



Yuki Kamiya

Institute of Theoretical Physics,
Chinese Academy of Sciences

Femtoscopic study on the $N\Xi-\Lambda\Lambda$ interaction

Kenji Sasaki (CiDER, Osaka Univ.)

Tokuro Fukui (RIKEN)

Tetsuo Hyodo (Tokyo Metropolitan Univ.)

Kenji Morita (RIKEN)

Kazuyuki Ogata (RCNP, Osaka Univ.)

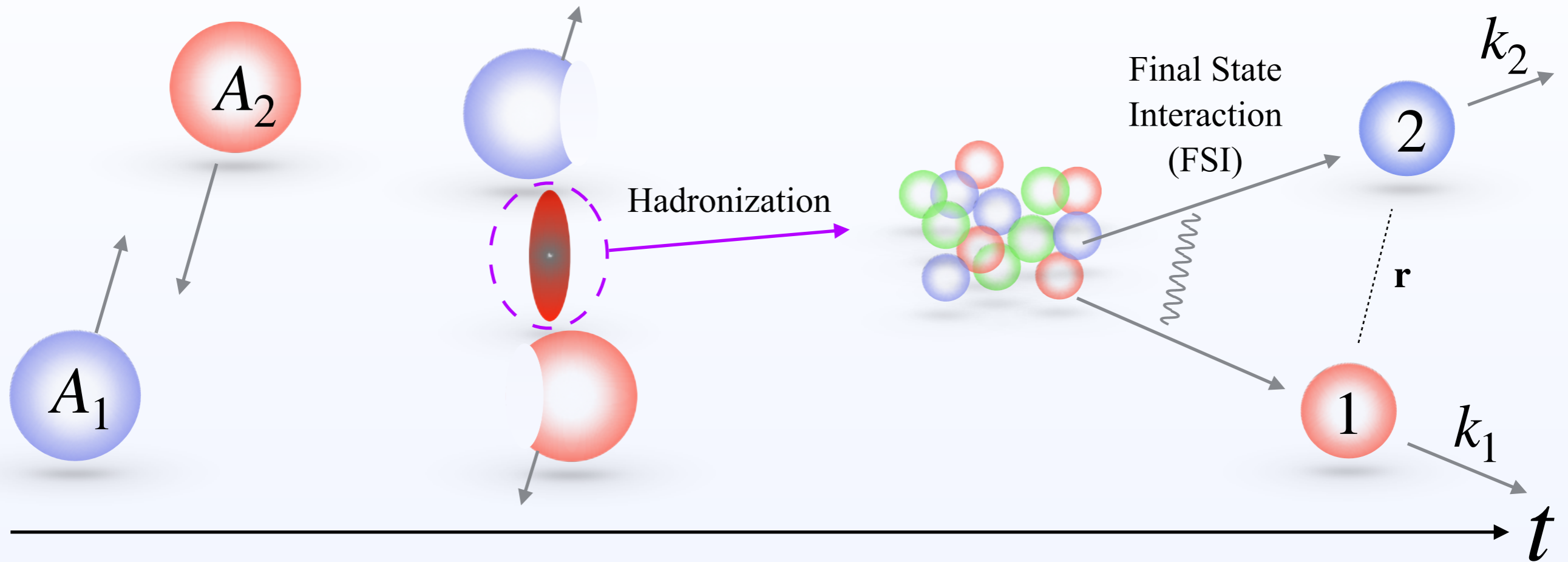
Tetsuo Hatsuda (RIKEN)

Akira Ohnishi (YITP, Kyoto Univ.)

HADRON 2021/07/28,
Zoom conference @ Mexico

Hadron correlation in high energy nuclear collision

- High energy nuclear collision and FSI

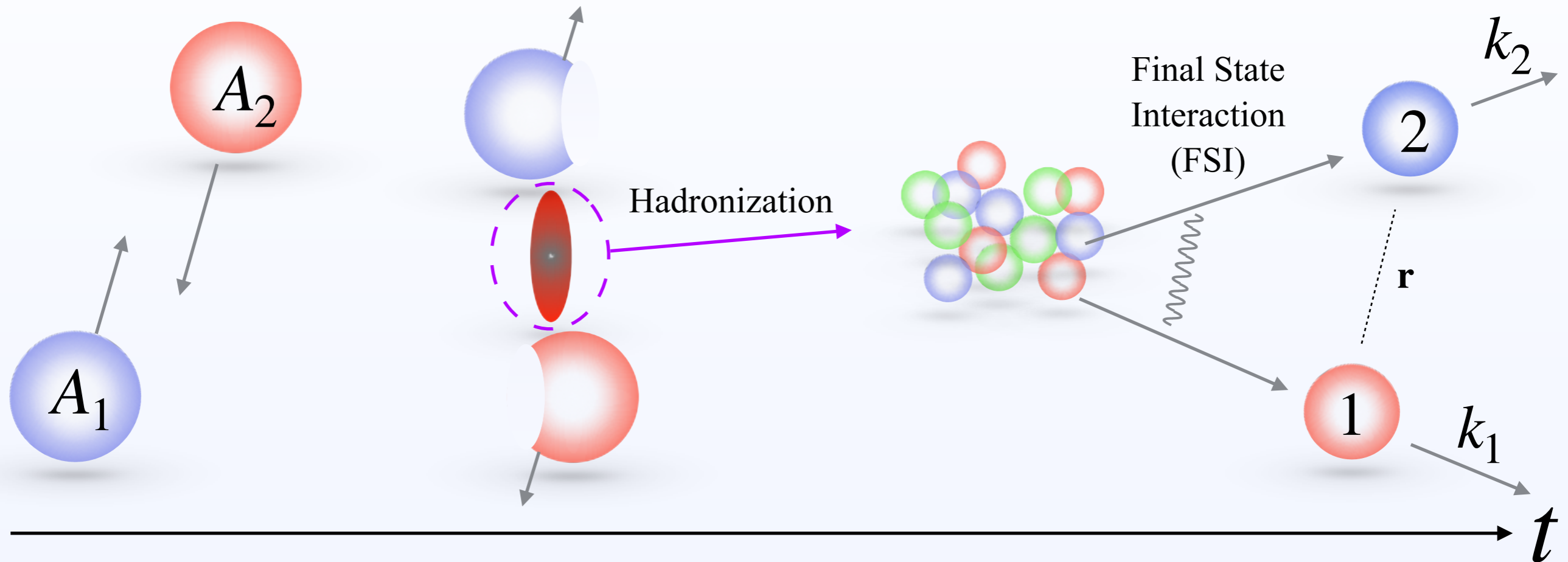


- Hadron-hadron correlation

$$C_{12}(k_1, k_2) = \frac{N_{12}(k_1, k_2)}{N_1(k_1)N_2(k_2)}$$
$$= \begin{cases} 1 & \text{(w/o correlation)} \\ \text{Others (w/ correlation)} \end{cases}$$

Hadron correlation in high energy nuclear collision

- High energy nuclear collision and FSI



- Hadron-hadron correlation

- Koonin-Pratt formula : S.E. Koonin, PLB 70 (1977)
S. Pratt et. al. PRC 42 (1990)

$$C(\mathbf{q}) \simeq \int d^3\mathbf{r} S(\mathbf{r}) |\varphi^{(-)}(\mathbf{q}, \mathbf{r})|^2$$

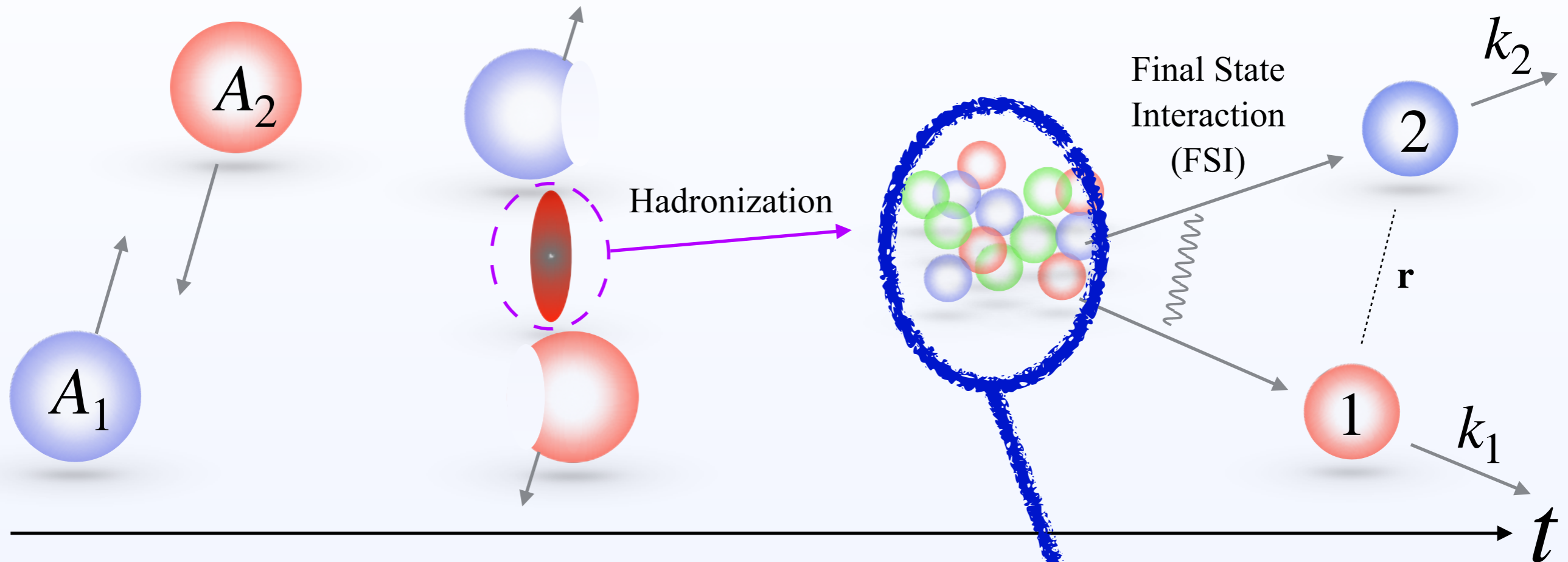
$$\mathbf{q} = (m_2\mathbf{k}_1 - m_1\mathbf{k}_2)/(m_1 + m_2)$$

$S(\mathbf{r})$: Source function

$\varphi^{(-)}(\mathbf{q}, \mathbf{r})$: Relative wave function

Hadron correlation in high energy nuclear collision

High energy nuclear collision and FSI



Hadron-hadron correlation

- Koonin-Pratt formula : S.E. Koonin, PLB 70 (1977)
S. Pratt et. al. PRC 42 (1990)

$$C(\mathbf{q}) \simeq \int d^3\mathbf{r} S(\mathbf{r}) |\varphi^{(-)}(\mathbf{q}, \mathbf{r})|^2$$

$$\mathbf{q} = (m_2\mathbf{k}_1 - m_1\mathbf{k}_2)/(m_1 + m_2)$$

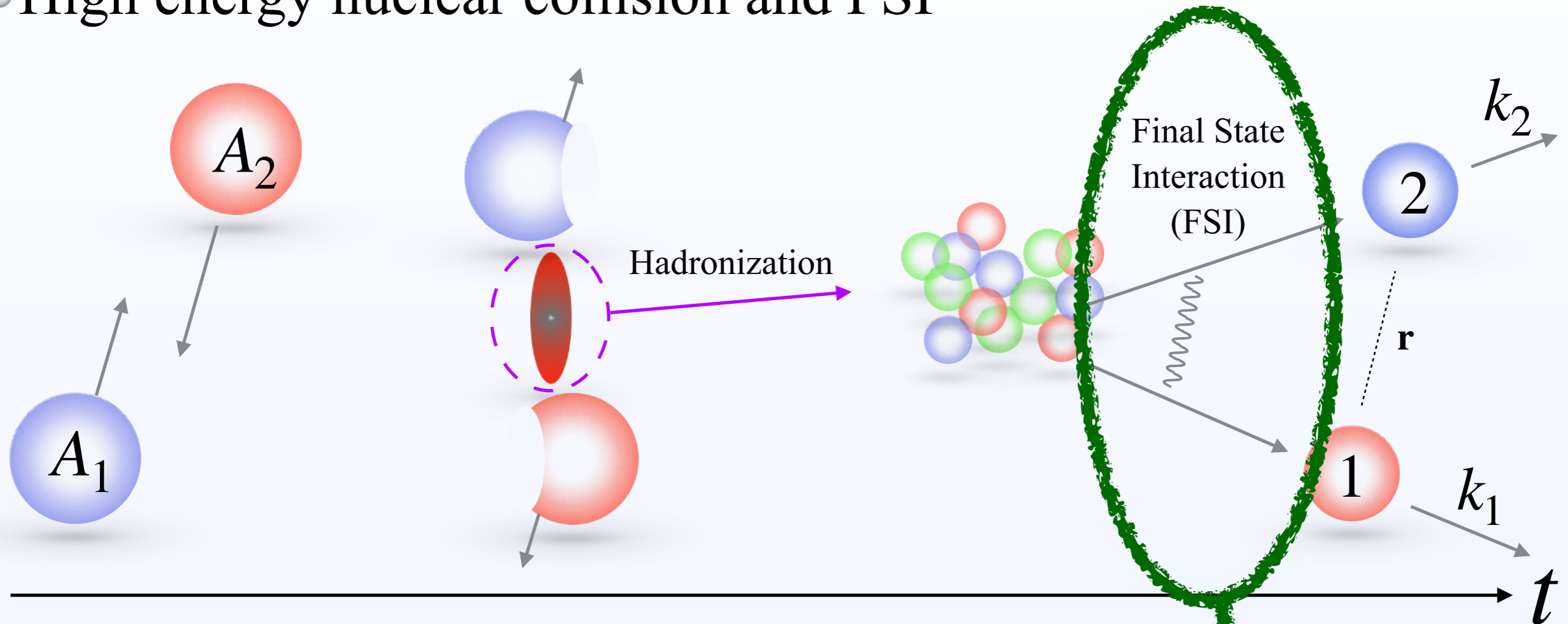
$S(\mathbf{r})$: Source function

$\varphi^{(-)}(\mathbf{q}, \mathbf{r})$: Relative wave function

- Depends on ...
Collision detail (A_i , energy, centrality)
- Including information of...
size of hadron source,
momentum dependence, weight...

