Three marticle scattering amplitudes from Lattice QCD

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Introduction Finite-Volume Spectrum Three particles in finite volume Applications to lattice QCD Summary and Outlook









Quantum chromodynamics is conceptually simple. Its realization in nature, however, is usually very complex.

Frank Wilczek

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 $\mathscr{L}_{QCD} = \sum_{i}^{N_f} \bar{q}_i \left(D_{\mu} \gamma^{\mu} + m_i \right) q_i + \frac{1}{4} G^a_{\mu\nu} G^{\mu\nu}_a$

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$\left< \mathcal{O}(t) \mathcal{O}(0) \right> = \frac{1}{\mathcal{Z}} \left[D\psi D\bar{\psi} DA \mathcal{O}(t) \mathcal{O}(0) e^{-S_E(\psi,\bar{\psi},A_\mu)} \right]$









$\left< \mathcal{O}(t) \mathcal{O}(0) \right> = \frac{1}{\mathcal{Z}} \int D\psi D\bar{\psi} DA \, \mathcal{O}(t) \mathcal{O}(0) e^{-S_E(\psi,\bar{\psi},A_\mu)}$ Euclidea



Euclidean action





$$\left< \mathcal{O}(t) \mathcal{O}(0) \right> = \frac{1}{\mathcal{Z}} \int D\psi D\bar{\psi} DA \, \mathcal{O}(t) \mathcal{O}(t)$$

It is systematically improvable: finite volume, discretization effects 0







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It is systematically improvable: finite volume, discretization effects 0

Can we obtain scattering amplitudes from Euclidean correlation functions?









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See plenary by D. Wilson, Today





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Volume Dependence of the Energy Spectrum in Massive Quantum Field Theories

II. Scattering States

M. Lüscher

Theory Division, Deutsches Elektronen-Synchrotron DESY, D-2000 Hamburg 52, Federal Republic of Germany





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Resonance	$I_{\pi\pi\pi}$	J^P
$\omega(782)$	0	1-
$h_1(1170)$	0	1+
$\omega_3(1670)$	0	3-
$\pi(1300)$	1	0-
$a_1(1260)$	1	1^{+}
$\pi_1(1400)$	1	1-
$\pi_2(1670)$	1	2^{-}
$a_2(1320)$	1	2^{+}
$a_4(1970)$	1	4^{+}

(with $\geq 3\pi$ decay modes)





Many-body nuclear physics, 3N force

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C CP violation in D and K decays

$$K \rightarrow 3\pi, D \rightarrow 4\pi$$

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$$\overrightarrow{p} = \frac{2\pi}{L}(n_x, n_y, n_z)$$

Two particles: E =

$$2\sqrt{m^2 + \frac{4\pi^2}{L^2}} \vec{n}^2$$

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Ground state to leading order $\underline{E_2} - 2m = \langle \phi(\vec{0})\phi(\vec{0}) | \mathbf{H_I} | \phi(\vec{0})\phi(\vec{0}) \rangle$ $\frac{E_2 - 2m}{8m^2L^3} + O(L^{-4})$ [Huang, Yang, 1958]







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The energy shift is related to the scattering amplitude







Interactions change the spectrum!

state to leading order $\mathbf{M} = \langle \phi(\vec{0})\phi(\vec{0}) | \mathbf{H}_{\mathbf{I}} | \phi(\vec{0})\phi(\vec{0}) \rangle$ $E_2 - 2m = \frac{\mathcal{M}_2(E = 2m)}{\Omega - 2\pi^3} + O(L^{-4})$

[Huang, Yang, 1958]

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 $C(t) = \langle \mathcal{O}(t) \mathcal{O}(0) \rangle =$

$$\sum_{n} \left| \langle 0 | \mathcal{O}(0) | n \rangle \right|^2 e^{-E_n t}$$

The energy levels of the theory are measured from Euclidean correlation functions 0

 $C(t) = \langle \mathcal{O}(t) \mathcal{O}(0) \rangle =$

$$\sum_{n} \left| \langle 0 | \mathcal{O}(0) | n \rangle \right|^2 e^{-E_n t}$$

Current techniques allow the determination of many energy levels!

