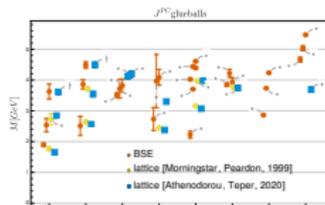
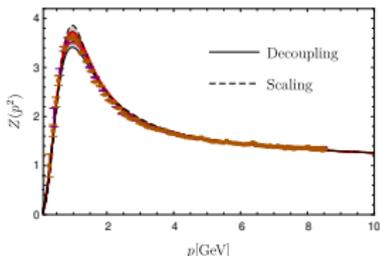
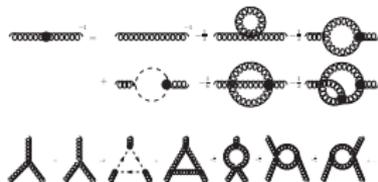


Glueballs with even charge parity from functional equations



Markus Q. Huber

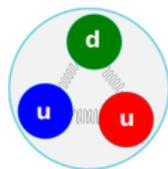
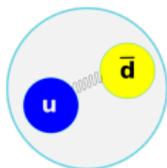
Institute of Theoretical Physics
Giessen University

MQH, Phys.Rev.D 101, arXiv:2003.13703

MQH, Fischer, Sanchis-Alepuz, Eur.Phys.J.C 80,
arXiv:2004.00415

19th International Conference on Hadron Spectroscopy and Structure:
HADRON2021, virtually in Mexico City, July 26, 2021

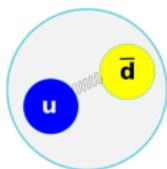
Bound states in QCD



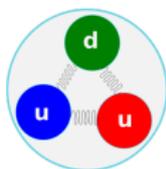
Mesons

Baryons

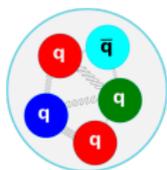
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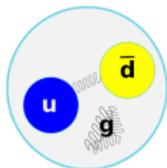
Pentaquarks

First observations 2015 (LHCb)



Tetraquarks

Increasing number of confirmed states. Bound state equations perspective: [Eichmann, Fischer, Heupel, Santowsky, Wallbott '20]



Hybrids



Glueballs

States of pure 'radiation'

Glueball observations

Experimental candidates, but situation not conclusive.

Scalar glueball: 0^{++} , mixing with scalar isoscalar mesons

Candidate reaction: $J/\psi \rightarrow \gamma + 2g$

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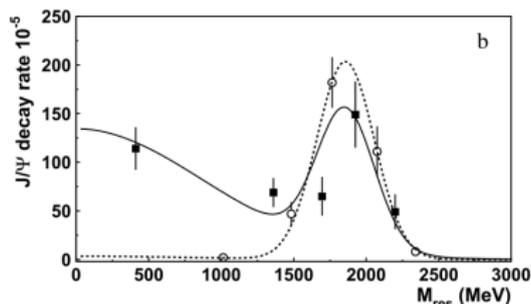
Candidate reaction: $J/\psi \rightarrow \gamma + 2g$

Recent analysis of BESIII data [Sarantsev, Denisenko, Thoma, Klempt '21]:

$$M = 1865 \pm 25^{+10}_{-30} \text{ MeV},$$

$$\Gamma = 370 \pm 50^{+30}_{-20} \text{ MeV}$$

→ Talks by Sarantsev, Klempt



Glueball calculations

Yang-Mills theory

- “Isolated” problem: only gluons
- Clean picture: well-established lattice results

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QCD glueballs: mixing with quarks

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QCD glueballs: mixing with quarks

Unquenching on the lattice [Gregory et al. '12]:

- Much higher statistics required (poor signal-to-noise ratio)
- Continuum extrapolation and inclusion of fermionic operators still to be done
- Mixing with $\bar{q}q$ challenging
- Tiny (e.g., 0^{++} , 2^{++}) to moderate unquenching effects (e.g., 0^{-+}) found
- $m_{\pi} = 360 \text{ MeV}$

Glueball calculations

Yang-Mills theory

- “Isolated” problem: only gluons
- Clean picture: well-established lattice results
- **Functional methods: High quality input available for bound state equations**

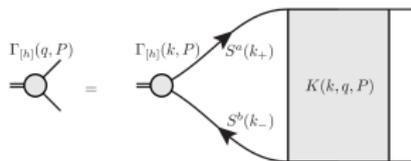
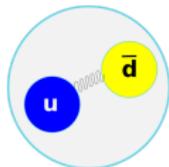
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Hadrons from bound state equations

Example: Meson

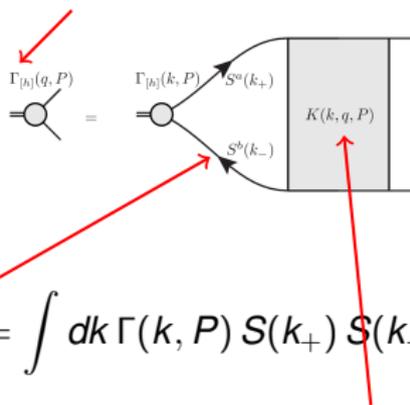
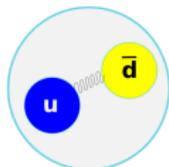


$$\text{Integral equation: } \Gamma(q, P) = \int dk \Gamma(k, P) S(k_+) S(k_-) K(k, q, P)$$

Hadrons from bound state equations

Bethe-Salpeter amplitude

Example: Meson



$$\text{Integral equation: } \Gamma(q, P) = \int dk \Gamma(k, P) S(k_+) S(k_-) K(k, q, P)$$

Ingredients:

- Quark propagator S

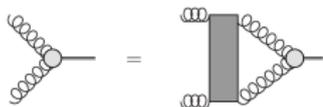
$$\text{Diagram: } \text{Quark propagator with dressing}^{-1} = \text{tree-level propagator}^{-1} + \text{dressed propagator}^{-1}$$

The diagram shows a quark propagator with a dressing loop. The left side is a circle with an arrow, labeled $S(p)$. The right side is a tree-level propagator (a straight line with an arrow) labeled $S_0(p)$, plus a term where a quark line with a gluon loop (represented by a wavy line) is attached to the propagator. The loop is labeled $D_{\mu\nu}(p-q)$ and the vertices are γ_μ and $\Gamma_\mu(p, q)$. The propagator part of the loop is labeled $S(q)$.

