



**Faculty
of Physics**

WARSAW UNIVERSITY OF TECHNOLOGY



ALICE



Studying the kaon- proton strong interactions with **ALICE** at the LHC

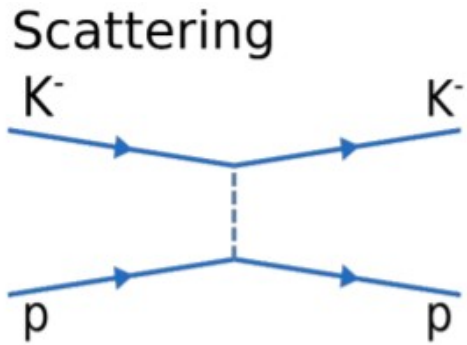
Łukasz Graczykowski
on behalf of the ALICE Collaboration

HADRON 2021

29 July 2021



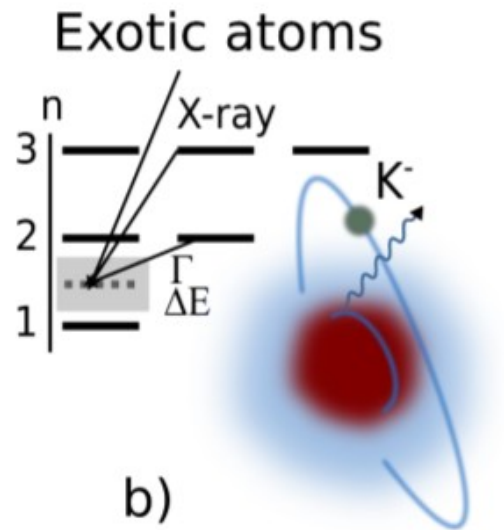
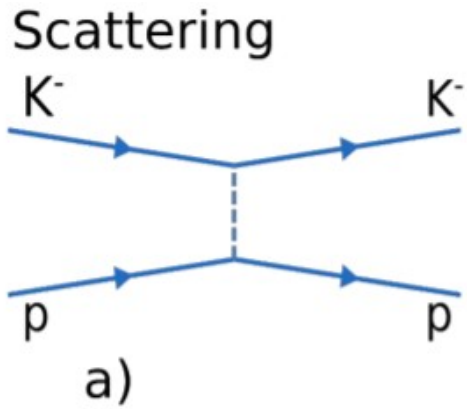
Experimental study of Kp interaction



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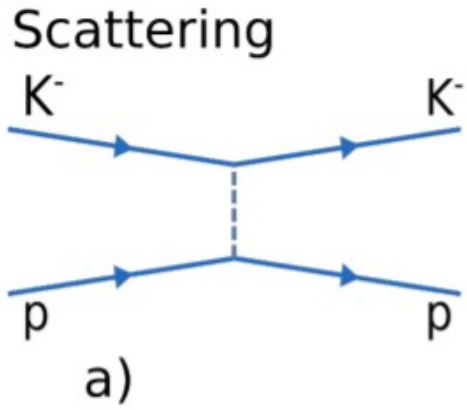
Experimental study of Kp interaction



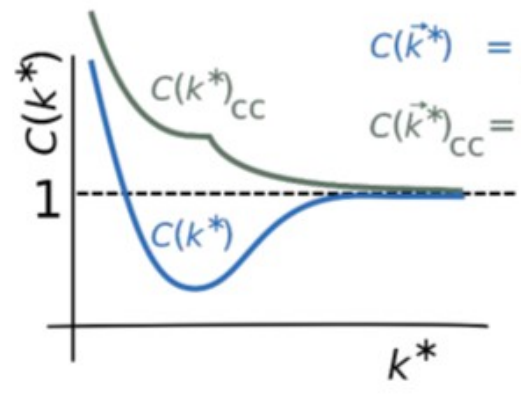
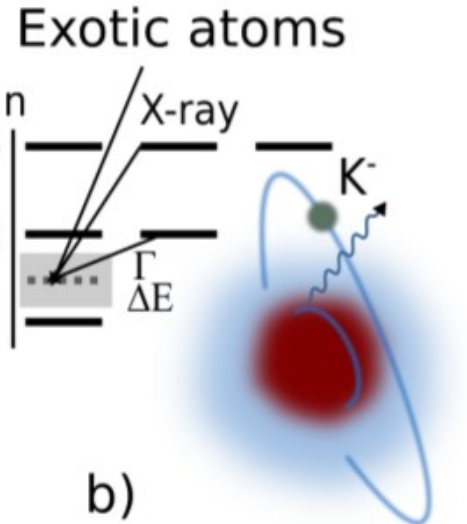
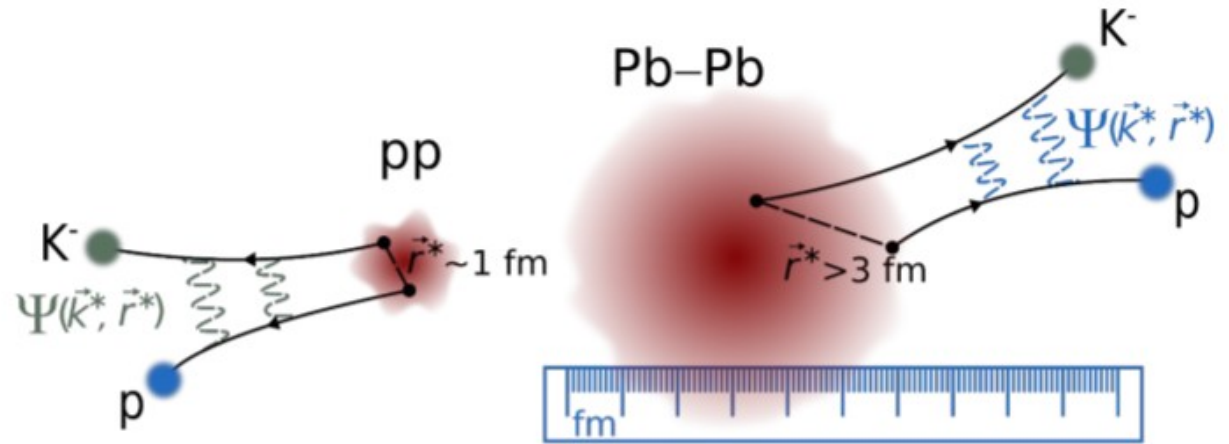
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Experimental study of Kp interaction



Femtoscscopy



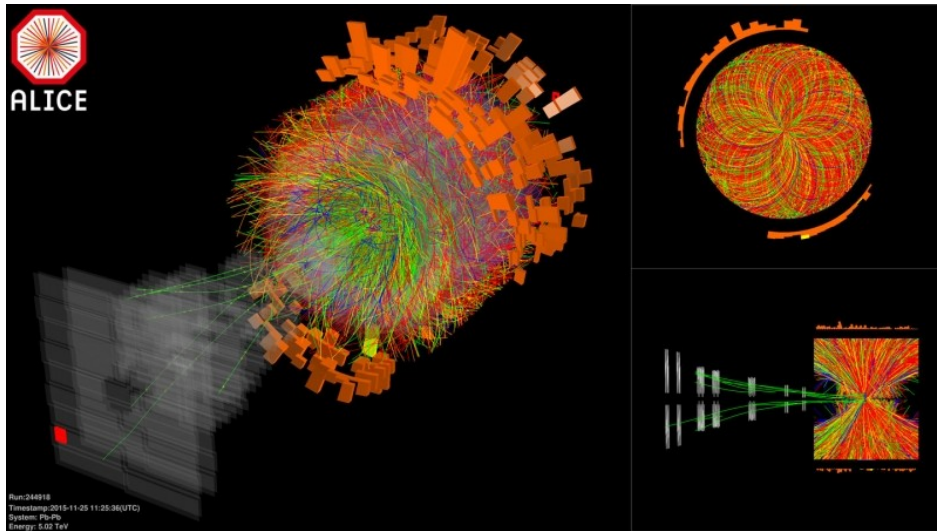
$$C(\vec{k}^*) = \int s(\vec{r}^*) |\Psi(\vec{k}^*, \vec{r}^*)|^2 d^3 r^*$$

$$C(\vec{k}^*)_{cc} = C(\vec{k}^*) + \sum_j \omega_j \int s_j(\vec{r}^*) |\Psi_j(\vec{k}^*, \vec{r}^*)|^2 d^3 r^*$$

