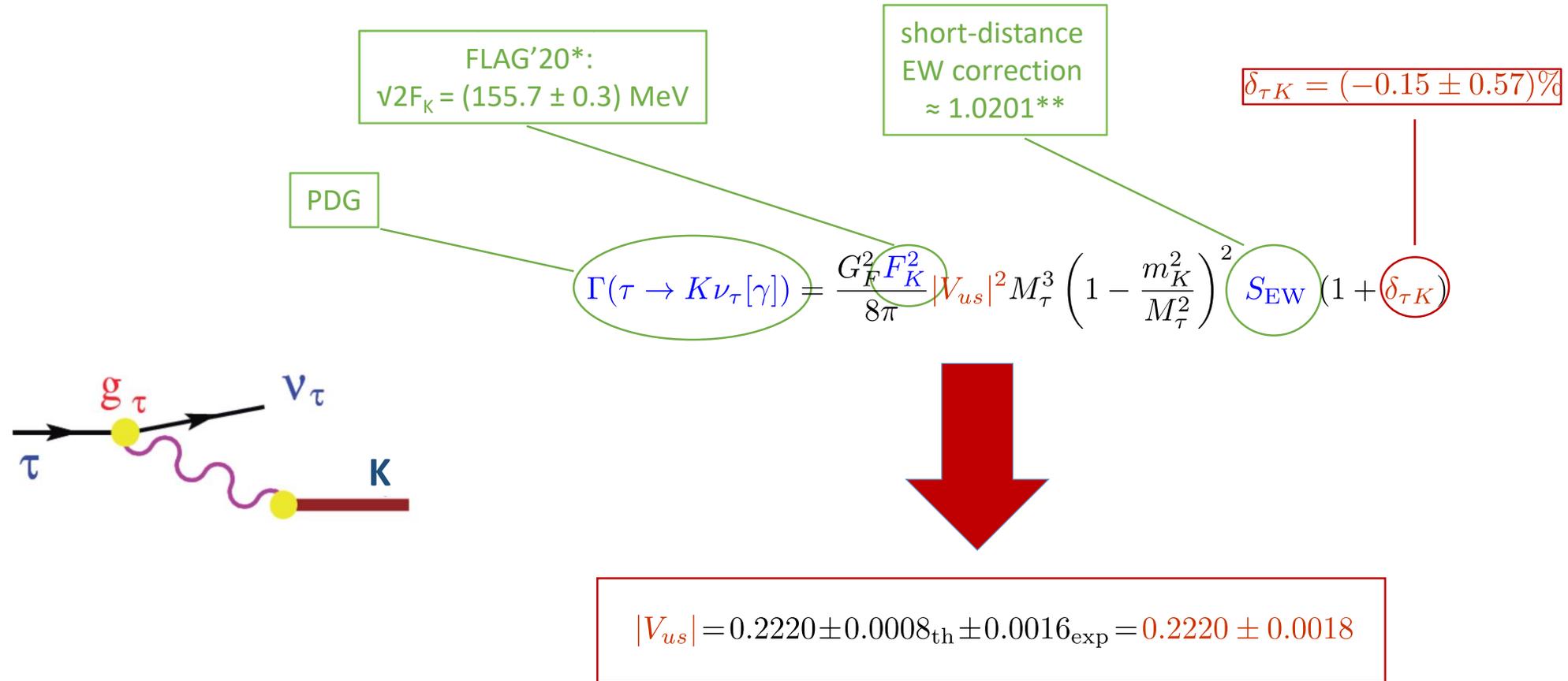


$$\Gamma(\tau \rightarrow K\nu_\tau[\gamma]) = \frac{G_F^2 F_K^2}{8\pi} |V_{us}|^2 M_\tau^3 \left(1 - \frac{m_K^2}{M_\tau^2}\right)^2 S_{EW} (1 + \delta_{\tau K})$$



$$|V_{us}| = 0.2220 \pm 0.0008_{\text{th}} \pm 0.0016_{\text{exp}} = 0.2220 \pm 0.0018$$

## 6. Application IV: CKM unitarity test in $\Gamma(\tau \rightarrow K\nu_\tau[\gamma])$

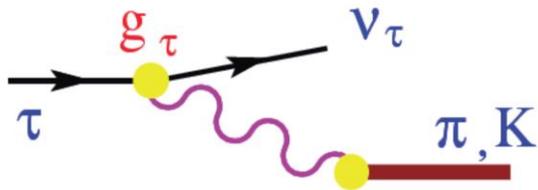


✓ 2.6 $\sigma$  away from CKM unitarity, considering  $|V_{ud}| = 0.97373 \pm 0.00031^{***}$ .

✓ To be compared with  $|V_{us}| = 0.2234 \pm 0.0015^\wedge$  or  $|V_{us}| = 0.2231 \pm 0.0006^\dagger$ , obtained this last one with kaon semileptonic decays. Our error does not reach this level due to lack of statistics in  $\tau$  decays.

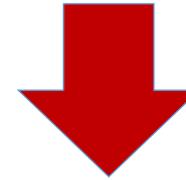
\* FLAG'20  
 \*\* Erler'02  
 \*\*\* Hardy & Towner'20  
 $^\wedge$  HFLAV'21  
 $^\dagger$  Seng et al.'21

## 6. Application V: constraining non-standard interactions in $\Gamma(\tau \rightarrow P\nu_\tau[\gamma])$



$$\Gamma(\tau \rightarrow P\nu_\tau[\gamma]) = \frac{G_F^2 |\tilde{V}_{uD}|^2 F_P^2}{8\pi} M_\tau^3 \left(1 - \frac{m_P^2}{M_\tau^2}\right)^2 S_{EW} (1 + \delta_{\tau P} + 2\Delta^{\tau P})$$

Values of  $\Delta^{\tau P}$  reported in the MS-scheme  
and at a scale of  $\mu=2$  GeV.



$$\Delta^{\tau P} = \epsilon_L^\tau - \epsilon_L^e - \epsilon_R^\tau - \epsilon_R^e - \frac{m_P^2}{M_\tau(m_u + m_D)} \epsilon_P^\tau = \begin{cases} \Delta^{\tau\pi} = -(0.15 \pm 0.72) \cdot 10^{-2} \\ \Delta^{\tau K} = -(0.36 \pm 1.18) \cdot 10^{-2} \end{cases}$$

## 6. Application V: constraining non-standard interactions in $\Gamma(\tau \rightarrow P\nu_\tau[\gamma])$

$|V_{ud}| = 0.97373 \pm 0.00031^*$   
 $|V_{us}/V_{ud}| = 0.2288 \pm 0.0020$

**FLAG'20\*:**  
 $\sqrt{2}F_\pi = (130.2 \pm 0.8) \text{ MeV}$   
 $\sqrt{2}F_K = (155.7 \pm 0.3) \text{ MeV}$

**short-distance  
EW correction**  
 $\approx 1.0201^{**}$

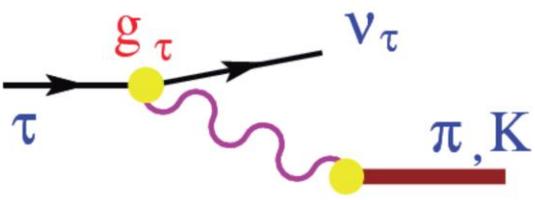
$\delta_{\tau\pi} = (-0.24 \pm 0.56)\%$   
 $\delta_{\tau K} = (-0.15 \pm 0.57)\%$

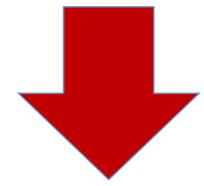
PDG

$\Gamma(\tau \rightarrow P\nu_\tau[\gamma])$

 $= \frac{G_F^2 |\tilde{V}_{uD}|^2 F_P^2}{8\pi} M_\tau^3 \left(1 - \frac{m_P^2}{M_\tau^2}\right)^2 S_{EW} (1 + \delta_{\tau P} + 2\Delta^{\tau P})$



Values of  $\Delta^{\tau P}$  reported in the MS-scheme and at a scale of  $\mu=2 \text{ GeV}$ .



