

Molecular interpretation of the LHCb P_c states from an analysis of $J/\psi p$ spectrum

Meng-Lin Du

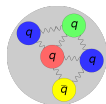
IFIC, University of Valencia
HISKP, Universität Bonn

In collaboration with V. Baru, F.-K. Guo, C. Hanhart,
U.-G. Meißner, J. A. Oller, and Q. Wang

Based on PRL124(2020)072001 and arXiv: 2102.07159 (2021)

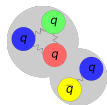
19th International Conference on Hadron Spectroscopy and Structure

Penta-Quark States



Pentaquark

↪ Compact object formed from q and \bar{q}



Hadronic-Molecule

↪ Extended object made of **Baryon** and **Meson**

▶ $\Lambda(1405)$

↪ $\bar{K}N$ predicted by Dalitz and Tuan, 1959

PRL2,425

↪ $\Lambda(1405) \rightarrow \Sigma\pi$ observed by Alston et al.,

PRL6,698

▶ “ $\theta(1540)$ ”

predicted by Diakonov et al., 1997 ($Z(1530)^+$)

ZPA359, 305

↪ NOT supported by many high statistics experiments

▶ N^* and Λ^* with hidden charm

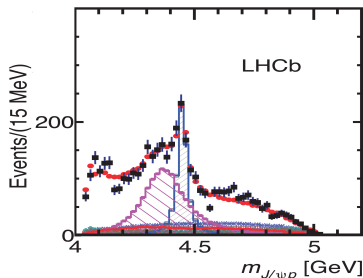
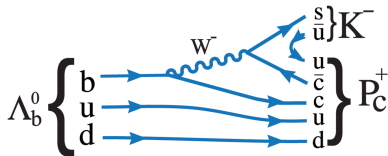
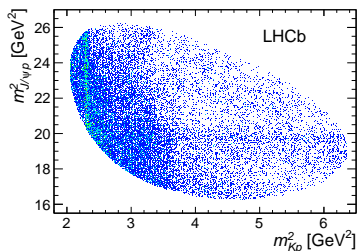
Wu et al., PRL 105, 232001 (2011)

↪ and many more works...

Charmonium-pentaquark states (I)

Observation of exotic structures (P_c) in $\Lambda_b^0 \rightarrow J/\psi p K^-$

LHCb, PRL 115, 072001 (2015)



$P_c(4380)^+$: $M = 4380 \pm 8 \pm 29$ MeV

$\Gamma = 205 \pm 18 \pm 86$ MeV

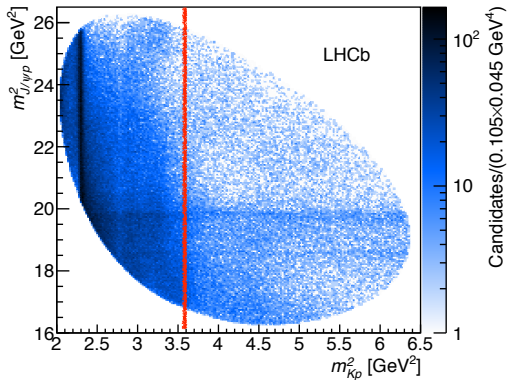
$P_c(4450)^+$: $M = 4449.8 \pm 1.7 \pm 2.5$ MeV

$\Gamma = 39 \pm 5 \pm 19$ MeV

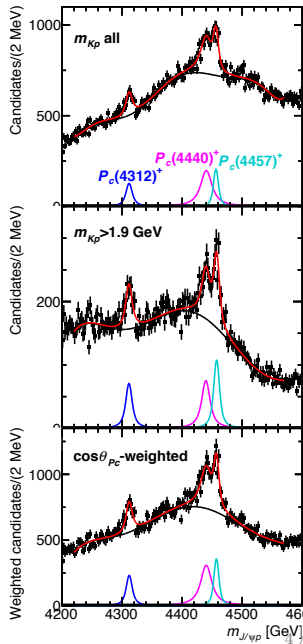
Preferred Parity: Opposite

Charmonium-pentaquark states (II)

LHCb, PRL 122, 222001 (2019)



| State | M [MeV] | Γ [MeV] |
|---------------|--------------------------------|-------------------------------|
| $P_c(4312)^+$ | $4311.9 \pm 0.7^{+6.8}_{-0.6}$ | $9.8 \pm 2.7^{+3.7}_{-4.5}$ |
| $P_c(4440)^+$ | $4440.3 \pm 1.3^{+4.1}_{-4.7}$ | $20.6 \pm 4.9^{+8.7}_{-10.1}$ |
| $P_c(4457)^+$ | $4457.3 \pm 0.6^{+4.1}_{-1.7}$ | $6.4 \pm 2.0^{+5.7}_{-1.9}$ |



