

Charm 2021

X(3872) Production and Suppression

(in pp collisions)

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General overview

• Brief review of *X*(3872)

• Production of X(3872) at hadron colliders

• Suppression of X(3872)

Brief Review of X(3872)

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Discovered at e^+e^- collider in B^+ \to K^+X, X \to J/\psi \pi^+\pi^- [Belle (2003)]
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Confirmed at $p\bar{p}$ collider [CDF (2003)]

Observed at pp collider [LHCb (2011), CMS (2011), ATLAS (2016)]

Most precise mass and first width [LHCb (2020)]:

$$M_X = 3871.695 \pm 0.096 \text{ MeV}, \ \Gamma_X^{BW} = 1.19 \pm 0.19 \text{ MeV}$$

7 observed decay channels

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J/\psi \pi^+ \pi^- [Belle (2003)] J/\psi \pi^+ \pi^- \pi^0 [BaBar (2010)] J/\psi \gamma [BaBar (2006)] \psi(2S)\gamma [BaBar (2009)]
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$$D^0 \overline{D}{}^0 \pi^0$$
 [Belle (2006)] $D^0 \overline{D}{}^0 \gamma$ [Belle (2010)] $\chi_{c1} \pi^0$ [BESIII (2019)]

Brief Review of X(3872)

Decays into $J/\psi\pi^+\pi^- = J/\psi\rho$ and $J/\psi\pi^+\pi^-\pi^0 = J/\psi\omega$ indicate severe isospin violation

Tiny binding energy [LHCb (2020)]:

$$E_X = M_X - (M_{D^{*0}} + M_{\overline{D}^0}) = -0.07 \pm 0.12 \text{ MeV}$$

Quantum numbers: $J^{PC} = 1^{++}$ [LHCb (2013)]

Imply X(3872) is S-wave loosely-bound charm-meson molecule

$$X = \frac{1}{\sqrt{2}} (|D^{*0}\overline{D}^{0}\rangle + |\overline{D}^{*0}D^{0}\rangle)$$

Universal properties determined by binding energy E_X (or scattering length $a_X = 1/\gamma_X$) [Braaten, Kusunoki (2003)]

$$|E_X| < 0.22 \text{ MeV}$$
 at 90% C.L.

$$\gamma_X = \sqrt{2\mu_{D^*\overline{D}}|E_X|} < 21 \text{ MeV}$$

Huge mean separation:

Wavefunction:

$$\langle r \rangle_X = \frac{1}{2v_X} > 4.5 \text{ fm}$$

$$\psi_{X(r)} = \frac{1}{\sqrt{8\pi\gamma_X}} \, \frac{e^{-\gamma_X r}}{r}$$

Brief Review of X(3872)

Other possibilities for *X*:

- Hybrid: combination of heavy quarks and a constituent gluon
- <u>Hadroquarkonia</u>: heavy charmonium core $c\bar{c}$ surrounded by light meson $q\bar{q}$ bound by QCD analog of van der Waals force
- Compact tetraquark: diquark and anti-diquark bound by color interactions
- Charmonium: $\chi_{c1}(2P)$
- Cusp: discontinuity in differential cross section across threshold

Regardless, the coupling of X to $D^{*0}\overline{D}{}^{0}$ transforms it into a large charm-meson molecule

Much past research has tried to use decay-channel predictions to determine the nature of X but without much success

More recently production and suppression mechanisms have been used to help determine X nature

Two contributions to inclusive production

- Prompt production at primary collision vertex
- Production by b hadron decay at secondary vertex

