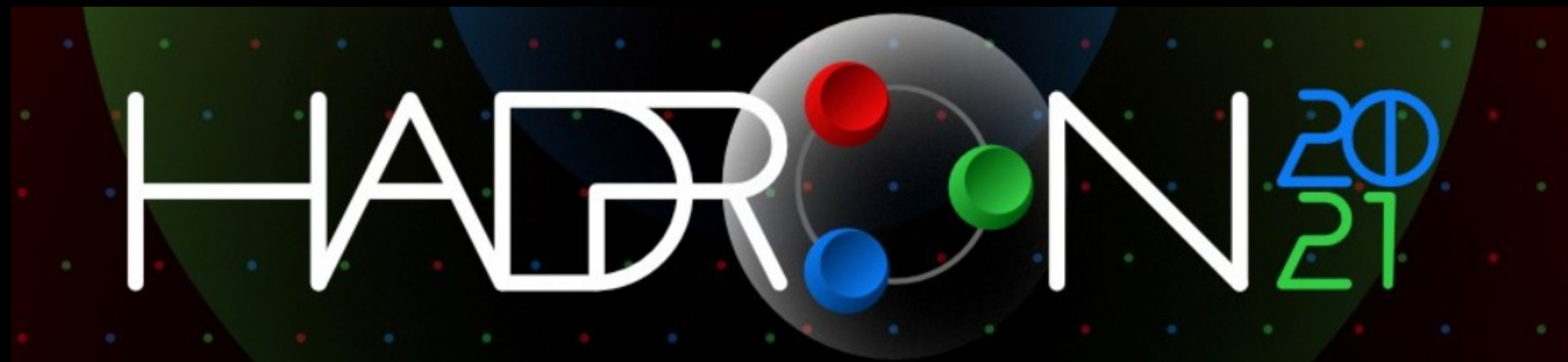
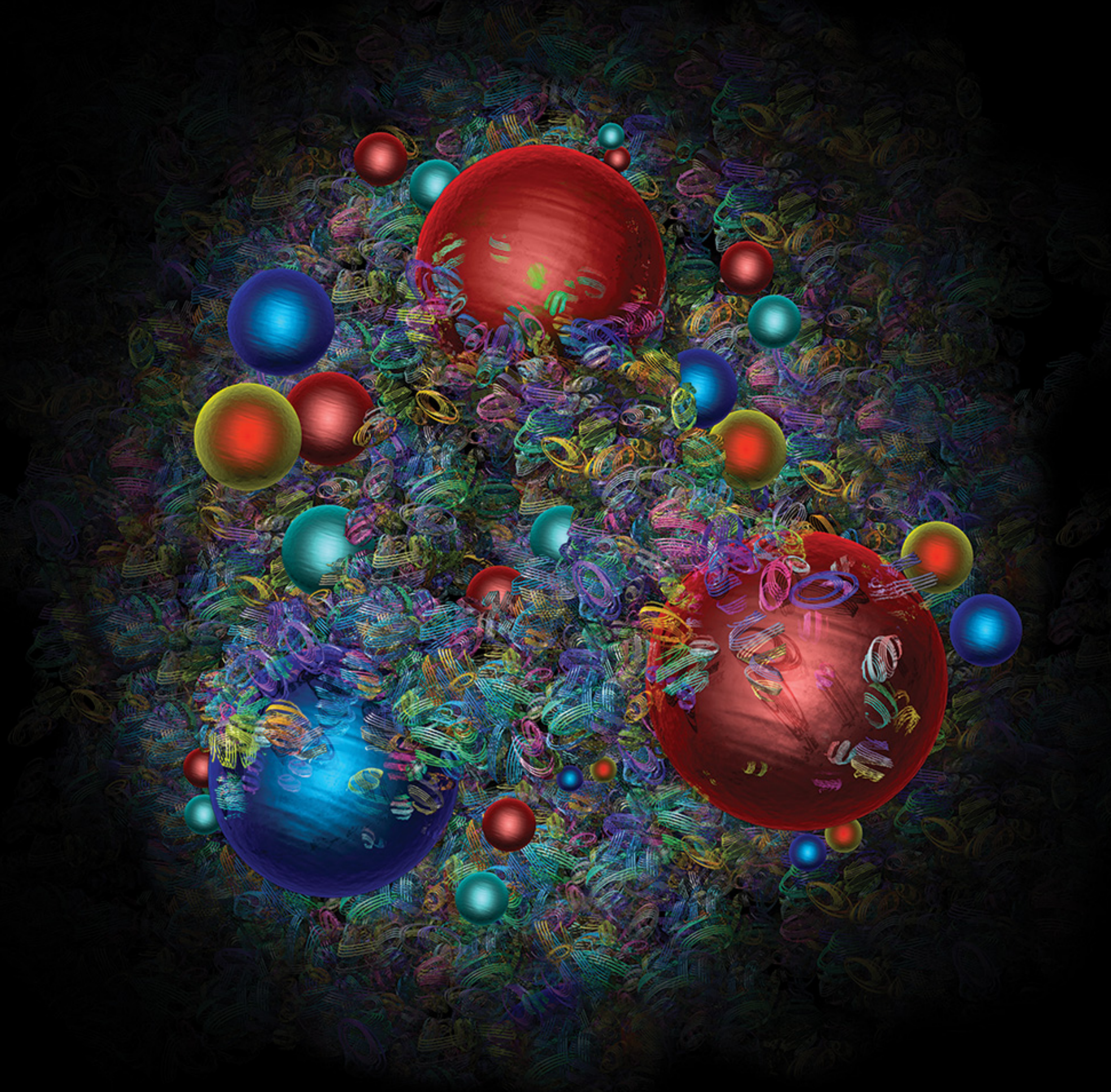


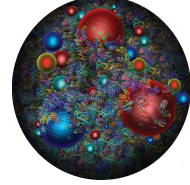
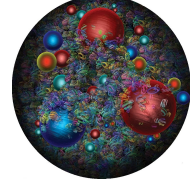
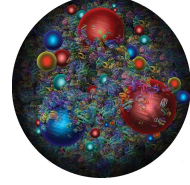
Hadron structure at small- x via unintegrated gluon densities

Michael Fucilla

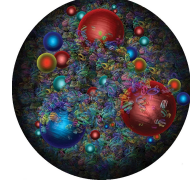
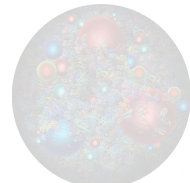
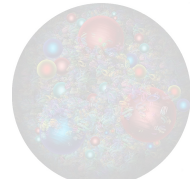
**Università della Calabria & INFN-gruppo collegato di
Cosenza**



Outline

-  **BFKL and unintegrated gluon densities**
-  **Exclusive forward meson lepton production**
-  **Conclusions and outlook**

Outline

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-  Conclusions and outlook

High-energy factorization and the UGD

Balitsky-Fadin-Kuraev-Lipatov (BFKL) resummation

* Leading-Logarithm-Approximation (LLA): $(\alpha_s \ln s)^n$

* Next-to-Leading-Logarithm-Approximation (NLLA): $\alpha_s(\alpha_s \ln s)^n$

Unintegrated gluon densities

* Definition $\mathcal{F}(x, \vec{k}), \quad f_g(x, Q^2) = \int \frac{d^2 \vec{k}}{\pi \vec{k}^2} \mathcal{F}(x, \vec{k}) \theta(Q^2 - \vec{k}^2)$

* Evolution equation as a function of $\ln(s/Q^2) = \ln(1/x)$

$$\frac{\partial \mathcal{F}}{\partial \ln(1/x)} = \mathcal{F} \otimes \mathcal{K}$$

High-energy factorization and the UGD

● Deep inelastic scattering

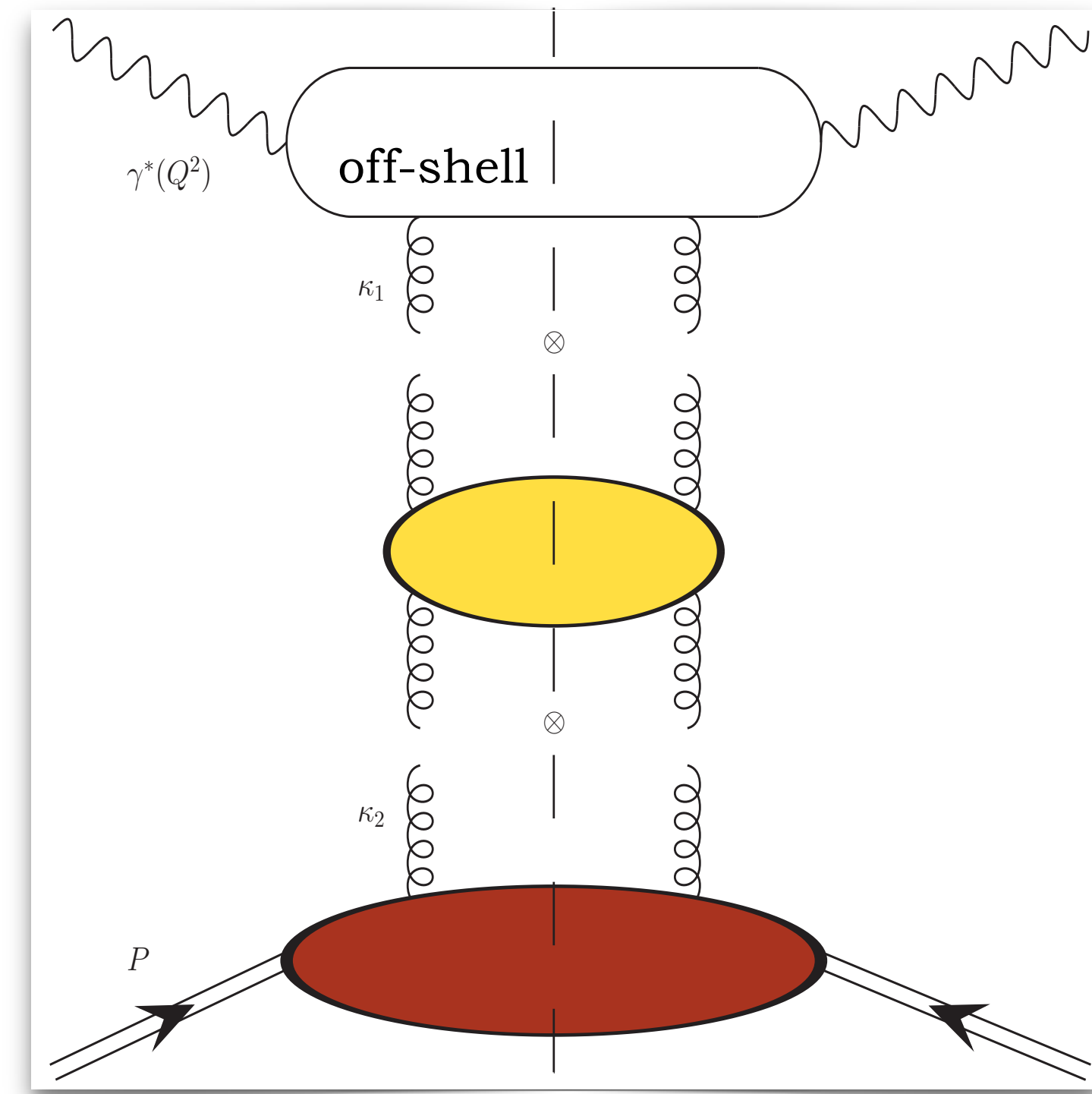
* Total cross section

$$\sigma_\lambda(x, Q^2) = \frac{\mathcal{G}}{(2\pi)^4} \int \frac{d^2\vec{k}_1}{k_1^2} \int \frac{d^2\vec{k}_2}{k_2^2} \Phi_\lambda(\vec{k}_1) F(x, \vec{k}_1, \vec{k}_2) \Phi_p(\vec{k}_2)$$

$$F(x, \vec{k}_1, \vec{k}_2) = \sum_{n=0}^{\infty} \int_{-\infty}^{\infty} d\nu \left(\frac{\vec{k}_1^2}{\vec{k}_2^2} \right)^{i\nu} \frac{e^{in(\theta_1 - \theta_2)}}{2\pi |\vec{k}_1| |\vec{k}_2|} e^{\bar{\alpha}_s \chi_n(\nu) \ln\left(\frac{1}{x}\right)}$$

