# Medal of the Mexican Society of Physics, Division of Particles and fields









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- Workshops in Guanajuato, Hermosillo, etc
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Physics with intact protons at the LHC: from the odderon discovery to the sensitivity to beyond standard model physics



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- Elastic interactions and introduction to the Odderon
- D0 pp̄ and TOTEM pp data
- The odderon discovery
- Study of quartic anomalous couplings and search for axion-like particles
- Ultra Fast Silicon detectors



#### What is elastic scattering? The pool game...



- We want to study "elastic" collisions between protons and proton-antiprotons
- In high energy physics:  $pp \rightarrow pp$ and  $p\bar{p} \rightarrow p\bar{p}$
- In these interactions, each proton/antiproton remains intact after interaction but are scattered at some angles and can lose/gain some momentum as in the pool game

#### What do we want to study?





- We want to study elastic interactions: pp o pp or par p o par p
- These are very clean events, where nothing is produced outside the two protons
- How to detect/measure these events? We need to detect the intact protons after interaction!
- Interactions explained by the exchange of a colorless object ( $\geq$  2 gluons, photon, etc...) between the two protons

#### How to explain the fact that protons can be intact?



- Quarks/gluons radiate lots of gluons when one tries to separate them (confinement)
- Gluons exchange color, interact with other gluons in the proton and in that case protons are destroyed in the final state
- In order to explain how protons can remain intact: we need colorless exchanges, or at least 2 gluons to be exchanged

#### $p\bar{p}$ interactions: the Tevatron



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#### pp interactions: The Large Hadron Collider at CERN

- Large Hadron Collider at CERN: proton proton collider with 2.76, 7, 8 and 13 TeV center-of-mass energy
- Circonference: 27 km; Underground: 50-100 m



#### Which tools do we have? Roman Pot detectors



- We use special detectors to detect intact protons/ anti-protons called Roman Pots
- These detectors can move very close to the beam (up to 3σ) when beam are stable so that protons scattered at very small angles can be measured



#### The odderon in a nutshell



- Let us assume that elastic scattering can be due to exchange of colorless objects: Pomeron and Odderon
- Charge parity C: Charge conjugation changes the sign of all quantum charges

- Pomeron and Odderon correspond to positive and negative C parity: Pomeron is made of two gluons which leads to a +1 parity whereas the odderon is made of 3 gluons corresponding to a -1 parity
- Scattering amplitudes can be written as:

 $A_{pp} = Even + Odd$  $A_{p\bar{p}} = Even - Odd$ 

 From the equations above, it is clear that observing a difference between *pp* and *pp̄* interactions would be a clear way to observe the odderon

### What is the odderon? The QCD picture



- Multi-gluon exchanges in hadron-hadron interactions in elastic *pp* interactions (Bartels-Kwiecinski-Praszalowicz)
- From B. Nicolescu: The Odderon is defined as a singularity in the complex plane, located at J = 1 when t = 0 and which contributes to the odd crossing amplitude



- Leads to contributions on 3,... gluon exchanges in terms of QCD for the perturbative odderon
- Colorless C-odd 3-gluon state (odderon) predicts differences in elastic dσ/dt for pp and pp̄ interactions since it corresponds to different amplitudes/ interferences

#### Measurement of elastic scattering at Tevatron and LHC



- Study of elastic pp → pp reaction: exchange of momentum between the two protons which remain intact
- Measure intact protons scattered close to the beam using Roman Pots installed both by D0 and TOTEM collaborations
- From counting the number of events as a function of |t| (4-momentum transferred square at the proton vertex measured by tracking the protons), we get  $d\sigma/dt$



- The situation is not that simple: elastic scattering at low energies can be due to exchanges of additional particles to pomeron/odderon: ρ, ω, φ, reggeons...
- How to distinguish between all these exchanges? Not easy...
- At ISR energies, there was already some indication of a possible difference between pp and  $p\bar{p}$  interactions, differences of about  $3\sigma$  between pp and  $p\bar{p}$  interactions but this was not considered to be a clean proof of the odderon because of these additional reggeon, meson exchanges at low  $\sqrt{s}$

#### What is the expected situation at the LHC?



- Expected elastic  $d\sigma/dt$  before LHC measurements
- Many different predictions including many possible contributions at high |t|, such as pomeron, reggeon, mesons (ω, φ) whereas other predictions mentioned that, at high energies, we should be more asymptotical and pomeron dominated
- Almost nobody thought about the odderon (except a few theorists such as Martynov, Nicolescu...)

