



Ghost dynamics from Schwinger-Dyson equations

Mauricio N. Ferreira

University of Campinas, Brazil.

Hadron 2021
July 30th, 2021.

Based on:

A. C. Aguilar, C. O. Ambrósio, F. De Soto, M. N. F., B. M. Oliveira, J. Papavassiliou and J. Rodríguez-Quintero, [arXiv:2107.00768](https://arxiv.org/abs/2107.00768).

Motivation

- Ghost sector Green's functions, in particular the **ghost propagator** and **ghost-gluon vertex**, appear as ingredients in several Schwinger-Dyson equations (SDEs), *e.g.* the gluon gap equation

The image contains two Schwinger-Dyson equations for ghost sector Green's functions. The first equation is for the ghost propagator inverse, showing a tree-level propagator with a ghost loop (a1) and a ghost-gluon loop (a2). The second equation is for the ghost-gluon vertex inverse, showing a tree-level vertex and three loop diagrams (a3, a4, a5) involving ghost and gluon lines.

- Studied using a variety of nonperturbative approaches, including Lattice QCD, SDEs, Functional Renormalization Group.

I. L. Bogolubsky, E. M. Ilgenfritz, M. Muller-Preussker and A. Sternbeck, *PoS LATTICE2007*, 290 (2007).
 P. Boucaud, J. P. Leroy, A. L. Yaouanc, J. Micheli, O. Pene and J. Rodriguez-Quintero, *JHEP* **06**, 012 (2008).
 M. Q. Huber and L. von Smekal, *JHEP* **04**, 149 (2013).
 A. C. Aguilar, D. Ibáñez and J. Papavassiliou, *Phys. Rev. D* **87**, no.11, 114020 (2013).
 A. K. Cyrol, L. Fister, M. Mitter, J. M. Pawłowski and N. Strodthoff, *Phys. Rev. D* **94**, no.5, 054005 (2016).

- ★ Ghost sector SDEs are the simplest: Clean environment for testing agreement with lattice and probe the effect of other Green's functions.

This work: Coupled dynamics of the ghost propagator and ghost-gluon vertex, consistency between SDE and lattice and impact of the three-gluon vertex.

Coupled SDEs

We solve the system of equations for the **ghost propagator** and **the ghost-gluon vertex**

$$\left(\text{---} \circ \text{---} \right)^{-1} = \left(\text{---} \xrightarrow{p} \text{---} \right)^{-1} + \dots$$

$$\text{Diagram} = \text{Diagram} + (d_1) + (d_2) + (d_3)$$

(d1) (d2) (d3)

Coupled SDEs

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$$\left(\text{---} \circ \text{---} \right)^{-1} = \left(\text{---} \right)^{-1} + \dots$$

$$\text{Vertex} = \text{Tree} + (d_1) + (d_2) + (d_3)$$

To close the system we need to provide

- Gluon propagator [all loop diagrams];

Coupled SDEs

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To close the system we need to provide

- Gluon propagator [all loop diagrams];
- **three-gluon vertex** [diagram (d₁)];

Coupled SDEs

We solve the system of equations for the **ghost propagator** and the **ghost-gluon vertex**

$$\left(\text{---} \circ \text{---} \right)_p^{-1} = \left(\text{---} \right)_p^{-1} + \text{---} \circ \text{---} \circ \text{---} \circ \text{---} \circ \text{---}$$

(d1) (d2) (d3)

$\approx 2\%$

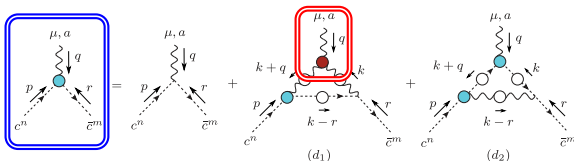
Huber, M. Q. *Eur. Phys. J. C* **77**, 733 (2017).

To close the system we need to provide

- Gluon propagator [all loop diagrams];
- **three-gluon vertex** [diagram (d1)];
- **ghost-ghost-gluon-gluon vertex** [diagram (d3)].

This we will omit \Rightarrow “**one-loop dressed**” truncation.

- Main remaining uncertainty is in the **three-gluon** vertex,

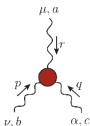


Sensitive to the three-gluon vertex.

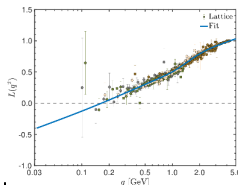
- Because this SDE has little contamination of more higher order functions

Clean probe for the three-gluon vertex.

