# Light Baryon Spectroscopy

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# **Motivation**

### **Theoretical Predictions**



[U. Loering, et al., Eur.Phys.J.A10:395 (2001)]

[R. Edwards et al., Phys.Rev.D 84 (2011) 07450]

[Eichmann, Fischer, Few Body Syst. 60 (2019) 1,2]

Discrepancies between measurement and calculations: "missing resonances" and level ordering

### **Theoretical Predictions**



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 $\rightarrow$  What are the relevant degrees of freedom?



### **Theoretical Predictions**



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 $\rightarrow$  What are the relevant degrees of freedom?

Most resonances observed in  $\pi N$  scattering:  $\rightarrow$  Experimental bias?



#### Resonances



Total cross section sensitive to dominant resonance contributions:

$$\sigma \sim |E_{0+}|^2 + |E_{1+}|^2 + |M_{1+}|^2 + |M_{1-}|^2 + \dots$$

Huge experimental effort from different experiments: Measurement of a wide range of final states



### Polarization Observables in photoproduction of pseudoscalar meson



Polarization Observables are a tool to access weak resonance contributions, sensitive to interference terms:

$$\Sigma \sim -2E_{0+}^*E_{2+}+2E_{0+}^*E_{2-}-2E_{0+}^*M_{2+}+\dots$$

			Target			Recoil		Target+Recoil				
		_	_	-	×'	y'	z'	×'	×'	z'	z'	
Photon		×	У	z	-	-	-	х	z	х	z	
unpolarized	$\sigma$	-	Т	-	-	Р	-	$T_{x'}$	$-L_{x'}$	$T_{z'}$	$L_{z'}$	
linearly pol.	Σ	Н	(-P)	-G	$O_{x'}$	(-T)	$O_{z'}$	-	-	-	-	
circularly pol.	-	F	-	-E	$-C_{x'}$	-	$-C_{z'}$	-	-	-	-	

#### **Complete Experiment:**

Extraction of the amplitudes without model dependence

For a single pseudoscalar meson at least **well-defined** 8 observables necessary

[Chiang and Tabakin, Phys.Rev.C 55 (1997) 2054-2066]

### **Complete Experiments**

#### Extraction of complete sets using graph theory:

#### **Complete Experiment:**

Extraction of the amplitudes without model dependence

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[Chiang and Tabakin, Phys.Rev.C 55 (1997) 2054-2066]



#### Electroproduction: 13 Observables

[Y. Wunderlich et al., Phys.Rev.C 102 (2020) 3, 034605]

Two meson photoproduction: 16 Observables [P. Kroenert et al., Phys.Rev.C 103 (2021) 1, 014607]

See presentations by Y. Wunderlich and P. Kroenert on Wednesday at 9:00h and 9:25! 5

### Examples of Important Experiments in the Last Years

A2 experiment at MAMI Mainz, Germany CBELSA/TAPS experiment Bonn, Germany CLAS experiment at JLAB Newport News, US



Common features:

- Good angular coverage of detector systems
- Polarized photons and polarized targets

Important differences:

- Different sensitivities (for charged or neutral particles)
- Different photon energies
- $\rightarrow$  Different physical foci

### Examples of Important Experiments in the Last Years

A2 experiment at MAMI Mainz, Germany

25000

1960 1970 1980

2000 2010 2020

Year

CBELSA/TAPS experiment Bonn, Germany CLAS experiment at JLAB Newport News, US



[D. Ireland et al., Prog.Part.Nucl.Phys. 111 (2020) 103752]

## **Measurement of Observables**

#### Cross Section Measurements at A2



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### Cusp Effect visible in $\eta$ Photoproduction by CBELSA/TAPS



High precision measurement of the Beam Asymmetry with high angular coverage

Cusp effect of the  $\eta'$  threshold visible in the Legendre coefficients



[F. Afzal et al, Phys.Rev.Lett. 125 (2020) 15, 152002]

### **Double Polarization Observable E** ( $\gamma p \rightarrow p \pi^0$ ): CBELSA/TAPS



E is a helicity asymmetry: Two spin configurations possible



Fits to the data: BnGa14-02 SAID 2015 JüBo16-1 [M. Gottschall et al., Phys. Rev. Lett. 112, 012003 (2014) Eur. Phys. J. A 57.1 (2021), p. 40]

### Measurements off (polarized) Neutrons with A2



Narrow peak observed in  $\eta$  photoproduction

Polarization observables used to shed further light on this structure

[D. Werthmüller et al., Phys.Rev. C90 (2014) no.1, 015205]

[L. Witthauer *et al.*, Phys. Rev. Lett. **117**, no. 13, 132502 (2016)]

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### Strangeness Production with CLAS: $\gamma p \rightarrow K \Lambda$

