



Are the $Z_{CS}(3985)$ and $Z_{CS}(4000)$ the same state?

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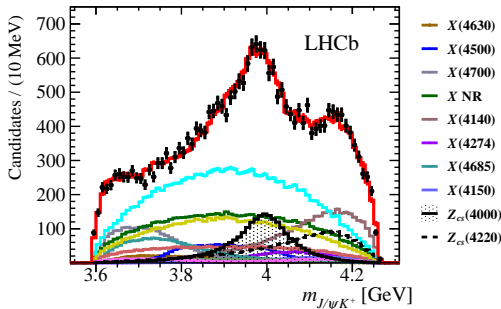
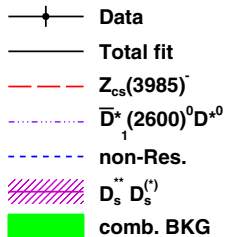
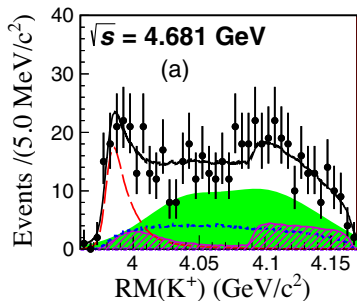
Together with B. Wang (HBU), G.J. Wang(JAEA) and S.L. Zhu (PKU)

Based on [arXiv:2104.08469](https://arxiv.org/abs/2104.08469), [PRD103, L021501](#), [PRD102,111502\(R\)](#)

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Experimental results from BESIII and LHCb



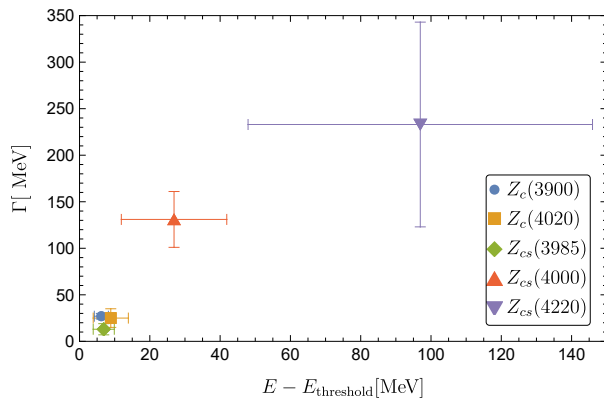
- BESIII: $e^+e^- \rightarrow K^+(D_s^- D^{*0} + D_s^{*-} D^0)$ at $\sqrt{s} = 4.681$ GeV
- LHCb: $Z_{cs}(4000)^+$ and $Z_{cs}(4220)^+$ in the $J/\psi K^+$ of the $B^+ \rightarrow J/\psi \phi K^+$

2103.01803

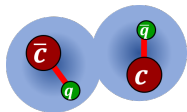
$(M, \Gamma) : (3982.5, 12.8)$ MeV VS $(4003, 131)$ MeV

- Width difference: 10 times. **Whether they are the same states?**
- Molecular scheme: symmetries (and their possible breaking)+ effective field theory (EFT)

Theoretical interpretation of $Z_{CS}(3985)$ state



Composed of two color singlet hadrons
(during most of its lifetime)



- Near-threshold resonances

$Z_c(3900)$	$D^* \bar{D} / D \bar{D}^*$
$Z_c(4020)$	$D^* \bar{D}^*$
$Z_{cs}(3985)$	$D_s \bar{D}^* / D_s^* \bar{D}$

- A natural explanation:

$\Rightarrow Z_{cs}(3985)$: SU(3) partner of $Z_c(3900)$

PRD103,074029; 2011.09404; 2011.10495; 2011.09225...

