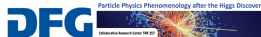


# Revisiting the Charm-meson lifetimes

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10th International Workshop on Charm Physics (online),  
31 May - 04 June 2021

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# Lifetimes and ratios: experiment

- Lifetimes of **charmed** hadrons are measured precisely in **experiments**  
[PDG-2021], [BESIII, 2104.07311]\*

	$D^0$	$D^+$	$D_s^+$
$\tau$ [ps]	0.4101(15)	1.040(7)	0.504(4)
$\Gamma_{\text{tot}}$ [ps $^{-1}$ ]	2.44(1)	0.96(1)	1.98(2)
$\frac{\tau(D_X)}{\tau(D^0)}$	1	2.54(2)	1.20(1)
$\text{Br}(D_X \rightarrow X e^+ \nu_e)$ [%]	6.49(11)	16.07(30)	6.30(16)*
$\frac{\Gamma(D_X \rightarrow X e^+ \nu_e)}{\Gamma(D^0 \rightarrow X e^+ \nu_e)}$	1	0.977(26)	0.790(26)*

# Lifetimes and ratios: theory

- The total width of a hadron  $H$  is given by

$$\Gamma(H) = \frac{1}{2m_H} \sum_X \int_{\text{PS}} (2\pi)^4 \delta^{(4)}(p_H - p_X) |\langle X(p_X) | \mathcal{H}_{\text{eff}} | H(p_H) \rangle|^2$$

Optical Theorem

$$= \frac{1}{2m_B} \text{Im} \langle H(p_H) | i \int d^4x T \{ \mathcal{H}_{\text{eff}}(x), \mathcal{H}_{\text{eff}}(0) \} | H(p_H) \rangle$$

- The lifetimes ratio  $\tau(H_1)/\tau(H_2)$

$$\frac{\tau(H_1)}{\tau(H_2)} = \frac{\Gamma_Q + \delta\Gamma_{H_2}}{\Gamma_Q + \delta\Gamma_{H_1}} = 1 + (\delta\Gamma_{H_2} - \delta\Gamma_{H_1}) \tau(H_1)$$

- May be sensitive to **New Physics** contributions

$$\frac{\tau(H_1)}{\tau(H_2)} = \frac{\Gamma_Q + \delta\Gamma_{H_2}}{\Gamma_Q + \delta\Gamma_{H_1}} = 1 + (\delta\Gamma_{H_2}^{\text{SM}} - \delta\Gamma_{H_1}^{\text{SM}}) \tau(H_1) + (\delta\Gamma_{H_2}^{\text{NP}} - \delta\Gamma_{H_1}^{\text{NP}}) \tau(H_1)$$

# HQE: Bottom vs Charm

- HQE:  $p_Q = m_Q v + k, \quad k \sim \Lambda \quad v = p_H/m_H$

- HQE in the **bottom** sector

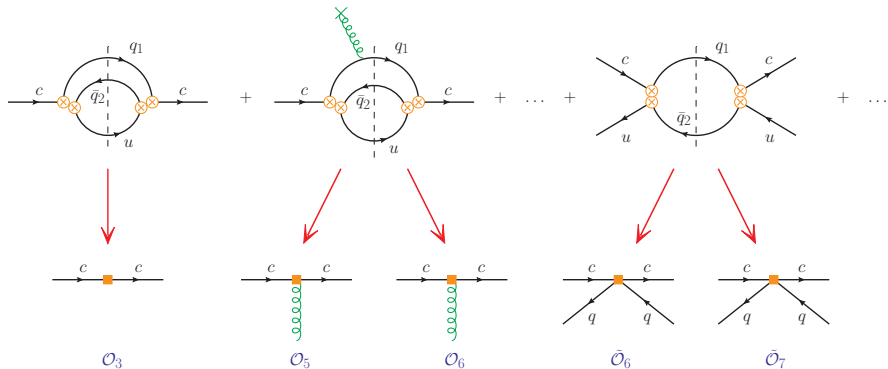
$$\frac{\Lambda}{m_b} \sim 0.12 \quad \alpha_s(m_b) \sim 0.22, \quad \frac{\alpha_s(m_b)}{\pi} \sim 0.07$$

- HQE in the **charm** sector

$$\frac{\Lambda}{m_c} \sim 0.39 \quad \alpha_s(m_c) \sim 0.38, \quad \frac{\alpha_s(m_c)}{\pi} \sim 0.12$$

- Convergence of HQE for charm-hadron lifetimes ?