- In the soft gluon limit, the **three-gluon vertex** is completely determined by the **lattice**:



Soft gluon

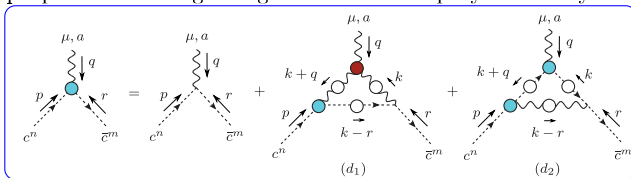


- Provides a nontrivial consistency check!

A. C. Aguilar, F. De Soto, M. N. F., J. Papavassiliou and J. Rodríguez-Quintero, *Phys. Lett. B* **818**, 136352 (2021).

Coupled SDEs: ghost-gluon vertex equation

Special properties of the ghost-gluon vertex simplify the analysis further:



- The most general structure of the ghost-gluon vertex is

$$\Gamma_\mu(r, p, q) = B_1(r, p, q) r_\mu + B_2(r, p, q) q_\mu.$$

- At tree level, $B_1^{(0)} = 1$ and $B_2^{(0)} = 0$.
- Now, in Landau gauge, by virtue of

$$\Delta_{\mu\nu}(q) = -iP_{\mu\nu}(q)\Delta(q^2), \quad P_{\mu\nu}(q) = g_{\mu\nu} - q_\mu q_\nu / q^2,$$

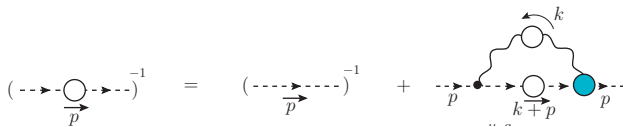
the form factor B_2 decouples from B_1 ,

$$B_1(r, p, q) = Z_1 + \int_\ell \mathcal{K}_1[B_1] + \int_\ell \mathcal{K}_2[B_1].$$

- Also, in Landau gauge this vertex is **finite**: $Z_1 = 1$ (Taylor scheme).
 \Rightarrow Further simplifies analysis.

Coupled SDEs: ghost propagator equation

We define the ghost dressing function, $F(p^2)$, as usual $D(p^2) = F(p^2)/p^2$.
The SDE for the ghost propagator also simplifies in Landau gauge,


$$\left(\text{---} \circ \text{---} \right)_p^{-1} = \left(\text{---} \text{---} \right)_p^{-1} + \text{---} \circ \text{---} \circ \text{---} \text{---} \text{---} \text{---}$$

$$F^{-1}(p^2) = Z_c + \Sigma(p^2),$$

with

$$\Sigma(p^2) = ig^2 C_A Z_1 \int_k \left[1 - \frac{(k \cdot p)^2}{k^2 p^2} \right] B_1(-p, k+p, -k) \Delta(k) D(k+p).$$

- Only depends on the form factor B_1 (Landau gauge).
- Renormalization in the Taylor scheme, where $F(\mu^2) = 1$ and $Z_1 = 1$ entails,

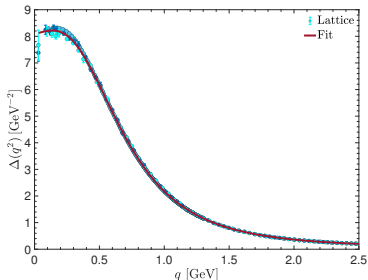
$$F^{-1}(p^2) = 1 + \Sigma(p^2) - \Sigma(\mu^2).$$

Inputs: gluon propagator

For the gluon propagator, we can use a fit to lattice data for $\Delta(q^2)$.

P. Boucaud, F. De Soto, K. Raya, J. Rodríguez-Quintero and S. Zafeiropoulos, *Phys. Rev. D* **98**, no.11, 114515 (2018).

I. L. Bogolubsky, E. M. Ilgenfritz, M. Müller-Preussker and A. Sternbeck, *PoS LATTICE2007*, 290 (2007).



- Most notable feature: IR saturation.
- Associated with the dynamical generation of a gluon mass scale.

A. C. Aguilar and J. Papavassiliou, *JHEP* **12**, 012 (2006).

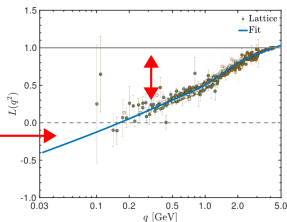
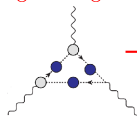
A. C. Aguilar, D. Binosi, C. T. Figueiredo and J. Papavassiliou, *Phys. Rev. D* **94**, no.4, 045002 (2016).