$\Rightarrow Z_c, Z'_c, Z_{cs}$: molecular resonances

- Hadronic molecule:

PRL67,556; RMP90,015004

\Rightarrow loosely bound states

\Rightarrow near-threshold di-hadron resonances

- $\bar{P}V/\bar{V}P$ di-meson wave functions

$$|C = -\eta\rangle = |G = \eta\rangle = \frac{1}{\sqrt{2}} (|\bar{D}D^*\rangle + \eta|\bar{D}^*D\rangle)^{I=1}$$

- $Z_c(3900) : G = +1 \Rightarrow \eta = +1$

$$\hat{G}_U = \hat{C}e^{i\hat{U}_2\pi}, \hat{G}_V = \hat{C}e^{i\hat{V}_2\pi}$$

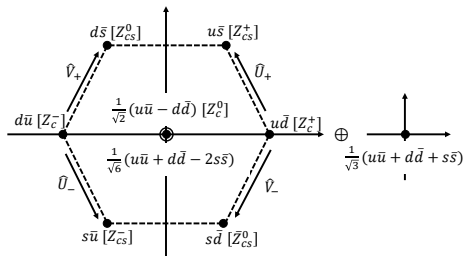
$$|G_U = +1\rangle = \frac{1}{\sqrt{2}} (|D_s^- D^{*+}\rangle + |D_s^{*-} D^+\rangle)$$

- $Z_{cs}(3985)^- : U\text{-spin partner of } Z_c(3900)^-$

- General notation:

$$|\bar{P}V/\bar{V}P, \pm\rangle = \frac{1}{\sqrt{2}} (|\bar{P}V\rangle \pm |\bar{V}P\rangle)$$

- Mixing: odd and even $G_{U/V}$ when ~~SU(3)_f~~



SU(2) subgroups		
$u, d (I)$	$d, s (U)$	$u, s (V)$

Heavy quark spin symmetry (HQSS)

- $V^{\text{spin-space}} = v_1 + v_2 \mathbf{s}_q \cdot \mathbf{s}_{\bar{q}} + \text{HQSS breaking terms}$

⇒ The heavy part: spectator

⇒ Heavy spin and light spin are conserved

- $\langle V \rangle$: HQSS partner states

$$\langle V_{q\bar{q}}^s \rangle_{\{\bar{V}V\}}^{1+} = \langle V_{q\bar{q}}^s \rangle_{\{\bar{P}V/\bar{V}P, +\}}^{1+}$$

$$\langle V_{q\bar{q}}^s \rangle_{\{\bar{V}V\}}^{2+} = \langle V_{q\bar{q}}^s \rangle_{\{\bar{P}V/\bar{V}P, -\}}^{1+}$$

- Rearrangement and selection rules: $|(\bar{c}q_1)_{j_1} (c\bar{q}_2)_{j_2}; J^P\rangle \Rightarrow |(\bar{c}c)_h (\bar{q}_1 q_2)_l; J^P\rangle$

- Example

$$|\bar{P}V/\bar{V}P, +\rangle = \frac{1}{\sqrt{2}} (|0_{\bar{c}c}^-, 1_{q_1\bar{q}_2}^-; 1^+\rangle - |1_{\bar{c}c}^-, 0_{q_1\bar{q}_2}^-; 1^+\rangle)$$

$$Z_c(3900) \sim |\bar{D}D^*/\bar{D}^*D, +\rangle \sim \frac{1}{\sqrt{2}} (|\eta_c\rho\rangle - |J/\psi\pi\rangle)$$

Calculation: Masses and widths

- Coupled-channel Lippmann-Schwinger equations (LSEs): $T = V + VGT$

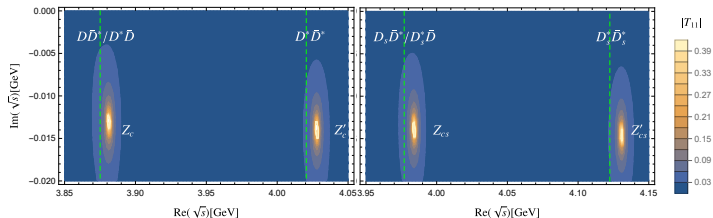
Channel	1	2	3
Z_c/Z'_c	$J/\psi\pi$	$\bar{D}D^*/\bar{D}^*D$	\bar{D}^*D^*
Z_{cs}/Z'_{cs}	$J/\psi K$	$\bar{D}_sD^*/\bar{D}_s^*D$	$\bar{D}_s^*D^*$

