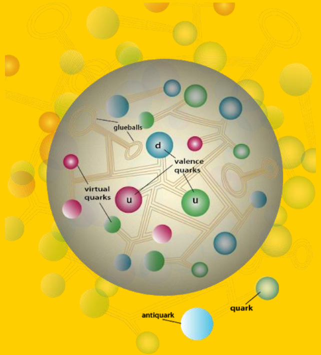
Double parton scattering via photon-proton interactions: a new light on the transverse proton structure



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in collaboration with

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Marco Traini
Sergio Scopetta
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STRONG 2020

INFN Istituto Nazionale di Fisica Nucleare

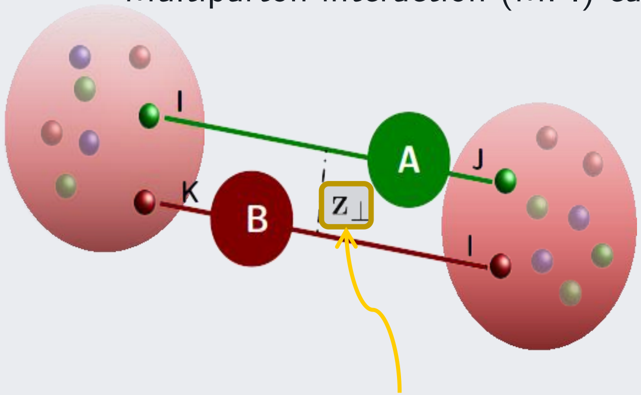
Roadmap

1. INTRODUCTION
2. 3D PROTON STRUCTURE VIA p-p DPS
3. INTERPRETATION OF LHC DATA
4. DPS VIA PHOTON-PROTON INTERACTIONS
5. PREDICTIONS AND OPPORTUNITIES
6. CONCLUSIONS

1 Double Parton Scattering @LHC

*See T. Kasemets's TALK

Multiparton interaction (MPI) can contribute to the, pp and pA, cross section @ the LHC:



Transverse distance between partons

The cross section for a double parton scattering (DPS) event can be written in the following way:

N. Paver, D. Treleani, Nuovo Cimento 70A, 215 (1982)

$$d\sigma \propto \int d^2 z_\perp \overbrace{F_{ik}(x_1, x_2, \vec{z}_\perp; \mu_A, \mu_B) \cdot F_{jl}(\underbrace{x_3, x_4}_{}, \vec{z}_\perp; \underbrace{\mu_A, \mu_B}_{}) }^{\text{double PDF (dPDF)}}$$

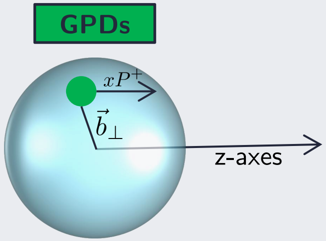
Momentum scales

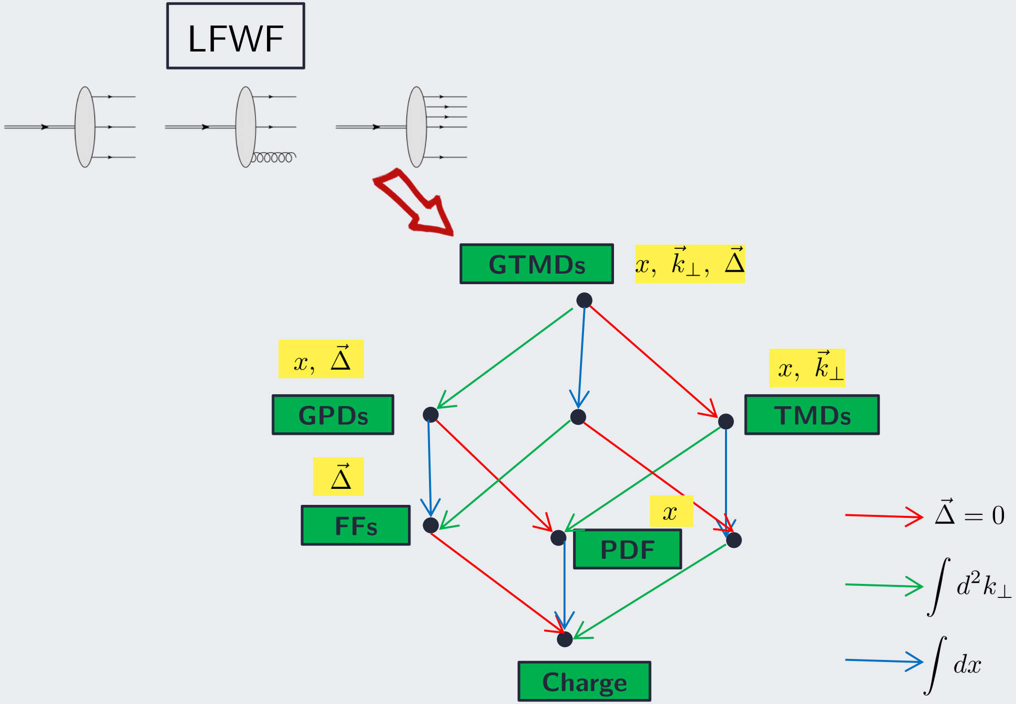
Momentum fractions carried by the parton inside the proton

DPS processes are important for fundamental studies, e.g. the background for the research of new physics and to grasp information on the **3D PARTONIC STRUCTURE OF THE PROTON**

1 Multidimensional Pictures of Hadron

1-body





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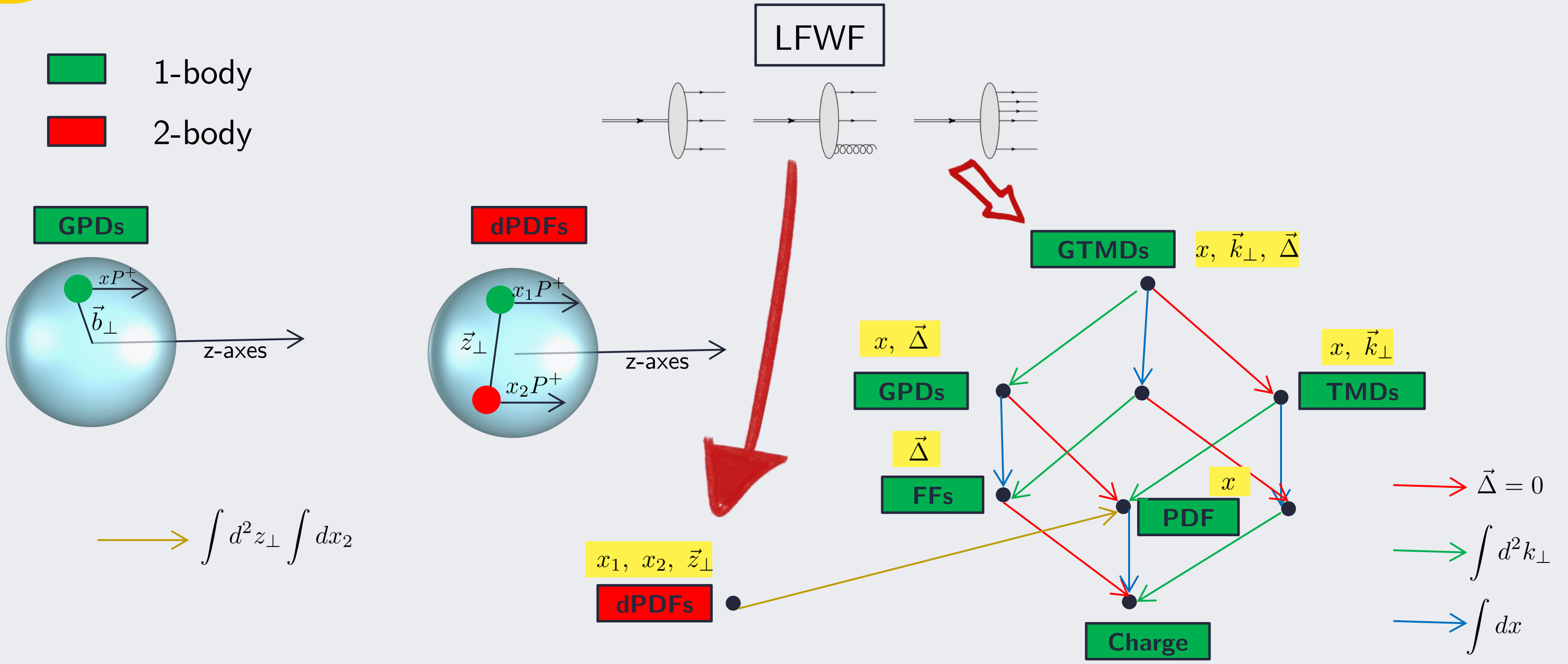

1 Double PDFs of the proton

$F_{ik}(x_1, x_2, \vec{z}_\perp)$ is unknown. However @LHC kinematics (small x and many partons produced)

$$F_{ik}(x_1, x_2, \vec{z}_\perp) \sim \underbrace{g(x_1, x_2)\tilde{T}(\vec{z}_\perp)}^{\text{1st uncorrelated scenario}}$$

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MODELS BASED ON SUM RULES

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→ PDF(x_1)*PDF(x_2)

→ [PDF(x_1)*PDF(x_2)] ⊗ pQCD EVOLUTION

2nd uncorrelated scenario

J. R. Gaunt and W. J. Stirling, JHEP 03, 005 (2010)

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Constituent quark models used to grasp basic **NON PERTURBATIVE** features

M.R., S. Scopetta et *al*, PRD 87 (2013) 114021
M.R., S, Scopetta et *al*, JHEP 12 (2014) 028

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PROBABILITY DISTRIBUTION
OF FINDING TWO PARTONS WITH
GIVEN TRANSVERSE DISTANCE

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M.R., S. Scopetta et *al*, PRD 87 (2013) 114021
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UNKNOWN: ONLY MODELS

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MODELS BASED ON SUM RULES

PDF(x_1)*PDF(x_2)

[PDF(x_1)*PDF(x_2)] $\otimes$ pQCD EVOLUTION

2nd uncorrelated scenario

PERTURBATIVE CORRELATIONS

Constituent quark models
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non perturbative feature

M.R., S. Scopetta et *al*, PRD 87 (2013) 114021
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1 Double PDFs of the proton

Some questions arose:

1) HOW CAN WE BE SURE OF THE ACCURACY OF SUCH APPROXIMATION?
dPDFs are non perturbative in QCD ➡ DPCs cannot be accessed from QCD

2) WHICH INFORMATION ON THE PROTON STRUCTURE COULD BE ACCESSED FROM DPS?

PERTURBATIVE CORRELATIONS

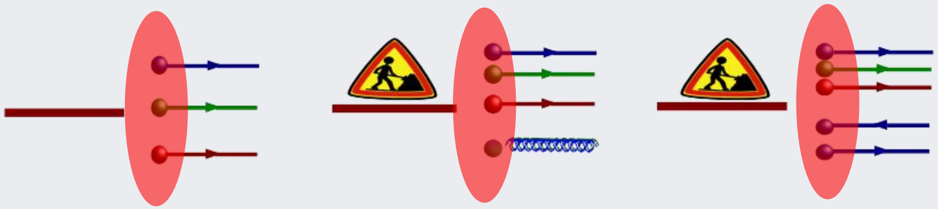
2 Double PDFs within the Light-Front

Extending the procedure developed in **S. Boffi, B. Pasquini and M. Traini, Nucl. Phys. B 649, 243 (2003)** for GPDs, we obtained the following expression of the dPDF in momentum space, often called ${}_2$GPDs:

$$F_{ij}(x_1, x_2, k_\perp) = 3(\sqrt{3})^3 \int \prod_{i=1}^{3} d\vec{k}_i \delta\left(\sum_{i=1}^{3} \vec{k}_i\right) \Phi^*(\{\vec{k}_i\}, k_\perp)\underbrace{\Phi(\{\vec{k}_i\}, -k_\perp)}_{\text{LF wave-function}}$$

$$\times \quad \delta\left(x_1 - \frac{k_1^+}{P_+}\right)\delta\left(x_2 - \frac{k_2^+}{P_+}\right)$$