Hadron correlation in high energy nuclear collision

High energy nuclear collision and FSI



Hadron-hadron correlation

- Koonin-Pratt formula : S.E. Koonin, PLB 70 (1977)
S. Pratt et. al. PRC 42 (1990)

$$C(\mathbf{q}) \simeq \int d^3\mathbf{r} S(\mathbf{r}) |\varphi^{(-)}(\mathbf{q}, \mathbf{r})|^2$$

$$\mathbf{q} = (m_2\mathbf{k}_1 - m_1\mathbf{k}_2)/(m_1 + m_2)$$

$S(\mathbf{r})$: Source function

$\varphi^{(-)}(\mathbf{q}, \mathbf{r})$: Relative wave function

- Depends on ...

Interaction (strong and Coulomb)

quantum statistics (Fermion, boson)

Hadron correlation in high energy nuclear collision

- Analytic model for ideal cases

$$C(\mathbf{q}) \simeq \int d^3\mathbf{r} S(\mathbf{r}) |\varphi^{(-)}(\mathbf{q}, \mathbf{r})|^2$$

$S(\mathbf{r})$: Source function

$\varphi^{(-)}(\mathbf{q}, \mathbf{r})$: Relative wave function

- Gaussian source with radius R
- Approximate φ by asymptotic wave func.
- $\mathcal{F}(q) = [-1/a_0 - iq]^{-1}$ with scat. length a_0

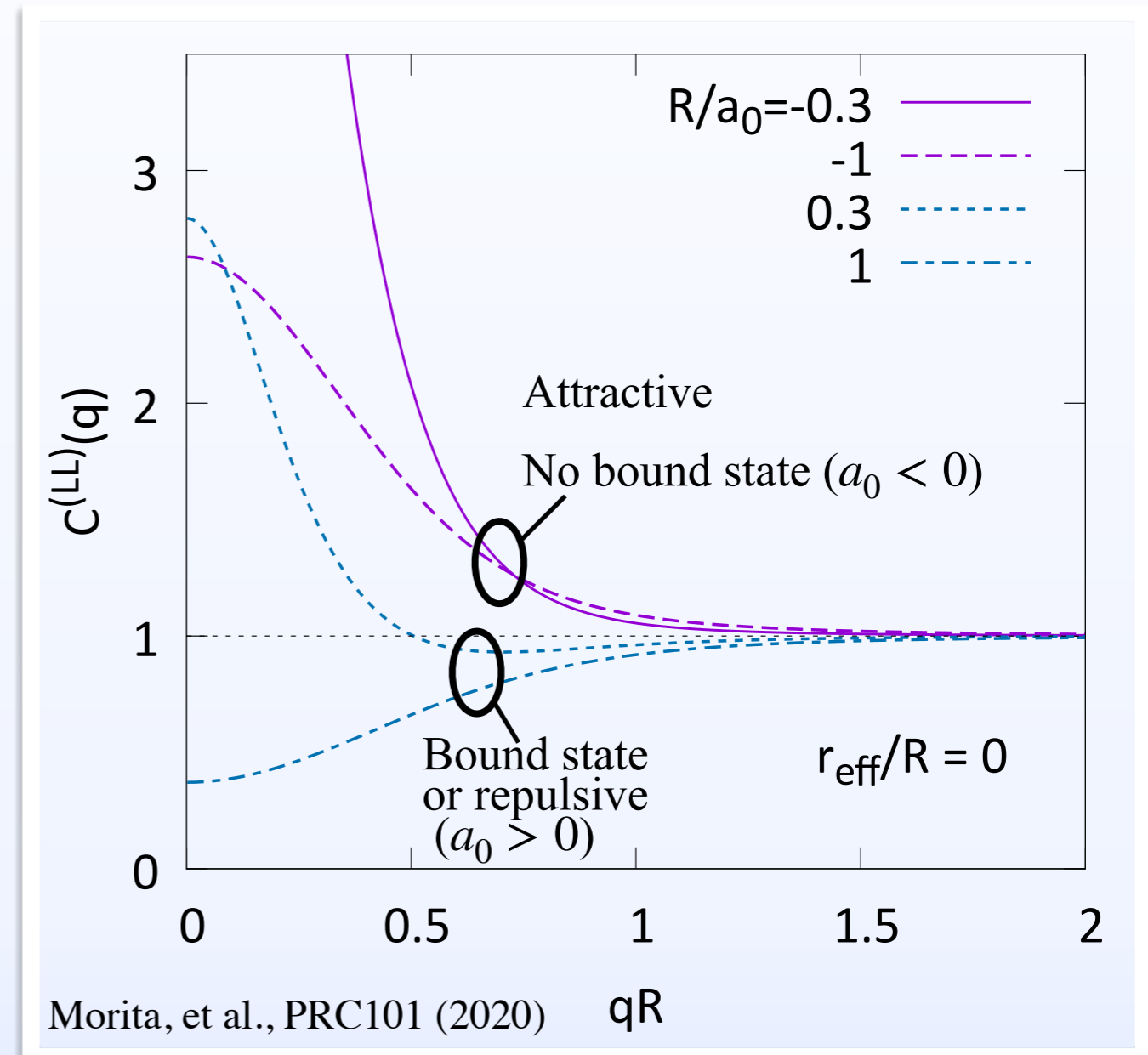
R. Lednicky, et al. Sov. J. Nucl. Phys. 35(1982).

➔ $C = C(qR, R/a_0)$

- $C(q)$ is sensitive to R/a_0 at $qR \lesssim 1$

Sgn(a_0)	Interaction
-	Attraction w/o bound state
+	Attraction w/ bound state or Repulsion

- ➔
- Clear relation between $C(q)$ and $\mathcal{F}(q)$
 - Sensitive to (non)existence of bound state



Hadron correlation in high energy nuclear collision

- Source size dependence of $C(q)$

$$C(\mathbf{q}) \simeq \int d^3\mathbf{r} S(\mathbf{r}) |\varphi^{(-)}(\mathbf{q}, \mathbf{r})|^2$$

$$\rightarrow C = C(qR, R/a_0)$$

- For the specific channel (or a_0)

we can get the different ratio R/a_0 from the different collision experiments

$$R \sim \begin{cases} 1 \text{ fm} & (\text{for } pp \text{ collisions}) \\ 2-5 \text{ fm} & (\text{for } AA \text{ collisions}) \end{cases}$$

- Unbound ($a_0 = -3 \text{ fm}$) case

small source \rightarrow Enhancement

Large source \rightarrow Enhancement

- Bound ($a_0 = 3 \text{ fm}$) case

small source \rightarrow Enhancement

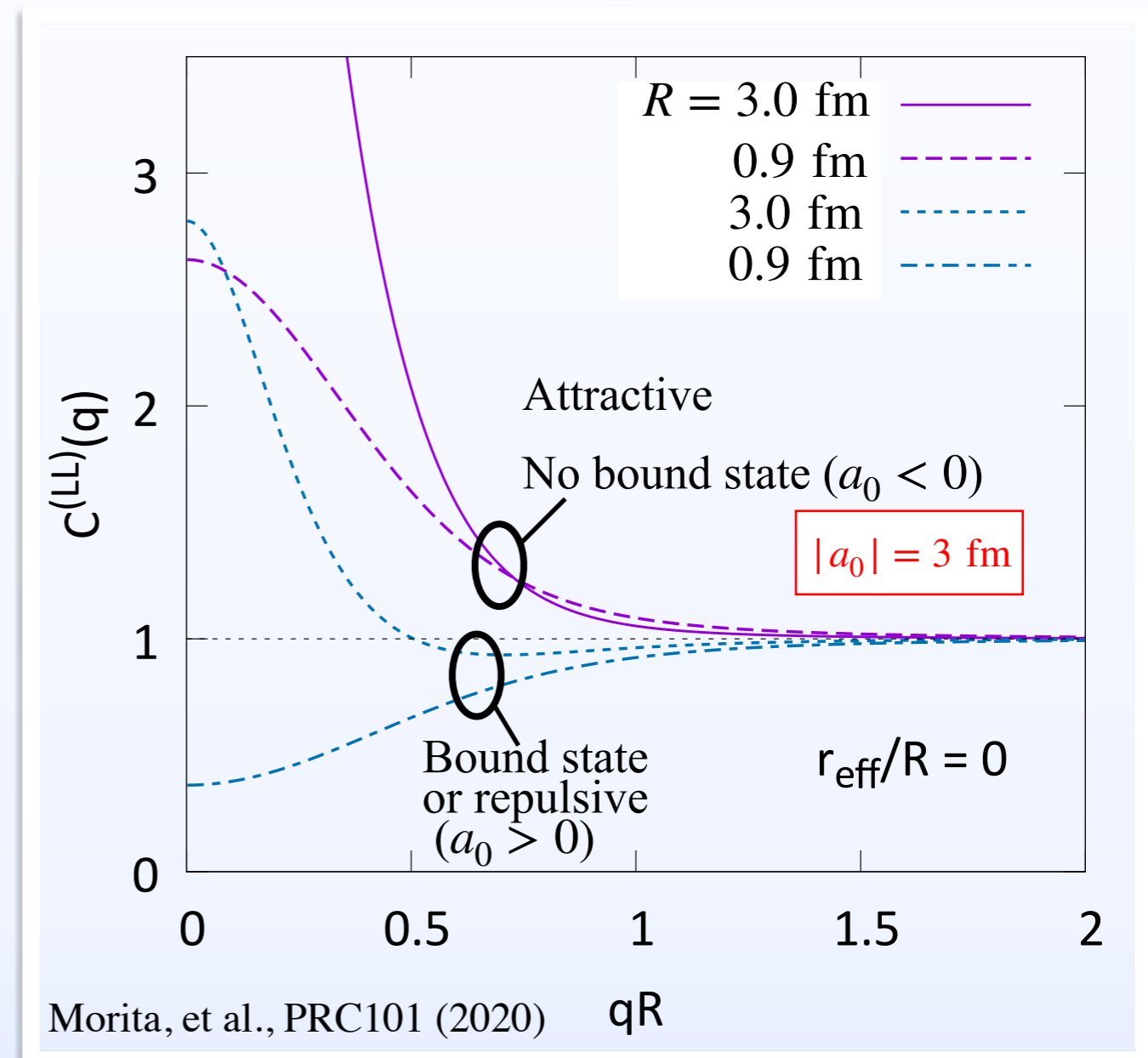
Large source \rightarrow Suppression

Correlation function in different collisions

are very important for the detailed study.

$S(\mathbf{r})$: Source function

$\varphi^{(-)}(\mathbf{q}, \mathbf{r})$: Relative wave function



$\Lambda\Lambda$ - $N\Xi$ interaction and H -dibaryon state

- $\Lambda\Lambda$ - $N\Xi$ interaction ($S = -2$) and H -dibaryon

- $J = 0$: Unique sector in flavor Octet-Octet baryon int.

$$8 \otimes 8 = \underline{1} \oplus 8_A \oplus 8_S \oplus 10 \oplus \bar{10} \oplus 27$$

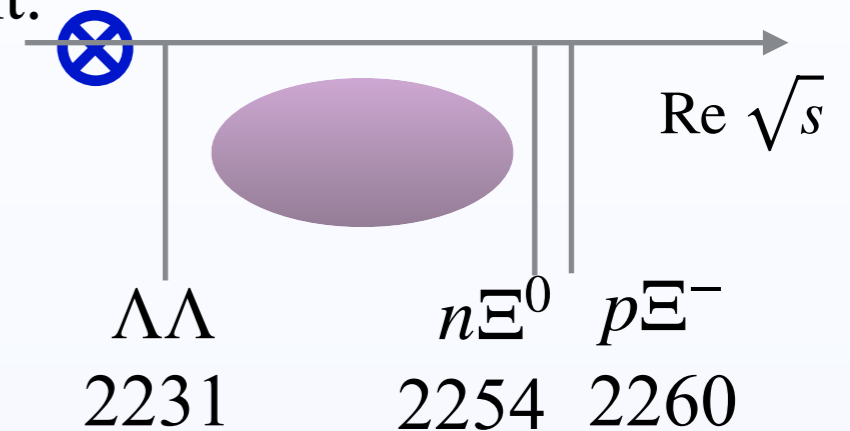
- Pauli arrowed
- Attractive color-magnetic int.

- Flavor-singlet dihyperon “H” R. L. Jaffe, PRL 38 (1977), 195.

Predicted as “single hadron” below $\Lambda\Lambda$

- Binding energy of double Λ hypernucleus
Takahashi et al., PRL87 (2001) 212502

→ $\Lambda\Lambda$ does NOT form (deep) bound state



$\Lambda\Lambda$ - $N\Xi$ interaction and H -dibaryon state

- $\Lambda\Lambda$ - $N\Xi$ interaction ($S = -2$) and H -dibaryon

- $J = 0$: Unique sector in flavor Octet-Octet baryon int.

$$8 \otimes 8 = \underline{1} \oplus 8_A \oplus 8_S \oplus 10 \oplus \bar{10} \oplus 27$$

- Pauli arrowed
- Attractive color-magnetic int.

- Flavor-singlet dihyperon “H” R. L. Jaffe, PRL 38 (1977), 195.

