

# Good and bad diquark properties and spatial correlations in lattice QCD

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*based on arXiv:2106.09080*

## Diquarks - an attractive concept

"The concept of diquarks is almost as old as the quark model, and actually predates QCD [1]"

↪ Snowmass '20, [1] PR 155, 1601 (1967)

- Successful for low-lying baryons and exotic hadrons.
  - But, experimental evidence has been elusive.
- Well founded in QCD with many predictions. Light quarks:
  - special "good" ( $\bar{3}_F, \bar{3}_C, J^P = 0^+$ ) configuration
  - "good" diquarks experience attraction effect
  - large mass splitting in good, bad and not-even-bad
  - non-vanishing size or compact?
- For heavy quarks, with HQSS, diquarks can act as single antiquarks  $[QQ] \leftrightarrow \bar{Q}$ .
  - ↪ opportunities for exotic hadrons, see B. Colquhoun on Tue.

good, bad and not-even

Diquark operator:

$$D_\Gamma = q^c C \Gamma q'$$

↪  $c, C$  = charge conjugation

↪  $\Gamma$  acts on Dirac space

$J^P$	C	F	Op: $\Gamma$
$0^+$	$\bar{3}$	$\bar{3}$	$\gamma_5, \gamma_0 \gamma_5$
$1^+$	$\bar{3}$	6	$\gamma_i, \sigma_{i0}$
$0^-$	$\bar{3}$	6	$\mathbb{1}, \gamma_0$
$1^-$	$\bar{3}$	$\bar{3}$	$\gamma_i \gamma_5, \sigma_{ij}$

towards a clearer understanding and footing in QCD

Goal: Measure diquark properties in QCD non-perturbatively

- **spectrum:** [diquark] mass differences are fundamental characteristics of QCD (Jaffe '05, arXiv:hep-ph/0409065)
- **spatial correlations:** study attraction and special status of the "good" diquark
- **structure:** estimate size and shape of the "good" diquark

## A gauge invariant probe - static quark as spectator

- A problem for the lattice is that diquarks are colored, i.e. not-gauge invariant.
  - Could fix a gauge, but then properties are gauge-dependent (masses, sizes,...)

↪ lattice and Dyson-Schwinger, see e.g. [15-20] in 2106.09080

- **Alternative:** Static spectator quark  $Q$  ( $m_Q \rightarrow \infty$ ) cancels in mass differences.
  - Diquark properties exposed in a gauge-invariant way.

↪ hep-lat/0510082, hep-lat/0509113, hep-lat/0609004, arxiv:1012.2353

$$C_\Gamma(t) \sim \exp \left[ -t \left( m_{D_\Gamma} + m_Q + \mathcal{O}(m_Q^{-1}) \right) \right]$$

↪  $t \rightarrow$  large,  $m_Q \rightarrow$  large

- **Lattice correlator:** Diquark embedded in a static-light-light baryon

$$C_\Gamma(t) = \sum_{\vec{x}} \langle [D_\Gamma Q](\vec{x}, t) [D_\Gamma Q]^\dagger(\vec{0}, 0) \rangle$$

↪ static quark=Q and  $D_\Gamma = q^c C \Gamma q$

↪ flavor combinations  $ud, \ell s, ss'$

↪ static-light mesons  $[\bar{Q} \Gamma q]$

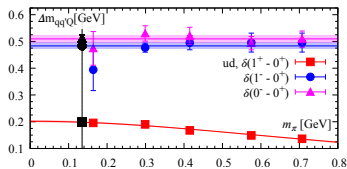
setting up on the lattice - we recycle

- $n_f = 2 + 1$  full QCD,  $32^3 \times 64$ ,  $a = 0.090\text{fm}$ ,  $a^{-1} = 2.194\text{GeV}$  (PACS-CS gauges)
- $m_\pi = 164, 299, 415, 575, 707 \text{ MeV}$ ,  $m_s \simeq m_s^{\text{phys}}$ , propagators re-used from before
- Quenched gauge  $a \simeq 0.1\text{fm}$ ,  $m_\pi^{\text{valence}} = 909 \text{ MeV}$ , to match hep-lat/0509113