#### Are we in the asymptotic regime at the LHC?



- Contrary to what some models expected before LHC, the elastic cross section is smooth: we do not see reggeons, mesons...!
- Effects of reggeon, meson exchanges are negligible at LHC energies: we can concentrate on pomeron/odderon studies!
- We can directly look for the existence of the odderon by comparing *pp* and *pp̄* elastic cross sections at very high energies: 1.96 TeV (Tevatron), 2.76, 7, 8, 13 (LHC)

### D0 elastic $p\bar{p} \ d\sigma/dt$ cross section measurements



- D0 collected elastic pp̄ data with intact p and p̄ detected in the Forward Proton Detector with 31 nb<sup>-1</sup> Phys. Rev. D 86 (2012) 012009
- Measurement of elastic  $p\bar{p} \ d\sigma/dt$  at 1.96 TeV for 0.26 <|t| < 1.2 GeV<sup>2</sup>

### Elastic cross section at the LHC: Forward coverage in CMS-TOTEM



Roman Pots: elastic & diffractive protons close to outgoing beams → Proton Trigger



#### TOTEM cross section measurements



### TOTEM elastic $pp \ d\sigma/dt$ cross section measurements

- Elastic *pp*  $d\sigma/dt$  measurements: tag both intact protons in TOTEM Roman Pots 2.76, 7, 8 and 13 TeV
- Very precise measurements at 2.76, 7, 8 and 13 TeV: Eur. Phys. J. C 80 (2020) no.2, 91; EPL 95 (2011) no. 41004; Nucl. Phys. B 899 (2015) 527; Eur. Phys. J. C79 (2019) no.10, 861



#### Strategy to compare pp and $p\bar{p}$ data sets



- In order to identify differences between pp and pp̄ elastic dσ/dt data, we need to compare TOTEM measurements at 2.76, 7, 8, 13 TeV and D0 measurements at 1.96 TeV
- All TOTEM dσ/dt measurements show the same features, namely the presence of a dip and a bump in data, whereas D0 data do not show this feature

### Reference points of elastic $d\sigma/dt$



• Define 8 characteristic points of elastic pp $d\sigma/dt$  cross sections (dip, bump...) that are feature of elastic pp interactions

- Determine how the values of |t| and  $d\sigma/dt$  of characteristic points vary as a function of  $\sqrt{s}$  in order to predict their values at 1.96 TeV
- We use data points closest to those characteristic points (avoiding model-dependent fits)
- Data bins are merged in case there are two adjacent dip or bump points of about equal value
- This gives a distribution of t and  $d\sigma/dt$  values as a function of  $\sqrt{s}$  for all characteristic points



- Bump over dip ratio measured for pp interactions at ISR and LHC energies
- Bump over dip ratio in *pp* elastic collisions: decreasing as a function of  $\sqrt{s}$  up to  $\sim 100$  GeV and flat above
- D0  $p\bar{p}$  shows a ratio of  $1.00\pm0.21$  given the fact that no bump/dip is observed in  $p\bar{p}$  data within uncertainties: more than  $3\sigma$  difference between pp and  $p\bar{p}$  elastic data (assuming flat behavior above  $\sqrt{s} = 100 \, GeV$ )

#### Variation of t and $d\sigma/dt$ values for reference points



$$|t| = a \log(\sqrt{s} [\text{TeV}]) + b$$
  $(d\sigma/dt) = c\sqrt{s} [\text{TeV}] + d$ 

