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Energy calibration of data recorded with the surface detectors of the Pierre Auger Observatory

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Abstract: The energy of the primary particles of the air showers recorded using the water-Cherenkov detectors of the Pierre Auger Observatory is inferred from simultaneous measurement of the showers with the fluorescence telescopes. The signal on the ground at 1000 m from the shower axis obtained using the water-Cherenkov detectors is related directly to the calorimetric energy measured with the telescopes. The energy assignment is therefore independent of air-shower simulations except for the assumptions that must be made about the energy carried into the ground by neutrinos and muons. The correlation between the signal at ground and the calorimetric energy is used to derive a calibration curve. A detailed description of the method used to determine the energy scale is presented. The systematic uncertainties on the calibration procedure are discussed.

Introduction

The Pierre Auger Observatory [1] detect the air showers with the surface deterctor array composed by water-Cherenkov detectors with a 100% duty cycle [2]. The interpolated signal at a fixed optimal distance from the shower core, $S(1000)$ for the surface detector, is a good energy estimator in the sense that it is well correlated with the energy of the primary cosmic ray [3]. A subsample of the air showers is detected using simultaneously the fluorescence telescopes. They provides a nearly calorimetric energy measurement E_{FD}, because the fluorescence light is produced in proportion to energy dissipation by a shower in the atmosphere [4, 5]. This method can be used only when the sky is moonless and dark, and thus has about a 10% duty cycle [6]. For this subsample of air showers, called "hybrid events", it possible to relate the shower energy E_{FD} to the ground parameter $S(1000)$. The energy scale obtained studying this data sample is applyed at the full sample of shower detected by the array of the water-Cherenkov detectors.

Data Analysis

In this analysis hybrid events collected by the Pierre Auger Observatory between the 1st of December 2004 and the 31st of May 2008 are used. To ensure that the shower is sampled to make an $S(1000)$ measurement with the surface array, the rejection of accidental triggers and the core of the shower contained inside the array are requested. The selection criteria used is that all six nearest neighbours of the station with the highest signal must be active.

A subset of high-quality hybrid events are selected requiring that, only events with the reconstructed zenith angle less than 60° are selected [7], the geometry of an event must be determined from the times recorded at a fluorescence telescope, supplemented by the time at the water-Cherenkov detector with the highest signal and with the core of the shower within 750 m from the shower axis [8]. It is also required that a reduced χ^2 is less than 2.5 for the fit of the longitudinal profile by Gaisser-Hillas function [9] and that the depth of shower

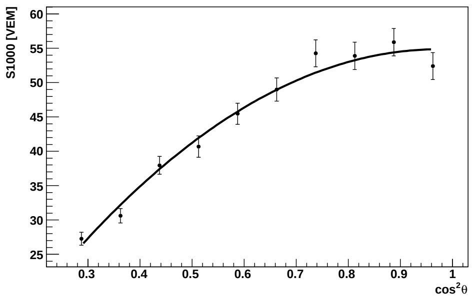

Figure 1: Derived attenuation curve, $CIC(\theta)$, fitted with a quadratic function.

maximum $Xmax$ be within the field of view of the telescopes. The fraction of the signal attributed to Cherenkov light must be less than 50%. The uncertainties on E_{FD} lower than 20% and on $Xmax$ lower than 40 g/cm^2 are also requested. The selection criteria include a measurement of the vertical aerosol optical depth profile (VAOD(h)) [10] using laser shots generated by the central laser facility (CLF) [11] and observed by the fluorescence telescopes in the same hour of each selected hybrid event. cite linear - ghChi2 and FD DB good periods cuts??

For a given energy the value of $S(1000)$ decreases with zenith angle, θ, due to attenuation of the shower particles and geometrical effects. Assuming an isotropic flux for the whole energy range considered, we extract the shape of the attenuation curve from the data [12]. The fitted attenuation curve, $CIC(\theta) = 1+a\ x+b\ x^2$, is a quadratic function of $x = \cos^2\theta - \cos^2 38^\circ$ as displayed in Figure 1 for a particular constant intensity cut, $I_0 = 330$ events, with $a = 0.91 \pm 0.05$ and $b = -1.28 \pm 0.23$. The average angle is $\langle\theta\rangle \simeq 38^\circ$ and we take this angle as reference and convert $S(1000)$ into S_{38° by $S_{38^\circ} \equiv S(1000)/CIC(\theta)$. It may be regarded as the signal $S(1000)$ the shower would have produced had it arrived at $\theta = 38^\circ$.

