

# The QCD Phase Diagram

Alfredo Raya

IFM-UMSNH

MexNICA Collaboration Winter Meeting 2020

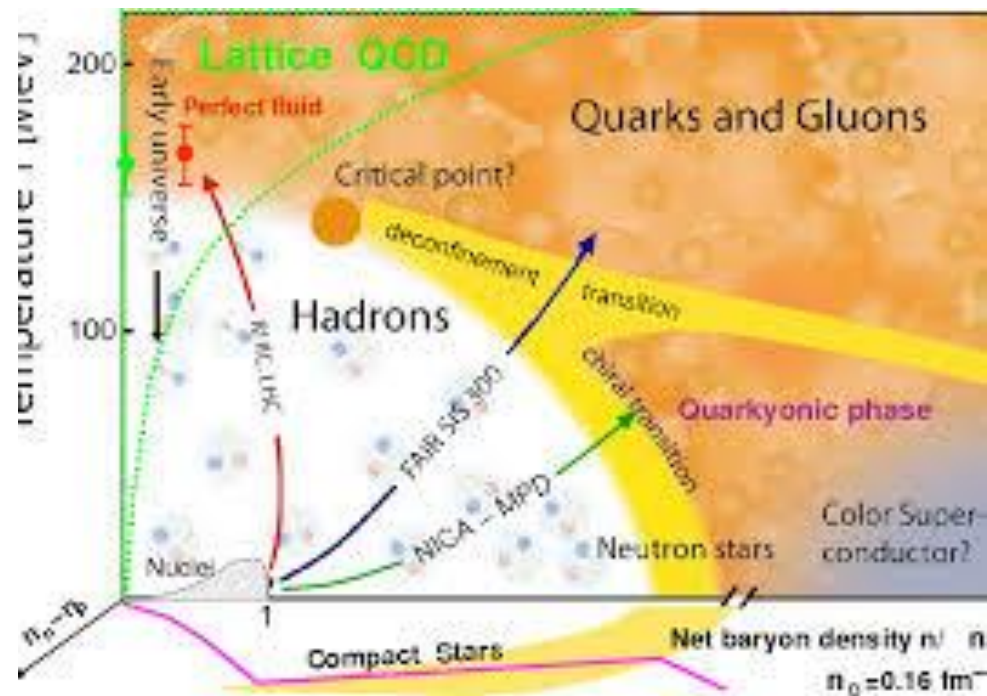
$\pi$

# Outlook

- › General features of the QCD Phase Diagram
- › Different approaches
  - Lattice QCD
  - Functional Methods
  - Effective Models
  - Sum Rules, etc
- › Physics interests of MPD Collaboration
- › Theory Meets Experiment
- › Final Remarks