Nonperturbative diagram: full momentum dependent dressings
 → numerical solution

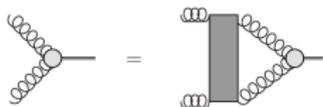
- Interaction kernel K
- Constrained by symmetries

Glueball BSE



Need  and , solve for . \rightarrow Mass

Glueball BSE



Need gluon self-energy and gluon propagator , solve for vertex . \rightarrow Mass
 Not quite...

Glueball BSE



Gluons couple to ghosts \rightarrow Include 'ghostball'-part.

(First step: no quarks
 \rightarrow Yang-Mills theory)

Glueball BSE



Gluons couple to ghosts \rightarrow Include 'ghostball'-part. (First step: no quarks
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Need ,  and $4 \times$ , solve for  and . \rightarrow Mass

Construction of kernel

Consistency with input: Apply same construction principle.

Glueball BSE



Gluons couple to ghosts \rightarrow Include 'ghostball'-part. (First step: no quarks
 \rightarrow Yang-Mills theory)

Need triple gluon , ghost-gluon and $4 \times \text{ghost loop}$, solve for three-gluon and three-ghost . \rightarrow Mass

Construction of kernel

Consistency with input: Apply same construction principle.

Previous BSE calculations for glueballs:

- ▶ [Meyers, Swanson '13]
- ▶ [Sanchis-Alepuz, Fischer, Kellermann, von Smekal '15]
- ▶ [Souza et al. '20]
- ▶ [Kaptari, Kämpfer '20]

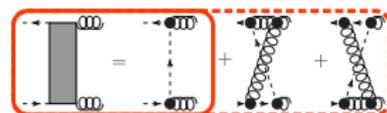
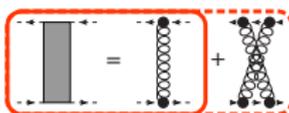
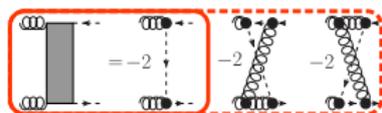
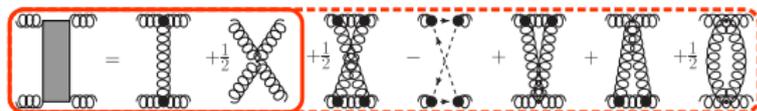
\Rightarrow **Input is important** for
 quantitative predictive
 power!

[MQH, Fischer, Sanchis-Alepuz '20]

Kernel construction

From 3PI effective action truncated to three-loops:

[Fukuda '87; McKay, Munczek '89; Sanchis-Alepuz, Williams '15; MQH, Fischer, Sanchis-Alepuz '20]

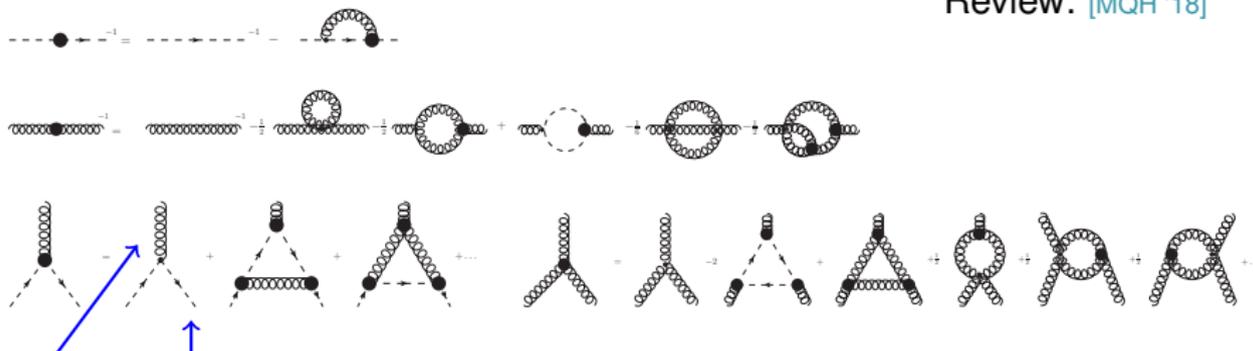


→ Need , , , .

- Some diagrams vanish for certain quantum numbers.
- Full QCD: Same for quarks → Mixing with mesons.

