



Flavourful Axion and the Peccei-Quinn Symmetry

arXiv:2102:03631

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May 11, 2021

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Outline

- 1 Axions and Flavour Symmetries
- 2 A Flavourful Axion Model
- 3 Numerical Results
- 4 Conclusions



Axions and the Peccei-Quinn Symmetry $U(1)_{PQ}$

- CP-violating term in QCD Lagrangian $\theta \frac{g_s^2}{32\pi^2} G \tilde{G}$,
 where $\tilde{G}^{a\mu\nu} = \frac{1}{2} \epsilon^{\mu\nu\rho\sigma} G_{\rho\sigma}^a$.
- $\bar{\theta} = \theta + \text{Arg}(\text{Det}(M^u M^d)) < 10^{-10}$ by neutron EDM¹.
- Unnaturally small?
- Introduce $U(1)_{PQ}$ and scalar field to SM².
- **Axion** appears as **pseudo-Goldstone boson of $U(1)_{PQ}$** ³, and
 $\mathcal{L}_{QCD+axion} \supset \frac{a}{f_a} \frac{g_s^2}{32\pi^2} G \tilde{G} + \bar{\theta} \frac{g_s^2}{32\pi^2} G \tilde{G}$.
- Quasi shift-symmetry of axion, $a \rightarrow a - f_a \bar{\theta}$, **solves strong CP problem**.

¹Abel et al. 2020.

²R. D. Peccei and Quinn 1977.

³Weinberg 1978; Wilczek 1978.

$U(1)_{PQ}$ as a Flavour Symmetry

- $U(1)_{PQ}$ can be a flavour symmetry.
- Identified with the Froggatt-Nielsen symmetry $U(1)_{FN}$.
- This gives rise to a *flavourful axion* or *axión sabroso*.
- FN symmetry can explain fermion mass hierarchy⁴.
- Flavon field σ couples to fermions like

$$\frac{C_{ij}^f}{\Lambda^n} \overline{F}_L^i H f_R^j \sigma^n$$
 or $\frac{C_{ij}^f}{\Lambda^n} \overline{F}_L^i \tilde{H} f_R^j \sigma^n$.
- After FN symmetry breaking $y_{ij}^f \varepsilon^n \overline{F}_L^i H f_R^j$ are the SM Yukawa terms, and $\varepsilon = \frac{v_\sigma}{\Lambda} \sim 0.2$, **hierarchy comes from powers of n** .

⁴Froggatt and Nielsen 1979.

Quark Sector

| Field/Symmetry | Q_{iL} | u_{iR} | d_{iR} | H_u | H_d | σ |
|-------------------------|------------------|-------------------|-------------------|-----------|----------|----------|
| $SU(2)_L \times U(1)_Y$ | (2, 1/6) | (1, 2/3) | (1, -1/3) | (2, -1/2) | (2, 1/2) | (1, 0) |
| $U(1)_{PQ}$ | (9/2, -5/2, 1/2) | (-9/2, 5/2, -1/2) | (-9/2, 5/2, -1/2) | 1 | 1 | 1 |

- Flavon σ is coupled to quark sector.
- Couplings up to dimension 5 are:

$$\mathcal{L} \supset \frac{C_{12}^u}{\Lambda} \bar{Q}_{1L} H_u u_{2R} \sigma + \frac{C_{21}^u}{\Lambda} \bar{Q}_{2L} H_u u_{1R} \sigma + \frac{C_{23}^u}{\Lambda} \bar{Q}_{2L} \tilde{H}_d u_{3R} \sigma^*$$

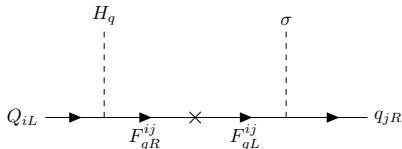
$$+ \frac{C_{32}^u}{\Lambda} \bar{Q}_{3L} \tilde{H}_d u_{2R} \sigma^* + y_{33}^u \bar{Q}_{3L} H_u u_{3R} + (u \rightarrow d).$$
- Next order is dimension 8.
- **NNI texture, with the (3, 3) entry generated at dimension 4**

$$M^{u/d} = \begin{pmatrix} 0 & \varepsilon v_{u/d} C_{12}^{u/d} & 0 \\ \varepsilon v_{u/d} C_{21}^{u/d} & 0 & \varepsilon v_{d/u} C_{23}^{u/d} \\ 0 & \varepsilon v_{d/u} C_{32}^{u/d} & y_{33}^{u/d} v_{u/d} \end{pmatrix}.$$

