



DPC-SMF



XVIII Mexican Workshop on Particles and Fields
November 21st - 25th

Reduction of couplings in two-Higgs-doublet models with natural flavour conservation

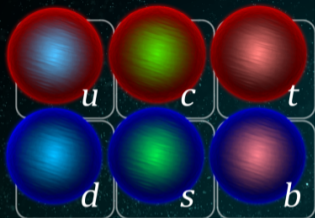
Puebla, México

Miguel Angel May Pech

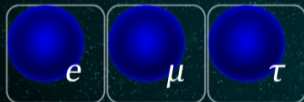
Dra. Myriam Mondragón Ceballos

November 25th, 2022

I. The Standard Model of Particle Physics



Quarks



Leptons

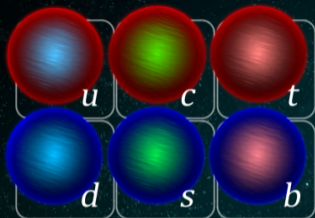


Higgs boson

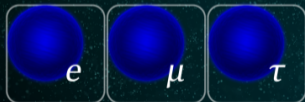


Forces

I. The Standard Model of Particle Physics



Quarks



Leptons

$$\mathcal{L}_{SM} = \mathcal{L}_f + \mathcal{L}_g + \mathcal{L}_H + \mathcal{L}_Y$$

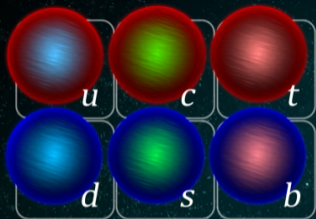


Higgs boson

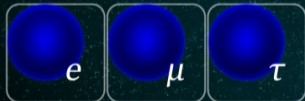


Forces

I. The Standard Model of Particle Physics



Quarks



Leptons

$$\mathcal{L}_{SM} = \mathcal{L}_f + \mathcal{L}_g + \mathcal{L}_H + \mathcal{L}_Y \rightarrow \text{External phenomena}$$

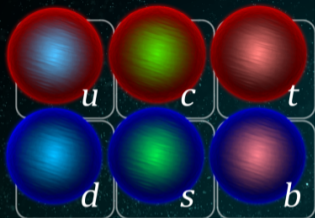


Higgs boson

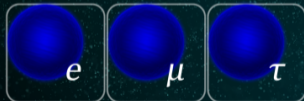


Forces

I. The Standard Model of Particle Physics

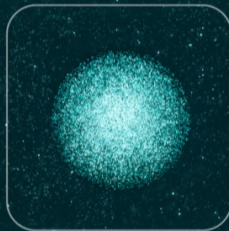


Quarks



Leptons

$$\mathcal{L}_{SM} = \mathcal{L}_f + \mathcal{L}_g + \mathcal{L}_H + \mathcal{L}_Y \begin{cases} \text{External phenomena} \\ \text{Internal issues} \end{cases}$$

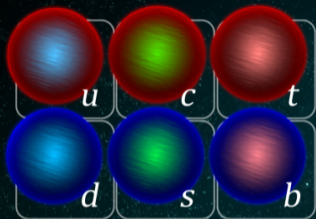


Higgs boson

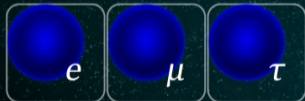


Forces

I. The Standard Model of Particle Physics



Quarks



Leptons

$$\mathcal{L}_{SM} = \mathcal{L}_f + \mathcal{L}_g + \mathcal{L}_H + \mathcal{L}_Y \rightarrow \text{External phenomena}$$

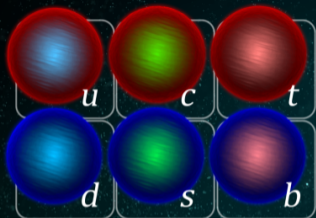


Higgs boson



Forces

I. The Standard Model of Particle Physics



Quarks



Leptons



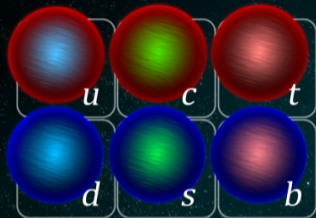
$$\mathcal{L}_{SM} = \mathcal{L}_f + \mathcal{L}_g + \mathcal{L}_H + \mathcal{L}_Y \rightarrow \text{External phenomena}$$



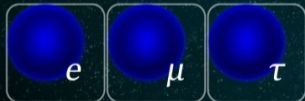
Higgs boson



I. The Standard Model of Particle Physics



Quarks



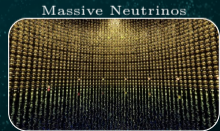
Leptons



$$\mathcal{L}_{SM} = \mathcal{L}_f + \mathcal{L}_g + \mathcal{L}_H + \mathcal{L}_Y \rightarrow \text{External phenomena}$$



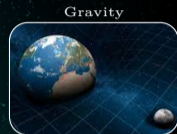
Higgs boson



Massive Neutrinos

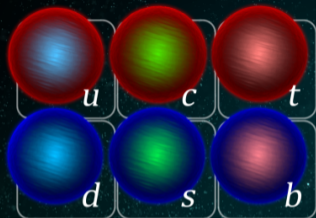
 γ  W^+  Z^0  g  W^-

Forces

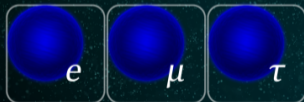


Gravity

I. The Standard Model of Particle Physics



Quarks



Leptons



$$\mathcal{L}_{SM} = \mathcal{L}_f + \mathcal{L}_g + \mathcal{L}_H + \mathcal{L}_Y \rightarrow \text{External phenomena}$$



Higgs boson



γ

W⁺Z⁰

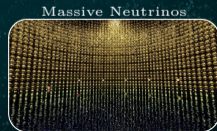
g

W⁻

Forces



Gravity

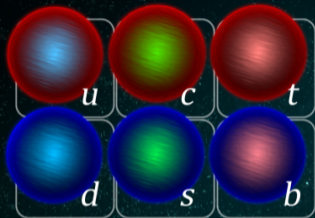


Massive Neutrinos

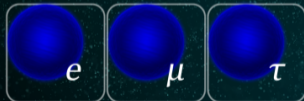


Dark Matter

I. The Standard Model of Particle Physics



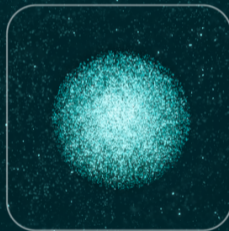
Quarks



Leptons



$$\mathcal{L}_{SM} = \mathcal{L}_f + \mathcal{L}_g + \mathcal{L}_H + \mathcal{L}_Y \rightarrow \text{Internal issues}$$

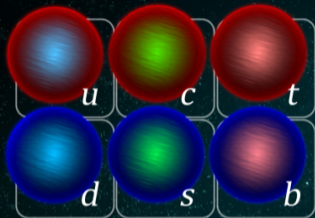


Higgs boson

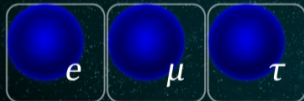


Forces

I. The Standard Model of Particle Physics



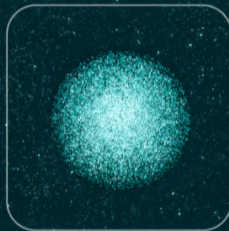
Quarks



Leptons



$$\mathcal{L}_{SM} = \mathcal{L}_f + \mathcal{L}_g + \mathcal{L}_H + \mathcal{L}_Y \rightarrow \text{Free parameters problem}$$

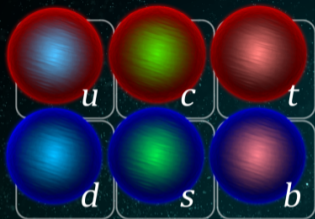


Higgs boson

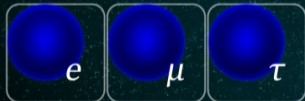


Forces

I. The Standard Model of Particle Physics



Quarks



Leptons

$$\mathcal{L}_{SM} = \mathcal{L}_f + \mathcal{L}_g + \mathcal{L}_H + \mathcal{L}_Y \rightarrow \text{Free parameters problem}$$



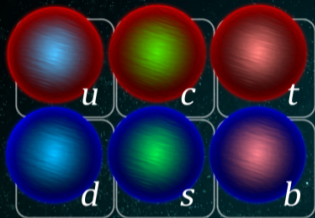
Higgs boson



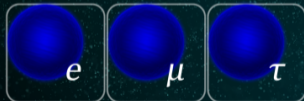
Forces

Even up to 19 free parameters!

I. The Standard Model of Particle Physics



Quarks



Leptons



$$\mathcal{L}_{SM} = \mathcal{L}_f + \mathcal{L}_g + \mathcal{L}_H + \mathcal{L}_Y \rightarrow \text{Free parameters problem}$$

Reduction of Couplings

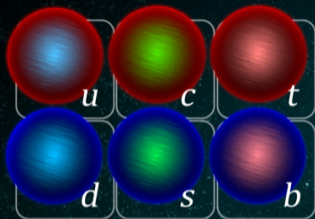


Higgs boson

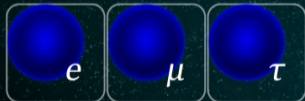


Forces

I. The Standard Model of Particle Physics



Quarks



Leptons



$$\mathcal{L}_{SM} = \mathcal{L}_f + \mathcal{L}_g + \mathcal{L}_H + \mathcal{L}_Y$$

External phenomena
Free parameters problem



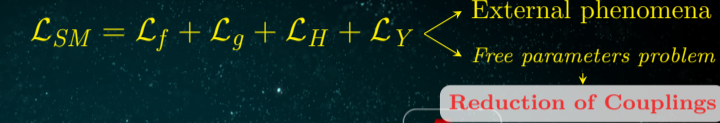
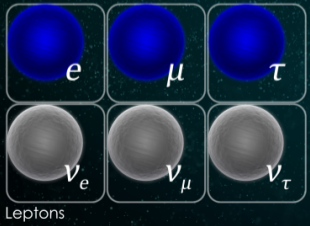
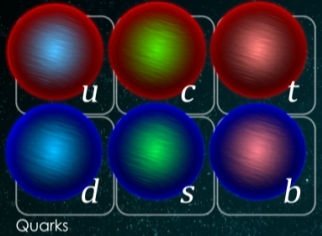
Higgs boson

Reduction of Couplings



Forces

I. The Standard Model of Particle Physics



II. Reduction of Couplings Method

Given a model with $n + 1$ coupling parameters: $\lambda_0, \lambda_1, \dots, \lambda_n \dots$

$$\lambda_j = \lambda_j(\lambda_0) \quad (1)$$

1st Condition: *Renormalization groups equations* (RGEs) or β_j functions on the *reduced system* must fulfil

$$\left[k \frac{\partial}{\partial k} + \sum_j \beta_j \frac{\partial}{\partial \lambda_j} + \gamma \right] \tau = 0 \rightarrow \left[k \frac{\partial}{\partial k} + \beta' \frac{\partial}{\partial \lambda_0} + \gamma' \right] \tau' = 0$$

Reducibility condition

$$\beta' \frac{d\lambda_j}{d\lambda_0} = \beta_j \quad (2)$$

Zimmermann W. (1985) *Reduction in the number of coupling parameters*. Commun. Math. Phys. 97,211-225

II. Reduction of Couplings Method

Given a model with $n + 1$ coupling parameters: $\lambda_0, \lambda_1, \dots, \lambda_n \dots$

$$\lambda_j = \lambda_j(\lambda_0) \quad (1)$$

1st Condition: *Renormalization groups equations* (RGEs) or β_j functions on the *reduced system* must fulfil

$$\left[k \frac{\partial}{\partial k} + \sum_j \beta_j \frac{\partial}{\partial \lambda_j} + \gamma \right] \tau = 0 \rightarrow \left[k \frac{\partial}{\partial k} + \beta' \frac{\partial}{\partial \lambda_0} + \gamma' \right] \tau' = 0$$

Reducibility condition

$$\beta' \frac{d\lambda_j}{d\lambda_0} = \beta_j \quad (2)$$

Zimmermann W. (1985) *Reduction in the number of coupling parameters*. Commun. Math. Phys. 97,211-225

II. Reduction of Couplings Method

Given a model with $n + 1$ coupling parameters: $\lambda_0, \lambda_1, \dots, \lambda_n \dots$

$$\lambda_j = \lambda_j(\lambda_0) \quad (1)$$

1st Condition: *Renormalization groups equations* (RGEs) or β_j functions on the *reduced system* must fulfil

$$\left[k \frac{\partial}{\partial k} + \sum_j \beta_j \frac{\partial}{\partial \lambda_j} + \gamma \right] \tau = 0 \rightarrow \left[k \frac{\partial}{\partial k} + \beta' \frac{\partial}{\partial \lambda_0} + \gamma' \right] \tau' = 0$$

Reducibility condition

$$\beta' \frac{d\lambda_j}{d\lambda_0} = \beta_j \quad (2)$$

Zimmermann W. (1985) *Reduction in the number of coupling parameters*. Commun. Math. Phys. 97,211-225

II. Reduction of Couplings Method

Given a model with $n + 1$ coupling parameters: $\lambda_0, \lambda_1, \dots, \lambda_n \dots$

$$\lambda_j = \lambda_j(\lambda_0) \quad (1)$$

1st Condition: *Renormalization groups equations* (RGEs) or β_j functions on the *reduced system* must fulfil

$$\left[k \frac{\partial}{\partial k} + \sum_j \beta_j \frac{\partial}{\partial \lambda_j} + \gamma \right] \tau = 0 \rightarrow \left[k \frac{\partial}{\partial k} + \beta' \frac{\partial}{\partial \lambda_0} + \gamma' \right] \tau' = 0$$

Reducibility condition

$$\beta' \frac{d\lambda_j}{d\lambda_0} = \beta_j \quad (2)$$

Zimmermann W. (1985) *Reduction in the number of coupling parameters*. Commun. Math. Phys. 97,211-225