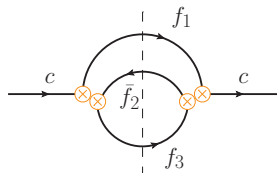
# HQE: diagrams



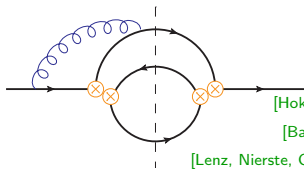
$$\Gamma(D) = \Gamma_3 + \Gamma_5 \frac{\langle \mathcal{O}_5 \rangle}{m_c^2} + \Gamma_6 \frac{\langle \mathcal{O}_6 \rangle}{m_c^3} + \dots + 16\pi^2 \left[ \tilde{\Gamma}_6 \frac{\langle \tilde{\mathcal{O}}_6 \rangle}{m_c^3} + \tilde{\Gamma}_7 \frac{\langle \tilde{\mathcal{O}}_7 \rangle}{m_c^4} + \dots \right]$$

$$\Gamma_i = \Gamma_i^{(0)} + \frac{\alpha_s}{4\pi} \Gamma_i^{(1)} + \dots$$

# Dimension-3 contribution



+



[Hokim, Pham (1984)]

[Bagan et al. (1994)]

[Lenz, Nierste, Ostermaier (1997)]

[Greub, Liniger (2000, 2001)]

[Krunner, Lenz, Rauh (2013)]

- Determine dimension-3 contribution  $\Gamma_3 = \Gamma_0 c_3$

$$c_3 = c_3^{s\bar{d}u} + c_3^{s\bar{s}u} + c_3^{d\bar{d}u} + c_3^{d\bar{s}u} + c_3^{se\bar{\nu}_e} + c_3^{s\mu\bar{\nu}_\mu}$$

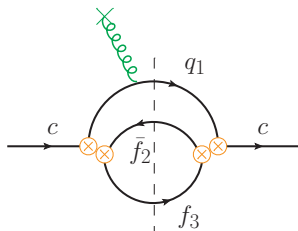
$$\Gamma_0 = \frac{G_F^2 m_c^5}{192\pi^2}$$

- Perturbative expansion in  $\alpha_s$

$$c_3 = c_3^{(0)} + \frac{\alpha_s}{4\pi} c_3^{(1)} + \dots$$

- **Universal** for all heavy hadrons with a c-quark
- There is no non-perturbative parameter at dimension-3

# Dimension-5 two-quark contributions



[Bigi et al. (1992)]

[Blok, Shifman (1993)]

- Determine power corrections for two-quark operators  $\Gamma_5$
- Non-perturbative parameters  $\langle \mathcal{O}_5 \rangle$  arise:
  - ▷ Kinetic  $\mu_\pi^2(D)$  and chromo-magnetic  $\mu_G^2(D)$
- Now dependent on the spectator quark in  $D$ -meson
- $\Gamma_5$  for semi-leptonic modes known at NLO
- $\Gamma_5$  for non-leptonic modes known at LO

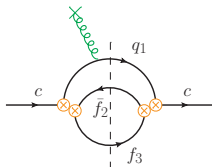
[Becher, Boos, Lunghi (2007), Alberti et al. (2013, 2014),  
Mannel, Pivovarov, Rosenthal (2015)]

[Blok, Shifman (1993)]

# Dimension-6 Darwin operator contribution

- Darwin parameter

$$2m_H \rho_D^3(H) = \langle H(p_H) | \bar{Q}_v(iD_\mu)(iv \cdot D)(iD^\mu)Q_v | H(p_H) \rangle$$



- In the **bottom** decays:  $m_b \sim m_c \gg \Lambda_{\text{QCD}}$

▷ Known for **semi-leptonic** modes upto **NLO**

[Gremm, Kapustin (1996)], [Mannel, Pivovarov (2019)]

▷ Recently determined for **non-leptonic** modes at **LO**

[Lenz, Piscopo, AR (2020)], [Mannel, Moreno, Pivovarov (2020)]

- In the **charm** decays:  $m_c \gg m_s \sim \Lambda_{\text{QCD}}$

▷ One cannot simply adopt results from **bottom** to **charm** decays (!)

