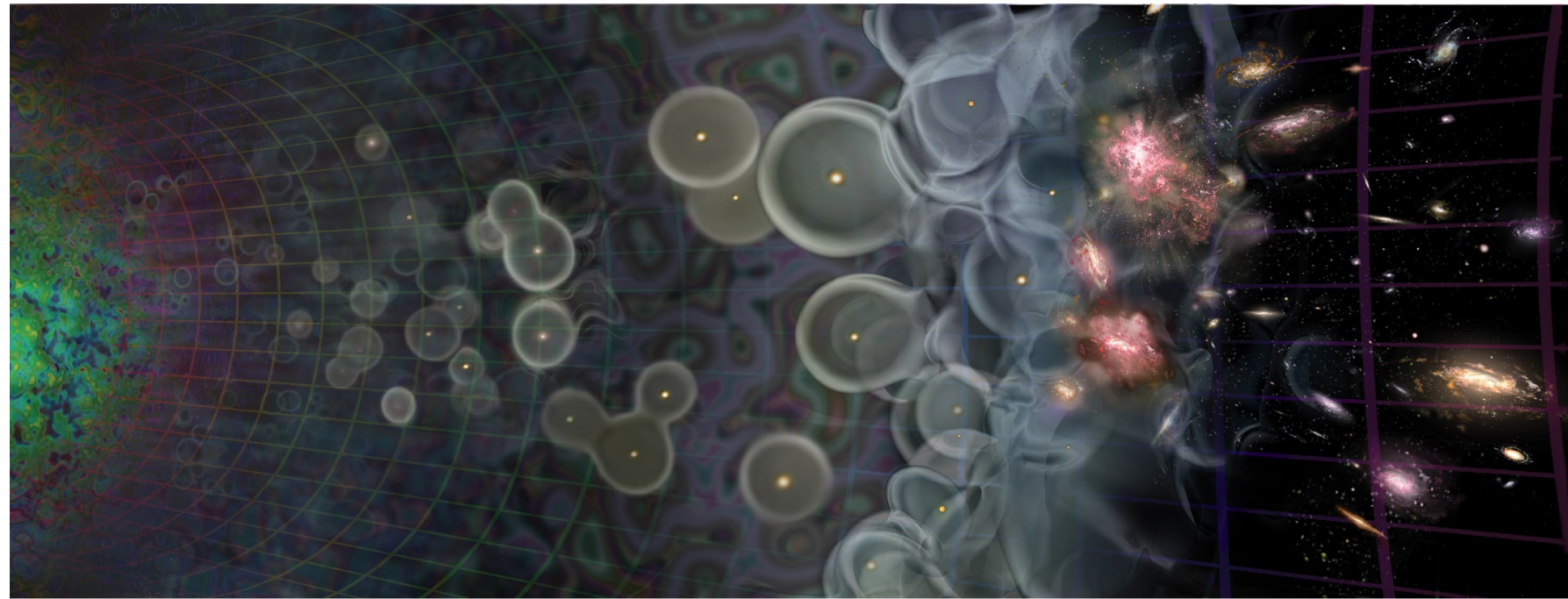


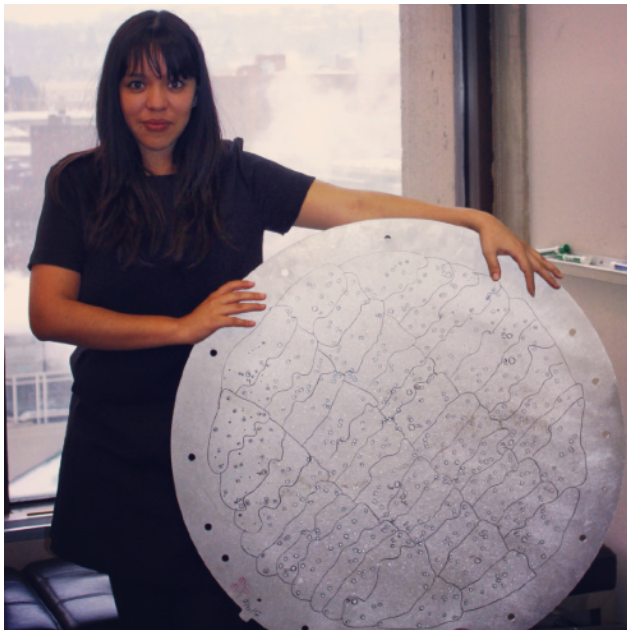


Proposal of local bias estimation for forward modeling Reconstruction in the BAO Peak analysis.

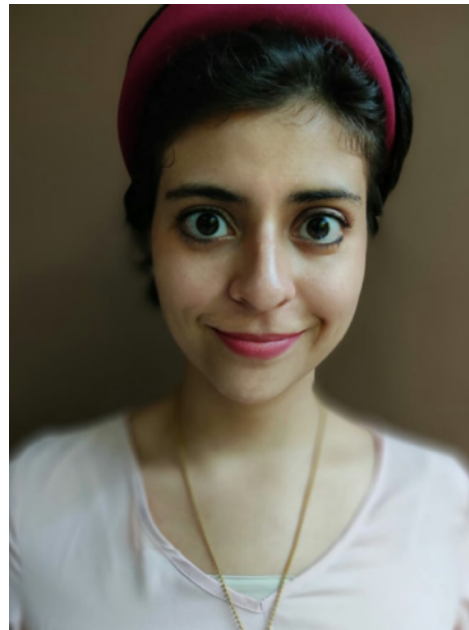
SEBASTIEN FROMENTEAU
ICF-UNAM



GROUP PARTICIPATING ON THIS WORK



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RODRIGUEZ OTERO
UNAM



MARCOS
TOLEDO ORTIZ
UAEM

Summary

- **Cosmology**

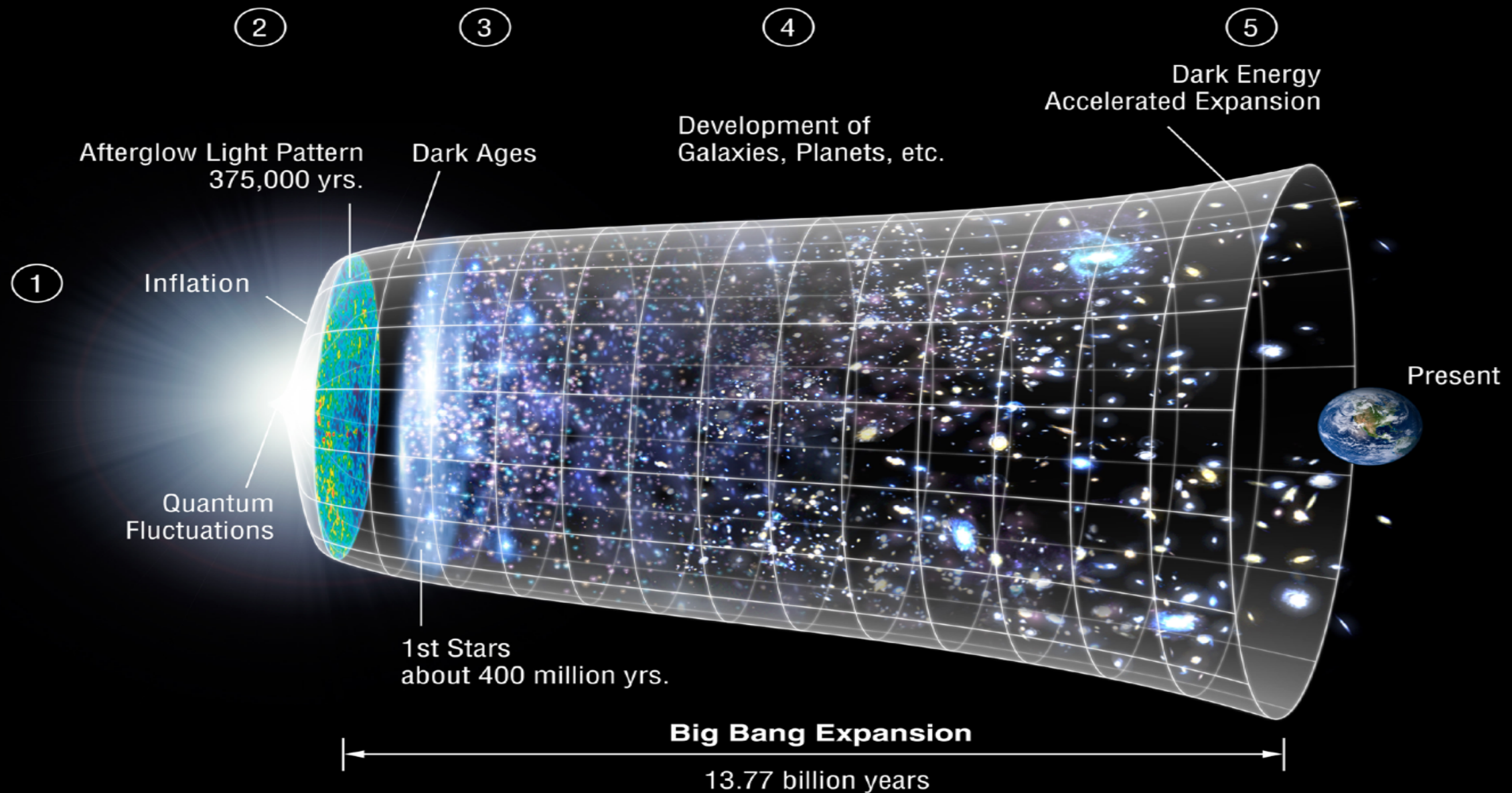
- Perturbations description and evolution
- Halo bias and Galaxy bias

- **Use of astrophysical local properties to enhance bias determination**

- Galaxy luminosity as a degeneracy break information
- A direct use for the CDM mapping - BAO Reconstruction
- Emission line galaxies

Cosmology Standard Model

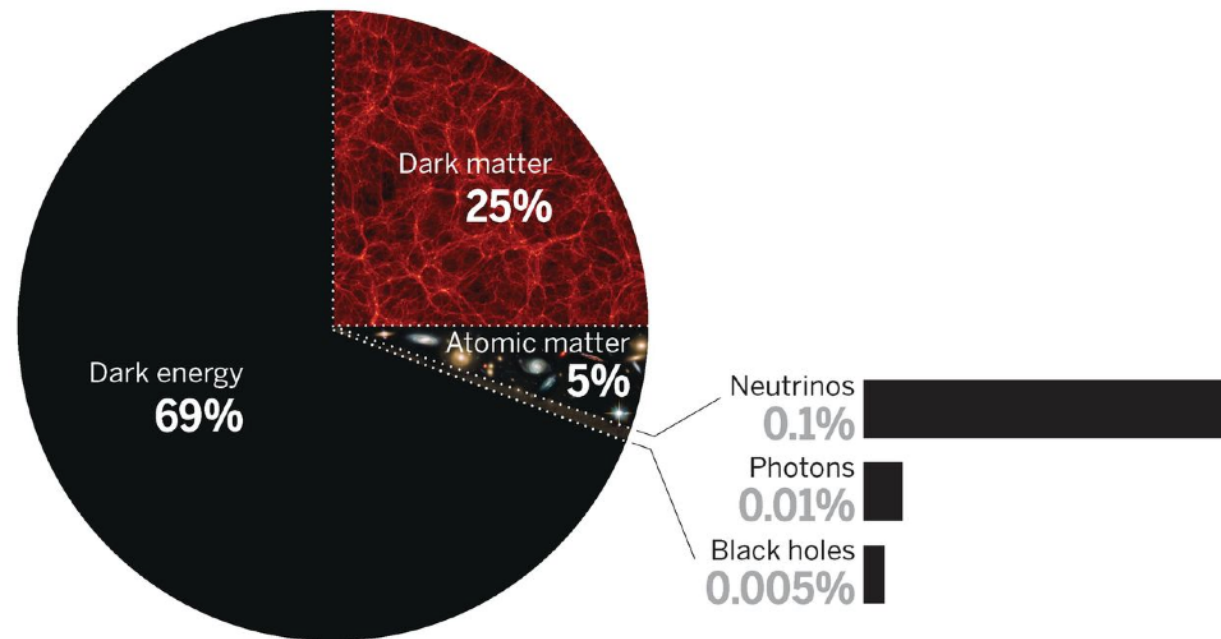
$$\Lambda - CDM$$



Cosmological Parameters

The multiple components that compose our universe

Current composition (as the fractions evolve with time)



$$\Omega_i, n_s, A_s, \tau, H_0$$

Perturbations

67.4 ± 0.5

73.9 ± 1.8

EARLY ROUTE

Cosmic Microwave Background
Galaxy clustering:
BAO
RSD
Galaxy weak lensing

LATE ROUTE

Supernovae Ia and all other
standardizable
observable based on
Cosmic Ladder

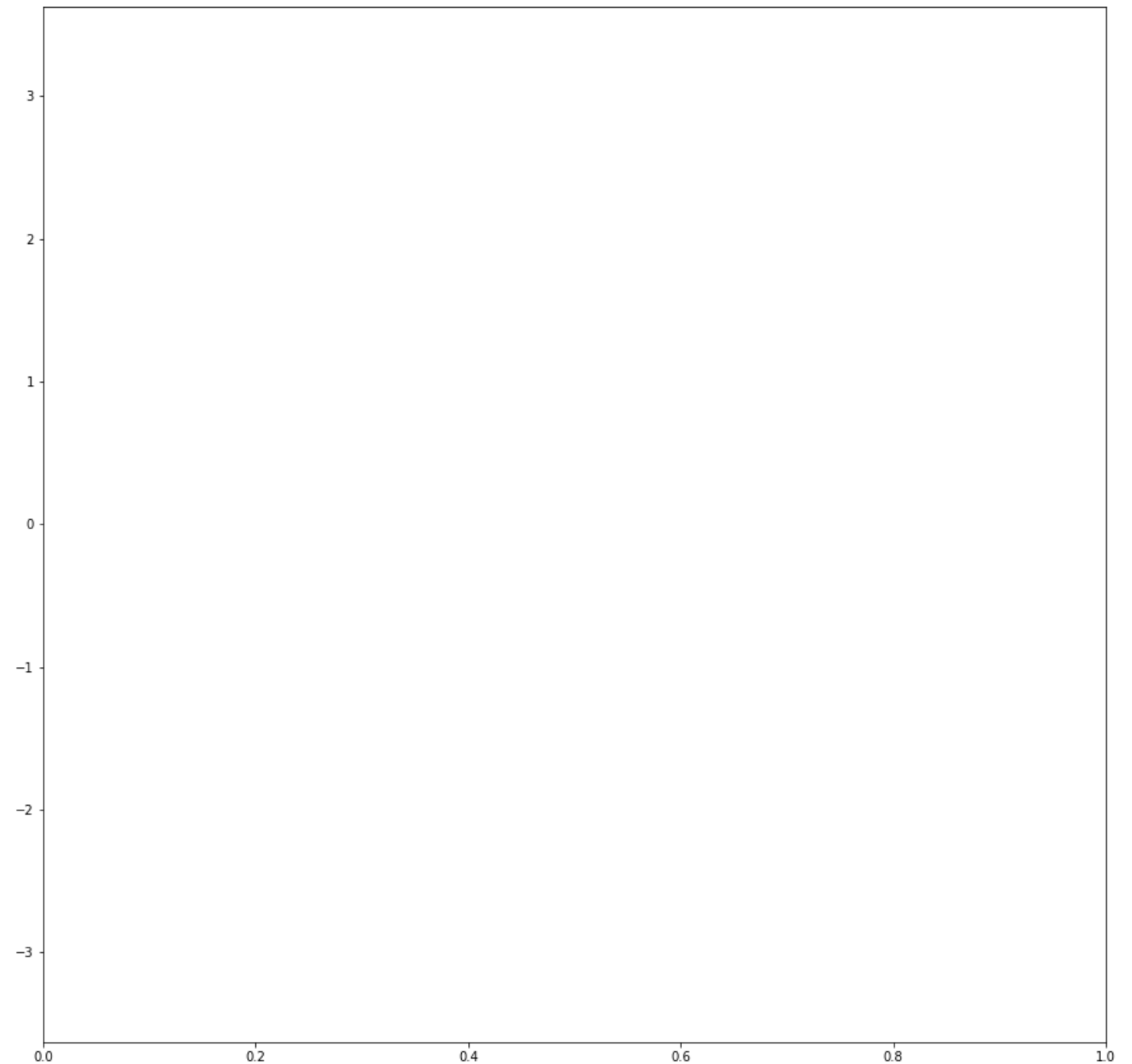
Quantum Fluctuations during inflation

Heisenberg uncertainty

$$\Delta x \times \Delta p \geq \hbar$$

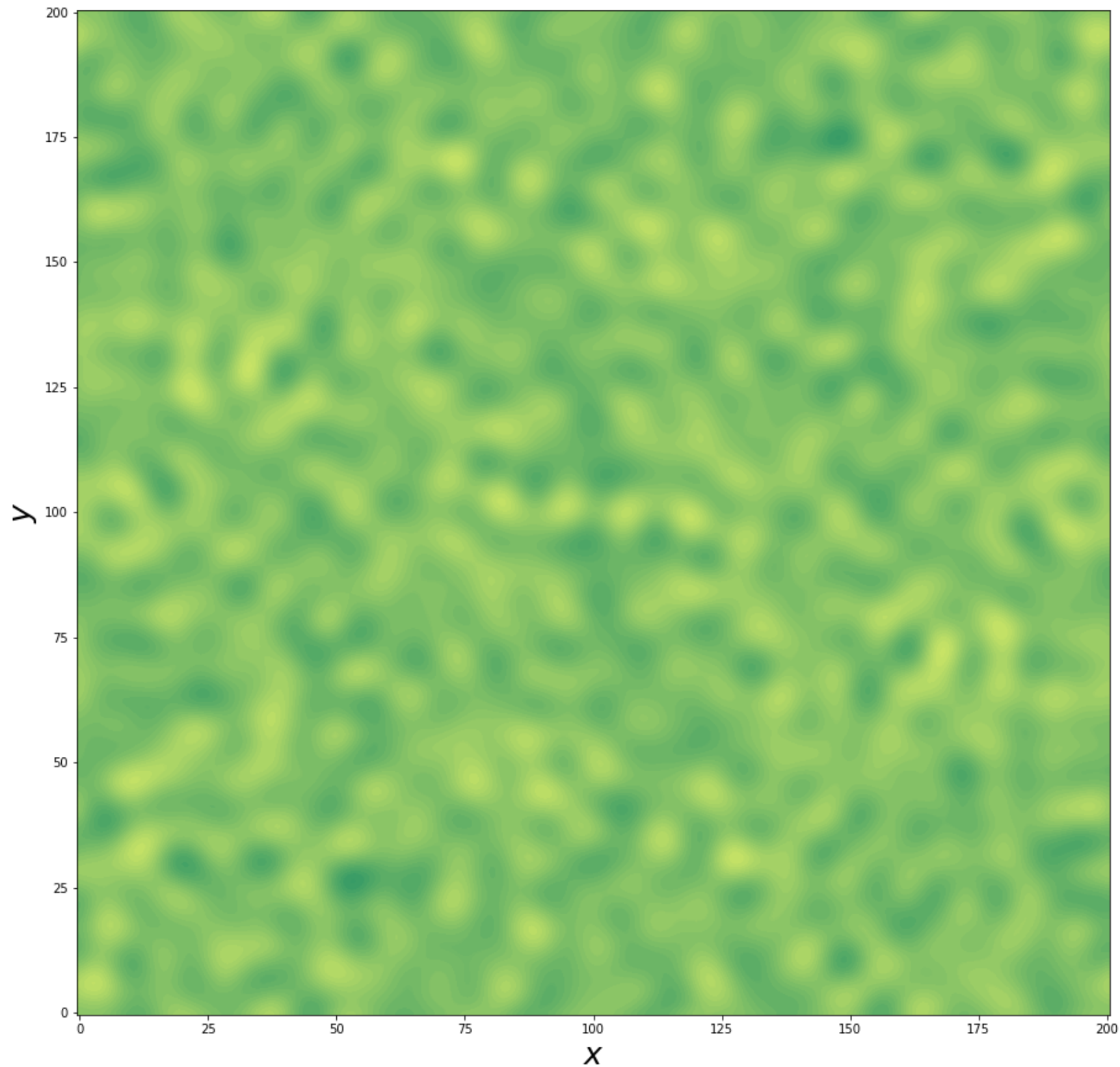


$$\Delta T \times \Delta E \geq \hbar$$

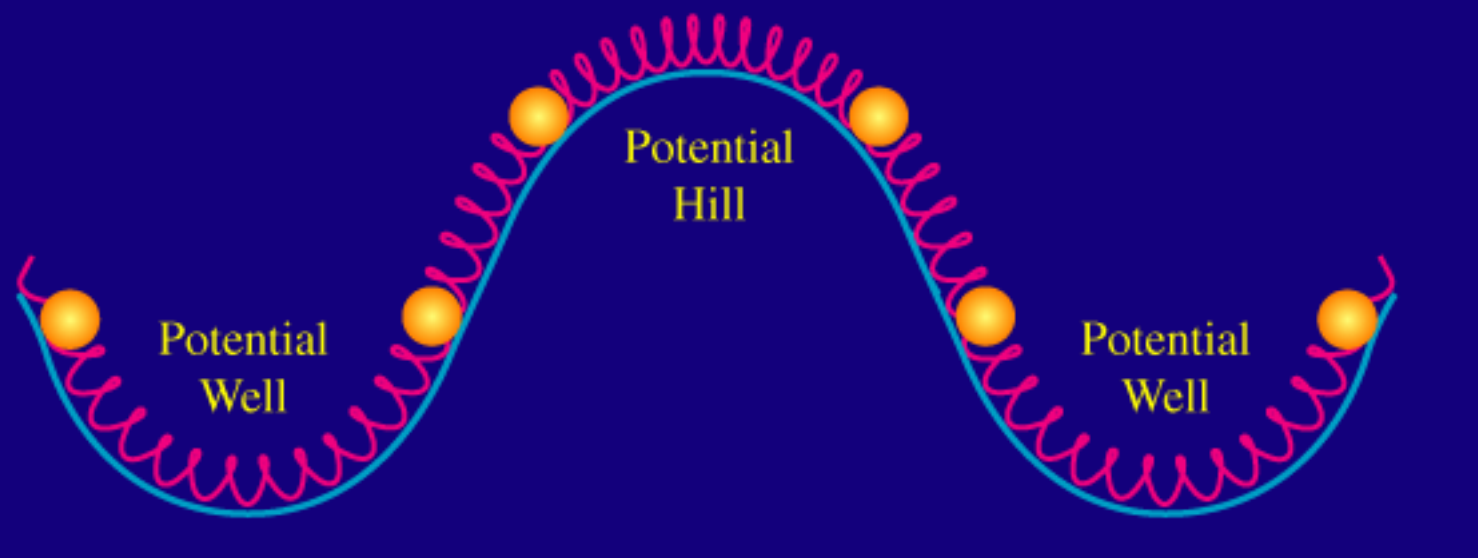


Freeze perturbations !!!!!

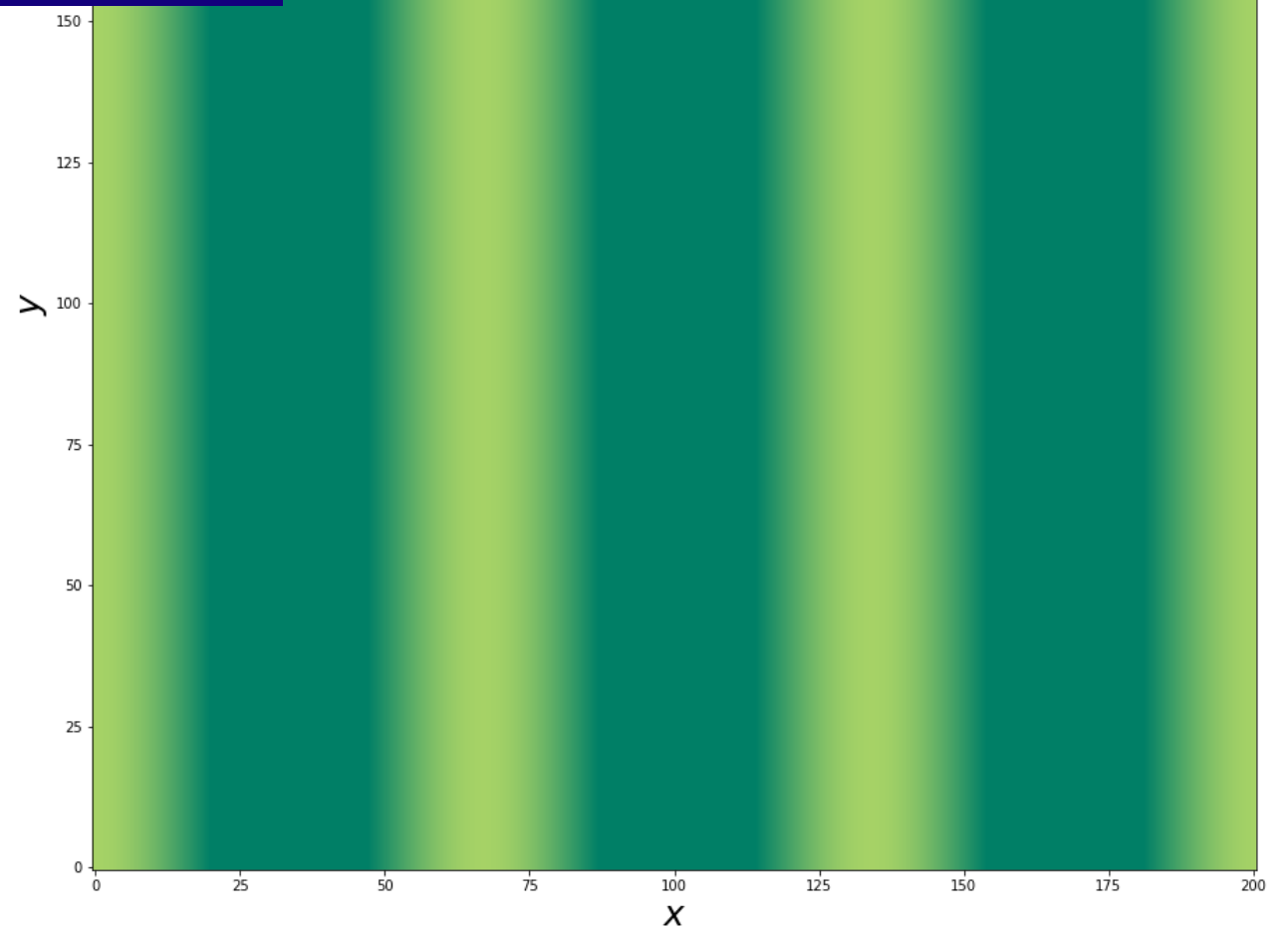
From Reheating to CMB

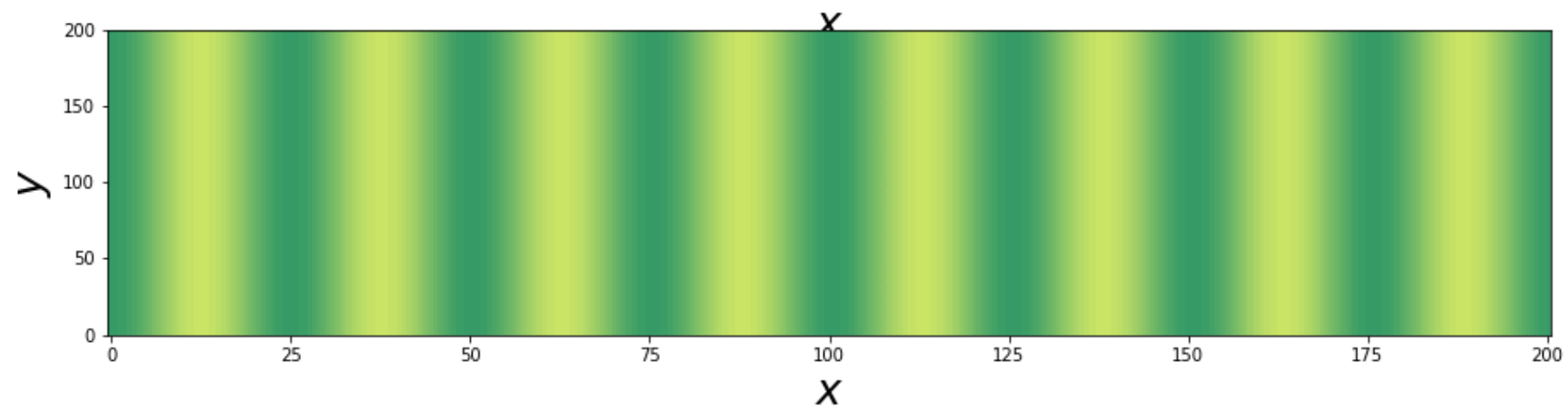
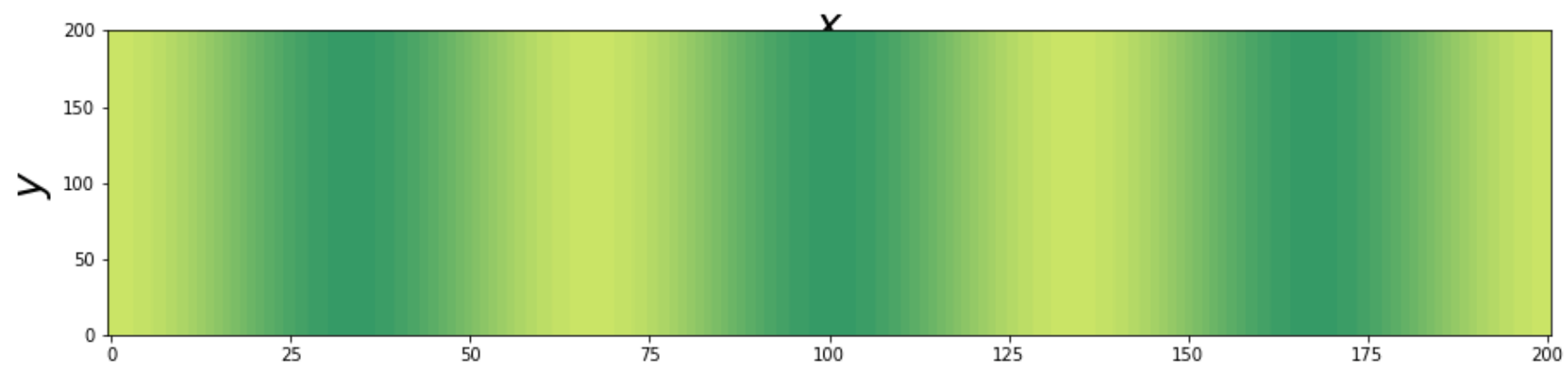
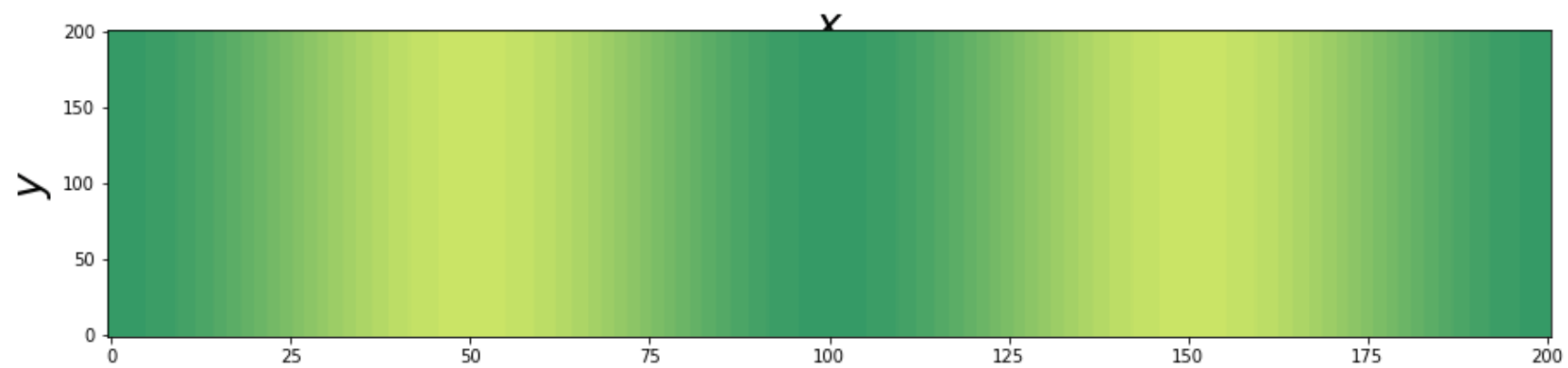
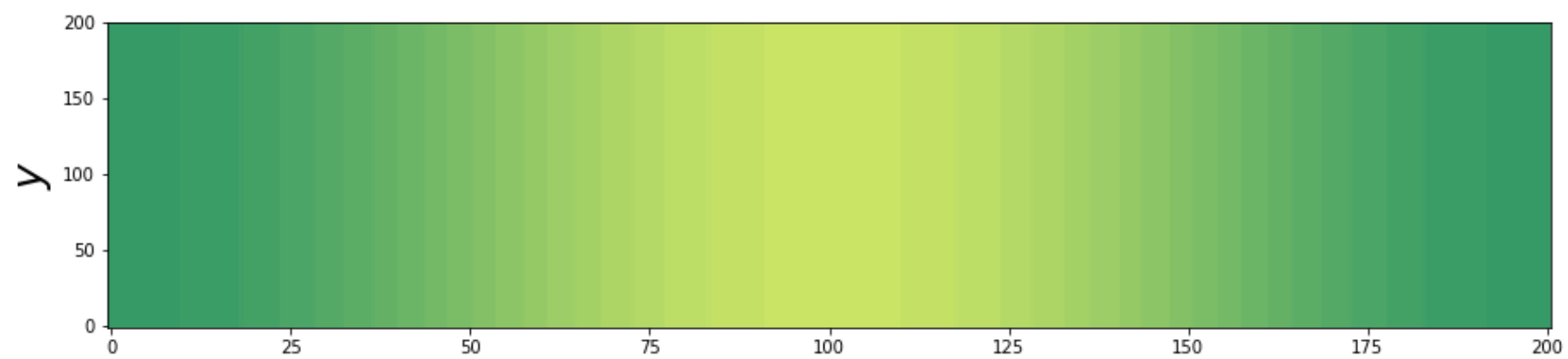