[Blanton, Hanlon, Hörz, Morningstar, FRL, Sharpe]

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Three-particle scattering amplitudes can be divergent for specific kinematics.
 They depend also on two-to-two interactions.
 But any separation between "two-particle" and "three-particle" effects is not well-defined
 However, the three-particle spectrum depends on S-matrix elements! [Polejaeva, Rusetsky]







Three-particle formalism(s)

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Relativistic, model-independent, three-particle quantization condition

Maxwell T. Hansen¹, * and Stephen R. Sharpe¹, †

Also [Blanton, Briceño, Hansen, Jackura, FRL, Szczepaniak, Sharpe]







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Three-particle quantization condition in a finite volume: 1. The role of the three-particle force

Hans-Werner Hammer^a, Jin-Yi Pang^b and Akaki Rusetsky^b Also [Döring, Geng, Hammer, Mai, Meißner, Müller, Pang, <u>FRL</u>, Rusetsky, Wu]





Non-Relativistic EFT (NREFT)

Finite-Volume Unitarity (FVU)

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Three-body Unitarity in the Finite Volume

M. Mai^{1, *} and M. Döring^{1, 2, †}



- **Higher partial waves**
- Nondegenerate and nonidentical scalars
- **Two-to-three transitions**
- **Three-particle decays**
- Analysis of lattice QCD data

Non-Relativistic EFT (NREFT)

- Nondegenerate (DDK systems)
- Perturbative expansions for three pions and excited states
- **Three-particle decays**
- **Relativistic kinematics can be included**

Finite-Volume Unitarity (FVU)

- **Chiral extrapolations**
- Analysis of lattice QCD data
- Study of resonant channels: $a_1(1260)$



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Hansen, Jackura, <u>FRL</u>, Szczepaniak, Sharpe]

ation condition in a finite volume: of the three-particle force

All three formalisms should be equivalent. Hans-Werner Hammer^a, Jin-Yi Pang^b and Akaki Rusetsky^b

Explicitly shown for FVU and RFT! [Blanton, Sharpe] Also [Döring, Geng, Hammer, Mai, Meißner, Müller, Pang, FRL, Rusetsky, Wu]

Three-body Unitarity in the Finite Volume

M. Mai^{1, *} and M. Döring^{1, 2, †}





O The determination of three-particle scattering amplitudes on the lattice is a two-step process!





$$\left[M(E,L)\right]\Big|_{E=E_n} = 0$$

$$_{3}, H_{0}, C_{0}$$



$$\left[M(E,L)\right]\Big|_{E=E_n} = 0$$























The RET Formalism





Solve integral equations to obtain the physical three-to-three amplitude

Solved in [Briceño et al], [Hansen et al.], [Jackura et al.]

The RET Formalism

Physical 3->3 amplitude $\mathcal{K}_2, \mathcal{K}_{df.3}$ Integral equations





 $a_0, \mathcal{K}_{df,3} \dots$



































• From the lattice, one can get the one-to-three finite-volume matrix element:

 $\langle E_n, \boldsymbol{P}, \Lambda \mu, L | \mathcal{H}_W(0) | K, \boldsymbol{P}, L \rangle$.





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How to relate that to the physical infinite-volume decay amplitude? 0

$$T_{K3\pi} = \langle 3\pi, \operatorname{out} | \mathcal{H}_W(0) | K, \boldsymbol{P} \rangle$$
,

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 $\operatorname{out} | \mathcal{H}_W(0) | K, \boldsymbol{P} \rangle$,

NREFT in [Müller, Rusetsky] RFT in [Hansen, <u>FRL</u>, Sharpe]

See talk by S. Sharpe, Wednesday 28th













Tyler D. Blanton,^{1, *} Fernando Romero-López,^{2, †} and Stephen R. Sharpe^{1, ‡} ¹Physics Department, University of Washington, Seattle, WA 98195-1560, USA ²Instituto de Física Corpuscular, Universitat de València and CSIC, 46980 Paterna, Spain

[see also earlier work using the ground state by Mai et al., and Beane et al.]







$$\frac{q}{M}\cot\delta_0 = \frac{\sqrt{sM}}{s - z_2^2} \left(B_0 + B_1 q^2 + \cdots\right)$$
$$\mathcal{K}_{df,3} = \mathcal{K}_{df,3}^{iso,0} + \mathcal{K}_{df,3}^{iso,1} \left(\frac{s - 9M^2}{9M^2}\right)$$











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Some tension with ChPT.




O Parametrize K-matrices with only s-wave interactions:







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1. 2 σ evidence for $\mathscr{K}_{df,3} \neq 0$.

Some tension with ChPT.

Same spectrum has been analyzed by [Mai, Döring, Culver, Alexandru]





0 [Fischer, Kostrzewa, Liu, <u>FRL</u>, Ueding, Urbach (ETMC)]



In a later article, the chiral dependence of $\mathscr{K}_{df,3}$ has been studied, including physical pions.