Predicted as “single hadron” below $\Lambda\Lambda$

- Binding energy of double Λ hypernucleus
Takahashi et al., PRL87 (2001) 212502

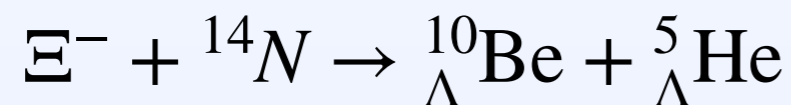
→ $\Lambda\Lambda$ does NOT form (deep) bound state

- Ξ - ^{14}N binding energy in twin-hyper nuclei

K. Nakazawa et al., PTEP 2015 (2015), 033D02

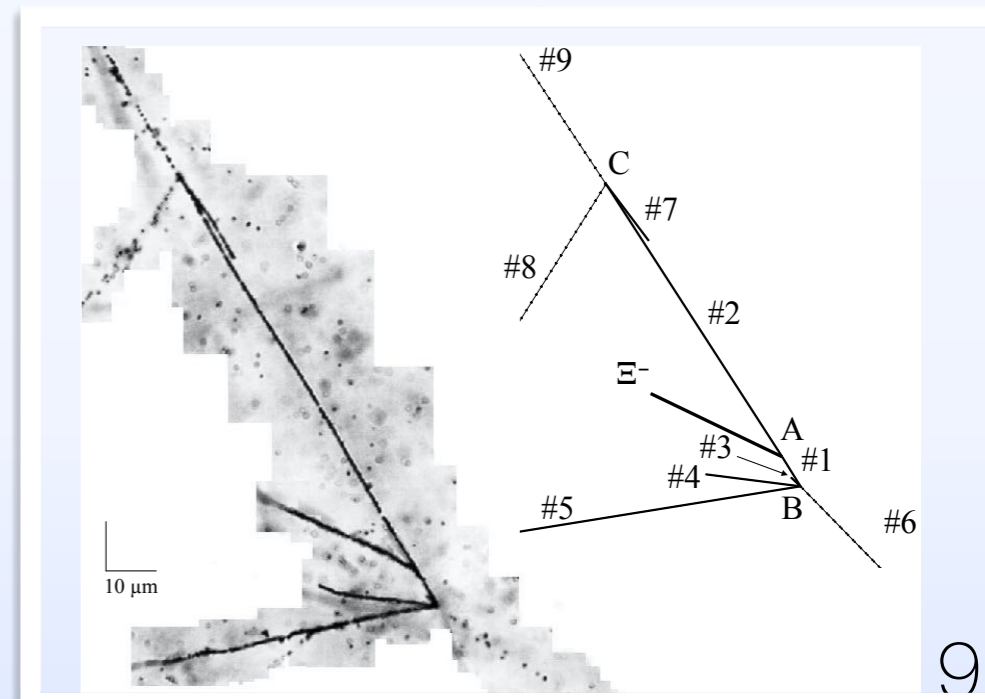
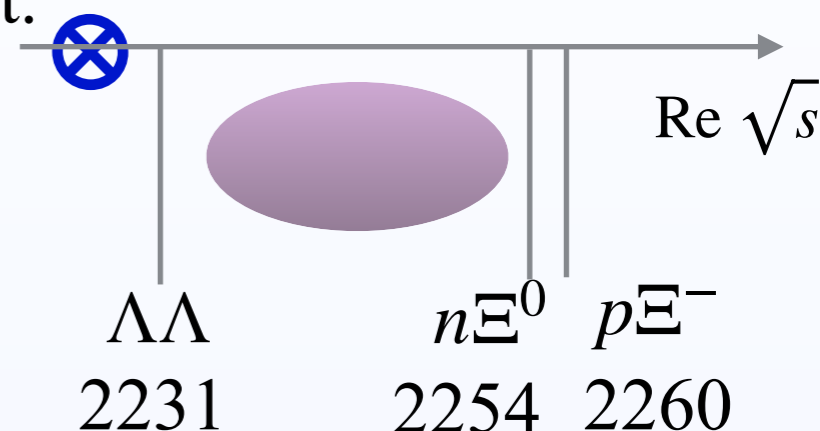
S. H. Hayakawa et al. [J-PARC E07], PRL 126 (2021) 062501

Ibuki event



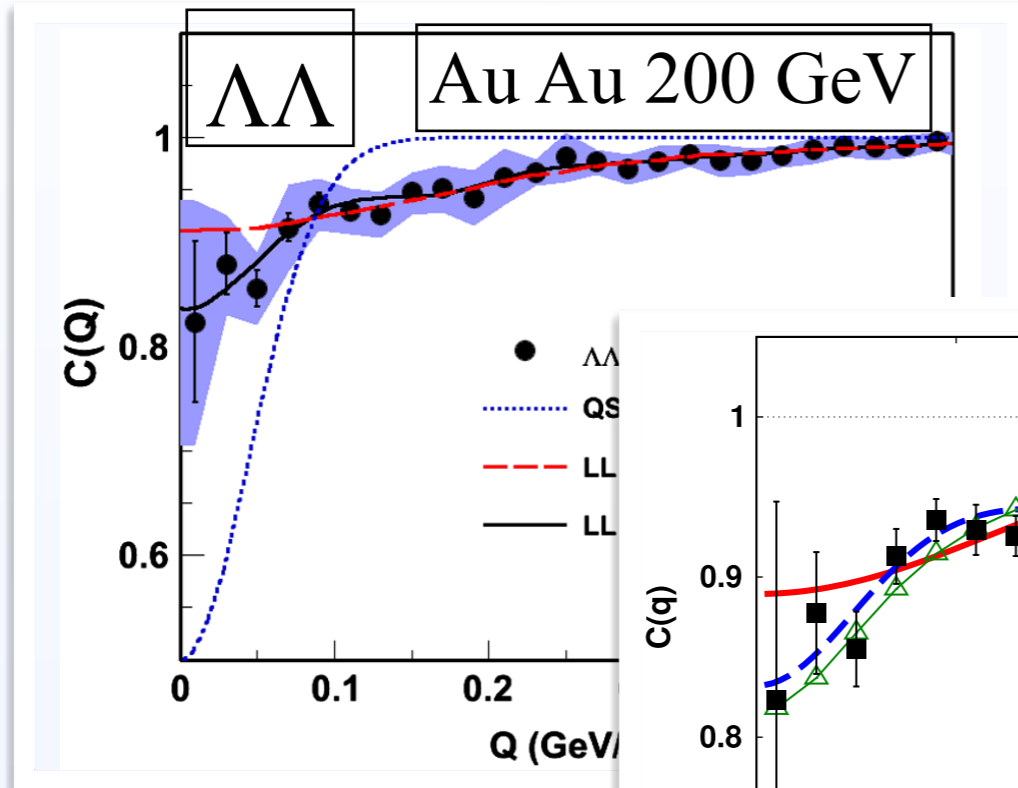
- • Consistent with attractive $N\Xi$ int.

Can dibaryon state emerge as $N\Xi$ quasibound?

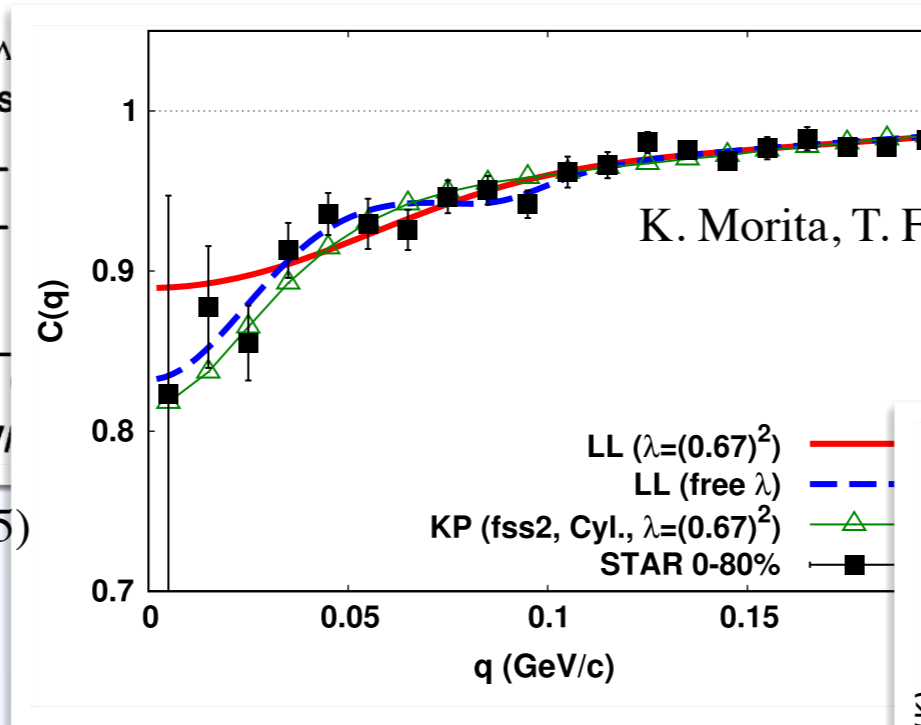


$\Lambda\Lambda$ correlation function

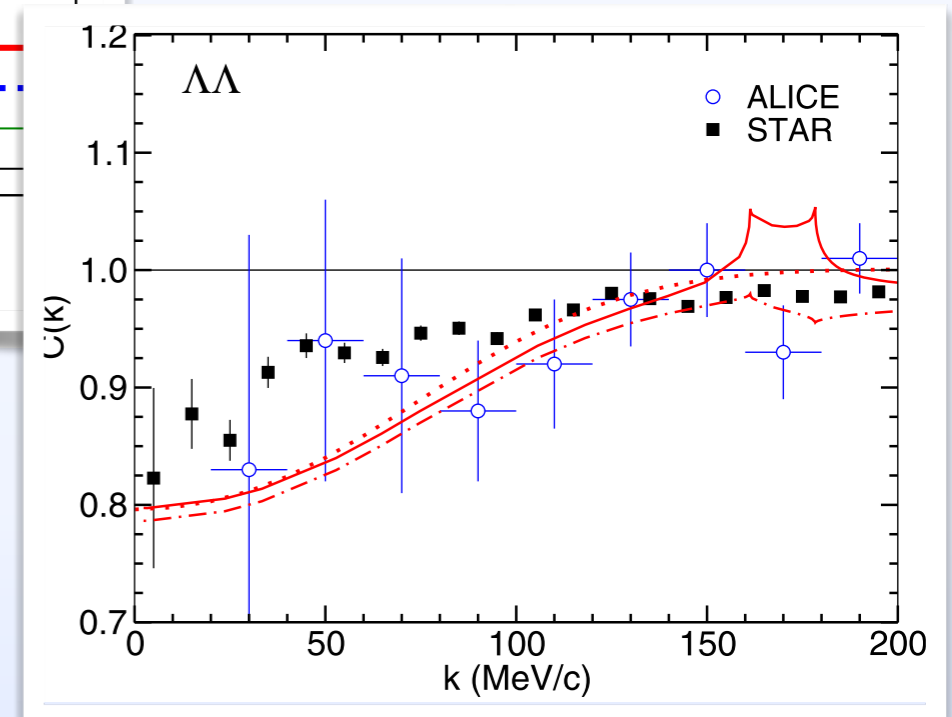
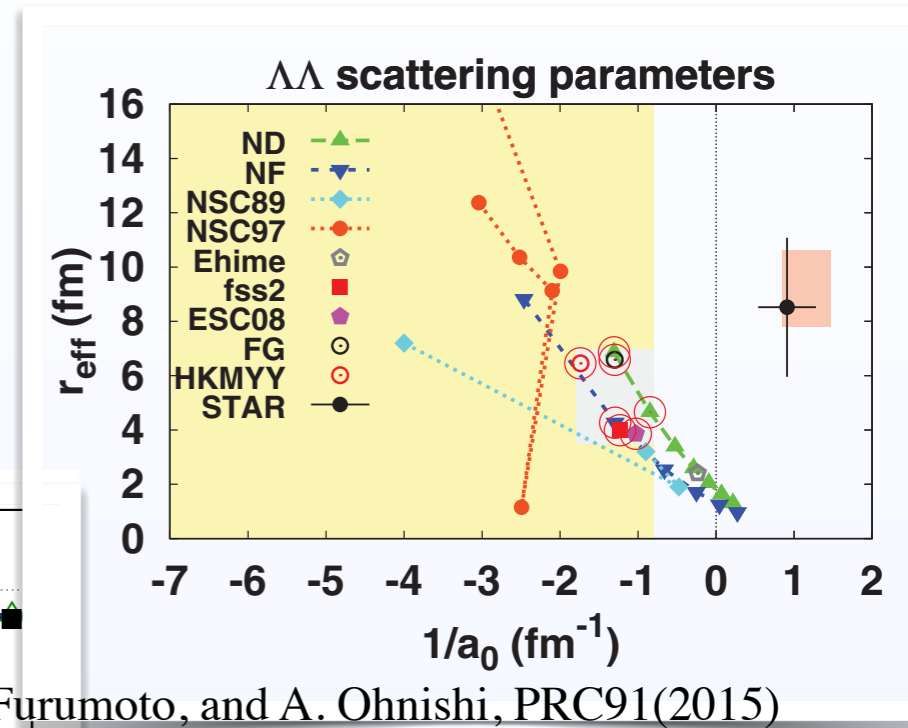
$\Lambda\Lambda$ correlation from STAR



L. Adamczyk, et al. [STAR], PRL114(2015)



K. Morita, T. Furumoto, and A. Ohnishi, PRC91(2015)



- Consistent with weakly attractive $\Lambda\Lambda$ interaction



No H dibaryon state below $\Lambda\Lambda$

$\Lambda\Lambda$ - $N\Xi$ HAL QCD potential

$N\Xi$ - $\Lambda\Lambda$ HAL QCD potential

K. Sasaki et al. [HAL QCD], NPA 998 (2020), 121737.

• HAL QCD method

Ishii, Aoki, Hatsuda, PRL99 (2007) 022001
N. Ishii et al Phys. Lett. B712(2012)437

$$\langle 0 | B_1 B_2(t, \vec{r}) \vec{I}(0) | 0 \rangle = A_0 \underbrace{\Psi(\vec{r}, E_0)}_{V(r)} e^{-E_0 t} + \dots$$

• Nearly physical mass calculation

$$m_\pi = 146 \text{ MeV} \quad m_K = 525 \text{ MeV}$$

$N\Xi$ - $\Lambda\Lambda$ $J = 0$ channel

• Strong attraction for $N\Xi$ ($I = 0$)