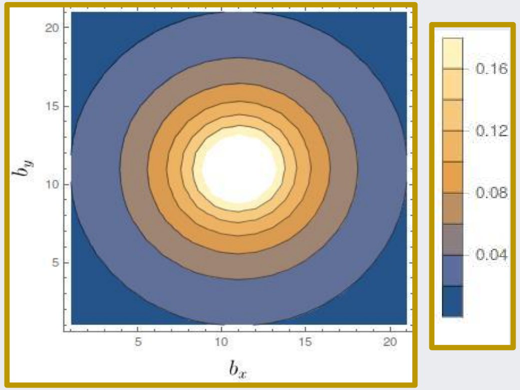
Conjugate to $\boxed{z_\perp}$



$$\Phi(\{\vec{k}_i\}, \pm k_\perp) = \Phi\left(\vec{k}_1 \pm \frac{\vec{k}_\perp}{2}, \vec{k}_2 \mp \frac{\vec{k}_\perp}{2}, \vec{k}_3\right)$$

2 Information from Quark Models

$z_\perp = b_\perp$

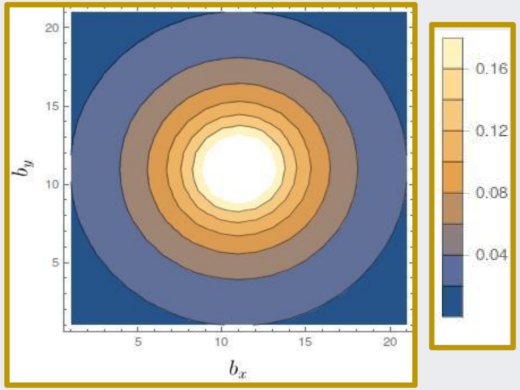


M.R. and F. A. Ceccopieri, JHEP 09 (2019) 097

1) e.g. the distance distribution of **two gluons** in the proton

$$\langle z_\perp^2 \rangle^{ij}_{x_1,x_2} = \frac{\int d^2 z_\perp \; z_\perp^2 F_{ij}(x_1, x_2, z_\perp)}{\int d^2 z_\perp \; F_{ij}(x_1, x_2, z_\perp)}$$

2 Information from Quark Models



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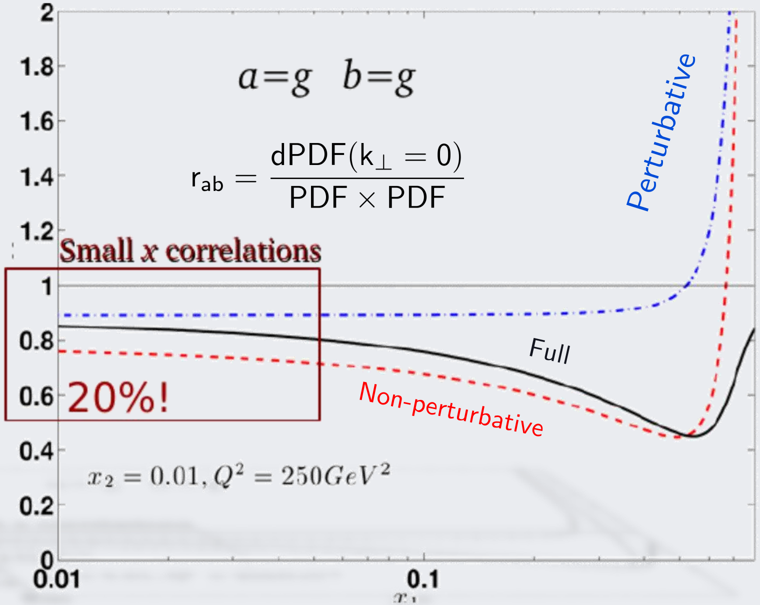

2) Correlations are important

M.R., S. Scopetta et al, JHEP 10 (2016) 063

M.R. and F. A. Ceccopieri PRD 95 (2017) 034040

3 Data and Effective Cross Section

A tool for the comprehension of the role of DPS in hadron-hadron collisions is the so called "effective X-section".

$$\underbrace{\sigma_{\text{eff}}^{pp} = \frac{m}{2} \frac{\sigma_A^{pp} \boxed{\sigma_B^{pp}}}{\sigma_{DPS}^{pp}}}_{\text{POCKET FORMULA}}$$

Differential cross section for the process:
$pp \longrightarrow A(B) + X$

Differential cross section for a DPS event:
$pp \longrightarrow A + B + X$

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color factors **PDF**

$$\sigma_{eff}(x_1,x_2,x_3,x_4) = \frac{\sum_{i,j,k,l} \boxed{C_{ik}C_{jl}} F_i(x_1)F_j(x_2)F_k(x_3)\boxed{F_l(x_4)}}{\sum_{i,j,k,l} C_{ik}C_{jl}\int d^2z_\perp\, F_{ij}(x_1,x_3,z_\perp)F_{kl}(x_3,x_4,z_\perp)}$$

M.R., S. Scopetta et al, PLB 752

M. Traini, M.R., S. Scopetta and V. Vento, PLB 768 (2017)

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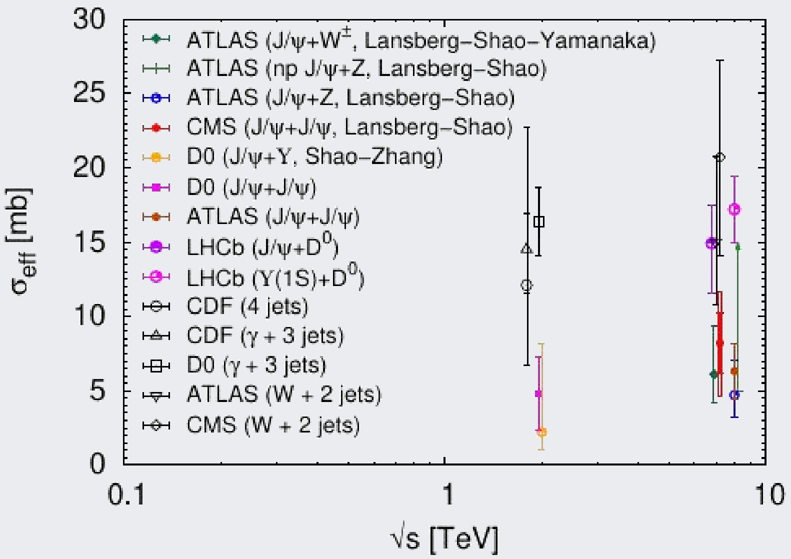
$$\sigma_{\text{eff}}^{\text{pp}} = \frac{\text{m}}{2}\frac{\sigma_{\text{A}}^{\text{pp}}\sigma_{\text{B}}^{\text{pp}}}{\sigma_{\text{DPS}}^{\text{pp}}}$$



J.P. Lansberg's slide
MPI-2019 workshop

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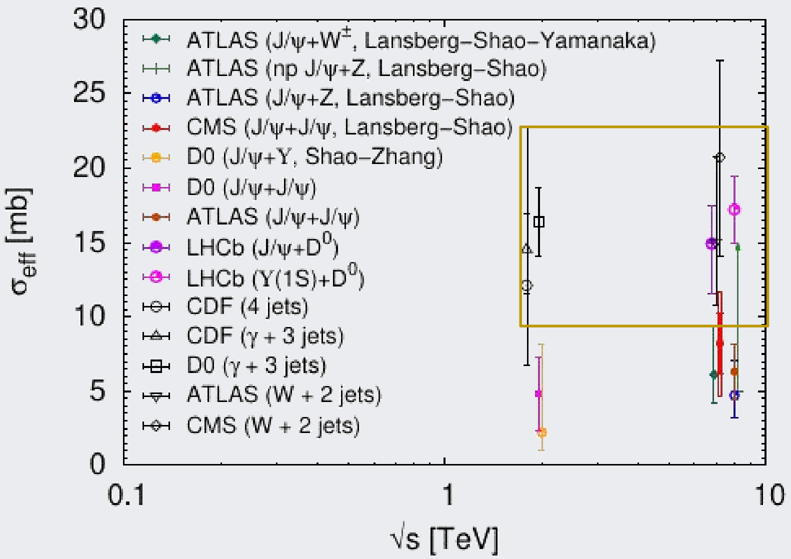

J.P. Lansberg's slide
MPI-2019 workshop

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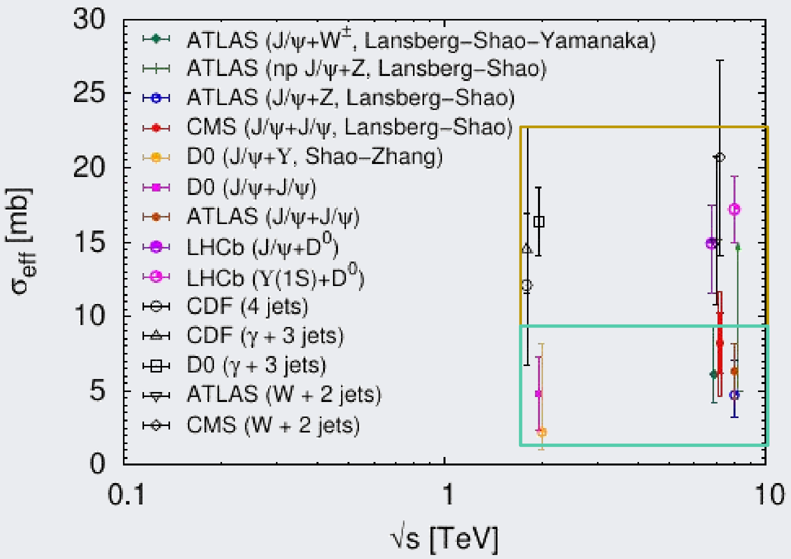
$$\sigma_{\rm eff}^{pp} = \frac{m}{2}\frac{\sigma_A^{pp}\sigma_B^{pp}}{\sigma_{\rm DPS}^{pp}}$$

- SENSITIVE TO CORRELATIONS
- PROCESS DEPENDENT?
- SENSITIVE TO INFORMATION ON THE PROTON STRUCTURE?
  As predicted by quark models

M.R. et al PLB 752,40 (2016)
M. Traini, M. R. et al, PLB 768, 270 (2017)
M. R. et al, Phys.Rev. D95 (2017) no.11, 114030



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- SENSITIVE TO CORRELATIONS
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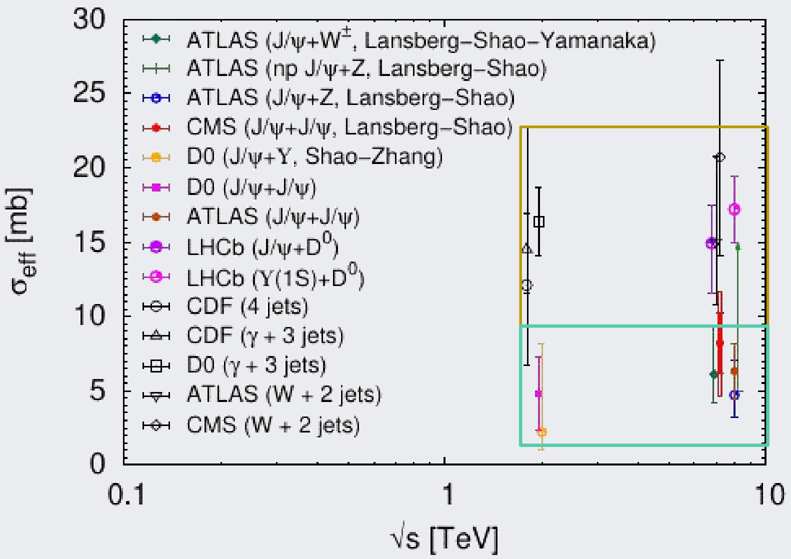
- SENSITIVE TO INFORMATION ON THE PROTON STRUCTURE?
  and phenomenological analyses
  T. Kasemets et al, JHEP 10 (2020) 214
  ...