Charmonium-pentaquark (theoretical)

- ▶ Compact pentaquark Cheng et al., PRD100(2019)054002

$P_c(4312), P_c(4440), P_c(4457): J^P = 3/2^-, 1/2^-, 3/2^-$

- ▶ Compact diquark model Ali et al., JHEP1910(2019)256

| | |
|---------|---------------|
| $3/2^-$ | 4240 ± 29 |
| $3/2^+$ | 4440 ± 35 |
| $5/2^+$ | 4457 ± 35 |

- ▶ P_c s as double triangle cusps Nakamura, PRD103(2020)L111503

- ▶ $P_c(4312)$: virtual state Fernández-Ramírez et al., PRL123(2019)092001

- ▶ K -matrix: $J/\psi p - \Sigma_c \bar{D} - \Sigma_c \bar{D}^*$ Kuang et al., EPJC80(2020)433
- $\hookrightarrow P_c$ s have same J^P . $P_c(4312)$: $\Sigma_c \bar{D}$, $P_c(4457)$: ? cusp effect

- ▶ Molecule (HQSS) Liu et al., PRL122,242001 (2019)

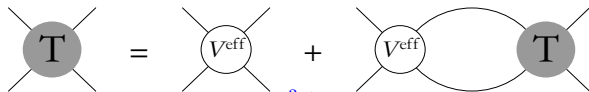
| | Molecule | J^P | M (MeV) | | Molecule | J^P | M (MeV) |
|---|-----------------------|-----------------|-----------------|---|-----------------------|-----------------|-----------------|
| A | $\bar{D}\Sigma_c$ | $\frac{1}{2}^-$ | 4311.8 – 4313.0 | B | $\bar{D}\Sigma_c$ | $\frac{1}{2}^-$ | 4306.3 – 4307.7 |
| A | $\bar{D}\Sigma_c^*$ | $\frac{3}{2}^-$ | 4376.1 – 4377.0 | B | $\bar{D}\Sigma_c^*$ | $\frac{3}{2}^-$ | 4370.5 – 4371.7 |
| A | $\bar{D}^*\Sigma_c$ | $\frac{1}{2}^-$ | 4440.3* | B | $\bar{D}^*\Sigma_c$ | $\frac{1}{2}^-$ | 4457.3* |
| A | $\bar{D}^*\Sigma_c$ | $\frac{3}{2}^-$ | 4457.3* | B | $\bar{D}^*\Sigma_c$ | $\frac{3}{2}^-$ | 4440.3* |
| A | $\bar{D}^*\Sigma_c^*$ | $\frac{1}{2}^-$ | 4500.2 – 4501.0 | B | $\bar{D}^*\Sigma_c^*$ | $\frac{1}{2}^-$ | 4523.2 – 4523.6 |
| A | $\bar{D}^*\Sigma_c^*$ | $\frac{3}{2}^-$ | 4510.6 – 4510.8 | B | $\bar{D}^*\Sigma_c^*$ | $\frac{3}{2}^-$ | 4516.5 – 4516.6 |
| A | $\bar{D}^*\Sigma_c^*$ | $\frac{5}{2}^-$ | 4523.3 – 4523.6 | B | $\bar{D}^*\Sigma_c^*$ | $\frac{5}{2}^-$ | 4500.2 – 4501.0 |

and many more works...

- ▶ quantum numbers? line shape? the existence of $P_c(4380)$?

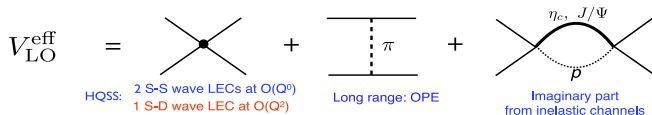
EFT approach at low energies: $\Sigma_c^{(*)} \bar{D}^{(*)}$ ($\Lambda_c \bar{D}^{(*)}$)

Lippmann-Schwinger Equation:



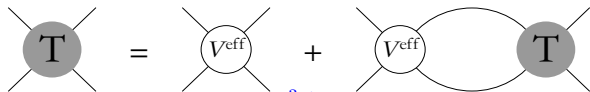
$$T_{\alpha\gamma}^J(E, p, p') = V_{\alpha\gamma}^J(E, p, p') - \sum_{\beta} \int \frac{d^3\vec{q}}{(2\pi)^3} V_{\alpha\beta}^J(E, p, q) G_{\beta}(E, q) T_{\beta\gamma}^J(E, q, p')$$