* Growth at small-x

$$F \sim \frac{x^{-\omega_0}}{\sqrt{\ln(1/x)}} \quad \omega_0 = 4\bar{\alpha}_s \ln 2$$



$\Phi_{\gamma^* \rightarrow \gamma^*}$

\otimes

$\mathcal{G}_{\text{BFKL}}$

\otimes

$\Phi_{[\text{NP}]}$

High-energy factorization and the UGD

IR-safe colorless $\{\Phi^{i \rightarrow 0}\}$

(Fadin-Martin theorem)

[V.S. Fadin, A.D. Martin (1999)]

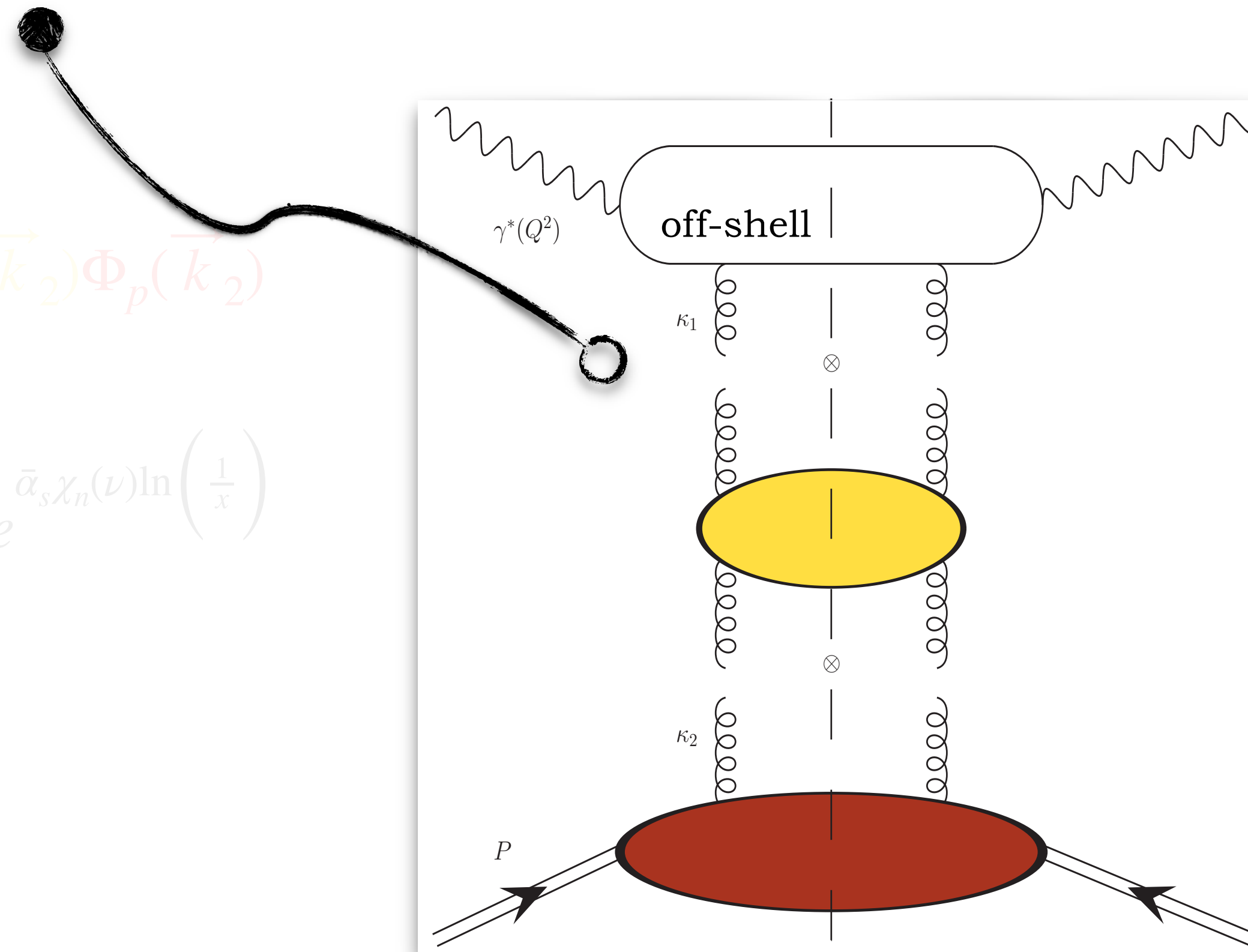
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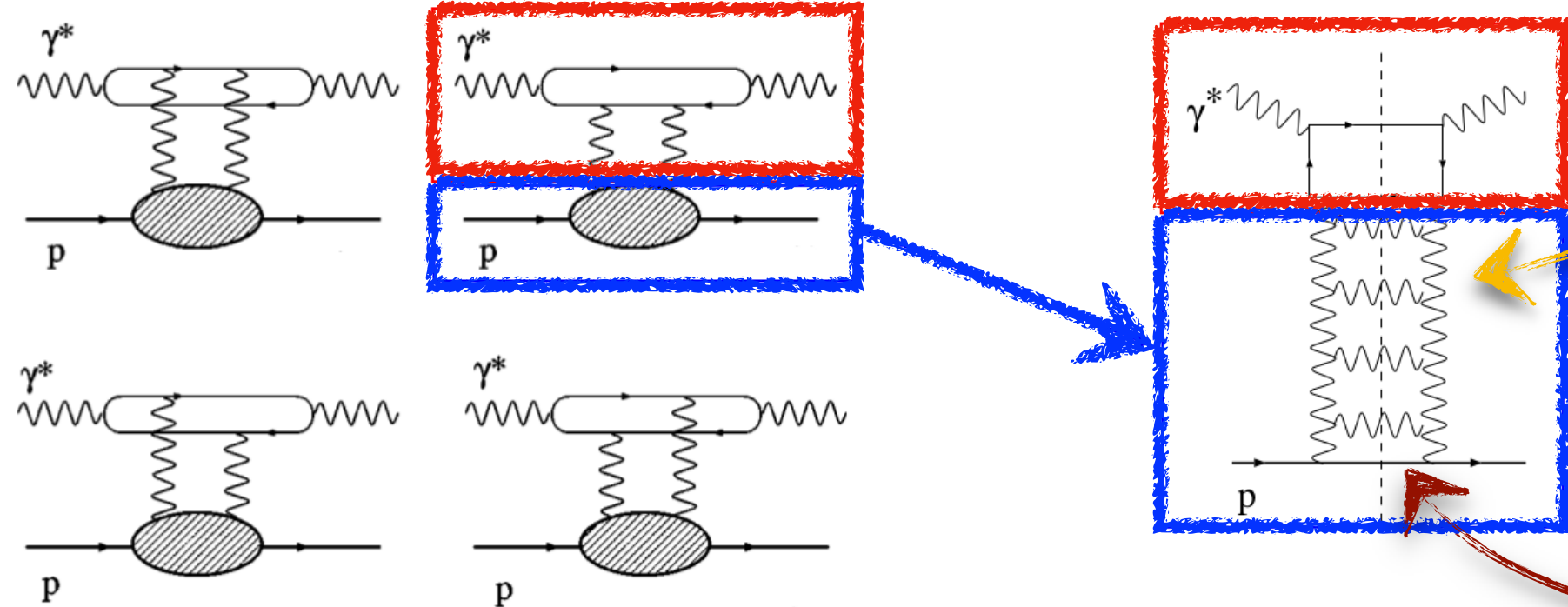
$\Phi^P_{[\text{NP}]}$

High-energy factorization and the UGD

Unintegrated gluon density

$$\sigma_\lambda(x, Q^2) = \frac{\mathcal{G}}{(2\pi)^4} \int \frac{d^2\vec{k}_1}{k_1^2} \int \frac{d^2\vec{k}_2}{k_2^2} \Phi_\lambda(\vec{k}_1) F(x, \vec{k}_1, \vec{k}_2) \Phi_p(\vec{k}_2) = \frac{\mathcal{G}}{(2\pi)^4} \int \frac{d^2\vec{k}_1}{k_1^4} \Phi_\lambda(\vec{k}_1) \mathcal{F}(x, \vec{k}_1)$$

$$\mathcal{F}(x, \vec{k}_1) \equiv \frac{k_1^2}{(2\pi)^3} \int \frac{d^2\vec{k}_2}{k_2^2} \Phi_p(\vec{k}_2) F(x, \vec{k}_1, \vec{k}_2)$$



- example: **virtual photoabsorption** in **high-energy factorization**

$$\sigma_{\text{tot}}(\gamma^* p \rightarrow X) \propto \text{Im}_s \{ \mathcal{A}(\gamma^* p \rightarrow \gamma^* p) \} \equiv \Phi_{\gamma^* \rightarrow \gamma^*} \otimes \mathcal{F}(x, \kappa^2)$$

- ◇ $\mathcal{F}(x, \kappa^2)$ is the **unintegrated gluon distribution (UGD)** in the proton

- ▶ Small- x limit: **UGD** = [**BFKL gluon ladder**] \otimes [**proton impact factor**]

- ◇ Takes into account the **resummation** of **high-energy logs**
- ◇ Describes the **coupling** of the gluon Green's function to the **proton**

- ▶ Proton impact factor is non-perturbative \implies UGD needs to be modeled!