J.P. Lansberg's slide
MPI-2019 workshop

M. R. and F. A. Ceccopieri, PRD 97, no. 7, 071501 (2018)

5 Clues from data?

If dPDFs factorize in terms of PDFs then $$\sigma_{\text{eff}}^{-1} = \int \frac{d^2 k_\perp}{(2\pi)^2} \boxed{T(k_\perp)}^2$$ → Effective form factor (EFF)

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$$T(k_\perp) \propto \int dx_1 dx_2 \, \tilde{F}(x_1, x_2, k_\perp)$$

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$$\langle z_\perp^2 \rangle \propto \frac{d}{k_\perp dk_\perp} T(k_\perp) \Big|_{k_\perp = 0}$$

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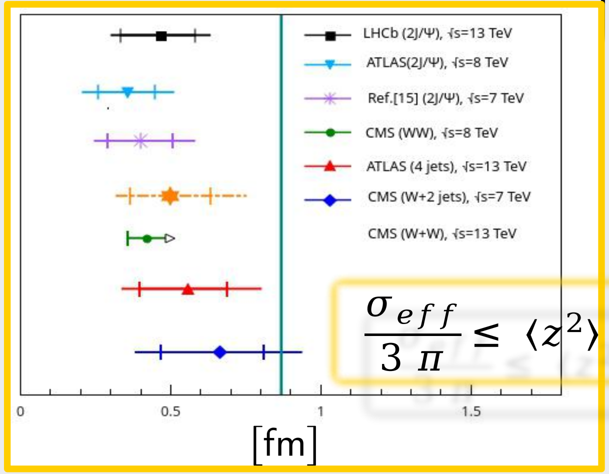
DPS processes:

The vertical line stands for the transverse proton radius



$$\frac{\sigma_{eff}}{3\pi} \leq \langle z^2 \rangle \leq \frac{\sigma_{eff}}{\pi}$$

M. R. and F. A. Ceccopieri, PRD 97, no. 7, 071501 (2018)

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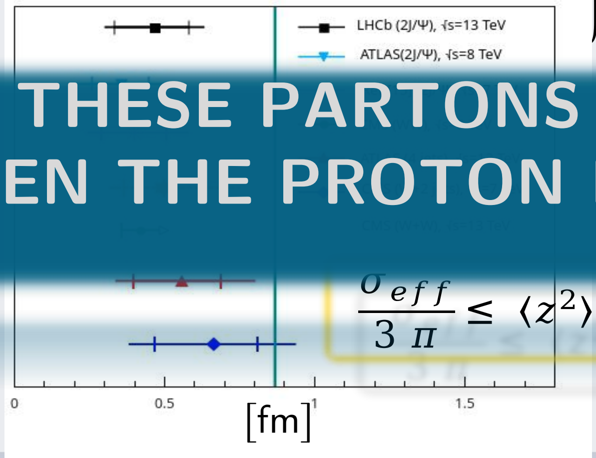
THE DISTANCE OF THESE PARTONS SEEMS TO BE SMALLER THEN THE PROTON RADIUS

the transverse proton radius



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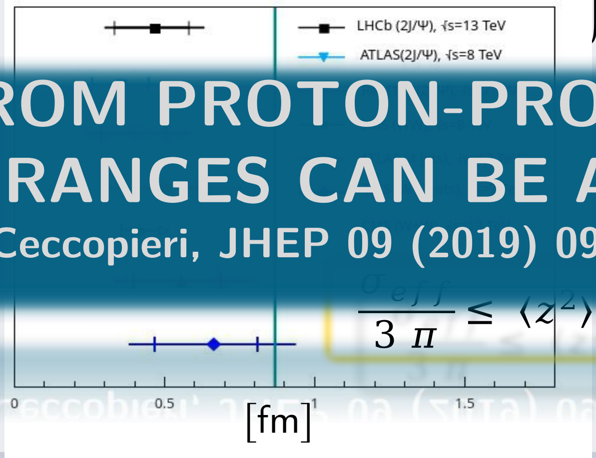
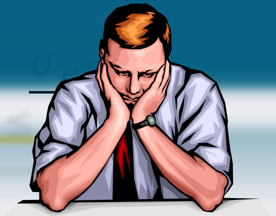


HOWEVER FROM PROTON-PROTON COLLISIONS ONLY RANGES CAN BE ACCESSED

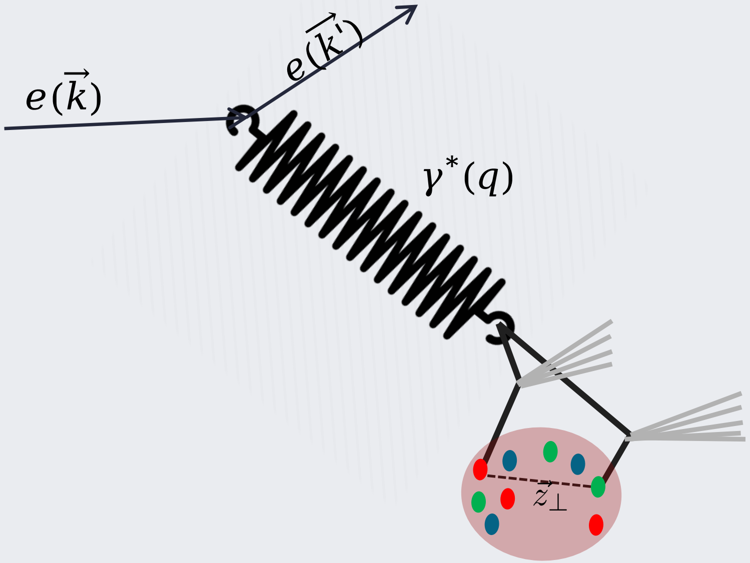
M.R. and F. A. Ceccopieri, JHEP 09 (2019) 097

$$\frac{\sigma_{eff}}{3\pi} \le \langle z^2 \rangle \le$$



6 New Idea: DPS via γ-p interaction

We consider the possibility offered by a DPS process involving a photon FLACTUATING in a quark-antiquark pair interacting with a proton:



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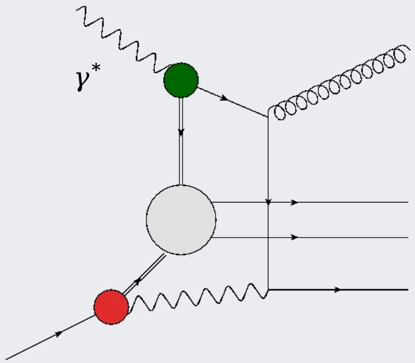
In order to study the impact of the DPS contribution to a process initiated via photon-proton interactions we evaluated the 4-JET photoproduction at HERA **(S. Checkanov et al. (ZEUS), Nucl. Phys B792, 1 (2008))**



In

1) **G. Abbiend et al, Phys. Commun 67, 465 (1992)**
2) **J.R. Forshaw et al, Z. Phys. C 72, 637 (1992)**

It has been shown that the agreement with data improves if MPI are included in the Monte Carlo

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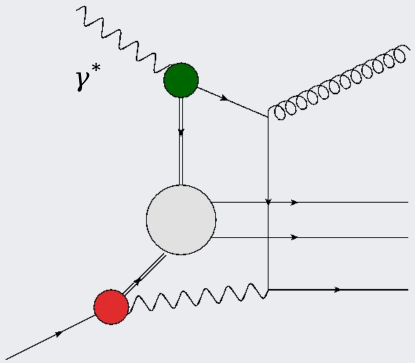

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WE EVALUATE THE DPS CONTRIBUTION TO THIS PROCESS

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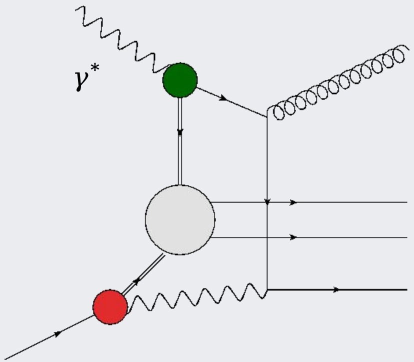

For this first investigation, we make use of the POCKET FORMULA:

$$d\sigma^{4j}_{DPS} = \frac{1}{2}\sum_{ab,cd}\int dy\, dQ^2 \frac{\boxed{f_{\gamma/e}(y,Q^2)}}{\sigma_{eff}^{\gamma p}(Q^2)} \times$$

$$\left.\begin{aligned} &\times \int dx_{p_a} dx_{\gamma_b} f_{a/p}(x_{p_a}) f_{b/\gamma}(x_{\gamma_b}) d\hat{\sigma}^{2j}_{ab}(x_{p_a}, x_{\gamma_b}) \\ &\times \int dx_{p_c} dx_{\gamma_d} \underbrace{f_{c/p}(x_{p_c})}_{\text{p-PDF}} \underbrace{f_{d/\gamma}(x_{\gamma_d})}_{\gamma\text{-PDF}} d\hat{\sigma}^{2j}_{cd}(x_{p_c}, x_{\gamma_d}) \end{aligned}\right\} \begin{matrix} \text{SPS*} \\ \times \\ \text{SPS} \end{matrix}$$

Flux Factor
P. Nason et al,PLB319 339 (1993)

p-PDF **(J. Pumplin et al. JHEP 07, 012 (2002))**

γ -PDF **(M. Gluck et al. PRD46, 1973 (1992))**

*Single Parton Scattering (SPS)

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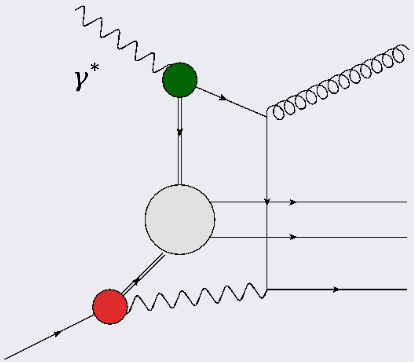

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$$d\sigma^{4j}_{DPS} = \frac{1}{\ldots} \sum \int dy\, dQ^2 \frac{\boxed{f_{\gamma/e}(y, Q^2)}}{\sigma_{eff}^{\gamma p}(Q^2)} \times$$

Flux Factor

P. Nason et al,PLB319 339 (1993)

$$\times \ldots) d\hat{\sigma}^{2j}_{ab}(x_{p_a}, x_{\gamma_b}) \quad \text{SPS*}$$

$$\times$$

$$\times \ldots) d\hat{\sigma}^{2j}_{cd}(x_{p_c}, x_{\gamma_d}) \quad \text{SPS}$$

PDF **(M. Gluck et al. PRD46, 1973 (1992))**

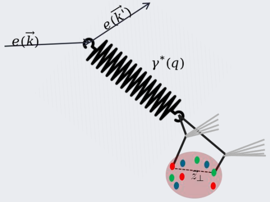
2002))

The main quantity we have to evaluate is:

$$\sigma_{eff}^{\gamma p}(Q^2)$$

*Single Parton Scattering (SPS)

6 The γ-p effective cross section



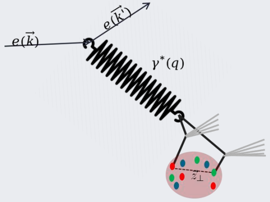
The expression of this quantity is very similar to the proton-proton collision case and can be formally derived by comparing the product of SPS cross sections and the DPS one obtained in **Gaunt, JHEP 01, 042 (2013)** and describing a DPS from a vector bosons splitting with given Q^2 virtuality

Proton EFF

$$\left[\sigma_{\rm eff}^{\gamma p}(Q^2)\right]^{-1} = \int \frac{d^2k_\perp}{(2\pi^2)} \boxed{T_p(k_\perp)}\, T_\gamma(k_\perp;Q^2)$$

M. R. and F. A. Ceccopieri, arXiv:2103.13480

6 The γ-p effective cross section



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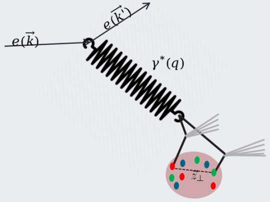
M. R. and F. A. Ceccopieri, arXiv:2103.13480

6 The γ-p effective cross section



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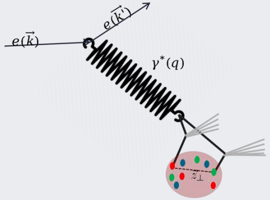
Proton EFF

This quantity is similar to an EFF

The full DPS cross section depends on the amplitude of the splitting photon in a $q\,\bar{q}$ pair. The latter can be formally described within a Light-Front (LF) approach in terms of LF wave functions.