UV Completion of the Quark Sector

Type-I seesaw:

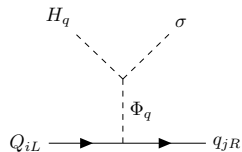
Heavy vector-like fermions F_q^{ij}



$$\mathcal{Y}_{ij}^u \overline{Q}_{iL} H_q F_{qR}^{ij} + \mathcal{M}_{ij}^q \overline{F_{qR}^{ij}} F_{qL}^{ij} + \mathcal{Y}'_{ij}{}^q \overline{F_{qL}^{ij}} \sigma q_{jR}$$

Type-II seesaw:

Heavy $SU(2)_L$ doublet scalars Φ_q



$$\mathcal{Y}_{ij}^q \overline{Q}_{iL} \Phi_q q_{jR} + \kappa_q H_q \Phi_q^\dagger \sigma$$

Can limit mass-generating terms!

Lepton Sector

| Field/Symmetry | L_{iL} | ℓ_{iR} | N_i | σ' |
|-------------------------|------------|-------------|-------------|-----------|
| $SU(2)_L \times U(1)_Y$ | (2, -1/2) | (1, -1) | (1, 0) | (1, 0) |
| $U(1)_{PQ}$ | (1, -3, 0) | (0, -2, -1) | (0, -2, -1) | 2 |

- Second flavon σ' is introduced.
- Neutrino masses generated in a type-I seesaw scenario, from $\mathcal{L} \supset y_1^\nu \bar{L}_{eL} H_u N_1 + y_2^\nu \bar{L}_{\mu L} \tilde{H}_d N_2 + y_3^\nu \bar{L}_{\tau} H_u N_3 + \frac{M_1}{2} \bar{N}_1^c N_1 + \frac{y_{12}^N}{2} \bar{N}_1^c N_2 \sigma' + \frac{y_{13}^N}{2} \bar{N}_1^c N_3 \sigma + \frac{y_{33}^N}{2} \bar{N}_3^c N_3 \sigma'$.

- The neutrino mass matrix gets an A_2 structure $\begin{pmatrix} 0 & a & 0 \\ a & b & c \\ 0 & c & d \end{pmatrix}$.
- Charged leptons mass matrix is diagonal at leading order.
- $0\nu\beta\beta$ at tree-level, and normal ordering.

Scalar Sector and the Axion

- The 4 or 6 scalar fields present lead to a complex scalar potential.
- Axion is the PQ Goldstone, orthogonal to the Z-boson Goldstone, i.e.

$$a = A_{PQ} - \left[\frac{\sum_i X_i Y_i v_i^2}{\sqrt{\sum_i Y_i^2 v_i^2} \sqrt{\sum_i X_i^2 v_i^2}} \right] A_Z.$$

- Axion mass and axion-photon coupling depend on f_a^5 , as
 $m_a \approx 5.70 \left(\frac{10^{12} \text{ GeV}}{f_a} \right) \mu\text{eV}$ and $g_{a\gamma} \approx \frac{\alpha}{2\pi f_a} \left(\frac{E}{N} - 1.92 \right)$.
- **Strong hierarchy between the scalars' vevs** leads to $f_a \approx \frac{v\sigma}{\sqrt{2N}}$.
- Here $N = 5$ and $E = \frac{28}{3}$.
- Compared to the $SU(5)$ GUT⁶, where $\frac{E}{N} = \frac{8}{3}$, $\frac{|g_{a\gamma}^{SU(5)}|}{|g_{a\gamma}^{Flaxion}|} \approx 14$.

⁵Grilli di Cortona et al. 2016.

⁶Georgi and Glashow 1974.