Inputs: Three-gluon vertex

The most complicated ingredient is the general kinematics **three-gluon vertex** appearing in the **ghost-gluon SDE**.

- Complicated tensor structure with 14 independent tensors.
- Each accompanied with a scalar form factor that is a function of three variables.
- Rich nonperturbative structure:
 - ◆ **IR suppression.**
 - ◆ **Zero-crossing**
 - ◆ IR divergences due to massless ghost loops.

log divergence



A. C. Aguilar, F. De Soto, M. N. F., J. Papavassiliou and J. Rodríguez-Quintero, *Phys. Lett. B* **818**, 136352 (2021).

A. Athenodorou, D. Binosi, Ph. Boucaud, et al. *Phys.Lett.B* 761 (2016) 444-449.

A. C. Aguilar, D. Binosi, D. Ibañez and J. Papavassiliou, *Phys. Rev. D* **89**, no.8, 085008 (2014).

In this work we resort to the Gauge Technique construction for the **three-gluon vertex**.

Gauge Technique three-gluon vertex

The three-gluon vertex satisfies a Slavnov-Taylor identity

$$q \cdot \Gamma^{\text{STI}} = F[\Delta^{-1}H - \Delta^{-1}H], \quad q \cdot \Gamma^{\text{Tr}} = 0, \quad \Gamma = \Gamma^{\text{STI}} + \Gamma^{\text{Tr}},$$

with H the ghost-gluon kernel [$q^\mu H_{\nu\mu}(q, p, r) = \Gamma_\mu(q, p, r)$].

A. C. Aguilar, M.N.F., C. T. Figueiredo and J. Papavassiliou, *Phys.Rev.D* 99 (2019) 9, 094010.

- The **transverse part** cannot be determined by the STI.
- ★ **But we can determine Γ^{STI} .**

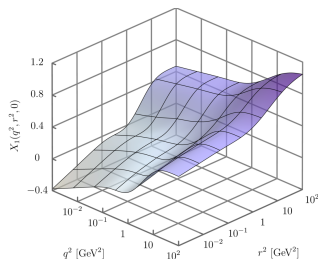
For simplicity, we retain only the “classical” tensor structure,

$$\Gamma_{\alpha\mu\nu}^{\text{STI}}(q, r, p) \approx (q - r)_\nu g_{\alpha\mu} X_1(q, r, p) + (r - p)_\alpha g_{\mu\nu} X_1(r, p, q) + (p - q)_\mu g_{\nu\alpha} X_1(p, q, r),$$

which reduces to the tree-level vertex with $X_1^{(0)} = 1$.

The form factor $X_1(q, r, p)$ captures the main features of the **three-gluon vertex**:

- IR suppression;
- logarithmic IR divergences;
- perturbative anomalous dimension.



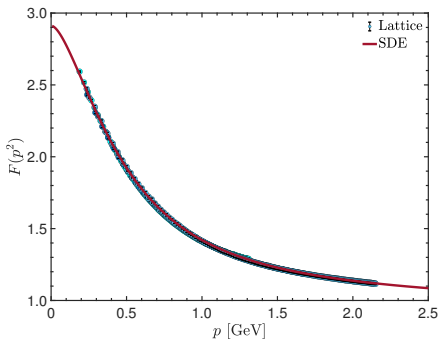
Results: Ghost dressing function

We first compare the resulting ghost dressing function to lattice data of P. Boucaud, F. De Soto, K. Raya, J. Rodriguez-Quintero and S. Zafeiropoulos, Phys. Rev. D **98**, no.11, 114515 (2018).

- Value of $\alpha_s = 0.244$ adjusted for best agreement.
- ★ **Excellent agreement!**
- ★ Careful treatment of the continuum limit of the lattice is key to this agreement.
- In previous works, with lattice not cured from discretization the agreement was not so perfect, *e.g.*

A. C. Aguilar, M. N. F., C. T. Figueiredo and J. Papavassiliou, Phys. Rev. D **99**, no.3, 034026 (2019).

A. C. Aguilar, D. Ibáñez and J. Papavassiliou, Phys. Rev. D **87**, no.11, 114020 (2013)

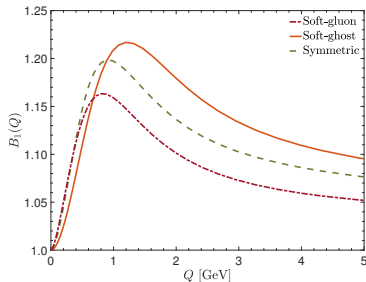
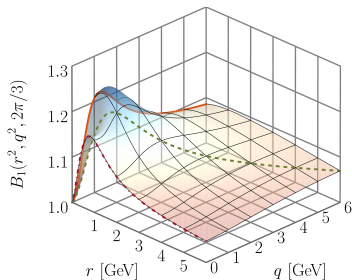


Results: ghost-gluon vertex

We can parametrize $B_1(r, p, q)$ in terms of the gluon and antighost momenta, q and r , respectively, and the angle θ between them.