M. R. and F. A. Ceccopieri, arXiv:2103.13480

6 The γ-p effective cross section



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$\boxed{T_p(k_\perp) \quad \text{proton EFF}}$

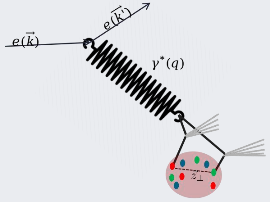
2

3 ψ^γ Photon WF

For the proton EFF use has been made of three choices:

M. R. and F. A. Ceccopieri, arXiv:2103.13480

6 The γ-p effective cross section



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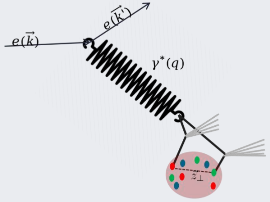
For the proton EFF use has been made of three choices:

1) G1: $\boxed{e^{-\alpha_1 k_\perp^2}}$, $\alpha_1 = 1.53 \text{ GeV}^{-2} \implies \sigma_{\rm eff}^{pp} = 15 \text{ mb}$

M. R. and F. A. Ceccopieri, arXiv:2103.13480

6 The γ-p effective cross section



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M. R. and F. A. Ceccopieri, arXiv:2103.13480

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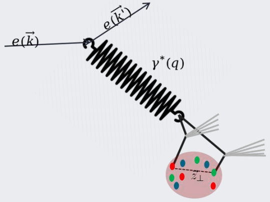

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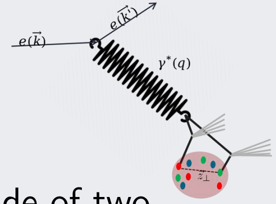
3) S: $\boxed{\left(1 + \frac{k_\perp^2}{m_g^2}\right)^{-4}}$, $m_g^2 = 1.1\ \mathrm{GeV}^2 \Longrightarrow \sigma_{\rm eff}^{pp} = 30$ mb

B. Blok et al, EPJC74, 2926 (2014)

M. R. and F. A. Ceccopieri, arXiv:2103.13480

6 The γ-p effective cross section



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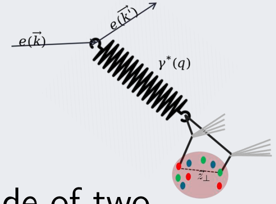
1) QED at LO **(S.J. Brodsky et al. PRD50, 3134 (1994))**:

$$\boxed{\psi^{\lambda=\pm}_{q,\bar{q}}(x, k_{1\perp}; Q^2) = -e_f \frac{\bar{u}_q(k)\ \gamma \cdot \varepsilon^\lambda\ v_{\bar{q}}(q-k)}{\sqrt{x(1-x)}\left[Q^2 + \frac{k_{1\perp}^2 + m^2}{x(1-x)}\right]}}$$

M. R. and F. A. Ceccopieri, arXiv:2103.13480

6 The γ-p effective cross section



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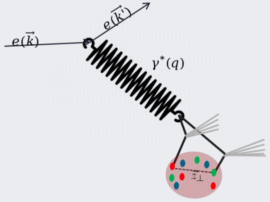
2) Non-Pertubative (NP) effects (E.R.Arriola et al, PRD74,054023 (2006))

$$\boxed{\psi_A^\gamma(x, k_{\perp 1}; Q^2) = \frac{6(1+Q^2/m_\rho^2)}{m_\rho^2\left(1 + 4\frac{k_{\perp 1}^2 + Q^2 x(1-x)}{m_\rho^2}\right)^{5/2}}}$$

M. R. and F. A. Ceccopieri, arXiv:2103.13480

6 The γ-p effective cross section



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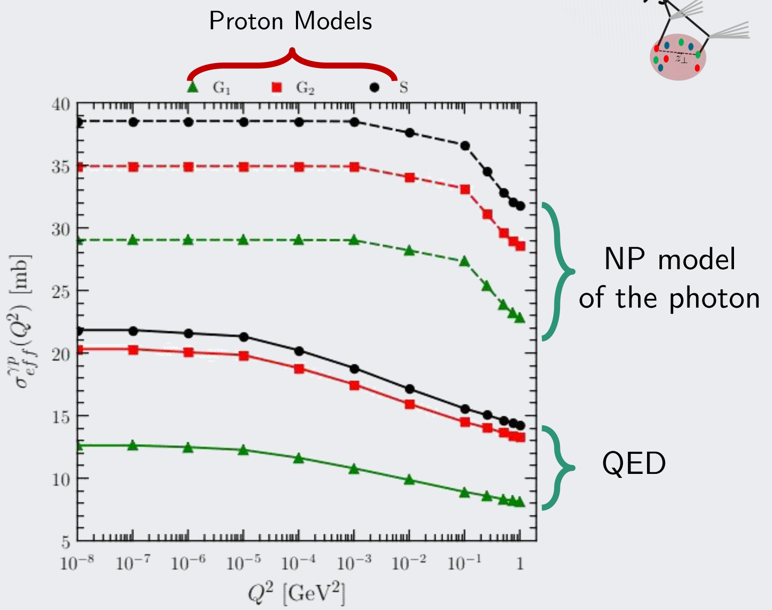
3 $\boxed{\psi_\gamma \quad \text{Photon WF}}$



M. R. and F. A. Ceccopieri, arXiv:2103.13480

6 The 4 jet DPS cross section

The HERA KINEMATICS:

S. Checkanov et al. (ZEUS), Nucl. Phys B792, 1 (2008)

$E_T^{\text{jet}} > 6\ \mathrm{GeV}$	Transverse energy of the jets
$\lvert\eta_{\text{jet}}\rvert < 2.4$	Pseudorapidity
$Q^2 < 1\ \mathrm{GeV}^2$	Photon virtuality
$0.2 \leqslant y \leqslant 0.85$	Inelasticity

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The ZEUS collaboration quoted an integrated total 4-jet cross section of 136 pb
S. Checkanov et al. (ZEUS), Nucl. Phys B792, 1 (2008)

6 The 4 jet DPS cross section

KINEMATICS:

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$Q^2 < 1$ GeV^2

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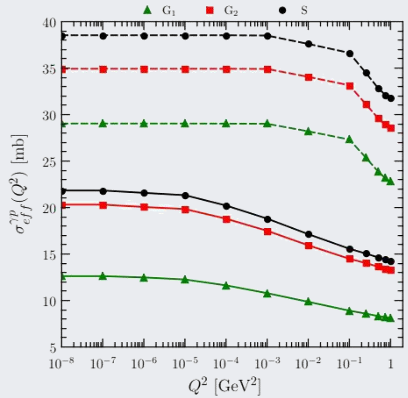
$$d\sigma^{4j}_{\rm DPS} = \frac{1}{2}\sum_{ab,cd}\int dy\, dQ^2\, \frac{f_{\gamma/e}(y, Q^2)}{\boxed{\sigma_{\rm eff}^{\gamma p}(Q^2)}} \times$$
$$\times \int dx_{p_a} dx_{\gamma_b} f_{a/p}(x_{p_a}) f_{b/\gamma}(x_{\gamma_b}) d\hat{\sigma}^{2j}_{ab}(x_{p_a}, x_{\gamma_b})$$
$$\times \int dx_{p_c} dx_{\gamma_d} f_{c/p}(x_{p_c}) f_{d/\gamma}(x_{\gamma_d}) d\hat{\sigma}^{2j}_{cd}(x_{p_c}, x_{\gamma_d})$$



M. R. and F. A. Ceccopieri, arXiv:2103.13480

6 The 4 jet DPS cross section

KINEMATICS

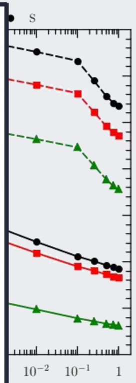
$E_T^{jet} > 6$ GeV

$|\eta_{jet}| < 2.4$

$Q^2 < 1$ GeV

$0.2 \leqslant y \leqslant 0.$

photon	proton	$Q^2 \leq 10^{-2}$ [GeV^2]	$10^{-2} \leq Q^2 \leq 1$ [GeV^2]	$Q^2 \leq 1$ [GeV^2]	$\frac{\sigma_{DPS}}{\sigma_{tot}}$ [%]
		σ_{DPS} [pb]			
NP model	G_1	35.1	18.6	53.7	40
NP model	G_2	29.1	15.2	44.3	33
NP model	S	26.4	13.7	40.1	30
QED	G_1	87.8	54.3	142.1	101
QED	G_2	54.3	33.4	87.7	65
QED	S	50.5	31.1	81.6	60



6 The 4 jet DPS cross section

KINEMATICS

$E_T^{jet} > 6$ Ge[...]

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		σ_{DPS} [pb] $Q^2 \leq 10^{-2}$ [GeV^2]	$10^{-2} \leq Q^2 \leq 1$ [GeV^2]	$Q^2 \leq 1$ [GeV^2]	$\frac{\sigma_{DPS}}{\sigma_{tot}}$ [%]
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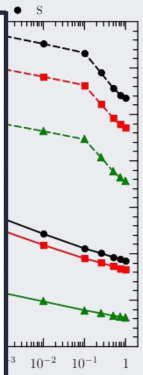
photon

proton



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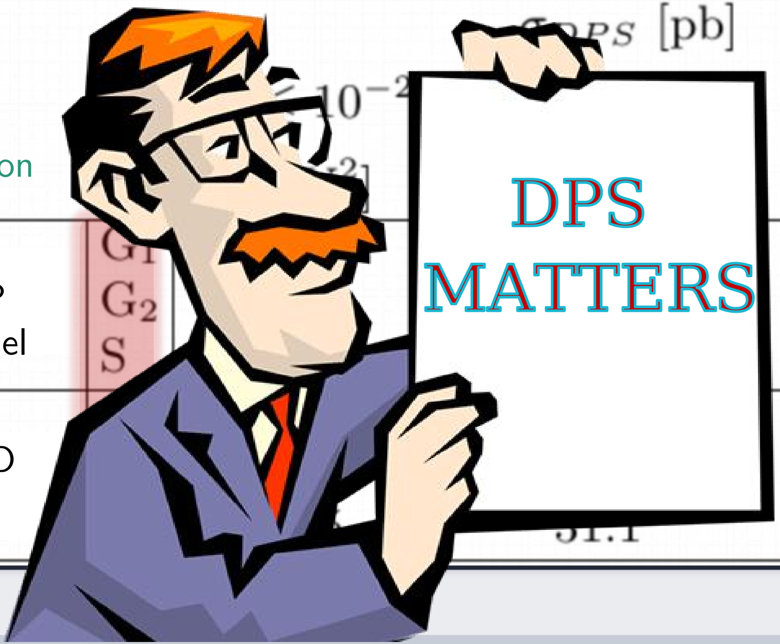

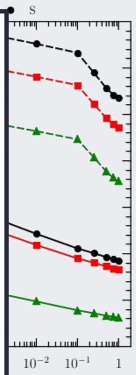

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photon	proton	σ_{DPS} [pb]		$Q^2 \le 1$ [GeV^2]	$\frac{\sigma_{DPS}}{\sigma_{tot}}$ [%]
NP model	G_1			3.7	40
	G_2			44.3	33
	S			40.1	30
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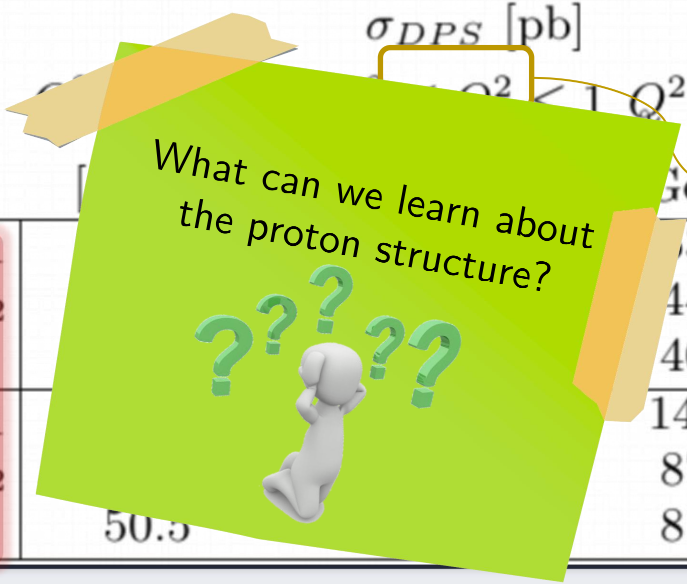

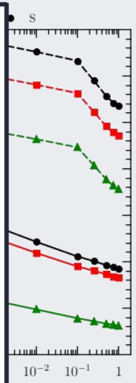

6 The effective cross section: a key for the proton structure

The effective cross section can be also written in terms of Fourier Transform of the EFF:

$$\tilde{F}(z_\perp)$$

The probability of finding a parton pair at distance

$$z_\perp$$

M. R. and F. A. Ceccopieri, arXiv:2103.13480

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If we can measure the dependence of the effective-cross section on the photon VIRTUALITY