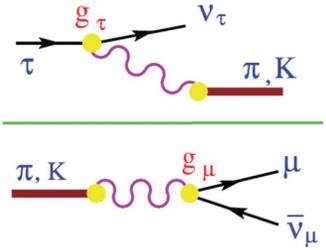
$$\Delta^{\tau P} = \epsilon_L^\tau - \epsilon_L^e - \epsilon_R^\tau - \epsilon_R^e - \frac{m_P^2}{M_\tau(m_u + m_D)} \epsilon_P^\tau = \begin{cases} \Delta^{\tau\pi} = -(0.15 \pm 0.72) \cdot 10^{-2} \\ \Delta^{\tau K} = -(0.36 \pm 1.18) \cdot 10^{-2} \end{cases}$$

- ✓ To be compared with  $\Delta^{\tau\pi} = -(0.15 \pm 0.67) \cdot 10^{-2}$  of Cirigliano et al.'19<sup>^</sup>.
- ✓ To be compared with  $\Delta^{\tau\pi} = -(0.12 \pm 0.68) \cdot 10^{-2}$  and  $\Delta^{\tau K} = (-0.41 \pm 0.93) \cdot 10^{-2}$  of González-Solís et al.'20<sup>†</sup>.

\* Hardy & Towner'20  
 \*\* FLAG'20  
 \*\*\* Erler'02

<sup>^</sup> Cirigliano et al.'19  
<sup>†</sup> González-Solís et al. '20

## 7. Conclusions



- ✓ The **observable** and **our result**:

$$R_{\tau/P} \equiv \frac{\Gamma(\tau \rightarrow P\nu_\tau[\gamma])}{\Gamma(P \rightarrow \mu\nu_\mu[\gamma])} = \left| \frac{g_\tau}{g_\mu} \right|_P^2 R_{\tau/P}^{(0)} (1 + \delta R_{\tau/P}) \quad \longrightarrow \quad \begin{cases} \delta R_{\tau/\pi} = (0.18 \pm 0.57)\% \\ \delta R_{\tau/K} = (0.97 \pm 0.58)\% \end{cases}$$

- ✓ **Framework**: ChPT for  $\pi$  decays and a **resonance extension of ChPT** for  $\tau$  decays.
- ✓ Consistent with **DF'95\***, but with more **robust assumptions** and yielding a **reliable uncertainty**.
- ✓ Applications:
  - ✓ Theoretical determination of **radiative corrections** in  $\Gamma(\tau \rightarrow P\nu_\tau[\gamma])$ .
  - ✓  $|g_\tau/g_\mu|_P$  at **0.9 $\sigma$**  ( $\pi$ ) and **1.8 $\sigma$**  (K) of LU, reducing **HFLAV'21\*\*** disagreement with LU.
  - ✓ **CKM unitarity** in  $\Gamma(\tau \rightarrow K\nu_\tau[\gamma])/\Gamma(\tau \rightarrow \pi\nu_\tau[\gamma])$ :  $|V_{us}/V_{ud}| = 0.2288 \pm 0.0020$ , at **2.1 $\sigma$**  from unitarity.
  - ✓ **CKM unitarity** in  $\Gamma(\tau \rightarrow K\nu_\tau[\gamma])$ :  $|V_{us}| = 0.2220 \pm 0.0018$ , at **2.6 $\sigma$**  from unitarity.
  - ✓ Constraining **non-standard interactions** in  $\Gamma(\tau \rightarrow P\nu_\tau[\gamma])$ : update of  $\Delta^{\tau P}$ .
- ✓ Our results have been **incorporated in the very recent HFLAV'22**.

\* Decker & Finkemeier'95

\*\* HFLAV'21

## 7. Conclusions

# Reliable NP tests for present & future exps.

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$$R_{\tau/P} \equiv \frac{\Gamma(\tau \rightarrow P\nu_\tau[\gamma])}{\Gamma(P \rightarrow \mu\nu_\mu[\gamma])} = \left| \frac{g_\tau}{g_\mu} \right|_P^2 R_{\tau/P}^{(0)} (1 + \delta R_{\tau/P}) \rightarrow \begin{cases} \delta R_{\tau/\pi} = (0.18 \pm 0.57)\% \\ \delta R_{\tau/K} = (0.97 \pm 0.58)\% \end{cases}$$

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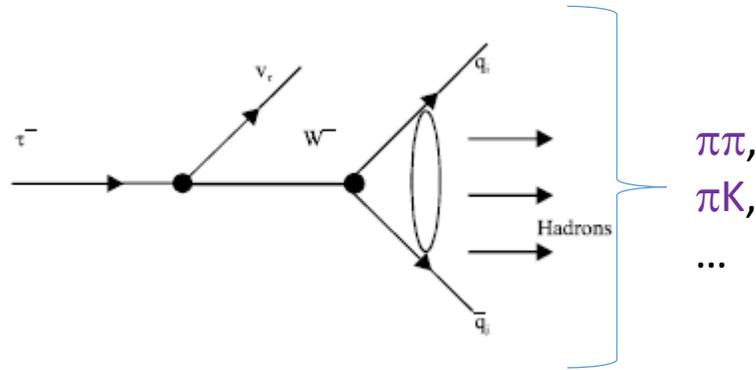
\* Decker & Finkemeier'95

\*\* HFLAV'21

## Comparison with Decker & Finkemeier'95 (DF'95) in the $\pi$ case

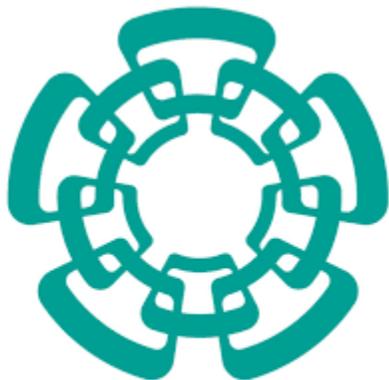
Contribution	$\delta R_{\tau\pi}$ by DF'95 [ $\mu_{\text{cut}} = 1.5 \text{ GeV}$ ]	our $\delta R_{\tau\pi}$
SI	+0.84%*	+1.05%
rSD	+0.05%	+0.15%
vSD	-0.49%*	-(1.02 ± 0.57)%
short-distance	-0.25%*	0
Total	+(0.16 ± 0.14)%*	+(0.18 ± 0.57)%

- ✓ Virtual corrections by DF'95 are  $\mu_{\text{cut}}$ -dependent, since long- and short-distance photonic contributions were separated unphysically: figures with an asterisk are cutoff-dependent.
- ✓ The quoted error in the radiative correction of DF'95 arises from **uncertainties in hadronic parameters** of SD contributions and from **variations in the cutoff parameter**,  $\mu_{\text{cut}}$ .
- ✓ For the SI contribution in DF'95 we have added to the result obtained in the point-like approximation (1.05%) the term coming from cutting off the loops at  $\mu_{\text{cut}}$  (-0.21%).
- ✓ Different contributions of  $\delta R_{\tau/K}$  are not provided in DF'95, which prevents a comparison.
- ✓ Although central values for the sum of all the corrections agree remarkably, this is a coincidence, since central values for the SD corrections are largely different within both approaches.