II. Reduction of Couplings Method

Given a model with $n + 1$ coupling parameters: $\lambda_0, \lambda_1, \dots, \lambda_n \dots$

$$\lambda_j = \lambda_j(\lambda_0) \quad (1)$$

1st Condition: *Renormalization groups equations* (RGEs) or β_j functions on the *reduced system* must fulfil

$$\left[k \frac{\partial}{\partial k} + \sum_j \beta_j \frac{\partial}{\partial \lambda_j} + \gamma \right] \tau = 0 \rightarrow \left[k \frac{\partial}{\partial k} + \beta' \frac{\partial}{\partial \lambda_0} + \gamma' \right] \tau' = 0$$

Reducibility condition

$$\beta_0 \frac{d\lambda_j}{d\lambda_0} = \beta_j \quad (2)$$

Zimmermann W. (1985) *Reduction in the number of coupling parameters*. Commun. Math. Phys. 97,211-225

II. Reduction of Couplings Method

Given a model with $n + 1$ coupling parameters: $\lambda_0, \lambda_1, \dots, \lambda_n \dots$

$$\lambda_j = \lambda_j(\lambda_0) \quad (1)$$

1st Condition: *Renormalization groups equations* (RGEs) or β_j functions on the *reduced system* must fulfil

$$\left[k \frac{\partial}{\partial k} + \sum_j \beta_j \frac{\partial}{\partial \lambda_j} + \gamma \right] \tau = 0 \rightarrow \left[k \frac{\partial}{\partial k} + \beta' \frac{\partial}{\partial \lambda_0} + \gamma' \right] \tau' = 0$$

Reducibility condition

$$\beta_0 \frac{d\lambda_j}{d\lambda_0} = \beta_j \rightarrow n \text{ reduction equations} \quad (2)$$

Zimmermann W. (1985) *Reduction in the number of coupling parameters*. Commun. Math. Phys. 97,211-225

2nd Condition: Weak coupling limit

$$\lim_{\lambda_0 \rightarrow 0} \lambda_j(\lambda_0) = 0$$

Complete solutions of the reduction equations

$$\lambda_i(\lambda_0) = \lambda_i^{(0)} \lambda_0 + \lambda_i^{(1)} \lambda_0^2 + \lambda_i^{(2)} \lambda_0^3 + \dots, \quad (3)$$

Reduction of Couplings on the Standard Model:

2nd Condition: Weak coupling limit

$$\lim_{\lambda_0 \rightarrow 0} \lambda_j(\lambda_0) = 0$$

Complete solutions of the reduction equations

$$\lambda_i(\lambda_0) = \lambda_i^{(0)} \lambda_0 + \lambda_i^{(1)} \lambda_0^2 + \lambda_i^{(2)} \lambda_0^3 + \dots, \quad (3)$$

Reduction of Couplings on the Standard Model:

2nd Condition: Weak coupling limit

$$\lim_{\lambda_0 \rightarrow 0} \lambda_j(\lambda_0) = 0$$

Complete solutions of the reduction equations

$$\lambda_i(\lambda_0) = \lambda_i^{(0)} \lambda_0 + \lambda_i^{(1)} \lambda_0^2 + \lambda_i^{(2)} \lambda_0^3 + \dots, \quad (3)$$

Reduction of Couplings on the Standard Model:

2nd Condition: Weak coupling limit

$$\lim_{\lambda_0 \rightarrow 0} \lambda_j(\lambda_0) = 0$$

Complete solutions of the reduction equations

$$\lambda_i(\lambda_0) = \lambda_i^{(0)} \lambda_0 + \lambda_i^{(1)} \lambda_0^2 + \lambda_i^{(2)} \lambda_0^3 + \dots, \quad (3)$$

Reduction of Couplings on the Standard Model:

2nd Condition: Weak coupling limit

$$\lim_{\lambda_0 \rightarrow 0} \lambda_j(\lambda_0) = 0$$

Complete solutions of the reduction equations

$$\lambda_i(\lambda_0) = \lambda_i^{(0)} \lambda_0 + \lambda_i^{(1)} \lambda_0^2 + \lambda_i^{(2)} \lambda_0^3 + \dots, \quad (3)$$

Reduction of Couplings on the Standard Model: Kubo, Sibold & Zimmermann

2nd Condition: Weak coupling limit

$$\lim_{\lambda_0 \rightarrow 0} \lambda_j(\lambda_0) = 0$$

Complete solutions of the reduction equations

$$\lambda_i(\lambda_0) = \lambda_i^{(0)} \lambda_0 + \lambda_i^{(1)} \lambda_0^2 + \lambda_i^{(2)} \lambda_0^3 + \dots, \quad (3)$$

Reduction of Couplings on the Standard Model: Kubo, Sibold & Zimmermann

Nuclear Physics B 259 (1985) 331-350

Nuclear Physics B259 (1985) 331-350
© North-Holland Publishing Company

HIGGS AND TOP MASS FROM REDUCTION OF COUPLINGS

J. KUBO*, K. SIBOLD** and W. ZIMMERMANN

Max-Planck-Institut für Physik und Astrophysik Werner-Heisenberg-Institut für Physik
Föhringer Ring 6, 8000 München 40, Federal Republic of Germany

Received 15 April 1985

We reduce the couplings in the standard model with one Higgs doublet and n generations and obtain for three generations 61 GeV and 81 GeV for the mass of the Higgs particle and the top quark respectively. The error is estimated to be about 10-15%.

2nd Condition: Weak coupling limit

$$\lim_{\lambda_0 \rightarrow 0} \lambda_j(\lambda_0) = 0$$

Complete solutions of the reduction equations

$$\lambda_i(\lambda_0) = \lambda_i^{(0)} \lambda_0 + \lambda_i^{(1)} \lambda_0^2 + \lambda_i^{(2)} \lambda_0^3 + \dots, \quad (3)$$

Reduction of Couplings on the Standard Model: Kubo, Sibold & Zimmermann

Nuclear Physics B 259 (1985) 331-350

Nuclear Physics B259 (1985) 331-350
© North-Holland Publishing Company

HIGGS AND TOP MASS FROM REDUCTION OF COUPLINGS

J. KUBO*, K. SIBOLD** and W. ZIMMERMANN

Max-Planck-Institut für Physik und Astrophysik Werner-Heisenberg-Institut für Physik
Föhringer Ring 6, 8000 München 40, Federal Republic of Germany

Received 15 April 1985

We reduce the couplings in the standard model with one Higgs doublet and n generations and obtain for three generations 61 GeV and 81 GeV for the mass of the Higgs particle and the top quark respectively. The error is estimated to be about 10-15%.

They found non-trivial solutions for the top quark Yukawa coupling constant and Higgs self-coupling constant

2nd Condition: Weak coupling limit

$$\lim_{\lambda_0 \rightarrow 0} \lambda_j(\lambda_0) = 0$$

Complete solutions of the reduction equations

$$\lambda_i(\lambda_0) = \lambda_i^{(0)} \lambda_0 + \lambda_i^{(1)} \lambda_0^2 + \lambda_i^{(2)} \lambda_0^3 + \dots, \quad (3)$$

Reduction of Couplings on the Standard Model: Kubo, Sibold & Zimmermann *Nuclear Physics B 259 (1985) 331-350*

Nuclear Physics B259 (1985) 331-350
© North-Holland Publishing Company

HIGGS AND TOP MASS FROM REDUCTION OF COUPLINGS

J. KUBO*, K. SIBOLD** and W. ZIMMERMANN

Max-Planck-Institut für Physik und Astrophysik Werner-Heisenberg-Institut für Physik
Föhringer Ring 6, 8000 München 40, Federal Republic of Germany

Received 15 April 1985

We reduce the couplings in the standard model with one Higgs doublet and n generations and obtain for three generations 61 GeV and 81 GeV for the mass of the Higgs particle and the top quark respectively. The error is estimated to be about 10-15%.

They found non-trivial solutions for the top quark Yukawa coupling constant and Higgs self-coupling constant

$$m_t = 81 \text{ GeV}$$

$$m_H = 61 \text{ GeV}$$

III. multi-Higgs Models

4.1 Motivations

- ▶ It is one of the easiest extensions to implement...
- ▶ The notion of generations can be brought in a more natural way
- ▶ Allows for a large amount of phenomenology:
 - Several Higgs bosons, charged and neutral.
 - Flavour Changing Neutral currents (FCNC) at tree level.
 - Additional forms of CP violation.
 - Opportunities for cosmology, as are dark matter candidates.

Ivanov ([Prog. Part. Nucl. Phys. 95 \(2017\) 160-208](#)).

III. multi-Higgs Models

4.1 Motivations

- ▶ It is one of the easiest extensions to implement...
- ▶ The notion of generations can be brought in a more natural way
- ▶ Allows for a large amount of phenomenology:
 - Several Higgs bosons, charged and neutral.
 - Flavour Changing Neutral currents (FCNC) at tree level.
 - Additional forms of CP violation.
 - Opportunities for cosmology, as are dark matter candidates.

Ivanov ([Prog. Part. Nucl. Phys. 95 \(2017\) 160-208](#)).

III. multi-Higgs Models

4.1 Motivations

- ▶ It is one of the easiest extensions to implement...
- ▶ The notion of generations can be brought in a more natural way
- ▶ Allows for a large amount of phenomenology:
 - Several Higgs bosons, charged and neutral.
 - Flavour Changing Neutral currents (FCNC) at tree level.
 - Additional forms of CP violation.
 - Opportunities for cosmology, as are dark matter candidates.



Ivanov ([Prog. Part. Nucl. Phys. 95 \(2017\) 160-208](#)).

III. multi-Higgs Models

4.1 Motivations

- ▶ It is one of the easiest extensions to implement...
- ▶ The notion of generations can be brought in a more natural way
- ▶ Allows for a large amount of phenomenology:
 - Several Higgs bosons, charged and neutral.
 - Flavour Changing Neutral currents (FCNC) at tree level.
 - Additional forms of CP violation.
 - Opportunities for cosmology, as are dark matter candidates.



Ivanov ([Prog. Part. Nucl. Phys. 95 \(2017\) 160-208](#)).

III. multi-Higgs Models

4.1 Motivations



- ▶ It is one of the easiest extensions to implement...
- ▶ The notion of generations can be brought in a more natural way
- ▶ Allows for a large amount of phenomenology:
 - Several Higgs bosons, charged and neutral.
 - Flavour Changing Neutral currents (FCNC) at tree level.
 - Additional forms of CP violation.
 - Opportunities for cosmology, as are dark matter candidates.



Ivanov ([Prog. Part. Nucl. Phys. 95 \(2017\) 160-208](#)).

III. multi-Higgs Models

4.1 Motivations



- ▶ It is one of the easiest extensions to implement...
- ▶ The notion of generations can be brought in a more natural way
- ▶ Allows for a large amount of phenomenology:
 - Several Higgs bosons, charged and neutral.
 - Flavour Changing Neutral currents (FCNC) at tree level.
 - Additional forms of CP violation.
 - Opportunities for cosmology, as are dark matter candidates.



Ivanov ([Prog. Part. Nucl. Phys. 95 \(2017\) 160-208](#)).



III. multi-Higgs Models

4.1 Motivations



- ▶ It is one of the easiest extensions to implement...
- ▶ The notion of generations can be brought in a more natural way
- ▶ Allows for a large amount of phenomenology:
 - Several Higgs bosons, charged and neutral.
 - Flavour Changing Neutral currents (FCNC) at tree level.
 - Additional forms of CP violation.
 - Opportunities for cosmology, as are dark matter candidates.



Ivanov ([Prog. Part. Nucl. Phys. 95 \(2017\) 160-208](#)).



III. multi-Higgs Models

4.1 Motivations



- ▶ It is one of the easiest extensions to implement...
- ▶ The notion of generations can be brought in a more natural way
- ▶ Allows for a large amount of phenomenology:



- Several Higgs bosons, charged and neutral.
- Flavour Changing Neutral currents (FCNC) at tree level.
- Additional forms of CP violation.
- Opportunities for cosmology, as are dark matter candidates.



Ivanov (Prog. Part. Nucl. Phys. 95 (2017) 160-208).



III. multi-Higgs Models

4.1 Motivations



- ▶ It is one of the easiest extensions to implement...
- ▶ The notion of generations can be brought in a more natural way
- ▶ Allows for a large amount of phenomenology:



- Several Higgs bosons, charged and neutral.
- Flavour Changing Neutral currents (FCNC) at tree level.
- Additional forms of CP violation.
- Opportunities for cosmology, as are dark matter candidates.



Ivanov ([Prog. Part. Nucl. Phys. 95 \(2017\) 160-208](#)).



