

Quarkoniumlike Mesons in the Diabatic Approach

Roberto Bruschini

roberto.bruschini@ific.uv.es

Instituto de Física Corpuscular

University of Valencia - CSIC

19th International Conference on Hadron Spectroscopy
and Structure in memoriam Simon Eidelman
Mexico City, 28 July 2021

1 Introduction

- Motivation
- Born-Oppenheimer Approximation
- Energy Levels in Lattice QCD

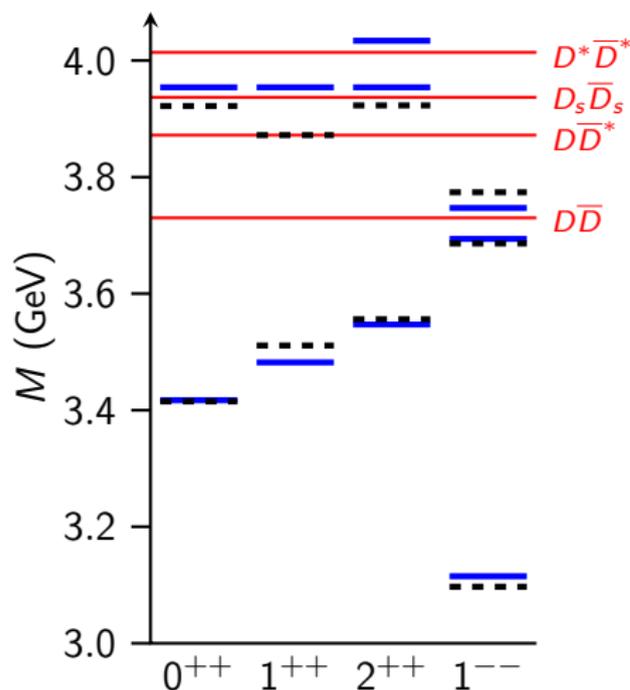
2 Diabatic Approach

- Diabatic Formalism
- String Breaking

3 Results

- Bottomoniumlike Mesons
- Charmoniumlike Mesons
- Summary and References

Unconventional Charmoniumlike Mesons

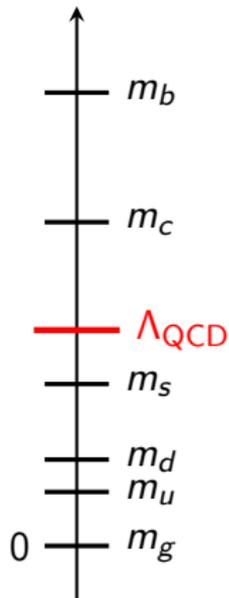


- - - experiment
 — quark model
 — thresholds

Threshold effects?

Unconventional mesons are often found near open-flavor thresholds.

Description in Heavy and Light Fields



Heavy ($m \gg \Lambda_{\text{QCD}}$)

- heavy quarks c, b

Light ($m \lesssim \Lambda_{\text{QCD}}$)

- gluons g
- light quarks u, d, s

Integrating the Light Fields

- 1 Separate kinetic energy of the heavy quarks

$$H |\psi\rangle = E |\psi\rangle, \quad H = K_{\text{heavy}} + H_{\text{light}}^{(\text{heavy})}.$$

- 2 Solve the light-field Hamiltonian for **static** heavy quarks

$$H_{\text{light}}^{(\text{heavy})} \rightarrow H_{\text{static}}(r), \quad H_{\text{static}}(r) |\zeta_i(r)\rangle = V_i(r) |\zeta_i(r)\rangle.$$

Static energy levels

The static energies $V_i(r)$ can be calculated *ab initio* in Lattice QCD.

Adiabatic Wave Function

Adiabatic expansion

$$|\psi\rangle = \sum_i \int dr \psi_i(r) |r\rangle |\zeta_i(r)\rangle$$

- Light field states calculated at **the same position** of the heavy quarks
- One wave function for each light-field energy

Adiabatic Schrödinger Equation

$$\sum_j \left[-\frac{\hbar^2}{2\mu} (\delta_{ij} \nabla + \tau_{ij}(r))^2 + \delta_{ij} (V_i(r) - E) \right] \psi_j(r) = 0$$