$\pi\pi,$   
 $\pi K,$   
...

# Radiative corrections in **two-meson** tau decays and reliable new physics tests



Cinvestav

Pablo Roig  
Cinvestav (Mexico)

In collaboration with:

J.A. **Miranda** (Cinvestav, Mexico & IFAE, Barcelona)

and also with R. **Escribano** (IFAE, Barcelona)

[PRD 102 \(2020\), 114117 \[arXiv:2007.11019\]](#)

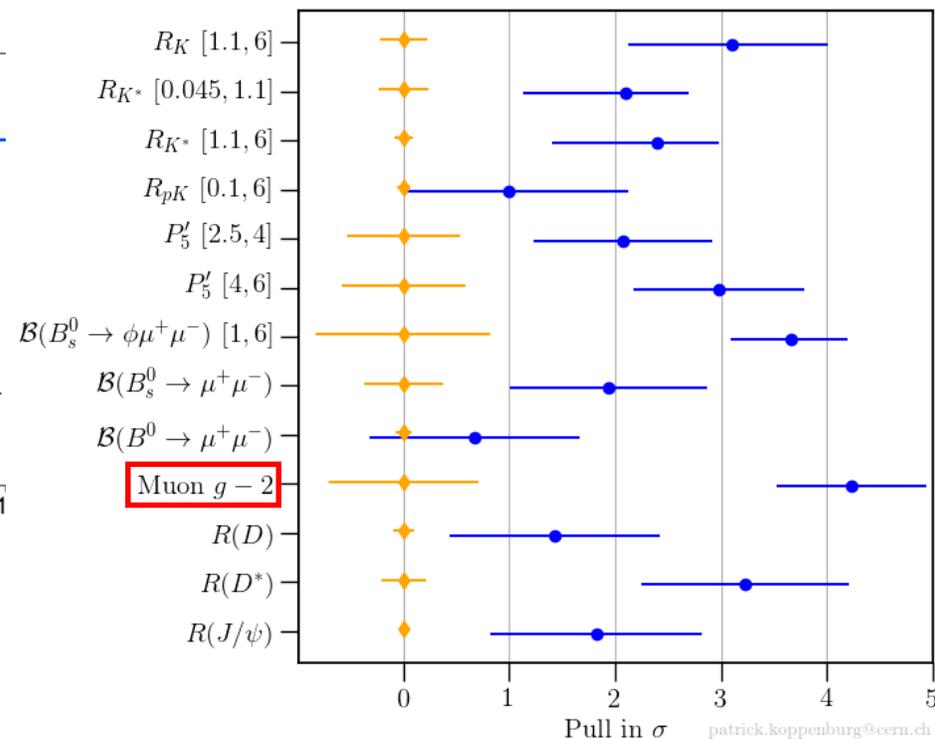
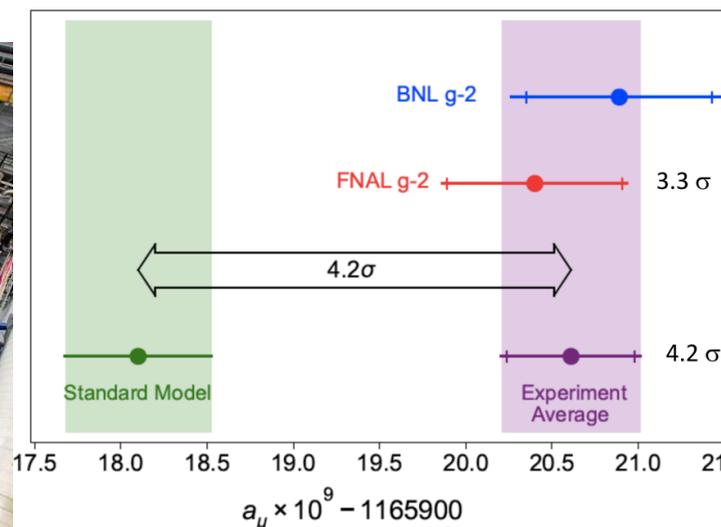
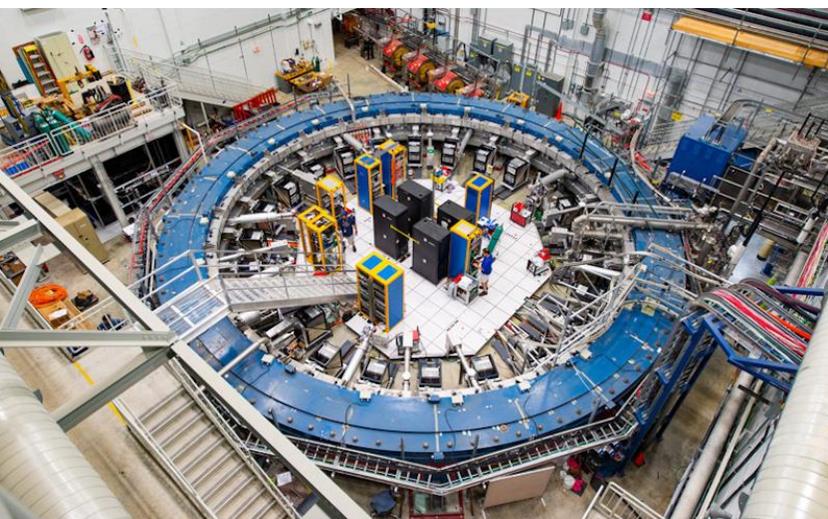
& to appear soon...