III. multi-Higgs Models

3.2 Implementation

Scalar Sector

Yukawa Sector

III. multi-Higgs Models

$$\mathcal{L}_{SM} = \mathcal{L}_f + \mathcal{L}_g + \mathcal{L}_H + \mathcal{L}_Y$$

3.2 Implementation

Scalar Sector

Yukawa Sector

III. multi-Higgs Models

$$\mathcal{L}_{SM} = \mathcal{L}_f + \mathcal{L}_g + \mathcal{L}_H + \mathcal{L}_Y$$

3.2 Implementation

Scalar Sector

$$\mathcal{L}_H = \sum_{i,j=1}^N (D_\mu \Phi_i)^\dagger (D_\mu \Phi_j) - V(\Phi),$$

Yukawa Sector

III. multi-Higgs Models

$$\mathcal{L}_{SM} = \mathcal{L}_f + \mathcal{L}_g + \mathcal{L}_H + \mathcal{L}_Y$$

3.2 Implementation

Scalar Sector

$$\mathcal{L}_H = \sum_{i,j=1}^N (D_\mu \Phi_i)^\dagger (D_\mu \Phi_j) - V(\Phi),$$

Yukawa Sector

$$\mathcal{L}_Y = - \left(G_i^{(\ell)} \bar{L}_L \Phi_i \ell_R + G_i^{(d)} \bar{Q}_L \Phi_i d_R + G_i^{(u)} \bar{Q}_L \tilde{\Phi}_i u_R \right) + \text{h. c.}$$

III. multi-Higgs Models

$$\mathcal{L}_{SM} = \mathcal{L}_f + \mathcal{L}_g + \mathcal{L}_H + \mathcal{L}_Y$$

3.2 Implementation

Scalar Sector

$$\mathcal{L}_H = \sum_{i,j=1}^N (D_\mu \Phi_i)^\dagger (D_\mu \Phi_j) - V(\Phi),$$

where

$$V(\Phi) = \mu_{ij} \Phi_i^\dagger \Phi_j + \Lambda_{ijkl} (\Phi_i^\dagger \Phi_j) (\Phi_k^\dagger \Phi_l).$$

Yukawa Sector

$$\mathcal{L}_Y = - \left(G_i^{(\ell)} \bar{L}_L \Phi_i \ell_R + G_i^{(d)} \bar{Q}_L \Phi_i d_R + G_i^{(u)} \bar{Q}_L \tilde{\Phi}_i u_R \right) + \text{h. c.}$$

III. multi-Higgs Models

$$\mathcal{L}_{SM} = \mathcal{L}_f + \mathcal{L}_g + \mathcal{L}_H + \mathcal{L}_Y$$

3.2 Implementation

Scalar Sector

$$\mathcal{L}_H = \sum_{i,j=1}^N (D_\mu \Phi_i)^\dagger (D_\mu \Phi_j) - V(\Phi),$$

where

$$V(\Phi) = \mu_{ij} \Phi_i^\dagger \Phi_j + \Lambda_{ijkl} (\Phi_i^\dagger \Phi_j) (\Phi_k^\dagger \Phi_l).$$

$$\mu_{ij} = \mu_{ji}^*,$$

$$\Lambda_{ijkl} = \Lambda_{klij} = \Lambda_{jilk}^*,$$

Yukawa Sector

$$\mathcal{L}_Y = - \left(G_i^{(\ell)} \bar{L}_L \Phi_i \ell_R + G_i^{(d)} \bar{Q}_L \Phi_i d_R + G_i^{(u)} \bar{Q}_L \tilde{\Phi}_i u_R \right) + \text{h. c.}$$

IV. Two-Higgs-doublet models (2HDM)



IV. Two-Higgs-doublet models (2HDM)



$$\begin{aligned}
 V(\Phi) = & \mu_{11}\Phi_1^\dagger\Phi_1 + \mu_{12}\Phi_1^\dagger\Phi_2 + \mu_{21}\Phi_2^\dagger\Phi_1 + \mu_{22}\Phi_2^\dagger\Phi_2 \\
 & + \Phi_1\Phi_1 \left(\Lambda_{1111}\Phi_1^\dagger\Phi_1 + \Lambda_{1112}\Phi_1^\dagger\Phi_2 + \Lambda_{1121}\Phi_2^\dagger\Phi_1 + \Lambda_{1122}\Phi_2^\dagger\Phi_2 \right) \\
 & + \Phi_1\Phi_2 \left(\Lambda_{1211}\Phi_1^\dagger\Phi_1 + \Lambda_{1212}\Phi_1^\dagger\Phi_2 + \Lambda_{1221}\Phi_2^\dagger\Phi_1 + \Lambda_{1222}\Phi_2^\dagger\Phi_2 \right) \\
 & + \Phi_2\Phi_1 \left(\Lambda_{2111}\Phi_1^\dagger\Phi_1 + \Lambda_{2112}\Phi_1^\dagger\Phi_2 + \Lambda_{2121}\Phi_2^\dagger\Phi_1 + \Lambda_{2122}\Phi_2^\dagger\Phi_2 \right) \\
 & + \Phi_2\Phi_2 \left(\Lambda_{2211}\Phi_1^\dagger\Phi_1 + \Lambda_{2212}\Phi_1^\dagger\Phi_2 + \Lambda_{2221}\Phi_2^\dagger\Phi_1 + \Lambda_{2222}\Phi_2^\dagger\Phi_2 \right).
 \end{aligned}$$

IV. Two-Higgs-doublet models (2HDM)



$$\begin{aligned}
 V(\Phi) = & \mu_{11}\Phi_1^\dagger\Phi_1 + \mu_{12}\Phi_1^\dagger\Phi_2 + \mu_{21}\Phi_2^\dagger\Phi_1 + \mu_{22}\Phi_2^\dagger\Phi_2 \\
 & + \Phi_1\Phi_1\left(\Lambda_{1111}\Phi_1^\dagger\Phi_1 + \Lambda_{1112}\Phi_1^\dagger\Phi_2 + \Lambda_{1121}\Phi_2^\dagger\Phi_1 + \Lambda_{1122}\Phi_2^\dagger\Phi_2\right) \\
 & + \Phi_1\Phi_2\left(\Lambda_{1211}\Phi_1^\dagger\Phi_1 + \Lambda_{1212}\Phi_1^\dagger\Phi_2 + \Lambda_{1221}\Phi_2^\dagger\Phi_1 + \Lambda_{1222}\Phi_2^\dagger\Phi_2\right) \\
 & + \Phi_2\Phi_1\left(\Lambda_{2111}\Phi_1^\dagger\Phi_1 + \Lambda_{2112}\Phi_1^\dagger\Phi_2 + \Lambda_{2121}\Phi_2^\dagger\Phi_1 + \Lambda_{2122}\Phi_2^\dagger\Phi_2\right) \\
 & + \Phi_2\Phi_2\left(\Lambda_{2211}\Phi_1^\dagger\Phi_1 + \Lambda_{2212}\Phi_1^\dagger\Phi_2 + \Lambda_{2221}\Phi_2^\dagger\Phi_1 + \Lambda_{2222}\Phi_2^\dagger\Phi_2\right).
 \end{aligned}$$

$$\begin{aligned}
 \mathcal{L}_Y = & -\sum_{i=1}^3 \left(G^{(\ell)_i} \bar{L}_{L,i}(\Phi_1 + \Phi_2)l_{R,i} + G_i^{(d)} \bar{Q}_{L,i}(\Phi_1 + \Phi_2)d_{R,i} \right. \\
 & \left. + G_i^{(u)} \bar{Q}_{L,i}(\tilde{\Phi}_1 + \tilde{\Phi}_2)u_{R,i} + \text{h. c.} \right)
 \end{aligned}$$

IV. Two-Higgs-doublet models (2HDM)



$$\begin{aligned}
 V(\Phi) = & \mu_{11}\Phi_1^\dagger\Phi_1 + \mu_{12}\Phi_1^\dagger\Phi_2 + \mu_{21}\Phi_2^\dagger\Phi_1 + \mu_{22}\Phi_2^\dagger\Phi_2 \\
 & + \Phi_1\Phi_1\left(\Lambda_{1111}\Phi_1^\dagger\Phi_1 + \Lambda_{1112}\Phi_1^\dagger\Phi_2 + \Lambda_{1121}\Phi_2^\dagger\Phi_1 + \Lambda_{1122}\Phi_2^\dagger\Phi_2\right) \\
 & + \Phi_1\Phi_2\left(\Lambda_{1211}\Phi_1^\dagger\Phi_1 + \Lambda_{1212}\Phi_1^\dagger\Phi_2 + \Lambda_{1221}\Phi_2^\dagger\Phi_1 + \Lambda_{1222}\Phi_2^\dagger\Phi_2\right) \\
 & + \Phi_2\Phi_1\left(\Lambda_{2111}\Phi_1^\dagger\Phi_1 + \Lambda_{2112}\Phi_1^\dagger\Phi_2 + \Lambda_{2121}\Phi_2^\dagger\Phi_1 + \Lambda_{2122}\Phi_2^\dagger\Phi_2\right) \\
 & + \Phi_2\Phi_2\left(\Lambda_{2211}\Phi_1^\dagger\Phi_1 + \Lambda_{2212}\Phi_1^\dagger\Phi_2 + \Lambda_{2221}\Phi_2^\dagger\Phi_1 + \Lambda_{2222}\Phi_2^\dagger\Phi_2\right).
 \end{aligned}$$

$$\begin{aligned}
 \mu_{ij} &= \mu_{ji}^*, \\
 \Lambda_{ijkl} &= \Lambda_{klij} = \Lambda_{jilk}^*,
 \end{aligned}$$

$$\begin{aligned}
 \mathcal{L}_Y = & - \sum_{i=1}^3 \left(G^{(\ell)i} \bar{L}_{L,i}(\Phi_1 + \Phi_2)l_{R,i} + G_i^{(d)} \bar{Q}_{L,i}(\Phi_1 + \Phi_2)d_{R,i} \right. \\
 & \left. + G_i^{(u)} \bar{Q}_{L,i}(\tilde{\Phi}_1 + \tilde{\Phi}_2)u_{R,i} + \text{h. c.} \right)
 \end{aligned}$$

IV. Two-Higgs-doublet models (2HDM)

$$\begin{aligned}
 V(\Phi) = & \mu_{11}\Phi_1^\dagger\Phi_1 + \mu_{12}\Phi_1^\dagger\Phi_2 + \mu_{21}\Phi_2^\dagger\Phi_1 + \mu_{22}\Phi_2^\dagger\Phi_2 \\
 & + \Phi_1\Phi_1\left(\Lambda_{1111}\Phi_1^\dagger\Phi_1 + \Lambda_{1112}\Phi_1^\dagger\Phi_2 + \Lambda_{1121}\Phi_2^\dagger\Phi_1 + \Lambda_{1122}\Phi_2^\dagger\Phi_2\right) \\
 & + \Phi_1\Phi_2\left(\Lambda_{1211}\Phi_1^\dagger\Phi_1 + \Lambda_{1212}\Phi_1^\dagger\Phi_2 + \Lambda_{1221}\Phi_2^\dagger\Phi_1 + \Lambda_{1222}\Phi_2^\dagger\Phi_2\right) \\
 & + \Phi_2\Phi_1\left(\Lambda_{2111}\Phi_1^\dagger\Phi_1 + \Lambda_{2112}\Phi_1^\dagger\Phi_2 + \Lambda_{2121}\Phi_2^\dagger\Phi_1 + \Lambda_{2122}\Phi_2^\dagger\Phi_2\right) \\
 & + \Phi_2\Phi_2\left(\Lambda_{2211}\Phi_1^\dagger\Phi_1 + \Lambda_{2212}\Phi_1^\dagger\Phi_2 + \Lambda_{2221}\Phi_2^\dagger\Phi_1 + \Lambda_{2222}\Phi_2^\dagger\Phi_2\right).
 \end{aligned}$$

$$\begin{aligned}
 \mathcal{L}_Y = & -\sum_{i=1}^3 \left(G^{(\ell)_i} \bar{L}_{L,i}(\Phi_1 + \Phi_2)l_{R,i} + G_i^{(d)} \bar{Q}_{L,i}(\Phi_1 + \Phi_2)d_{R,i} \right. \\
 & \left. + G_i^{(u)} \bar{Q}_{L,i}(\tilde{\Phi}_1 + \tilde{\Phi}_2)u_{R,i} + \text{h. c.} \right)
 \end{aligned}$$



$$\begin{aligned}
 \mu_{ij} &= \mu_{ji}^*, \\
 \Lambda_{ijkl} &= \Lambda_{klij} = \Lambda_{jilk}^*,
 \end{aligned}$$

IV. Two-Higgs-doublet models (2HDM)



$$\begin{aligned}
 V(\Phi) = & m_{11}^2 \Phi_1^\dagger \Phi_1 + m_{22}^2 \Phi_2^\dagger \Phi_2 - (m_{12}^2 \Phi_1^\dagger \Phi_2 + \text{h. c.}) \\
 & + \frac{1}{2} \lambda_1 (\Phi_1^\dagger \Phi_1)^2 + \frac{1}{2} \lambda_2 (\Phi_2^\dagger \Phi_2)^2 + \lambda_3 (\Phi_1^\dagger \Phi_1) (\Phi_2^\dagger \Phi_2) + \lambda_4 (\Phi_1^\dagger \Phi_2) (\Phi_2^\dagger \Phi_1) \\
 & + \left(\frac{1}{2} \lambda_5 (\Phi_1^\dagger \Phi_2)^2 + \lambda_6 (\Phi_1^\dagger \Phi_1) (\Phi_1^\dagger \Phi_2) + \lambda_7 (\Phi_2^\dagger \Phi_2) (\Phi_1^\dagger \Phi_2) + \text{h. c.} \right);
 \end{aligned}$$

$$\begin{aligned}
 \mathcal{L}_Y = - \sum_{i=1}^3 \left(G^{(\ell)_i} \bar{L}_{L,i} (\Phi_1 + \Phi_2) l_{R,i} + G_i^{(d)} \bar{Q}_{L,i} (\Phi_1 + \Phi_2) d_{R,i} \right. \\
 \left. + G_i^{(u)} \bar{Q}_{L,i} (\tilde{\Phi}_1 + \tilde{\Phi}_2) u_{R,i} + \text{h. c.} \right)
 \end{aligned}$$

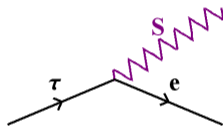
4.2 2HDM Types

One of the main features of the 2HDM is that allows Flavour Changing Neutral Currents (FCNC) at tree level:

Paschos-Glashow-Weinberg theorem allows to eliminate the FCNC at tree level

4.2 2HDM Types

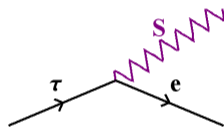
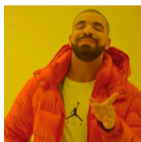
One of the main features of the 2HDM is that allows Flavour Changing Neutral Currents (FCNC) at tree level:



Paschos-Glashow-Weinberg theorem allows to eliminate the FCNC at tree level

4.2 2HDM Types

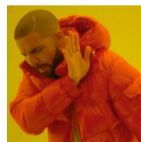
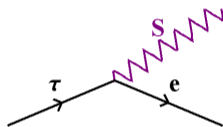
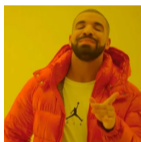
One of the main features of the 2HDM is that allows Flavour Changing Neutral Currents (FCNC) at tree level:



Paschos-Glashow-Weinberg theorem allows to eliminate the FCNC at tree level

4.2 2HDM Types

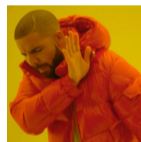
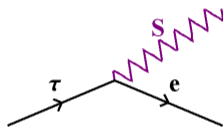
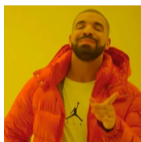
One of the main features of the 2HDM is that allows Flavour Changing Neutral Currents (FCNC) at tree level:



Paschos-Glashow-Weinberg theorem allows to eliminate the FCNC at tree level

4.2 2HDM Types

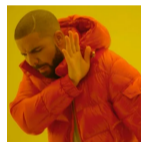
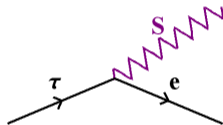
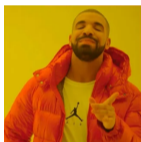
One of the main features of the 2HDM is that allows Flavour Changing Neutral Currents (FCNC) at tree level:



Paschos-Glashow-Weinberg theorem allows to eliminate the FCNC at tree level

4.2 2HDM Types

One of the main features of the 2HDM is that allows Flavour Changing Neutral Currents (FCNC) at tree level:

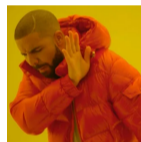
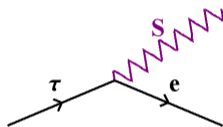
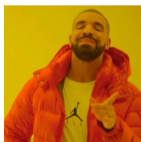


Paschos-Glashow-Weinberg theorem allows to eliminate the FCNC at tree level introducing a \mathbb{Z}_2 symmetry

$$\mathbb{Z}_2 : \Phi_1 \rightarrow -\Phi_1, \quad \Phi_2 \rightarrow \Phi_2$$

4.2 2HDM Types

One of the main features of the 2HDM is that allows Flavour Changing Neutral Currents (FCNC) at tree level:

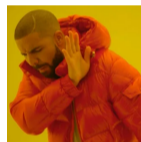
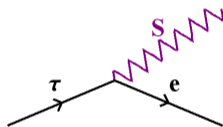
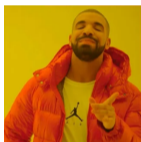


Paschos-Glashow-Weinberg theorem allows to eliminate the FCNC at tree level introducing a \mathbb{Z}_2 symmetry \rightarrow **this models are said to have Natural Flavour Conservation (NFC)**.

$$\mathbb{Z}_2 : \Phi_1 \rightarrow -\Phi_1, \quad \Phi_2 \rightarrow \Phi_2$$

4.2 2HDM Types

One of the main features of the 2HDM is that allows Flavour Changing Neutral Currents (FCNC) at tree level:



Paschos-Glashow-Weinberg theorem allows to eliminate the FCNC at tree level introducing a \mathbb{Z}_2 symmetry → **this models are said to have Natural Flavour Conservation (NFC)**.

$$\mathbb{Z}_2 : \Phi_1 \rightarrow -\Phi_1, \quad \Phi_2 \rightarrow \Phi_2$$

Model	u_R	d_R	ℓ_R
Type I	Φ_2	Φ_2	Φ_2
Type II	Φ_2	Φ_1	Φ_1
Lepton-specific (X)	Φ_2	Φ_2	Φ_1
Flipped (Y)	Φ_2	Φ_1	Φ_2

4.3 The Model

Denner (Nuclear Physics B, Volume 347, Issues 1–2, 24 December 1990, Pages 184-202).

Scalar sector

$$\begin{aligned}
 V(\Phi) = & m_{11}^2 \Phi_1^\dagger \Phi_1 + m_{22}^2 \Phi_2^\dagger \Phi_2 - (m_{12}^2 \Phi_1^\dagger \Phi_2 + \text{h. c.}) \\
 & + \frac{1}{2} \lambda_1 (\Phi_1^\dagger \Phi_1)^2 + \frac{1}{2} \lambda_2 (\Phi_2^\dagger \Phi_2)^2 + \lambda_3 (\Phi_1^\dagger \Phi_1) (\Phi_2^\dagger \Phi_2) + \lambda_4 (\Phi_1^\dagger \Phi_2) (\Phi_2^\dagger \Phi_1) \\
 & + \left(\frac{1}{2} \lambda_5 (\Phi_1^\dagger \Phi_2)^2 + \lambda_6 (\Phi_1^\dagger \Phi_1) (\Phi_1^\dagger \Phi_2) + \lambda_7 (\Phi_2^\dagger \Phi_2) (\Phi_1^\dagger \Phi_2) + \text{h. c.} \right);
 \end{aligned}$$

Higgs self-coupling conditions

$$\lambda_4 < 0, \quad \lambda_1 > 0, \quad \lambda_2 > 0, \quad \sqrt{\lambda_1 \lambda_2} + \lambda_3 + \lambda_4 > |\lambda_5|;$$

Yukawa sector

$$\mathcal{L}_Y = - \left(G^{(\ell)} \bar{L}_L \Phi_i \ell_R + G^{(d)} \bar{Q}_L \Phi_j d_R + G^{(u)} \bar{Q}_L \tilde{\Phi}_2 u_R \right) + \text{h. c.}$$

4.3 The Model

Denner ([Nuclear Physics B, Volume 347, Issues 1–2, 24 December 1990, Pages 184-202](#)).

Scalar sector

$$\begin{aligned}
 V(\Phi) = & m_{11}^2 \Phi_1^\dagger \Phi_1 + m_{22}^2 \Phi_2^\dagger \Phi_2 - (m_{12}^2 \Phi_1^\dagger \Phi_2 + \text{h. c.}) \\
 & + \frac{1}{2} \lambda_1 (\Phi_1^\dagger \Phi_1)^2 + \frac{1}{2} \lambda_2 (\Phi_2^\dagger \Phi_2)^2 + \lambda_3 (\Phi_1^\dagger \Phi_1) (\Phi_2^\dagger \Phi_2) + \lambda_4 (\Phi_1^\dagger \Phi_2) (\Phi_2^\dagger \Phi_1) \\
 & + \left(\frac{1}{2} \lambda_5 (\Phi_1^\dagger \Phi_2)^2 + \lambda_6 (\Phi_1^\dagger \Phi_1) (\Phi_1^\dagger \Phi_2) + \lambda_7 (\Phi_2^\dagger \Phi_2) (\Phi_1^\dagger \Phi_2) + \text{h. c.} \right);
 \end{aligned}$$

Higgs self-coupling conditions

$$\lambda_4 < 0, \quad \lambda_1 > 0, \quad \lambda_2 > 0, \quad \sqrt{\lambda_1 \lambda_2} + \lambda_3 + \lambda_4 > |\lambda_5|;$$

Yukawa sector

$$\mathcal{L}_Y = - \left(G^{(\ell)} \bar{L}_L \Phi_i \ell_R + G^{(d)} \bar{Q}_L \Phi_j d_R + G^{(u)} \bar{Q}_L \tilde{\Phi}_2 u_R \right) + \text{h. c.}$$

4.3 The Model

Denner ([Nuclear Physics B, Volume 347, Issues 1–2, 24 December 1990, Pages 184-202](#)).

Scalar sector

$$\begin{aligned}
 V(\Phi) = & m_{11}^2 \Phi_1^\dagger \Phi_1 + m_{22}^2 \Phi_2^\dagger \Phi_2 - (m_{12}^2 \Phi_1^\dagger \Phi_2 + \text{h. c.}) \\
 & + \frac{1}{2} \lambda_1 (\Phi_1^\dagger \Phi_1)^2 + \frac{1}{2} \lambda_2 (\Phi_2^\dagger \Phi_2)^2 + \lambda_3 (\Phi_1^\dagger \Phi_1) (\Phi_2^\dagger \Phi_2) + \lambda_4 (\Phi_1^\dagger \Phi_2) (\Phi_2^\dagger \Phi_1) \\
 & + \left(\frac{1}{2} \lambda_5 (\Phi_1^\dagger \Phi_2)^2 + \cancel{\lambda_6 (\Phi_1^\dagger \Phi_1) (\Phi_1^\dagger \Phi_2)} + \cancel{\lambda_7 (\Phi_2^\dagger \Phi_2) (\Phi_1^\dagger \Phi_2)} + \text{h. c.} \right);
 \end{aligned}$$

Higgs self-coupling conditions

$$\lambda_4 < 0, \quad \lambda_1 > 0, \quad \lambda_2 > 0, \quad \sqrt{\lambda_1 \lambda_2} + \lambda_3 + \lambda_4 > |\lambda_5|;$$

Yukawa sector

$$\mathcal{L}_Y = - \left(G^{(\ell)} \bar{L}_L \Phi_i \ell_R + G^{(d)} \bar{Q}_L \Phi_j d_R + G^{(u)} \bar{Q}_L \tilde{\Phi}_2 u_R \right) + \text{h. c.}$$

4.3 The Model

Denner ([Nuclear Physics B, Volume 347, Issues 1–2, 24 December 1990, Pages 184-202](#)).

Scalar sector

$$\begin{aligned}
 V(\Phi) = & m_{11}^2 \Phi_1^\dagger \Phi_1 + m_{22}^2 \Phi_2^\dagger \Phi_2 - \cancel{(m_{12}^2 \Phi_1^\dagger \Phi_2 + \text{h. c.})} \\
 & + \frac{1}{2} \lambda_1 (\Phi_1^\dagger \Phi_1)^2 + \frac{1}{2} \lambda_2 (\Phi_2^\dagger \Phi_2)^2 + \lambda_3 (\Phi_1^\dagger \Phi_1) (\Phi_2^\dagger \Phi_2) + \lambda_4 (\Phi_1^\dagger \Phi_2) (\Phi_2^\dagger \Phi_1) \\
 & + \left(\frac{1}{2} \lambda_5 (\Phi_1^\dagger \Phi_2)^2 + \cancel{\lambda_6 (\Phi_1^\dagger \Phi_1) (\Phi_1^\dagger \Phi_2)} + \cancel{\lambda_7 (\Phi_2^\dagger \Phi_2) (\Phi_1^\dagger \Phi_2)} + \text{h. c.} \right);
 \end{aligned}$$

Higgs self-coupling conditions

$$\lambda_4 < 0, \quad \lambda_1 > 0, \quad \lambda_2 > 0, \quad \sqrt{\lambda_1 \lambda_2} + \lambda_3 + \lambda_4 > |\lambda_5|;$$

Yukawa sector

$$\mathcal{L}_Y = - \left(G^{(\ell)} \bar{L}_L \Phi_i \ell_R + G^{(d)} \bar{Q}_L \Phi_j d_R + G^{(u)} \bar{Q}_L \tilde{\Phi}_2 u_R \right) + \text{h. c.}$$

4.3 The Model

Denner ([Nuclear Physics B, Volume 347, Issues 1–2, 24 December 1990, Pages 184-202](#)).

Scalar sector

$$\begin{aligned}
 V(\Phi) = & m_{11}^2 \Phi_1^\dagger \Phi_1 + m_{22}^2 \Phi_2^\dagger \Phi_2 + \frac{1}{2} \lambda_1 (\Phi_1^\dagger \Phi_1)^2 + \frac{1}{2} \lambda_2 (\Phi_2^\dagger \Phi_2)^2 \\
 & + \lambda_3 (\Phi_1^\dagger \Phi_1) (\Phi_2^\dagger \Phi_2) + \lambda_4 (\Phi_1^\dagger \Phi_2) (\Phi_2^\dagger \Phi_1) + \left(\frac{1}{2} \lambda_5 (\Phi_1^\dagger \Phi_2)^2 + \text{h. c.} \right);
 \end{aligned}$$

Higgs self-coupling conditions

$$\lambda_4 < 0, \quad \lambda_1 > 0, \quad \lambda_2 > 0, \quad \sqrt{\lambda_1 \lambda_2} + \lambda_3 + \lambda_4 > |\lambda_5|;$$

Yukawa sector

$$\mathcal{L}_Y = - \left(G^{(\ell)} \bar{L}_L \Phi_i \ell_R + G^{(d)} \bar{Q}_L \Phi_j d_R + G^{(u)} \bar{Q}_L \tilde{\Phi}_2 u_R \right) + \text{h. c.}$$

4.3 The Model

Denner ([Nuclear Physics B, Volume 347, Issues 1–2, 24 December 1990, Pages 184-202](#)).

Scalar sector

$$\begin{aligned}
 V(\Phi) = & m_{11}^2 \Phi_1^\dagger \Phi_1 + m_{22}^2 \Phi_2^\dagger \Phi_2 + \frac{1}{2} \lambda_1 (\Phi_1^\dagger \Phi_1)^2 + \frac{1}{2} \lambda_2 (\Phi_2^\dagger \Phi_2)^2 \\
 & + \lambda_3 (\Phi_1^\dagger \Phi_1)(\Phi_2^\dagger \Phi_2) + \lambda_4 (\Phi_1^\dagger \Phi_2)(\Phi_2^\dagger \Phi_1) + \left(\frac{1}{2} \lambda_5 (\Phi_1^\dagger \Phi_2)^2 + \text{h. c.} \right);
 \end{aligned}$$

Higgs self-coupling conditions

$$\lambda_4 < 0, \quad \lambda_1 > 0, \quad \lambda_2 > 0, \quad \sqrt{\lambda_1 \lambda_2} + \lambda_3 + \lambda_4 > |\lambda_5|;$$

Yukawa sector

$$\mathcal{L}_Y = - \left(G^{(\ell)} \bar{L}_L \Phi_i \ell_R + G^{(d)} \bar{Q}_L \Phi_j d_R + G^{(u)} \bar{Q}_L \tilde{\Phi}_2 u_R \right) + \text{h. c.}$$

4.3 The Model

Denner ([Nuclear Physics B, Volume 347, Issues 1–2, 24 December 1990, Pages 184-202](#)).