Strangeness production self analyzing, allows the extraction of observables with recoil polarization

$$\frac{d\sigma}{d\Omega} = \left(\frac{d\sigma}{d\Omega}\right)_0 \{1 - P^{\gamma} \Sigma \cos 2\phi + \alpha \cos \theta_x P^{\gamma} O_x \sin 2 + \alpha \cos \theta_y P - \alpha \cos \theta_y P^{\gamma} T \cos 2\phi + \alpha \cos \theta_z P^{\gamma} O_z \sin 2\phi\},\$$



red: ANL-Osaka green: BnGa14 blue: BnGa refit (2016)

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## **Observables in Multi-Meson Final States**

- Multi-meson final states like  $\gamma p \to p \pi^0 \pi^0$  or  $\pi^0 \eta$  preferred at higher energies
- Probes the high mass region, where the missing resonances occur
- Can help to observe cascading decays



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## Polarization Observables T, P, H ( $\gamma p \rightarrow p \pi^0 \pi^0$ ): CBELSA/TAPS

Here:

only results shown in quasi two-body kinematics







Observables also extracted for different kinematic variables

Full three-body kinematics allows the measurement of further observables.

Talk by T. Seifen on Wednesday at 07:45h! 13

## First Indication of Triangle Singularities in $\gamma p \rightarrow p \pi^0 \eta$ by CBELSA/TAPS



Structure observed in the  $p\eta$  invariant mass

Triangle singularity can describe this structure

First observation of a triangle singularity in baryon spectroscopy?

See presentation by M. Nanova on Wednesday at 09:50!

## Interpretation

### **Multipoles and CGLN Amplitudes**

Multipoles give informations about the intermediate states, can be combined into four CGLN amplitudes:

$$egin{aligned} F_1(W,z) &= \sum_{\ell=0}^\infty [\ell M_{\ell+} + E_{\ell+}] \cdot P'_{\ell+1}(z) + [(\ell+1)M_{\ell-} + E_{\ell-}] \cdot P'_{\ell-1}(z) \ F_2(W,z) &= \sum_{\ell=0}^\infty \ldots \end{aligned}$$



with  $z = \cos \theta_{\pi}$  and the Legendre polynomials  $P_{\ell}(z)$ .

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All observables can be expressed in CGLN amplitudes, for example:

$$\hat{\Sigma} = \frac{\Sigma \cdot \sigma(\theta_{\pi})}{\rho_0} = -\sin^2 \theta_{\pi} \cdot Re\left[\frac{1}{2}|F_3|^2 + \frac{1}{2}|F_4|^2 + F_2^*F_3 + F_1^*F_4 + \cos \theta F_3^*F_4\right]\rho_0$$

with the density of states  $\rho_0 = k/q$ .

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#### Example of a Truncated Partial Wave Analysis

Observable described by

$$\check{T} = T \cdot \sigma = rac{q}{k} \sin heta \left[ \sum_{h=0}^{2L_{max}-1} A_h (\cos heta)^h 
ight]$$

• using S- and P-waves 
$$(L_{max} = 1)$$
:  
 $\check{T} = \frac{q}{k} \sin \theta \left[ A_0 + A_1 \cdot \cos \theta \right]$ 

• using S-, P- and D-waves 
$$(L_{max} = 2)$$
:

$$\check{\mathcal{T}} = \frac{q}{k}\sin\theta[A_0 + A_1 \cdot \cos\theta + A_2 \cdot \cos^2\theta + A_3 \cdot \cos^3\theta]$$

• using S-, P-, D- and F-waves ( $L_{max} = 3$ ):

$$\check{\mathcal{T}} = \frac{q}{k} \sin \theta [A_0 + A_1 \cdot \cos \theta + A_2 \cdot \cos^2 \theta + A_3 \cdot \cos^3 \theta + A_4 \cdot \cos^4 \theta + A_5 \cdot \cos^5 \theta]$$

#### First Interpretation with a Truncated Partial Wave Analysis



[Y. Wunderlich, F. Afzal, A. Thiel, R. Beck, Eur.Phys.J. A53 (2017) no.5, 86]

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### New Fits from different Analyses

New observables for  $p\pi^0$  have been included in the analyses of the groups:

- BnGa (black)
- JüBo (red)
- SAID (blue)



For all other multipoles see: [Anisovich et al., Eur.Phys.J. A52 (2016) no.9, 284]

### New Fits from different Analyses

New observables for  $p\pi^0$  have been included in the analyses of the groups:

- BnGa (black)
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Variance between the different analyses decreases!