$\pi$ 

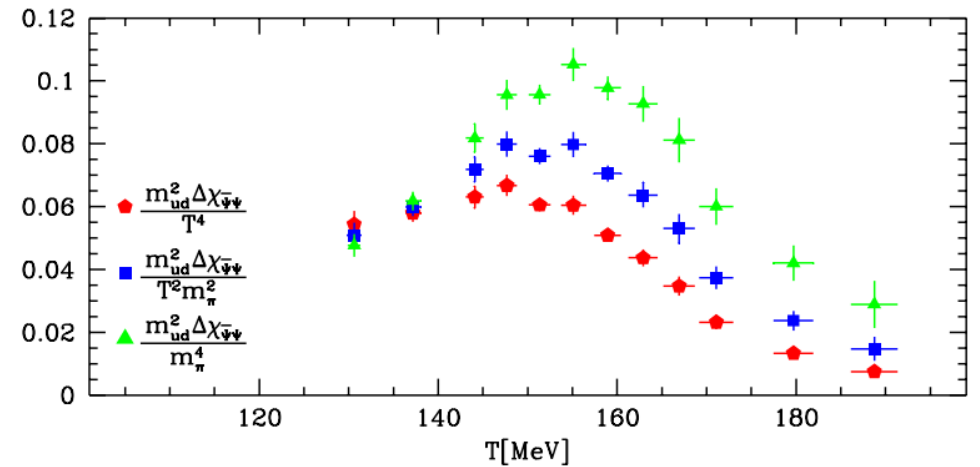
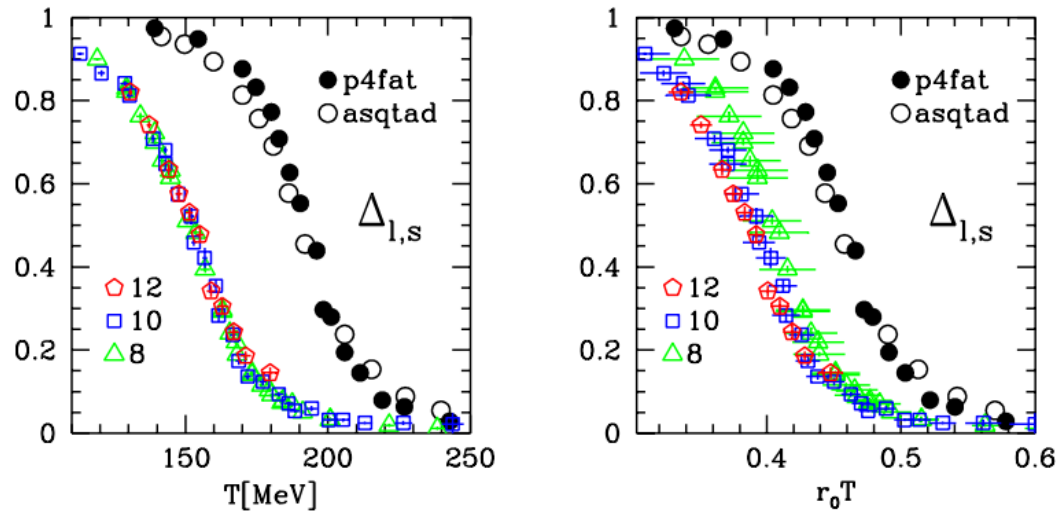
# The QCD Phase Diagram



$\pi$ 

# Different approaches

- › Lattice Simulations
  - Chiral Condensate
  - Chiral Susceptibilities



Adapted from J. High Energy Phys. 2009

# Different Approaches

- › Functional Methods:
- › Schwinger-Dyson equations
- › Functional Renormalization Group

The diagram shows the Schwinger-Dyson equation for a fermion propagator. On the left is a horizontal line with a single black dot, followed by a superscript  $-1$ . This is set equal to the sum of two terms. The first term is a horizontal line with no dots, followed by a superscript  $-1$ . The second term is a horizontal line with three black dots, with a loop of curly lines (representing a gluon) connecting the first and second dots from the left.

$$\langle \bar{\psi} \psi \rangle_f = -3 Z_2^f Z_m^f \sum_{\ell_k} \int \frac{d^3 k}{(2\pi)^3} \text{Tr}[S_f(k)]$$

## Different Approaches

- › Effective Models
- › NJL Model (Angelo's talk)
- › Sigma Model (Marcelo's talk)
- › Many, many others (Rest of the talks)

# Different Approaches

- › Sum Rules
  - Finite Energy Sum Rules

$$\frac{1}{\pi} \int_0^{s_0} ds K(s) \text{Im} \Pi(s)|_{\text{Had}} = \frac{-1}{2\pi i} \oint_{s_0} ds K(s) \Pi(s)|_{\text{QCD}},$$

$$\Pi(q) = \not{q} \Pi_1(q^2) + \Pi_2(q^2).$$

$$\begin{aligned} \Pi_1(s) = & -\frac{1}{64\pi^4} s^2 \ln(-s/\nu^2) - \frac{1}{32\pi^3} \langle \alpha_s G^2 \rangle \ln(-s/\nu^2) \\ & - \frac{2}{3} \frac{\langle \bar{q}q\bar{q}q \rangle}{s} + C_8 \frac{\langle \mathcal{O}_8 \rangle}{s^2} + C_{10} \frac{\langle \mathcal{O}_{10} \rangle}{s^3} + \dots, \quad (6) \end{aligned}$$

$$\begin{aligned} \Pi_2(s) = & \frac{1}{4\pi^2} \langle \bar{q}q \rangle s \ln(-s/\nu^2) - \frac{1}{12\pi} \frac{\langle \alpha_s G^2 \bar{q}q \rangle}{s} \\ & + C'_9 \frac{\langle \mathcal{O}_9 \rangle}{s^2} + C'_{11} \frac{\langle \mathcal{O}_{11} \rangle}{s^3} + \dots, \quad (7) \end{aligned}$$