High-energy factorization and the UGD

Features

- * Small- x and large k_t
- * Speaks the language of Reggeized gluon
- * Inclusive or exclusive processes
- * Double-log-approximation (DLA): $\alpha_s \ln(Q^2/Q_0^2) \ln(1/x)$

BFKL \longleftrightarrow *DGLAP*

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Unitarity violation and diffusion

- * Violation of the Martin-Froissart bound: $\sigma_{\text{tot}} \leq \frac{\pi \Delta^2}{2m_\pi^2} \ln^2 s$
- * Diffusion to the infrared: $l_\perp e^{-k\sqrt{\Delta Y/2}} \lesssim k_\perp \lesssim l_\perp e^{k\sqrt{\Delta Y/2}}$

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Non-linear evolution



BK/JIMWLK domain

High-energy factorization and the UGD

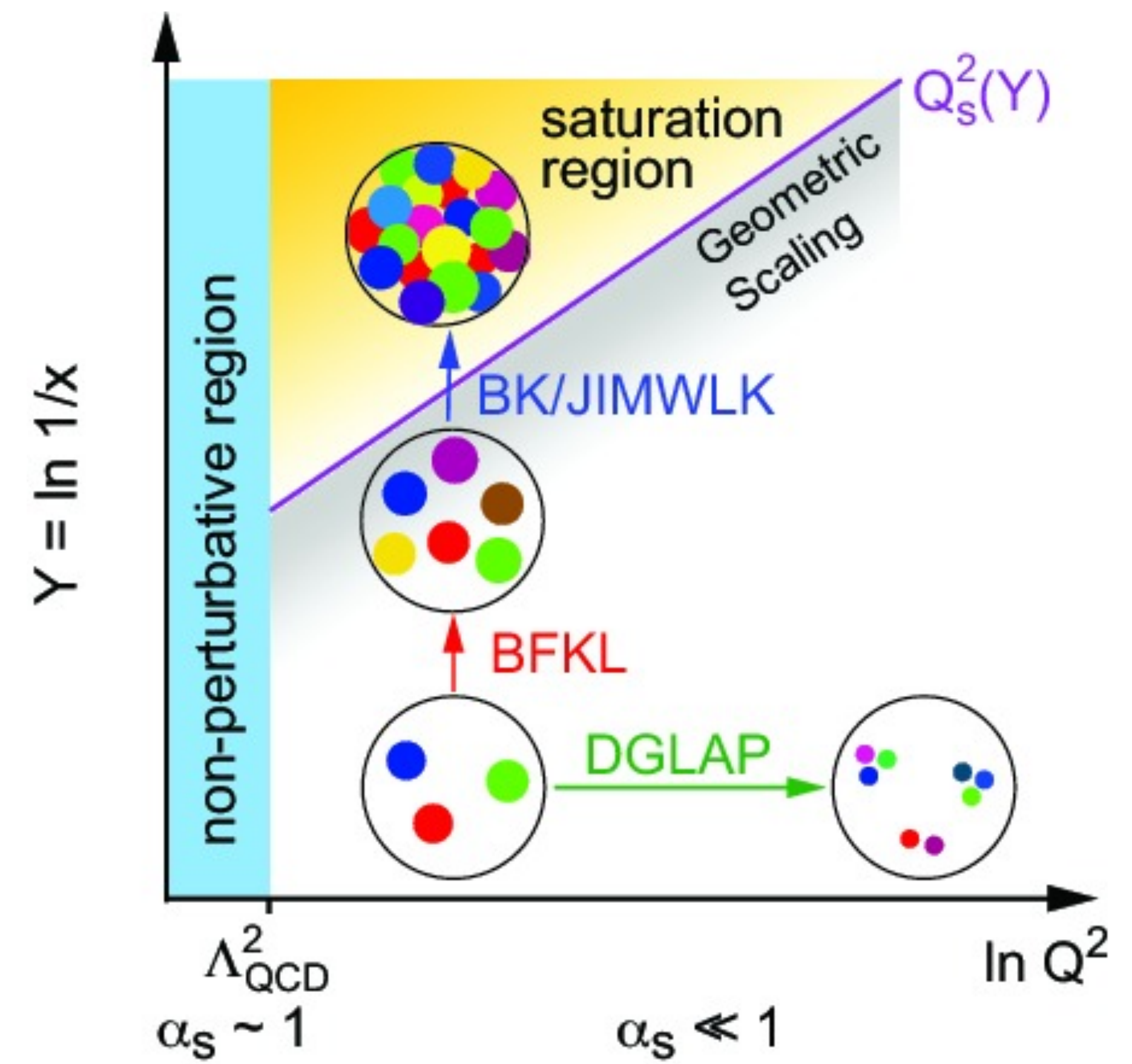
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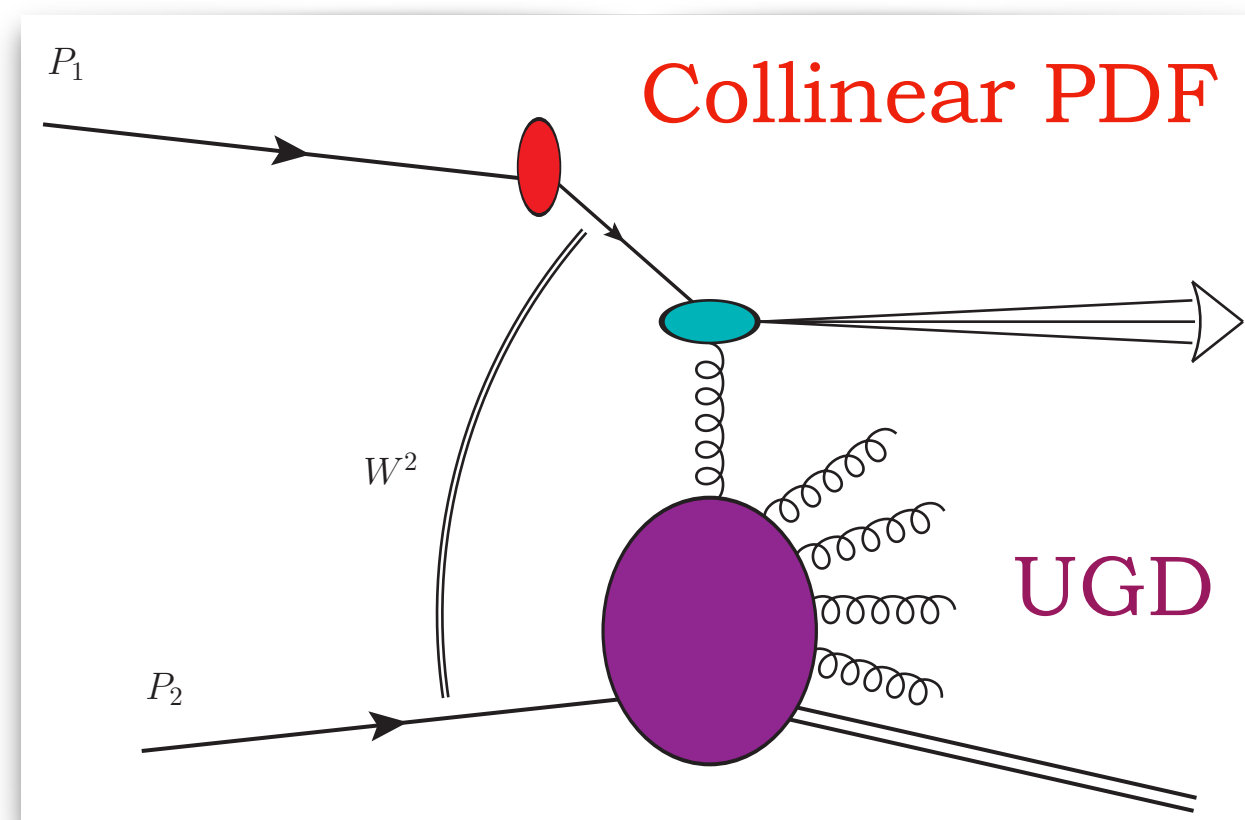


BK/JIMWLK domain

Hybrid or pure factorization?

Forward emissions

- * *Asymmetric* config. \leftrightarrow fast parton + small- x gluon
- * Hybrid **high-energy/collinear** factorization

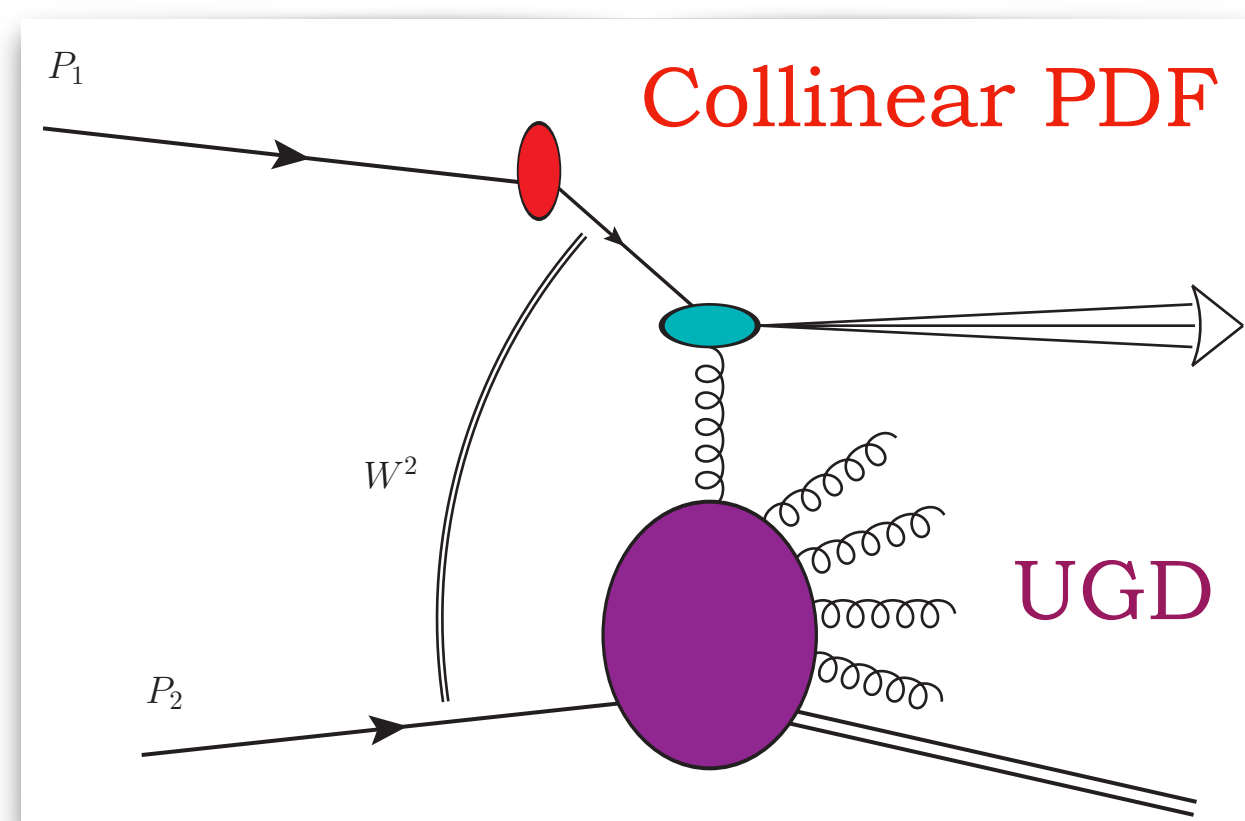


- * *Distinctive signals* of small- x dynamics **expected**
- * Phenomenology:
forward jet, Drell-Yan, Higgs or vector meson

Hybrid or pure factorization?