Equations of motion from 3-loop 3PI effective action

Review: [MQH '18]

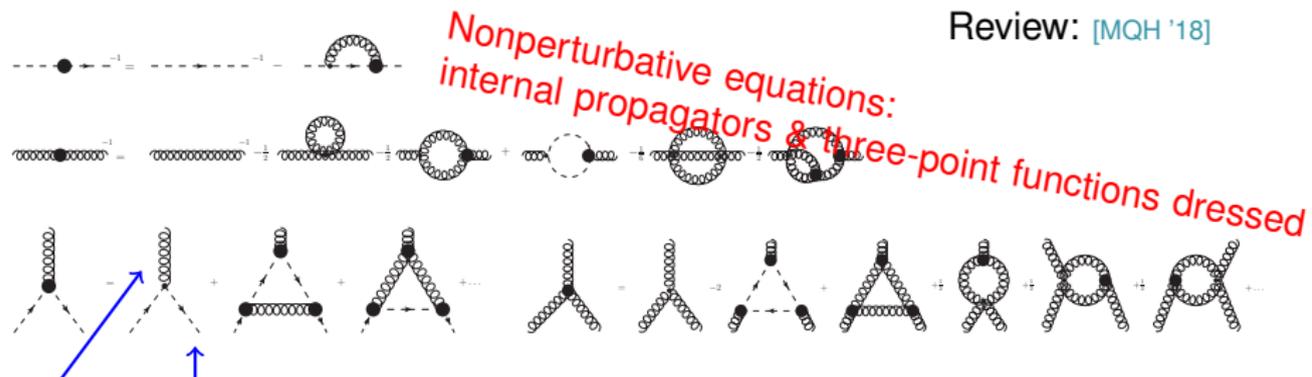


Gluon and ghost fields: Elementary fields of Yang-Mills theory in the Landau gauge

Self-contained system of equations with the scale as the only input.

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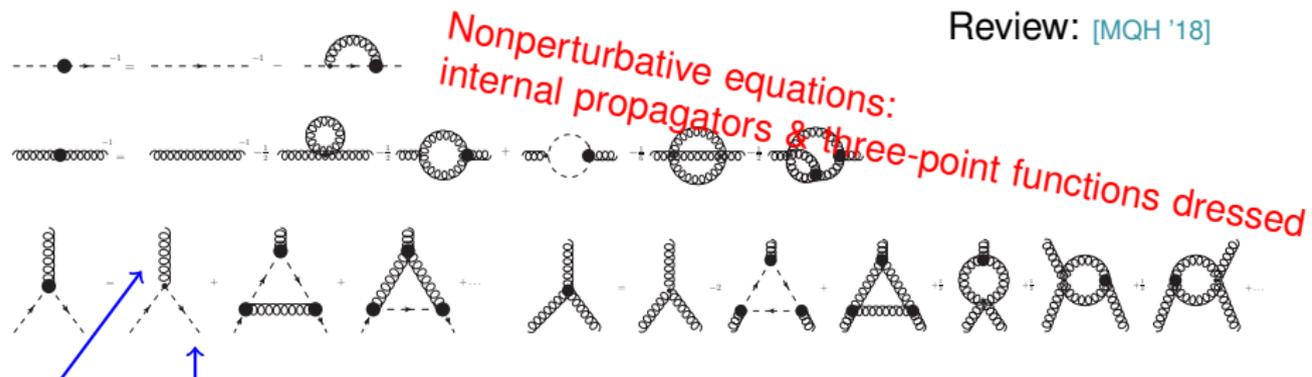


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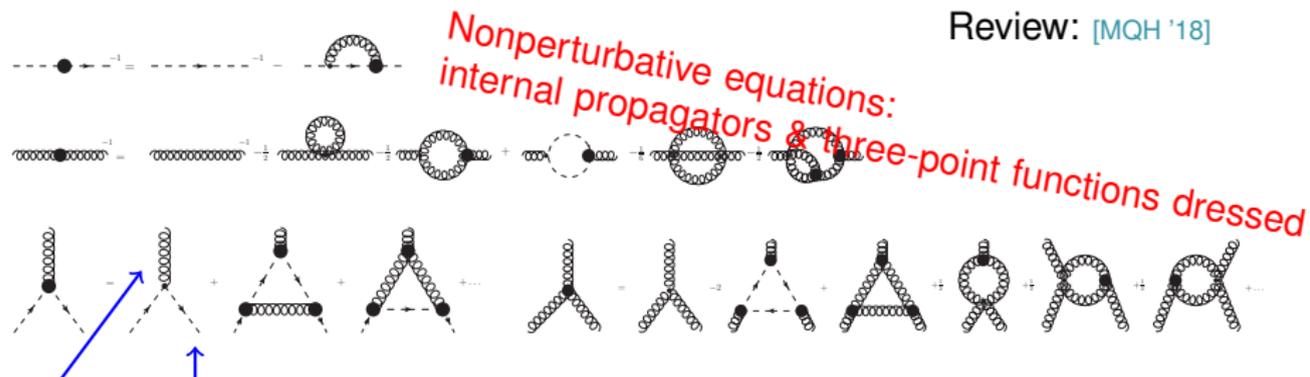
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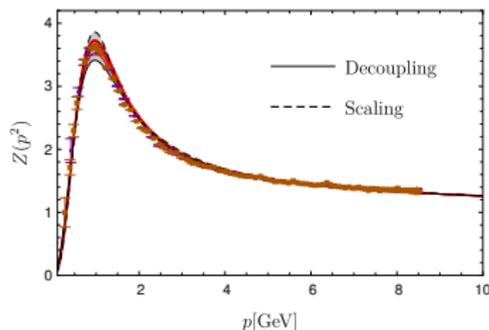
Self-contained system of equations with the scale as the only input.

Truncation \rightarrow 3-loop expansion of 3PI effective action [Berges '04]

- 4 coupled integral equations with full kinematic dependence.
- Sufficient numerical accuracy required for renormalization.
- One- and two-loop diagrams [Meyers, Swanson '14; MQH '17; Eichmann, Pawłowski, Silva '21].