# NICA

## PROGRAM:

- › Hadronic interactions and multiparticle production mechanisms at high barionic density abundances produced in Heavy Ion collisions
- › Even by evento basis
- › Global variables: Multiplicity and transverse energy
- › Quark confinement
- › Hadrons to partons
- › Critical End Point



# MPD Physics

- › Hadron abundances produced in Heavy Ion collisions
  - Different regions of the phase diagram correspond to different collision energies
  - Below (and above) a threshold, hadronic multiplicity will be different
  - Sizable change in the hadronic freeze-out pattern

# MPD Physics

- › Anisotropic Flow measurements
  - Strongly-coupled QGP
  - EoS, Sound velocity, viscosities
  - Fourier coefficients

# MPD Physics

## › Equation of State

- First order transition characterized by a softening in the EoS
- Dramatic drop of pressure
- Strongly connected with anisotropic flow

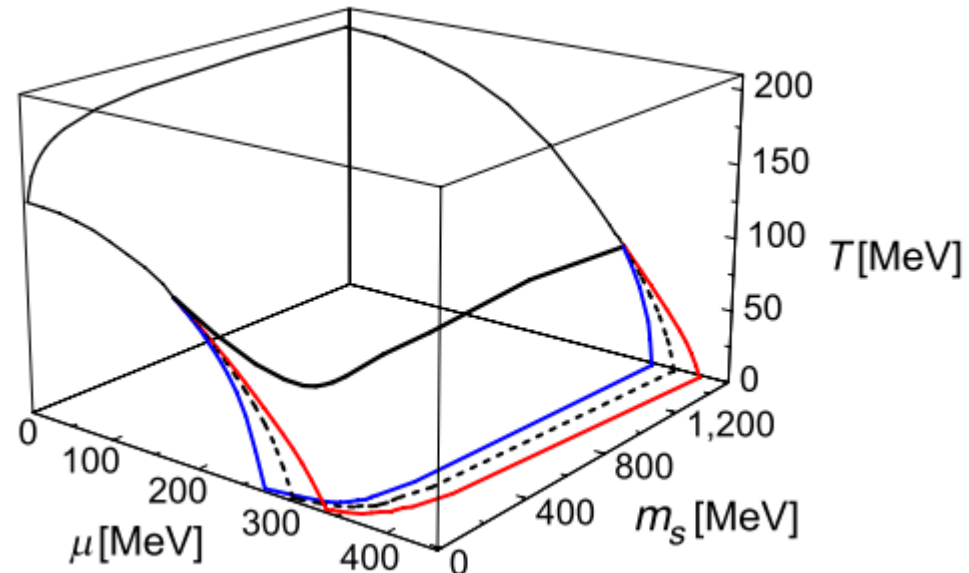
# MPD Physics

- › Femtoscopy: Pion interferometry
  - Spatio-temporal dynamics of the evolution
  - Two-pion measurements
  - Pion interferometry varying with the collision energy

# MPD Physics

## › Strangeness

- Deconfinement: Strangeness yield higher than that of hadron gas
- Novel effect on the phase diagram: Second critical point



# Theory Meets Experiment (EoS)

Thermodynamics

$$\Omega(T, \mu) = -\frac{T}{V} \log \mathcal{Z}(T, \mu)$$

$$p(T, \mu) = -(\Omega(T, \mu) - \Omega(0, 0)), \quad s(T, \mu) = -\frac{\partial \Omega(T, \mu)}{\partial T} \quad n(T, \mu) = -\frac{\partial \Omega(T, \mu)}{\partial \mu}$$

$$\varepsilon(T, \mu) = Ts(T, \mu) + \mu n(T, \mu) - p(T, \mu)$$

$$I(T, \mu) = \varepsilon(T, \mu) - 3p(T, \mu)$$

# Theory Meets Experiment (EoS)

$$\langle \bar{\psi}\psi \rangle(T, \mu; m) = \frac{\partial \Omega(T, \mu; m)}{\partial m}$$

$$\Omega(T, \mu; m_2) - \Omega(T, \mu; m_1) = \int_{m_1}^{m_2} dm' \langle \bar{\psi}\psi \rangle(T, \mu; m') .$$