[see also talk by Matteo Fael on Wednesday]

▷ Determined for **semi-leptonic** modes at **LO**

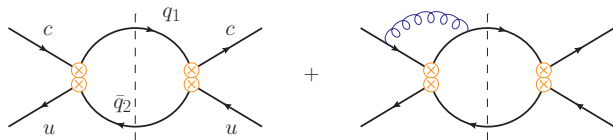
e.g. [Gambino, Kamenik (2010)]

[Fael, Mannel, Vos (2019)]

▷ Results for **non-leptonic** modes at **LO** are **new!**

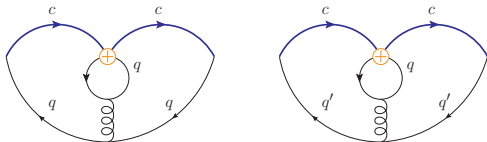


# Four-quark contributions



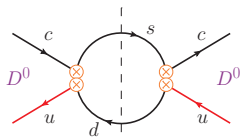
[Uraltsev (1996)]  
 [Neubert, Sachrajda (1996)]  
 [Beneke et al. (2002)]  
 [Franco et al. (2002)]

- Determine  $\tilde{\Gamma}_6, \tilde{\Gamma}_7$  starting at one-loop
- Non-perturbative input  $\langle \tilde{\mathcal{O}}_6 \rangle, \langle \tilde{\mathcal{O}}_7 \rangle$ 
  - ▷ Dim-6: 4 Bag parameters  $B_{1,2}, \epsilon_{1,2}$
  - ▷ Dim-7: 12 Bag parameters  $\rho_i, \sigma_i$  ( $i = 1, \dots, 6$ ) (in HQET)
- **New** determination of dim-6 Bag parameters from **HQET SR** both for  $D^{*+}$ - and  $D_s^{*+}$ -mesons (incl. so-called eye-contractions)

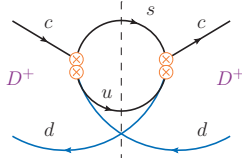


[King, Lenz, Rauh (in progress)]

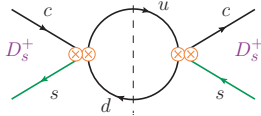
# Dominant four-quark contributions



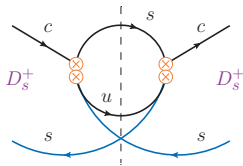
$\sim B_1 - B_2 \approx 0$  (helicity suppression)



Large! (no helicity suppression)



CKM-dominant,  
helicity suppressed



CKM-suppressed,  
no helicity suppression

# Lifetimes and their ratios: theory status

$$\Gamma(D) = \Gamma_3 + \Gamma_5 \frac{\langle \mathcal{O}_5 \rangle}{m_c^2} + \Gamma_6 \frac{\langle \mathcal{O}_6 \rangle}{m_c^3} + \dots + 16\pi^2 \left[ \tilde{\Gamma}_6 \frac{\langle \tilde{\mathcal{O}}_6 \rangle}{m_c^3} + \tilde{\Gamma}_7 \frac{\langle \tilde{\mathcal{O}}_7 \rangle}{m_c^4} + \dots \right]$$

$$\tau(D_{(s)}^+) / \tau(D^0) = 1 + \left[ \Gamma^{\text{HQE}}(D^0) - \Gamma^{\text{HQE}}(D_{(s)}^+) \right] \tau^{\text{exp}}(D_{(s)}^+)$$

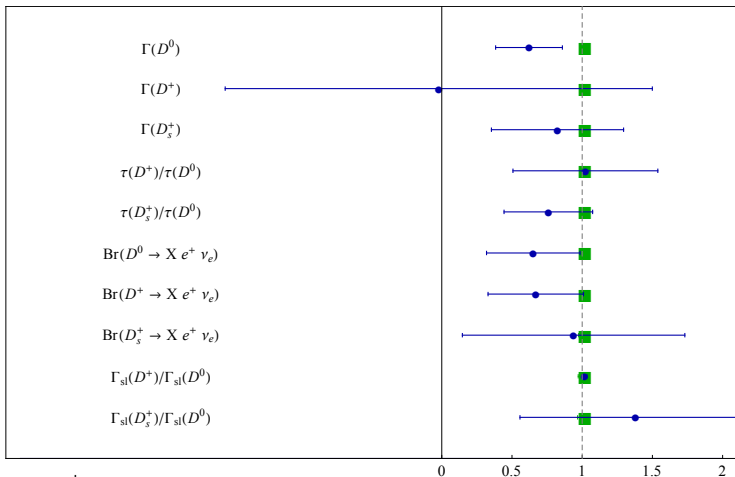
	SL	NL
$\Gamma_3$	NNNLO <sup>†</sup>	NLO
$\Gamma_5$	NLO	LO
$\Gamma_6$	LO	LO [New!]
$\Gamma_7$	LO	–
$\tilde{\Gamma}_6$	NLO	NLO*
$\tilde{\Gamma}_7$	LO	LO