$\pi\pi$

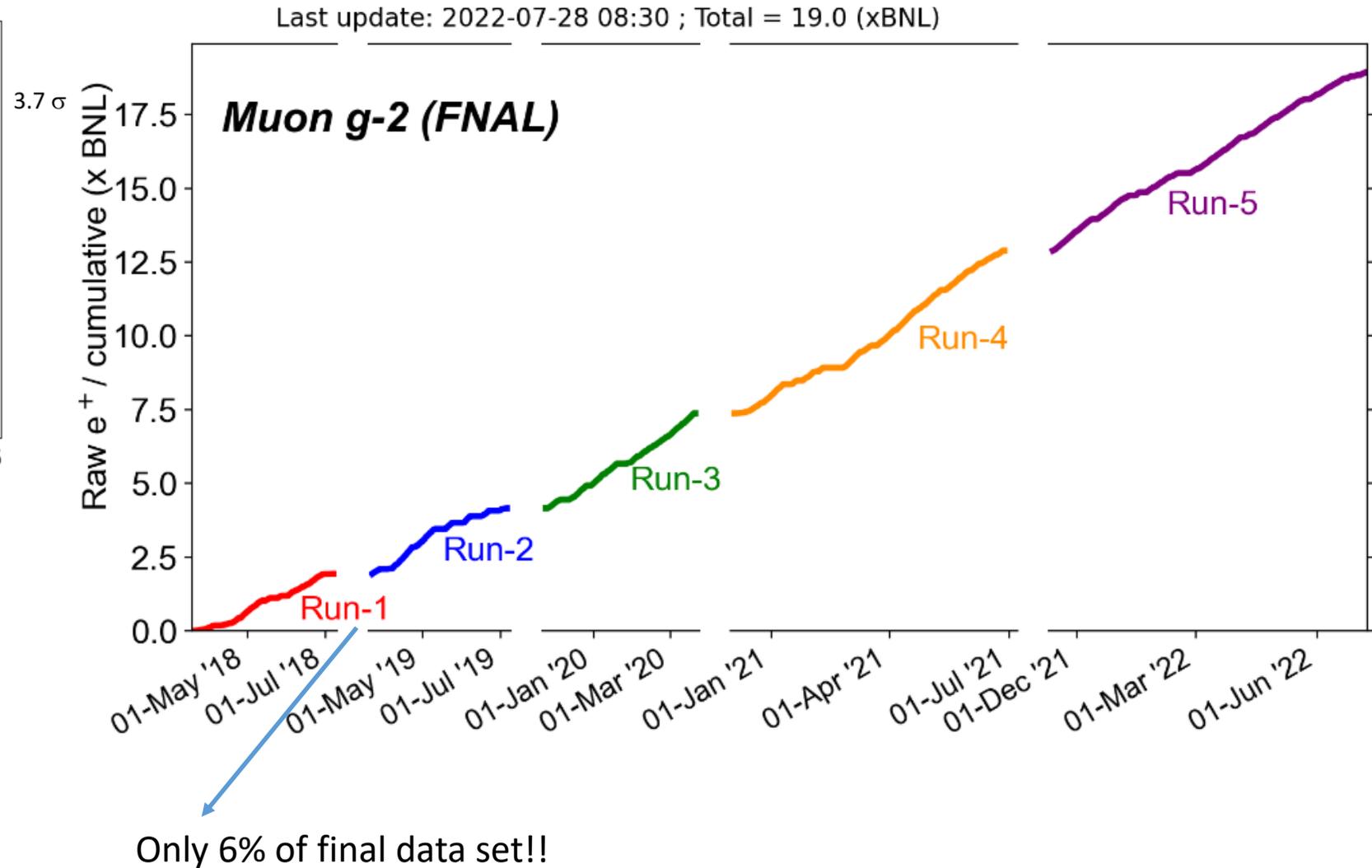
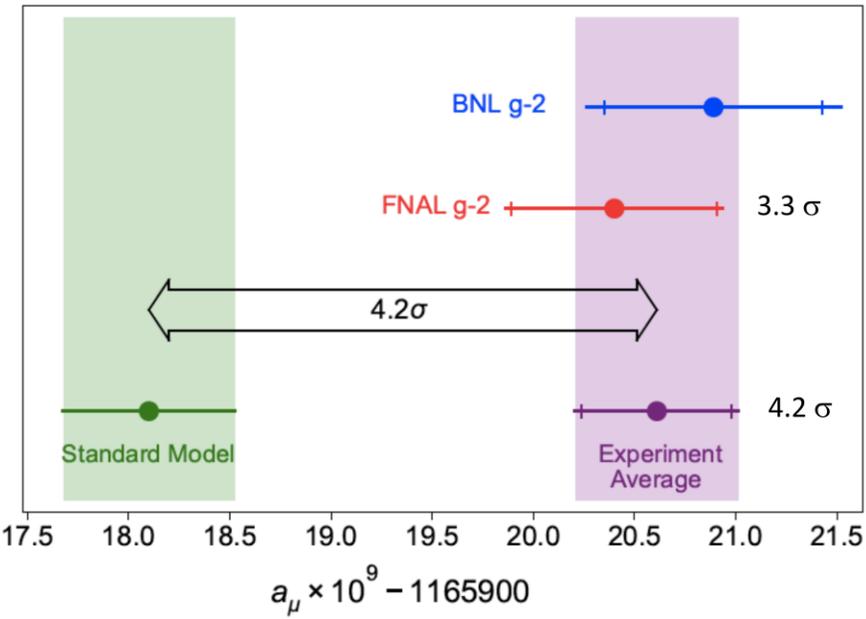
All other modes

# THE MUON ANOMALOUS MAGNETIC MOMENT & POSSIBLE NEW PHYSICS AFTER THE FNAL 1st MEASUREMENT

Support from Cátedra Marcos Moshinsky is acknowledged



# Waiting eagerly for next FNAL measurement in 2023...

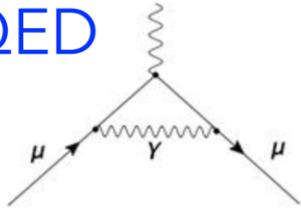


$$a_\mu(\text{SM}) = a_\mu(\text{QED}) + a_\mu(\text{Weak}) + a_\mu(\text{Hadronic})$$

<https://doi.org/10.1016/j.physrep.2020.07.006> (SM prediction, 'White Paper')

<https://muon-gm2-theory.illinois.edu>

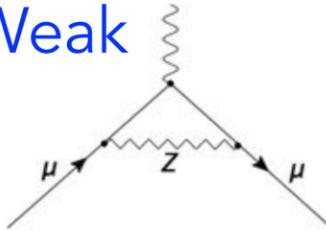
QED



+ ...

$$116\,584\,718.9(1) \times 10^{-11} \quad 0.001 \text{ ppm}$$

Weak



+ ...

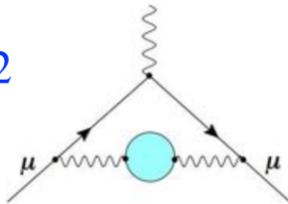
$$153.6(1.0) \times 10^{-11} \quad 0.01 \text{ ppm}$$

Hadronic...

**Saturate the error**

...Vacuum Polarization (HVP)

$\alpha^2$



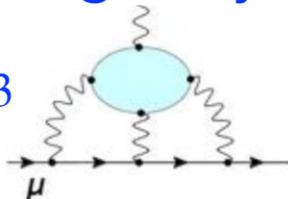
+ ...

$$6845(40) \times 10^{-11} \quad 0.37 \text{ ppm}$$

[0.6%]

...Light-by-Light (HLbL)