M. R. and F. A. Ceccopieri, arXiv:2103.13480

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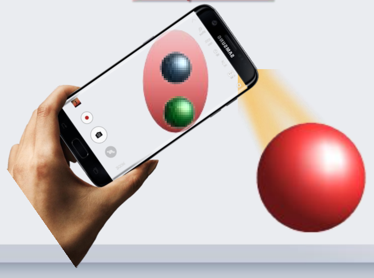
We could access for the first time the <u>mean transverse distance between partons in the proton</u>



M. R. and F. A. Ceccopieri, arXiv:2103.13480

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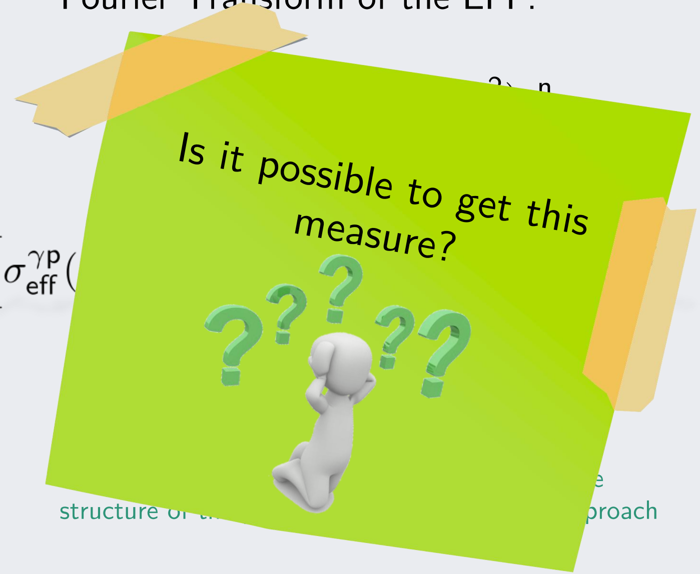

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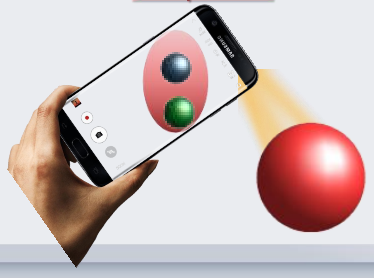

6 The effective cross section: a key for the proton structure

To test if in future a dependence of the effective cross section on the photon virtuality could be observed, we considered again the 4 JET photoproduction:

M. R. and F. A. Ceccopieri, arXiv:2103.13480

6 The effective cross section: a key for the proton structure

To test if in future a dependence of the effective cross section on the photon virtuality could be observed, we considered again the 4 JET photoproduction:

1) We divided the integral of the cross section on Q^2 in two intervals:

$$Q^2 \leqslant 10^{-2} \quad \text{and} \quad 10^{-2} \leqslant Q^2 \leqslant 1 \quad \text{GeV}^2$$

M. R. and F. A. Ceccopieri, arXiv:2103.13480

6 The effective cross section: a key for the proton structure

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2) We have estimated for each photon and proton models a constant effective cross section $\bar{\sigma}_{\text{eff}}^{\gamma p}$ (with respect to Q^2) such that the total integral of the cross section on Q^2 reproduce the full calculation obtained by means of $\sigma_{\text{eff}}^{\gamma p}(Q^2)$

M. R. and F. A. Ceccopieri, arXiv:2103.13480

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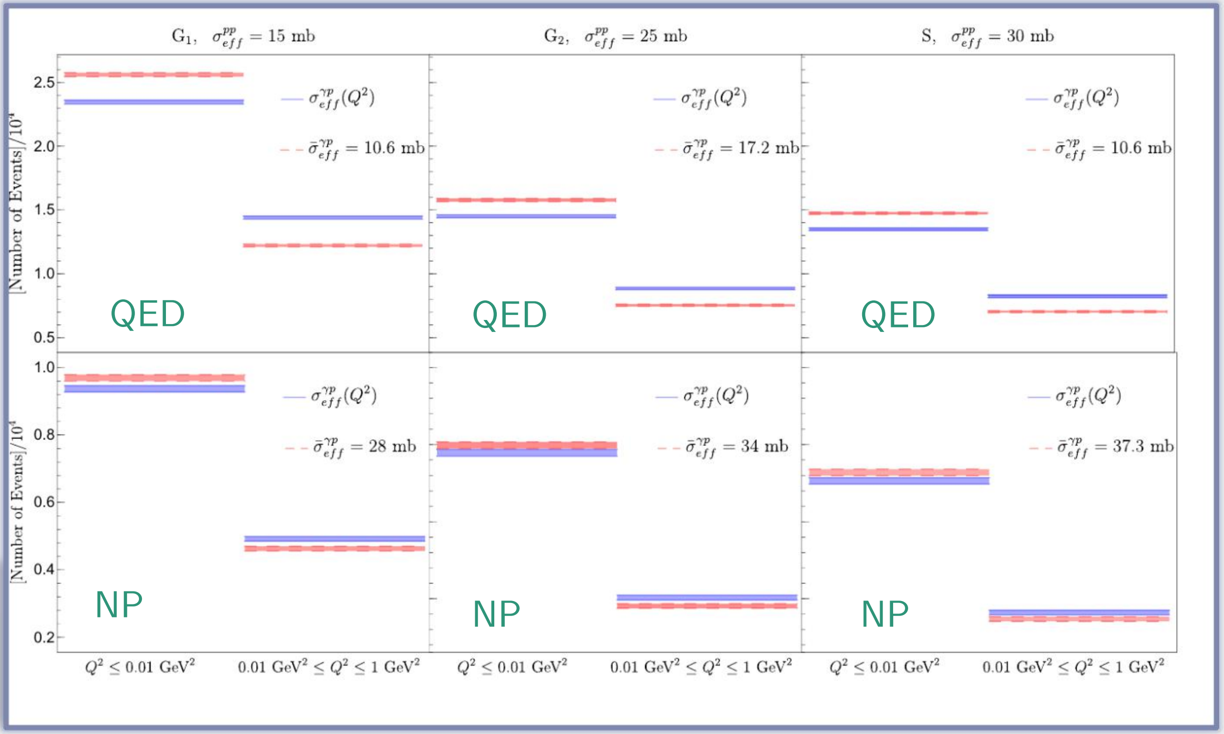
3) We estimate the minimum luminosity to distinguish the two cases

M. R. and F. A. Ceccopieri, arXiv:2103.13480

6 The effective cross section: a key for the proton structure

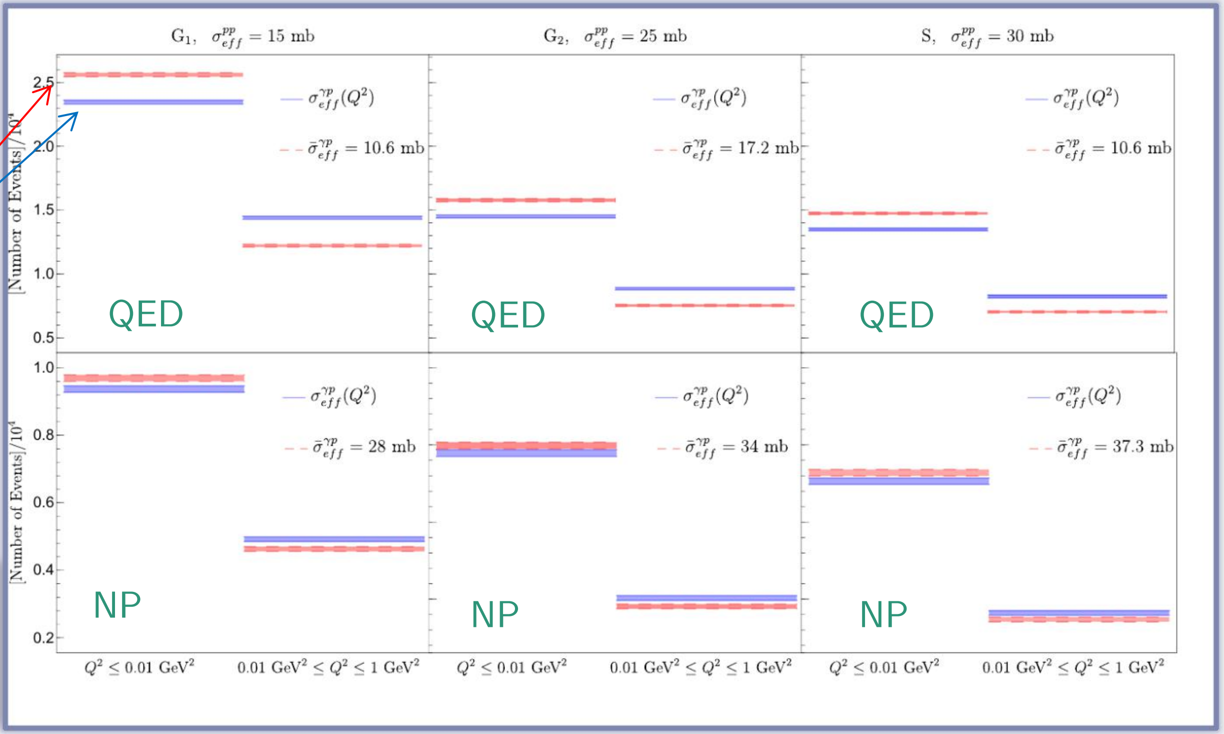
Proton

photon

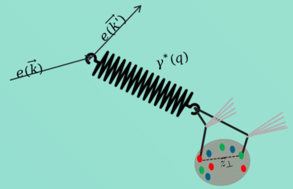


6 The effective cross section: a key for the proton structure

With an integrated luminosity of 200 pb^{-1} we can separate:



CONCLUSIONS



1) We investigated the impact of correlations in DPS proton-proton collisions to learn something new on the parton structure of the proton

2) We demonstrated that in p-p collisions only some limited information on the proton can be obtained

3) We proposed to consider DPS initiated via photon-proton interactions by showing that:

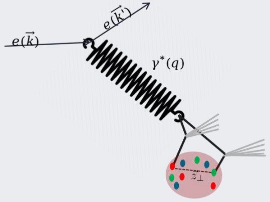
   * DPS can contribute also in this case. Cross section of the 4 jet photo production strongly affected

   * The dependence of $\sigma_{\text{eff}}^{\gamma p}(Q^2)$ on the Q^2 can unveil the mean distance of partons in the proton

   * We show that by increasing the luminosity such a dependence can be exposed in future facilities such as the **Electron Ion Collider**

   * In the future could be interesting to study other processes with different final states such as those associated to the **QUARKONIUM PRODUCTION**

6 The γ-p effective cross section

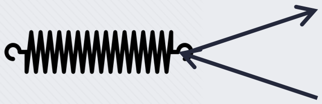


The expression of this quantity is very similar to the proton-proton collision case and can be formally derived by comparing the product of SPS cross sections and the DPS one obtained in **Gaunt, JHEP 01, 042 (2013)** and describing a DPS from a vector bosons splitting with given Q^2 virtuality

$$\left[\sigma_{\text{eff}}^{\gamma p}(Q^2)\right]^{-1} = \int \frac{d^2 k_\perp}{(2\pi^2)} T_p(k_\perp) T_\gamma(k_\perp; Q^2)$$

The full DPS cross section depends on the amplitude of the splitting photon in a $q\,\bar{q}$ pair. The latter can be formally described within a Light-Front (LF) approach in terms of LF wave functions (W.F.):

$$f^{\gamma}_{q,\bar{q}}(x, \tilde{k}_\perp; Q^2) = \int d^2 k_{\perp,1}\, \psi^{\dagger\gamma}_{q\bar{q}}(x, \vec{k}_{\perp,1}; Q^2) \times \psi^{\gamma}_{q\bar{q}}(x, \vec{k}_{\perp,1} + \vec{k}_\perp; Q^2)$$

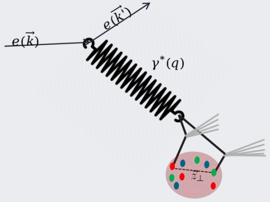


$q: (x, \vec{k}_{\perp,1})$

$\bar{q}: (1-x, -\vec{k}_{\perp,1})$

M. R. and F. A. Ceccopieri, arXiv:2103.13480

6 The γ-p effective cross section



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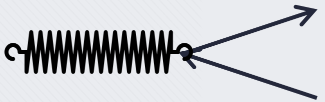

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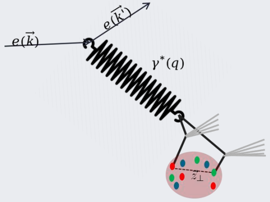
Similar definition of a meson dPDF