A representative case: $\theta = 2\pi/3$

Some special limits:



- For all kinematics we find moderate deviations only from the tree level value, $B_1^{(0)} = 1$.
- Qualitative agreement with previous works.

Iigenfritz, E.-M., Muller-Preussker, M., Sternbeck, et al. *Braz.J. Phys.* **37**, 193–200 (2007).

Cucchieri, A., Maas, A. and Mendes, T. *Phys. Rev. D* **77**, 094510 (2008).

A. C. Aguilar, D. Ibáñez and J. Papavassiliou, *Phys. Rev. D* **87**, no.11, 114020 (2013).

Huber, M. Q. and von Smekal, L. *JHEP* **04**, 149 (2013).

Cyrol, A. K., Fister, L., Mitter, M., et al. *Phys. Rev. D* **94**, 054005 (2016).

Boucaud, P., Dudal, D., Leroy, J., et al. *JHEP* **12**, 018 (2011).

Barrios, N., Pelaez, M., Reinosa, U., et al. *Phys.Rev.D* **102**, 11401 (2020).

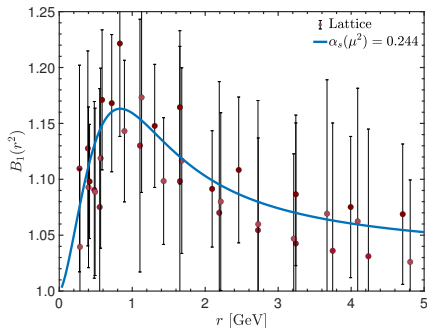
Results: ghost-gluon vertex

In the soft gluon limit there is lattice data available for the ghost-gluon vertex.

A. Sternbeck, (2006), arXiv:hep-lat/0609016 [hep-lat].

E.-M. Ilgenfritz, et al., Braz. J. Phys. 37, 193 (2007).

Again in agreement with the lattice:



Soft gluon limit: A nontrivial consistency check

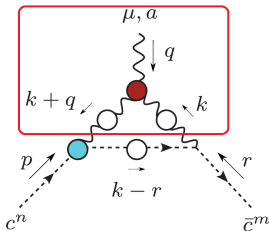
In the soft gluon limit, $q = 0$, the **three-gluon vertex** contribution to the **ghost-gluon SDE** simplifies dramatically.

It can be shown that

$$P_{\mu}^{\mu'}(r)P_{\nu}^{\nu'}(-r)\Gamma_{\mu'\nu'\alpha}(r, -r, 0) = L(r^2)\lambda_{\mu'\nu'\alpha}(r), \quad \lambda_{\mu'\nu'\alpha}(r) := 2r_{\alpha}P_{\mu\nu}(r),$$

recalling that $P_{\mu\nu}(r) = g_{\mu\nu} - r_{\mu}r_{\nu}/r^2$ is the transverse projector.

- **Only one tensor structure.**
- In the **ghost-gluon SDE**, for $q = 0$, the above is all we need:



Since in Landau gauge

$$\Delta_{\mu\nu}(q) = -iP_{\mu\nu}(q)\Delta(q^2),$$

the combination to the left becomes

$$-\Delta^2(k^2)L(k^2)\lambda_{\alpha\beta\mu}(k).$$

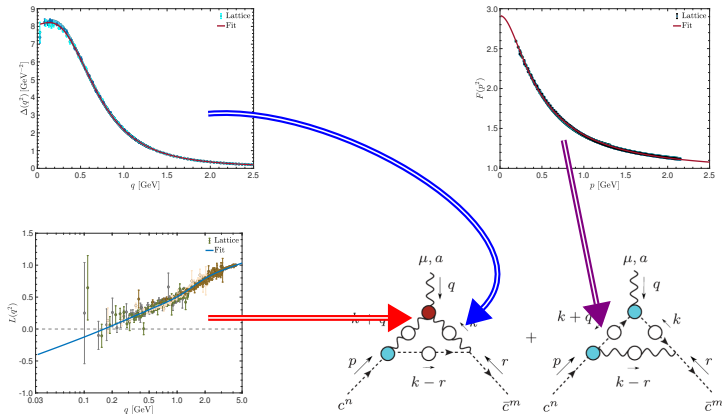
★ **Projection $L(q^2)$ accurately determined on the lattice.**