	Source
$\langle \mathcal{O}_5 \rangle$	Heavy quark symmetry; Spectroscopy relations
$\langle \mathcal{O}_6 \rangle$	EOM relation to $\langle \tilde{\mathcal{O}}_6 \rangle$
$\langle \tilde{\mathcal{O}}_6 \rangle$	HQET Sum Rules [New!]
$\langle \tilde{\mathcal{O}}_7 \rangle$	VIA

\* Penguin operator contribution is still missing

† see talk by Matteo Fael on Wednesday

# Preliminary numerical results



Green - exp. values

Blue - th. predictions

Note: All quantities are normalised to the corresponding experimental values

# Conclusion and Outlook

- Update of *D*-meson lifetimes and their ratios incl.
  - ▷ First computation of the Darwin operator contribution in NL decays
  - ▷ New dim-6 Bag parameters for both *D*- and *D<sub>s</sub>*-mesons from HQET SR
- In agreement with data though huge theoretical uncertainties
- To improve theory predictions one needs
  - ▷ NNLO corrections at dimension-3
  - ▷ Dimension-6 Bag parameters (Lattice QCD) [M. Black, O. Witzel]
  - ▷ Dimension-7 Bag parameters (Lattice QCD) [M. Black, O. Witzel]
  - ▷ NNLO dimension-6 with four-quark operators [U. Nierste, M. Steinhauser (CRC project)]
  - ▷ NLO dimension-7 with four-quark operators [U. Nierste, M. Steinhauser (CRC project)]
  - ▷ Dimension-7 with two-quark operators
  - ▷ Dimension-8 with four-quark operators
  - ▷ NLO for dimension-5 and dimension-6 with two-quark operators
  - ▷ Experimental data on various moments in semileptonic decays to extract HQE parameters ( $\mu_\pi^2(D)$ ,  $\mu_G^2(D)$ ,  $\rho_D^3(D)$ , etc. ) [BESIII]

# Backup

# Definition of non-perturbative parameters

- Dimension-5 two-quark

$$2m_D \mu_\pi^2(D) = -\langle D(p_D) | \bar{c}_v(iD_\mu)(iD^\mu)c_v | D(p_D) \rangle,$$

$$2m_D \mu_G^2(D) = \langle D(p_D) | \bar{c}_v(iD_\mu)(iD_\nu)(-i\sigma^{\mu\nu})c_v | D(p_D) \rangle,$$

- Dimension-6 two-quark

$$2m_D \rho_D^3(D) = \langle D(p_D) | \bar{c}_v(iD_\mu)(iv \cdot D)(iD^\mu)c_v | D(p_D) \rangle,$$

$$2m_D \rho_{LS}^3(D) = \langle D(p_D) | \bar{c}_v(iD_\mu)(iv \cdot D)(iD_\nu)(-i\sigma^{\mu\nu})c_v | D(p_D) \rangle,$$

- Dimension-6 four-quark

$$\langle D_q | (\bar{c}_v \gamma_\mu (1 - \gamma_5) q) (\bar{q} \gamma^\mu (1 - \gamma_5) c_v) | D_q \rangle = f_{D_q}^2 M_{D_q}^2 B_1(D_q)$$

$$\langle D_q | (\bar{c}_v (1 - \gamma_5) q) (\bar{q} (1 + \gamma_5) c_v) | D_q \rangle = f_{D_q}^2 M_{D_q}^2 B_2(D_q)$$

$$\langle D_q | (\bar{c}_v \gamma_\mu (1 - \gamma_5) t^a q) (\bar{q} \gamma^\mu (1 - \gamma_5) t^a c_v) | D_q \rangle = f_{D_q}^2 M_{D_q}^2 \varepsilon_1(D_q)$$

$$\langle D_q | (\bar{c}_v (1 - \gamma_5) t^a q) (\bar{q} (1 + \gamma_5) t^a c_v) | D_q \rangle = f_{D_q}^2 M_{D_q}^2 \varepsilon_2(D_q)$$