- Interaction

$$V_{ij} = \begin{bmatrix} 0 & v_{12} & v_{12} \\ v_{12} & C_d + \frac{C'_d}{2}(\mathbf{p}^2 + \mathbf{p}'^2) & v_{23} \\ v_{12} & v_{23} & C_d + \frac{C'_d}{2}(\mathbf{p}^2 + \mathbf{p}'^2) \end{bmatrix}$$

- Masses and widths of $Z_c(3900)$ and $Z_c(4020)$ as input

Calculation: Masses and widths



(M, Γ)	This work	Exp
Z_{CS}	(3984, 27)	(3982.5, 12.8)
Z'_{CS}	(4130, 29)	

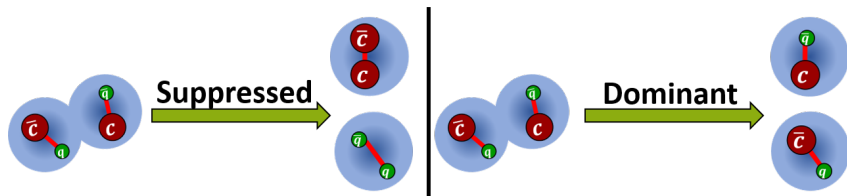
- $Z_{CS}(3985)$: agree with the experiment results well
- Predict $Z_{CS}(4130)$, $SU(3)_F$ partner of $Z_c(4020)$ and HQSS partner of $Z_{CS}(3985)$

PRD102,111502(R)

Calculation: Decays

Z_c/Z_{c_s}	$\frac{\Gamma_2}{\Gamma_1} > 10$	
Z'_c/Z'_{c_s}	$\frac{\Gamma_3}{\Gamma_1} > 10$	$\frac{\Gamma_3}{\Gamma_2} \sim 1$

- $Z_{c(s)} \rightarrow \bar{D}_{(s)} D^* / \bar{D}_{(s)}^* D$ is dominant: fall apart
- $Z_{c(s)} \rightarrow J/\psi \pi(K)$ is suppressed: recluster the heavy quarks
- $Z_{c(s)} : \bar{D}_{(s)} D^* / \bar{D}_{(s)}^* D$ interaction is not attractive enough but confine them for a finite time



- Coupled-channel effect from $J/\psi \pi(K)$ is tiny

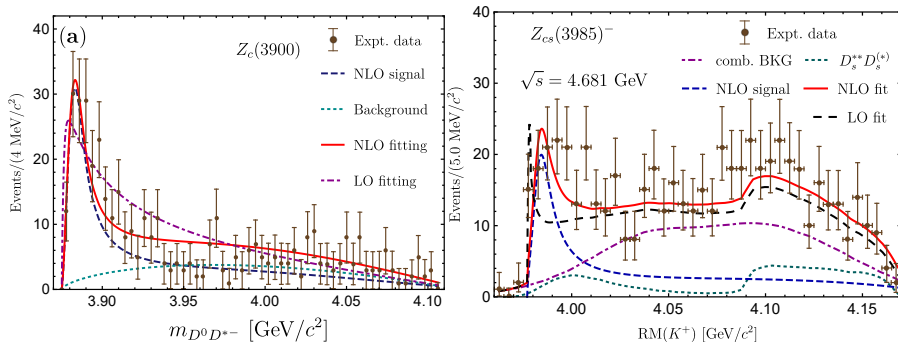
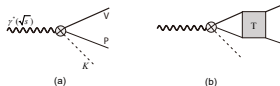
PRD102,111502(R)

- The parameters are extracted from the $Z_c(3900)$ data

PRD102,114019; PRD92,092006

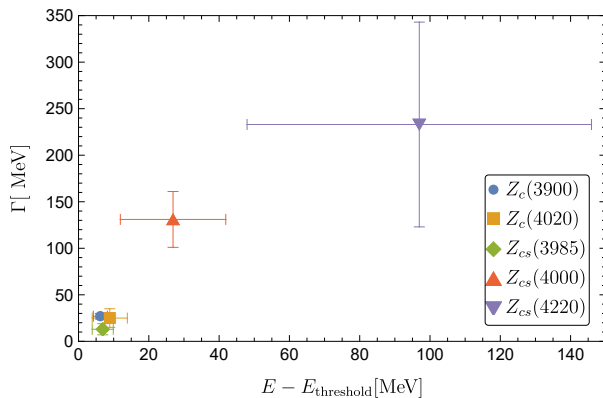
- The event distribution of $Z_{cs}(3985)$ are consistent with the experimental results