Scalar sector

$$\begin{aligned}
 V(\Phi) = & m_{11}^2 \Phi_1^\dagger \Phi_1 + m_{22}^2 \Phi_2^\dagger \Phi_2 + \frac{1}{2} \lambda_1 (\Phi_1^\dagger \Phi_1)^2 + \frac{1}{2} \lambda_2 (\Phi_2^\dagger \Phi_2)^2 \\
 & + \lambda_3 (\Phi_1^\dagger \Phi_1) (\Phi_2^\dagger \Phi_2) + \lambda_4 (\Phi_1^\dagger \Phi_2) (\Phi_2^\dagger \Phi_1) + \left(\frac{1}{2} \lambda_5 (\Phi_1^\dagger \Phi_2)^2 + \text{h. c.} \right);
 \end{aligned}$$

Higgs self-coupling conditions

$$\lambda_4 < 0, \quad \lambda_1 > 0, \quad \lambda_2 > 0, \quad \sqrt{\lambda_1 \lambda_2} + \lambda_3 + \lambda_4 > |\lambda_5|;$$

Yukawa sector

$$\mathcal{L}_Y = - \left(G^{(\ell)} \bar{L}_L \Phi_i \ell_R + G^{(d)} \bar{Q}_L \Phi_j d_R + G^{(u)} \bar{Q}_L \tilde{\Phi}_2 u_R \right) + \text{h. c.}$$

4.3 The Model

Masses of Particles

With the vacuum expectation values given as follows:

$$\langle \Phi_1 \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v_1 \end{pmatrix}, \quad \langle \Phi_2 \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v_2 \end{pmatrix};$$

four Higgs particles will be obtained, whose masses at the tree level are expressed as:

$$m_H^2 = \frac{1}{2}\eta_+, \quad m_h^2 = \frac{1}{2}\eta_-, \quad m_{H^\pm}^2 = -\frac{1}{2}(\lambda_4 + \lambda_5)v^2, \quad m_A^2 = -\lambda_5 v^2,$$

where

$$\eta_\pm = \lambda_1 v_1^2 + \lambda_2 v_2^2 \pm \sqrt{(\lambda_1 v_1^2 - \lambda_2 v_2^2)^2 + 4(\lambda_3 + \lambda_4 + \lambda_5)^2 v_1^2 v_2^2}.$$

For diagonal Yukawa matrices the fermion masses will be given as:

$$m_i^{(f)} = \frac{1}{\sqrt{2}} G_i^{(f)} v_j;$$

since fermions couples to only one of the Higgs.

4.3 The Model

Masses of Particles

With the vacuum expectation values given as follows:

$$\langle \Phi_1 \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v_1 \end{pmatrix}, \quad \langle \Phi_2 \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v_2 \end{pmatrix};$$

four Higgs particles will be obtained, whose masses at the tree level are expressed as:

$$m_H^2 = \frac{1}{2}\eta_+, \quad m_h^2 = \frac{1}{2}\eta_-, \quad m_{H^\pm}^2 = -\frac{1}{2}(\lambda_4 + \lambda_5)v^2, \quad m_A^2 = -\lambda_5 v^2,$$

where

$$\eta_\pm = \lambda_1 v_1^2 + \lambda_2 v_2^2 \pm \sqrt{(\lambda_1 v_1^2 - \lambda_2 v_2^2)^2 + 4(\lambda_3 + \lambda_4 + \lambda_5)^2 v_1^2 v_2^2}.$$

For diagonal Yukawa matrices the fermion masses will be given as:

$$m_i^{(f)} = \frac{1}{\sqrt{2}} G_i^{(f)} v_j;$$

since fermions couples to only one of the Higgs.

4.3 The Model

Masses of Particles

With the vacuum expectation values given as follows:

$$\langle \Phi_1 \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v_1 \end{pmatrix}, \quad \langle \Phi_2 \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v_2 \end{pmatrix};$$

four Higgs particles will be obtained, whose masses at the tree level are expressed as:

$$m_H^2 = \frac{1}{2}\eta_+, \quad m_h^2 = \frac{1}{2}\eta_-, \quad m_{H^\pm}^2 = -\frac{1}{2}(\lambda_4 + \lambda_5)v^2, \quad m_A^2 = -\lambda_5 v^2,$$

where

$$\eta_\pm = \lambda_1 v_1^2 + \lambda_2 v_2^2 \pm \sqrt{(\lambda_1 v_1^2 - \lambda_2 v_2^2)^2 + 4(\lambda_3 + \lambda_4 + \lambda_5)^2 v_1^2 v_2^2}.$$

For diagonal Yukawa matrices the fermion masses will be given as:

$$m_i^{(f)} = \frac{1}{\sqrt{2}} G_i^{(f)} v_j;$$

since fermions couples to only one of the Higgs.

4.3 The Model

Masses of Particles

With the vacuum expectation values given as follows:

$$\langle \Phi_1 \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v_1 \end{pmatrix}, \quad \langle \Phi_2 \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v_2 \end{pmatrix};$$

four Higgs particles will be obtained, whose masses at the tree level are expressed as:

$$m_H^2 = \frac{1}{2}\eta_+, \quad m_h^2 = \frac{1}{2}\eta_-, \quad m_{H^\pm}^2 = -\frac{1}{2}(\lambda_4 + \lambda_5)v^2, \quad m_A^2 = -\lambda_5 v^2,$$

where

$$\eta_\pm = \lambda_1 v_1^2 + \lambda_2 v_2^2 \pm \sqrt{(\lambda_1 v_1^2 - \lambda_2 v_2^2)^2 + 4(\lambda_3 + \lambda_4 + \lambda_5)^2 v_1^2 v_2^2}.$$

For diagonal Yukawa matrices the fermion masses will be given as:

$$m_i^{(f)} = \frac{1}{\sqrt{2}} G_i^{(f)} v_j;$$

since fermions couples to only one of the Higgs.

4.3 The Model

Masses of Particles

With the vacuum expectation values given as follows:

$$\langle \Phi_1 \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v_1 \end{pmatrix}, \quad \langle \Phi_2 \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v_2 \end{pmatrix};$$

four Higgs particles will be obtained, whose masses at the tree level are expressed as:

$$m_H^2 = \frac{1}{2}\eta_+, \quad m_h^2 = \frac{1}{2}\eta_-, \quad m_{H^\pm}^2 = -\frac{1}{2}(\lambda_4 + \lambda_5)v^2, \quad m_A^2 = -\lambda_5 v^2,$$

where

$$\eta_\pm = \lambda_1 v_1^2 + \lambda_2 v_2^2 \pm \sqrt{(\lambda_1 v_1^2 - \lambda_2 v_2^2)^2 + 4(\lambda_3 + \lambda_4 + \lambda_5)^2 v_1^2 v_2^2}.$$

For diagonal Yukawa matrices the fermion masses will be given as:

$$m_i^{(f)} = \frac{1}{\sqrt{2}} G_i^{(f)} v_j;$$

since fermions couples to only one of the Higgs.

4.3 The Model

Masses of Particles

With the vacuum expectation values given as follows:

$$\langle \Phi_1 \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v_1 \end{pmatrix}, \quad \langle \Phi_2 \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v_2 \end{pmatrix};$$

four Higgs particles will be obtained, whose masses at the tree level are expressed as:

$$m_H^2 = \frac{1}{2}\eta_+, \quad m_h^2 = \frac{1}{2}\eta_-, \quad m_{H^\pm}^2 = -\frac{1}{2}(\lambda_4 + \lambda_5)v^2, \quad m_A^2 = -\lambda_5 v^2,$$

where

$$\eta_\pm = \lambda_1 v_1^2 + \lambda_2 v_2^2 \pm \sqrt{(\lambda_1 v_1^2 - \lambda_2 v_2^2)^2 + 4(\lambda_3 + \lambda_4 + \lambda_5)^2 v_1^2 v_2^2}.$$

For diagonal Yukawa matrices the fermion masses will be given as:

$$m_i^{(f)} = \frac{1}{\sqrt{2}} G_i^{(f)} v_j;$$

since fermions couples to only one of the Higgs.

V. Reduction of Couplings in 2HDM with NFC

Renormalization Groups Equations

$$16\pi^2\beta_{g_1} = 7g_1^3,$$

$$16\pi^2\beta_{g_2} = -3g_2^3,$$

$$16\pi^2\beta_{g_3} = -7g_3^3.$$

$$16\pi^2\beta_{G_j^{(\ell)}} = G_j^{(\ell)} \left(- \left(\frac{15}{4}g_1^2 + \frac{9}{4}g_2^2 \right) + \frac{3}{2}G_j^{(\ell)2} + \sum_k \left(a_1 G_k^{(\ell)2} + 3a_2 G_k^{(d)2} + 3a_3 G_k^{(u)2} \right) \right),$$

$$16\pi^2\beta_{G_j^{(d)}} = G_j^{(d)} \left(- \left(\frac{5}{12}g_1^2 + \frac{9}{4}g_2^2 + 8g_3^2 \right) + \frac{3}{2}G_j^{(d)2} + a_4 G_j^{(u)2} + \sum_k \left(a_5 G_k^{(\ell)2} + 3a_6 G_k^{(d)2} + 3a_7 G_k^{(u)2} \right) \right),$$

$$16\pi^2\beta_{G_j^{(u)}} = G_j^{(u)} \left(- \left(\frac{17}{12}g_1^2 + \frac{9}{4}g_2^2 + 8g_3^2 \right) + \frac{3}{2}G_j^{(u)2} + a_8 G_j^{(d)2} + \sum_k \left(a_9 G_k^{(\ell)2} + 3a_{10} G_k^{(d)2} + 3a_{11} G_k^{(u)2} \right) \right);$$

$$16\pi^2\beta_{\lambda_1} = \frac{3}{4}g_1^4 + \frac{3}{2}g_1^2g_2^2 + \frac{9}{4}g_2^4 - 3g_1^2\lambda_1 - 9g_2^2\lambda_1 + 12\lambda_1^2 + 4\lambda_3^2 + 4\lambda_3\lambda_4 + 2\lambda_4^2 + 2\lambda_5^2$$

$$16\pi^2\beta_{\lambda_2} = \frac{3}{4}g_1^4 + \frac{3}{2}g_1^2g_2^2 + \frac{9}{4}g_2^4 - 3g_1^2\lambda_2 - 9g_2^2\lambda_2 + 12\lambda_2^2 + 4\lambda_3^2 + 4\lambda_3\lambda_4 + 2\lambda_4^2 + 2\lambda_5^2 + \left(-12G_b^4 - 12G_t^4 - 4G_\tau^4 + (12G_b^2 + 12G_t^2 + 4G_\tau^2)\lambda_2 \right)$$

$$16\pi^2\beta_{\lambda_3} = \frac{3}{4}g_1^4 - \frac{3}{2}g_1^2g_2^2 + \frac{9}{4}g_2^4 - 3g_1^2\lambda_3 - 9g_2^2\lambda_3 + (\lambda_1 + \lambda_2)(6\lambda_3 + 2\lambda_4) + 4\lambda_3^2 + 2\lambda_4^2 + 2\lambda_5^2 + (6G_b^2 + 6G_t^2 + 2G_\tau^2)\lambda_3$$

$$16\pi^2\beta_{\lambda_4} = 3g_1^2g_2^2 - (3g_1^2 + 9g_2^2)\lambda_4 + 2\lambda_1\lambda_4 + 2\lambda_2\lambda_4 + 8\lambda_3\lambda_4 + 4\lambda_4^2 + 8\lambda_5^2 + (6G_b^2 + 6G_t^2 + 2G_\tau^2)\lambda_4$$

$$16\pi^2\beta_{\lambda_5} = \lambda_5 \left(-3g_1^2 - 9g_2^2 + 2\lambda_1 + 2\lambda_2 + 8\lambda_3 + 12\lambda_4 + (6G_b^2 + 6G_t^2 + 2G_\tau^2) \right)$$

V. Reduction of Couplings in 2HDM with NFC

Renormalization Groups Equations

$$16\pi^2 \beta_{g_1} = 7g_1^3,$$

$$16\pi^2 \beta_{g_2} = -3g_2^3,$$

$$16\pi^2 \beta_{g_3} = -7g_3^3.$$

$$16\pi^2 \beta_{G_j^{(\ell)}} = G_j^{(\ell)} \left(- \left(\frac{15}{4} g_1^2 + \frac{9}{4} g_2^2 \right) + \frac{3}{2} G_j^{(\ell)2} + \sum_k \left(a_1 G_k^{(\ell)2} + 3a_2 G_k^{(d)2} + 3a_3 G_k^{(u)2} \right) \right),$$

$$16\pi^2 \beta_{G_j^{(d)}} = G_j^{(d)} \left(- \left(\frac{5}{12} g_1^2 + \frac{9}{4} g_2^2 + 8g_3^2 \right) + \frac{3}{2} G_j^{(d)2} + a_4 G_j^{(u)2} + \sum_k \left(a_5 G_k^{(\ell)2} + 3a_6 G_k^{(d)2} + 3a_7 G_k^{(u)2} \right) \right),$$

$$16\pi^2 \beta_{G_j^{(u)}} = G_j^{(u)} \left(- \left(\frac{17}{12} g_1^2 + \frac{9}{4} g_2^2 + 8g_3^2 \right) + \frac{3}{2} G_j^{(u)2} + a_8 G_j^{(d)2} + \sum_k \left(a_9 G_k^{(\ell)2} + 3a_{10} G_k^{(d)2} + 3a_{11} G_k^{(u)2} \right) \right);$$

$$16\pi^2 \beta_{\lambda_1} = \frac{3}{4} g_1^4 + \frac{3}{2} g_1^2 g_2^2 + \frac{9}{4} g_2^4 - 3g_1^2 \lambda_1 - 9g_2^2 \lambda_1 + 12\lambda_1^2 + 4\lambda_3^2 + 4\lambda_3 \lambda_4 + 2\lambda_4^2 + 2\lambda_5^2$$

$$16\pi^2 \beta_{\lambda_2} = \frac{3}{4} g_1^4 + \frac{3}{2} g_1^2 g_2^2 + \frac{9}{4} g_2^4 - 3g_1^2 \lambda_2 - 9g_2^2 \lambda_2 + 12\lambda_2^2 + 4\lambda_3^2 + 4\lambda_3 \lambda_4 + 2\lambda_4^2 + 2\lambda_5^2 + \left(-12G_b^4 - 12G_t^4 - 4G_\tau^4 + (12G_b^2 + 12G_t^2 + 4G_\tau^2) \lambda_2 \right)$$

$$16\pi^2 \beta_{\lambda_3} = \frac{3}{4} g_1^4 - \frac{3}{2} g_1^2 g_2^2 + \frac{9}{4} g_2^4 - 3g_1^2 \lambda_3 - 9g_2^2 \lambda_3 + (\lambda_1 + \lambda_2)(6\lambda_3 + 2\lambda_4) + 4\lambda_3^2 + 2\lambda_4^2 + 2\lambda_5^2 + (6G_b^2 + 6G_t^2 + 2G_\tau^2) \lambda_3$$

$$16\pi^2 \beta_{\lambda_4} = 3g_1^2 g_2^2 - (3g_1^2 + 9g_2^2) \lambda_4 + 2\lambda_1 \lambda_4 + 2\lambda_2 \lambda_4 + 8\lambda_3 \lambda_4 + 4\lambda_4^2 + 8\lambda_5^2 + (6G_b^2 + 6G_t^2 + 2G_\tau^2) \lambda_4$$

$$16\pi^2 \beta_{\lambda_5} = \lambda_5 \left(-3g_1^2 - 9g_2^2 + 2\lambda_1 + 2\lambda_2 + 8\lambda_3 + 12\lambda_4 + (6G_b^2 + 6G_t^2 + 2G_\tau^2) \right)$$