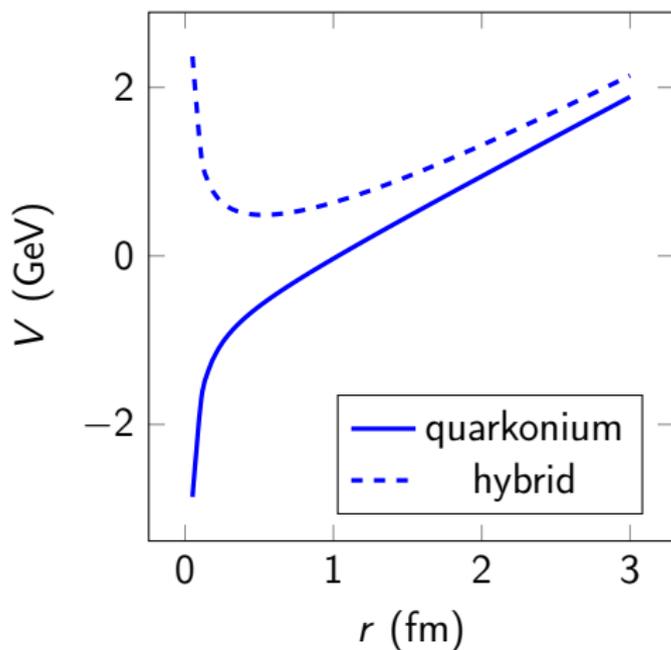
Non-adiabatic coupling terms

The kinetic energy term mixes different channels through the non-adiabatic couplings $\tau_{ij}(r) = \langle \zeta_i(r) | \nabla \zeta_j(r) \rangle$.

Adiabatic potentials

The potentials $V_i(r)$ in the Schrödinger equation are the energy levels calculated in Lattice QCD.

Adiabatic Potentials in Quenched Lattice QCD

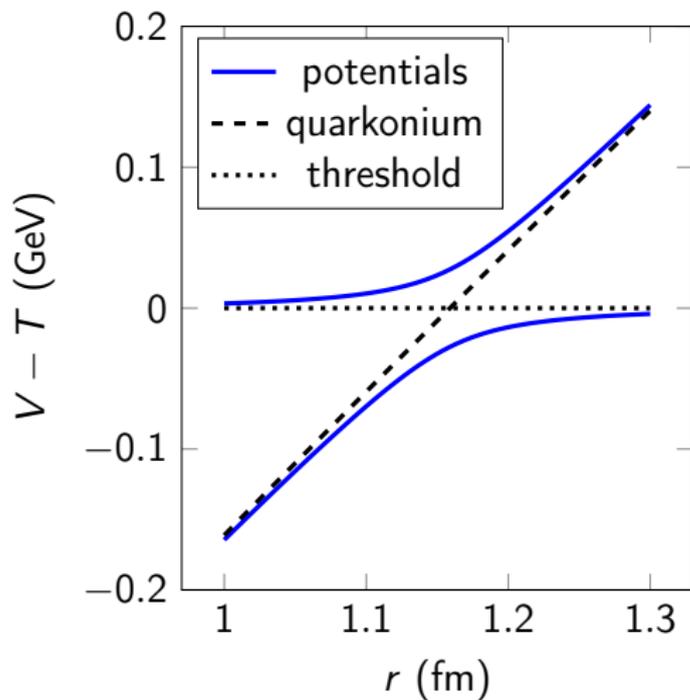


Quenched Lattice QCD

Gluons only,
without light quarks

- Ground state: quarkonium potential
- Excited states: hybrid potentials

Adiabatic Potentials in Unquenched Lattice QCD



Unquenched Lattice QCD

Gluons and light quarks

String breaking

Channel mixing significant near the avoided crossing

From Adiabatic to Diabatic

Diabatic expansion

$$|\psi\rangle = \sum_i \int dr \tilde{\psi}_i(r, r_0) |r\rangle |\zeta_i(r_0)\rangle$$

Diabatic channels

$$\tilde{\psi}_i(r, r_0) \rightarrow \psi_{Q\bar{Q}}(r), \psi_{M\bar{M}}(r)$$

- Light field states are calculated at a fixed position r_0 .
- For r_0 far from the avoided crossing, they correspond to **quark-antiquark** and **meson-meson**.