$\alpha^3$

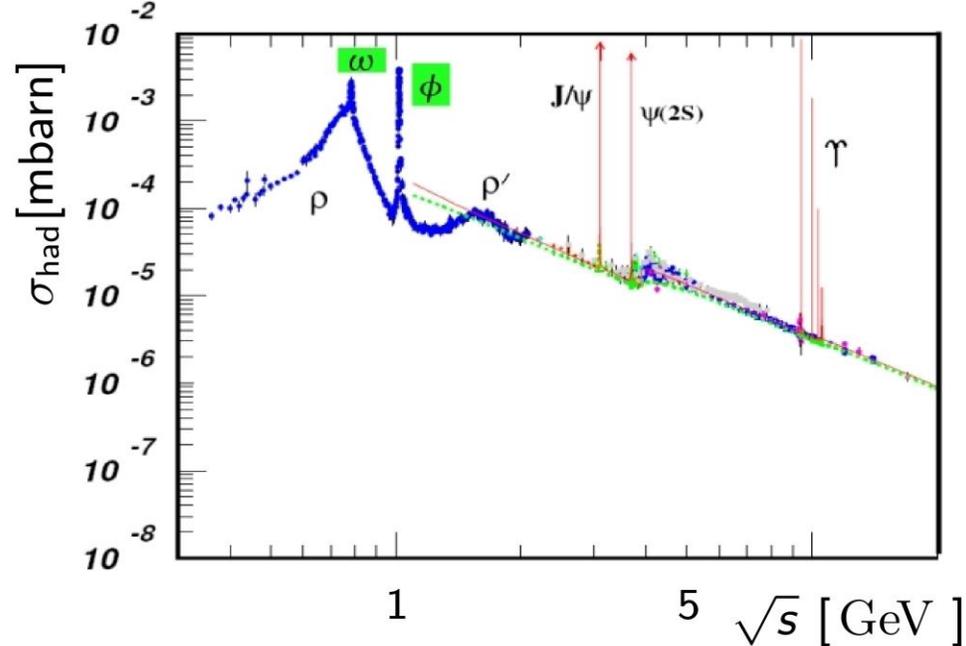
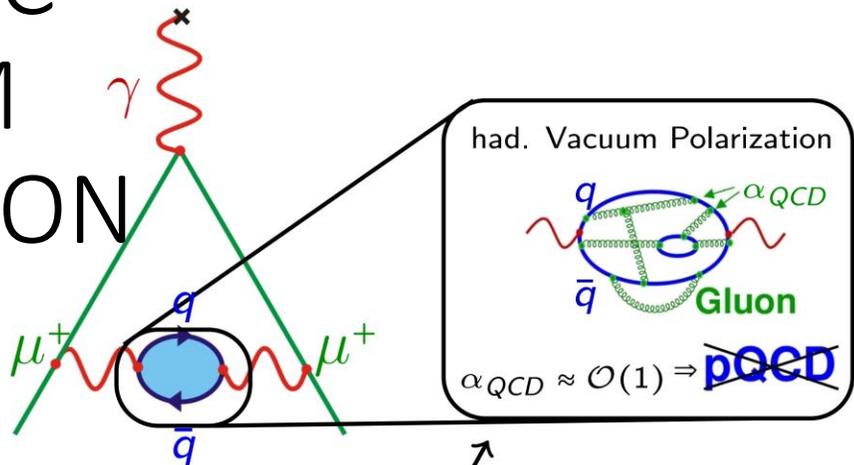


+ ...

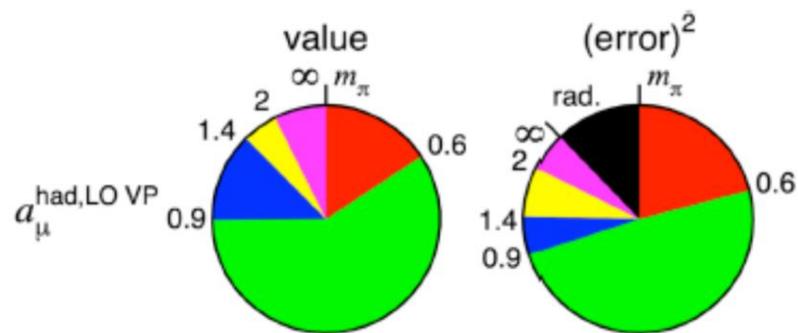
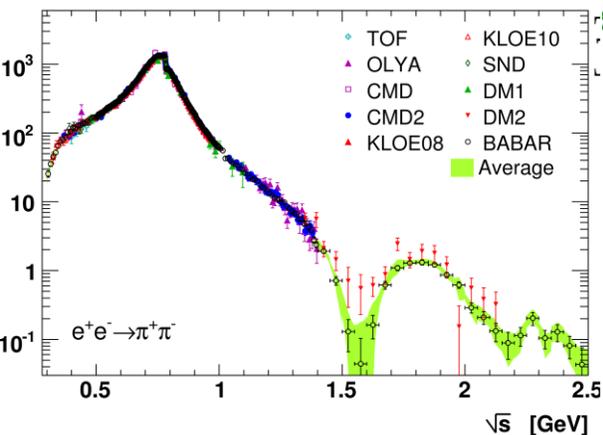
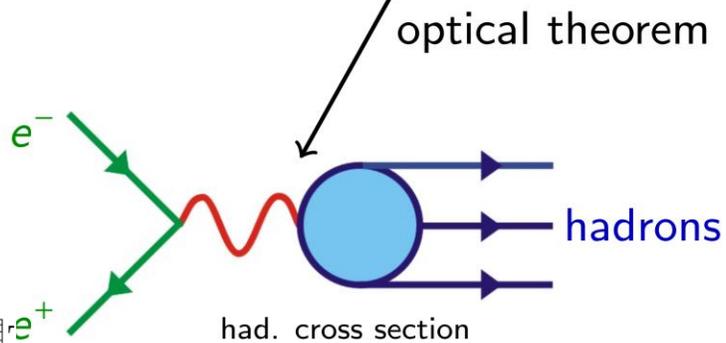
$$92(18) \times 10^{-11} \quad 0.15 \text{ ppm}$$

[20%]

# HADRONIC VACUUM POLARIZATION



## DATA-DRIVEN METHOD



$$a_{\mu,LO}^{\text{had}} = \frac{1}{4\pi^3} \int_{m_{\pi^0}^2}^{\infty} ds K(s) \sigma_{\text{had}}(s)$$

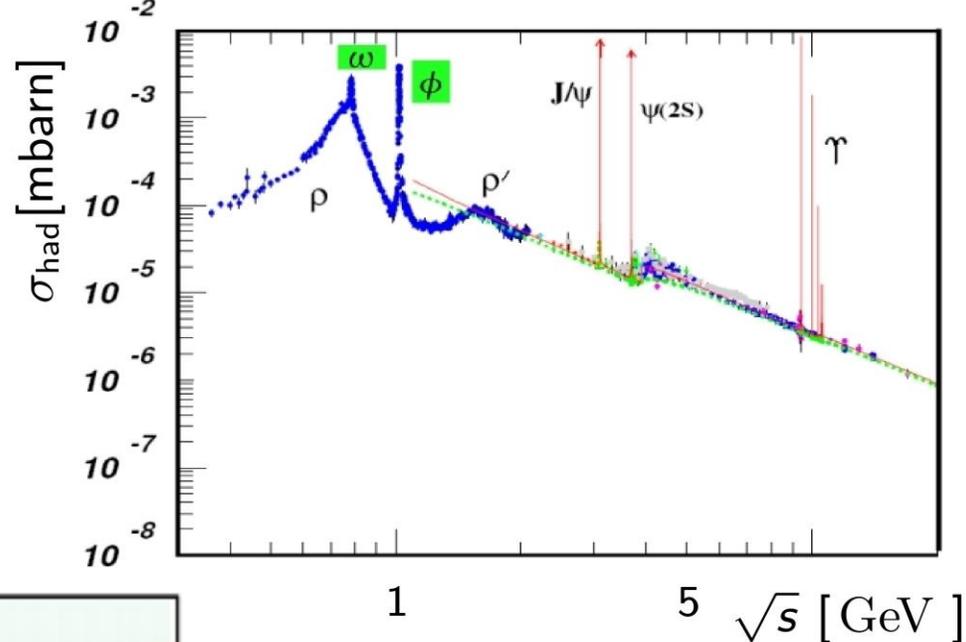
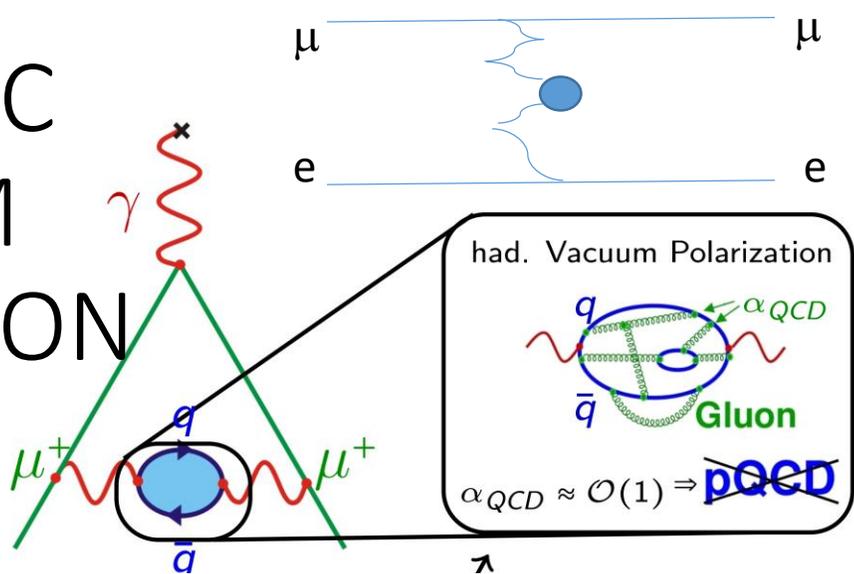
$$\sigma_{\text{had}}(s) \sim 1/s \quad \& \quad K(s) \sim 1/s$$