Baryon-photon oscillates

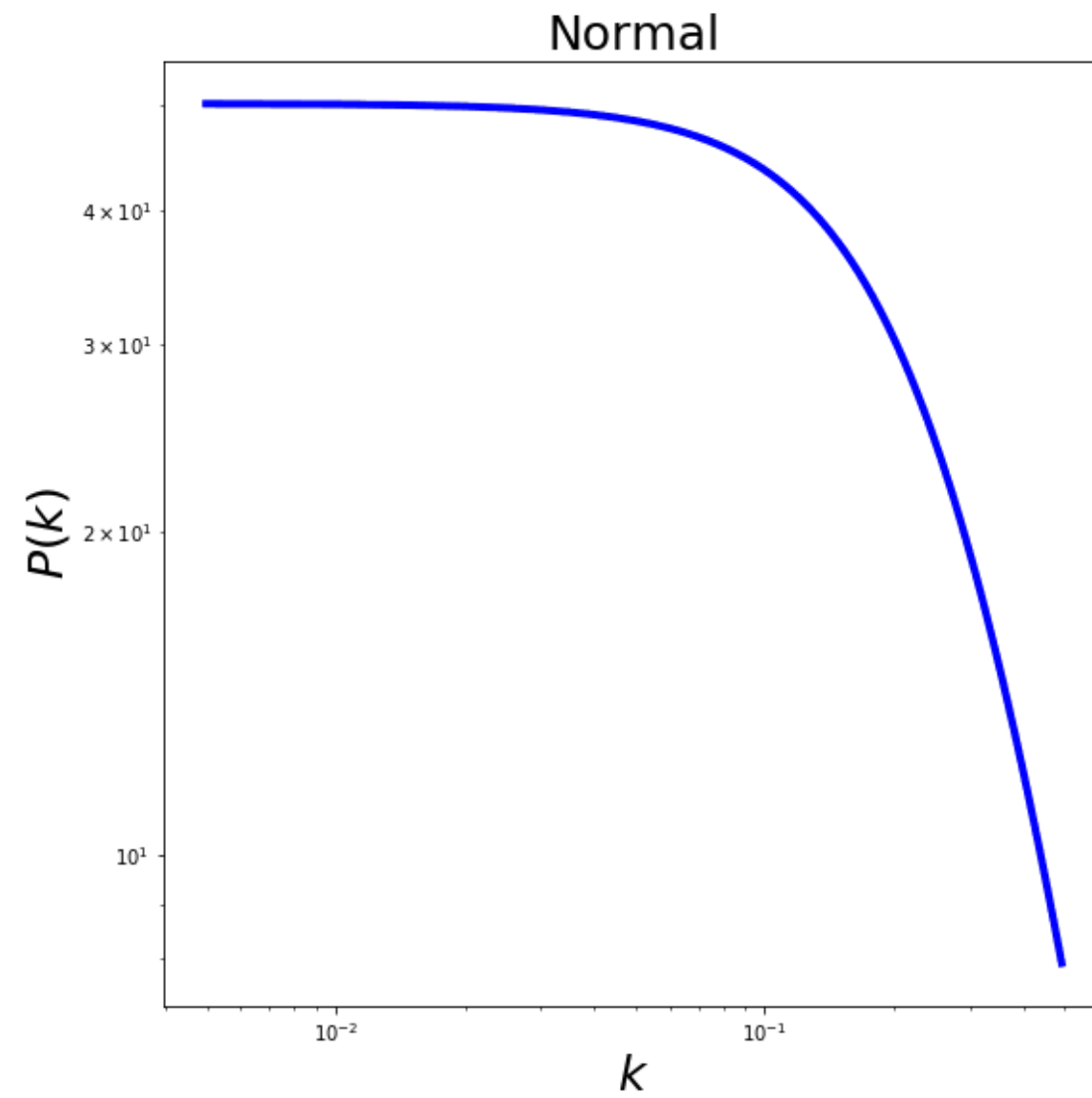
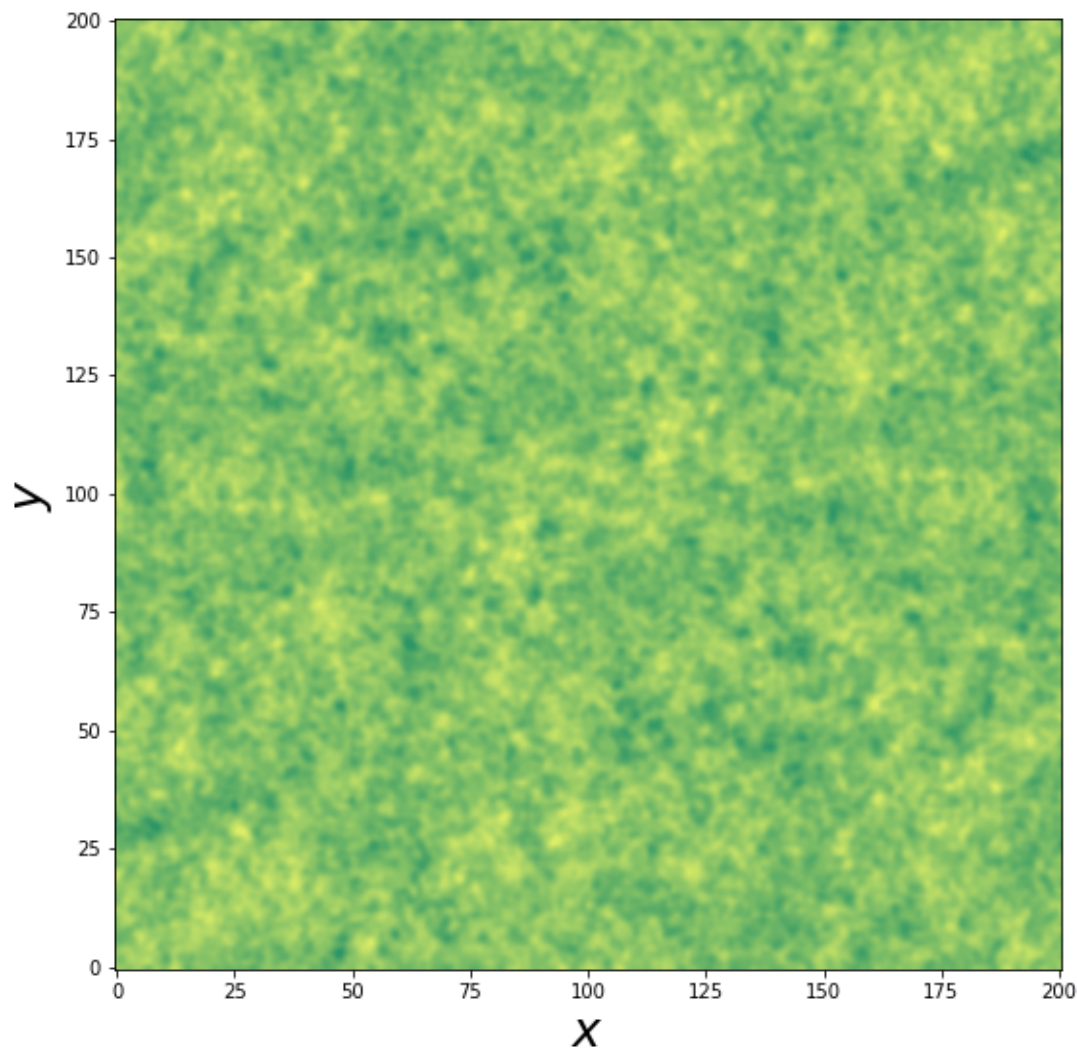


Wayne Hu's tutorial

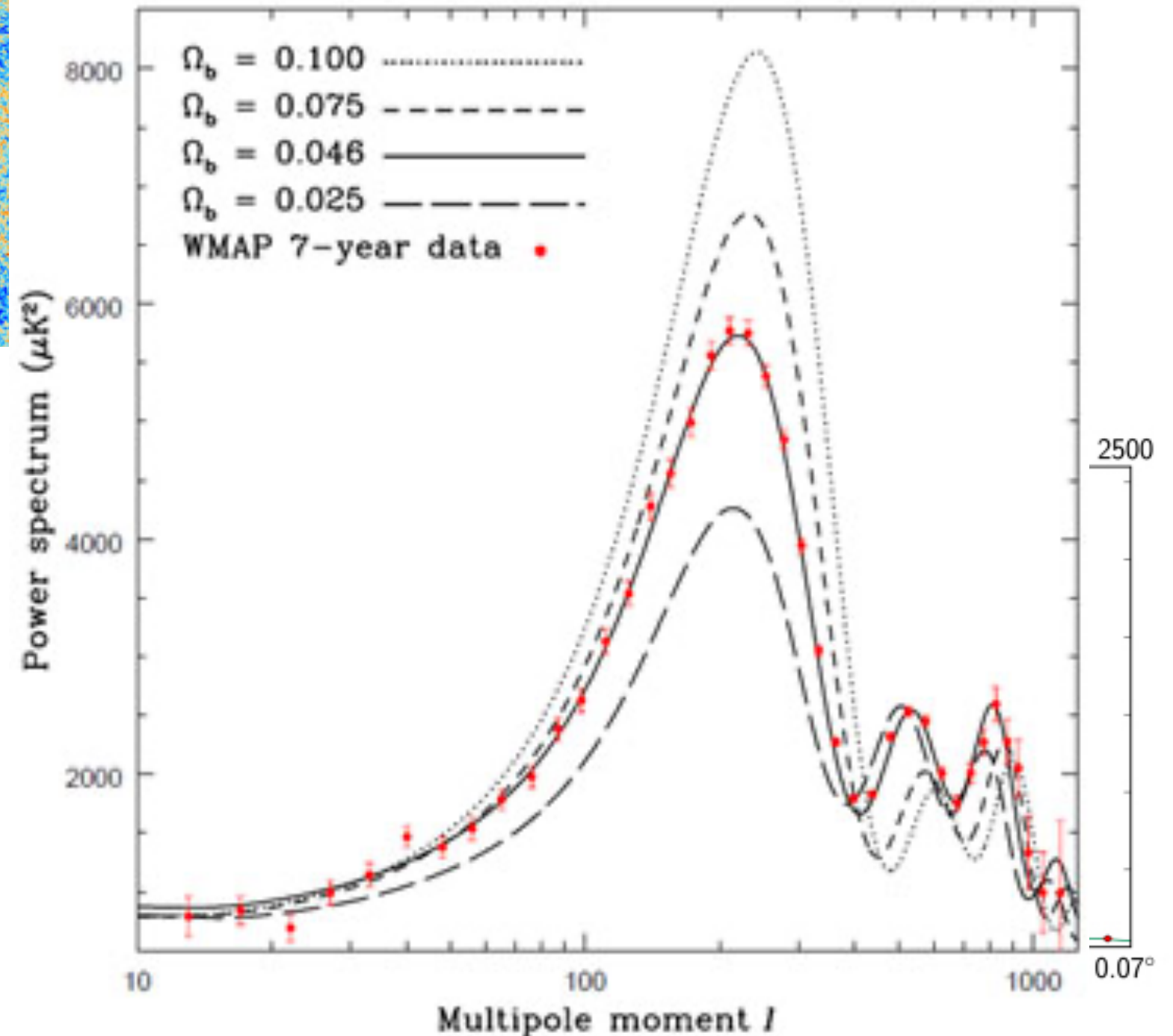
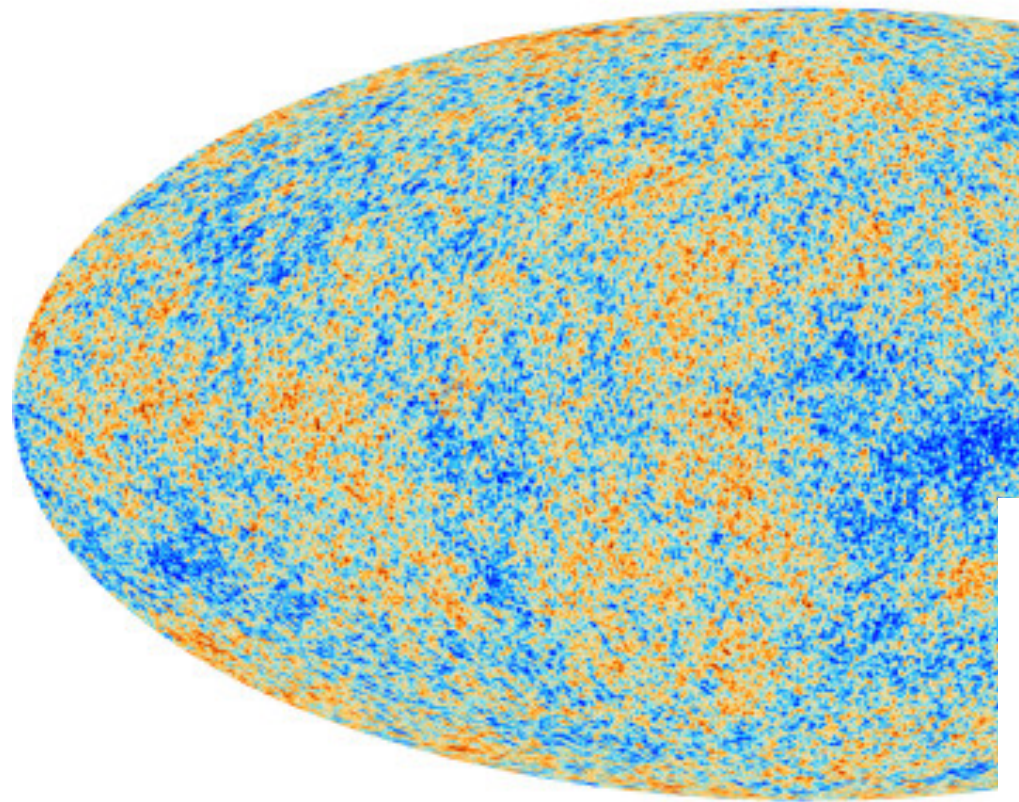




Fourier / Spherical Harmonics (CMB)



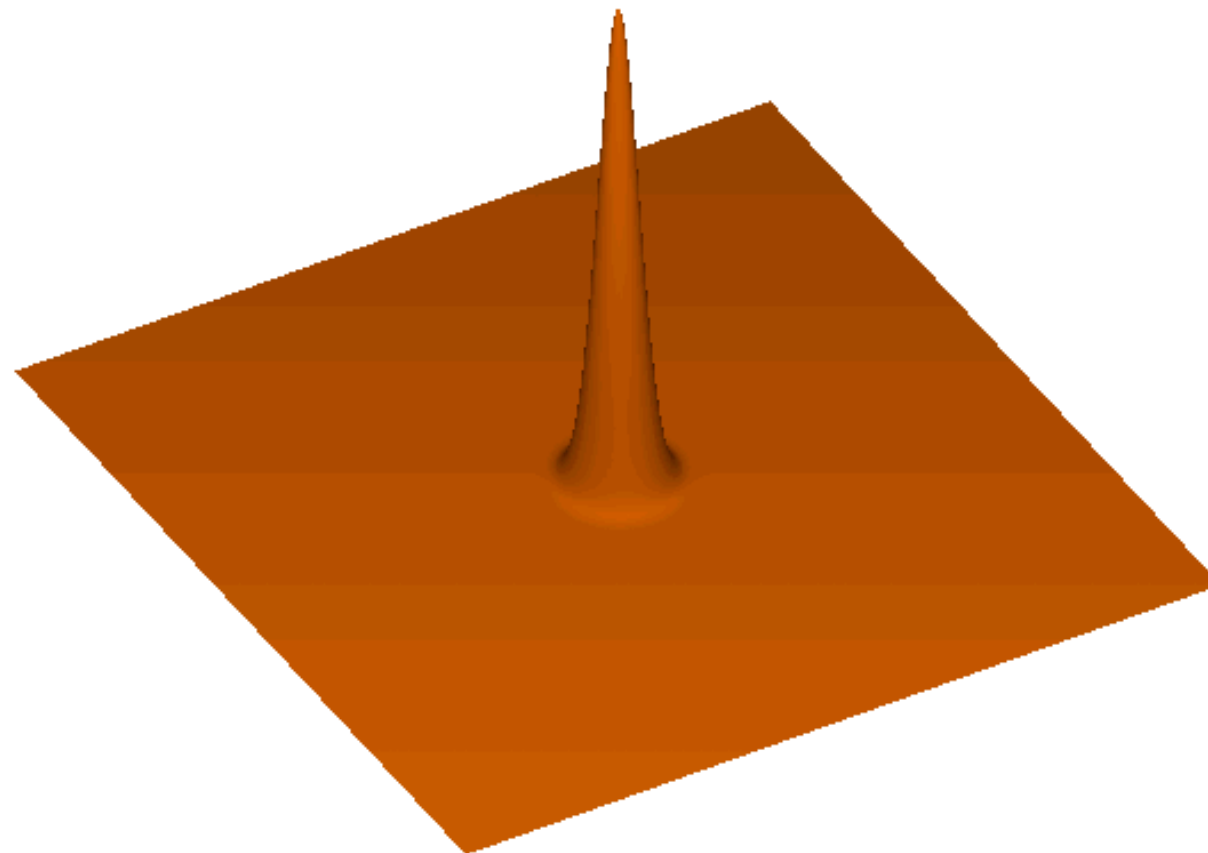
Cosmic Microwave Background (CMB)



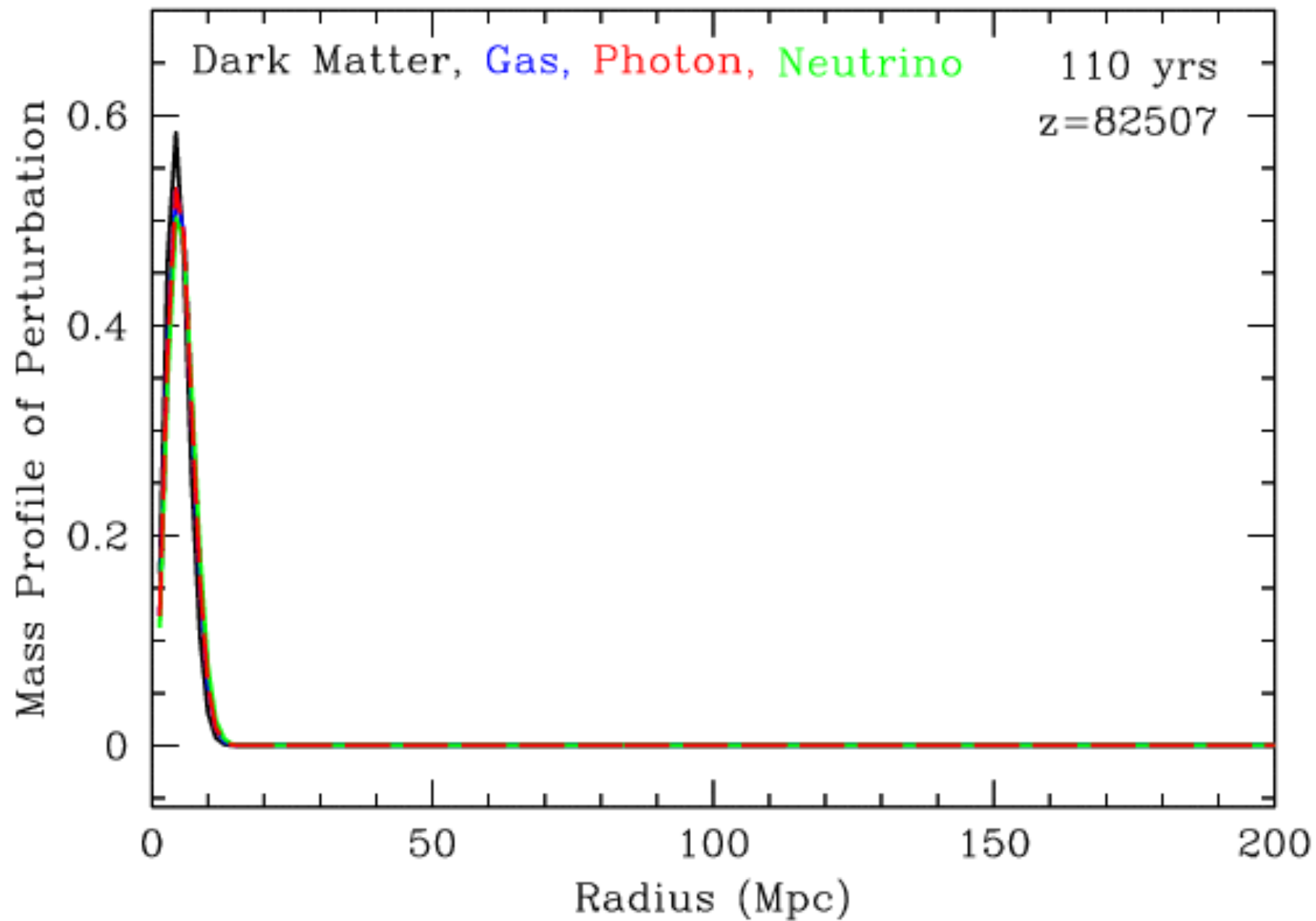
Baryonic Acoustic Oscillations (BAO) Real Space

Around all over-densities

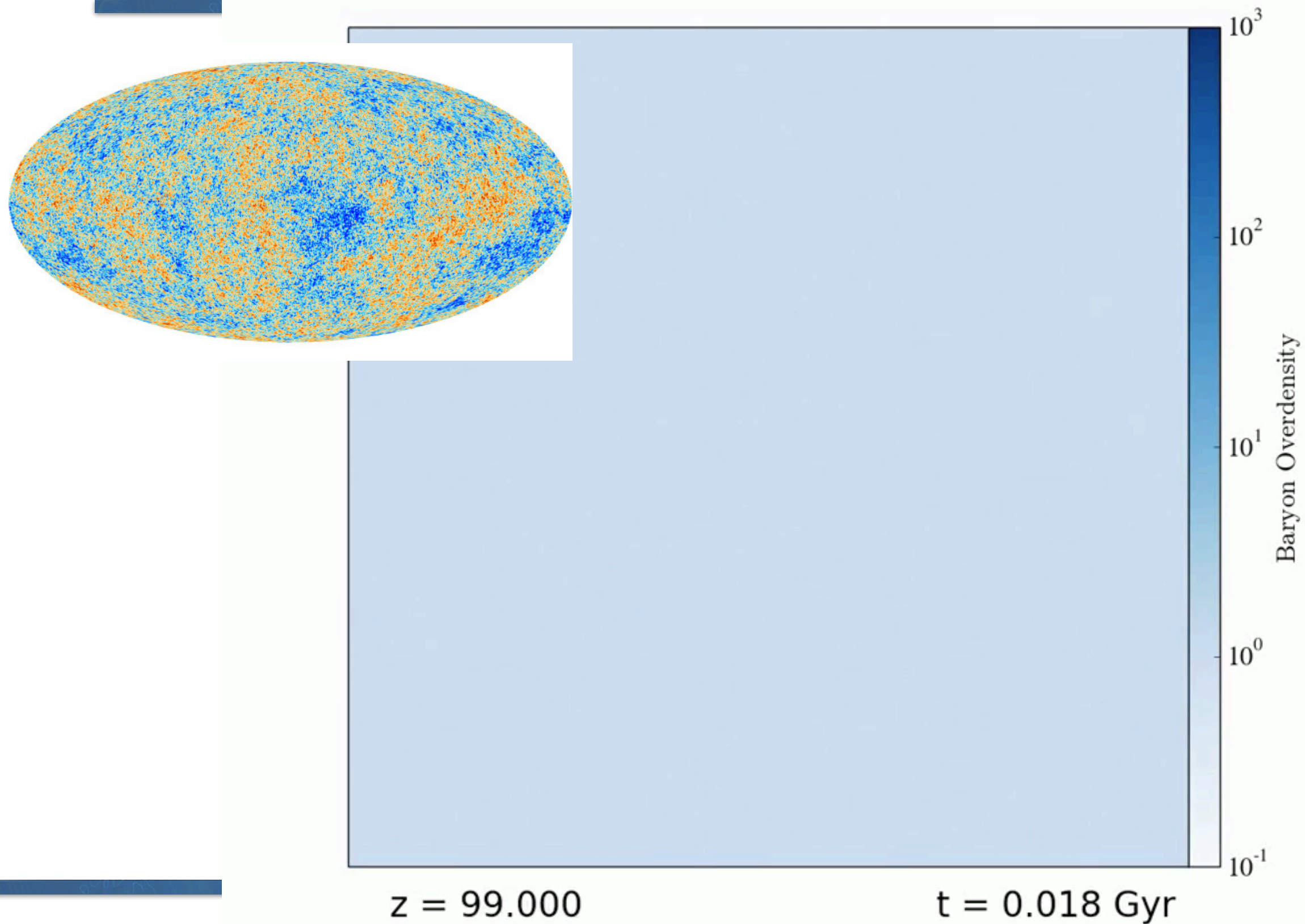
Stop to propagate when the plasma disappear



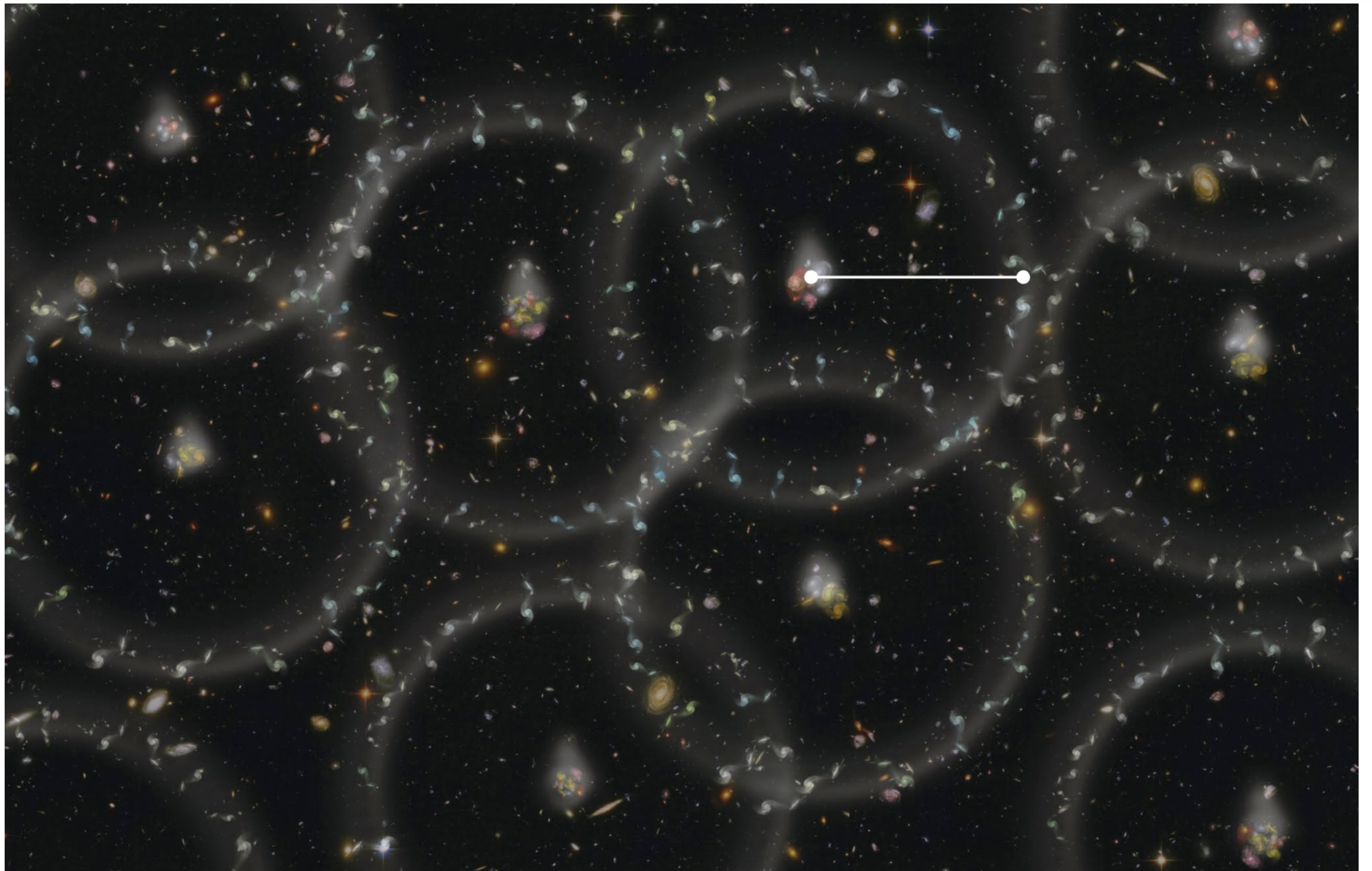
BAO in real space



Perturbations after CMB



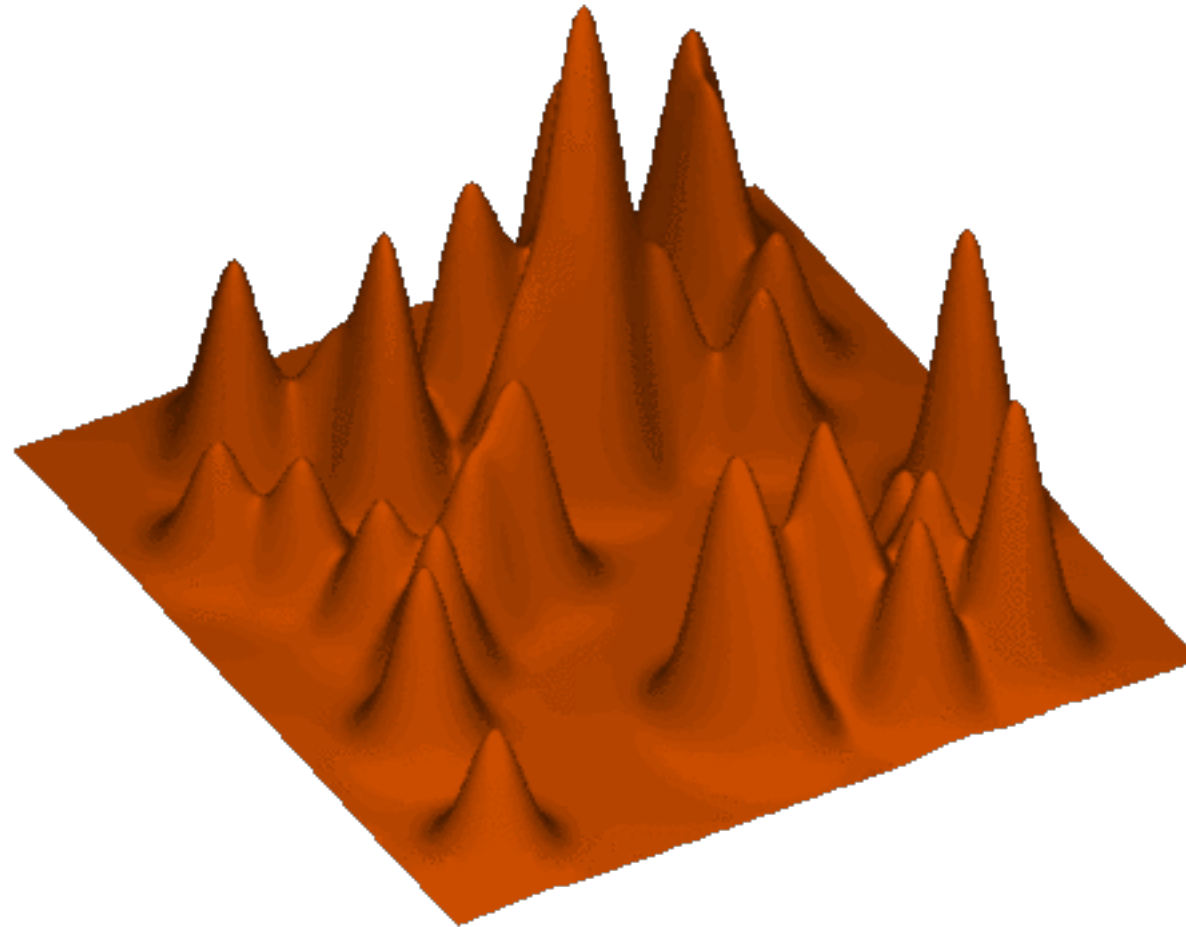
BAO in galaxies



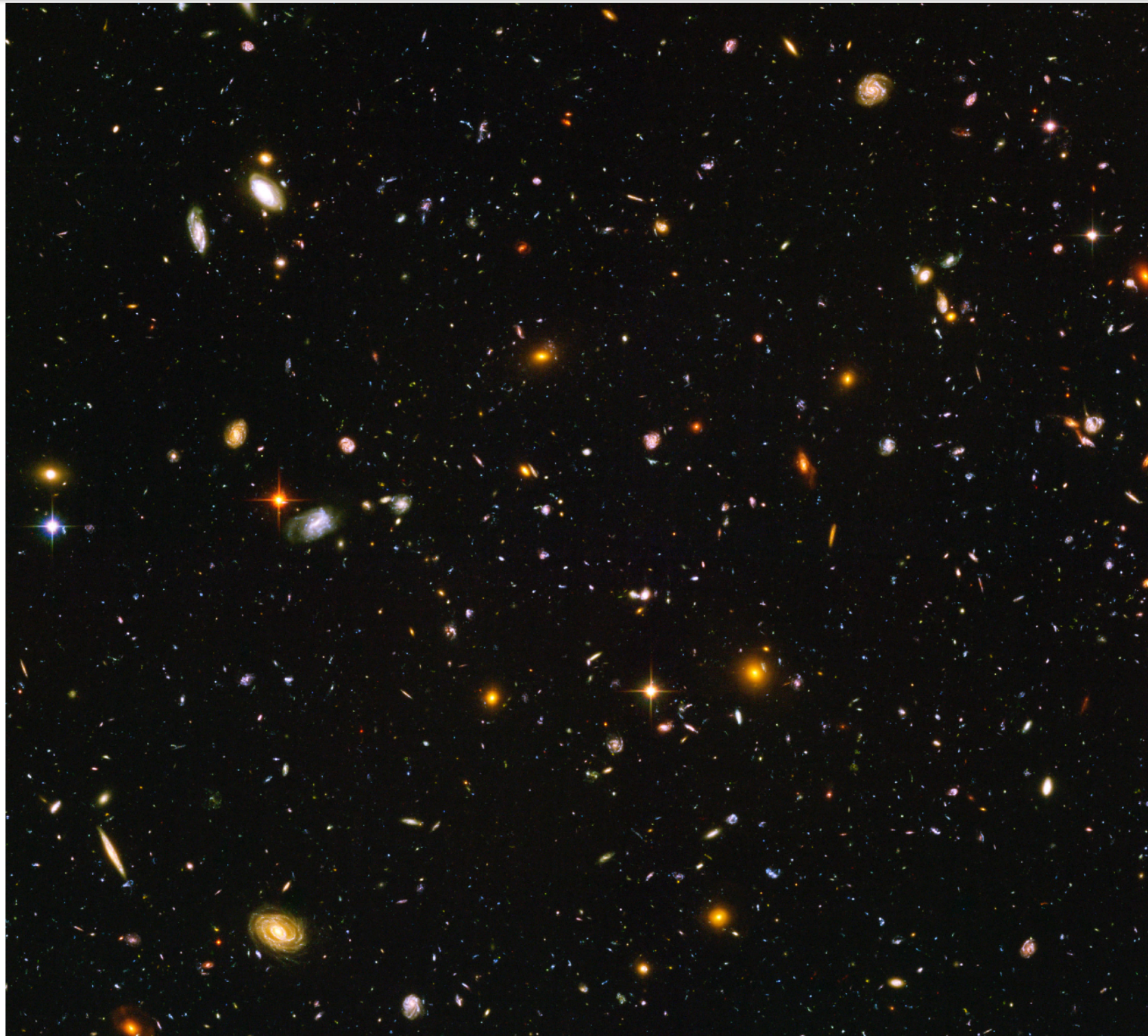
Baryonic Acoustic Oscillations (BAO) Real Space

Around all over-densities

Stop to propagate when the plasma disappear



Reality



Measuring Clustering

$$dP_{12} = \bar{n}_g^2 [1 + \xi(\vec{r}_{12})] dV_1 dV_2$$

We estimate it using galaxy pair counts



2-pt statistics

$$\xi(r) = \langle \delta(x)\delta(x+r) \rangle_{\text{volumen}}$$

$\delta(x)$ is the contrast density

$$\delta(x) = \frac{\rho(x) - \langle \rho \rangle}{\langle \rho \rangle}$$

Power Spectrum $P(k)$ is the FT of $\xi(r)$:

$$P(k) = \int d^3x e^{-ik \cdot r} \xi(r)$$

Power Spectrum and 2pt-correlation function

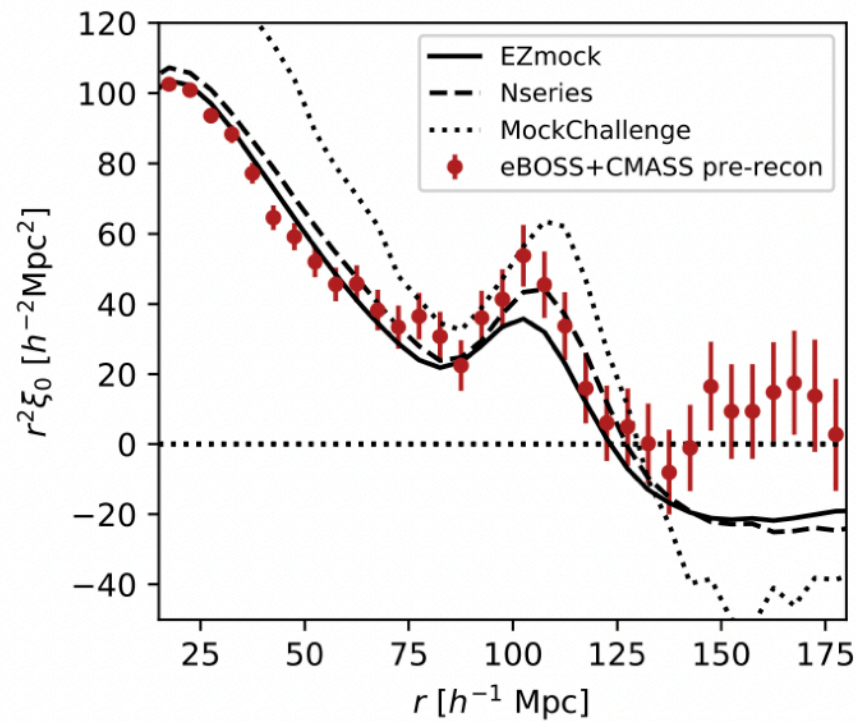
$$\xi(\vec{r}) = \langle \delta(\vec{x}) \cdot \delta(\vec{x} + \vec{r}) \rangle$$

$$P(\vec{k}) = \frac{1}{(2\pi)^3} \int d^3\vec{r} \xi(\vec{r}) e^{-i\vec{k} \cdot \vec{r}}$$