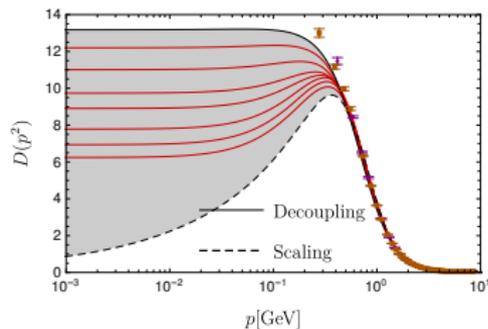
Landau gauge propagators

Gluon dressing function:

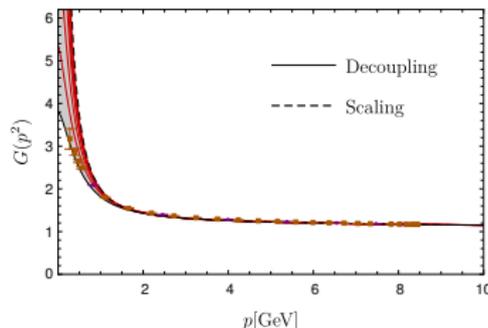


- Family of solutions:
Nonperturbative completions of Landau gauge [Maas '10]?
- Realized by condition on $G(0)$
[Fischer, Maas, Pawłowski '08; Alkofer, MQH, Schwenzer '08]
- Results here independent of $G(0)$

Gluon propagator:



Ghost dressing function:



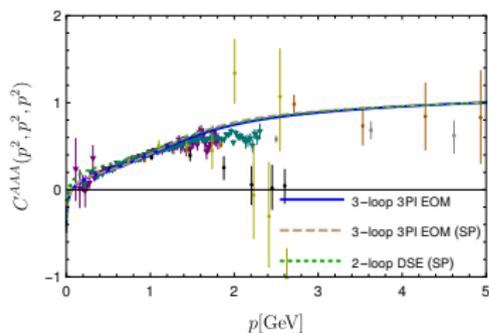
[Sternbeck '06; MQH '20]

Concurrence of functional methods

→ See also talks by Papavassiliou, Rodriguez-Quintero

Exemplified with three-gluon vertex.

3PI vs. 2-loop DSE:



[Cucchieri, Maas, Mendes '08; Sternbeck et al. '17;

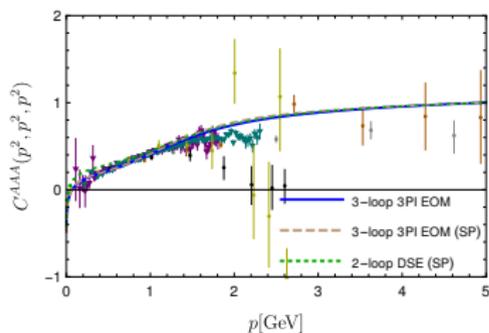
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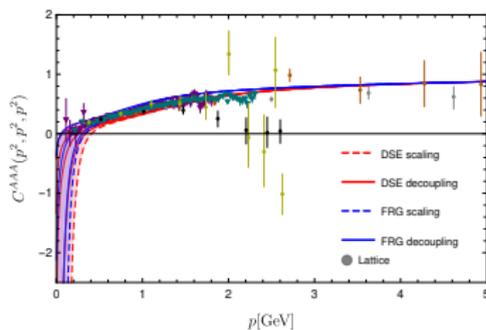
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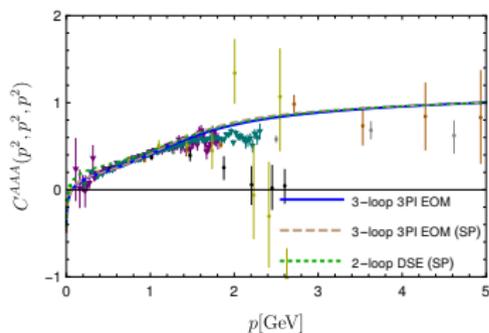
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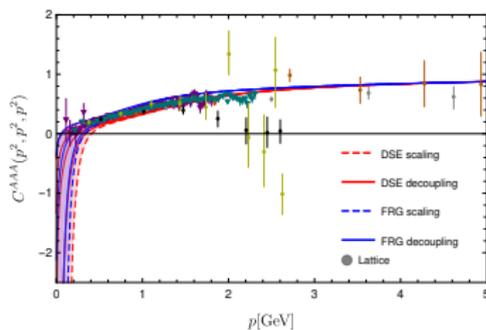
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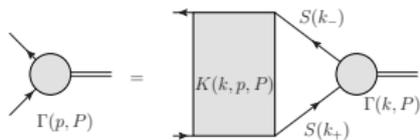


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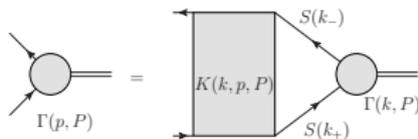
Beyond this truncation

- Further dressings of three-gluon vertex [Eichmann, Williams, Alkofer, Vujanovic '14]
- Effects of four-point functions [MQH '16, MQH '17, Corell et al. '18, MQH '18]

Solving a BSE



Solving a BSE

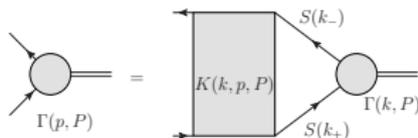


Consider the eigenvalue problem (Γ is the BSE amplitude)

$$\mathcal{K} \cdot \Gamma(P) = \lambda(P) \Gamma(P).$$

$\lambda(P^2) = 1$ is a solution to the BSE \Rightarrow Glueball mass $P^2 = -M^2$

Solving a BSE



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$\lambda(P^2) = 1$ is a solution to the BSE \Rightarrow Glueball mass $P^2 = -M^2$

Calculation requires quantities for

$$k_{\pm}^2 = P^2 + k^2 \pm 2\sqrt{P^2 k^2} \cos \theta = -M^2 + k^2 \pm 2iM\sqrt{k^2} \cos \theta.$$

\Rightarrow **Complex momentum** arguments.