Quark Sector

| Parameter | Best fit |
|-----------------------------|----------|
| $A_u/(10^{-2} \text{ GeV})$ | 1.519 |
| $B_u/(10^{-2} \text{ GeV})$ | -5.368 |
| C_u/GeV | -3.004 |
| $D_u/(10^1 \text{ GeV})$ | 3.562 |
| $E_u/(10^2 \text{ GeV})$ | 1.679 |
| $A_d/(10^{-2} \text{ GeV})$ | -1.233 |
| $B_d/(10^{-2} \text{ GeV})$ | 1.228 |
| $C_d/(10^{-1} \text{ GeV})$ | -3.074 |
| $D_d/(10^{-1} \text{ GeV})$ | -4.782 |
| E_d/GeV | -2.793 |
| $\alpha_u/^\circ$ | 96.56 |
| $\beta_u/^\circ$ | 98.23 |

| Observable | Global-fit value | | Model best-fit |
|-----------------------------|------------------|-----------------------------|----------------|
| | Best-fit value | 1σ range | |
| $\theta_{12}^q/^\circ$ | 13.09 | 13.06 \rightarrow 13.12 | 12.988 |
| $\theta_{13}^q/^\circ$ | 0.207 | 0.202 \rightarrow 0.213 | 0.2000 |
| $\theta_{23}^q/^\circ$ | 2.32 | 2.29 \rightarrow 2.37 | 2.381 |
| $\delta^q/^\circ$ | 68.53 | 66.06 \rightarrow 71.10 | 68.720 |
| $m_u/(10^{-3} \text{ GeV})$ | 1.288 | 0.766 \rightarrow 1.550 | 1.2743 |
| $m_c/(10^{-1} \text{ GeV})$ | 6.268 | 6.076 \rightarrow 6.459 | 6.2592 |
| m_t/GeV | 171.68 | 170.17 \rightarrow 173.18 | 171.687 |
| $m_d/(10^{-3} \text{ GeV})$ | 2.751 | 2.577 \rightarrow 3.151 | 2.7330 |
| $m_s/(10^{-2} \text{ GeV})$ | 5.432 | 5.153 \rightarrow 5.728 | 5.4311 |
| m_b/GeV | 2.854 | 2.827 \rightarrow 2.880 | 2.8501 |
| χ_q^2 | | | 1.0901 |

Lepton Sector

| Parameter | Best fit |
|--------------------------|----------|
| $a/(10^{-3} \text{ eV})$ | 9.933 |
| $b/(10^{-2} \text{ eV})$ | 2.646 |
| $c/(10^{-2} \text{ eV})$ | 2.475 |
| $d/(10^{-2} \text{ eV})$ | 2.264 |
| $\phi_a/^\circ$ | 29.87 |
| $\phi_b/^\circ$ | 91.88 |
| $\phi_c/^\circ$ | 3.03 |
| $\phi_d/^\circ$ | -109.97 |

| Observable | Global-fit value | | Model best-fit |
|--|------------------|-------------------------------------|----------------|
| | Best-fit value | 1σ range | |
| $\theta_{12}^l/^\circ$ | 34.5 | 33.5 \rightarrow 35.7 | 34.85 |
| $\theta_{13}^l/^\circ$ | 8.45 | 8.31 \rightarrow 8.61 | 8.432 |
| $\theta_{23}^l/^\circ$ | 47.7 | 46.0 \rightarrow 48.9 | 48.11 |
| $\delta^l/^\circ$ | 218 | 191 \rightarrow 256 | 258.8 |
| $\alpha^l/^\circ$ | | | 65.27 |
| $\beta^l/^\circ$ | | | 265.08 |
| $\Delta m_{21}^2/(10^{-5} \text{ eV}^2)$ | 7.55 | 7.39 \rightarrow 7.75 | 7.571 |
| $\Delta m_{32}^2/(10^{-3} \text{ eV}^2)$ | 2.424 | 2.394 \rightarrow 2.454 | 2.4221 |
| $\sum m_\nu/(10^{-2} \text{ eV})$ | | | 6.453 |
| m_e/MeV | 0.4865763 | 0.4865735 \rightarrow 0.4865789 | - |
| m_μ/GeV | 0.10271897 | 0.10271866 \rightarrow 0.10271931 | - |
| m_τ/GeV | 1.74618 | 1.74602 \rightarrow 1.74633 | - |
| χ^2_j | | | 2.0053 |