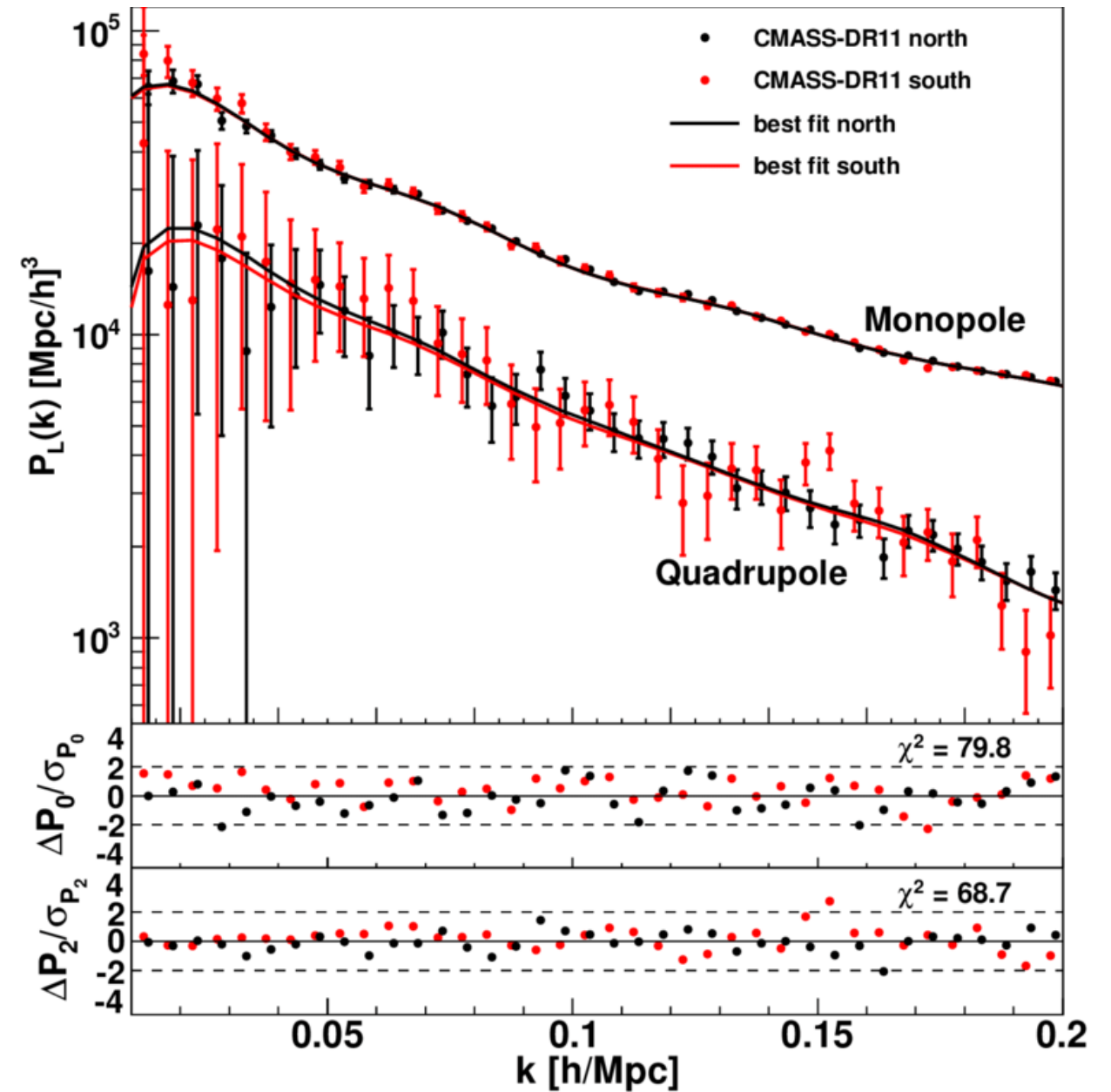
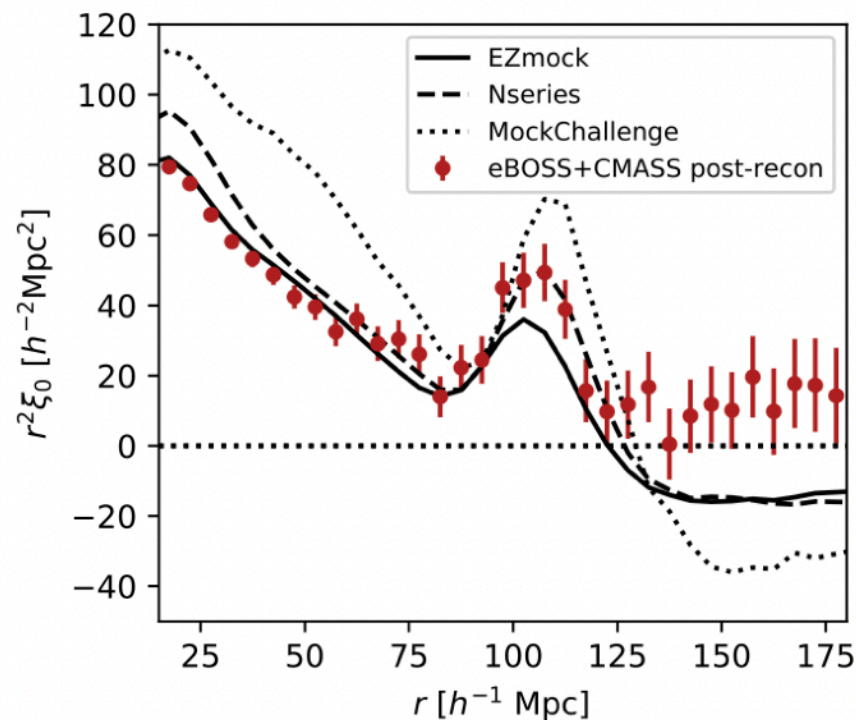
Real field condition : $\delta(-\vec{k}) = \delta^*(\vec{k})$

$$(2\pi)^3 P(\vec{k}) = \langle \delta(\vec{k}) \cdot \delta^*(\vec{k}) \rangle = \langle |\delta(\vec{k})|^2 \rangle$$

Galaxy 2pt-statistics measurement

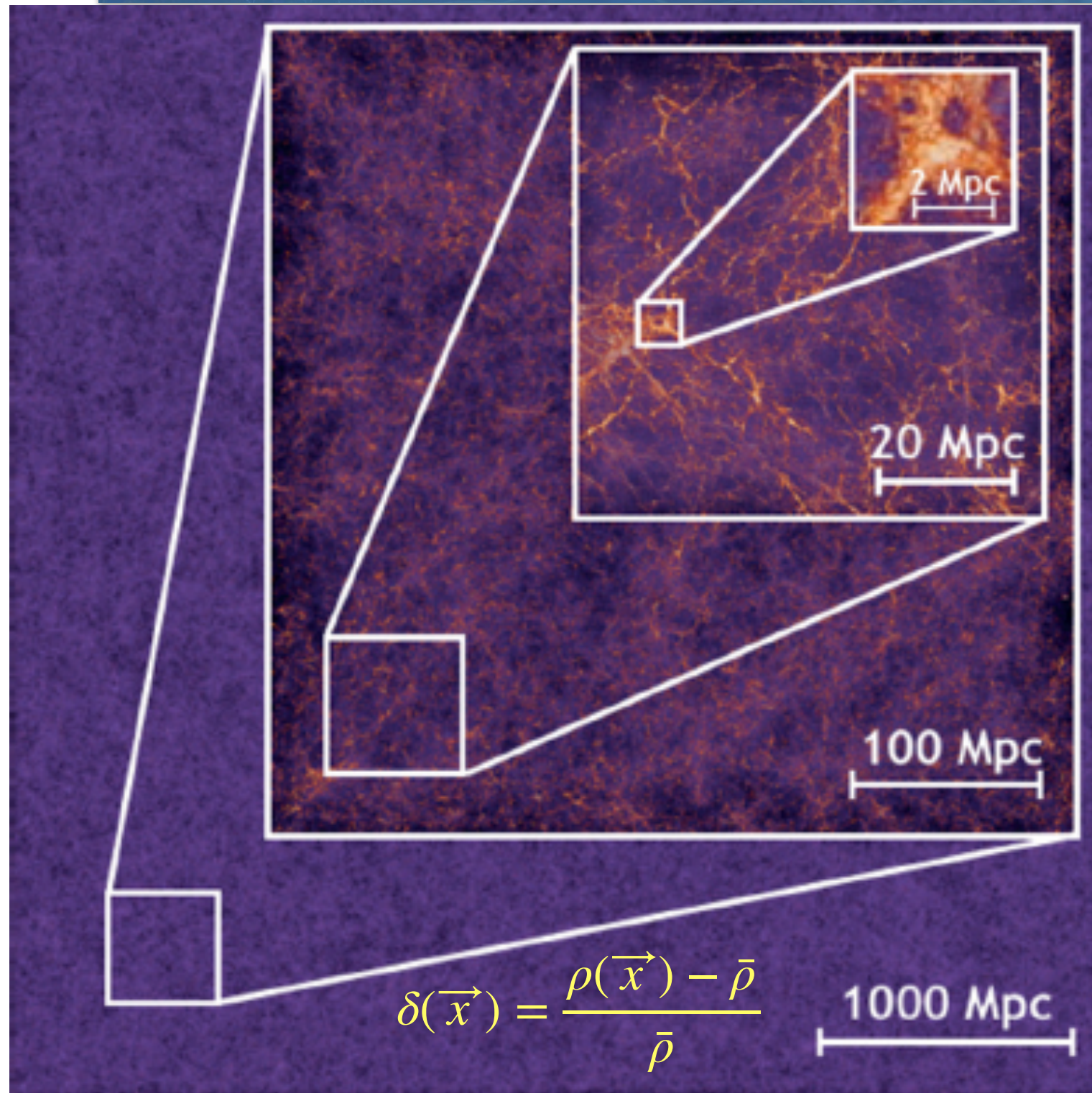


Bautista et al. 2021



Beutler et al. 2013

Scales in Cosmology



$\delta > 1$: Non-Linear

- **Simulations**

$\delta \sim 1$: Quasi-Linear

- **Simulations**
- **2-LPT**

$\delta \ll 1$: Linear

- **Simulations**
- **LPT** (Lagrangian Pert. Th)
- **Analytical approx.**

Linear evolution of Perturbations

$$\dot{\rho} + \nabla_{\vec{r}} \cdot (\rho \vec{u}) = 0,$$

: Continuity equation

$$\rho \left[\dot{\vec{u}} + (\vec{u} \cdot \nabla_{\vec{r}}) \vec{u} \right] = -\nabla_{\vec{r}} p - \rho \nabla_{\vec{r}} \Phi,$$

: Euler's equation

$$\nabla_{\vec{r}}^2 \Phi = 4\pi G(\rho + 3p) - \Lambda.$$

: Poisson's equation

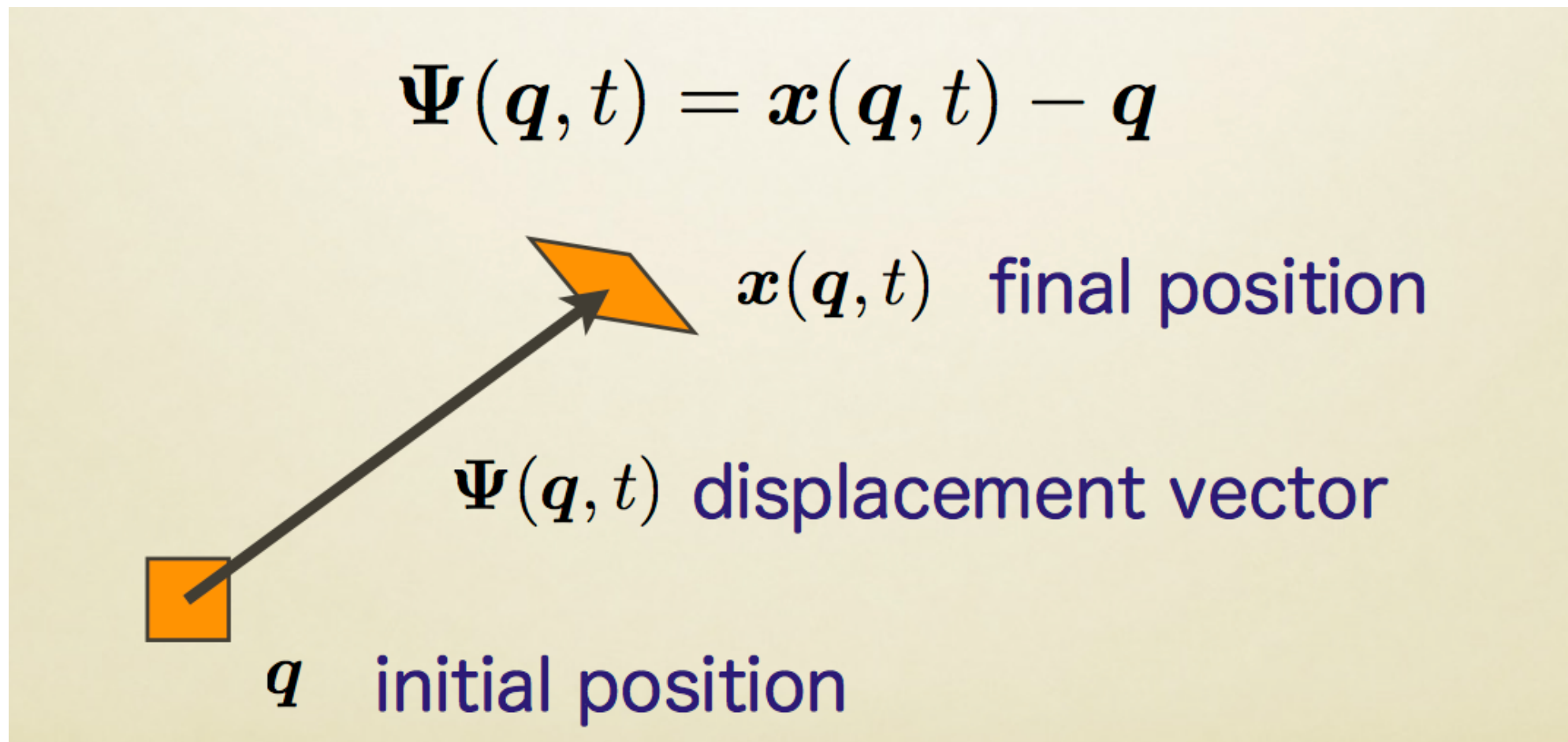
$$\ddot{\delta} + 2H\dot{\delta} = 4\pi G\rho_m\delta$$

**Solution for whole matter
as Dark Matter**

$$\delta = \frac{\rho - \bar{\rho}}{\bar{\rho}} \ll 1$$

linearity condition

Lagrangian Perturbation Theory



Credit to Matsubara's presentation

Lagrangian Perturbation Theory

Equation of motion & Poisson's Equation

$$\ddot{\Psi} + \frac{\dot{a}}{a} \dot{\Psi} = -\frac{1}{a^2} \nabla_x \phi$$
$$\Delta_x \phi = 4\pi G \bar{\rho} a^2 \delta(x, t)$$

Linearization & Zeldovich Approximation

$$\delta(x, t) = \left[\det \left(I + \frac{\partial \Psi}{\partial q} \right) \right]^{-1} - 1 \approx -\nabla_q \cdot \Psi$$

$$\Psi \approx -D(t) \nabla_q \rho_0(q)$$

Credit to Lile Wang's presentation

Lagrangian Perturbation Theory

- Taking into account the higher-order perturbations in the displacement

$$\Psi = \sum_{n=1}^{\infty} \Psi^{(n)} = \Psi^{(1)} + \Psi^{(2)} + \Psi^{(3)} + \dots$$



$$\Psi^{(1)} = -D(t) \nabla \varphi_0(\mathbf{q})$$

(First order: Zel'dovich approx.)

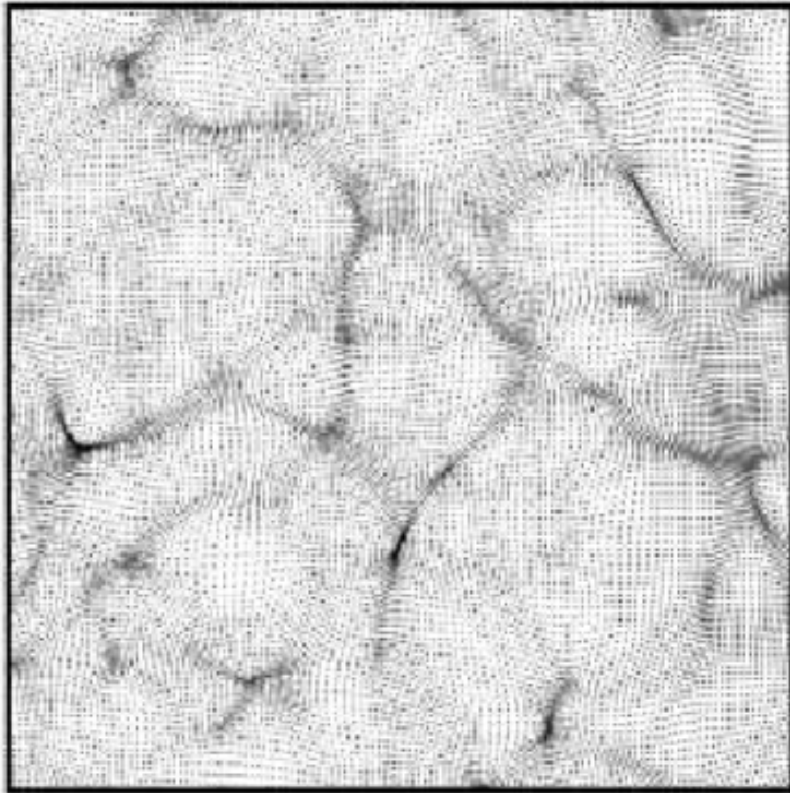
$$\Psi^{(2)} = -\frac{1}{2} D_2(t) \nabla \Delta^{-1} \left[\Psi_{i,i}^{(1)} \Psi_{j,j}^{(1)} - \Psi_{i,j}^{(1)} \Psi_{i,j}^{(1)} \right]$$

$$\begin{aligned} \Psi^{(3)} = -\frac{1}{3!} \left[D_{3a}(t) \nabla \Delta^{-1} \left(\Psi_{i,i}^{(1)} \Psi_{j,j}^{(2)} - \Psi_{i,j}^{(1)} \Psi_{i,j}^{(2)} \right) + D_{3b}(t) \nabla \Delta^{-1} \det \left(\Psi_{i,j}^{(1)} \right) \right. \\ \left. + D_{3c}(t) \Delta^{-1} \left(\Psi_{,j}^{(1)} \Psi_{i,j}^{(2)} - \Psi_{i,j}^{(1)} \Psi_{,j}^{(2)} \right)_{,i} \right] \end{aligned}$$

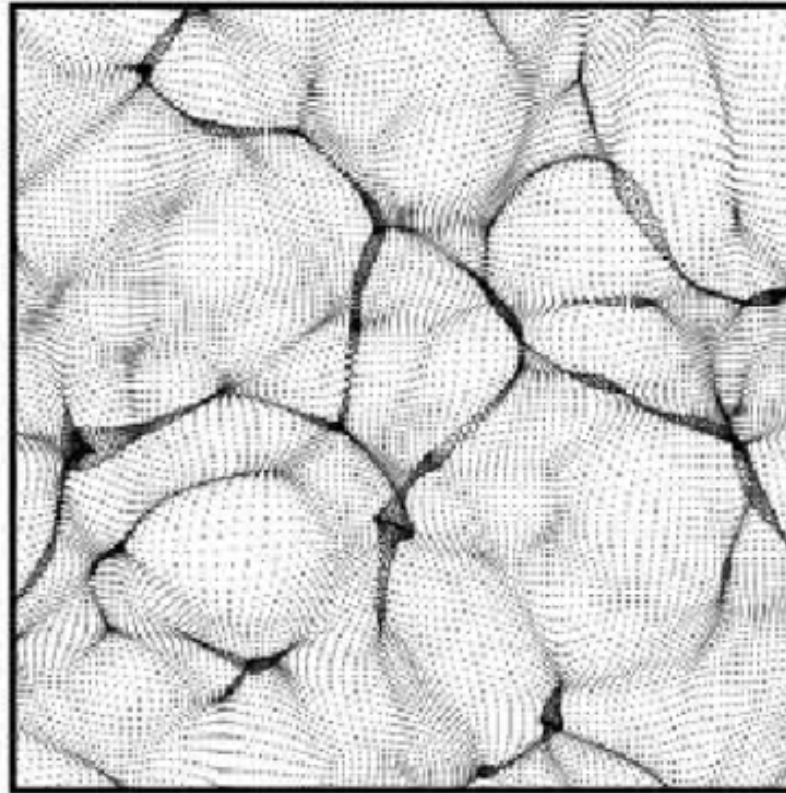
Recursive solutions

Credit to Matsubara's presentation

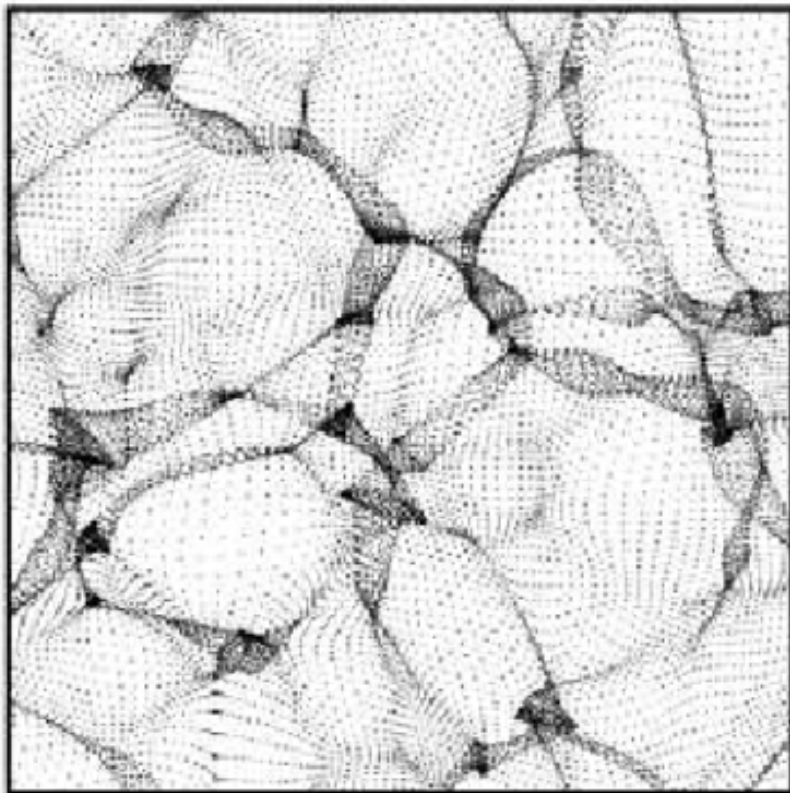
$\sigma = 0.5$



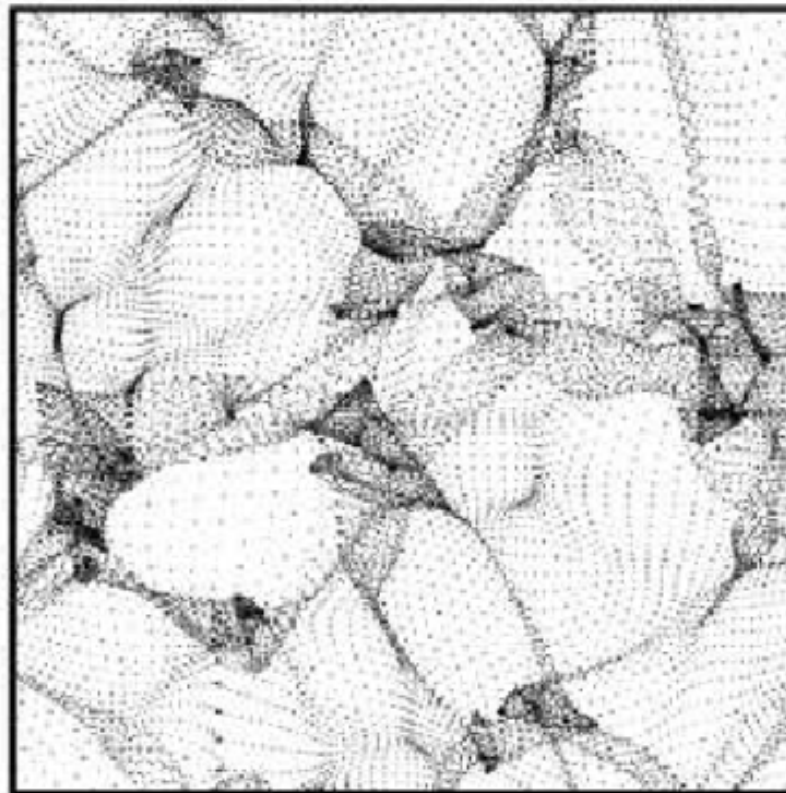
$\sigma = 1.$



$\sigma = 1.5$



$\sigma = 2.$



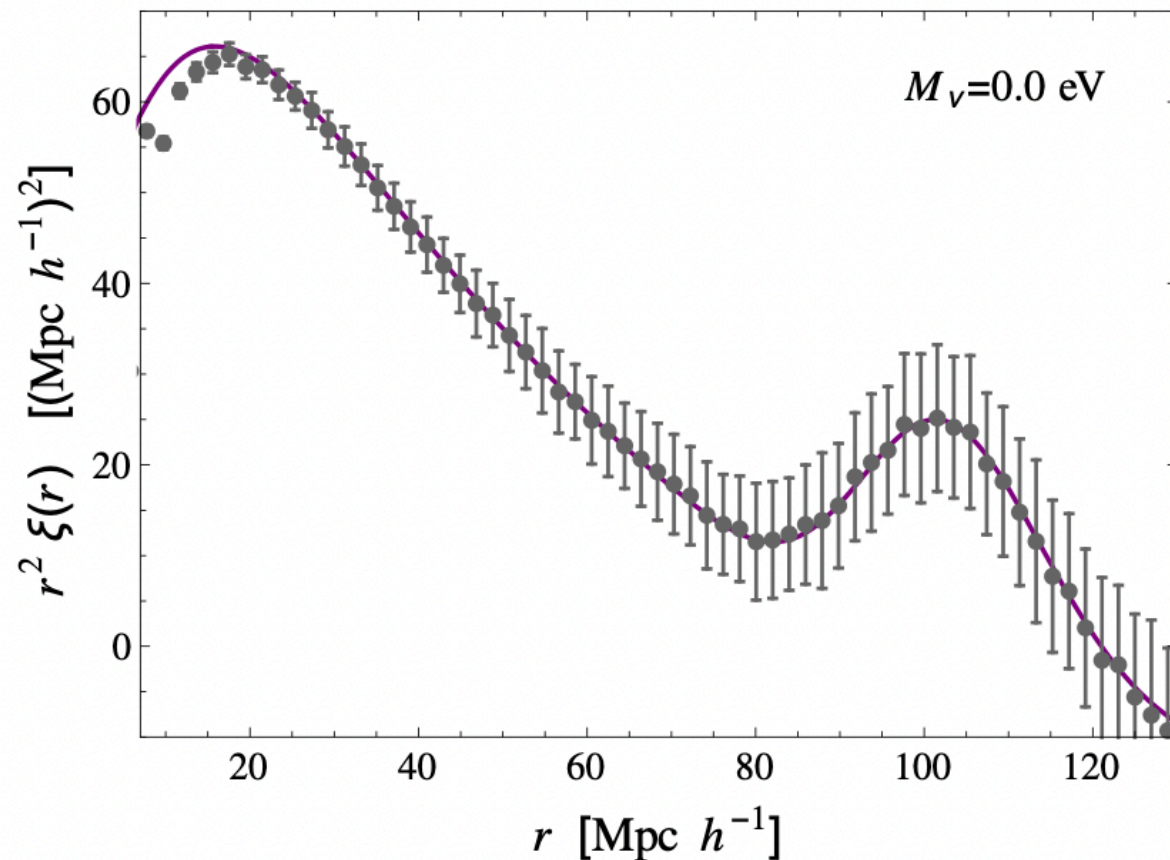
Applying the first order displacement field on initial particle positions.

Perturbation theory can evaluate in average how these displacements modify the correlation function and Power Spectrum

Perturbation Theory on stats

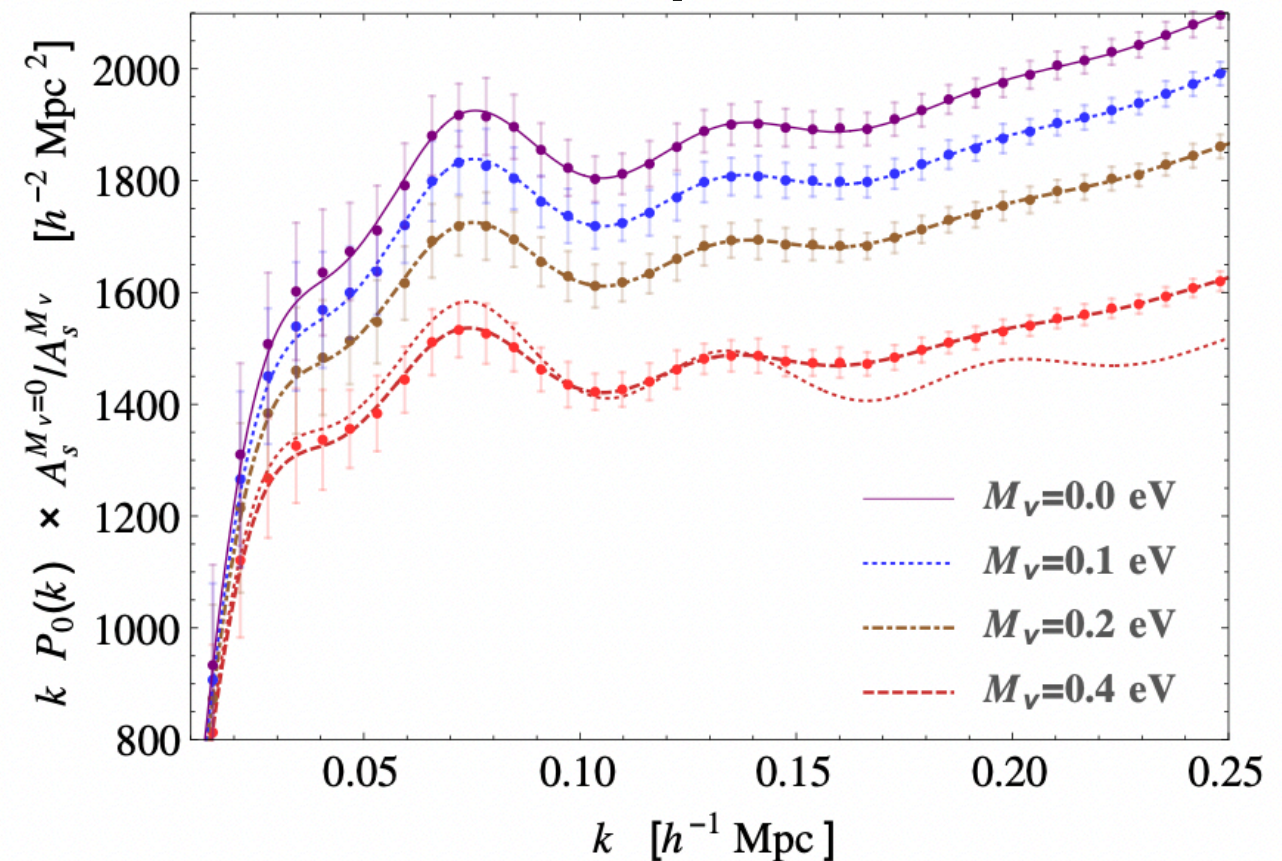
V_s CDM simulations

2pt-corr function



A, Aviles & A, Banerjee
2020

Power spectrum



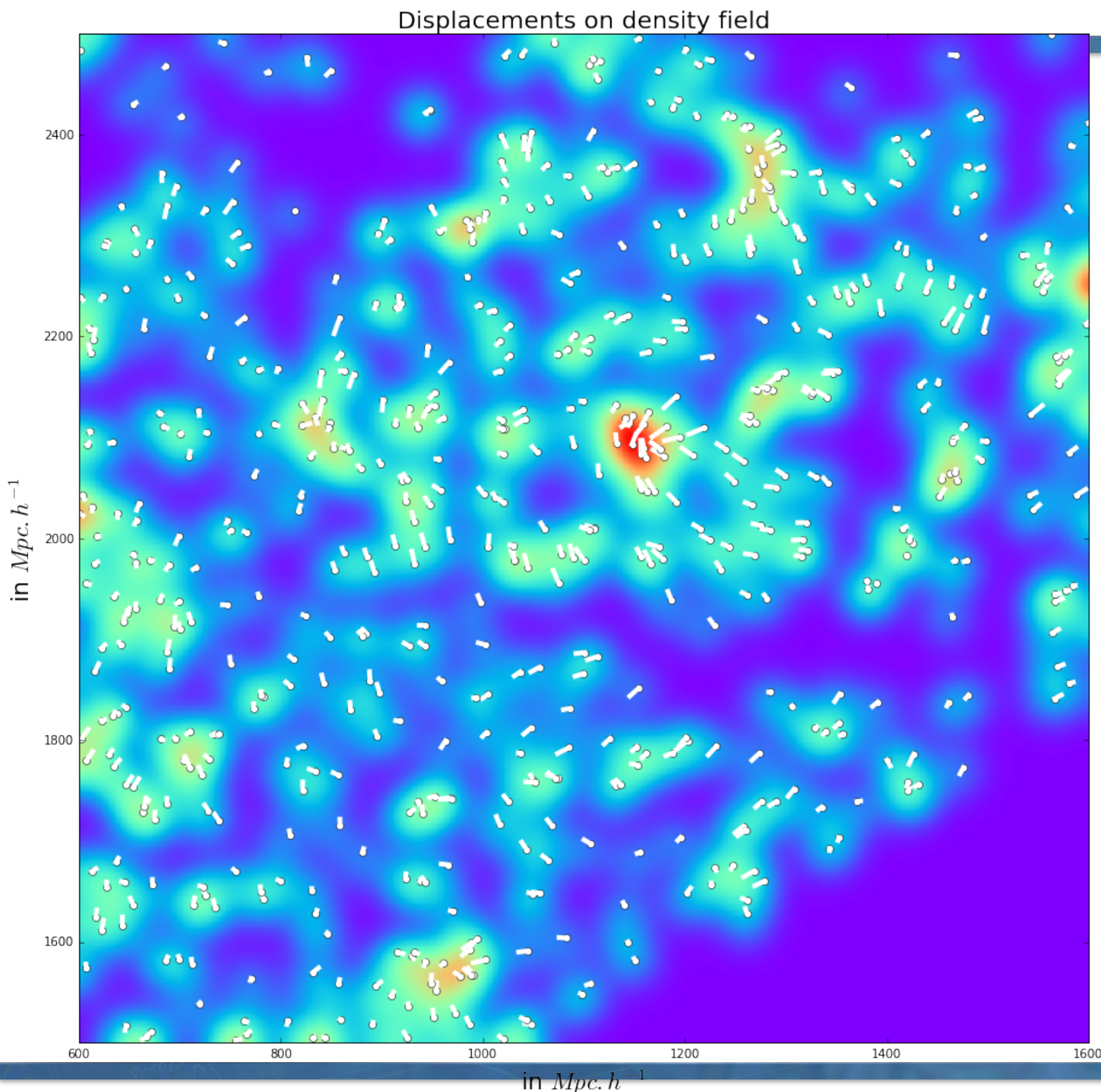
Alejandro Aviles et al 2021

Works very well up to :

$$r \sim 25 \text{ Mpc}/h$$

$$k \sim 0.3 \text{ h}/\text{Mpc}$$

Reconstruction to enhance BAO peak



Vargas, Ho, Fromenteau, Cuesta (2017)

We estimate the smoothed galaxy density field



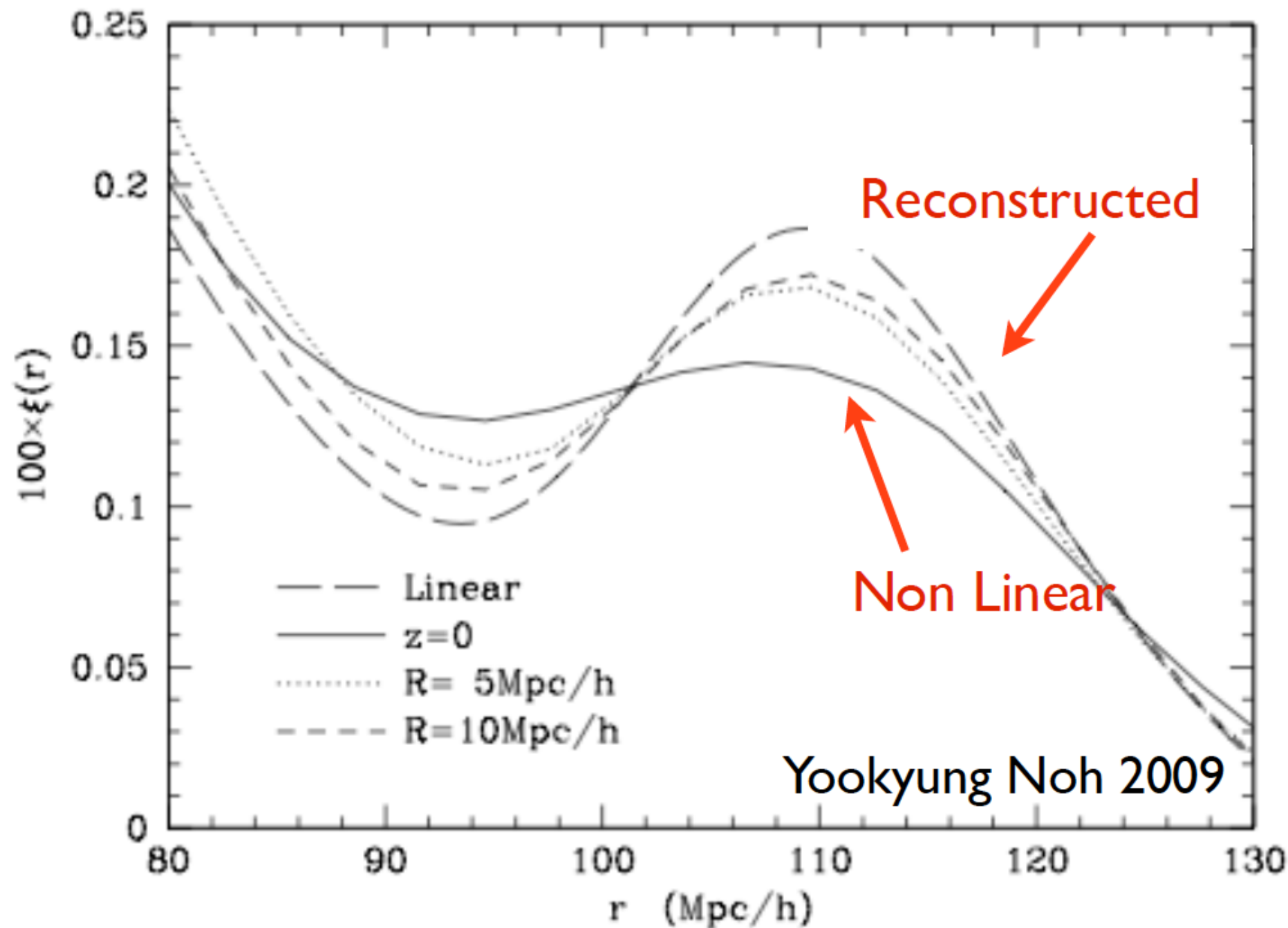
Using linear bias, we derive the matter density field