The Diabatic Schrödinger Equation

$$\left[\begin{pmatrix} -\frac{\nabla^2}{2\mu_{Q\bar{Q}}} & 0 \\ 0 & -\frac{\nabla^2}{2\mu_{M\bar{M}}} \end{pmatrix} + \begin{pmatrix} V_{Q\bar{Q}}(r) & V_{\text{mix}}(r) \\ V_{\text{mix}}(r) & T_{M\bar{M}} \end{pmatrix} - E \right] \begin{pmatrix} \psi_{Q\bar{Q}}(r) \\ \psi_{M\bar{M}}(r) \end{pmatrix} = 0$$

Diabatic potential matrix

The potential couples quark-antiquark and meson-meson.

Adiabatic-to-diabatic transformation

The eigenvalues of the diabatic potential matrix are the adiabatic potentials calculated in Lattice QCD.

Mixing Effects on the Quarkoniumlike Spectrum

Above threshold

- Meson states acquire **decay width**.
- Quarkonium masses are shifted.

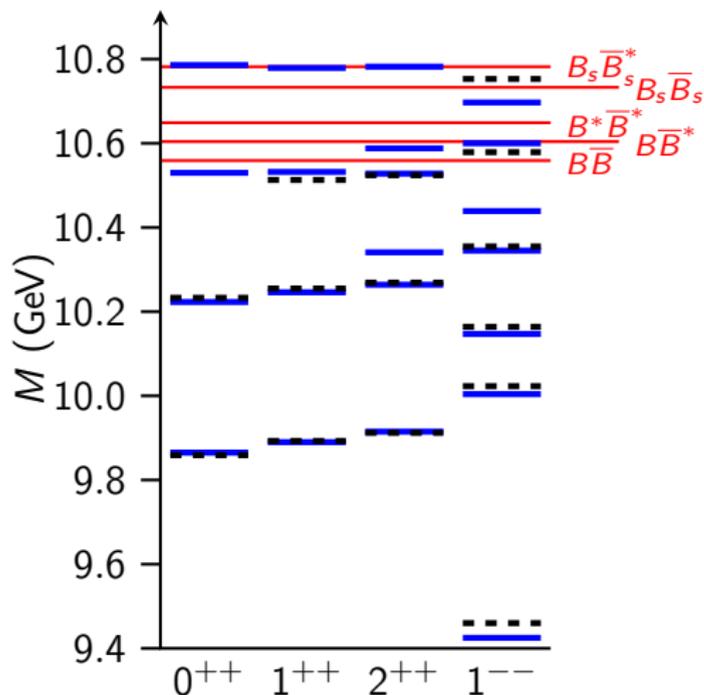
Below threshold

- Meson states acquire **molecular components**.
- Unconventional mesons may appear near threshold.

Unified description above and below threshold

Appearance of unconventional mesons and resonance decays are described by the same mixing potential.

Bottomoniumlike Spectrum



- - - experiment
 — diabatic
 — thresholds

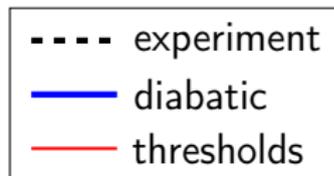
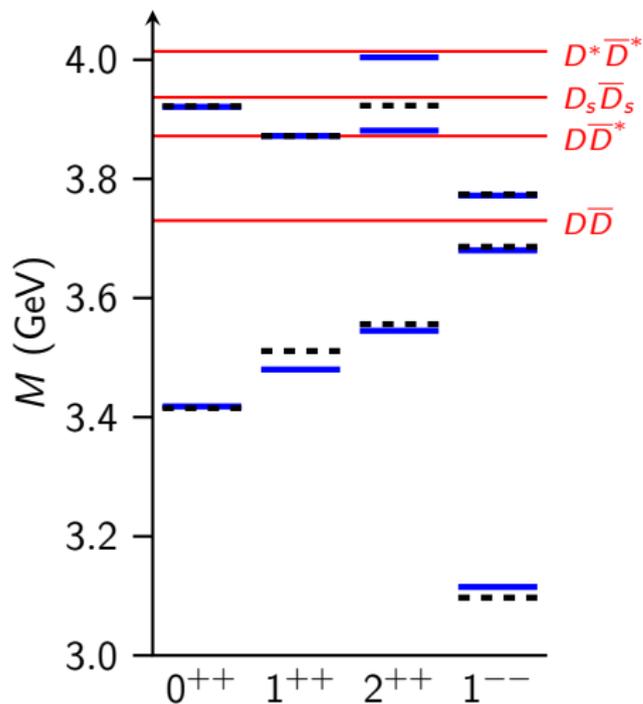
- From Lattice QCD
 $\Delta \approx 50$ MeV
- Relatively small threshold effects