#### Fits of TOTEM extrapolated characteristic points at 1.96 TeV

- Last step: predict the *pp* elastic cross sections at the same *t* values as measured by D0 in order to make a direct comparison
- Fit the reference points extrapolated to 1.96 TeV from TOTEM measurements using a double exponential fit ( $\chi^2 = 0.63$  per dof):  $h(t) = a_1 e^{-b_1|t|^2 c_1|t|} + d_1 e^{-f_1|t|^3 g_1|t|^2 h_1|t|}$



- Differences in normalization taken into account by adjusting TOTEM and D0 data sets to have the same cross sections at the optical point  $d\sigma/dt(t=0)$
- Predict the *pp* total cross section from extrapolated fit to TOTEM data  $(\chi^2 = 0.27)$ : *pp*  $\sigma_{tot} = 82.7 \pm 3.1$  mb at 1.96 TeV

# Relative normalization between D0 measurement and extrapolated TOTEM data: Rescaling TOTEM data

- Adjust 1.96 TeV  $d\sigma/dt(t=0)$  from extrapolated TOTEM data to D0 measurement
- From TOTEM *pp*  $\sigma_{tot}$ , obtain  $d\sigma/dt(t=0)$  :

$$\sigma_{tot}^2 = \frac{16\pi(\hbar c)^2}{1+\rho^2} \left(\frac{d\sigma}{dt}\right)_{t=0}$$

- Assuming  $\rho = 0.145$ , the ratio of the imaginary and the real part of the elastic amplitude, as taken from COMPETE extrapolation
- This leads to a TOTEM  $d\sigma/dt(t=0)$  at the OP of 357.1  $\pm$  26.4 mb/GeV<sup>2</sup>
- D0 measured the optical point of  $d\sigma/dt$  at small t:  $341\pm48 \text{ mb/GeV}^2$
- $\bullet$  TOTEM data rescaled by 0.954  $\pm$  0.071
- NB: We do not claim that we performed a measurement of  $d\sigma/dt$  at the OP at t = 0 (it would require additional measurements closer to t = 0), but we use the two extrapolations simply in order to obtain a common and somewhat arbitrary normalization point

# Predictions at $\sqrt{s} = 1.96$ TeV

- Reference points at 1.96 TeV (extrapolating TOTEM data) and  $1\sigma$  uncertainty band
- Comparison with D0 data: the  $\chi^2$  test with six degrees of freedom yields the *p*-value of 0.00061, corresponding to a significance of 3.4 $\sigma$



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#### Combination with additional TOTEM measurement: $\rho$ measurement



• Measure elastic scattering at very low t: Coulomb-Nuclear interference region

$$rac{d\sigma}{dt} \sim |A^{C} + A^{N}(1 - lpha G(t))|^{2}$$

- The differential cross section is sensitive to the phase of the nuclear amplitude
- In the CNI region, both the modulus and the phase of the nuclear amplitude can be used to determine  $\rho = \frac{Re(A^N(0))}{Im(A^N(0))}$  where the modulus is constrained by the measurement in the hadronic region and the phase by the t dependence

# A previous measurement by TOTEM: $\rho$ and $\sigma_{tot}$ measurements as an indication for odderon



- $\rho$  is the ratio of the real to imaginary part of the elastic amplitude at t = 0
- Using low |t| data in the Coulomb-nuclear interference region, measurement of  $\rho$  at 13 TeV:  $\rho = 0.09 \pm 0.01$  (EPJC 79 (2019) 785)
- Combination of the measured  $\rho$  and  $\sigma_{tot}$  values not compatible with any set of models without odderon exchange (COMPETE predictions above as an example)
- This result can be explained by the exchange of the Odderon in addition to the Pomeron

- Combination with the independent evidence of the odderon found by the TOTEM Collaboration using  $\rho$  and total cross section measurements at low t in a completely different kinematical domain
- For the models included in COMPETE, the TOTEM  $\rho$  measurement at 13 TeV provided a 3.4 to 4.6 $\sigma$  significance, to be combined with the D0/TOTEM result
- The combined significance ranges from 5.3 to 5.7 $\sigma$  depending on the model
- Models without colorless *C*-odd gluonic compound are excluded including the Durham model and different sets of COMPETE models (blue, magenta and green bands on the previous slide)

#### Searching for beyond standard model physics using intact protons



Physics with intact protons at the LHC: from the odderon discovery to the sensitivity to beyond stand.