Convenient to benchmark X(3872) against $\psi(2S) = \psi(3686)$

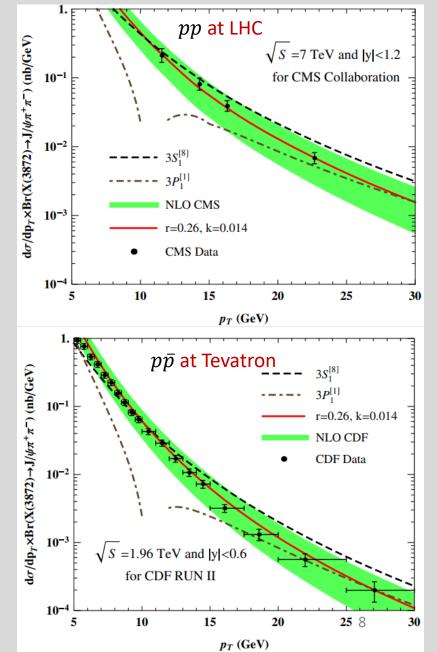
- Both are observed in $J/\psi \pi^+\pi^-$ channel
- They have similar masses

Cross section for producing X related to cross section for creating $c\bar{c}$ at short distances through Long Distance Matrix Elements (LDMEs)

[Meng, Han, Chao (2017)] calculate p_T -distribution assuming production of X dominated by $\chi_{c1}(2P)$ component

LDMEs at NLO in NRQCD:

- $\widehat{O}^{\chi'_{c_1}}\left(3S_1^{[8]}\right)$: from fits
- $\widehat{O}^{\chi'_{c1}}\left(3P_1^{[1]}\right)$: related to $\chi_{c1}(2P)$ wavefunction at origin
- Normalization factor $k = Z_{c\bar{c}} \mathrm{Br}(X \to J/\psi \pi^+ \pi^-)$, where $Z_{c\bar{c}}$ is probability $|\langle \chi'_{c1} | X \rangle|^2$



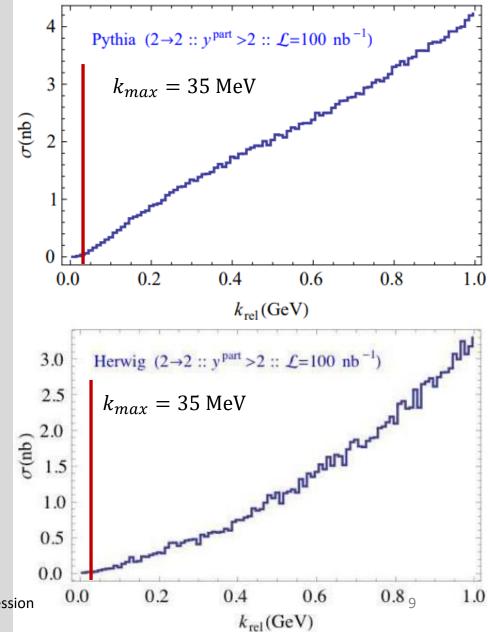
[Bignamini *et al.* (2009)] If X is a loosely bound charm-meson molecule its cross section can be approximated by the cross section for producing $D^{*0}\overline{D}^{0}$ with relative momentum k less than some k_{max}

$$\sigma[X] = \sigma[D^{*0}\overline{D}^{0}(k < k_{max})]$$

Assume that
$$k_{max} \approx \gamma_X$$
, $\gamma_X = \sqrt{2\mu_{D^{*0}\overline{D}^0}|E_X|}$

Calculate $\sigma[D^{*0}\overline{D}^{0}]$ using Pythia and Herwig event generator

Calculated cross section at Tevatron is orders of magnitudes smaller than observed cross section



$$\sigma[X] = \sigma[D^{*0}\overline{D}^{0}(k < k_{max})]$$

[Artoisenet, Braaten (2010)]

- Assume that $k_{max} \approx m_{\pi}$
- Calculated $\sigma[D^{*0}\overline{D}{}^{0}]$ using PYTHIA event generator

[Albaladejo et al. (2017)]

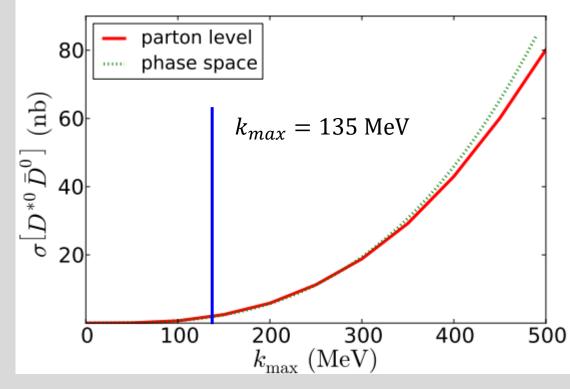
• Used deuteron to show that $k_{max} \sim m_{\pi}$, similar should be true for X(3872)