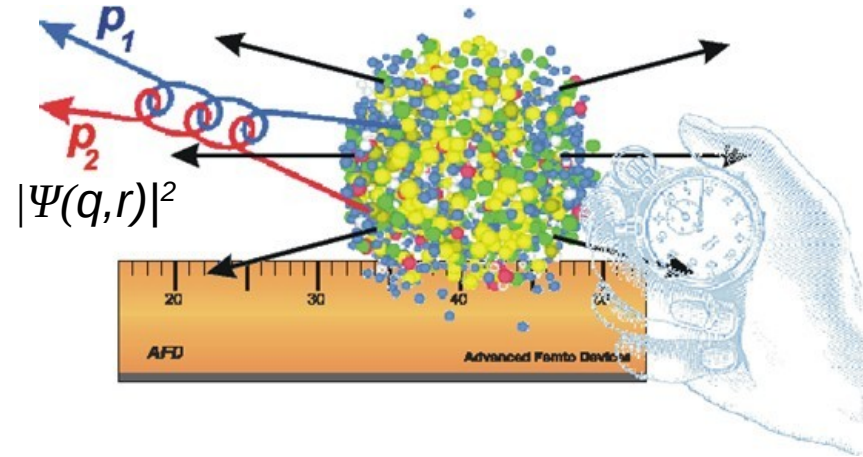
c)



Femtoscscopy technique



from M. Lisa and S. Pratt



- Femtoscscopy – measures space-time characteristics of the source using particle correlations in momentum space

$$C(\vec{p}_1, \vec{p}_2) = \frac{P_{12}(\vec{p}_1, \vec{p}_2)}{P_1(\vec{p}_1)P_2(\vec{p}_2)}$$

$$\vec{q} = \vec{p}_1 - \vec{p}_2$$

$$\vec{r} = \vec{x}_1 - \vec{x}_2$$

experiment

$$C(\vec{q}) = \frac{A(\vec{q})}{B(\vec{q})}$$

$A(\vec{q})$ - correlated pairs (“same events”)

$B(\vec{q})$ - uncorrelated pairs (“mixed events”)

theory (models)

$$C(\vec{q}) = \frac{\int d^3 r S_{12}(\vec{q}, \vec{r}) |\Psi(\vec{q}, \vec{r})|}{\int d^3 x_1 S_1(\vec{x}_1, \vec{p}_1) \int d^3 x_2 S_2(\vec{x}_2, \vec{p}_2)}$$

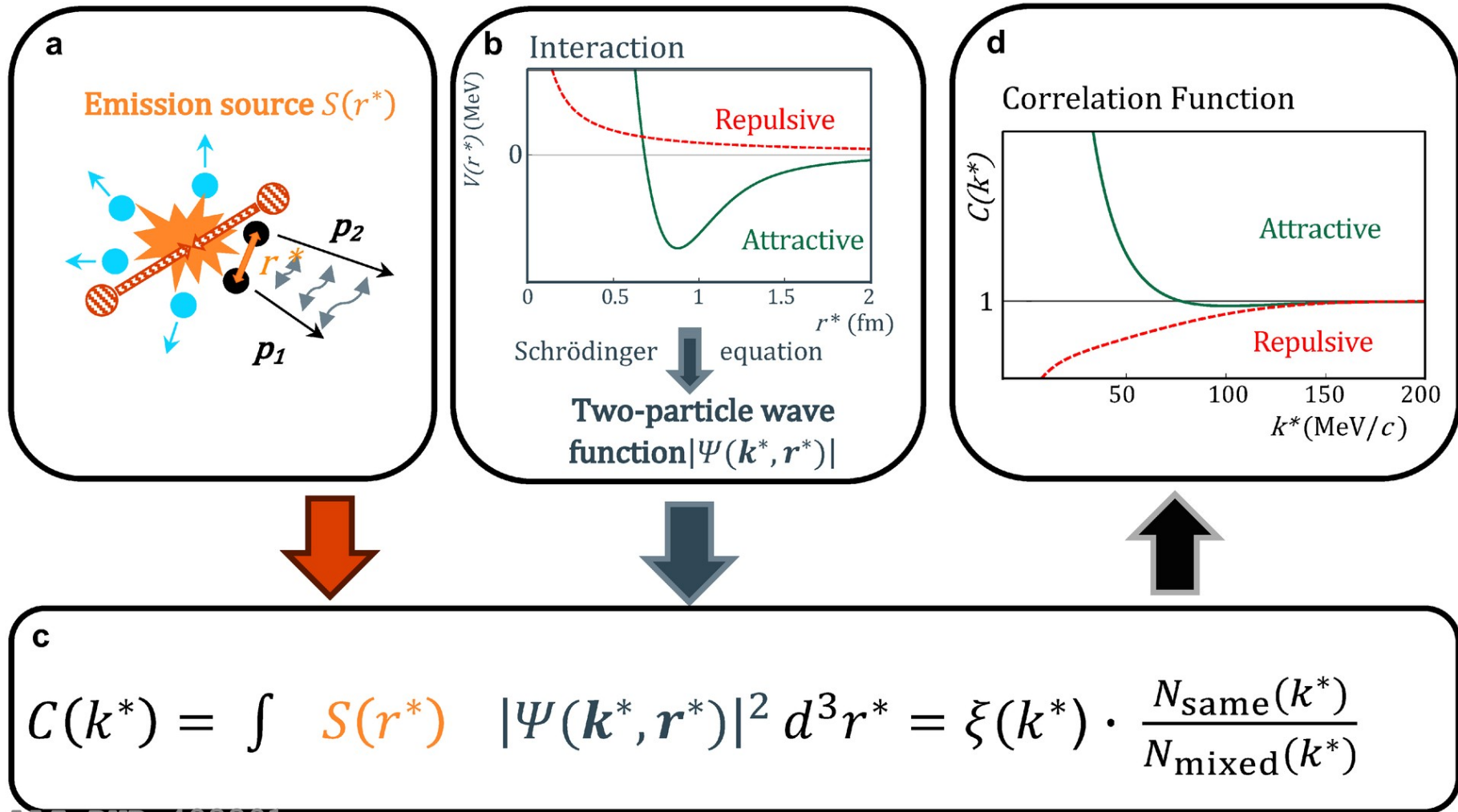
$$C(\vec{q}) = \int d^3 r S(\vec{q}, \vec{r}) |\Psi(\vec{q}, \vec{r})|$$



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Strong interaction from femtoscopy



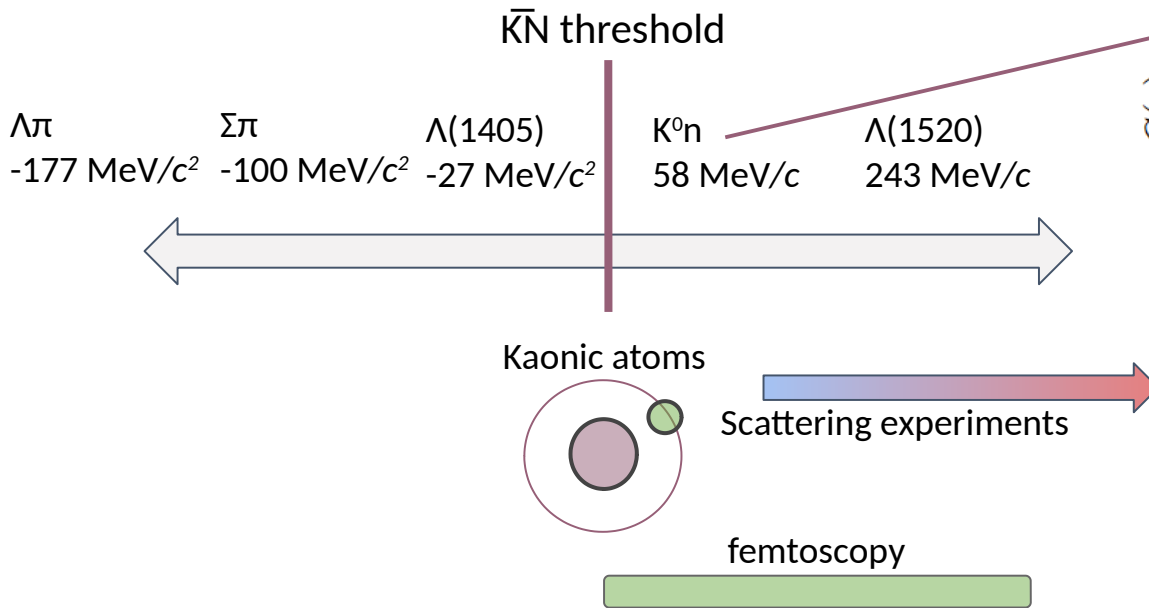
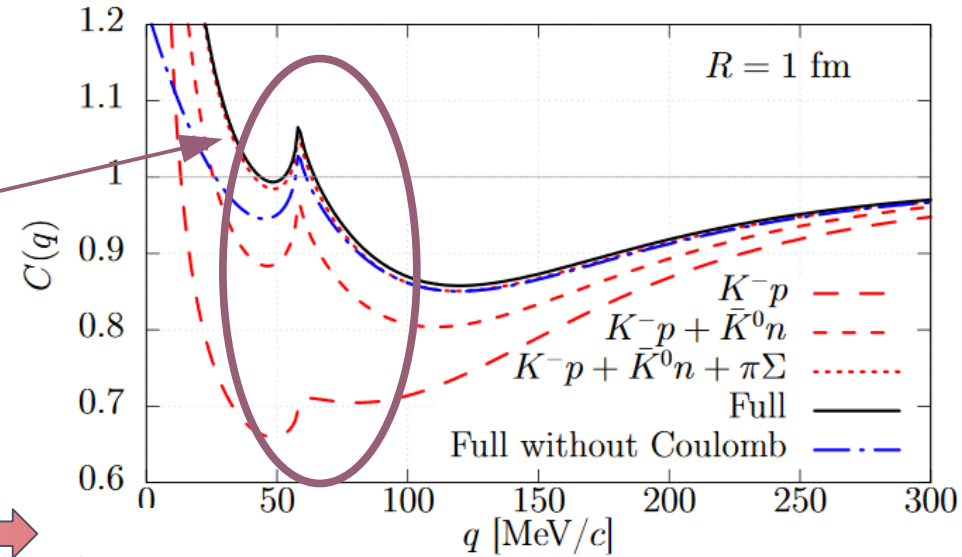
ALI-PUB-483391