We move back the galaxies to their original position

$$\vec{\psi}(\vec{k}) = \frac{-i\vec{k}}{k^2} \frac{\delta_g(\vec{k})}{b_g}$$

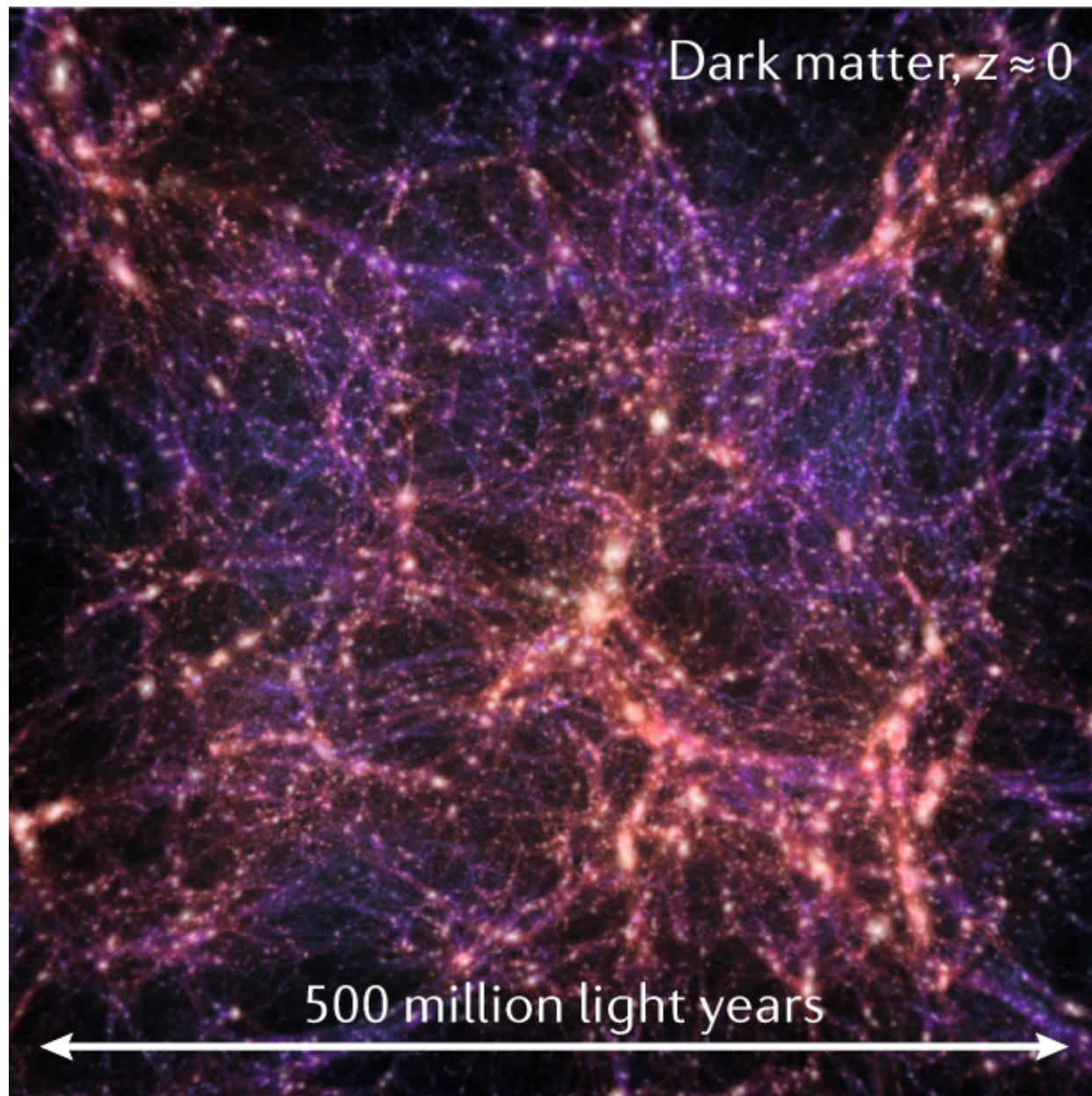
Reconstruction to enhance BAO peak



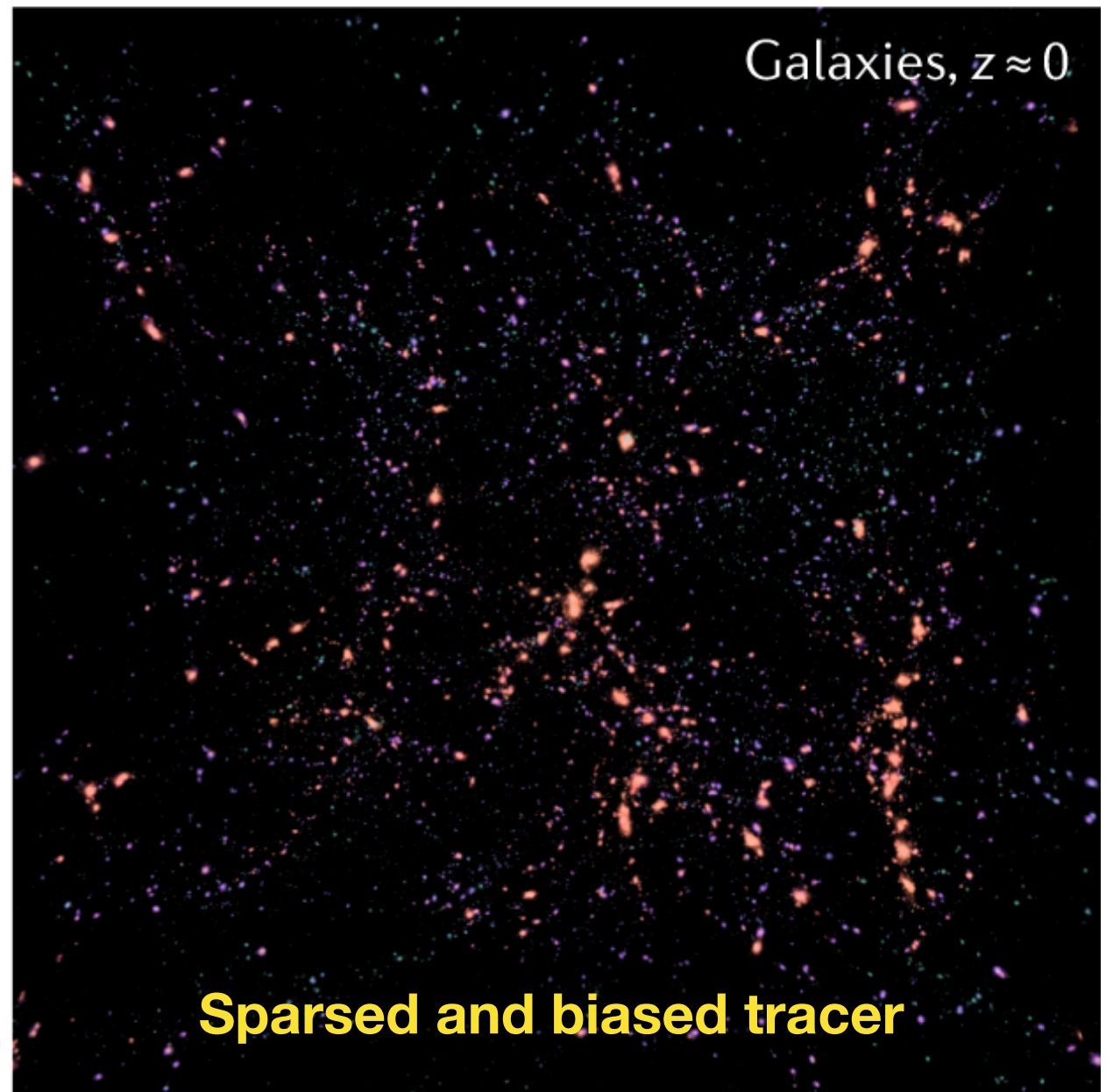
Is part of the standard
BAO analysis

What we do observe?

a That's what we need

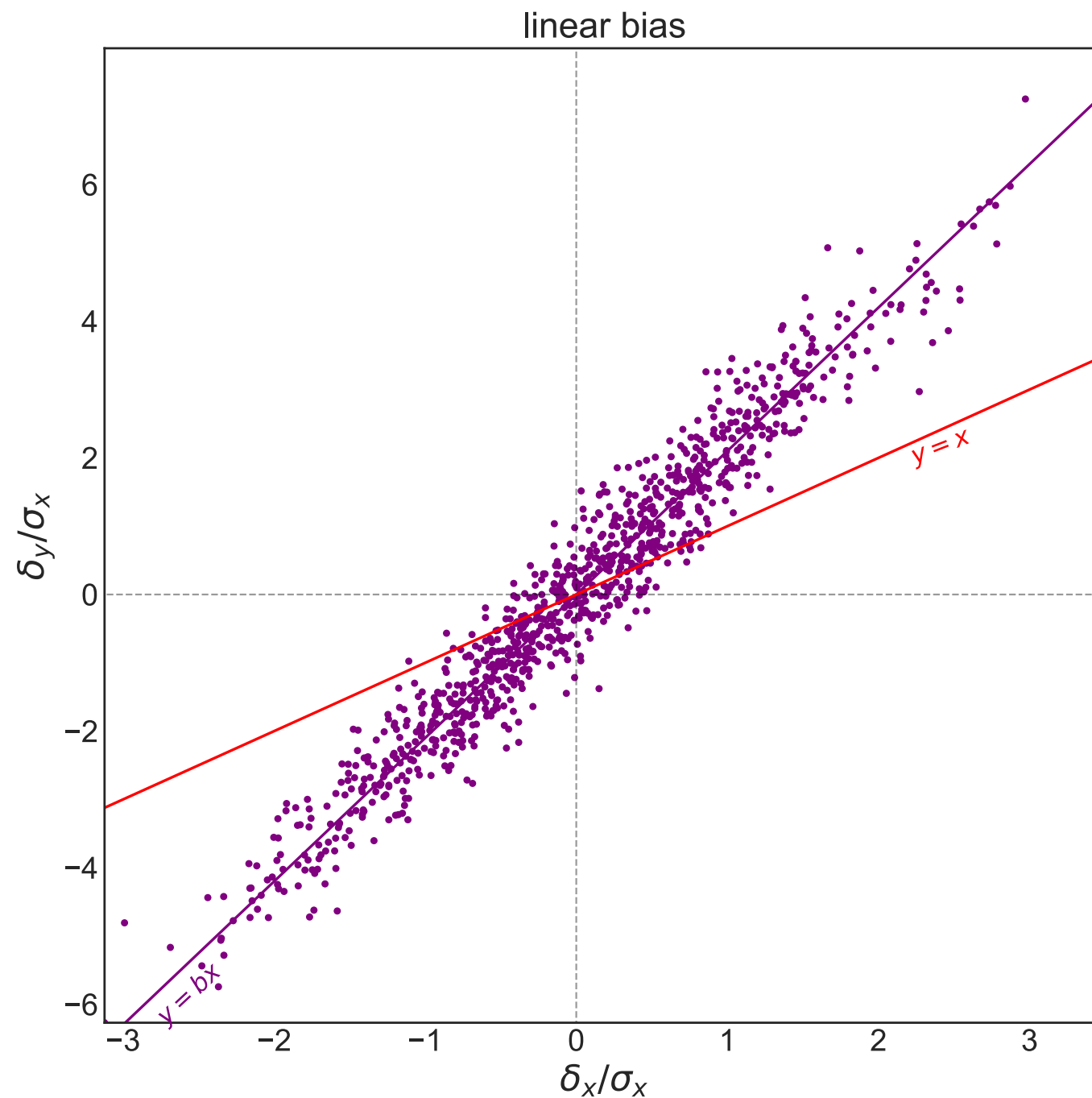


b That's what we have



(Image from Robertson et al. 2019)

Global bias (linear)



$$\delta_y = b_1(\delta_x - \bar{\delta}_x)$$

2pt-statistics and linear bias

$$\xi_{mm}(\vec{r}) = \left\langle \delta_m(\vec{x} + \vec{r}) \cdot \delta_m(\vec{x}) \right\rangle_{\vec{x}}$$

$$\xi_{gg}(\vec{r}) = \left\langle \delta_g(\vec{x} + \vec{r}) \cdot \delta_g(\vec{x}) \right\rangle_{\vec{x}}$$

$$\xi_{gg}(\vec{r}) = \left\langle b_g \delta_m(\vec{x} + \vec{r}) \cdot b_g \delta_m(\vec{x}) \right\rangle_{\vec{x}}$$

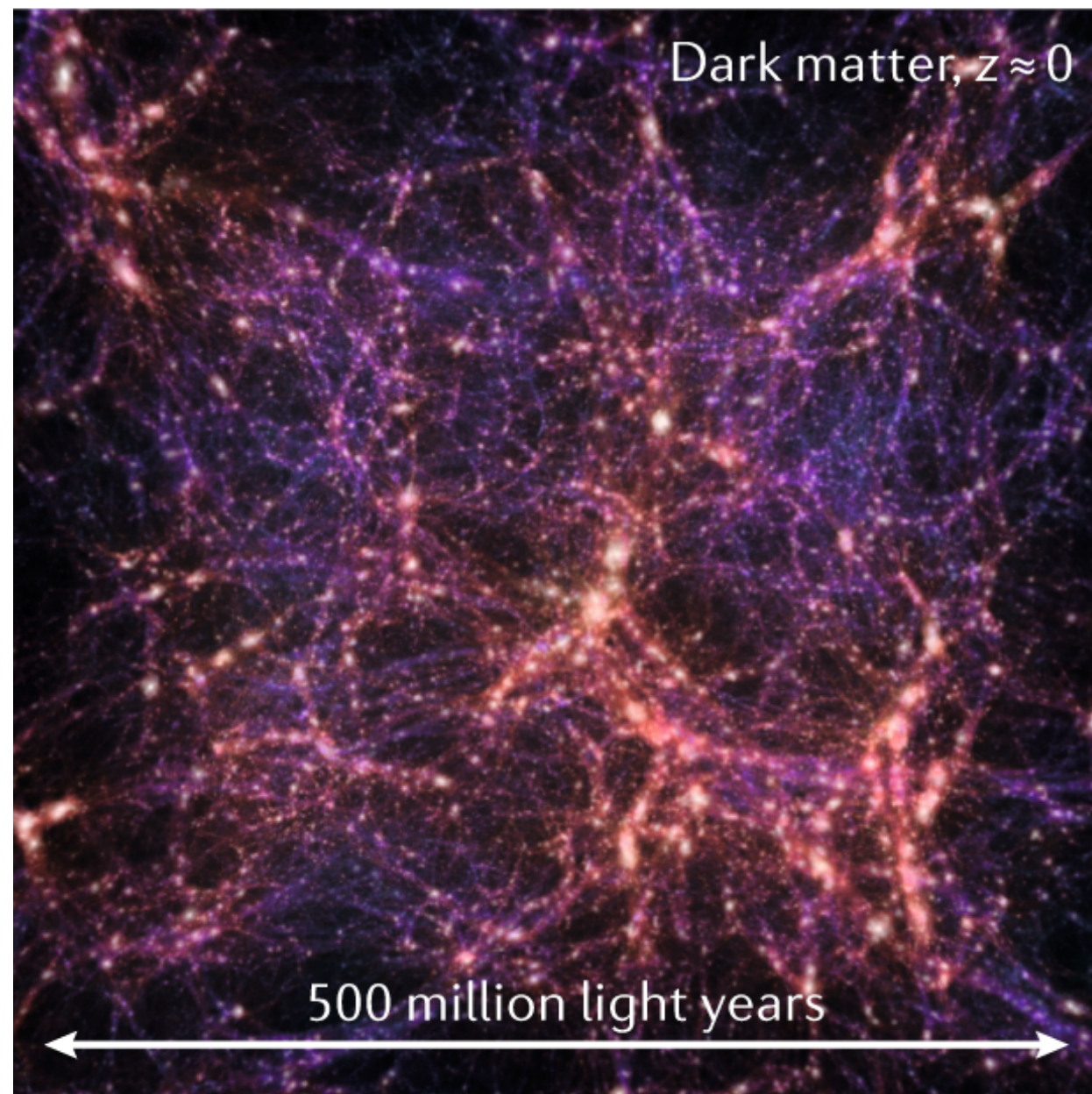
$$\delta_x = \frac{\rho_x - \bar{\rho}_x}{\bar{\rho}_x}$$

b = linear bias

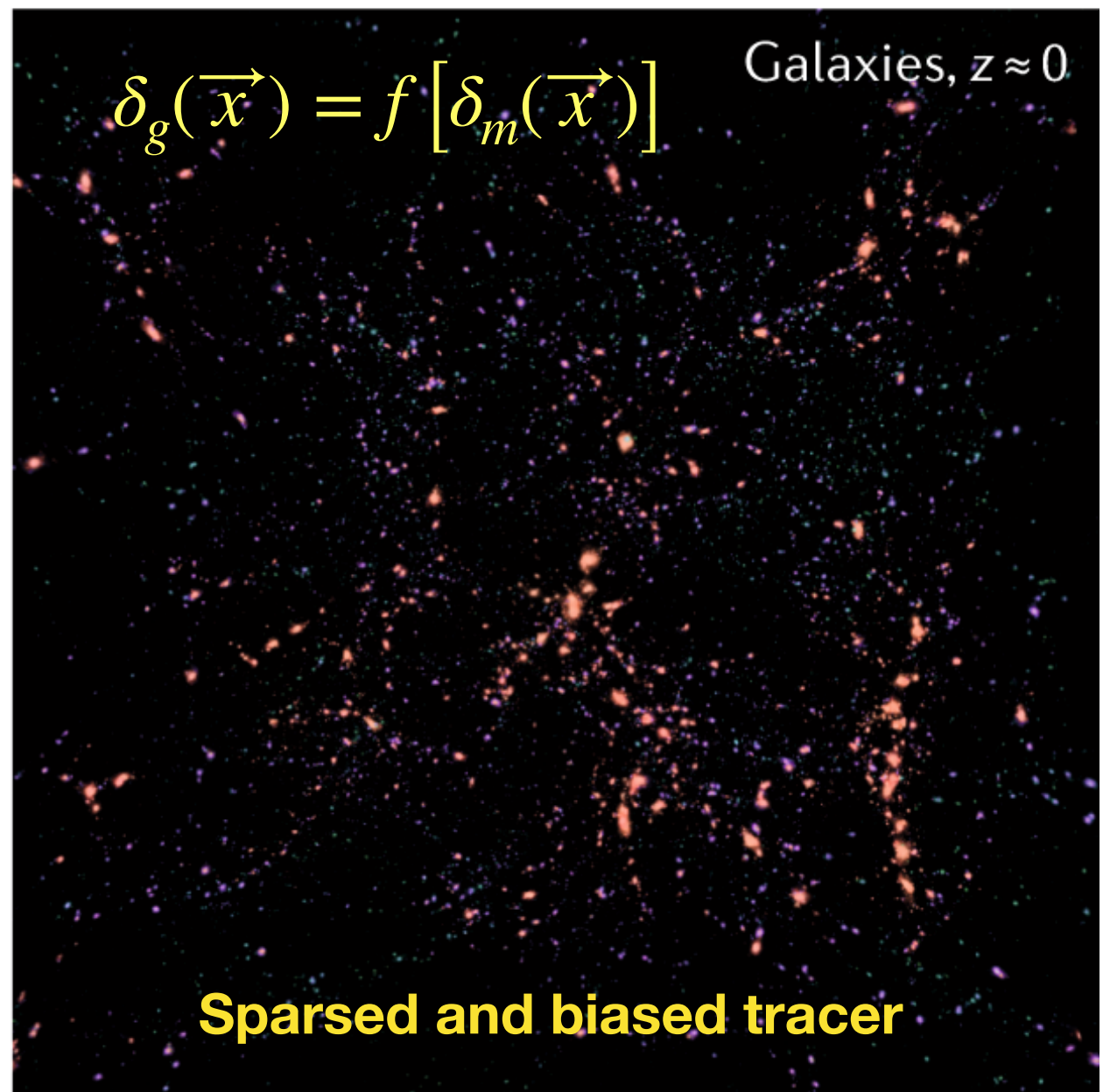
$$\xi_{gg}(\vec{r}) = b_g^2 \xi_{mm}(\vec{r})$$

Complicated because bias can be complex

a That's what we need



b That's what we have



(Image from Robertson et al. 2019)

But what to do?

$$\xi_{mm}(\vec{r}) = \left\langle \delta_m(\vec{x} + \vec{r}) \cdot \delta_m(\vec{x}) \right\rangle_{\vec{x}}$$

$$\xi_{gg}(\vec{r}) = \left\langle \delta_g(\vec{x} + \vec{r}) \cdot \delta_g(\vec{x}) \right\rangle_{\vec{x}}$$

$$\delta_x = \frac{\rho_x - \bar{\rho}_x}{\bar{\rho}_x}$$

$$\xi_{gg}(\vec{r}) = \left\langle f[\delta_m(\vec{x} + \vec{r})] \cdot f[\delta_m(\vec{x})] \right\rangle_{\vec{x}}$$

???

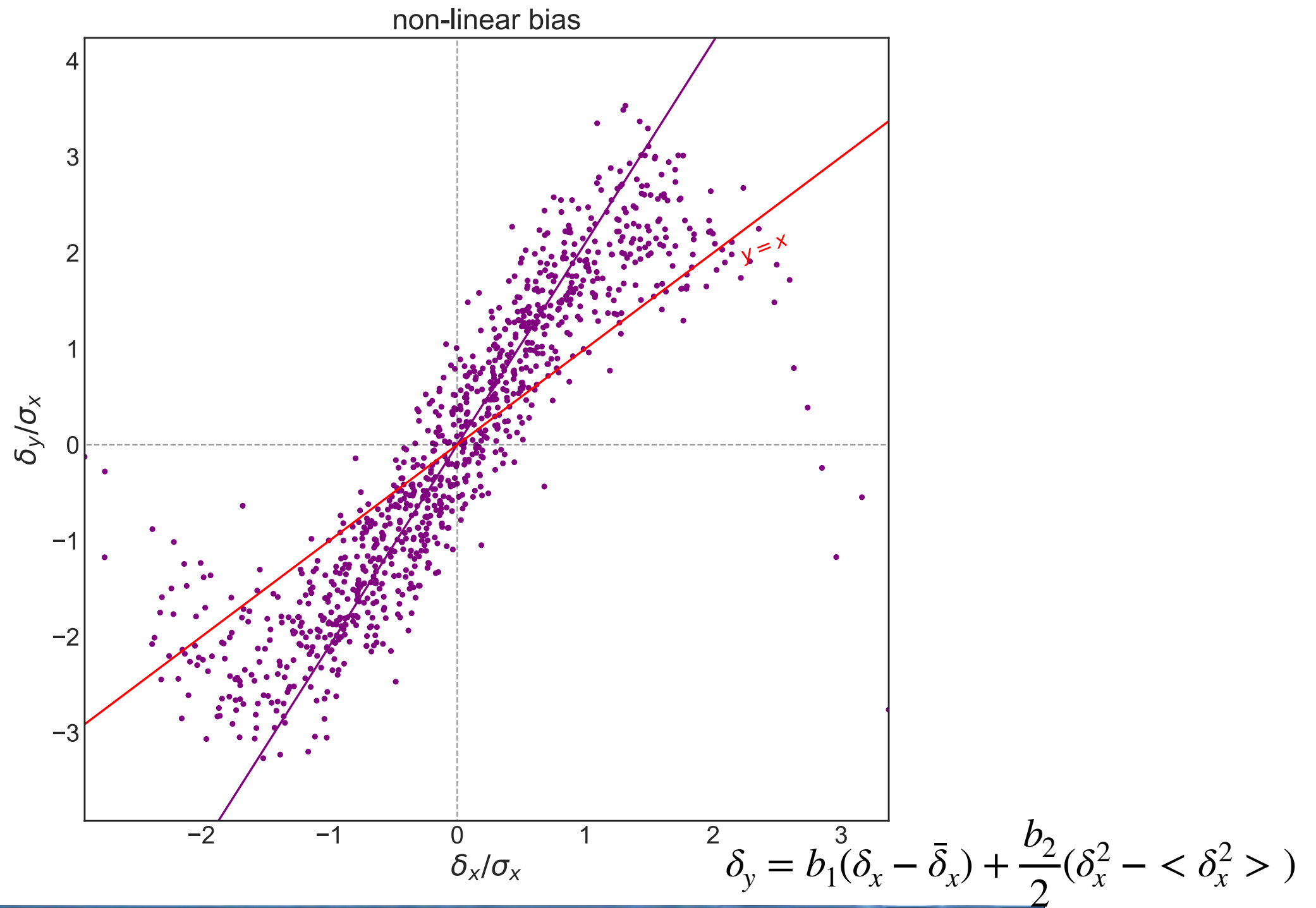
Standard bias scheme

$$\delta_g(\vec{x}) = f[\delta_m(\vec{x})] \longrightarrow \delta_g(\vec{x}) = \sum_{i=0}^{\infty} \frac{b_i}{i!} \delta_m^i(\vec{x})$$

Local
No specific form

non-Local
Taylor expansion

Non-local bias order 2



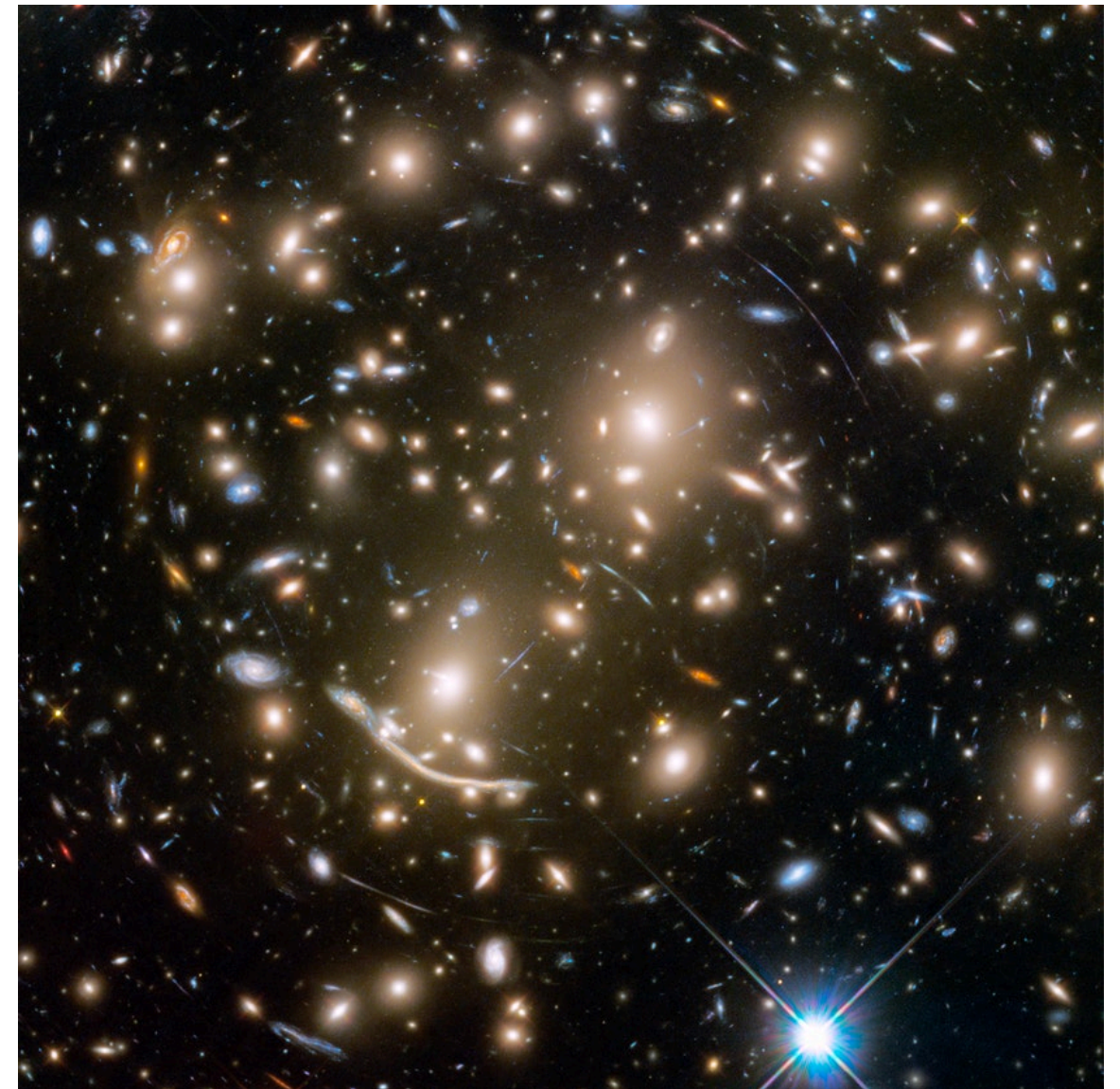
Real application : Hard!!

$$\begin{aligned}\delta_g = & c_\delta \delta + \frac{1}{2} c_{\delta^2} (\delta^2 - \sigma^2) + \frac{1}{2} c_{s^2} \left(s^2 - \frac{2}{3} \sigma^2 \right) + \frac{1}{3!} c_{\delta^3} \delta^3 + \frac{1}{2} c_{\delta s^2} \delta s^2 + c_\psi \psi + c_{st} st + \frac{1}{3!} c_{s^3} s^3 \\ & + c_\epsilon \epsilon + c_{\delta\epsilon} \delta\epsilon + \frac{1}{2} c_{\delta^2\epsilon} \delta^2\epsilon + \frac{1}{2} c_{s^2\epsilon} s^2\epsilon + \frac{1}{2} c_{\epsilon^2} (\epsilon^2 - \sigma_\epsilon^2) + \frac{1}{2} c_{\delta\epsilon^2} \delta\epsilon^2 + \frac{1}{3!} c_{\epsilon^3} \epsilon^3 + \dots\end{aligned}$$

$$\begin{aligned}P_{mg}(k) = & c_\delta P_{NL}(k) \\ & + c_{\delta^2} \int \frac{d^3\mathbf{q}}{(2\pi)^3} P(q) P(|\mathbf{k} - \mathbf{q}|) F_S^{(2)}(\mathbf{q}, \mathbf{k} - \mathbf{q}) + \frac{34}{21} c_{\delta^2} \sigma^2 P(k) \\ & + c_{s^2} \int \frac{d^3\mathbf{q}}{(2\pi)^3} P(q) P(|\mathbf{k} - \mathbf{q}|) F_S^{(2)}(\mathbf{q}, \mathbf{k} - \mathbf{q}) S(\mathbf{q}, \mathbf{k} - \mathbf{q}) \\ & + 2 c_{s^2} P(k) \int \frac{d^3\mathbf{q}}{(2\pi)^3} P(q) F_S^{(2)}(-\mathbf{q}, \mathbf{k}) S(\mathbf{q}, \mathbf{k} - \mathbf{q}) \\ & + \frac{1}{2} c_{\delta^3} \sigma^2 P(k) + \frac{1}{3} c_{\delta s^2} \sigma^2 P(k) \\ & + 2 c_\psi P(k) \int \frac{d^3\mathbf{q}}{(2\pi)^3} P(q) \left[\frac{3}{2} D_S^{(3)}(\mathbf{q}, -\mathbf{q}, -\mathbf{k}) - 2 F_S^{(2)}(-\mathbf{q}, \mathbf{k}) D_S^{(2)}(\mathbf{q}, \mathbf{k} - \mathbf{q}) \right] \\ & + 2 c_{st} P(k) \int \frac{d^3\mathbf{q}}{(2\pi)^3} P(q) D_S^{(2)}(-\mathbf{q}, \mathbf{k}) S(\mathbf{q}, \mathbf{k} - \mathbf{q}) \\ & + \frac{1}{2} c_{\delta\epsilon^2} \sigma_\epsilon^2 P(k) .\end{aligned}$$

McDonald & Roy 2009

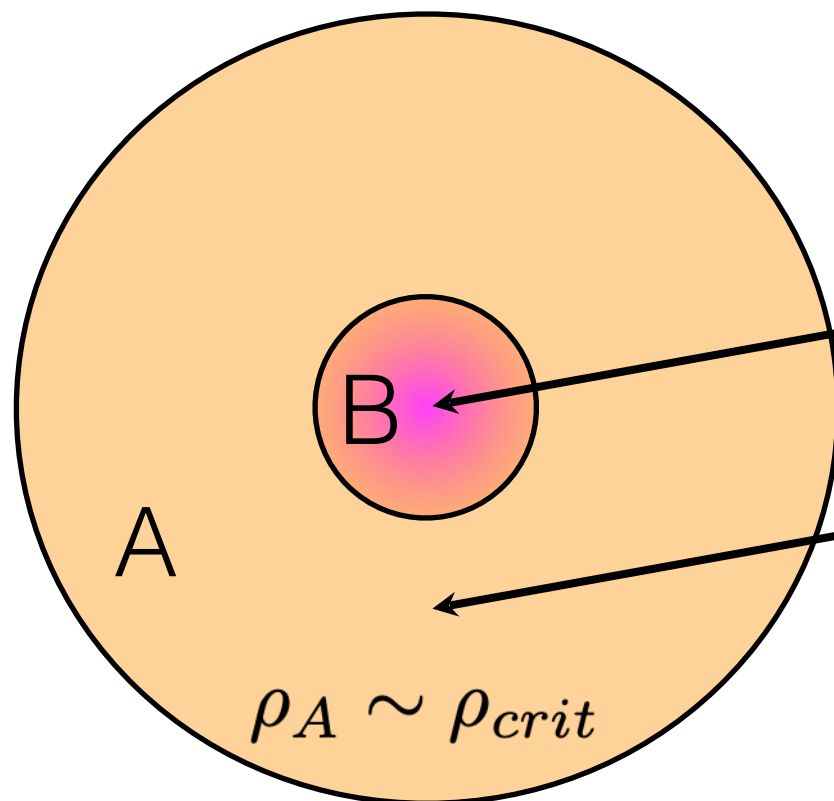
All galaxies are leaving in halos



Spherical halo model

The 2 areas A & B are considered as local universe with proper scale factor evolution

(Gunn & Gott 1972)

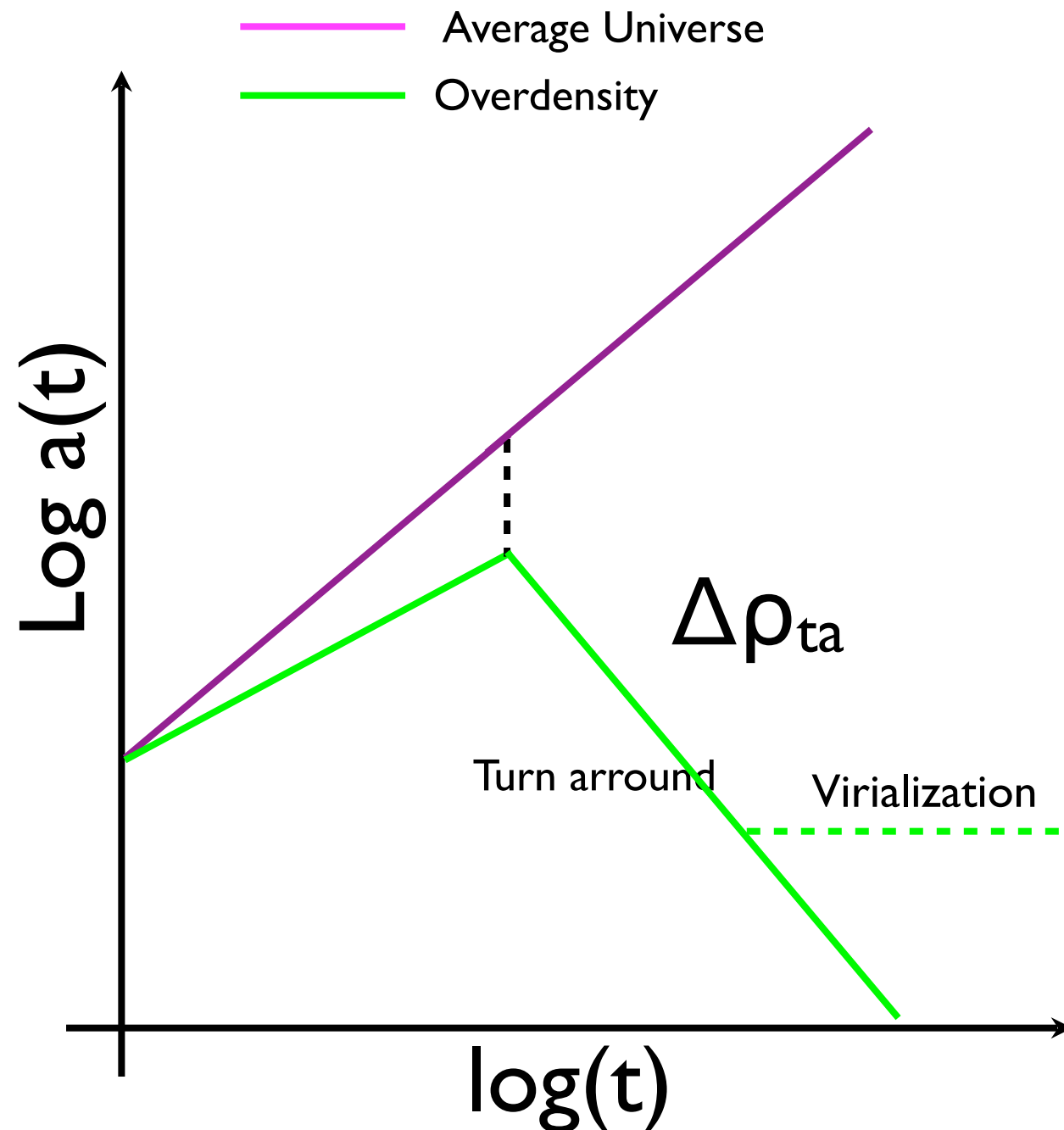


$$H_B^2 = \left(\frac{\dot{a}}{a} \right)^2 = \frac{8\pi G}{3} \rho_B + \frac{\Lambda}{3} - \frac{k}{a^2}$$

$$H_A^2 = \left(\frac{\dot{a}}{a} \right)^2 = \frac{8\pi G}{3} \rho_A + \frac{\Lambda}{3}$$