Low energy region important!  $\sim 1/s^2$

Sum of exclusive  $\sigma_{\text{had}}$

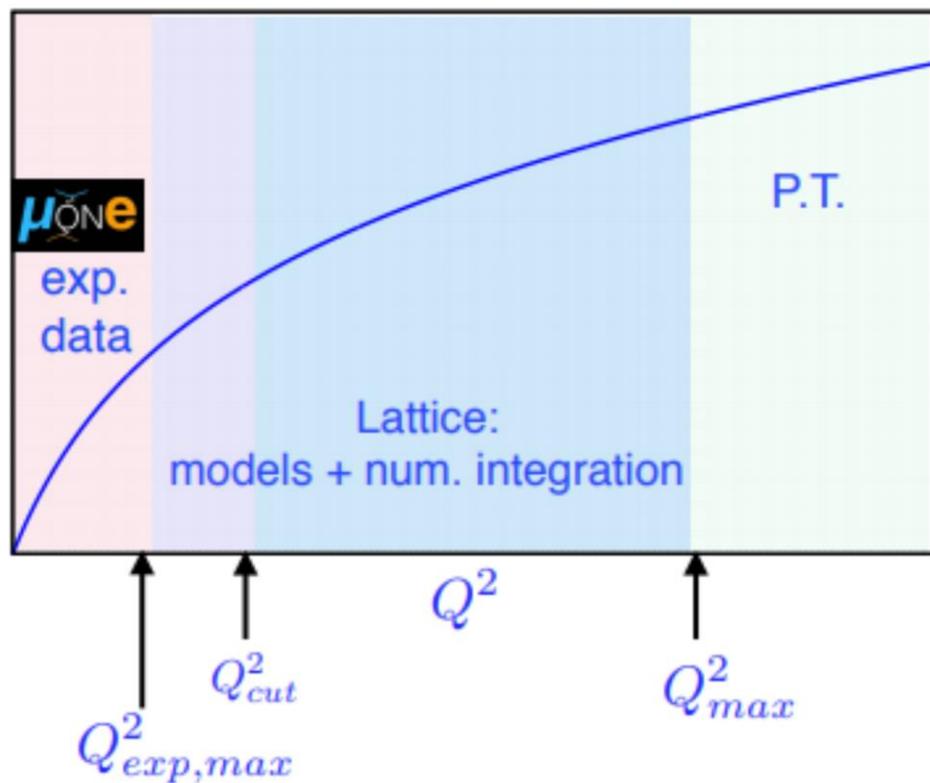
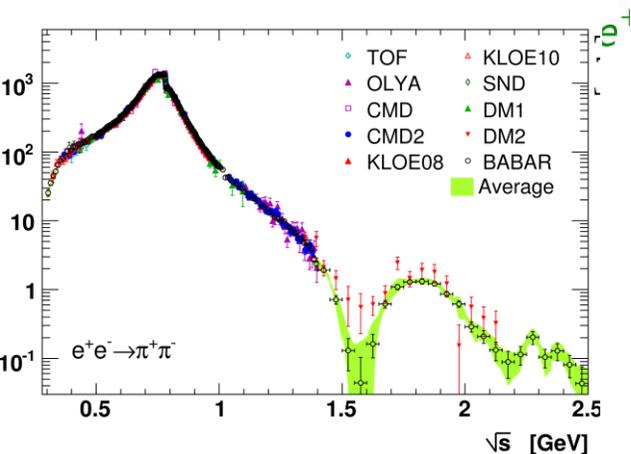
Hadronic contribution of  $a_{\mu}$

# HADRONIC VACUUM POLARIZATION



## DATA-DRIVEN METHOD

(In future also MUonE)



$$= \frac{1}{4\pi^3} \int_{m_{\pi^0}^2}^{\infty} ds K(s) \sigma_{\text{had}}(s)$$

$$d(s) \sim 1/s \quad \& \quad K(s) \sim 1/s$$

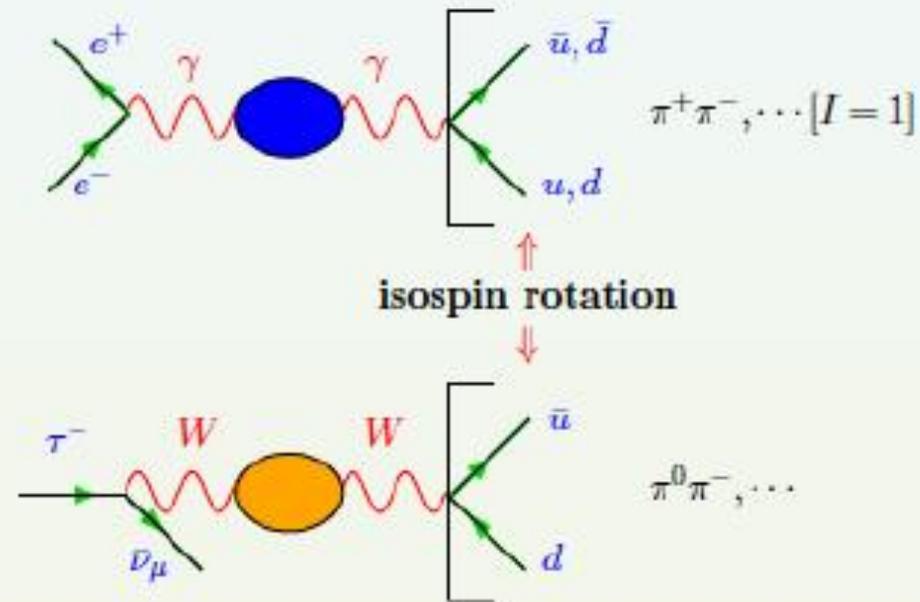
energy region important!  $\sim 1/s^2$

Sum of exclusive  $\sigma_{\text{had}}$

hadronic contribution of  $a_\mu$

Need precise data:

