

# Quark mass matrices in the physical basis

**Virendra Gupta**

Departamento de Física Aplicada  
CINVESTAV-IPN, Unidad Mérida

# Introduction

The quark mixing matrix  $V$  can be viewed in three basis.

1. Physical basis:  $M_u, M_d$  diagonal.
2.  $M_u$  diagonal
3.  $M_d$  diagonal

$V \rightarrow 3$  angles, 1 phase.

## Basic formulas

The  $3 \times 3$  hermitian quark mass matrix  $M_q$  is diagonalized by  $V_q$  so that  $M_q = V_q^\dagger \hat{M}_q V_q$ ,  $q=u,d$ .

The eigenvalues are denoted by  $(\lambda_u, \lambda_c, \lambda_t)$  and  $(\lambda_d, \lambda_s, \lambda_b)$  for the up and down quark mass matrices.

Each mass matrix can be expressed in terms of its projectors

$$M_u = \sum_{\alpha=u,c,t} \lambda_\alpha N_\alpha \quad \text{and} \quad M_d = \sum_{j=d,s,b} \lambda_j N_j. \quad (1)$$

## Basic formulas

$$|V_{\alpha j}|^2 = \text{Tr}[N_{\alpha}N_j], \quad (2)$$

$$N_{\alpha} = \frac{(\lambda_{\beta} - M_u)(\lambda_{\gamma} - M_u)}{(\lambda_{\beta} - \lambda_{\alpha})(\lambda_{\gamma} - \lambda_{\alpha})} \quad (3)$$

$$N_j = \frac{(\lambda_k - M_d)(\lambda_l - M_d)}{(\lambda_k - \lambda_j)(\lambda_l - \lambda_j)}, \quad (4)$$

with  $(\alpha, \beta, \gamma)$  and  $(j, k, l)$  any permutation of  $(u, c, t)$  and  $(d, s, b)$ .

## Basic formulas

The Jarlskog invariant  $J(V)$ , which is a measure of CP-violation can be directly expressed in terms of  $M_u$  and  $M_d$ .

$$\text{Det}([M_u, M_d]) = 2i D(\lambda_\alpha) D(\lambda_j) J(V),$$

$$D(\lambda_\alpha) = (\lambda_c - \lambda_u)(\lambda_t - \lambda_u)(\lambda_t - \lambda_c)$$

$$D(\lambda_j) = (\lambda_s - \lambda_d)(\lambda_b - \lambda_d)(\lambda_b - \lambda_s).$$

If one mass matrix is diagonal then

$$J(V) = \text{Im} (M_{12} M_{23} M_{13}^*) / D(\lambda)$$

## Fritzsch type Mass Matrices

$$M_u = \begin{pmatrix} 0 & A & 0 \\ A^* & 0 & B \\ 0 & B^* & C \end{pmatrix}, \quad M_d = \begin{pmatrix} 0 & A' & 0 \\ A'^* & 0 & B \\ 0 & B'^* & C' \end{pmatrix}$$

Without lack of generality we can take  $C$  and  $C'$  to be positive and  $A$  and  $B$  to be real and positive. Then,  $M_u$  and  $M_d$  have eight real parameters  $A, B, C, C', |A'|, |B'|$  and the phases  $\phi_{A'}$  and  $\phi_{B'}$ .

## Gupta-Rajpoot type mass matrices

For the three families of quarks the Gupta-Rajpoot mass matrices are of the form

$$M_u = \begin{pmatrix} 0 & A & 0 \\ A^* & |D| & B \\ 0 & B^* & |C| \end{pmatrix}, \quad M_d = \begin{pmatrix} 0 & A' & 0 \\ A'^* & |D'| & B \\ 0 & B'^* & |C'| \end{pmatrix}$$

Without lack of generality we can make  $M_u$  to be real and positive. Then,  $M_u$  and  $M_d$  have ten real parameters  $A$ ,  $B$ ,  $C$ ,  $D$ ,  $|A'|$ ,  $|B'|$ ,  $|C'|$ ,  $|D'|$ , and the phases  $\phi_{A'}$  and  $\phi_{B'}$ . To reduce the number of parameters to nine we shall take  $|D'| = 0$ .