Kyoto model with coupled channels



PRL 124 (2020) 132501

$$C(\mathbf{q}) = \int d^3r \sum_j \omega_j S_j(\mathbf{r}) |\Psi_j^{(-)}(\mathbf{q}, \mathbf{r})|^2$$



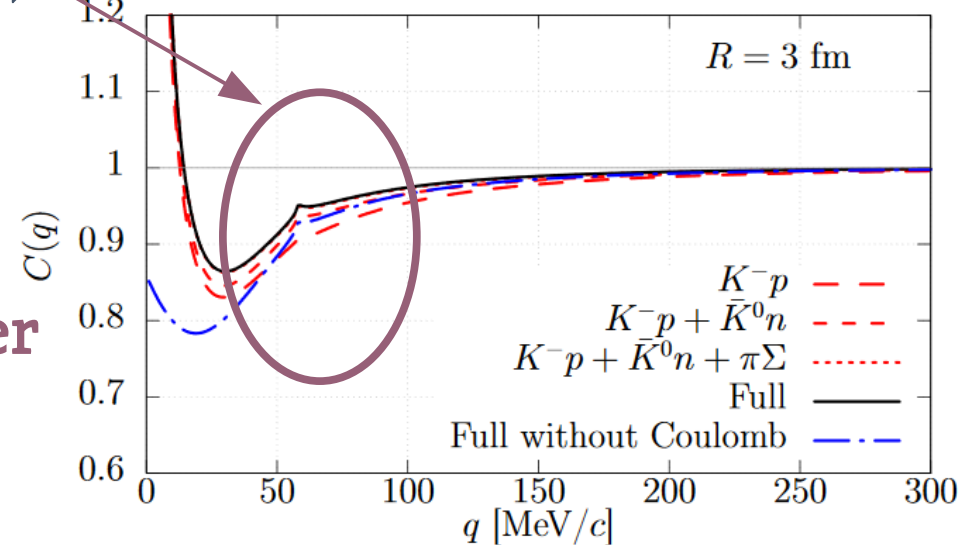
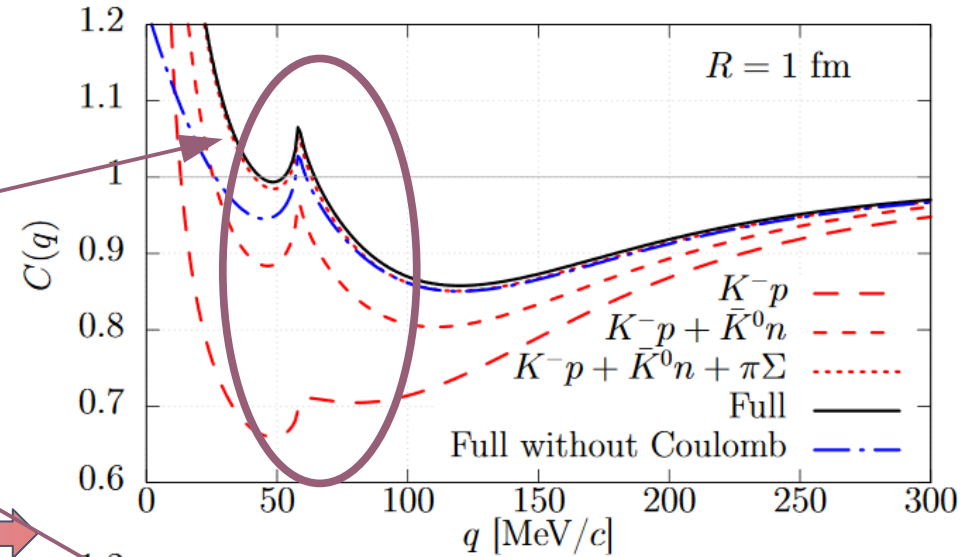
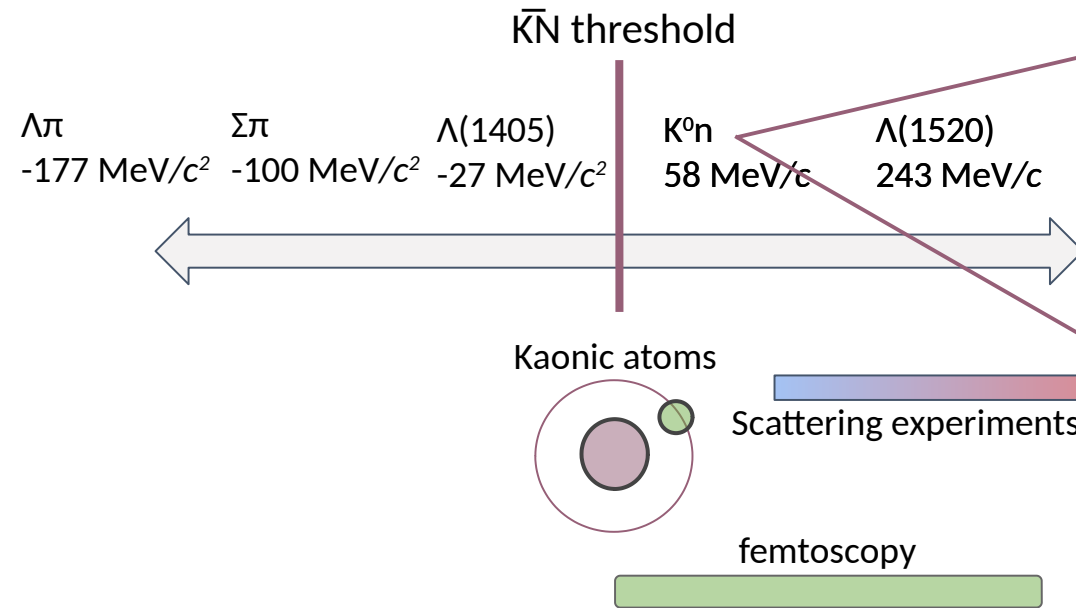
- Clear cusp structure visible for small radii

Kyoto model with coupled channels



$$C(\mathbf{q}) = \int d^3r \sum_j \omega_j S_j(\mathbf{r}) |\Psi_j^{(-)}(\mathbf{q}, \mathbf{r})|^2$$

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- Cusp structure fading out for larger radii

Lednický & Lyuboshitz model

$$C(q) = \int S(r) |\Psi(q, r)|^2 d^3 r \quad q = 2 \cdot k^* = p_1 - p_2$$

measured correlation emission function (source size/shape) pair wave function (includes cross section)

pair wave function \longrightarrow $\Psi_{-\vec{k}^*}^{(+)}(\vec{r}^*) = \sqrt{A_C(\varepsilon)} \frac{1}{\sqrt{2}} \left[e^{-i\vec{k}^* \times \vec{r}^*} F(-i\varepsilon, 1, i\zeta^+) + f_C(\vec{k}^*) \frac{\tilde{G}(\rho, \varepsilon)}{r^*} \right]$

scattering amplitude \longrightarrow $f_C^{-1}(k^*) = \frac{1}{f_0} + \frac{1}{2} d_0 k^{*2} - \frac{2}{a_C} h(k^* a_C) - i k^* a_C$