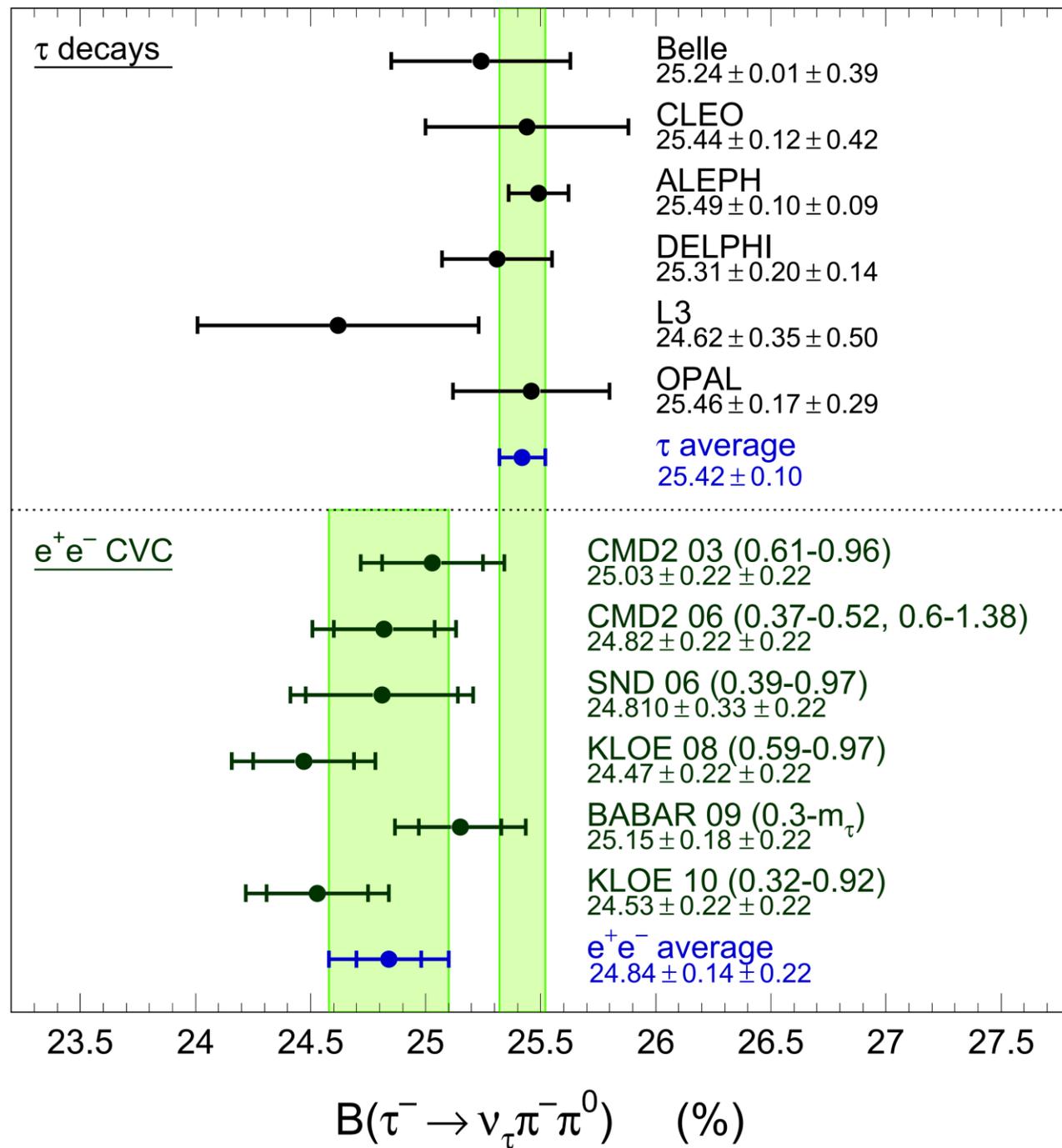
- ❖ Good old idea: use isospin symmetry to include existing high quality  $\tau$ -data (including isospin corrections)



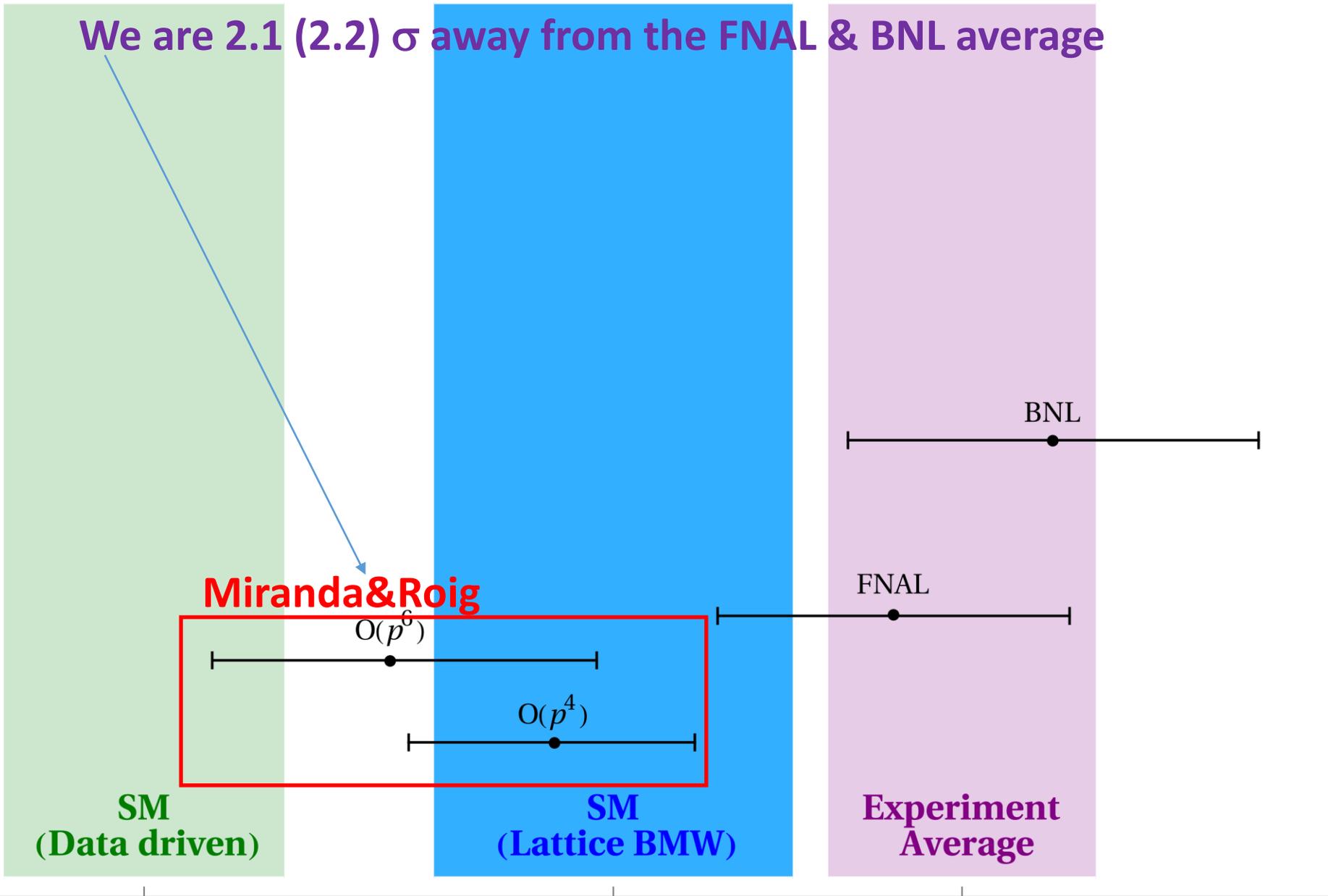
(Francisco Flores-Baez, Alain Flores-Tlalpa, Gabriel López Castro & Genaro Toledo'06&'07)