↳



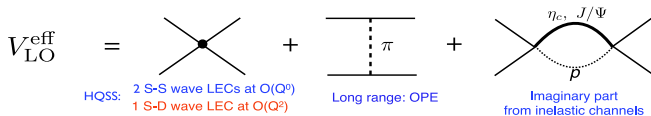
EFT approach at low energies: $\Sigma_c^{(*)} \bar{D}^{(*)} (\Lambda_c \bar{D}^{(*)})$

Lippmann-Schwinger Equation:



$$T_{\alpha\gamma}^J(E, p, p') = V_{\alpha\gamma}^J(E, p, p') - \sum_{\beta} \int \frac{d^3\vec{q}}{(2\pi)^3} V_{\alpha\beta}^J(E, p, q) G_{\beta}(E, q) T_{\beta\gamma}^J(E, q, p')$$

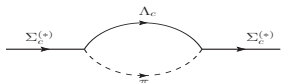
☞



☞ Green function: $\Gamma(\Sigma_c^* \rightarrow \Lambda_c \pi) = 15.0 \text{ MeV} \sim \Gamma(P_c)$

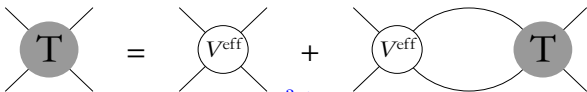
$$G_{\beta}(E, \mathbf{q}) = \frac{m_{\Sigma_c^{(*)}} m_{D^{(*)}}}{E_{\Sigma_c^{(*)}}(\mathbf{q}) E_{D^{(*)}}(\mathbf{q})} \frac{1}{E_{\Sigma_c^{(*)}}(\mathbf{q}) + E_{D^{(*)}}(\mathbf{q}) - E - \frac{\tilde{\Sigma}_R^{(*)}(s)}{2E_{\Sigma_c^{(*)}}(\mathbf{q})}},$$

↪ The self-energy: $\tilde{\Sigma}_R^{(*)}(s) \sim ig^2 \frac{p^3}{\sqrt{s}}$



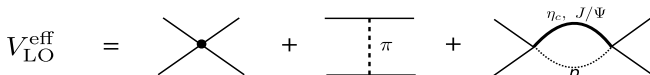
EFT approach at low energies: $\Sigma_c^{(*)} \bar{D}^{(*)}$ ($\Lambda_c \bar{D}^{(*)}$)

Lippmann-Schwinger Equation:



$$T_{\alpha\gamma}^J(E, p, p') = V_{\alpha\gamma}^J(E, p, p') - \sum_{\beta} \int \frac{d^3\vec{q}}{(2\pi)^3} V_{\alpha\beta}^J(E, p, q) G_{\beta}(E, q) T_{\beta\gamma}^J(E, q, p')$$

☞



$V_{\text{LO}}^{\text{eff}} =$

HQSS: 2 S-S wave LECs at $\mathcal{O}(Q^0)$
 1 S-D wave LEC at $\mathcal{O}(Q^2)$

Long range: OPE

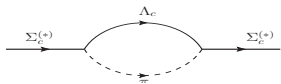
Imaginary part from inelastic channels

☞ Green function: $\Gamma(\Sigma_c^* \rightarrow \Lambda_c \pi) = 15.0 \text{ MeV} \sim \Gamma(P_c)$

$$G_{\beta}(E, \mathbf{q}) = \frac{m_{\Sigma_c^{(*)}} m_{D^{(*)}}}{E_{\Sigma_c^{(*)}}(\mathbf{q}) E_{D^{(*)}}(\mathbf{q})} \frac{1}{E_{\Sigma_c^{(*)}}(\mathbf{q}) + E_{D^{(*)}}(\mathbf{q}) - E - \frac{\tilde{\Sigma}_R^{(*)}(s)}{2E_{\Sigma_c^{(*)}}(\mathbf{q})}}$$

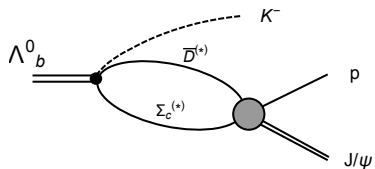
↪ The self-energy: $\tilde{\Sigma}_R^{(*)}(s) \sim ig^2 \frac{p^3}{\sqrt{s}}$