Curvature

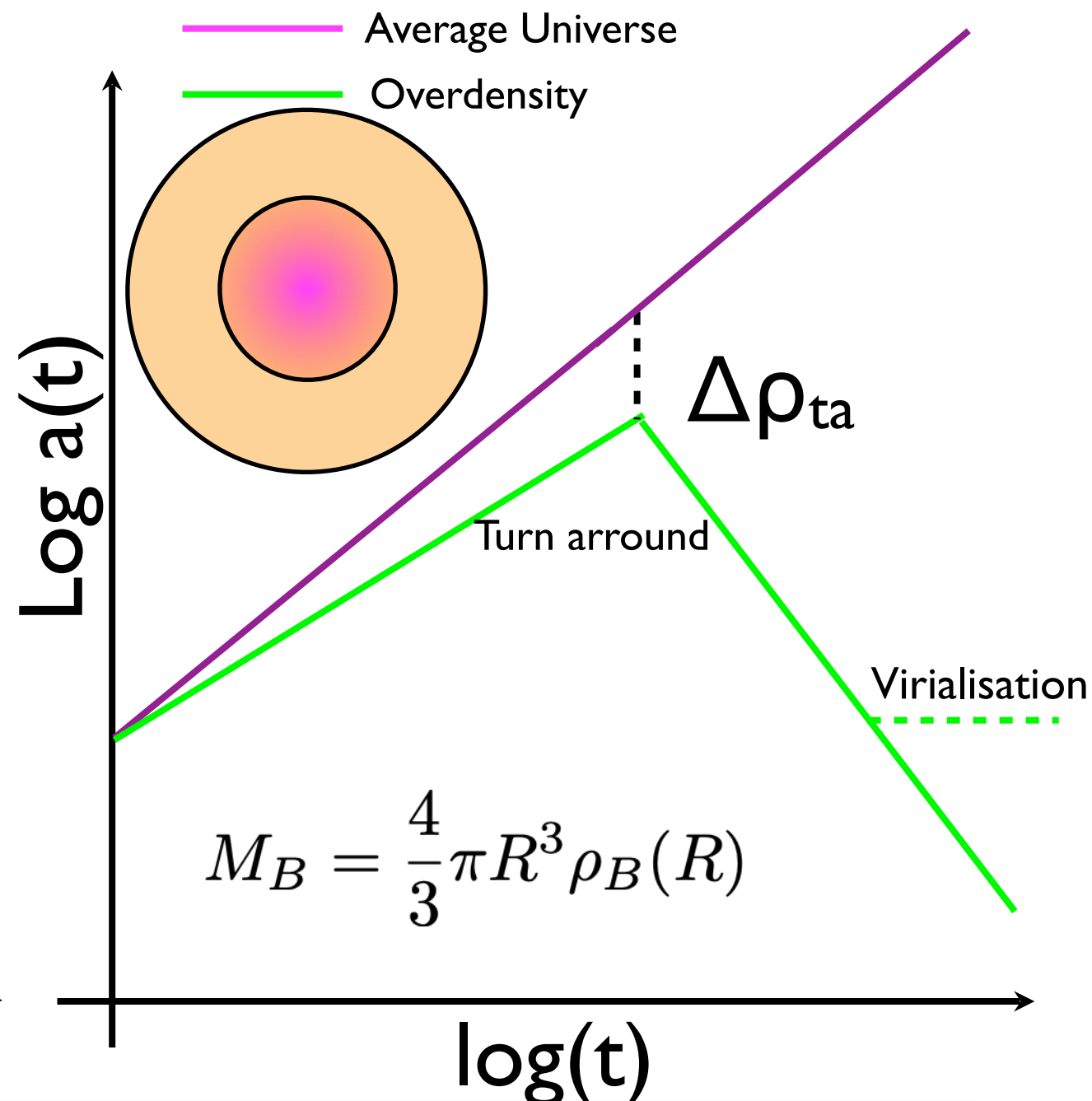
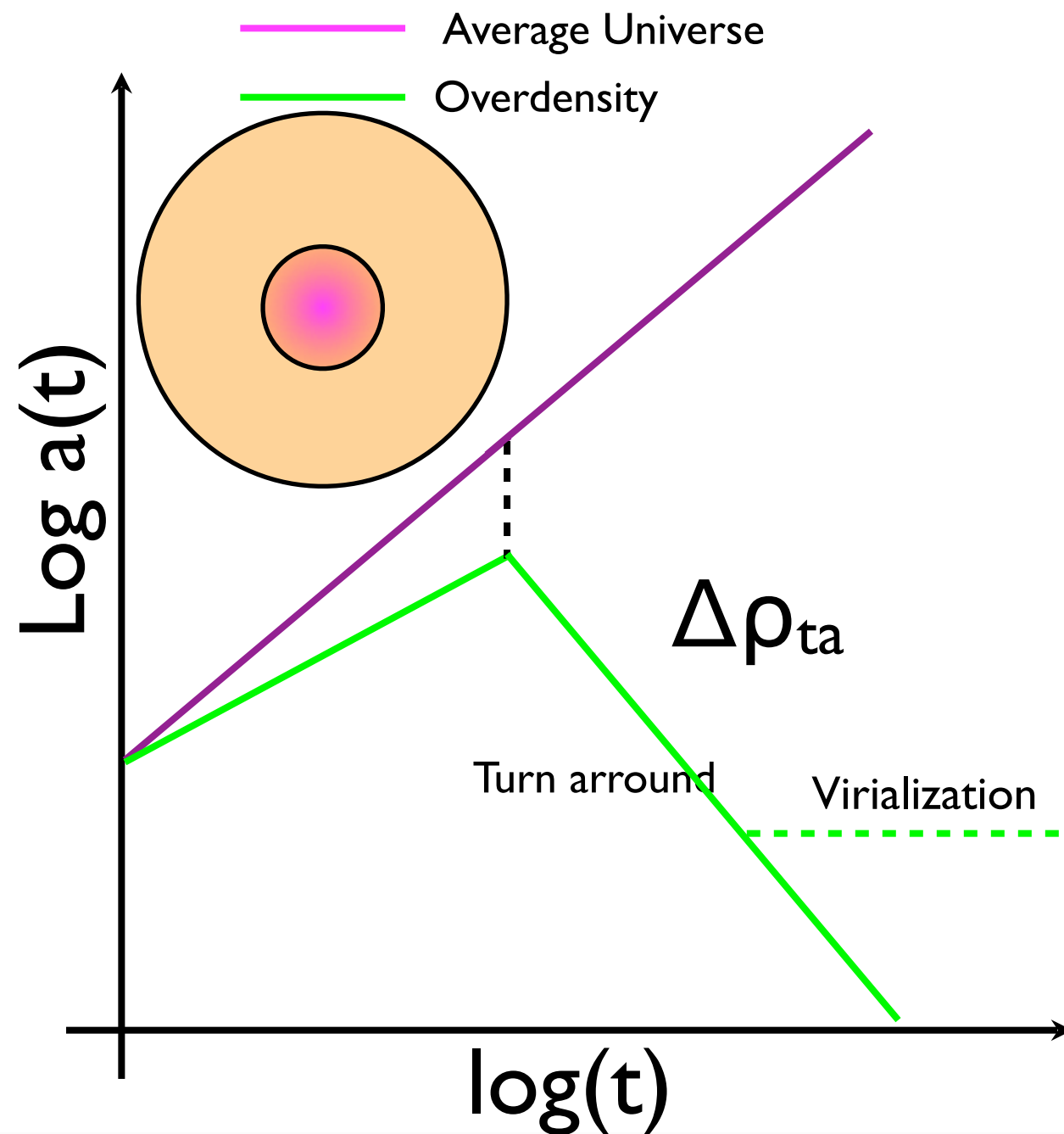
Galaxy clusters



$$M_B = \frac{4}{3}\pi R^3 \rho_B(R)$$

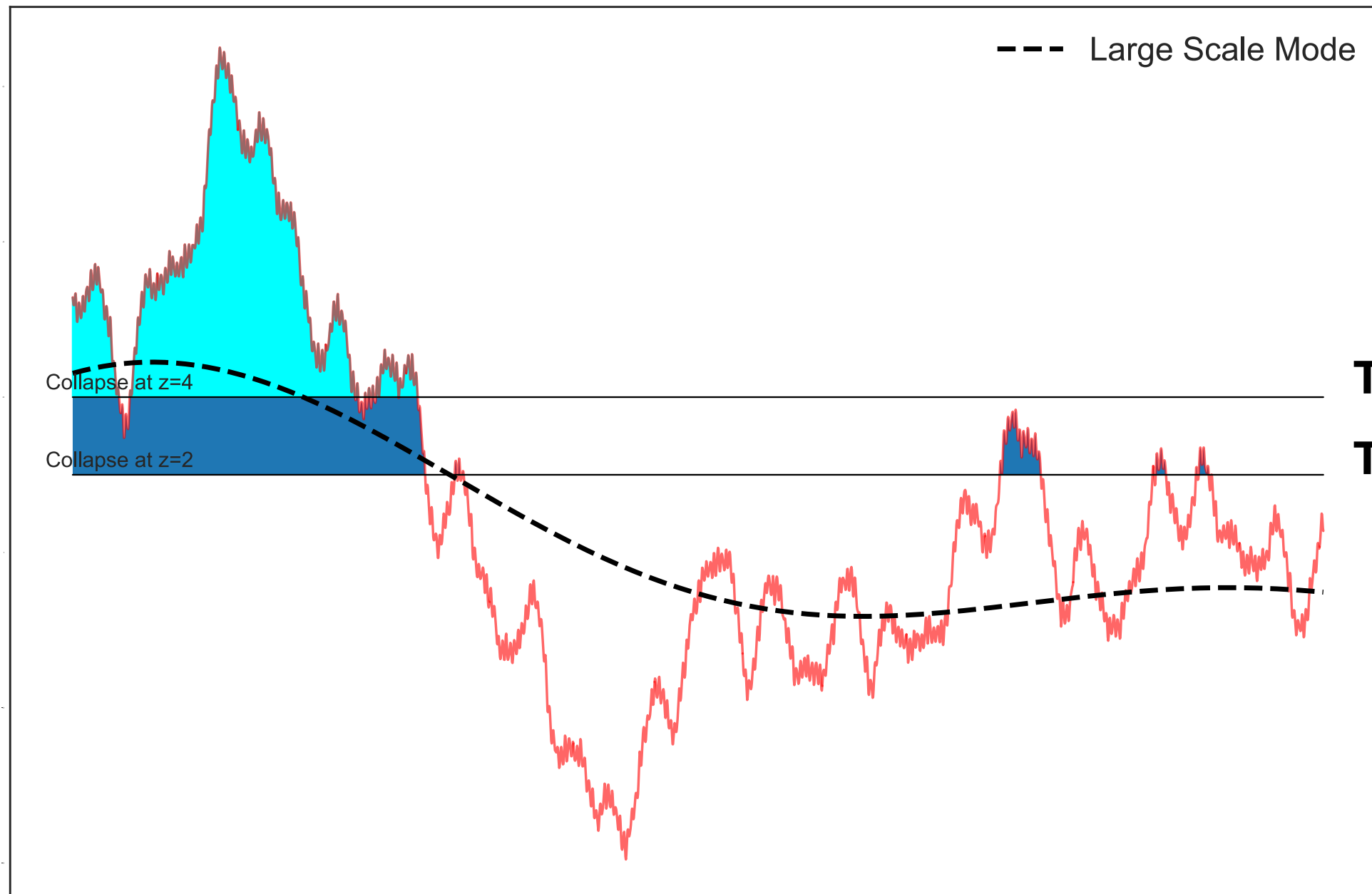
The number of virialized halos of a given mass at a given redshift

Galaxy clusters



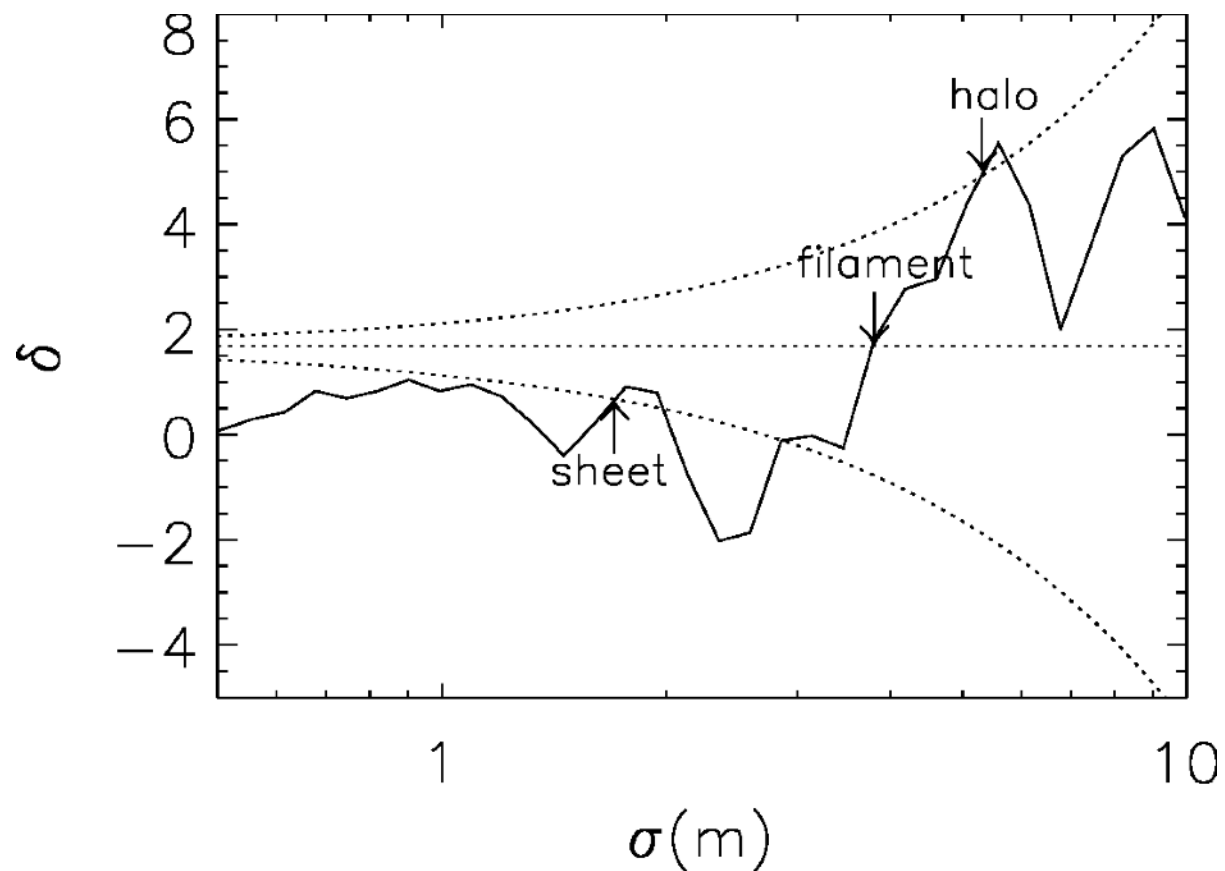
The number of virialized halos of a given mass at a given redshift

More massive halos formed in large scale overdensities



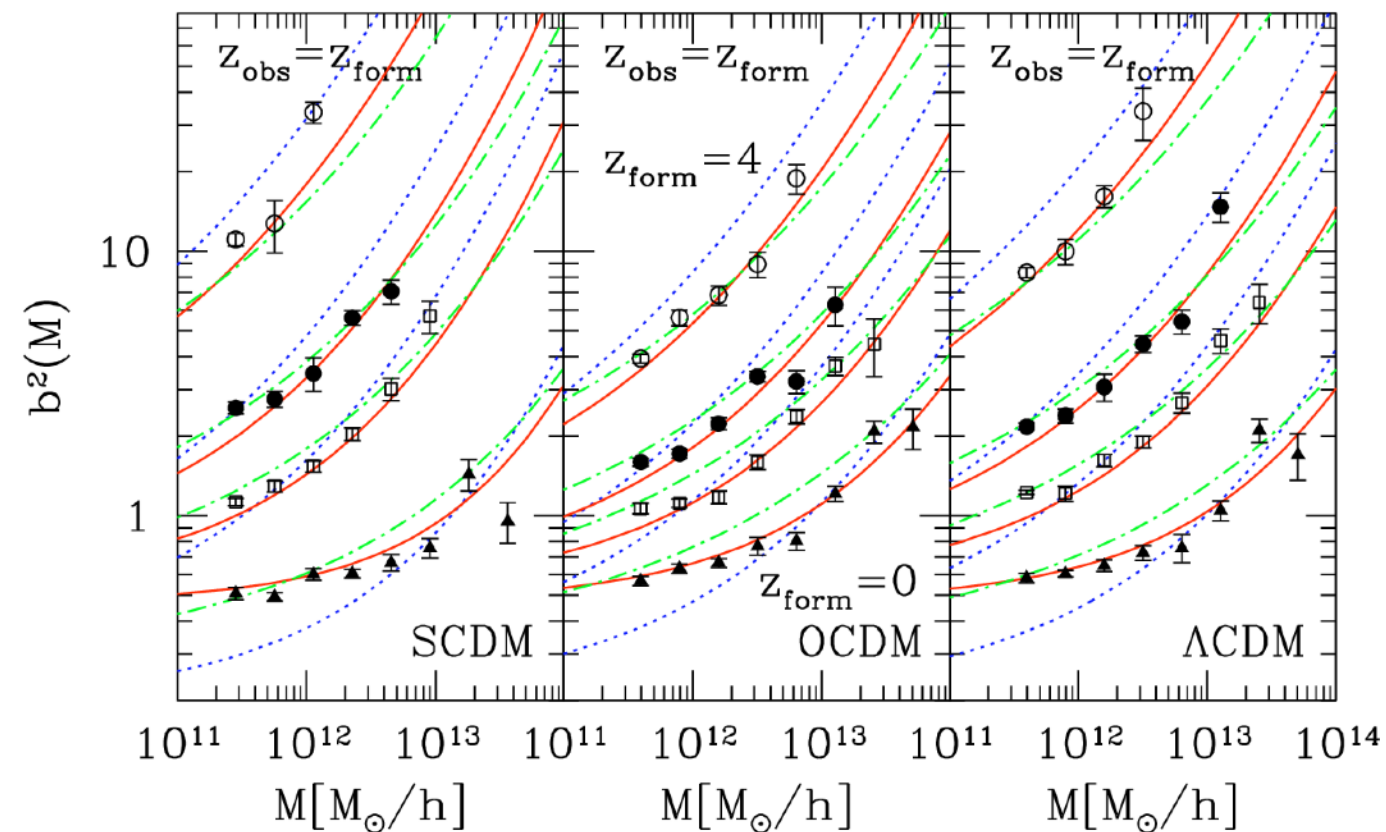
Conexión halo - materia

Excursion Set



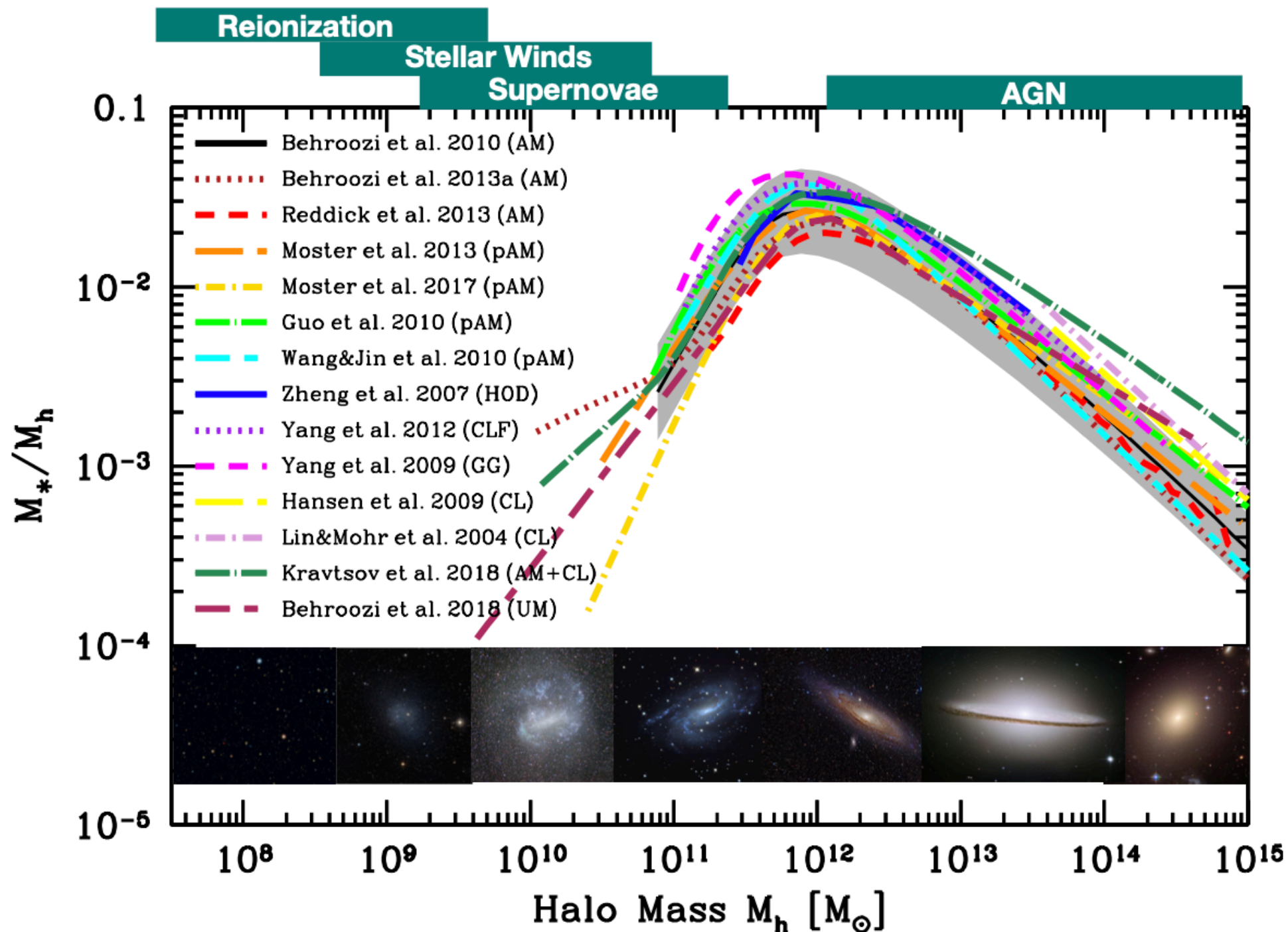
Bond et. 1991

Peak Background Split



Sheth & Tormen 1999

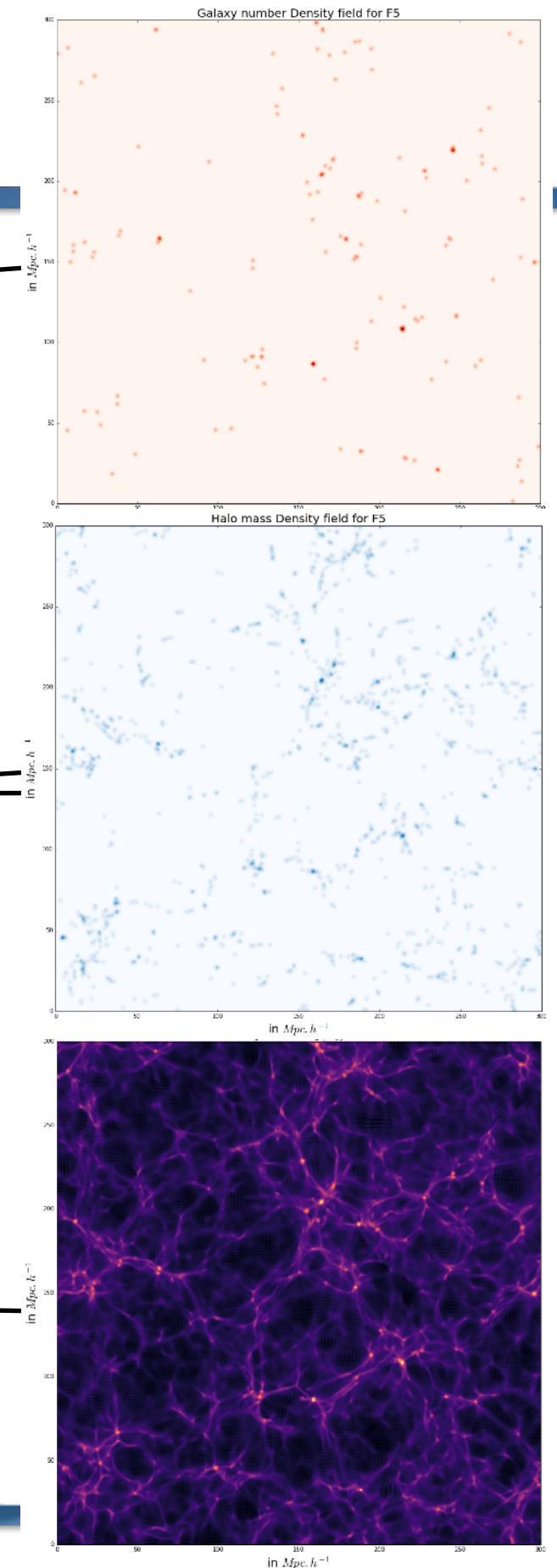
Stellar formation and galaxy evolution is much more complex

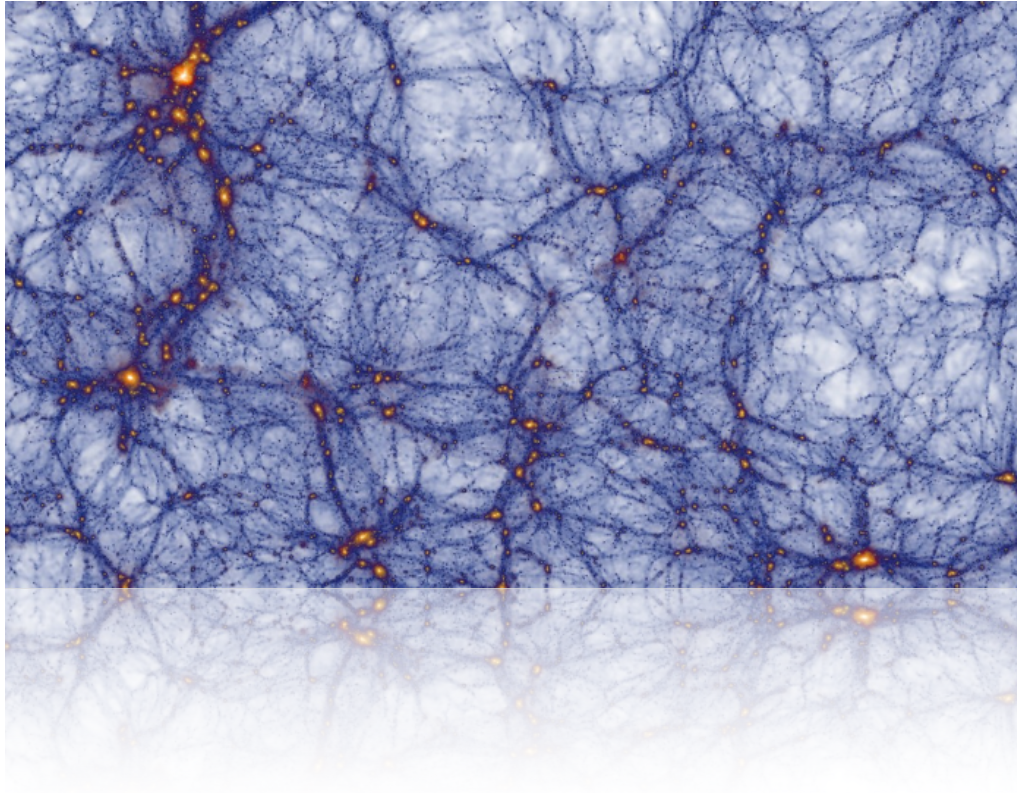


Idea

**Use local astrophysical information
to infer the hosting halo mass**

**Use the halo mass bias to connect
with the matter density field**

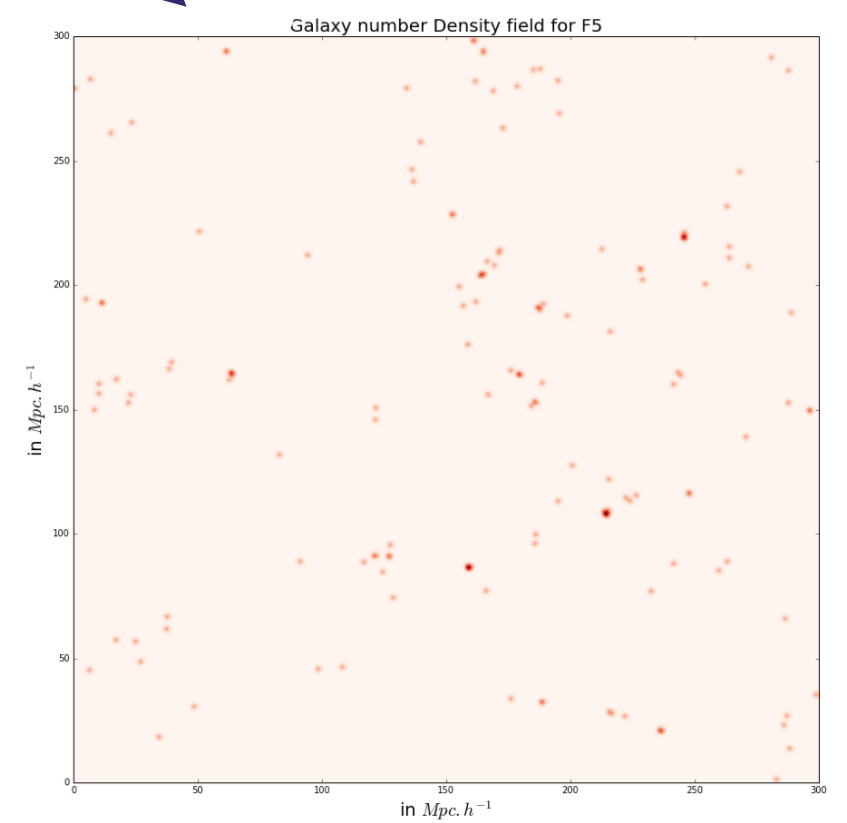
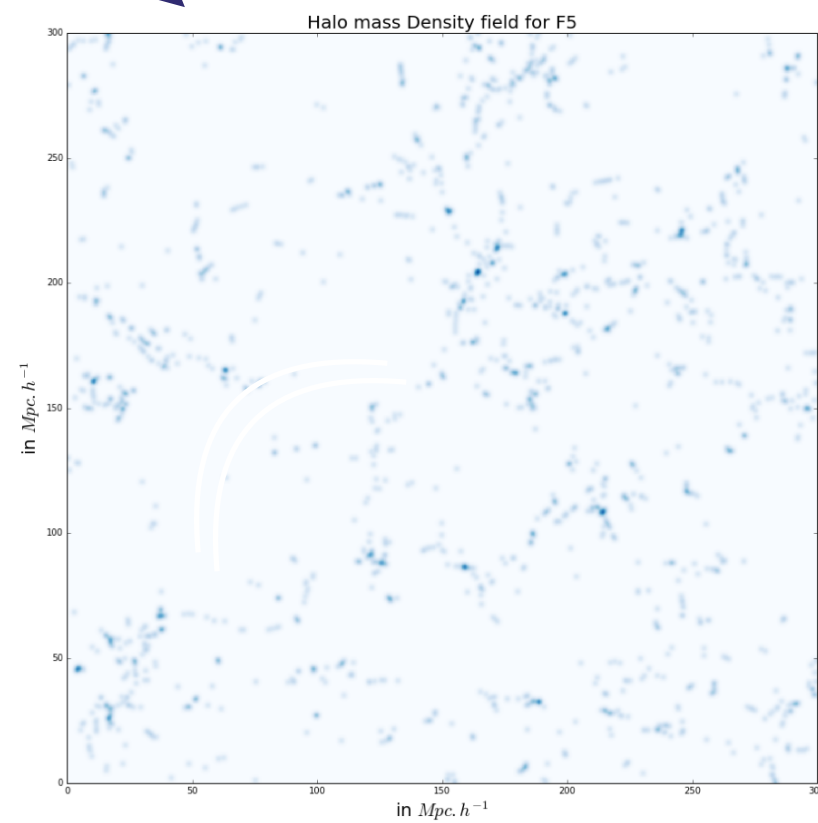
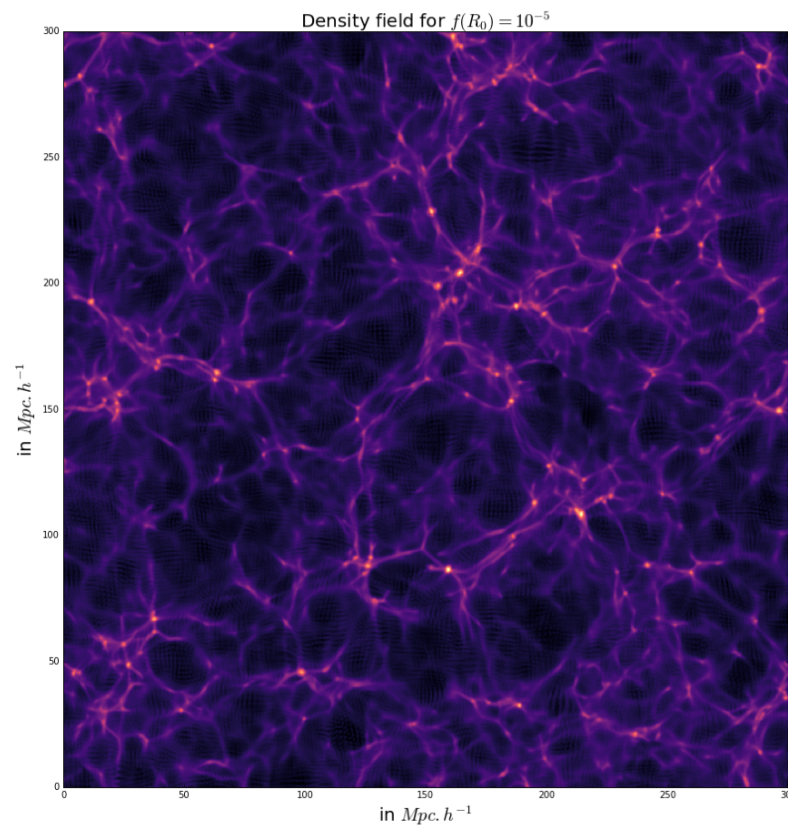




Invert HOD statistics

ROCKSTAR

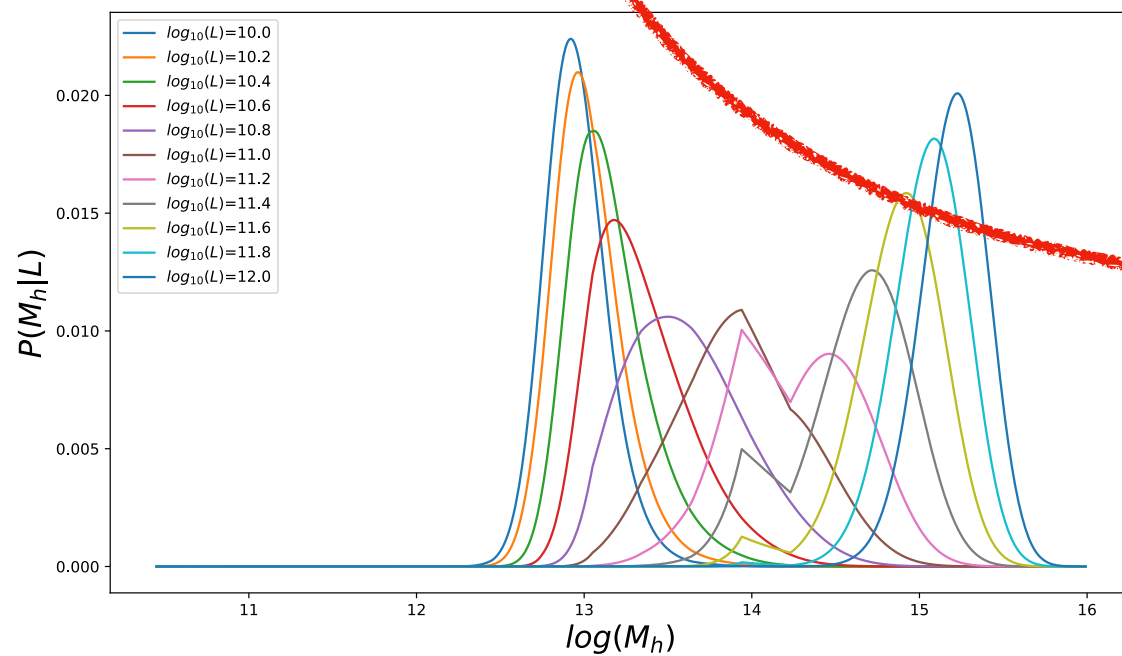
HOD + CLF



Dark Matter
particles

Dark Matter
Halos

Galaxies with
Luminosity



Infer the underlying CDM field

DATA

Reverting the information

$$\mathcal{P}(M_h | L) = \frac{\mathcal{P}(L | M_h) \cdot \mathcal{P}(M_h)}{\mathcal{P}(L)}$$