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The FVU formalism has also been applied to three-pion systems 0 [Brett et al.]



three-particle "contact" term

See talk by A. Alexandru, Wednesday 28th





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[Blanton, <u>FRL</u>, Sharpe]

$$\mathcal{K}_{\mathrm{df},3} = \mathcal{K}_{\mathrm{df},3}^{\mathrm{iso},0} + \mathcal{K}_{\mathrm{df},3}^{\mathrm{iso},1}\Delta + \mathcal{K}_{\mathrm{df},3}^{\mathrm{iso},2}\Delta^2 + \mathcal{K}_A\Delta_A + \mathcal{K}_B\Delta_B,$$





[Blanton, <u>FRL</u>, Sharpe]

Swave

 $\mathcal{K}_{\mathrm{df},3} = \mathcal{K}_{\mathrm{df},3}^{\mathrm{iso},0} + \mathcal{K}_{\mathrm{df},3}^{\mathrm{iso},1}\Delta + \mathcal{K}_{\mathrm{df},3}^{\mathrm{iso},2}\Delta^2 + \mathcal{K}_A\Delta_A + \mathcal{K}_B\Delta_B,$ $\Delta \equiv \frac{s - 9m^2}{2}$

$$\Delta \equiv \frac{s - sm}{9m^2}$$





[Blanton, <u>FRL</u>, Sharpe]

d-wave s-wave $\mathcal{K}_{\mathrm{df},3} = \mathcal{K}_{\mathrm{df},3}^{\mathrm{iso},0} + \mathcal{K}_{\mathrm{df},3}^{\mathrm{iso},1}\Delta + \mathcal{K}_{\mathrm{df},3}^{\mathrm{iso},2}\Delta^2 + \mathcal{K}_A\Delta_A + \mathcal{K}_B\Delta_B,$ $\Delta \equiv \frac{s - 9m^2}{9m^2}$





[Blanton, FRL, Sharpe]

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 $t_{ij} = (p_i - k_j)^2$





[Blanton, <u>FRL</u>, Sharpe]

d-wave s-wave $\mathcal{K}_{\mathrm{df},3} = \mathcal{K}_{\mathrm{df},3}^{\mathrm{iso},0} + \mathcal{K}_{\mathrm{df},3}^{\mathrm{iso},1}\Delta + \mathcal{K}_{\mathrm{df},3}^{\mathrm{iso},2}\Delta^2 + \mathcal{K}_A\Delta_A + \mathcal{K}_B\Delta_B,$ $\Delta \equiv \frac{s - 9m^2}{9m^2}$ $\Delta_B = \sum_{i,j=1} t_{ij}^2 - \Delta^2,$

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See talk by A. Hanlon, Wednesday 28th



[Blanton, Hanlon, Hörz, Morningstar, FRL, Sharpe]



Integral equations (RFT Contraction and marked in section in an animal the second second Final step Physical 3->3 amplitude M_{z} $\mathcal{K}_2, \mathcal{K}_{df,3}$ Integral equations ĸĸĸŎŢŦĸſĊĊſĸŖŊĊŎŎĔĸĸĊĊŶĊŦĔĿĿŢŶĊĹĬĬſĊĬĿĿĊŴŢĬŢſŎŢŦĸſĊĊŦŖŊĊŎĔŔĸĿĊŶŶĔŦĔĿĿŖ









Integral equations (RF Final slep Physical 3->3 amplitude $\mathcal{K}_2, \mathcal{K}_{df,3}$ M_{2} Integral equations

Particle-Dimer phase shift [Jackura et al.]









Final slep $\mathscr{K}_2, \mathscr{K}_{df,3}$ M2 Integral equations





Integral equations Final step Physical 3->3 amplitude M_{2} $\mathcal{K}_2, C_{\ell\ell'}$ Integral equations









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- Three versions of finite-volume formalism for identical particles.
- Various generalizations for nonidentical scalar particles.
- Formalism for three-body decays.

Many recent formal developments for three-particle systems.





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- O Many recent formal developments for three-particle systems.
 - Three versions of finite-volume formalism for identical particles.
 - Various generalizations for nonidentical scalar particles.
 - Formalism for three-body decays.
- O Applications to simple systems successfully undertaken
 - Some lattice studies of three charged pions and kaons
 - It is possible to even study d-wave interactions [Blanton et al.]
 - Recent study in FVU formalism of the $a_1(1260)$ [Mai et al.]