[Braaten, He, Ingles (2019)]

• Quantitative estimate on $k_{max} = 7.7\gamma_X$, where

$$\gamma_X = \sqrt{2\mu_{D^{*0}\bar{D}^0}} |E_X|$$
Note: $7.7^3 \approx 500$

Cross section of X(3872) as a molecule reproduces order of magnitude of experiments



LHCb:
$$2.5 < y < 4.5$$
, $5 < p_T < 20 \text{ GeV/c}$
 $\sigma[X]\text{Br}[X \to J\psi\pi^+\pi^-] = 5.4 \pm 1.3 \pm 0.8 \text{ nb}$

CMS:
$$y < |1.2|$$
, $10 < p_T < 20 \text{ GeV/c}$
 $\sigma[X] \text{Br}[X \to J\psi \pi^+ \pi^-] = 1.06 \pm 0.11 \pm 0.15 \text{ nb}$

Production of X can come from creation of $\overline{D}{}^0D^{*0}$, $D^0\overline{D}{}^{*0}$ at short distances

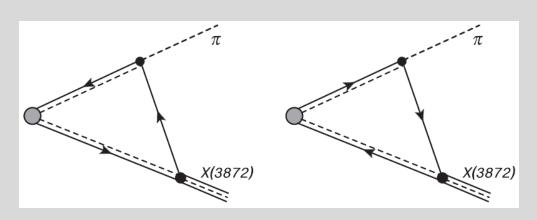
Production of $X\pi^+$ with soft π can come from creation of $D^{*+}\overline{D}^{*0}$ at short distances

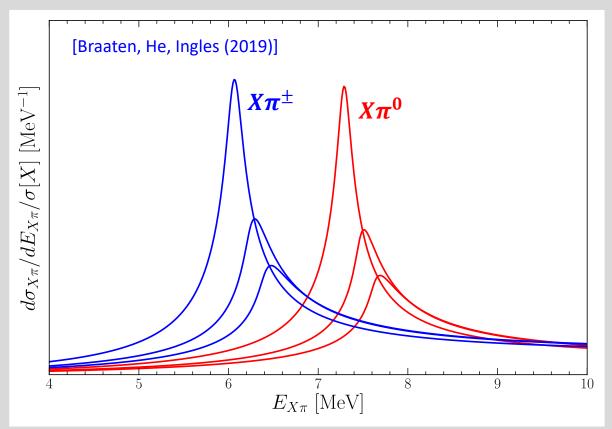
• Triangle singularity in process $D^{*+}\overline{D}^{*0} \to X\pi^{+}$ gives peak about 6 MeV above $X\pi^{+}$

threshold with width < 1 MeV

Charm-meson triangle loop ⇒ triangle singularity

Decay width of D^* and binding energy of X reduce \log^2 -divergence to narrow peak





Suppression of X(3872)

Suppression of X(3872)

Proton-proton collision:

Interactions with comoving gluons (or pions)
 Comover Interaction Model (CIM)

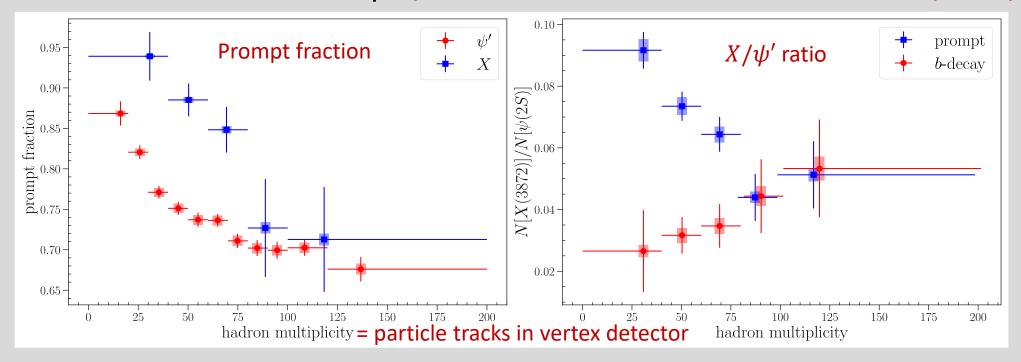
Proton-nucleus collision:

- Interactions with comoving gluons (or pions)
- Cold nuclear matter effects: PDFs of p and n, nuclear shadowing, absorption by nucleons etc.