POSTERIOR WE WANT TO EVALUATE

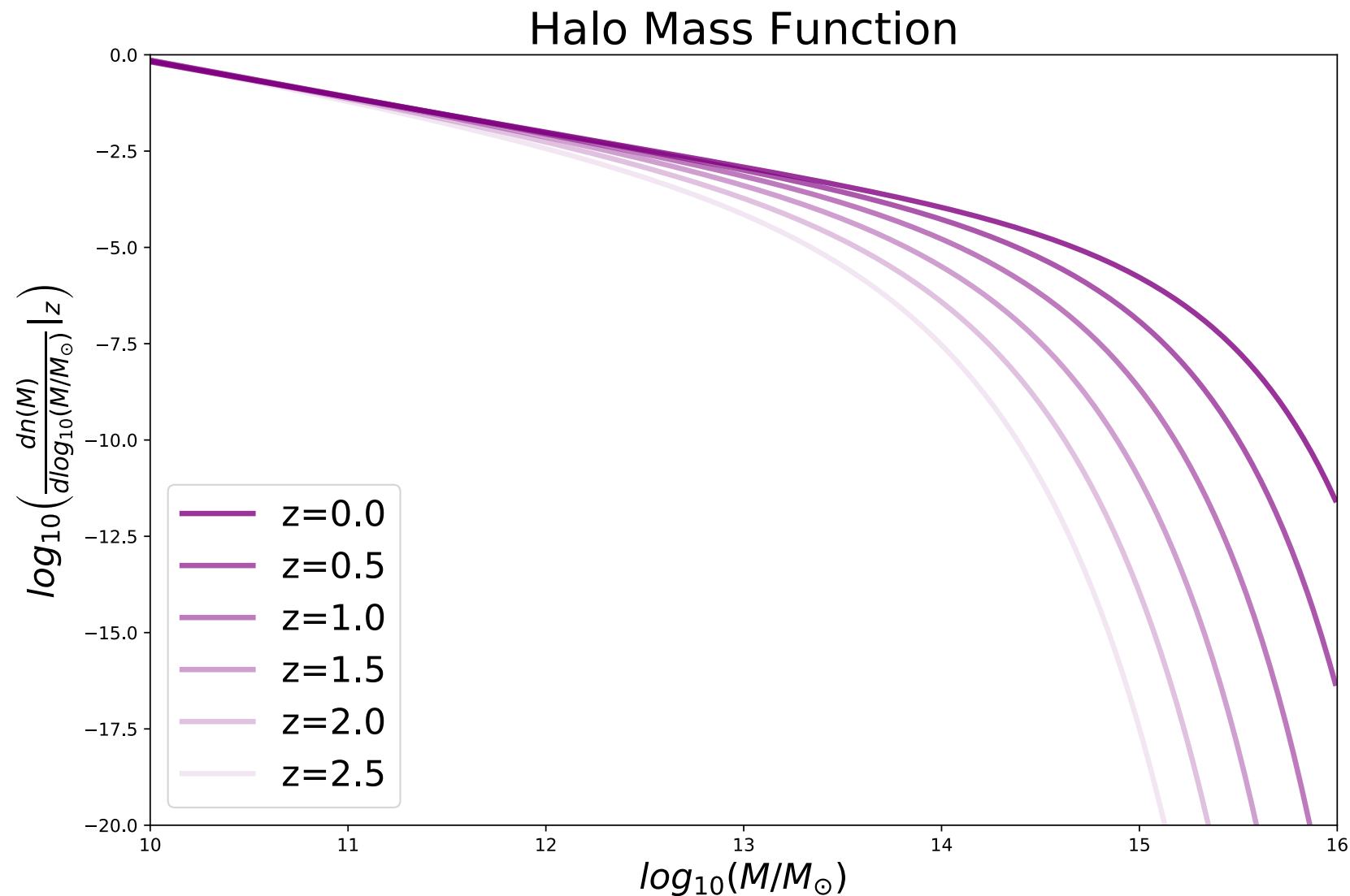
CONDITIONAL LUMINOSITY FUNCTION

CONVOLUTION OF HOD AND HMF

NORMALIZATION WHICH HAVE NOT IMPACT HERE

Halo Mass Function

1pt-statistics based on Press-Schechter and spherical collapse (Gunn & Gott) formalisms
+ Semi-analytical corrections



Halo Occupation Distribution

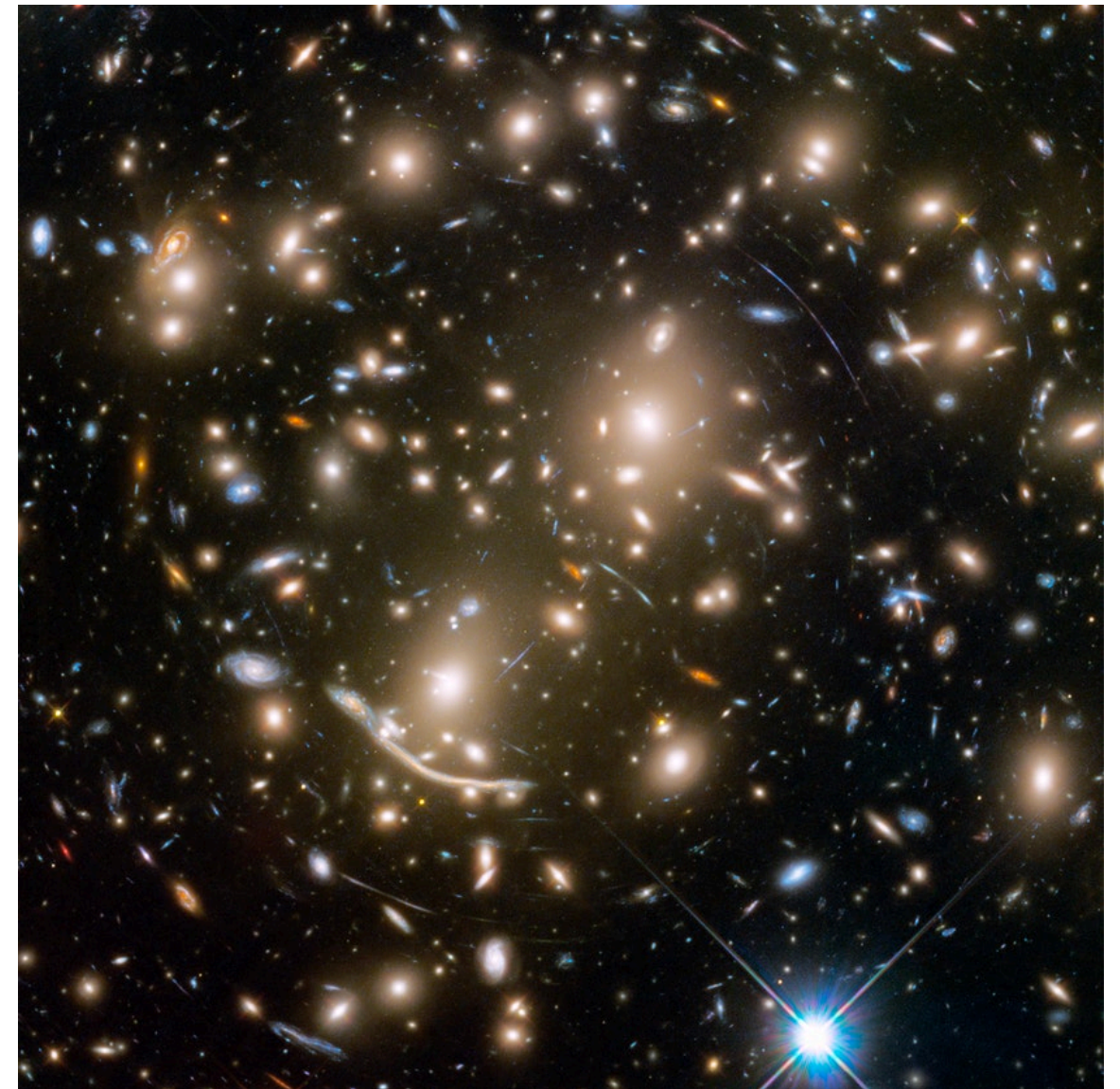
Link between the clustering of observed galaxies and the theoretical correlation function and the Halo profile

Associate a mean number of galaxies at a given halo mass

Each sample of galaxies has its proper HOD

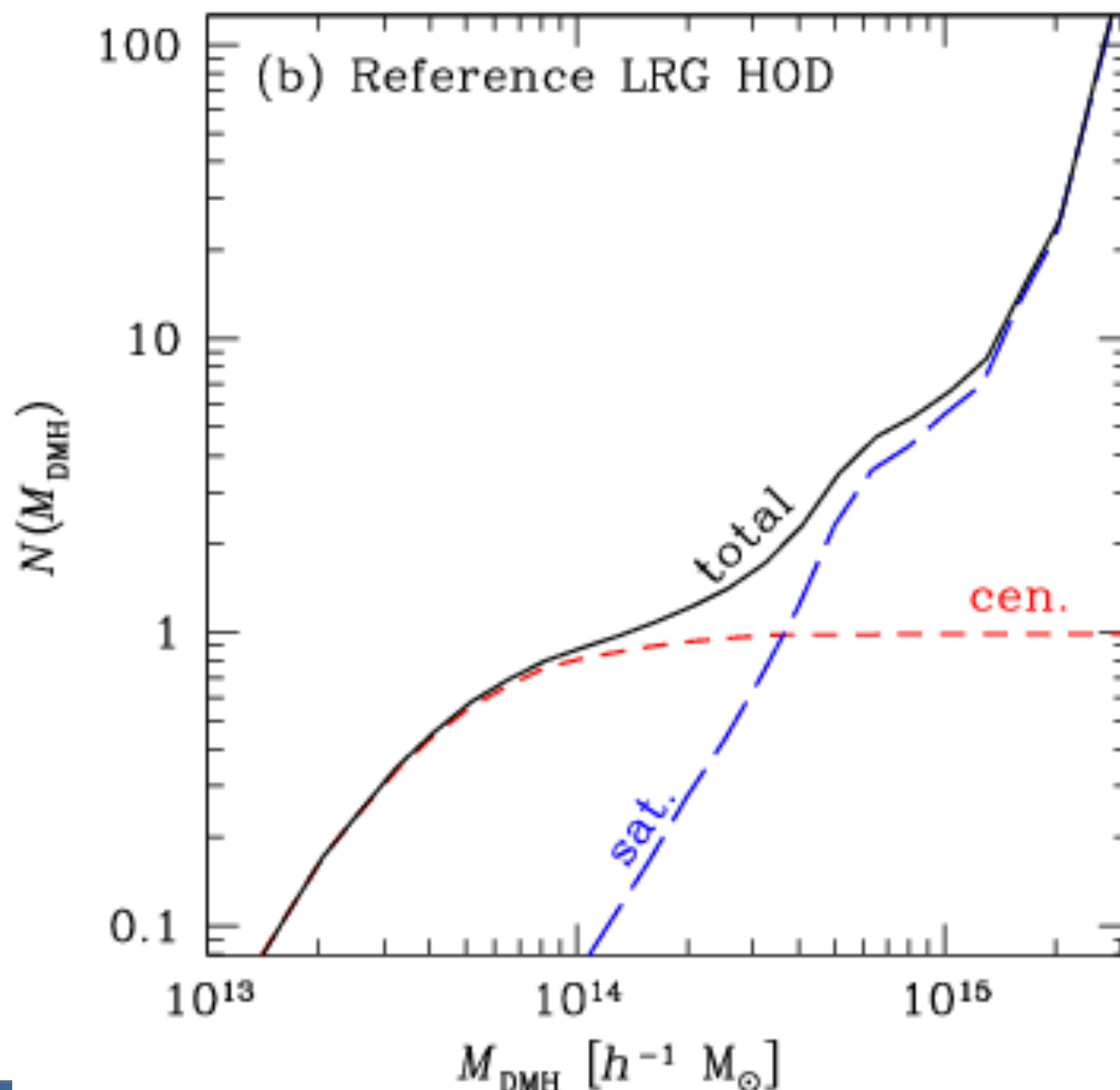
$$\langle N(M) \rangle = \langle N_{cen}(M) \rangle + \langle N_{sat}(M) \rangle$$

Halo Occupation Distribution



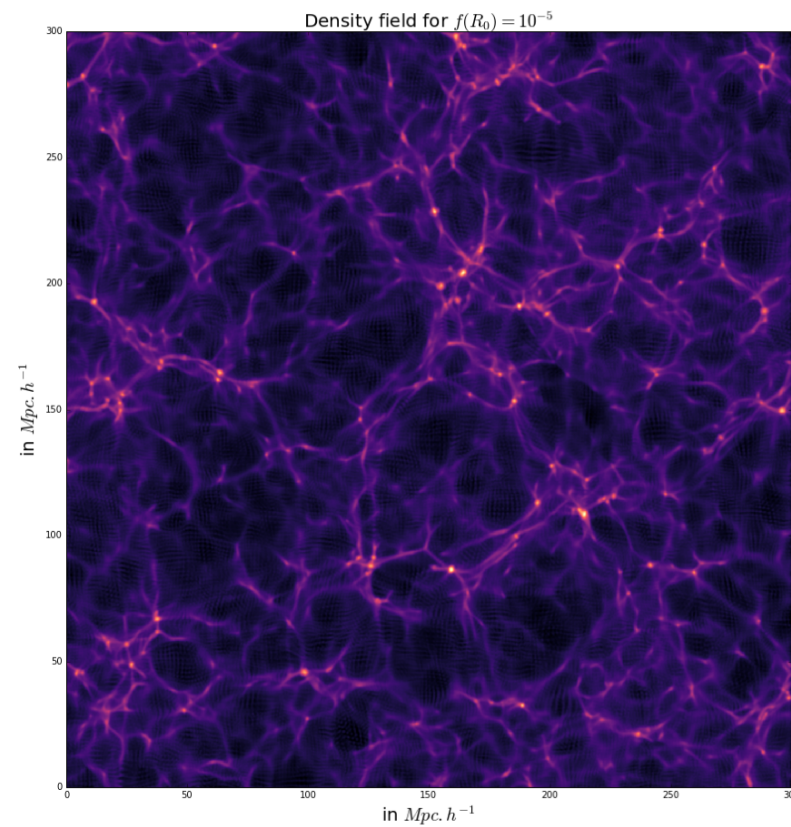
Example of HOD parametrization

$$\langle N_{cen}(M) \rangle = \frac{1}{2} \left[1 + \operatorname{erf} \left(\frac{\log M - \log M_{min}}{\sigma_{\log M}} \right) \right] \quad \langle N_{sat}(M) \rangle = \left(\frac{M - M_0}{M_1} \right)^\alpha$$

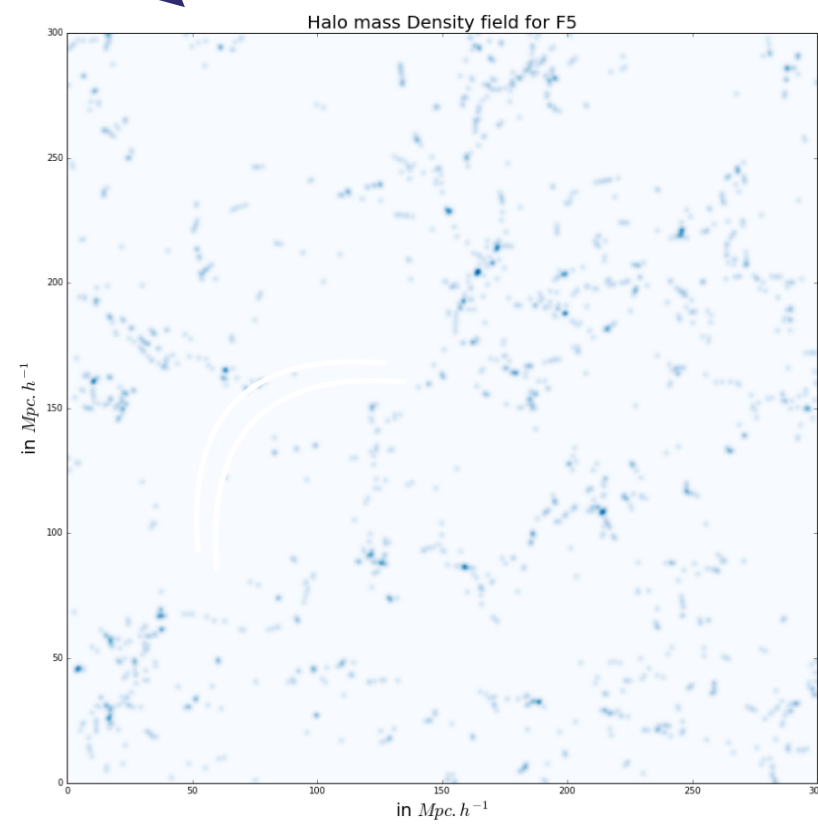


ROCKSTAR

HOD + CLF



Dark Matter
particles



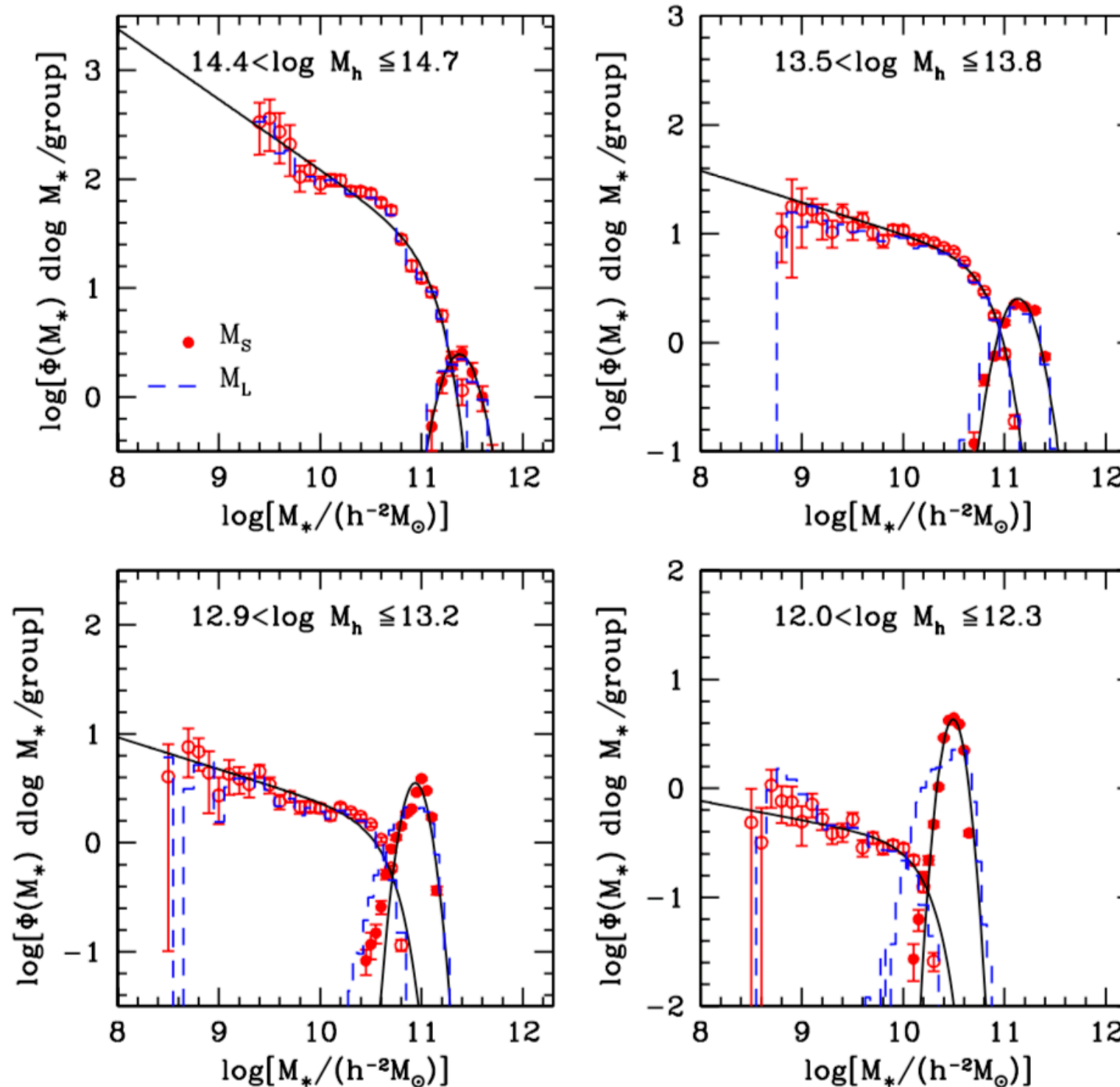
Dark Matter
Halos



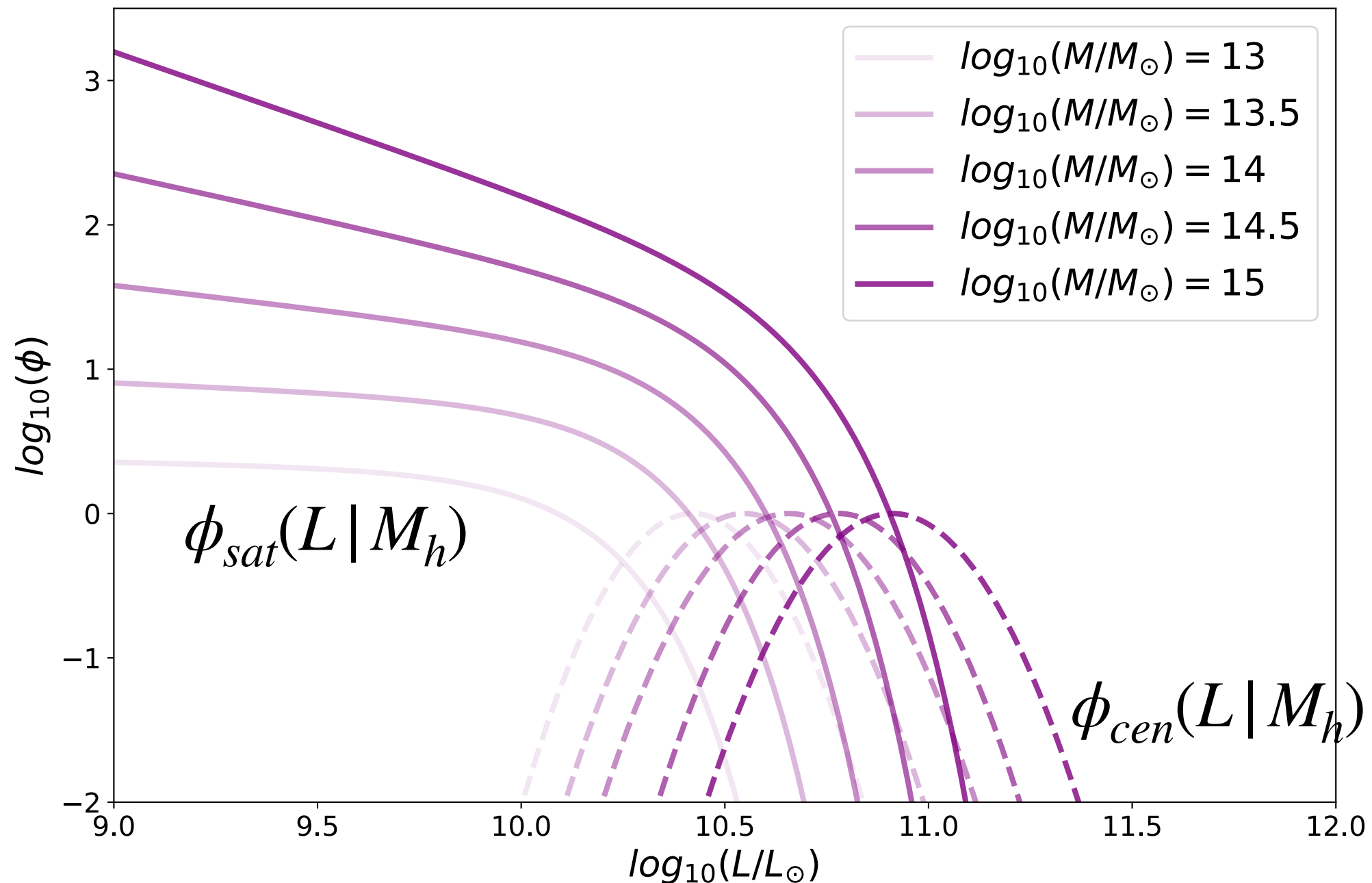
Galaxies with
Luminosity

Conditional Luminosity Function

(Yang, Mo, Van den Bosch 2009)



Conditional Luminosity Function



Conditional Luminosity Function

Satellite density as modified Schechter function:

$$\phi_{sat}(L | M_h) = \phi_s^* \left(\frac{L}{L_s^*} \right)^{(\alpha_s^*)} \exp \left\{ - \left(\frac{L}{L_s^*} \right)^2 \right\}$$

Follow the sub-halo mass function

Central density as modified log-normal function:

$$\phi_{cen}(L | M_h) = \frac{1}{\sqrt{2\pi}\sigma_c} \exp \left\{ - \frac{(\log L - \log L_c)^2}{2\sigma_c^2} \right\}$$

Canibal evol.
Can be at the most one.

Reverting the information

$$\mathcal{P}(M_h | L) = \frac{\mathcal{P}(L | M_h) \cdot \mathcal{P}(M_h)}{\mathcal{P}(L)}$$

POSTERIOR WE WANT TO EVALUATE

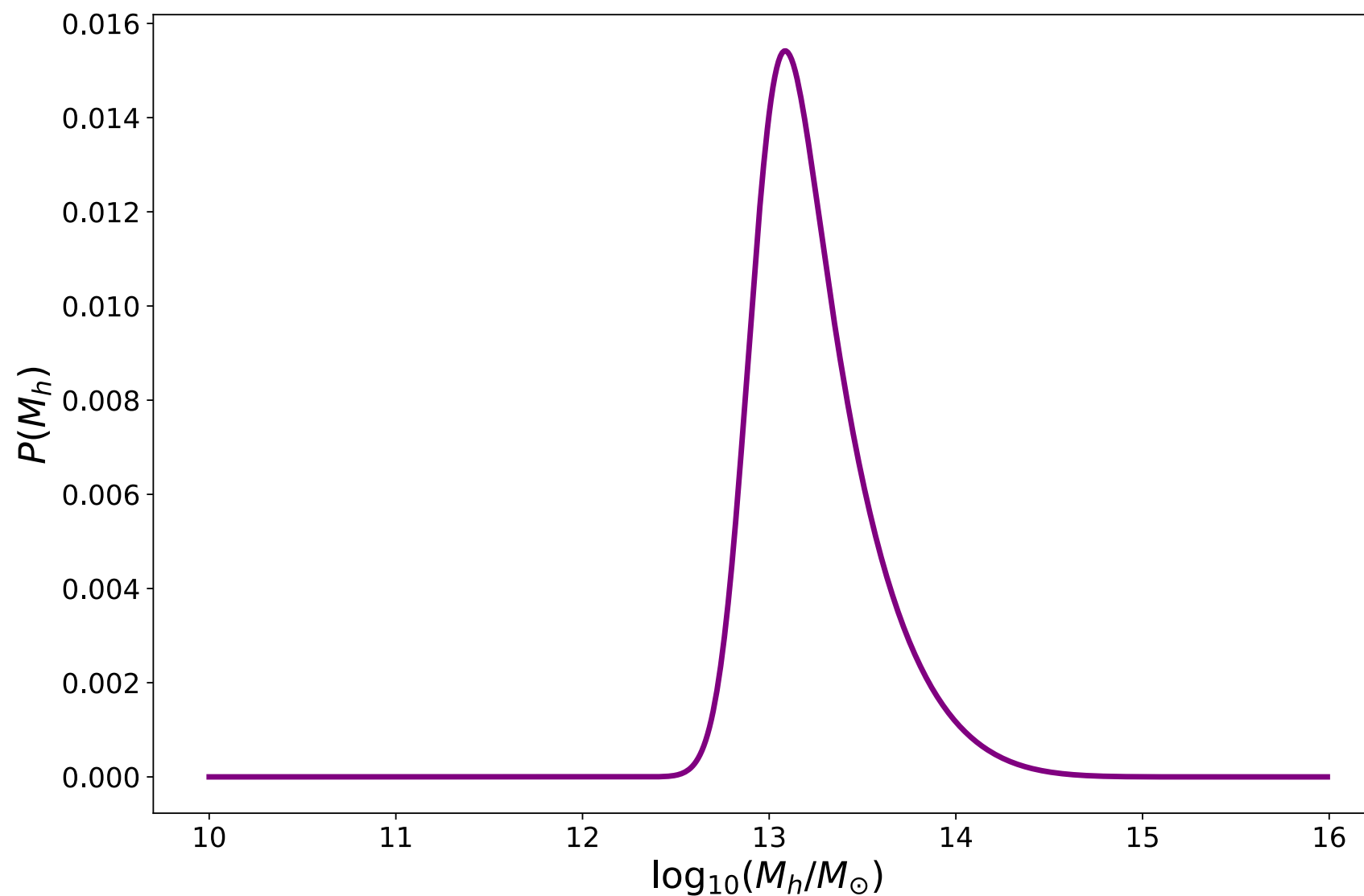
CONDITIONAL LUMINOSITY FUNCTION

CONVOLUTION OF HOD AND HMF

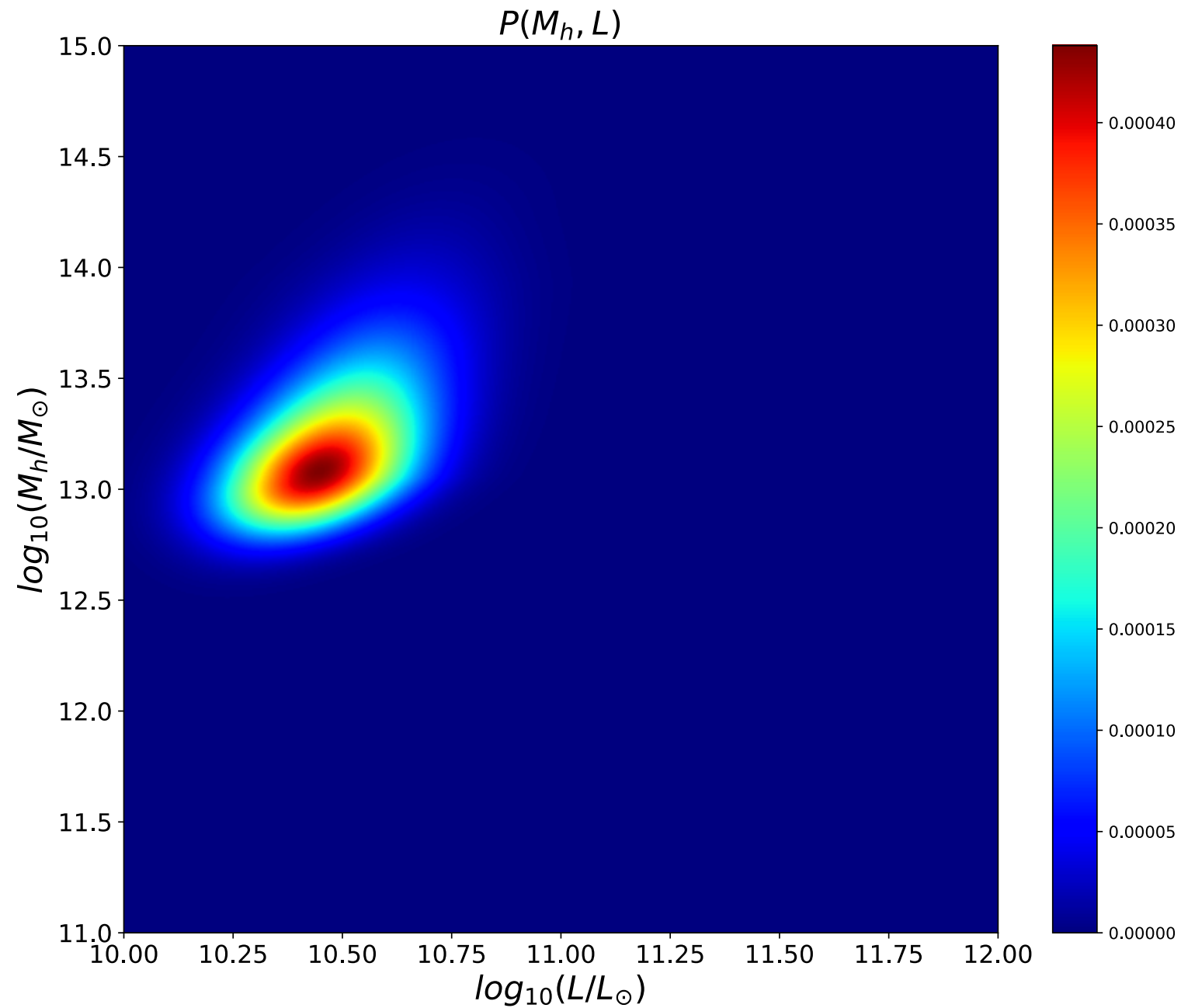
NORMALIZATION WHICH HAVE NOT IMPACT HERE

Prior estimation

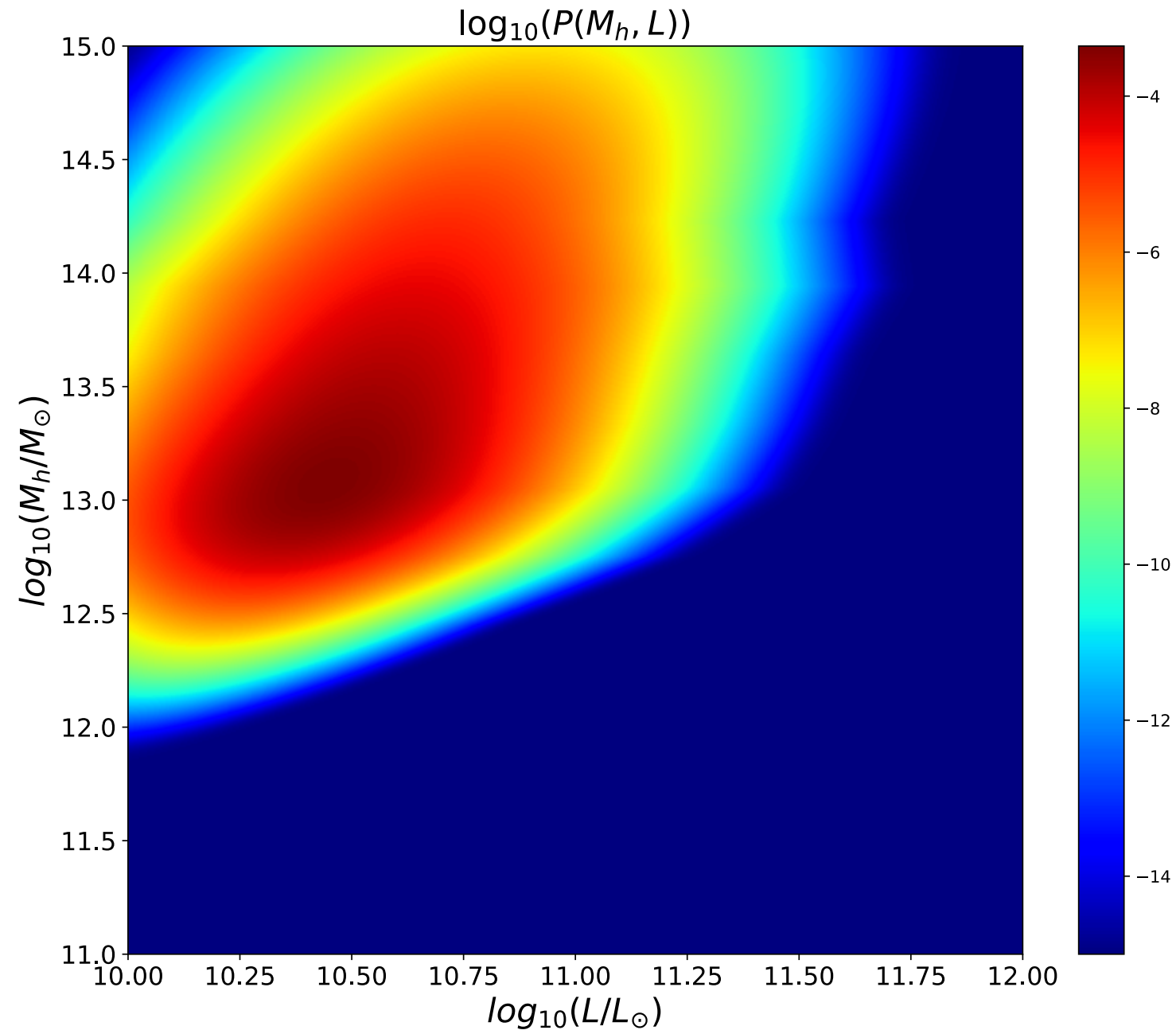
$$\mathcal{P}(M_h) = \frac{\mathcal{P}(N_{gal} \geq 1 | M_h) \times \frac{dn(M_h)}{d \log M}}{\int_0^\infty d \log M \mathcal{P}(N_{gal} \geq 1 | M_h) \times \frac{dn(M)}{d \log M}},$$