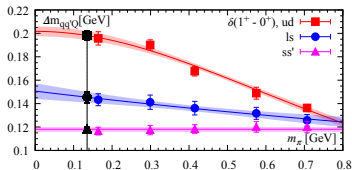
## *I. Diquark spectroscopy*

# Lattice spectroscopy - diquark-diquark differences

$ud$   $0^+$  versus  $1^+$ ,  $0^-$  and  $1^-$



$(1^+ - 0^+)_{qq'}$  splitting



We consider differences of  $qq'Q$  baryons:

$$C_{\Gamma}^{qq'Q}(t) - C_{\gamma_5}^{qq'Q}(t)$$

$\rightsquigarrow Q$  drops out

$\rightsquigarrow$  measures diquark-diquark mass difference

Bad-good diquark splitting:

- o Special status of good diquark observed
- o Good  $0^+$   $ud$  diquark lies lowest in the spectrum
- o Bad  $1^+$   $ud$  diquark 100-200 MeV above
- o  $0^-$  and  $1^-$   $ud$  diquarks  $\sim 0.5$  GeV above
- o Pattern repeated in  $ls$  and  $ss'$

$\Delta m_{qq'Q}(m_\pi)$  dependence:

- o Chiral limit:  $\sim \text{const}$
- o Heavy-quark limit: decreases  $\sim 1/(m_{q_1} m_{q_2})$ , with  $m_\pi \sim (m_{q_1} + m_{q_2})$

$$\delta(1^+ - 0^+)_{q_1 q_2} = A / \left[ 1 + (m_\pi / B)^{n \in \{0, 1, 2\}} \right]$$

## Lattice spectroscopy - diquark-quark differences

We consider differences of a  $qq'Q$  baryon and a light-static meson:

$$C_{\Gamma=\gamma_5}^{qq'Q}(t) - C_{\gamma_5}^{q'\bar{Q}}(t)$$

$\rightsquigarrow Q$  drops out  
 $\rightsquigarrow$  diquark-quark mass difference

$\Delta m_{qq'Q}(m_\pi)$  dependence:

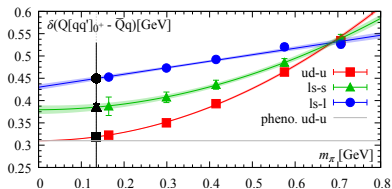
- Chiral vs. heavy-quark limiting behaviours, as before

$$\delta(Q[q_1q_2]_{0^+} - \bar{Q}q_2) = C [1 + (m_\pi/D)^{n \in \{0,1,2\}}]$$

Diquark-quark splitting:

- Established relative masses between a good diquark and an [anti]quark
- May prove useful in identifying favourable tetra-, pentaquark channels
- Omits possible distortions through additional light quarks, Pauli-blocking, spin-spin interactions ...

### $Qqq' - \bar{Q}q'$ splittings



## Diquark spectroscopy - phenomenological estimates

We want to compare our results with phenomenology

- Key resource: (Jaffe '05, arXiv:hep-ph/0409065), updated with PDG 2021 input
- For pheno estimates use charm and bottom hadron masses where leading  $\mathcal{O}(1/m_Q)$  ( $Q = c, b$ ) can be cancelled

Four estimates considered:

- $\delta(1^+ - 0^+)_{ud}$ : 
$$\frac{1}{3} (2M(\Sigma_Q^*) + M(\Sigma_Q)) - M(\Lambda_Q)$$

- $\delta(1^+ - 0^+)_{us}$ : 
$$\frac{2}{3} (M(\Xi_Q^*) + M(\Sigma_Q) + M(\Omega_Q)) - M(\Xi_Q) - M(\Xi_Q')$$

- $\delta(Q[ud]_{0^+} - \bar{Q}u)$ : 
$$M(\Lambda_Q) - \frac{1}{4} (M(P_{Qu}) + 3M(V_{Qu}))$$