The reconstruction accuracy of the parameter $S(1000)$, $\sigma_{S(1000)}$ is composed by 3 contributions: a statistical uncertainty due to the finite size of the detector and the limited dynamic range of the signal detection, a systematic uncertainty due to the assumptions of the shape of the lateral distribution and finally due to the shower-to-shower fluctuations [13]. These and the uncertainty on to the attenuation curve parameters, are taken into account in inferring S_{38° and its uncertainty $\sigma_{S_{38^\circ}}$ and the relative uncertainty is about $\sigma_{S_{38^\circ}}/S_{38^\circ} = 16\%$ as shown in figure 2 and it is energy dependent.

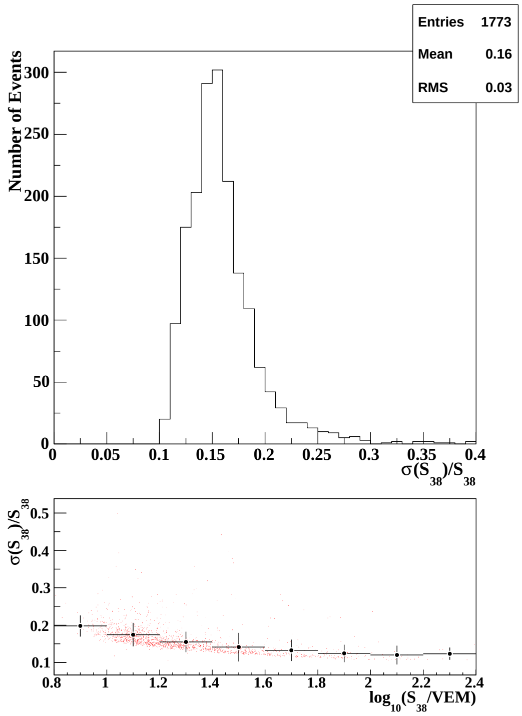

Figure 2: Upper panel: S_{38° resolution. Lower panel $\sigma_{S_{38^\circ}}/S_{38^\circ}$ vs $log_{10}(S_{38^\circ}/VEM)$ scatter plot with mean profile.

Not all the energy of a primary cosmic ray particle ends up in the electromagnetic part of an air shower detected by fluorescence telescopes. Neutrinos escape undetected and muons need long path lengths to release their energy. This is usually accounted for by multiplying the electromagnetic energy by a correction factor f_{inv} determined from shower simulations to obtain the total primary energy. Due to the energy dependence of the meson decay probabilities in the atmosphere, and thus the neutrino and muon production probabilities, the correction depends on the energy for different

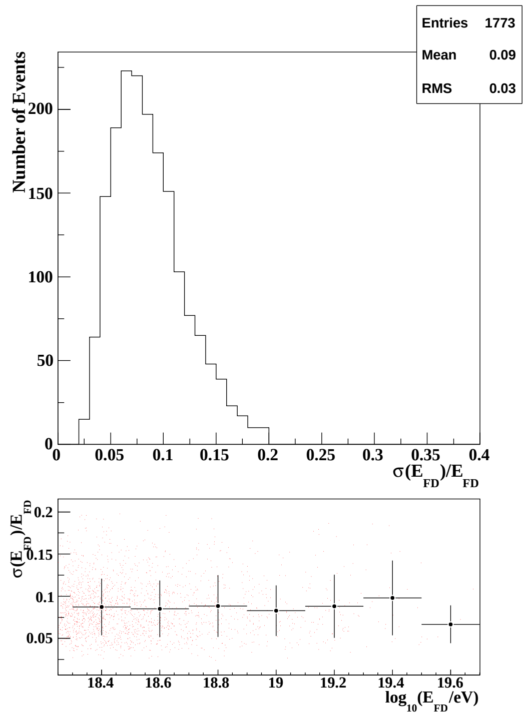

Figure 3: Upper panel: E_{FD} resolution. Lower panel $\sigma_{E_{FD}}/E_{FD}$ *vs* $log_{10}(E_{FD}/eV)$ scatter plot with mean profile.

hadronic interaction model, and is also subject to shower-to-shower fluctuations [14]. The so-called *invisible energy* correction is based on the average for proton and iron showers simulated with the QGSJet model and sums up to about 10% and its systematic uncertainty contributes 4% to the total systematic uncertainty in the energy obtained by the fluorescence telescopes.