M. R. and F. A. Ceccopieri, arXiv:2103.13480

M. R. et al., EPJC78, 781 (2018)

6 The γ-p effective cross section



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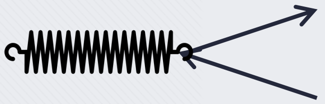
$$f_{q,\bar{q}}^{\gamma}(x, \tilde{k}_\perp; Q^2) = \int d^2k_{\perp,1}\, \psi_{q\bar{q}}^{\dagger\gamma}(x, \vec{k}_{\perp,1}; Q^2) \times \psi_{q\bar{q}}^{\gamma}(x, \vec{k}_{\perp,1} + \vec{k}_\perp; Q^2)$$

$$T_\gamma(k_\perp; Q^2) = \frac{\sum_q \int dx\, f_{q,\bar{q}}^{\gamma}(x, k_\perp; Q^2)}{\sum_q \int dx\, f_{q,\bar{q}}^{\gamma}(x, k_\perp = 0; Q^2)}$$



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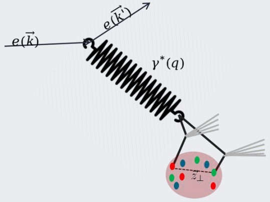


The expression of this quantity is very similar t... collision
case and can be formall... ...ss sections
and the DPS on... ...DPS from a

$$\left[\sigma_{\text{eff}}^{\gamma p}(Q^2)\right]^{-1} = \int \frac{d\ldots}{(2\ldots}$$

$$\ldots \int dx\, f^{\gamma}_{q,\bar{q}}(x, k_\perp; Q^2)$$

$$\ldots dx\, f^{\gamma}_{q,\bar{q}}(x, k_\perp = 0; Q^2)$$

$$f^{\gamma}_{q,\bar{q}}(x, \tilde{k}_\perp; Q^2) = \int d^2\ldots$$

INGREDIENTS OF THE CALCULATIONS

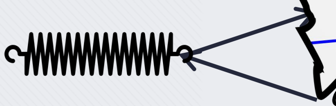
1

$T_p(k_\perp)$ proton EFF

2

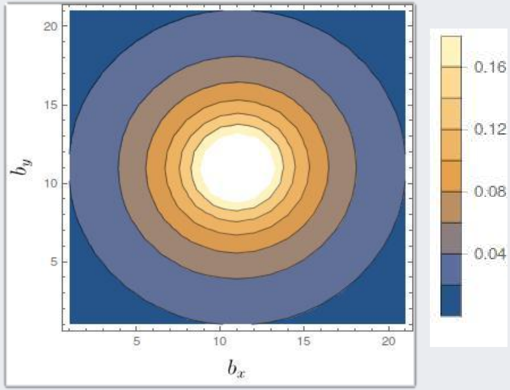
ψ^γ Photon WF

3



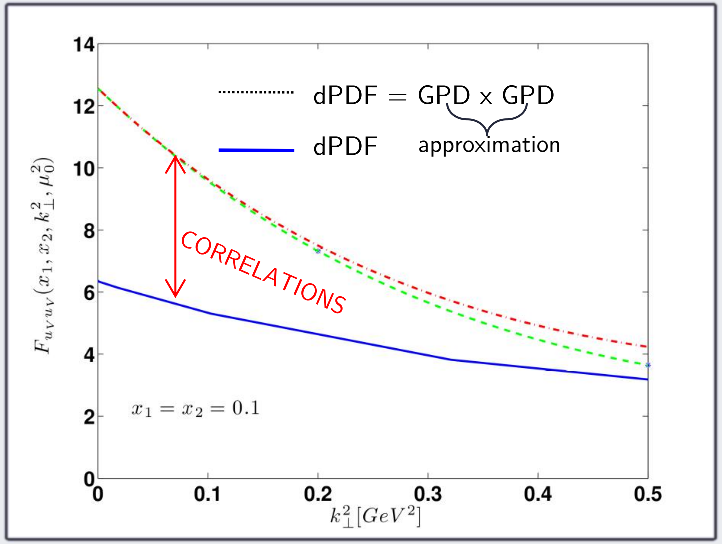
M. R. and F. A. Ceccopieri, arXiv:2103.13480

2 Information from Quark Models



M.R. and F. A. Ceccopieri, JHEP 09 (2019) 097

1) e.g. the distance distribution of two gluons in the proton



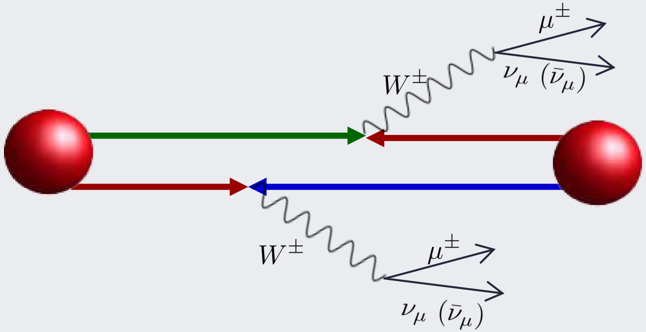
2) Correlations are important

M.R., S. Scopetta et al, JHEP 10 (2016) 063

M.R. and F. A. Ceccopieri PRD 95 (2017) 034040

4 Same sign W's production at the LHC

M. R. et al, Phys.Rev. D95 (2017) no.11, 114030

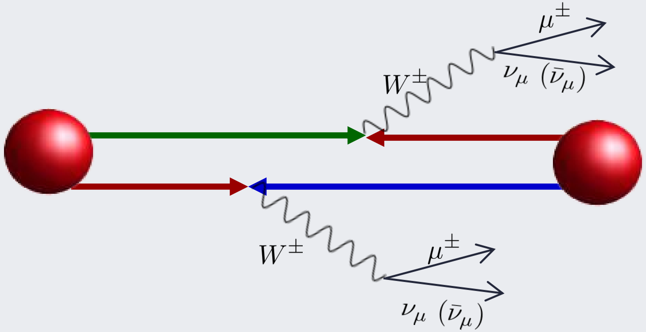


In this channel, the single parton scattering (usually dominant w.r.t to the double one) starts to contribute to higher order in strong coupling constant.

"Same-sign W boson pairs production is a promising channel to look for signature of double Parton interactions at the LHC."

4 Same sign W's production at the LHC

M. R. et al, Phys.Rev. D95 (2017) no.11, 114030



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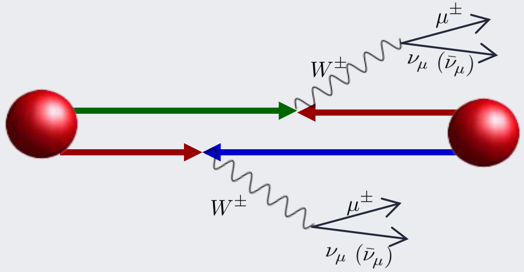
Can double parton correlations be observed for the first time in the next LHC run ?

4 Same sign W's production at the LHC

M. R. et al, Phys.Rev. D95 (2017) no.11, 114030



Kinematical cuts

$$pp,\ \sqrt{s} = 13\ \text{TeV}$$
$$p_{T,\mu}^{leading} > 20\ \text{GeV}, \quad p_{T,\mu}^{subleading} > 10\ \text{GeV}$$
$$|p_{T,\mu}^{leading}| + |p_{T,\mu}^{subleading}| > 45\ \text{GeV}$$
$$|\eta_\mu| < 2.4$$
$$20\ \text{GeV} < M_{inv} < 75\ \text{GeV or } M_{inv} > 105\ \text{GeV}$$

DPS cross section:

$$\frac{d^4\sigma^{pp\to\mu^\pm\mu^\pm X}}{d\eta_1 dp_{T,1} d\eta_2 dp_{T,2}} = \sum_{i,k,j,l} \frac{1}{2} \int d^2\vec{b}_\perp F_{ij}(x_1, x_2, \vec{b}_\perp, M_W) F_{kl}(x_3, x_4, \vec{b}_\perp, M_W) \frac{d^2\sigma_{ik}^{pp\to\mu^\pm X}}{d\eta_1 dp_{T,1}} \frac{d^2\sigma_{jl}^{pp\to\mu^\pm X}}{d\eta_2 dp_{T,2}} \mathcal{I}(\eta_i, p_{T,i})$$

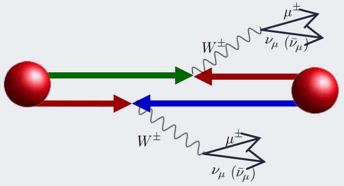
In order to estimate the role of double parton correlations we have used as input of dPDFs:

1) Longitudinal and transverse correlations arise from the relativistic CQM model describing three valence quarks

2) These correlations propagate to sea quarks and gluons through pQCD evolution

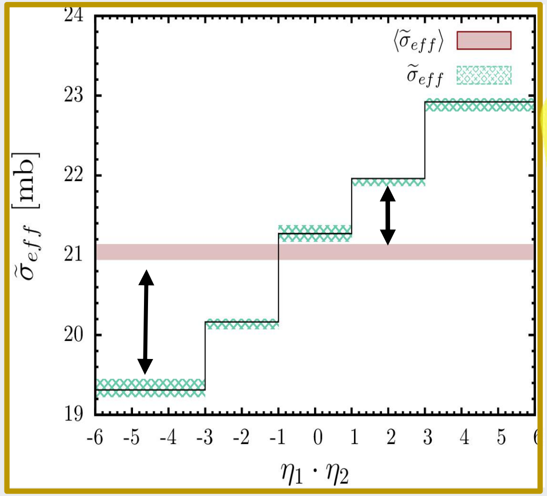
4 Same sign W's production at the LHC

M. R. et al, Phys.Rev. D95 (2017) no.11, 114030



$$\eta_1 \cdot \eta_2 \simeq \frac{1}{4} ln\frac{x_1}{x_3} ln\frac{x_2}{x_4}$$

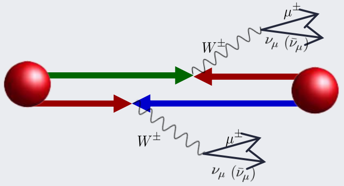
$$\boxed{\langle\tilde\sigma_{eff}\rangle = 21.04\,^{+0.07}_{-0.07}\,(\delta Q_0)\,^{+0.06}_{-0.07}(\delta\mu_F)\text{ mb}\,.}$$



Difference between green and red line is due to correlations effects

4 Same sign W's production at the LHC

M. R. et al, Phys.Rev. D95 (2017) no.11, 114030



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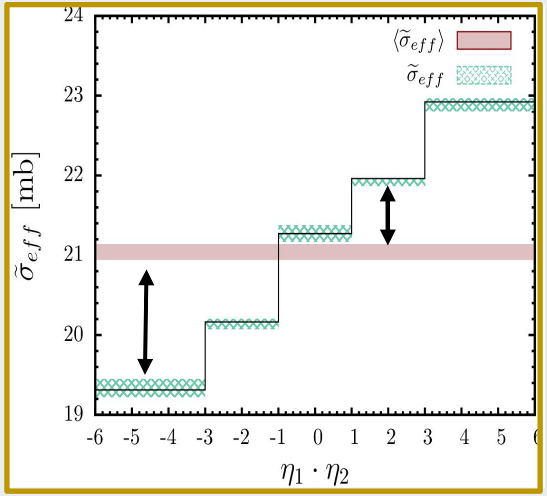


<u>*x*- dependence of effective x-section</u>

M.Rinaldi et al PLB 752,40 (2016)