$\rightsquigarrow P_{Qu}, V_{Qu}$  are the ground-state, heavy-light mesons

- $\delta(Q[us]_{0^+} - \bar{Q}s)$ :

$$M(\Xi_Q) + M(\Xi_Q') - \frac{1}{2} (M(\Sigma_Q) + M(\Omega_Q)) - \frac{1}{4} (M(P_{Qs}) + 3M(V_{Qs}))$$

$\rightsquigarrow P_{Qs}, V_{Qs}$  are the ground-state, heavy-strange mesons

## Diquark spectroscopy - comparing results

- We summarise the main spectroscopy results as:

All in [MeV]	$\delta E_{\text{lat}}(m_{\pi}^{\text{phys}})$	$\delta E_{\text{pheno}}$	$\delta E_{\text{pheno}}^{\text{bottom}}$	$\delta E_{\text{pheno}}^{\text{charm}}$
$\delta(1^+ - 0^+)_{ud}$	198(4)	206(4)	206	210
$\delta(1^+ - 0^+)_{\ell s}$	145(5)	145(3)	145	148
$\delta(1^+ - 0^+)_{ss'}$	118(2)			
$\delta(Q[ud]_{0^+} - \bar{Q}u)$	319(1)	306(7)	306	313
$\delta(Q[\ell s]_{0^+} - \bar{Q}s)$	385(9)	397(1)	397	398
$\delta(Q[\ell s]_{0^+} - \bar{Q}\ell)$	450(6)			

↪ updated pheno using PDG '21

↪ use the bottom estimate for static

↪ use charm-bottom difference as estimate for deviation from static

⇒  $\lesssim \mathcal{O}(7)\text{MeV}$  deviation

- Overall, very good agreement observed.



## *II. Diquark structure*

## Diquarks - spatial correlations

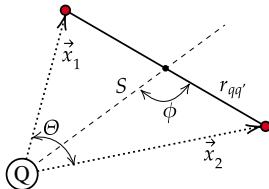
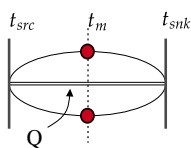
We access (good) diquark structure information through density-density correlations:

$$C_{\Gamma}^{dd}(\vec{x}_1, \vec{x}_2, t) = \left\langle \mathcal{O}_{\Gamma}(\vec{0}, 2t) \rho(\vec{x}_1, t) \rho(\vec{x}_2, t) \mathcal{O}_{\Gamma}^{\dagger}(\vec{0}, 0) \right\rangle$$

$$\rightsquigarrow \mathcal{O}_{\Gamma} = q^c C \Gamma q \text{ and } \rho(\vec{x}, t) = \bar{q}(\vec{x}, t) \gamma_0 q(\vec{x}, t)$$

$$\rightsquigarrow t_m = (t_{snk} + t_{src})/2 \text{ to minimize excited states}$$

Main tool: Correlations between two light quarks' relative positions to the static quark



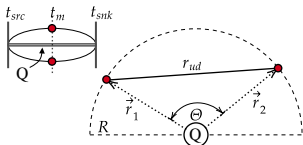
$$\rightsquigarrow \vec{r}_{ud} = \vec{x}_2 - \vec{x}_1 \text{ and } \vec{S} = (\vec{x}_1 - \vec{x}_2)/2$$

$$\rho_2(r_{ud}, S, \phi; \Gamma) = C_{\Gamma}^{dd}(\vec{x}_1, \vec{x}_2, t_m)$$

Note, when  $S$  and  $r_{ud}$  fixed, distance between static quark  $Q$  and light quarks  $q, q'$  is

- Minimized for  $\phi = \pi$ , possible disruption due to  $Q$  is largest
- Maximized for  $\phi = \pi/2$ , possible disruption due to  $Q$  is smallest

# Good diquark attraction



Setting  $\phi = \pi/2$ :

- $|\vec{x}_1| = |\vec{x}_2| = R$ , use  $R, \Theta$ :

$$\rho_2^\perp(R, \Theta) = \rho_2(r_{ud}, S, \pi/2)$$

- Attraction visible through increase in  $\rho_2^\perp$  for small  $\Theta$  at any fixed  $R$