↪ Nonrelativistic limit:



$$G_{\beta}(E, \mathbf{q}) = \frac{1}{\frac{\mathbf{q}^2}{2\mu} + m_{D^{(*)}} + m_{\Sigma_c^{(*)}} - E}$$

$$\underline{\Lambda_b^0 \rightarrow K^- J/\psi p}$$

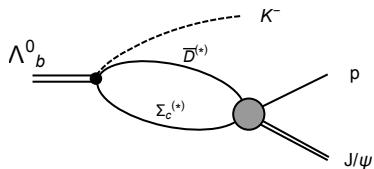


$$\Rightarrow m_{J/\psi p} \sim 4440 \text{ MeV}$$

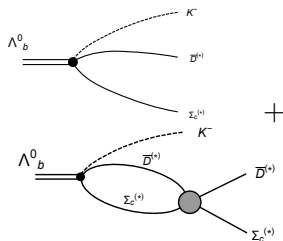
$$\hookrightarrow |\mathbf{p}| \sim 810 \text{ MeV}$$

$$\hookrightarrow J/\psi p(S), J/\psi p(D)$$

$$\underline{\Lambda_b^0 \rightarrow K^- J/\psi p}$$

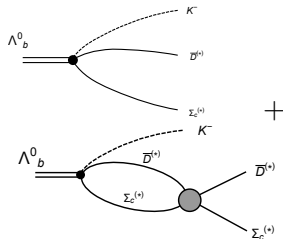
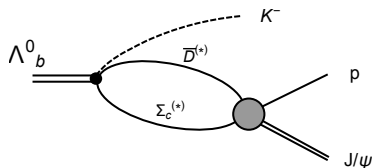


- ☞ $m_{J/\psi p} \sim 4440 \text{ MeV}$
- ↪ $|\mathbf{p}| \sim 810 \text{ MeV}$
- ↪ $J/\psi p(S), J/\psi p(D)$



- ☞ Weak production:
- ↪ S -wave $\Sigma_c^{(*)} \bar{D}^{(*)}$
- ↪ 7 parameters: P_α^J

$$\underline{\Lambda_b^0 \rightarrow K^- J/\psi p}$$



- ☞ $m_{J/\psi p} \sim 4440 \text{ MeV}$
- ☞ $|\mathbf{p}| \sim 810 \text{ MeV}$
- ☞ $J/\psi p(S), J/\psi p(D)$

- ☞ Weak production:
- ☞ S -wave $\Sigma_c^{(*)} \bar{D}^{(*)}$
- ☞ 7 parameters: P_α^J



$$\text{channels} \begin{cases} \Sigma_c^{(*)} \bar{D}^{(*)}(S/D), \Lambda_c \bar{D}^{(*)}(S/D) & \rightarrow \alpha, \beta, \gamma \\ J/\psi p(S/D), \eta_c p(S/D) & \rightarrow i, j, k \end{cases}$$



$$U_\alpha^J(E, p) = P_\alpha^J(E, p) - \sum_\beta \int \frac{d\mathbf{q}^3}{(2\pi)^3} V_{\alpha\beta}^J(E, p, q) G_\beta(E, q) U_\beta^J(q),$$

$$U_i^J(E, p) = \sum_\beta \int \frac{d\mathbf{q}^3}{(2\pi)^3} \mathcal{V}_{\beta i} G_\beta(E, q) U_\beta^J(q).$$

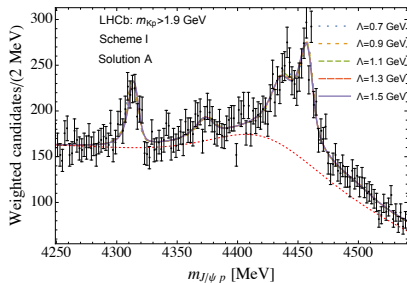
Fit Schemes

☞ Fit schemes:

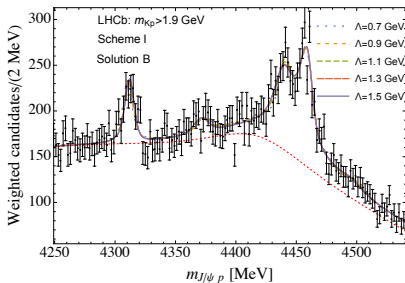
- Scheme I: pure contact potential w/o $\Lambda_c \bar{D}^{(*)}$
- Scheme II: Scheme I + OPE + S-D counter term w/o $\Lambda_c \bar{D}^{(*)}$
↪ coupled channel
- Scheme III: contact + OPE + S-D counter terms w/ $\Lambda_c \bar{D}^{(*)}$