PRD103, L021501



Theoretical analysis in two-state scheme

If the $Z_{cs}(3985)$ and $Z_{cs}(4000)$ are two different states



- Bad $SU_F(3)$ symmetry
 - \Rightarrow Only one Z_c state with $J^P = 1^+$ near the $\bar{D}D^*/\bar{D}^*D$ threshold
 - \Rightarrow Two Z_{cs} states (narrower one and broader one) with $J^P = 1^+$ near the $\bar{D}_sD^*/\bar{D}_s^*D$ threshold
- Good HQSS: $Z_c(3900)$ and $Z_c(4020)$

broad, $Z_{cs}(4000)$	$\xleftrightarrow{\text{Corresponding?}}$	$ \bar{P}V/\bar{V}P, +\rangle$	$\xleftrightarrow{\text{HQSS}}$	$ \bar{V}V, 1^+\rangle$
narrow, $Z_{cs}(3985)$	$\xleftrightarrow{\text{Mixing?}}$	$ \bar{P}V/\bar{V}P, -\rangle$	$\xleftrightarrow{\text{HQSS}}$	$ \bar{V}V, 2^+\rangle$

- Neglect the $J/\psi\pi$ and $J/\psi K$ channels

Heavy quark spin symmetry violation effect

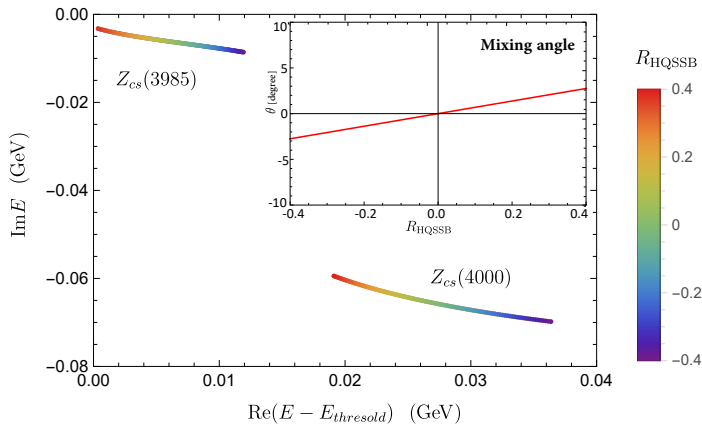
$$V_{\{\bar{P}V, \bar{V}P\}}^{1+} = \frac{1}{2\Lambda} \begin{bmatrix} c_a^+ + c_a^- & c_a^+ - c_a^- \\ c_a^+ - c_a^- & c_a^+ + c_a^- + 4\delta c_a \end{bmatrix} + \text{NLO term}$$

$$V_{\{\bar{P}V/\bar{V}P, +- \}}^{1+} = \begin{bmatrix} \frac{c_a^+ + \delta c_a}{\Lambda} & \frac{\delta c_a}{\Lambda} \\ \frac{\delta c_a}{\Lambda} & \frac{c_a^- + \delta c_a}{\Lambda} \end{bmatrix} + \begin{bmatrix} \frac{c_b^+ (\mathbf{p}^2 + \mathbf{p}'^2)}{\Lambda^3} & \\ & \frac{c_b^- (\mathbf{p}^2 + \mathbf{p}'^2)}{\Lambda^3} \end{bmatrix}$$