Forward emissions

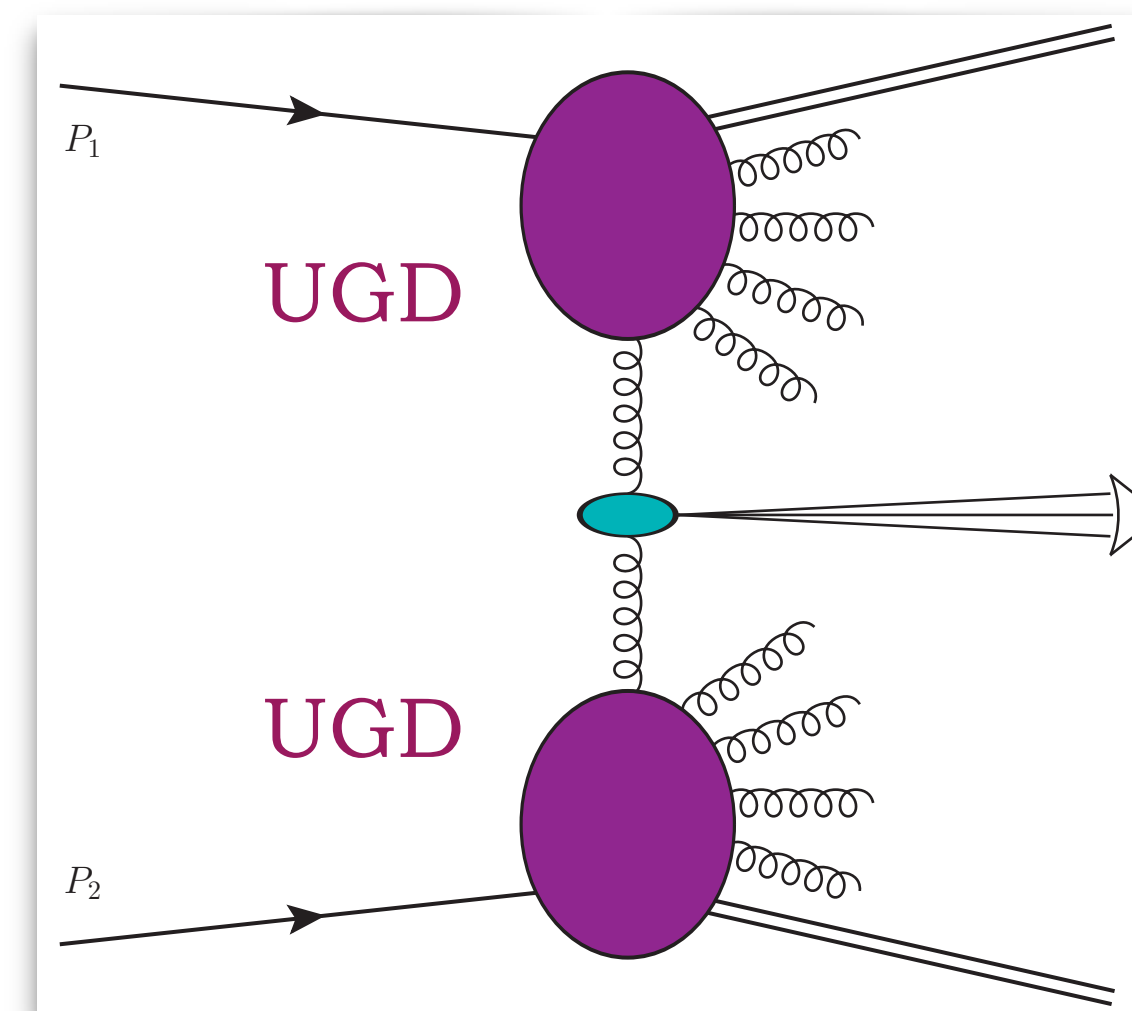
- * *Asymmetric* config. \leftrightarrow fast parton + small- x gluon
- * Hybrid **high-energy/collinear** factorization



- * *Distinctive signals* of small- x dynamics **expected**
- * Phenomenology:
forward jet, Drell-Yan, Higgs or vector meson

Central emissions

- * *Gluon induced* \leftrightarrow small- x gluons
- * Pure **high-energy** factorization



- * Small- x dynamics to **enhance** f.o. description
- * Phenomenology:
central jet, Higgs or vector meson

Outline

- BFKL and unintegrated gluon densities
- **Exclusive forward meson lepton production**
- Conclusions and outlook

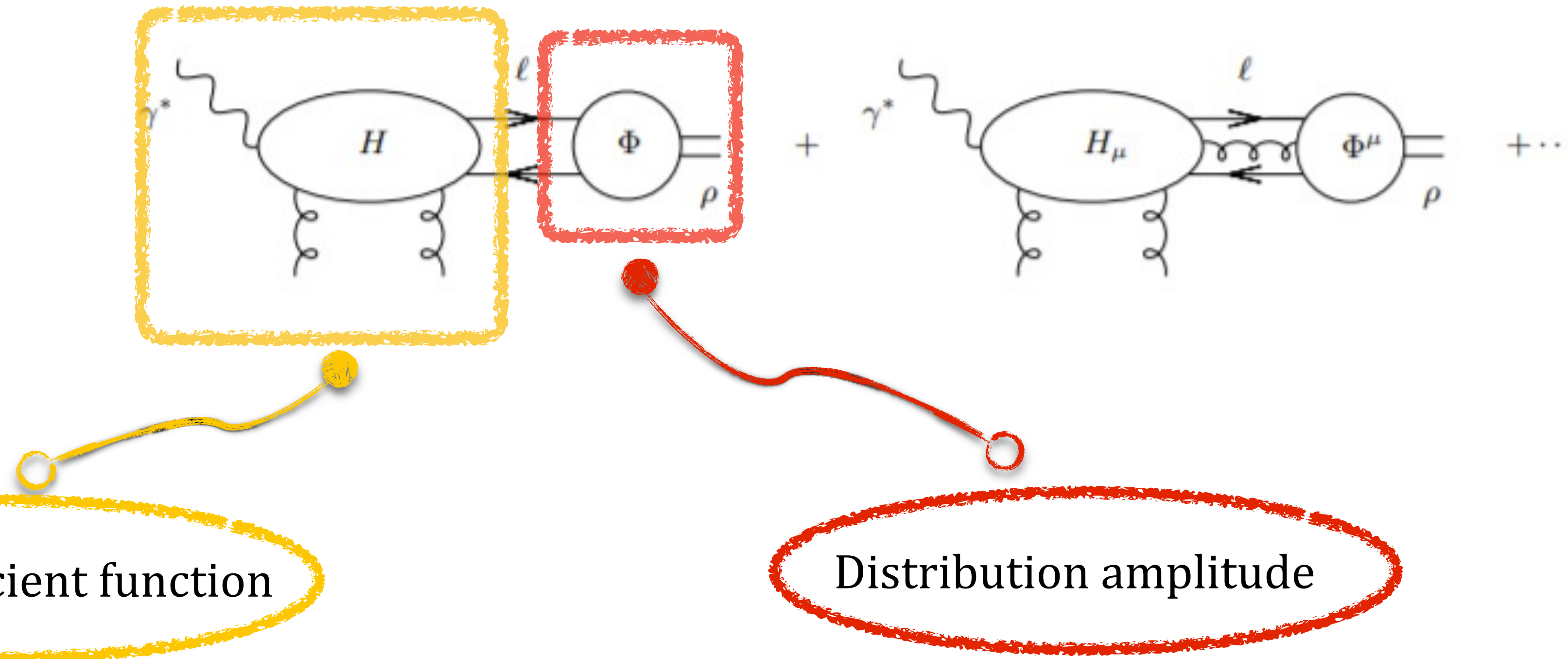
ρ -meson leptonproduction

[I. V. Anikin, D. Yu. Ivanov, B. Pire, L. Szymanowski and S. Wallon(2011)]

Impact factor

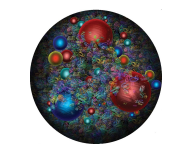
* Amplitude twist-expansion

$$A = \int d^4l \text{Tr}[H(l)\Phi(l)] + \int d^4l_1 \int d^4l_2 \text{Tr}[H_\mu(l_1, l_2)\Phi^\mu(l_1, l_2)] + \dots$$



* Amplitude factorization achieved by a Taylor expansion of the hard part

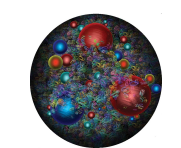
ρ -meson leptonproduction



Longitudinal case: $\gamma_L \rightarrow \rho_L$

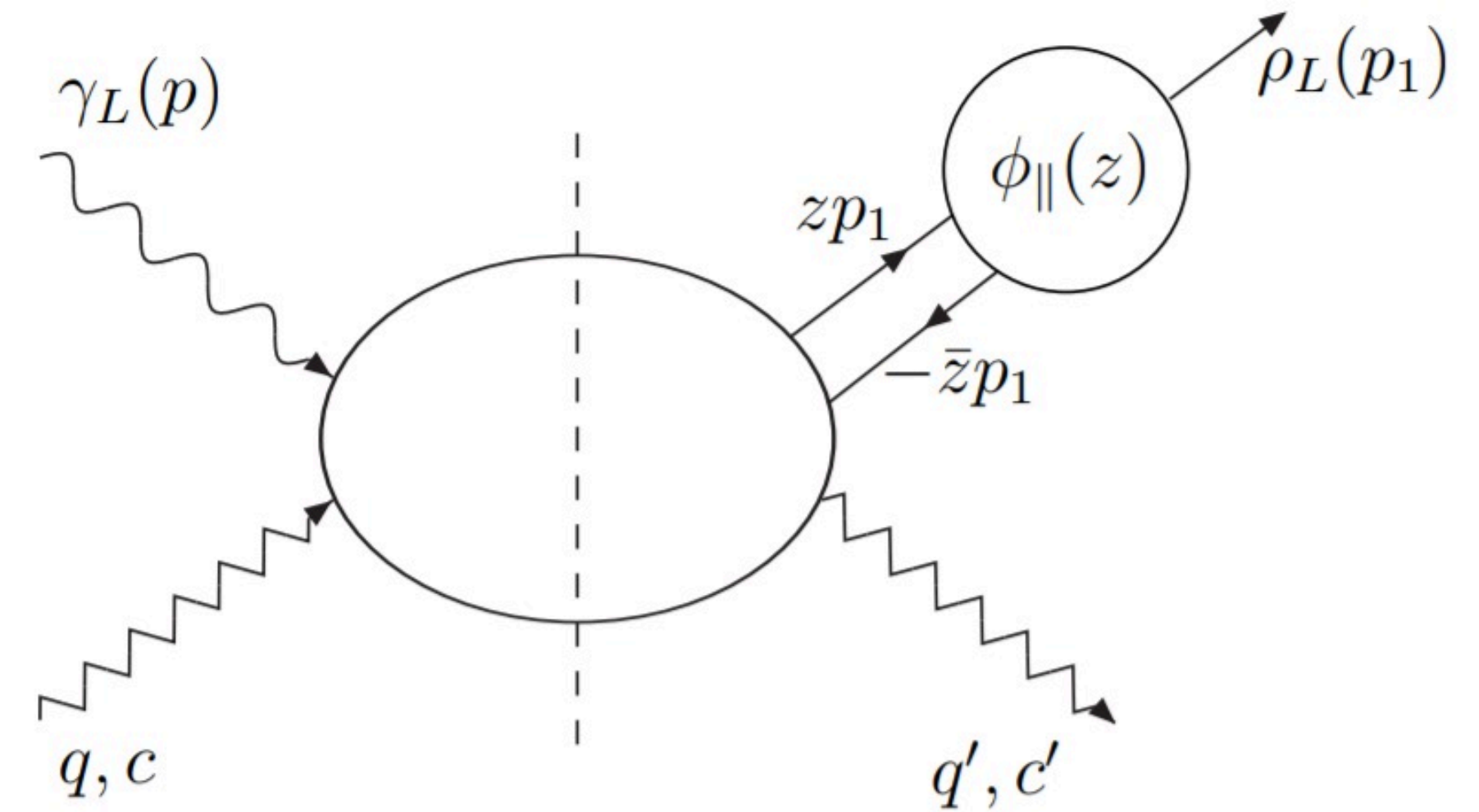
- * Starts the leading twist (twist two)
- * Known up to next-to-leading order
- * LO expression

$$\Phi_{\gamma_L \rightarrow \rho_L}(k, Q, \mu^2) = 2B \frac{\sqrt{N_c^2 - 1}}{QN_c} \int_0^1 dy \varphi_1(y; \mu^2) \left(\frac{\alpha}{\alpha + y\bar{y}} \right)$$



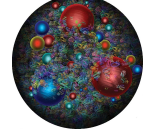
Transverse case: $\gamma_T \rightarrow \rho_T$