Scheme I: pure contact potential w/o $\Lambda_c \bar{D}^{(*)}$

Solution A



Solution B



☞ $\Lambda > \Lambda_{\text{soft}} \sim \sqrt{2\mu\delta} \sim 0.7 \text{ GeV}$

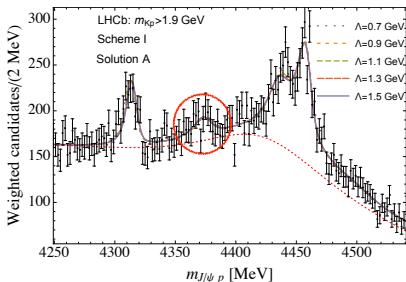
☞ Cutoff-independent for both solution A and B

| $\Sigma_c \bar{D}^*$ | $P_c(4440)$ | $P_c(4457)$ |
|----------------------|-----------------|-----------------|
| Fit A | $\frac{1}{2} -$ | $\frac{3}{2} -$ |
| Fit B | $\frac{3}{2} -$ | $\frac{1}{2} -$ |

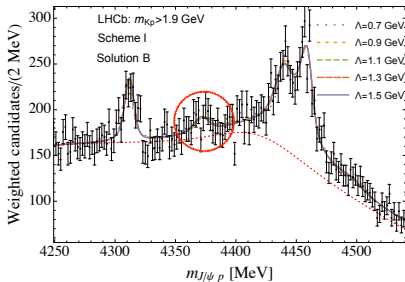
☞ No need for $\Lambda_c \bar{D}^{(*)}$

Scheme I: pure contact potential w/o $\Lambda_c \bar{D}^{(*)}$

Solution A



Solution B



☞ $\Lambda_{\text{soft}} \sim \sqrt{2\mu\delta} \sim 0.7$ GeV

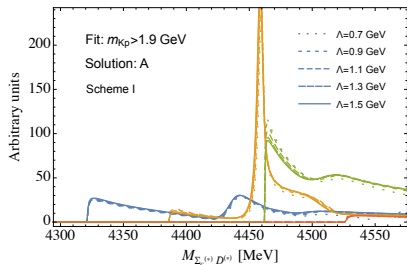
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| | $\Sigma_c \bar{D}^*$ | $P_c(4440)$ | $P_c(4457)$ |
|-------|----------------------|-------------|---------------|
| Fit A | $\frac{1}{2}$ | — | $\frac{3}{2}$ |
| Fit B | $\frac{1}{2}$ | — | $\frac{1}{2}$ |

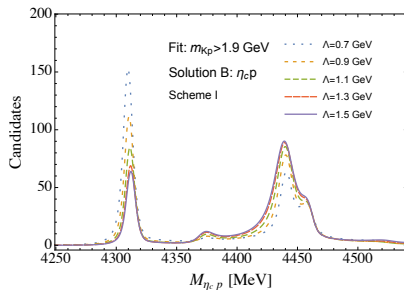
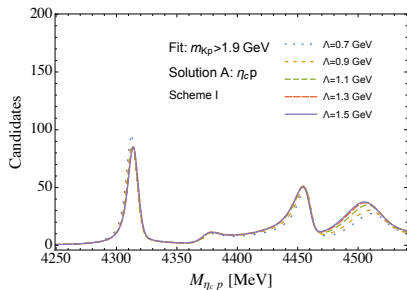
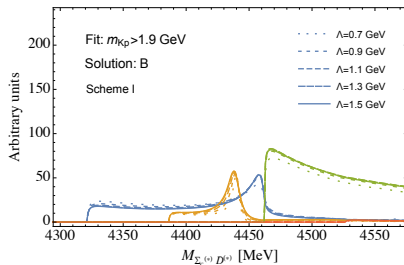
☞ No need for $\Lambda_c \bar{D}^{(*)}$