V. Reduction of Couplings in 2HDM with NFC

Renormalization Groups Equations

$$16\pi^2 \beta_{g_1} = 7g_1^3,$$

$$16\pi^2 \beta_{g_2} = -3g_2^3,$$

$$16\pi^2 \beta_{g_3} = -7g_3^3.$$

$$16\pi^2 \beta_{G_j^{(\ell)}} = G_j^{(\ell)} \left(- \left(\frac{15}{4} g_1^2 + \frac{9}{4} g_2^2 \right) + \frac{3}{2} G_j^{(\ell)2} + \sum_k \left(a_1 G_k^{(\ell)2} + 3a_2 G_k^{(d)2} + 3a_3 G_k^{(u)2} \right) \right),$$

$$16\pi^2 \beta_{G_j^{(d)}} = G_j^{(d)} \left(- \left(\frac{5}{12} g_1^2 + \frac{9}{4} g_2^2 + 8g_3^2 \right) + \frac{3}{2} G_j^{(d)2} + a_4 G_j^{(u)2} + \sum_k \left(a_5 G_k^{(\ell)2} + 3a_6 G_k^{(d)2} + 3a_7 G_k^{(u)2} \right) \right),$$

$$16\pi^2 \beta_{G_j^{(u)}} = G_j^{(u)} \left(- \left(\frac{17}{12} g_1^2 + \frac{9}{4} g_2^2 + 8g_3^2 \right) + \frac{3}{2} G_j^{(u)2} + a_8 G_j^{(d)2} + \sum_k \left(a_9 G_k^{(\ell)2} + 3a_{10} G_k^{(d)2} + 3a_{11} G_k^{(u)2} \right) \right);$$

$$16\pi^2 \beta_{\lambda_1} = \frac{3}{4} g_1^4 + \frac{3}{2} g_1^2 g_2^2 + \frac{9}{4} g_2^4 - 3g_1^2 \lambda_1 - 9g_2^2 \lambda_1 + 12\lambda_1^2 + 4\lambda_3^2 + 4\lambda_3 \lambda_4 + 2\lambda_4^2 + 2\lambda_5^2$$

$$16\pi^2 \beta_{\lambda_2} = \frac{3}{4} g_1^4 + \frac{3}{2} g_1^2 g_2^2 + \frac{9}{4} g_2^4 - 3g_1^2 \lambda_2 - 9g_2^2 \lambda_2 + 12\lambda_2^2 + 4\lambda_3^2 + 4\lambda_3 \lambda_4 + 2\lambda_4^2 + 2\lambda_5^2 + \left(-12G_b^4 - 12G_t^4 - 4G_\tau^4 + (12G_b^2 + 12G_t^2 + 4G_\tau^2) \lambda_2 \right)$$

$$16\pi^2 \beta_{\lambda_3} = \frac{3}{4} g_1^4 - \frac{3}{2} g_1^2 g_2^2 + \frac{9}{4} g_2^4 - 3g_1^2 \lambda_3 - 9g_2^2 \lambda_3 + (\lambda_1 + \lambda_2)(6\lambda_3 + 2\lambda_4) + 4\lambda_3^2 + 2\lambda_4^2 + 2\lambda_5^2 + (6G_b^2 + 6G_t^2 + 2G_\tau^2) \lambda_3$$

$$16\pi^2 \beta_{\lambda_4} = 3g_1^2 g_2^2 - (3g_1^2 + 9g_2^2) \lambda_4 + 2\lambda_1 \lambda_4 + 2\lambda_2 \lambda_4 + 8\lambda_3 \lambda_4 + 4\lambda_4^2 + 8\lambda_5^2 + (6G_b^2 + 6G_t^2 + 2G_\tau^2) \lambda_4$$

$$16\pi^2 \beta_{\lambda_5} = \lambda_5 \left(-3g_1^2 - 9g_2^2 + 2\lambda_1 + 2\lambda_2 + 8\lambda_3 + 12\lambda_4 + (6G_b^2 + 6G_t^2 + 2G_\tau^2) \right)$$

V. Reduction of Couplings in 2HDM with NFC

Renormalization Groups Equations

$$16\pi^2 \beta_{g_1} = 7g_1^3,$$

$$16\pi^2 \beta_{g_2} = -3g_2^3,$$

$$16\pi^2 \beta_{g_3} = -7g_3^3.$$

$$16\pi^2 \beta_{G_j^{(\ell)}} = G_j^{(\ell)} \left(- \left(\frac{15}{4} g_1^2 + \frac{9}{4} g_2^2 \right) + \frac{3}{2} G_j^{(\ell)2} + \sum_k \left(a_1 G_k^{(\ell)2} + 3a_2 G_k^{(d)2} + 3a_3 G_k^{(u)2} \right) \right),$$

$$16\pi^2 \beta_{G_j^{(d)}} = G_j^{(d)} \left(- \left(\frac{5}{12} g_1^2 + \frac{9}{4} g_2^2 + 8g_3^2 \right) + \frac{3}{2} G_j^{(d)2} + a_4 G_j^{(u)2} + \sum_k \left(a_5 G_k^{(\ell)2} + 3a_6 G_k^{(d)2} + 3a_7 G_k^{(u)2} \right) \right),$$

$$16\pi^2 \beta_{G_j^{(u)}} = G_j^{(u)} \left(- \left(\frac{17}{12} g_1^2 + \frac{9}{4} g_2^2 + 8g_3^2 \right) + \frac{3}{2} G_j^{(u)2} + a_8 G_j^{(d)2} + \sum_k \left(a_9 G_k^{(\ell)2} + 3a_{10} G_k^{(d)2} + 3a_{11} G_k^{(u)2} \right) \right);$$

$$16\pi^2 \beta_{\lambda_1} = \frac{3}{4} g_1^4 + \frac{3}{2} g_1^2 g_2^2 + \frac{9}{4} g_2^4 - 3g_1^2 \lambda_1 - 9g_2^2 \lambda_1 + 12\lambda_1^2 + 4\lambda_3^2 + 4\lambda_3 \lambda_4 + 2\lambda_4^2 + 2\lambda_5^2$$

$$16\pi^2 \beta_{\lambda_2} = \frac{3}{4} g_1^4 + \frac{3}{2} g_1^2 g_2^2 + \frac{9}{4} g_2^4 - 3g_1^2 \lambda_2 - 9g_2^2 \lambda_2 + 12\lambda_2^2 + 4\lambda_3^2 + 4\lambda_3 \lambda_4 + 2\lambda_4^2 + 2\lambda_5^2 + \left(-12G_b^4 - 12G_t^4 - 4G_\tau^4 + (12G_b^2 + 12G_t^2 + 4G_\tau^2) \lambda_2 \right)$$

$$16\pi^2 \beta_{\lambda_3} = \frac{3}{4} g_1^4 - \frac{3}{2} g_1^2 g_2^2 + \frac{9}{4} g_2^4 - 3g_1^2 \lambda_3 - 9g_2^2 \lambda_3 + (\lambda_1 + \lambda_2)(6\lambda_3 + 2\lambda_4) + 4\lambda_3^2 + 2\lambda_4^2 + 2\lambda_5^2 + (6G_b^2 + 6G_t^2 + 2G_\tau^2) \lambda_3$$

$$16\pi^2 \beta_{\lambda_4} = 3g_1^2 g_2^2 - (3g_1^2 + 9g_2^2) \lambda_4 + 2\lambda_1 \lambda_4 + 2\lambda_2 \lambda_4 + 8\lambda_3 \lambda_4 + 4\lambda_4^2 + 8\lambda_5^2 + (6G_b^2 + 6G_t^2 + 2G_\tau^2) \lambda_4$$

$$16\pi^2 \beta_{\lambda_5} = \lambda_5 \left(-3g_1^2 - 9g_2^2 + 2\lambda_1 + 2\lambda_2 + 8\lambda_3 + 12\lambda_4 + (6G_b^2 + 6G_t^2 + 2G_\tau^2) \right)$$

V. Reduction of Couplings in 2HDM with NFC

Renormalization Groups Equations

$$16\pi^2 \beta_{g_1} = 7g_1^3,$$

$$16\pi^2 \beta_{g_2} = -3g_2^3,$$

$$16\pi^2 \beta_{g_3} = -7g_3^3.$$

$$16\pi^2 \beta_{G_j^{(\ell)}} = G_j^{(\ell)} \left(- \left(\frac{15}{4} g_1^2 + \frac{9}{4} g_2^2 \right) + \frac{3}{2} G_j^{(\ell)2} + \sum_k \left(a_1 G_k^{(\ell)2} + 3a_2 G_k^{(d)2} + 3a_3 G_k^{(u)2} \right) \right),$$

$$16\pi^2 \beta_{G_j^{(d)}} = G_j^{(d)} \left(- \left(\frac{5}{12} g_1^2 + \frac{9}{4} g_2^2 + 8g_3^2 \right) + \frac{3}{2} G_j^{(d)2} + a_4 G_j^{(u)2} + \sum_k \left(a_5 G_k^{(\ell)2} + 3a_6 G_k^{(d)2} + 3a_7 G_k^{(u)2} \right) \right),$$

$$16\pi^2 \beta_{G_j^{(u)}} = G_j^{(u)} \left(- \left(\frac{17}{12} g_1^2 + \frac{9}{4} g_2^2 + 8g_3^2 \right) + \frac{3}{2} G_j^{(u)2} + a_8 G_j^{(d)2} + \sum_k \left(a_9 G_k^{(\ell)2} + 3a_{10} G_k^{(d)2} + 3a_{11} G_k^{(u)2} \right) \right);$$

$$16\pi^2 \beta_{\lambda_1} = \frac{3}{4} g_1^4 + \frac{3}{2} g_1^2 g_2^2 + \frac{9}{4} g_2^4 - 3g_1^2 \lambda_1 - 9g_2^2 \lambda_1 + 12\lambda_1^2 + 4\lambda_3^2 + 4\lambda_3 \lambda_4 + 2\lambda_4^2 + 2\lambda_5^2$$

$$16\pi^2 \beta_{\lambda_2} = \frac{3}{4} g_1^4 + \frac{3}{2} g_1^2 g_2^2 + \frac{9}{4} g_2^4 - 3g_1^2 \lambda_2 - 9g_2^2 \lambda_2 + 12\lambda_2^2 + 4\lambda_3^2 + 4\lambda_3 \lambda_4 + 2\lambda_4^2 + 2\lambda_5^2 + \left(-12G_b^4 - 12G_t^4 - 4G_\tau^4 + (12G_b^2 + 12G_t^2 + 4G_\tau^2) \lambda_2 \right)$$

$$16\pi^2 \beta_{\lambda_3} = \frac{3}{4} g_1^4 - \frac{3}{2} g_1^2 g_2^2 + \frac{9}{4} g_2^4 - 3g_1^2 \lambda_3 - 9g_2^2 \lambda_3 + (\lambda_1 + \lambda_2)(6\lambda_3 + 2\lambda_4) + 4\lambda_3^2 + 2\lambda_4^2 + 2\lambda_5^2 + (6G_b^2 + 6G_t^2 + 2G_\tau^2) \lambda_3$$

$$16\pi^2 \beta_{\lambda_4} = 3g_1^2 g_2^2 - (3g_1^2 + 9g_2^2) \lambda_4 + 2\lambda_1 \lambda_4 + 2\lambda_2 \lambda_4 + 8\lambda_3 \lambda_4 + 4\lambda_4^2 + 8\lambda_5^2 + (6G_b^2 + 6G_t^2 + 2G_\tau^2) \lambda_4$$

$$16\pi^2 \beta_{\lambda_5} = \lambda_5 \left(-3g_1^2 - 9g_2^2 + 2\lambda_1 + 2\lambda_2 + 8\lambda_3 + 12\lambda_4 + (6G_b^2 + 6G_t^2 + 2G_\tau^2) \right)$$