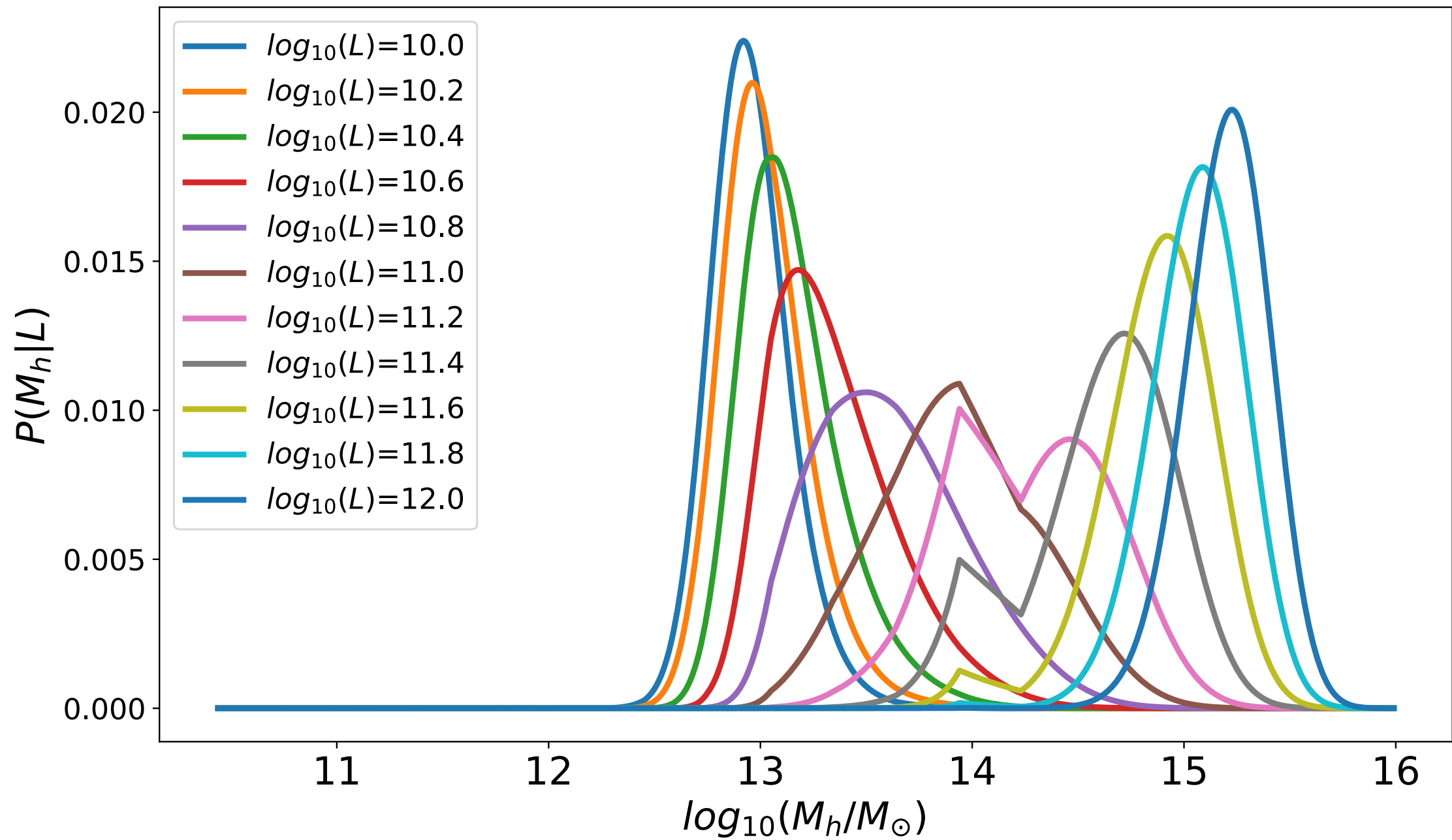
Joint Probability



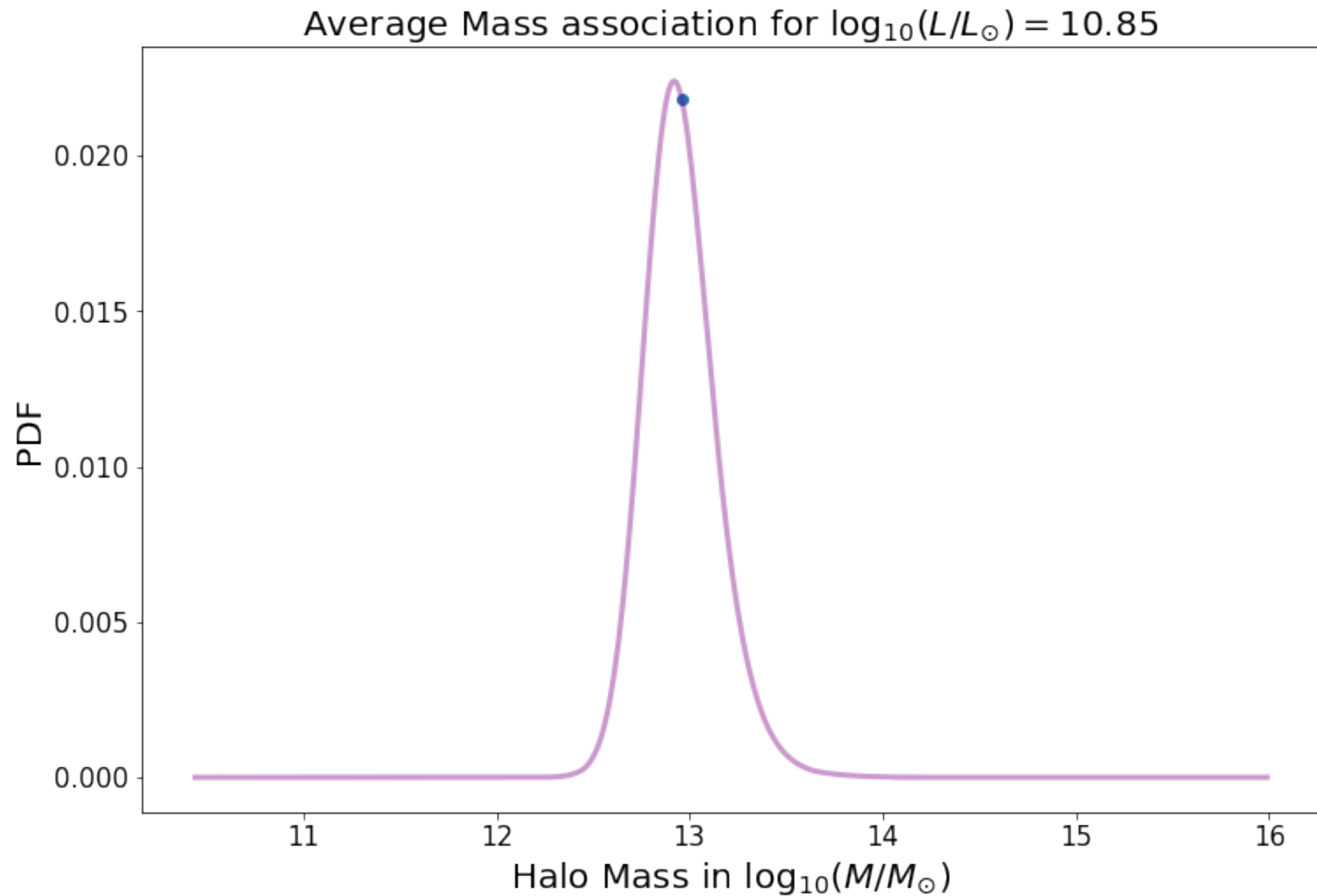
Joint Probability



Posterior Results

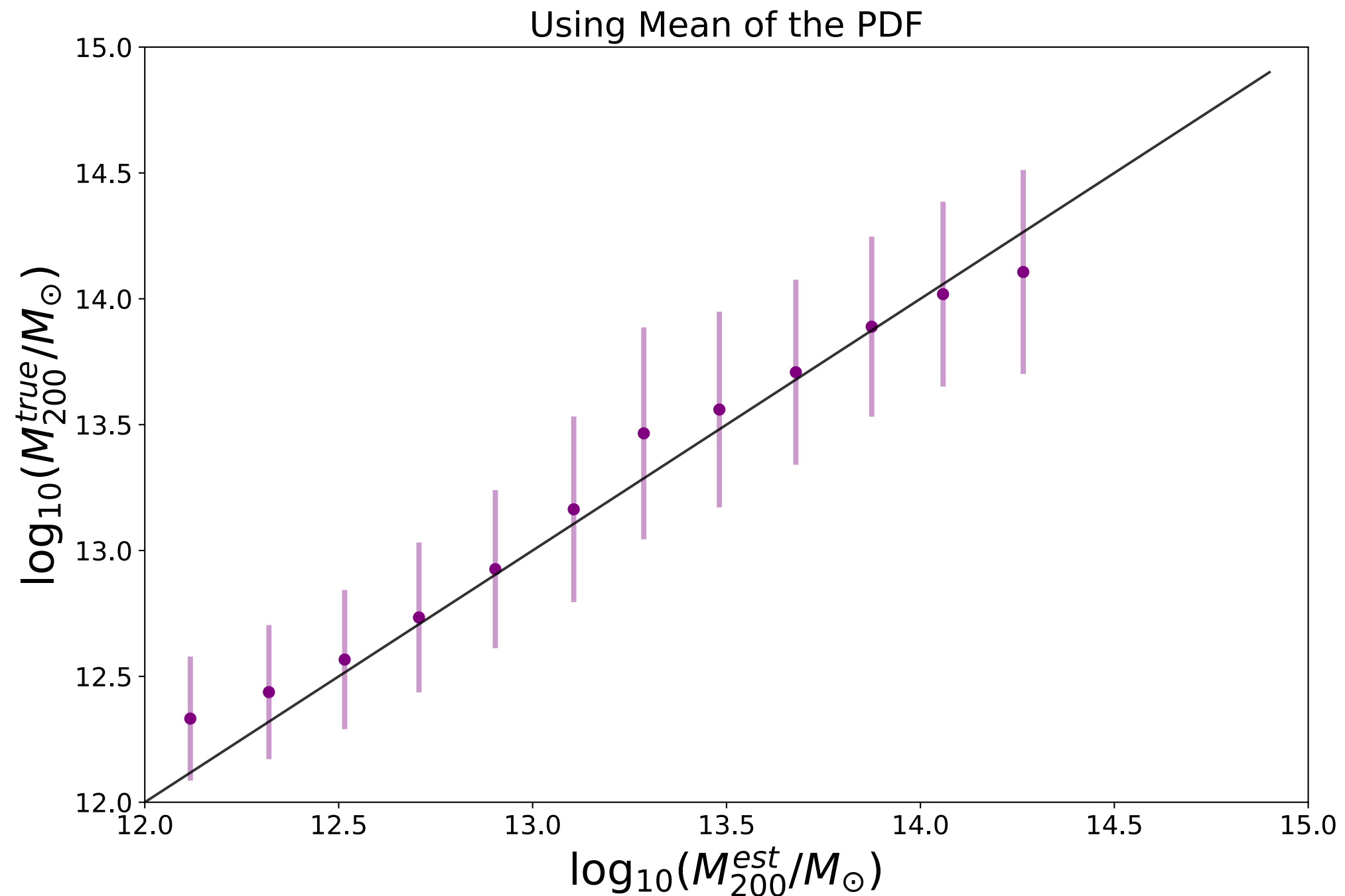


Average Mass



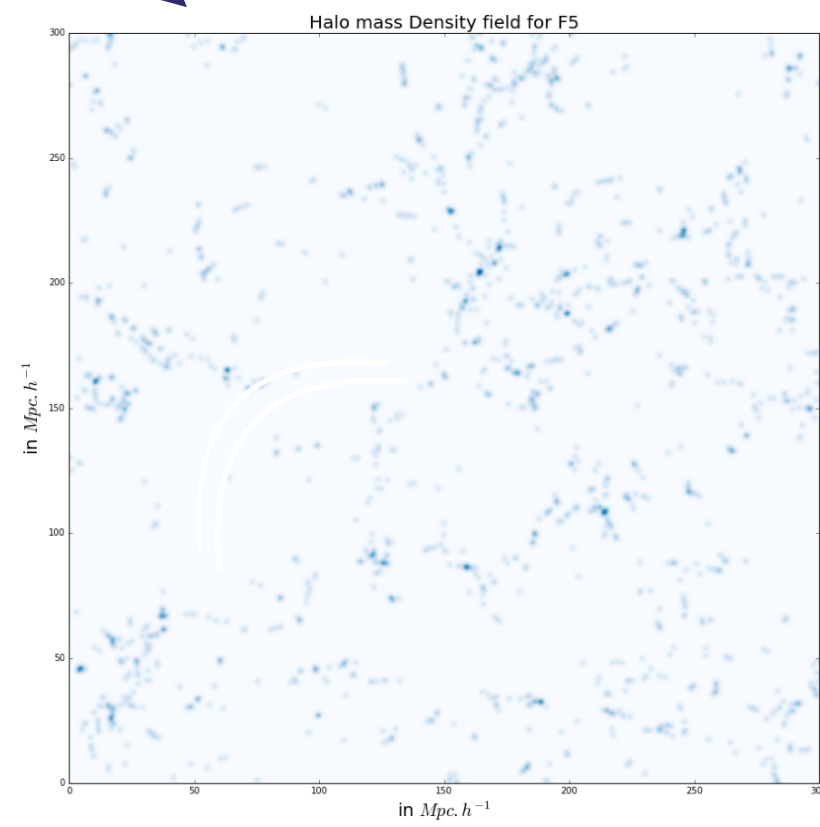
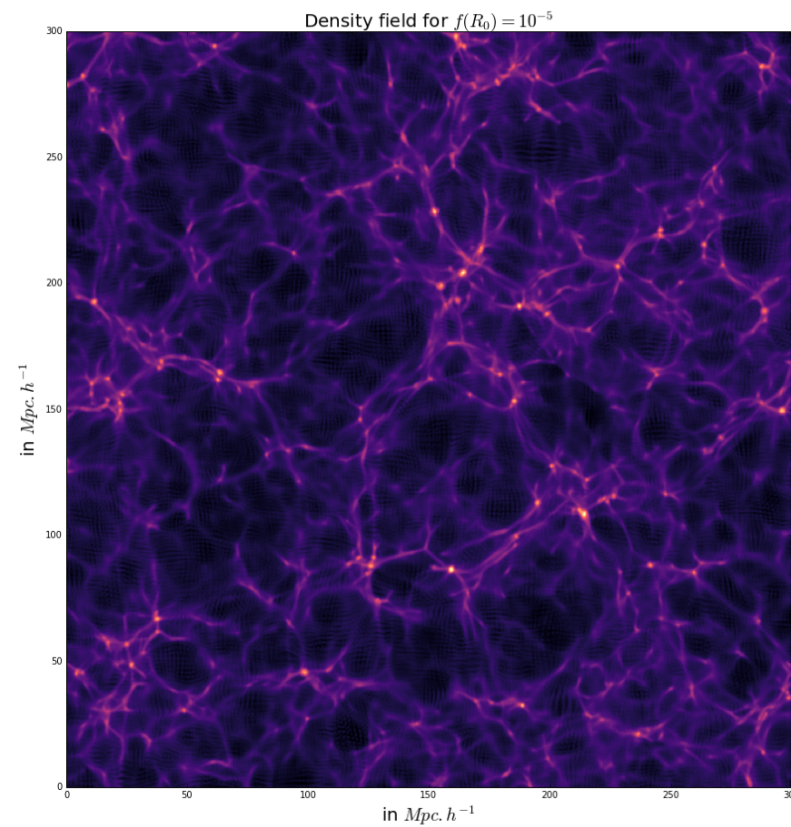
Estimation of the mass using L

In simulation



ROCKSTAR

HOD + CLF



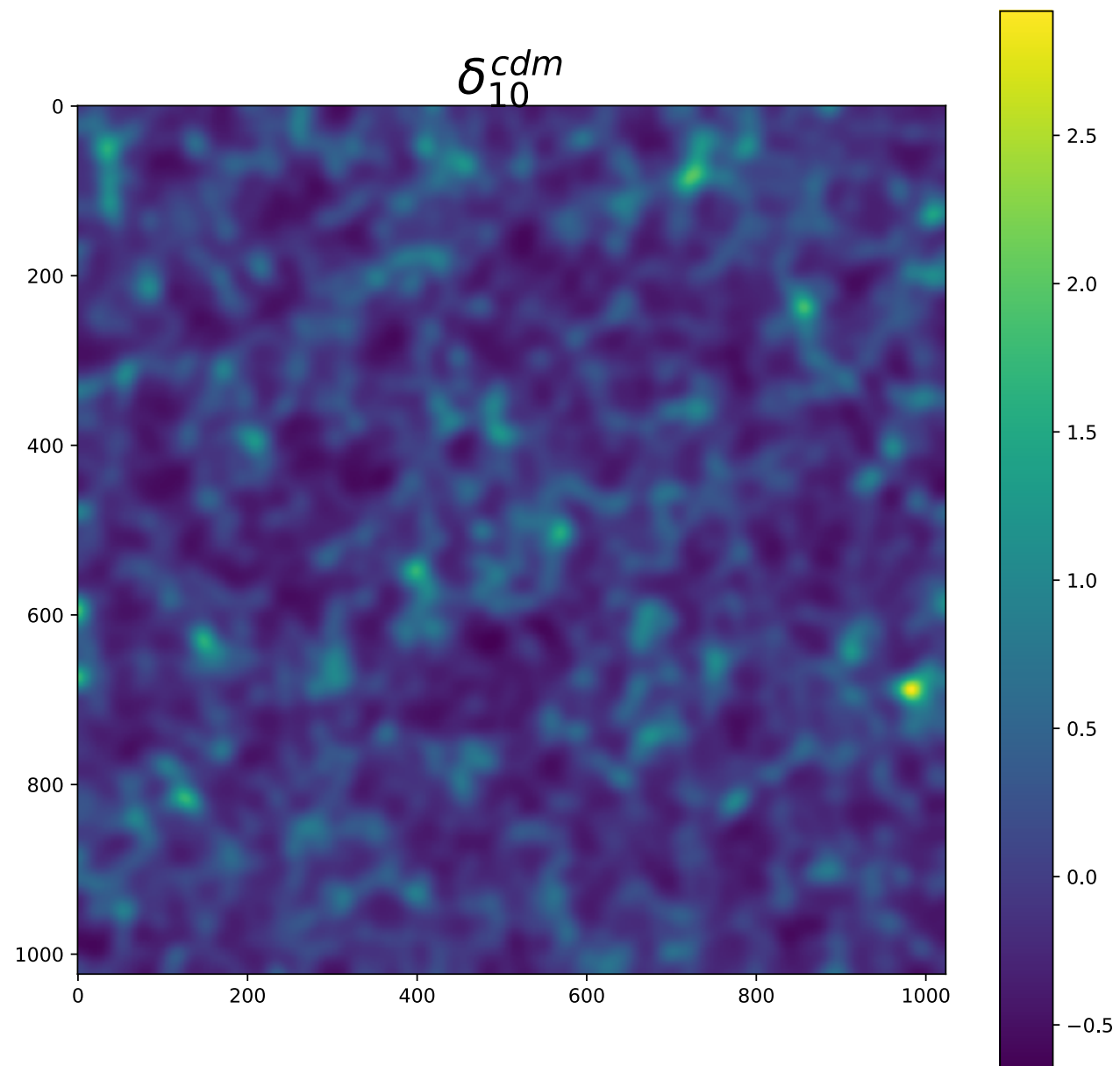
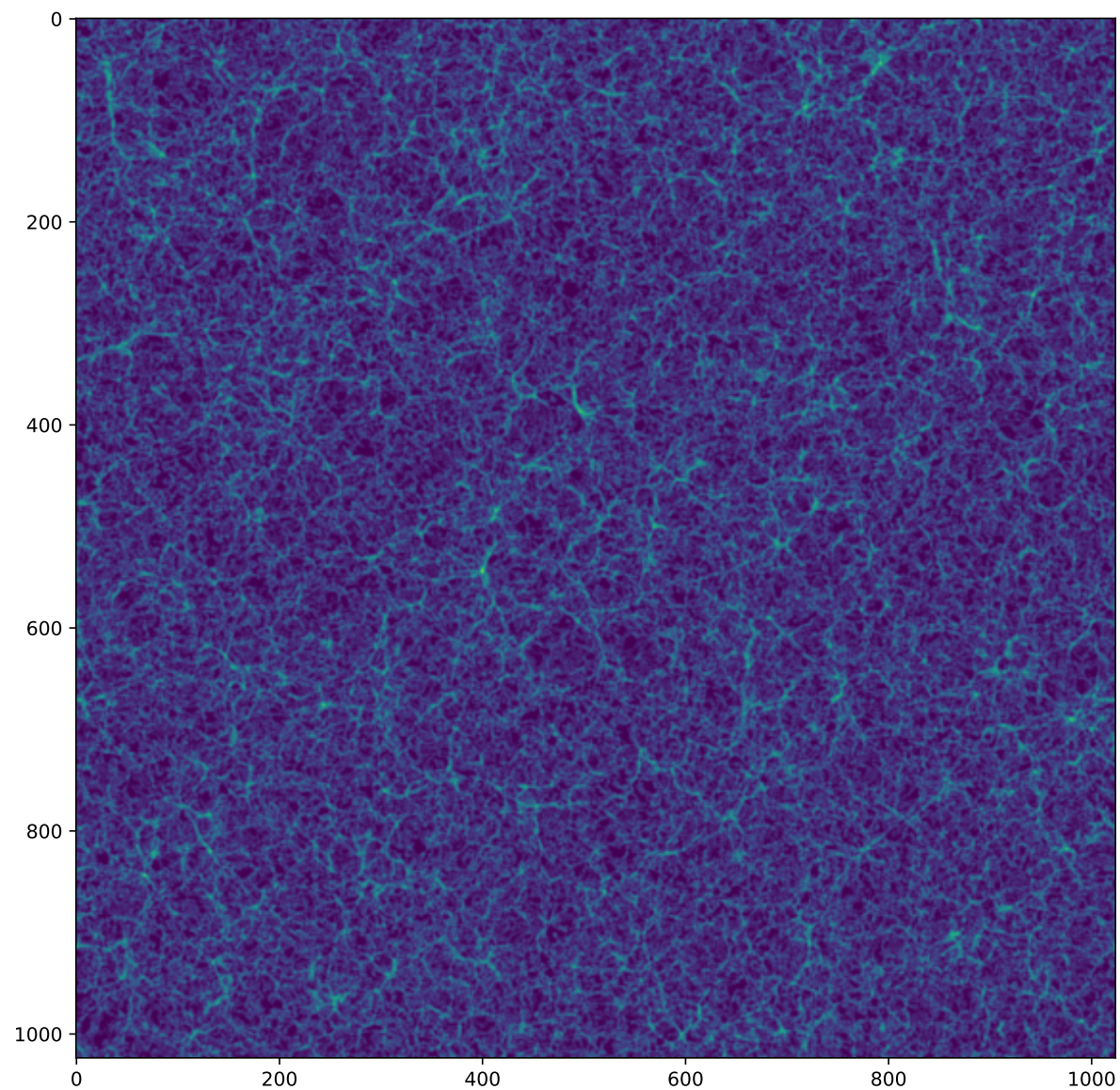
Dark Matter
particles

Dark Matter
Halos

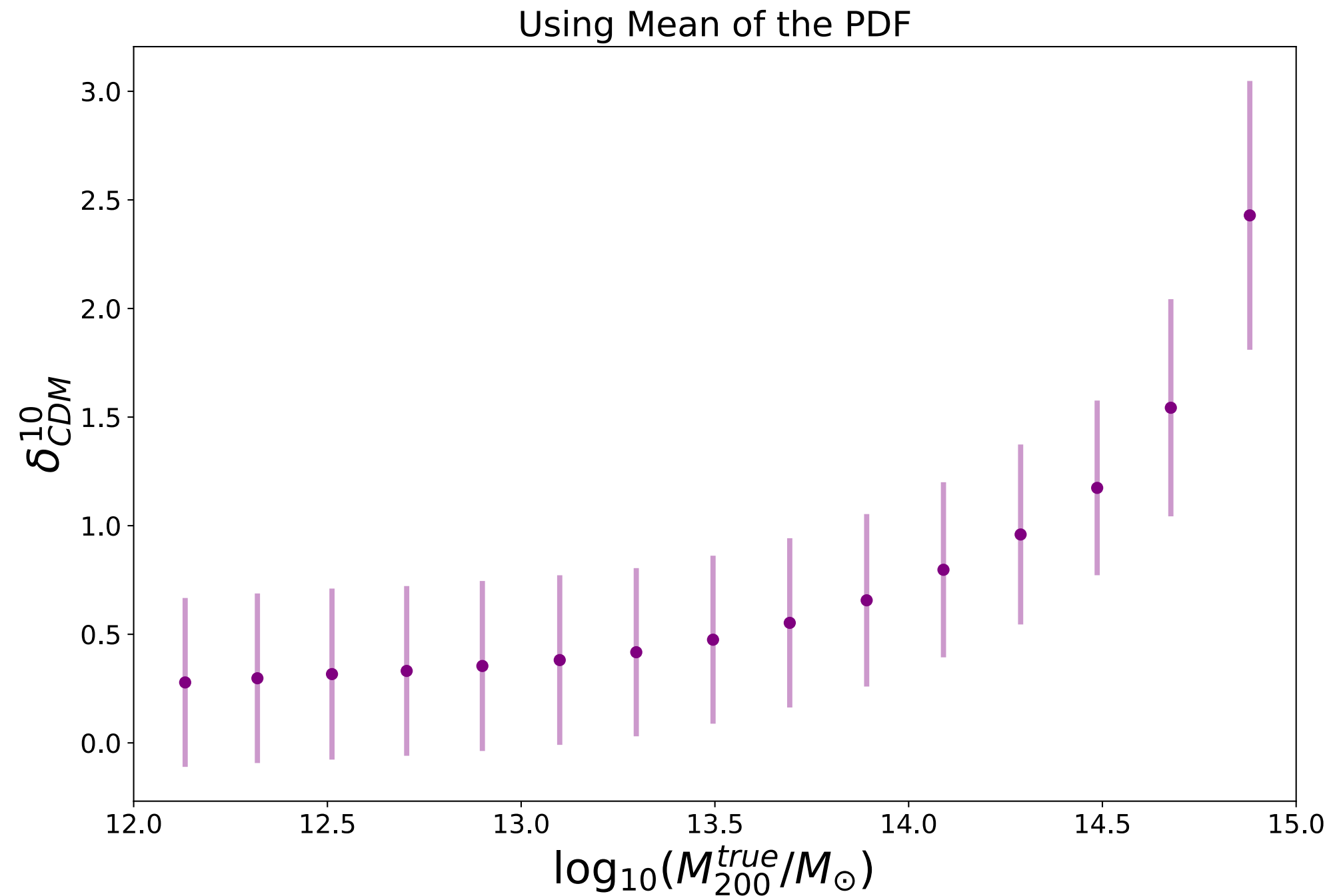
Galaxies with
Luminosity



Gaussian filter at 10 Mpc

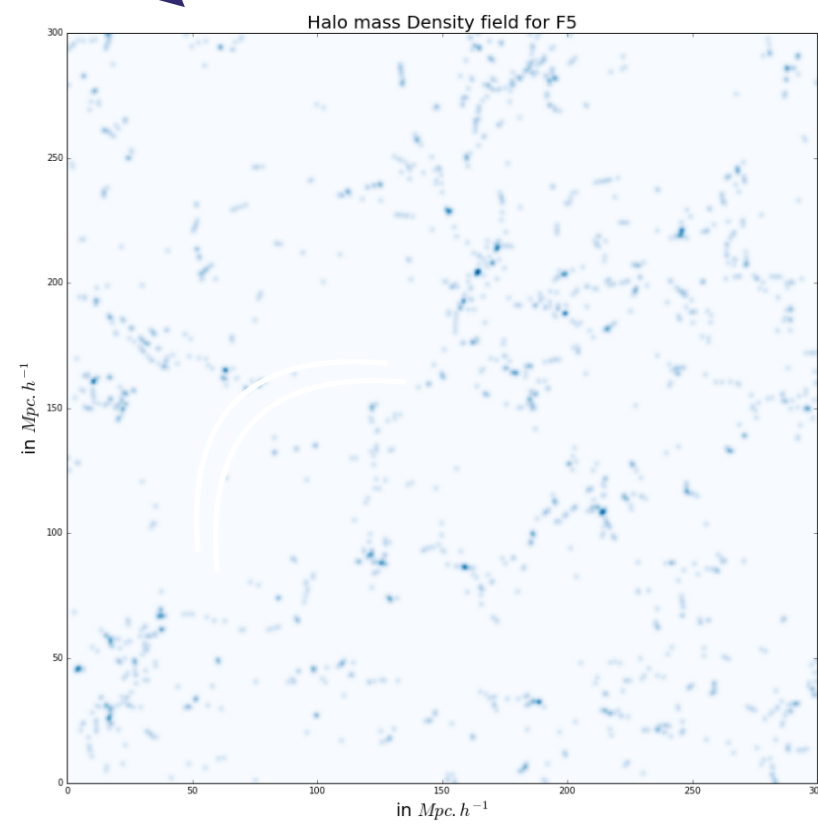
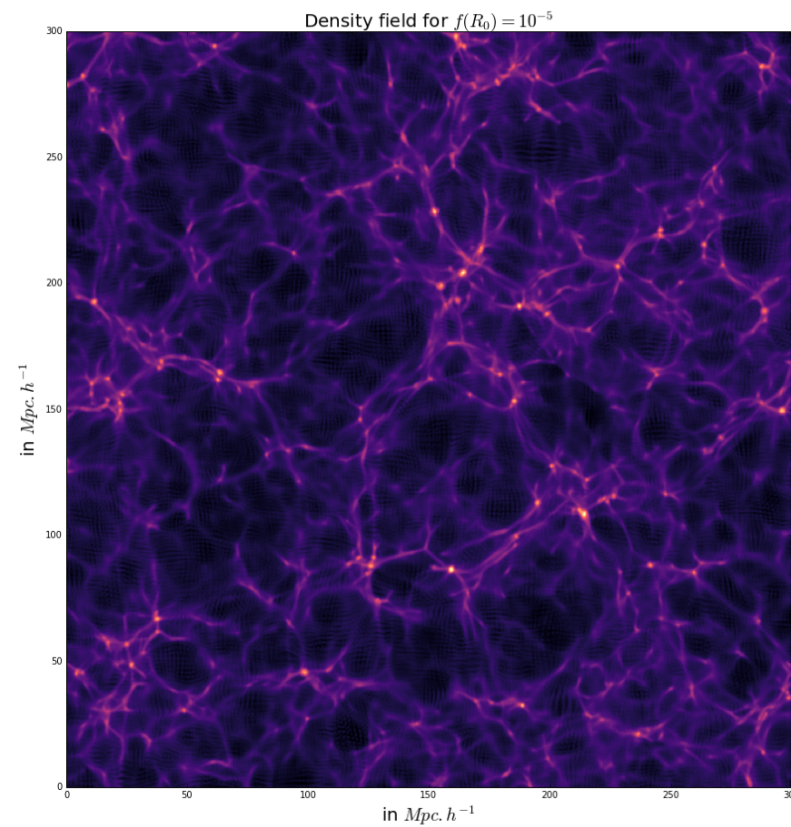


Gaussian filter at 10 Mpc



ROCKSTAR

HOD + CLF

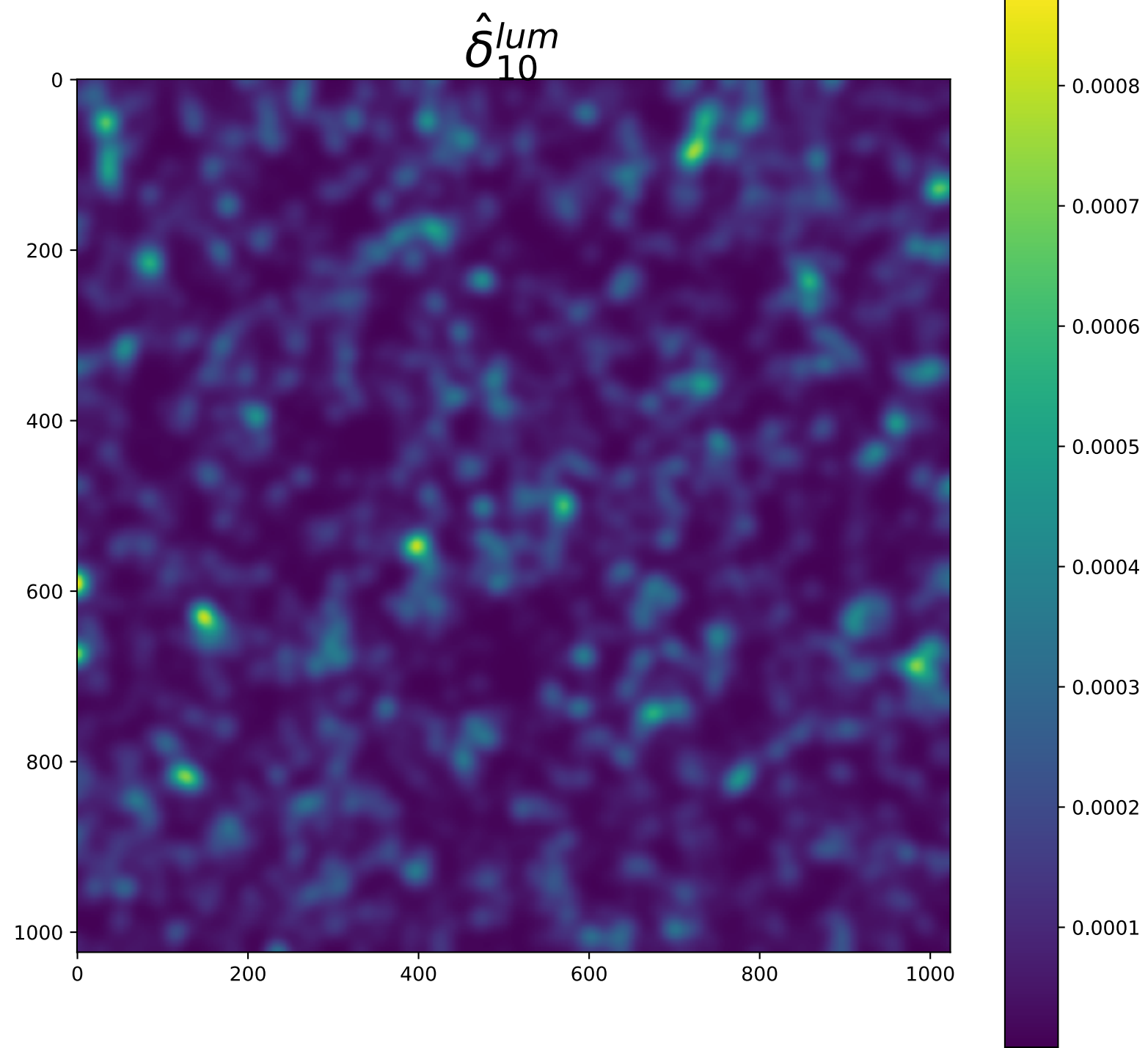


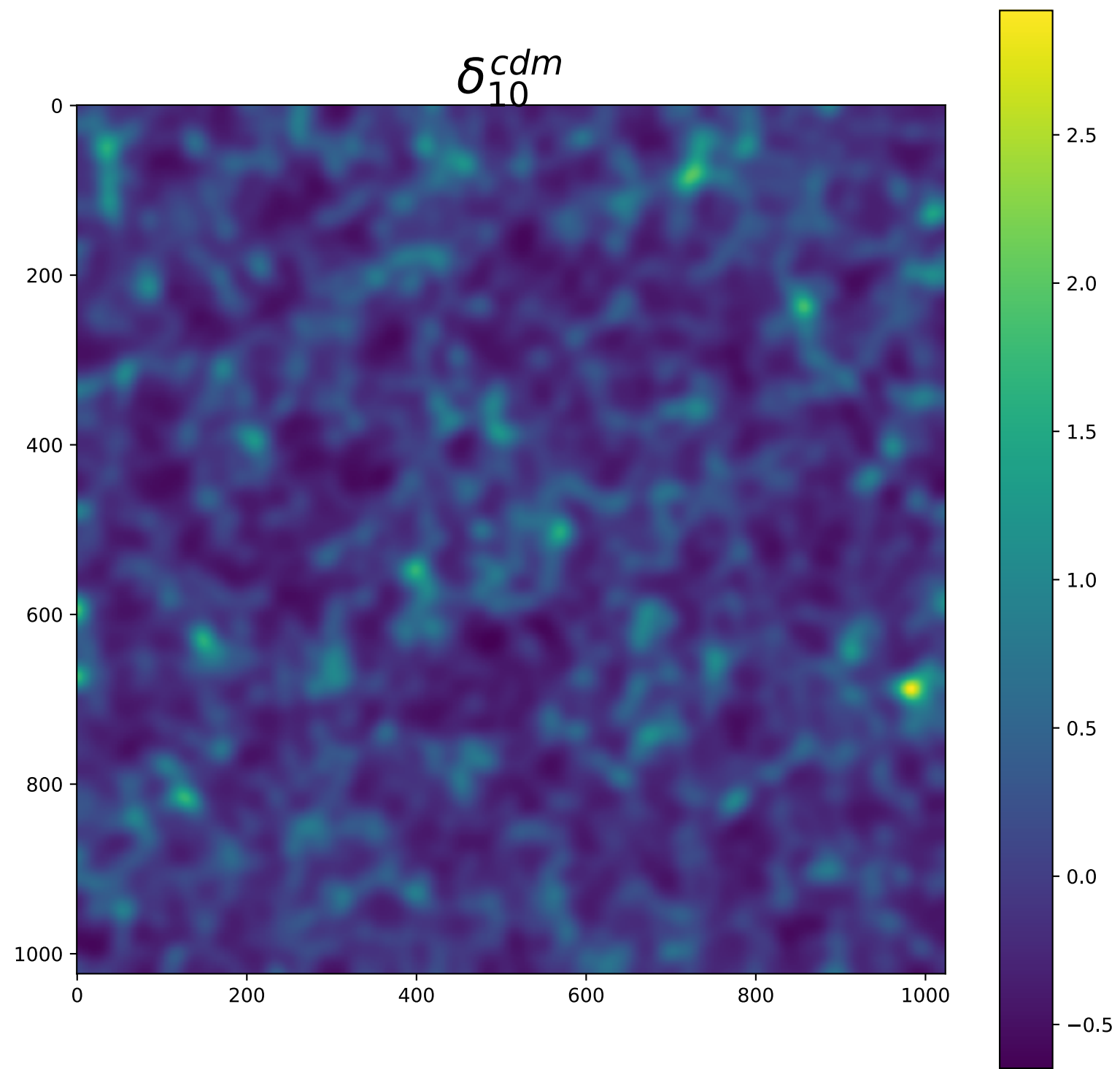
Dark Matter
particles

Dark Matter
Halos

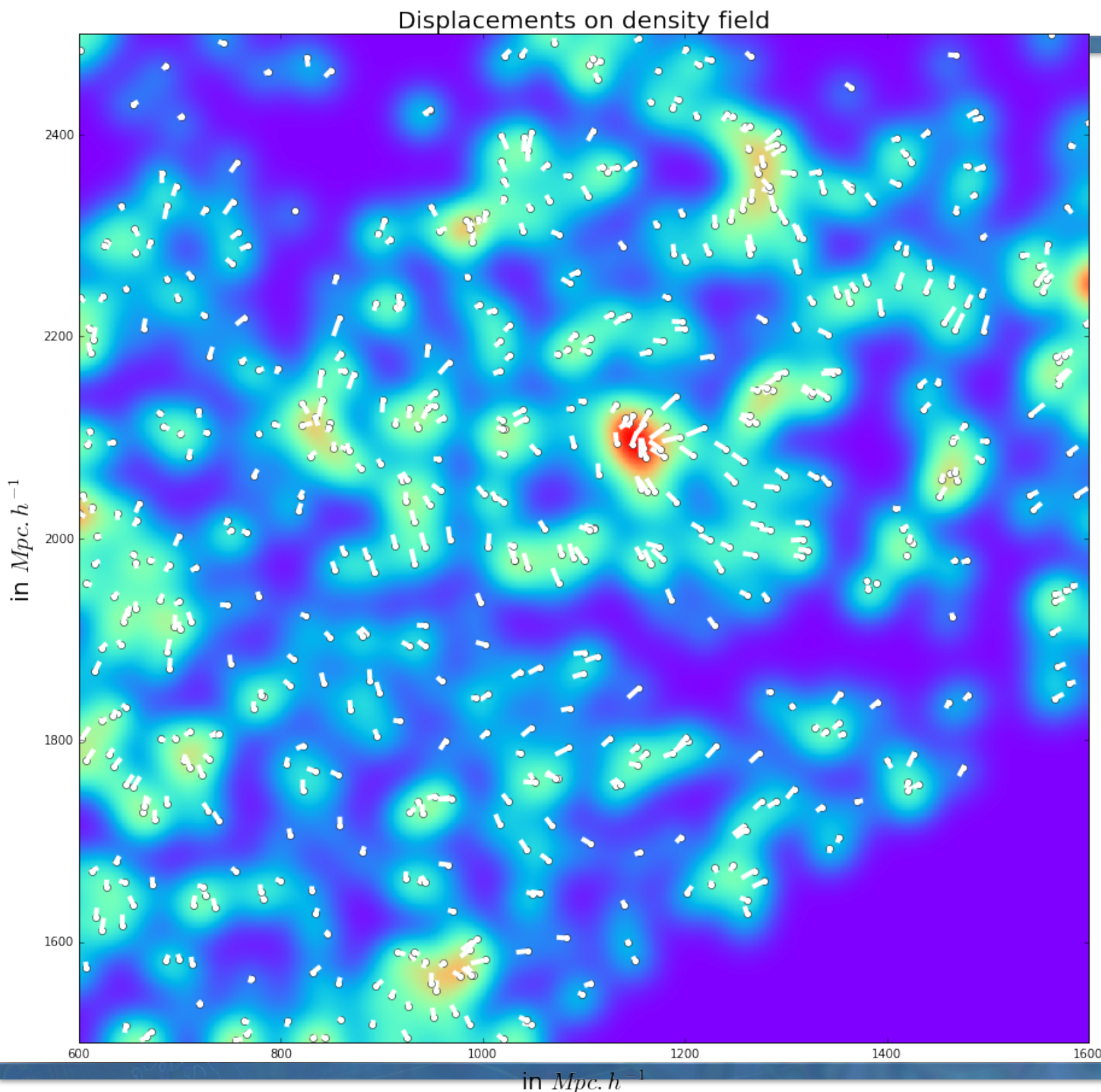
Galaxies with
Luminosity







Reconstruction to enhance BAO peak



Vargas, Ho, Fromenteau, Cuesta (2017)