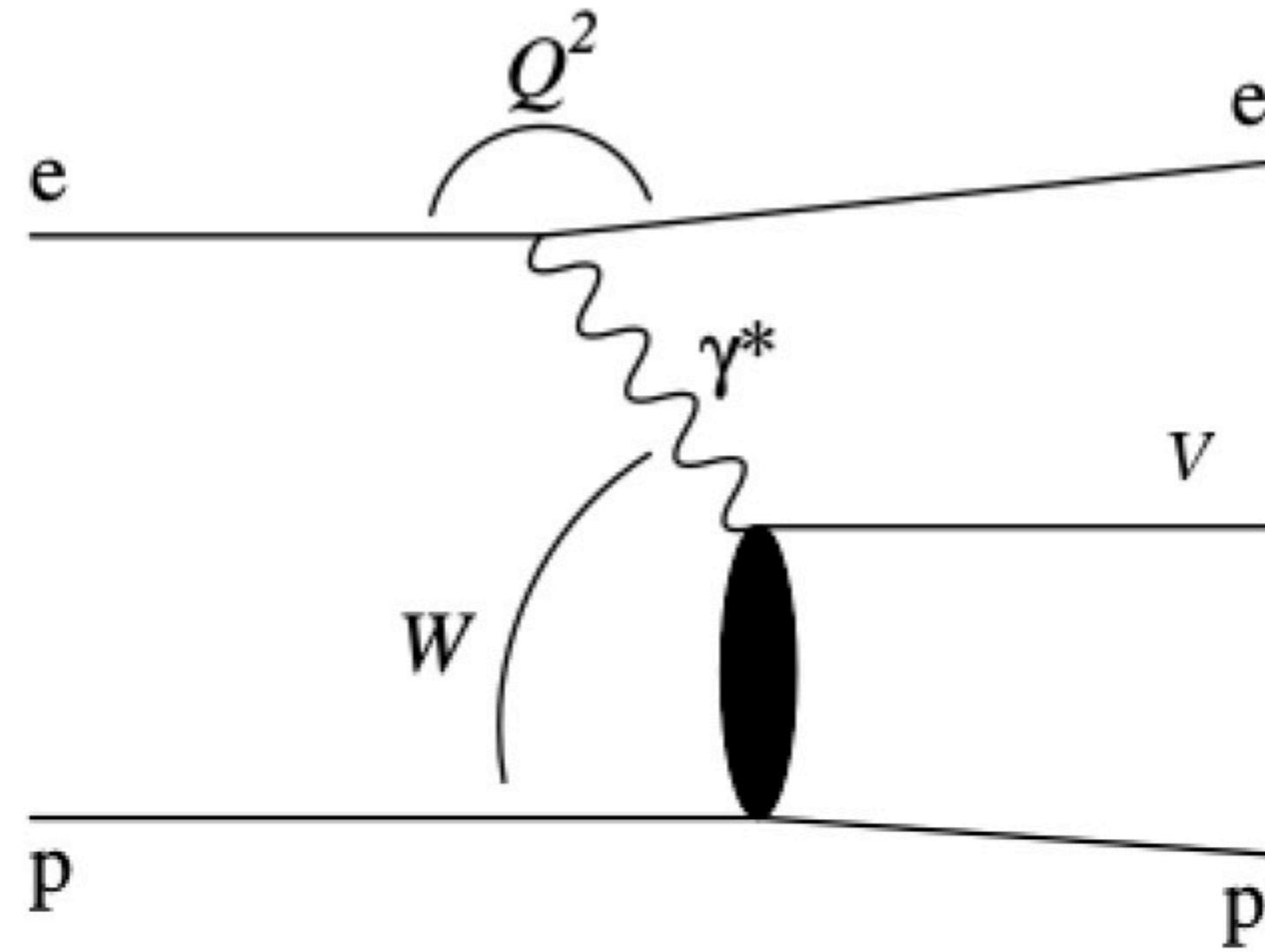
- * Starts at the next-to-leading twist (twist three)
- * Known up to leading order



$$\alpha = \frac{k^2}{Q^2}, \quad B = 2\pi\alpha_s \frac{e}{\sqrt{2}} f_\rho, \quad \varphi_1 \rightarrow DA$$

ρ -meson leptonproduction

 **Electron-proton collision** $e + p \longrightarrow e + \gamma^* + p \longrightarrow e + \rho + p$



* Exclusive reaction

* High-energy regime

$$s \equiv W^2 \gg Q^2 \gg \Lambda^2 \longrightarrow \text{small } x = \frac{Q^2}{W^2}$$

* Photon virtuality Q is the **hard scale**

* Process solved in helicity

 **Available data**

* Hera

$$2.5 \text{ GeV}^2 < Q^2 < 60 \text{ GeV}^2$$

$$35 \text{ GeV} < W < 180 \text{ GeV}$$

* Zeus

$$2 \text{ GeV}^2 < Q^2 < 60 \text{ GeV}^2$$

$$32 \text{ GeV} < W < 180 \text{ GeV}$$

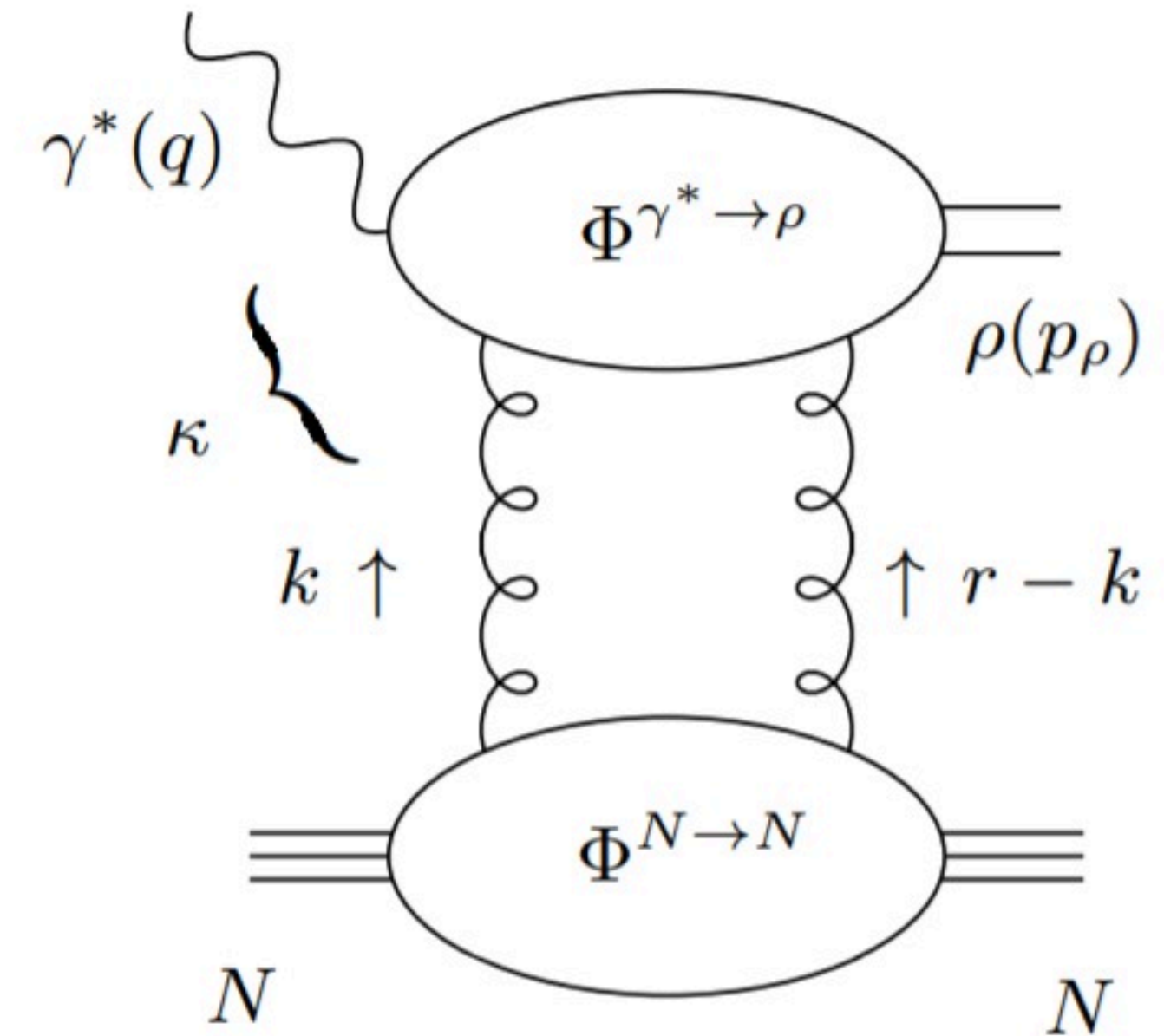
ρ -meson leptonproduction

● **Leading helicity amplitudes** $T_{\lambda_\rho \lambda_\gamma}(W^2, Q^2)$

* $Im_s\{A(\gamma^* p \rightarrow VP)\}$ dominates

* $T_{00} \gg T_{11} \gg T_{10} \gg T_{01} \gg T_{-11}$

* Small-size dipole mechanism $\longrightarrow k_T$ -factorization



$$T_{\lambda_V \lambda_\gamma}(s, Q^2) = is \int \frac{d^2k}{(k^2)^2} \Phi^{\gamma^*(\lambda_\gamma) \rightarrow V(\lambda_V)}(k^2, Q^2) \mathcal{F}(x, k^2), \quad x = \frac{Q^2}{s}$$

* $V = \rho, \phi$ via **distribution amplitude (DAs)**: $\varphi(y) = \varphi^{WW}(y) + \varphi^{gen}(y)$

ρ -meson leptonproduction

- Wandzura-Wilczek (WW) approximation \rightarrow genuine terms neglected

$$T_{11} = is \frac{2BC}{Q^2} \int \frac{d^2k}{(k^2)^2} \mathcal{F}(x, k^2) \int_0^1 \frac{dy}{(y\bar{y} + \tau)} \varphi_+^{WW}(y, \mu^2) \frac{\alpha(\alpha + 2y\bar{y} + 2\tau)}{(\alpha + y\bar{y} + \tau)^2} + o(\tau^2)$$

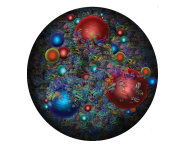
$$T_{00} = is \frac{4BC}{Q} \int \frac{d^2k}{(k^2)^2} \mathcal{F}(x, k^2) \int_0^1 dy \frac{\bar{y}y}{(y\bar{y} + \tau)} \varphi_+^{as}(y, \mu^2) \frac{\alpha}{(\alpha + y\bar{y} + \tau)}$$

$$\tau = \frac{m_q^2}{Q^2}, \quad C = \sqrt{4\pi\alpha_{em}}$$

- Generalized massive formula:
 - * $\tau = 0 \rightarrow$ no quark mass $\rightarrow \rho$ -production
 - * $\tau \neq 0 \rightarrow$ with quark mass $\rightarrow \phi$ -production
- [A. D. Bolognino, A. Szczurek, W. Schafär]

- Vector meson-DAs employed:
 - * asymptotic $\varphi_1^{as}(y) \rightarrow a_2(\mu^2) = 0$
 - * $\varphi_+^{WW}(y, \mu^2) = (2y - 1)\varphi_{1T}^{WW}(y, \mu^2) + \varphi_{AT}^{WW}(y, \mu^2)$

ρ -meson leptonproduction



Models of unintegrated gluon density (UGD)

* **ABIPSW:** x-independent model $\mathcal{F}(x, k^2) = \frac{A}{(2\pi)^2 M^2} \left[\frac{k^2}{k^2 + M^2} \right]$

[\[I. V. Anikin et al. \(2011\)\]](#)

* **Toy model:** gluon momentum derivative $\mathcal{F}(x, k^2) = \frac{d(xg(x, k^2))}{d \ln k^2}$

* **IN:** soft-hard model $\mathcal{F}(x, k^2) = \mathcal{F}_{soft}(x, k^2) + \mathcal{F}_{hard}(x, k^2)$