Nucleus-Nucleus collision:

- Interactions with comoving gluons (or pions)
- Cold nuclear matter effects
- Thermal effects in quark-gluon plasma
- Thermal effects in expanding, cooling hadron gas

[LHCb (2021)] measured X and ψ' yields as functions of hadron multiplicity



Prompt fractions for X and ψ' decrease with multiplicity Prompt fraction for ψ' seems to saturates at large multiplicity

Survival probability in Comover Interaction Model [Armesto, Capella (1998)]

$$S = \exp\left[-\frac{\langle v\sigma \rangle}{\sigma_0} \frac{dN}{dy} \log\left(\frac{1}{N_0} \frac{dN}{dy}\right)\right]$$

 N_0 : multiplicity at which interactions stop

 σ_0 : parameter that depends center-of-mass energy

Model for breakup reaction rate and momentum distribution for comovers

[Ferreiro, Lansberg (2018)]

$$\langle v\sigma \rangle = \pi \langle r^2 \rangle;$$

 r^2 : separation squared of constituents

 E^{thr} : energy required to break X apart

E: gluon (or pion) relativistic energy

$$f(E_{co}) = \left(e^{E/T_{eff}} - 1\right)^{-1}$$

$$T_{eff} = (250 \pm 50) \text{ MeV}$$

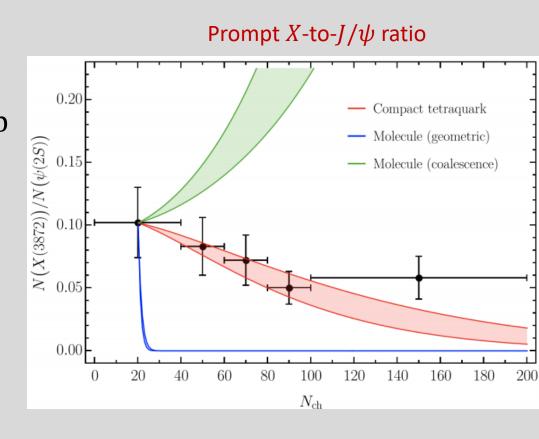
[Esposito, et al. (2020)] estimated X/ψ' ratio assuming CIM and

- X as a tightly bound tetraquark $\langle v\sigma \rangle \sim 10 \text{ mb}$
- X as charm-meson molecule $\langle v\sigma \rangle \sim 1200 \text{ mb}$
- X as charm-meson molecule and with processes $\pi \overline{D} D^* \to X \pi$

Glauber Monte Carlo modeling

- Generate distribution of comovers' momentum and spacetime distributions
- X can only interact with comovers within range r given by geometric cross section $\pi\langle r^2\rangle$

Authors conclude that CIM favors tetraquark interpretation

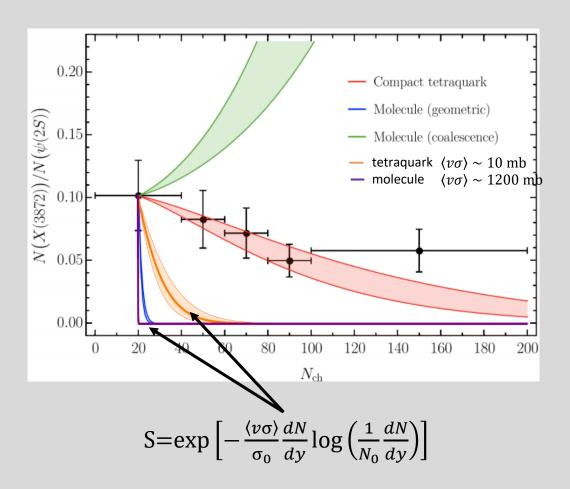


Problem with few-body physics from Esposito *et al.* (2020): their breakup reaction rate given by geometric cross section $\pi \langle r^2 \rangle$ where $\langle r^2 \rangle \propto E_X^{-1}$