#### Roman pot detectors from PPS installed in the tunnel



- Good acceptance at high mass in standard runs (PPS in CMS, AFP in ATLAS)
- $\bullet \ {>}100 \ fb^{-1}$  collected in Run II





# Search for quartic $\gamma\gamma\gamma\gamma\gamma$ anomalous coupling



• Search for production of two photons and two intact protons in the final state:  $pp \rightarrow p\gamma\gamma p$ 

- Additional channels: WW, ZZ,  $\gamma Z$ ,  $t\bar{t}$
- Possible larger number of events than expected in SM due to extra-dimensions, composite Higgs models, axion-like particles
- Anomalous couplings can appear via loops of new particles coupling to photons or via resonances decaying into two photons
- JHEP 1806 (2018) 131; JHEP 1502 (2015) 165; Phys.Rev. D89 (2014) 114004; Phys.Rev. D81 (2010) 074003; Phys.Rev. D78 (2008) 073005

#### Motivations to look for quartic $\gamma\gamma$ anomalous couplings



- Two effective operators and two different couplings at low energies  $\zeta$
- $\gamma\gamma\gamma\gamma$  couplings can be modified in a model independent way by loops of heavy charge particles

$$\zeta_1 = \alpha_{em}^2 Q^4 m^{-4} N c_{1,s}$$

where the coupling depends only on  $Q^4 m^{-4}$  (charge and mass of the charged particle) and on spin,  $c_{1,s}$  depends on the spin of the particle This leads to  $\zeta_1$  of the order of  $10^{-14}$ - $10^{-13}$ 

#### Motivations to look for quartic $\gamma\gamma$ anomalous couplings



- Two effective operators at low energies
- $\zeta_1$  can also be modified by neutral particles at tree level (extensions of the SM including scalar, pseudo-scalar, and spin-2 resonances that couple to the photon)  $\zeta_1 = (f_s m)^{-2} d_{1,s}$  where  $f_s$  is the  $\gamma \gamma X$  coupling of the new particle to the photon, and  $d_{1,s}$  depends on the spin of the particle; for instance, 2 TeV dilatons lead to  $\zeta_1 \sim 10^{-13}$

#### $\gamma\gamma$ exclusive production: SM contribution



- QCD production dominates at low  $m_{\gamma\gamma}$ , QED at high  $m_{\gamma\gamma}$
- Important to consider W loops at high  $m_{\gamma\gamma}$
- At high masses (> 200 GeV), the photon induced processes are dominant
- Conclusion: Two photons and two tagged protons means photon-induced process

#### Removing pile up at the LHC

- Advantage of tagging protons: negligible background after matching mass/rapidity of photon and proton systems (JHEP 1502 (2015) 165; Phys.Rev. D89 (2014) 114004)
- Possibility to use fast timing detectors to measure proton time of flights



# First search for high mass exclusive $\gamma\gamma$ production (CMS/TOTEM)



- Search for exclusive diphoton production: back-to-back, high diphoton mass ( $m_{\gamma\gamma} > 350$  GeV), matching in rapidity and mass between diphoton and proton information
- First limits on quartic photon anomalous couplings:  $|\zeta_1| < 2.9 \ 10^{-13} \ \text{GeV}^{-4}$ ,  $|\zeta_2| < 6. \ 10^{-13} \ \text{GeV}^{-4}$  with about 10 fb<sup>-1</sup>, accepted by PRL (2110.05916)
- Limit updates with 102.7 fb<sup>-1</sup>:  $|\zeta_1| < 7.3 \ 10^{-14} \ \text{GeV}^{-4}$ ,  $|\zeta_2| < 1.5 \ 10^{-13} \ \text{GeV}^{-4}$