Landau gauge propagators in the complex plane

Propagators for complex momenta

- **Reconstruction** from Euclidean results: mathematically ill-defined, bias in solution
- **Direct calculation** from functional methods possible, e.g., contour deformation or spectral DSEs [Horak, Pawłowski, Wink '20; Horak, Papavassiliou, Pawłowski, Wink '21]
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Contour deformation: Special technique to respect analyticity (avoid branch cuts in the integrand)

- ▶ QED3 [Maris '95 (QED)]
- ▶ Quark propagator [Alkofer, Fischer, Detmold, Maris '04]
- ▶ Self-consistent solution: **Ray technique**, YM propagators [Strauss, Fischer, Kellermann '12; Fischer, MQH '20]
- ▶ Glueball correlators [Windisch, Alkofer, Haase, Liebmann '13; Windisch, MQH, Alkofer '13]
- ▶ Meson decays [Weil, Eichmann, Fischer, Williams '17; Williams '18]
- ▶ Spectral functions at $T > 0$ [Pawłowski, Strodthoff, Wink '18]
- ▶ Quark-photon vertex [Miramontes, Sanchis-Alepuz '19]
- ▶ Scalar scattering amplitude [Eichmann, Duarte, Pena, Stadler '19]
- ▶ **Talk by Eichmann**

Landau gauge propagators in the complex plane

Simpler truncation:

$$\text{---} \bullet \text{---}^{-1} = \text{---} \text{---}^{-1} - \frac{1}{2} \text{---} \circ \text{---} + \text{---} \circ \text{---}$$

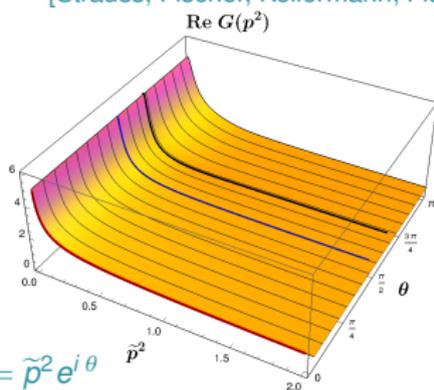
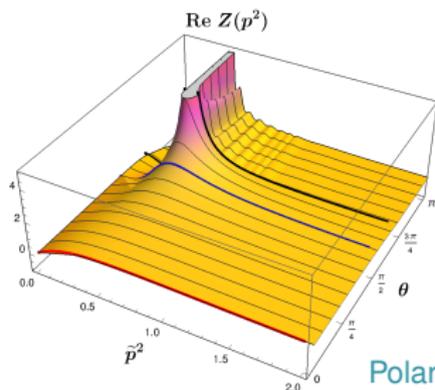
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Ray technique for self-consistent solution of a DSE:

[Strauss, Fischer, Kellermann; Fischer, MQH '20].



Polar coordinates: $p^2 = \tilde{p}^2 e^{i\theta}$

- Current truncation leads to a pole-like structure in the gluon propagator.
- Analyticity up to 'pole' confirmed by various tests (Cauchy-Riemann, Schlessinger, reconstruction)

[Fischer, MQH '20]

Extrapolation of $\lambda(P^2)$

Extrapolation method

- Extrapolation to time-like P^2 using Schlessinger's continued fraction method (proven superior to default Padé approximants) [\[Schlessinger '68\]](#)
- Average over extrapolations using subsets of points for error estimate

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$$f(x) = \frac{f(x_1)}{1 + \frac{a_1(x-x_1)}{1 + \frac{a_2(x-x_2)}{1 + \frac{a_3(x-x_3)}{\dots}}}}$$

Coefficients a_i can
determined such that
 $f(x)$ exact at x_i .

Extrapolation of $\lambda(P^2)$

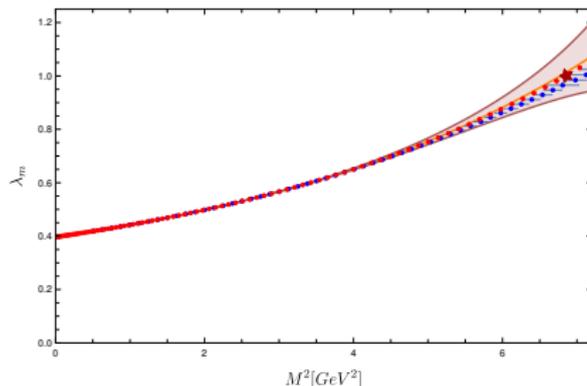
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Test extrapolation for solvable system:
Heavy meson [MQH, Sanchis-Alepuz, Fischer '20]

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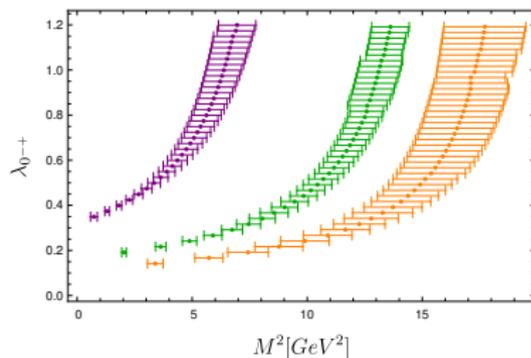
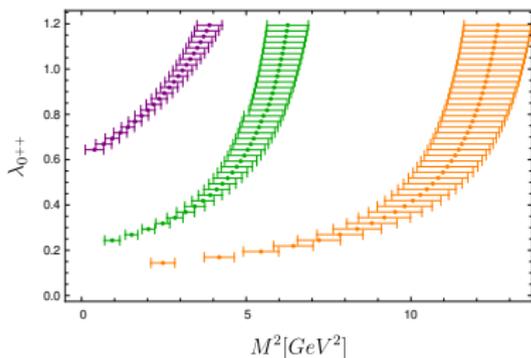


Extrapolation of $\lambda(P^2)$ for glueballs

Higher eigenvalues: Excited states.