## CGS type mass matrices

$$M = \begin{pmatrix} 0 & a & d \\ a^* & 0 & b \\ d^* & b^* & c \end{pmatrix}$$

In the physical basis with the CGS-type mass matrix it is enough to consider the case where  $M_u$  is Fritzsch-type while  $M_d$  is CGS-type

$$M_u = \begin{pmatrix} 0 & a & 0 \\ a & 0 & b \\ 0 & b & c \end{pmatrix}, \quad M_d = \begin{pmatrix} 0 & a' & i|d'| \\ a' & 0 & b' \\ -i|d'| & b' & c' \end{pmatrix}$$

## Fits

1. For a given choice of the mass matrices we form a  $\chi^2$ -function which contains eleven summands.
2. The first five compare the experimental values the four best measured moduli ( $|V_{ud}|$ ,  $|V_{us}|$ ,  $|V_{cd}|$ ,  $|V_{cs}|$ ) and  $J(V)$  with their theoretical expressions.
3. The last six summands constrain the eigenvalues of quark mass matrices to the experimentally deduced quark masses at a specified energy scale.

# Quark masses at various energy scales.

Quark mass	Scale		
	2 GeV	$M_Z = 91,1876 \text{ GeV}$	PDG(06)
$m_u$	$2,2^{+0,8}_{-0,7}$	$1,28^{+0,50}_{-0,43}$	$2,25 \pm 0,75$
$m_d$	$5,0 \pm 2,0$	$2,91^{+1,24}_{-1,20}$	$5,0 \pm 2,0$
$m_s$	$95 \pm 25$	$55^{+16}_{-15}$	$95 \pm 25$
$m_c$	$1,07 \pm 0,12$	$0,624 \pm 0,083$	$1,25 \pm 0,09$
$m_b$	$5,04^{+0,16}_{-0,15}$	$2,89 \pm 0,09$	$4,20 \pm 0,07$
$m_t$	$318,9^{+13,1}_{-12,3}$	$172,5 \pm 3,0$	$174,2 \pm 3,3$

## Observed values of CKM parameters.

Observable	Exp. value
$ V_{ud} $	$0,97383 \pm 0,00024$
$ V_{us} $	$0,2272 \pm 0,0010$
$ V_{cd} $	$0,2271 \pm 0,0010$
$ V_{cs} $	$0,97296 \pm 0,00024$
$J(V)$	$(3,08 \pm 0,18) \times 10^{-5}$

## Fit for quark masses at $M_Z$ energy scale.

Type of mass matrix	Basis	Number of parameters	$\chi^2/(\text{dof})$
Fritzsch	Physical ( $\phi_{A'}$ and $\phi_{B'}$ free)	8	4.23/3 = 1.41
	Physical ( $\phi_{A'} = -\pi/2$ and $\phi_{B'} = 0$ )	6	4.84/5 = 0.97
Gupta-Rajpoot	Physical ( $\phi_{A'}$ and $\phi_{B'}$ free)	9	0.80/2 = 0.40
	Physical ( $\phi_{A'} = -\pi/2$ and $\phi_{B'} = 0$ )	7	2.62/4 = 0.66
$M_u$ Fritzsch-type and $M_d$ CGS-type	Physical	7	1.89/4 = 0.47

## Fit for quark at 2 GeV energy scale.

Type of mass matrix	Basis	Number of parameters	$\chi^2/(\text{dof})$
Fritzsch	Physical ( $\phi_{A'}$ and $\phi_{B'}$ free)	8	4.80/3 = 1.60
	Physical ( $\phi_{A'} = -\pi/2$ and $\phi_{B'} = 0$ )	6	5.49/5 = 1.10
Gupta-Rajpoot	Physical ( $\phi_{A'}$ and $\phi_{B'}$ free)	9	0.87/2 = 0.44
	Physical ( $\phi_{A'} = -\pi/2$ and $\phi_{B'} = 0$ )	7	2.76/4 = 0.69
$M_u$ Fritzsch-type and $M_d$ CGS-type	Physical	7	2.47/4 = 0.62

## Fit of PDG quark masses.

Type of mass matrix	Basis	Number of parameters	$\chi^2/(\text{dof})$
Fritzsch	Physical ( $\phi_{A'}$ and $\phi_{B'}$ free)	8	3.32/3 = 1.11
	Physical ( $\phi_{A'} = -\pi/2$ and $\phi_{B'} = 0$ )	6	4.27/5 = 0.85
Gupta-Rajpoot	Physical ( $\phi_{A'}$ and $\phi_{B'}$ free)	9	0.78/2 = 0.39
	Physical ( $\phi_{A'} = -\pi/2$ and $\phi_{B'} = 0$ )	7	3.77/4 = 0.94
$M_u$ Fritzsch-type and $M_d$ CGS-type	Physical	7	0.80/4 = 0.20

## Conclusion

1. The stability of this type of analysis with respect to evolution of the quark masses is important. As can be seen the results for  $\chi^2/\text{dof}$  for the different cases at 2 GeV scale are very similar to results at  $M_Z$  scale.
2. The CGS type of mass matrix with 7 parameters is favored over the F type or the choice of the G-R type.