Not useful

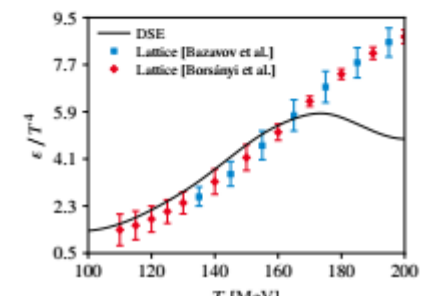
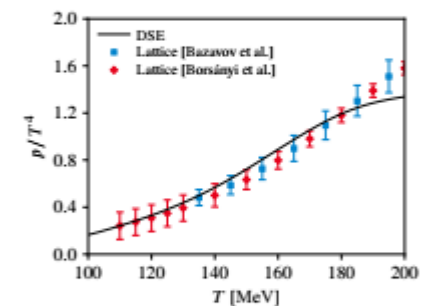
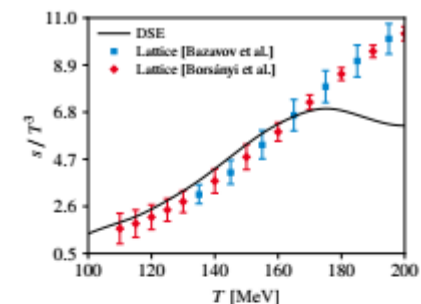
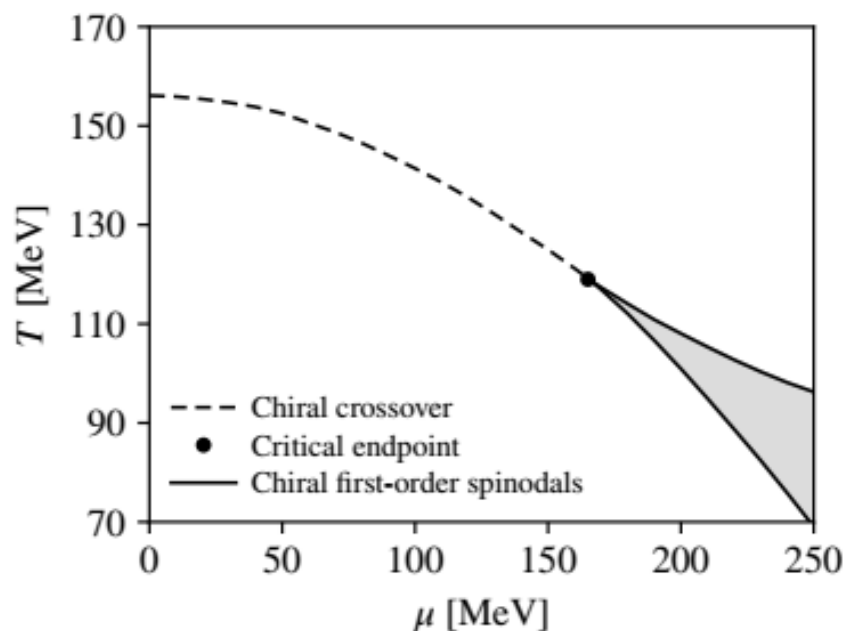
$$s(T, \mu; m_2) - s(T, \mu; m_1) = - \int_{m_1}^{m_2} dm' \frac{\partial \langle \bar{\psi}\psi \rangle}{\partial T}(T, \mu; m')$$

Divergence free

$$s(T, \mu; m) = s_{\text{YM}}(T) + \int_m^\infty dm' \frac{\partial \langle \bar{\psi}\psi \rangle}{\partial T}(T, \mu; m')$$

# Theory Meets Experiment (EoS)

$$p(T, \mu) = p(T_0, 0) + \int_{T_0}^T dT' s(T', 0) + \int_0^\mu d\mu' n(T, \mu').$$





## Final Remarks

- › It is possible to obtain thermodynamics information from theoretical quantities obtained in different frameworks
- › This information is constrained by the observed EoS
- › Physics com MPD would help to refine these approaches