- δc_a HQSS breaking effect, inducing the mixing of $|\bar{P}V/\bar{V}P, +\rangle$ and $|\bar{P}V/\bar{V}P, -\rangle$
- 5 LECs, 4 input: masses and widths of $Z_{cs}(3985)$ and $Z_{cs}(4000)$
- Set $\delta c_a = 0$ and determine the other 4 LECs
- Varying δc_a to make $-0.4 < R_{\text{HQSSB}} < 0.4$

$$R_{\text{HQSSB}} = \frac{4\delta c_a}{|c_a^+ + c_a^-|}$$

Mixing effect is not significant



- Relative orders for M and Γ do not change
- Mixing angle of $|PV/VP, +\rangle$ and $|PV/VP, -\rangle$ is tiny

Dimeson components of two states

broad, $Z_{cs}(4000)$	$\xleftrightarrow{\text{Corresponding}}$	$ \bar{P}V/\bar{V}P, +\rangle$	$\xleftrightarrow{\text{HQSS}}$	$ \bar{V}V, 1^+\rangle$	\leftarrow	$Z_{cs}(4220)$, broad state
narrow, $Z_{cs}(3985)$	$\xleftrightarrow{\text{tiny mixing}}$	$ \bar{P}V/\bar{V}P, -\rangle$	$\xleftrightarrow{\text{HQSS}}$	$ \bar{V}V, 2^+\rangle$	\leftarrow	Prediction

- $|Z_{cs}(4000)\rangle = |\bar{D}_s^* D / \bar{D}_s D^*, +\rangle$ and $Z_{cs}(3985) = |\bar{D}_s^* D / \bar{D}_s D^*, -\rangle$

$$\mathcal{R}(Z_{cs} \rightarrow \bar{D}_s^* D / Z_{cs} \rightarrow \bar{D}_s D^*) \approx 0.5$$

- Tensor $\bar{D}_s^* D^*$ state as the HQSS partner of $Z_{cs}(3985)$.

$$M = 4126 \pm 3 \text{ MeV}, \quad \Gamma = 13 \pm 6 \text{ MeV}.$$

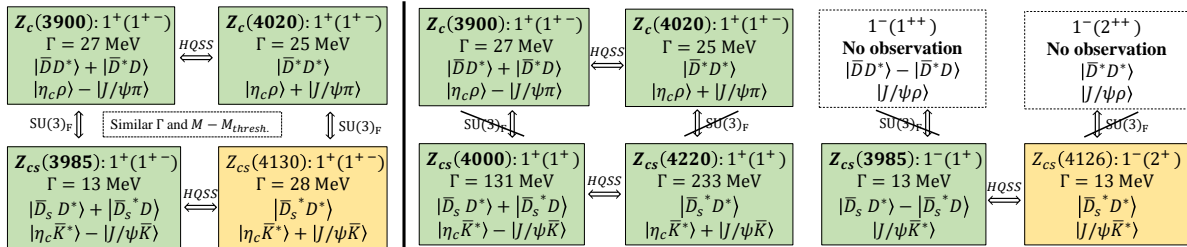
- $Z_{cs}(3985) \rightarrow J/\psi K$ is suppressed compared with $Z_{cs}(4000) \rightarrow J/\psi K$ in the HQSS limit.

$$|\bar{P}V/\bar{V}P, +\rangle = \frac{1}{\sqrt{2}} (|0_{\bar{c}c}^-, 1_{q_1 \bar{q}_2}^-; 1^+\rangle - |1_{\bar{c}c}^-, 0_{q_1 \bar{q}_2}^-; 1^+\rangle)$$

$$|\bar{P}V/\bar{V}P, -\rangle = |1_{\bar{c}c}^-, 1_{q_1 \bar{q}_2}^-; 1^+\rangle$$

Summary

- Based on molecule scheme



(a) $Z_{cs}(3985)$ as the $SU(3)_F$ partner of $Z_c(3900)$.

(b) Implications $Z_{cs}(3985)$ and $Z_{cs}(4000)$ as two different states.

- Take home messages

$\Rightarrow Z_{cs}(3985) \rightarrow J/\psi K$ in suppressed in the two states schemes

\Rightarrow Tensor $Z_{cs}(4126)$ in two states schemes

- Compact tetraquark scheme: predict more states (two nonets)

Thanks for your attention!