Strong Decay Widths of Bottomoniumlike States

J^{PC}	M	$\Gamma_{B\bar{B}}$	$\Gamma_{B\bar{B}^*}$	$\Gamma_{B^*\bar{B}^*}$	$\Gamma_{B_s\bar{B}_s}$	$\Gamma_{\text{total}}^{\text{Theor}}$	$\Gamma_{\text{total}}^{\text{Expt}}$
0^{++}	10785.8	1.6		5.3	0.7	7.6	
1^{++}	10778.9		0.2	1.7		1.9	
2^{++}	10588.4	4.3				4.3	
2^{++}	10782.3	5.4	1.5	21.0	10.4	38.3	
1^{--}	10599.8	21.9				21.9	20.5 ± 2.5
1^{--}	10697.0	2.0	1.0	38.0		41.0	

Masses and widths in MeV units

Charmoniumlike Spectrum



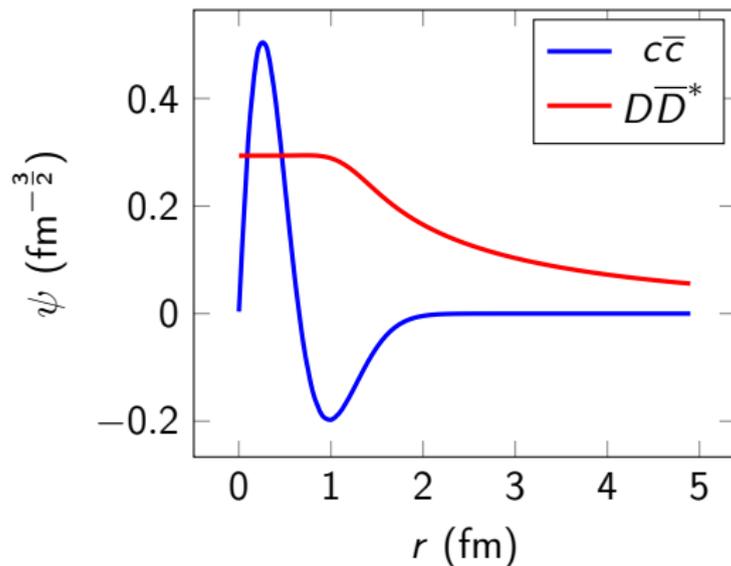
- Fitting $\chi_{c1}(3872)$
 $\Delta \approx 130$ MeV
- More prominent threshold effects

Strong Decay Widths of Charmoniumlike States

J^{PC}	M	$\Gamma_{D\bar{D}}$	$\Gamma_{D\bar{D}^*}$	$\Gamma_{D_s\bar{D}_s}$	$\Gamma_{\text{total}}^{\text{Theor}}$	$\Gamma_{\text{total}}^{\text{Expt}}$
0^{++}	3920.9	0.6			0.6	
2^{++}	3881.1	49.5	0.4		49.9	35.3 ± 2.8
2^{++}	4003.9	4.8	6.3	3.5	14.5	
1^{--}	3771.7	20.2			20.2	27.2 ± 1.0

Masses and widths in MeV units

Diabatic Wave Function of $\chi_{c1}(3872)$



- about 3% of $c\bar{c}$
- about 97% of $D\bar{D}^*$
- $\sqrt{\langle r^2 \rangle} \approx 11$ fm

Diabatic $\chi_{c1}(3872)$

It can be described as a $D\bar{D}^*$ molecule created by the mixing between $D\bar{D}^*$ and $c\bar{c}$.

Summary

- The Born-Oppenheimer approximation gives a description of quarkonium firmly based on Lattice QCD.
- The diabatic framework allows to include open-flavor mesons and string breaking into this description nonperturbatively.
- The diabatic potential matrix, calculable from Lattice QCD, can give account of the spectrum as well as strong decays of quarkoniumlike mesons.

For Further Reading

-  R. Bruschini and P. González.
Diabatic description of charmoniumlike mesons.
[Phys. Rev. D 102, 074002 \(2020\)](#).
-  R. Bruschini and P. González.
Diabatic description of charmoniumlike mesons. II.
Mass corrections and strong decay widths.
[Phys. Rev. D 103, 074009 \(2021\)](#).
-  R. Bruschini and P. González.
Diabatic description of bottomoniumlike mesons.
[Phys. Rev. D 103, 114016 \(2021\)](#).