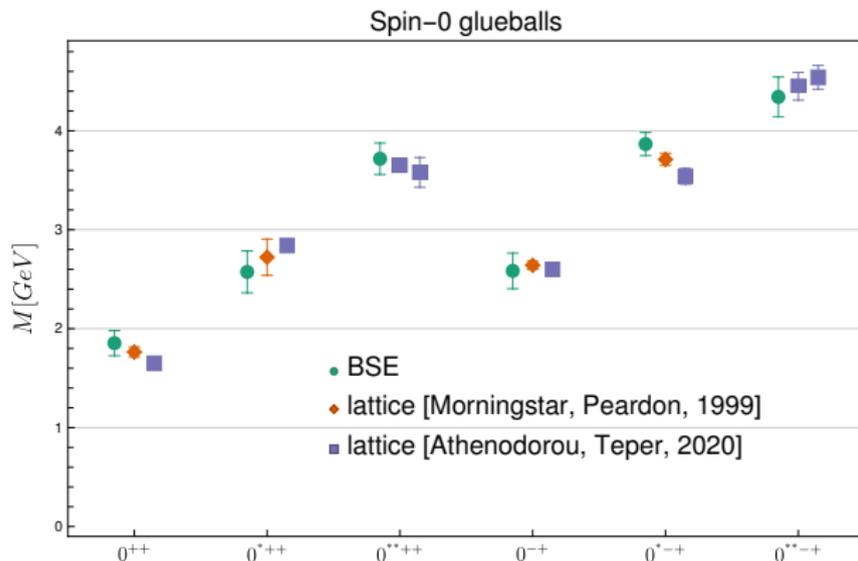
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Higher eigenvalues: Excited states.



Physical solutions for $\lambda(P^2) = 1$.

Glueballs masses for $0^{\pm+}$

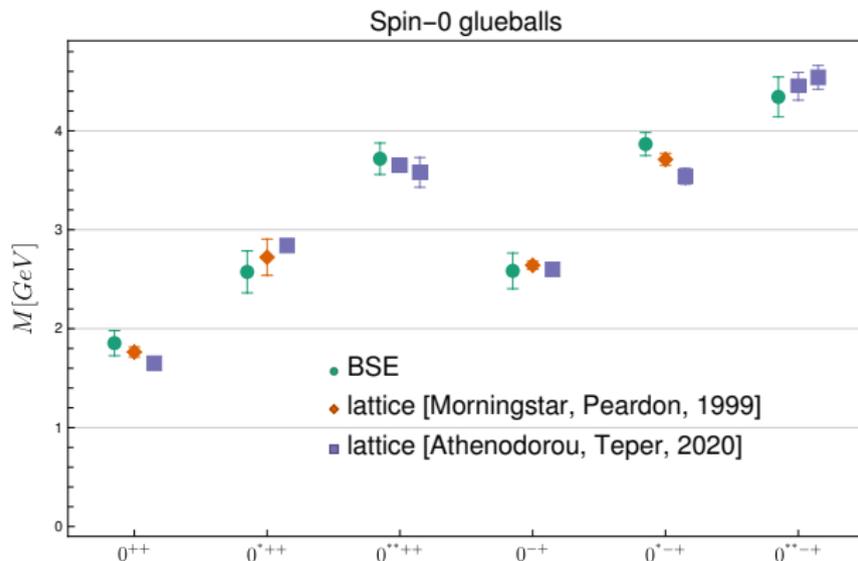


Lattice 0^{**++} :
 Conjectured based on
 irred. rep. of octahedral
 group

All results for $r_0 = 1/418(5)$ MeV.

[MQH, Fischer, Sanchis-Alepuz '20]

Glueballs masses for $0^{\pm+}$



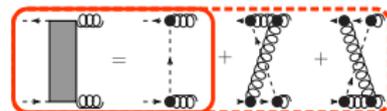
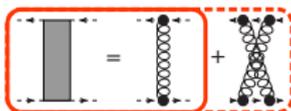
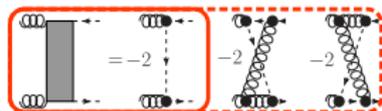
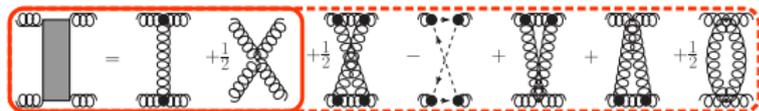
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Two-loop diagrams

Results from [MQH, Fischer, Sanchis-Alepuz '20] were from one-loop terms only:

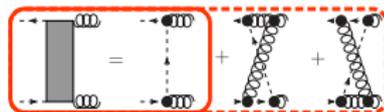
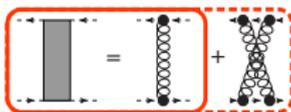
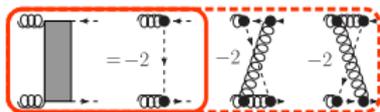
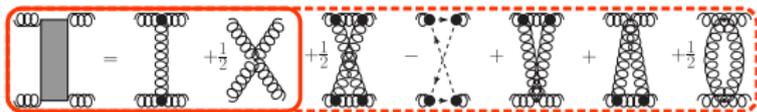


For a fully **self-consistent DSE/BSE truncation**, the two-loop terms are necessary.