In this limit all external ingredients are known from lattice simulations:

A. C. Aguilar, F. De Soto, M. N. F., J. Papavassiliou and J. Rodríguez-Quintero, Phys. Lett. B **818**, 136352 (2021).

A. Athenodorou, D. Binosi, Ph. Boucaud, et al. Phys.Lett.B 761 (2016) 444-449.

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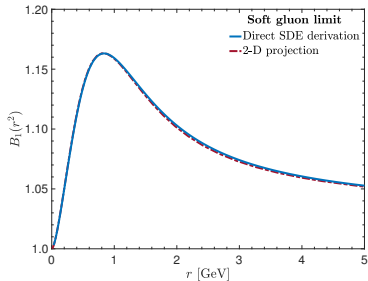
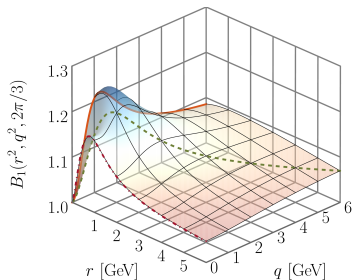
- All from the same lattice setups;
- and renormalized at the same point, $\mu = 4.3$ GeV.

Comparison to general kinematics

Comparing the soft gluon limit of the coupled system with the result of the soft-gluon SDE with lattice inputs only provides a consistency check:

Soft gluon SDE:

Coupled system:

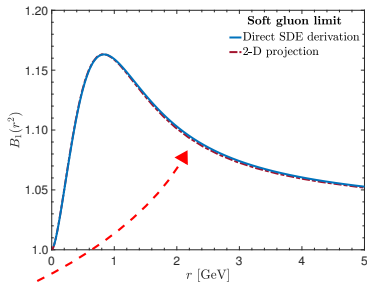
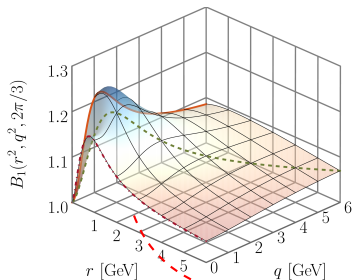


Comparison to general kinematics

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Soft gluon SDE:

Coupled system:



soft gluon slice

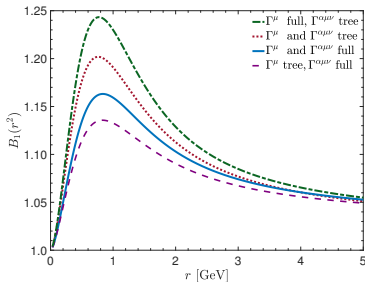
★ Results agree at the 2% level in quantum correction ($B_1 - 1$).

Impact of the vertex dressings

We now consider the effect of the the vertex dressings in the result for B_1 .

Results exhibit the expected hierarchy:

- Significant **suppression** due to the **three-gluon vertex**;
- and **enhancement** due to the **ghost-gluon vertex** itself.
- **Effects are most prominent** around ~ 1 GeV.

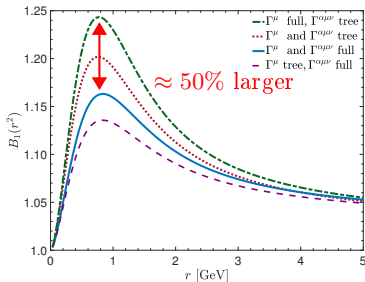


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- ★ Most importantly, neglecting the **three-gluon vertex** yields 50% too big quantum correction for B_1 .

Conclusions

- We solved the coupled SDEs for the ghost propagator and general kinematics ghost-gluon vertex with most up-to-date lattice and gauge technique inputs.
- Results exhibit **excellent agreement with lattice** data, after careful treatment of discretization artifacts.
- In the soft gluon limit, a consistency check is possible using lattice inputs for all external ingredients.
 - ★ Consistency of **Lattice**, **SDE** and **Gauge Technique**.
- The characteristic **IR suppression of the three-gluon vertex** has an important quantitative effect on the ghost-gluon vertex.