The statistical uncertainties of the total energy (E_{FD}) measured by the fluorescence telescopes ($\sigma_{E_{FD}}$) is composed by the statistical uncertainty of the light flux (σ_{flux}), the uncertainty due to the core location and shower direction (σ_{geo}), the uncertainty on the invisible energy correction (σ_{inv}) and the uncertainty related to the measured VAOD profile (σ_{atm}). The total relative uncertainty is about $\sigma_{E_{FD}}/E_{FD} = 8\%$ as shown in figure 3 and do not depend strongly on the energy measured in this energy range. Check CIC values with

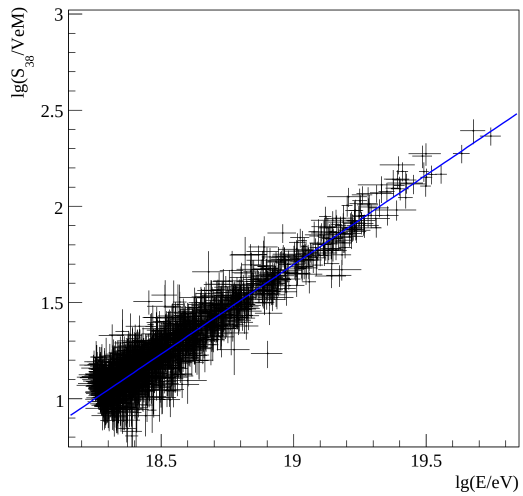

Figure 4: Correlation between $\lg(S_{38})$ and $\lg(E_{FD})$ for the 1773 hybrid events used in the fit. The line is the best fit

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Calibration Curve

The full efficienty of the surface detectors array is for $S_{38^\circ} = 15VEM$, this cut value (S_{38}^{cut}) is related to a cut value for the fluorescence telescopes (E_{FD}^{cut}). For any given pair of ($S_{38}^{cut}, E_{FD}^{cut}$), in the hypothesis of gaussian distributed measurements, we expect the data points to be distributed according to the standard error ellipse with center in ($S_{38}^{cut}, E_{FD}^{cut}$) Rejecting events below any line cutting the ellipse, a bias it is introduced. In fact, when averaging over the points of the ellipse above the cut, $(< S_{38} >, < E_{FD} >) \neq (S_{38}^{cut}, E_{FD}^{cut})$. A selection criteria it is necessary to select event for which $(< S_{38} >, < E_{FD} >) = (S_{38}^{cut}, E_{FD}^{cut})$ in the correlation between E_{FD} and S_{38°. This is obtained by selecting events which land in the 90% C.L. ellipse centered in ($S_{38}^{cut}, E_{FD}^{cut}$). With this selection is possible to take into account the resolution of both detectors in the cut values ($S_{38}^{cut}, E_{FD}^{cut}$). The 1773 hybrid selected events appear to be well described by a power-law: $E = a \cdot {S_{38}}^b$ as shown in figure 4.

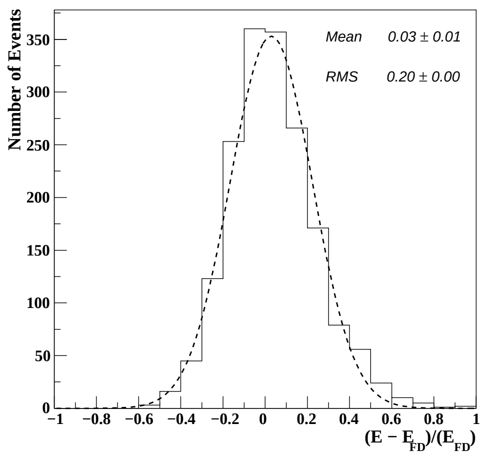

Figure 5: Fractional difference between the FD and SD energy for the 1773 selcted events.

The results of the fit are:
$a = (1.50 \pm 0.03(stat) \pm 0.12(syst)) \times 10^{17}\ eV$,
$b = 1.07 \pm 0.01(stat) \pm 0.04(syst)$,
with a reduced χ^2 of 1.03. S_{38} grows approximately linearly with energy. The root-mean-square deviation of the distribution is 20% as shown in figure 5, in good agreement with the quadratic sum of the S_{38° and E_{FD} statistical uncertainties of 18%. The calibration accuracy at the highest energies is limited by the number of events: the most energetic is about 6×10^{19} eV. The calibration at low energies extends below the range of interest.

Systematic Uncertainties

The systematics uncertainty on the curve parameters due to the calibration procedure are quoted changing the ellipse cut C.L. (from 68% up to 95%), appling the cut horizzontally and vertically respect to the ellipse and changing the fit function with a linear function in log. scale.

At this uncertainty the systematic uncertainty due to the fluorescence telescope energy meareaments must be considered.

The individual systematic uncertainties in determining E_{SD} coming from the FD sum up to 22%. The largest uncertainties are given by the absolute fluorescence yield (14%)[15], the absolute calibration of the FD (9.5%) and the reconstruction method (10%).

The uncertainty due to the dependence of the fluorescence spectrum on pressure (1%), humidity (5%) and temperature (5%) are take into account as well as the wavelength dependent response of the fluorescence telescopes, the aerosol phase function, invisible energy and others, which are well below 4%. The *invisible energy* correction introduce a systematic uncertainty contributes 4% at the total systematic of 22% [6].

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