M. Traini, M. R. et al, PLB 768, 270 (2017)

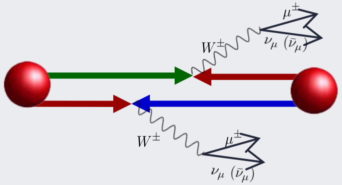
Assuming that the results of the first and the last bins can be distinguished if they differ by 1 sigma, we estimated that:

$$\boxed{\mathcal{L} = 1000\ \mathrm{fb}^{-1}}$$

is necessary to observe correlations

* to be updated to new CMS cuts

4 Same sign W's production at the LHC

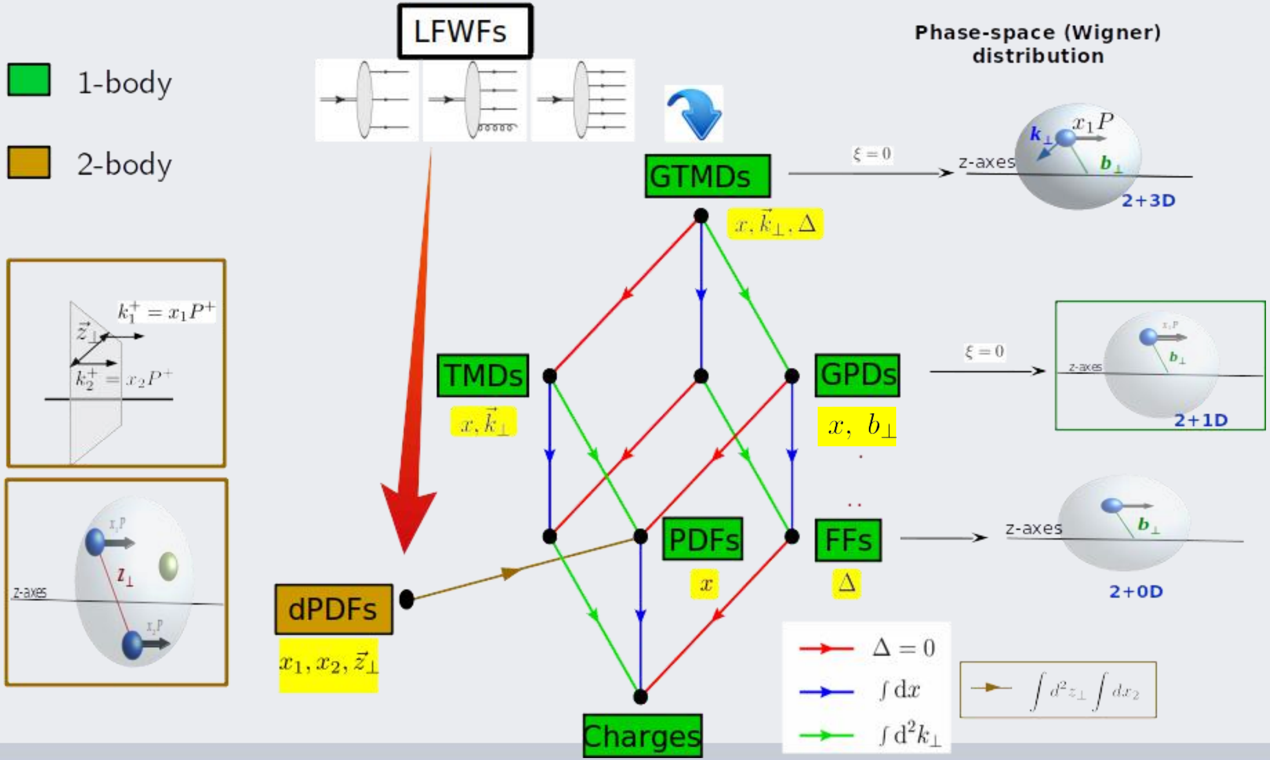


In Ref. **S. Cotogno et al, JHEP 10 (2020) 214**, it has been shown that several experimental observable are sensitive to **double spin correlations**.

The LHC has the potential to access these new information!

IN THIS CHANNEL, WE ESTABLISHED THE POSSIBILITY TO OBSERVE, FOR THE FIRST TIME, TWO-PARTON CORRELATIONS IN THE NEXT LHC RUN!

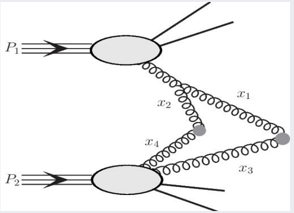
2 Multidimensional Pictures of Hadron



4 Further implementations

Relevant for processes involving heavy particles is the splitting term (1v2 mechanism) i.e.:

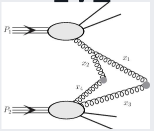
$^{*}D_{j_1,j_2}(x_1,x_2) = \int d^2b_\perp \tilde{F}_{j_1,j_2}(x_1,x_2,b_\perp)$



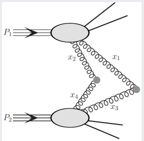
In pQCD evolution: $$\frac{dD_{j_1 j_2}(x_1,x_2;t)}{dt} = \begin{cases} \text{Homogeneous term (double DGLAP)} \\ + \\ \sum_{j'} F_{j'}(x_1+x_2;t)\underbrace{\frac{1}{x_1+x_2}P_{j'\to j_1 j_2}}_{\text{SPLITTING TERM}}\left(\frac{x_1}{x_1+x_2}\right) \end{cases}$$

Gaunt J.R. and Stirling W. J., JHEP 03 (2010)

1v2



+



2v2

$$\frac{\sigma_{eff}}{3\pi}\left(1+\frac{3}{2}r_v\right) \le \langle b^2\rangle \le \frac{\sigma_{eff}}{\pi}\left(1+2\boxed{r_v}\right)$$

$$r_v \sim \frac{F^{splitting}_{j_1 j_2}(x_1,x_2,k_\perp=0;t)}{F_{j_1 j_2}(x_1,x_2,k_\perp=0;t)}$$

with:

$0 \le r_v \le 1$

J.R. Gaunt, R. Maciula and A. Szczurek,
PRD 90 (2014) 054017

Due to the difficulty in the estimate of the 2 contributions:

Absolute minimum $r_v = 0$ $$\boxed{\frac{\sigma_{eff}}{3\pi}} \le \langle b^2\rangle \le \boxed{\frac{3\ \sigma_{eff}}{\pi}}$$ Absolute maximum $\mathbf{r_v = 1}$

4 Further implementations

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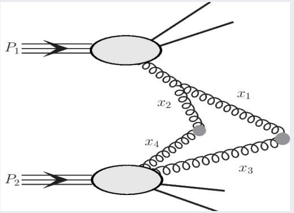
$) = \int d^2b_\perp \tilde{F}_{j_1,j_2}(x_1,x_2,b_\perp)$



In p

Gaus

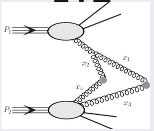
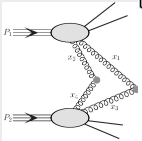
e DGLAP)

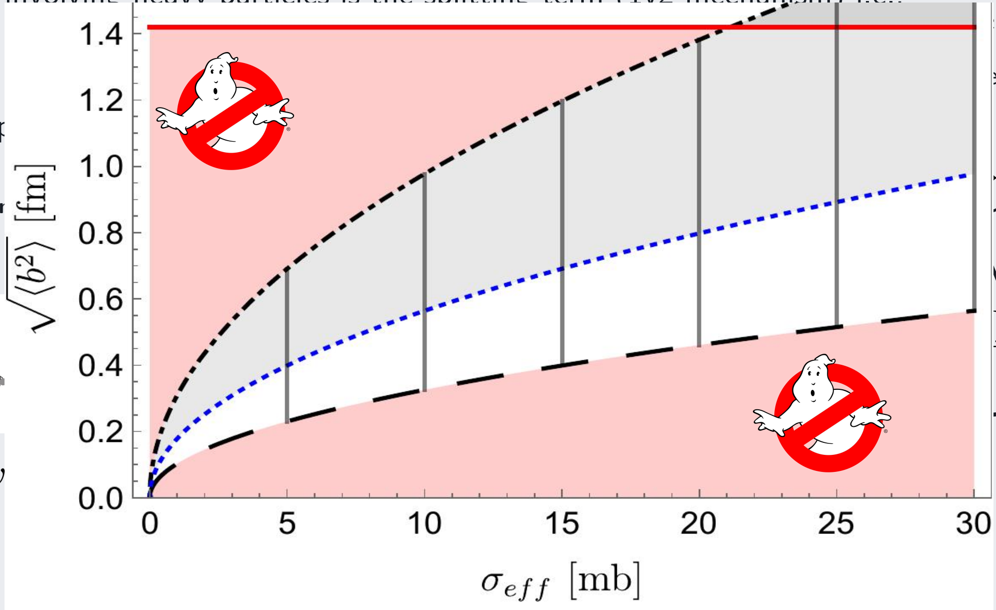
$j_1 j_2 \left(\frac{x_1}{x_1+x_2}\right)$

M

$\frac{{}^{splitting}_{j_1 j_2}(x_1,x_2,k_\perp=0;t)}{F_{j_1 j_2}(x_1,x_2,k_\perp=0;t)}$

1v2

 + 



with:
$0 \leq r_v \leq 1$

1) Minimum as function of r_v

$m(r_v)$

2) Maximum as function of r_v

$M(r_v)$

Absolute minimum $\boxed{\frac{\sigma_{eff}}{3\pi}} \leq \langle b^2 \rangle \leq \boxed{\frac{3\,\sigma_{eff}}{\pi}}$ Absolute maximum

$r_v = 0$ $\qquad$ $\mathbf{r_v = 1}$

4 Further implementations

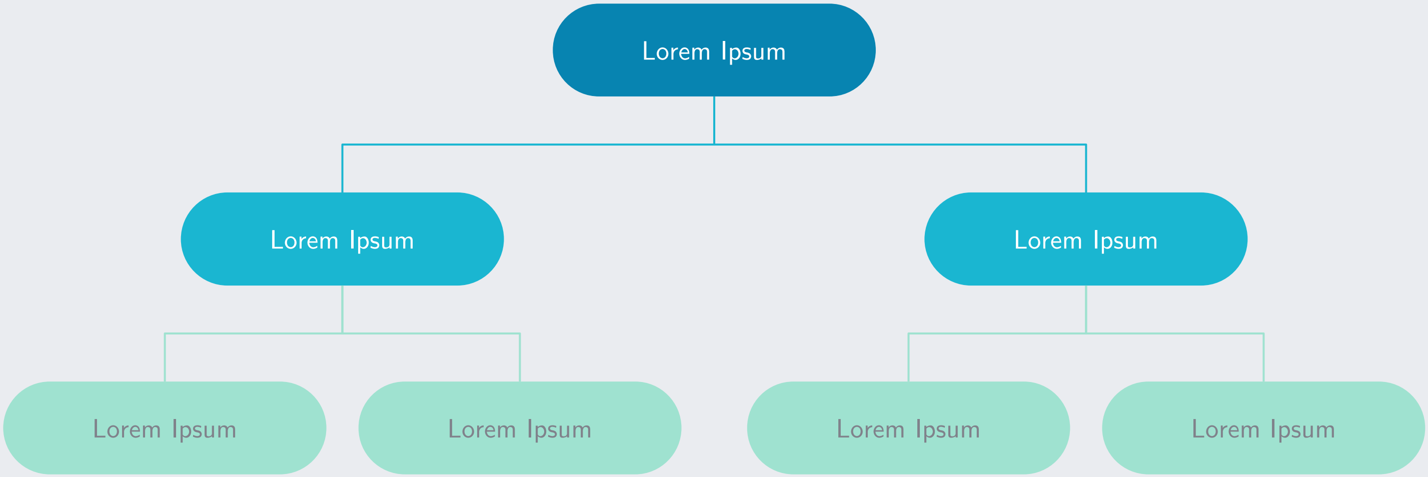
IF WE DO NOT CONSIDER ANY FACTORIZATION ANSATZ IN DOUBLE PDFs:

$$\frac{\sigma_{eff}(x_1,x_2)}{3\pi}\left[r^{2v2}(x_1,x_2)^2+\frac{3}{2}r^{2v1}(x_1,x_2)^2\ r_v\right]\leq\langle b^2\rangle_{x_1,x_2}\leq\frac{\sigma_{eff}(x_1,x_2)}{\pi}\left[r^{2v2}(x_1,x_2)^2+2r^{2v1}(x_1,x_2)^2\ r_v\right]$$