• Weak attraction for $\Lambda\Lambda$ channel

• Weak $\Lambda\Lambda$ - $N\Xi$ coupling

• Solving Schrödinger eq. with physical masses

Y. Kamiya, et al. in preparation

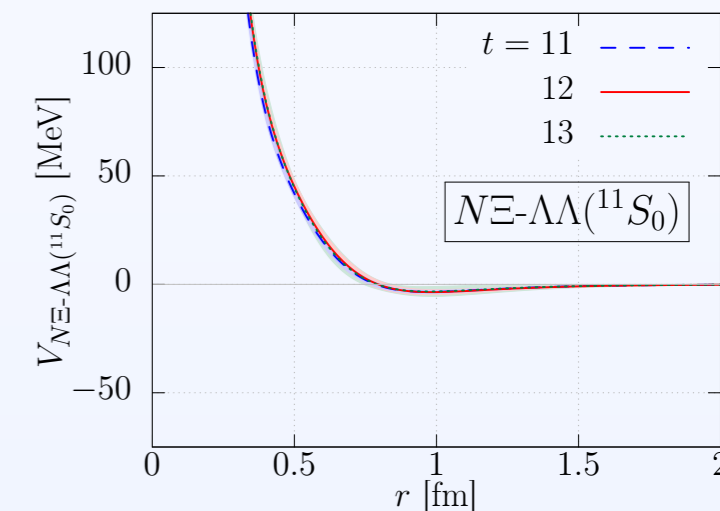
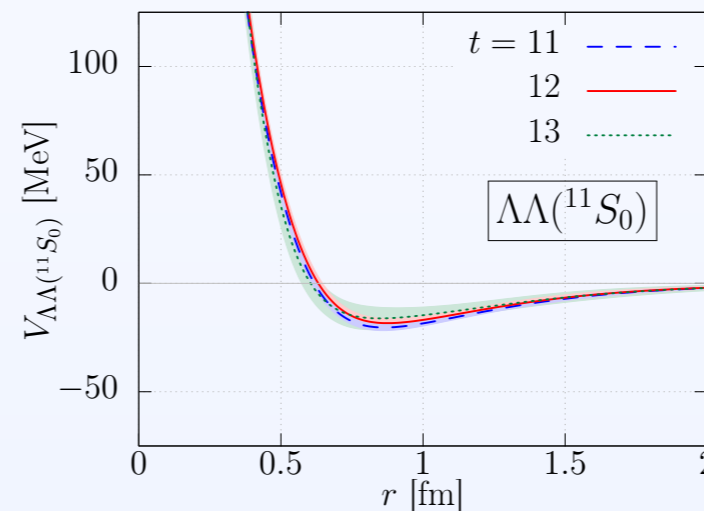
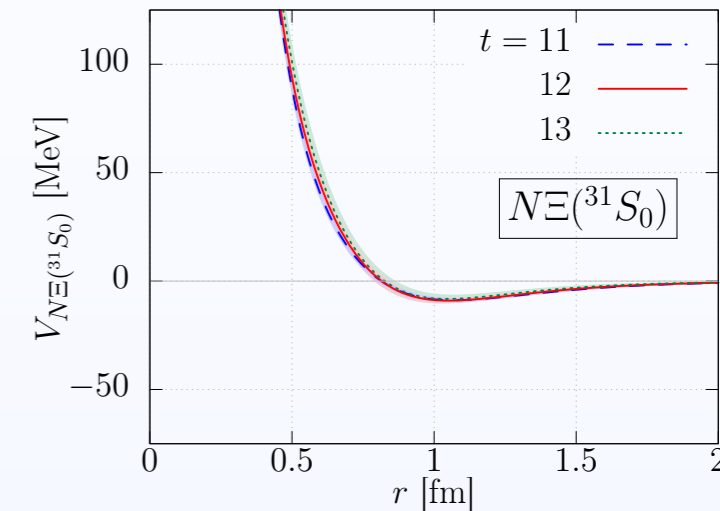
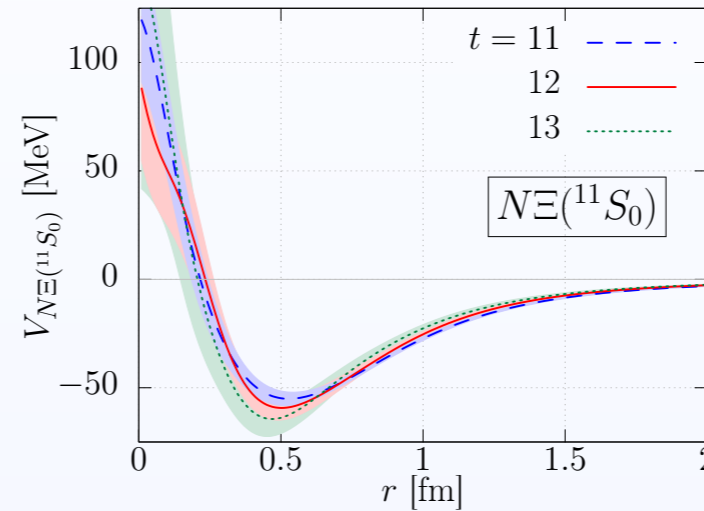
$$\text{Scat. length : } a_0 \equiv -\mathcal{F}(E_{\text{th}})$$

$$\text{Virtual pole : } -3.9 - i0.3 \text{ MeV (from } n\Xi^0 \text{ thr.)}$$



H dibaryon state: merely unbound

$2I+1, 2s+1 L_J$



channel	a_0 [fm]
$J = 0 \quad p\Xi^-$	$-1.22 \pm 0.13_{-0.00}^{+0.08} - i1.57 \pm 0.35_{-0.23}^{+0.18}$
$n\Xi^0$	$-2.07 \pm 0.39_{-0.35}^{+0.28} - i0.14 \pm 0.08_{-0.01}^{+0.00}$
$\Lambda\Lambda$	$-0.78 \pm 0.22_{-0.13}^{+0.00}$

$\Lambda\Lambda$ - $N\Xi$ HAL QCD potential

● $N\Xi$ $J = 1$ channel K. Sasaki et al. [HAL QCD], NPA 998 (2020), 121737.
Y. Kamiya, et al. in preparation

- Weekly attractive
- Similar potential for both I components

Negligible coupling between $n\Xi^0$ and $p\Xi^-$

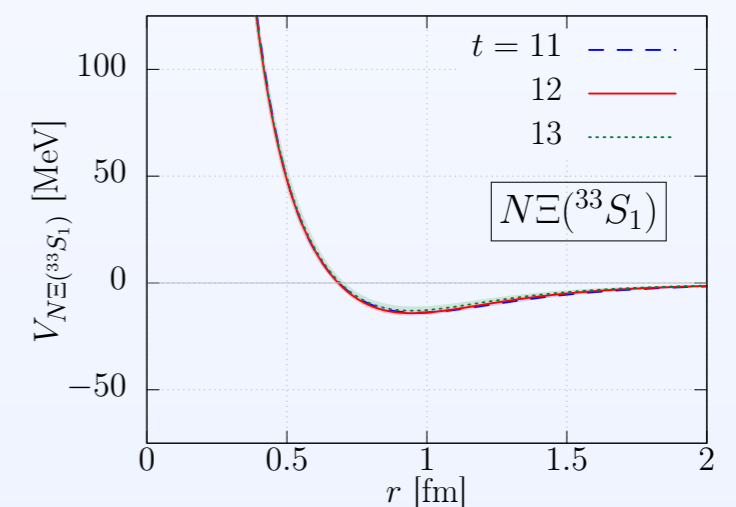
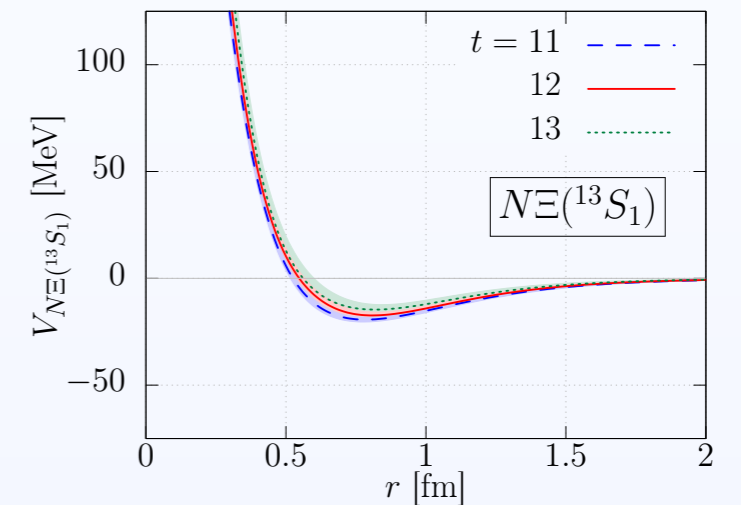


small negative a_0

channel	a_0 [fm]
$J = 1$ $p\Xi^-$	$-0.35 \pm 0.06^{+0.09}_{-0.07} - i0.00$
$n\Xi^0$	$-0.35 \pm 0.06^{+0.09}_{-0.07}$

- For calculation of $C_{p\Xi^-}$,
both ($J = 0$ and 1) contribution must be summed up.

$2I+1, 2s+1 L_J$



$\Lambda\Lambda$ and $p\Xi^-$ correlation function

$p\Xi^-$ and $\Lambda\Lambda$ correlation from ALICE

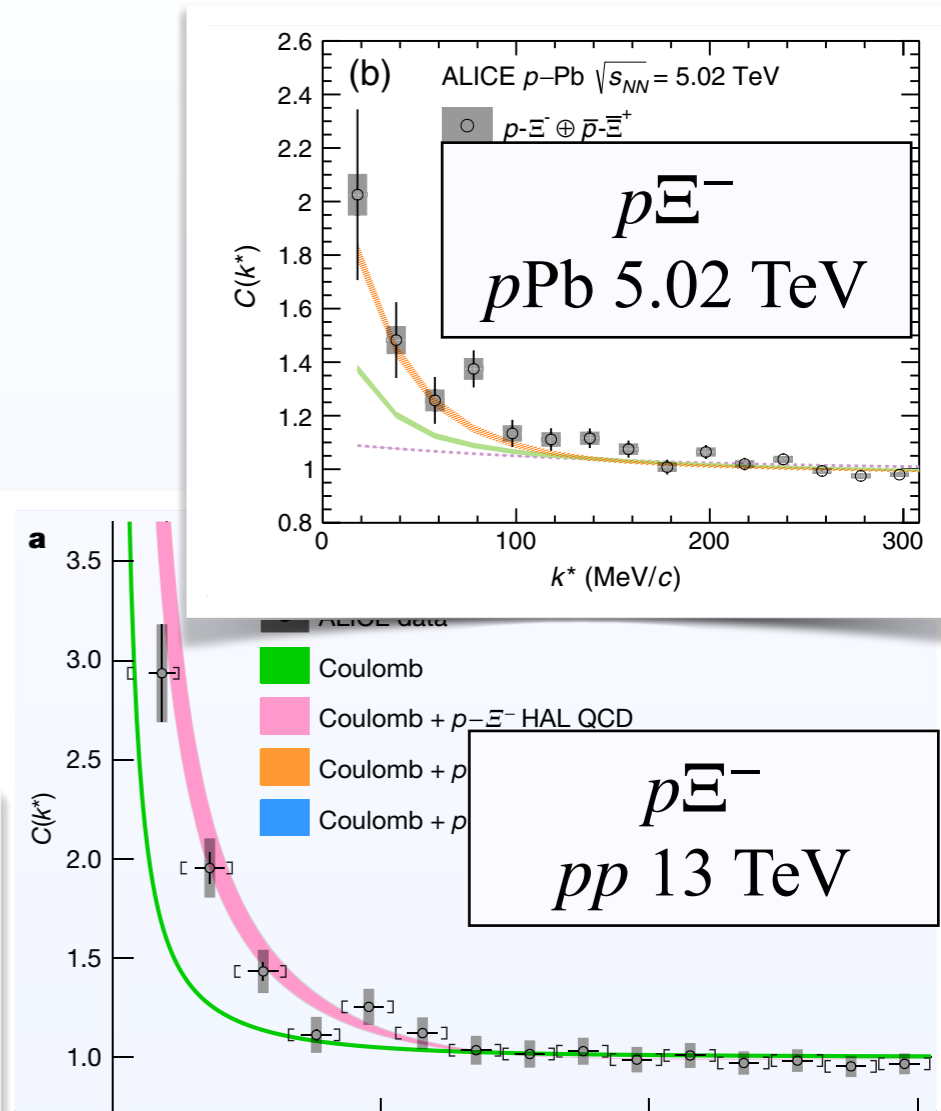
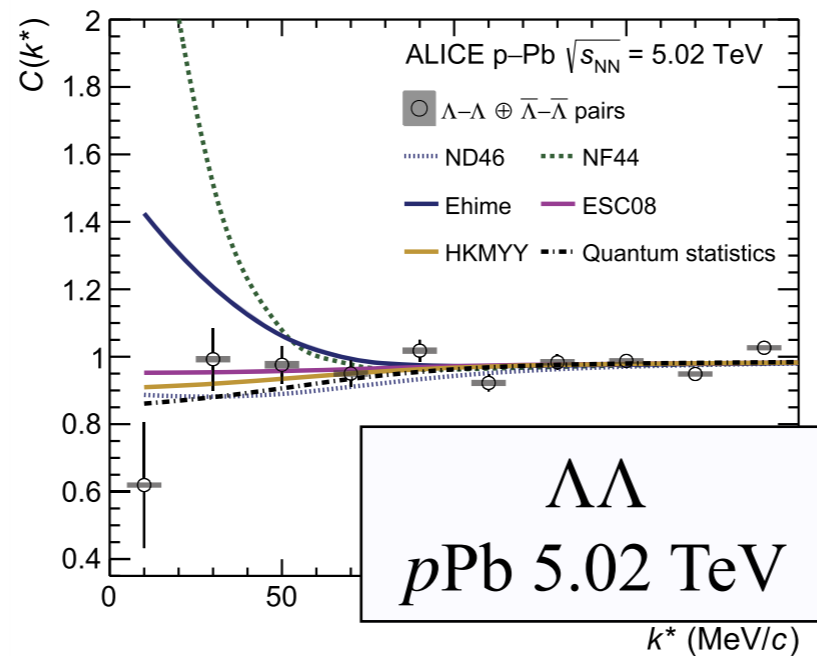
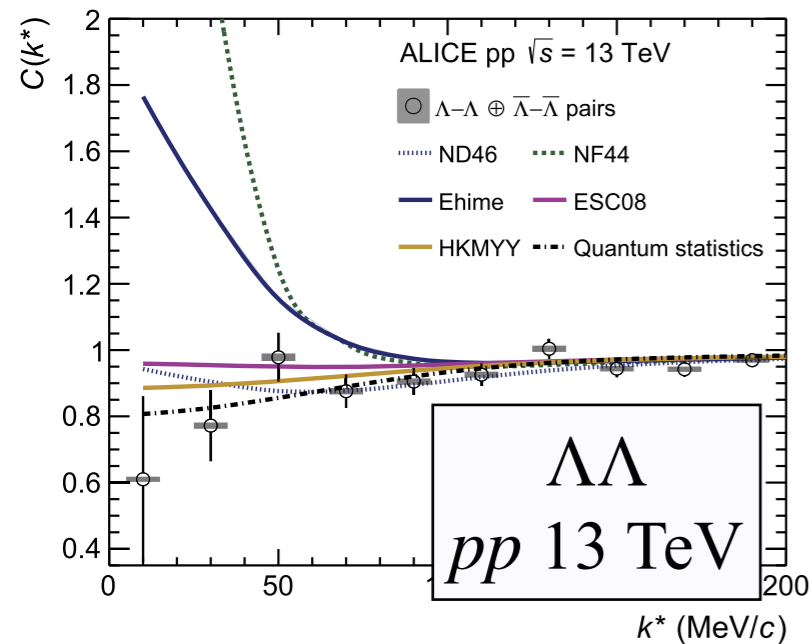
$p\Xi^-$

pPb 5.02 TeV : ALICE, PRL, 123 (2019), 112002

pp 13 TeV : ALICE, Nature 588 (2020), 232-238

$\Lambda\Lambda$

pPb 5.02 TeV S. Acharya et al. [ALICE], PLB 797 (2019).
pp 13 TeV



$\Lambda\Lambda$ and $p\Xi^-$ correlation function

- $p\Xi^-$ and $\Lambda\Lambda$ correlation from ALICE

- $p\Xi^-$

pPb 5.02 TeV : ALICE, PRL, 123 (2019), 112002
 pp 13 TeV : ALICE, Nature 588 (2020), 232-238

- $\Lambda\Lambda$

pPb 5.02 TeV
 pp 13 TeV S. Acharya et al. [ALICE], PLB 797 (2019).