But breakup reaction rate in the high energy limit should be given by cross section for scattering with constituents $\sim (m_{\pi}/f_{\pi}^2)^2$

Survival probability from CIM gives exponential suppression as function of multiplicity

LHCb data indicates a saturation



From LHCb data, prompt fraction for ψ' seems to saturates at large multiplicity

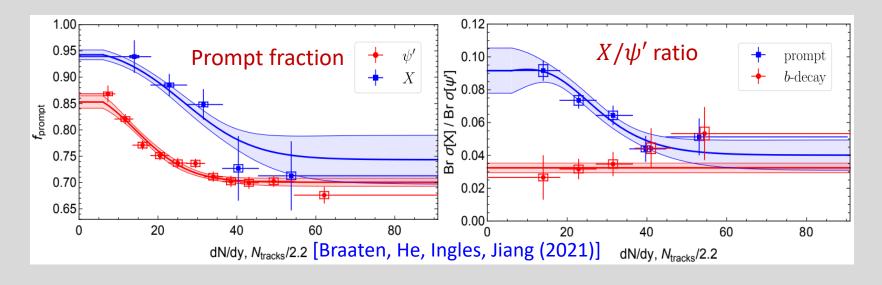
Assumptions of arXiv:2012.13499:

- 1. prompt cross section is sum of
 - term with survival probability $S = \exp\left[-\frac{\langle v\sigma\rangle(dN/dy)}{\sigma_0}\log\left(\frac{dN/dy}{N_0}\right)\right]$
 - term with survival probability 1
- 2. b-decay cross section independent of dN/dy

26 data points 7 fitting parameters $\chi^2/\text{dof} = 0.99$ Fitted reaction rates

4/20/2021

$$\langle v\sigma \rangle_{\psi}$$
, = 3.9 ± 0.8 mb
 $\langle v\sigma \rangle_X$ = 2.6 ± 0.7 mb



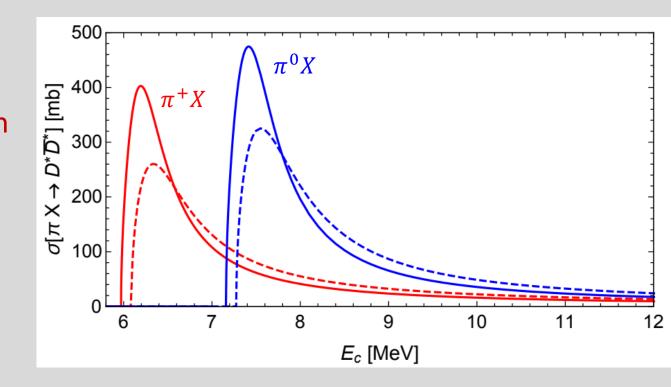
Cross section for πX scattering in XEFT E_c - comover energy relative to πX threshold

Narrow peak in $\pi^+ X$ 6 MeV above threshold Narrow peak in $\pi^0 X$ 8 MeV above threshold

Order of magnitude at peak is geometric mean of geometric cross section $\pi \langle r^2 \rangle$ and cross section for scattering with constituents $(m_\pi/f_\pi^2)^2$

Cross section for πX scattering decreases to cross section for scattering with constituents $\pi D^{(*)}$ at higher energies

$$\sigma(\pi X) \sim \left(\frac{m_{\pi}}{f_{\pi}^2}\right)^2 = 25 \text{ mb}$$



Summary

Still no consensus on nature of X(3872)

Studying production and suppression mechanisms will help reveal X(3872) nature

Important to use correct few-body physics for X(3872)

Fitted results of multiplicity dependence from LHCb suggest breakup reaction rate for X(3872) from collisions with comoving pions should be of order 3 mb