\tilde{G} is the combination of the regular and singular s-wave Coulomb functions $h(\varepsilon) = \varepsilon^2 \sum_{n=1}^{\infty} [n(n^2 + \varepsilon^2)]^{-1} - \gamma - \ln|\varepsilon|$ ($\gamma = 0.5772$ is the Euler constant)
 F is the confluent hypergeometric function A_C is the Gamow factor, $\zeta^{\pm} = k^* r^* (1 \pm \cos \theta^*)$, $\varepsilon = 1/(k^* a_C)$
 Bohr radius of the pair ($a_C = -83.59$ fm for $K^- p$ pairs)

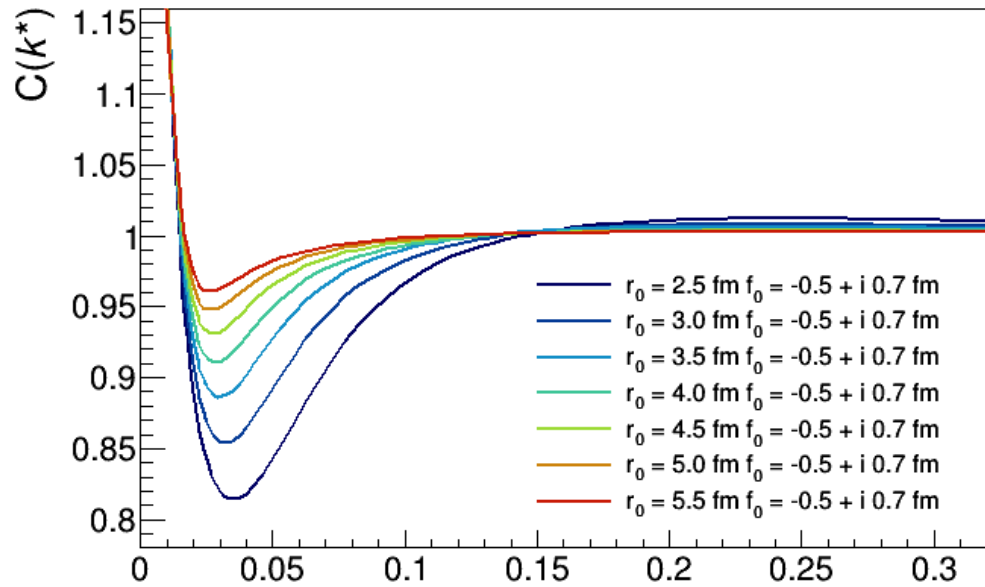
- Numerically solvable (Coulomb + strong)

• The correlation function is characterized by **three parameters**:

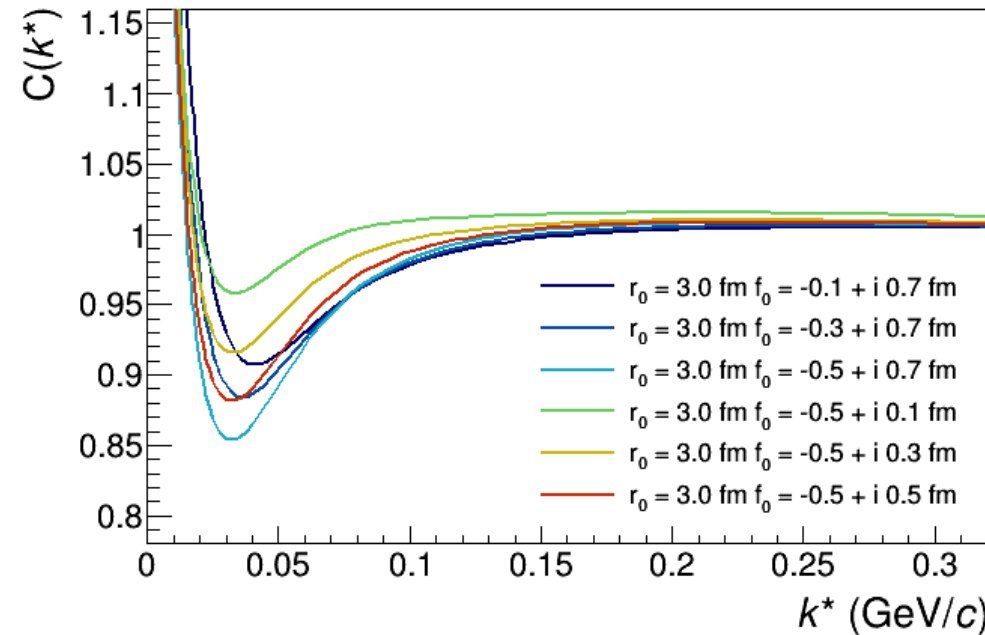
- radius R , scattering length f_0 , and effective radius d_0
- $d_0=0$ effective range approximation
- cross section σ (at low k^*) is simply: $\sigma = 4 \pi |f|^2$