- Our study: Systematic analysis with including

- Coulomb interaction
- Coupled-channel effect
- Threshold difference

using latest HAL QCD coupled-channel potential

- Systematic comparison to

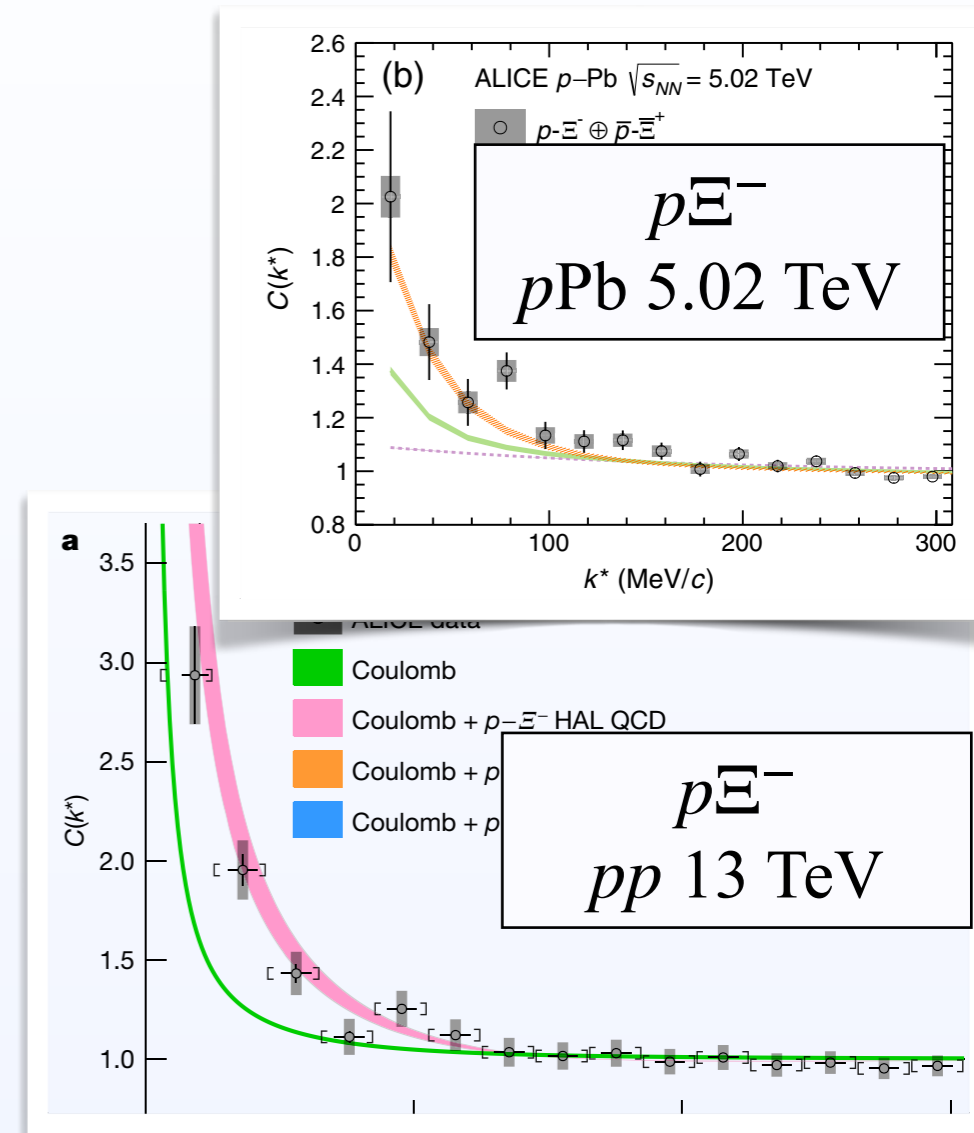
$\Lambda\Lambda$ and $p\Xi^-$ correlation from pp and $p\text{Pb}$ collisions data from ALICE

- Static spherical Gaussian with $R_{N\Xi} \sim R_{\Lambda\Lambda}$

- Fitting for comparison

$$C_{\text{fit}}(q) = \frac{A_{\text{non-femt}}(q)}{a + bq} \times [1 + \frac{\lambda(C_{\text{Theor}}(q) - 1)}{< 1}]$$

- Miss identification
- feed-down



KP Formula for Coupled-channel systems

- Koonin-Pratt formula for coupled-channel systems

$$\text{Koonin-Pratt formula : } C(\mathbf{q}) \simeq \int d^3\mathbf{r} S(\mathbf{r}) |\psi^{(-)}(\mathbf{q}; \mathbf{r})|^2$$

S.E. Koonin, PLB 70 (1977)
S. Pratt et. al. PRC 42 (1990)



- For coupled-channel systems

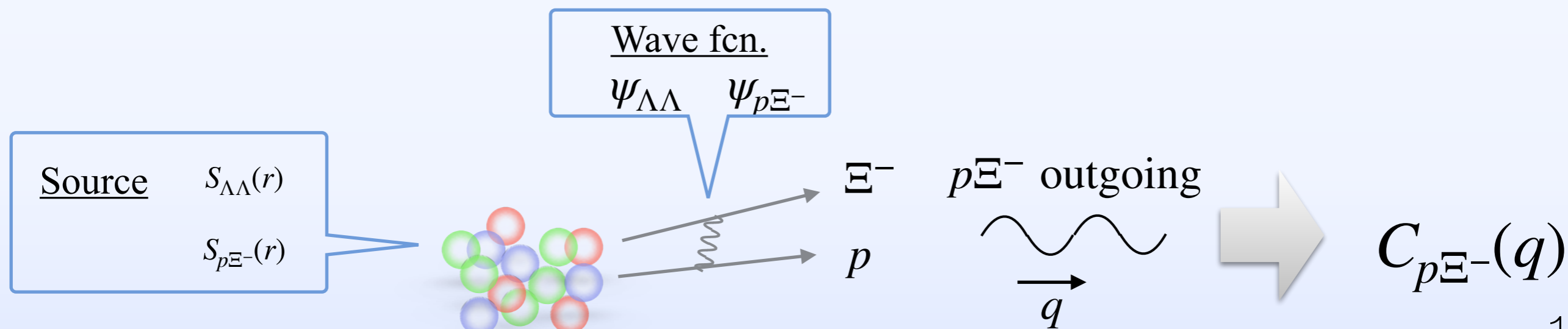
R. Lednicky, et. al. Phys. At. Nucl. 61 (1998)
Haidenbauer NPA 981 (2018)

$$\text{KPLLL formula : } C_i(\mathbf{q}) = \int d^3\mathbf{r} S_i(\mathbf{r}) |\psi_i^{C,(-)}(q; r)|^2 + \sum_{j \neq i} \omega_j \int d^3\mathbf{r} S_j(\mathbf{r}) |\psi_j^{C,(-)}(q; r)|^2$$

(Koonin-Pratt-Lednicky
-Lyuboshits-Lyuboshits)

Coupled-channel
wave function

- Contribution from coupled-channel source (for $p\Xi^-$)



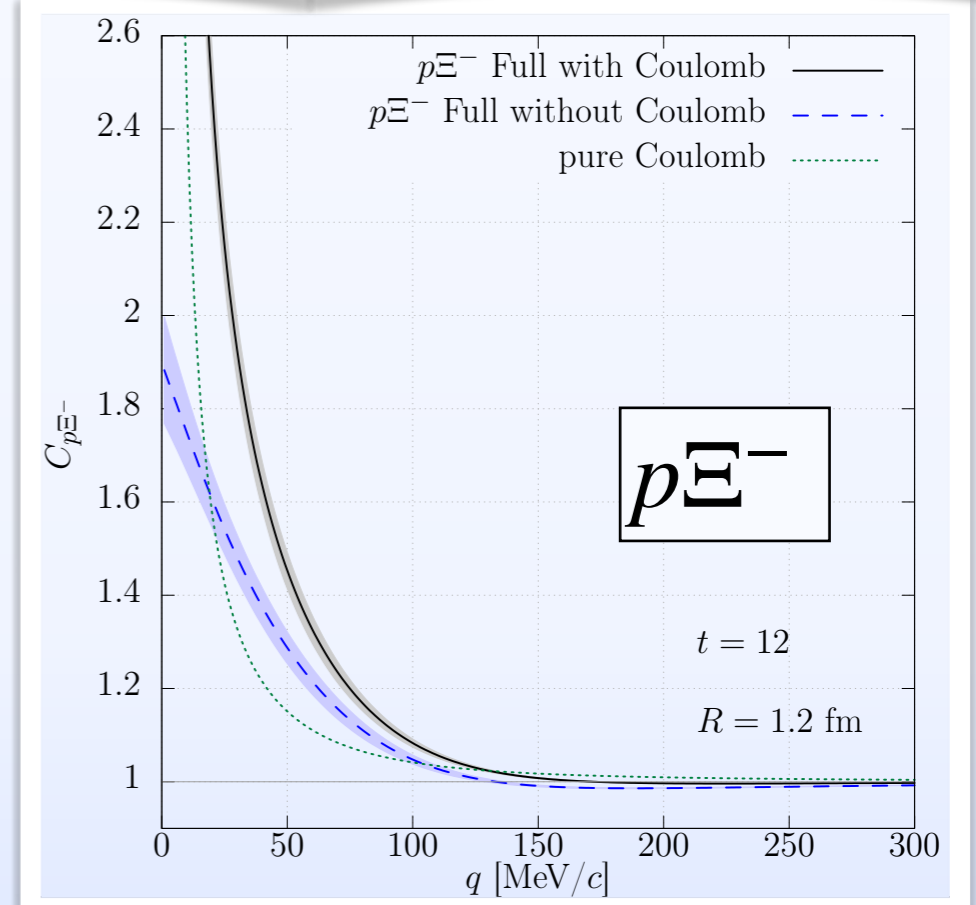
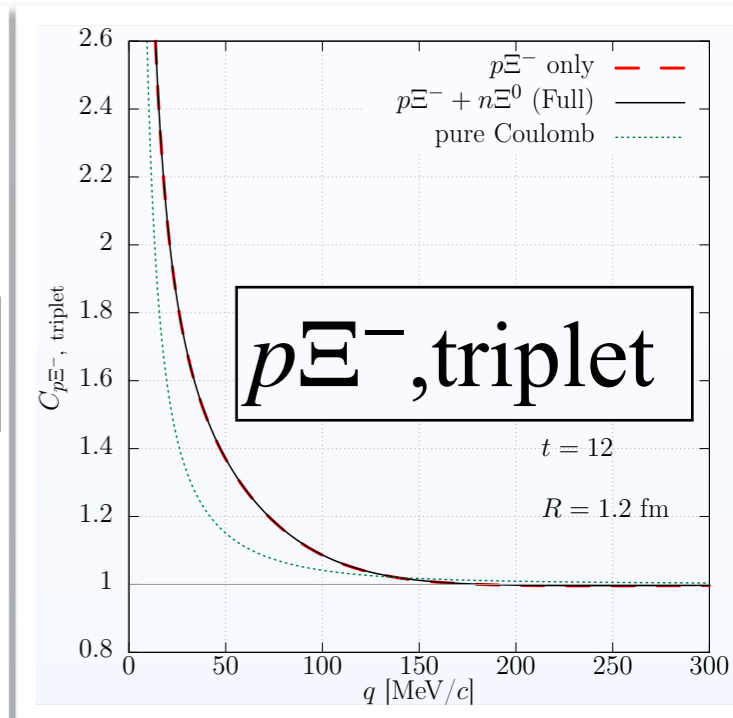
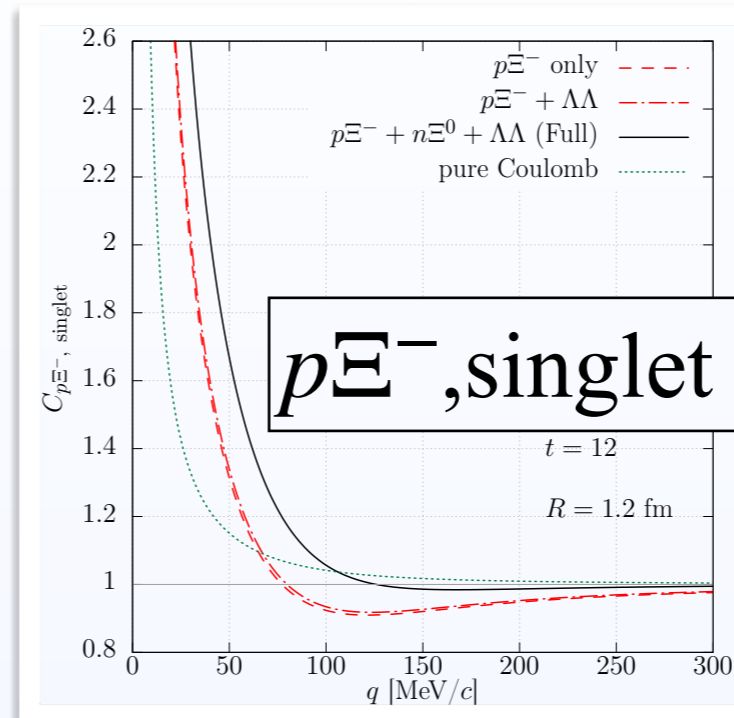
$\Lambda\Lambda$ - $N\Xi$ interaction and $\Lambda\Lambda$ and $p\Xi^-$ correlation function

$p\Xi^-$ correlation function

$$C_{p\Xi^-} = \frac{1}{4} C_{p\Xi^-, \text{singlet}} + \frac{3}{4} C_{p\Xi^-, \text{triplet}}$$

Couples to $\Lambda\Lambda$
(H-dibaryon channel)

- Enhancement from **pure Coulomb** case
- $n\Xi^0$ source contribution
Singlet (J=0) : sizable enhancement
Triplet (J=1) : negligible
- $\Lambda\Lambda$ source contribution : Negligible



$\Lambda\Lambda$ - $N\Xi$ interaction and $\Lambda\Lambda$ and $p\Xi^-$ correlation function

- $p\Xi^-$ correlation function

$$C_{p\Xi^-} = \frac{1}{4} C_{p\Xi^-, \text{singlet}} + \frac{3}{4} C_{p\Xi^-, \text{triplet}}$$