# First search for high mass production of axion-like particles (CMS/TOTEM)



- First limits on ALPs at high mass (CMS-PAS-EXO-21-007)
- Sensivities projected with 300 fb<sup>-1</sup> (C. Baldenegro, S. Fichet, G. von Gersdorff, C. Royon, JHEP 1806 (2018) 13)

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- Production of ALPs via photon exchanges in heavy ion runs: Complementarity to *pp* running
- Sensitivity to low mass ALPs: low luminosity but cross section increased by Z<sup>4</sup>, C. Baldenegro, S. Hassani, C.R., L. Schoeffel, ArXiv:1903.04151
- Similar gain of three orders of magnitude on sensitivity for γγγZ couplings in pp collisions:
   C. Baldenegro, S. Fichet, G. von Gersdorff, C. R., JHEP 1706 (2017) 142

# Exclusive production of W boson pairs (CMS/TOTEM)



• Search with fully hadronic decays of *W* bosons: anomalous production of *WW* events dominates at high mass with a rather low cross section

- 2 "fat" jets (radius 0.8), jet  $p_T > 200$ GeV, 1126<  $m_{jj} < 2500$  GeV, jets back-to-back ( $|1 - \phi_{jj}/\pi| < 0.01$ )
- Signal region defined by the correlation between central *WW* system and proton information



# WW and ZZ exclusive productions (CMS/TOTEM)



- Searches performed in full hadronic decays of *W* bosons (high cross section) with AK8 jets
- SM cross section is low
- Limits on SM cross section  $\sigma_{WW} < 67 {\rm fb}, \ \sigma_{ZZ} < 43 {\rm fb}$  for 0.04  $< \xi <$  0.2 (CMS-PAS-EXO-21-014)
- New limits on quartic anomalous couplings:  $a_0^W/\Lambda^2 < 4.3 \ 10^{-6} \ \text{GeV}^{-2}$ ,  $a_C^W/\Lambda^2 < 1.6 \ 10^{-5} \ \text{GeV}^{-2}$ ,  $a_0^Z/\Lambda^2 < 0.9 \ 10^{-5} \ \text{GeV}^{-2}$ ,  $a_C^Z/\Lambda^2 < 4. \ 10^{-5} \ \text{GeV}^{-2}$  with 52.9 fb<sup>-1</sup>

#### The future: Observation of exclusive WW production



- SM contributiona appears at lower WW masses compared to anomalous couplings
- Use purely leptonic channels for *W* decays (the dijet background is too high at low masses for hadronic channels)
- SM prediction on exclusive WW (leptonic decays) after selection: about 50 events for 300 fb<sup>-1</sup> (2 background)
- JHEP 2012 (2020) 165, C. Baldenegro, G. Biagi, G. Legras, C.R.

 $\gamma\gamma\gamma\gamma Z$  quartic anomalous coupling: leptonic and hadronic decays of Z boson



- C. Baldenegro, S. Fichet, G. von Gersdorff, C. Royon, JHEP 1706 (2017) 142
- Best expected reach at the LHC by about three orders of magnitude
- Sensitivity to wide/narrow resonances, loops of new particles

# Exclusive $t\bar{t}$ production (CMS/TOTEM)



dilep channel ( $\bar{t}t \rightarrow l\nu b + l\nu \bar{b}$ )	Semilep channel ( $\bar{t}t \rightarrow l\nu b + jj\bar{b}$ )				
Object selection					
Leptons: pT>30(20)GeV,  η <2.1 Jets: pT>30GeV,  η <2.4, ΔR(j,l)>0.4	Leptons: pT>30GeV,  η <2.1(2.4) for e(μ) Jets: pT>25GeV,  η <2.4, ΔR(j,l)>0.4				
Event selection					
≥2 leptons (OS pair), m(ll)-m(Z) >15GeV ≥2 b-jets 1 proton / side	=1 lepton ≥2 b-jets, ≥2 non b-jets 1 proton / side				

