

Interacciones efectivas de Yukawa y texturas universales

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Introduction

		fermions (3 générations de la matière)			bosons (forces)	
		I	II	III		
masse →		2.4 MeV	1.27 GeV	171.2 GeV	0	électromagnétisme
charge →		$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	0	
spin →		$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	
nom →		u up	c charm	t top	γ photon	
	Quarks	4.8 MeV	104 MeV	4.2 GeV	0	interaction forte
		$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$	0	
		$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	
		d down	s strange	b bottom	g gluon	
	Leptons	<2.2 eV	<0.17 MeV	<15.5 MeV	91.2 GeV	interaction faible
		0	0	0	0	
		$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	
		ν_e neutrino électronique	ν_μ neutrino muonique	ν_τ neutrino tauique	Z⁰ boson Z ⁰	
		0.511 MeV	105.7 MeV	1.777 GeV	80.4 GeV	interaction faible
		-1	-1	-1	±1	
		$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	
		e électron	μ muon	τ tau	W[±] boson W	

Figure: Particles of the SM

- Gauge theory that unifies the fundamental forces... most of them.
- Able to predict many experimental data.

$$\mathcal{L}_{ME} = \mathcal{L}_G + \mathcal{L}_H + \mathcal{L}_D + \mathcal{L}_Y \quad (1)$$

- \mathcal{L}_G (Gauge sector)
- \mathcal{L}_H (Scalar Sector)
- \mathcal{L}_D (Dirac Sector)
- \mathcal{L}_Y (Yukawa Sector)

Neutrino massess

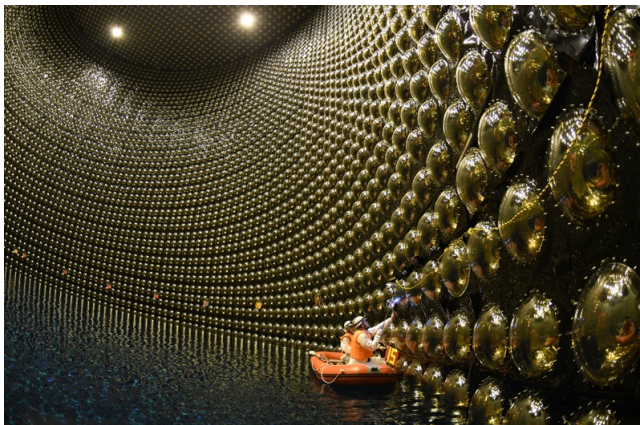


Figure: SuperKamiokande

Dark Matter

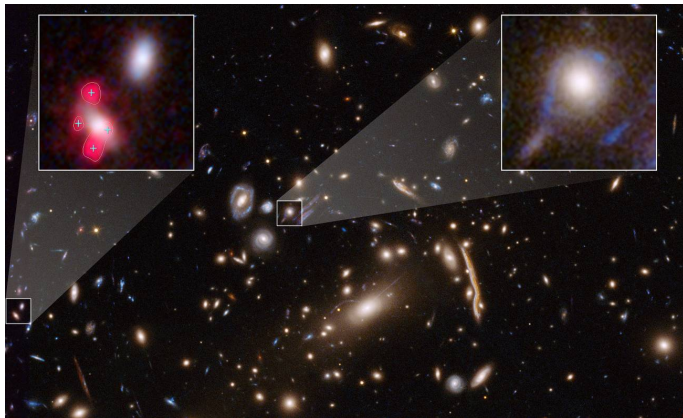


Figure: Gravitational lens

Hierarchy of masses

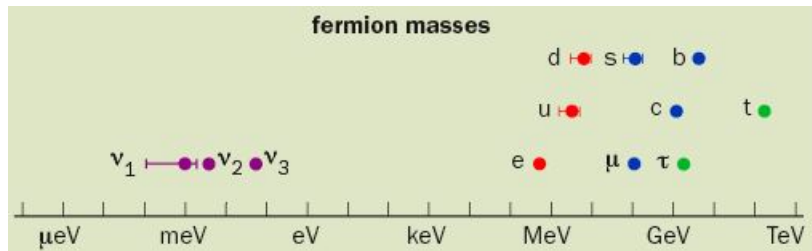


Figure: Hierarchy of masses

SM Questions

- Why we have this hierarchy on the fermions masses?
- All the parameters on the flavour sector are arbitrary? Are connected?