Two limiting cases for the two quarks:

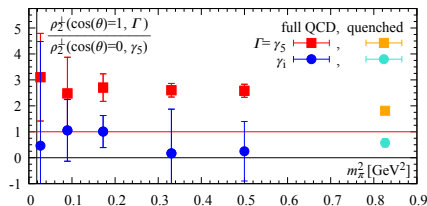
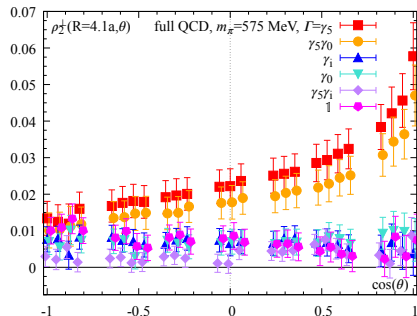
- $\cos(\Theta) = 1$  on top of each other
- $\cos(\Theta) = -1$  opposite each other

"Lift" as qualitative criterion:

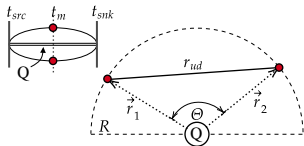
$$\frac{\rho_2^\perp(R, \Theta = 0, \Gamma)}{\rho_2^\perp(R, \Theta = \pi/2, \gamma_5)}$$

Increase observed in good diquark only

## Spatial correlation over $\Theta$



## Good diquark size



- Distance between quarks:

$$r_{ud} = R\sqrt{2(1 - \cos(\Theta))}$$

~> different visualisation

- $\rho_2^\perp(R, r_{ud}) \sim \exp(-r_{ud}/r_0)$   
~> "characteristic size"  $r_0$

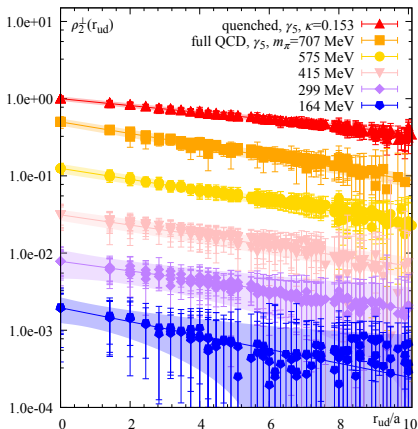
- Need to control:

- interference from  $Q$   
~> we limit analysis to  $r_{ud} < R$
- periodicity effects  
~> in practice we find  $L = 5r_0$

- Further checks:

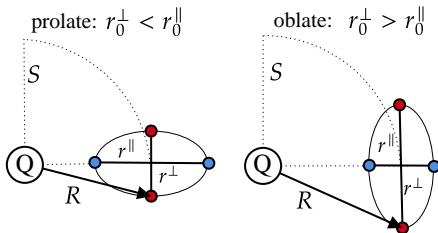
$$A(R, r_{ud} = 0) \sim \exp(-R/R_0)$$

## Spatial correlation over $r_{ud}$



- $r_{ud} = 0$  normalised, offset for each  $m_\pi$
- all  $R$  shown simultaneously
- combined fits over  $\forall R$  with shared  $r_0$

## Shape of good diquarks - studying oblateness



### Tangential and radial spatial correlation decay

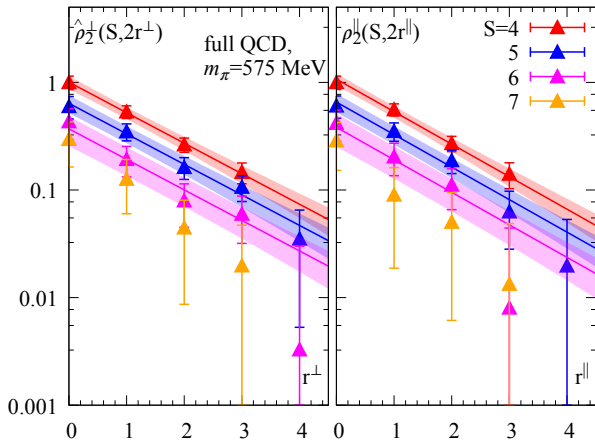
As opposed to before  $R \neq$  fixed:

- $\phi = \pi$ : radial correlation,  
size  $\rightsquigarrow r_0^\parallel$
- $\phi = \pi/2$ : tangential correlation,  
size  $\rightsquigarrow r_0^\perp$

- $r_0^\perp / r_0^\parallel$  gives information on shape:  
= 1, spherical  
 $\neq 1$ , prolate/oblate

- Probe  $J = 0$  nature of good diquark
- Diquark polarisation through static quark?