Lednický & Lyuboshitz model



Change with R
Constant f_0



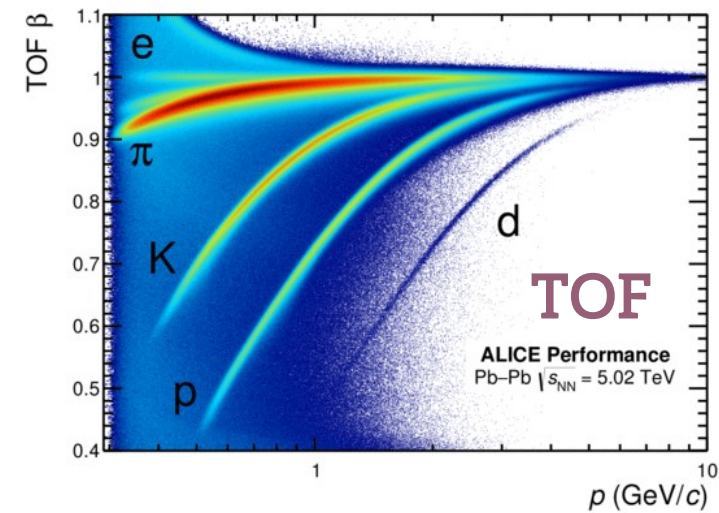
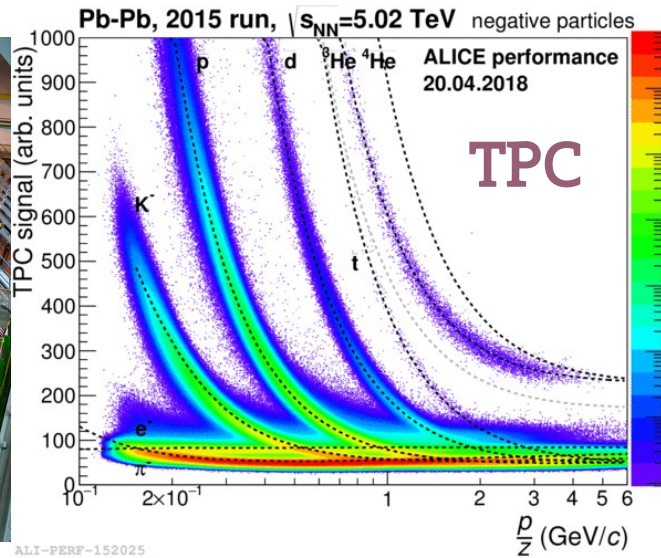
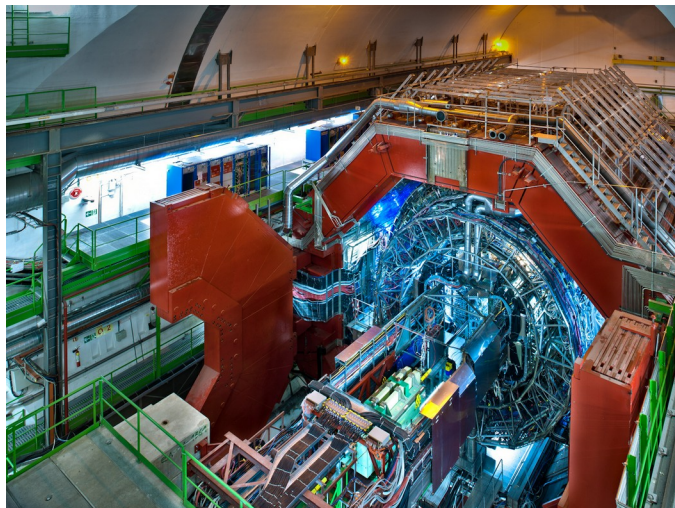
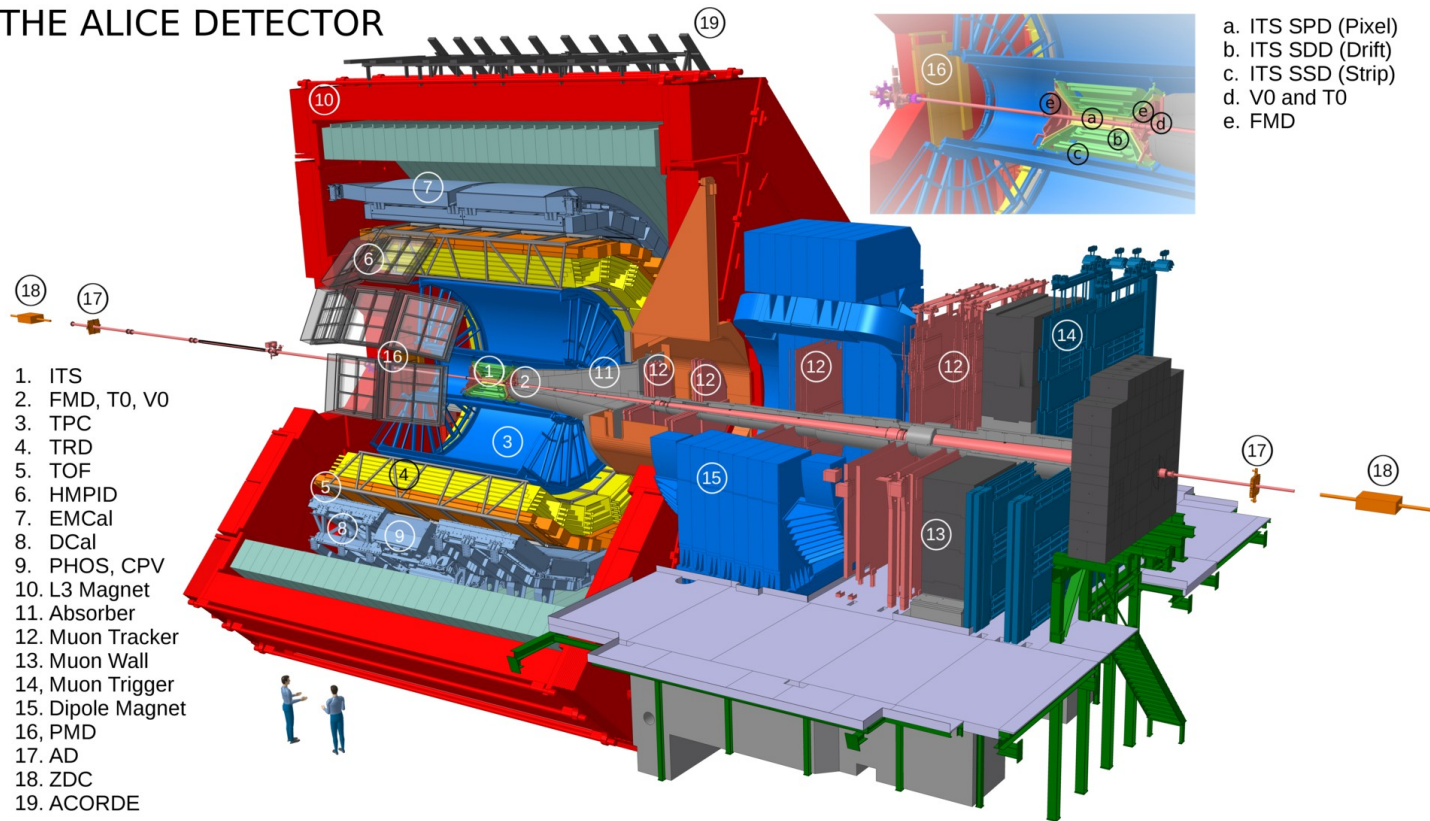
K^-p

Change with f_0
Constant R

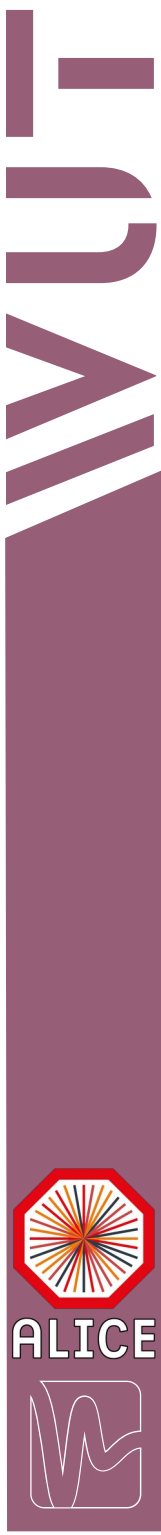
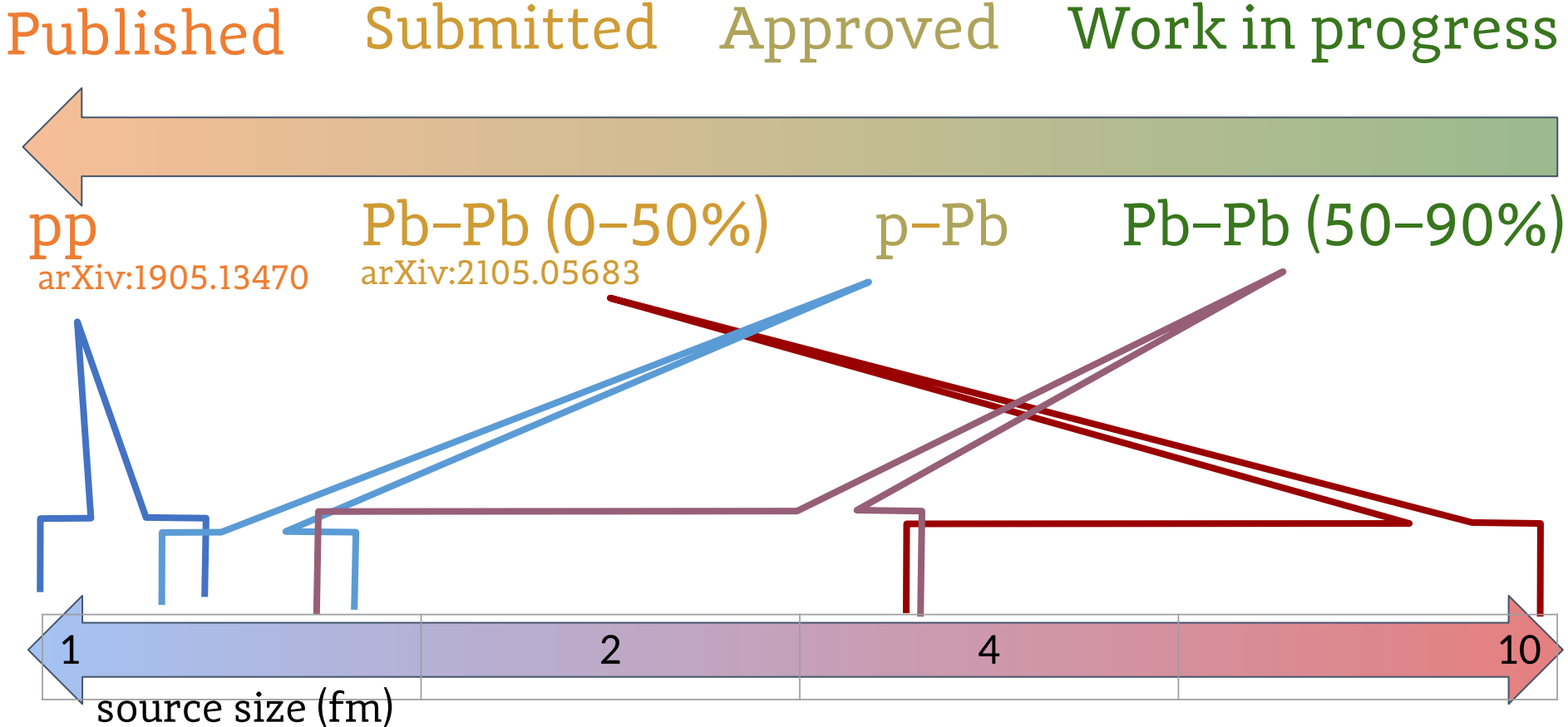