Scheme I: pure contact potential $w/o \Lambda_c \bar{D}^{(*)}$

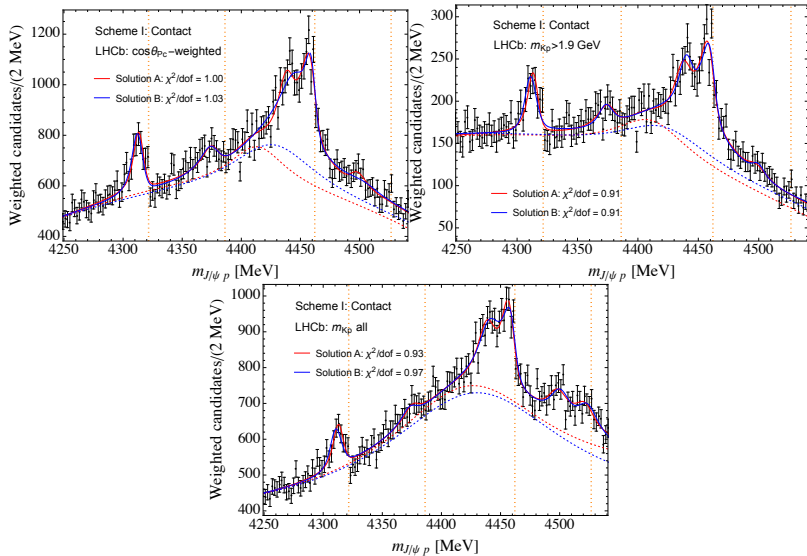
Solution A



Solution B



Scheme I: Contact fits to three sets of LHCb data



Scheme I: pole positions

| | DC ([MeV]) | Solution A | | Solution B | |
|---------------|---------------------------------|-----------------|---------------------|-----------------|---------------------|
| | | J^P | Pole [MeV] | J^P | Pole [MeV] |
| $P_c(4312)$ | $\Sigma_c \bar{D}$ (4321.6) | $\frac{1}{2}^-$ | 4314(1) - 4(1) i | $\frac{1}{2}^-$ | 4312(2) - 4(2) i |
| $P_c(4380)^*$ | $\Sigma_c^* \bar{D}$ (4386.2) | $\frac{3}{2}^-$ | 4377(1) - 7(1) i | $\frac{3}{2}^-$ | 4375(2) - 6(1) i |
| $P_c(4440)$ | $\Sigma_c \bar{D}^*$ (4462.1) | $\frac{1}{2}^-$ | 4440(1) - 9(2) i | $\frac{3}{2}^-$ | 4441(3) - 5(2) i |
| $P_c(4457)$ | $\Sigma_c \bar{D}^*$ (4462.1) | $\frac{3}{2}^-$ | 4458(2) - 3(1) i | $\frac{1}{2}^-$ | 4462(4) - 5(3) i |
| P_c | $\Sigma_c^* \bar{D}^*$ (4526.7) | $\frac{1}{2}^-$ | 4498(2) - 9(3) i | $\frac{1}{2}^-$ | 4526(3) - 9(2) i |
| P_c | $\Sigma_c^* \bar{D}^*$ (4526.7) | $\frac{3}{2}^-$ | 4510(2) - 14(3) i | $\frac{3}{2}^-$ | 4521(2) - 12(3) i |
| P_c | $\Sigma_c^* \bar{D}^*$ (4526.7) | $\frac{5}{2}^-$ | 4525(2) - 9(3) i | $\frac{5}{2}^-$ | 4501(3) - 6(4) i |

✎ * NOT the broad $P_c(4380)$ reported by LHCb in 2015

✎ Bound states with respect to the dominant channel (DC)

✎ $\Sigma_c^* \bar{D}^*$ states are not seen yet, production rate suppressed?

↪ prompt production in the pp collision in the LHC

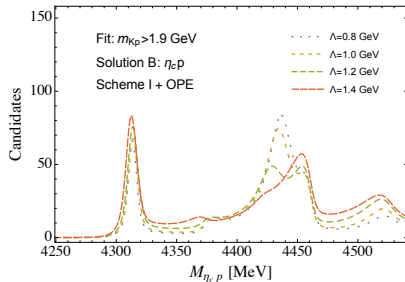
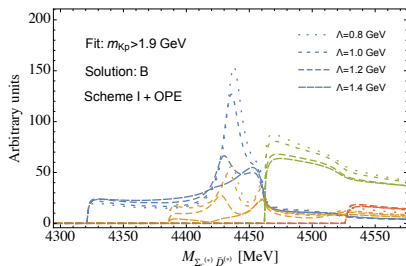
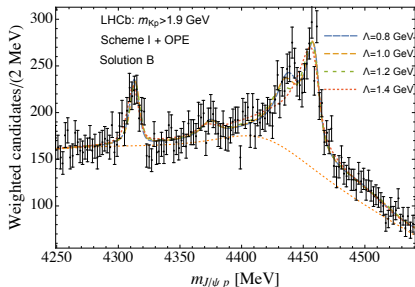
P. Ling, X.-H. Dai, MLD and Q. Wang, arXiv:2104.11133