# Oblateness - results



Goal:

- $r_0^\perp, r_0^\parallel$  at fixed  $S$

Technical issue:

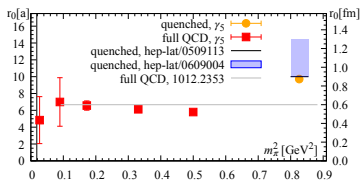
- ( $\parallel$ ) as before:  
 $R = S$
- ( $\perp$ ) different:  
 $R = \sqrt{(r^\perp)^2 + S^2}$

Solution:

- Introduce "nuisance" parameter  $R_0$
- Adjusted in figure
- Parallel lines  $\rightsquigarrow r_0^\perp = r_0^\parallel$

# Diquark structure - overview

## Size dependence $r_0(m_\pi)$



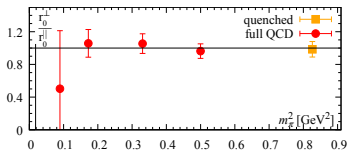
Good diquark size:

- Agreement w/ prev. quenched and dynamical
- Refinement through our results
- $r_0 \simeq \mathcal{O}(0.6)\text{fm}$  weak  $m_\pi$  dependence  
 $\rightsquigarrow \sim r_{\text{meson, baryon}}$ , arXiv:1604.02891

$r_0(m_\pi)$  dependence:

- $m_{q,q'} \uparrow$  should produce more compact object
- But, diquark attraction  $\downarrow$  works opposite
- Former effect dominates at large  $m_\pi$ ?
- But, in quenched diquarks definitely larger...

## Shape dependence $r_0^\perp / r_0^\parallel(m_\pi)$



$r_0^\perp / r_0^\parallel(m_\pi)$  dependence:

- Ratio  $\simeq 1$  for all  $m_\pi$
- Consistent w/ scalar,  $J = 0$ , shape
- No diquark polarisation through  $Q$  observed

# Summary

## Gauge invariant approach to diquarks in $n_f = 2 + 1$ lattice QCD

- Lattice setup with short chiral extrapolations, continuum limit still required

## Diquark spectroscopy

- Special status of "good" diquark confirmed, attraction of 198(4)MeV over "bad"
- Chiral and flavor dependence modelled through simple Ansatz
- Very good agreement with phenomenological estimates

## Diquark structure

- $q - q$  attraction in good diquark induces compact spatial correlation
- Good diquark size  $r_0 \simeq \mathcal{O}(0.6)\text{fm} \sim r_{\text{meson, baryon}}$ , weakly  $m_\pi$  dependent
- Good diquark shape appears nearly spherical

## Outlook

- Results provide clear, quantitative support for the good diquark picture
- Hope to refine diquark model parameters
- Insights for studies of exotic tetraquarks (esp. doubly heavy), heavy-baryons, etc.
- In future, extend towards tetraquarks, finite temperature ...



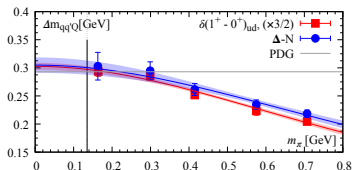
*Thank you for your attention.*



*Further material*

## $\Delta$ -Nucleon mass difference

$[\Delta - N](m_\pi)$



Measured the mass difference of  $\Delta - N$

- Prediction:  $\delta(\Delta - N) = 3/2 \times \delta(1^+ - 0^+)_{ud}$
- Same Ansatz as before
- Prediction holds well, even at fairly large  $m_\pi$