→ full system of equations from 3-loop truncated 3PI effective action

Two-loop diagrams

Results from [MQH, Fischer, Sanchis-Alepuz '20] were from one-loop terms only:



For a fully **self-consistent DSE/BSE truncation**, the two-loop terms are necessary.

→ full system of equations from 3-loop truncated 3PI effective action

Drastic increase in computational resources, hence lower precision used.

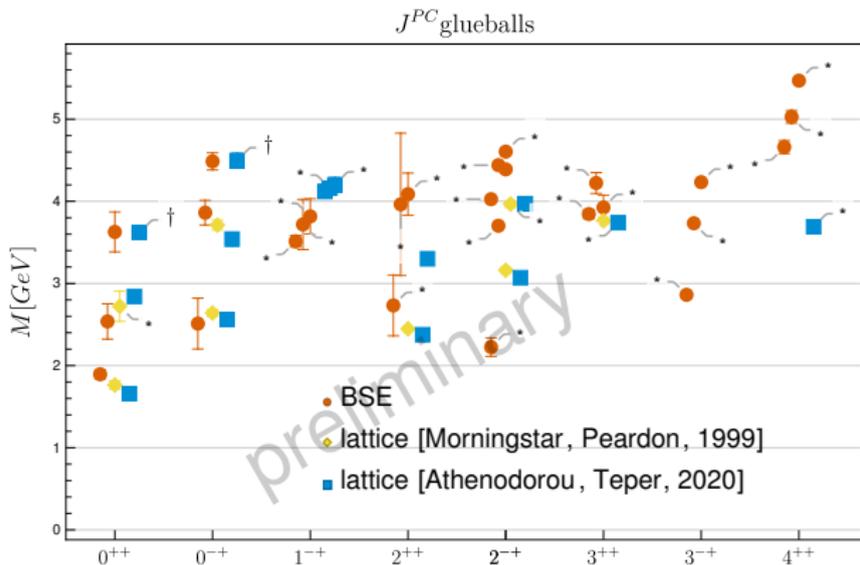
Preliminary result: **No effect on mass.**

Glueball masses for $J^{\pm+}$

For higher spin, larger tensor bases: more tensors, more indices

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Lattice:

*: identification with some uncertainty

†: conjecture based on irred. rep of octahedral group

[MQH, Fischer, Sanchis-Alepuz, in preparation]

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Parameter-free determination of glueball masses from functional methods.

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Thank you for your attention.

Glueballs as bound states

Hadron masses from correlation functions of **color singlet operators**.

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Example: For $J^{PC} = 0^{++}$ glueball take $O(x) = F_{\mu\nu}(x)F^{\mu\nu}(x)$:

$$D(x - y) = \langle O(x)O(y) \rangle$$

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Hadron masses from correlation functions of **color singlet operators**.

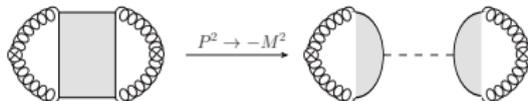
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- Complicated object in a diagrammatic language: 2-, 3- and 4-gluon contributions

Put total momentum on-shell and consider individual 2-, 3- and 4-gluon contributions. → Each can have a pole at the glueball mass.

A^4 -part of $D(x - y)$, total momentum on-shell:



Charge parity

Transformation of gluon field under charge conjugation:

$$A_{\mu}^a \rightarrow -\eta(a)A_{\mu}^a$$

where

$$\eta(a) = \begin{cases} +1 & a = 1, 3, 4, 6, 8 \\ -1 & a = 2, 5, 7 \end{cases}$$

Color neutral operator with two gluon fields:

$$A_{\mu}^a A_{\nu}^a \rightarrow \eta(a)^2 A_{\mu}^a A_{\nu}^a = A_{\mu}^a A_{\nu}^a.$$

$$\Rightarrow C = +1$$

Negative charge parity, e.g.:

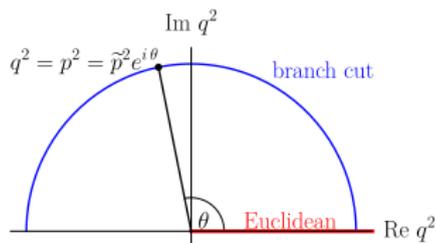
$$\begin{aligned} d^{abc} A_{\mu}^a A_{\nu}^b A_{\rho}^c &\rightarrow -d^{abc} \eta(a)\eta(b)\eta(c) A_{\mu}^a A_{\nu}^b A_{\rho}^c = \\ &= -d^{abc} A_{\mu}^a A_{\nu}^b A_{\rho}^c. \end{aligned}$$

Only nonvanishing elements of the symmetric structure constant d^{abc} : zero or two indices equal to 2, 5 or 7.

Landau gauge propagators in the complex plane

Simpler truncation:

$$\text{gluon propagator}^{-1} = \text{gluon propagator}^{-1} - \frac{1}{2} \text{ghost loop} + \text{ghost loop}$$

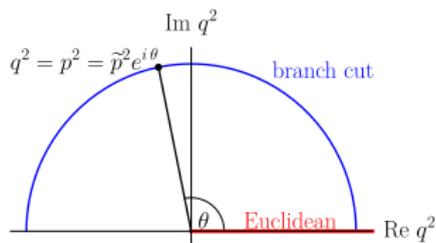


→ Opening at $q^2 = p^2$.