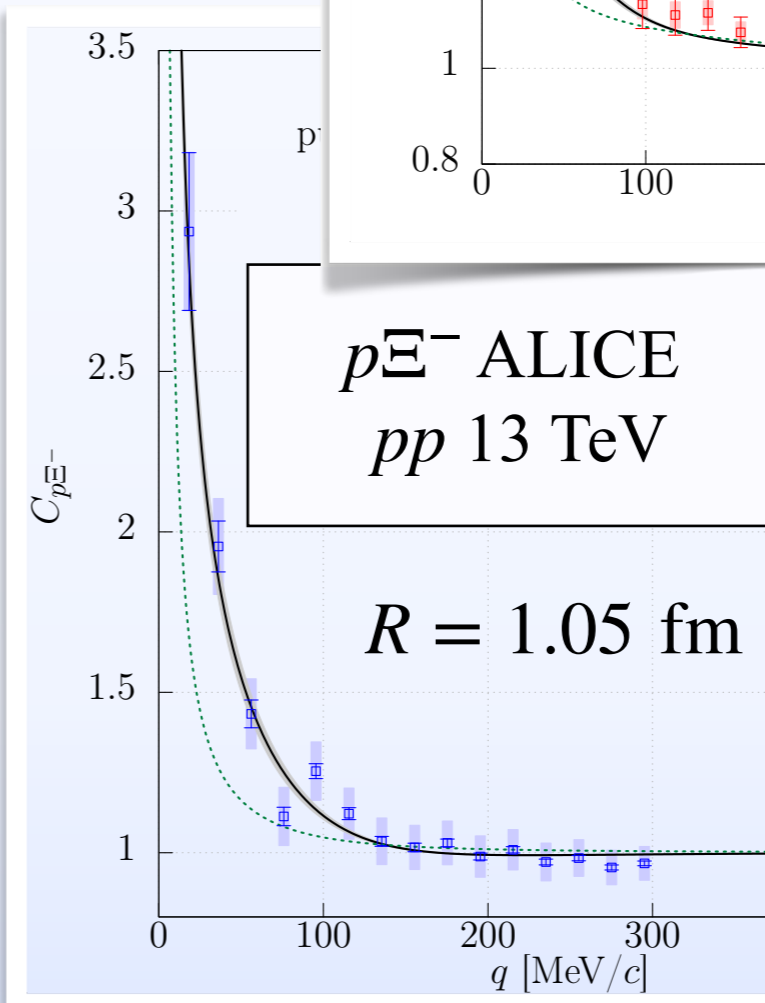
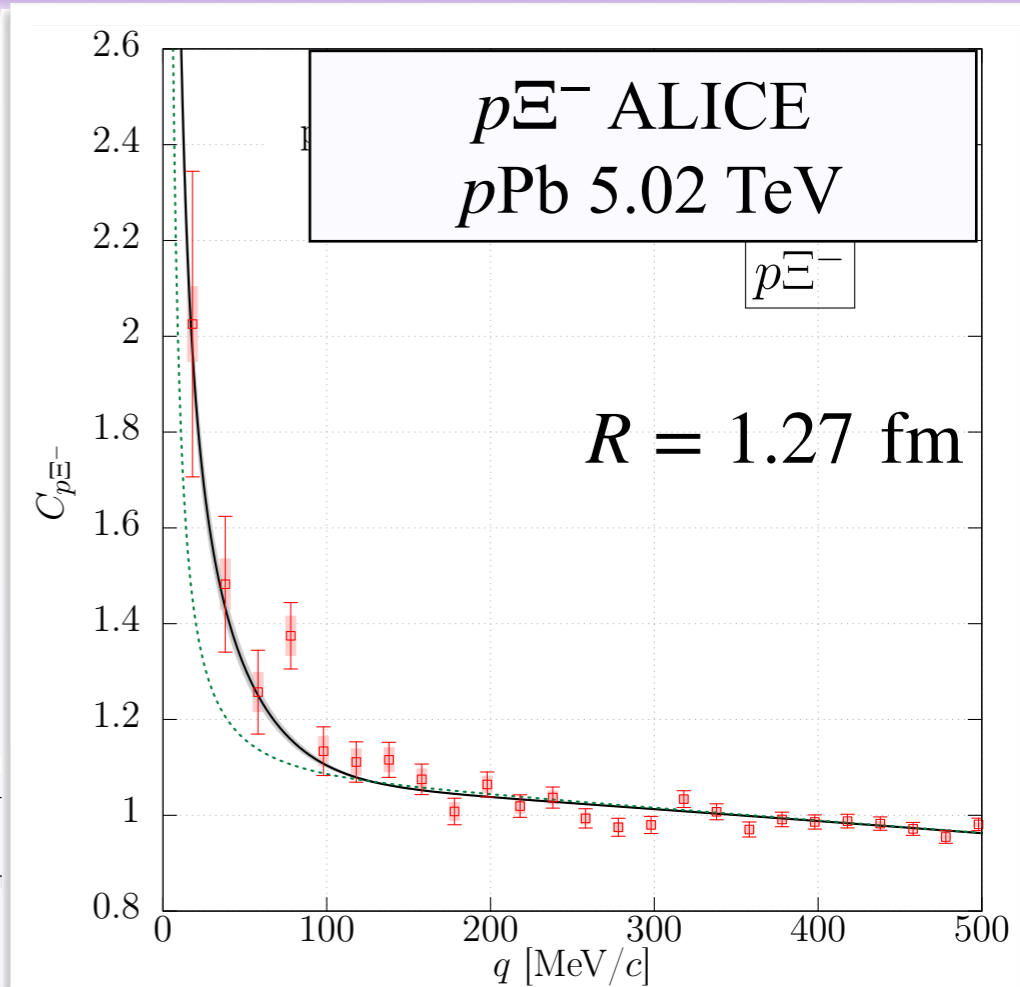
Couples to $\Lambda\Lambda$
(H-dibaryon channel)

- Enhancement from **pure Coulomb** case
- $n\Xi^0$ source contribution
Singlet (J=0) : sizable enhancement
Triplet (J=1) : negligible
- $\Lambda\Lambda$ source contribution : Negligible
- Comparison with ALICE data

pPb 5.02 TeV : ALICE, PRL, 123 (2019), 112002
pp 13 TeV : ALICE, Nature 588 (2020), 232-238

$$C_{\text{fit}}(q) = C_{\text{non-femt}}(q) \times C_{\text{theor}}(q)$$

Well reproduced with enhancement
from pure Coulomb



$\Lambda\Lambda$ - $N\Xi$ interaction and $\Lambda\Lambda$ and $p\Xi^-$ correlation function

- $\Lambda\Lambda$ correlation function

$$C_{\Lambda\Lambda} = 1 - \frac{1}{2} \frac{\exp(-4q^2 R^2)}{\cancel{\quad}} + \frac{\Delta C_{\Lambda\Lambda}}{\text{Strong int.}}$$

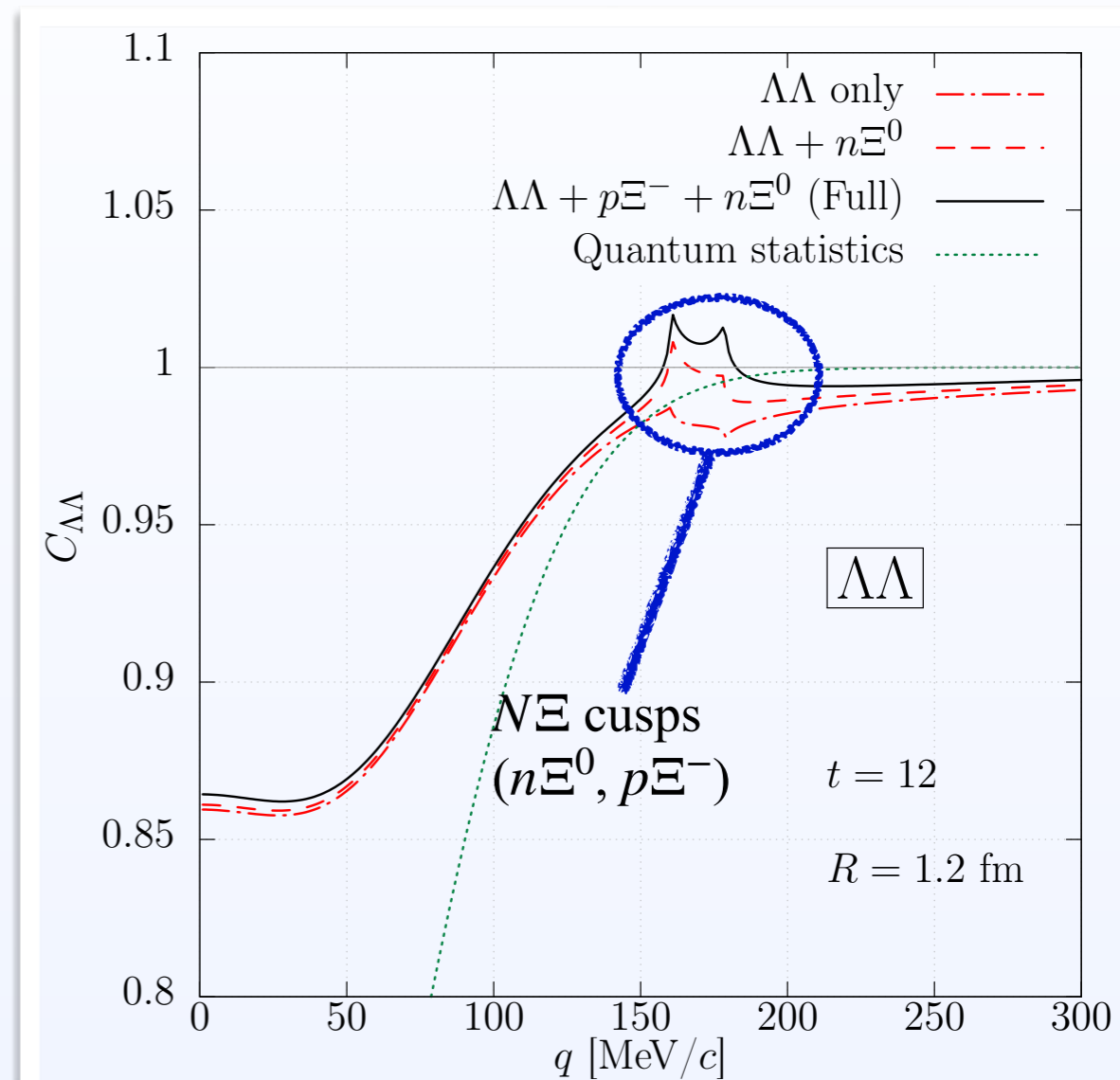
Quantum statistics

- Enhancement from quantum statistics weak attractive interaction

- $N\Xi$ cusps related to the coupling and existence of H-dibaryon

J. Haidenbauer, Nucl. Phys. A 981 (2019),

→ Almost invisible with HAL-QCD potential (small coupling)



$\Lambda\Lambda$ - $N\Xi$ interaction and $\Lambda\Lambda$ and $p\Xi^-$ correlation function

- $\Lambda\Lambda$ correlation function

$$C_{\Lambda\Lambda} = 1 - \frac{1}{2} \frac{\exp(-4q^2 R^2)}{\cancel{\quad}} + \frac{\Delta C_{\Lambda\Lambda}}{\text{Strong int.}}$$

Quantum statistics

- Enhancement from quantum statistics weak attractive interaction

- $N\Xi$ cusps related to the coupling and existence of H-dibaryon

J. Haidenbauer, Nucl. Phys. A 981 (2019),

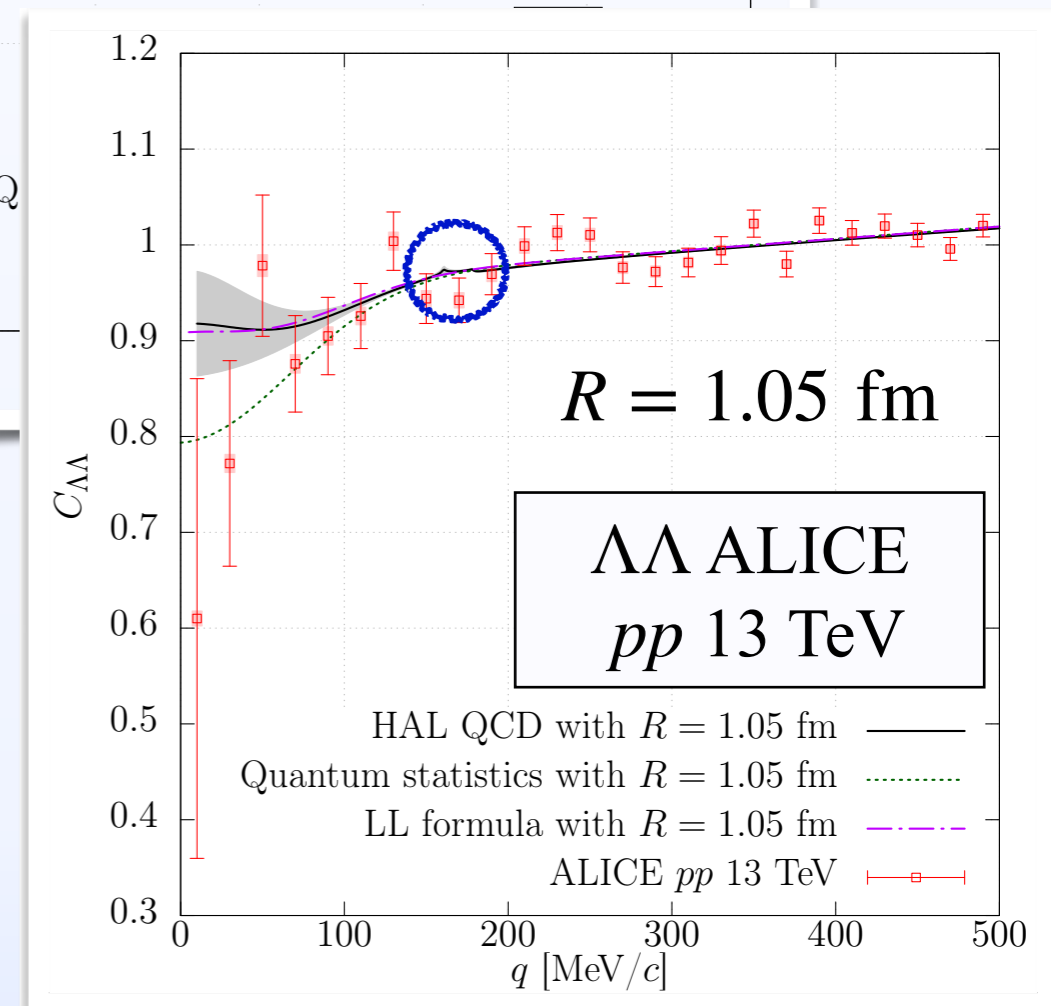
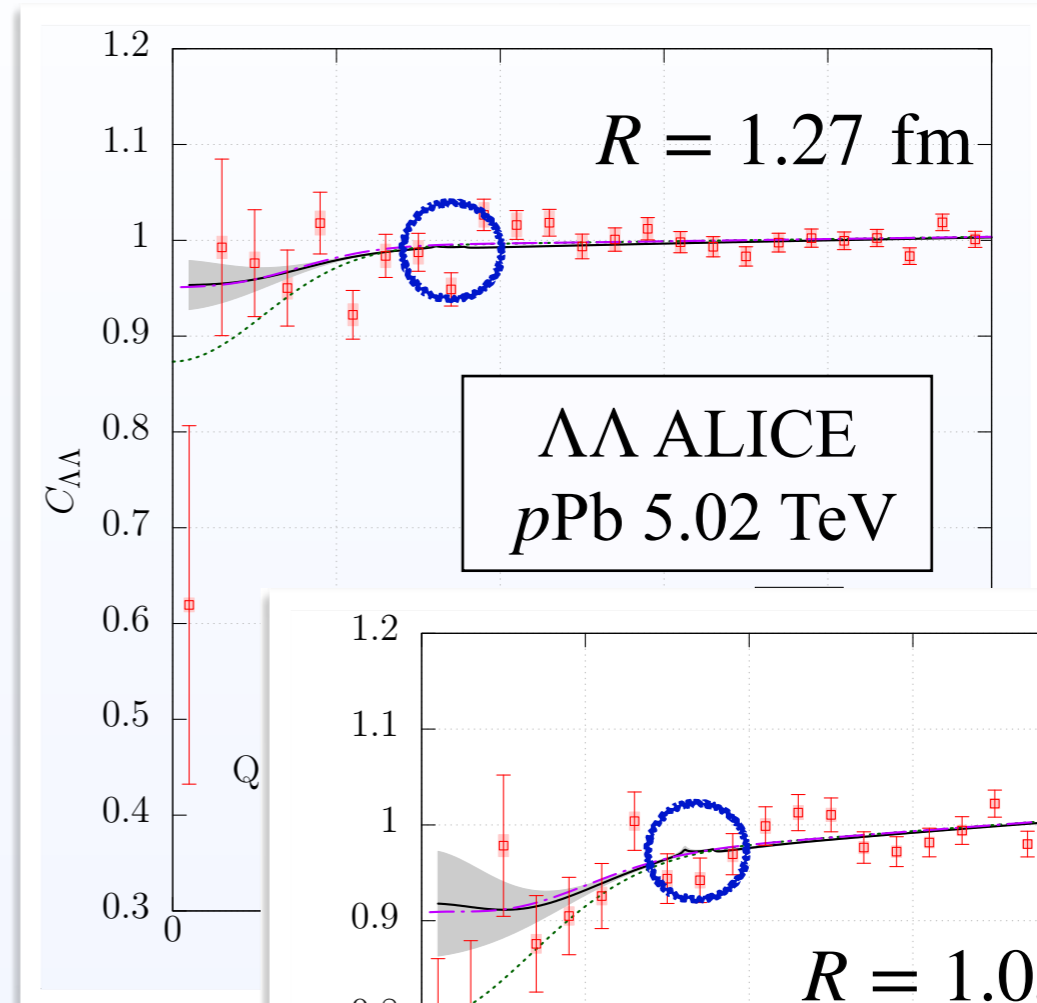
→ Almost invisible with HAL-QCD potential (small coupling)