[\[I. P. Ivanov and N. N. Nikolaev \(2002\)\]](#)

* **HSS:** $\mathcal{F}(x, k^2) = \Phi_P \otimes \mathcal{G}_{BFKL}$

[\[M. Hentschinski, A. Sabio Vera, C. Salas \(2013\)\]](#)

* **WMR:** angular ordering of gluon emissions

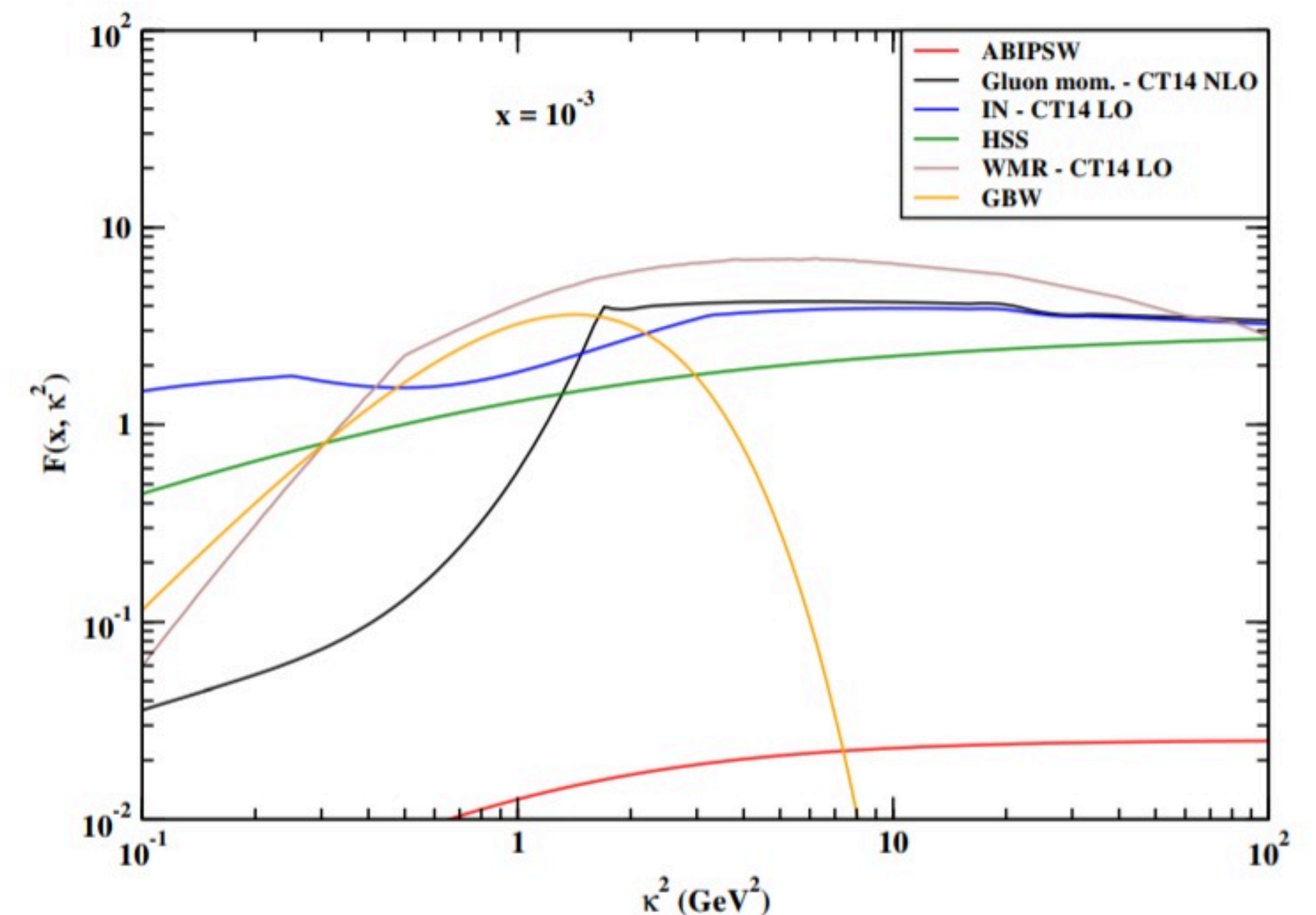
[\[G. Watt, A. D. Martin, M. G. Ryskin \(2003\)\]](#)

* **GBW:** FT of dipole cross section

[\[K. J. Golec-Biernat, M. Wüsthoff \(1998\)\]](#)

* **BCRT:** small-x improved unpolarized gluon TMD

[\[A. Bacchetta, F.G. Celiberto, M. Radici, P. Taelis \(2020\)\]](#)



ρ -meson leptonproduction at HERA

$$\sigma_L(\gamma^* p \rightarrow Vp) = \frac{1}{16\pi b(Q^2)} \frac{|T_{00}(s, Q^2)|^2}{W^2}$$

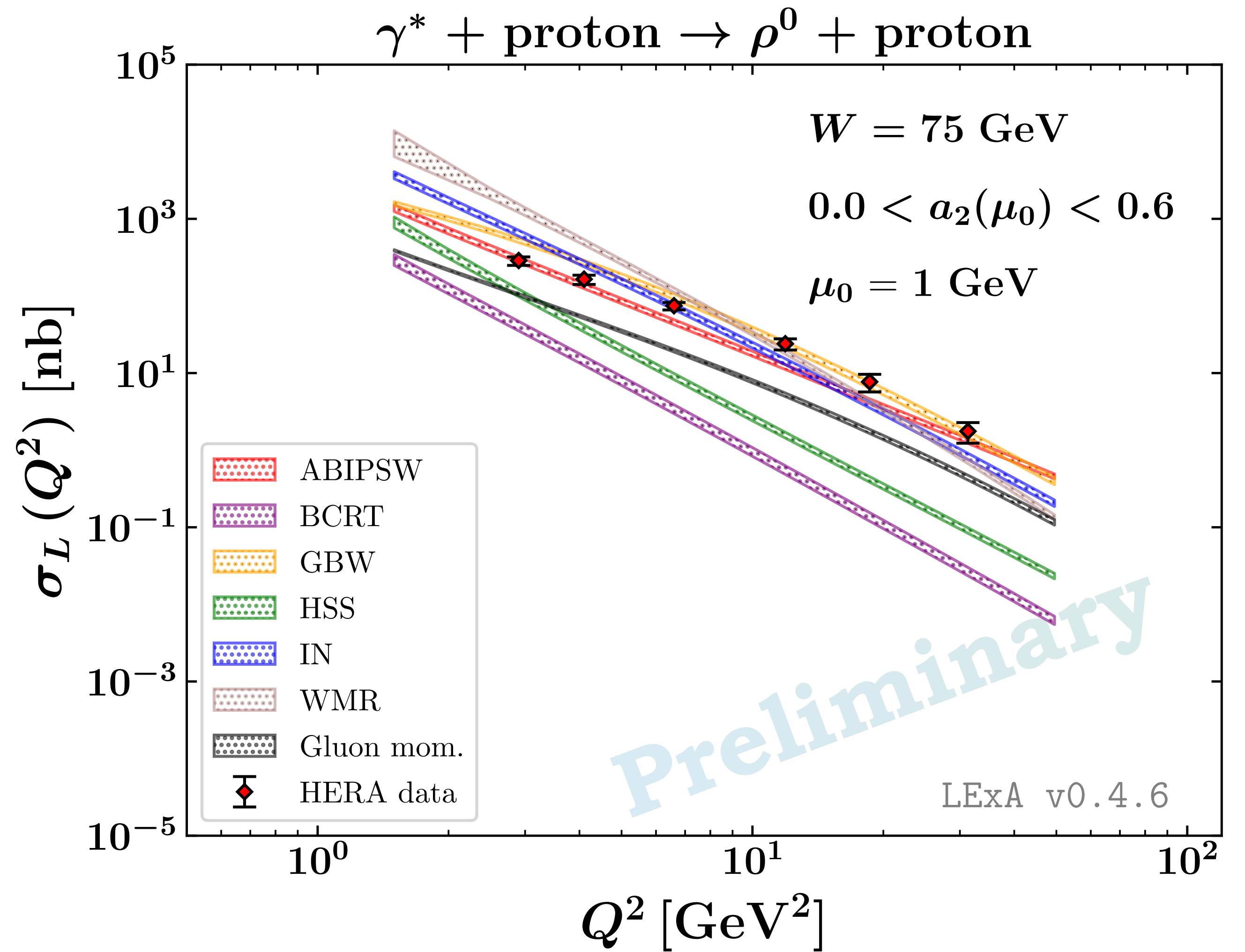
$$\sigma_T(\gamma^* p \rightarrow Vp) = \frac{1}{16\pi b(Q^2)} \frac{|T_{11}(s, Q^2)|^2}{W^2}$$

* $b(Q^2)$ -slope for light vector mesons

$$b(Q^2) \approx \beta_0 - \beta_1 \ln \left[\frac{Q^2 + m_V^2}{m_{J/\Psi}^2} \right] + \frac{\beta_2}{Q^2 + m_V^2}$$

* For ρ -meson:

$$\beta_0 = 6.5 \text{ GeV}^{-2}, \beta_1 = 1.2 \text{ GeV}^{-2}, \beta_2 = 1.1 \text{ GeV}^{-2}$$



[A. D. Bolognino, F. G. Celiberto, D. Yu. Ivanov, A. Papa, A. Szczurek, W. Schaf er]