Scheme I + OPE w/o $\Lambda_c \bar{D}^{(*)}$

👉 No solution A

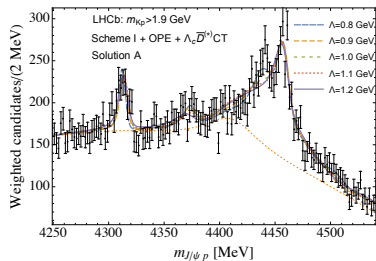
👉 Solution B:
Cut-off dependent

👉 $\Lambda_{\text{soft}} \sim 700$ MeV

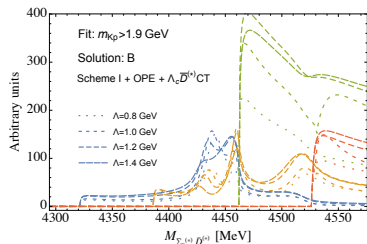
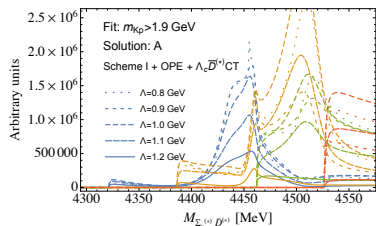
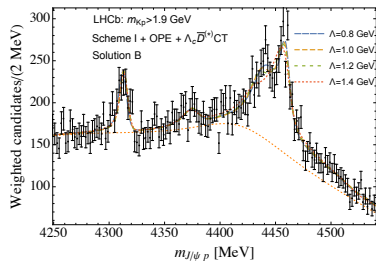


Scheme I + OPE + CT for $\Lambda_c \bar{D}^{(*)}$

Solution A



Solution B

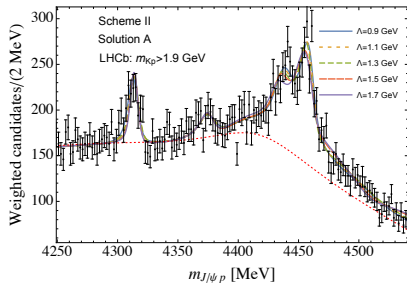


• Cut-off dependent

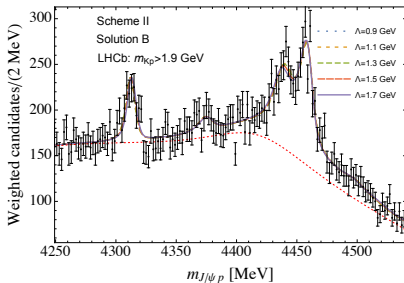
• $\Lambda_{\text{soft}} \sim 900$ MeV $\Lambda_c \bar{D}^{(*)}$

Scheme II: contact + OPE + S-D w/o $\Lambda_c \bar{D}^{(*)}$

Solution A



Solution B



☞ $\Lambda_{\text{soft}} \sim 0.7$ GeV

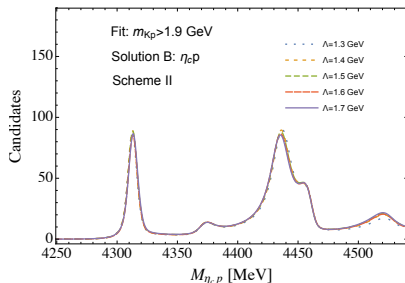
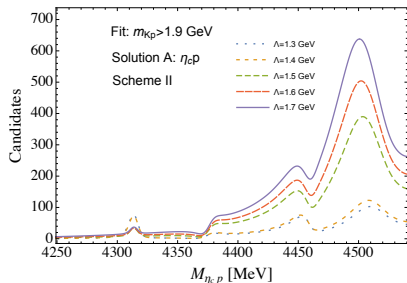
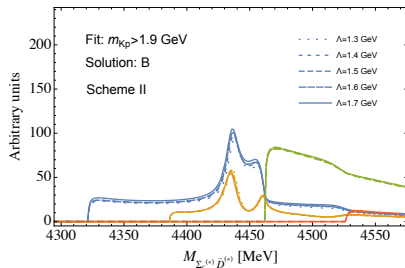
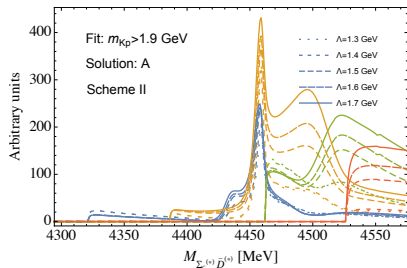
☞ Cutoff-independent only for solution B