- Comparison with ALICE data

pPb 5.02 TeV, pp 13 TeV collisions :

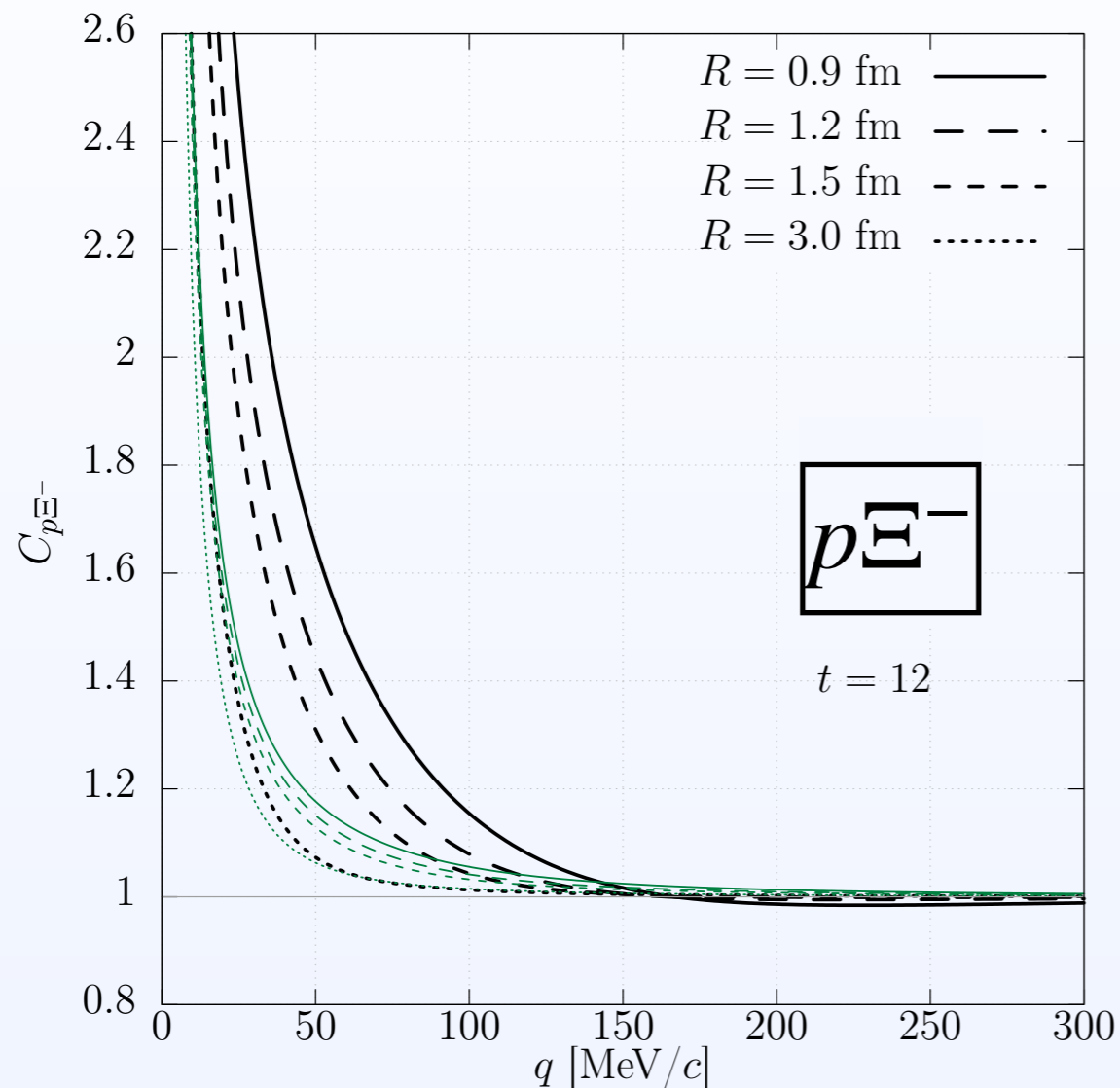
S. Acharya et al. [ALICE], PLB 797 (2019).

- Weak attraction of $\Lambda\Lambda$ int.
- There is no signal of H-dibaryon



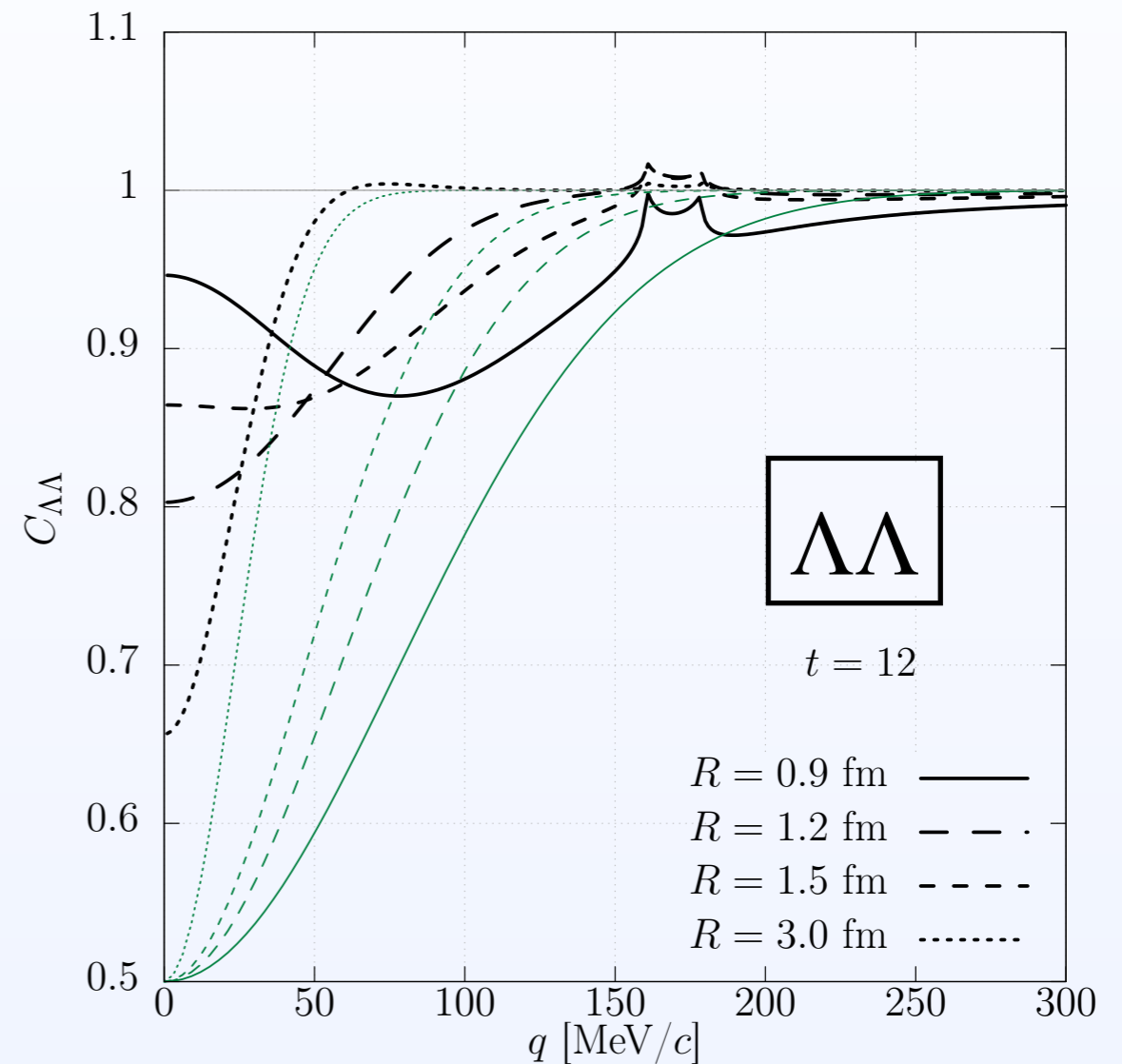
$\Lambda\Lambda$ - $N\Xi$ interaction and $\Lambda\Lambda$ and $p\Xi^-$ correlation function