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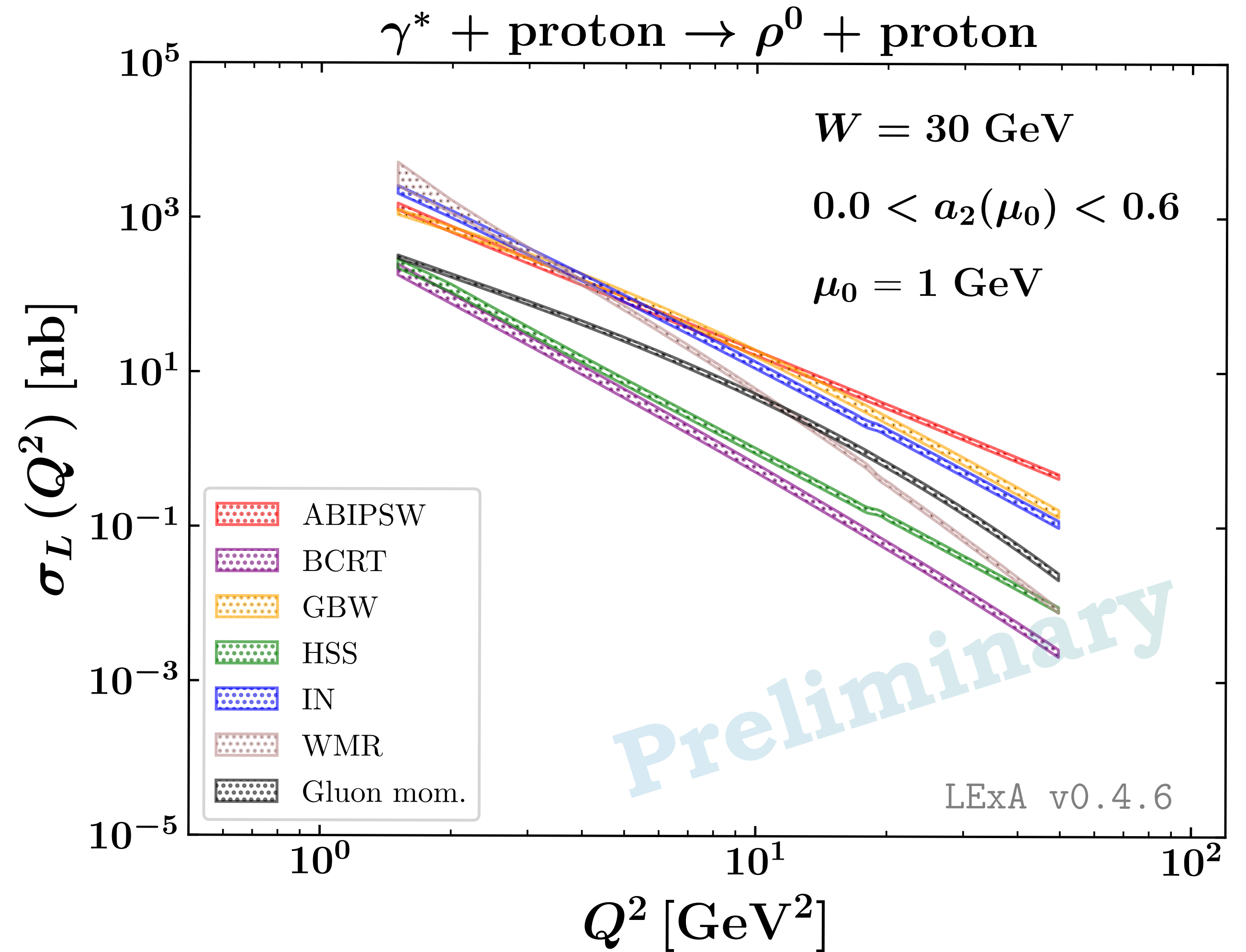
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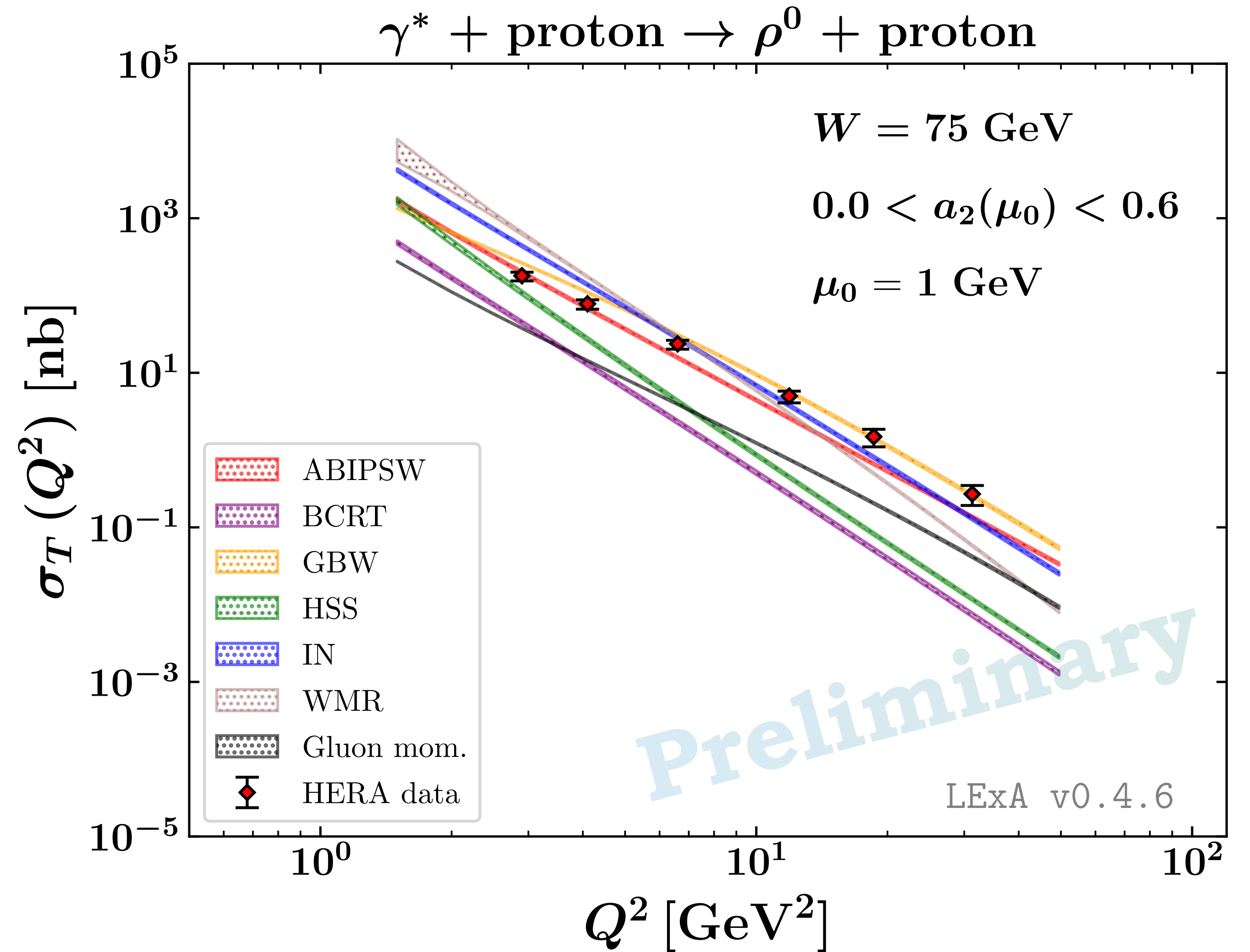
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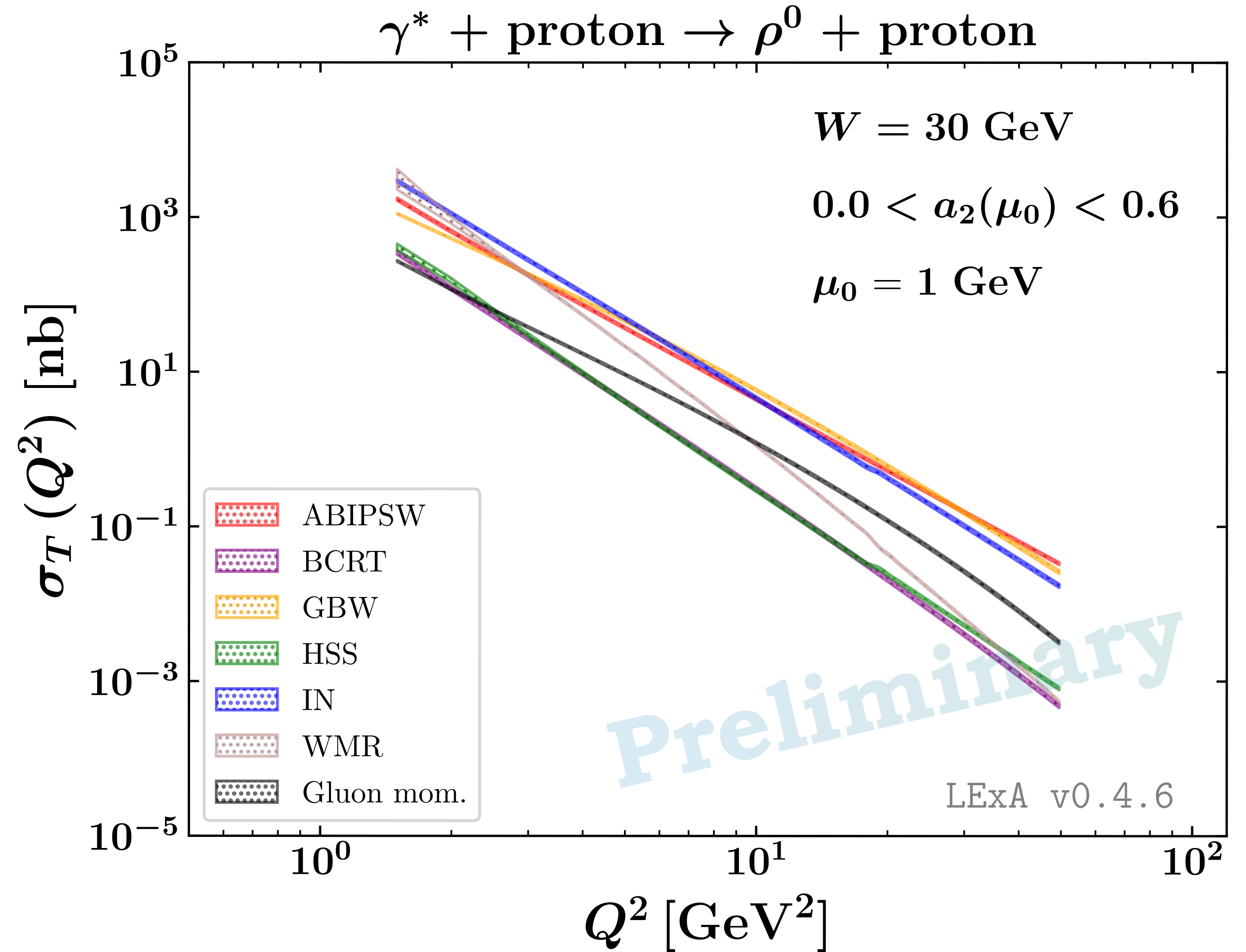
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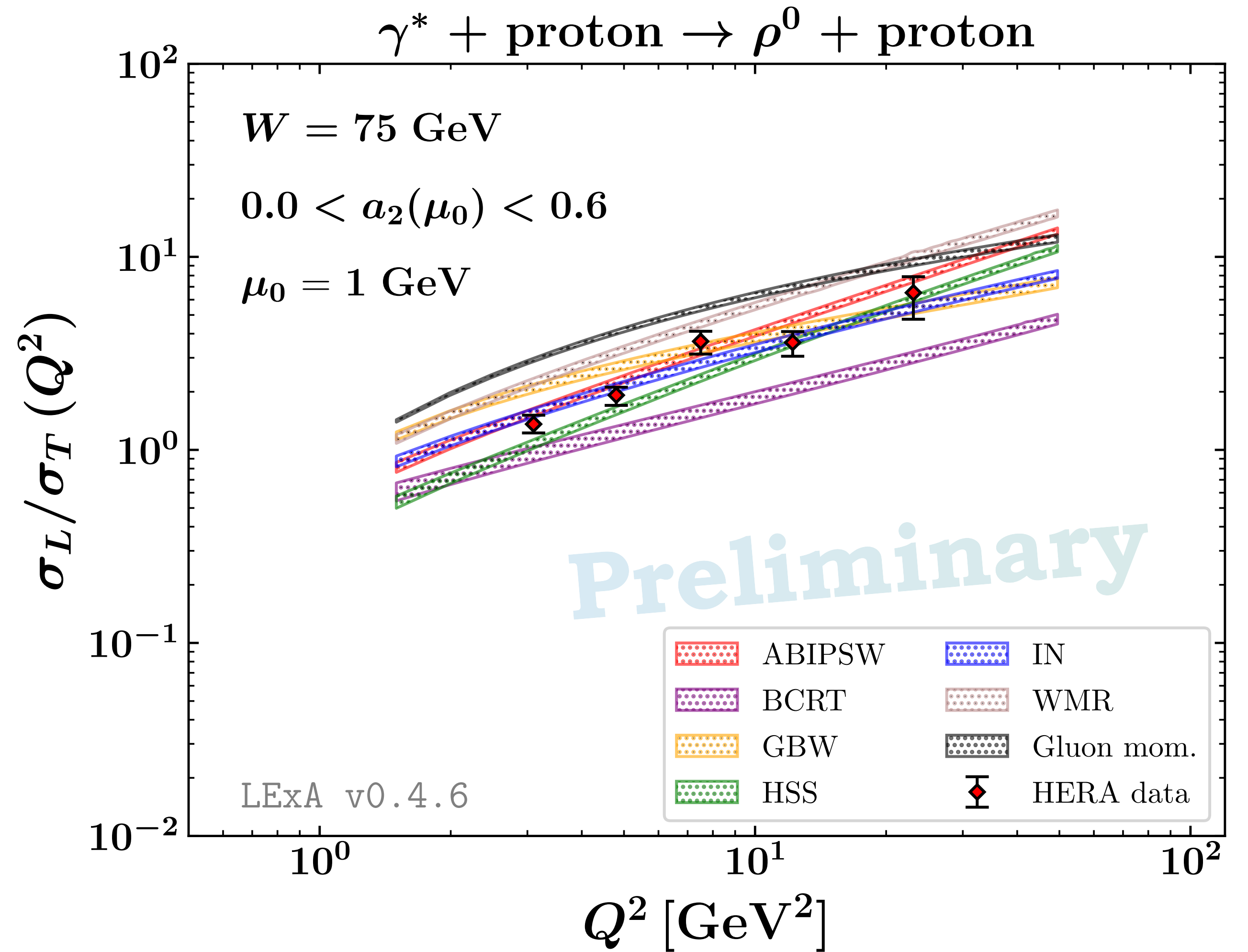
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* For ρ -meson:

$$\beta_0 = 6.5 \text{ GeV}^{-2}, \beta_1 = 1.2 \text{ GeV}^{-2}, \beta_2 = 1.1 \text{ GeV}^{-2}$$