Flavour Violating Decays with Axions

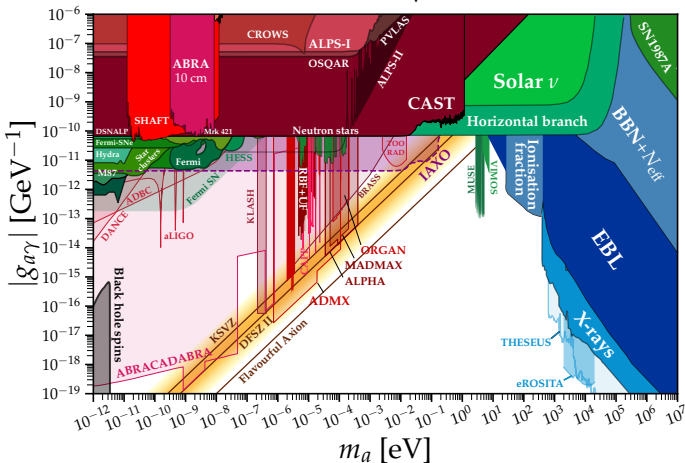
- Flavour violating Yukawa couplings can lead to $q_i \rightarrow q_j a$.
- Can be probed through meson decays.
- $\Gamma(K^+ \rightarrow \pi^+ a) \approx \frac{1.9 \times 10^{-9} \text{ GeV}^3}{v_\sigma^2}$.
- E949 Collaboration constrain $\frac{\Gamma(K^+ \rightarrow \pi^+ a)}{\Gamma_{\text{Total}}(K^+)} < 7.3 \times 10^{-117}$.
- Relationship $f_a \approx \sqrt{2} v_\sigma N$ implies $f_a \geq 7 \times 10^9 \text{ GeV}$.
- $m_a < 0.7 \times 10^{-3} \text{ eV}$ and $|g_{a\gamma}(\text{GeV}^{-1})| < 0.8 \times 10^{-14}$.
- $B^+ \rightarrow K^+ a$ (Belle-II⁸) provides $f_a \geq 6 \times (10^6 - 10^7) \text{ GeV}$.

⁷Adler et al. 2008.

⁸Calibbi et al. 2017.

Flavour Violating Decays with Axions

Flavourful axion and experiments⁹

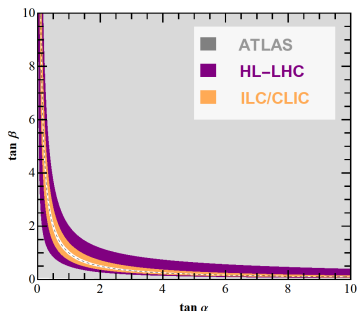


⁹O'Hare 2020.

Flavour Violating Higgs Couplings

- $v_{u/d} \ll v_{\sigma/\sigma'}$.
- Two Higgs bosons:
 $h \approx h_0^u \cos \alpha + h_0^d \sin \alpha$ and
 $H \approx -h_0^u \sin \alpha + h_0^d \cos \alpha$.
- $v_{SM} = v_u \sin \beta + v_d \cos \beta$.
- Flavour conservation if $\tan \alpha = \cot \beta$.
- $t \rightarrow hc$ sets bounds¹⁰.
- Effective Lagrangian and χ^2 fit:

$$\left| \frac{\cos \alpha}{\sin \beta} (1 - \tan \alpha \tan \beta) \right| \leq 17 \frac{\Gamma_{t \rightarrow hc}^{Exp}}{[GeV]}.$$



¹⁰ATLAS 2016; Vos et al. 2016.

Flavourful Axion as a Dark Matter Candidate

- Through the misalignment mechanism¹¹ cold axions can be produced, with $\Omega_a h^2 \approx 2 \times 10^4 \left(\frac{f_a}{10^{16} \text{GeV}} \right)^{7/6} \langle \theta_i^2 \rangle$.
- Relic density can be matched to $\Omega_{DM} h^2 \sim 0.12$.
- For $\theta_i \in (0, 2\pi)$ and $5 \times 10^{10} \text{ GeV} < f_a < 1 \times 10^{15} \text{ GeV}$.
- $N > 1$ implies formation of domain walls, which have to be destabilized or inflated away.

¹¹Fox, Pierce, and Thomas 2004.

Conclusions

- $U(1)_{PQ}$ can be identified with $U(1)_{FN}$.
- Mass and vev hierarchies.
- At leading order mass quark mass matrix gets NNI structure and neutrino mass matrix gets A_2 structure.
- Theoretical predictions match χ^2 fit.
- Flavourful axion arises after PQ symmetry breaking.
- Flavour violating decays with the axion set bounds on f_a , m_a , and $g_{a\gamma}$.
- Possibility of FCNCs being present.
- Flavourful axion can account for dark matter.
- Future possibilities: KSVZ-like model, gauge symmetries.

Thank you for listening!