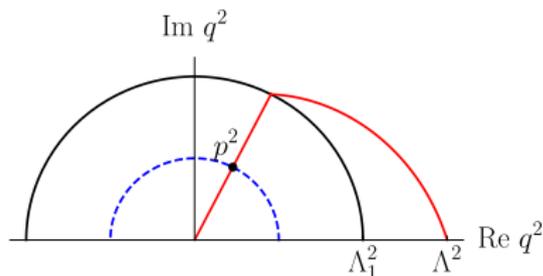
Landau gauge propagators in the complex plane

Simpler truncation:

$$\text{wavy line with dot}^{-1} = \text{wavy line}^{-1} - \frac{1}{2} \text{wavy line with loop} + \text{wavy line with dashed loop}$$



→ Opening at $q^2 = p^2$.



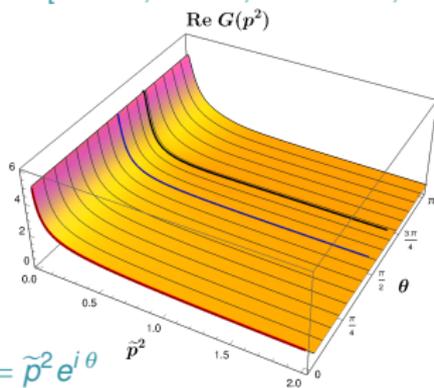
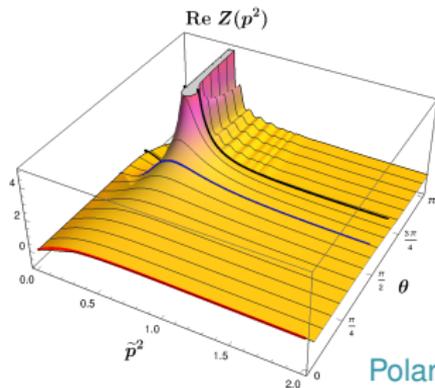
Appearance of branch cuts for complex momenta forbids integration directly to cutoff.

Deformation of integration contour necessary [Maris '95]. Recent resurgence: [Alkofer et al. '04; Windisch, MQH, Alkofer, '13; Williams '19; Miramontes, Sanchis-Alepuz '19; Eichmann et al. '19, ...]

Landau gauge propagators in the complex plane

Ray technique for self-consistent solution of a DSE:

[Strauss, Fischer, Kellermann; Fischer, MQH '20].



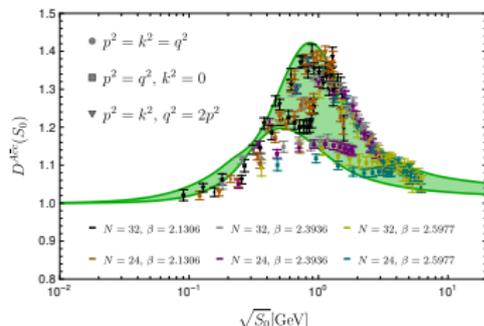
Polar coordinates: $p^2 = \tilde{p}^2 e^{i\theta}$

- Current truncation leads to a pole-like structure in the gluon propagator.
- Analyticity up to 'pole' confirmed by various tests (Cauchy-Riemann, Schlessinger, reconstruction)
- No proof of existence of complex conjugate poles due to simple truncation.

[Fischer, MQH '20]

Landau gauge vertices

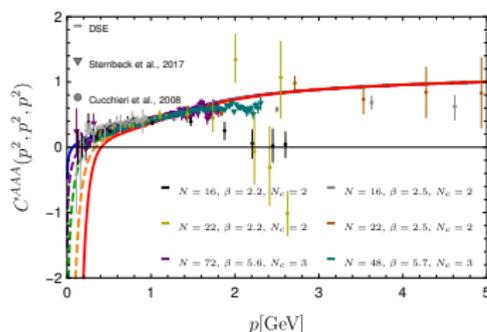
Ghost-gluon vertex:



[Maas '19; MQH '20]

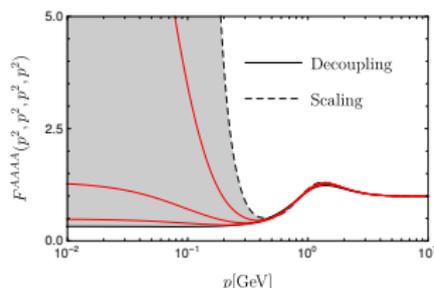
- Nontrivial kinematic dependence of ghost-gluon vertex
- Simple kinematic dependence of three-gluon vertex
- Four-gluon vertex from solution

Three-gluon vertex:



[Cucchieri, Maas, Mendes '08; Sternbeck et al. '17; MQH '20]

Four-gluon vertex:

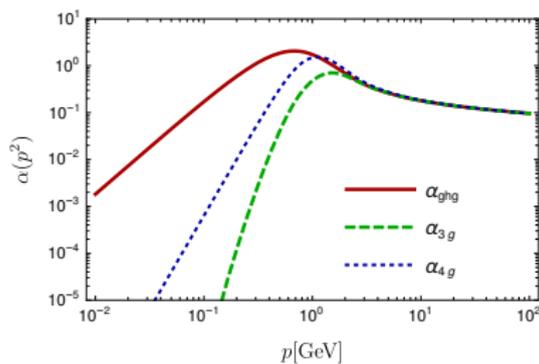


[MQH '20]

Some properties of the Landau gauge solution

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- Slavnov-Taylor identities (gauge invariance): Vertex couplings agree down to GeV regime



Some properties of the Landau gauge solution

[MQH '20]

- Slavnov-Taylor identities (gauge invariance): Vertex couplings agree down to GeV regime
- Renormalization: First parameter-free subtraction of quadratic divergences
 \Rightarrow **One unique free parameter** (family of solutions)