The ALICE experiment

THE ALICE DETECTOR



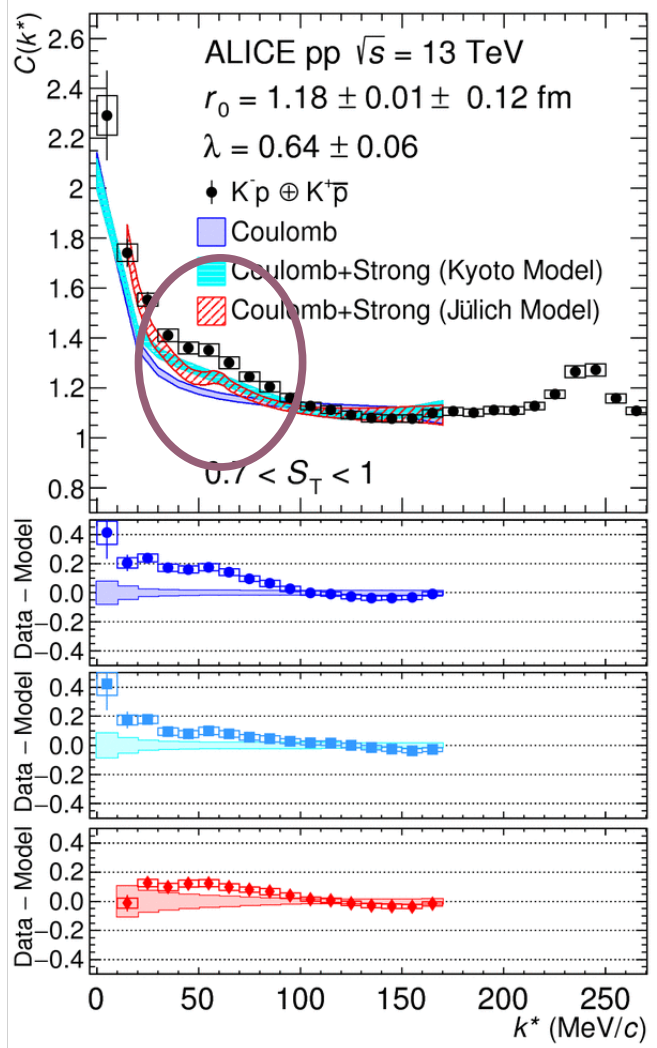
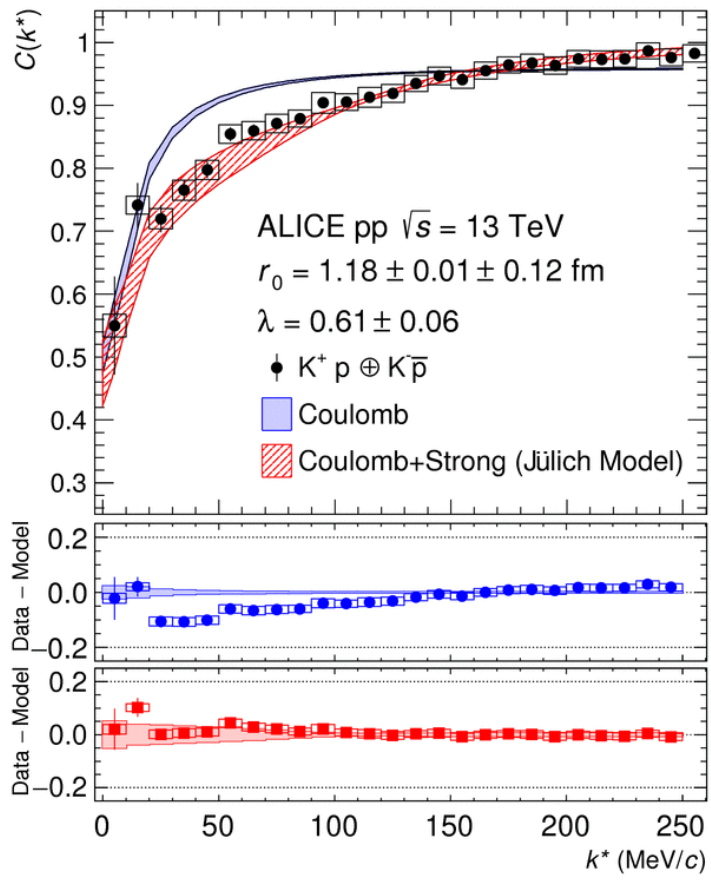
Status of K^-p femtoscopy in ALICE



Proton-kaon correlations in pp

K^+p

K^-p



- First experimental evidence for the opening of the K^0n isospin breaking channel.
- Constraints for low-energy QCD chiral models.

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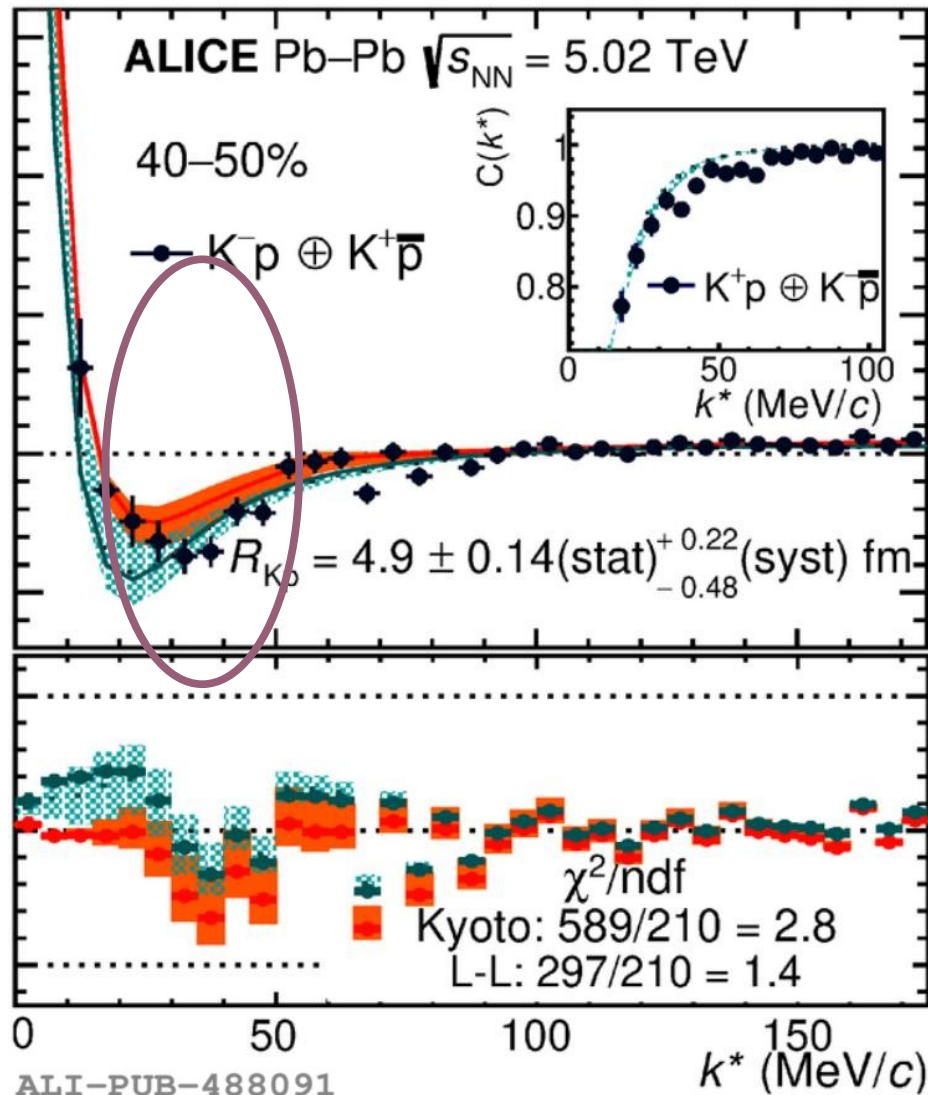
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Proton-kaon correlations in Pb–Pb

<https://arxiv.org/abs/2105.05683>

NEW!