For all other multipoles see: [Anisovich et al., Eur.Phys.J. A52 (2016) no.9, 284]

### Comparison between PDG values

- Until 2010: almost only results from pion nucleon scattering used in the PDG, only few pion photoproduction data used
- PWA groups include photoproduction data with different final states from several experiments
- Now: new values from the fits are entering the PDG

Particle	$J^{P}$	overall	$N\gamma$	$N\pi$	$\Delta \pi$	$N\sigma$	$N\eta$	$\Lambda K$	$\Sigma K$	$N\rho$	$N\omega$	$N\eta'$
N	$1/2^+$	****										
N(1440)	$1/2^+$	****	***	****	***	***	-			-		
N(1520)	3/2-	****	****	****	****	**	****					
N(1535)	$1/2^{-}$	****	****	****	***	*	****					
N(1650)	$1/2^{-}$	****	***	****	***	*	****	*				
N(1675)	$5/2^{-}$	****	****	****	****	***	*	*	*	2.00		
N(1680)	$5/2^+$	****	****	****	****	***	*	*	*			
N(1700)	$3/2^{-}$	***	**	***	***	*	*		2.1	2.00		
N(1710)	$1/2^+$	***	***	***	*_		***	**	*	*	*	
N(1720)	$3/2^+$	****	****	****	***	*	*	****	*	*_	*	
N(1860)	$5/2^{+}$	**	*	**		*	*					
N(1875)	$3/2^{-}$	***	**	**	*	**	*	*	*	*	*	
N(1880)	$1/2^+$	***	**	*	**	*	*	**	**		**	
N(1895)	$1/2^{-}$	****	****	*	*	*	****	**	**	*	*	****
N(1900)	$3/2^+$	****	****	**	**	*	*	**	**	2.00	*	**
N(1990)	$7/2^+$	**	**	**			*	*	*			
N(2000)	$5/2^+$	**	**	*_	**	*	*	-			*	
N(2040)	$3/2^{+}$	*		*								
N(2060)	$5/2^{-}$	***	***	**	*	*	*	*	*	*	*	
N(2100)	$1/2^{+}$	***	**	***	**	**	*	*		*	*	**
N(2120)	$3/2^{-}$	***	***	**	**	**		**	*		*	*
N(2190)	$7/2^{-}$	****	****	****	****	**	*	**	*	*	*	
N(2220)	$9/2^+$	****	**	****			*	*	*			
N(2250)	$9/2^{-}$	****	**	****			*	*	*			
N(2300)	$1/2^{+}$	**		**								
N(2570)	$5/2^{-}$	**		**								
N(2600)	$11/2^{-}$	***		***								
N(2700)	$13/2^+$	**		**								

Large improvement, but still lot of work to be done!

### Still Many Open Questions...

• Do parity doublets exist? Is chiral symmetry restored at high energies? Do they exist for all high mass states?

Not predicted by the current lattice QCD calculations nor by constituent quark models.



Parity Doublets at high masses:

Partner of the  $\Delta(1950)7/2^+$  seems to be missing Search in different final states revealed state with  $7/2^-$  at much higher masses (2200 MeV)  $\rightarrow$  No parity partner found



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Fit with and without  $\Delta(2200)7/2^-$  reveals high sensitivity of the data sets Further identification of weak resonance contributions possible!



## Still Many Open Questions...

• Do parity doublets exist? Is chiral symmetry restored at high energies? Do they exist for all high mass states?

Not predicted by the current lattice QCD calculations nor by constituent quark models.



• Is it possible to do a complete experiment? How many observables and which precision is needed?

Recently: Correction of the decay parameter  $\alpha$  by BESIII [M. Ablikim et al., Nature Phys.15, 631 (2019)] Parameter has a substantial influence on the polarization observable for  $\Lambda$  production

Fierz identities of the measured (double) polarization observables

 $O_x^2 + O_z^2 + C_x^2 + C_z^2 + \Sigma^2 - T^2 + P^2 = 1$  $\Sigma P - C_x O_z + C_z O_x - T = 0$  Recently: Correction of the decay parameter  $\alpha$  by BESIII [M. Ablikim et al., Nature Phys.15, 631 (2019)] Parameter has a substantial influence on the polarization observable for  $\Lambda$  production

Fierz identities of the measured (double) polarization observables

$$O_x^2 + O_z^2 + C_x^2 + C_z^2 + \Sigma^2 - T^2 + P^2 = 1$$
  
$$\Sigma P - C_x O_z + C_z O_x - T = 0$$



Measured (double) polarization observables for  $\gamma p \rightarrow K \Lambda$  can be used to give an additional view on the value the decay parameter

[D. Ireland et al., Phys.Rev.Lett. 123 (2019) no.18, 182301]

# Summary

### **Conclusion and Outlook**

- New era of experiments allows precise measurements of (polarization) observables for various reactions
- Data is included in the different partial wave analyses and the multipoles are converging
- New polarization data will help to understand the resonance spectrum and will provide an experimental basis for comparison with constituent quark models, lattice QCD or other methods
- $\rightarrow~$  Baryon Spectroscopy has reached the high precision area!

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Thank you for your attention.