Scheme II: contact + OPE + S-D w/o $\Lambda_c \bar{D}^{(*)}$

Solution A

$\Lambda_{\text{soft}} \sim 0.7 \text{ GeV}$

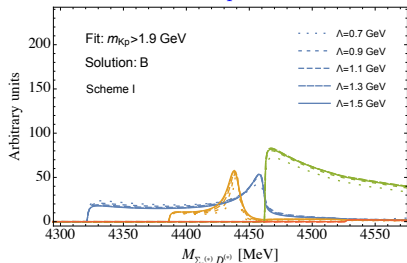
Solution B



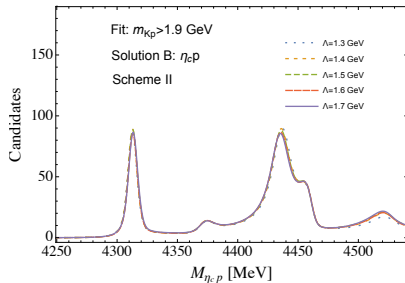
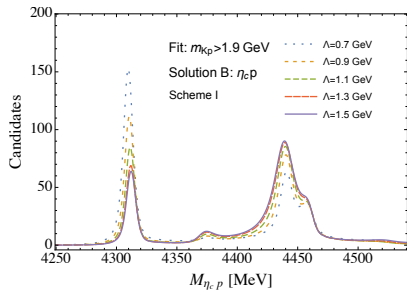
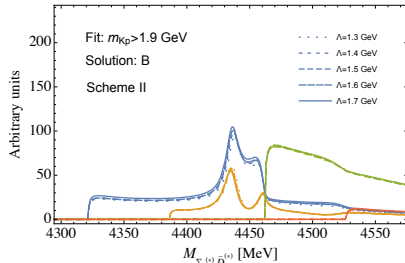
Scheme I vs Scheme II w/o $\Lambda_c \bar{D}^{(*)}$ $\Lambda_{\text{soft}} \sim 0.7 \text{ GeV}$

Solution B

Scheme I: pure contact

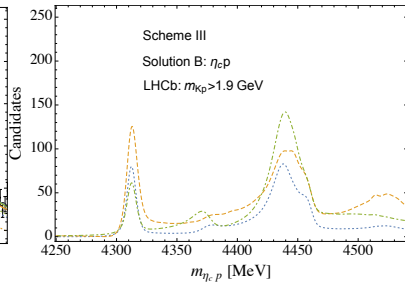
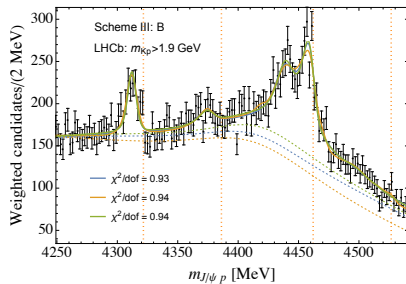


Scheme II: contact + OPE + SD



Scheme III: CT + OPE + SD w/ $\Lambda_c \bar{D}^{(*)}$ $\Lambda_{\text{soft}} \sim 0.9$ GeV

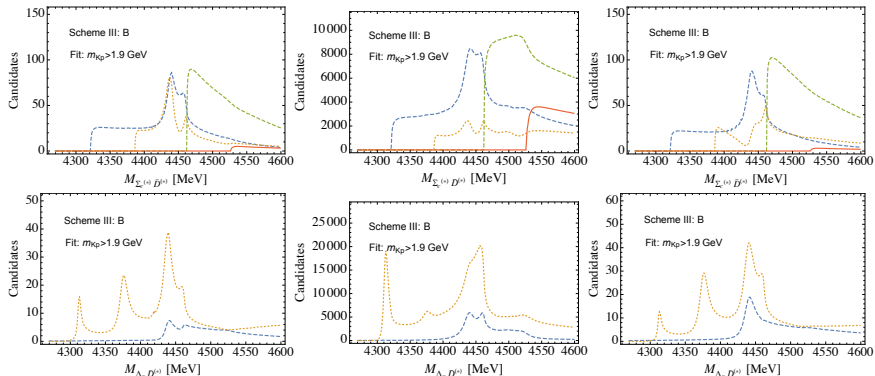
Solution B



👉 $\Lambda = 1.3$ GeV.

👉 Overdetermined.

Solution B



👉 J/ψ data alone are not enough to constrain $\Lambda_c \bar{D}^{(*)}$ interactions.

Summary & Outlook

- ☞ Solving Lippmann-Schwinger equation with respect to
 - ▶ Unitarity, three-body cut
↪ width of $\Sigma_c^{(*)}$
 - ▶ Coupled-channels
↪ cut-off independence: OPE \rightarrow SD counter term
 - ▶ Heavy quark spin symmetry
↪ 7 $\Sigma_c^{(*)} \bar{D}^{(*)}$ molecular states
- ☞ Preferred spin assignment (Solution B)

$$P_c(4312) \quad : \quad 1/2^-, \quad (\bar{D}\Sigma_c)$$

$$P_c(4440) \quad : \quad 3/2^-, \quad (\bar{D}^*\Sigma_c)$$

$$P_c(4457) \quad : \quad 1/2^-, \quad (\bar{D}^*\Sigma_c)$$

- ☞ We can not say much about $\Lambda_c \bar{D}^{(*)}$ interaction without data in this channel.
- ☞ A narrow $P_c(4380)$, different from the broad one reported by LHCb in 2015.

Thank you very much for your attention!