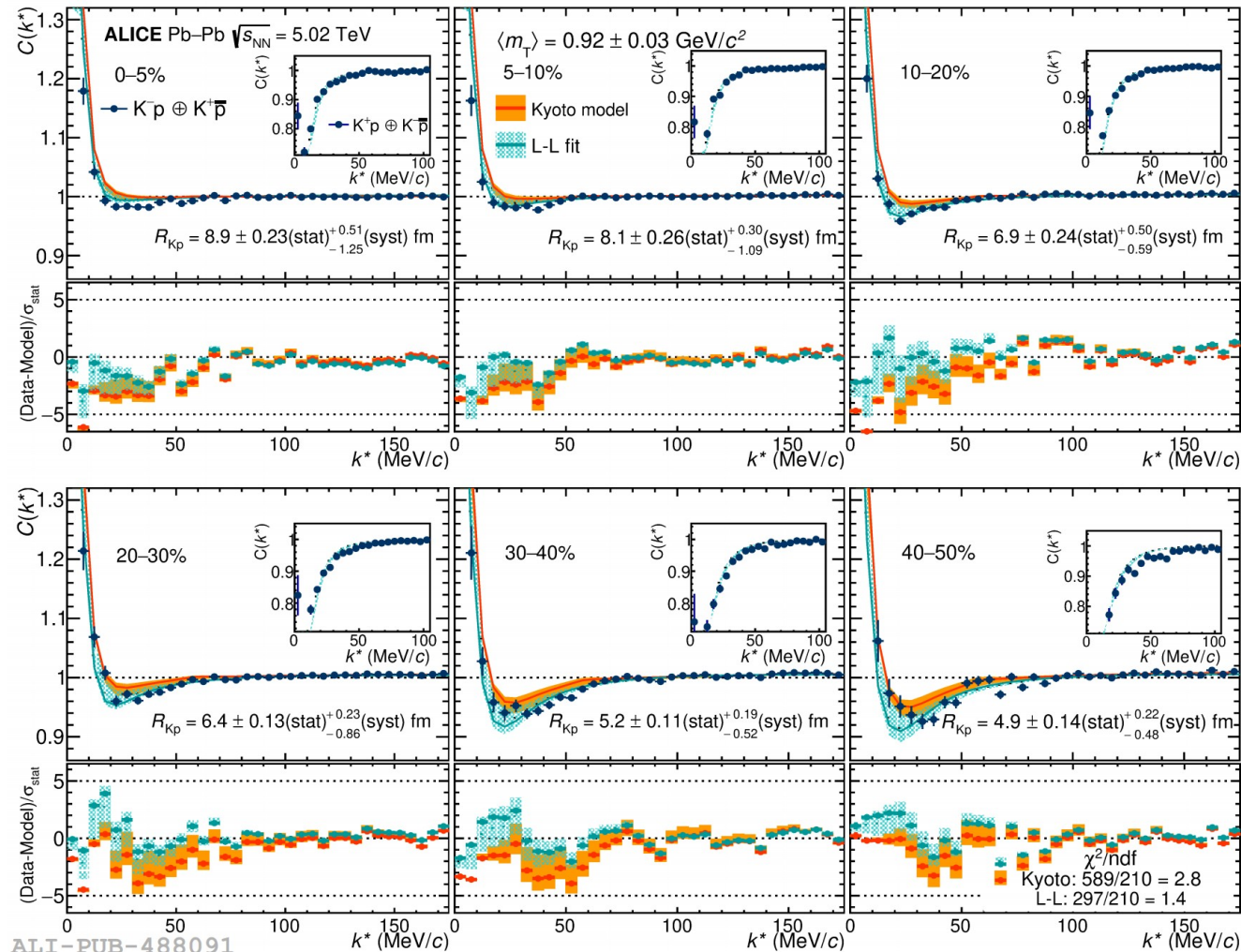


- No K^0n structure visible
- Radii constrained from K^+p pairs

Proton-kaon correlations in Pb-Pb

<https://arxiv.org/abs/2105.05683>

NEW!



- No K^0n structure visible
- Radii constrained from K^+p pairs
- Simultaneous description (and fit) of the CFs for 6 centralities with two parameters and 6 radii



ALICE



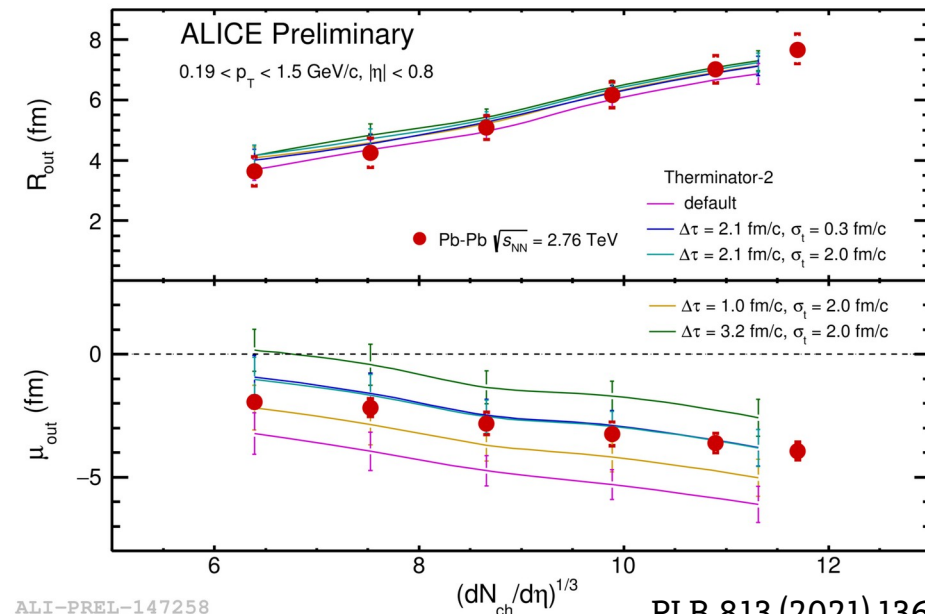
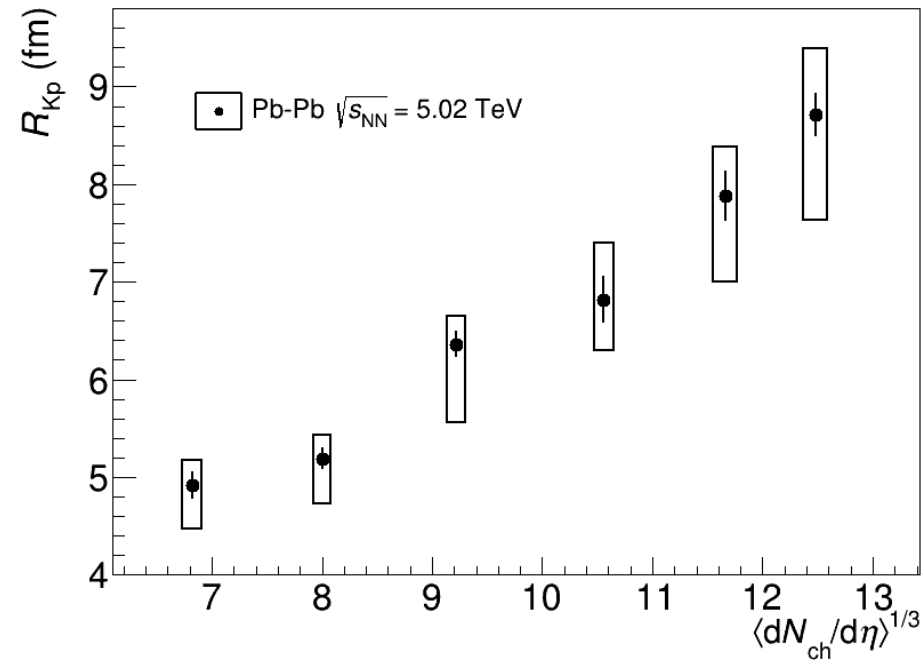
Proton-kaon correlations in Pb-Pb



Kp

- Radii well constrained from scalings
- Radii follow scaling with cube root of charged particle multiplicity density and $\langle m_T \rangle$ as predicted by hydrodynamics calculations