V. Reduction of Couplings in 2HDM with NFC

Renormalization Groups Equations

$$\begin{aligned}
 16\pi^2 \beta_{g_1} &= 7g_1^3, \\
 16\pi^2 \beta_{g_2} &= -3g_2^3, \\
 16\pi^2 \beta_{g_3} &= -7g_3^3.
 \end{aligned}
 \qquad
 \beta_0 \frac{d\lambda_j}{d\lambda_0} = \beta_j
 \qquad
 \begin{aligned}
 16\pi^2 \beta_{\lambda_1} &= \frac{3}{4}g_1^4 + \frac{3}{2}g_1^2g_2^2 + \frac{9}{4}g_2^4 - 3g_1^2\lambda_1 - 9g_2^2\lambda_1 + 12\lambda_1^2 + 4\lambda_3^2 + 4\lambda_3\lambda_4 \\
 &\quad + 2\lambda_4^2 + 2\lambda_5^2 \\
 16\pi^2 \beta_{\lambda_2} &= \frac{3}{4}g_1^4 + \frac{3}{2}g_1^2g_2^2 + \frac{9}{4}g_2^4 - 3g_1^2\lambda_2 - 9g_2^2\lambda_2 + 12\lambda_2^2 + 4\lambda_3^2 + 4\lambda_3\lambda_4 \\
 &\quad + 2\lambda_4^2 + 2\lambda_5^2 + (-12G_b^4 - 12G_t^4 - 4G_\tau^4 + (12G_b^2 + 12G_t^2 + 4G_\tau^2)\lambda_2) \\
 16\pi^2 \beta_{\lambda_3} &= \frac{3}{4}g_1^4 - \frac{3}{2}g_1^2g_2^2 + \frac{9}{4}g_2^4 - 3g_1^2\lambda_3 - 9g_2^2\lambda_3 + (\lambda_1 + \lambda_2)(6\lambda_3 + 2\lambda_4) \\
 &\quad + 4\lambda_3^2 + 2\lambda_4^2 + 2\lambda_5^2 + (6G_b^2 + 6G_t^2 + 2G_\tau^2)\lambda_3 \\
 16\pi^2 \beta_{\lambda_4} &= 3g_1^2g_2^2 - (3g_1^2 + 9g_2^2)\lambda_4 + 2\lambda_1\lambda_4 + 2\lambda_2\lambda_4 + 8\lambda_3\lambda_4 + 4\lambda_4^2 + 8\lambda_5^2 \\
 &\quad + (6G_b^2 + 6G_t^2 + 2G_\tau^2)\lambda_4 \\
 16\pi^2 \beta_{\lambda_5} &= \lambda_5(-3g_1^2 - 9g_2^2 + 2\lambda_1 + 2\lambda_2 + 8\lambda_3 + 12\lambda_4 \\
 &\quad + (6G_b^2 + 6G_t^2 + 2G_\tau^2))
 \end{aligned}$$

$$\begin{aligned}
 16\pi^2 \beta_{G_j^{(\ell)}} &= G_j^{(\ell)} \left(- \left(\frac{15}{4}g_1^2 + \frac{9}{4}g_2^2 \right) + \frac{3}{2}G_j^{(\ell)2} \right. \\
 &\quad \left. + \sum_k (a_1 G_k^{(\ell)2} + 3a_2 G_k^{(d)2} + 3a_3 G_k^{(u)2}) \right), \\
 16\pi^2 \beta_{G_j^{(d)}} &= G_j^{(d)} \left(- \left(\frac{5}{12}g_1^2 + \frac{9}{4}g_2^2 + 8g_3^2 \right) + \frac{3}{2}G_j^{(d)2} \right. \\
 &\quad \left. + a_4 G_j^{(u)2} + \sum_k (a_5 G_k^{(\ell)2} + 3a_6 G_k^{(d)2} + 3a_7 G_k^{(u)2}) \right), \\
 16\pi^2 \beta_{G_j^{(u)}} &= G_j^{(u)} \left(- \left(\frac{17}{12}g_1^2 + \frac{9}{4}g_2^2 + 8g_3^2 \right) + \frac{3}{2}G_j^{(u)2} \right. \\
 &\quad \left. + a_8 G_j^{(d)2} + \sum_k (a_9 G_k^{(\ell)2} + 3a_{10} G_k^{(d)2} + 3a_{11} G_k^{(u)2}) \right);
 \end{aligned}$$

V. Reduction of Couplings in 2HDM with NFC

Renormalization Groups Equations

$$16\pi^2 \beta_{g_1} = 7g_1^3,$$

$$16\pi^2 \beta_{g_2} = -3g_2^3,$$

$$16\pi^2 \beta_{g_3} = -7g_3^3.$$

$$\beta_0 \frac{d\lambda_j}{d\lambda_0} = \beta_j$$

$$16\pi^2 \beta_{G_j^{(\ell)}} = G_j^{(\ell)} \left(- \left(\frac{15}{4} g_1^2 + \frac{9}{4} g_2^2 \right) + \frac{3}{2} G_j^{(\ell)2} + \sum_k \left(a_1 G_k^{(\ell)2} + 3a_2 G_k^{(d)2} + 3a_3 G_k^{(u)2} \right) \right) + \left(12G_b^4 - 12G_t^4 - 4G_\tau^4 + (12G_b^2 + 12G_t^2 + 4G_\tau^2)\lambda_2 \right)$$

$$16\pi^2 \beta_{G_j^{(d)}} = G_j^{(d)} \left(- \left(\frac{5}{12} g_1^2 + \frac{9}{4} g_2^2 + 8g_3^2 \right) + \frac{3}{2} G_j^{(d)2} + a_4 G_j^{(u)2} + \sum_k \left(a_5 G_k^{(\ell)2} + 3a_6 G_k^{(d)2} + 3a_7 G_k^{(u)2} \right) \right),$$

$$16\pi^2 \beta_{G_j^{(u)}} = G_j^{(u)} \left(- \left(\frac{17}{12} g_1^2 + \frac{9}{4} g_2^2 + 8g_3^2 \right) + \frac{3}{2} G_j^{(u)2} + a_8 G_j^{(d)2} + \sum_k \left(a_9 G_k^{(\ell)2} + 3a_{10} G_k^{(d)2} + 3a_{11} G_k^{(u)2} \right) \right);$$

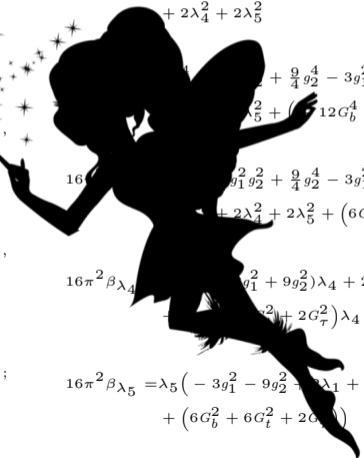
$$16\pi^2 \beta_{\lambda_1} = \frac{3}{4} g_1^4 + \frac{3}{2} g_1^2 g_2^2 + \frac{9}{4} g_2^4 - 3g_1^2 \lambda_1 - 9g_2^2 \lambda_1 + 12\lambda_1^2 + 4\lambda_3^2 + 4\lambda_3 \lambda_4 + 2\lambda_4^2 + 2\lambda_5^2$$

$$16\pi^2 \beta_{\lambda_2} = \frac{3}{4} g_1^4 + \frac{3}{2} g_1^2 g_2^2 + \frac{9}{4} g_2^4 - 3g_1^2 \lambda_2 - 9g_2^2 \lambda_2 + 12\lambda_2^2 + 4\lambda_3^2 + 4\lambda_3 \lambda_4 + 2\lambda_4^2 + 2\lambda_5^2 + (6G_b^2 + 6G_t^2 + 2G_\tau^2)\lambda_3$$

$$16\pi^2 \beta_{\lambda_3} = \frac{3}{4} g_1^4 + \frac{3}{2} g_1^2 g_2^2 + \frac{9}{4} g_2^4 - 3g_1^2 \lambda_3 - 9g_2^2 \lambda_3 + (\lambda_1 + \lambda_2)(6\lambda_3 + 2\lambda_4) + 2\lambda_4^2 + 2\lambda_5^2 + (6G_b^2 + 6G_t^2 + 2G_\tau^2)\lambda_3$$

$$16\pi^2 \beta_{\lambda_4} = \frac{3}{4} g_1^4 + \frac{3}{2} g_1^2 g_2^2 + \frac{9}{4} g_2^4 - 3g_1^2 \lambda_4 - 9g_2^2 \lambda_4 + 2\lambda_1 \lambda_4 + 2\lambda_2 \lambda_4 + 8\lambda_3 \lambda_4 + 4\lambda_4^2 + 8\lambda_5^2 + (6G_b^2 + 6G_t^2 + 2G_\tau^2)\lambda_4$$

$$16\pi^2 \beta_{\lambda_5} = \lambda_5 \left(-3g_1^2 - 9g_2^2 - 2\lambda_1 + 2\lambda_2 + 8\lambda_3 + 12\lambda_4 + (6G_b^2 + 6G_t^2 + 2G_\tau^2)\lambda_4 \right)$$



V. Reduction of Couplings in 2HDM with NFC

Reduction Equations

$$-14x \frac{du}{dx} = -6u^2 + 14u,$$

$$-14x \frac{dv}{dx} = 14v^2 + 14v,$$

$$-14x \frac{d\rho_{\ell_j}}{dx} = 2\rho_{\ell_j} \left(-\left(\frac{15}{4}v + \frac{9}{4}u\right) + \frac{3}{2}\rho_{\ell_j} + \sum_k (\rho_{\ell_k} + 3\rho_{d_k} + 3\rho_{u_k}) + 7 \right),$$

$$-14x \frac{d\rho_{d_j}}{dx} = 2\rho_{d_j} \left(-\left(\frac{5}{12}v + \frac{9}{4}u\right) + \frac{3}{2}\rho_{d_j} - \frac{3}{2}\rho_{u_j} + \sum_k (\rho_{\ell_k} + 3\rho_{d_k} + 3\rho_{u_k}) - 1 \right),$$

$$-14x \frac{d\rho_{u_j}}{dx} = 2\rho_{u_j} \left(-\left(\frac{17}{12}v + \frac{9}{4}u\right) + \frac{3}{2}\rho_{u_j} - \frac{3}{2}\rho_{d_j} + \sum_k (\rho_{\ell_k} + 3\rho_{d_k} + 3\rho_{u_k}) - 1 \right),$$

$$-14x \frac{d\rho_1}{dx} = \frac{3}{4}v^2 + \frac{3}{2}vu + \frac{9}{4}u^2 - 3v\rho_1 - 9u\rho_1 + 12\rho_1^2 + 4\rho_3^2 + 4\rho_3\rho_4 + 2\rho_4^2 + 2\rho_5^2 + 14\rho_1,$$

$$-14x \frac{d\rho_2}{dx} = \frac{3}{4}v^2 + \frac{3}{2}vu + \frac{9}{4}u^2 - 3v\rho_2 - 9u\rho_2 + 12\rho_2^2 + 4\rho_3^2 + 4\rho_3\rho_4 + 2\rho_4^2 + 2\rho_5^2 + 14\rho_2 + 4\rho_2 \sum_k (\rho_{\ell_k} + 3\rho_{d_k} + 3\rho_{u_k}) - 4 \sum_k (\rho_{\ell_k}^2 + 3\rho_{d_k}^2 + 3\rho_{u_k}^2),$$

$$-14x \frac{d\rho_3}{dx} = (\rho_1 + \rho_2)(6\rho_3 + 2\rho_4) + 4\rho_3^2 + 2\rho_4^2 + 2\rho_5^2 + 14\rho_3 + 2\rho_3 \sum_k (\rho_{\ell_k} + 3\rho_{d_k} + 3\rho_{u_k}),$$

$$-14x \frac{d\rho_4}{dx} = 2\rho_1\rho_4 + 2\rho_2\rho_4 + 8\rho_3\rho_4 + 4\rho_4^2 + 8\rho_5^2 + 14\rho_4 + 2\rho_4 \sum_k (\rho_{\ell_k} + 3\rho_{d_k} + 3\rho_{u_k}),$$

$$-14x \frac{d\rho_5}{dx} = (2\rho_1 + 2\rho_2 + 8\rho_3 + 12\rho_4 + 14 + 2 \sum_k (\rho_{\ell_k} + 3\rho_{d_k} + 3\rho_{u_k}))\rho_5.$$

V. Reduction of Couplings in 2HDM with NFC

Reduction Equations

$$-14x \frac{du}{dx} = -6u^2 + 14u,$$

$$-14x \frac{dv}{dx} = 14v^2 + 14v,$$

$$-14x \frac{d\rho_{\ell_j}}{dx} = 2\rho_{\ell_j} \left(-\left(\frac{15}{4}v + \frac{9}{4}u\right) + \frac{3}{2}\rho_{\ell_j} + \sum_k (\rho_{\ell_k} + 3\rho_{d_k} + 3\rho_{u_k}) + 7 \right),$$

$$-14x \frac{d\rho_{d_j}}{dx} = 2\rho_{d_j} \left(-\left(\frac{5}{12}v + \frac{9}{4}u\right) + \frac{3}{2}\rho_{d_j} - \frac{3}{2}\rho_{u_j} + \sum_k (\rho_{\ell_k} + 3\rho_{d_k} + 3\rho_{u_k}) - 1 \right),$$

$$-14x \frac{d\rho_{u_j}}{dx} = 2\rho_{u_j} \left(-\left(\frac{17}{12}v + \frac{9}{4}u\right) + \frac{3}{2}\rho_{u_j} - \frac{3}{2}\rho_{d_j} + \sum_k (\rho_{\ell_k} + 3\rho_{d_k} + 3\rho_{u_k}) - 1 \right),$$

$$x = \frac{g_2^2}{4\pi},$$

$$u = \frac{g_2^2}{g_3^2}$$

$$v = \frac{g_1^2}{g_3^2},$$

$$\rho_i^{(f)} = \frac{G_i^{(f)2}}{g_3^2},$$

$$\rho_i = \frac{\lambda_i}{g_3^2};$$

$$-14x \frac{d\rho_1}{dx} = \frac{3}{4}v^2 + \frac{3}{2}vu + \frac{9}{4}u^2 - 3v\rho_1 - 9u\rho_1 + 12\rho_1^2 + 4\rho_3^2 + 4\rho_3\rho_4 + 2\rho_4^2 + 2\rho_5^2 + 14\rho_1,$$

$$-14x \frac{d\rho_2}{dx} = \frac{3}{4}v^2 + \frac{3}{2}vu + \frac{9}{4}u^2 - 3v\rho_2 - 9u\rho_2 + 12\rho_2^2 + 4\rho_3^2 + 4\rho_3\rho_4 + 2\rho_4^2 + 2\rho_5^2 + 14\rho_2 + 4\rho_2 \sum_k (\rho_{\ell_k} + 3\rho_{d_k} + 3\rho_{u_k}) - 4 \sum_k (\rho_{\ell_k}^2 + 3\rho_{d_k}^2 + 3\rho_{u_k}^2),$$

$$-14x \frac{d\rho_3}{dx} = (\rho_1 + \rho_2)(6\rho_3 + 2\rho_4) + 4\rho_3^2 + 2\rho_4^2 + 2\rho_5^2 + 14\rho_3 + 2\rho_3 \sum_k (\rho_{\ell_k} + 3\rho_{d_k} + 3\rho_{u_k}),$$

$$-14x \frac{d\rho_4}{dx} = 2\rho_1\rho_4 + 2\rho_2\rho_4 + 8\rho_3\rho_4 + 4\rho_4^2 + 8\rho_5^2 + 14\rho_4 + 2\rho_4 \sum_k (\rho_{\ell_k} + 3\rho_{d_k} + 3\rho_{u_k}),$$

$$-14x \frac{d\rho_5}{dx} = (2\rho_1 + 2\rho_2 + 8\rho_3 + 12\rho_4 + 14 + 2 \sum_k (\rho_{\ell_k} + 3\rho_{d_k} + 3\rho_{u_k}))\rho_5.$$