$$r^{2v2}(x_1,x_2)=\frac{F(x_1,x_2,k_\perp=0;t)}{F(x_1;t)F(x_2;t)}$$

$$r^{2v1}(x_1,x_2)=\frac{F^{splitting}(x_1,x_2,k_\perp=0;t)}{F(x_1;t)F(x_2;t)}$$

Use diagrams to explain your ideas



Our process is easy

Vestibulum congue tempus

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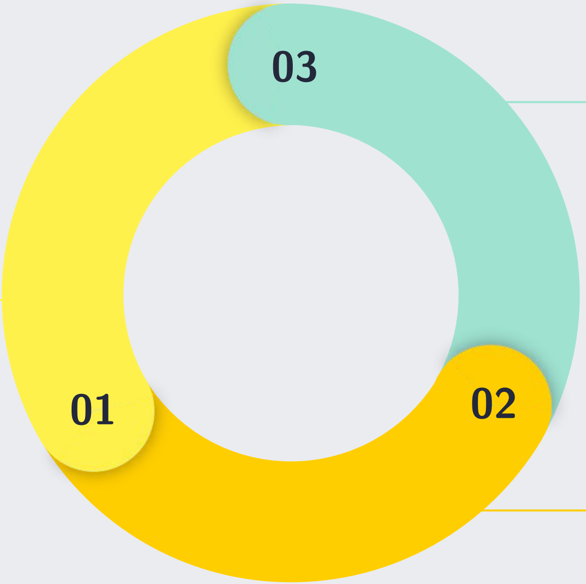
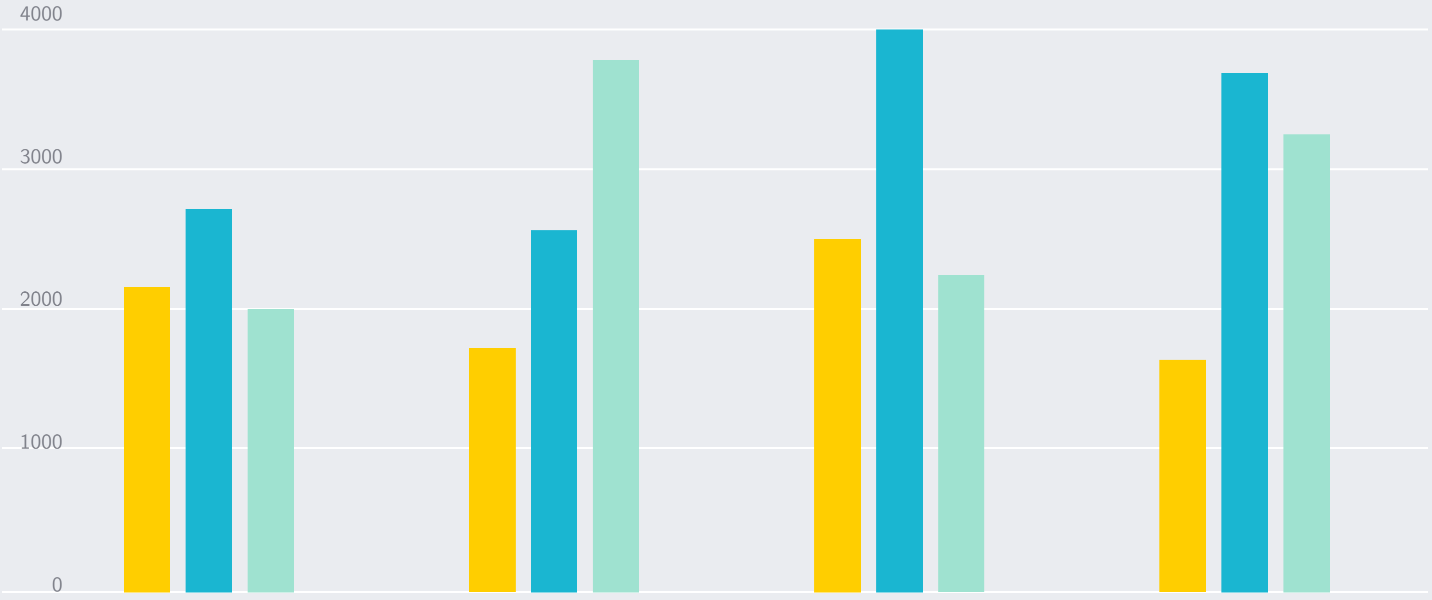


Vestibulum congue tempus

Lorem ipsum dolor sit amet, consectetur adipiscing elit, sed do eiusmod tempor. Donec facilisis lacus eget mauris.

Vestibulum congue tempus

Lorem ipsum dolor sit amet, consectetur adipiscing elit, sed do eiusmod tempor. Donec facilisis lacus eget mauris.



You can insert graphs from Excel or Google Sheets

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Show and explain your web, app or software projects using these gadget templates.

Timeline

JAN — Blue is the colour of the clear sky and the deep sea

FEB — Yellow is the color of gold, butter and ripe lemons

MAR — Red is the colour of danger and courage

APR — White is the color of milk and fresh snow

MAY — Black is the color of ebony and of outer space

JUN — Blue is the colour of the clear sky and the deep sea

JUL — Yellow is the color of gold, butter and ripe lemons

AUG — Red is the colour of danger and courage

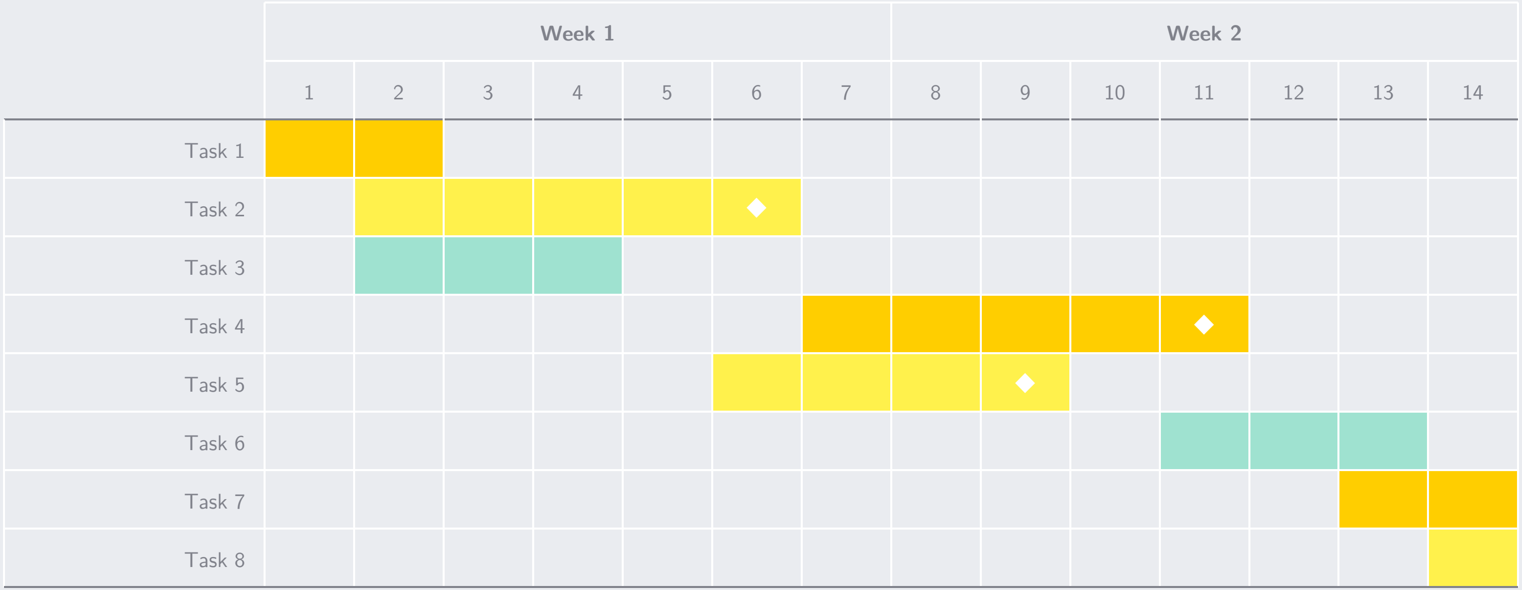
SEP — White is the color of milk and fresh snow

OCT — Black is the color of ebony and of outer space

NOV — Blue is the colour of the clear sky and the deep sea

DEC — Yellow is the color of gold, butter and ripe lemons

Gantt chart


Week 1
Week 2
1
2
3
4
5
6
7
8
9
10
11
12
13
14
Task 1
Task 2
Task 3
Task 4
Task 5
Task 6
Task 7
Task 8

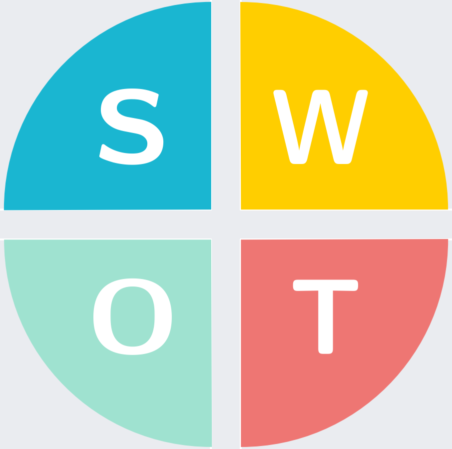
SWOT Analysis

STRENGTHS

Blue is the colour of the clear sky and the deep sea

WEAKNESSES

Yellow is the color of gold, butter and ripe lemons



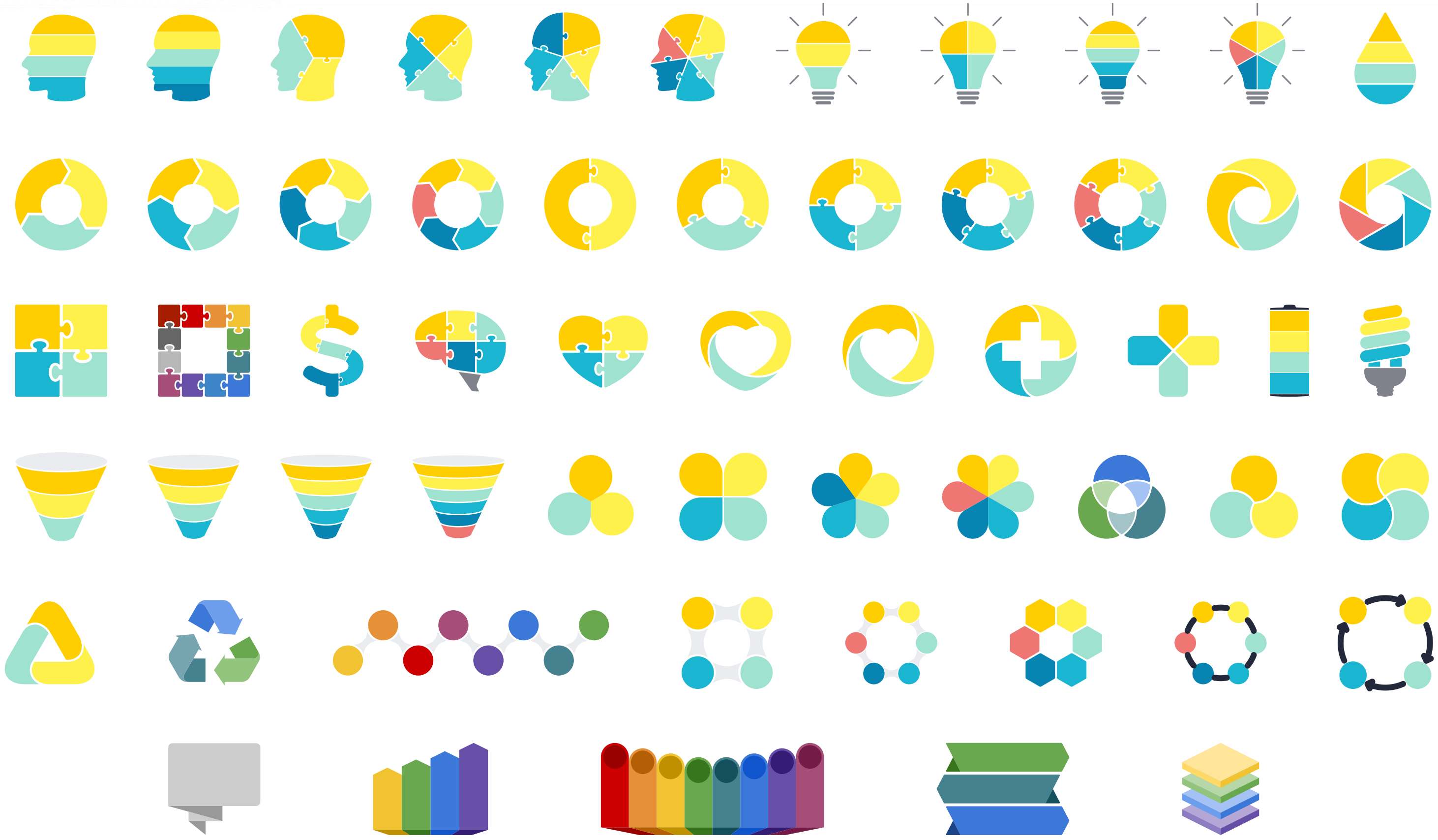
Black is the color of ebony and of outer space

OPPORTUNITIES

White is the color of milk and fresh snow

THREATS

Diagrams and infographics



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