We estimate the smoothed galaxy density field



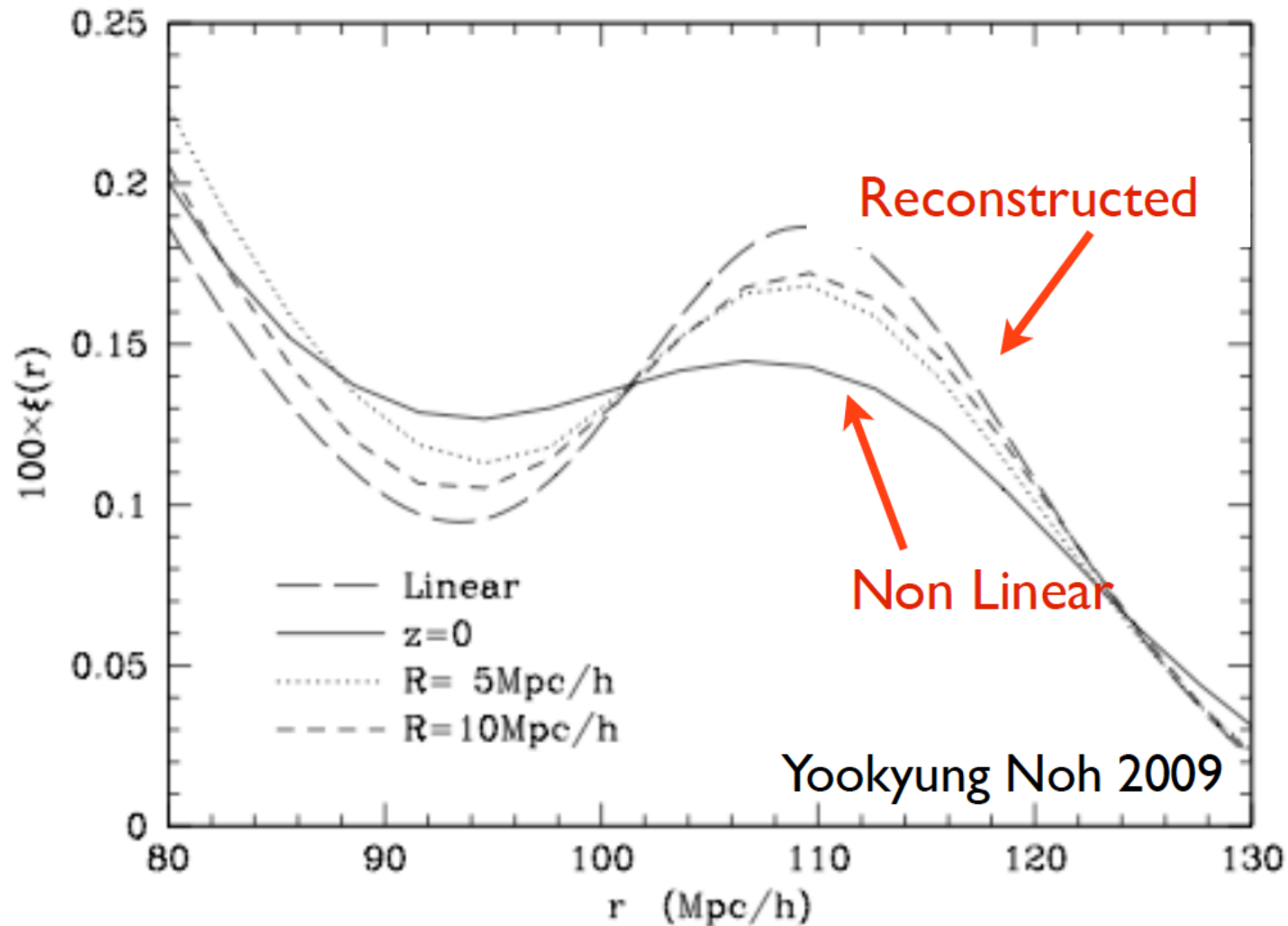
Using linear bias, we derive the matter density field



We move back the galaxies to their original position

$$\vec{\psi}(\vec{k}) = \frac{-i\vec{k}}{k^2} \frac{\delta_g(\vec{k})}{b_g}$$

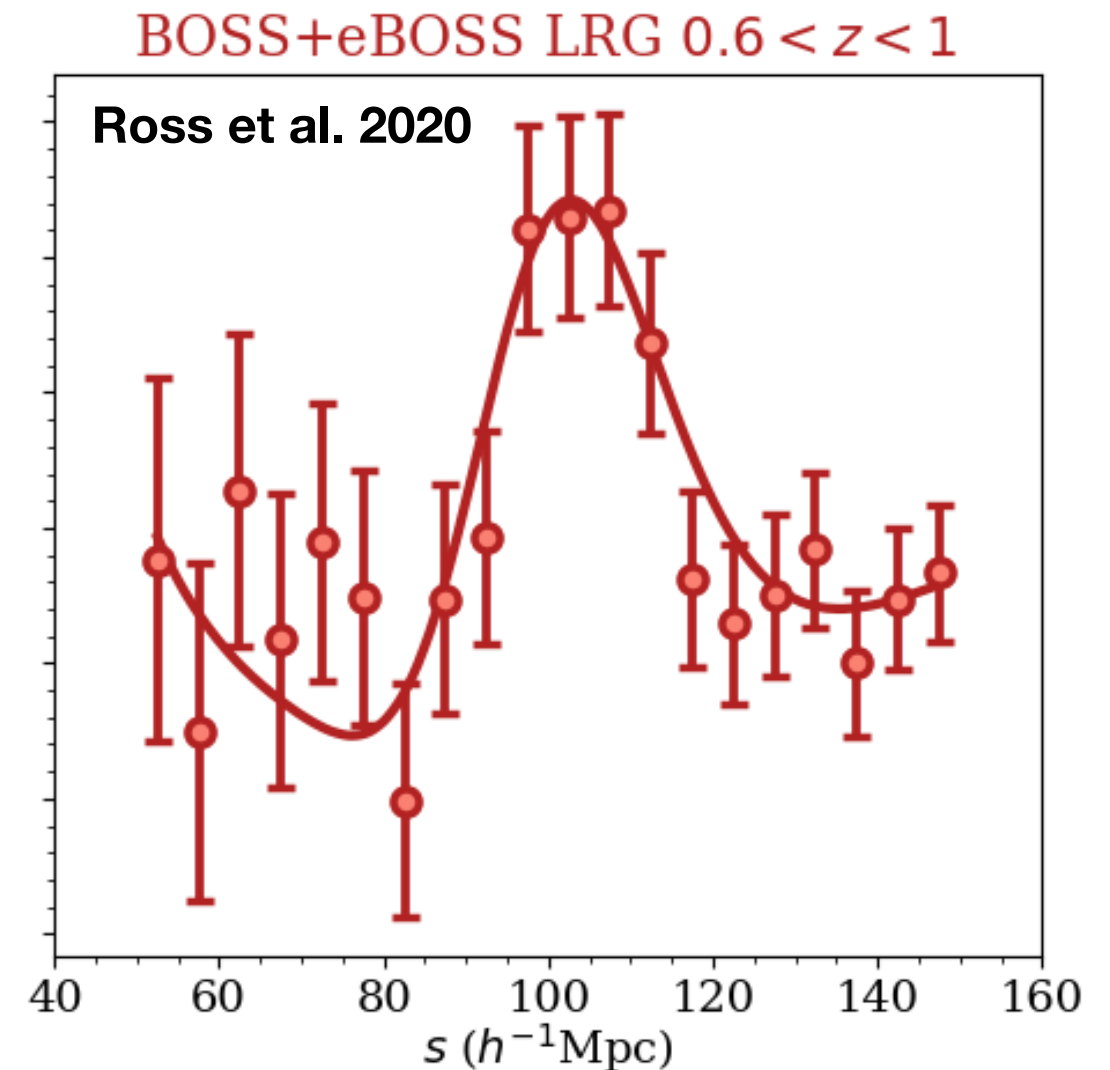
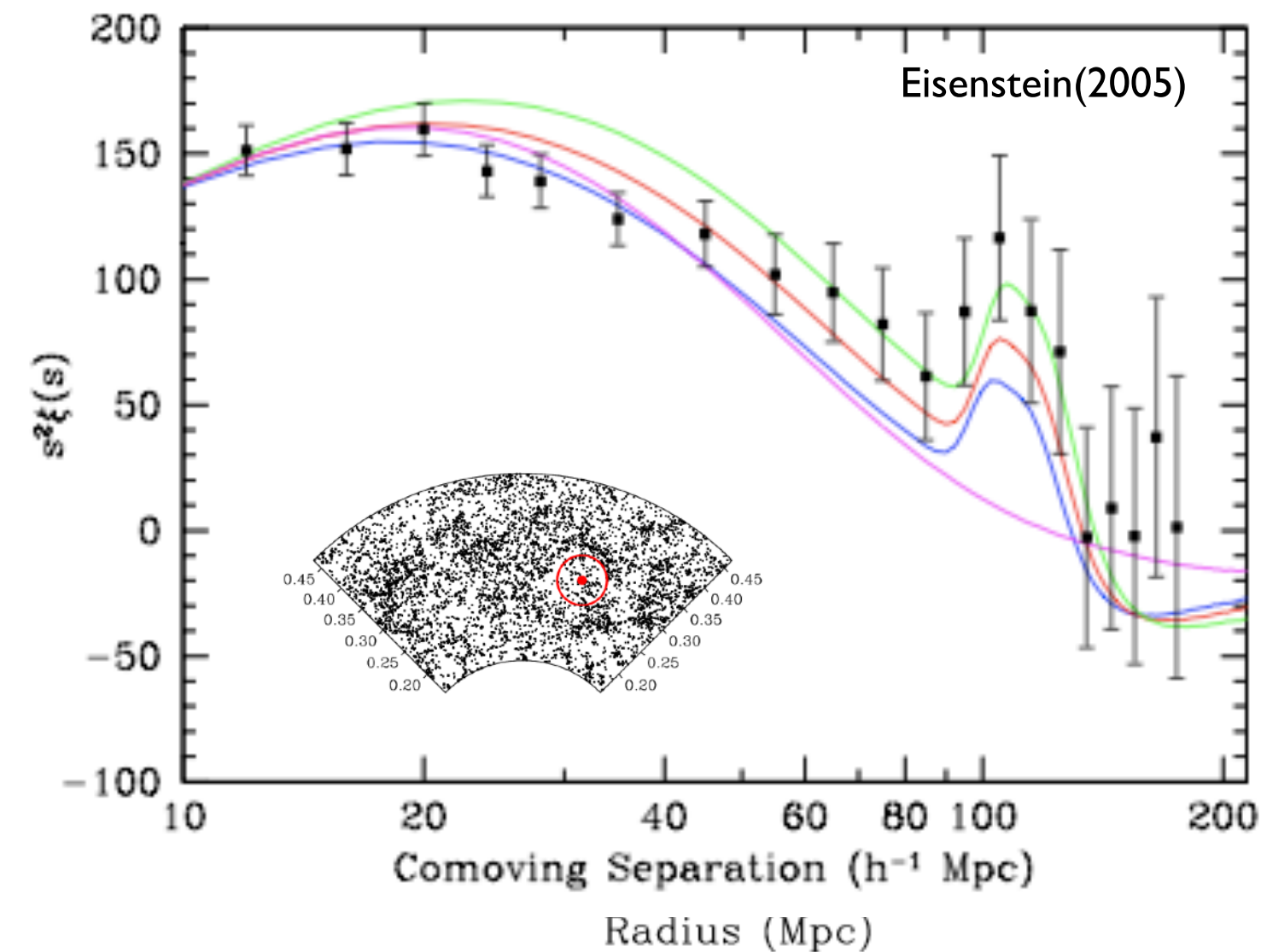
Reconstruction to enhance BAO peak



Is part of the standard
BAO analysis

Baryonic Acoustic Oscillations Peak

BAO Detection in the correlation function of Luminous Red Galaxies



Mariana's talk on Saturday

Conclusions

Obtain non-linear local bias information using luminosity information in a probabilistic way

This method can be generalized to any statistical information connecting the tracer to the halo mass (for example density)

We can use this method in order to implement a redshift dependent reconstruction inside the large surveys

Different tracers



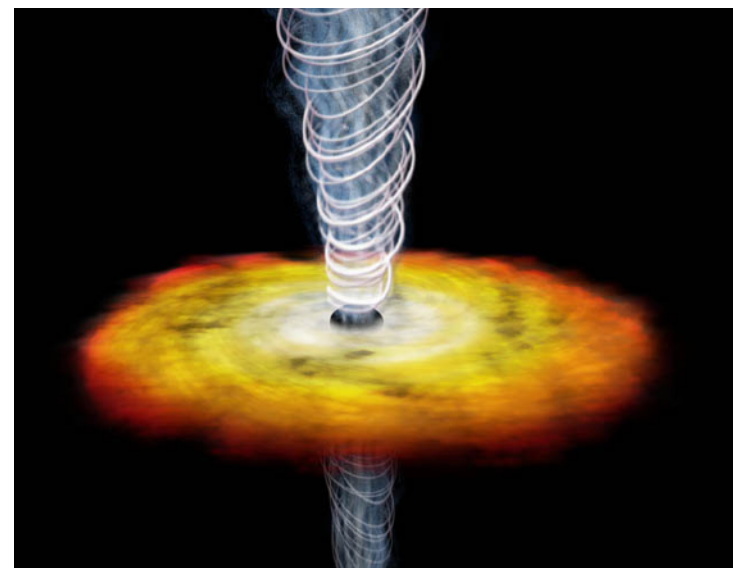
**Luminous Red
Galaxy**

LRG



**Emission Line
Galaxy**

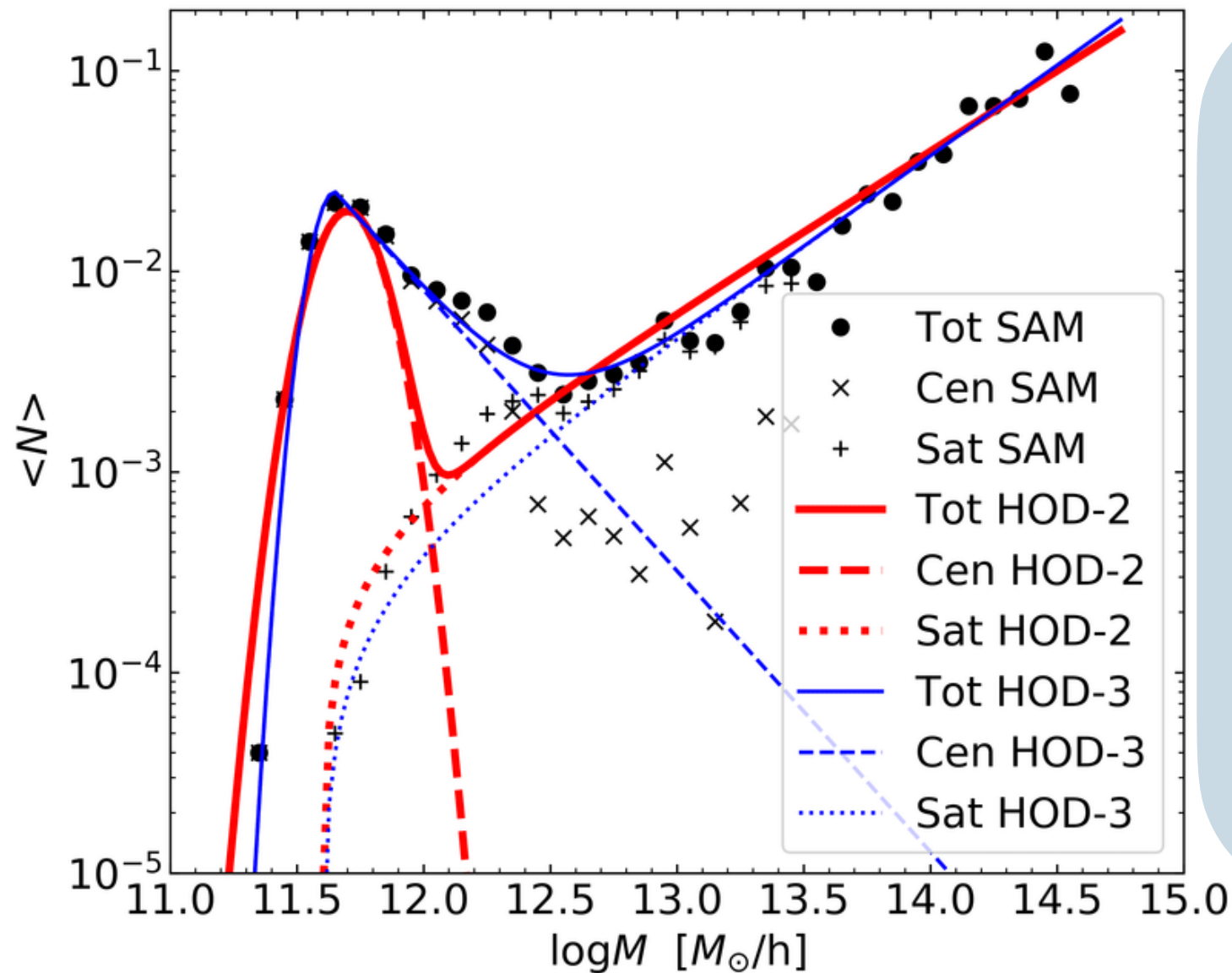
ELG



Quasar

QSO/AGN

Emission Line galaxies



Avila et al 2020

**Important new tracer for
high redshift**

**The galaxy Halo connection
is very complex**

So galaxy bias is highly non-trivial

The background of the slide is a Cosmic Microwave Background (CMB) radiation map, showing a dense field of orange and red spots against a dark purple background, representing the early universe's temperature fluctuations.

THANKS !

Halo Occupation Distribution

1 HALO TERM

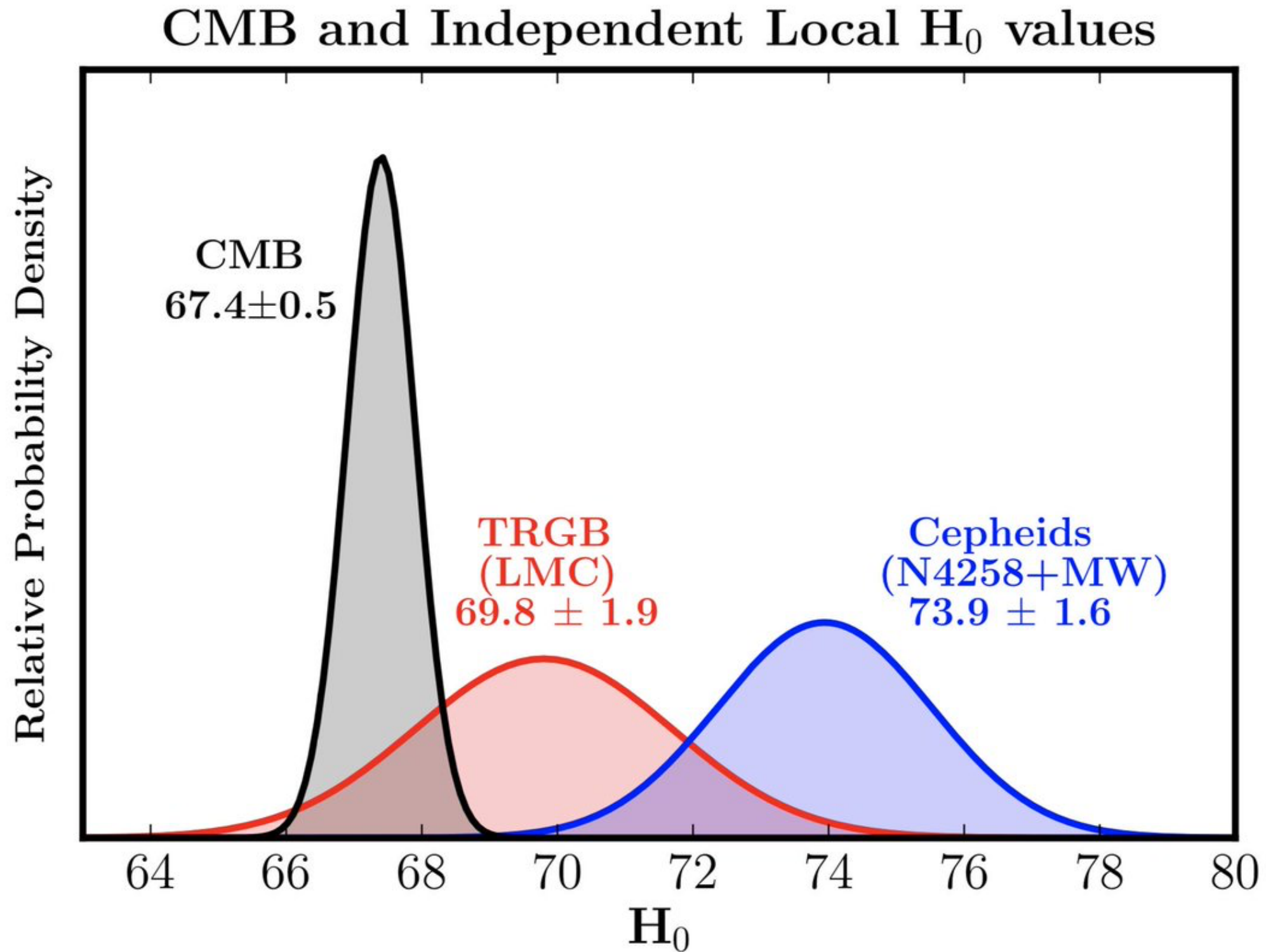
$$\xi_{gg}(r) \propto \int d \log(M) \frac{dn}{d \log(M)} \langle N(M) \rangle (\langle N(M) \rangle - 1) P_{NFW}(r | M)$$

$$+ \int d \log(M_1) \int d \log(M_2) \langle N(M_1) \rangle \cdot \langle N(M_2) \rangle b^2 \xi_{mm}(r) \times \\ \times \frac{P_{NFW}(r_1 | M_1) dn}{d \log(M_1)} \frac{P_{NFW}(r_2 | M_2) dn}{d \log(M_2)}$$

2 HALO TERM

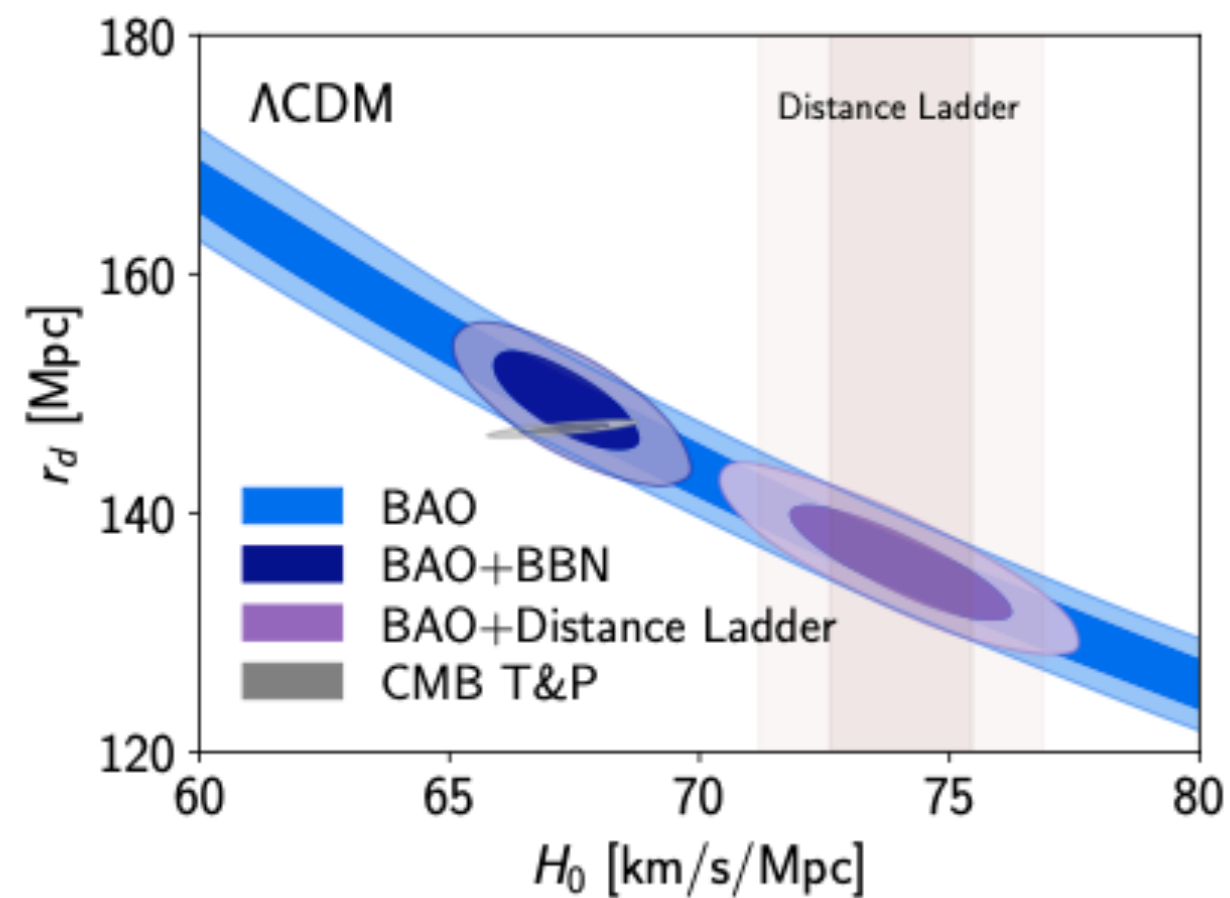
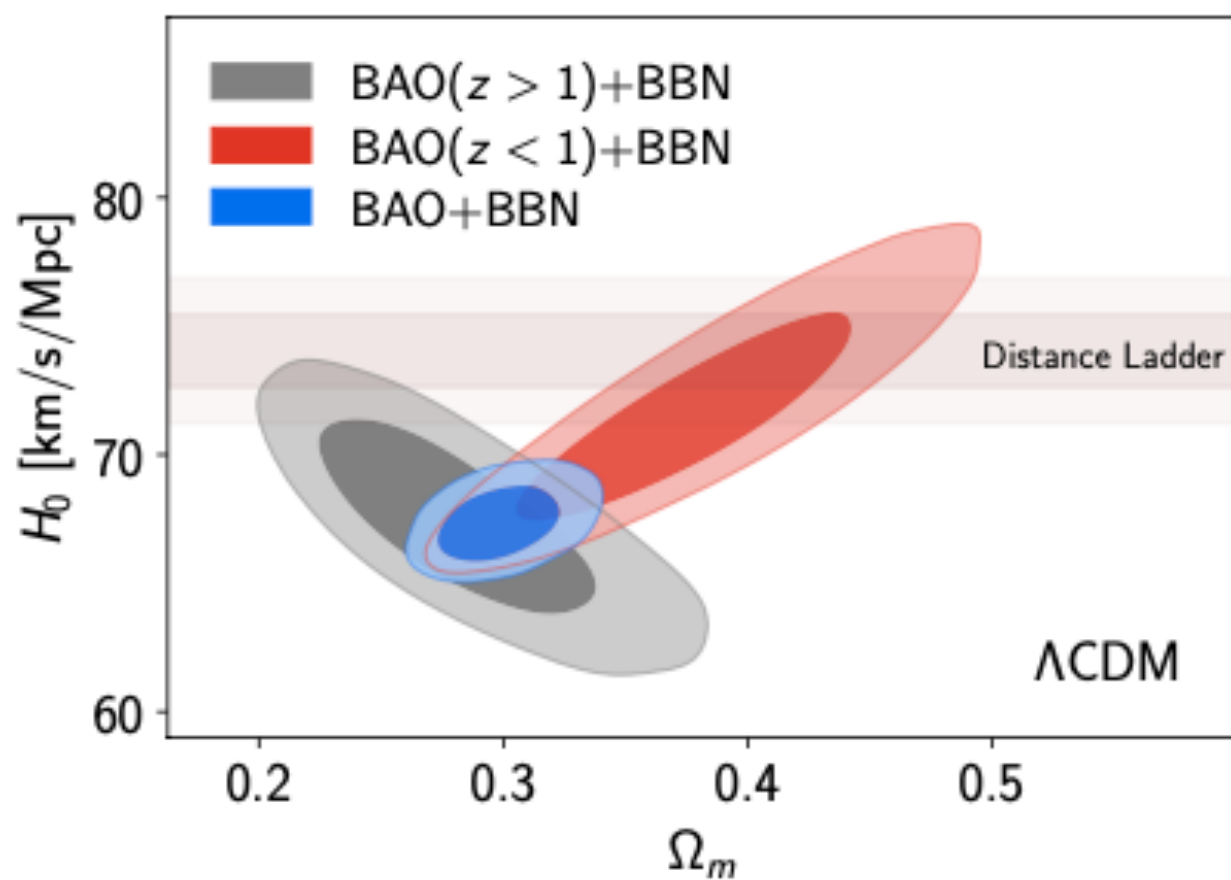
$$P_{NFW}(r | M) = \frac{\rho_{NFW}(r | M) \times 4\pi r^2 dr}{M}$$

Increase H_0 crisis



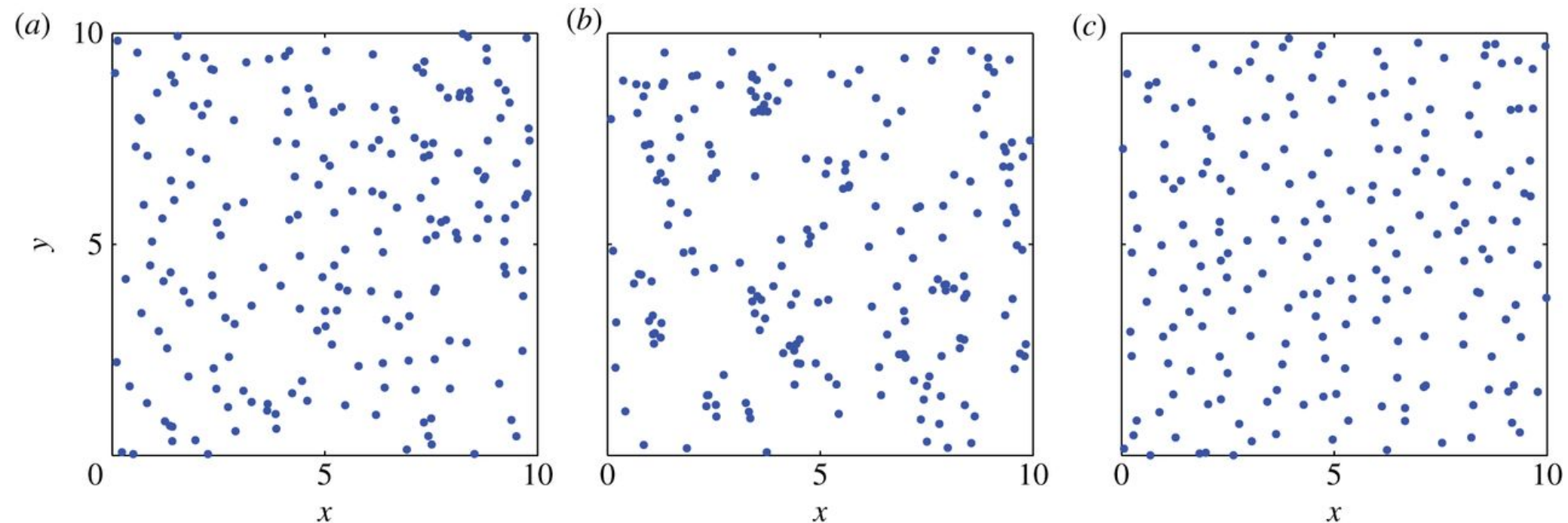


Increase H_0 crisis



eBOSS collaboration 2020 (in prep.)

2-point correlation function



We count the number of pairs of galaxies for each distance r

We compare with the number obtain for random distribution

Estimator

$$\xi(r) = \frac{N_{pairs}(r) - N_{pairs_{rand}}(r)}{N_{pairs_{rand}}(r)}$$

Quantity

$$P(r) = 4\pi r^2 dr \cdot \bar{n} [1 + \xi(r)]$$

2-point correlation function