(Gabriel López Castro, Genaro Toledo & Orsay, CERN & IHEP collaborators'10)

**Corrected data: large discrepancy [ $\sim 10\%$ ] persists!  $\tau$  vs.  $e^+e^-$  problem! [manifest since 2002]**



We are 2.1 (2.2)  $\sigma$  away from the FNAL & BNL average



Miranda&Roig

$O(p^6)$

$O(p^4)$

SM  
(Data driven)

SM  
(Lattice BMW)

Experiment  
Average

0.00116591810

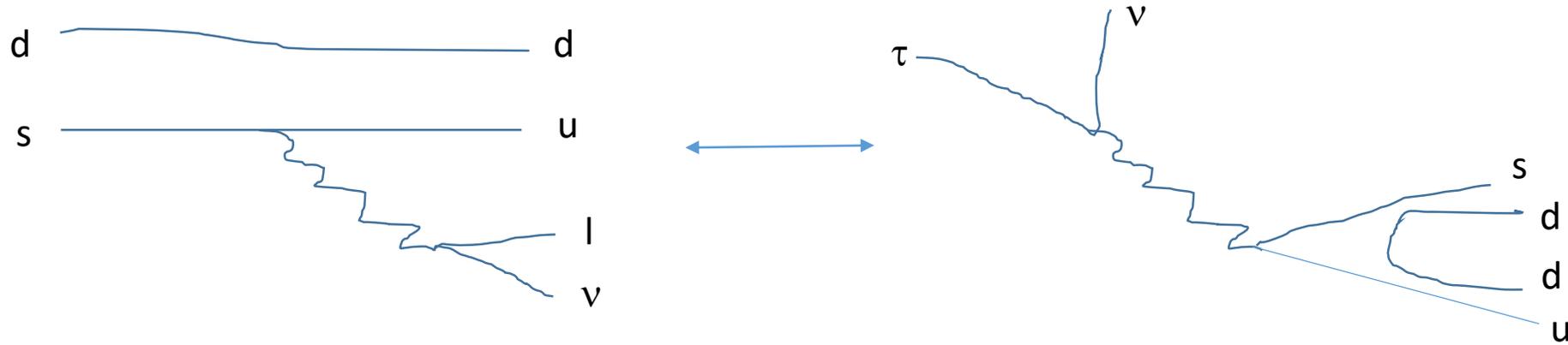
0.00116591954

0.00116592061

$a_\mu$

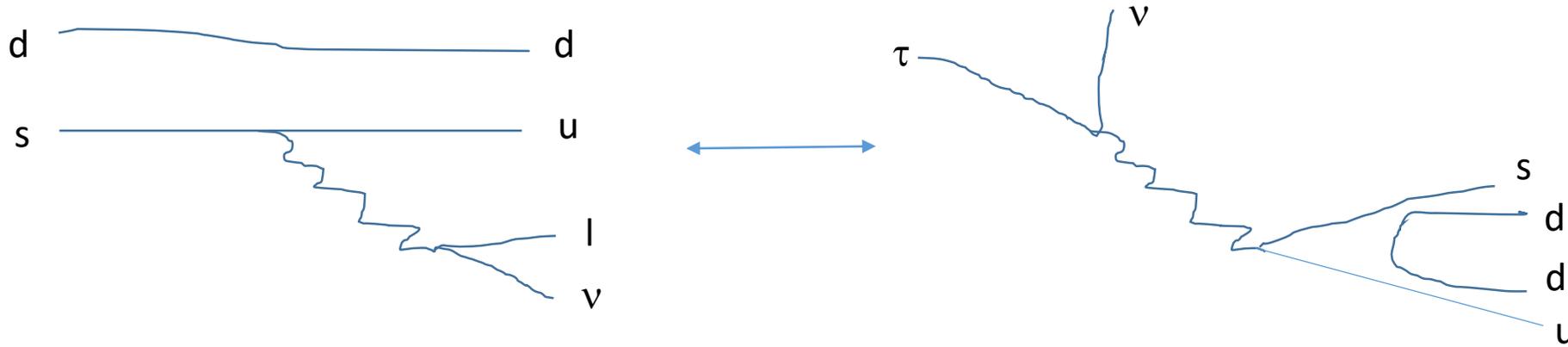
# Radiative corrections to other two-meson channels

Antonelli-Cirigliano-Lusiani-Passemar'13 reduced the CKM unitarity violation by using hadron input from Kaon semileptonic decays in strangeness-changing tau decays



# Radiative corrections to other two-meson channels

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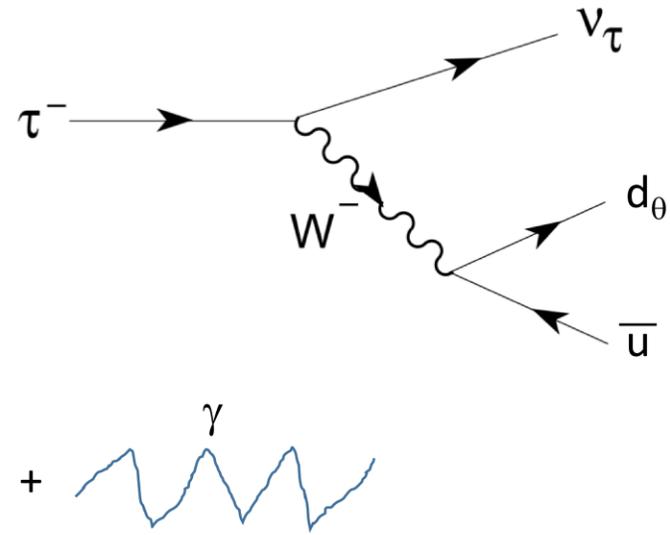


In the one-meson tau decays they used the old Decker-Finkemeier RadCors & for the two-meson channels they computed only the SI part and estimated the size of the model-dependent corrections, **which we have calculated now.**

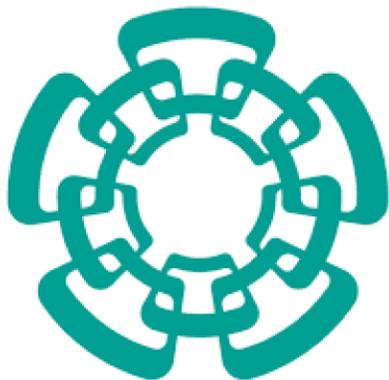
$$\delta^{K^- \pi^0} = \left( -0.009_{-0.118}^{+0.008} \right) \%, \quad \delta^{\bar{K}^0 \pi^-} = \left( -0.166_{-0.122}^{+0.010} \right) \%$$

$$\delta^{K^- K^0} = \left( -0.030_{-0.179}^{+0.026} \right) \%, \quad \delta^{\pi^- \pi^0} = \left( -0.186_{-0.169}^{+0.024} \right) \%$$

**The first two were -0.20(20) and -0.15(20) in Antonelli et al. Our RadCors (also for one-meson case) will enable improved NP tests: most notably CKM unitarity.**



## Radiative corrections in semileptonic tau decays and reliable new physics tests



**Cinvestav**

Pablo Roig  
Cinvestav (Mexico)