1. Generalizing the formalism for generic two- and three- particle systems (e.g. nucleons, Roper)





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- 2. Apply the formalism to more complex systems (e.g. $\pi\pi K$, $h_1(1170)...$)





- 1 Generalizing the formalism for generic two- and three- particle systems (e.g. nucleons, Roper)
- 2. Apply the formalism to more complex systems (e.g. $\pi\pi K$, $h_1(1170)...$)
- Study three-particle decays with the available framework. (e.g. $K \rightarrow 3\pi$)





- Generalizing the formalism for generic two- and three- particle systems (e.g. nucleons, Roper)
- 2. Apply the formalism to more complex systems (e.g. $\pi\pi K$, $h_1(1170)...$)
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19:00	Multibody decay analyses tool: A phenomenological model for meson-meson subamplitudes	Dr. Patricia MAGALHAES
	Progress in relativistic three-hadron scattering from lattice QCD	Dr. Andrew JACKURA
	Mexico City	18:40 - 19:05
	Unitarity in hadronic three-body decay and application to physics beyond the Standard Model	Mr. Mehmet Hakan AKDAG
	The six-pion amplitude	Tomas HUSEK
	Mexico City	19:30 - 19:55

Wednesday 28th

	Three-pion scattering from lattice QCD	Maxwell HANSEN
	Mexico City	18:15 - 18:40
	Three pion and three kaon scattering from lattice QCD	Andrei ALEXANDRU
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19:00	Beyond s-wave interactions of two- and three-meson systems with maximal isospin from lattice QCD	Dr. Andrew HANLON
	Generalizing the Lellouch-Luscher formula to three-particle decays	Prof. Stephen SHARPE 📄
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	Study of scalar meson production in three body \$\eta_c\$ decays at BABAR	Dr. Alessandro PILLONI
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Threshold expansion:

$\mathcal{K}_{\mathrm{df},3} = \mathcal{K}_{\mathrm{df},3}^{\mathrm{iso},0} + \mathcal{K}_{\mathrm{df},3}^{\mathrm{iso},1}\Delta + \mathcal{K}_{\mathrm{df},3}^{\mathrm{iso},2}\Delta^2 + \mathcal{K}_A\Delta_A + \mathcal{K}_B\Delta_B,$

[Blanton, <u>FRL</u>, Sharpe]



Threshold expansion:

$$\mathcal{K}_{\mathrm{df},3} = \mathcal{K}_{\mathrm{df},3}^{\mathrm{iso},0} + \mathcal{K}_{\mathrm{df},3}^{\mathrm{iso},1}\Delta + \mathcal{K}$$

[Blanton, <u>FRL</u>, Sharpe]

 $\Delta \equiv \frac{s - 9m^2}{9m^2}$



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d-wave

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d-wave

 $\mathcal{K}^{\mathrm{iso},2}_{\mathrm{df},3}\Delta^2 + \mathcal{K}_A\Delta_A + \mathcal{K}_B\Delta_B$ $\Delta_B = \sum_{i,j=1}^{0} \widetilde{t}_{ij}^2 - \Delta^2,$



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[Blanton, <u>FRL</u>, Sharpe]





[Blanton, Hanlon, Hörz, Morningstar, FRL, Sharpe]



d-wave

 $\mathcal{K}^{\mathrm{iso},2}_{\mathrm{df},3}\Delta^2 + \mathcal{K}_A\Delta_A + \mathcal{K}_B\Delta_B$ $\Delta_B = \sum_{i,j=1}^{0} \widetilde{t}_{ij}^2 - \Delta^2,$



Higher partial waves

Threshold expansion:

$$\mathcal{K}_{\mathrm{df},3} = \mathcal{K}_{\mathrm{df},3}^{\mathrm{iso},0} + \mathcal{K}_{\mathrm{df},3}^{\mathrm{iso},1}\Delta + \mathcal{K}$$

[Blanton, <u>FRL</u>, Sharpe]





[Blanton, Hanlon, Hörz, Morningstar, FRL, Sharpe]

d-wave

$$\begin{split} \mathcal{K}_{\mathrm{df},3}^{\mathrm{iso},2} \Delta^2 + \mathcal{K}_A \Delta_A + \mathcal{K}_B \Delta_B \\ \Delta_B \ = \sum_{i,j=1}^3 \widetilde{t}_{ij}^2 - \Delta^2 \,, \end{split}$$