● Source size dependence



● $p\Xi^-$

- small enhancement for large ($R \sim 3$ fm) case w/o any dip

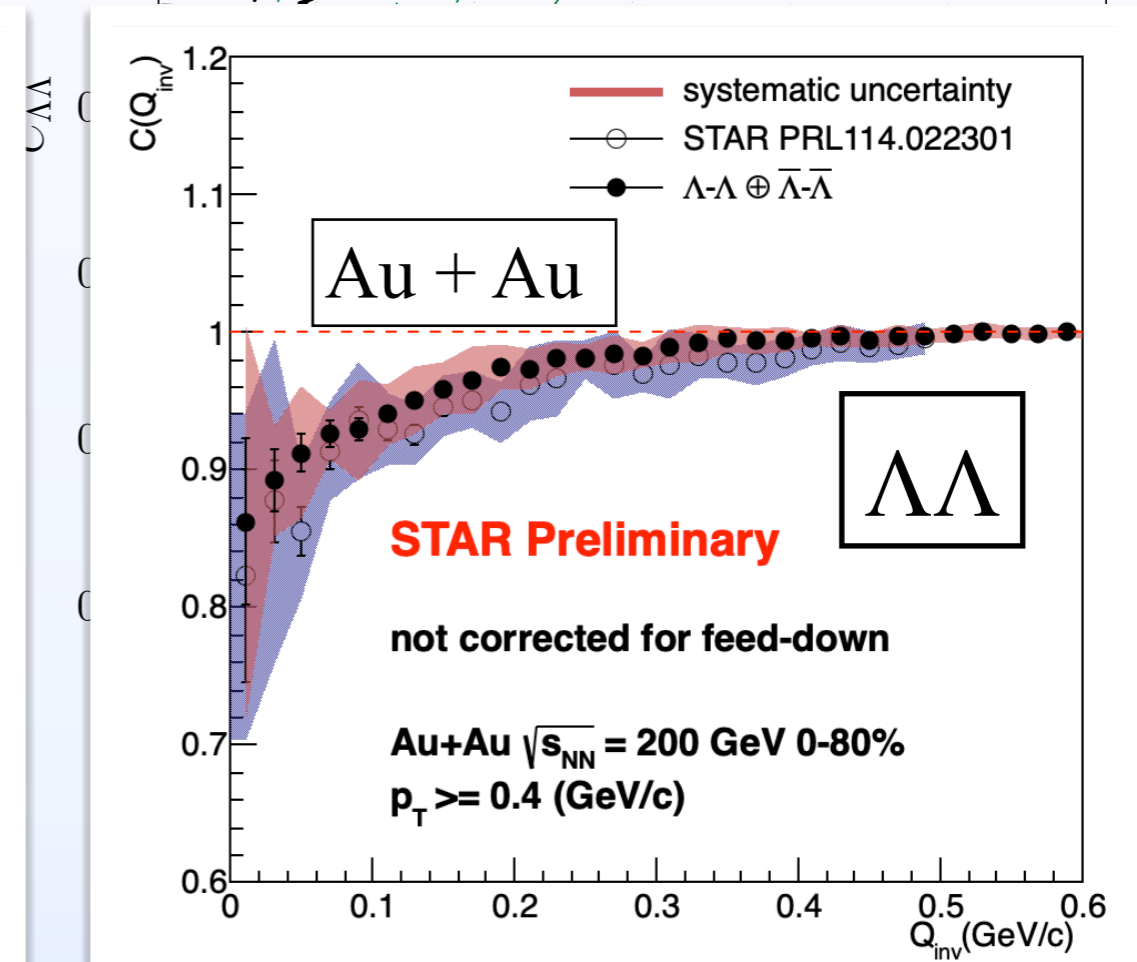
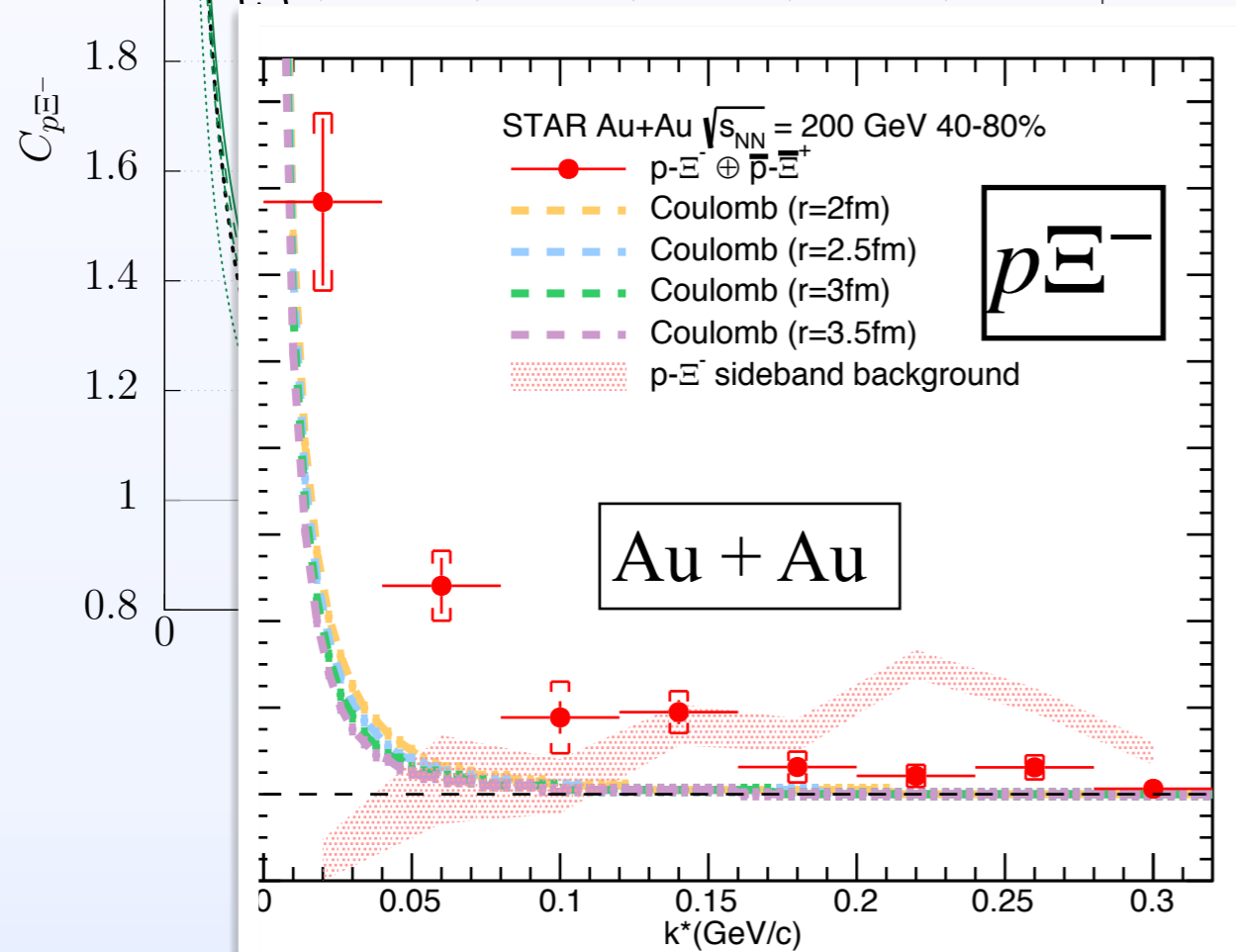
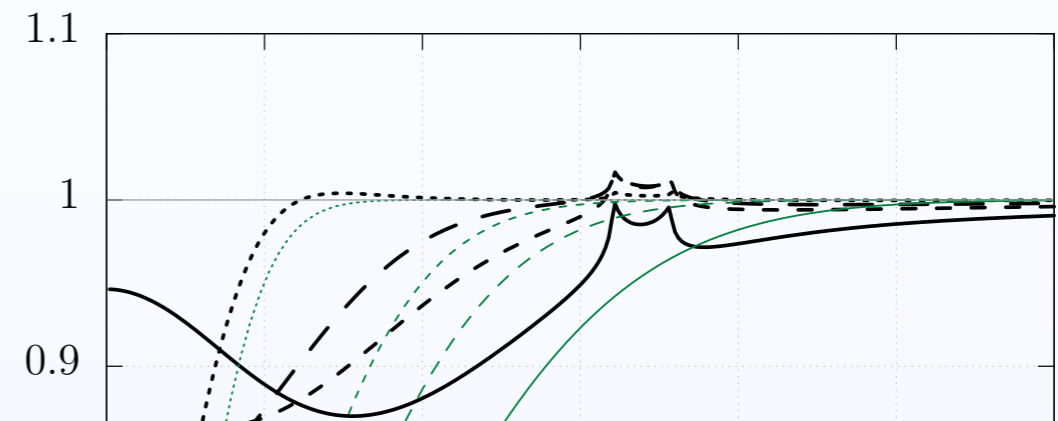
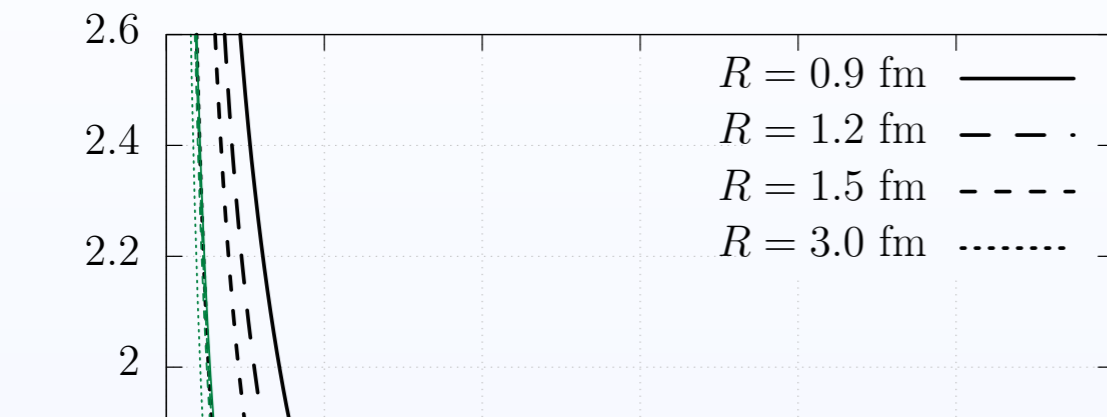


● $\Lambda\Lambda$

- To see source size dependence at small q is mandatory to see the weak interaction

$\Lambda\Lambda$ - $N\Xi$ interaction and $\Lambda\Lambda$ and $p\Xi^-$ correlation function

Source size dependence



$d\Xi^-$ correlation function

- Three body system of $d(np)\Xi$

K. Ogata, T. Fukui, Y. Kamiya, and A. Ohnishi, arXiv:2103.00100

- Three body problem :
continuum-discretized
coupled-channels method (CDCC)

N. Austern, M. Yahiro, and M. Kawai, PRL 63 2649(1989)

N. Austern, M. Kawai, and M. Yahiro, PRC 53 314 (1996)

$\implies d\text{-}\Xi$ relative wave function

- $N\Xi$ interaction : HAL QCD potential

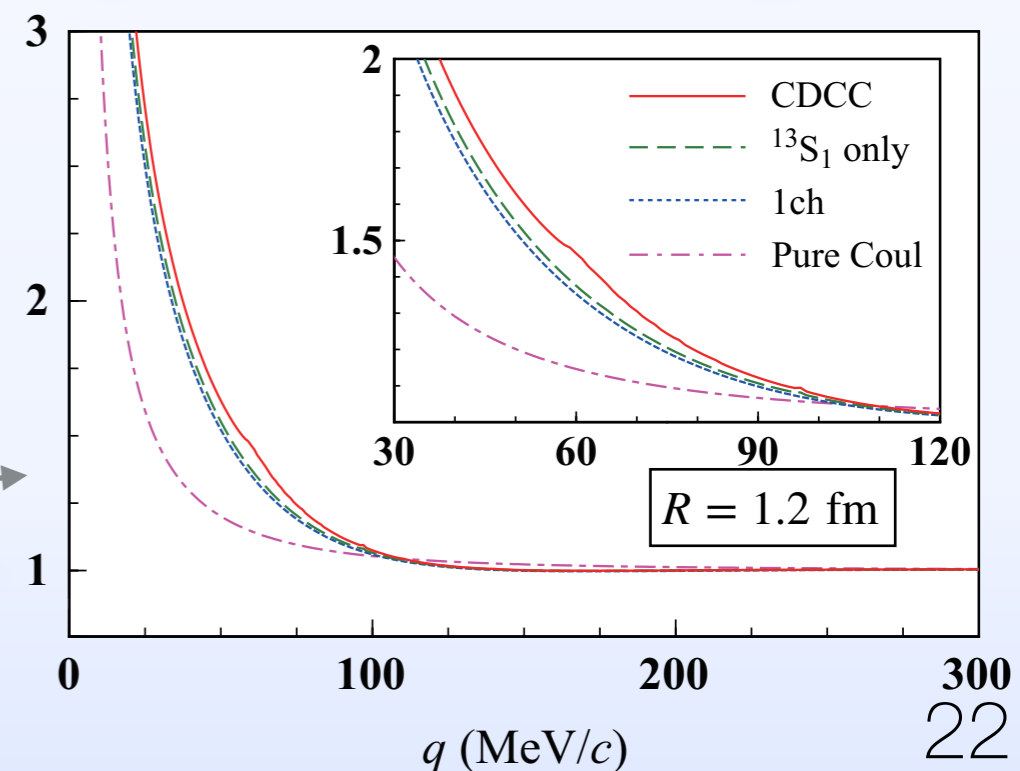
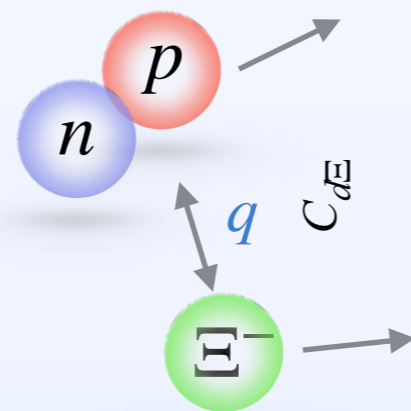
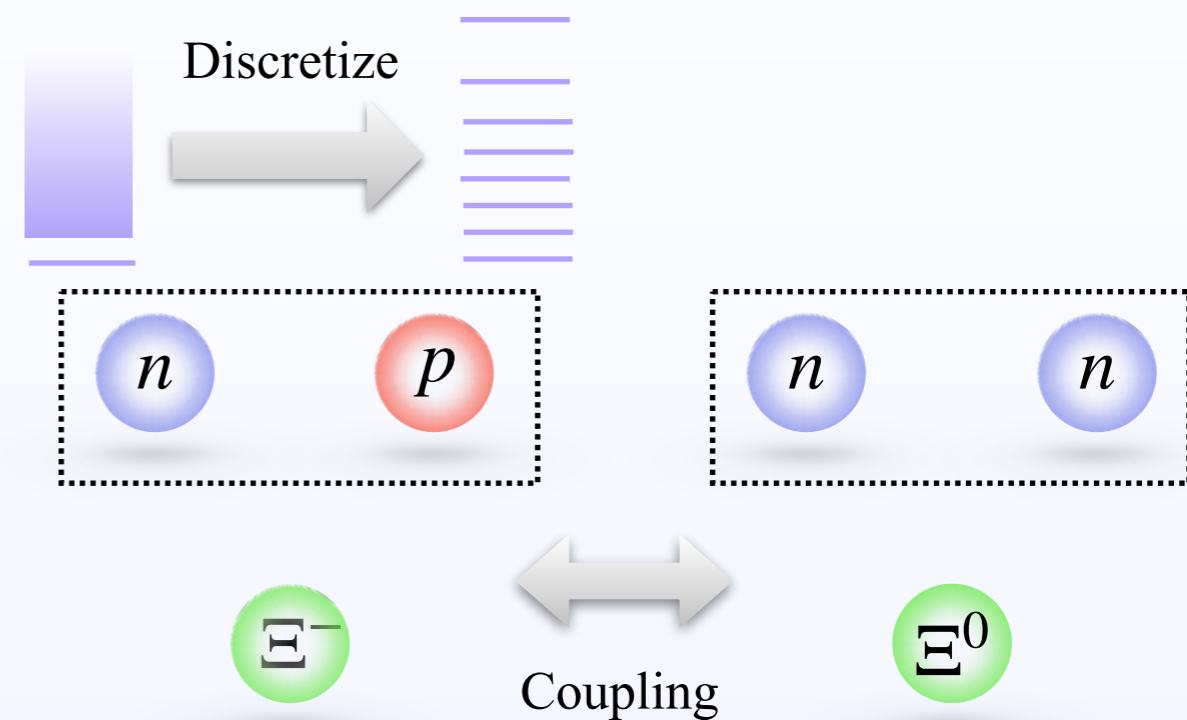
K. Sasaki et al. [HAL QCD], NPA 998 (2020), 121737.

- Coupling between $np\Xi^-$ - $nn\Xi^0$ included

- Theoretical model for $C_{d\Xi^-}$

- Strong enhancement compared to pure Coulomb case

- Coupling effect by $np\Xi^-$ - $nn\Xi^0$ is estimated to be 6–8 %



Summary

Summary

Y. Kamiya, K. Sasaki, T. Fukui, T. Hyodo, K. Morita, K. Ogata, A. Ohnishi, T. Hatsuda in prep.

- Femtoscopic correlation function in high energy nuclear collisions is a powerful tool to investigate the hadron-hadron interaction.
- With the latest HAL QCD potential, the H state does not exist as physical state but emerges as virtual pole in amplitude.
- $\Lambda\Lambda$ and $p\Xi^-$ correlation function is studied with including full coupled-channel effect and Coulomb interaction in the consistent manner. The result shows good agreement with ALICE data (pp and pPb collisions).
- For the further determination of interactions, to extract the source size dependence is important using other collisions systems.

Thank you for your attention!