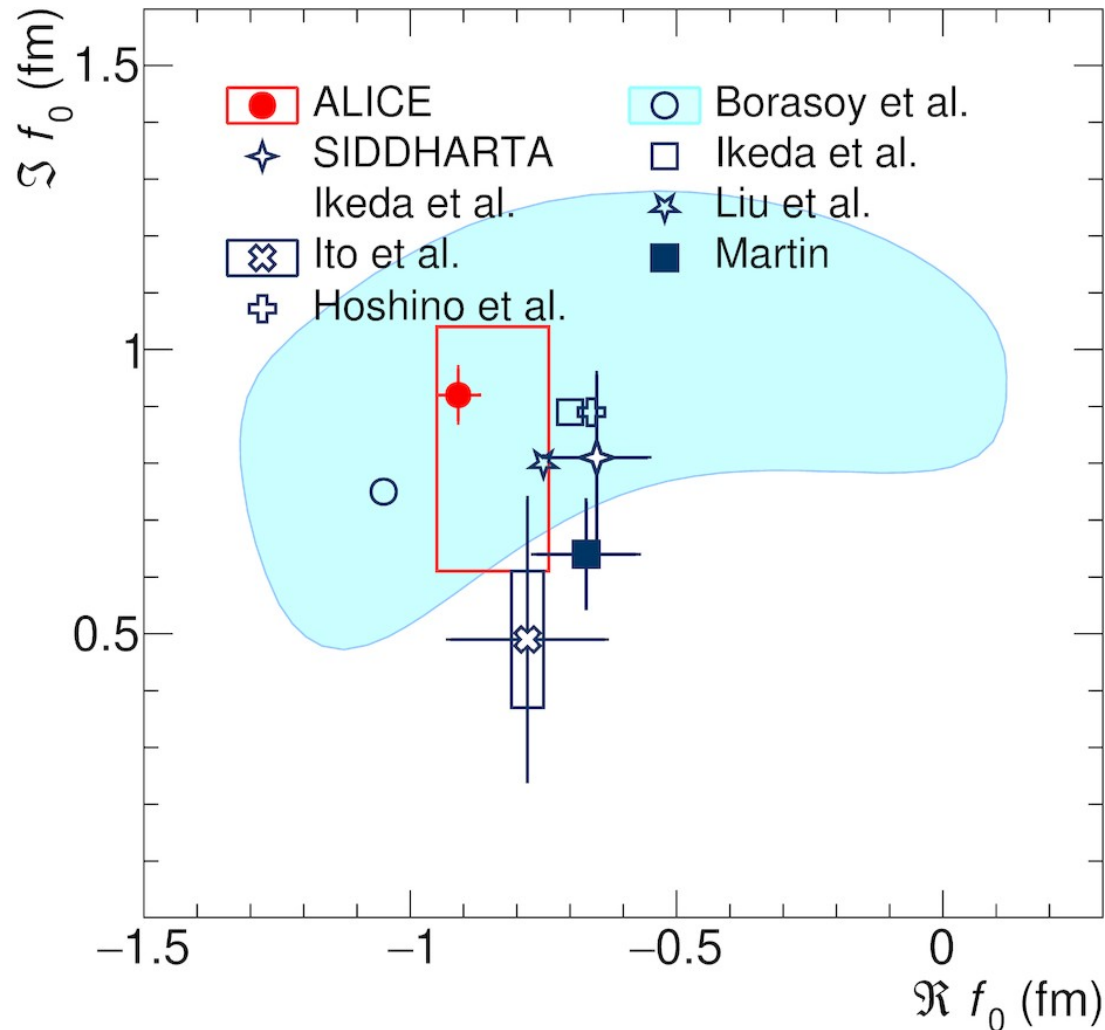
πK



K^-p scattering length

<https://arxiv.org/abs/2105.05683>

NEW!



- $\text{Re } f_0$ and $\text{Im } f_0$ in agreement with available data and calculations!
- Complementary to exotic atoms and scattering experiments



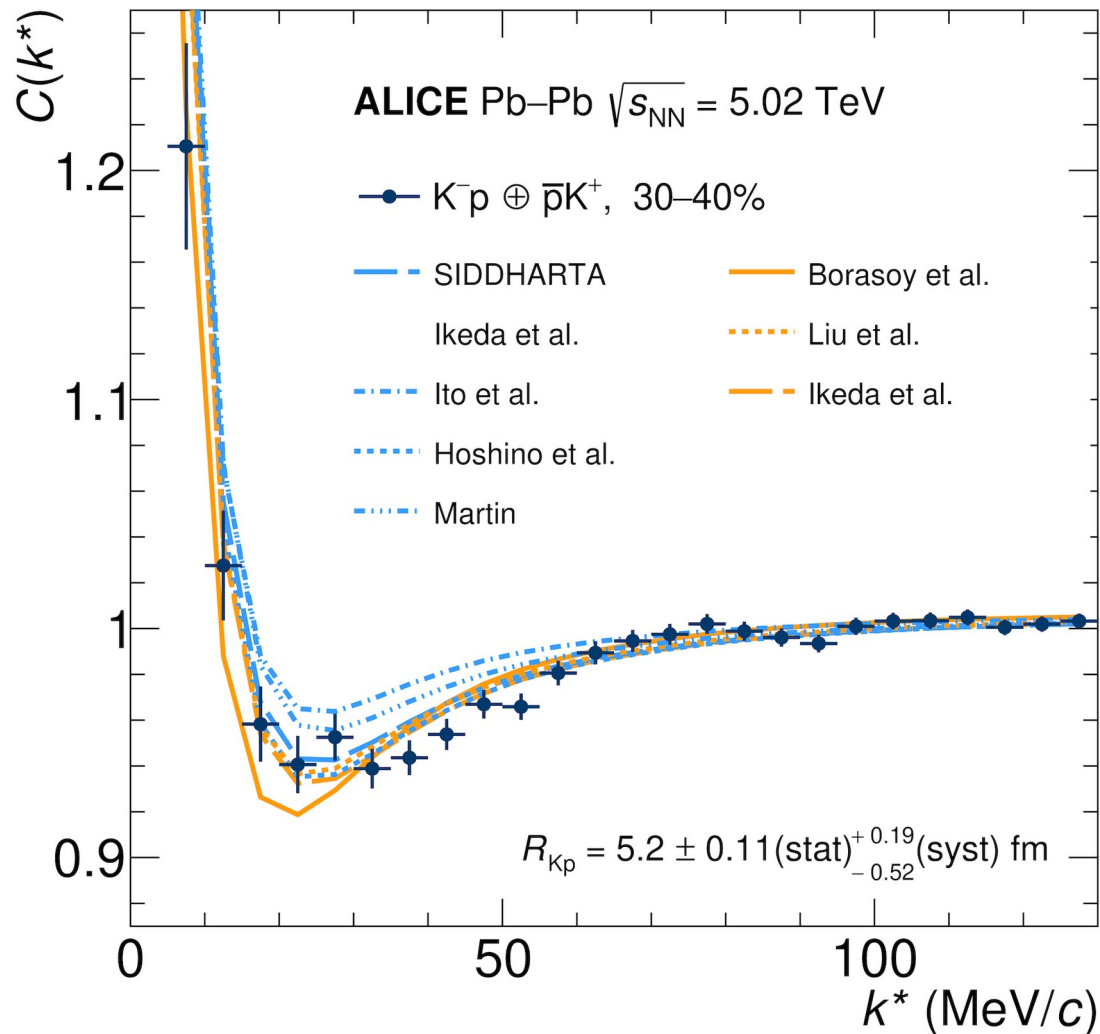
ALICE



K⁻p scattering length

<https://arxiv.org/abs/2105.05683>

NEW!



- $\text{Re } f_0$ and $\text{Im } f_0$ in agreement with available data and calculations!
- Complementary to exotic atoms and scattering experiments
- Plotted with L&L model for comparison



ALICE



K⁻p scattering length

<https://arxiv.org/abs/2105.05683>

NEW!

Model calculation:	$\Re f_0$ (fm)	$\Im f_0$ (fm)	χ^2/ndf
Lednický–Lyuboshitz fit to data	$-0.91 \pm 0.03(\text{stat})_{-0.03}^{+0.17}(\text{syst})$	$0.92 \pm 0.05(\text{stat})_{-0.33}^{+0.12}(\text{syst})$	1.4
Kyoto [39, 76]	–	–	2.8
Lednický–Lyuboshitz with fixed parameters from:			
Kaonic deuterium (Hoshino et al.) [74]	–0.66	0.89	2.0
Scattering experiments (Martin) [71]	-0.67 ± 0.1	0.64 ± 0.1	3.3
Chiral SU(3) (Ikeda et al.) [17, 18]	–0.7	0.89	1.9
SIDDHARTA chiral SU(3) [17, 18]	-0.65 ± 0.1	0.81 ± 0.15	2.3
Hamiltonian EFT (Liu et al.) [73]	–0.75	0.80	1.9
Kaonic hydrogen (Ito et al.) [72]	-0.78 ± 0.15	0.49 ± 0.25	4.2
Chiral SU(3) (Borasoy et al.) [75]	-1.05 ± 0.5	0.75 ± 0.4	1.6

- Chi-square values show Kyoto model is slightly worse than L&L calculations, requiring further developments



Summary

- ALICE has measured K^+p correlations in various collision systems (pp, Pb–Pb) and source sizes
- At very small ranges (1 fm in pp) the coupled channel effects are dominant
- For larger sources neither the cusp from $K^0n \rightarrow K^-p$ isospin breaking channel nor enhancement due to below threshold coupled-channel effects are visible, providing access to the $K^-p \rightarrow K^-p$ process
- Radii are constrained with K^+p pairs and they follow hydrodynamic scalings



ALICE



THANK YOU!



Teotihuacán
ISMD 2017