# Exclusive $t\bar{t}$ production (CMS/TOTEM)



• Kinematic fitter based on *W* and *t* mass constraints to reduce background



- Search for exclusive  $t\bar{t}$  production in leptonic and semi-leptonic modes
- $\sigma_{t\bar{t}}^{excl.} <$  0.6 pb (CMS-PAS-TOP-21-007)

#### Exclusive $t\bar{t}$ production: the future

- Search for  $\gamma\gamma t\bar{t}$  anomalous coupling in semi-leptonic decays with 300 fb<sup>-1</sup>
- Use similar selection: high  $t\bar{t}$  mass, matching between pp and  $t\bar{t}$  information
- Use fast timing detectors to suppress further the pile up background
- C. Baldenegro, A. Bellora, S. Fichet, G. von Gersdorff, M. Pitt, CR arXiv:2205.01173

Coupling $[10^{-11}  \text{GeV}^{-4}]$	$95\%~{ m CL}$	$5\sigma$	$95\%{ m CL}(60{ m ps})$	$5\sigma \ (60  \mathrm{ps})$	$95\%{ m CL}(20{ m ps})$	$5\sigma \ (20  \mathrm{ps})$
$\zeta_1$	1.5	2.5	1.1	1.9	0.74	1.5
$\zeta_2$	1.4	2.4	1.0	1.7	0.70	1.4
$\zeta_3$	1.4	2.4	1.0	1.7	0.70	1.4
$\zeta_4$	1.5	2.5	1.0	1.8	0.73	1.4
$\zeta_5$	1.2	2.0	0.84	1.5	0.60	1.2
$\zeta_6$	1.3	2.2	0.92	1.6	0.66	1.3

# Additional method to remove pile up: Measuring proton time-of-flight



- Measure the proton time-of-flight in order to determine if they originate from the same interaction as the selected photon
- Typical precision: 10 ps means 2.1 mm
- Idea: use ultra-fast Si detectors (signal duration of ~few ns and possibility to use fast sampling to reconstruct full signal)

#### Timing measurements in Particle Physics



- Proton going through a detector (for instance scintillator, Silicon) emits a signal
- Measure this signal using an oscilloscope, or some electronics



- Amplify the signal
- Very fast digitization of the signal: measure many points on the fast increasing signal as an example
- Allows reconstructing both the shape and amplitude of signal
- Leads to precise timing measurements (using for instance time when signal starts), and energy/type of particle measurements

#### Measuring cosmic ray in space: the AGILE project

- We want to measure the type of particles (*p*, *He*, *Fe*, *Pb*, ...) and at the same time their energies
- Analysis of cosmic ray particles: using a cube sat, cheap to be sent into space
- Use similar technics: measure the signal (Bragg peak) where the particle stops in a ultra-fast Si detector
- Allows extracting type/energy of particles: project in collaboration with NASA, to be launched in Spring 2022, https://arxiv.org/abs/2103.00613



# Tests performed at St Luke hospital, University of Dublin, Ireland



- Measurement of charge deposited in Si detector compared to standard measurement using an ion chamber: good correlation
- Our detectors see in addition the beam structure (periodicity of the beam of  $\sim$ 330 ps, contrary to a few seconds for the ion chamber): measure single particles from the beam
- Fundamental to measure instantaneous doses for high intensity proton therapy as example
- For more details: https://arxiv.org/abs/2101.07134

#### Conclusion

- pp and  $p\bar{p}$  cross sections differ with a significance of  $3.4\sigma$  in a model-independent way and thus provides evidence that the Colorless *C*-odd gluonic compound i.e. the odderon is needed to explain elastic scattering at high energies
- When combined with the  $\rho$  and total cross section result at 13 TeV, the significance is in the range 5.3 to 5.7 $\sigma$  and thus constitutes the first experimental observation of the odderon: Major discovery at CERN/Tevatron
- PPS allows probing quartic anomalous couplings with unprecedented precision: sensitivity to composite Higgs, extra-dimension models, axion-like particles
- Development of fast timing detectors for HEP and applications in medicine, cosmic-ray physics



#### We need to look everywhere! For instance using intact protons...



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