$$\mathcal{L}_Y = Y_{ij}^d \bar{Q}_{Li} \phi D_{Rj} + Y_{ij}^u \bar{Q}_{Li} \tilde{\phi} U_{Rj} + Y_{ij}^l \bar{L}_{Li} \phi l_{Rj} + H.C \quad (2)$$

$$\mathcal{L}_Y = Y_{ij}^d \bar{Q}_{Li} \phi D_{Rj} + Y_{ij}^u \bar{Q}_{Li} \tilde{\phi} U_{Rj} + Y_{ij}^l \bar{L}_{Li} \phi l_{Rj} + \underline{Y_{ij}^\nu \bar{L}_{Li} \tilde{\phi} \nu_{Rj}} + \mathcal{L}_m + H.C \quad (3)$$

Considering the leptonic terms on the Yukawa Lagrangian

$$\mathcal{L}_Y \supset Y_{ij}^l L_{Li}^- \phi l_{Rj} + Y_{ij}^\nu L_{Li}^- \tilde{\phi} \nu_{Rj} + \mathcal{L}_m + HC \quad (4)$$

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About the model

- Extension of the SM that introduces particles of "flavor" types and right handed neutrinos.
- Inspired by the Froggat Nielsen Mechanism

$$\mathcal{L}_Y \supset Y_{ij}^l \bar{L}_{Li} \phi l_{Rj} + Y_{ij}^\nu \bar{L}_{Li} \tilde{\phi} \nu_{Rj} + HC \quad (5)$$

L_L the doublet of $SU(2)$ that agroups the left fermionic fields , l_R y ν_R are singlets from $SU(2)$. And ϕ it is de Higgs doublet.

If we propose a "Yukawa Matrix at high energies"

$$Y'_{ij}{}^l = \frac{\lambda_l}{\Lambda} Z_{Li}^{l\dagger}(x) Z_{Rj}^l(x) \quad (6)$$

Such that, taking the VEV

$$\langle Y'_{ij}{}^l \rangle = \frac{\lambda_l}{\Lambda} \langle Z_{Li}^{l\dagger}(x) \rangle \langle Z_{Rj}^l(x) \rangle = Y_{ij}^l \quad (7)$$

Where, Z_L y Z_R are scalar fields transforming as triplets of flavor on the fundamental of $SU(3)_F$.

Where:

$$Z_i = e^{\frac{i\psi_a(x)T_a}{\Lambda\psi}} X_i(x) \quad (8)$$

And if we rewrite the lagrangian

$$\mathcal{L}_Y \supset L_{Li}^- Y_{ij}'^l \phi l_{Rj} + L_{Li}^- Y_{ij}'^\nu \tilde{\phi} \nu_{Rj} + HC \quad (9)$$

And taking his explicit form:

$$Y'_{ij} = \frac{\lambda_l}{\Lambda} e^{\frac{-i\psi_a(x)T_a}{\Lambda_\psi}} X_{Li}^\dagger(x) X_{Rj}(x) e^{\frac{i\psi_a(x)T_a}{\Lambda_\psi}} \quad (10)$$

Applying the Baker-Campbell-Hausdorff we have:

$$Y'_{ij} = \frac{\lambda_l}{\Lambda} (X_{Li}^\dagger X_{Rj} - i \frac{\psi_a(x)}{\Lambda_\psi} [T_a, X_{Li}^\dagger X_{Rj}]) \quad (11)$$

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- Using Python language we write a code where we generate the Yukawa matrix through the assignment of random parameters.
- χ^2 parameter to filter data.