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ρ -meson leptonproduction at the EIC

$$\sigma_L(\gamma^* p \rightarrow Vp) = \frac{1}{16\pi b(Q^2)} \frac{|T_{00}(s, Q^2)|^2}{W^2}$$

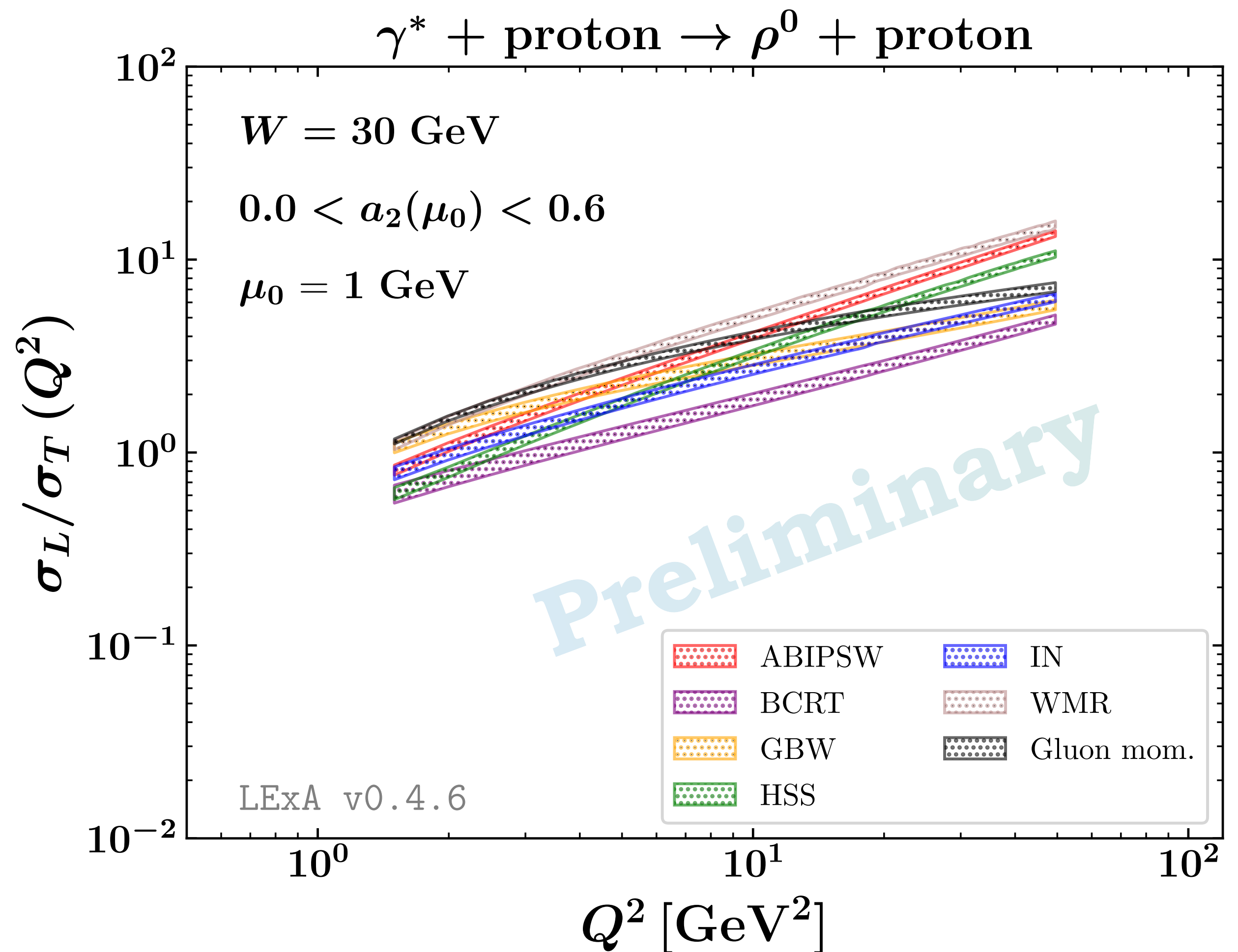
$$\sigma_T(\gamma^* p \rightarrow Vp) = \frac{1}{16\pi b(Q^2)} \frac{|T_{11}(s, Q^2)|^2}{W^2}$$

* $b(Q^2)$ -slope for light vector mesons

$$b(Q^2) \approx \beta_0 - \beta_1 \ln \left[\frac{Q^2 + m_V^2}{m_{J/\Psi}^2} \right] + \frac{\beta_2}{Q^2 + m_V^2}$$

* For ρ -meson:

$$\beta_0 = 6.5 \text{ GeV}^{-2}, \beta_1 = 1.2 \text{ GeV}^{-2}, \beta_2 = 1.1 \text{ GeV}^{-2}$$



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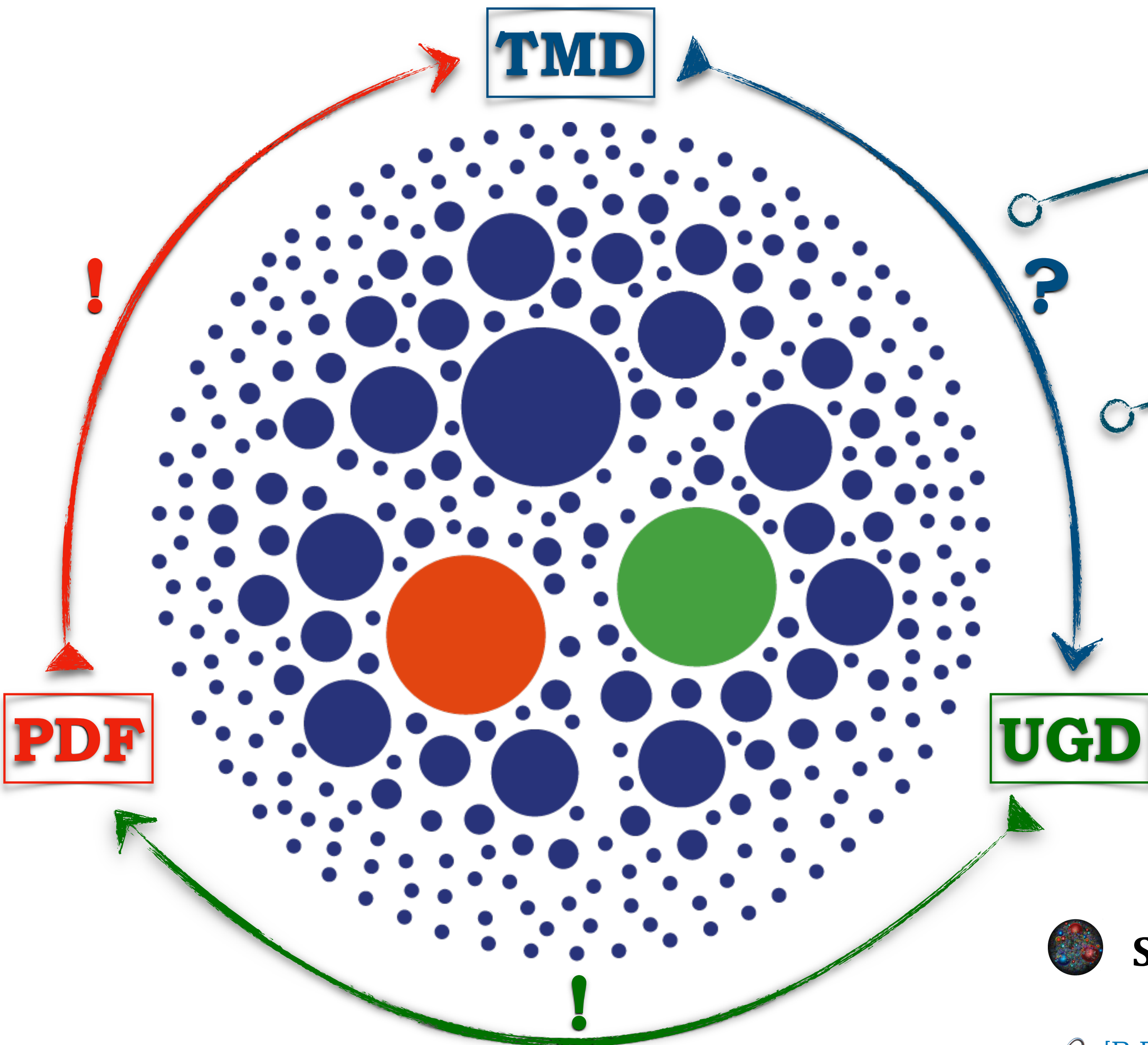
Outline

- BFKL and unintegrated gluon densities
- Exclusive forward meson leptonproduction
- **Conclusions and outlook**

Conclusions and summary

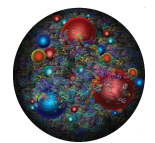



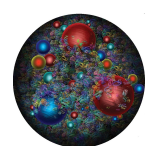

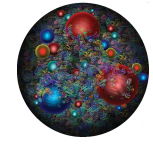

- **Unintegrated gluon densities are essential for the description of high-energy QCD**
- **Vector meson leptonproduction is a suitable tool for the investigation of the UGD**
 - * Impact factors for both longitudinally and transversely polarized ρ -meson are known
 - * Hera data are available and predictions for future studies at the EIC has been built
- **None of models is able to reproduce the entire HERA Q^2 -spectrum**
 - * UGD model extraction from fits
 - * Towards a unification of formalism

Towards a unification of formalisms



M. Nefedov's talk

F. Celiberto's talk

-  **TMD evolution vs BFKL evolution**
 -  [M. Hentschinski (2021)]
 -  [M. Nefedov (2021)]
 -  [M. Hentschinski, A. Kusina, K. Kutak, M. Serino (2018)]
-  **Small- x input to gluon TMDs**
 -  [A. Bacchetta, F.G. Celiberto, M. Radici, P. Taelis (2020)]
-  **Small- x resummed collinear PDFs**
 -  [R.D. Ball, V. Bertone, M. Bonvini, S. Marzani, J. Rojo, L. Rottoli (2018)]

Thanks for the attention!