V. Reduction of Couplings in 2HDM with NFC

Solutions: Complete Reduction

In the so-called **complete reduction**, the reduction equations must satisfy:

$$\rho_i(x) = \rho_i^{(0)} + \rho_i^{(1)}x + \rho_i^{(2)}x^2 + \dots$$

$U(1)$ coupling

$$-14x \frac{dv}{dx} = 14v^2 + 14v,$$

$$-14x \frac{d}{dx} \left(v^{(0)} \right) = 0 = 14(v^{(0)})^2 + 14v^{(0)},$$

\therefore

$$v_-^{(0)} = -1, \quad v_+^{(0)} = 0;$$

$SU(2)$ coupling

$$-14x \frac{du}{dx} = -6u^2 + 14u,$$

$$-14x \frac{d}{dx} \left(u^{(0)} \right) = 0 = -6(u^{(0)})^2 + 14u^{(0)},$$

\therefore

$$u_-^{(0)} = 0, \quad u_+^{(0)} = \frac{7}{3}$$

V. Reduction of Couplings in 2HDM with NFC

Solutions: Complete Reduction

In the so-called **complete reduction**, the reduction equations must satisfy:

$$\rho_i(x) = \rho_i^{(0)} + \rho_i^{(1)}x + \rho_i^{(2)}x^2 + \dots$$

$U(1)$ coupling

$$-14x \frac{dv}{dx} = 14v^2 + 14v,$$

$$-14x \frac{d}{dx} \left(v^{(0)} \right) = 0 = 14(v^{(0)})^2 + 14v^{(0)},$$

\therefore

$$v_-^{(0)} = -1, \quad v_+^{(0)} = 0;$$

$SU(2)$ coupling

$$-14x \frac{du}{dx} = -6u^2 + 14u,$$

$$-14x \frac{d}{dx} \left(u^{(0)} \right) = 0 = -6(u^{(0)})^2 + 14u^{(0)},$$

\therefore

$$u_-^{(0)} = 0, \quad u_+^{(0)} = \frac{7}{3}$$

V. Reduction of Couplings in 2HDM with NFC

Solutions: Complete Reduction

In the so-called **complete reduction**, the reduction equations must satisfy:

$$\rho_i(x) = \rho_i^{(0)} + \rho_i^{(1)}x + \rho_i^{(2)}x^2 + \dots$$

$U(1)$ coupling

$SU(2)$ coupling

$$-14x \frac{dv}{dx} = 14v^2 + 14v,$$

$$-14x \frac{du}{dx} = -6u^2 + 14u,$$

$$-14x \frac{d}{dx} \left(v^{(0)} \right) = 0 = 14(v^{(0)})^2 + 14v^{(0)},$$

$$-14x \frac{d}{dx} \left(u^{(0)} \right) = 0 = -6(u^{(0)})^2 + 14u^{(0)},$$

\therefore

\therefore

$$v_-^{(0)} = -1, \quad v_+^{(0)} = 0;$$

$$u_-^{(0)} = 0, \quad u_+^{(0)} = \frac{7}{3}$$

V. Reduction of Couplings in 2HDM with NFC

Solutions: Complete Reduction

In the so-called **complete reduction**, the reduction equations must satisfy:

$$\rho_i(x) = \rho_i^{(0)} + \rho_i^{(1)}x + \rho_i^{(2)}x^2 + \dots$$

$U(1)$ coupling

$$-14x \frac{dv}{dx} = 14v^2 + 14v,$$

$$-14x \frac{d}{dx} \left(v^{(0)} \right) = 0 = 14(v^{(0)})^2 + 14v^{(0)},$$

\therefore

$$v_-^{(0)} = -1, \quad v_+^{(0)} = 0;$$

$SU(2)$ coupling

$$-14x \frac{du}{dx} = -6u^2 + 14u,$$

$$-14x \frac{d}{dx} \left(u^{(0)} \right) = 0 = -6(u^{(0)})^2 + 14u^{(0)},$$

\therefore

$$u_-^{(0)} = 0, \quad u_+^{(0)} = \frac{7}{3} > 1$$

Results

Solutions for all types of 2HDM with NFC

- For gauge couplings, only trivial solutions are physically acceptable:

$$v(x) = 0 \rightarrow g_1(g_3) = 0$$

$$u(x) = 0 \rightarrow g_2(g_3) = 0$$

- Given the previous result, the only solution that satisfies the physical conditions is also the trivial one

$$\rho_{\ell_j}(x) = 0 \rightarrow G_j^{(\ell)}(g_3) = 0$$

Results

Solutions for all types of 2HDM with NFC

- For gauge couplings, only trivial solutions are physically acceptable:

$$v(x) = 0 \rightarrow g_1(g_3) = 0$$

$$u(x) = 0 \rightarrow g_2(g_3) = 0$$

- Given the previous result, the only solution that satisfies the physical conditions is also the trivial one

$$\rho_{\ell_j}(x) = 0 \rightarrow G_j^{(\ell)}(g_3) = 0$$

Results

Solutions for all types of 2HDM with NFC

- For gauge couplings, only trivial solutions are physically acceptable:

$$v(x) = 0 \rightarrow g_1(g_3) = 0$$

$$u(x) = 0 \rightarrow g_2(g_3) = 0$$

- Given the previous result, the only solution that satisfies the physical conditions is also the trivial one

$$\rho_{\ell_j}(x) = 0 \rightarrow G_j^{(\ell)}(g_3) = 0$$

Results

Solutions for Yukawa and Higgs couplings

Type I and X

$$G_t(g_3) \approx \sqrt{\frac{2}{9}} g_3 + \dots$$

$$\lambda_1(g_3) = 0$$

$$\lambda_2(g_3) \approx 0.03469 g_3 + \dots$$

$$\lambda_3(g_3) = 0$$

$$\lambda_4(g_3) = 0$$

$$\lambda_5(g_3) = 0$$

Type II and Y

$$G_t(g_3) \approx \sqrt{\frac{1}{5}} g_3 + \dots$$

$$\lambda_1(g_3) \approx 0.0285714 g_3 + \dots$$

$$\lambda_2(g_3) \approx 0.0285714 g_3 + \dots$$

$$\lambda_3(g_3) \approx 0.0285714 g_3 + \dots$$

$$\lambda_4(g_3) \approx -0.028868 g_3 + \dots$$

$$\lambda_5(g_3) = 0$$

Results

Solutions for Yukawa and Higgs couplings

Type I and X

$$G_t(g_3) \approx \sqrt{\frac{2}{9}} g_3 + \dots$$

$$\lambda_1(g_3) = 0$$

$$\lambda_2(g_3) \approx 0.03469 g_3 + \dots$$

$$\lambda_3(g_3) = 0$$

$$\lambda_4(g_3) = 0$$

$$\lambda_5(g_3) = 0$$

Type II and Y

$$G_t(g_3) \approx \sqrt{\frac{1}{5}} g_3 + \dots$$

$$\lambda_1(g_3) \approx 0.0285714 g_3 + \dots$$

$$\lambda_2(g_3) \approx 0.0285714 g_3 + \dots$$

$$\lambda_3(g_3) \approx 0.0285714 g_3 + \dots$$

$$\lambda_4(g_3) \approx -0.028868 g_3 + \dots$$

$$\lambda_5(g_3) = 0$$

Results

Solutions for Yukawa and Higgs couplings

Type I and X

$$G_t(g_3) \approx \sqrt{\frac{2}{9}} g_3 + \dots$$

$$\lambda_1(g_3) = 0$$

$$\lambda_2(g_3) \approx 0.03469 g_3 + \dots$$

$$\lambda_3(g_3) = 0$$

$$\lambda_4(g_3) = 0$$

$$\lambda_5(g_3) = 0$$

Type II and Y

$$G_t(g_3) \approx \sqrt{\frac{1}{5}} g_3 + \dots$$

$$\lambda_1(g_3) \approx 0.0285714 g_3 + \dots$$

$$\lambda_2(g_3) \approx 0.0285714 g_3 + \dots$$

$$\lambda_3(g_3) \approx 0.0285714 g_3 + \dots$$

$$\lambda_4(g_3) \approx -0.028868 g_3 + \dots$$

$$\lambda_5(g_3) = 0$$

Results

Mass predictions: Type I y X

At the scale of the m_Z the strong coupling constant is given as $x = \frac{g_3^2}{4\pi} = 0.1179(9)$ ([PDG, 2021](#)), with this the constants will take the values:

$$G_t \approx 0.573793$$

$$\lambda_2 \approx 0.513947$$

$$\lambda_1 = \lambda_3 = \lambda_4 = \lambda_5 = 0$$

$$m_t \leq 100 \text{ GeV}$$

$$m_H \leq 56 \text{ GeV}$$

Results

Mass predictions: Type II y Y

At the scale of the M_Z the strong coupling constant is given as $x = \frac{g_3^2}{4\pi} = 0.1179(9)$ ([PDG, 2021](#)), with this the constants will take the values:

$$G_t \approx 0.544348$$

$$\lambda_1 \approx 0.0423307$$

$$\lambda_2 \approx 0.0423307$$

$$\lambda_3 \approx 0.0423307$$

$$\lambda_4 \approx -0.0427701$$

$$\lambda_5 = 0$$

$$m_t \approx 94.7 \text{ GeV}$$

$$m_H \approx 50.6 \text{ GeV}$$

$$m_h \approx 1.7 \text{ GeV}$$

$$m_{H^\pm} \approx 36 \text{ GeV}$$

$$m_A = 0$$

Conclusions

- The **Reduction of Couplings Method** (RCM) is presented as an attractive tool to deal with the ***free parameters problem***.
- Although the preliminary results obtained seem not to agree with the phenomenology, this doesn't discard the method
- Dr. M. Mondragón et al. searches in the MSSM has given more accurated results
- Given the physical possibilities that multi-Higgs models offer, we will continue the exploration of these models using the RCM.

Conclusions

- The **Reduction of Couplings Method** (RCM) is presented as an attractive tool to deal with the ***free parameters problem***.
- Although the preliminary results obtained seem not to agree with the phenomenology, this doesn't discard the method
- Dr. M. Mondragón et al. searches in the MSSM has given more accurated results
- Given the physical possibilities that multi-Higgs models offer, we will continue the exploration of these models using the RCM.

Conclusions

- The **Reduction of Couplings Method** (RCM) is presented as an attractive tool to deal with the ***free parameters problem***.
- Although the preliminary results obtained seem not to agree with the phenomenology, this doesn't discard the method → **it's known that the tested models are not actually the final physical theories.**
- **Dr. M. Mondragón et al.** searches in the MSSM has given more accurated results
- Given the physical possibilities that multi-Higgs models offer, we will continue the exploration of these models using the RCM.

Conclusions

- The **Reduction of Couplings Method** (RCM) is presented as an attractive tool to deal with the ***free parameters problem***.
- Although the preliminary results obtained seem not to agree with the phenomenology, this doesn't discard the method → **it's known that the tested models are not actually the final physical theories.**
- **Dr. M. Mondragón et al.** searches in the MSSM has given more accurated results
- Given the physical possibilities that multi-Higgs models offer, we will continue the exploration of these models using the RCM.

Conclusions

- The **Reduction of Couplings Method (RCM)** is presented as an attractive tool to deal with the ***free parameters problem***.
- Although the preliminary results obtained seem not to agree with the phenomenology, this doesn't discard the method → **it's known that the tested models are not actually the final physical theories.**
- **Dr. M. Mondragón et al.** searches in the MSSM has given more accurated results → **this suggest that RCM can be used as a testing tool for QFT models symmetries.**
- Given the physical possibilities that multi-Higgs models offer, we will continue the exploration of these models using the RCM.



Conclusions

- The **Reduction of Couplings Method** (RCM) is presented as an attractive tool to deal with the ***free parameters problem***.
- Although the preliminary results obtained seem not to agree with the phenomenology, this doesn't discard the method → **it's known that the tested models are not actually the final physical theories.**
- **Dr. M. Mondragón et al.** searches in the MSSM has given more accurated results → **this suggest that RCM can be used as a testing tool for QFT models symmetries.**
- Given the physical possibilities that multi-Higgs models offer, we will continue the exploration of these models using the RCM.



Thanks for your attention!