We want to replicate the observables that we have

- $\Delta m_{32}^2 \approx 2.515 \pm 0.027 \times 10^{-3}$
- $\Delta m_{21}^2 \approx 7.42 \pm 0.21 \times 10^{-5}$

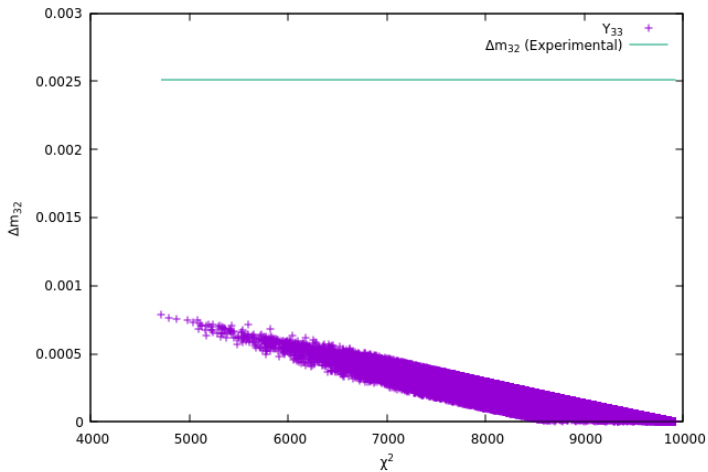


Figure: Δm_{32}^2 vs χ^2

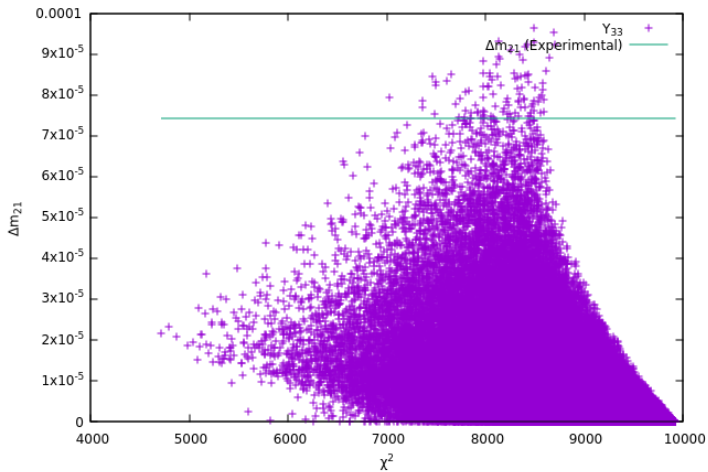


Figure: Δm_{21}^2 vs χ^2

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Conclusions and perspective

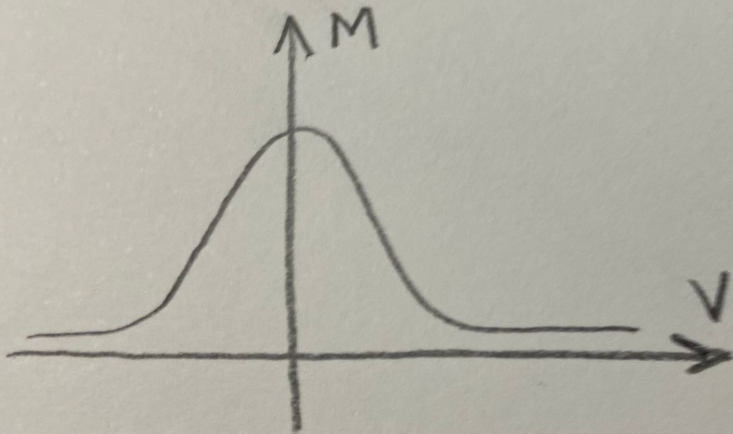
Conclusions

Introducing this structure of flavon type to generate the Yukawa Matrix it is possible to demonstrate that exists a common mechanism to generate the textures whose eigenvalues are the masses of the fermions.

Perspective

It is necessary to improve our model. The next scenario is including Majorana mass terms.

Thanks!



References

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